

PART V

**what do we learn
from all that ?**

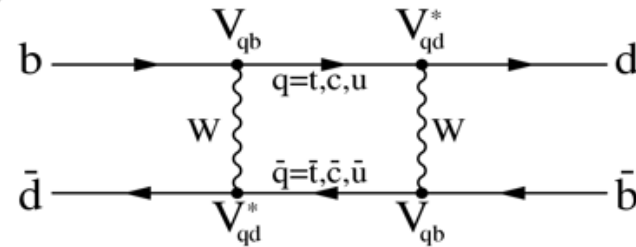
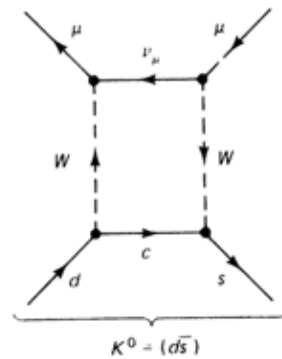
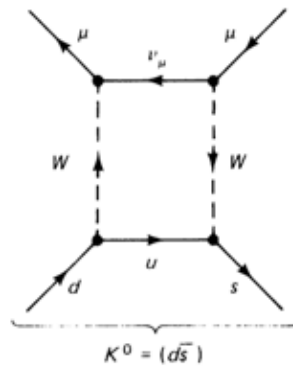
The quarks **c**, **b** and **t** were discovered in an indirect way by using **rare decays**, and more precisely **FCNC**, such as

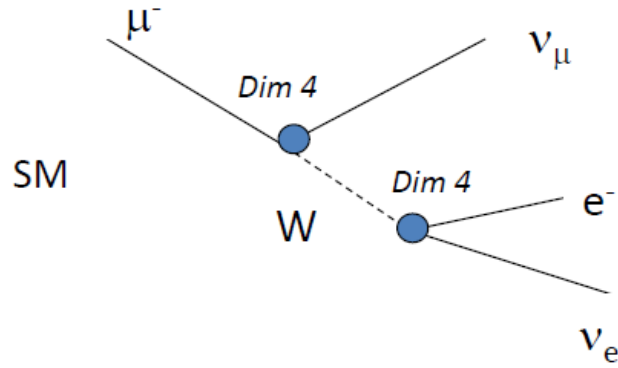
$$K^0 \rightarrow \mu\mu, \quad K_L \rightarrow \pi\pi \quad \text{and} \quad \mathbf{B} \text{ oscillations}$$

- ① ~1970 charm quark from FCNC and GIM-mechanism $K^0 \rightarrow \mu\mu$
- ② ~1973 3rd generation from CP violation in kaon (ϵ_K) KM-mechanism
- ③ ~1990 heavy top from B oscillations Δm_B

The Quantum path

*The indirect searches
look for “New Physics”
through virtual effects from new particles in loop
corrections*



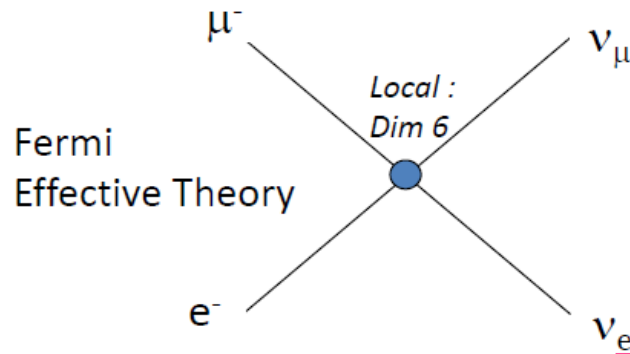


$$L_{CC} = \frac{g}{\sqrt{2}} (J_{\mu}^{-} W_{\mu}^{-} + J_{\mu}^{+} W_{\mu}^{+})$$

$$M = \left(\frac{g}{\sqrt{2}} \bar{u}_{\nu_{\mu}} \gamma^{\mu} \frac{1}{2} (1 - \gamma_5) u_{\mu} \right) \frac{1}{M_W^2 - q^2} \left(\frac{g}{\sqrt{2}} \bar{u}_e \gamma_{\mu} \frac{1}{2} (1 - \gamma_5) u_{\nu_e} \right)$$

if $q^2 \ll M_W^2$ (case of beta decay)

$$M = \frac{g^2}{8M_W^2} (\bar{u}_{\nu_{\mu}} \gamma^{\mu} (1 - \gamma_5) u_{\mu}) (\bar{u}_e \gamma_{\mu} (1 - \gamma_5) u_{\nu_e})$$



$$N = \frac{G_F}{\sqrt{2}} (\bar{u}_{\nu_{\mu}} \gamma^{\mu} (1 - \gamma_5) u_{\mu}) (\bar{u}_e \gamma_{\mu} (1 - \gamma_5) u_{\nu_e})$$

Experimentally from muon decay $G_F = 1.16 \times 10^{-5} \text{ GeV}^{-2}$

→

$$\frac{G_F}{\sqrt{2}} = \frac{g^2}{8M_W^2}$$

The effective coupling constant G_F is expressed as the SM « fundamental » coupling constant g divided by the mass of the propagator M_W squared (consequence of Dim=6 operators [4legs])

In this specific case we know

- from SM $e = g \sin(\theta_W)$
- Experimentally $M_W \sim 80 \text{ GeV}$

The weak interaction is not weak because of $g \ll e$ but because of the large value for the W mass

Effective Flavour Theory to New Flavour Physics :

A game of scale and coupling

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_{k=1} \left(\sum_i C_i^k Q_i^{(k+4)} \right) / \Lambda^k$$

New operators which are of dimension >4, in principle the theory is not Renormalizable...as Fermi theory was not..!

