Functional QCD and the QCD phase structure

PhD School XQCD 2024

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GEFÖRDERT VOM



Bundesministerium für Bildung und Forschung



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Phase structure of QCD















Phase structure of QCD







Collection of reviews & lecture notes on the FRG & DSE

Structure of the FRG: Aspects of the FRG

The nonperturbative functional renormalization group and its applications

QCD at finite temperature and density within the fRG approach: An overview

Material

Topical reviews

JMP, Annals Phys. 322 (2007) 2831-2915

Dupuis et al, Phys.Rep. 910 (2021) 1-114

Fu, Commun.Theor.Phys. 74 (2022) 9, 097304

Outline

I) Functional Renormalisation group

(II) Functional QCD and the QCD phase structure

(I) Functional Renormalisation Group for QCD

Introduction to the functional renormalisation group

- Derivation of the flow equation
- Spontaneous symmetry breaking
- Systematic error control & optimisation

• Functional flows for QCD

- Flows for correlation functions & chiral symmetry breaking
- Getting dynamical: emergent hadrons & diquarks
- Dynamical hadronisation at work

(II) Functional QCD and the QCD phase structure

QCD at finite temperature and density

- Benchmarks in the vacuum
- Correlation functions at finite temperature
- Polyakov loop from functional approaches

QCD phase structure

- Locating the QCD phase boundary and the critical end point
- Fluctuations of conserved charges: Ripples of the critical end point

(I) Functional Renormalisation Group for QCD

Introduction to the functional renormalisation group

- Derivation of the flow equation
- Spontaneous symmetry breaking
- Systematic error control & optimisation

• Functional flows for QCD

- Flows for correlation functions & chiral symmetry breaking
- Getting dynamical: emergent hadrons & diquarks
- Dynamical hadronisation at work

Introduction to the functional renormalisation group

Derivation of the flow equation

Generating functional Z

$$Z[J] = \frac{1}{\mathcal{N}} \int d\varphi \, e^{-S[\varphi] + \int_x \, J\varphi}$$

partition function

$$S[\varphi] = \frac{1}{2} \int_{x} \left[\partial_{\mu} \varphi \partial_{\mu} \varphi + \mathbf{m}^{2} \varphi \right]$$

classical action

zero-dimensional example: 'Functional' flows for integrals





Generating functional Z

$$Z[J] = \frac{1}{\mathcal{N}} \int d\varphi \, e^{-S[\varphi] + \int_x \, J\varphi}$$

partition function

Effective action Γ

$$\Gamma[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\hat{\varphi} + \phi]}$$

free energy





$$\begin{split} \varphi &= \hat{\varphi} + \phi \\ \langle \hat{\varphi} \rangle_{\frac{\delta\Gamma}{\delta\phi}} &= 0 \\ J &= \frac{\delta\Gamma}{\delta\phi} \end{split}$$

Generating functional Z

$$Z[J] = \frac{1}{\mathcal{N}} \int d\varphi \, e^{-S[\varphi] + \int_x \, J\varphi}$$

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Effective action Γ

$$\Gamma[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\hat{\varphi} + \phi]}$$

free energy

$$\Gamma[\phi] = \sup_{J} \left(\int_{x} J \cdot \phi - \log Z \right)$$

Legendre transform





$$\begin{split} \varphi &= \hat{\varphi} + \phi \\ \langle \hat{\varphi} \rangle_{\frac{\delta\Gamma}{\delta\phi}} &= 0 \\ J &= \frac{\delta\Gamma}{\delta\phi} \end{split}$$



Generating functional Z

$$Z[J] = \frac{1}{\mathcal{N}} \int d\varphi \, e^{-S[\varphi] + \int_x \, J\varphi}$$

partition function

Effective action Γ

$$\Gamma[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\hat{\varphi} + \phi] + \int_x \, \hat{\varphi} \, \frac{\delta \Gamma[\phi]}{\delta \phi}}$$

free energy

Dyson-Schwinger equation



quantum equation of motion

$$\langle \varphi \rangle_J = \phi$$

$$egin{aligned} & arphi &= \hat{arphi} + \phi \ & \langle \hat{arphi}
angle_{rac{\delta\Gamma}{\delta\phi}} &= 0 \ & J = rac{\delta\Gamma}{\delta\phi} \end{aligned}$$

Dyson-Schwinger equation

$$\frac{\delta\Gamma[\phi]}{\delta\phi(x)} = \left\langle \frac{\delta S[\hat{\varphi} + \phi]}{\delta\phi(x)} \right\rangle$$

Diagrammatics

$$S[\phi] = \frac{1}{2} \int_{x} \left[\partial_{\mu} \phi \partial_{\mu} \phi + m^{2} \phi^{2} + \frac{\lambda}{4} \phi^{4} \right]$$

Dyson-Schwinger equation

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$$\frac{\delta\Gamma[\phi]}{\delta\phi(x)} = \left\langle \frac{\delta S[\hat{\varphi} + \phi]}{\delta\phi(x)} \right\rangle$$

Diagrammatics

$$S[\phi] = \frac{1}{2} \int_{x} \left[\partial_{\mu} \phi \partial_{\mu} \phi + m^{2} \phi^{2} + \frac{\lambda}{4} \phi^{4} \right]$$











