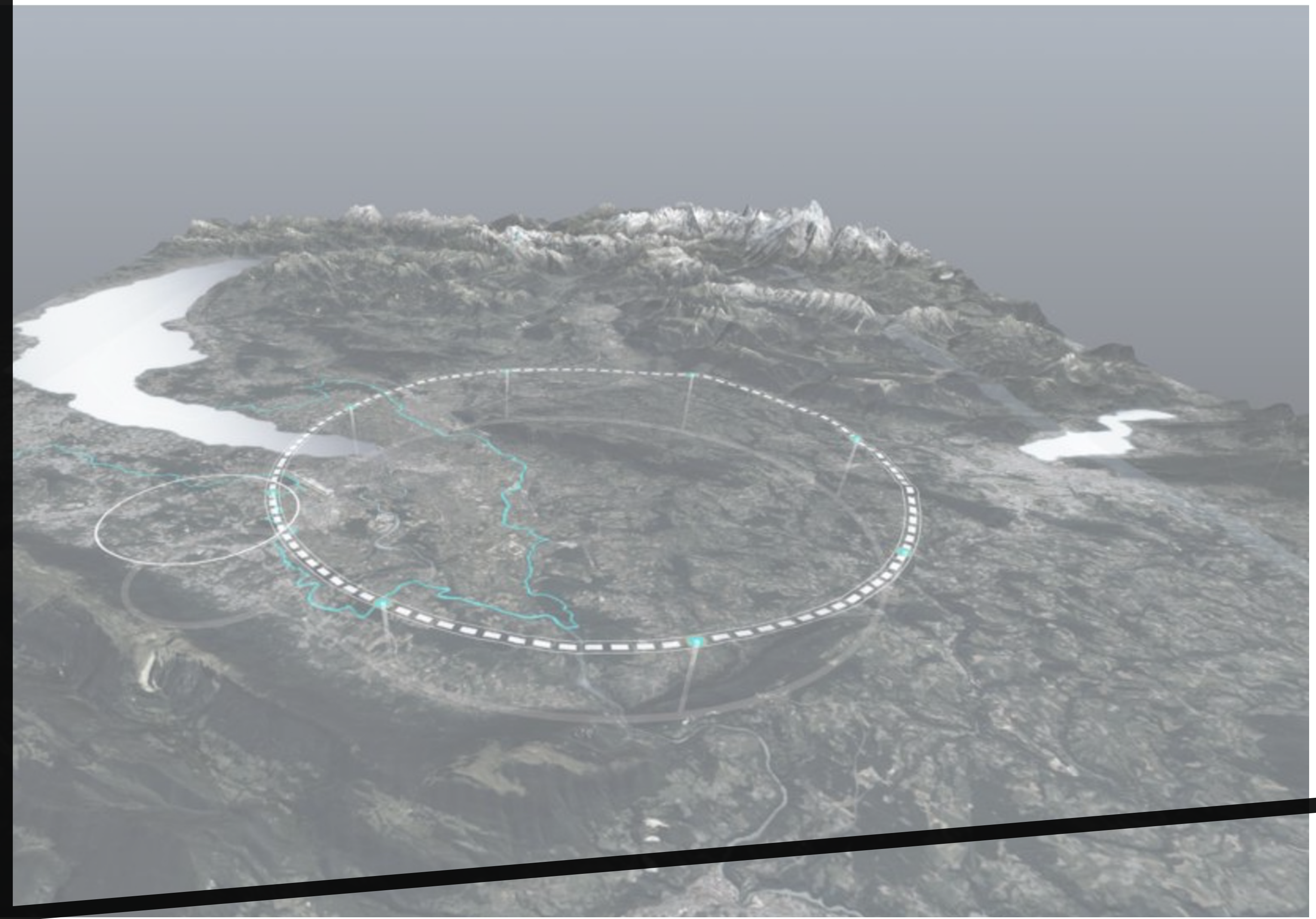
Future Circular Collider
— Physics Case —

CERN, Feb. 13th 2024



Christophe Grojean

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Humboldt University (Berlin)
CERN
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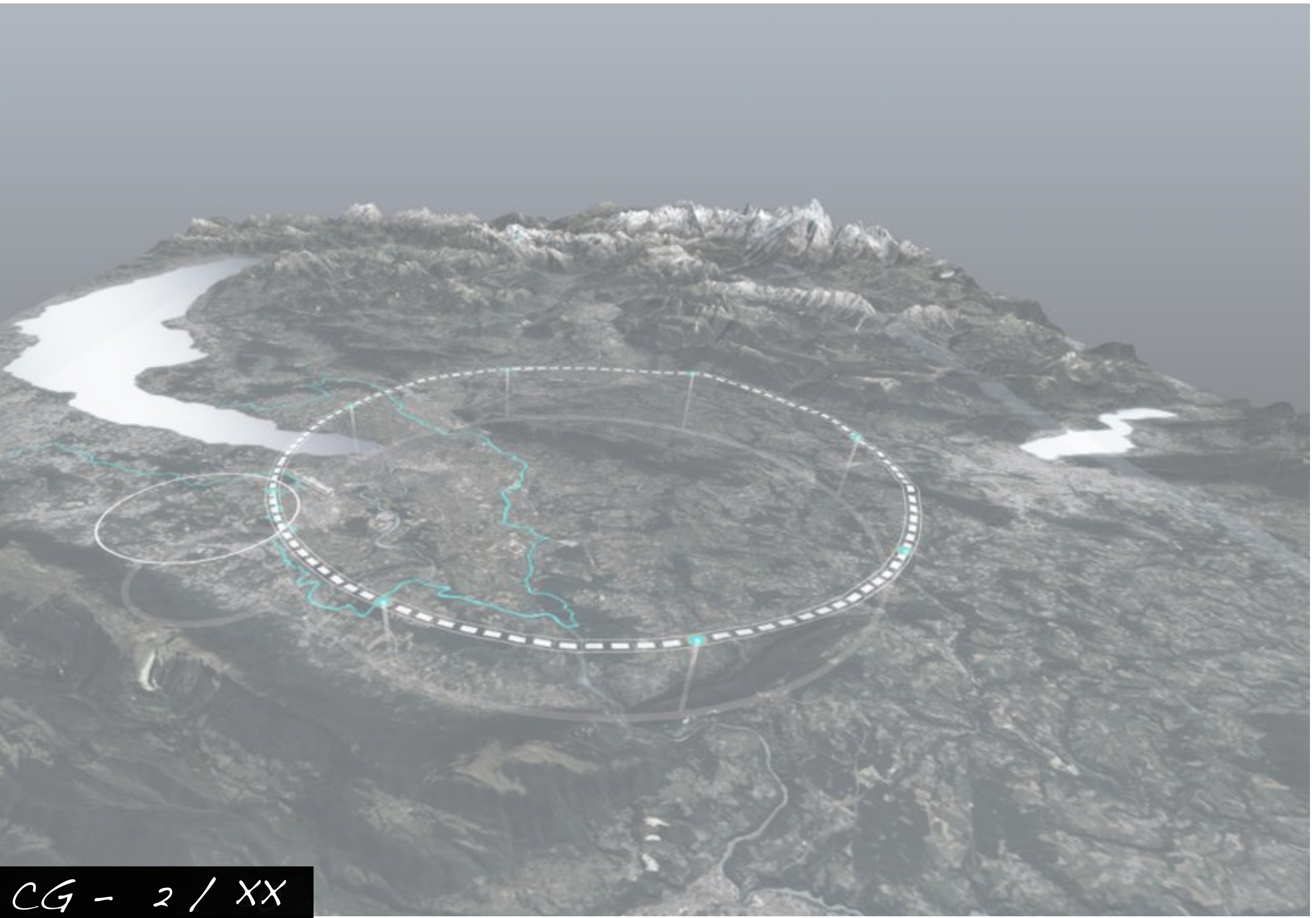
— on behalf of the FCC team—

Future Circular Collider

- A versatile particle collider housed in a 91km underground ring
- Implemented in several stages:
  - an e^+e^- “Higgs/EW/Flavour/top/QCD” factory running at 90-365 GeV ▶ **FCC-ee**
  - followed by a high-energy pp collider reaching 100 TeV ▶ **FCC-hh**



Outline

1. Why do we need a new collider?
2. FCC-ee: much more than a Higgs factory
3. FCC-ee/hh as a Higgs/electroweak factory
4. FCC-ee as a flavour factory
5. FCC-hh: the broadest exploration potential at high-energy
6. FCC-ee↔FCC-hh: complementarity and synergy

1.
Why do we need a new collider?

The LHC Legacy (so far).

- **Standard Model (SM) confirmed to high accuracy up to energies of several TeV**
  (thanks to a firm control of exp. & th. syst. uncertainties, the LHC became a precision machine)

- **Higgs boson discovered** at the mass predicted* by LEP precision EW measurements
  *within the Standard Model

- **Absence of new physics** ⟶
  - Traditional New Physics models are under siege
  - New approaches: relaxion, Nnaturalness, clockwork…

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We need a broad, versatile and ambitious programme that
1. sharpens our knowledge of already discovered physics
2. pushes the frontiers of the unknown at high and low scales
— together FCC-ee & FCC-hh combine these 2 aspects —
more PRECISION and more ENERGY, for more SENSITIVITY to New Physics

Precision as a discovery tool.

Many historical examples

- Uranus anomalous trajectory ⇢ Neptune
- Mercury perihelion ⇢ General Relativity
- Z/W interactions to quarks and leptons ⇢ Higgs boson
- …

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(e.g. Rayleigh-Jeans UV catastrophe for QM, unitarity breakdown for the Higgs)
but precision gave valuable additional clues.
In any case, experimentalists shouldn't lean too heavily on theorist priors/prejudices
(remember discovery of CP violation).
At times when we don't have a precise theoretical guidance that we need powerful experimental tools to make progress.

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The Standard Model structure has seemingly "accidental" aspects (e.g. B, L number conservations) that should be probed to form a deeper understanding of Nature.

The Higgs requires more precision.

"The Higgs isn't everything !"

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"The Higgs isn't everything; it's the only thing!"*

The scalar discovery in 2012 has been an important milestone for HEP.

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lifetime of stars

(why $t_{Sun} \sim t_{life\ evolution}$?)

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$m_e, m_u, m_d \leftrightarrow$ Higgs couplings ?

m_e → size of atoms; m_u, m_d → nuclei stability

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m_e, m_u, m_d ↔ Higgs couplings ?
size of atoms (m_e) — nuclei stability (m_u, m_d)

EWSB @ $t \sim 10^{-10}$s ↔ ? Higgs self-coupling(s) / Higgs(es) potential

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matter/anti-matter ↔ ? CPV in Higgs sector

The Higgs requires more precision.

(HL)-LHC will make remarkable progress.
But it won't be enough.
A new collider is needed!

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The launch of the feasibility study.

"An **electron-positron** Higgs factory is the highest-priority next collider. For the longer term, the European particle physics community has the ambition to operate a **proton-proton** collider at the highest achievable energy."

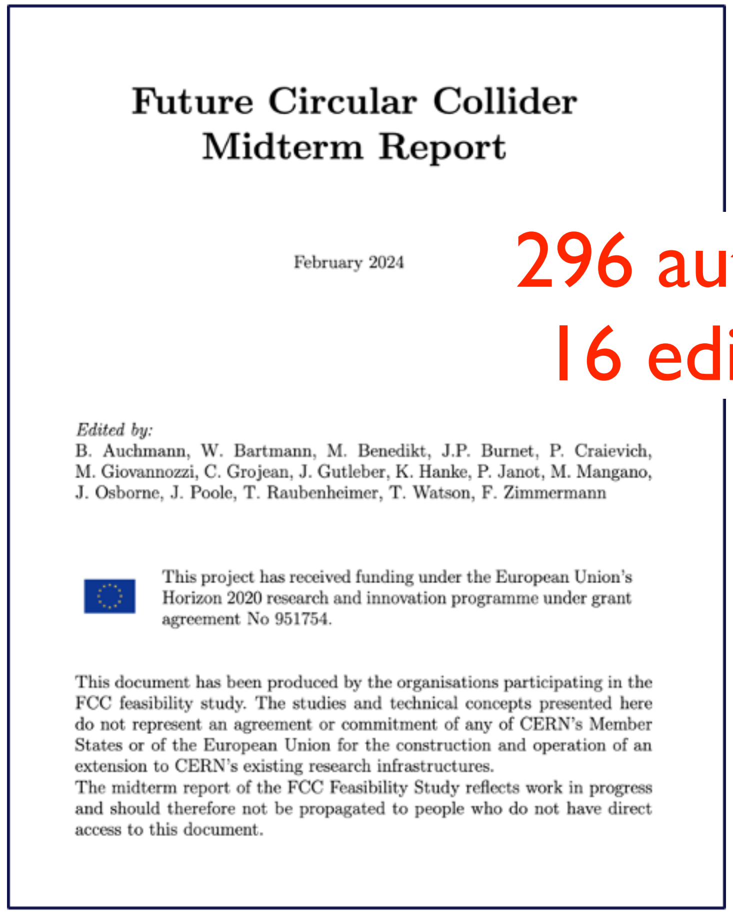
— CERN council approved the Strategy and CERN management implemented it —

FCC Feasibility Study (FS) started in 2021 and will be completed in 2025.

Mid-term review in 2023.

FCC feasibility mid-term report.

- **703 pages**: 7 chapters (cost and financial feasibility is a separate document) + refs.
  - Placement scenario (75 pages)
  - Civil engineering (50 pages)
  - Implementation with the host states (45 pages)
  - Technical infrastructure (110 pages)
  - FCC-ee collider design and performance (170 pages)
  - FCC-hh accelerator (60 pages)
  - (Cost and financial feasibility)
  - Physics and experiments (110 pages)
  - References (70 pages)
- **Executive summary**: 44 pages
- Reviewed by
  - Scientific Advisory Committee and Cost Review Panel on Oct. 16-18
  - Scientific Policy Committee and Financial Committee on Nov. 21-22
  - CERN Council Feb. 2


Future Circular Collider
Midterm Report

February 2024

Edited by:
B. Auchmann, W. Bartmann, M. Benedikt, J.P. Burnet, P. Craievich, M. Giovannozzi, C. Grojean, J. Gutleber, K. Hanke, P. Janot, M. Mangano, J. Osborne, J. Poole, T. Raubenheimer, T. Watson, F. Zimmermann

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This document has been produced by the organisations participating in the FCC feasibility study. The studies and technical concepts presented here do not represent an agreement or commitment of any of CERN's Member States or of the European Union for the construction and operation of an extension to CERN's existing research infrastructures.
The midterm report of the FCC Feasibility Study reflects work in progress and should therefore not be propagated to people who do not have direct access to this document.

296 authors
16 editors

confidential documents (work in progress) available to CERN personnel

Physics, Experiments, Detectors.

- FCC Feasibility Study PED deliverables for mid-term review

8. Physics & Experiments	C. Grojean, P. Janot, M. Mangano	8.1 Overview	
		8.2. Documentation of the specificities of the FCC-ee and FCC-hh physics cases.	deliverables explicitly requested from SPC & Council
		8.3 Strategic plans for the improved theoretical calculations.	
		8.4 FCC-ee Detector Requirements.	

- Content of the mid-term PED chapter (60 pages were expected → 110 pages delivered)

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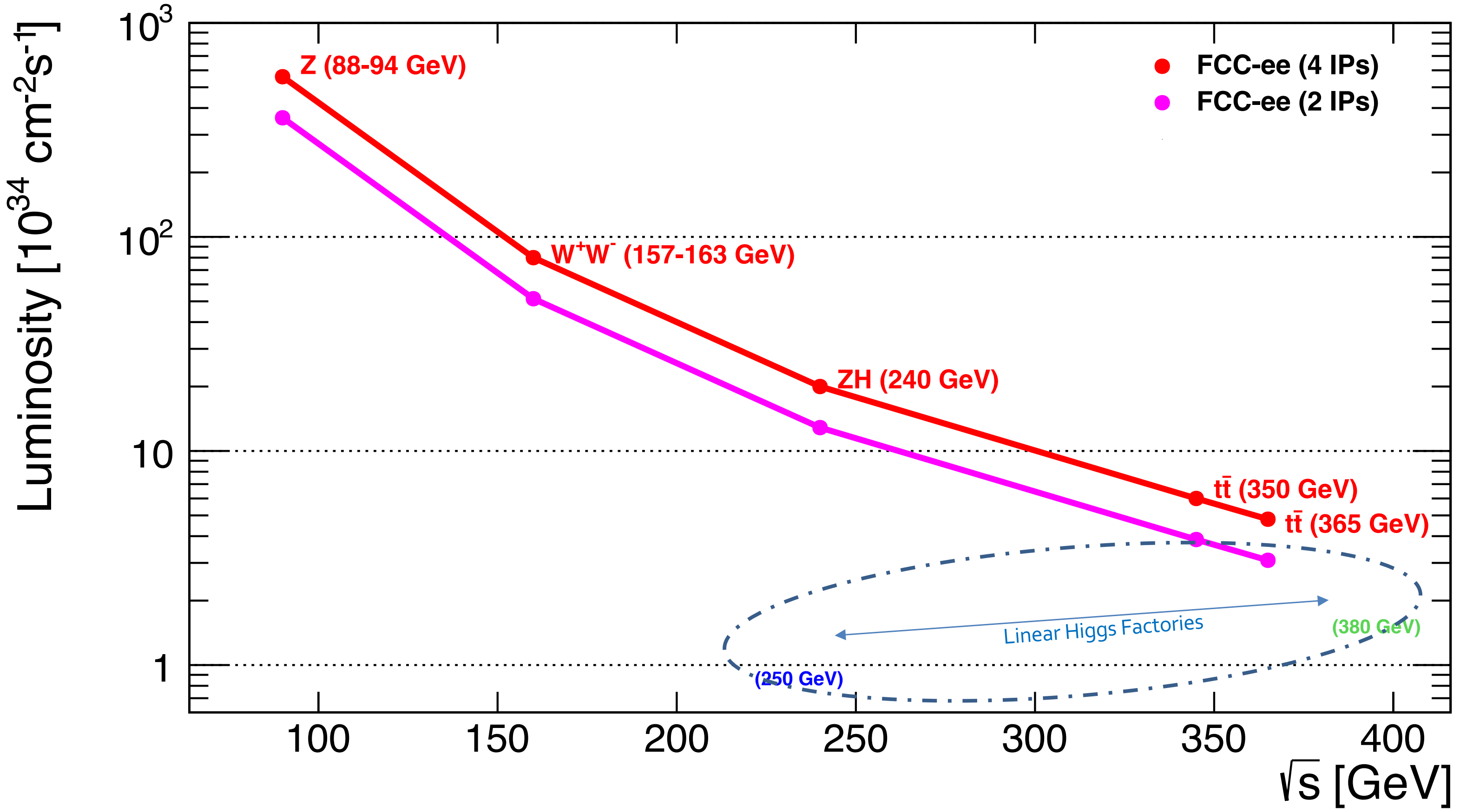
I am not going to present the details of this report.
The limited time of this talk wouldn't do justice to all the work accomplished,
and my personal limitations wouldn't allow me to touch upon the breadth of the studies.
I limit myself to a few highlights

2.
FCC-ee: much more than a Higgs factory

FCC-ee Run Plan.

LEP1 data accumulated in **every 2 mn**. Exciting & diverse programme with different priorities every few years.

(order of the different stages still subject to discussion/optimisation)



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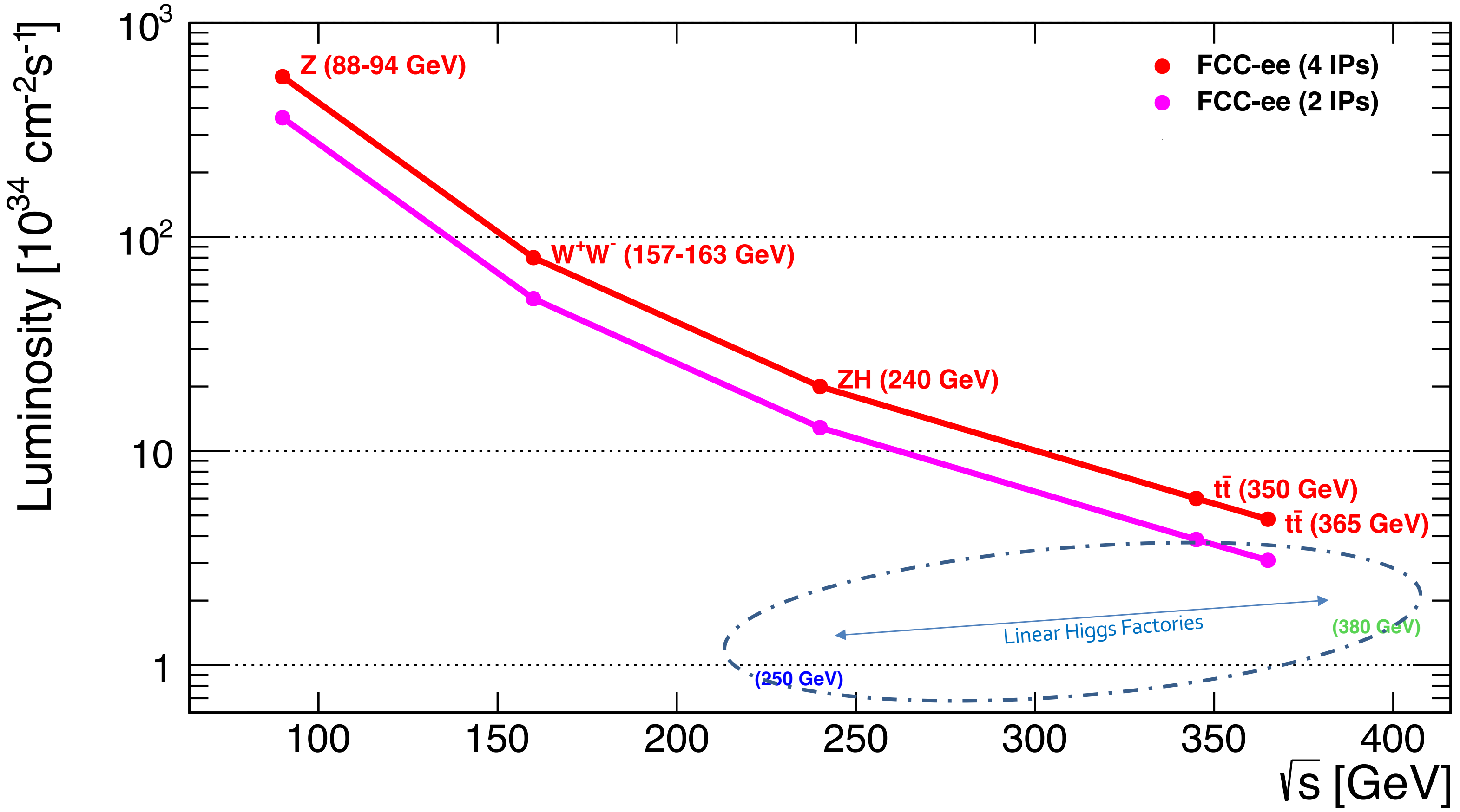

Working point	Z, years 1-2	Z, later	WW, years 1-2	WW, later	ZH	$t\bar{t}$	
$\sqrt{s}$ (GeV)	88, 91, 94		157, 163		240	340–350	365
Lumi/IP ($10^{34}\,\mathrm{cm}^{-2}\mathrm{s}^{-1}$)	70	140	10	20	5.0	0.75	1.20
Lumi/year (ab^{-1})	34	68	4.8	9.6	2.4	0.36	0.58
Run time (year)	2	2	2	–	3	1	4
Number of events	6×10^{12} Z		2.4×10^8 WW		1.45×10^6 ZH + 45k WW $\rightarrow$ H	1.9×10^6 $t\bar{t}$ +330k ZH +80k WW $\rightarrow$ H	

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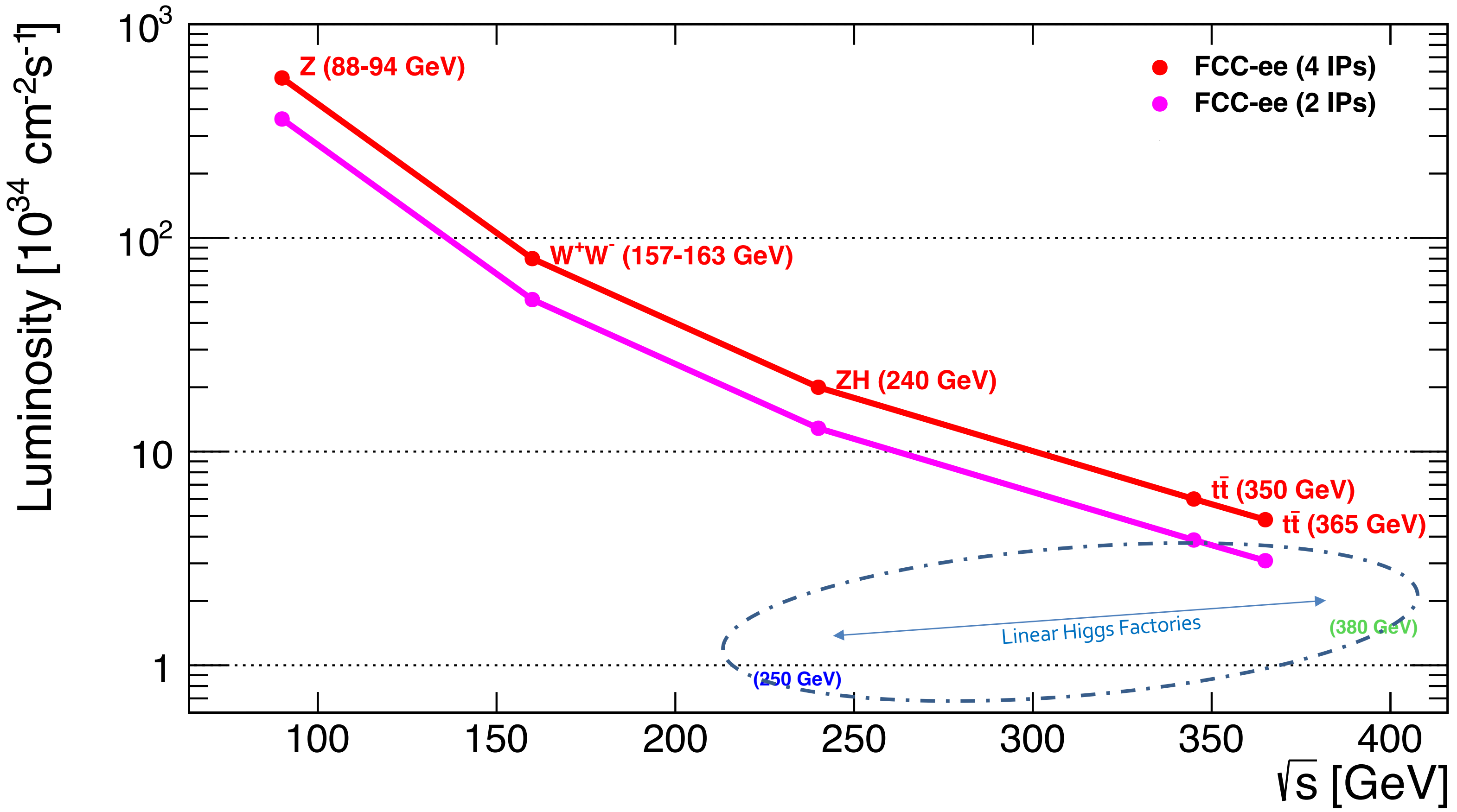

— Superb statistics achieved in only 15 years —

in each detector:
10^5 Z/sec, 10^4 W/hour,
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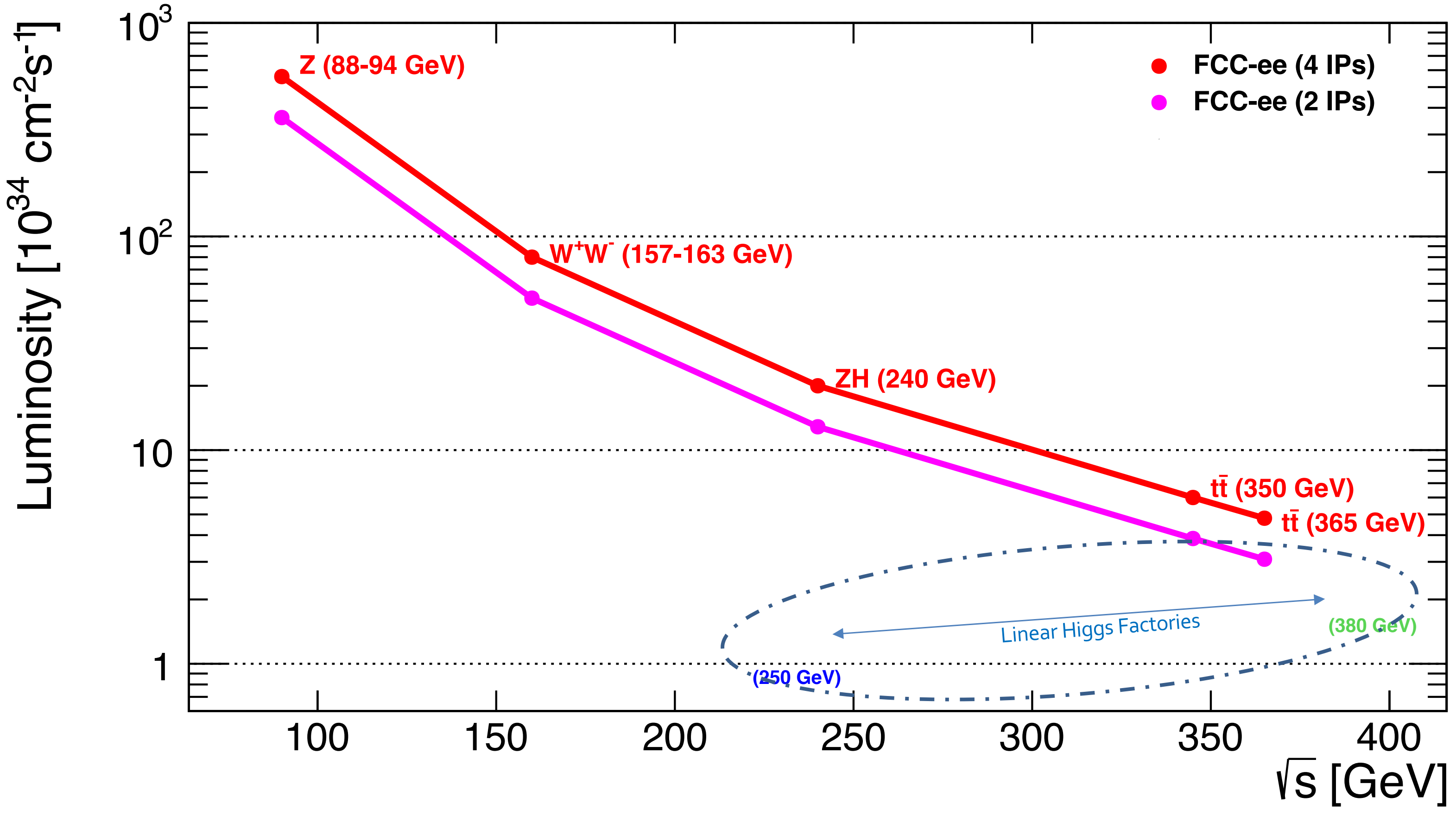

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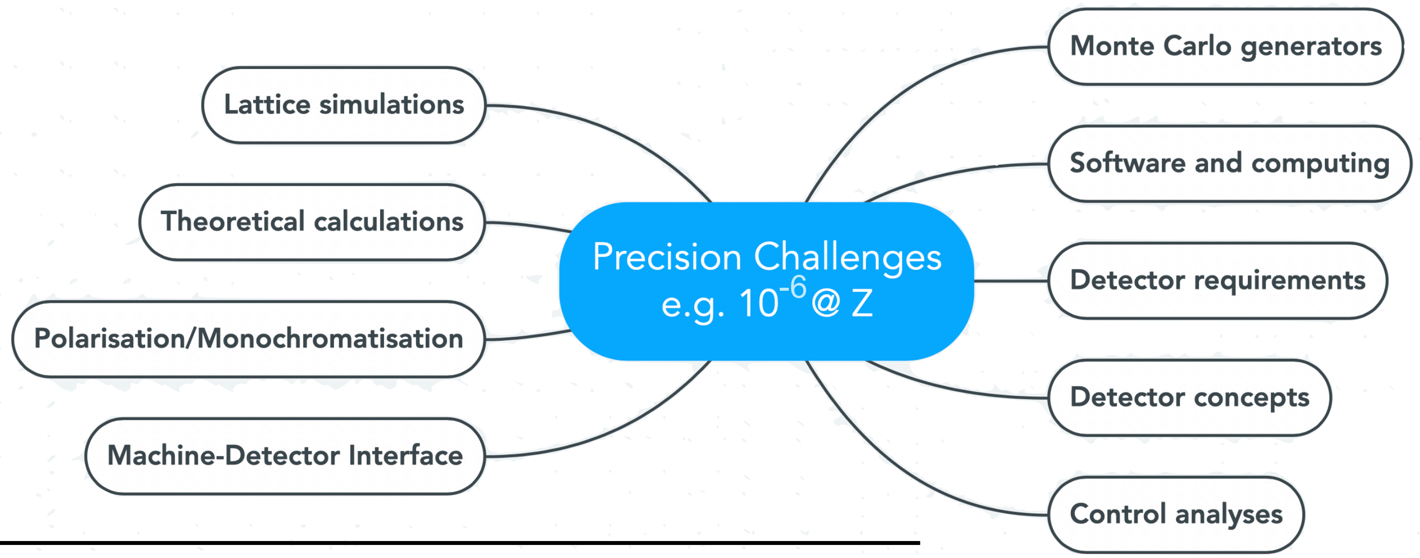

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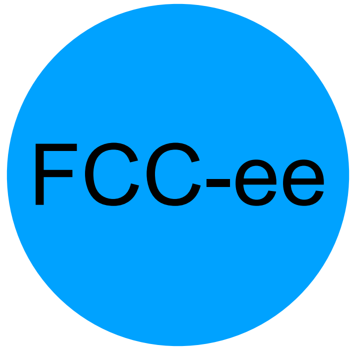

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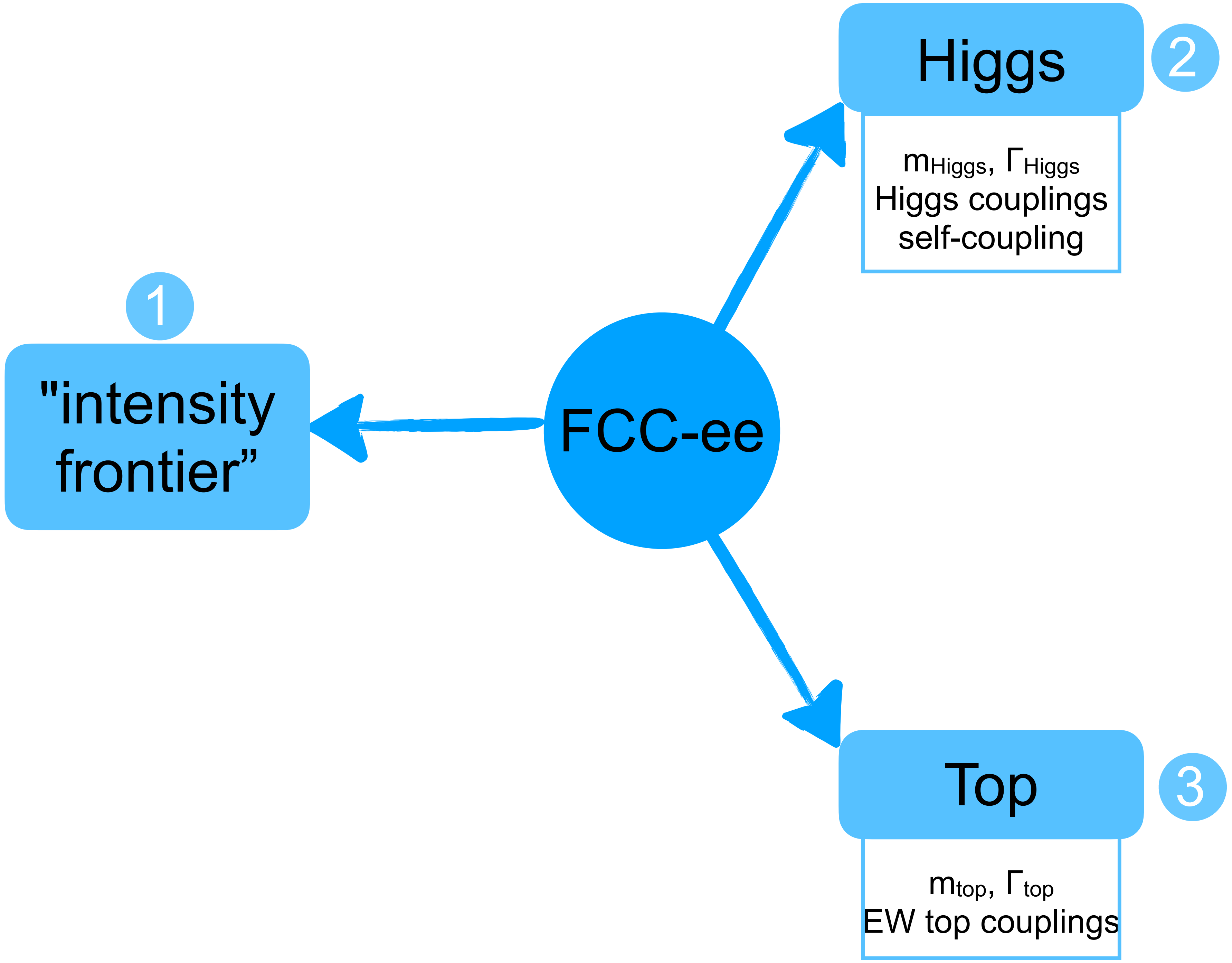

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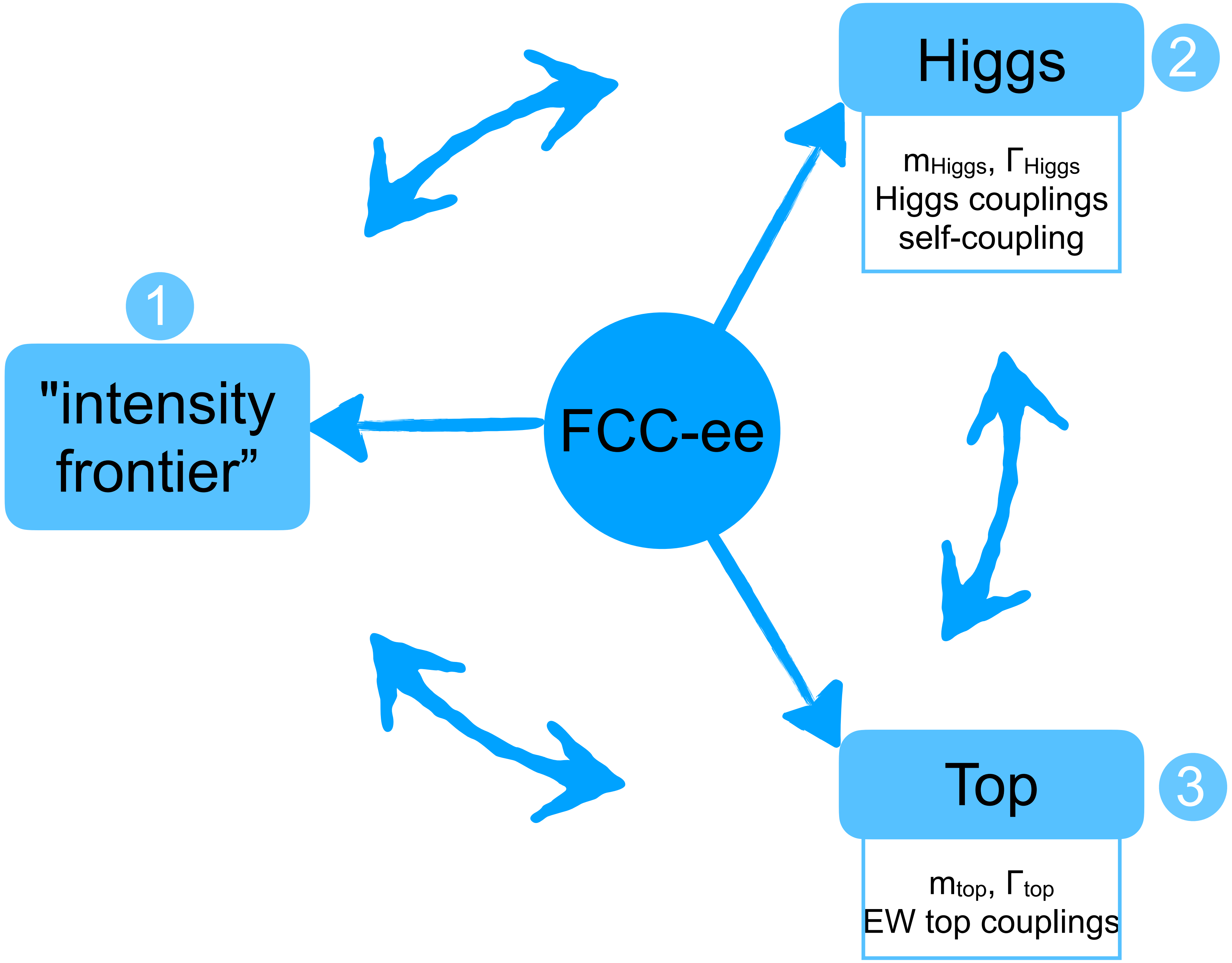

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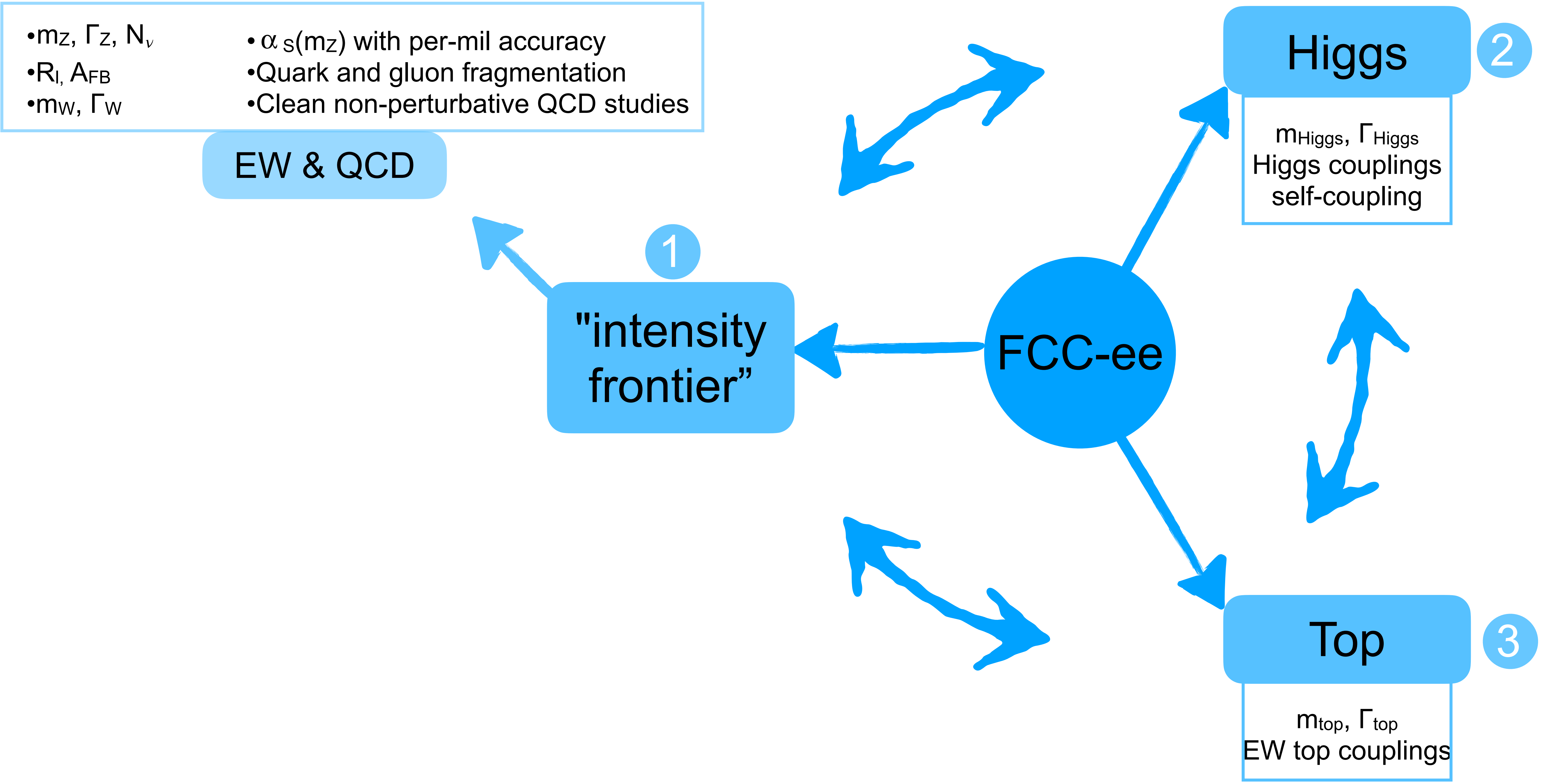

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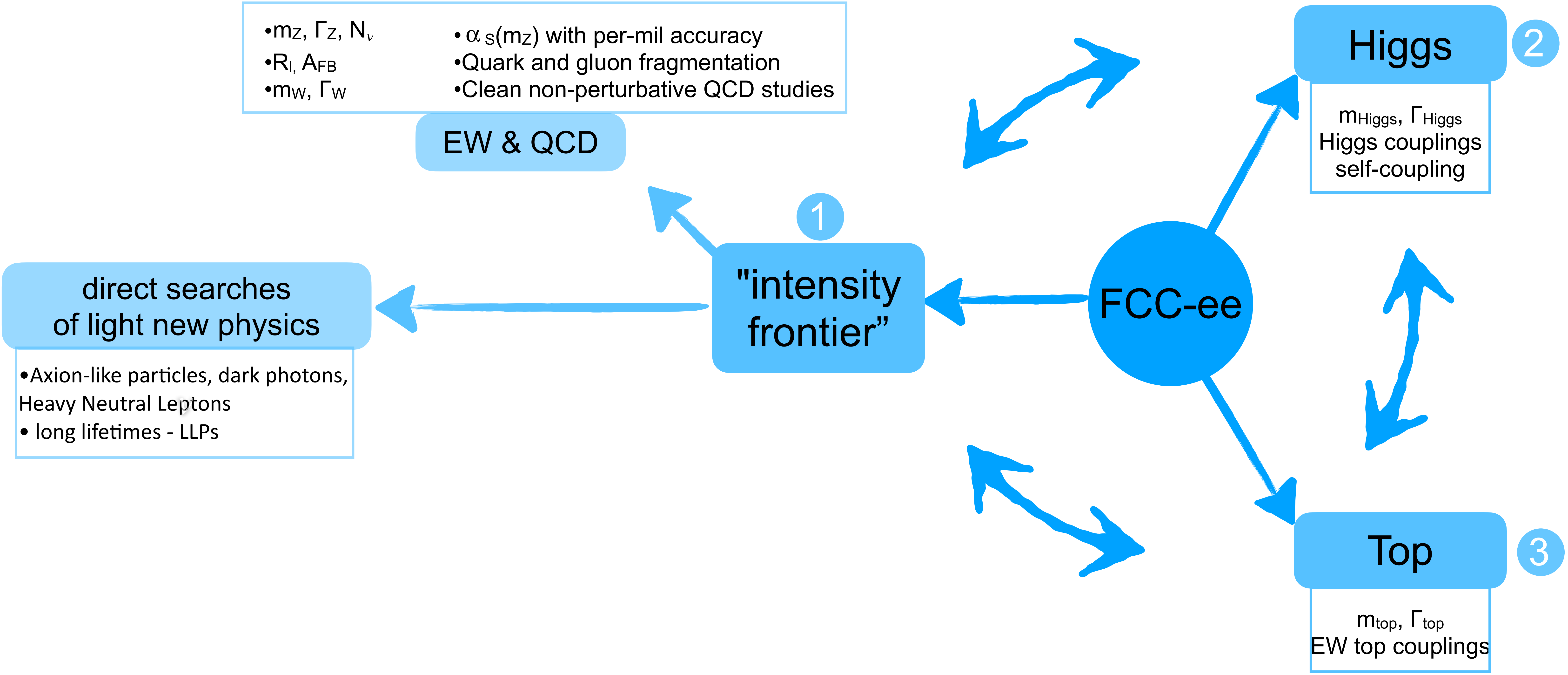