[You can show that in B physics the new operators have dimension 6]

NP flavour effects are governed by two players

→ the value of the new physics scale Λ

→ the effective flavour-violating coupling C 's

In explicit models

$\Lambda \sim$ mass of virtual particle

(Fermi theory : M_W)

$C \sim$ loop coupling \times flavour coupling

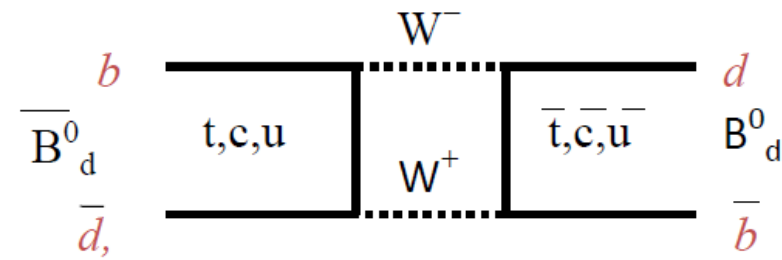
(SM/MFV $\alpha_w \times$ CKM)

Example for B oscillations (FCNC- $\Delta B=2$)

FCNC processes are ideal place to look for NP effects because they are suppressed in SM

Precise measurements are needed. Effects goes $1/\Lambda^2$

SM

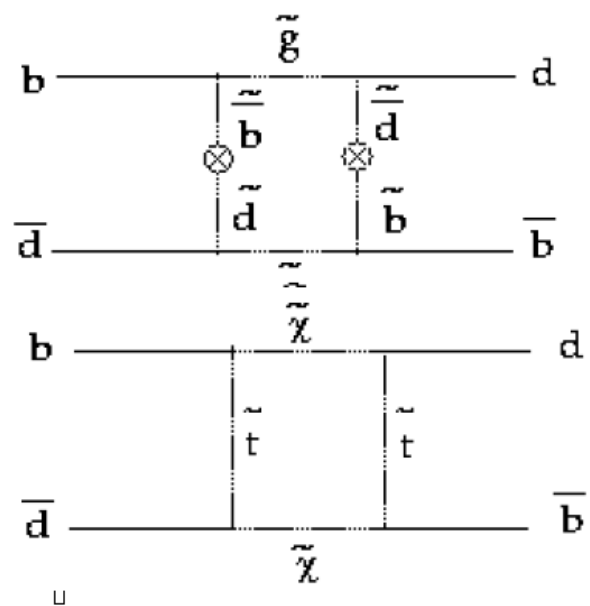


$$f(m) \times \frac{|V_{tb}^* V_{tq}|^2}{M_W^2}$$

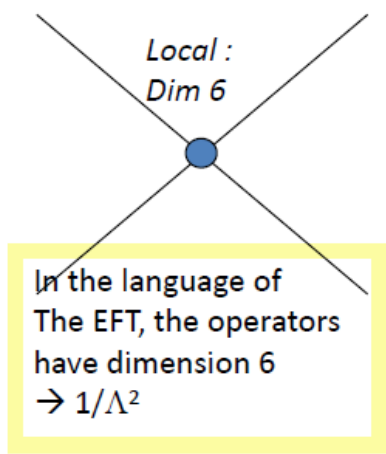
GIM mechanism

When I squared I have the m_t^2 at the numerator and M_W^4 at denominator. IMPORTANT at the end the dimension is $1/E^2$

BSM



$$\frac{|\delta_{bq}|^2}{\Lambda_{eff}^2}$$



In the language of The EFT, the operators have dimension 6 $\rightarrow 1/\Lambda^2$

The measurements (in this case Δm_d) are modified wrt the predictions of the SM by the presence of BSM particles.

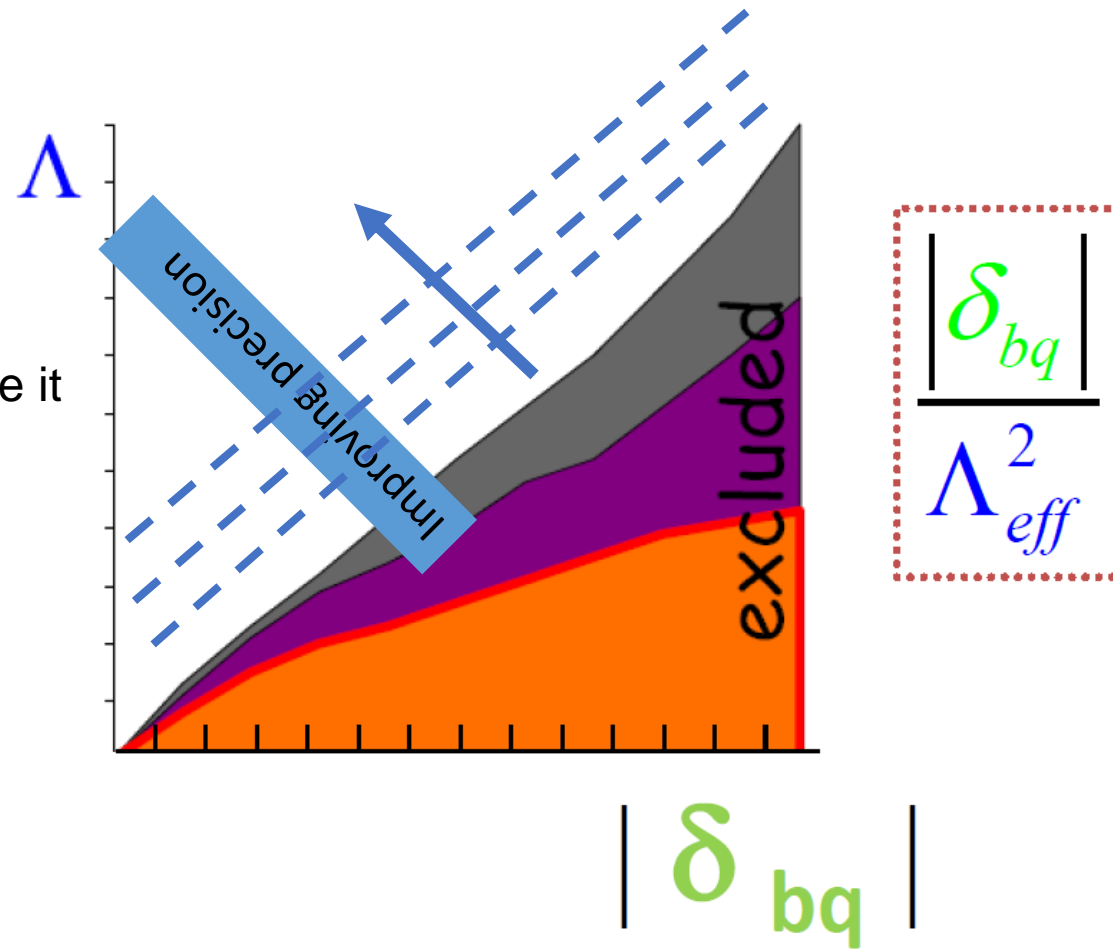
modifications are important if couplings are larger and/or NP masses are lighter