Effective action Γ

$$\Gamma[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\hat{\varphi}+\phi] + \int_x \, \hat{\varphi} \, \frac{\delta\Gamma[\phi]}{\delta\phi}}$$

No quantum fluctuations

$$\Gamma[\phi] = -\log e^{-S[\phi]} = S[\phi]$$



Effective action Γ

$$\Gamma[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\hat{\varphi} + \phi]}$$



Effective action Γ

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\phi + \hat{\varphi}] + \frac{1}{2} \int_p}$$



 $\hat{\varphi}(p)R_k(p^2)\hat{\varphi}(-p) + \int_x \hat{\varphi}(x) \frac{\delta\Gamma_k[\phi]}{\delta\phi(x)}$

DSE $\frac{\delta \Gamma_k[\phi]}{\delta \phi(x)}$ $\delta S[\hat{\varphi} +$



Effective action Γ

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\phi + \hat{\varphi}] + \frac{1}{2} \int_p}$$



 $\hat{\varphi}(p)R_k(p^2)\hat{\varphi}(-p) + \int_x \hat{\varphi}(x) \frac{\delta\Gamma_k[\phi]}{\delta\phi(x)}$

 p^2

$$\frac{R_k(p^2)}{2k^2}$$

$$t = \log \frac{k}{\Lambda}$$

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Effective action Γ

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\phi + \hat{\varphi}] + \frac{1}{2} \int_p}$$



$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \left\langle \hat{\varphi}(p) \hat{\varphi}(-p) \right\rangle \partial_t R_k(p^2)$$

Flow

 $\hat{\varphi}(p)R_k(p^2)\hat{\varphi}(-p) + \int_x \hat{\varphi}(x) \frac{\delta\Gamma_k[\phi]}{\delta\phi(x)}$

$$t = \log \frac{k}{\Lambda}$$

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \left\langle \hat{\varphi}(p) \hat{\varphi}(-p) \right\rangle \partial_t R_k(p^2)$$

Propagator

Flow



$$G^{-1}[\phi] = \Gamma_k^{(2)}[\phi] + R_k$$

$$t = \log \frac{k}{\Lambda}$$

$$- = \langle \hat{\varphi}(x)\hat{\varphi}(y)\rangle_c$$



 $\Gamma_k^{(2)}[\phi]$ $S^{(2)}[\phi]$

$$\hat{\varphi}(p)\hat{\varphi}(-p)\rangle \partial_t R_k(p^2)$$

$$t = \log \frac{k}{\Lambda}$$

$$- = \langle \hat{\varphi}(x)\hat{\varphi}(y)\rangle_c$$

$$\Gamma_k^{(2)}[\phi] + R_k$$

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} \, e^{-S[\hat{\varphi}+\phi] + \frac{1}{2} \int_p \hat{\varphi}(p) R_k(p^2) \hat{\varphi}(-p) + \int_x \hat{\varphi} \, \frac{\delta \Gamma_k[\phi]}{\delta \phi}}$$

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Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \operatorname{Tr} \frac{1}{\Gamma_k^{(2)}}$$

Diagrammatics

(Inverse) propagator



Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \operatorname{Tr} \frac{1}{\Gamma_k^{(2)}}$$

(Inverse) propagator

fRG



DSE

 $\partial_t \Gamma^{(n)} = \operatorname{Flow}_n[\Gamma^{(m)}; m = 2, ..., n+2]$



$$\Gamma^{(n)} = \text{DSE}_n[S^{(m)}, \Gamma^{(m)}; m = 2, ..., n+2]$$

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Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \operatorname{Tr} \frac{1}{\Gamma_k^{(2)}}$$

Properties

- 1-loop exact
- closed
- **RG-scaling**
- energy/particle-number conservation

automatic







only in specific approximation schemes



energy/particle-number conservation

automatic

only in specific approximation schemes

$$\partial_t \Gamma_k[\phi] = \frac{1}{2}$$

Derivative expansion

- Expansion in powers of momenta
- Controlled in the presence of a mass gap $m_{
 m gap}$
- **Expansion parameter**



Vertex expansion

- Expansion in number n of external fields
- **Controlled in perturbation theory/presence of symmetries**
- Expansion parameter n

Mixtures, exact resummation schemes,

$$\operatorname{Fr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

$$\partial_t \Gamma_k[\phi] = \frac{1}{2}$$

Derivative expansion

- Expansion in powers of momenta
- Controlled in the presence of a mass gap $m_{
 m gap}$
- **Expansion parameter**



Lowest order: 0th order

$$\Gamma_k[\phi] = \frac{1}{2} \int_p \phi \, p^2 \phi + \int_x V_k(\phi) + O(p^2)$$

$$\operatorname{Fr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

$$\partial_t \Gamma^{(n)} = \operatorname{Flow}_n[\Gamma^{(m)}; m = 2, ..., n]$$



$$\Gamma_k^{(2)}[\phi](p,q) = \left(p^2 + V_k''(\phi)\right) (2\pi)^d \delta(p)$$





$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \operatorname{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