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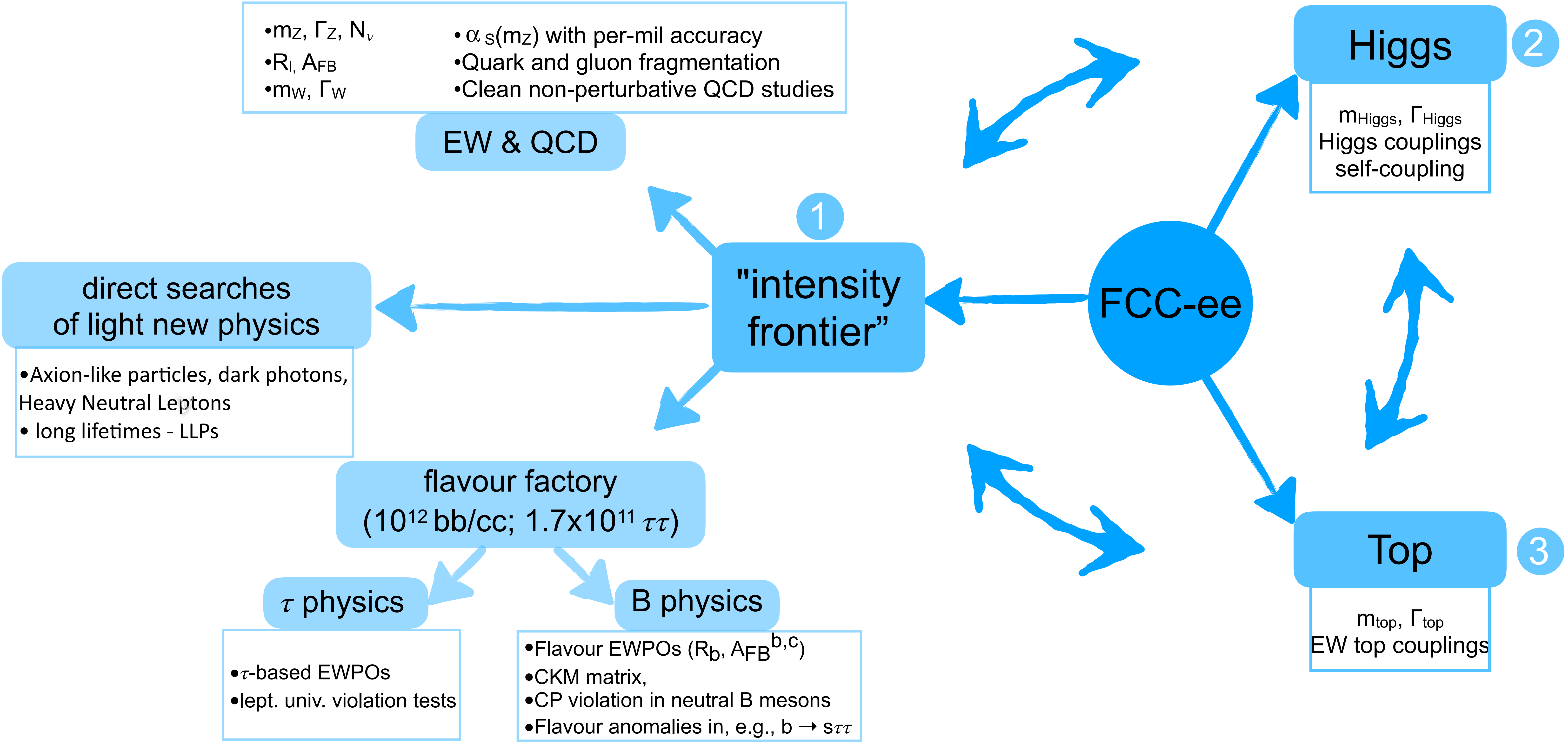

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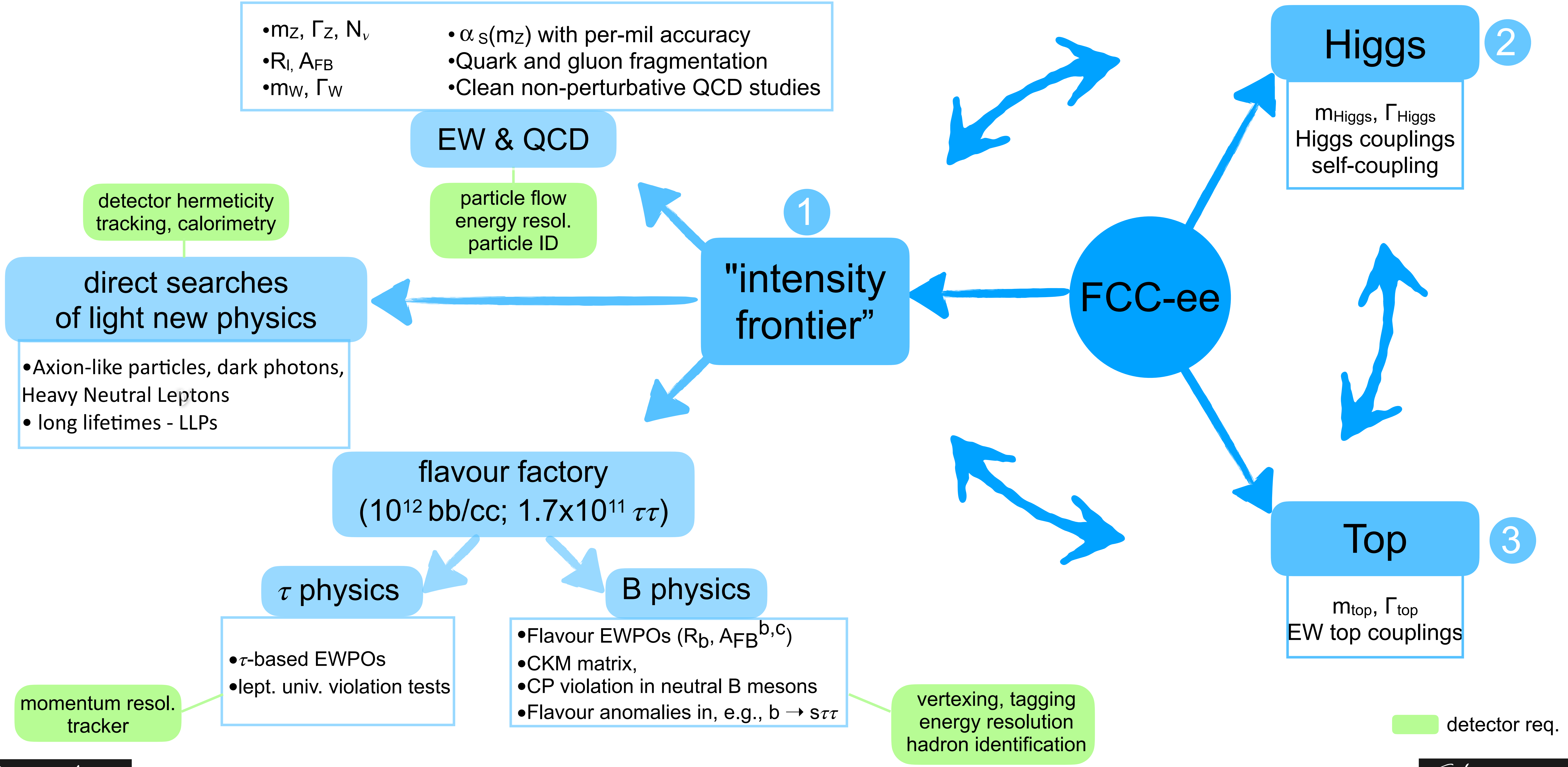
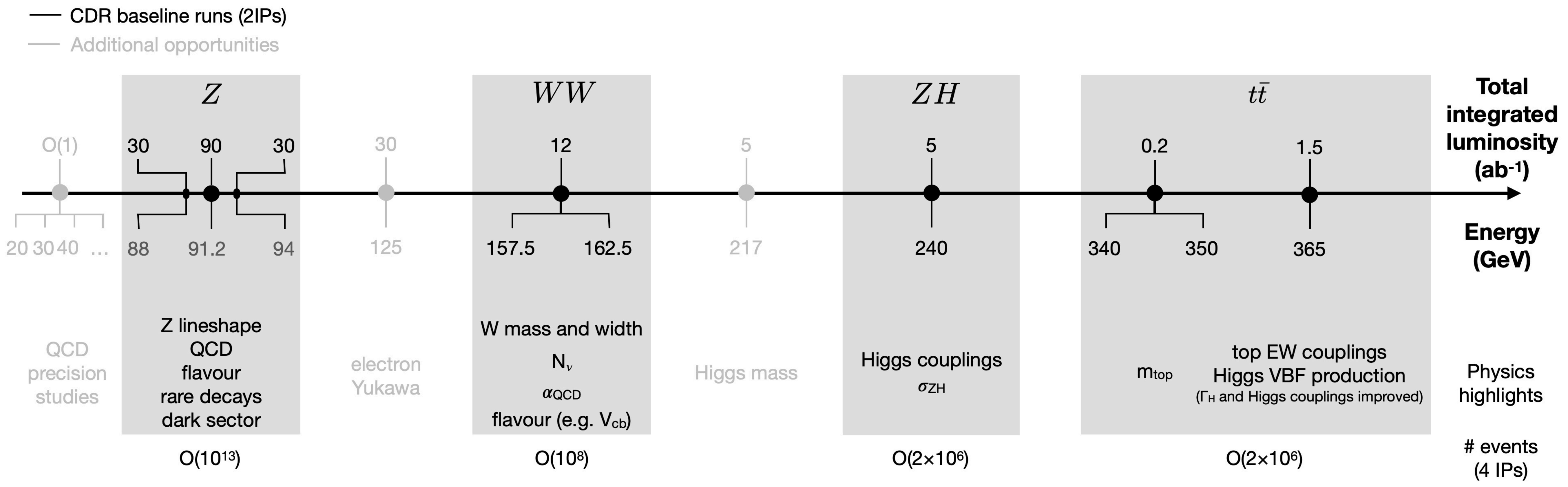


FCC-ee Physics Programme.

baseline FCC-ee detector performance

track momentum	$\frac{\sigma_p}{p} = 0.02 \cdot 10^{-3} \cdot p_T(\text{GeV}) \oplus 1 \cdot 10^{-3}$
track impact parameter	$\sigma_{d_0} = \frac{15\,\mu\text{m}}{\sin^{3/2}\theta} \oplus 5\,\mu\text{m}$
electromagnetic energy	$\frac{\sigma_{E_\gamma}}{E_\gamma} = \frac{15\%}{E_\gamma} \oplus 1\%$
electromagnetic energy xy position	$\sigma_{\gamma,xy} = \frac{6\,\text{mm}}{E(\text{GeV})} \oplus 2\,\text{mm}$

Collider Programme (and beyond).



- **Opportunities** beyond the baseline plan (√s below Z, 125GeV, 217GeV; larger integrated lumi…)
- **Opportunities** to exploit FCC facility differently (to be studied more carefully):
  - using the electrons from the injectors for beam-dump experiments,
  - extracting electron beams from the booster,
  - reusing the synchrotron radiation photons.

FCC-ee: Explore & Discover.

PED @ CERN-SPC 2022

- **EXPLORE INDIRECTLY** the 10-100 TeV energy scale with precision measurements
  - From the correlated properties of the Z , b, c, τ, W, Higgs, and top particles
    - Up to 20-50-fold improved precision on ALL electroweak observables (EWPO)
      → m_Z , m_W , m_{top} , Γ_Z , $\sin^2\theta_W^{eff}$, R_b , $\alpha_{QED}(m_Z)$, $\alpha_s(m_Z\ m_W\ m_t)$, top EW couplings …
    - Up to 10 × more precise and model-independent Higgs couplings (width, mass) measurements
      → Access the Higgs potential and infer the vacuum structure of the Universe
      → Reveals the dynamics of the EW phase transition and infer the fate of the EW vacuum

- **DISCOVER** that the Standard Model does not fit
  - New Physics! → Pattern of deviations may point to the source.

See Bonus Slides for examples and plots

- **DISCOVER** a violation of flavour conservation / universality
  - $Z\to\tau\mu$ in 5×10^{12} Z decays; $\tau\to\mu\nu/e\nu$ in 2×10^{11} τ decays; $B^0\to K^{*0}\tau^+\tau^-$ or $B_S\to\tau^+\tau^-$ in 10^{12} bb evts

- **DISCOVER** dark matter, e.g., as invisible decays of Higgs or Z

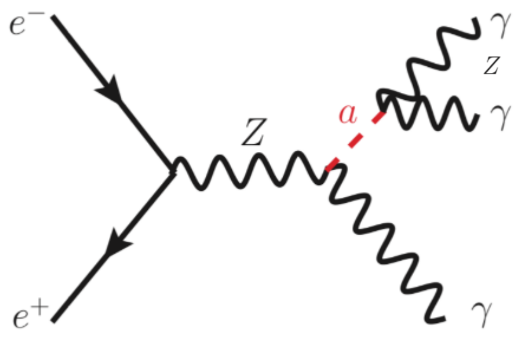
- **DISCOVER DIRECTLY** elusive (aka feebly-coupled) particles
  - in the 5-100 GeV mass range, such as right-handed neutrinos, dark photons, light Higgs-like scalars, dilaton, ALPs, relaxions…

FCC-ee: Explore & Discover.

- **ALPs@ colliders**

e.g. $e^+e^- \to \gamma a$

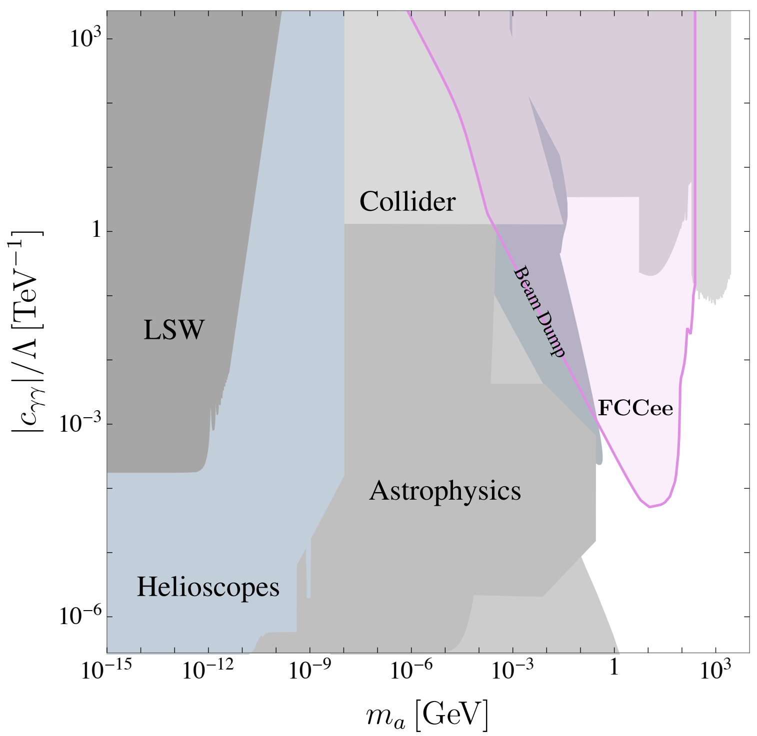
$e^+e^- \to ha$



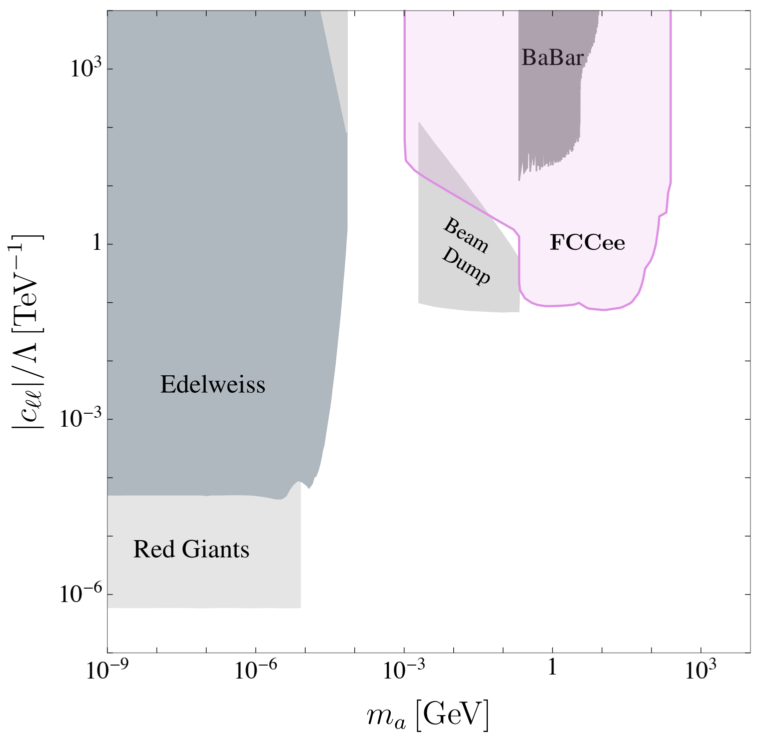
Knapen, Thamm arXiv:2108.08949

Astro/Cosmo → long-lived ALPs
colliders → short-lived ALPs MeV+

ALP coupling to photons

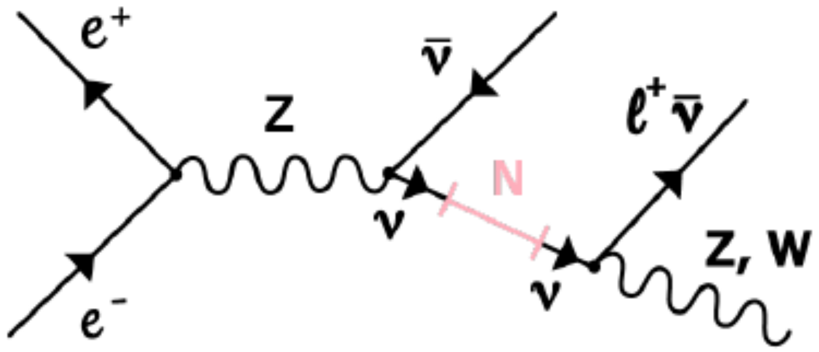


ALP coupling to electrons



- **Search for ν_{RH}.**

Direct observation
in Z decays
from LH-RH mixing



mixing active-sterile neutrinos

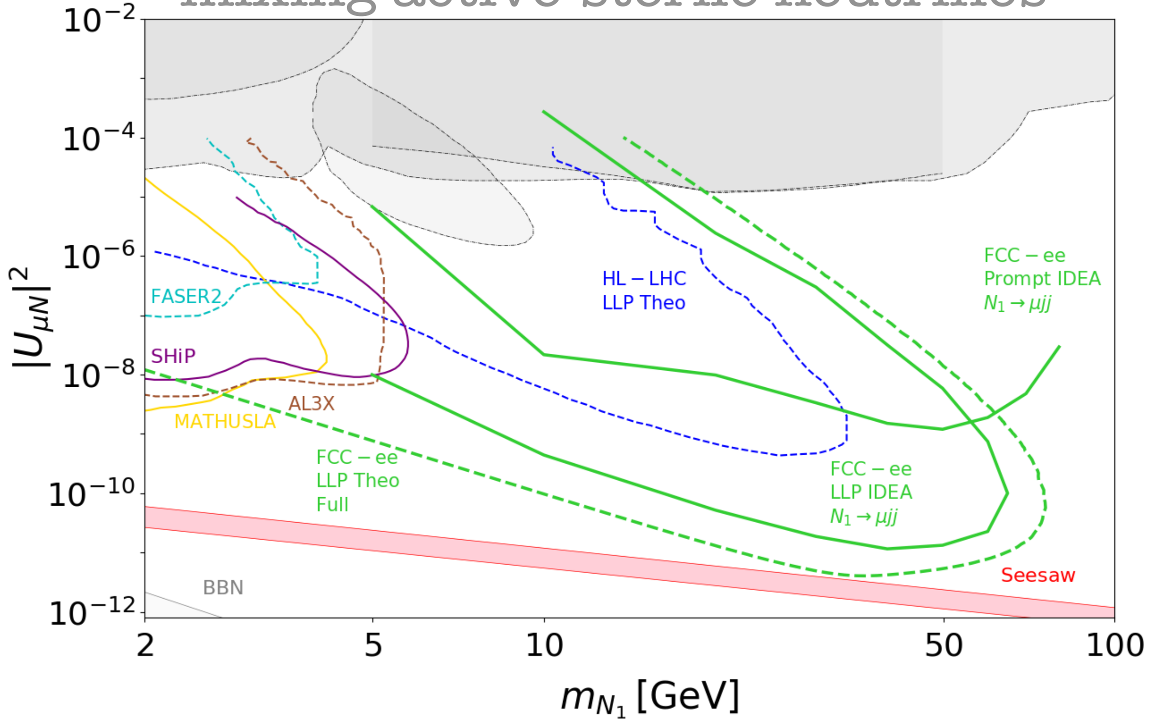


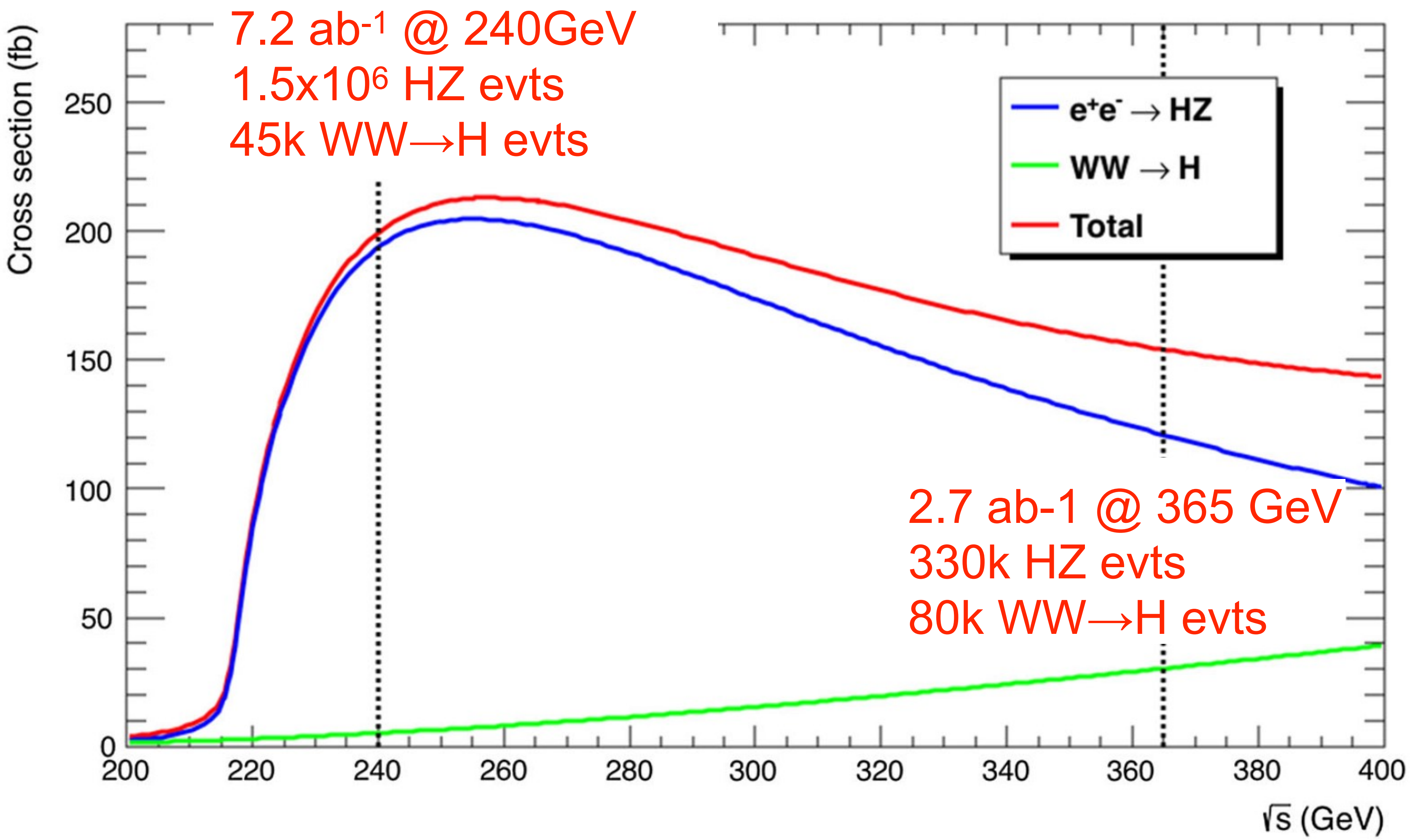
Fig. from mid-term report

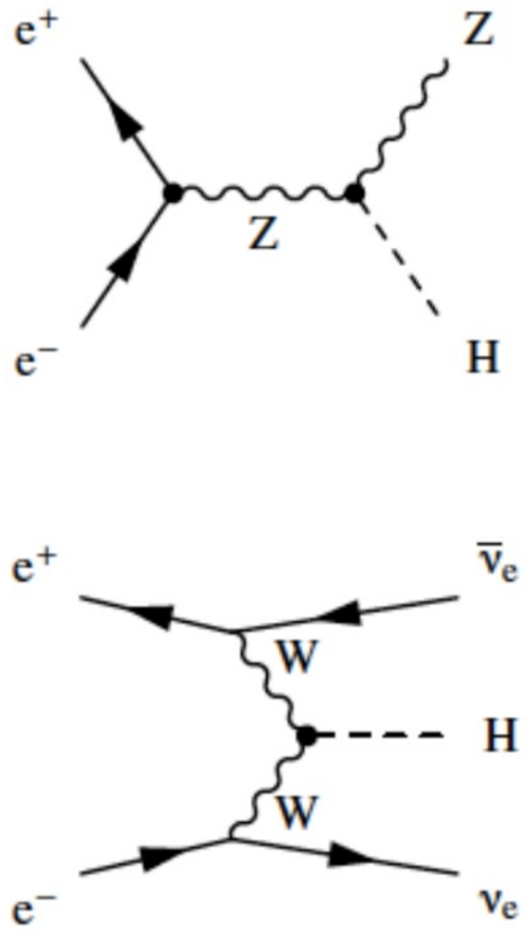
3.
FCC-ee/hh as a Higgs/electroweak factory

Higgs @ FCC-ee.

Central goal of FCC-ee: model-independent measurement of Higgs width and couplings with (<)% precision. Achieved through operation at two energy points.

G. Wilkinson, FCC Physics WS '22





Sensitivity to both processes very helpful in improving precision on couplings.

Complementarity with 365GeV on top of 240GeV improvement factor: ∞/3/2/1.5/1.2 on $\kappa_\lambda / \kappa_W / \kappa_b / \kappa_g, \kappa_c / \kappa_\gamma$ (plot in bonus)

Higgs @ FCC-ee.

- Absolute normalisation of couplings (by recoil method)
- Measurement of width (from ZH>ZZZ* and WW>H)
- $\delta\Gamma_H \sim 1\%, \delta m_H \sim 3\,\mathrm{MeV}$ (resp. 25%, 30 MeV @ HL-LHC)
- Model-independent coupling determination and improvement factor up to 10 compared to LHC
- (Indirect) sensitivity to new physics up to 70 TeV (for maximally strongly coupled models)

$$(\delta\kappa_X = v^2/f^2 \quad \& \quad m_{\mathrm{NP}} = g_{\mathrm{NP}} f)$$

Higgs coupling sensitivity

Coupling	HL-LHC	FCC-ee (240–365 GeV) 2 IPs / 4 IPs
κ_W [%]	1.5*	0.43 / 0.33
κ_Z [%]	1.3*	0.17 / 0.14
κ_g [%]	2*	0.90 / 0.77
κ_γ [%]	1.6*	1.3 / 1.2
$\kappa_{Z\gamma}$ [%]	10*	10 / 10
κ_c [%]	–	1.3 / 1.1
κ_t [%]	3.2*	3.1 / 3.1
κ_b [%]	2.5*	0.64 / 0.56
κ_μ [%]	4.4*	3.9 / 3.7
κ_τ [%]	1.6*	0.66 / 0.55
$\mathrm{BR}_{\mathrm{inv}}$ (<%, 95% CL)	1.9*	0.20 / 0.15
$\mathrm{BR}_{\mathrm{unt}}$ (<%, 95% CL)	4*	1.0 / 0.88

Table from mid-term report

$$\kappa_X = \frac{g_{hXX}}{g_{hXX}^{\mathrm{SM}}}$$

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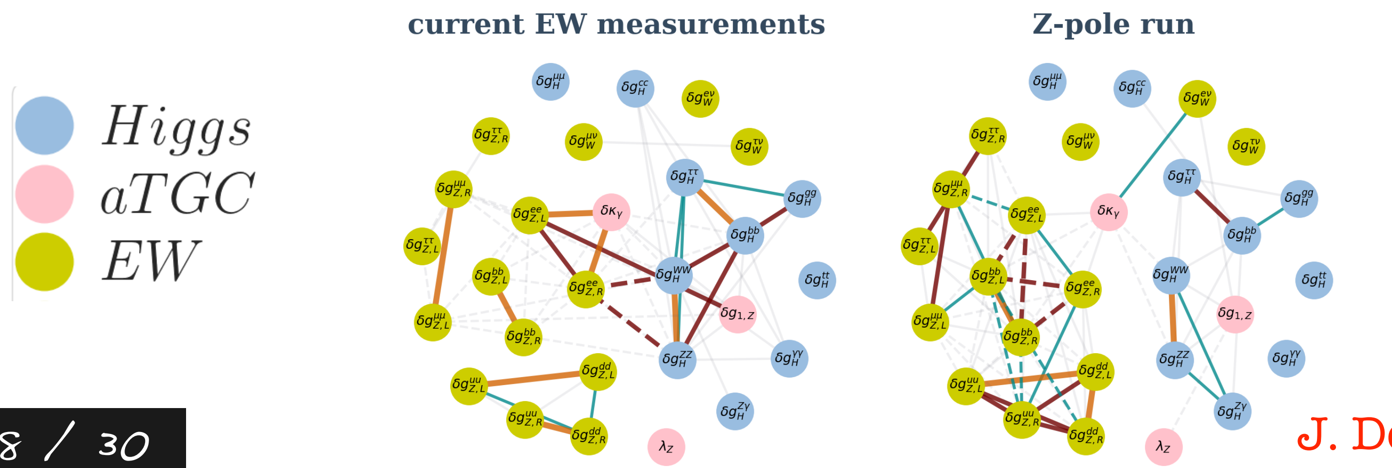
— Higgs programme needs Z-pole —



J. De Blas et al. 1907.04311

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Table from mid-term report

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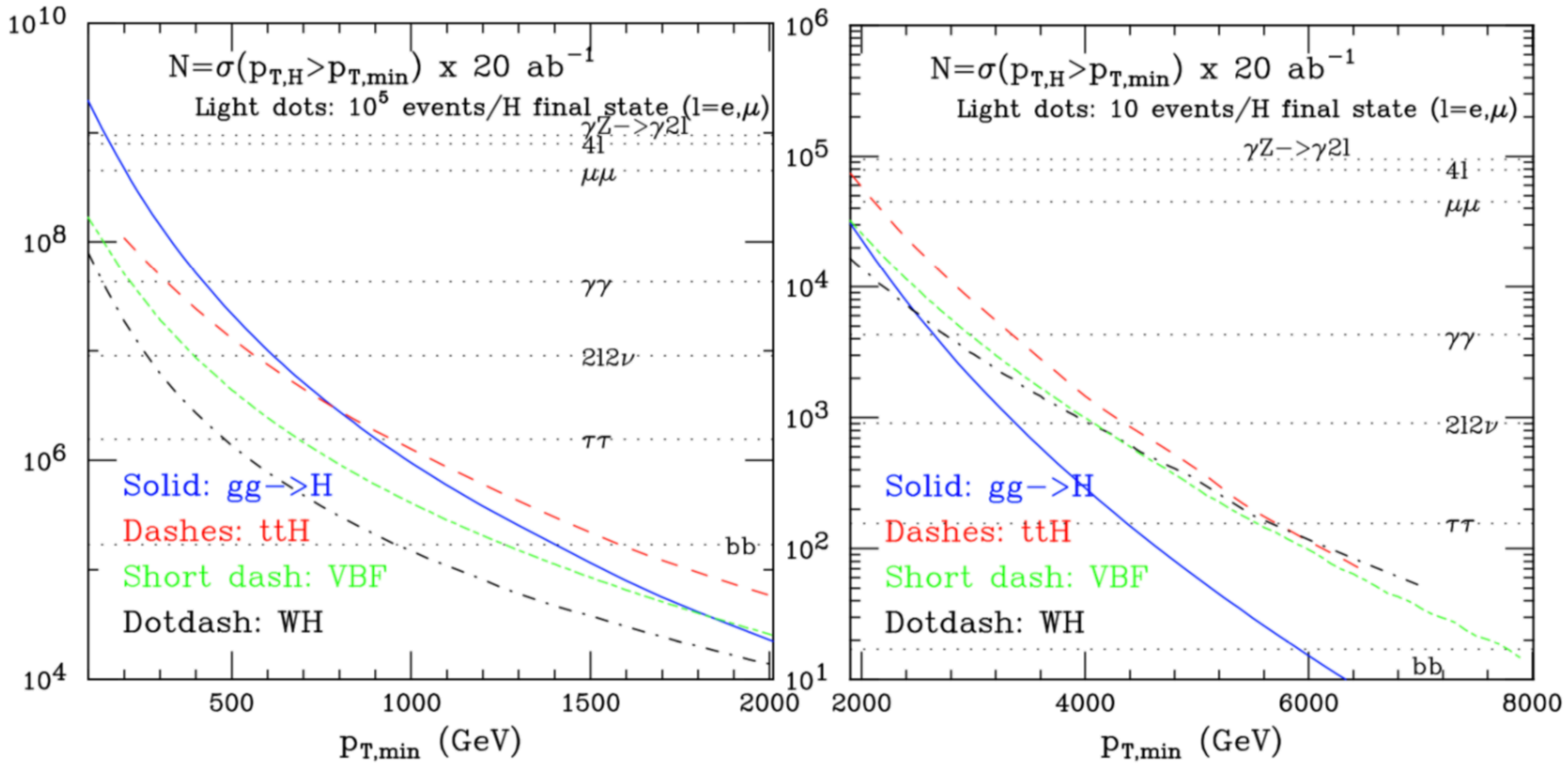
Higgs @ FCC-hh.

	ggH (N^3LO)	VBF (N^2LO)	WH (N^2LO)	ZH (N^2LO)	$t\bar{t}H$ (N^2LO)	HH (NLO)
N100	24×10^9	2.1×10^9	4.6×10^8	3.3×10^8	9.6×10^8	3.6×10^7
N100/N14	180	170	100	110	530	390

$(\mathrm{N100} = \sigma_{100\,\mathrm{TeV}} \times 30\,\mathrm{ab}^{-1} \;\;\&\;\; \mathrm{N14} = \sigma_{14\,\mathrm{TeV}} \times 3\,\mathrm{ab}^{-1})$



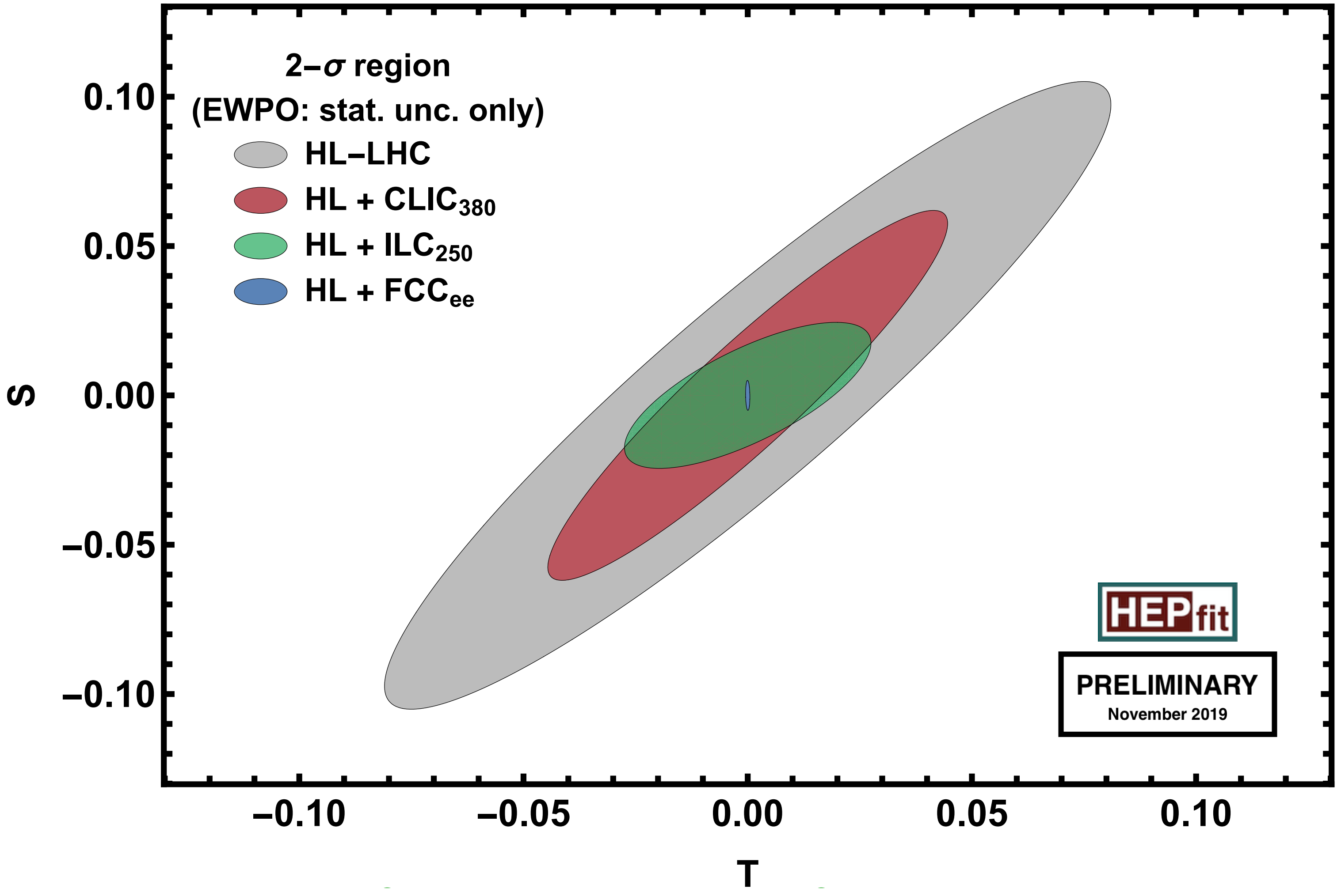
- Large rate (> 10^{10}H, > 10^7 HH)
  - unique sensitivity to rare decays
  - few % sensitivity to self-coupling
- Explore extreme phase space:
  - e.g. 10^6 H w/ pT>1 TeV
  - clean samples with high S/B
  - small systematics

Tera-Z EW precision measurements.

- The target is to reduce syst. uncertainties to the level of stat. uncertainties.
  (exploit the large samples and innovative control analyses)
- Exquisite √s precision (100keV@Z, 300keV@WW) reduces beam uncertainties (EPOL)

➡ ~50 times better precision than LEP/LSD on EW precision observables

w/ stat. and param. only



Indirect sensitivity to 70TeV-scale sector connected to EW/Higgs

(For the impact of the theory uncertainties on the EW fit, see bonus slides)

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Need TH results to fully exploit Tera-Z

Table from mid-term report

Quantity	Current precision	FCC-ee stat. (syst.) precision	Required theory input	Available calc. in 2019	Needed theory improvement†
m_Z Γ_Z $\sin^2\theta^\ell_{\rm eff}$	2.1 MeV 2.3 MeV 1.6×10^{-4}	0.004 (0.1) MeV 0.004 (0.025) MeV $2(2.4)\times10^{-6}$	non-resonant $e^+e^-\to f\bar{f}$, initial-state radiation (ISR)	NLO, ISR logarithms up to 6th order	NNLO for $e^+e^-\to f\bar{f}$
m_W	12 MeV	0.25 (0.3) MeV	lineshape of $e^+e^-\to WW$ near threshold	NLO (ee → 4f or EFT framework)	NNLO for ee → WW, W → ff in EFT setup
HZZ coupling	—	0.2%	cross-sect. for $e^+e^-\to ZH$	NLO + NNLO QCD	NNLO electroweak
$m_{\rm top}$	100 MeV	17 MeV	threshold scan $e^+e^-\to t\bar{t}$	N³LO QCD, NNLO EW, resummations up to NNLL	Matching fixed orders with resummations, merging with MC, α_s (input)

†The listed needed theory calculations constitute a minimum baseline; additional partial higher-order contributions may also be required.

Indirect sensitivity to 70TeV-scale sector connected to EW/Higgs

4.
FCC-ee as a flavour factory

Flavour potential.

At present (Z/h/NewPhysics) FCNCs mostly constrained by low energy observables.

The large statistics of FCC will open on-shell opportunities.

See S. Monteil, Flavour@FCC '22

Particle production (10^9)	B^0 / $\overline{B}^0$	B^+ / B^-	B^0_s / $\overline{B}^0_s$	Λ_b / $\overline{\Lambda}_b$	$c\bar{c}$	τ^-/τ^+
Belle II	27.5	27.5	n/a	n/a	65	45
FCC-*ee*	300	300	80	80	600	150

FCC-ee = 10 x Belle II

Decay mode/Experiment	Belle II (50/ab)	LHCb Run I	LHCb Upgr. (50/fb)	FCC-*ee*
EW/H penguins				
$B^0 \to K^*(892)e^+e^-$	~ 2000	~ 150	~ 5000	~ 200000
$\mathcal{B}(B^0 \to K^*(892)\tau^+\tau^-)$	~ 10	–	–	~ 1000
$B_s \to \mu^+\mu^-$	n/a	~ 15	~ 500	~ 800
$B^0 \to \mu^+\mu^-$	~ 5	–	~ 50	~ 100
$\mathcal{B}(B_s \to \tau^+\tau^-)$				
Leptonic decays				
$B^+ \to \mu^+\nu_{mu}$	5%	–	–	3%
$B^+ \to \tau^+\nu_{tau}$	7%	–	–	2%
$B_c^+ \to \tau^+\nu_{tau}$	n/a	–	–	5%
CP / hadronic decays				
$B^0 \to J/\Psi K_S$ ($\sigma_{\sin(2\phi_d)}$)	$\sim 2. * 10^6$ (0.008)	41500 (0.04)	$\sim 0.8 \cdot 10^6$ (0.01)	$\sim 35 \cdot 10^6$ (0.006)
$B_s \to D_s^{\pm} K^{\mp}$	n/a	6000	~ 200000	$\sim 30 \cdot 10^6$
$B_s(B^0) \to J/\Psi\phi$ (σ_{ϕ_s} rad)	n/a	96000 (0.049)	$\sim 2.10^6$ (0.008)	$16 \cdot 10^6$ (0.003)

out of reach at LHCb/Belle

boosted b's/τ's at FCC-ee

Makes possible a topological rec. of the decays w/ miss. energy

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						45
						150

FCC-ee = 10 x Belle II

Flavour @ FCC vs Belle/pp

Attribute	$\Upsilon(4S)$	pp	Z^0
All hadron species		✓	✓
High boost		✓	✓
Enormous production cross-section		✓	
Negligible trigger losses	✓		✓
Low backgrounds	✓		✓
Initial energy constraint	✓		(✓)

Decay ...				...C-ee
EW/H...				
$B^0 \to$ ...				...0000
$\mathcal{B}(B^0 \to$ ...				...000
$B_s \to \mu$ ...				...800
$B^0 \to \mu$ ...				...100
$\mathcal{B}(B_s \to$ ...				
Leptoni...				
$B^+ \to$ ...				...%
$B^+ \to \tau^+\nu_{tau}$	...%			...%
$B_c^+ \to \tau^+\nu_{tau}$	n/a	–	–	5%
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Low backgrounds	✓		✓
Initial energy constraint	✓		(✓)

Decay				C-ee
EW/H				
$B^0 \to$				0000
$\mathcal{B}(B^0 -$				000
$B_s \to \mu$				800
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$\mathcal{B}(B_s -$				
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Makes possible a topological rec. of the decays w/ miss. energy

Flavour defines shared (vertexing, tracking, calorimetry) and specific (hadronic PID) detector requirements.

FCC-ee flavour opportunities.