Pictorially

Flavour Physics

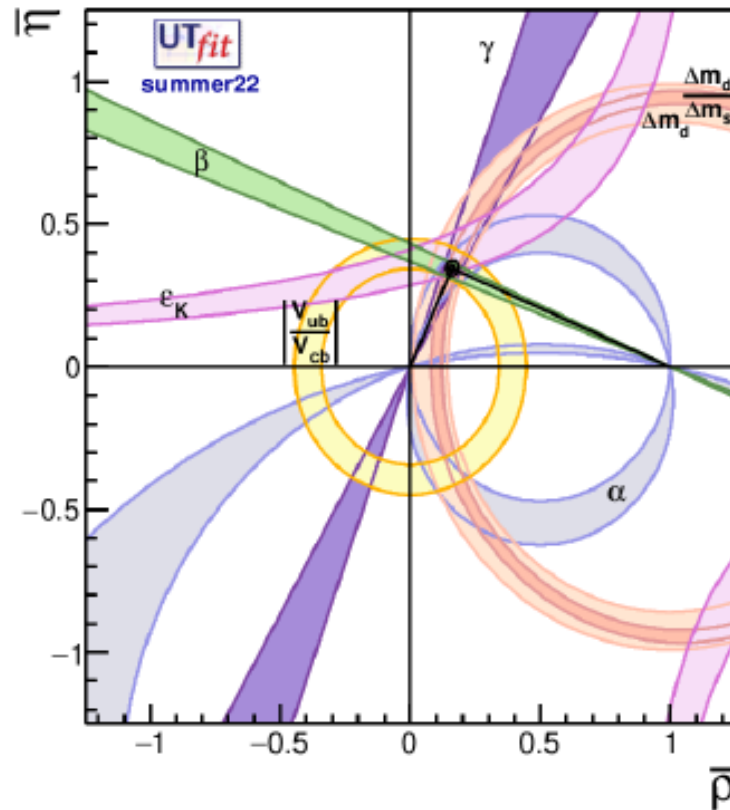
It is a game of couplings and scales

This is the basic region because it is important to have precise measurements



Global Fit within the SM

All the constraints
Look compatibles !



*Coherent picture of
FCNC and CPV
processes in SM*

CKM matrix is the dominant source of
flavour mixing and CP violation

Out of these measurement there a general agreement that we have limited
the contributions of New Physics amplitudes (A_{NP}) wrt to SM ones (A_{SM})
at the the level of

$$R = \frac{A_{NP}}{A_{SM}} < 20\%$$

What does it imply ?

UT analysis including new physics

fit simultaneously for the CKM and the NP parameters (generalized UT fit)

- add most general loop NP to all sectors
- use all available experimental info
- find out NP contributions to $\Delta F=2$ transitions

B_d and B_s mixing amplitudes
(2+2 real parameters):

$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}} = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})} \right) A_q^{SM} e^{2i\phi_q^{SM}}$$