Derivative expansion

Lowest order: 0th order

$$\Gamma_k[\phi] = \frac{1}{2} \int_p \phi \, p^2 \phi + \int_x V_k(\phi) + O(p^2)$$

$$\partial_t V_k(\phi) = \frac{1}{d} \frac{\Omega_d}{(2\pi)^d} \frac{k^{2+d}}{k^2 + V_k''(\phi)}$$

$$\Gamma_k^{(2)}[\phi](p,q) = \left(p^2 + V_k''(\phi)\right) (2\pi)^d \delta(p)$$

$$R_{k,\text{opt}}(p^2) = (k^2 - p^2)\theta(k^2)$$

$$\partial_t R_{k,\text{opt}}(p^2) = 2k^2\theta(k^2)$$

 $\Gamma^{(2)}[\phi](p) + R_{k,\text{opt}}(p^2) = \left[k^2 + V''(\phi)\right]\theta(k^2 - p^2) + (p^2 + V''(\phi))\theta(p^2 - k^2)$



$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \operatorname{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

Derivative expansion

Lowest order: 0th order

$$\Gamma_k[\phi] = \frac{1}{2} \int_p \phi \, p^2 \phi + \int_x V_k(\phi) + O(p^2)$$

Flow

$$\partial_t V_k(\phi) = \frac{1}{d} \frac{\Omega_d}{(2\pi)^d} \frac{k^{2+d}}{k^2 + V_k''(\phi)}$$

$$\Omega_d = \frac{2\pi^{d/2}}{\Gamma(d/2)}$$

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$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \operatorname{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

Derivative expansion

Lowest order: 0th order

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Spontaneous symmetry breaking

Approximation schemes & phase structure



- bosonic flow is symmetry-restoring
- flow guarantees convexity of effective action





Approximation schemes & phase structure



- bosonic flow is symmetry-restoring
- flow guarantees convexity of effective action




Approximation schemes & phase structure



Example: 3d critical exponents with fRG

Simple approximation: LPA' $\Gamma_k[\phi] = \frac{1}{2} \int_p Z_\phi \,\phi(p) \, p^2 \,\phi(-p) + \int_x V_k(\phi)$

Taylor exp

A simple program to compute critical exponents in O(N)-models with the Wetterich equation

pansion
$$V_k(\phi) = \sum_{n=1}^{N_{\max}} \frac{\lambda_n}{n!} (\phi^2 - \phi_{0,k}^2)^n$$

Michael Scherer

Ising universality

$$N = 1: \ \nu_{\rm Ising} = 0.630$$

fRG: LPA'

 $N = 1: \ \nu_{\text{Ising}} = 0.637...$





 $\partial_t V_k(\phi) = -$



- bosonic flow is symmetry-restoring
- fermionic flow is symmetry-breaking
- competing dynamics decides about fate of symmetries
- flow guarantees convexity



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'governs general phase structures'



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'governs general phase structures'





 $\partial_t V_k(\phi) = -$



- bosonic flow is symmetry-restoring
- fermionic flow is symmetry-breaking
- competing dynamics decides about fate of symmetries
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Grossi, Wink, SciPost Phys. Core 6 (2023)

State of the art time steppers Ihssen, Sattler, Wink, CPC 300 (2024) 109182

'governs general phase structures'









$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \Gamma_k[\phi]$$

g_3 $\Gamma_{\Lambda} = S$ $A_k^{(3)}$ $R_{1}^{(2)}$ $R_k^{(1)}$ ${\stackrel{\bullet}{\bullet}} \{g_i\}$ $\Gamma_0 = \Gamma$ $\searrow g_2$ $\downarrow g_1$

full flow



approximated flow

Optimisation: find $R_k^{(2)}$!

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \Gamma_k[\phi]$$

g_3 $\Gamma_{\Lambda} = S$ $R_{k}^{(3)}$ (2) $R_k^{(1)}$ • $\{g_i\}$ $\Gamma_0 = \Gamma$ $\blacktriangleright g_2$ $\checkmark g_1$

full flow

Principle of minimal sensitivity

Liao, Polonyi, Strickland, NPB 567 (2000) 493-514 eg. Canet, Delamotte, Mouhanna, Vidal, PRD 67 (2003) 065004



approximated flow

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Principle of minimal sensitivity

Liao, Polonyi, Strickland, NPB 567 (2000) 493-514 eg. Canet, Delamotte, Mouhanna, Vidal, PRD 67 (2003) 065004



approximated flow

Optimisation: find $R_{k}^{(2)}$!