- **CKM element V_{cb}** (critical for normalising the Unitarity Triangle) from WW decays
- **Tau physics** (>10^{11} pairs of tau's produced in Z decays)
  - test of lepton flavour universality: G_F from tau decays @ 10 ppm @ FCC-ee (0.5 ppm from muon decays)
  - lepton flavour violation:
    - $\tau \to \mu\gamma$: 4×10^{-8} @Belle2021→10^{-9} @ FCC-ee
    - $\tau \to 3\mu$: 2×10^{-8} @Belle → 3×10^{-10} @BelleII → 10^{-11} @ FCC-ee
  - tau lifetime uncertainty:
    - 2000 ppm → 10 ppm
  - tau mass uncertainty:
    - 70 ppm → 14 ppm
- **Semi-leptonic mixing asymmetries** a^s_{sl} and a^d_{sl}
- …

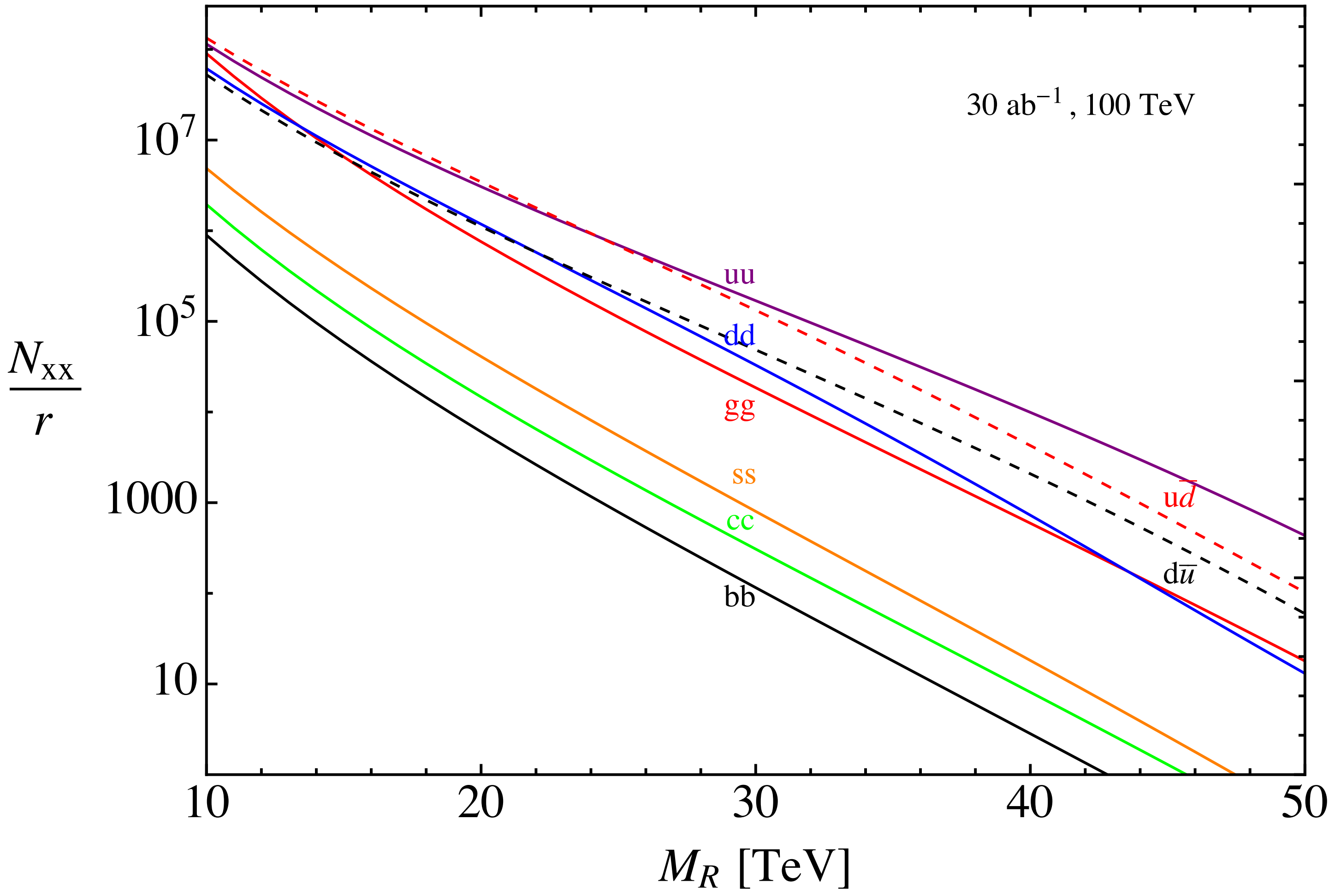
5.
FCC-hh: the broadest exploration potential at high-energy

Resonance production.

Protons are made of 5 quarks, gluons, photons, W/Z

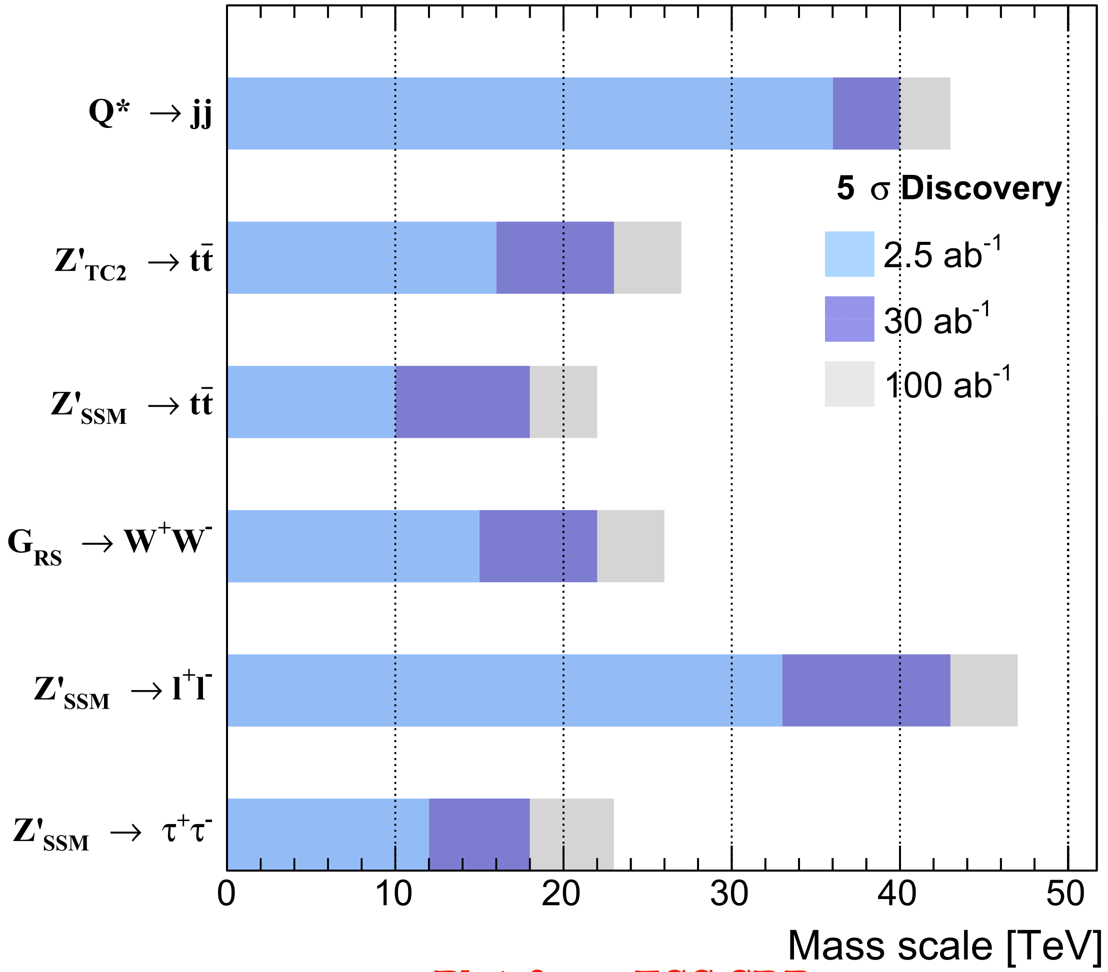
FCC-hh effectively collides 196 different initial states, perfect exploratory machine

resonances produced



Plot from mid-term report

FCC-hh mass reach



Plot from FCC CDR

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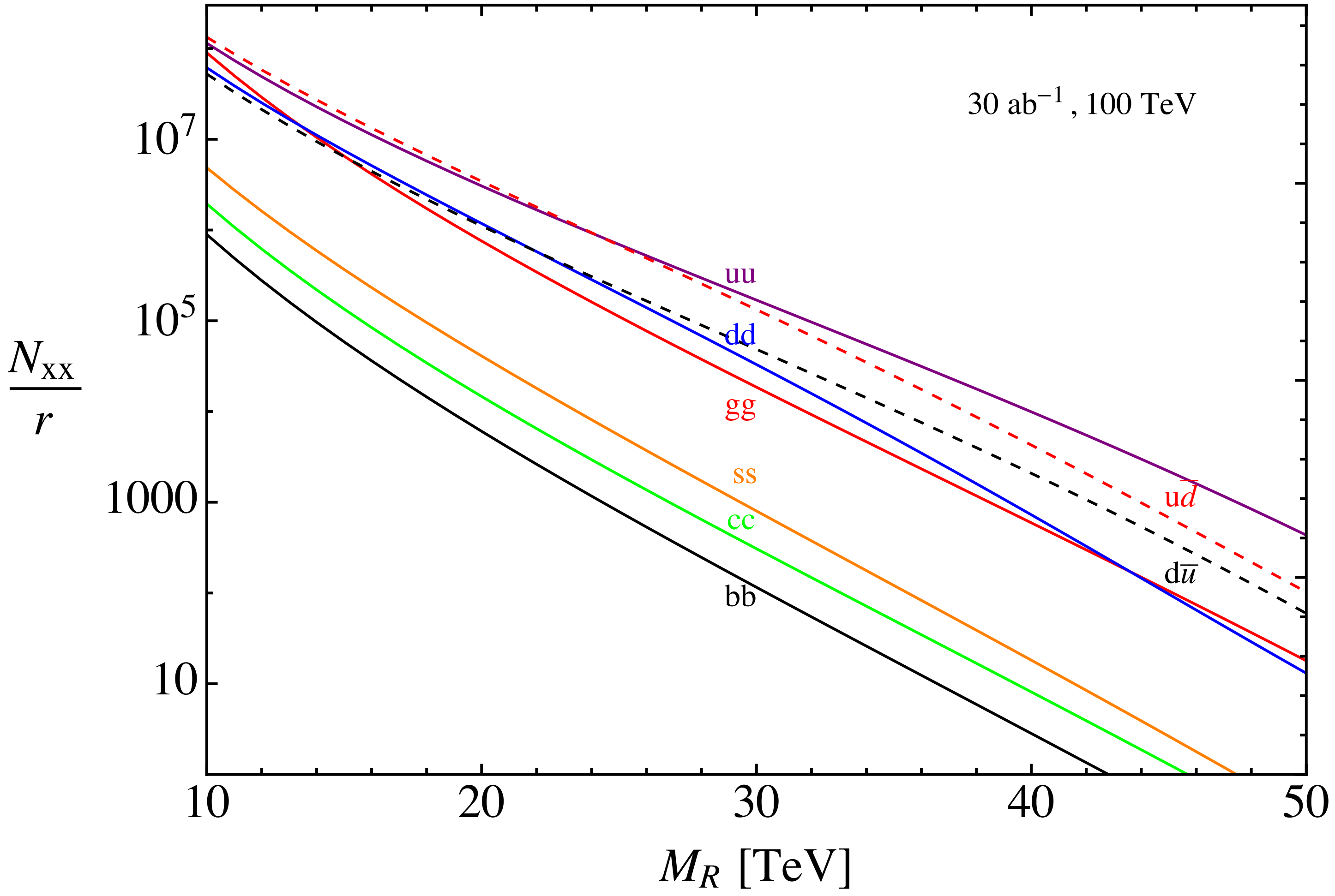
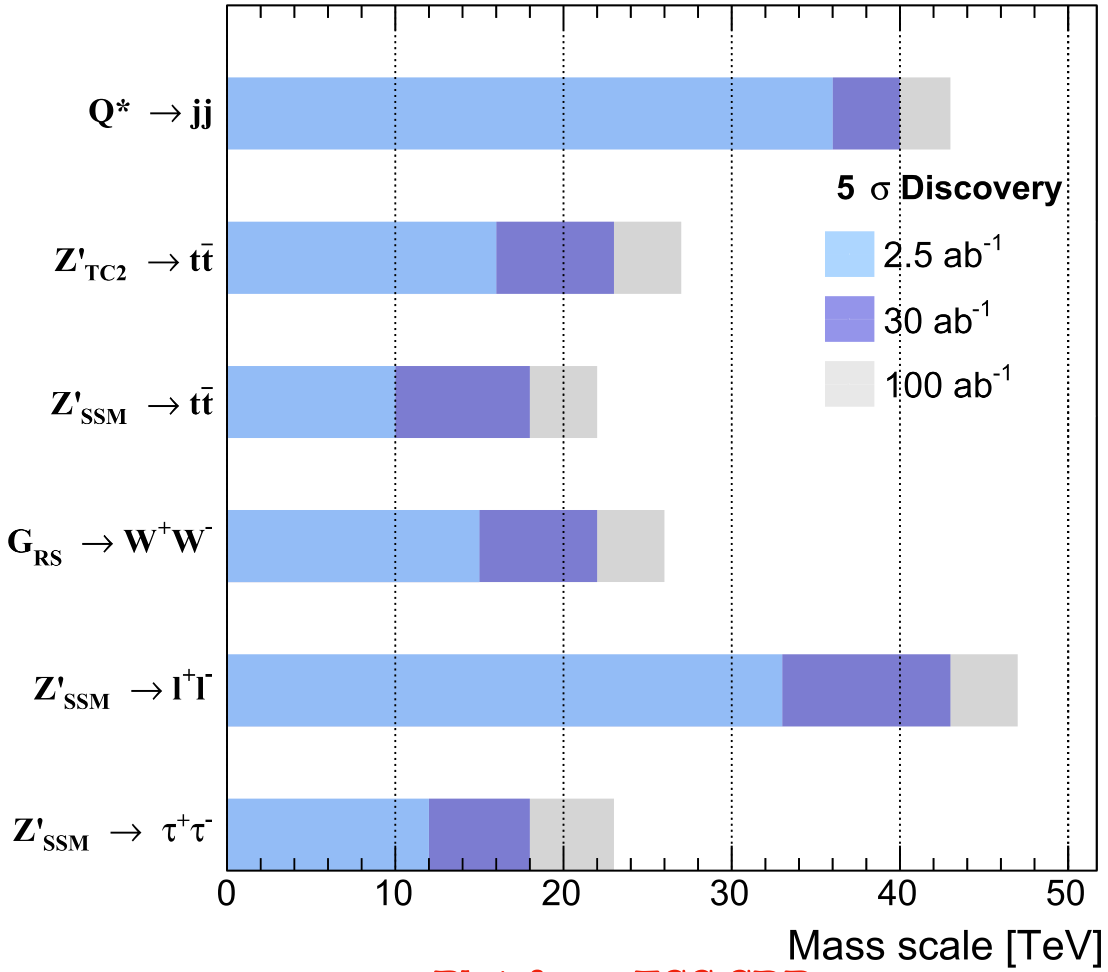


Plot from mid-term report

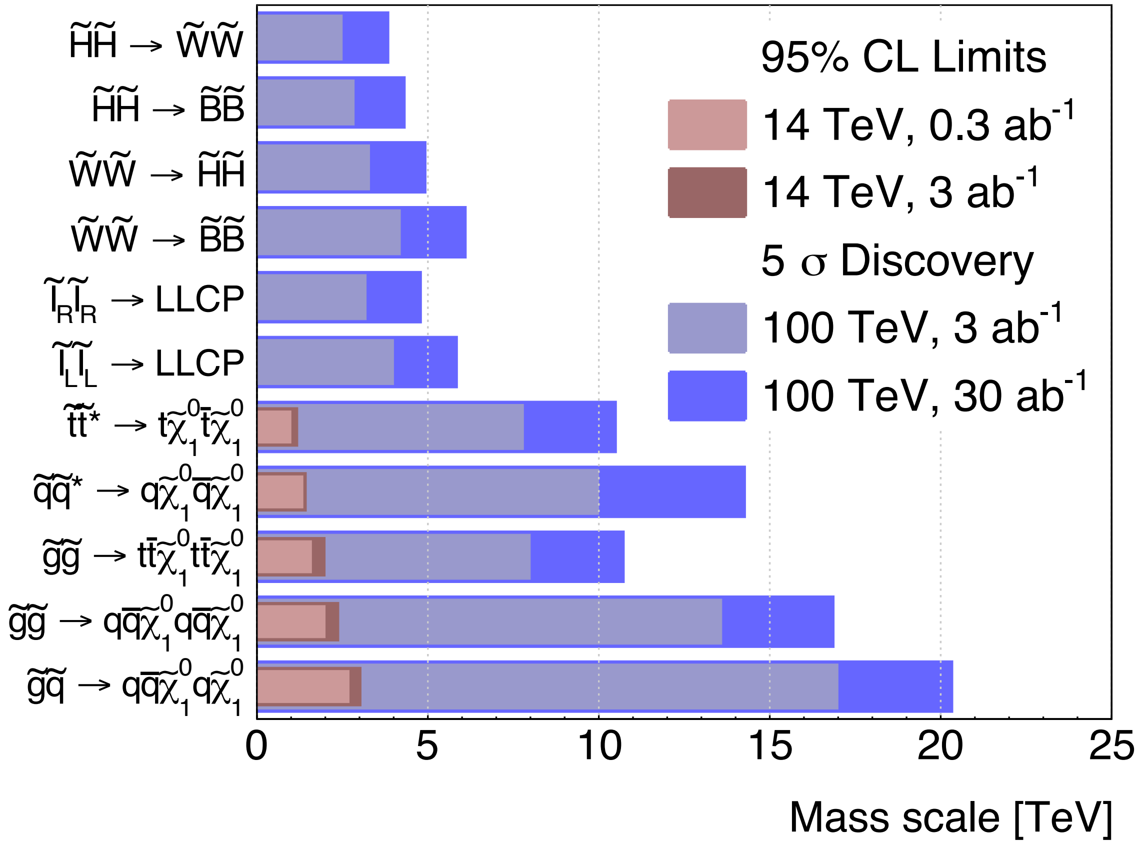
FCC-hh mass reach



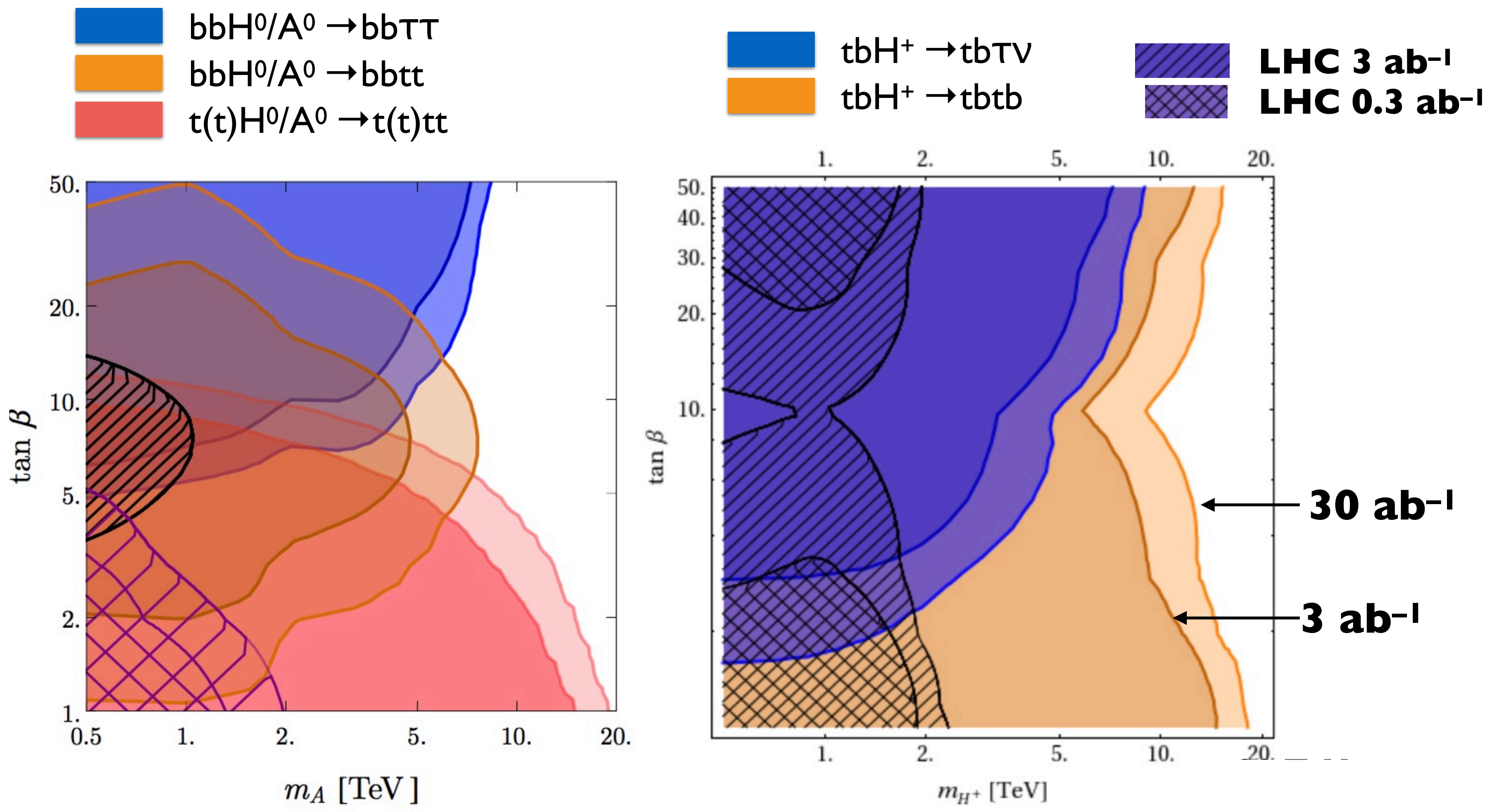
Plot from FCC CDR

FCC-hh allows the direct exploration of new physics at energy scales up to 40 TeV, including any physics that may be indirectly indicated by precision Higgs and EW measurements at FCC-ee.

Pushing limits of SUSY.



Plot from arXiv:1606.00947



Plot from arXiv:1605.08744 and arXiv:1504.07617

15-20TeV squarks/gluinos

require kinematic threshold 30-40TeV:

FCC-hh is more than a $\sqrt{\hat{s}}$~10TeV factory

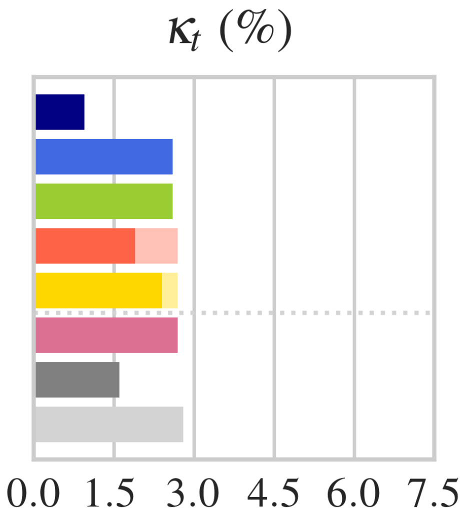
Factor 10 increase on the HL-LHC limits.

6.
Examples of complementarity & synergy FCC-ee↔FCC-hh

Synergy ee↔hh.

1 FCC-hh without ee could bound BR_{inv} but it could say nothing about $BR_{untagged}$ (FCC-ee needed for absolute normalisation of Higgs couplings)

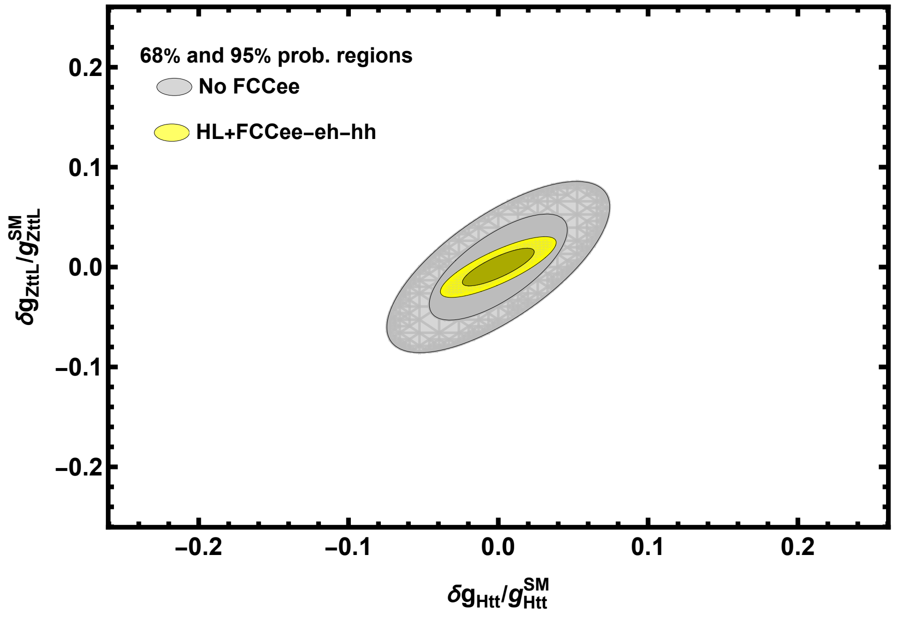
2 FCC-hh is determining top Yukawa through ratio tth/ttZ
So the extraction of top Yukawa heavily relies on the knowledge of ttZ from FCC-ee



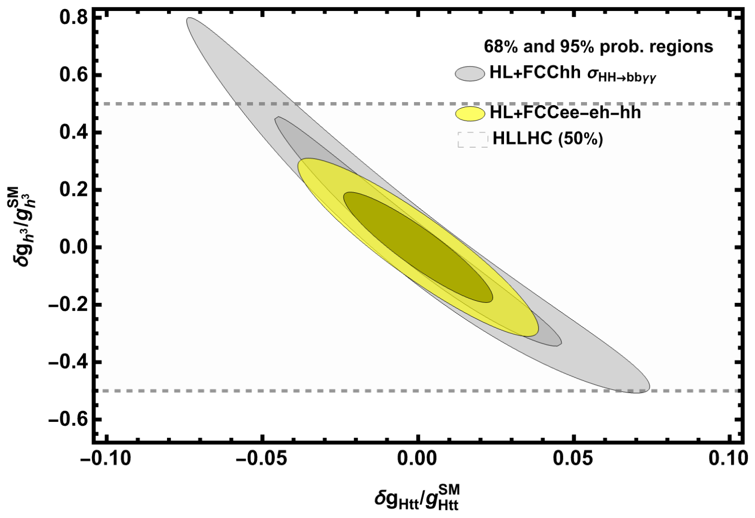
Mangano+ '15

	$\sigma(t\bar{t}H)$[pb]	$\sigma(t\bar{t}Z)$[pb]	$\frac{\sigma(t\bar{t}H)}{\sigma(t\bar{t}Z)}$
13 TeV	$0.475^{+5.79\%+3.33\%}_{-9.04\%-3.08\%}$	$0.785^{+9.81\%+3.27\%}_{-11.2\%-3.12\%}$	$0.606^{+2.45\%+0.525\%}_{-3.66\%-0.319\%}$
100 TeV	$33.9^{+7.06\%+2.17\%}_{-8.29\%-2.18\%}$	$57.9^{+8.93\%+2.24\%}_{-9.46\%-2.43\%}$	$0.585^{+1.29\%+0.314\%}_{-2.02\%-0.147\%}$

(uncertainty drops in ratio)



3 Subsequently, the 1% sensitivity on tth is essential to determine h^3 at O(5%) at FCC-hh



Plots from mid-term report

FCC-hh tunnel is great for FCC-ee.

- **80-100 km is needed to accelerate pp up to 100 TeV**
- **80-100 km is also exactly what is needed**
  - to get enough luminosity (5 times more than in 27 km) to maybe get sensitivity to the Higgs self coupling, the electron Yukawa coupling, or sterile neutrinos,
  - to make TeraZ a useful flavour factory,
  - for transverse polarisation to be available all the way to the WW threshold (allowing a precise W mass measurement)
  - for the top threshold to be reached and exceeded.

Conclusions & Outlook

A **circular "Higgs factory"** like FCC-ee has a rich potential:

* Search directly and indirectly for New Physics
* Establish new organising principles of Nature (LEP→ gauge symmetries, FCC→??)
* Probe the **HEP-Cosmo connections** thanks to the high statistics of the **Z-pole run**

(omitting this exploration would be ignoring the outcome of LHC).

And FCC-ee is an essential part of an **integrated** programme to probe the energy frontier.

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FCC will for sure contribute.

We can learn a lot from nice pictures/**observations**
but **experiments** remain the driver of physics.

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Colliders are the most powerful microscopes we have to study Nature at the smallest scales and also from the early moments of the Universe. They provide a quantitative understanding to progress forward.

Acknowledgement.

This project is supported from the European Union's Horizon 2020 research and innovation programme under grant agreement No 951754.

BONUS

Some work ahead of us.

- Development of a common software and the estimate of the computing needs
- Evaluation of the physics performance and requirements for detectors
- Conceptualisation of detectors capable of delivering these requirements
- Mitigation of the interaction region constraints on detectors and vice versa
- Design of methods and tools for centre-of-mass energy calibration, beam polarisation, and monochromatization
- Understanding and optimisation of the physics programm
- Exploration of the physics opportunities
- Development of the theoretical tools and observables needed to meet the measurement targets

Higgs and EW measurements

Experimental Inputs.

A circular ee Higgs factory starts as a Z/EW factory (**TeraZ**)

A linear ee Higgs factory operating above Z-pole can also preform EW measurements via **Z-radiative** return

A linear ee Higgs factory could also operate on the Z-pole though at lower lumi (**GigaZ**)

	Higgs	aTGC	EWPO	Top EW
FCC-ee	Yes (μ, σ_{ZH}) (Complete with HL-LHC)	Yes (aTGC dom.) Warning	Yes	Yes (365 GeV, Ztt)
ILC	Yes (μ, σ_{ZH}) (Complete with HL-LHC)	Yes (HE limit) **Warning**	LEP/SLD (Z-pole) + HL-LHC + W (ILC)	Yes (500 GeV, Ztt)
CEPC	Yes (μ, σ_{ZH}) (Complete with HL-LHC)	Yes (aTGC dom) Warning	Yes	No
CLIC	Yes (μ, σ_{ZH})	Yes (Full EFT parameterization)	LEP/SLD (Z-pole) + HL-LHC + W (CLIC)	Yes
HE-LHC	Extrapolated from HL-LHC	N/A → LEP2	LEP/SLD + HL-LHC (M_W, $\sin^2\theta_W$)	-
FCC-hh	Yes (μ, BR_i/BR_j) Used in combination with FCCee/eh	From FCC-ee	From FCC-ee	-
LHeC	Yes (μ)	N/A → LEP2	LEP/SLD + HL-LHC (M_W, $\sin^2\theta_W$)	-
FCC-eh	Yes (μ) Used in combination with FCCee/hh	From FCC-ee	From FCC-ee + Zuu, Zdd	-

EW Precision Measurements at FCC-ee

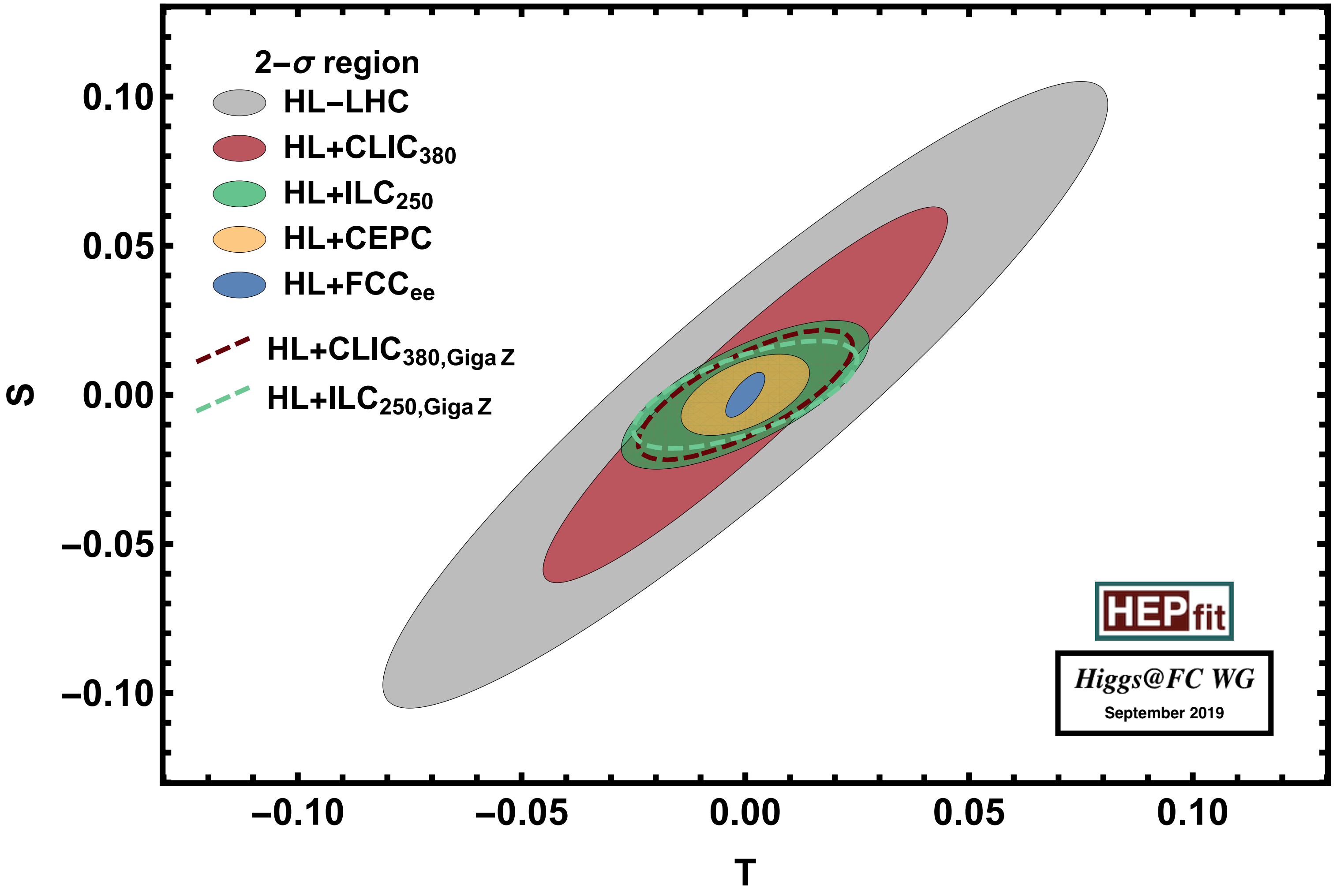
Observable	present value	±	error	FCC-ee **Stat.**	FCC-ee Syst.	Comment and leading error
m_Z (keV)	91186700	±	2200	**4**	100	From Z line shape scan Beam energy calibration
Γ_Z (keV)	2495200	±	2300	**4**	25	From Z line shape scan Beam energy calibration
$\sin^2\theta_W^{eff}(\times 10^6)$	231480	±	160	**2**	2.4	From $A_{FB}^{\mu\mu}$ at Z peak Beam energy calibration
$1/\alpha_{QED}(m_Z^2)(\times 10^3)$	128952	±	14	**3**	small	From $A_{FB}^{\mu\mu}$ off peak QED&EW errors dominate
R_ℓ^Z $(\times 10^3)$	20767	±	25	**0.06**	0.2-1	Ratio of hadrons to leptons Acceptance for leptons
$\alpha_s(m_Z^2)$ $(\times 10^4)$	1196	±	30	**0.1**	0.4-1.6	From R_ℓ^Z
σ_{had}^0 $(\times 10^3)$ (nb)	41541	±	37	**0.1**	4	Peak hadronic cross-section Luminosity measurement
$N_\nu(\times 10^3)$	2996	±	7	**0.005**	1	Z peak cross-sections Luminosity measurement
R_b $(\times 10^6)$	216290	±	660	**0.3**	< 60	Ratio of $b\bar{b}$ to hadrons Stat. extrapol. from SLD
$A_{FB}^b, 0$ $(\times 10^4)$	992	±	16	**0.02**	1-3	b-quark asymmetry at Z pole From jet charge
$A_{FB}^{pol,\tau}$ $(\times 10^4)$	1498	±	49	**0.15**	<2	τ polarization asymmetry τ decay physics
τ lifetime (fs)	290.3	±	0.5	**0.001**	0.04	Radial alignment
τ mass (MeV)	1776.86	±	0.12	**0.004**	0.04	Momentum scale
τ leptonic $(\mu\nu_\mu\nu_\tau)$ B.R. (%)	17.38	±	0.04	**0.0001**	0.003	e/μ/hadron separation
m_W (MeV)	80350	±	15	**0.25**	0.3	From WW threshold scan Beam energy calibration
Γ_W (MeV)	2085	±	42	**1.2**	0.3	From WW threshold scan Beam energy calibration
$\alpha_s(m_W^2)(\times 10^4)$	1010	±	270	**3**	small	From R_ℓ^W
$N_\nu(\times 10^3)$	2920	±	50	**0.8**	small	Ratio of invis. to leptonic in radiative Z returns
m_{top} (MeV)	172740	±	500	**17**	small	From $t\bar{t}$ threshold scan QCD errors dominate
Γ_{top} (MeV)	1410	±	190	**45**	small	From $t\bar{t}$ threshold scan QCD errors dominate
$\lambda_{top}/\lambda_{top}^{SM}$	1.2	±	0.3	**0.10**	small	From $t\bar{t}$ threshold scan QCD errors dominate
ttZ couplings		±	30%	**0.5 – 1.5** %	small	From $\sqrt{s} = 365$ GeV run

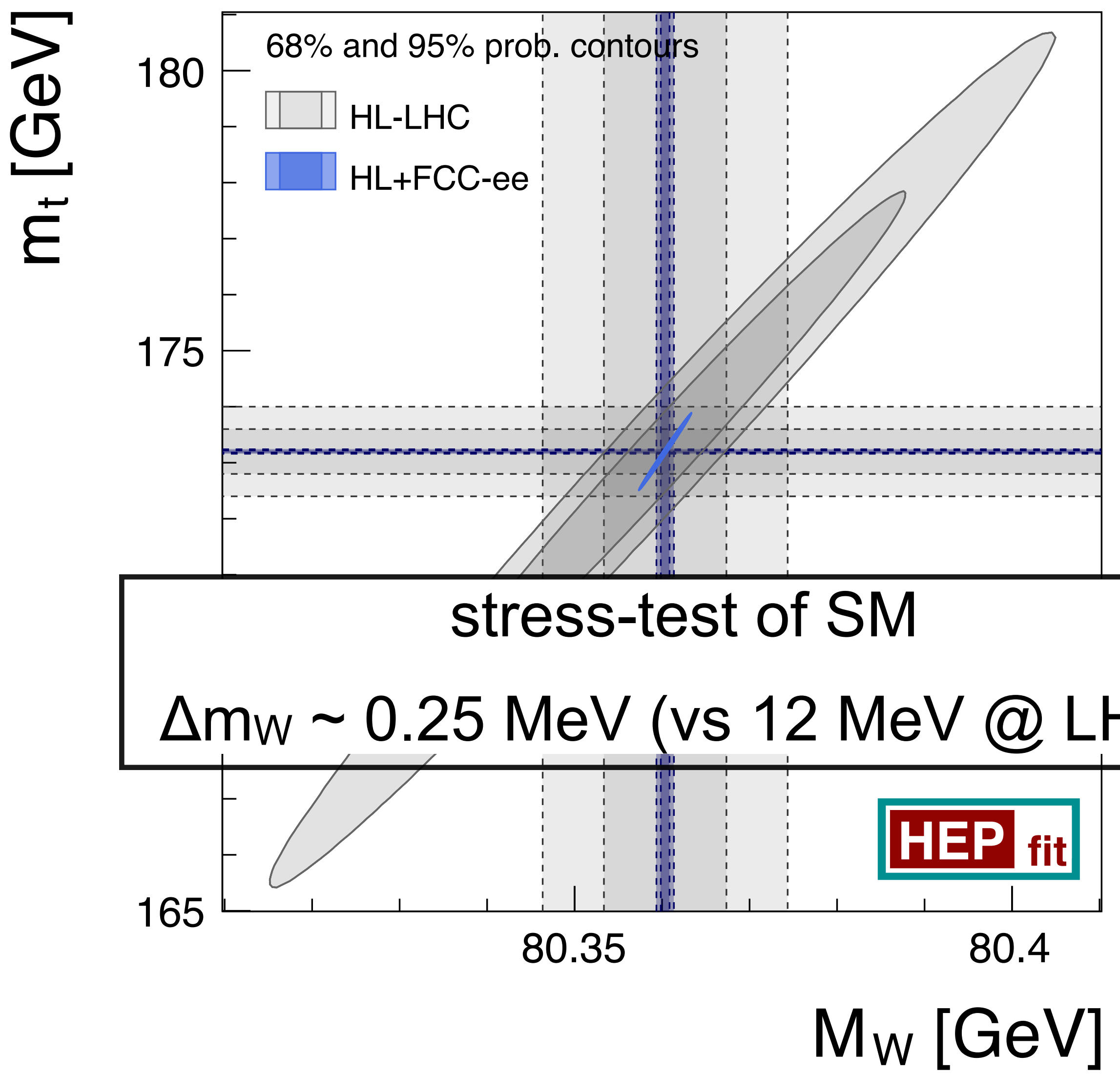
Table from mid-term report

Improvements of EW measurements

Exquisite measurements of m_Z (100 keV) , Γ_Z (25 keV), m_W (<500 keV), $\alpha_{QED}(m_Z)$ (3.10^{-5}) (all unique to FCC-ee)

w/. stat.+ param. + th-exp syst.





stress-test of SM

Δm_W ~ 0.25 MeV (vs 12 MeV @ LHC)

The importance of improved EW measurements is threefold:

1) improve mass reach in indirect search for NP ($S\sim10^{-2} \rightarrow M\sim70$ TeV)
2) reduced parametric uncertainties for other measurements
3) reduced degeneracies in a global fit for Higgs couplings

Improvements of EW measurements

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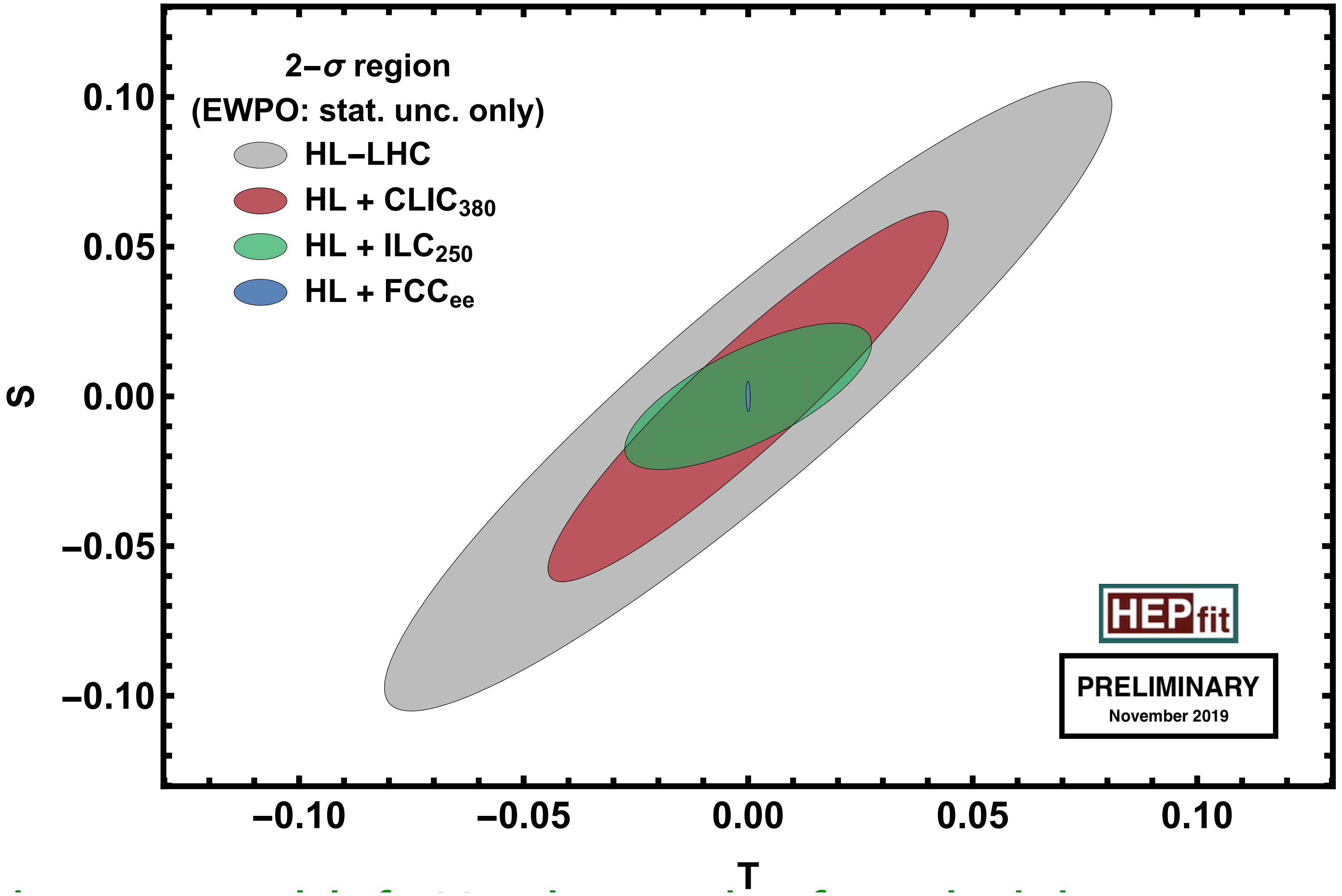
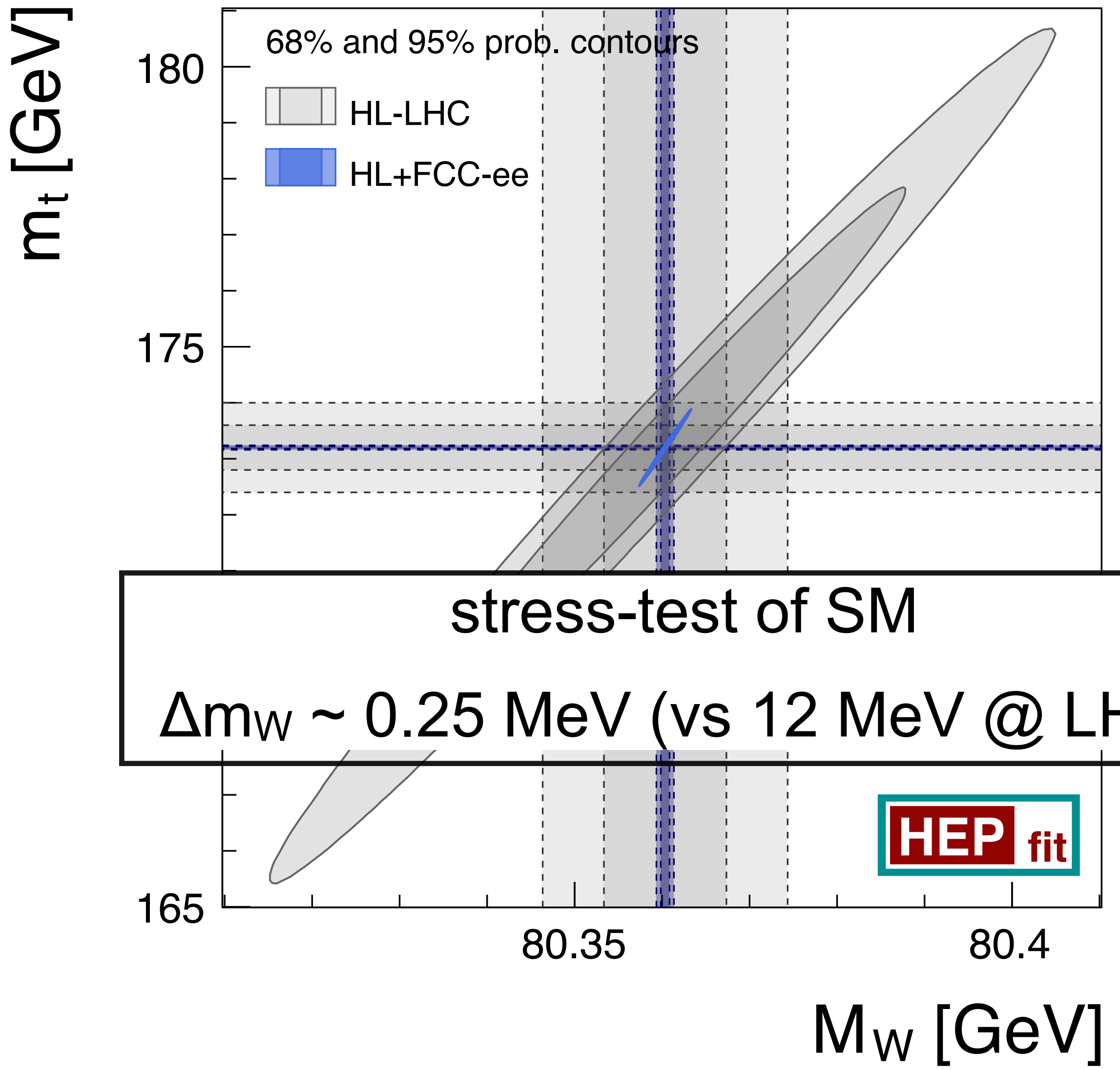
w/ stat. and param. only





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Systematics vs. Statistics.

