

Controlling Jets with SCET

Jonathan Walsh
UC Berkeley / LBL

work with

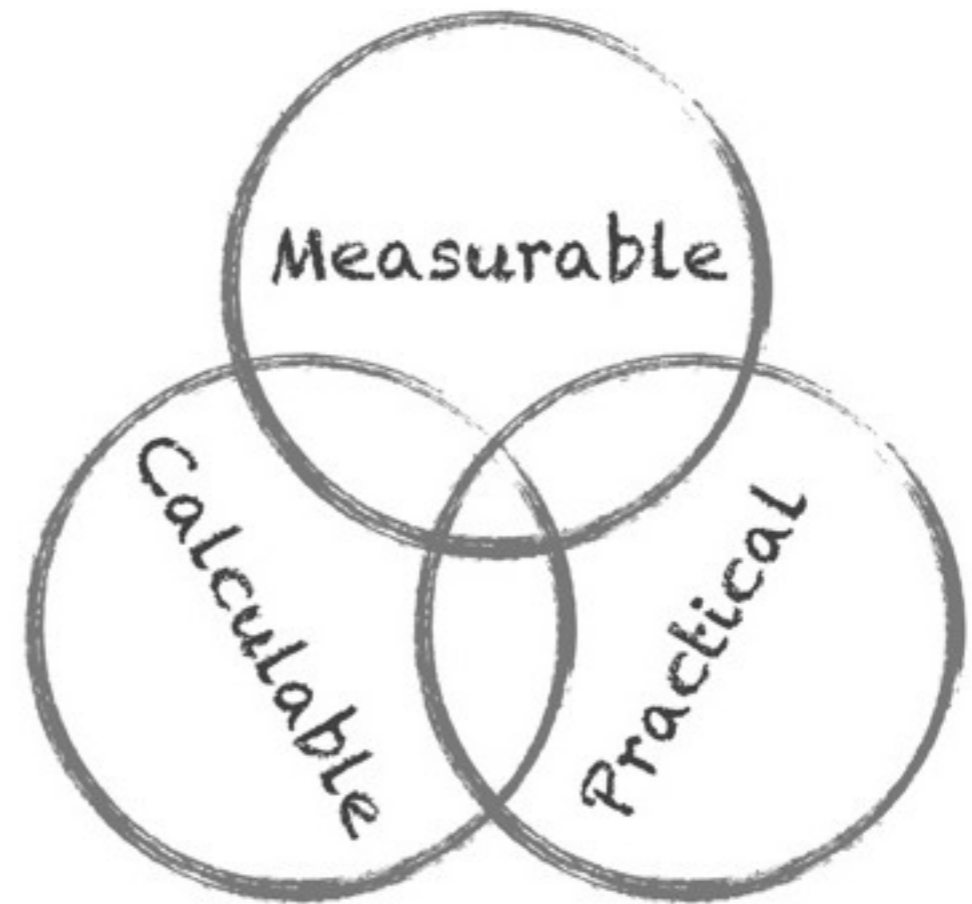
Christian Bauer, Frank Tackmann, Saba Zuberi

Outline

- Accessing jet structure
- Constructing an effective theory for jet structure
- Future Outlook

Calculable, Measurable, Predictable

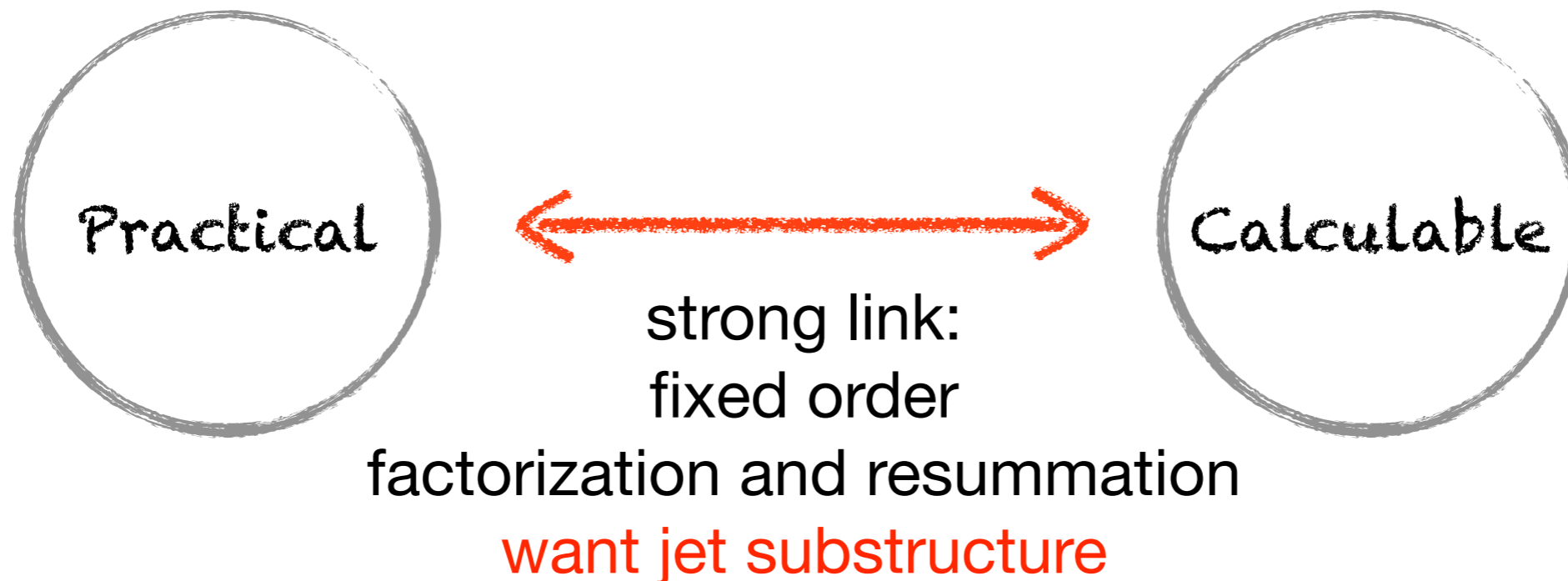
- Is jet substructure measurable?
 - Yes!
- Is jet substructure practical?
 - Yes!
- Is jet substructure calculable?
 - Kind of...