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Most rapid convergence at fixed points

Litim, PLB 486 (2000) 92-99

Functional optimisation: Integrability

JMP, AP 322 (2007) 2831 JMP, Scherer, Schmidt, Wetzel, AP 384 (2017) 165



full flow



Functional optimisation: Integrability

JMP, AP 322 (2007) 2831 JMP, Scherer, Schmidt, Wetzel, AP 384 (2017) 165

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Flows for correlation functions & chiral symmetry breaking



Dupuis et al, Phys.Rept. (2021) Fu, Commun.Theor.Phys. 74 (2022) 9, 097304

ab initio







 $\Phi = (A_{\mu}, c, \bar{c}, q, \bar{q})$

functional RG:

Dupuis et al, Phys.Rept. (2021) Fu, Commun.Theor.Phys. 74 (2022) 9, 097304

free energy at momentum



quarks & gluons

ab initio



quark quantum fluctuations (RG-scale k: $t = \ln k$)

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$$functional RG: \left(\partial_t + \int_x \dot{\Phi} \frac{\delta}{\delta \Phi}\right) \Gamma_k[\Phi] = \begin{bmatrix} \frac{1}{2} & \frac{\delta}{\delta \Phi} \\ \frac{1}{2} & \frac{\delta}{\delta \Phi} \end{bmatrix} = \begin{bmatrix} 1 & \frac{1}{2} & \frac{\delta}{\delta \Phi} \\ \frac{1}{2} & \frac{\delta}{\delta \Phi} \end{bmatrix} = \begin{bmatrix} 1 & \frac{1}{2} & \frac{\delta}{\delta \Phi} \\ \frac{1}{2} & \frac{\delta}{\delta \Phi} \end{bmatrix} = \begin{bmatrix} 1 & \frac{\delta}{\delta \Phi} \\ \frac{1}{2} & \frac{\delta}{\delta \Phi} \end{bmatrix} = \begin{bmatrix} 1 & \frac{\delta}{\delta \Phi} \\ \frac{1}{2} & \frac{\delta}{\delta \Phi} \end{bmatrix} = \begin{bmatrix} 1 & \frac{\delta}{\delta \Phi} \\ \frac{1}{2} & \frac{\delta}{\delta \Phi} \end{bmatrix} = \begin{bmatrix} 1 & \frac{\delta}{\delta \Phi} \\ \frac{1}{2} & \frac{\delta}{\delta \Phi} \end{bmatrix} = \begin{bmatrix} 1 & \frac{\delta}{\delta \Phi} \\ \frac{1}{2} & \frac{\delta}{\delta \Phi} \end{bmatrix} = 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free energy/ grand potential

$$\Phi = (A_{\mu}, c, \bar{c}, q, \bar{q})$$

Dupuis et al, Phys.Rept. (2021) Fu, Commun.Theor.Phys. 74 (2022) 9, 097304

ab initio

fRG approach with emergent composites/dynamical hadronisation











functional RG:

 $\partial_t \Gamma_k[\Phi]$ $=\frac{-}{2}$

free energy/ grand potential

Correlation functions

glue quantum fluctuations



quark quantum fluctuations

functional RG:



gluon propagator $\langle A_\mu A_
u
angle(p)$

Pure glue

Correlation functions

glue quantum fluctuations

> quark quantum fluctuations

 $\partial_t \Gamma_k[\Phi]$ functional RG: = $\overline{2}$ free energy/ grand potential gluon propagator Pure glue $\partial_t \dots = \frac{2}{2}$ $\partial_t \quad \text{mmOulline} = \frac{2}{2}$ HO HE \bigcirc $\hat{\mathbf{Q}}$ $\oplus \frac{4}{2}$ ∂_t Oum \otimes $\partial_{t} \int_{e^{-5}}^{2} \partial_{t} \int_{e^{-5}}^{2} \int_{e^{-5}}^{e^{-5}} \partial_{t} \int_{e^{-5}}^{e^{-5}} \partial_{t}$

glue quantum fluctuations



quark quantum fluctuations

Correlation functions



--- 1-loop exact

functional RG:



Correlation functions

quark propagator \mathcal{D} $\langle qq \rangle$

gluon propagator (p) $A_{\mu}A_{\nu}$

no hadronic composites

glue quantum fluctuations

> quark quantum fluctuations

functional RG:



gluon propagator (\mathcal{D}) $A_{\mu}A_{\nu}$

quark propagator $\langle q\bar{q}\rangle$ (\mathcal{D})

no hadronic composites

glue quantum fluctuations

> quark quantum fluctuations

Correlation functions

quark-gluon vertex $, p_2)$ $\langle q\bar{q}A$ Eight tranverse tensor structures

functional RG:



gluon propagator (p) $A_{\mu}A_{\nu}$

quark-gluon vertex quark propagator $\langle q\bar{q}A_{\mu}\rangle$ (p_1, p_2) $\langle q\bar{q}\rangle(p)$ Eight tranverse tensor structures

no hadronic composites



quark quantum fluctuations

Correlation functions

glue

quark—anti-quark scattering

 $\langle q\bar{q}q\bar{q}\rangle(p_1,p_2,p_3)$

functional RG:

quark propagator

 $\langle q \bar{q} \rangle(p)$



gluon propagator (p) $A_{\mu}A_{\nu}$



no hadronic composites



quark quantum fluctuations

Correlation functions

glue

functional RG:





2 tensor structures

$$\Gamma_{\text{mat}} = \int_{p} \bar{q}(-p) \left[Z_{q}(p) i \not p + M_{q}(p) \right] q(p) + \sum_{i=1}^{10} \int_{\mathbf{p}} \lambda_{\bar{q}^{2}q^{2}}^{(i)}(\mathbf{p}) \left(\bar{q}^{2} \mathcal{T}_{\bar{q}^{2}q^{2}}^{(i)} q^{2} \right) (\mathbf{p}) + \cdots$$

$$\phi$$
 =

glue quantum fluctuations

> quark quantum fluctuations



Two flavours: 10 momentum-independent tensor structures

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Chiral symmetry breaking in a nutshell