- We often hear that more Z pole statistics is useless, because they are systematics-limited
  - This is a passive attitude, which leads to pessimistic expectations and wrong conclusions/planning
    - Experience shows that a careful experimental systematic analysis boils down to a statistical problem
    - If well prepared, theory will go as far as deemed useful : this preparation starts today (and needs SUPPORT)
    - We are working in the spirit of matching systematic errors to expected statistics for all precision measurements
  - Take the Z lineshape

FCC-ee special

$\alpha_{QED}(m_Z)$: Stat. 3×10^{-5}

Obtained at FCC-ee from off-peak asymmetries (87.9 & 94.3 GeV): for the first time, it is a direct measurement of this quantity (game changer)

- Enters as a limiting parametric uncertainties in the new physics interpretation many past and future measurements.
- Is statistics limited and will directly benefit from more luminosity
- No useful impact on $\alpha_{QED}(m_Z)$ with five times less luminosity

$\sin^2\theta_W^{eff}$ and Γ_Z (also m_W vs m_Z) : Stat. 2×10^{-6} and 4 keV

Error dominated by point-to-point energy uncertainties.
Based on in-situ comparisons between $\sqrt{s}$ (e.g. with muon pairs), with measurements made every few minutes (100's times per day)
Boils down to

- statistics (the more data the better, scales down as $1/\sqrt{L}$)
- detector systematics (uncorrelated between experiments, scales down a $1/\sqrt{N_{experiments}}$)

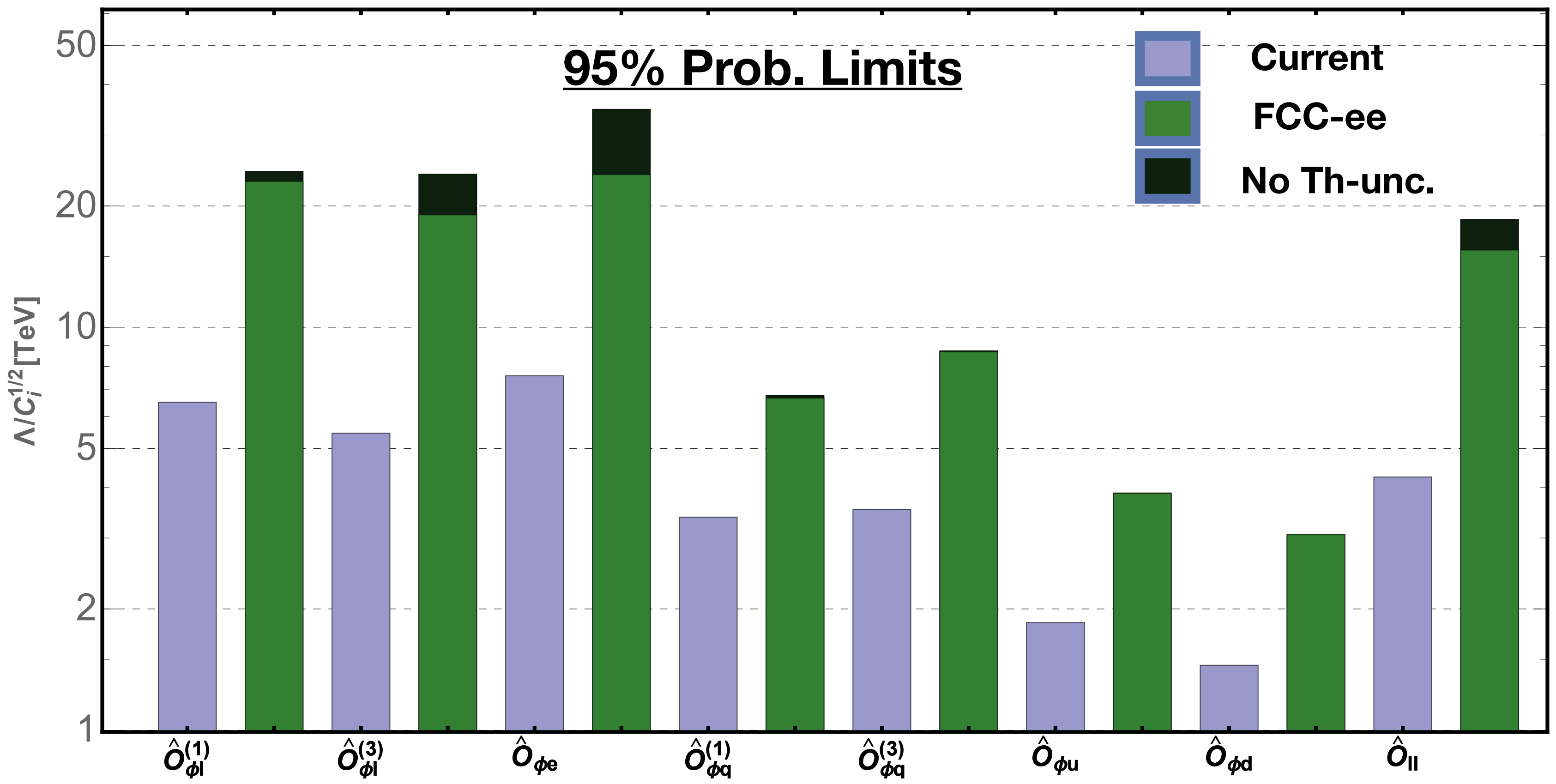
Z (and W) mass: Stat. 4 keV (250 keV)

Error dominated by $\sqrt{s}$ determination with resonant depolarization.
As more understanding is gained, progress are made at a constant pace, and this error approaches regularly the statistical limit

- Most of the work is (will be) on systematics
  - But huge statistics will turn into better precision
    - ➔ A real chance for discovery

Impact of TH uncertainties.

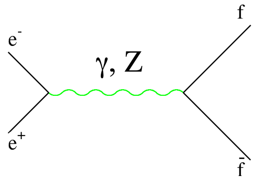
J. de Blas, FCC CDR overview '19



	Current		FCCee		
	Exp.	SM	Exp.	SM (par.)	SM (th.)
δM_W [MeV]	± 15	± 8	± 1	$\pm 0.6/\pm 1$	± 1
$\delta \Gamma_Z$ [MeV]	± 2.3	± 0.73	± 0.1	± 0.1	± 0.2
$\delta \mathcal{A}_\ell$ [$\times 10^{-5}$]	± 210	± 93	± 2.1	$\pm 8/\pm 14$	± 11.8
δR_b^0 [$\times 10^{-5}$]	± 66	± 3	± 6	± 0.3	± 5

Some EW measurements @ Tera

measure $\sigma(e^+e^- \to \mu^+\mu^-)$ and $A_{FB}^{\mu\mu}$ at (a) judicious $\sqrt{s}$



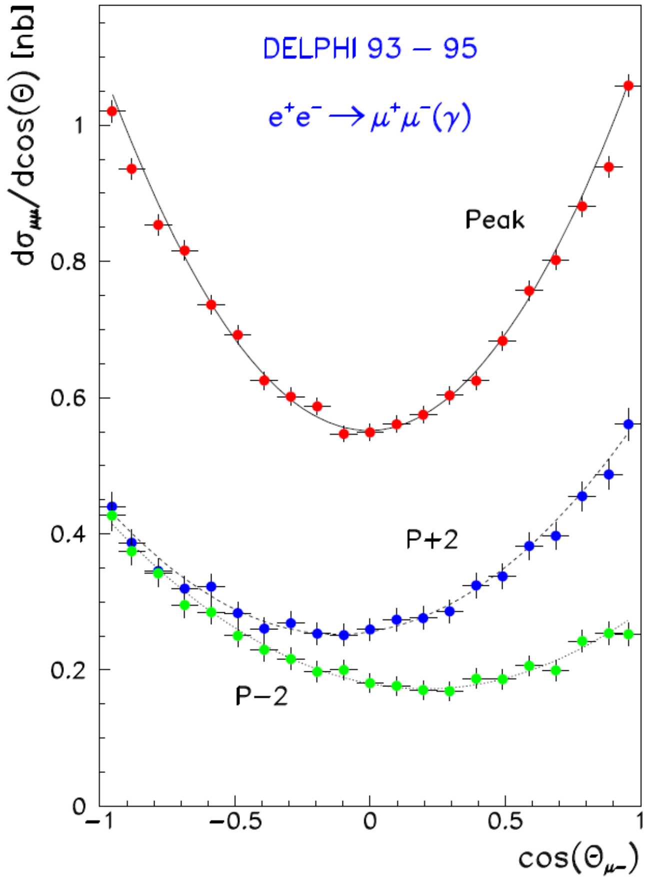
The γ exchange term is proportional to $\alpha^2_{QED}(\sqrt{s})$
The Z exchange term is proportional to G^2_F, hence independent of α_{QED}
The γZ interference is proportional to $\alpha_{QED}(\sqrt{s}) \times G_F$

Excellent experimental control of off-peak di-muon asymmetry motivates campaign to collect 50-80 ab^{-1} off peak to gain highest sensitivity to Z-γ interference

strongly depends on $\sqrt{s}$ ←
direct measurement of $\alpha_{QED}(s)$ at $\sqrt{s}$!= m_Z ←
measure $\sin^2\theta_W$ to high precision ←

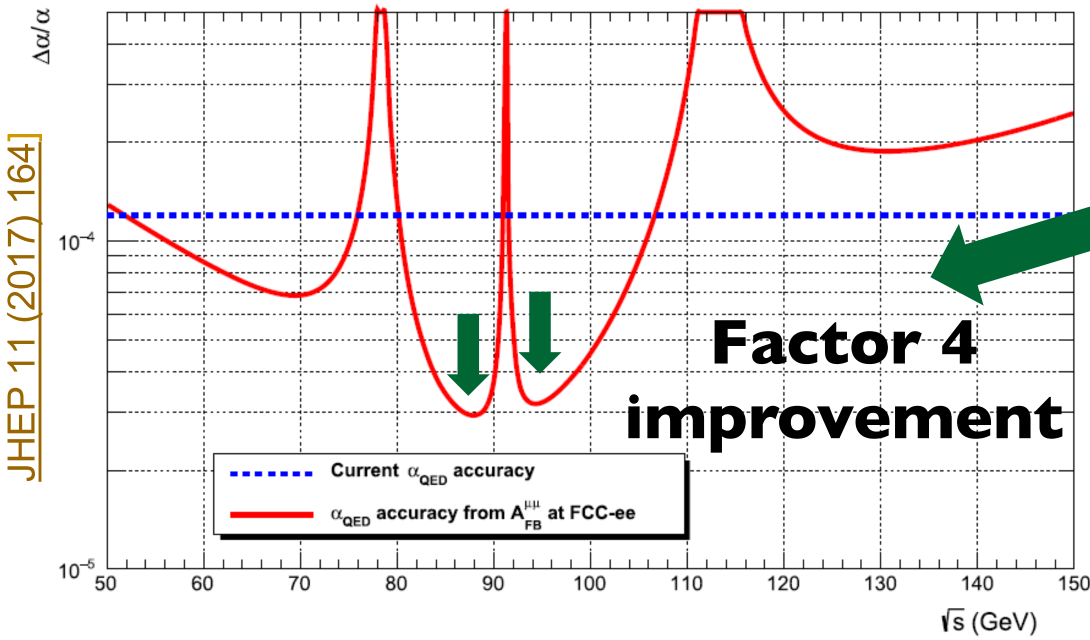
$$A_{FB}^{\mu\mu}(s) \simeq \frac{3}{4} \mathcal{A}_e \mathcal{A}_\mu \times \left[1 + \frac{8\pi\sqrt{2}\alpha_{QED}(s)}{m_Z^2 G_F \left(1 - 4\sin^2\theta_W^{eff}\right)^2} \frac{s - m_Z^2}{2s} \right]$$

Allows for clean determination of $\alpha_{QED}(m_Z^2)$, which is a *critical* input for m_W closure tests (see later).



relative α_{QED} uncertainty with 80 ab^{-1}

P.Janot '16
[JHEP 02 (2016) 053, JHEP 11 (2017) 164]



This dependence, & location of half-integer spin tunes, guides the choice of off-peak energies: 87.8 & 93.9 GeV.

→ Measure $\alpha_{QED}(m_Z^2)$ to 3×10^{-5} rel. precision (currently 1.1×10^{-4})
→ Stat. dominated; syst. uncertainties < 10^{-5} (dominated by $\sqrt{s}$ calib)
→ Theoretical uncertainties ~ 10^{-4}, higher order calcs needed

G. Wilkinson @ FCC week 2021

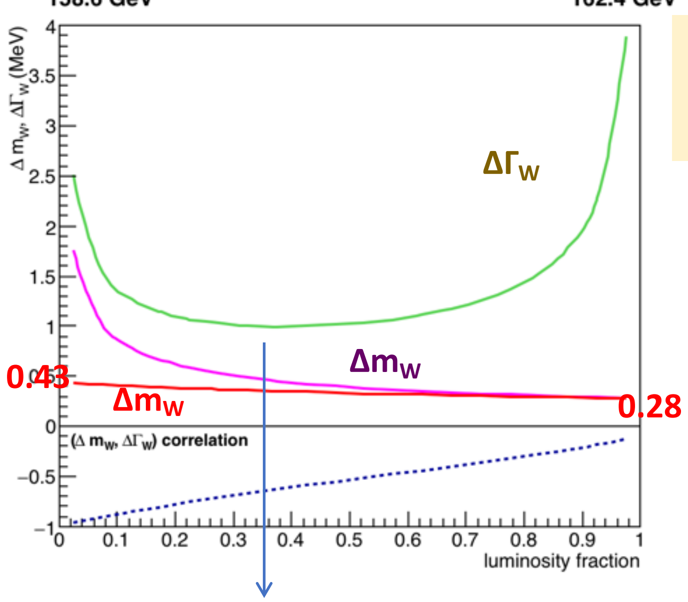
Mw.

- Two independent W mass and width measurements @FCCee :

1. **The** m_W and Γ_W determinations from the WW threshold cross section lineshape, with 12/ab at $E_{CM} \simeq$ 157.5-162.5 GeV **Δm_W=0.4 MeV $\Delta\Gamma_W$=1 MeV**

2. Other measurements of m_W and Γ_W from the decay products kinematics at $E_{CM} \simeq$ 162.5-240-365 GeV **Δm_W , $\Delta\Gamma_W$= 2-5 MeV ?**

Scans of possible E_1 E_2 data taking energies and luminosity fractions f (at the E_2 point)



Δm_W =0.45 MeV , $\Delta\Gamma_W$=1 MeV (r=-0.6)
Δm_W=0.35 MeV

A -minimum of **$\Delta\Gamma_W$=0.91 MeV** with **Δm_W=0.55 MeV**
taking data at E_1=**156.6** GeV E_2=**162.4** GeV f=**0.25**
yields **Δm_W=0.47 MeV** (as single par)

B- minimum of **Δm_W=0.28 MeV** **$\Delta\Gamma_W$=3.3 MeV** with
E_1=**155.5** GeV E_2=**162.4** GeV f=**0.95**
yields **Δm_W=0.28 MeV** (as single par)

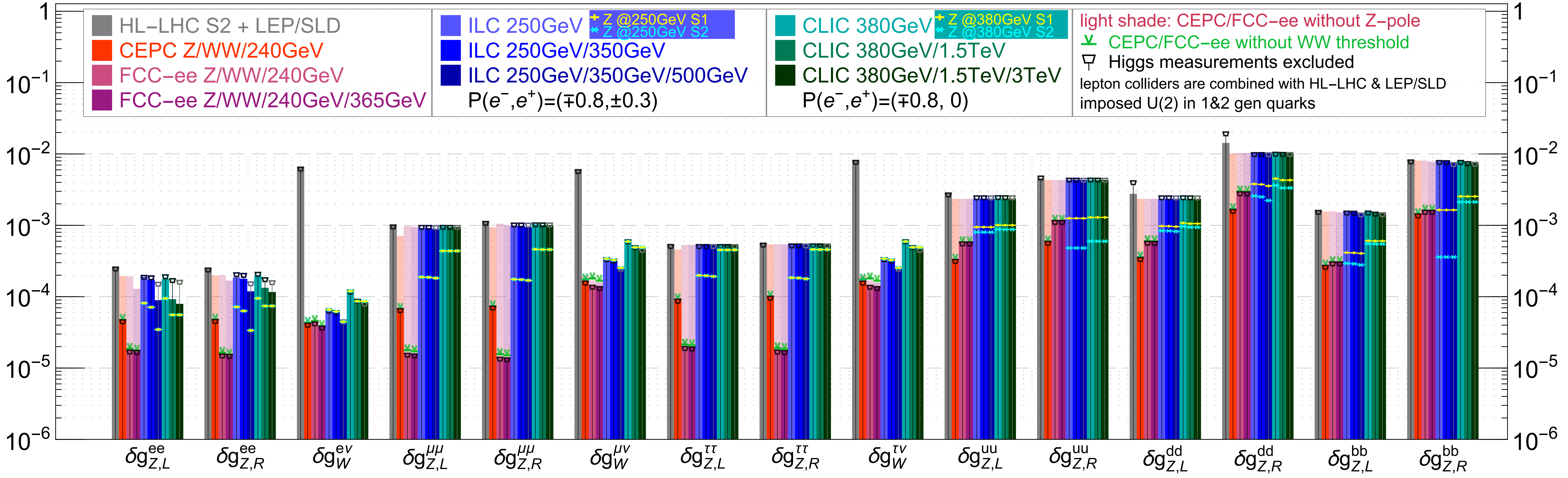
C- minimum of **$\Delta\Gamma_W$=0.96 MeV** +**Δm_W=0.41 MeV** with
E_1=**157.5** GeV E_2=**162.4** GeV f=**0.45**
yields and **Δm_W=0.37 MeV** (as single par)

Δm_W , $\Delta\Gamma_W$: error on W mass and width from fitting both
Δm_W : error on W mass from fitting only m_W

FCC workshop - 27 Jan 2023 P.Azzurri - W mass and width

Sensitivity on EW couplings.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



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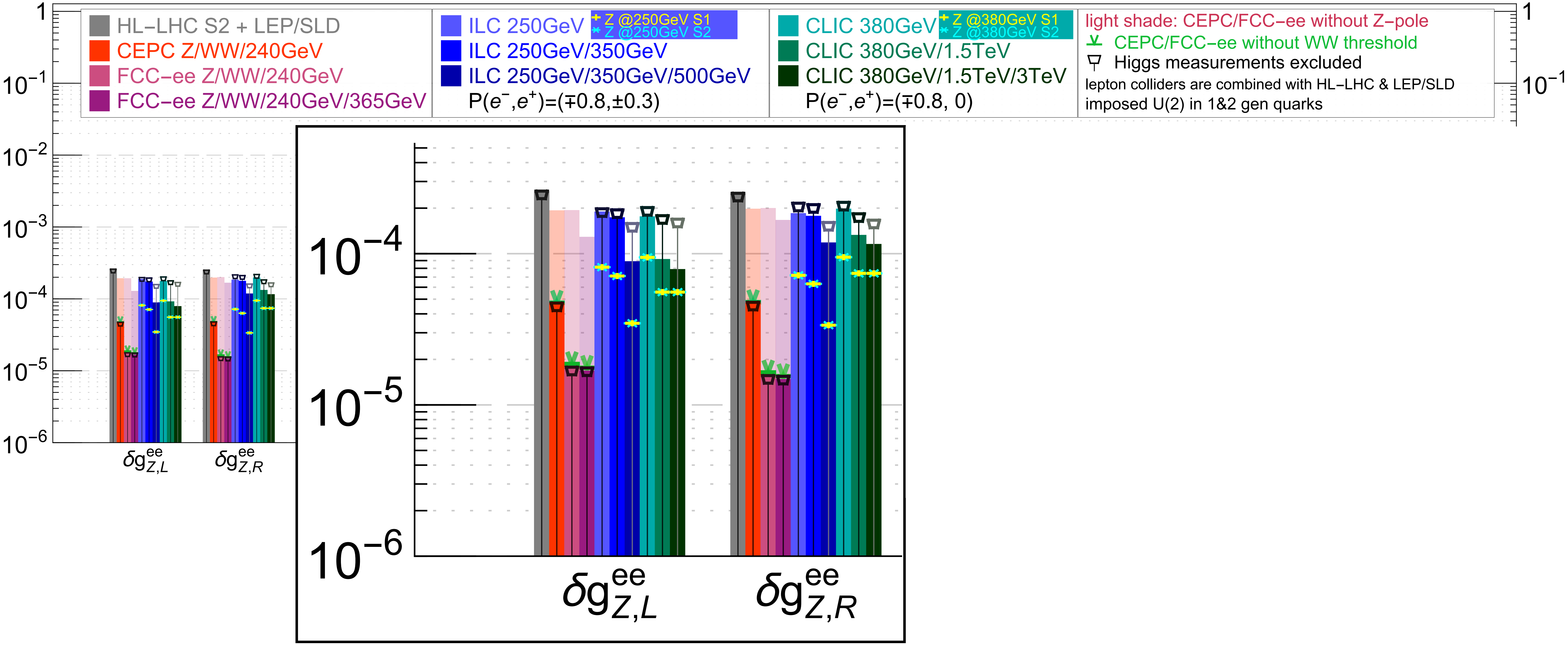

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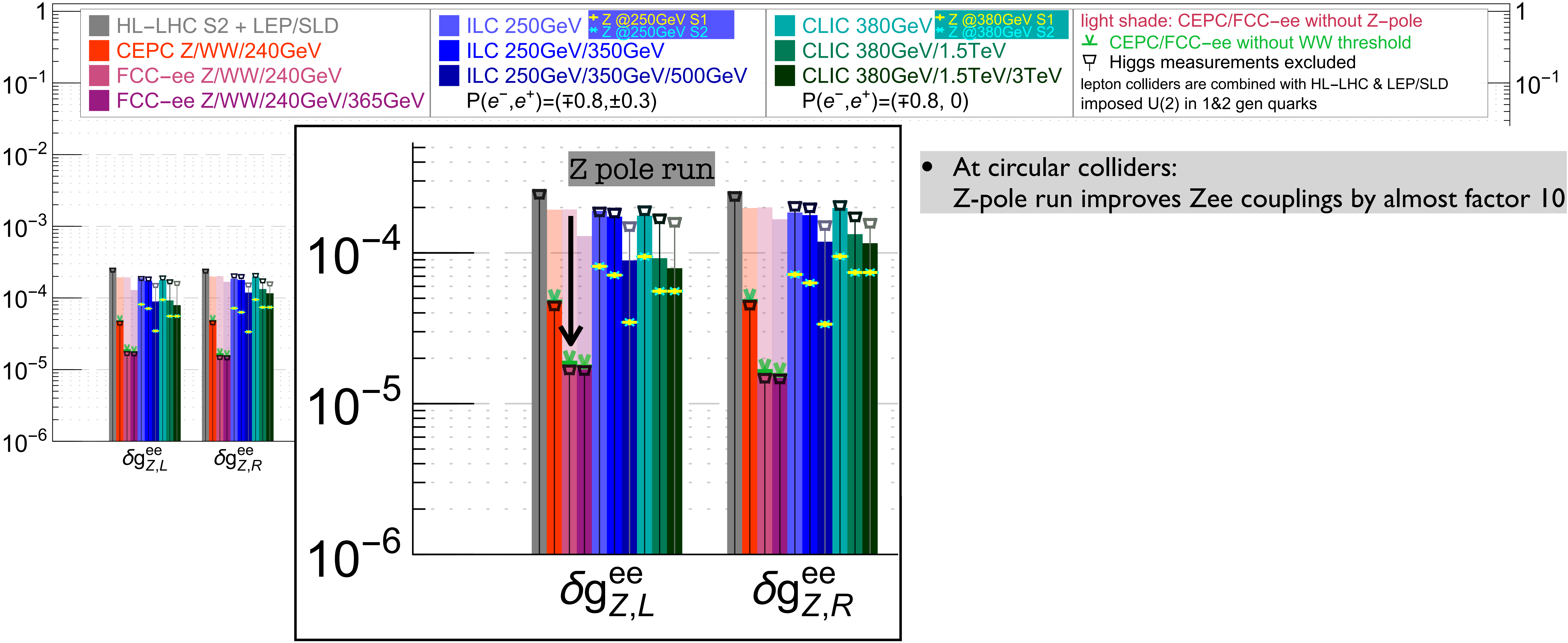

- At circular colliders:
  Z-pole run improves Zee couplings by almost factor 10

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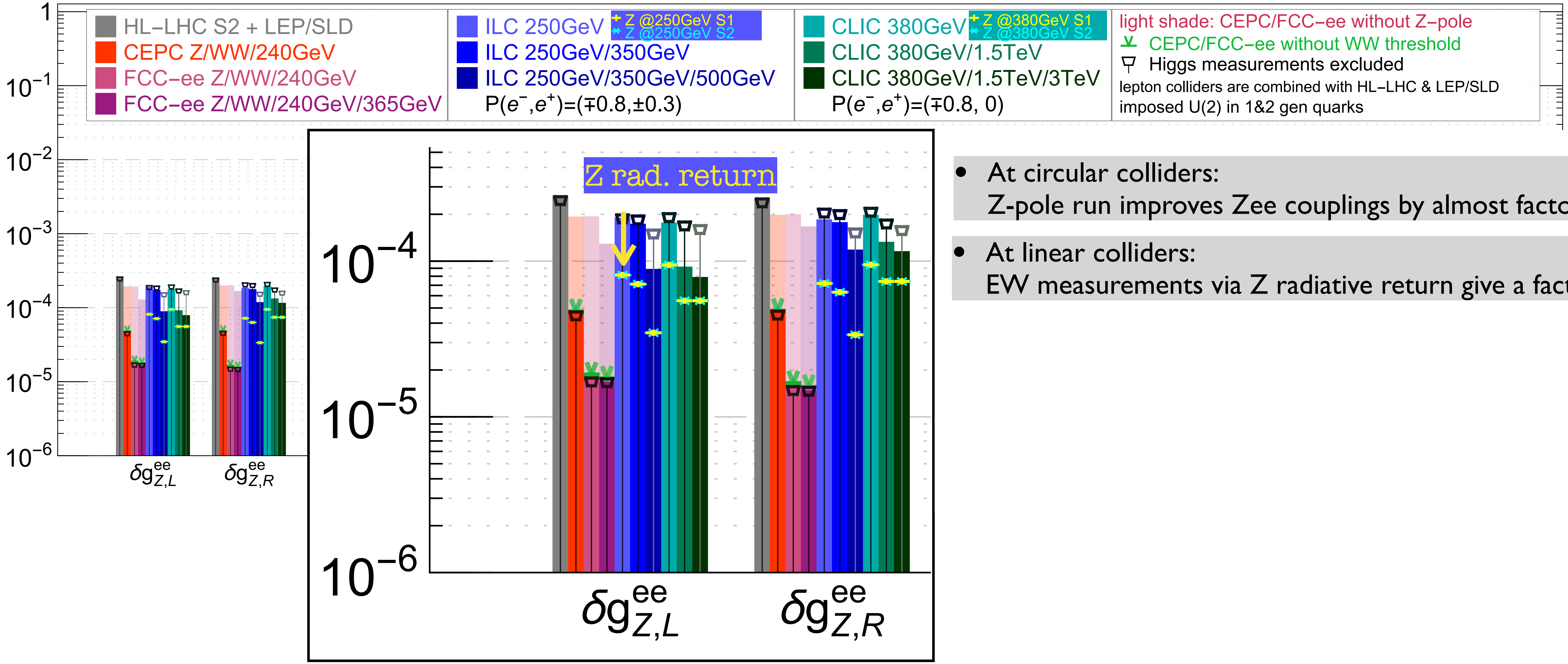

- At circular colliders:
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- At linear colliders:
  EW measurements via Z radiative return give a factor 3

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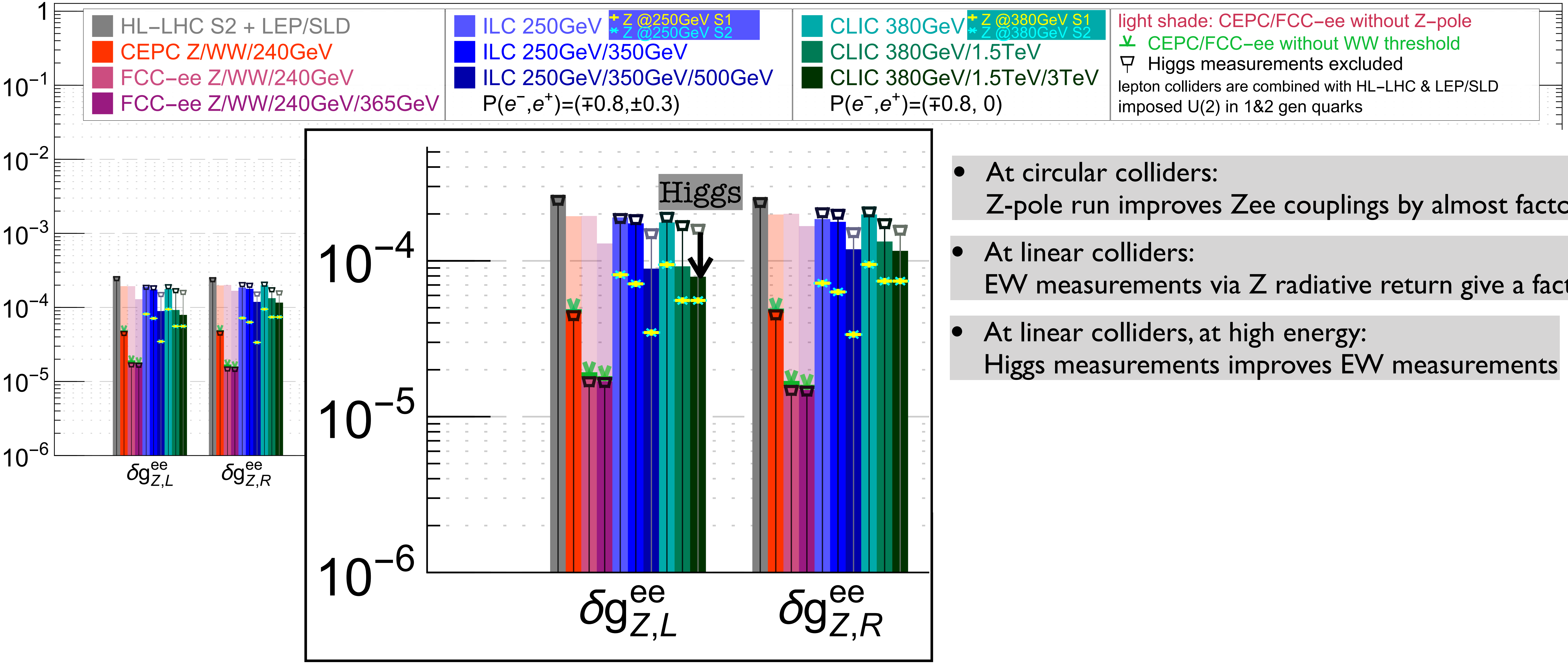
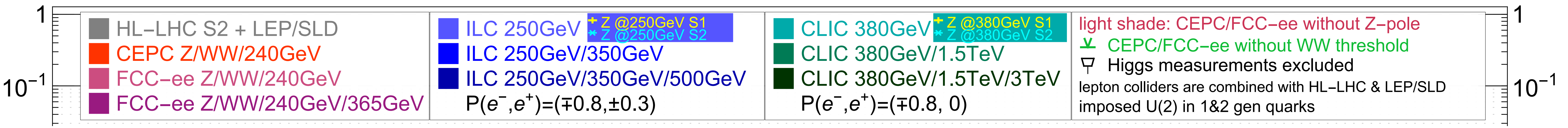


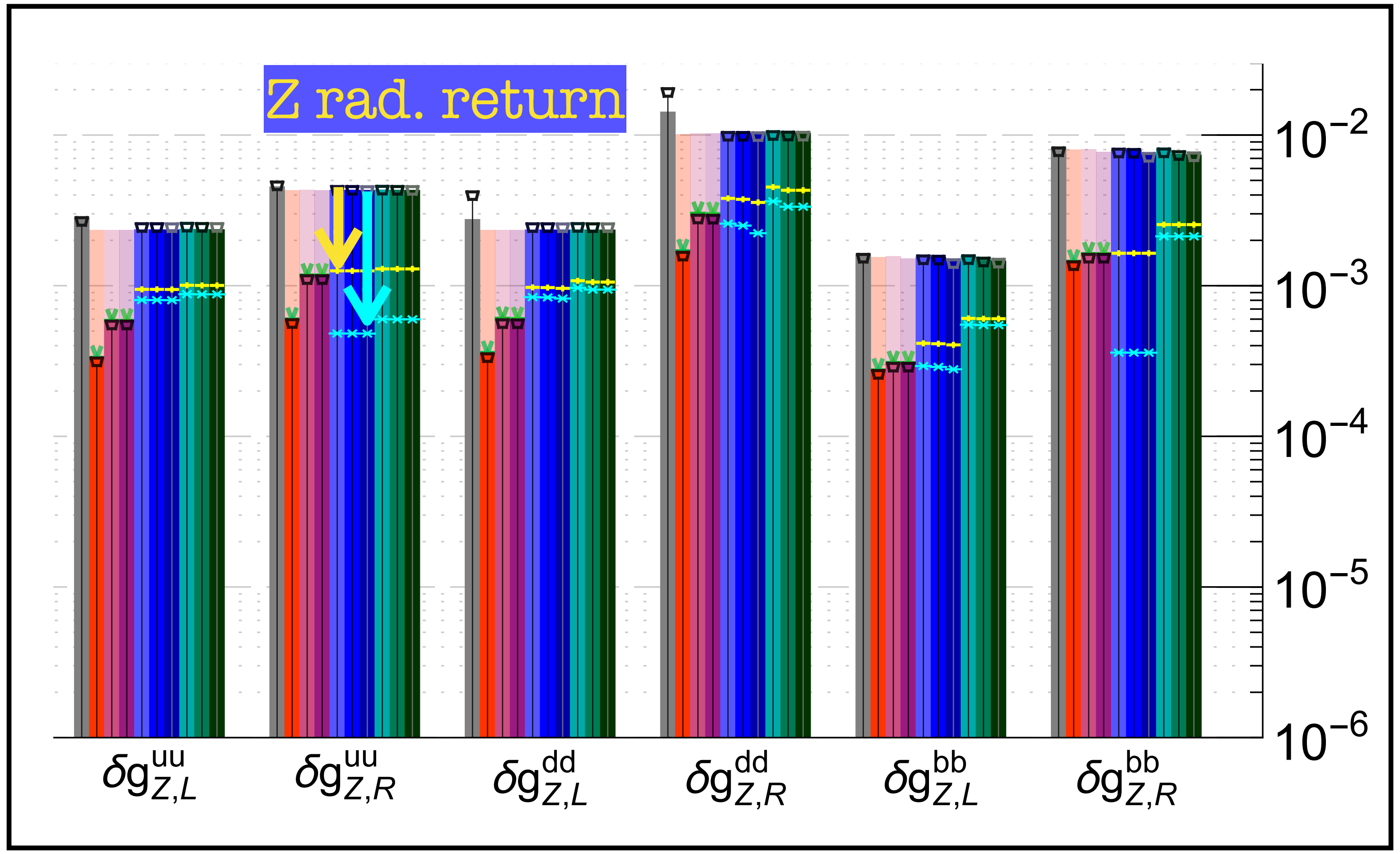
- At circular colliders:
  Z-pole run improves Zee couplings by almost factor 10
- At linear colliders:
  EW measurements via Z radiative return give a factor 3
- At linear colliders, at high energy:
  Higgs measurements improves EW measurements

Sensitivity on EW couplings.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311

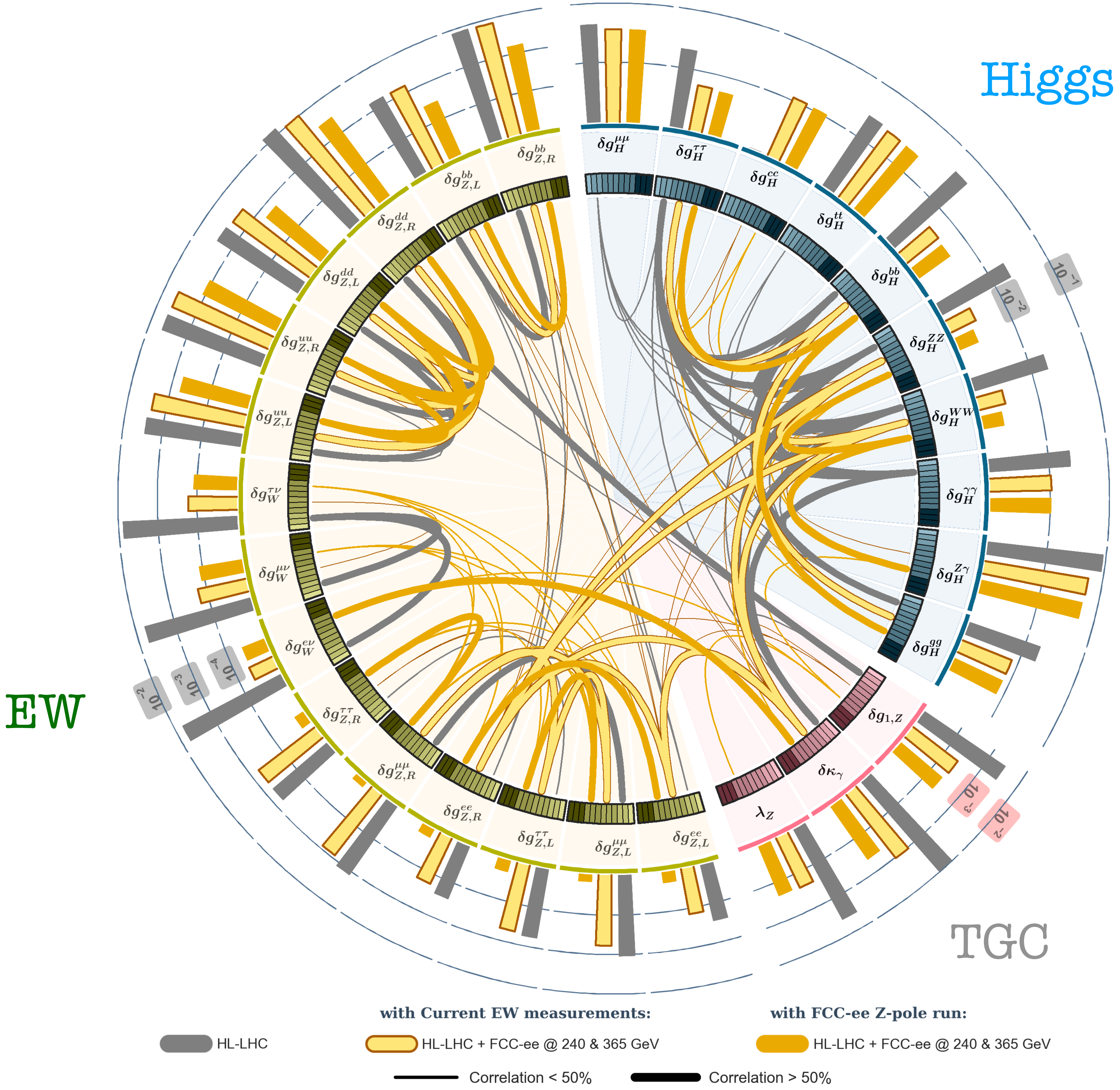


- At linear colliders, at high energy:
  EW measurements via Z-radiative return has a large impact on Zqq couplings
- Improvements depend a lot on hypothesis on systematic uncertainties
  - Yellow: LEP/SLD systematics / 2
  - Blue: small EXP and TH systematics



Why Z-pole for Higgs?

J. De Blas et al. 1907.04311



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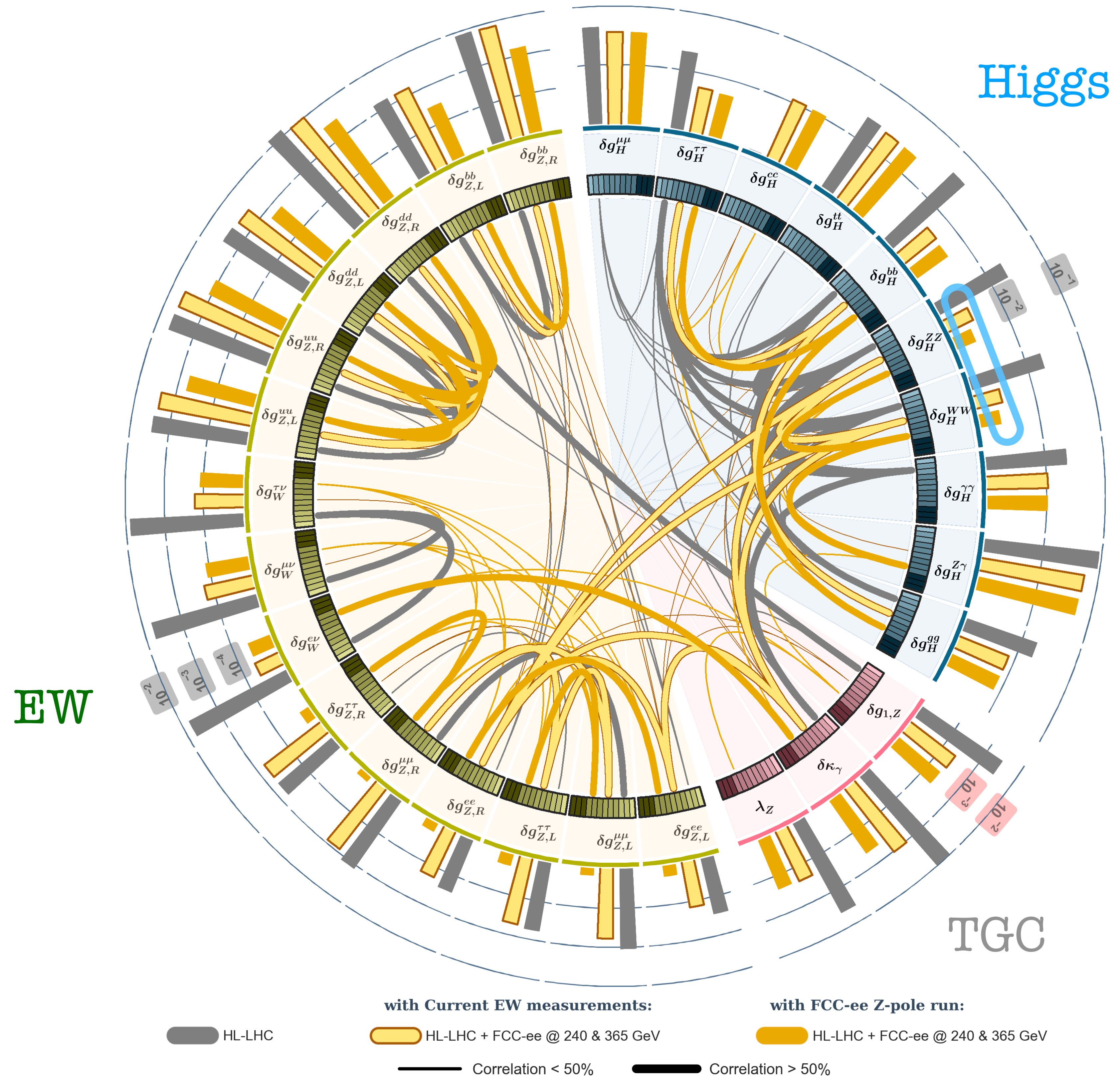

J. De Blas et al. 1907.04311

With Z-pole measurements,
Higgs coupling determination improves
by up to 50%

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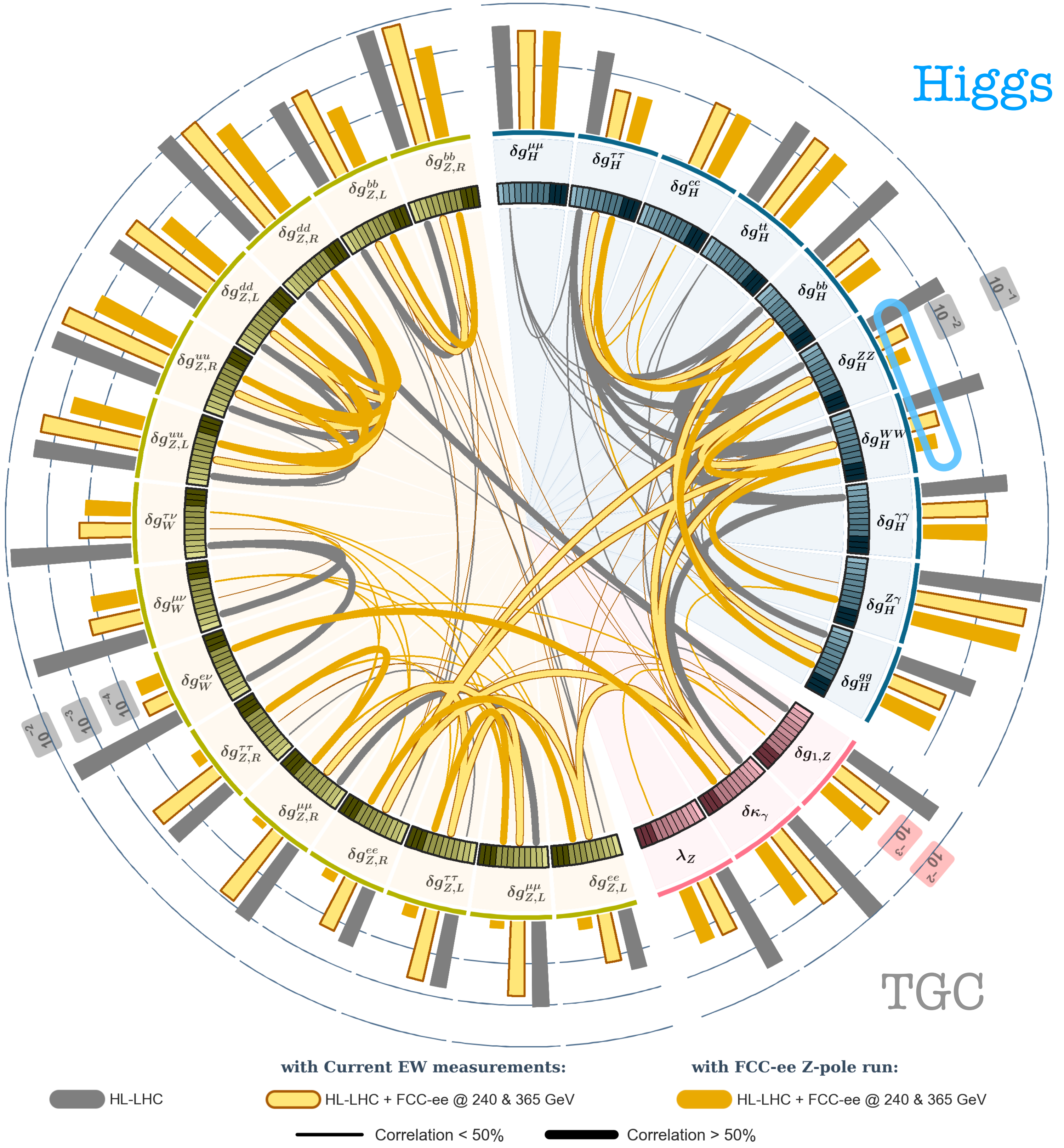