Calculable Approaches

Analytic approaches (pQCD & SCET, factorization and resummation)

- Pros: well controlled, provides accuracy, helps improve MC
- Cons: inclusive observables, currently limited to simple configurations, difficult to add a full range of physics



Calculable Framework for Jet Substructure

- Calculation-friendly algorithms, observables

Catani, Dokshitzer, Olsson, Turnock, Webber; Ellis, Soper;
Dokshitzer, Leder, Moretti, Webber; Cacciari, Salam, Soyez; Bauer, Hornig, Tackmann

- Hard scattering for the LHC

Everyone ever

- Merging between multijet processes

Catani, Seymour; Catani, Krauss, Kuhn, Webber; Mangano + many MC authors

- Jet shapes and multijet events

Ellis; Seymour; Berger, Sterman; Almeida, Lee, Perez, Sterman, Sung, Virzi;
Ellis, Hornig, Lee, Vermilion, JW; Jouttenus, Stewart, Tackmann, Waalewijn + many event shapes authors

- Non-global logarithms

Dasgupta, Salam, Rubin, Banfi, Khelifa-Kerfa, Marzani,
Kelley, Schabinger, Schwartz, Zhu; Hornig, Lee, Stewart, JW, Zuberi

- Non-perturbative corrections and uncertainties

Lee, Sterman; Dasgupta, Magnea, Salam; Hoang, Stewart; Bauer, Lee, Manohar, Wise

N-jettiness

Stewart, Tackmann, Waalewijn
Jouttenus, Stewart, Tackmann, Waalewijn

for each jet:

$$\mathcal{T}_j = \sum_i n_j \cdot k_i$$

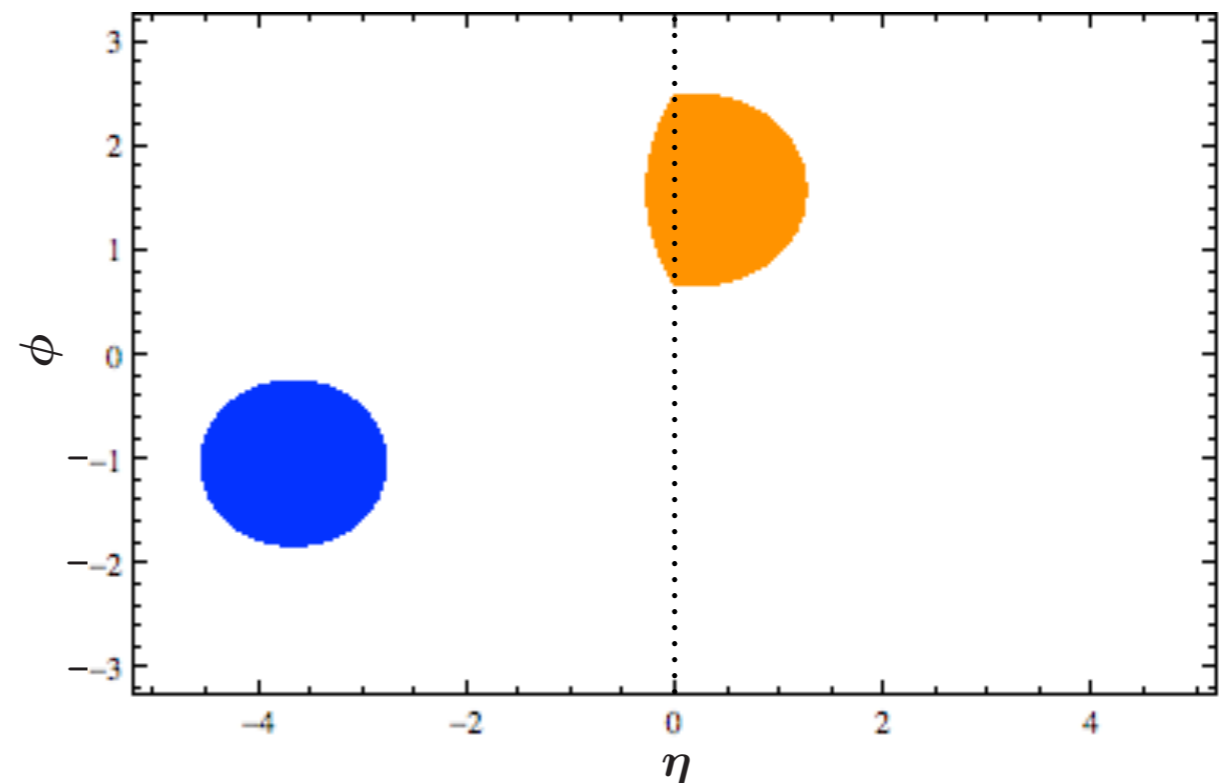
advantages:

inclusive over phase space
calculable
no boundary parameters

We will study a specific multijet configuration using N-jettiness

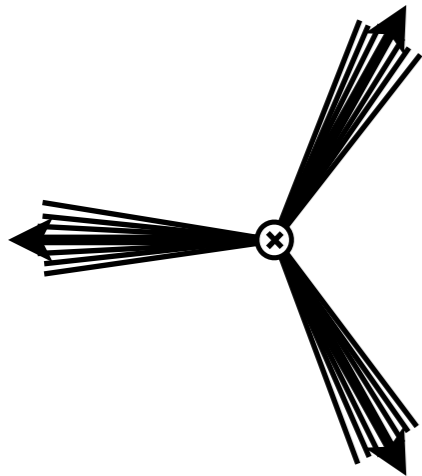
But the framework we use applies to other jet definitions and observables

boundary regions for jets, beam

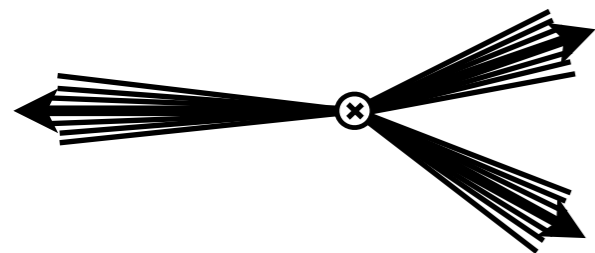


jet assignment depends only
on particle direction

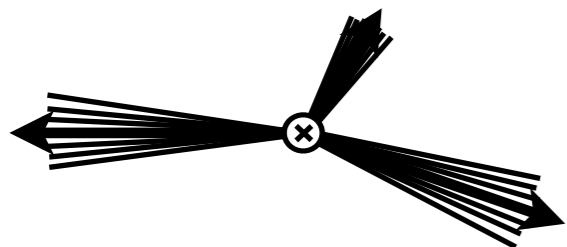
Multi-jet and Multi-subjet Events



uncommon
well-separated
energetic
all scales $\sim Q$

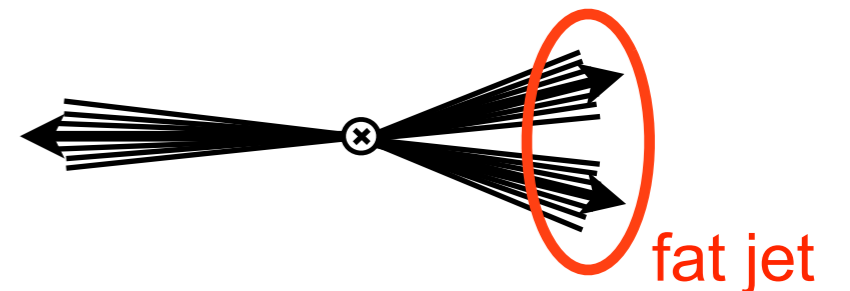


common
nearby jets
energetic
small dijet invariant mass t



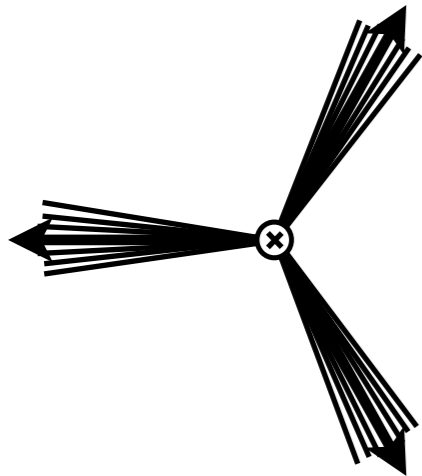
common
well-separated
hierarchy of jet energies
small dijet invariant masses

substructure limit:
2 jets merge



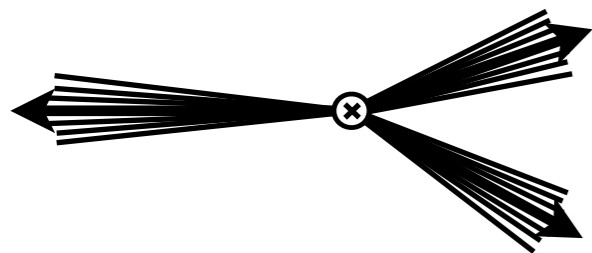
Can think about the
subjects as their own jets