Chiral symmetry breaking in a nutshell

$$= 2\lambda - A(k, M_q) \lambda^2$$







Chiral symmetry breaking in a nutshell

$$= 2\lambda - A(k, M_q) \lambda^2 - B(k, M_q, M_{\rm gap}) \lambda \alpha_s$$

Chiral symmetry breaking in a nutshell







$$= \left[2 - B(k, M_q, M_{gap})\alpha_s\right]\lambda - A(k, M_q)\lambda^2 - C(k, M_q, M_q)\lambda^2$$



Chiral symmetry breaking in a nutshell







$$= \left[2 - B(k, M_q, M_{gap})\alpha_s\right]\lambda - A(k, M_q)\lambda^2 - C(k, M_q, M_q)\lambda^2$$



Getting dynamical: emergent hadrons & diquarks

Gies, Wetterich, PRD 65 (2002) 065001 PRD 69 (2004) 025001

JMP, AP 322 (2007) 2831-2915 Floerchinger, Wetterich, PLB 680 (2009) 371

Fu, JMP, Rennecke, PRD 101, (2020) 054032 Fukushima, JMP, Strodthoff, 2103.01129

functional RG:
$$\left(\partial_t + \int_x \dot{\Phi} \frac{\delta}{\delta \Phi}\right) \Gamma_k[\Phi] = \frac{1}{2}$$

'DynHad for mesons & diquarks is BSE-DSE for QCD in a 'unified' effective action approach'





Dynamical hadronisation



Implementation:

functional RG:
$$\left(\partial_t + \int_x \dot{\Phi} \frac{\delta}{\delta \Phi}\right) \Gamma_k[\Phi] = \frac{1}{2}$$



Implementation:

functional RG:
$$\left(\partial_t + \int_x \dot{\Phi} \frac{\delta}{\delta \Phi}\right) \Gamma_k[\Phi] = \frac{1}{2}$$

Consider path integral in the presence of sources for composite operators

$$Z[J_q, J_{\bar{q}}, J_{\mathcal{O}}] = \int dq d\bar{q}$$



JMP, AP 322 (2007) 2831-2915

 $J_e e^{-S[q,\bar{q}] + \int J_q q - \bar{q} J_{\bar{q}} + \int J_{\mathcal{O}} \mathcal{O}[q,\bar{q}]}$

Implementation:

functional RG:
$$\left(\partial_t + \int_x \dot{\Phi} \frac{\delta}{\delta \Phi}\right) \Gamma_k[\Phi] = \frac{1}{2}$$

Consider path integral in the presence of sources for composite operators

$$Z[J_q, J_{\bar{q}}, J_{\mathcal{O}}] = \int dq d\bar{q} \, e^{-S[q,\bar{q}] + \int J_q q - \bar{q} J_{\bar{q}} + \int J_{\mathcal{O}} \, \mathcal{O}[q,\bar{q}]}$$

Choose scale-dependent $\mathcal{O}_k[q, \overline{q}]$ 'to optimise dynamics'!

$$\partial_t \Gamma_k[A_\mu, q, \bar{q}]$$



JMP, AP 322 (2007) 2831-2915

$$\partial_t \Gamma_k[\Phi] + \partial_t \mathcal{O}_k^{(i)}[\Phi] \frac{\delta \Gamma_k}{\delta \Phi_i}$$

$$\Phi = (A_{\mu}, q, \bar{q}, \langle \mathcal{O}^{(1)} \rangle, \dots)$$

Implementation:

$$\frac{\lambda_{\psi}}{2} \left[(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2 \right] = \left[ih\,\bar{\psi}(\tau\cdot\Phi)\psi + \frac{1}{2}m_{\phi}^2\Phi^2 \right]_{\mathrm{EoM}(\Phi)}$$



Consider path integral in the presence of sources for composite operators

$$Z[J_q, J_{\bar{q}}, J_{\mathcal{O}}] = \int dq d\bar{q}$$

Common choices

$$T^i = (1, \gamma_5 \vec{\sigma})$$

Scalar-pseudoscalar channel

2001 - : Braun, Flörchinger, Fu Gies, JMP, Rennecke, Wetterich, ...

Hubbard-Stratonovich



JMP, AP 322 (2007) 2831-2915

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Scalar-pseudoscalar channel



2001 - : Braun, Flörchinger, Fu Gies, JMP, Rennecke, Wetterich, ...