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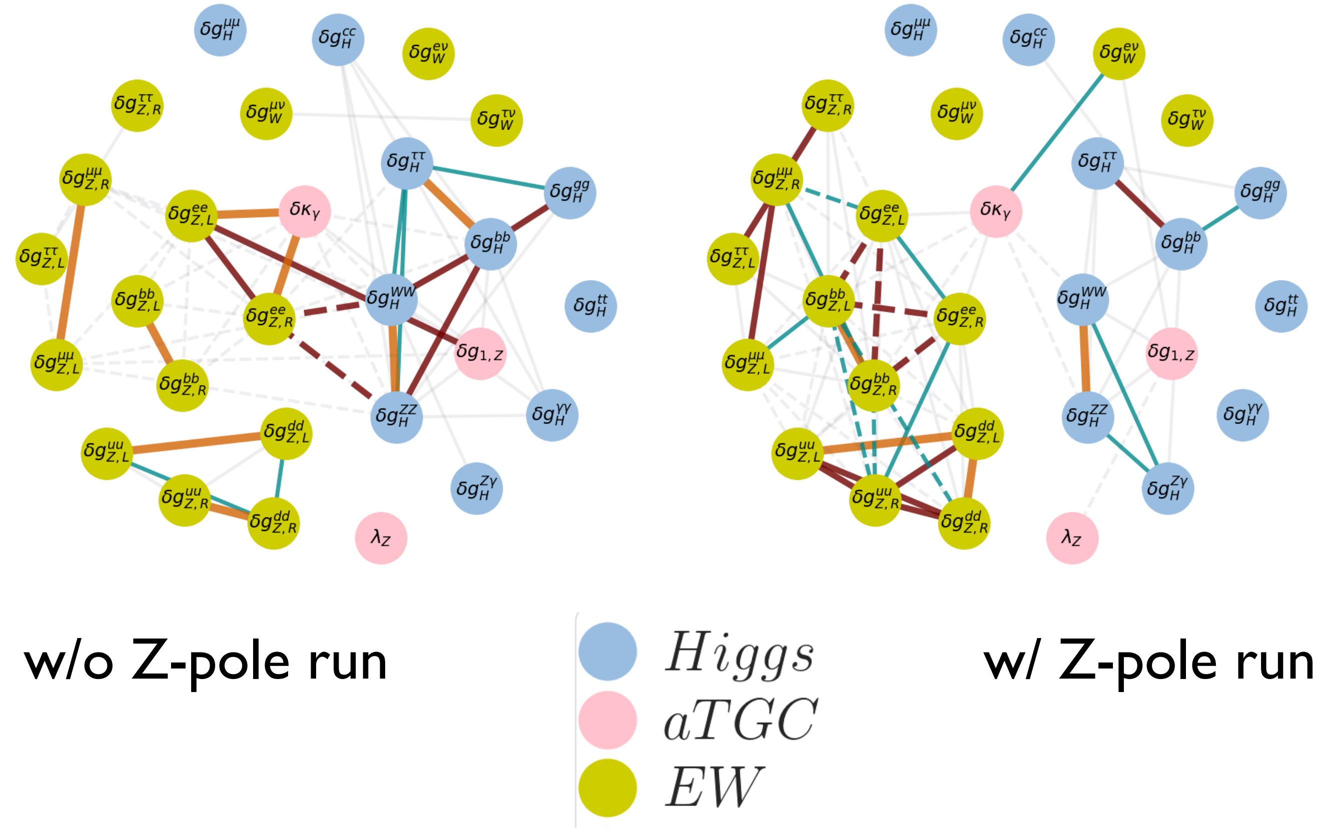
With Z-pole measurements, Higgs coupling determination improves by up to 50%

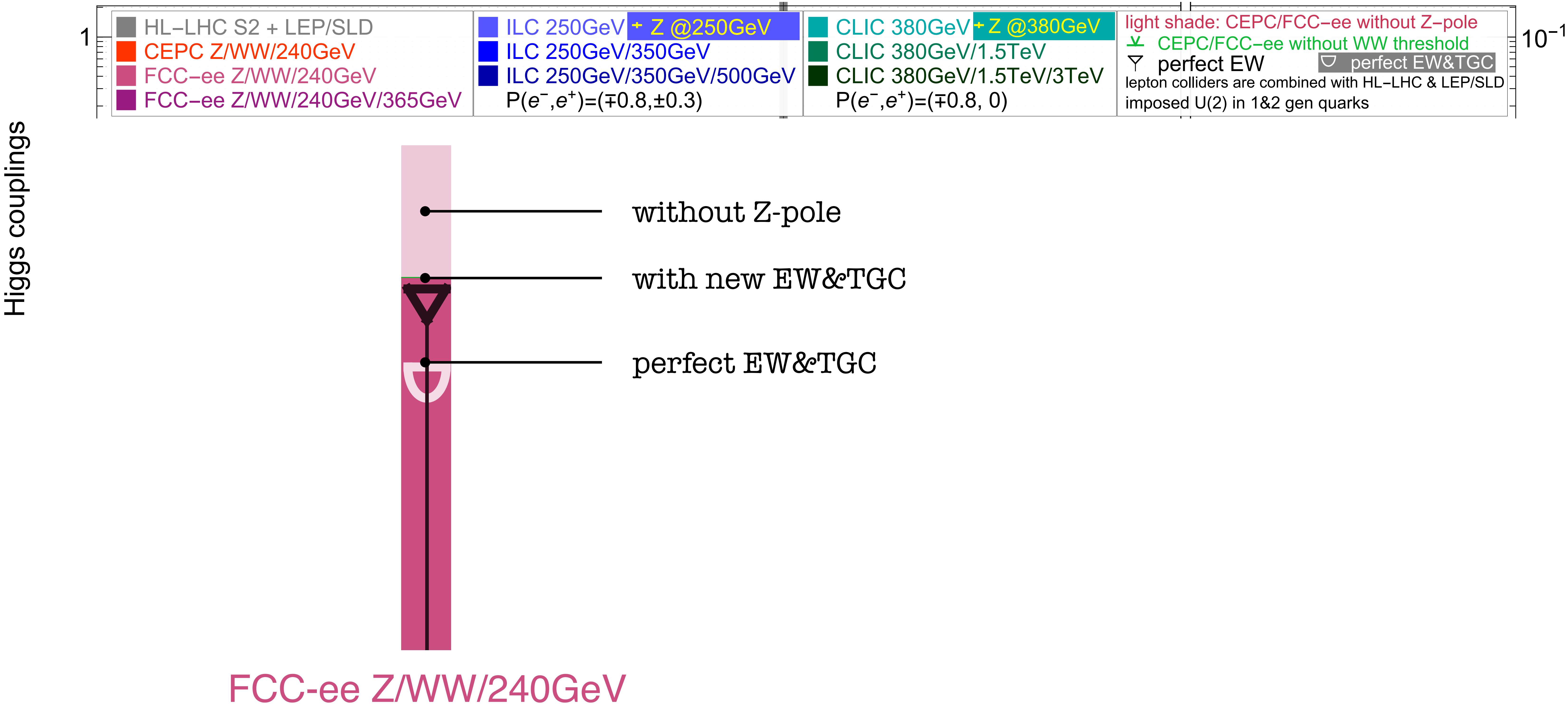
Z-pole run at circular colliders decorrelates EW and Higgs sectors from each other



Impact of Z-pole on Higgs.

J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements



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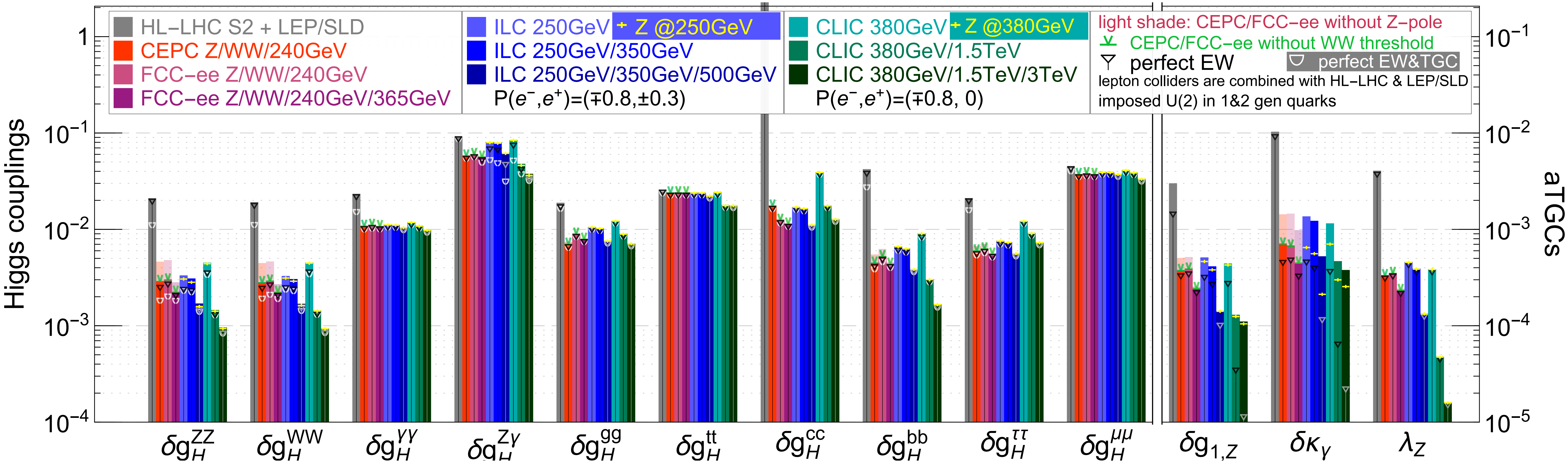

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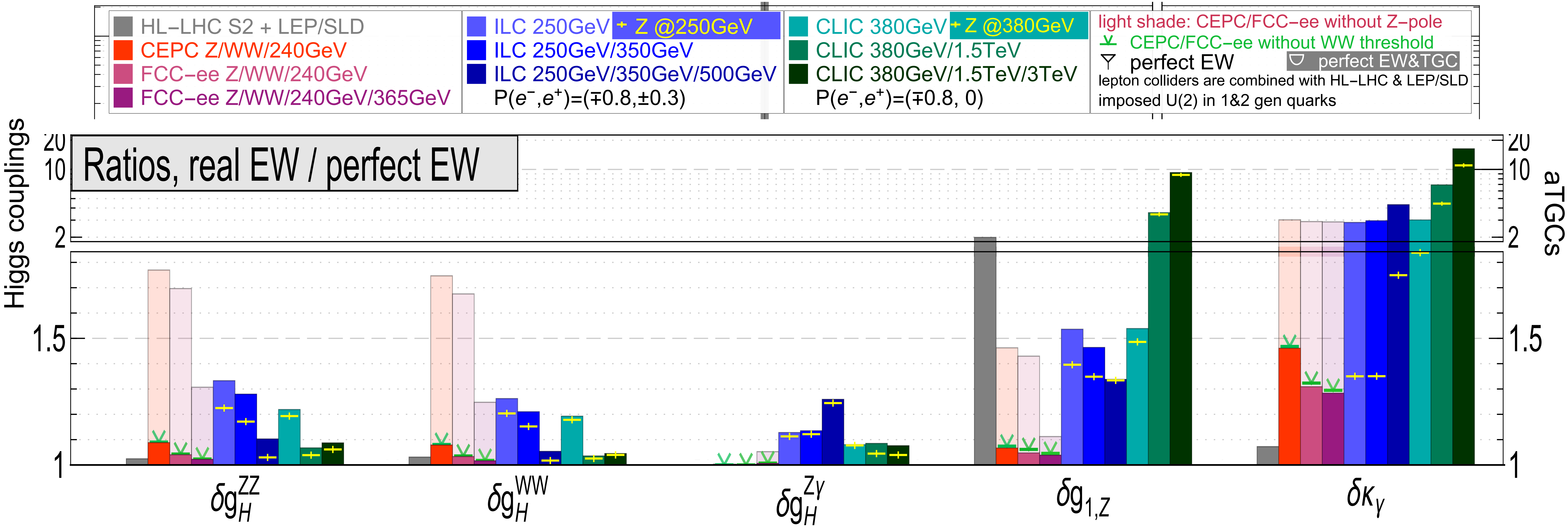

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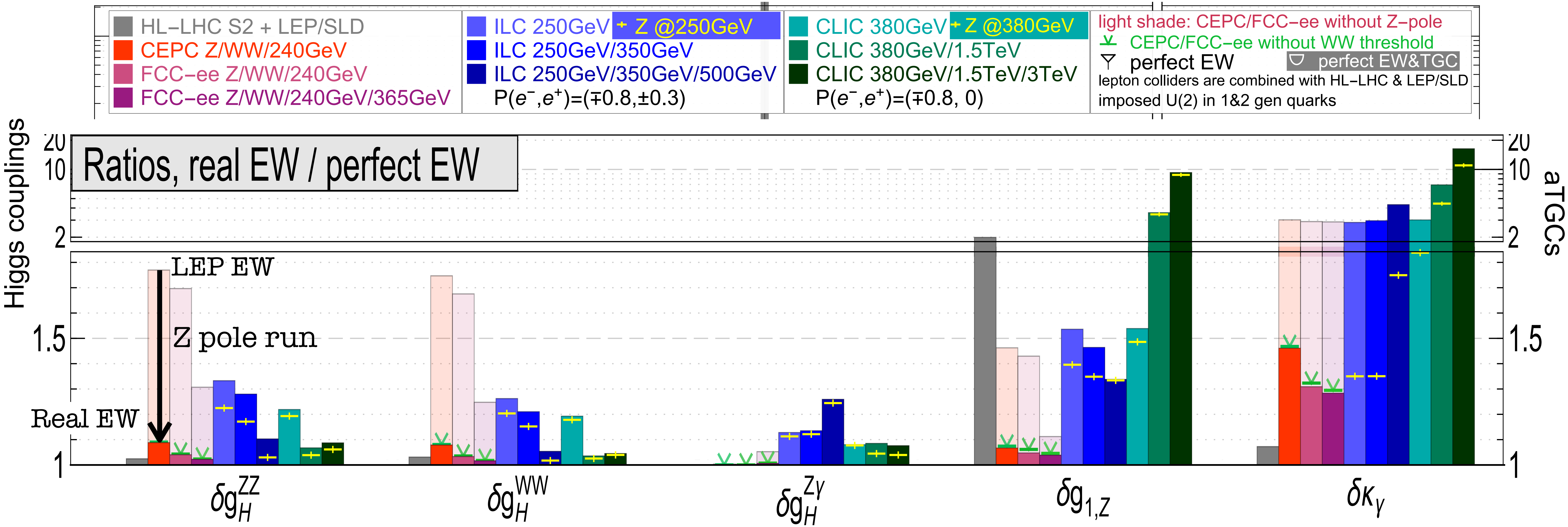

- FCC-ee and CEPC benefit a lot (>50% on HVV) from Z-pole run
- FCC-ee and CEPC EW measurements are almost perfect for what concerns Higgs physics (<10%).

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J. De Blas et al. 1907.04311

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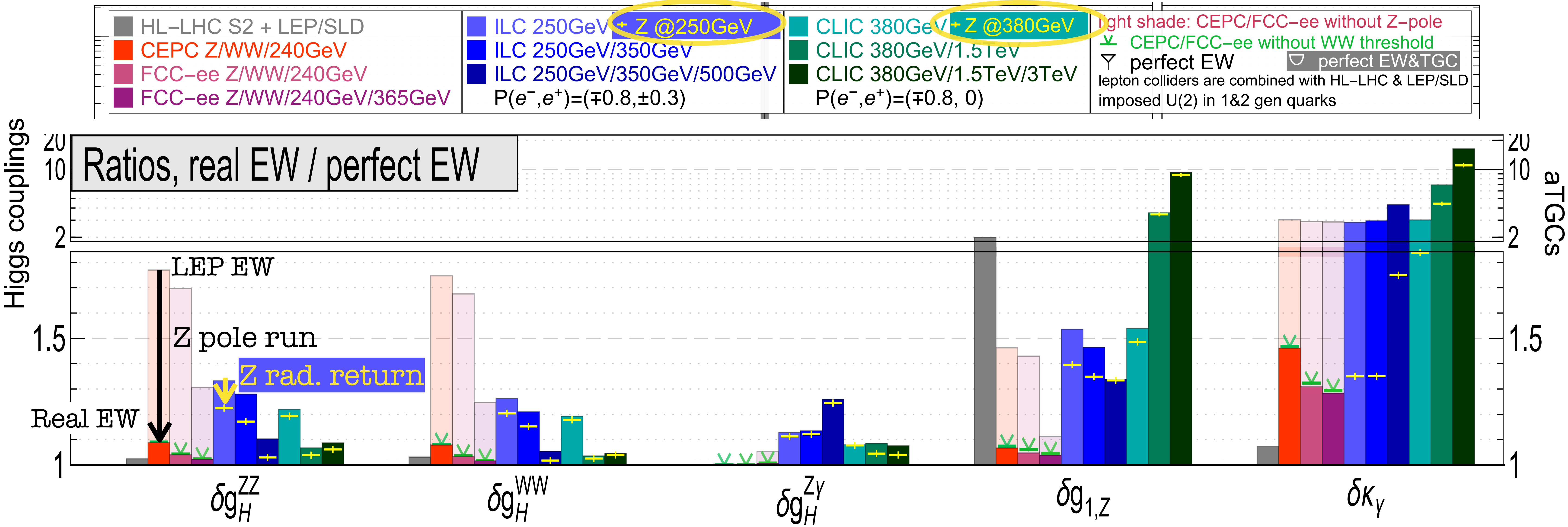
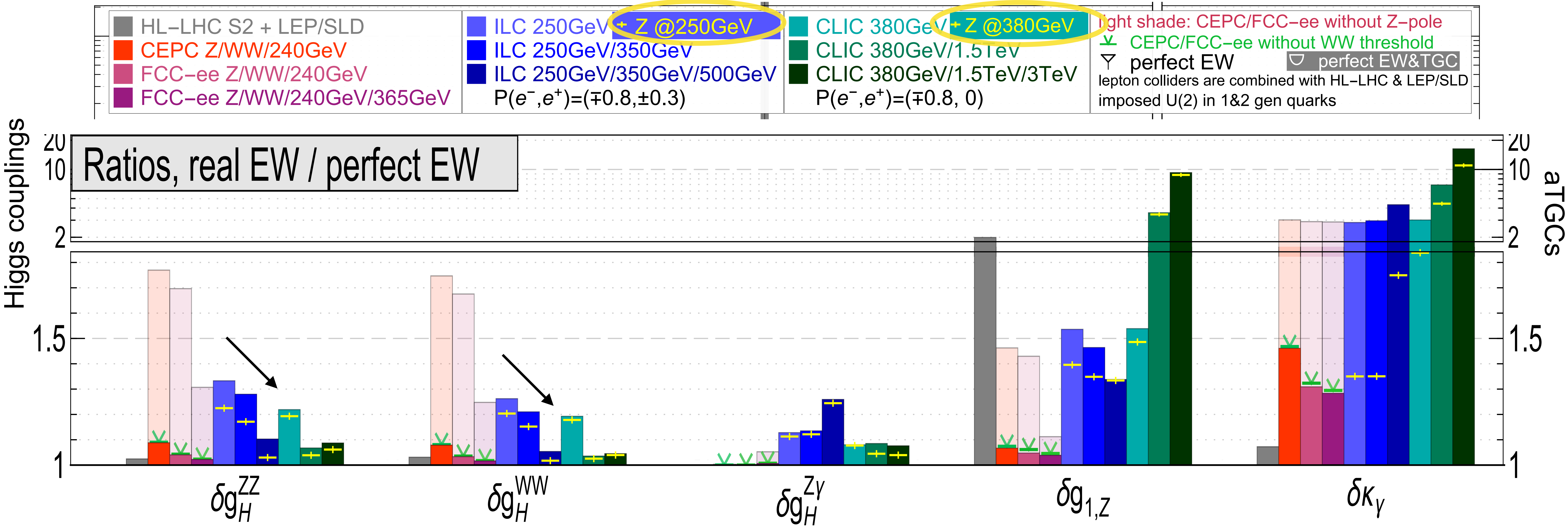


- FCC-ee and CEPC benefit a lot (>50% on HVV) from Z-pole run
- FCC-ee and CEPC EW measurements are almost perfect for what concerns Higgs physics (<10%).
- LEP EW measurements are a limiting factor (~30%) to Higgs precision at ILC, especially for the first runs But EW measurements at high energy (via Z-radiative return) help mitigating this issue

Impact of Z-pole on Higgs.

J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements

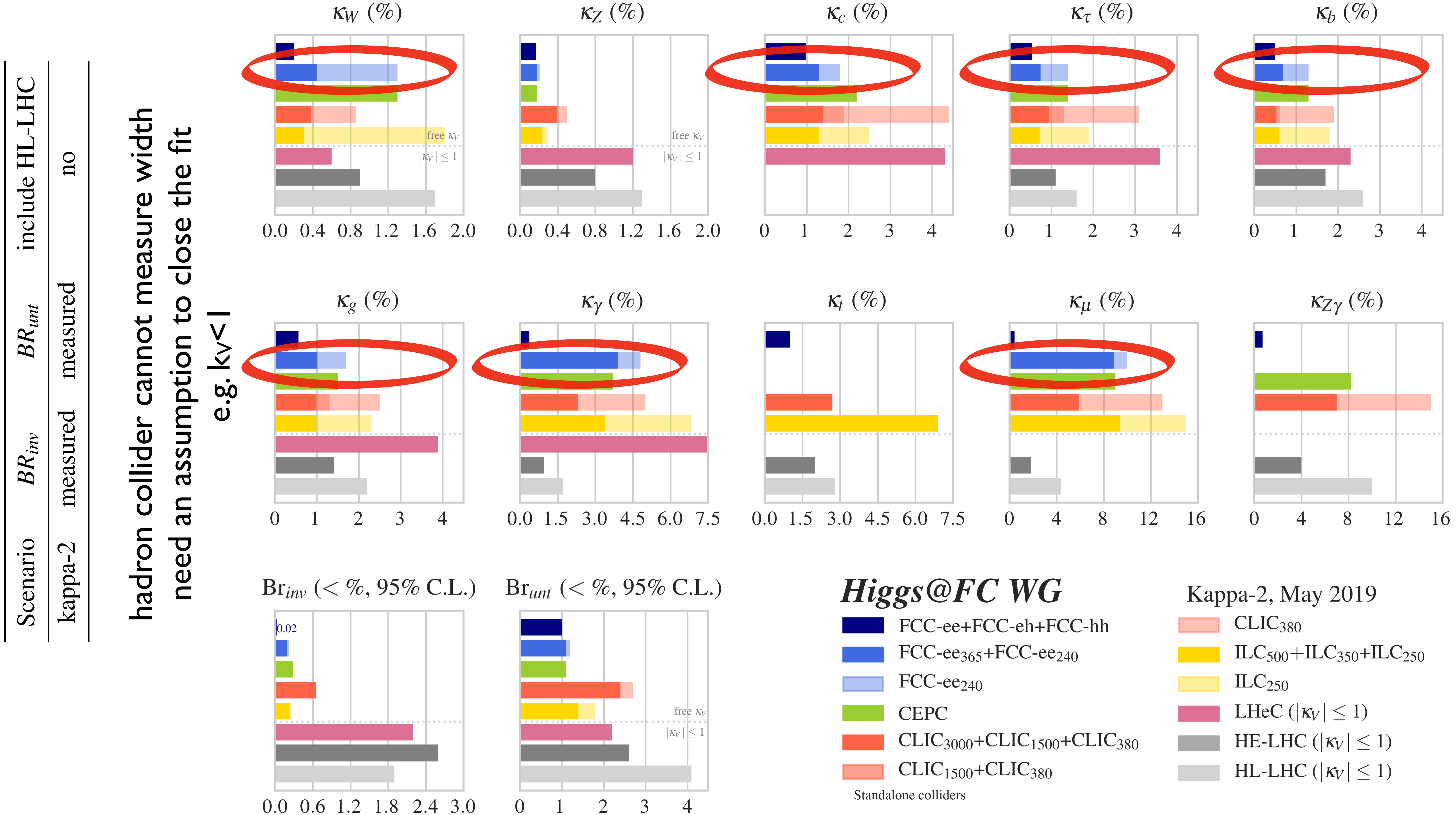


- Higher energy runs reduce the EW contamination in Higgs coupling extraction

Complementarity 240↔365 GeV.

ECFA Higgs study group '19

hadron collider cannot measure width
need an assumption to close the fit
e.g. $\kappa_V<1$



Complementarity FCC-ee↔HL-LHC.

ECFA Higgs study group '19

Scenario	BR_{inv}	BR_{unt}	include HL-LHC
kappa-2	measured	measured	no

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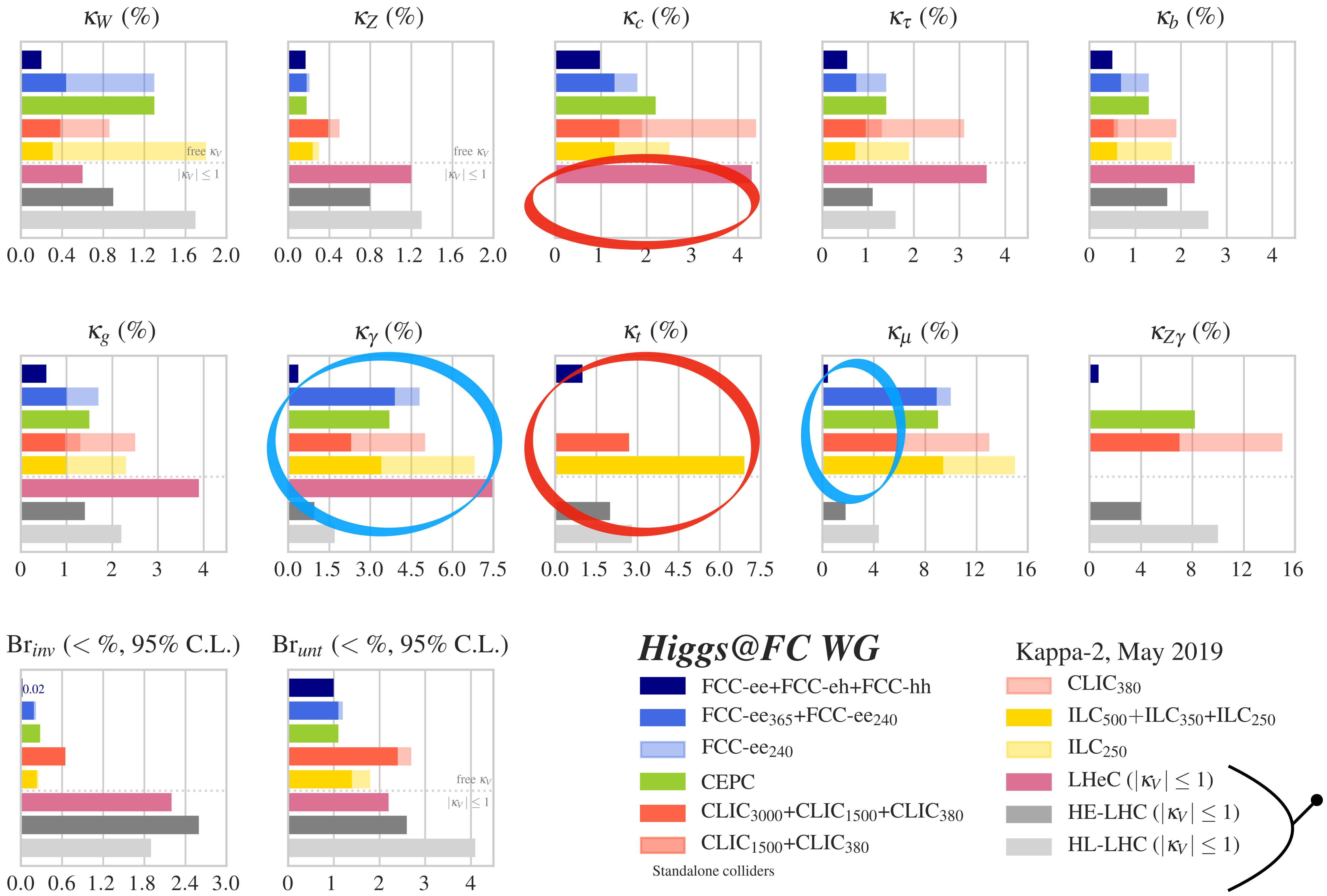

assumption needed for the fit to close at hadron machines

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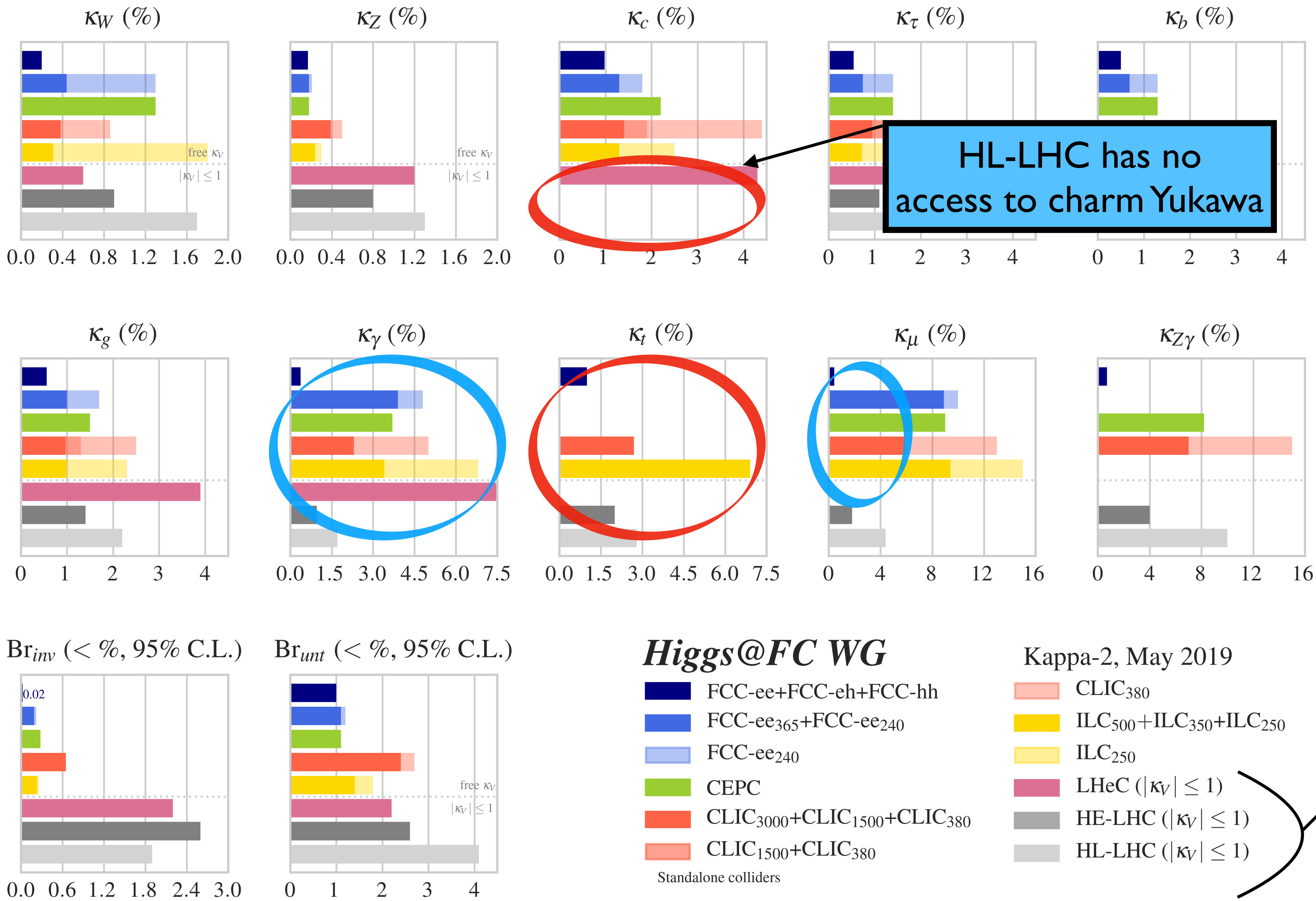

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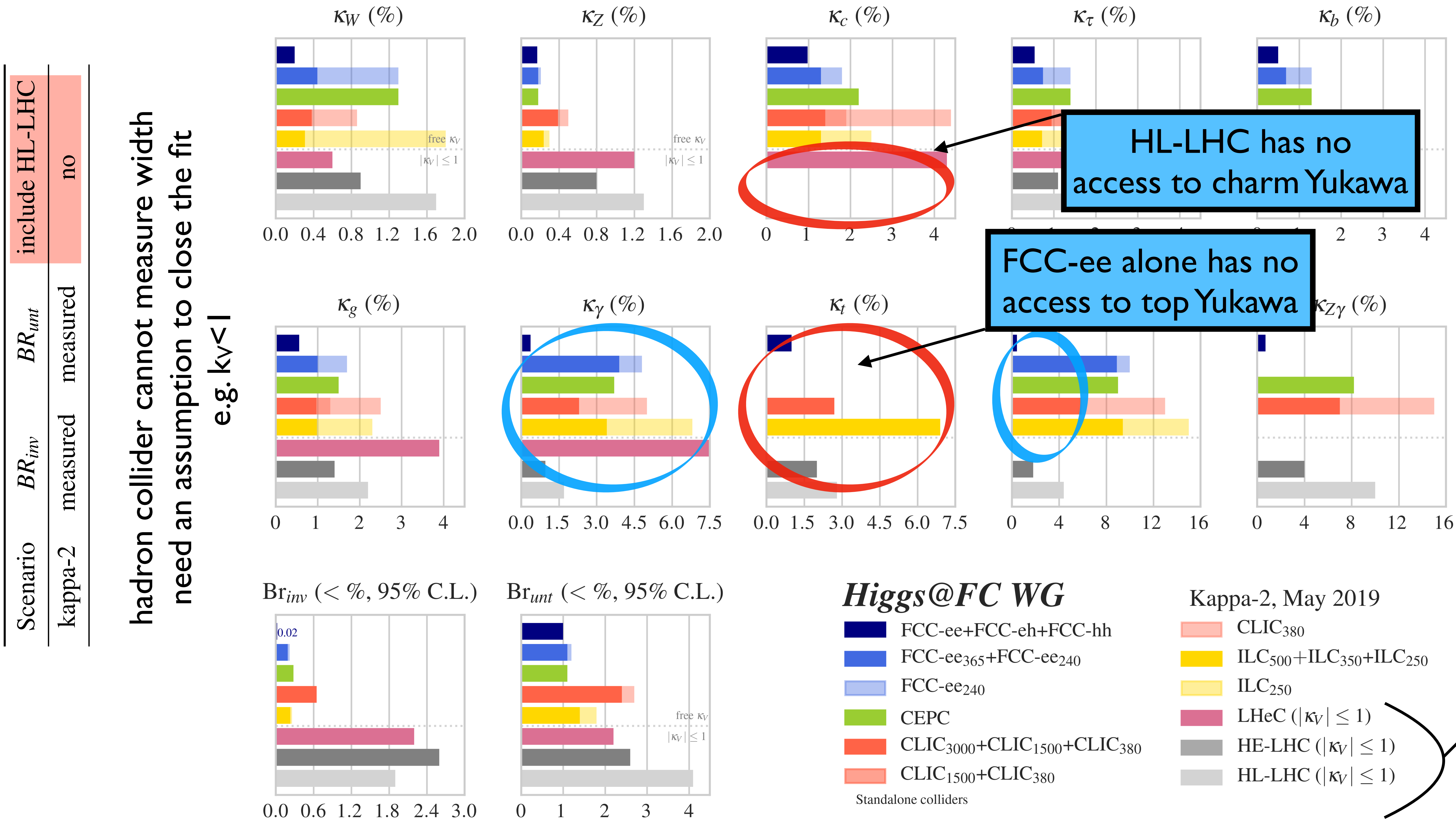

hadron collider cannot measure width
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e.g. κ_V<1

HL-LHC has no access to charm Yukawa

FCC-ee alone has no access to top Yukawa

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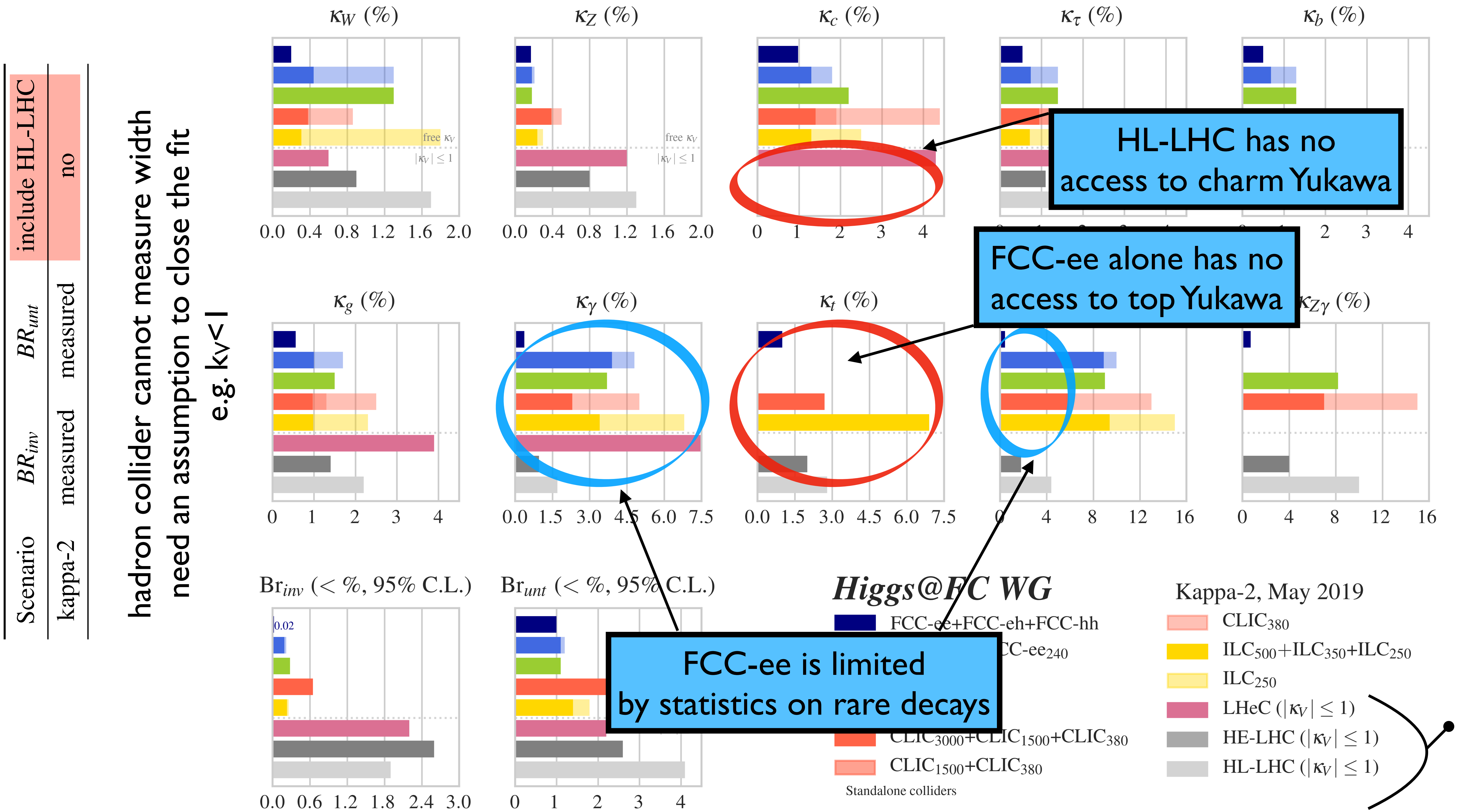


HL-LHC has no access to charm Yukawa

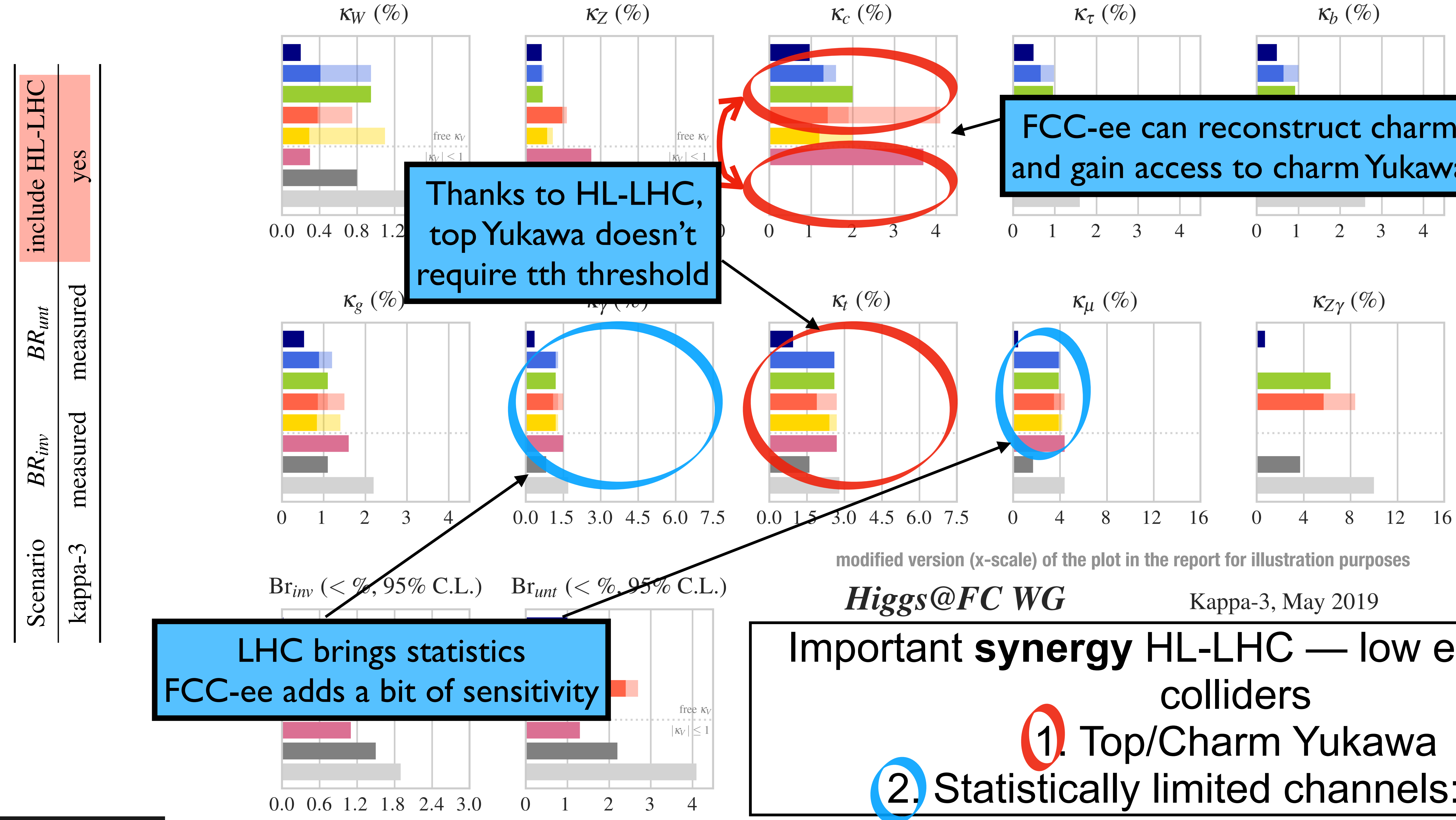
FCC-ee alone has no access to top Yukawa

FCC-ee is limited by statistics on rare decays

assumption needed for the fit to close at hadron machines

Complementarity FCC-ee↔HL-LHC.

ECFA Higgs study group '19



FCC-ee can reconstruct charm and gain access to charm Yukawa

Thanks to HL-LHC, top Yukawa doesn't require tth threshold

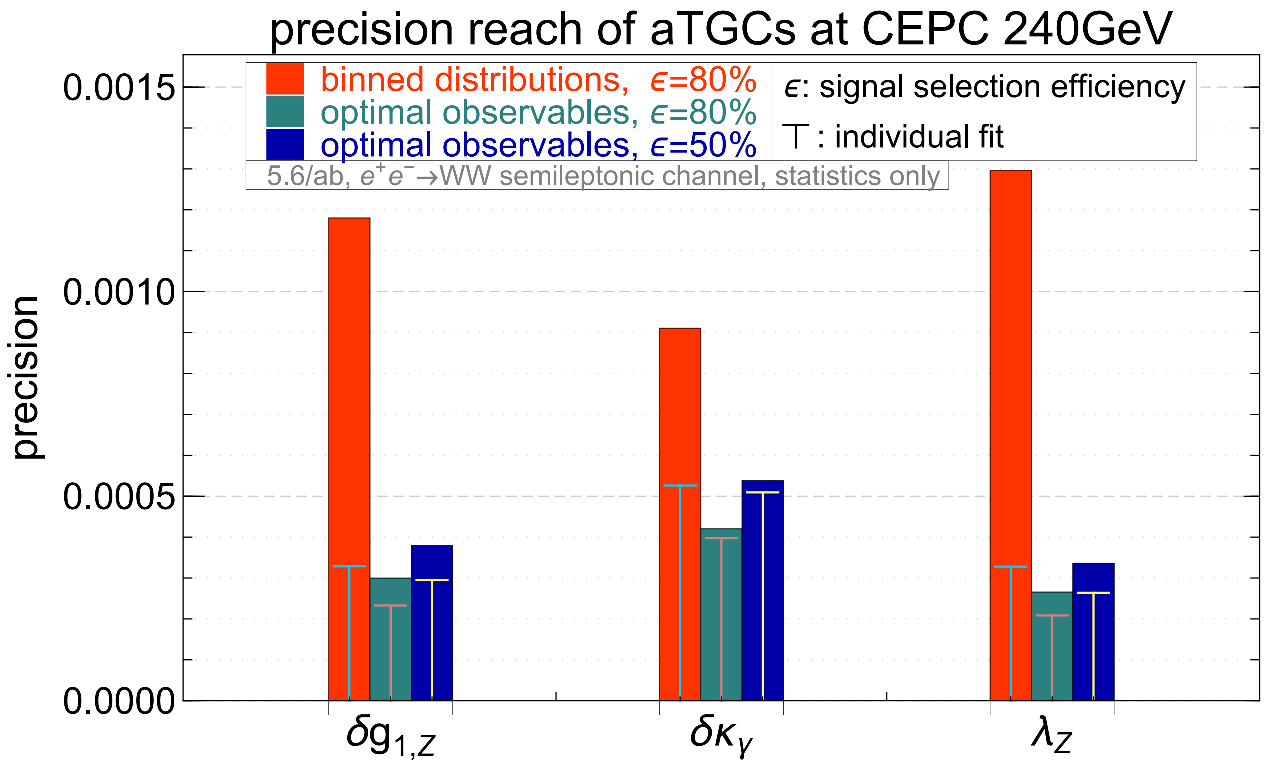
LHC brings statistics
FCC-ee adds a bit of sensitivity

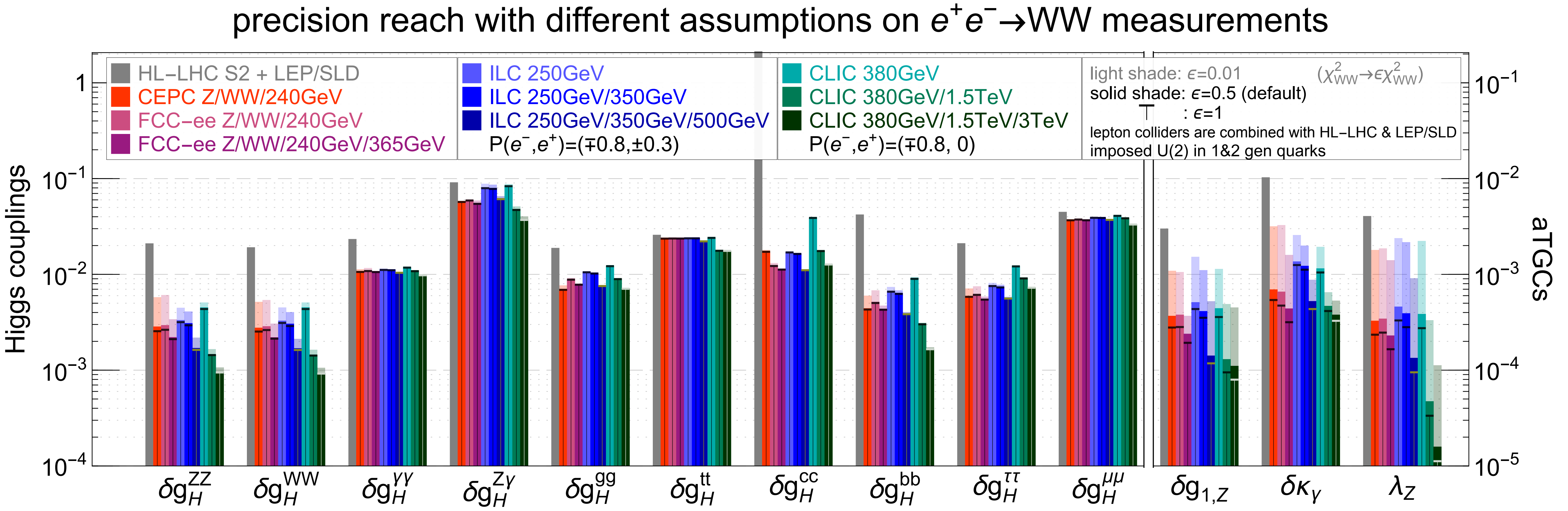
Important **synergy** HL-LHC — low energy lepton colliders

1. Top/Charm Yukawa
2. Statistically limited channels: γγ, μμ

Impact of Diboson Systematics.