$$\Delta m_{q/K} = C_{B_q/\Delta m_K} (\Delta m_{q/K})^{SM}$$

$$A_{CP}^{B_d \rightarrow J/\psi K_S} = \sin 2(\beta + \phi_{B_d})$$

$$A_{SL}^q = \text{Im}(\Gamma_{12}^q/A_q)$$

$$\varepsilon_K = C_\varepsilon \varepsilon_K^{SM}$$

$$A_{CP}^{B_s \rightarrow J/\psi \phi} \sim \sin 2(-\beta_s + \phi_{B_s})$$

$$\Delta \Gamma^q / \Delta m_q = \text{Re}(\Gamma_{12}^q/A_q)$$

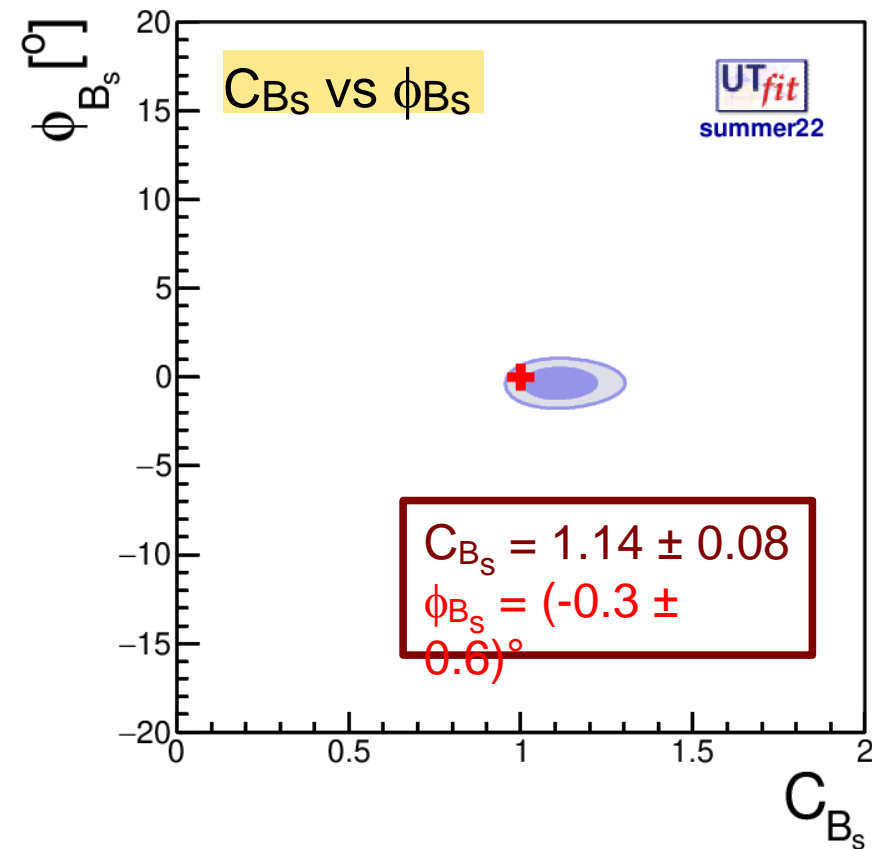
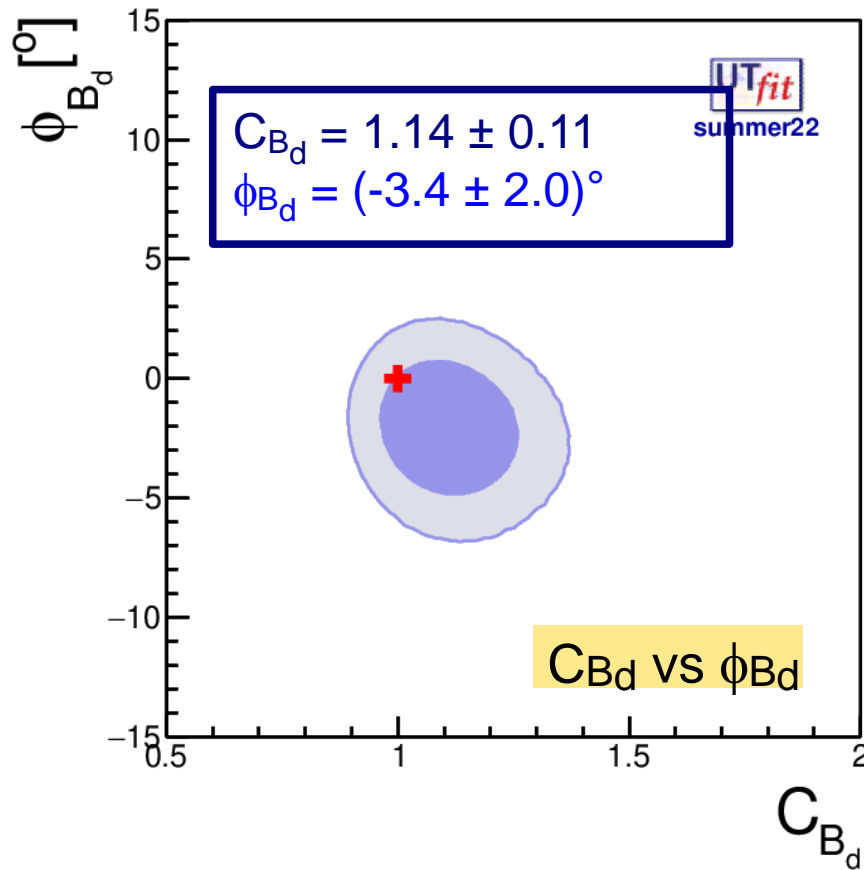
NP parameter results