Multi-jet and Multi-subjet Events



uncommon
 well-separated
 energetic
 all scales $\sim Q$

scales:
 center of mass energy
 dijet invariant masses
 jet masses
 soft scales



common
 nearby jets
 energetic
 small dijet invariant mass t

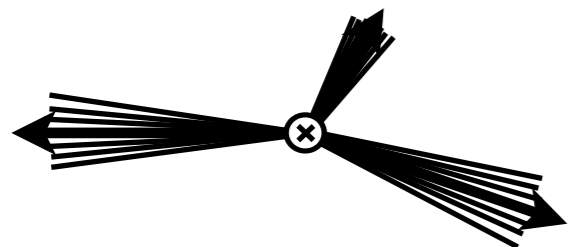
$$\sqrt{\hat{s}}$$

$$\sqrt{t}$$

$$m_J$$

$$m_J^2 / \sqrt{t}$$

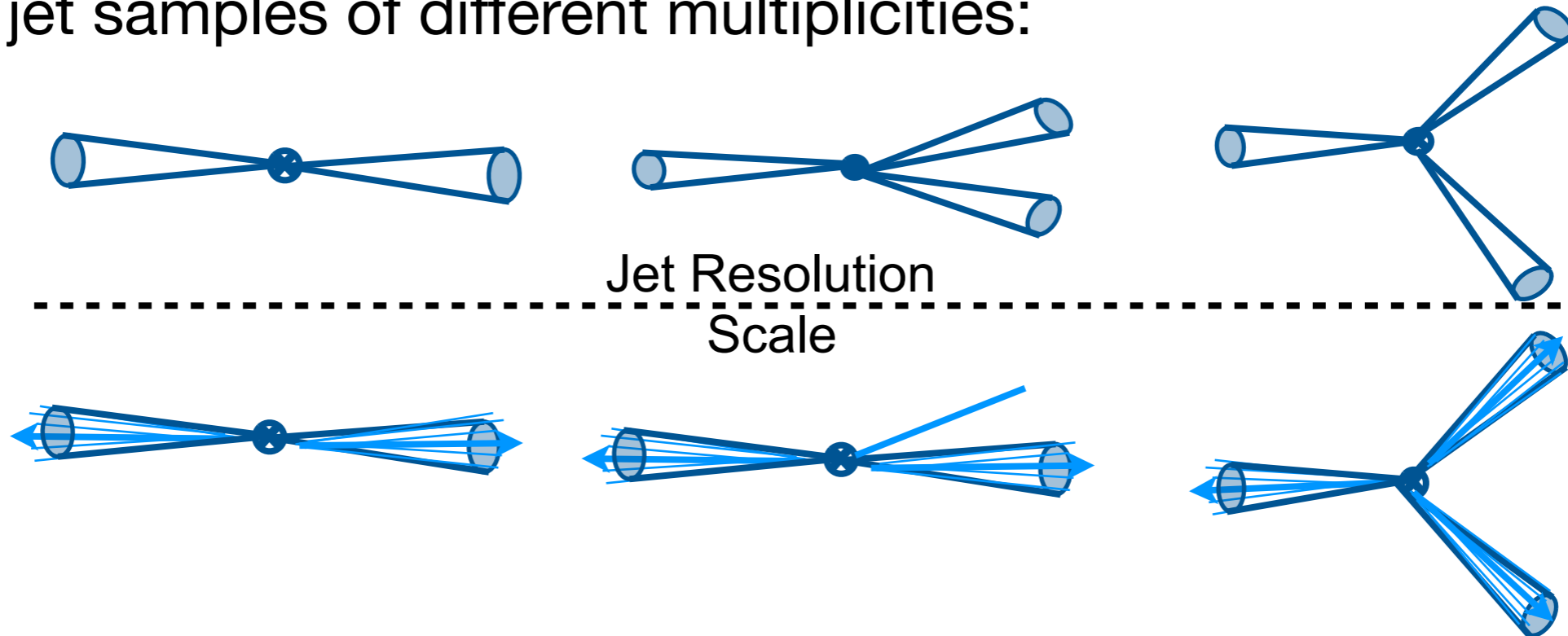
$$m_J^2 / \sqrt{\hat{s}}$$



common
 well-separated
 hierarchy of jet energies
 small dijet invariant masses

Multijets and Monte Carlo

- Merge jet samples of different multiplicities:



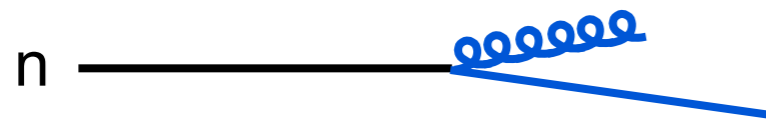
- Current state of the art: NLO/LL for 1 multiplicity + LO/LL for additional multiplicities
- Goal: NLO/LL for many multiplicities



Christian Bauer, Calvin Berggren, Nicholas Dunn, Andrew Hornig, Frank Tackmann, Jesse Thaler, Christopher Vermilion, Jonathan Walsh, Saba Zuberi

Modes for Jet Physics

collinear $p_c \sim Q(\lambda^2, 1, \lambda)$



light-cone
coords: $n_j = (1, \hat{n}_j)$
 $\bar{n}_j = (1, -\hat{n}_j)$

soft $p_s \sim Q(\lambda^2, \lambda^2, \lambda^2)$



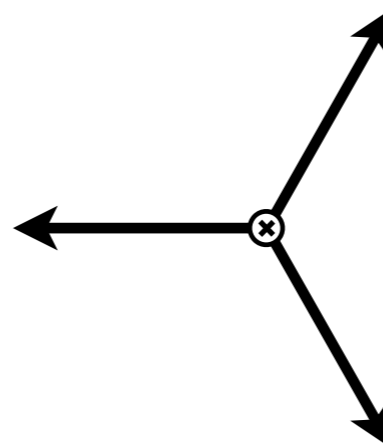
$k = (\bar{n} \cdot k, n \cdot k, k_\perp)$
 $= (k^-, k^+, k^\perp)$

$\lambda \equiv \frac{p_\perp}{Q} \ll 1$: defines the power counting parameter in SCET

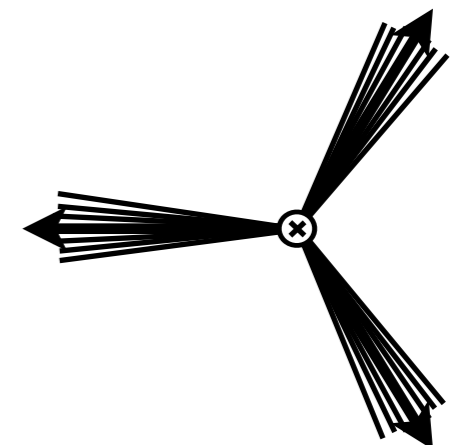
getting from QCD to SCET:

$\sqrt{\hat{s}}$ $\xrightarrow{\text{QCD}}$ $C_N(q_i)$
 $\xrightarrow{\text{SCET}}$
N-jet operator

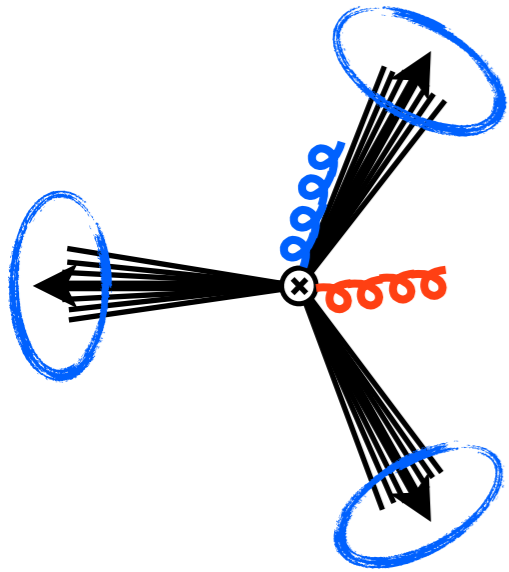
hard parton
configuration



jet, soft
evolution



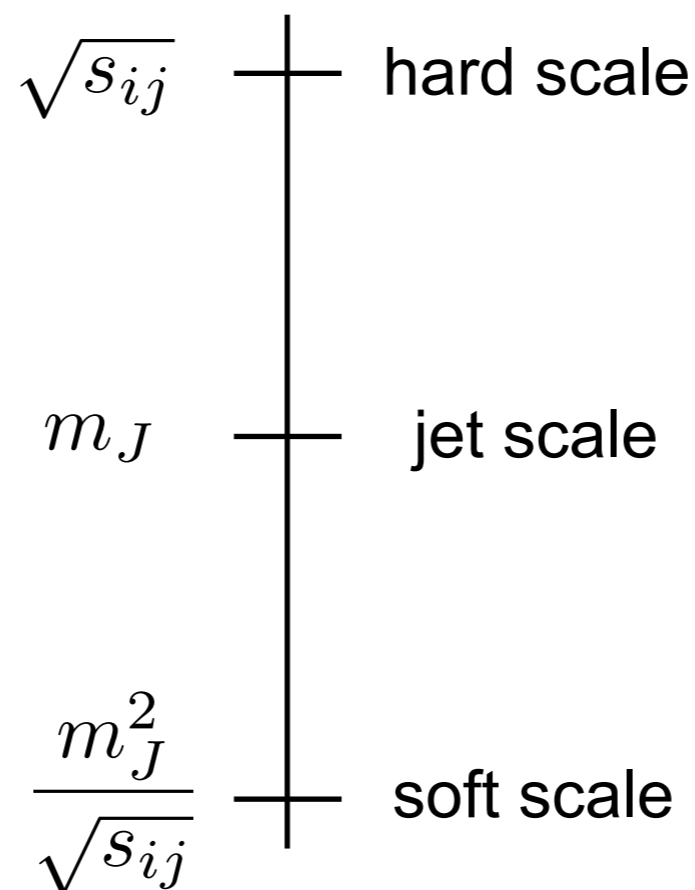
Factorization and Scales in Multijet Events



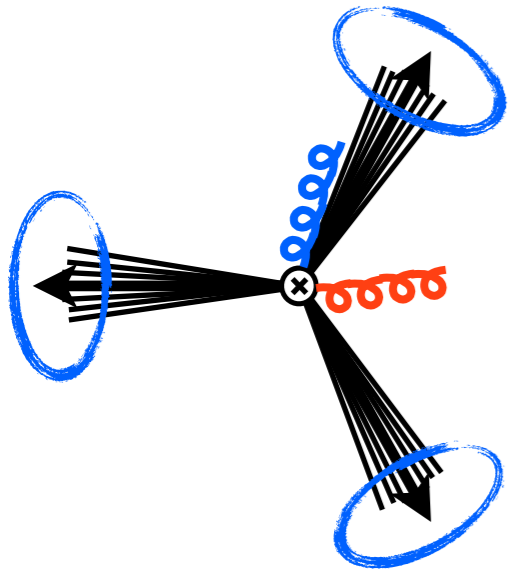
hard:	$p_h \sim \sqrt{s_{ij}}(1, 1, 1)$	$p_h^2 \sim s_{ij}$
collinear:	$p_c \sim E_J(1, \lambda^2, \lambda)$	$p_c^2 \sim E_J^2 \lambda^2 \sim m_J^2$
soft:	$p_s \sim E_J(\lambda^2, \lambda^2, \lambda^2)$	$p_s^2 \sim E_J^2 \lambda^4 \sim m_J^4 / E_J^2$

factorization theorem:

$$\frac{d\sigma}{d\mathcal{T}_i} = \frac{d\sigma^0}{d\Phi_3} H_3 J_i \otimes S_3$$