Hubbard-Stratonovich



JMP, AP 322 (2007) 2831-2915

$$=\gamma_0$$

Density channel (part of vector multiplet)


Implementation:

$$\frac{\lambda_{\psi}}{2} \left[(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2 \right] = \left[ih\,\bar{\psi}(\tau\cdot\Phi)\psi + \frac{1}{2}m_{\phi}^2\Phi^2 \right]_{\mathrm{EoM}(\Phi)}$$



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Hubbard-Stratonovich



JMP, AP 322 (2007) 2831-2915

$$=(\gamma_0\,,\,ec\gamma)$$

Density channel (part of vector multiplet)



Implementation:

$$\frac{\lambda_{\psi}}{2} \left[(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2 \right] = \left[ih\,\bar{\psi}(\tau\cdot\Phi)\psi + \frac{1}{2}m_{\phi}^2\Phi^2 \right]_{\mathrm{EoM}(\Phi)}$$



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JMP, AP 322 (2007) 2831-2915

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Diquark channels



Implementation:

$$\frac{\lambda_{\psi}}{2} \left[(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2 \right] = \left[ih\,\bar{\psi}(\tau\cdot\Phi)\psi + \frac{1}{2}m_{\phi}^2\Phi^2 \right]_{\text{EoM}(\Phi)}$$



Common choices

$$T^i = (1, \gamma_5 \vec{\sigma})$$

Scalar-pseudoscalar channel



Density channel (part of vector multiplet)

$$\mathbf{N_f}=\mathbf{2:}\mathbf{10}$$

Momentum-independent tensor structures

Hubbard-Stratonovich

$$\lambda_{\psi} = \frac{h^2}{m_{\phi}^2}$$
$$\Phi = (\sigma, \vec{\pi})$$
$$\tau \cdot \Phi = \sigma + i\gamma_5$$

 $T^i = (\gamma_0, \vec{\gamma})$

Diquark channels

Complete basis

$$\mathbf{N_f} = \mathbf{3}:$$
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Implementation:

$$\frac{\lambda_{\psi}}{2} \left[(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2 \right] = \left[ih\,\bar{\psi}(\tau\cdot\Phi)\psi + \frac{1}{2}m_{\phi}^2\Phi^2 \right]_{\text{EoM}(\Phi)}$$



Common choices

$$T^i = (1, \gamma_5 \vec{\sigma})$$

Scalar-pseudoscalar channel



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 $T^i = (\gamma_0, \vec{\gamma})$

Diquark channels

Complete basis

$$\mathbf{N_f}=\mathbf{3:26}$$

All tensor structures for $\, N_f = 2:256 \,$



Implementation:

$$\frac{\lambda_{\psi}}{2} \left[(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5 \vec{\tau}\psi)^2 \right] = \left[i h \bar{\psi}(\tau \cdot \Phi)\psi + \frac{1}{2}m_{\phi}^2 \Phi^2 \right]_{\text{EoM}(\Phi)}$$

Hubbard-Stratonovich



General dynamical hadronisation

hadronised Flow

$$\frac{\partial}{\partial t}\Big|_{\phi}\Gamma_{k}[\phi] = \frac{1}{2}G_{k,\phi}\dot{R}_{k,\phi} + R_{k}G_{k,\phi}\frac{\delta\dot{\phi}}{\delta\phi} - \frac{\delta\Gamma}{\delta\phi}\dot{\phi}$$

$$\phi = (A_{\mu}, C, \bar{C}, q, \bar{q}, \Phi, ..., n, \bar{n}, ...)$$

How to fix
$$\phi_k$$
&

$$\lambda_{\psi} = \frac{h^2}{m_{\phi}^2}$$
$$\Phi = (\sigma, \vec{\pi})$$
$$\tau \cdot \Phi = \sigma + i\gamma_5$$

baryons mesons



 $\dot{\Phi}_k \simeq \dot{A}_k \bar{\psi} \tau \psi + \dot{B}_k \Phi_k + \dot{C}_k$



Implementation:





Implementation:





Flow for four-fermion coupling $\hat{\lambda}_{\psi} = \lambda_{\psi} k^2$ with infrared scale k



+

 $\partial_t \Phi$ - terms

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Implementation:





Flow for four-fermion coupling $\hat{\lambda}_{\psi} = \lambda_{\psi} k^2$ with infrared scale k

Implementation:

Full bosonisation $\hat{\lambda}_{\psi}=0$





=0

! Reminder !

Full bosonisation $\hat{\lambda}_{\psi} = 0$ Really?





(i) Complete dynamical hadronisation of one tensor channel removes one momentum channel!

(ii) Residual four-quark vertex left!



Stability & decoupling





Cutoff scale of dynamical chiral symmetry breaking



Stability & decoupling



Stability & decoupling



Stability & decoupling



quarks & gluons

Stability & decoupling



Pions: Chiral perturbation theory

quarks & gluons

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Mesons & diquarks:



Mesons & diquarks:





Mesons & diquarks:







Schematical flow

Fukushima, JMP, Strodthoff, 2103.01129

Mesons & diquarks:







Schematical flow

Fukushima, JMP, Strodthoff, 2103.01129

Mesons & diquarks:







Schematical flow

baryons:

Dominant UV-process:





baryons:

Dominant UV-process:











baryons:

Baryon formation processes



Baryonisation



baryons:

three-quark scattering



Baryonisation

quark-diquark scattering



baryons:

three-quark scattering



Yukawa-flows with baryonisation





nucleon-nucleon — ω_{μ} scattering:



Baryonisation

quark-diquark scattering





'DynHad for mesons, diquarks & baryons is Faddeev-BSE-DSE for QCD in a 'unified' effective action approach'

$$\left(\partial_t + \partial_t \Phi_{i,k}[\Phi] \frac{\delta}{\delta \Phi_i}\right)$$



(II) Functional QCD and the QCD phase structure

QCD at finite temperature and density

- Benchmarks in the vacuum
- Correlation functions at finite temperature
- Polyakov loop from functional approaches

QCD phase structure

- Locating the QCD phase boundary and the critical end point
- Fluctuations of conserved charges: Ripples of the critical end point

The unreasonable effectiveness of low energy effective theories and how to use them

Dalian, Beijing, Darmstadt, Heidelberg, Gießen

Braun, Chen, Fu, Gao, Geissel, Huang, Lu, Ihssen, Pawlowski, Rennecke, Sattler, Schallmo, Stoll, Tan, Töpfel, Turnwald, Wessely, Wen, Wink, Yin, Zheng, Zorbach



Functional flows for QCD











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Example: 4-quark scattering vertex









Example: 4-quark scattering vertex







The unreasonable effectiveness of low energy effective theories






fQCD: workflow

European Research Council

Established by the European Commission

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QCD at finite temperature and density

Benchmarks in the vacuum

Current set of correlation functions



Extension, work in progress:

Fu, Huang, Ihssen, JMP, Rennecke, Sattler, Tan

Cyrol, Mitter, JMP, Strodthoff, PRD 97 (2018) 054006, PRD 97 (2018) 054015

Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 94 (2016) 054005

Mitter, JMP, Strodthoff, PRD 91 (2015) 054035 67

Current set of correlation functions



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Fu, Huang, Ihssen, JMP, Rennecke, Sattler, Tan

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Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 94 (2016) 054005

Mitter, JMP, Strodthoff, PRD 91 (2015) 054035 67

Euclidean propagators





Two-flavour QCD





Euclidean propagators





simple correlations

Two-flavour QCD





Vertices



Aiming at apparent convergence



Vertices



Aiming at apparent convergence



Quark-gluon vertex

$$\left[\Gamma_{\bar{q}qA}^{(3)}\right]_{\mu}^{a}(p,q) = 1_{2\times 2}^{\text{flav}} T^{a} \sum_{i=1}^{8} \lambda_{i}(p,q) \left[\mathcal{T}_{\bar{q}qA}^{(i)}\right]_{\mu}(p,q)\right]$$

$$\begin{split} \bar{q} \not{D} q : & \left[\mathcal{T}_{\bar{q}qA}^{(1)} \right]_{\mu} (p,q) = -i \gamma_{\mu} \\ \bar{q} \not{D}^{3} q : & \left[\mathcal{T}_{\bar{q}qA}^{(5)} \right]_{\mu} (p,q) = i \left(\not{p} + \not{q} \right) (p-q)_{\mu} \\ & \left[\mathcal{T}_{\bar{q}qA}^{(6)} \right]_{\mu} (p,q) = i \left(\not{p} - \not{q} \right) (p-q)_{\mu} \\ & \left[\mathcal{T}_{\bar{q}qA}^{(7)} \right]_{\mu} (p,q) = \frac{i}{2} [\not{p}, \not{q}] \gamma_{\mu} \end{split}$$

Aiming at apparent convergence

covariant expansion scheme

$$\bar{q} \not{D}^2 q : \left[\mathcal{T}_{\bar{q}qA}^{(2)} \right]_{\mu} (p,q) = (p-q)_{\mu} \mathbf{1}_{4 \times 4}$$
$$\left[\mathcal{T}_{\bar{q}qA}^{(3)} \right]_{\mu} (p,q) = (\not{p} - \not{q}) \gamma_{\mu}$$
$$\left[\mathcal{T}_{\bar{q}qA}^{(4)} \right]_{\mu} (p,q) = (\not{p} + \not{q}) \gamma_{\mu}$$

quenched: Mitter, JMP, Strodthoff, PRD 91 (2015) 054035 70 Cyrol, Mitter, JMP, Strodthoff, PRD 97 (2018) 054006

Quark-gluon vertex



Quark-gluon vertex





p,q in MeV

 $p^2 q^2$

 $\theta = -$

Quark-gluon vertex



up-to-date 1st principles works:

FunMethods: Williams, EPJ A51 (2015) 57

Sanchis-Alepuz, Williams, PLB 749 (2015) 592 Williams, Fischer, Heupel, PRD 93 (2016) 034026

Aguilar, Binosi, Ibanez, Papavassiliou, PRD 89 (2014) 065027 Binosi, Chang, Papavassiliou, Qin, Roberts, PRD 95 (2017) 031501 Aguilar, Cardona, Ferreira, Papavassiliou, arXiv:1610.06158

Mitter, JMP, Strodthoff, PRD 91 (2015) 054035

Pelaez, Tissier, Wschebor, PRD 92 (2015) 045012

Eichmann, Sanchis-Alepuz, Williams, Alkofer, Fischer, PPNP 91 (2016) 1

lattice: Oliveira, Kizilersü, Silva, Skullerud, Sternbeck, Williams, APP Suppl. 9 (2016) 363