$$A_q = C_{Bq} e^{2i\phi_{Bq}} A_q^{SM} e^{2i\phi_q^{SM}}$$

dark: 68%
light: 95%
SM: red cross

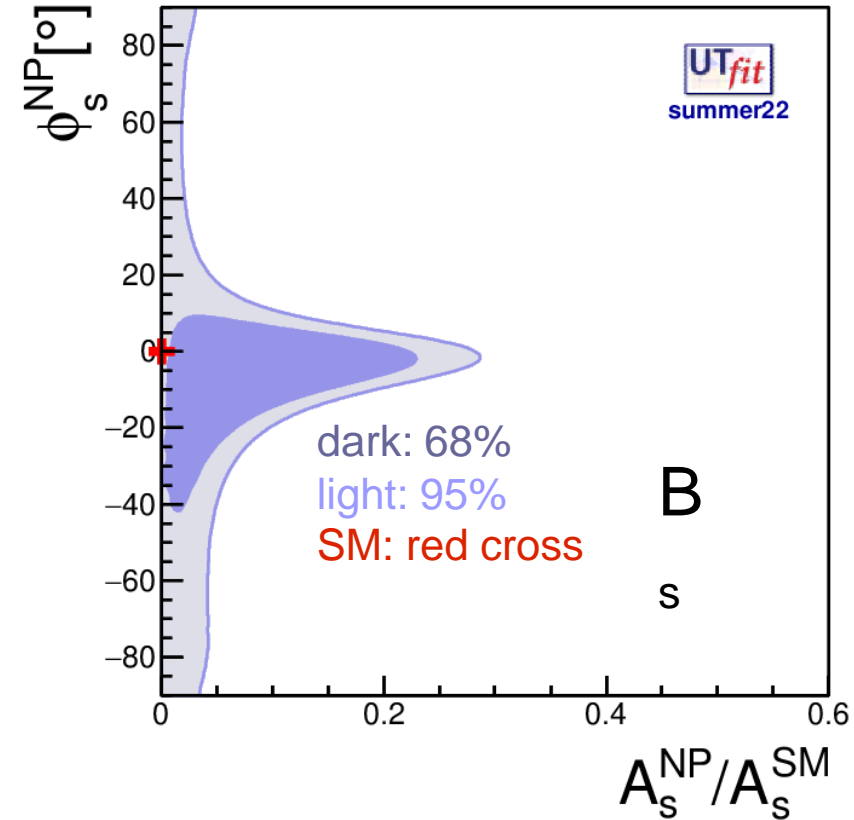
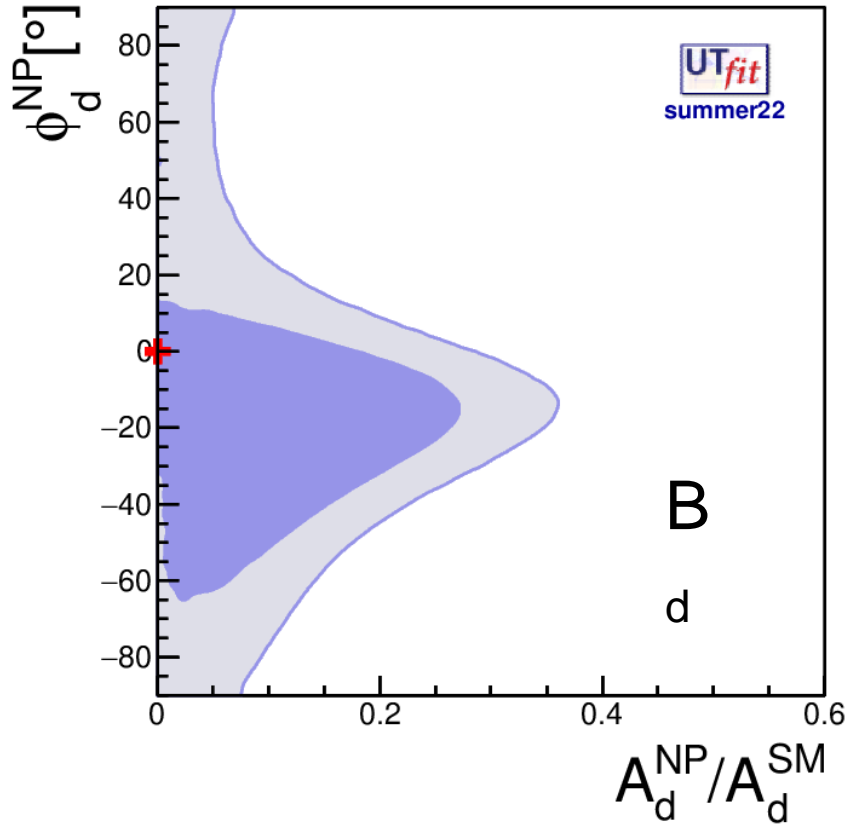
K system

$$C_{eK} = 1.12 \pm 0.12$$



NP parameter results

$$A_q = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})} \right) A_q^{SM} e^{2i\phi_q^{SM}}$$

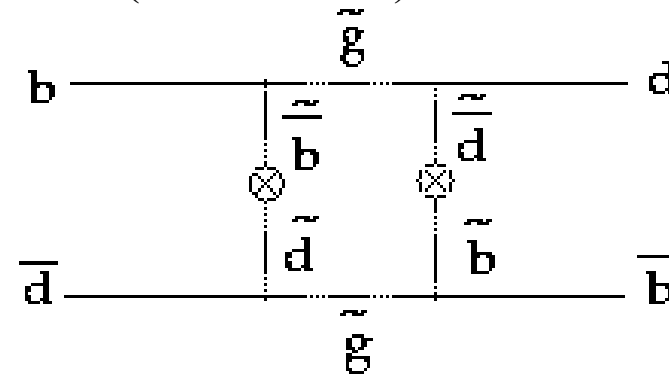
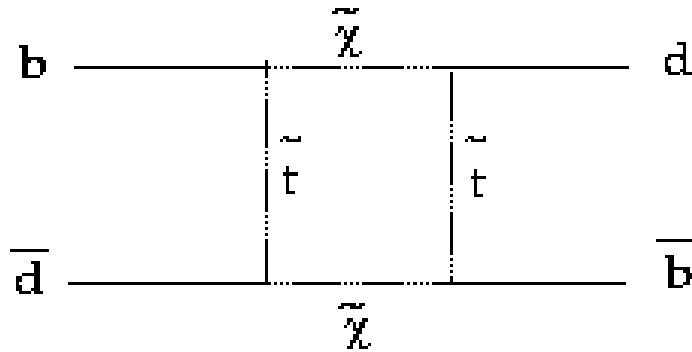


The ratio of NP/SM amplitudes is:

< 25% @68% prob. (35% @95%) in B_d mixing

< 25% @68% prob. (30% @95%) in B_s mixing

Simplify the discussion with one example : B oscillations (FCNC- $\Delta B=2$) :



$$\left| \frac{Q_{\Delta B=2}^{NP}}{Q_{\Delta B=2}^{SM}} \right| \leq r \rightarrow \frac{|\delta_{bq}|}{\Lambda_{eff}} \leq \sqrt{r} \frac{|V_{tb}^* V_{tq}|}{M_W}$$

r upper limit of the relative NP contribution
 δ_{bd} NP physics coupling
 Λ_{eff} NP scale (masses of new particles)

If couplings ~ 1

$$\delta_{bq} \sim 1 \quad \Lambda_{eff} \sim 10/\sqrt{r} \text{ TeV}$$

$$\delta_{bs} \sim 1 \quad \Lambda_{eff} \sim 2/\sqrt{r} \text{ TeV}$$

Minimal Flavour Violation

\equiv no new sources of flavour and CP violation NP contributions governed by SM Yukawa couplings.

$$\delta_{q'd} \approx V_{tq'}^* V_{td}$$

(couplings small as CKM elements)