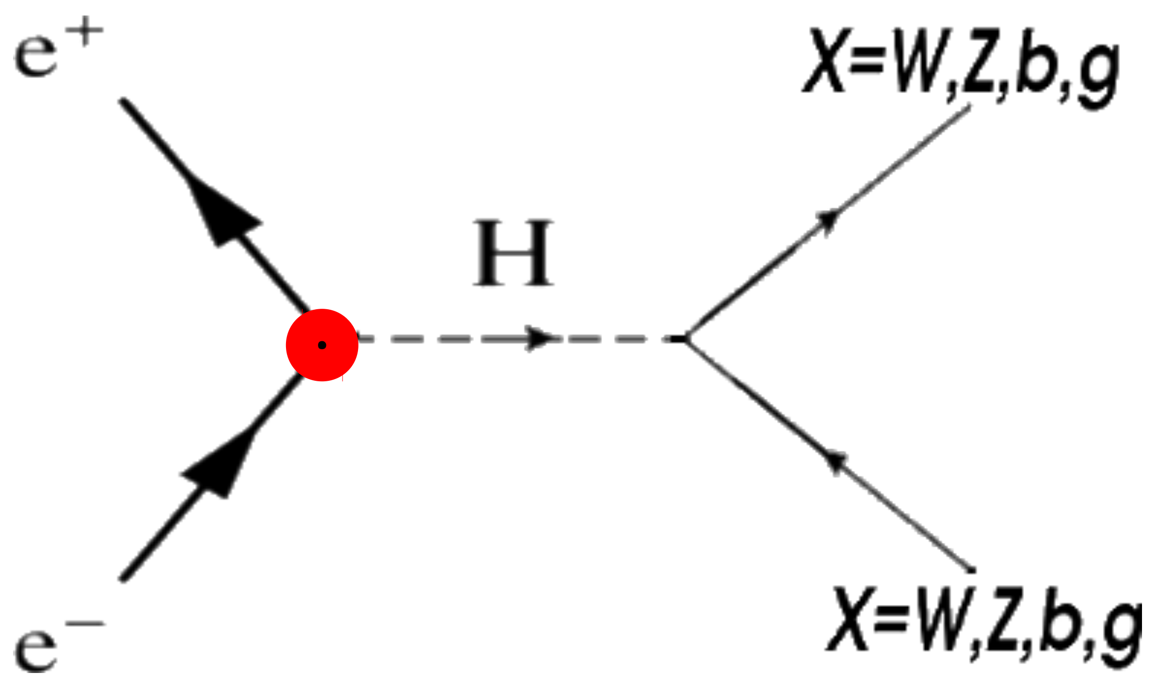
J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



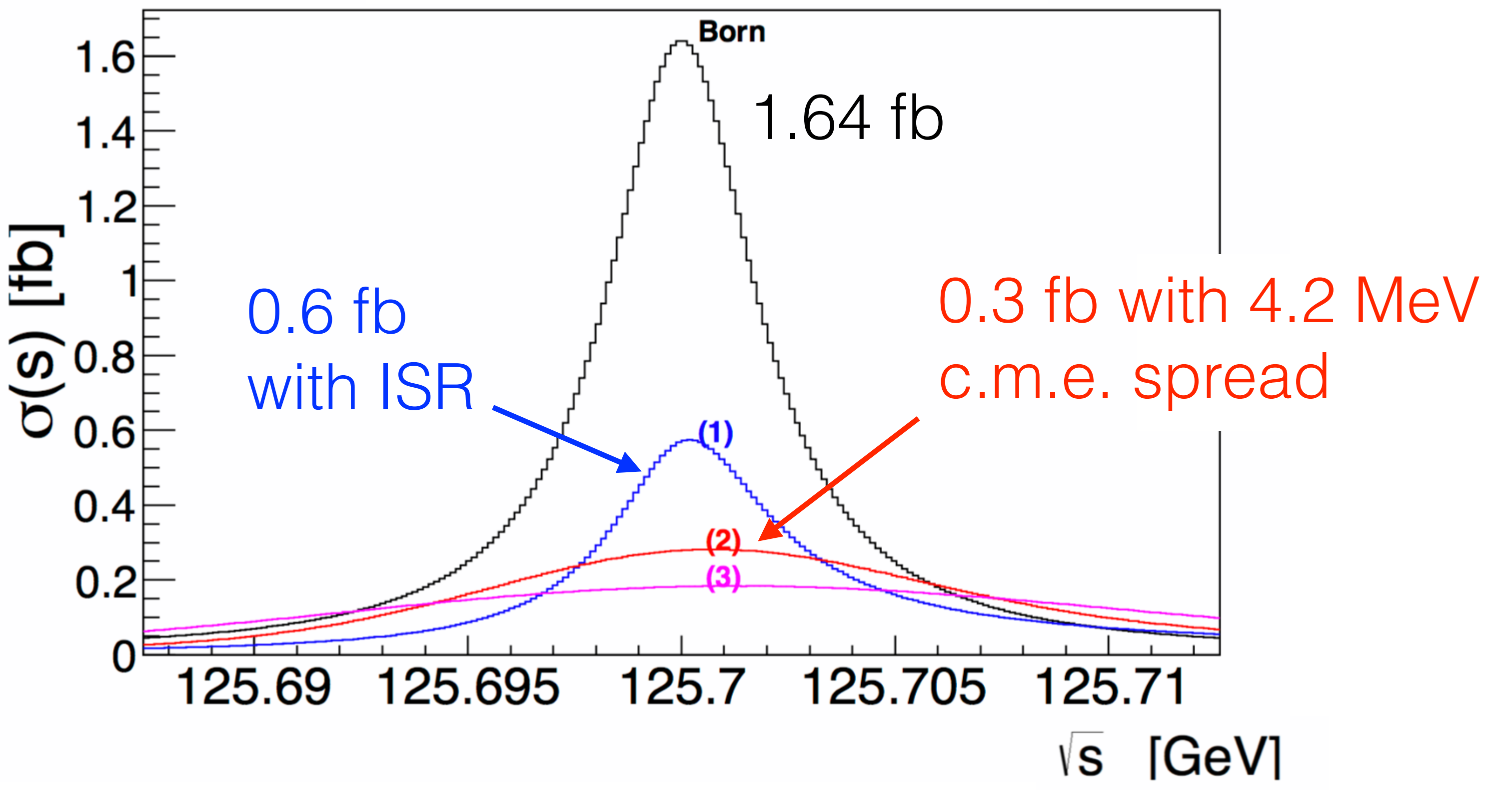


The stuff we are made of.

The high luminosity, the precise control of the beam √s, the clean reconstruction of final states make it possible to observe



Jadach+, arXiv: 1509.02406



$\sigma(e^+e^- \to H) = 1.64$ fb

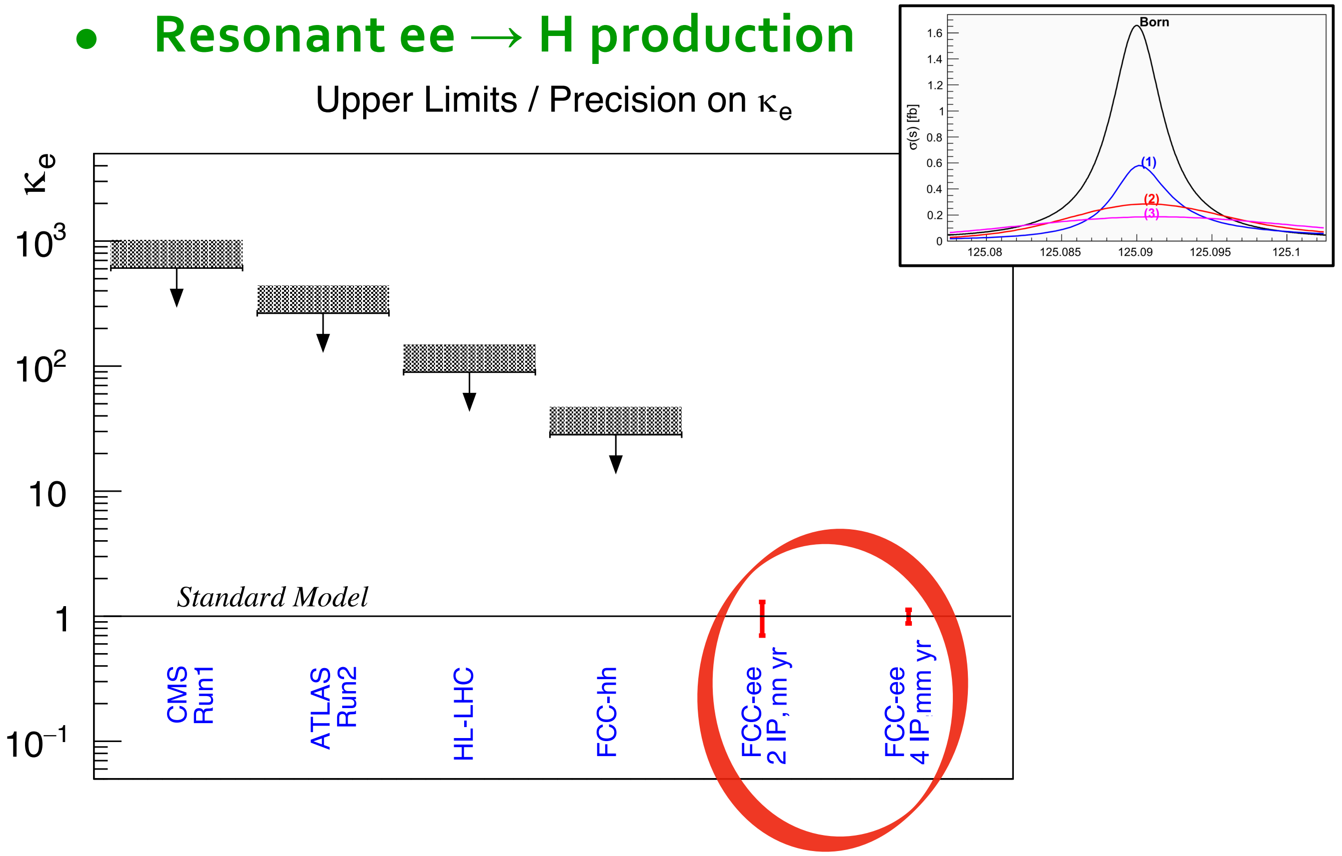
$\sigma_{spread+ISR}(e^+e^- \to H) = 0.17 \times \sigma(e^+e^- \to H) = 290$ ab

that the Higgs is responsible for
the mass of the stable elementary particles ordinary matter is made of

The stuff we are made of.

The high luminosity, the precise control of the beam √s, the clean reconstruction of final states make it possible to observe

- **20 ab^{-1} / year at $\sqrt{s}$ = 125 GeV** *(not in baseline FCC-ee)*
- **Monochromatization $\sigma_{\sqrt{s}} \sim 1\text{-}2 \times \Gamma_H \sim$ 6 to 10 MeV**
  - **Resonant ee → H production**



The stuff we are made of.

The high luminosity, the precise control of the beam √s, the clean reconstruction of final states make it possible to observe

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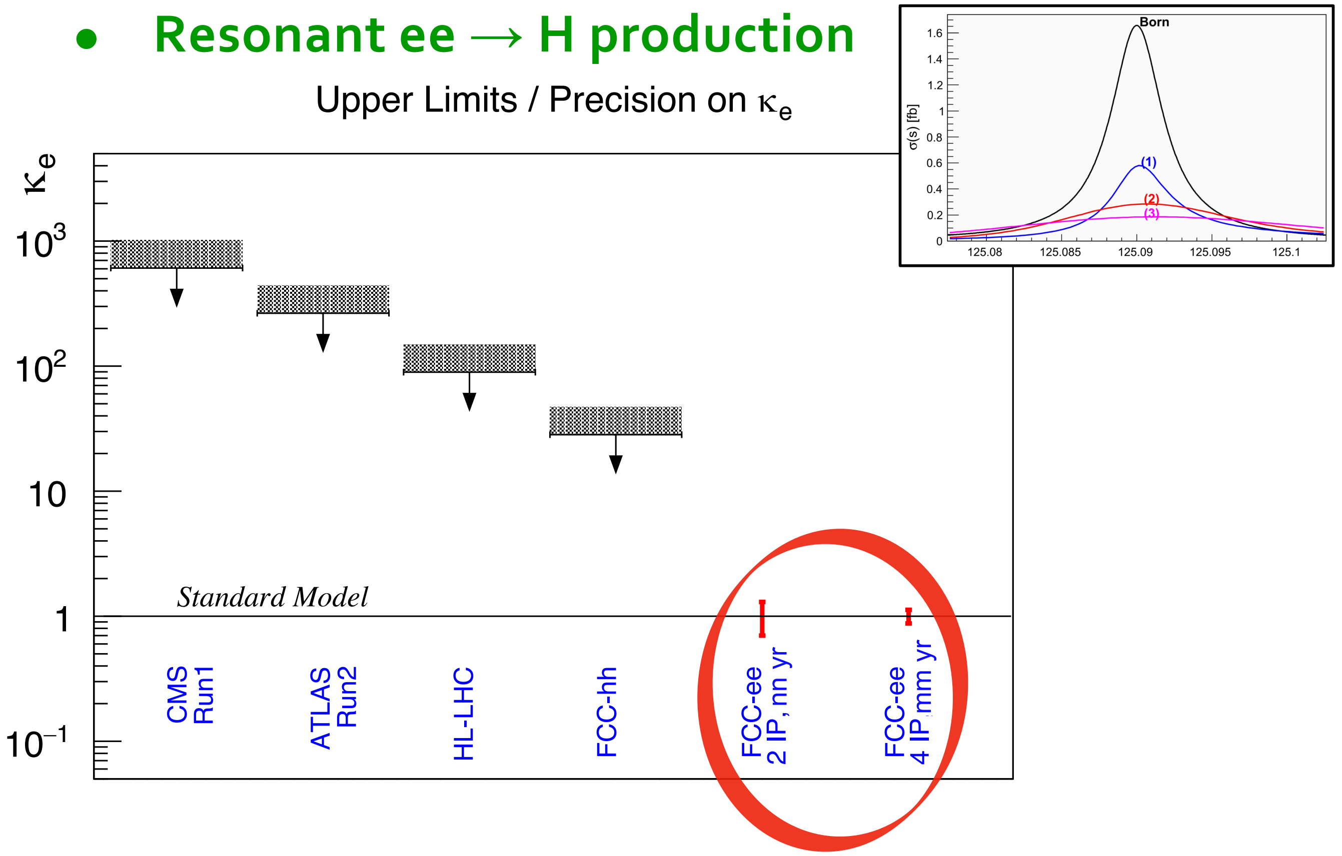
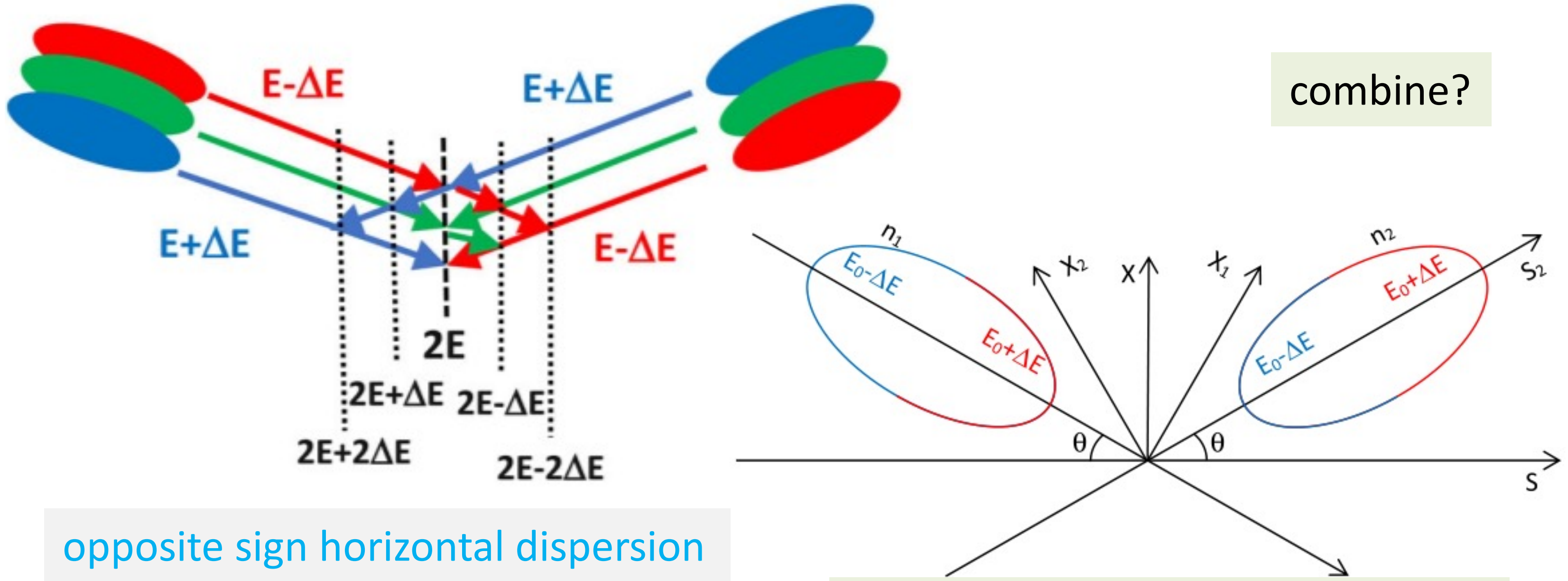


Monochromatisation

Monochromatization: **UNDER STUDY**
taking advantage of the separate e+ and e- rings, one can distribute in opposite way high and low energies in the beam (in x, z time)



opposite sign horizontal dispersion

opposite difference in arrival time

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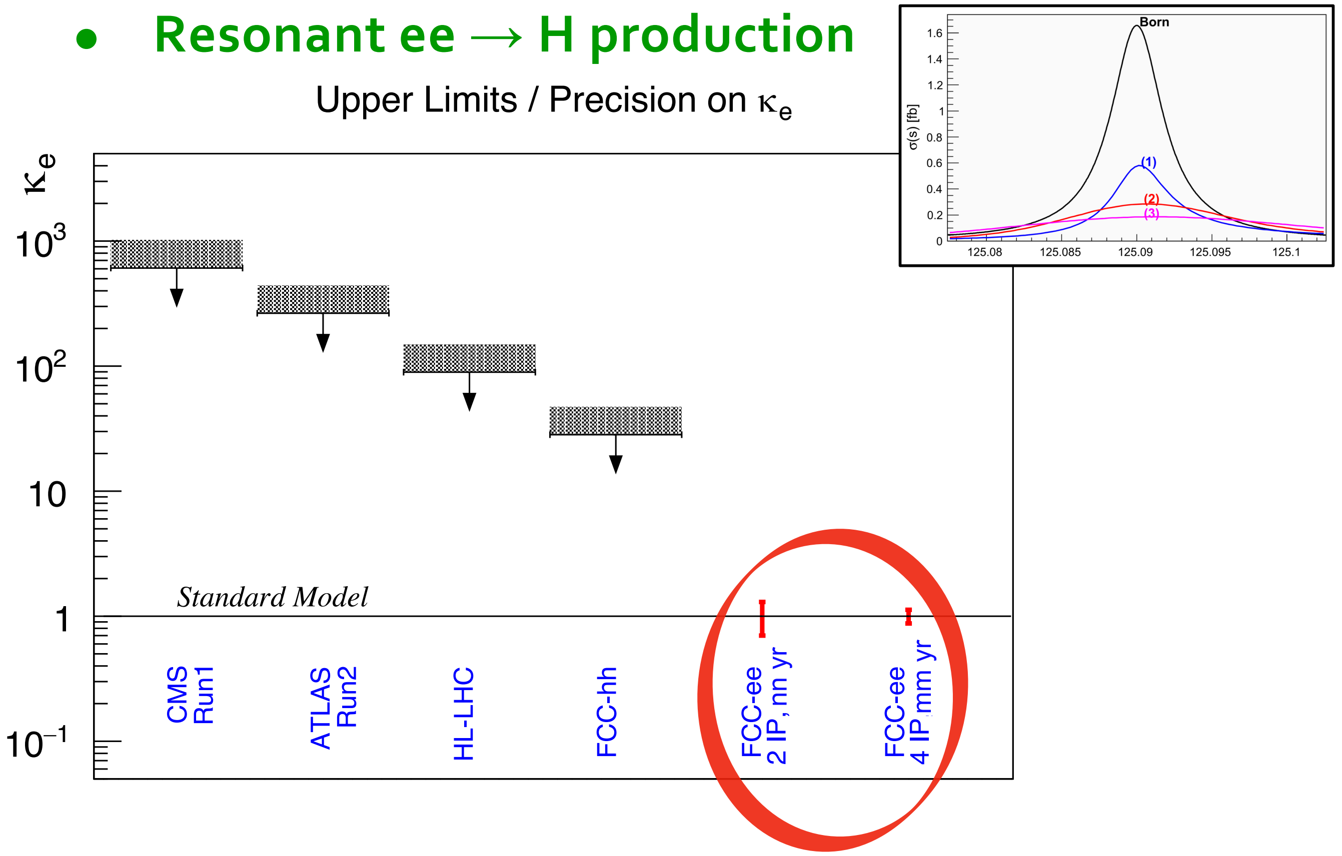

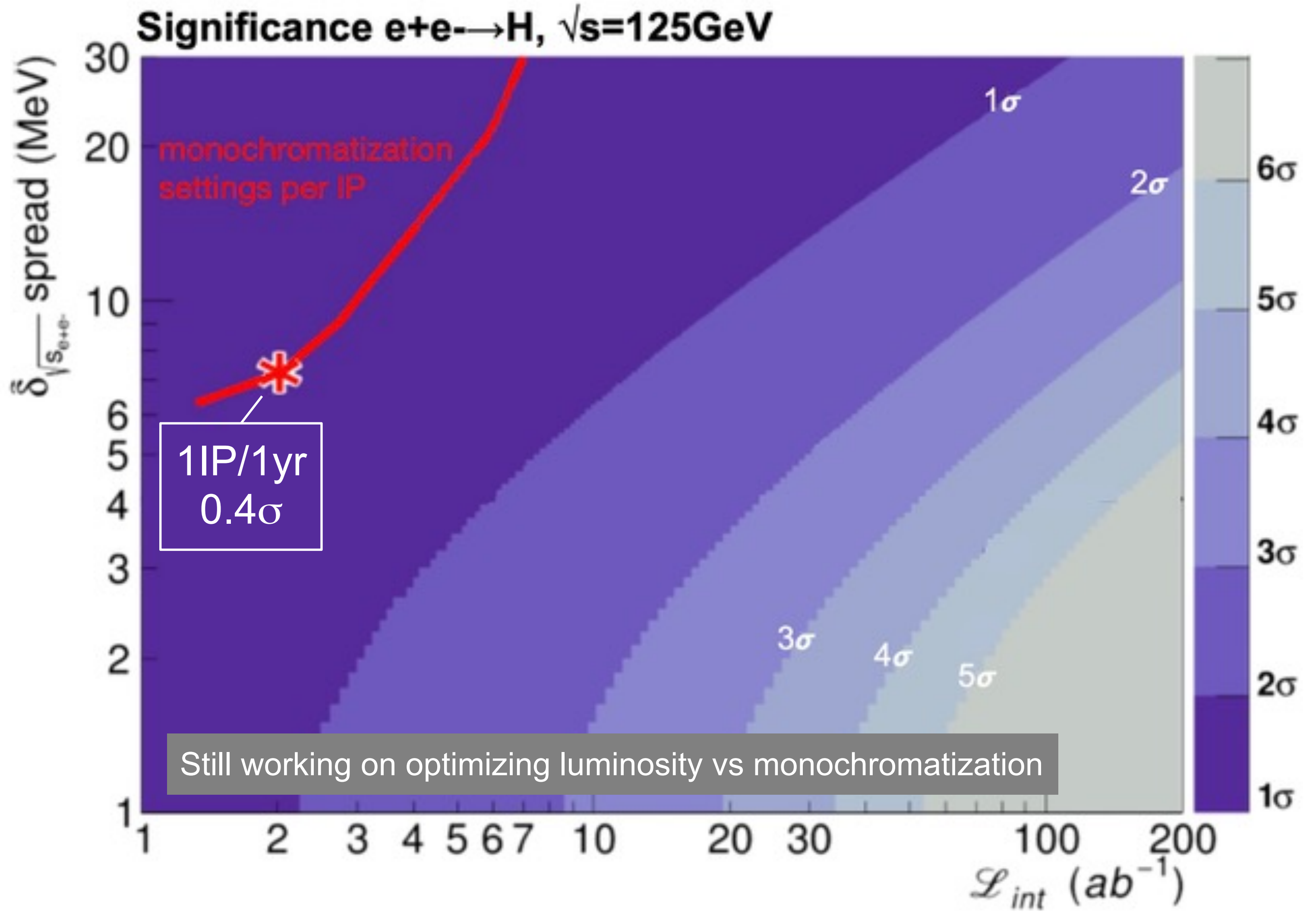

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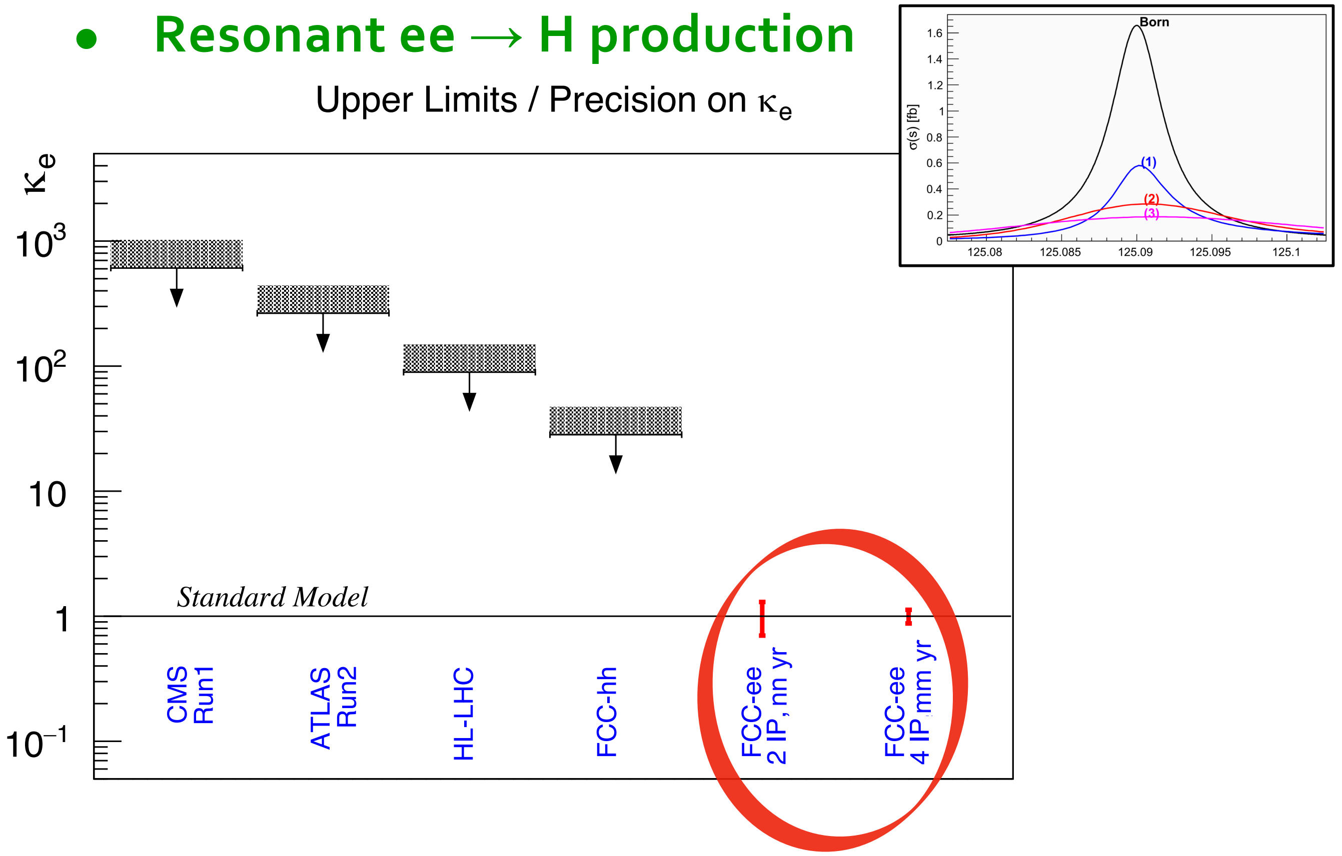

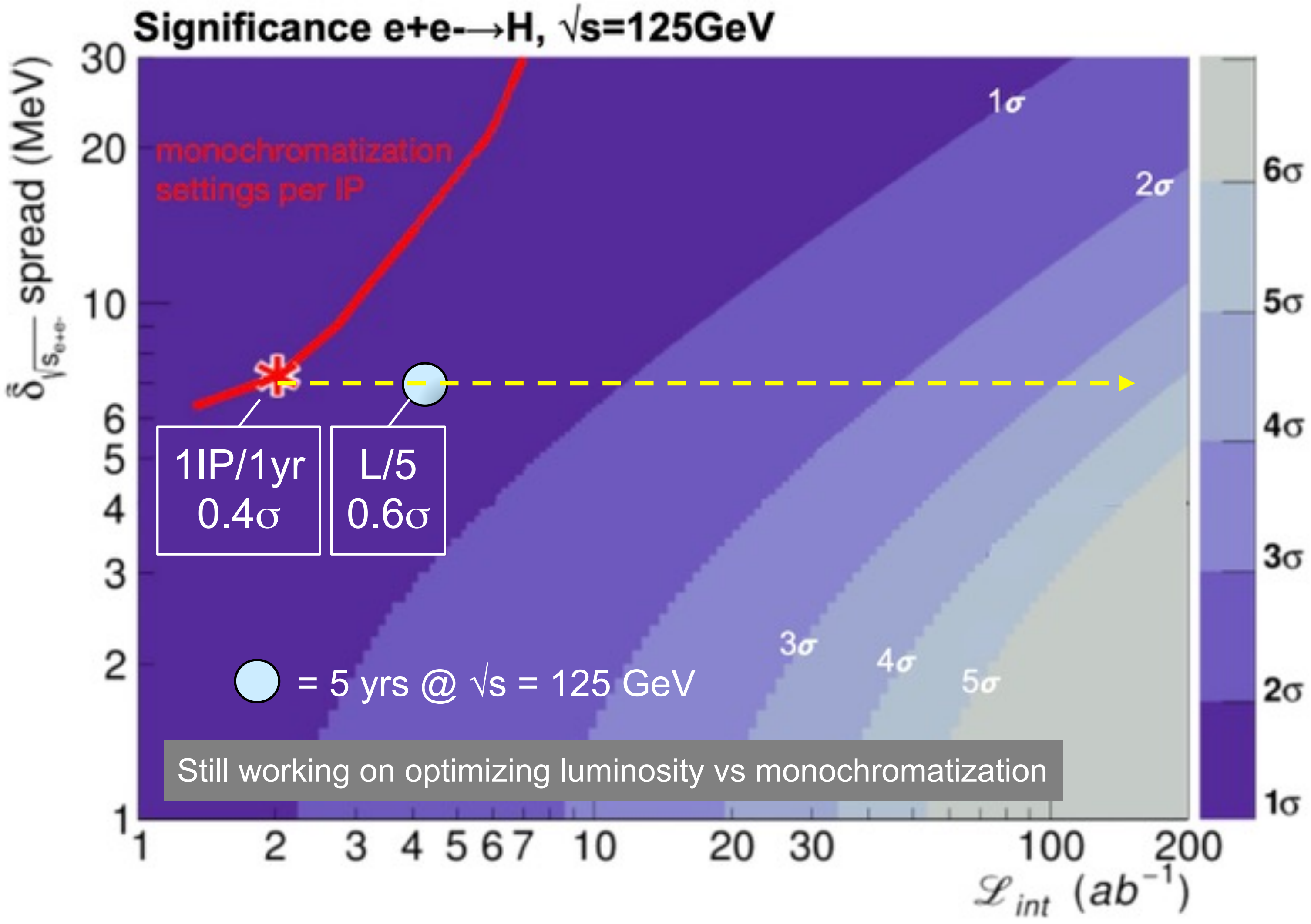

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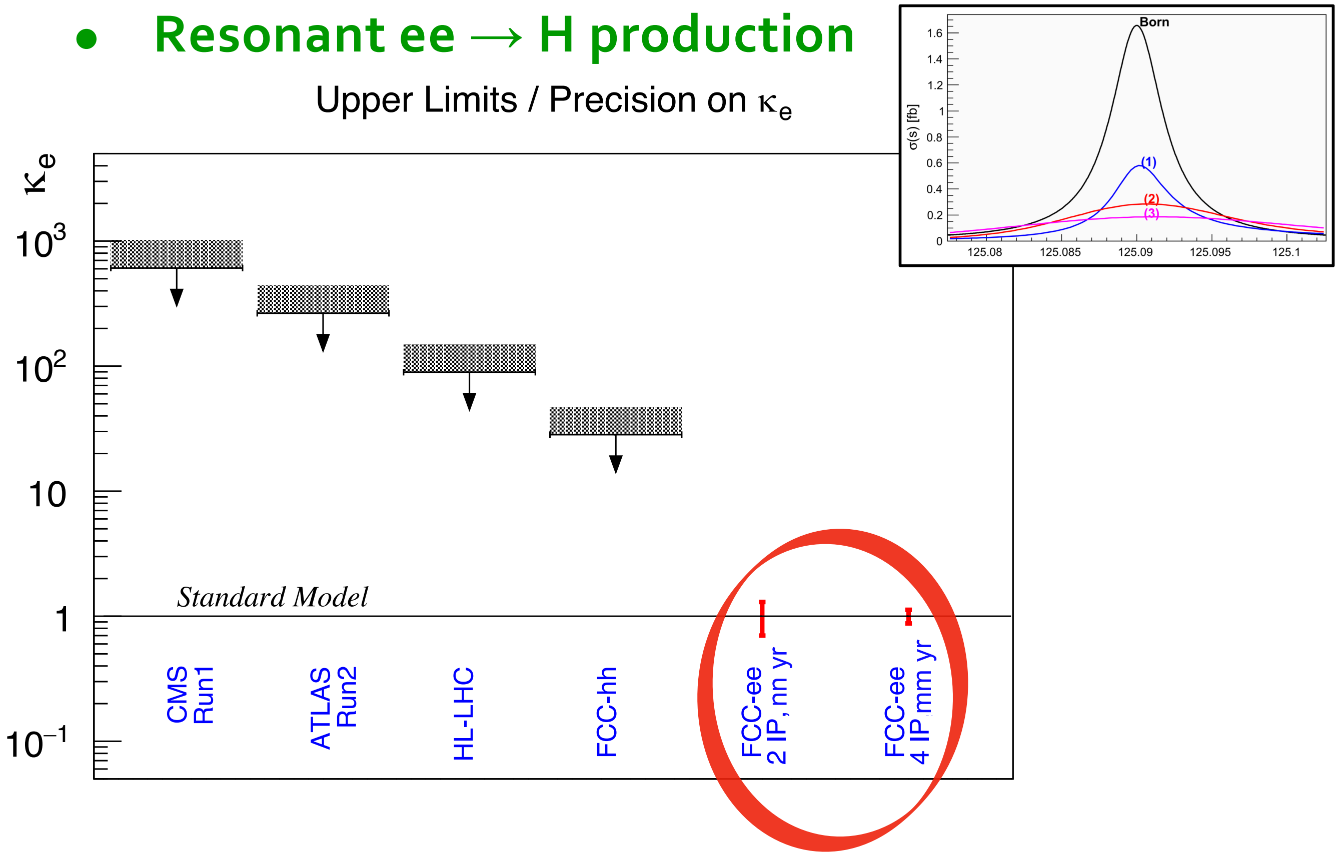

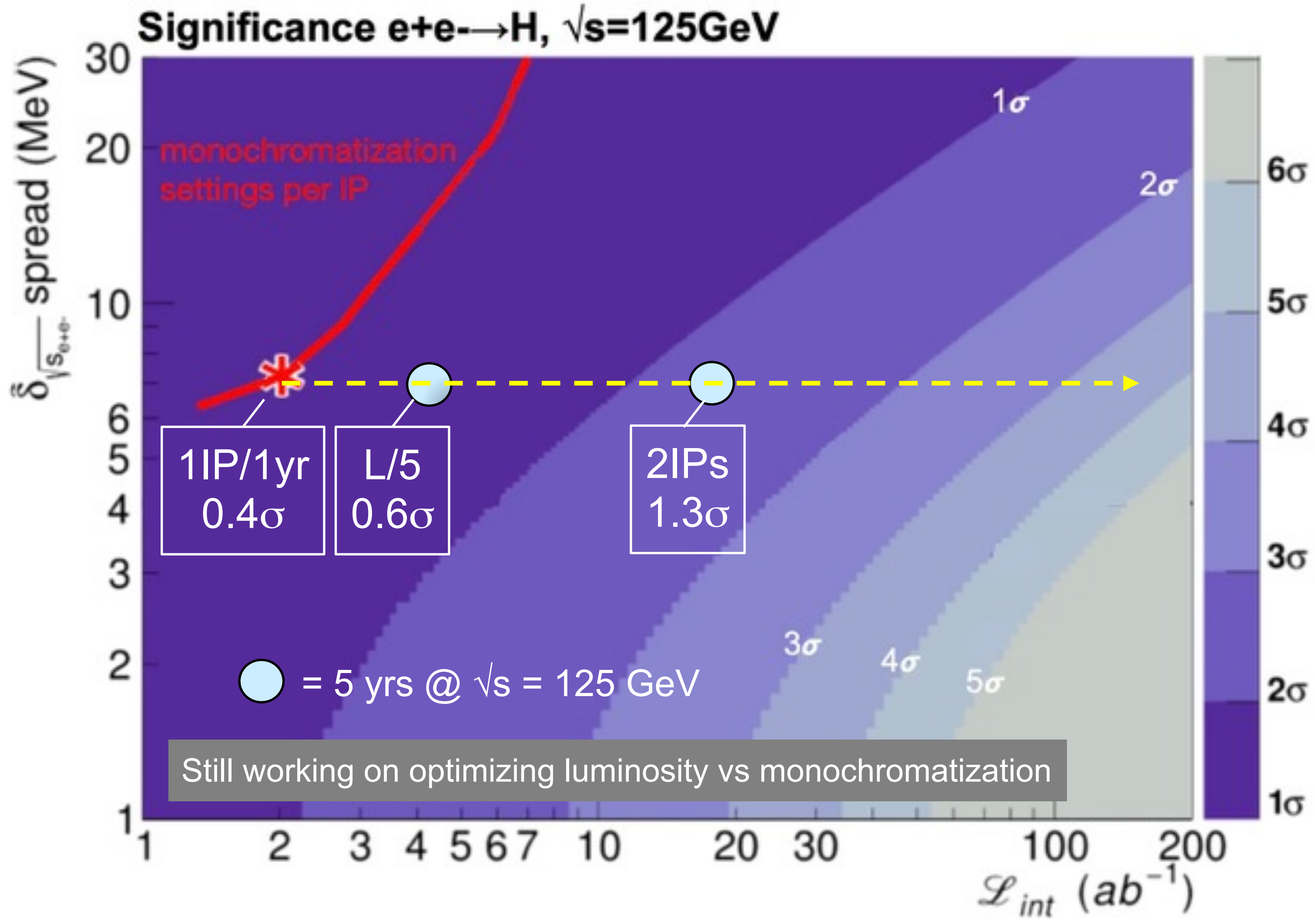

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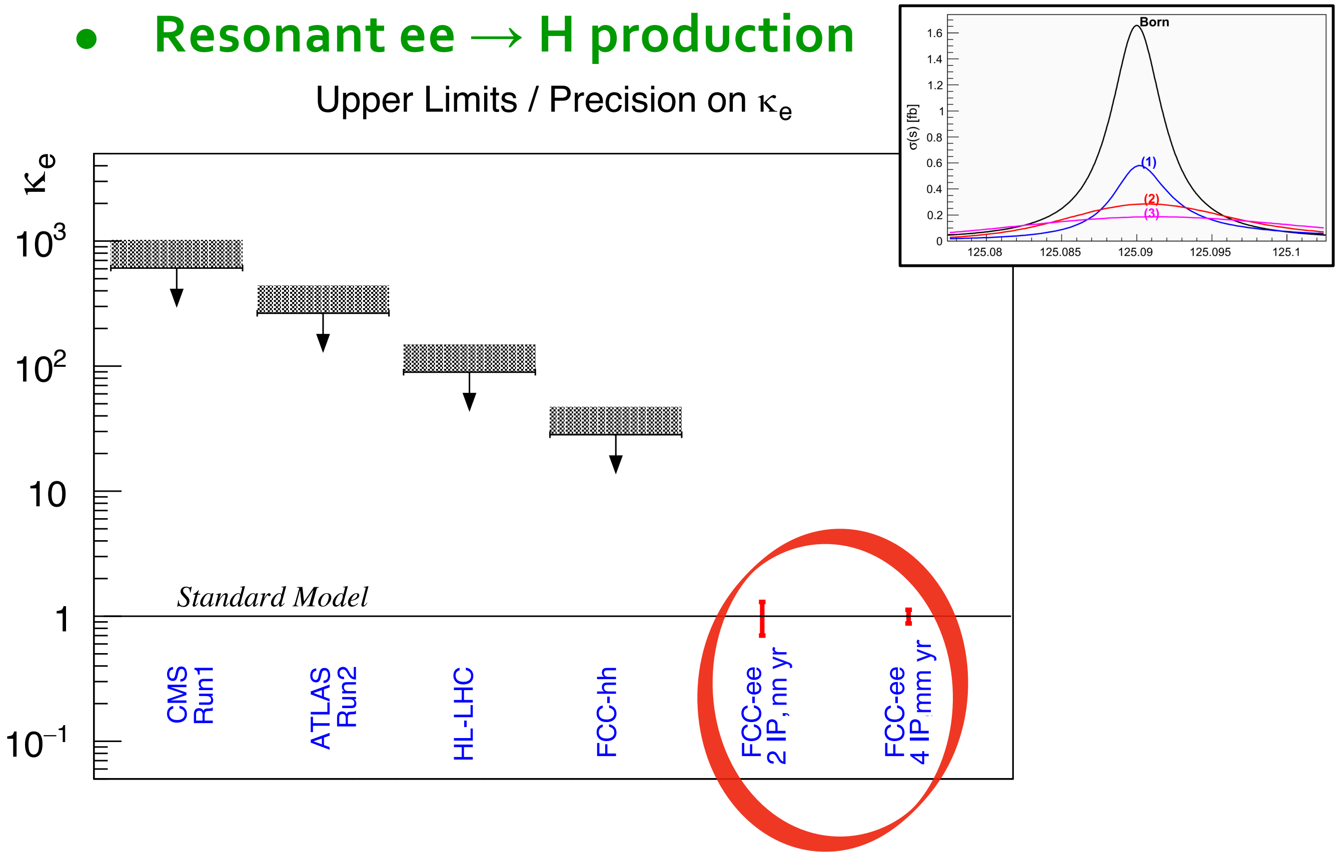