Scales in Multijet Events



hard: $p_h \sim \sqrt{s_{ij}}(1, 1, 1)$

$$p_h^2 \sim s_{ij}$$

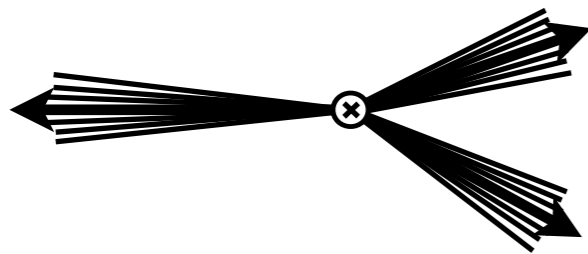
collinear: $p_c \sim E_J(1, \lambda^2, \lambda)$

$$p_c^2 \sim E_J^2 \lambda^2 \sim m_J^2$$

soft: $p_s \sim E_J(\lambda^2, \lambda^2, \lambda^2)$

$$p_s^2 \sim E_J^2 \lambda^4 \sim m_J^4 / E_J^2$$

jet configuration
changes



Hard scales become widely separated

Cannot sum large logarithms in the hard function
- same problem in the soft function

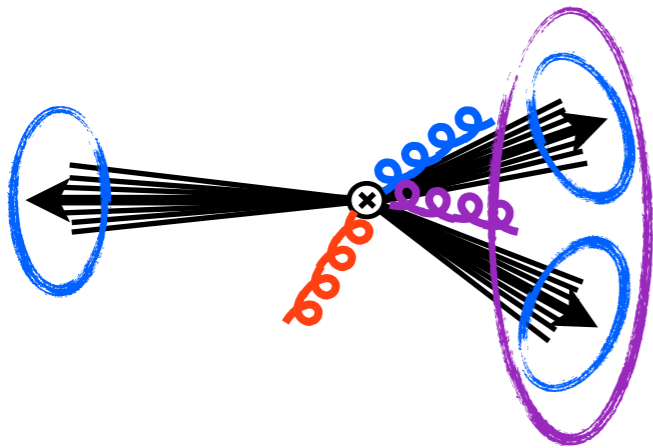
Obtain large logs of small angles: $\ln n_i \cdot n_j$

For nearby jets, need a
new mode for the dijet system
and a new factorization theorem

$\ln n_i \cdot n_j : \text{ninja}$



Solution: Add a New Mode



culprits:

need additional hard scattering factorization
soft radiation between the dijets lives at a different scale

Hard scattering factorization solved by

[Bauer, Schwartz](#)

[Baumgart, Marcantonini, Stewart](#)

We will add a new collinear-soft (**csoft**) mode which
contributes to the dijet system

Build this new mode into a new version of SCET

SCET₊: an EFT for multijets with small dijet invariant masses

Also useful for jet substructure: nearby subjects

Why Collinear and Soft Modes?

hard: $p_h \sim \sqrt{s_{ij}} (1, 1, 1)$

collinear: $p_c \sim E_J(1, \lambda^2, \lambda)$

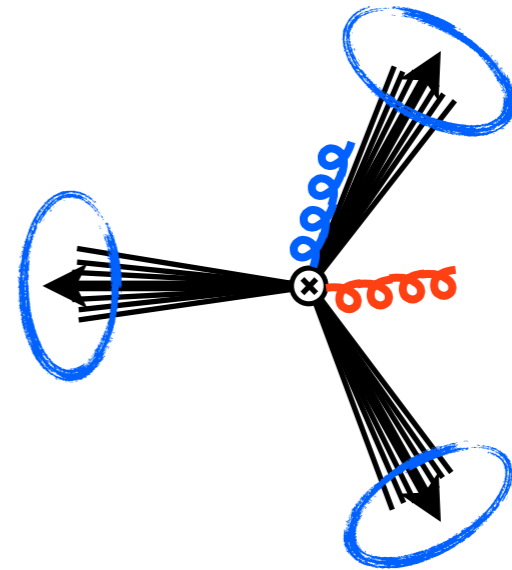
soft: $p_s \sim E_J(\lambda^2, \lambda^2, \lambda^2)$

relate them to the observable:

$$\mathcal{T}_j(p) = n_j \cdot p$$

$$\mathcal{T}_j(p_c) = n_j \cdot p_c \sim E_J \lambda^2 + p_c^2 \sim E_J^2 \lambda^2 \Rightarrow p_c \sim E_J(1, \lambda^2, \lambda)$$

$$\mathcal{T}_j(p_s) = n_j \cdot p_s \sim E_J \lambda^2 + p_s^2 \sim E_J^2 \lambda^4 \Rightarrow p_s \sim E_J(\lambda^2, \lambda^2, \lambda^2)$$



asking for a given sized contribution to the observable
and certain kinematics (collinear, soft) fixes the scaling of the modes

Modes with Nearby Jets: Hard Modes

Bauer, Schwartz
Baumgart, Marcantonini, Stewart

QCD

$$\sqrt{s_{ij}} \text{ ————— } C_2(q_i) \quad \text{hard: } p_h \sim \sqrt{s_{ij}} (1, 1, 1)$$

O_2
2-jet operator



resolve 2 jets

Modes with Nearby Jets: Hard Modes

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QCD

$$\sqrt{s_{ij}} \text{ ————— } C_2(q_i) \quad \text{hard: } p_h \sim \sqrt{s_{ij}} (1, 1, 1)$$

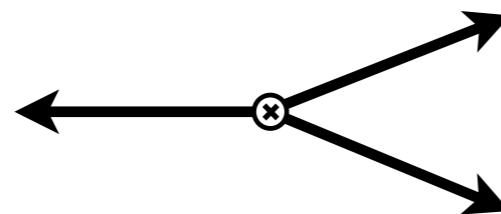
O_2
2-jet operator



resolve 2 jets

$$\sqrt{t} \text{ ————— } C_3(q_i) \quad \text{collinear: } p_c \sim E_J (1, \lambda_t^2, \lambda_t)$$