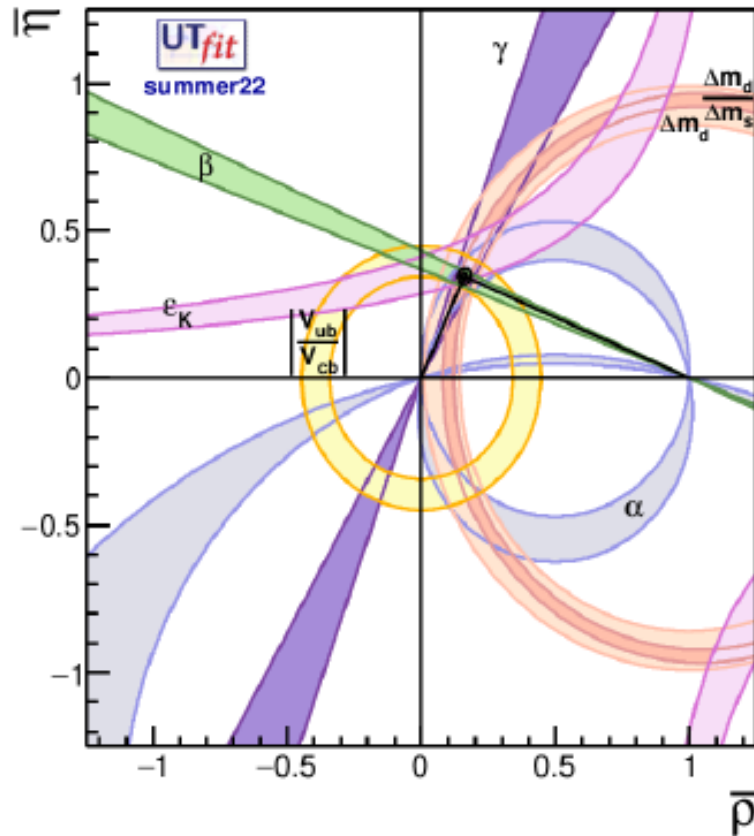
$$\Lambda_{eff} \sim 0.08/\sqrt{r} \text{ TeV}$$

Coupling	r=20%
δ	today
Order 1	$\Lambda_{eff} \sim 20 \text{ TeV}$
MFV	$\Lambda_{eff} \sim 180 \text{ GeV}$

$r < 0.2 \rightarrow \Lambda_{eff} > 180 \text{ GeV}$. Particles below 180 GeV circulating in the loop would have given visible effects within the present level of precision

The test of the SM (in fermion sector)

1990-now → a huge number of precise measurements



Discovery : absence of New Particles up to the $\sim 2 \times$ Electroweak Scale !

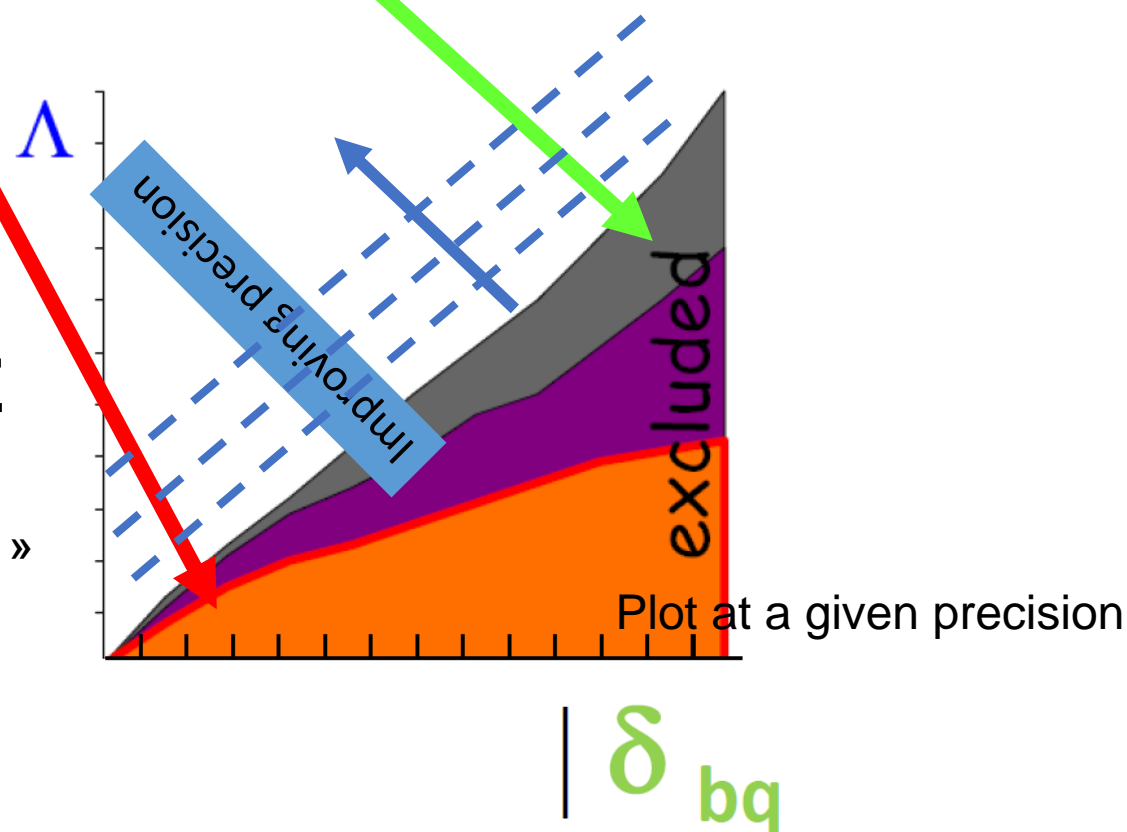
....and at much larger scale ($> \text{TeV}$) for some New physics model