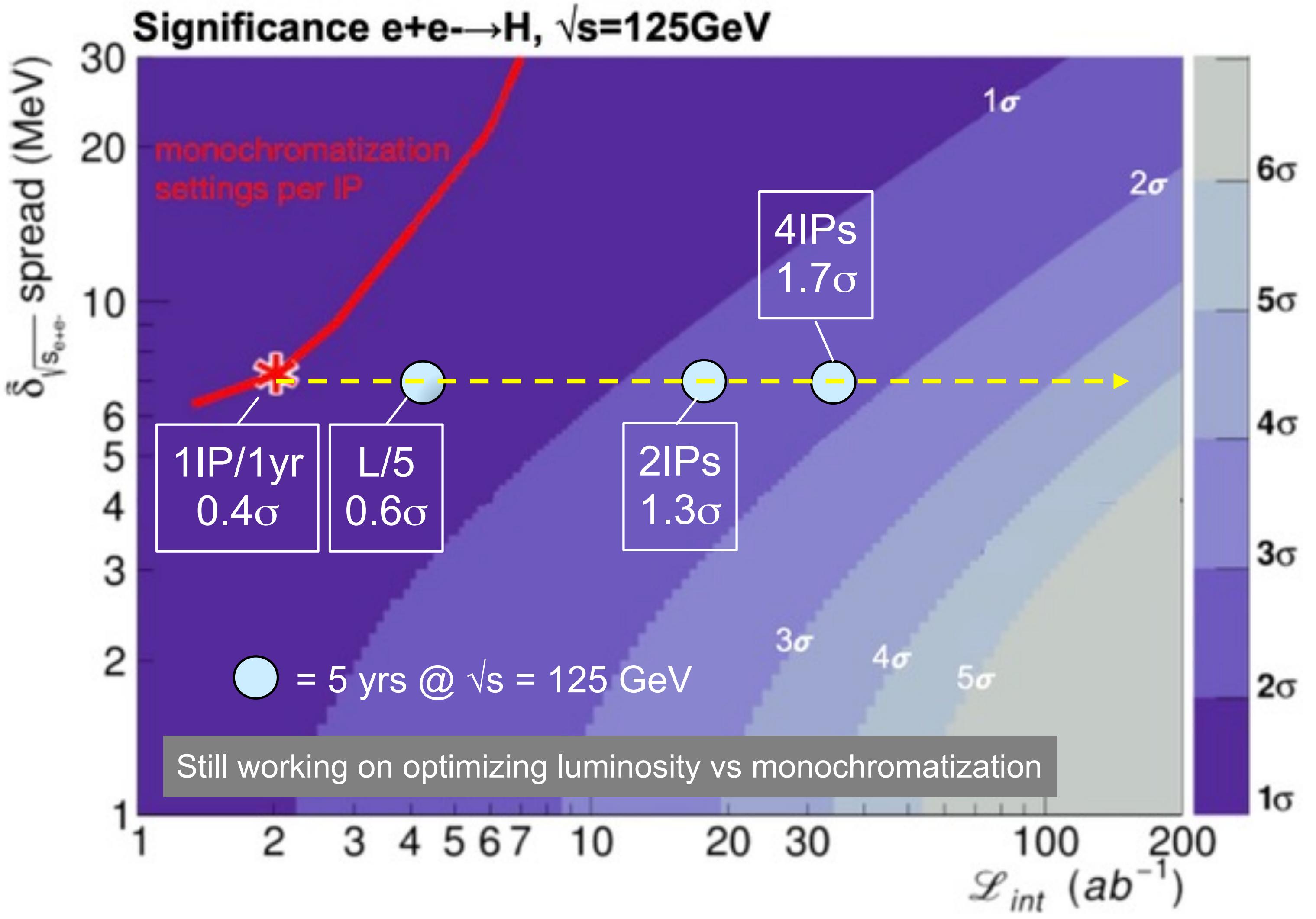

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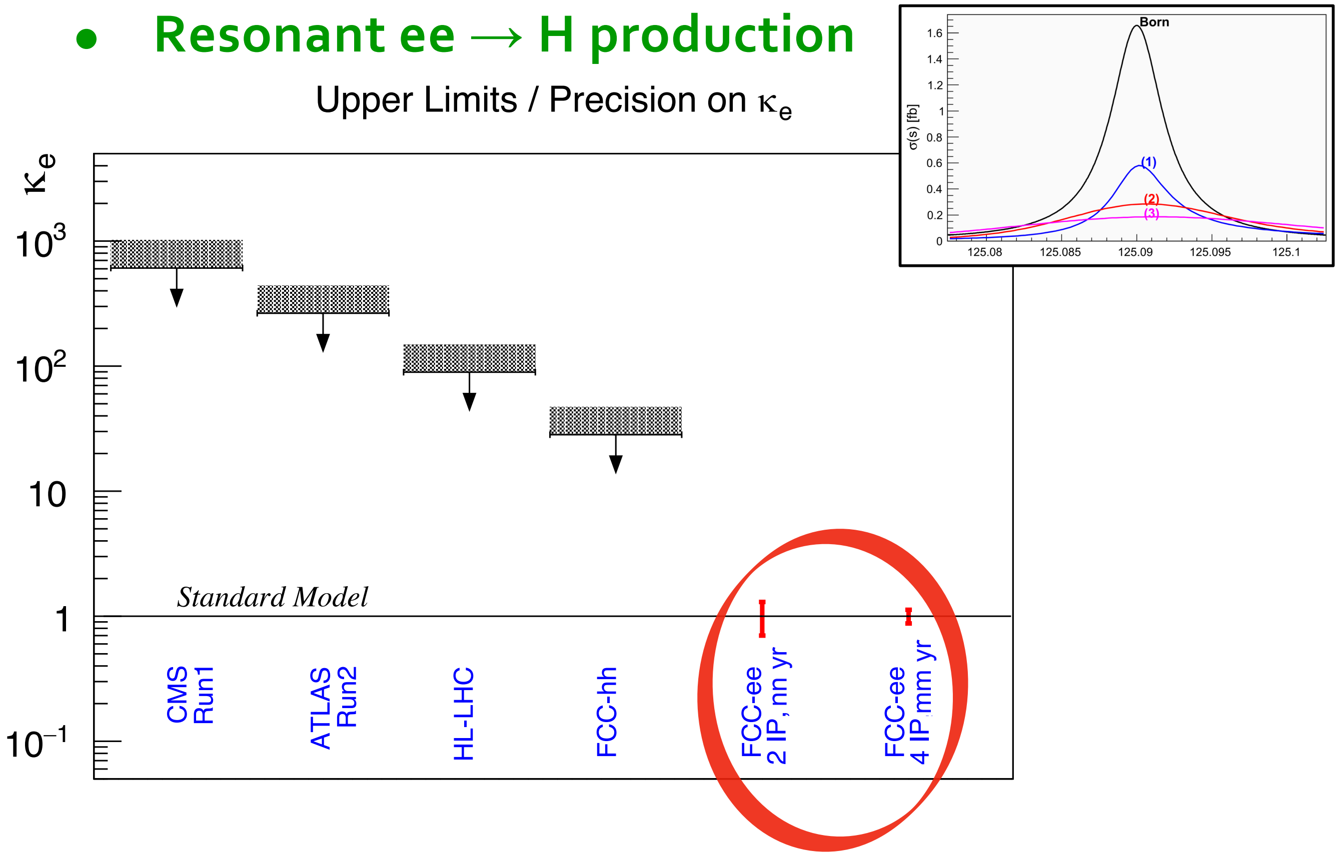
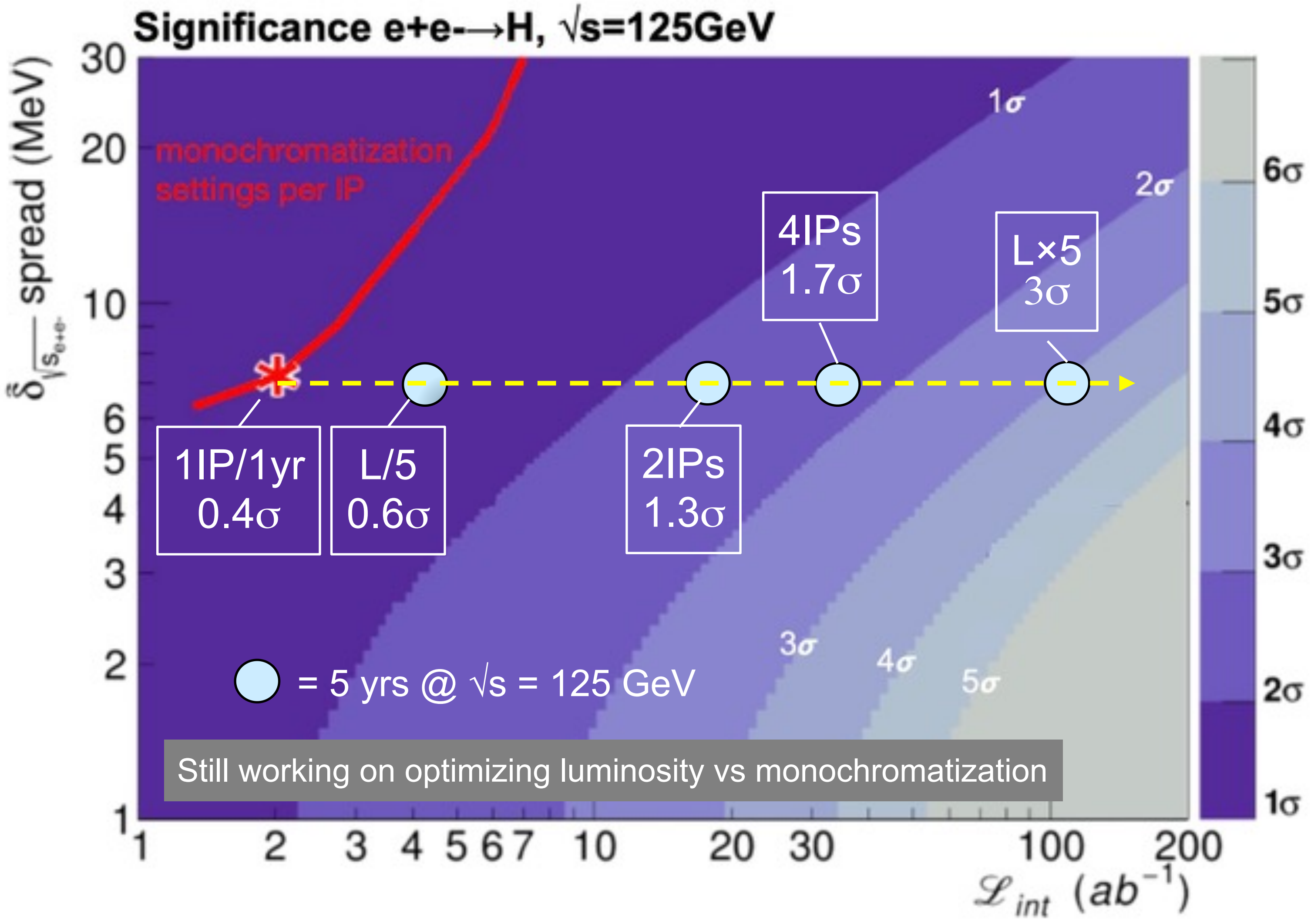
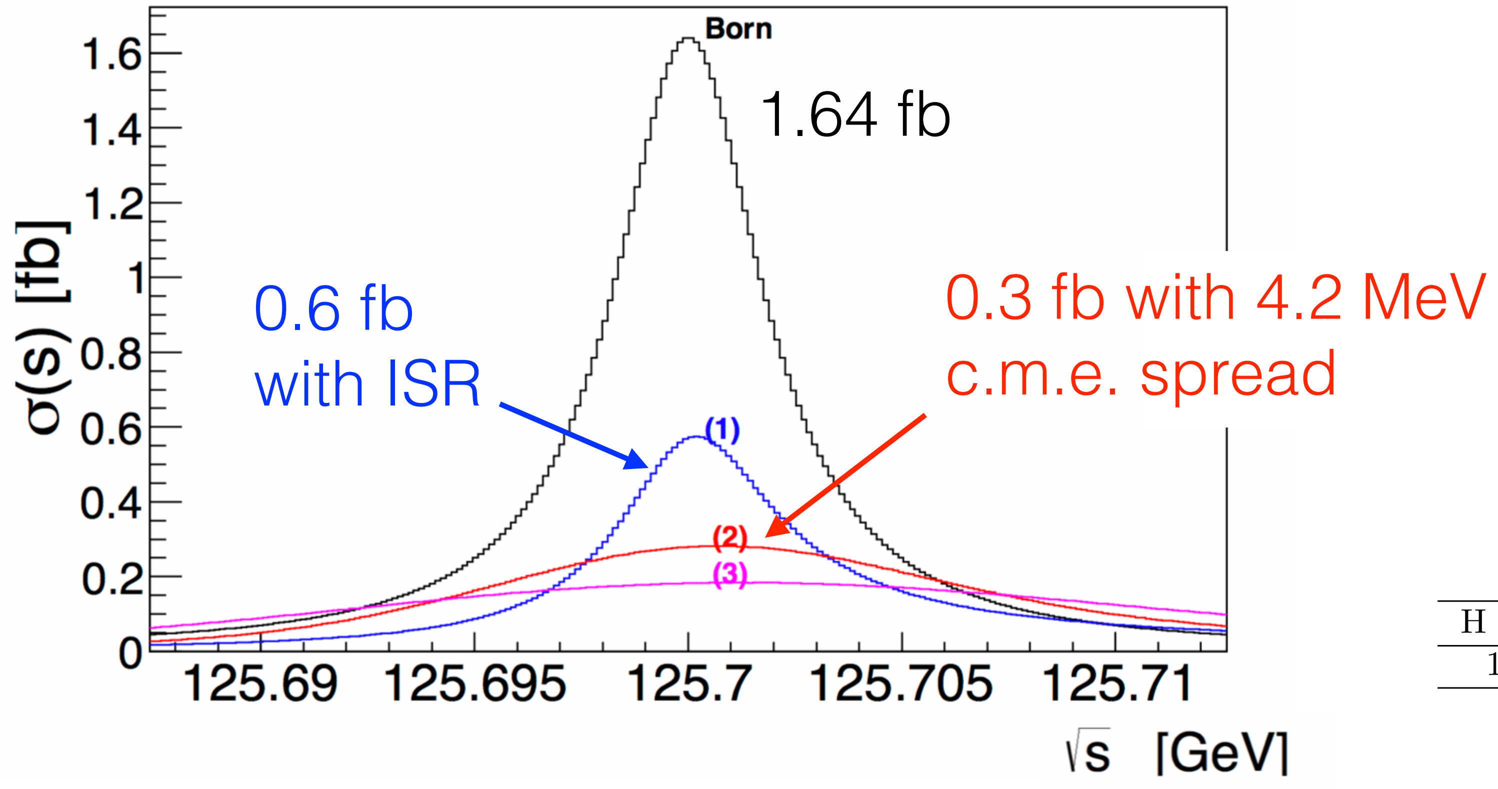




Electron Yukawa Coupling.

Jadach+, arXiv: 1509.02406



d'Enterria+, arXiv: 2107.02686

Higgs decay channel	$\mathcal{B}$	$\sigma \times \mathcal{B}$	Irreducible background	σ	S/B
$e^+e^- \to H \to b\bar{b}$	58.2%	164 ab	$e^+e^- \to b\bar{b}$	19 pb	$\mathcal{O}(10^{-5})$
$e^+e^- \to H \to gg$	8.2%	23 ab	$e^+e^- \to q\bar{q}$	61 pb	$\mathcal{O}(10^{-3})$
$e^+e^- \to H \to \tau\tau$	6.3%	18 ab	$e^+e^- \to \tau\tau$	10 pb	$\mathcal{O}(10^{-6})$
$e^+e^- \to H \to c\bar{c}$	2.9%	8.2 ab	$e^+e^- \to c\bar{c}$	22 pb	$\mathcal{O}(10^{-7})$
$e^+e^- \to H \to WW^* \to \ell\nu\, 2j$	21.4%×67.6%×32.4%×2	26.5 ab	$e^+e^- \to WW^* \to \ell\nu\, 2j$	23 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \to H \to WW^* \to 2\ell\, 2\nu$	21.4%×32.4%×32.4%	6.4 ab	$e^+e^- \to WW^* \to 2\ell\, 2\nu$	5.6 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \to H \to WW^* \to 4j$	21.4%×67.6%×67.6%	27.6 ab	$e^+e^- \to WW^* \to 4j$	24 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \to H \to ZZ^* \to 2j\, 2\nu$	2.6%×70%×20%×2	2 ab	$e^+e^- \to ZZ^* \to 2j\, 2\nu$	273 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \to H \to ZZ^* \to 2\ell\, 2j$	2.6%×70%×10%×2	1 ab	$e^+e^- \to ZZ^* \to 2\ell\, 2j$	136 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \to H \to ZZ^* \to 2\ell\, 2\nu$	2.6%×20%×10%×2	0.3 ab	$e^+e^- \to ZZ^* \to 2\ell\, 2\nu$	39 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \to H \to \gamma\gamma$	0.23%	0.65 ab	$e^+e^- \to \gamma\gamma$	79 pb	$\mathcal{O}(10^{-8})$

w. 10/ab

$H \to gg$	$H \to WW^* \to \ell\nu\, 2j;\ 2\ell\, 2\nu;\ 4j$	$H \to ZZ^* \to 2j\, 2\nu;\ 2\ell\, 2j;\ 2\ell\, 2\nu$	$H \to b\bar{b}$	$H \to \tau_{had}\tau_{had};\ c\bar{c};\ \gamma\gamma$	Combined
1.1σ	$(0.53 \otimes 0.34 \otimes 0.13)\sigma$	$(0.32 \otimes 0.18 \otimes 0.05)\sigma$	0.13σ	$< 0.02\sigma$	1.3σ

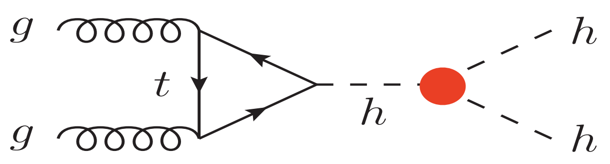
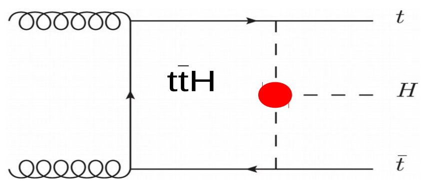
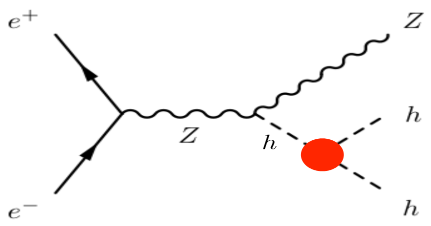
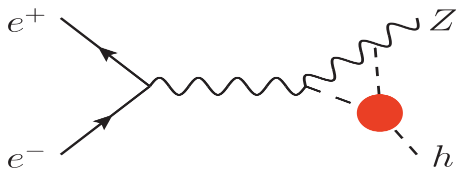
w/ 10/ab: S~55, B~2400 → 1.1σ

Electron Yukawa Coupling.

Higgs self-coupling.

Higgs self-couplings is very interesting for a multitude of reasons
(vacuum stability, hierarchy, baryogenesis, GW, EFT probe…).

How much can it deviate from SM given the tight constraints on other Higgs couplings?
Do you need to reach HH production threshold to constrain h^3 coupling?

	Directly: Higgs-pair prod	**Indirectly: via single Higgs**
Hadron Colliders		
Lepton Colliders		

ECFA Higgs study group '19

	di-Higgs	single-H
exclusive	**1. di-H, excl.** • Use of σ(HH) • only deformation of κλ	**3. single-H, excl.** • single Higgs processes at higher order • only deformation of κλ
global	**2. di-H, glob.** • Use of σ(HH) • deformation of κλ + of the single-H couplings (a) do not consider the effects at higher order of κλ to single H production and decays (b) these higher order effects are included	**4. single-H, glob.** • single Higgs processes at higher order • deformation of κλ + of the single Higgs couplings

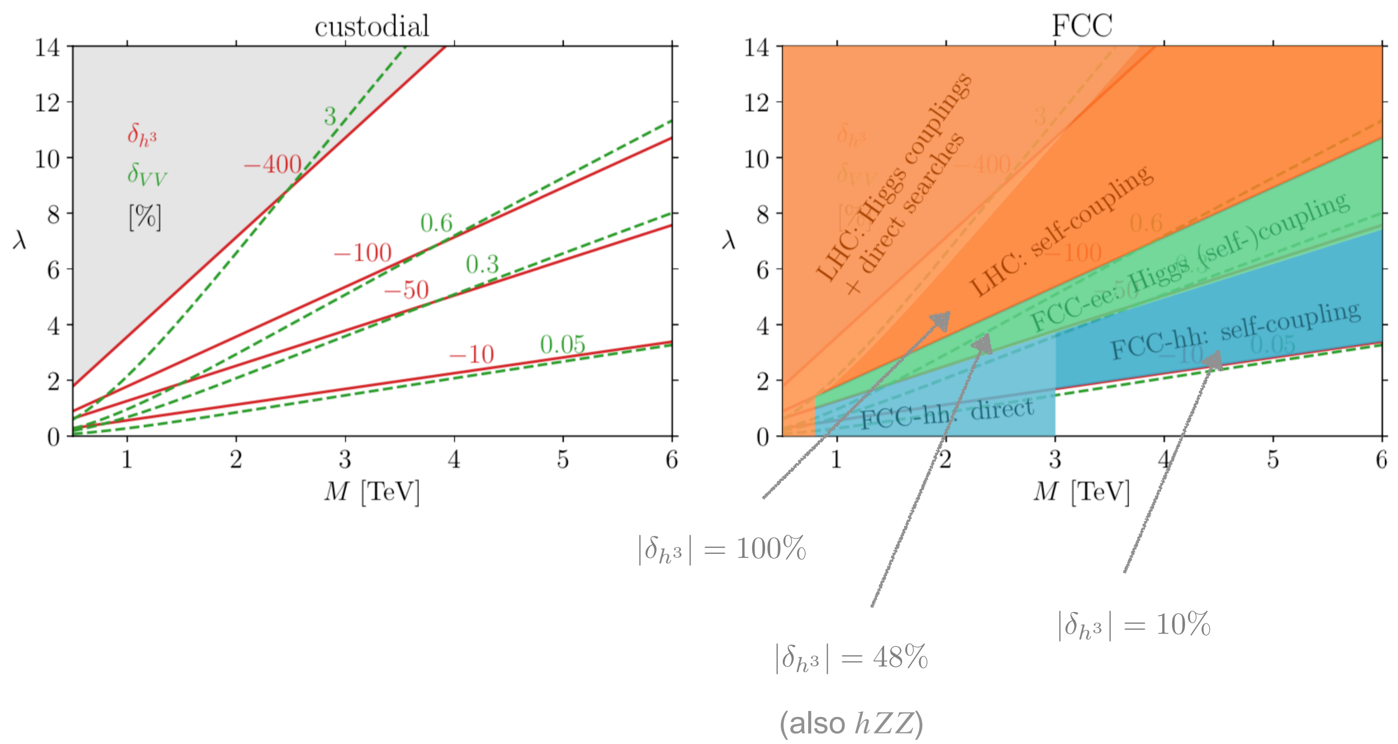
Large self-coupling scenarios.

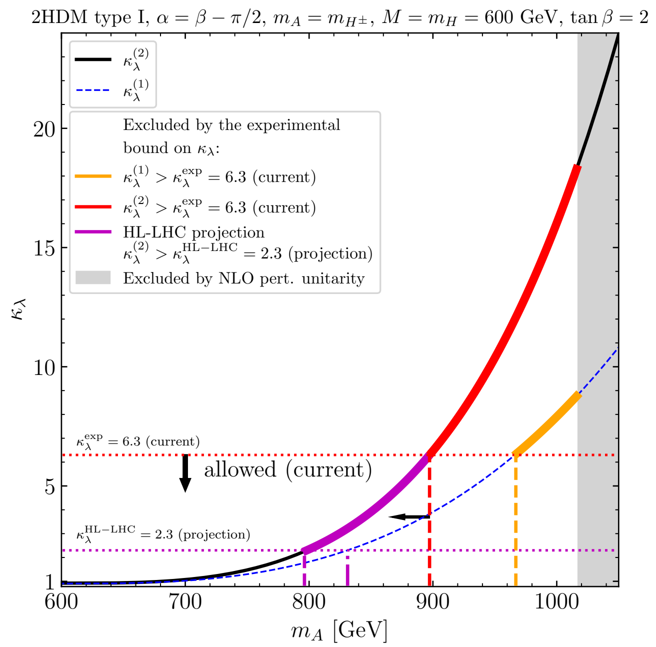
Generically: $\left|\frac{\delta_{h^3}}{\delta_{\text{single } h}}\right| \sim O(1)$ (composite Higgs/susy)

Particular exceptions: Higgs DM-portal models or custodial EW quadruplet

DiVita et al,: 1704.01953 Falkowski, Rattazzi: 1902.05936 Durieux, McCullough, Salvioni: 2209.00666

h^3 generically is not a tool to discover BSM but exceptions exist.



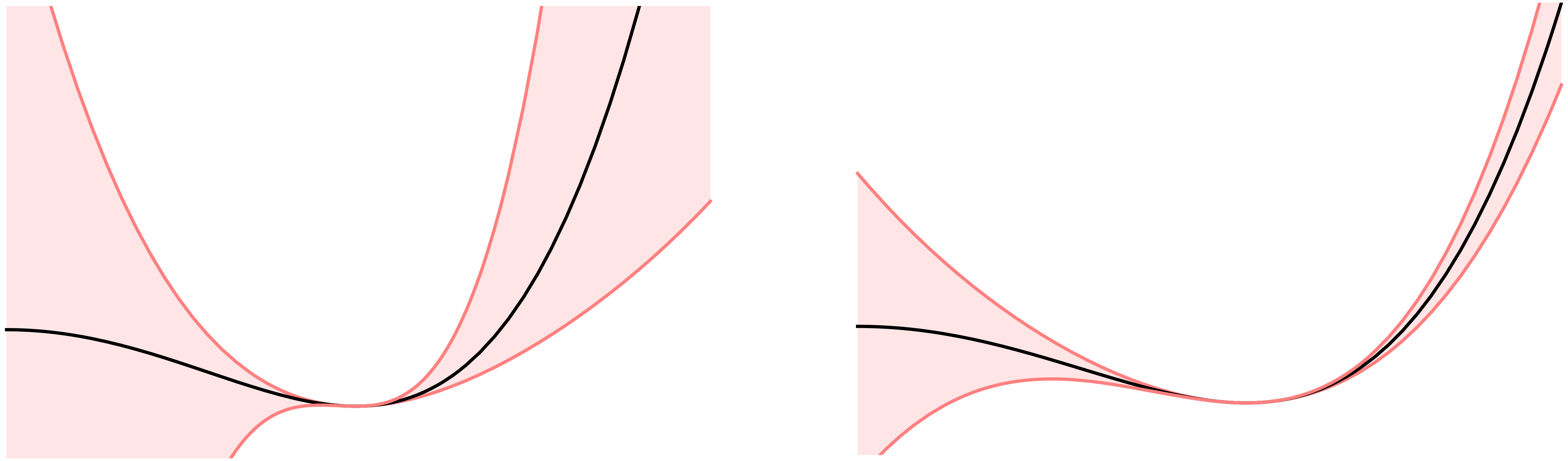


Other exceptions: non-decoupled/fine-tuned spectra

Bahl, Braathen, Weiglein: 2202.03453

Large self-coupling scenarios.

It is true that we haven't "measured" the Higgs potential but there are only peculiar physics scenarios that produce large deviations in the shape of the potential without leaving imprints elsewhere.



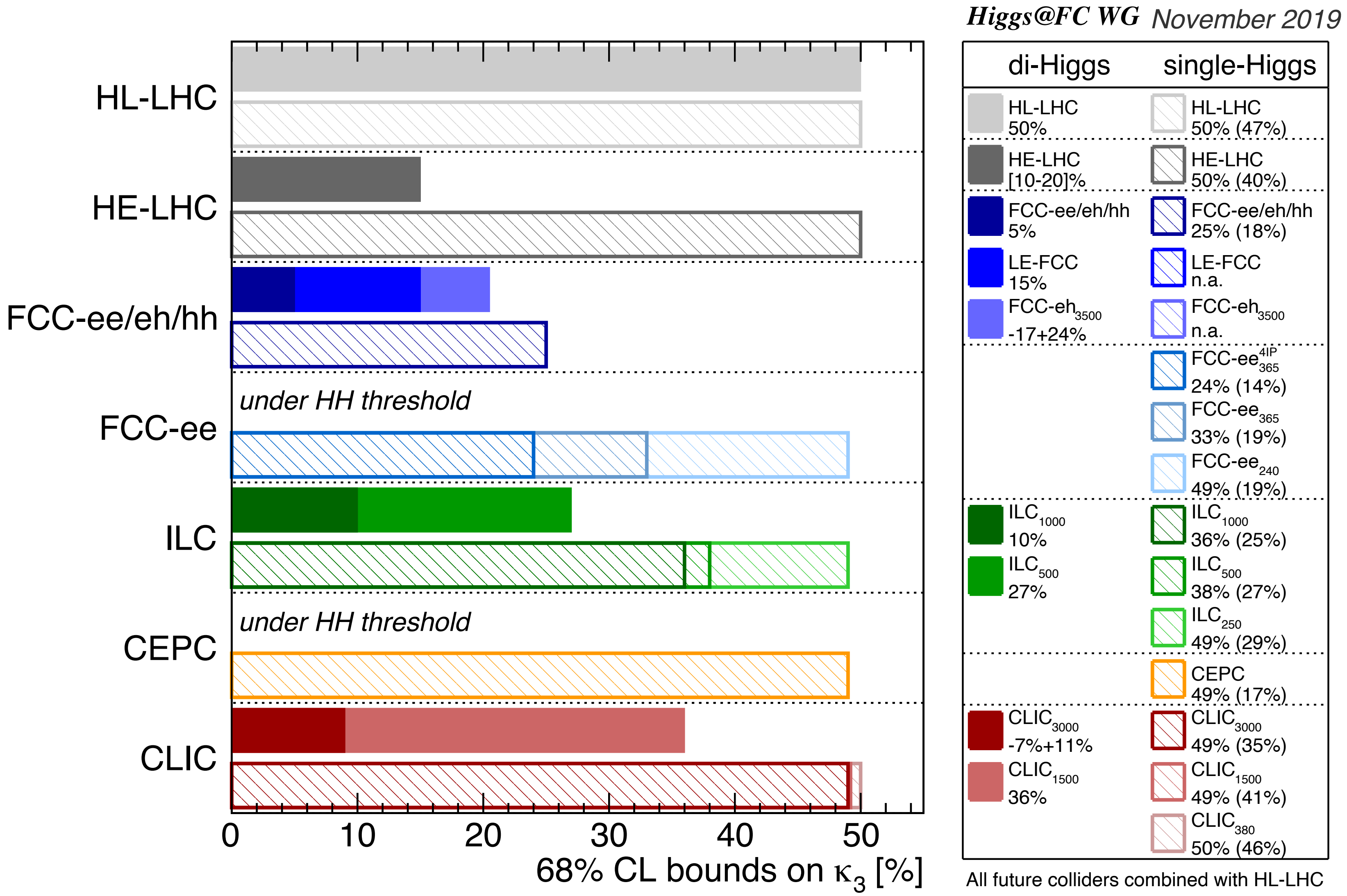
Current LHC → **HL-LHC**

R. Petrossian-Byrne/N. Craig @ LCWS'23

Important to understand which dynamics is really probed when embarking into challenging measurements. Actually, double Higgs production is also interesting to probe new physics in its tail rather than near threshold (where the sensitivity to Higgs self-coupling comes from).

Higgs self-coupling.

ECFA Higgs study group '19



1

Don't need to reach HH threshold
to have access to h^3.
Z-pole run is very important
if the HH threshold cannot be reached

2

The determination of h^3 at FCC-hh
relies on HH channel,
for which FCC-ee is of little direct help.
But the extraction of h^3
requires precise knowledge of y_t.
1% $y_t \leftrightarrow$ 5% h^3
Precision measurement of y_t needs ee

50% sensitivity: establish that $h^3 \neq 0$ at 95%CL
20% sensitivity: 5σ discovery of the SM h^3 coupling
5% sensitivity: getting sensitive to quantum corrections to Higgs potential

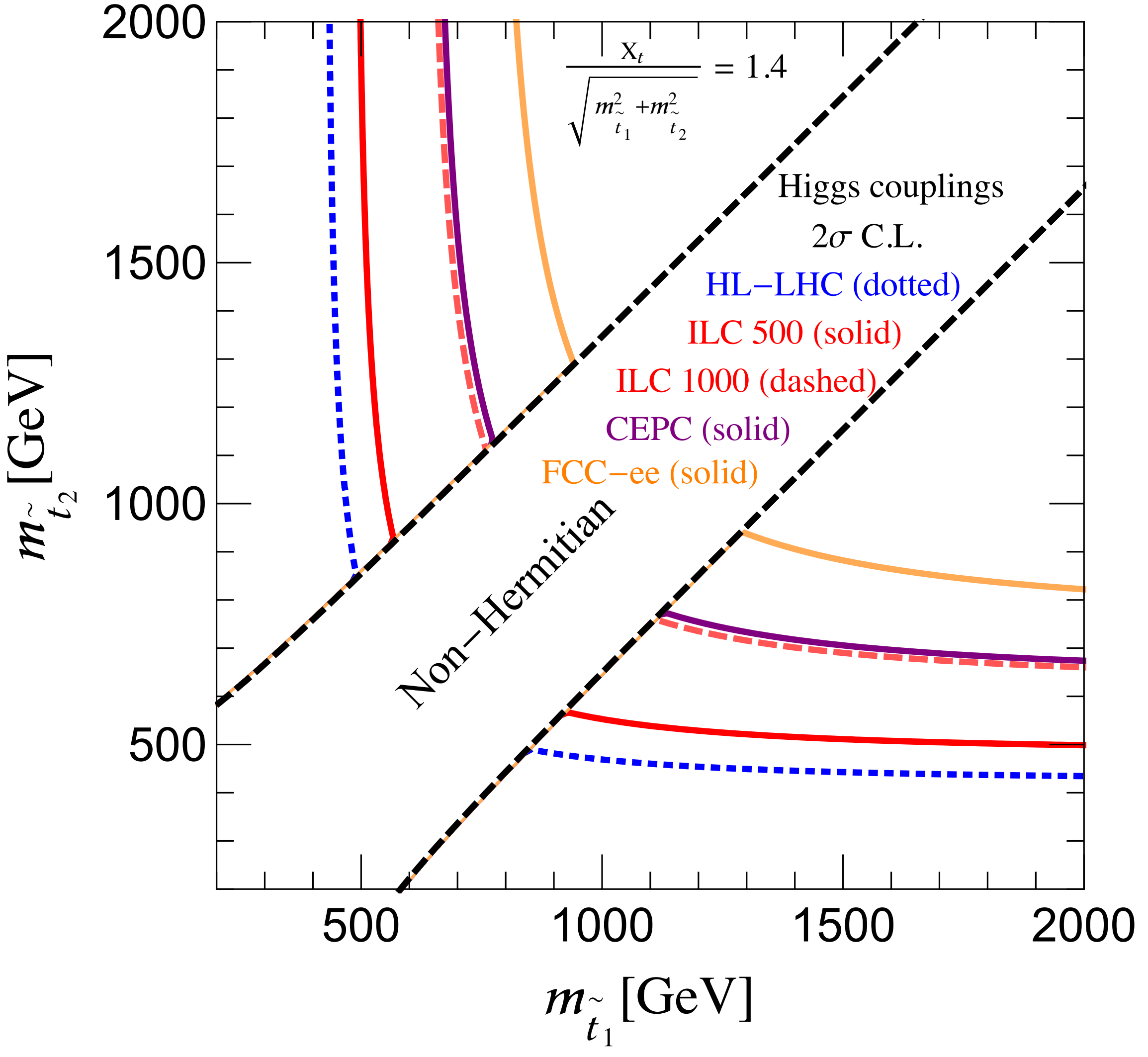
Discovery potential beyond LHC

Discovery Potential Beyond LHC.

Precisely measured EW and Higgs observables are sensitive to heavy New Physics

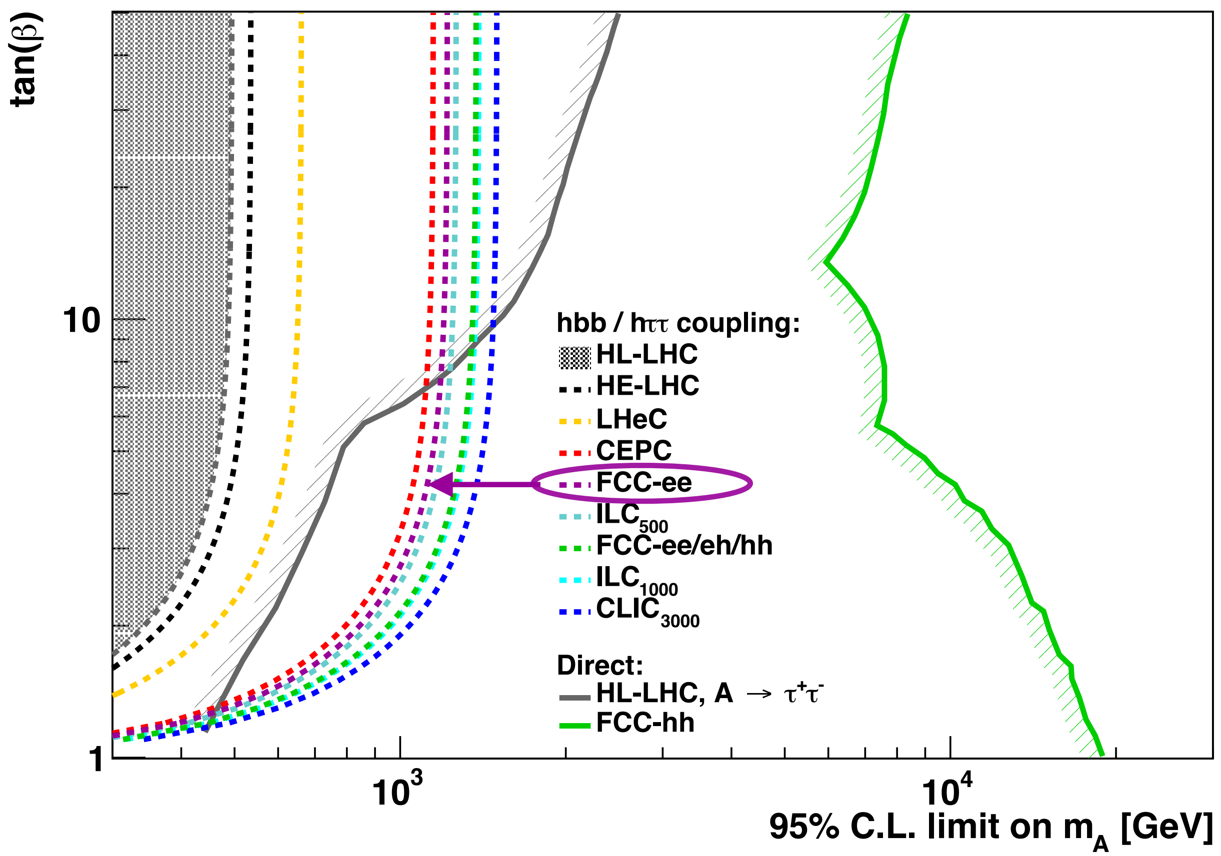
Examples of improved sensitivity wrt direct reach @ HL-LHC: SUSY

stops



Fan, Reece, Wang '14

Heavy neutral Higgses

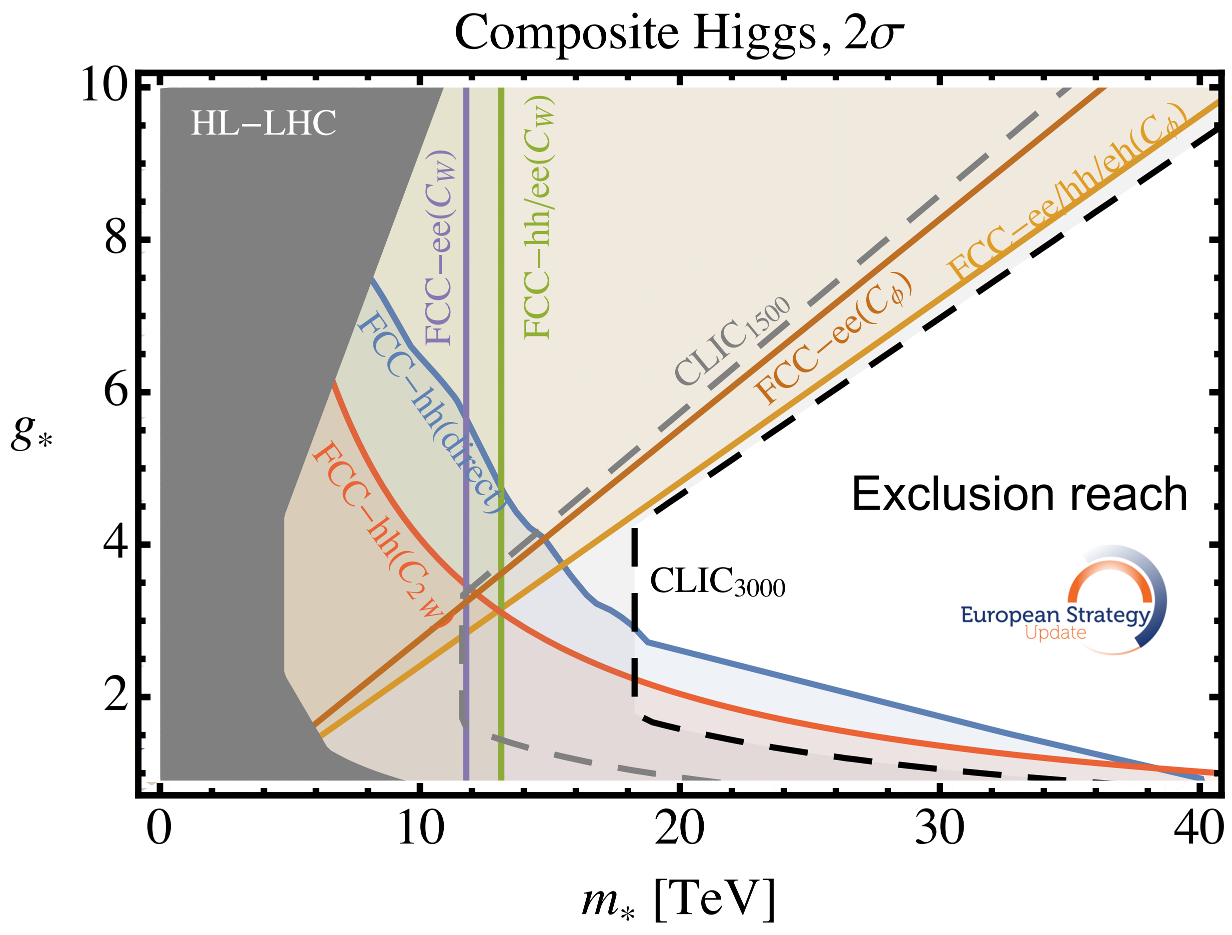


ESU Physics BB '19

Discovery Potential Beyond LHC.

Precisely measured EW and Higgs observables are sensitive to heavy New Physics

Examples of improved sensitivity wrt direct reach @ HL-LHC: Composite Higgs



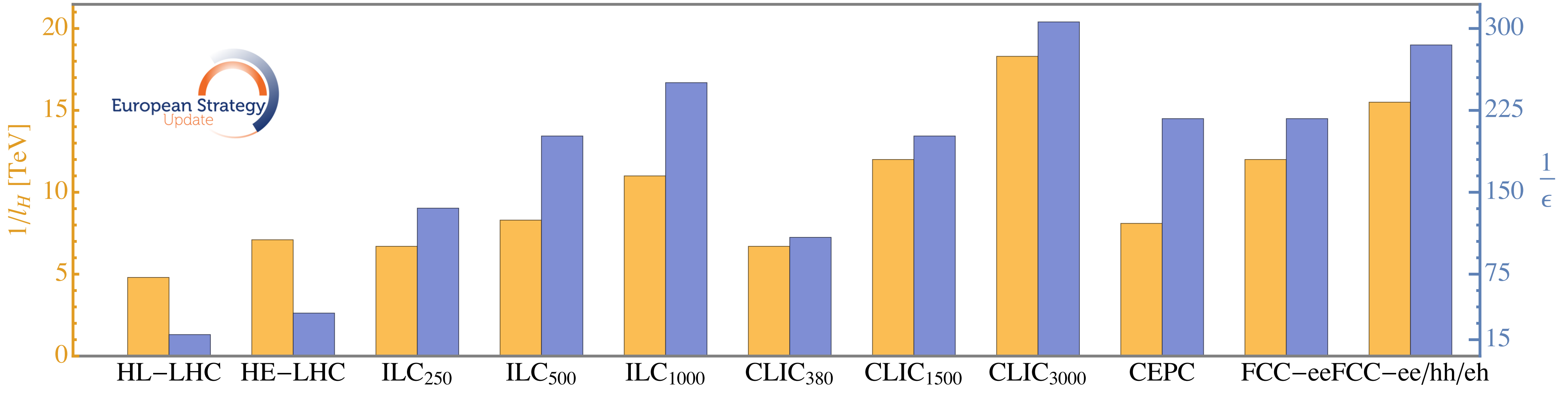


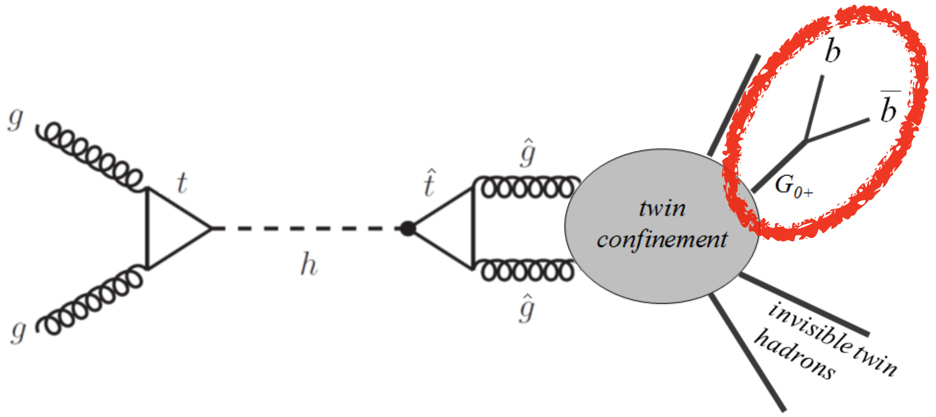
Fig. 8.5: Exclusion reach of different colliders on the inverse Higgs length $1/\ell_H = m_*$ (orange bars, left axis) and the tuning parameter $1/\varepsilon$ (blue bars, right axis), obtained by choosing the weakest bound valid for any value of the coupling constant g_*.