O_3
3-jet operator



resolve 3 jets

$$\lambda = \frac{m_J^2}{Q^2}$$

$$\lambda_t = \frac{t}{Q^2}$$

Modes with Nearby Jets: Hard Modes

Bauer, Schwartz
Baumgart, Marcantonini, Stewart

QCD

$$\sqrt{s_{ij}} \text{ ————— } C_2(q_i) \quad \text{hard: } p_h \sim \sqrt{s_{ij}} (1, 1, 1)$$

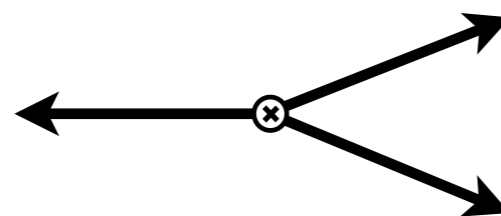
O_2
2-jet operator



resolve 2 jets

$$\sqrt{t} \text{ ————— } C_3(q_i) \quad \text{collinear: } p_c \sim E_J(1, \lambda_t^2, \lambda_t)$$

O_3
3-jet operator

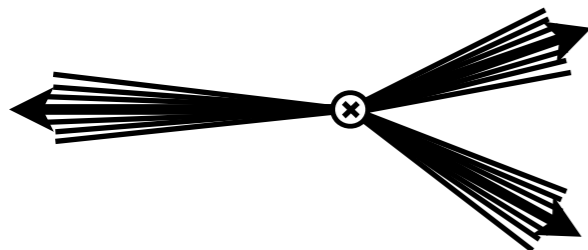


resolve 3 jets

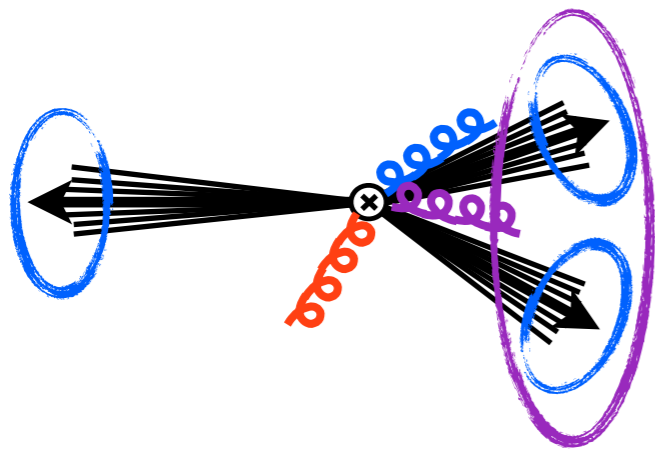
$$\lambda = \frac{m_J^2}{Q^2}$$

$$\lambda_t = \frac{t}{Q^2}$$

jet evolution
↓



Modes with Nearby Jets: Collinear and Soft Modes



collinear: $p_c \sim E_J(1, \lambda^2, \lambda)$

csoft: $p_{cs} \sim E_J \frac{\lambda^2}{\lambda_t^2} (1, \lambda_t^2, \lambda_t)$

soft: $p_s \sim E_J(\lambda^2, \lambda^2, \lambda^2)$

$$\lambda = \frac{m_J^2}{Q^2}$$

$$\lambda_t = \frac{t}{Q^2}$$

$$\mathcal{T}_j(p) = n_j \cdot p$$

$$\mathcal{T}_j(p_c) = n_j \cdot p_c \sim E_J \lambda^2 \quad \& \quad p_c^2 \sim E_J^2 \lambda^2 \quad \Rightarrow \quad p_c \sim E_J(1, \lambda^2, \lambda)$$

$$\mathcal{T}_j(p_{cs}) = n_j \cdot p_{cs} \sim E_J \lambda^2 \quad \& \quad \frac{p_{cs}^+}{p_{cs}^-} \sim \lambda_t^2 \quad \Rightarrow \quad p_{cs} \sim E_J \frac{\lambda^2}{\lambda_t^2} (1, \lambda_t^2, \lambda_t)$$

$$\mathcal{T}_j(p_s) = n_j \cdot p_s \sim E_J \lambda^2 \quad \& \quad p_s^2 \sim E_J^2 \lambda^4 \quad \Rightarrow \quad p_s \sim E_J(\lambda^2, \lambda^2, \lambda^2)$$

Scales and the Factorization Theorem

Complete factorization in SCET₊

QCD	$\sqrt{s_{ij}}$	hard 2 jet	\longrightarrow	$\ln \frac{\mu}{\sqrt{s_{ij}}}$
SCET O_2	\sqrt{t}	hard 3 jet	\longrightarrow	$\ln \frac{\mu}{\sqrt{t}}$
SCET ₊ O_3	m_J	jet	\longrightarrow	$\ln \frac{\mu}{m_J}$
soft ₊ S_3	$\frac{m_J^2}{\sqrt{t}}$	csoft	\longrightarrow	$\ln \frac{\mu}{m_J^2/\sqrt{t}}$
soft S_2	$\frac{m_J^2}{\sqrt{s_{ij}}}$	soft	\longrightarrow	$\ln \frac{\mu}{m_J^2/\sqrt{s_{ij}}}$