- 1 ~1970 charm quark from FCNC and GIM-mechanism $K^0 \rightarrow \mu\mu$
- 2 ~1973 3rd generation from CP violation in kaon (ϵ_K) KM-mechanism
- 3 ~1990 heavy top from B oscillations Δm_B
- 4 >2010 success of the description of FCNC and CPV in SM

Coupling δ	r=20% today	r=10% tomorrow	r=1% after tomorrow
Order 1	$\Lambda_{\text{eff}} \sim 20 \text{ TeV}$	$\Lambda_{\text{eff}} \sim 30 \text{ TeV}$	$\Lambda_{\text{eff}} \sim 100 \text{ TeV}$
MFV	$\Lambda_{\text{eff}} \sim 180 \text{ GeV}$	$\Lambda_{\text{eff}} \sim 250 \text{ GeV}$	$\Lambda_{\text{eff}} \sim 800 \text{ GeV}$

FLAVOUR PHYSICS IS CENTRAL IN THE NP SEARCH

because ALLOW TO « DISCOVER »
VERY HEAVY PARTICLE NOT
ACCESSIBLE TO DIRECT
SEARCH DEPENDING ON
COUPLINGS

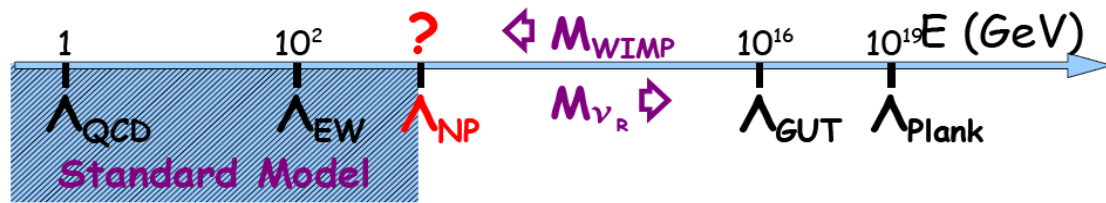


FLAVOUR PHYSICS IS AN EXPLORATORY PHYSICS !
IT ALLOWS TO EXPLORE SCALE FAR BEYOND THE
REACH OF THE DIRECT SEARCH

SPAN FROM ELECTROWEAK SCALE TO
SEVERAL TEV (up to 100TEV) !

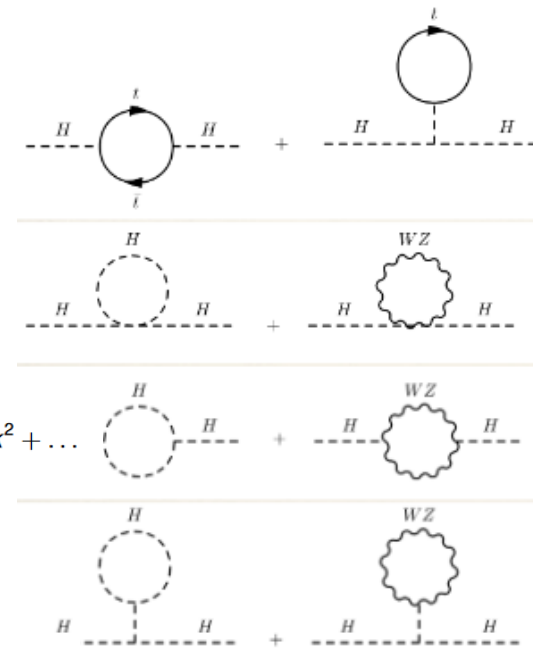
**Where do we expect
the scale of New Physics??
This IS THE QUESTION (for me) !**

Rephrasing : The problem of particle physics today is :
 where is the NP scale $\Lambda \sim 0.5, 1 \dots 10^{16}$ TeV



$$m_H^2 \rightarrow m_{\text{bare}}^2 + \delta m_H^2$$

$$M_H^2(p^2) = M_H^2(\Lambda^2) + Cg^2 \int_{p^2}^{\Lambda^2} dk^2 + \dots$$



$$\delta M_H^2 = \frac{G_F \Lambda^2}{4\pi^2 \sqrt{2}} (6 M_W^2 + 3 M_Z^2 + M_H^2 - 12 m_t^2) \approx -\frac{3 G_F}{\pi^2 \sqrt{2}} m_t^2 \Lambda^2 \approx -0.075 \Lambda^2.$$

The corrections to the **Higgs mass are quadratically divergent*** and Λ is the Energy cut-off to the divergent loop integrals. (* ΔM_H increases depending of the choosed cut off Λ .)

Λ is the energy scale where we decide to cut off..

Λ Set the scale of the « validity » of the theory

hierarchy problem

Stabilizing the mass $M_H = 125$ GeV requires
 DELICATE BALANCE of TWO NUMBERS.

$\Lambda = 1 \text{ TeV}$
 $M_H^2(p^2) = (125 \text{ GeV})^2 = 1.56 \cdot 10^4 \text{ GeV}^2 = M_H^2(\Lambda^2) - 7.5 \cdot 10^4 \text{ GeV}^2$

$\Lambda = 10 \text{ TeV}$
 $M_H^2(p^2) = (125 \text{ GeV})^2 = 1.56 \cdot 10^4 \text{ GeV}^2 = M_H^2(\Lambda^2) - 7.5 \cdot 10^6 \text{ GeV}^2$

This delicate balance become more and more delicate as much as Λ increase ... up to the Plank scale