ESU Physics BB '19

Direct Searches for Elusive New Physics

- **LLP searches with displaced vertices**
  e.g. in twin Higgs models glueballs that mix with the Higgs and decay back to b-quarks

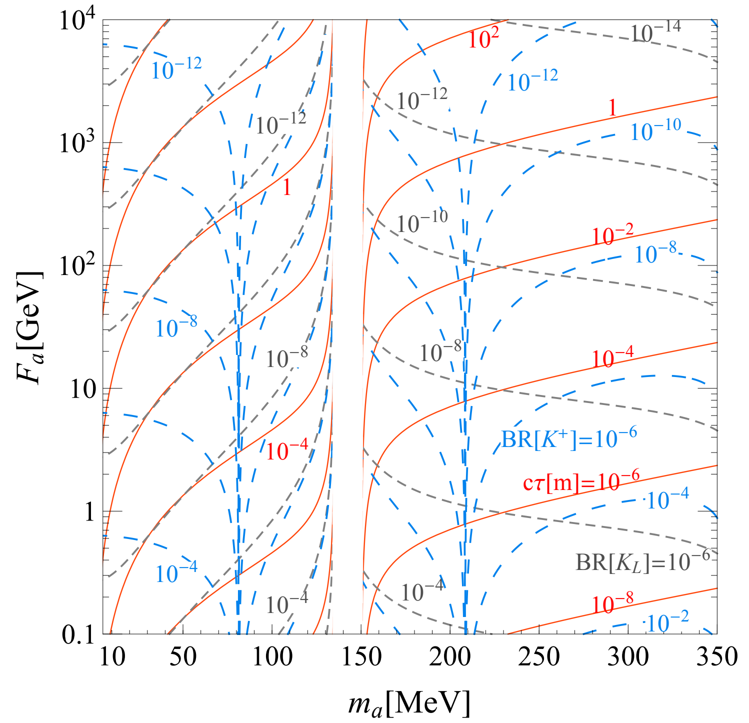
Craig et al, arXiv:1501.05310



- **Rare decays** Gori et al arXiv:2005.05170
  e.g. ALP mixing w/ SM mesons:

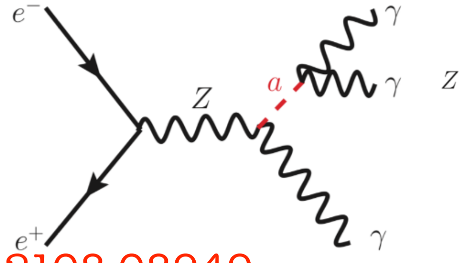
$$K_L \to \pi^0 a \to \pi^0 \gamma\gamma \text{ (KOTO)}$$
$$K^+ \to \pi^+ a \to \pi^+ \gamma\gamma \text{ (NA62)}$$

$$\mathcal{L} = \frac{\alpha_s}{8\pi F_a} a G_{\mu\nu} \tilde{G}^{\mu\nu}$$

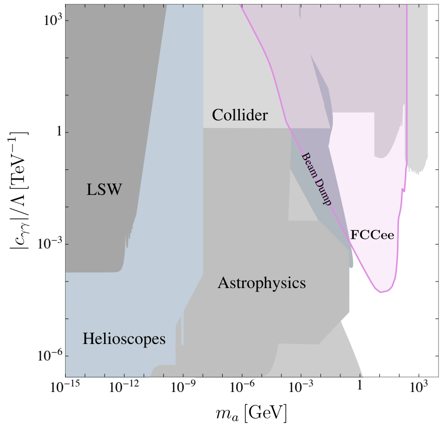


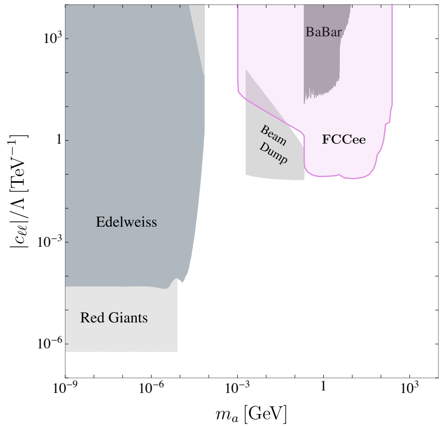
- **ALPs@ colliders**
  e.g. $e^+e^- \to \gamma a$

  $e^+e^- \to ha$



Knapen, Thamm arXiv:2108.08949

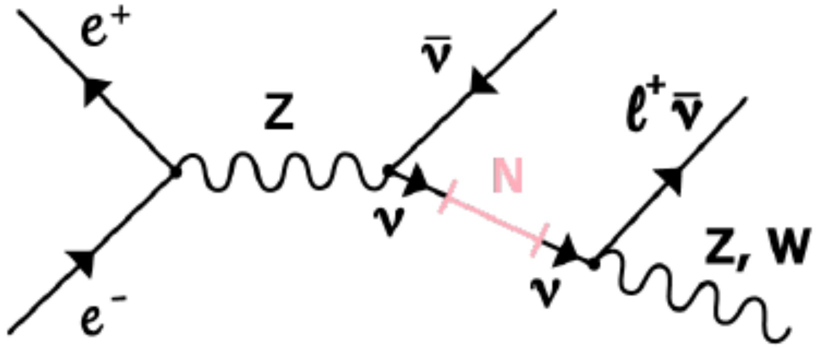




Astro/Cosmo → long-lived ALPs
colliders → short-lived ALPs MeV+

Search for ν_{RH}.

Direct observation
in Z decays
from LH-RH mixing



Important to understand
1. how neutrinos acquired mass
2. if lepton number is conserved
3. if leptogenesis is realised

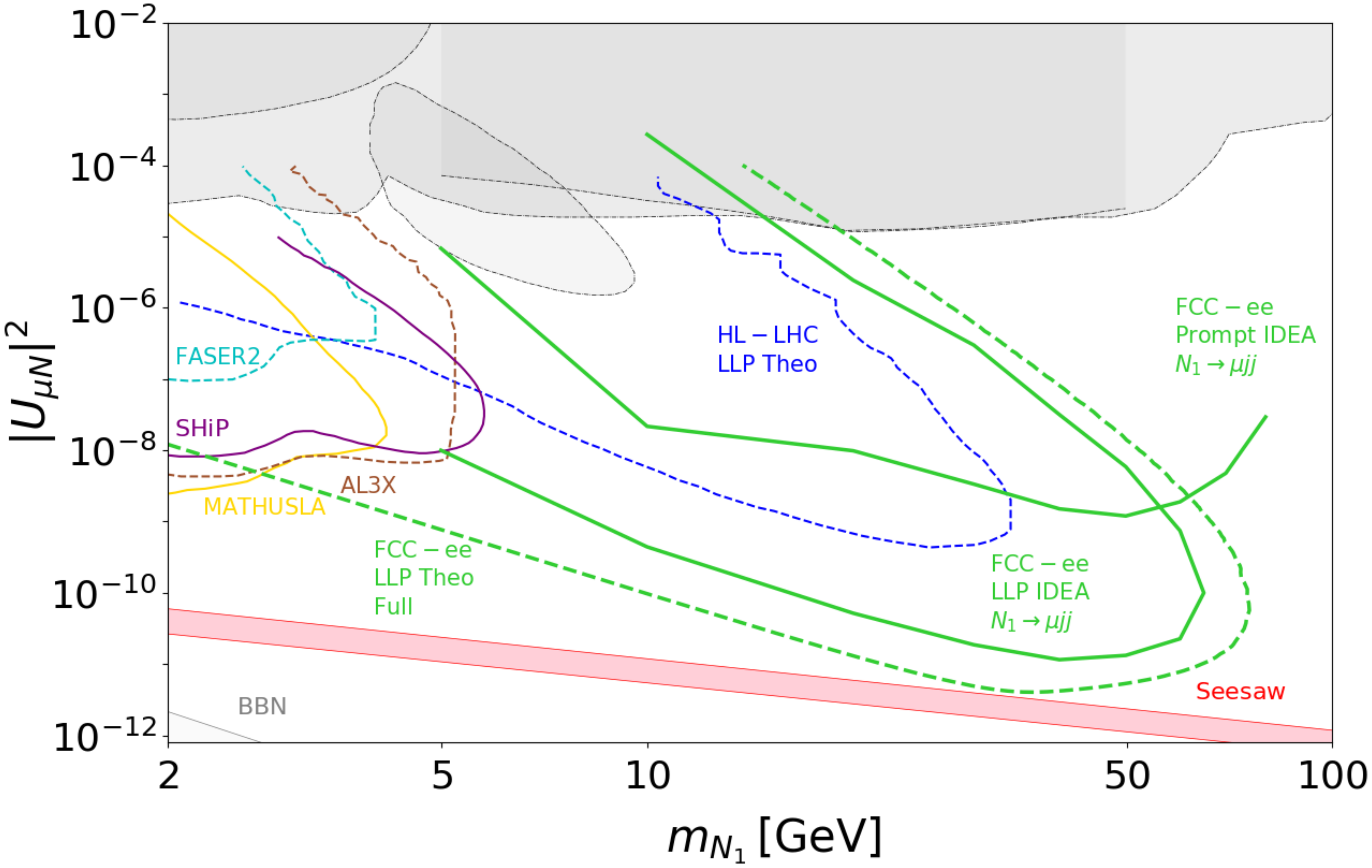


Fig. from mid-term report

Search for ν_{RH}.

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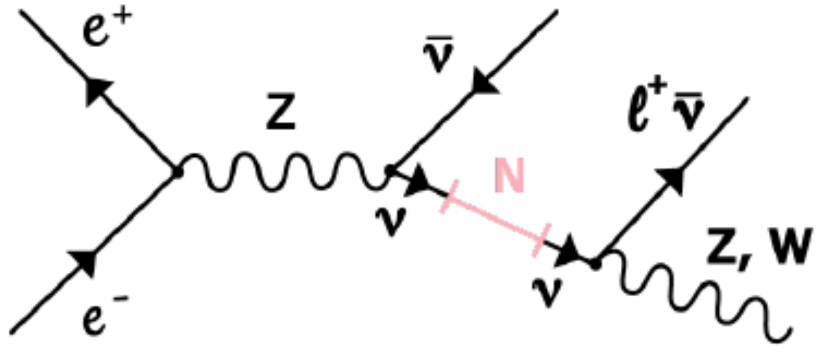

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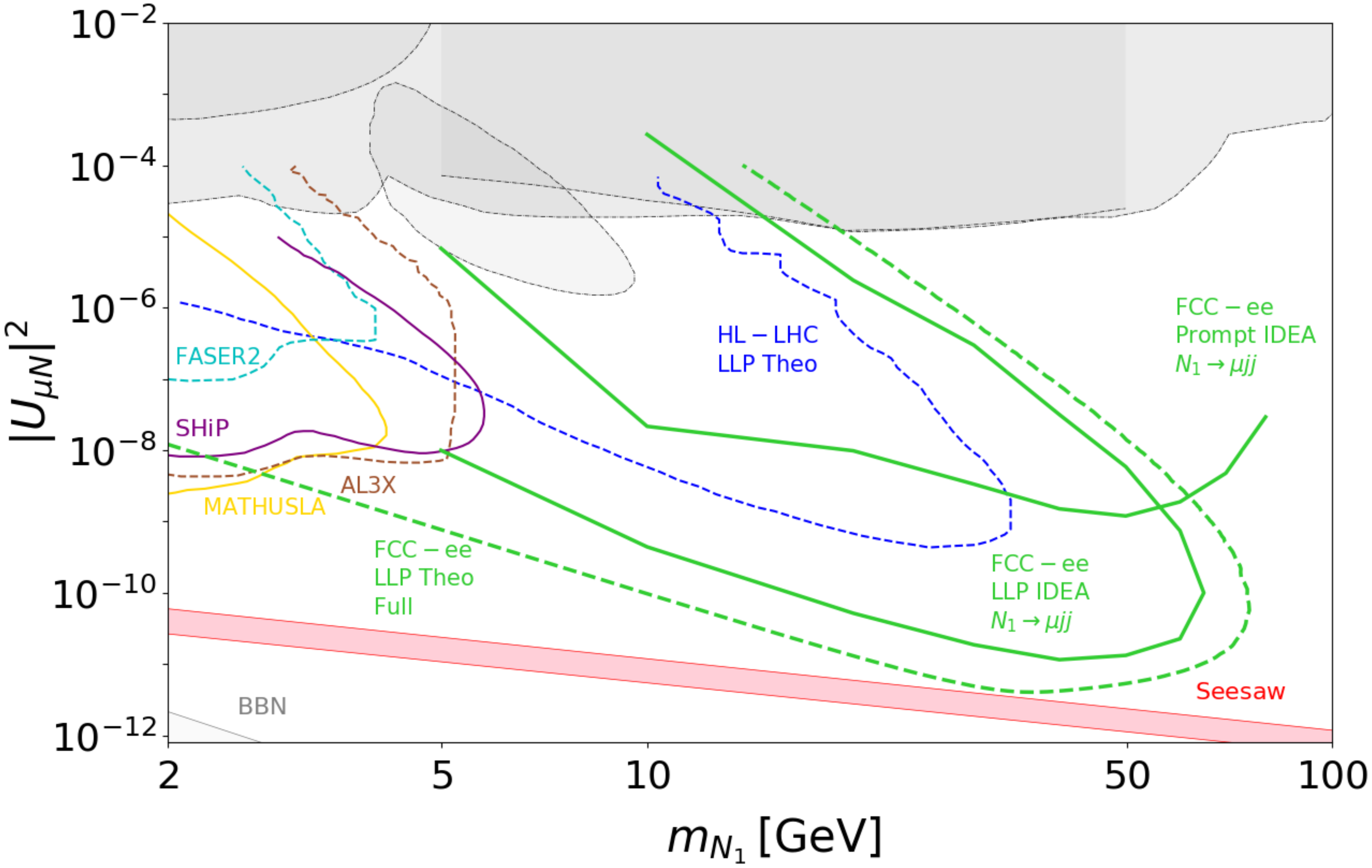
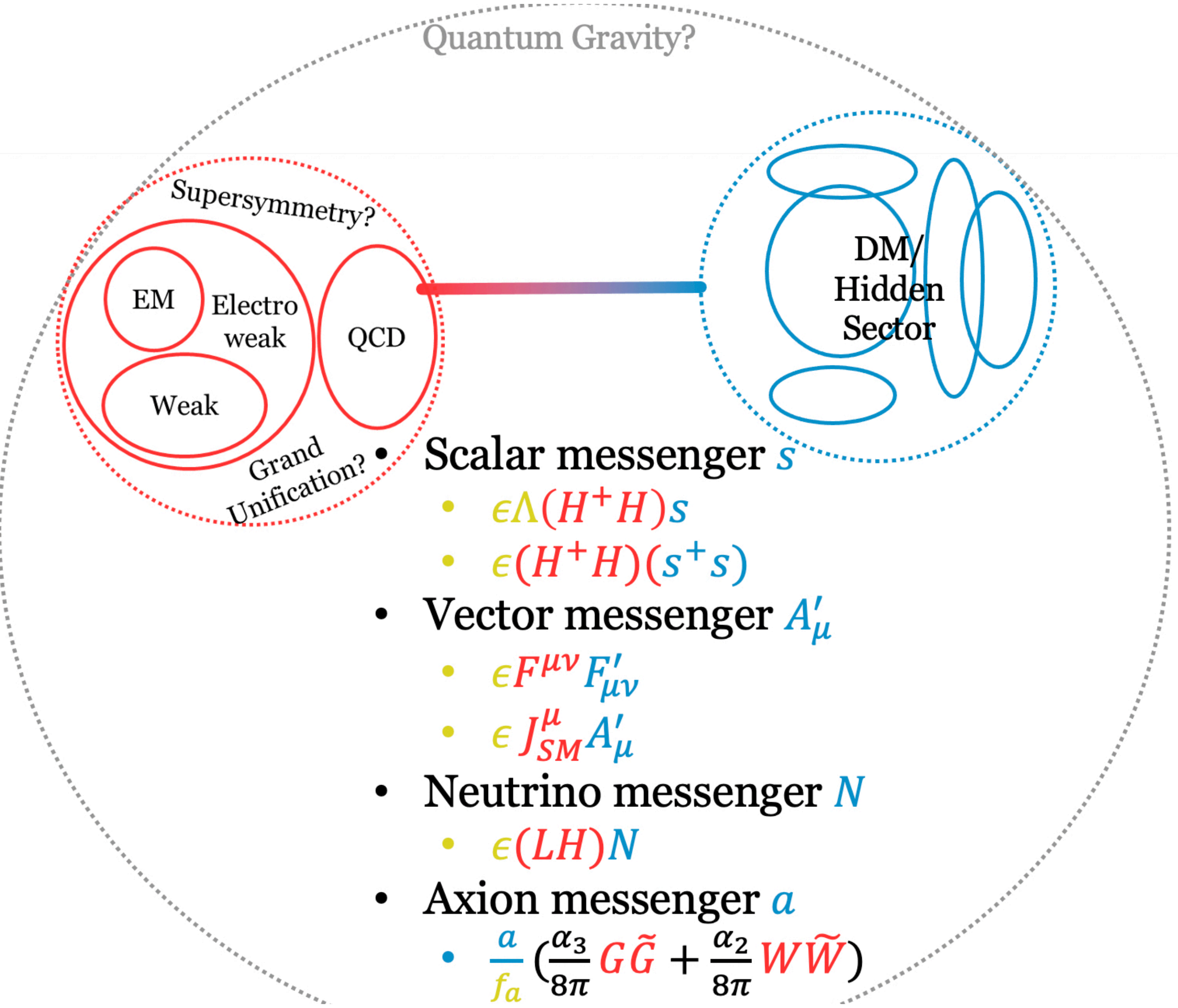


Fig. from mid-term report

Exotics/Long Lived Particles.

Z. Liu @ CEPC 2020



- Scalar messenger s
  - $\epsilon\Lambda(H^+H)s$
  - $\epsilon(H^+H)(s^+s)$
- Vector messenger A'_μ
  - $\epsilon F^{\mu\nu}F'_{\mu\nu}$
  - $\epsilon\, J^\mu_{SM} A'_\mu$
- Neutrino messenger N
  - $\epsilon(LH)N$
- Axion messenger a
  - $\frac{a}{f_a}(\frac{\alpha_3}{8\pi}G\tilde{G} + \frac{\alpha_2}{8\pi}W\tilde{W})$

The Higgs could be a good portal to Dark Sector

— rich exotic signatures —

Decay Topologies	Decay mode $\mathcal{F}_i$
$h \to 2$	$h \to \not{E}_T$
$h \to 2 \to 3$	$h \to \gamma + \not{E}_T$
	$h \to (b\bar{b}) + \not{E}_T$
	$h \to (jj) + \not{E}_T$
	$h \to (\tau^+\tau^-) + \not{E}_T$
	$h \to (\gamma\gamma) + \not{E}_T$
	$h \to (\ell^+\ell^-) + \not{E}_T$
$h \to 2 \to 3 \to 4$	$h \to (b\bar{b}) + \not{E}_T$
	$h \to (jj) + \not{E}_T$
	$h \to (\tau^+\tau^-) + \not{E}_T$
	$h \to (\gamma\gamma) + \not{E}_T$
	$h \to (\ell^+\ell^-) + \not{E}_T$
	$h \to (\mu^+\mu^-) + \not{E}_T$
$h \to 2 \to (1+3)$	$h \to b\bar{b} + \not{E}_T$
	$h \to jj + \not{E}_T$
	$h \to \tau^+\tau^- + \not{E}_T$
	$h \to \gamma\gamma + \not{E}_T$
	$h \to \ell^+\ell^- + \not{E}_T$

Decay Topologies	Decay mode $\mathcal{F}_i$
$h \to 2 \to 4$	$h \to (b\bar{b})(b\bar{b})$
	$h \to (b\bar{b})(\tau^+\tau^-)$
	$h \to (b\bar{b})(\mu^+\mu^-)$
	$h \to (\tau^+\tau^-)(\tau^+\tau^-)$
	$h \to (\tau^+\tau^-)(\mu^+\mu^-)$
	$h \to (jj)(jj)$
	$h \to (jj)(\gamma\gamma)$
	$h \to (jj)(\mu^+\mu^-)$
	$h \to (\ell^+\ell^-)(\ell^+\ell^-)$
	$h \to (\ell^+\ell^-)(\mu^+\mu^-)$
	$h \to (\mu^+\mu^-)(\mu^+\mu^-)$
	$h \to (\gamma\gamma)(\gamma\gamma)$
	$h \to \gamma\gamma + \not{E}_T$
$h \to 2 \to 4 \to 6$	$h \to (\ell^+\ell^-)(\ell^+\ell^-) + \not{E}_T$
	$h \to (\ell^+\ell^-) + \not{E}_T + X$
$h \to 2 \to 6$	$h \to \ell^+\ell^-\ell^+\ell^- + \not{E}_T$
	$h \to \ell^+\ell^- + \not{E}_T + X$

LHC's strength

Hard at LHC due to missing energy

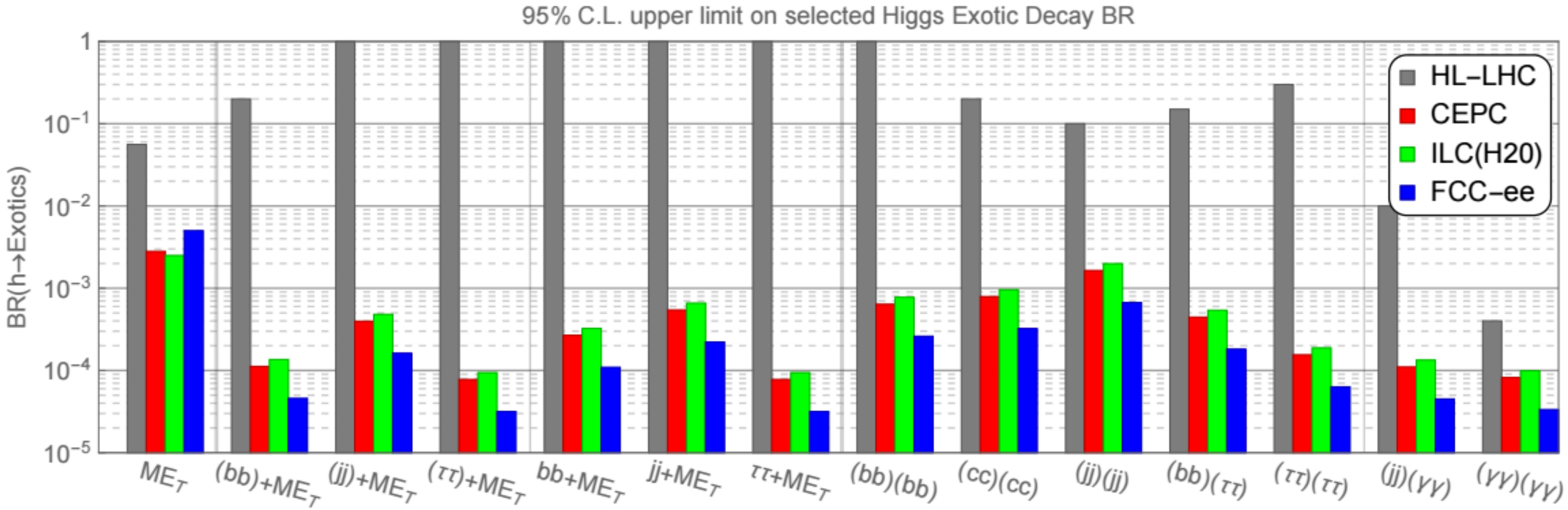
Hard at LHC due to hadronic background

Lepton colliders' strength

Exotics/Long Lived Particles.

Z. Liu @ CEPC 2020

The Higgs could be a good portal to Dark Sector
— rich exotic signatures —



How to improve?
> Dedicated detectors, see e.g. talk by R. Gonzalez Suarez @ FCC week 2021

Cost of Operation

Energy and carbon footprint.

- FCC-ee total instantaneous power demand at each centre-of-mass energies

J.-P. Burnet, FCC Week'22

		Z	W	H	TT
Beam energy (GeV)		45.6	80	120	182.5
Magnet current		25%	44%	66%	100%
Power ratio		6%	19%	43%	100%
PRF EL (MW)	Storage	146	146	146	146
PRFb EL (MW)	Booster	2	2	2	2
Pcryo (MW)	all	1,3	12,6	15,8	47,5
Pcv (MW)	all	33	34	36	40.2
PEL magnets (MW)	Stroage	6	17	39	89
PEL magnets (MW)	Booster	1	3	5	11
Experiments (MW)	Pt A & G	8	8	8	8
Data centers (MW)	Pt A & G	4	4	4	4
General services (MW)		36	36	36	36
Power during beam operation (MW)		237	262	291	384

Ongoing R&D

Ongoing R&D

Potential energy savings

- At 240 GeV, the instantaneous power of FCC-ee amounts to 291 MW
  - As a comparison, P(ILC_{250})=140 MW, P($CLIC_{380}$)=110 MW : less power hungry than FCC-ee?
    - Not clear: both produce (2 to 4 times) less Higgs than FCC-ee_{240}, with (3 to 6 times) longer running time

Energy and carbon footprint.

- Our first responsibility (as particle physicists) is to do the maximum of science
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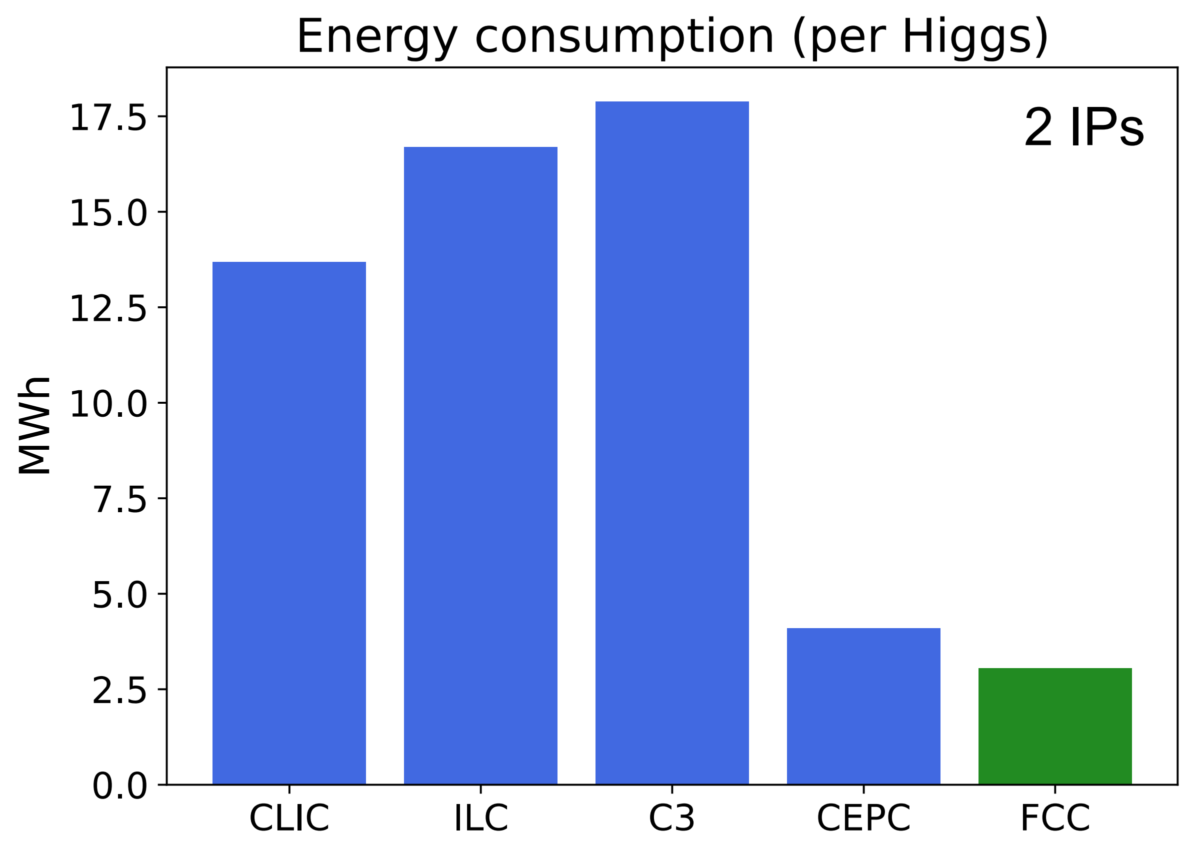

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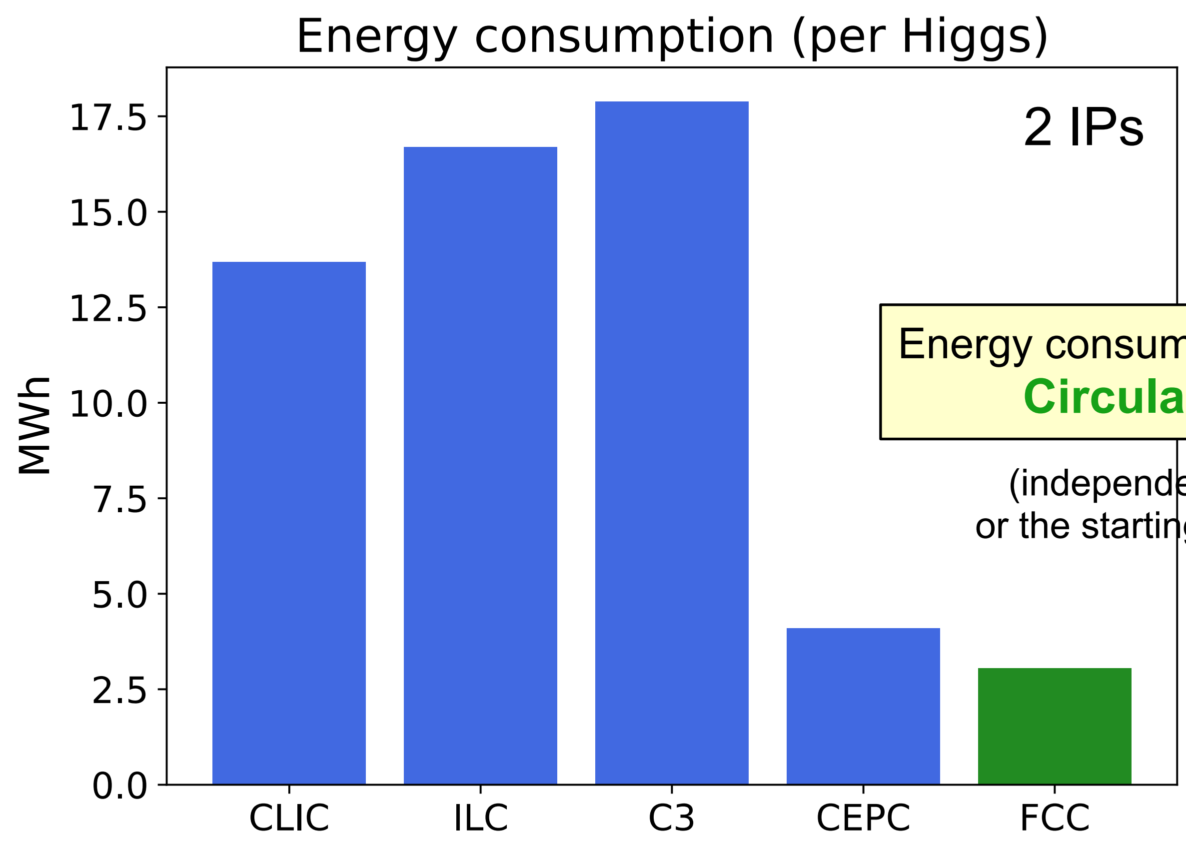

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Energy consumption / Higgs with 2IP
Circular ~ Linear / 5

(independently of the location
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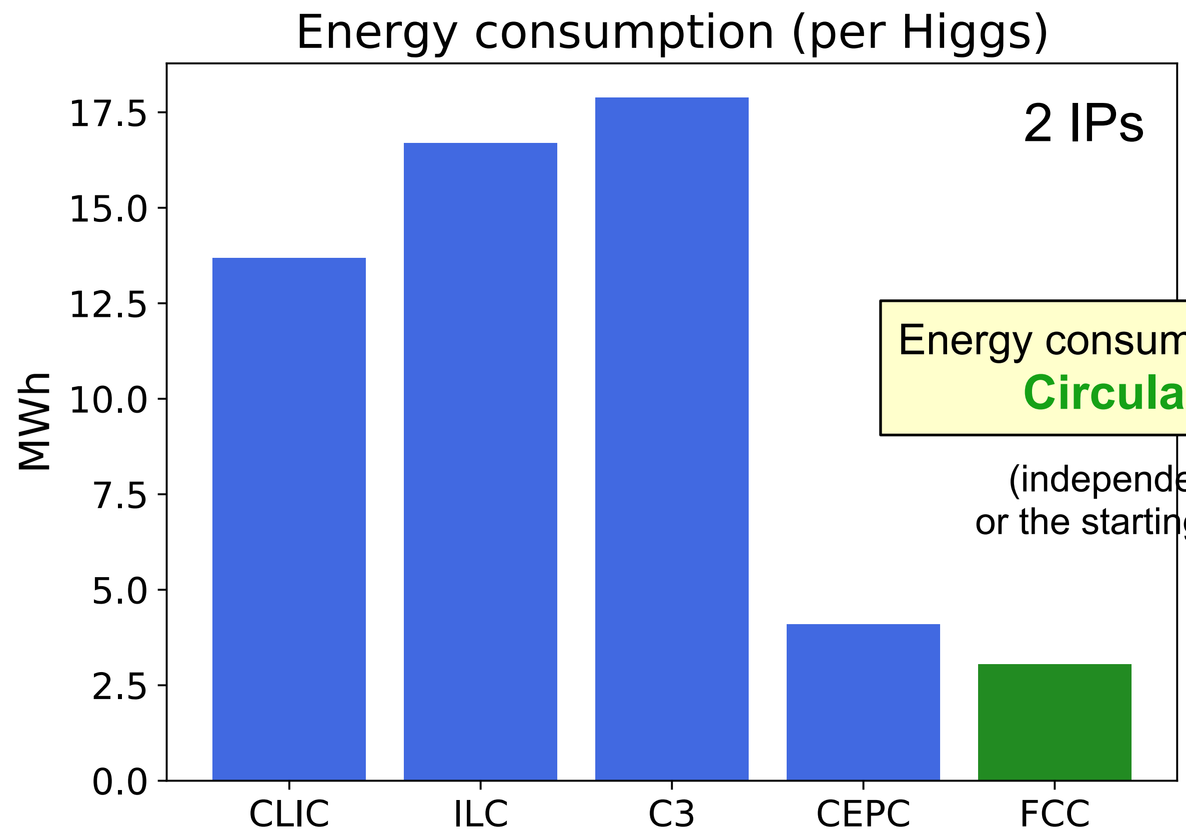

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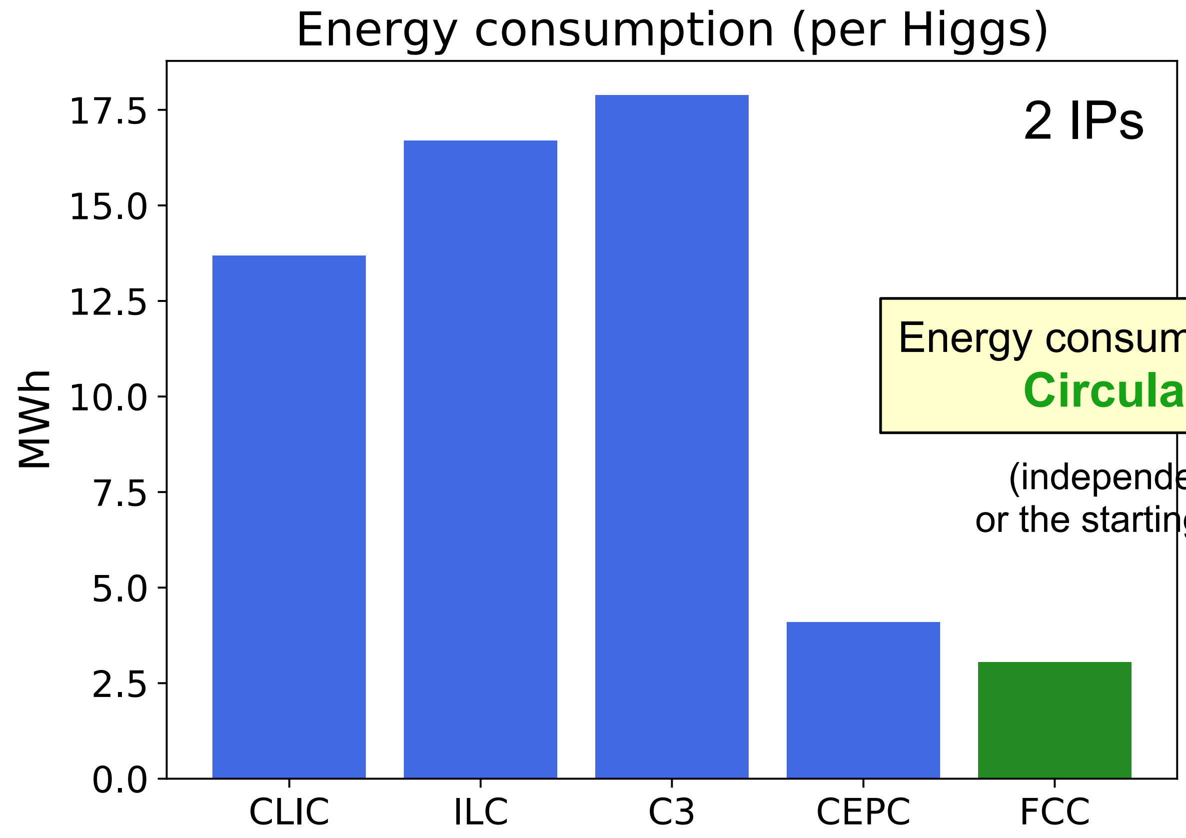

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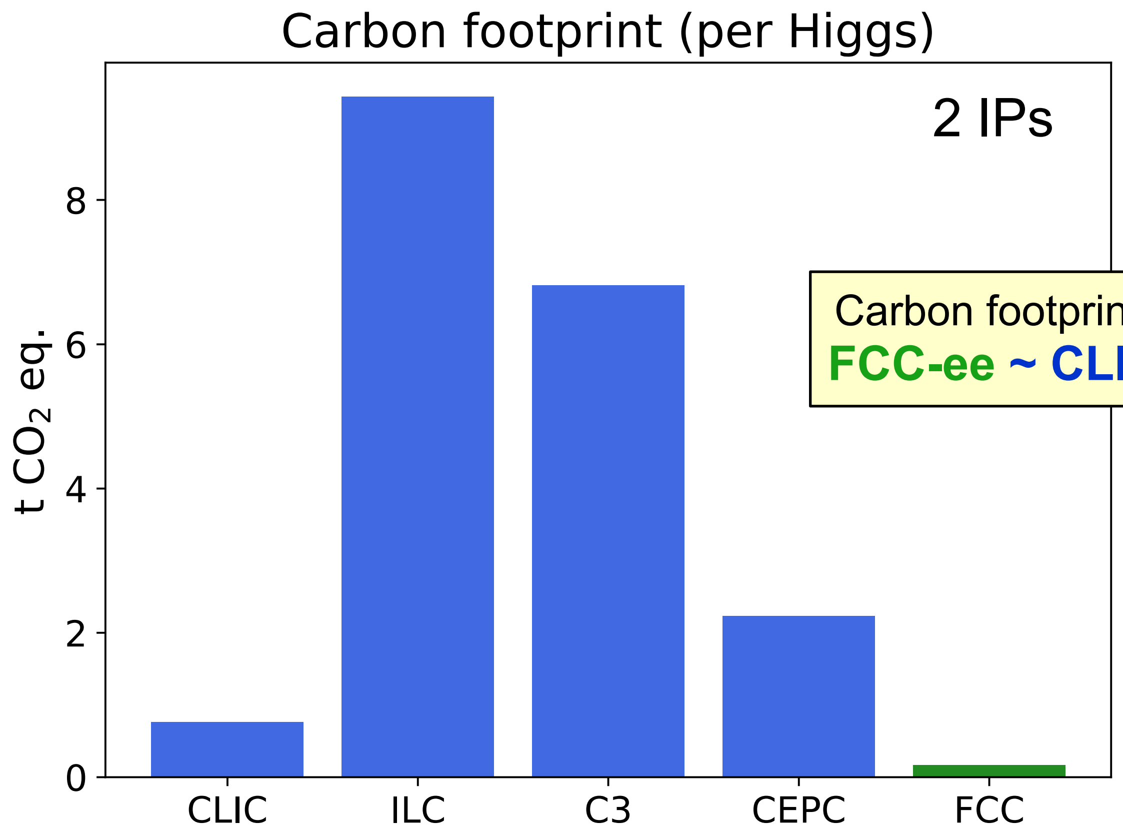

Carbon footprint / Higgs with 2IP
FCC-ee ~ CLIC / 5 ~ ILC / 50

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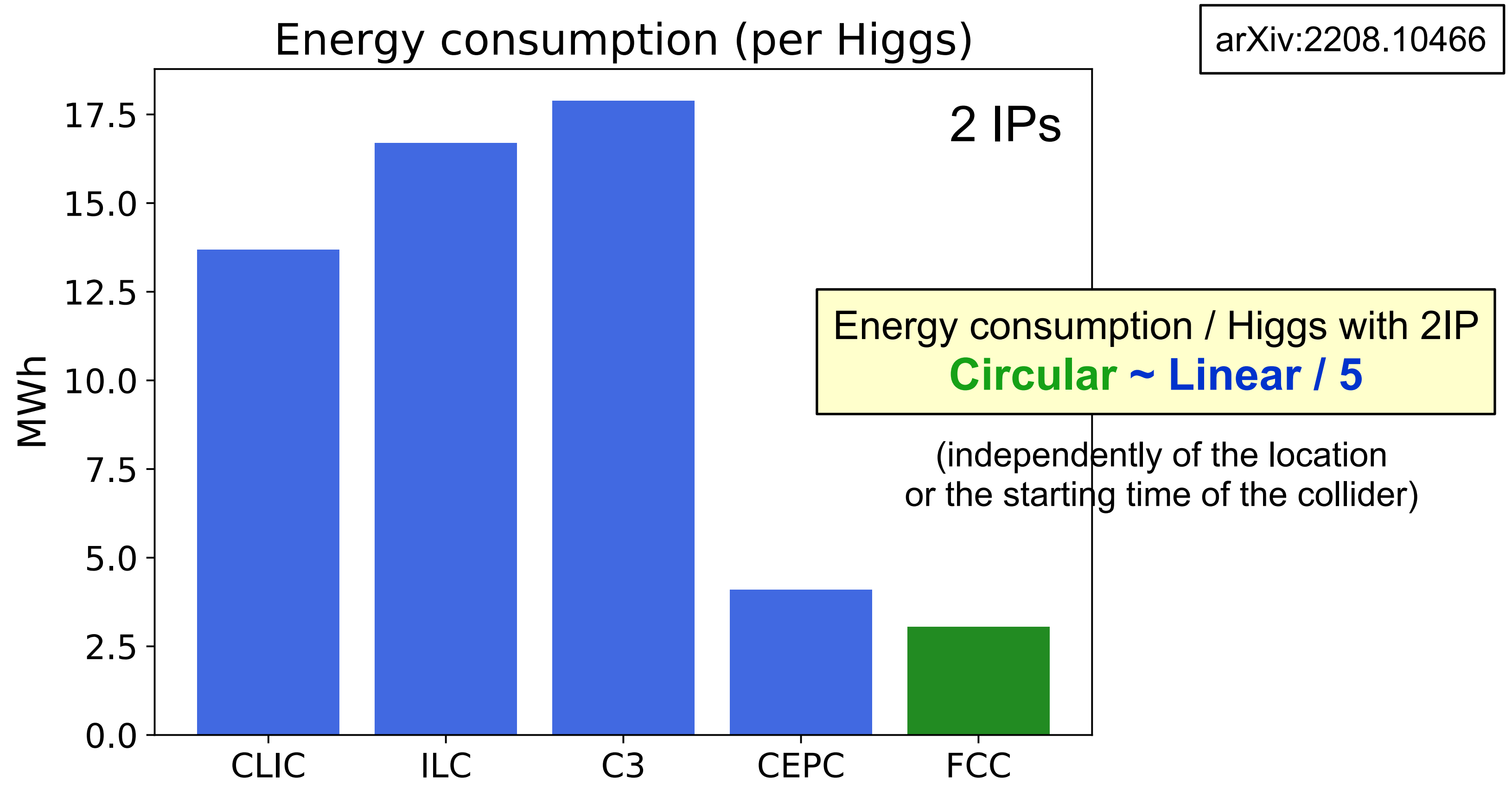

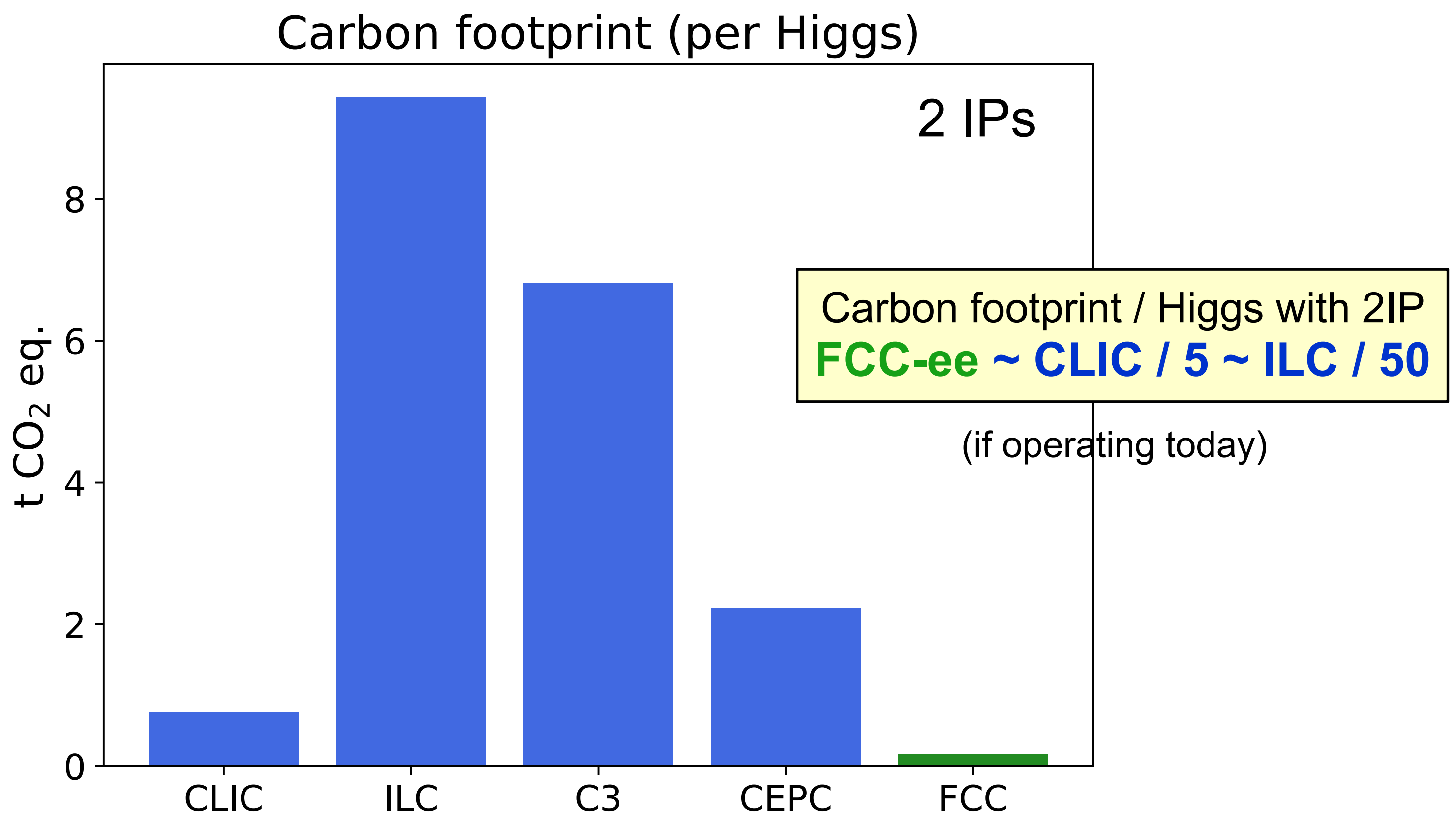

Cost of Operation.

The total electrical energy consumption over the fourteen years of the FCC-ee research programme is estimated to be around 27 TWh [58], corresponding to an average electricity consumption of 1.9 TWh/year over the entire operation programme, to be compared with the 1.2 TWh/year consumed by CERN today and the expected 1.4 TWh/year for HL-LHC[9]. At the CERN electricity prices from 2014/15, the electricity cost for FCC-ee collider operation would be about 85 MEuro per year. In the HZ running mode, about one million Higgs bosons are expected to be produced in three years, which sets the price of each FCC-ee Higgs boson at 255 Euros. A similar exercise can be done for the first stage of CLIC, expected to consume 0.8 TWh/year over 8 years at 380 GeV to produce about 150,000 Higgs bosons, which sets the price of a CLIC Higgs boson at about 2000 Euros. Finally, with the official ILC operation cost in Japan of 330 MEuro per year [10], its 11.5 to 18.5 years of operation (Section 5), and the 500,000 Higgs bosons produced in total, the price of an ILC Higgs boson is between 7,000 and 12,000 Euros, i.e., between 30 and 50 times more expensive than at FCC-ee. These operation costs are summarized in Table 8.

Table 8: Operation costs of low-energy Higgs factories, expressed in Euros per Higgs boson.

Collider	ILC_{250}	CLIC_{380}	FCC-ee_{240}
Cost (Euros/Higgs)	7,000 to 12,000	2,000	255

FCC-ee, 1906.02693