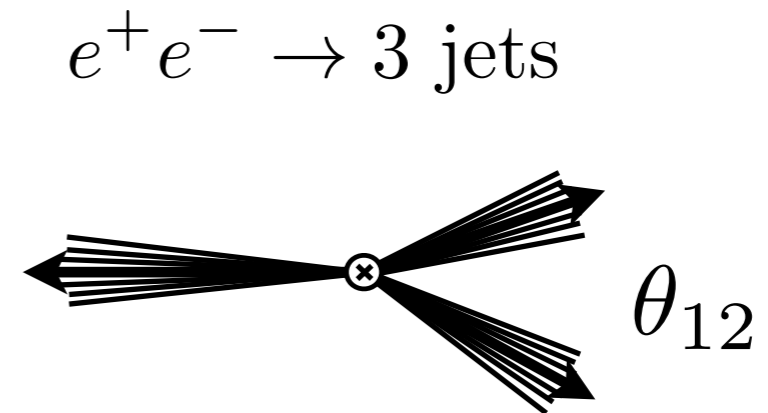
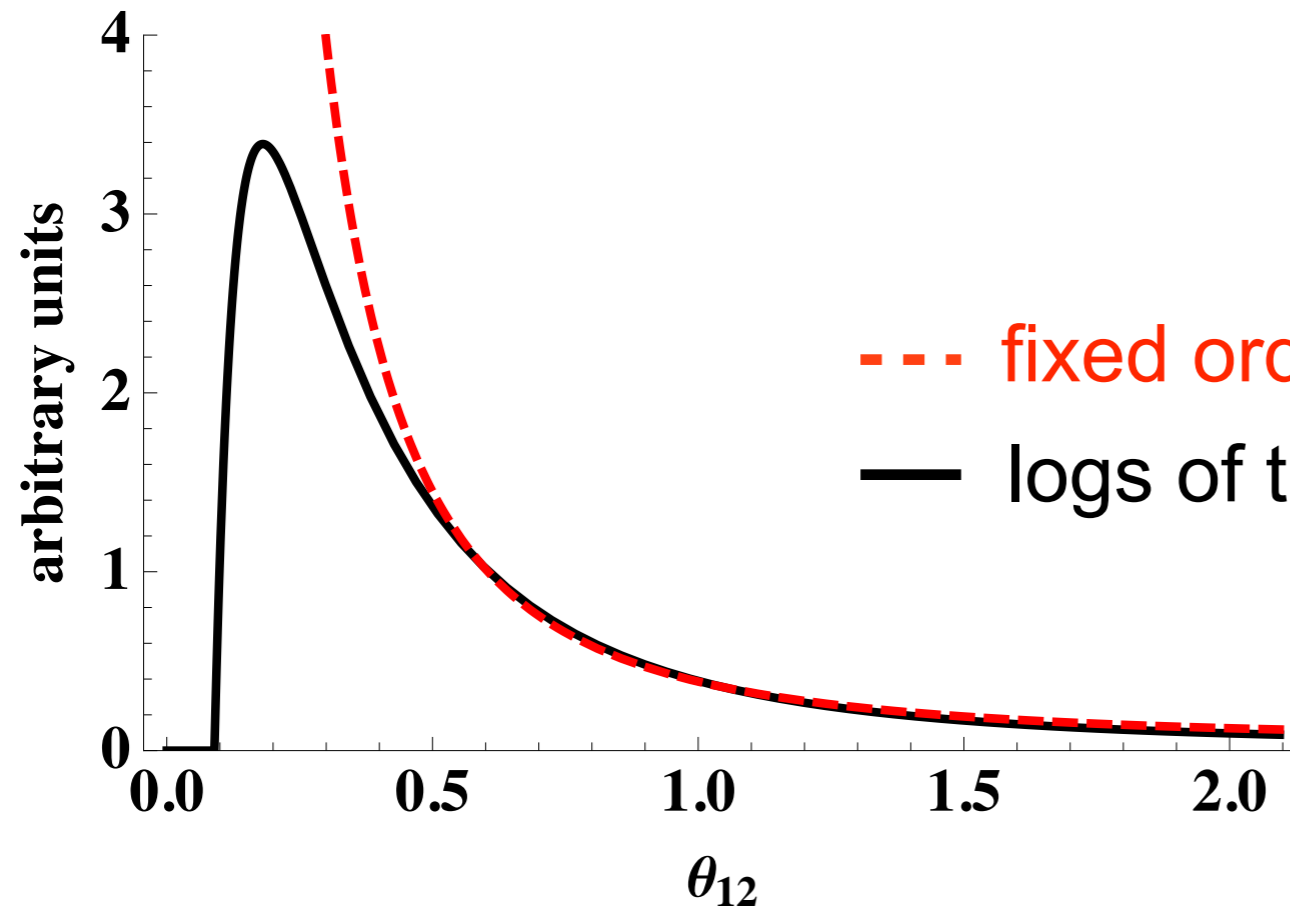
$$\frac{d\sigma}{d\mathcal{T}_i} = \frac{d\sigma^0}{d\Phi_3} H_2 H_3^+ J_i \otimes S_c \otimes S_2$$

all functions
fully factorized
(one scale per function)

Resumming Kinematic Logs

preliminary

$$\frac{d\sigma}{dt}$$



rate for jets with 3-jettiness less than a cutoff

$$Q\tau_3 < 0.1$$

Conclusions

- Worthwhile to access jet substructure through analytic approaches
- SCET developing more tools to describe realistic jets
- Have built an effective theory to describe nearby energetic jets
 - Applications for jet substructure
 - Factorization constraints on jet substructure algorithms,
JW & Zuberi, 1106.xxxx
- Goal is convergence with experimental measurements