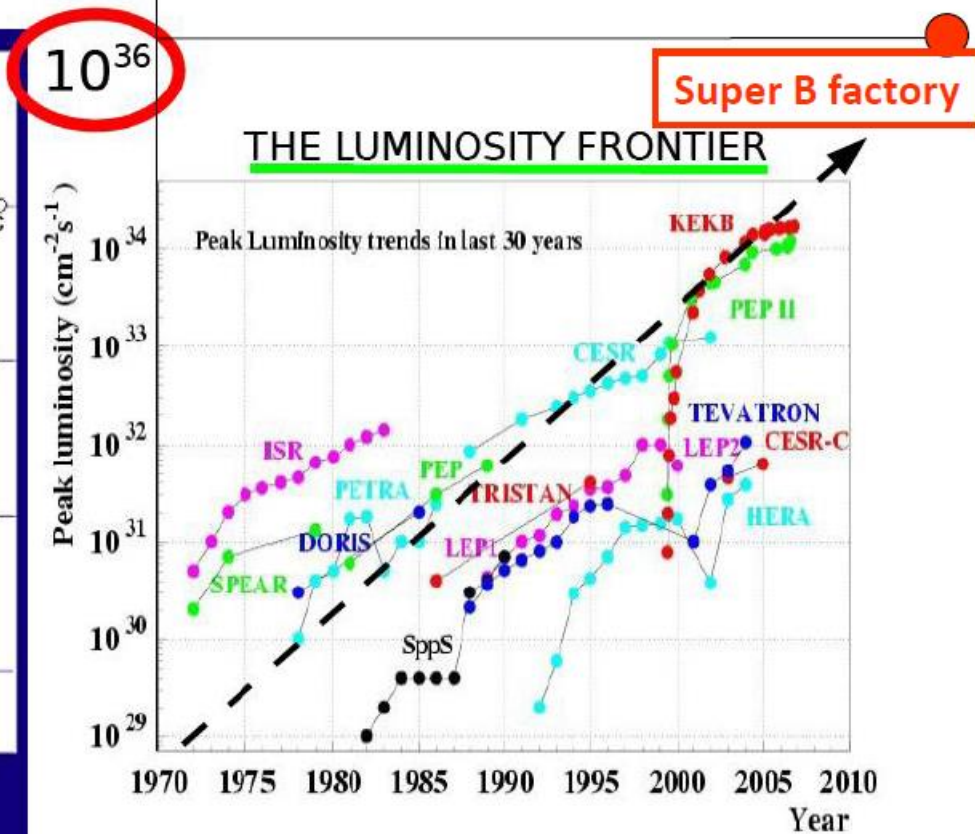
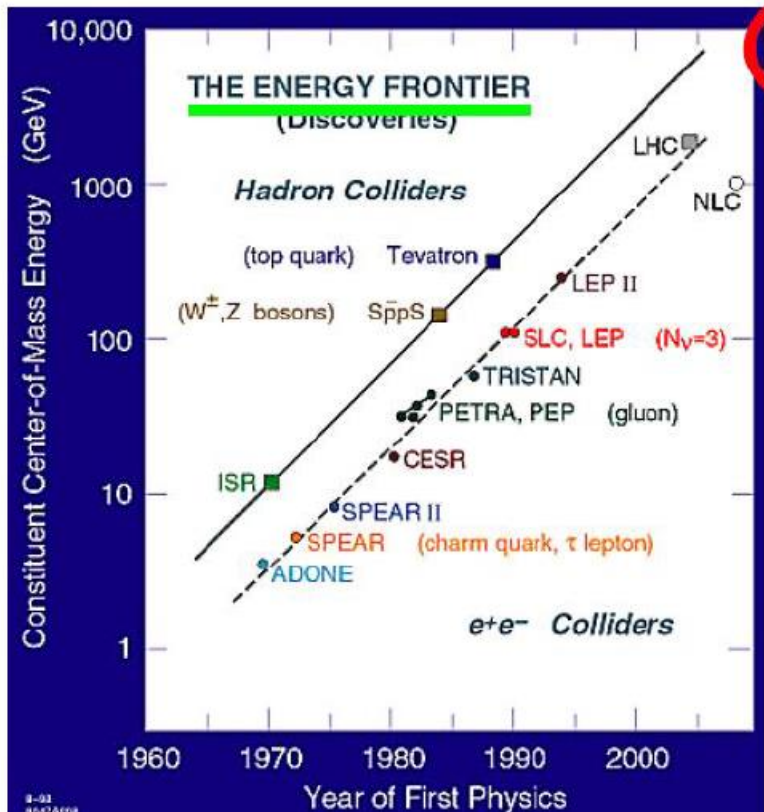
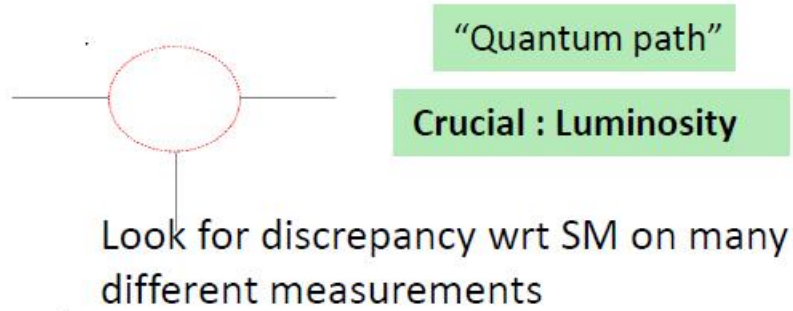
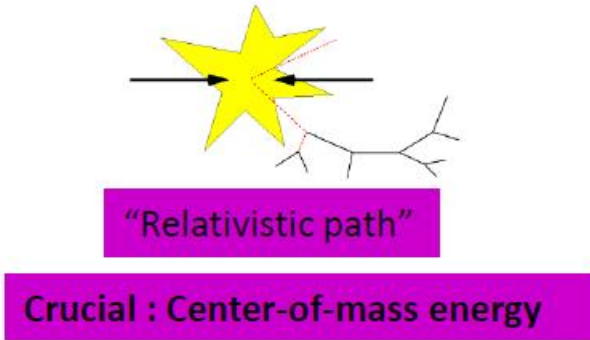
MH(bare) 300 GeV

MH(bare) 2,5 TeV

Another way of saying , **we need new particles in the loop to regularise this divergence**

Explore New Energy scales

...Indeed historically we have always followed the two ways...



10^{36}