More generally: X-assisted Y

X=fRG, DSE, nPI, lattice, exp. data Y=fRG, DSE, nPI

Example: use

(a) 2+1 fRG-assisted gluon

(b) optional: two-flavour fRG quark-gluon vertex

More generally: X-assisted Y

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Example: use

(a) 2+1 fRG-assisted gluon

(b) optional: two-flavour fRG quark-gluon vertex



Further example, e.g. lattice-assisted 2+1 flavour DSE Aguilar et al, EPC 80 (2020) 2, 154

More generally: X-assisted Y

X=fRG, DSE, nPI, lattice, exp. data Y=fRG, DSE, nPI



Example: use

(a) 2+1 fRG-assisted gluon

(b) optional: two-flavour fRG quark-gluon vertex



in 2+1 flavour DSE quark gap eq.



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Example: use

(a) 2+1 fRG-assisted gluon

(b) optional: two-flavour fRG quark-gluon vertex











Example: use

(a) 2+1 fRG-assisted gluon

(b) optional: two-flavour fRG quark-gluon vertex







Correlation functions at finite temperature

YM-theory: gluonic correlation functions



Aiming at apparent convergence



YM-theory: gluonic correlation functions



Aiming at apparent convergence



Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 97 (2018) 5, 054015



Euclidean gluon propagator at finite T

chromo-magnetic propagator



Fister, JMP, arXiv:1112.5440

Lattice: Maas, JMP, Smekal, Spielmann, PRD 85 (2012) 034037

CF model: Reinosa, Serreau, Tissier, Tresmontant, PRD 95 (2017) 045014

Aiming at apparent convergence





Euclidean gluon propagator at finite T



Debye mass (chromo-electric)

chromo-electric propagator

Lattice: Silva, Oliveira, Bicudo, Cardoso, PRD89 (2014) 7, 074503

Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 97 (2018) 5, 054015



Euclidean gluon propagator at finite T



Debye mass (chromo-electric)

 $\langle A_0 \rangle \neq 0$

chromo-electric propagator

Lattice: Silva, Oliveira, Bicudo, Cardoso, PRD89 (2014) 7, 074503

Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 97 (2018) 5, 054015



Polyakov loop from functional approaches

$$\left(L[A_0] = \frac{1}{\mathbf{N}_{\mathbf{c}}} \mathbf{tr} \, \mathcal{P} \mathbf{e}^{\mathbf{i} \, \mathbf{g}} \int_{\mathbf{0}}^{\beta} \mathbf{A}_{\mathbf{0}}(\mathbf{x})\right)$$

FRG: Braun, Gies, JMP, PLB 684 (2010) 262 FRG, DSE, 2PI: Fister, JMP, PRD 88 (2013) 045010



$$\left(L[A_0] = \frac{1}{\mathbf{N}_{\mathbf{c}}} \operatorname{tr} \mathcal{P} \mathbf{e}^{\mathbf{i} \mathbf{g}} \int_{\mathbf{0}}^{\beta} \mathbf{A}_{\mathbf{0}}(\mathbf{x})\right)$$



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FRG: Braun, Gies, JMP, PLB 684 (2010) 262 FRG, DSE, 2PI: Fister, JMP, PRD 88 (2013) 045010



$T_c/\sqrt{\sigma} = 0.658 \pm 0.023$

lattice : $T_c/\sqrt{\sigma} = 0.646$

$$\left(L[A_0] = \frac{1}{\mathbf{N}_{\mathbf{c}}} \operatorname{tr} \mathcal{P} \mathbf{e}^{\mathbf{i} \mathbf{g}} \int_{\mathbf{0}}^{\beta} \mathbf{A}_{\mathbf{0}}(\mathbf{x})\right)$$



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Flow equation for the Polyakov loop expectation value



Herbst, Luecker, JMP, arXiv:1510.03830

quation for composite operators

JMP, AP 322 (2007) 2831

Igarashi, Itoh, Sonoda, PTP Suppl. 181 (2010) 1

Pagani, PRD 94 (2016) 045001





Flow equation for the Polyakov loop expectation value



Parameterisation

$$\langle L[A_0]
angle = Z_L[ar{A},\phi]\cdot L[A_0]$$
 with $\phi=(a_\mu,c,ar{c})$

Herbst, Luecker, JMP, arXiv:1510.03830

quation for composite operators

JMP, AP 322 (2007) 2831

Igarashi, Itoh, Sonoda, PTP Suppl. 181 (2010) 1

Pagani, PRD 94 (2016) 045001




Confinement

Flow equation for the Polyakov loop expectation value



Parameterisation

$$\langle L[A_0]
angle = Z_L[ar{A},\phi]\cdot L[A_0]$$
 with $\phi=(a_\mu,c,ar{c})$

Flow for Polyakov loop wave function

$$\partial_t Z_L[\bar{A}, \phi] = \operatorname{Flow}_{Z_L}[\bar{A}; Z_L, G_A, G_c, L[A_0]]$$

Herbst, Luecker, JMP, arXiv:1510.03830

quation for composite operators

JMP, AP 322 (2007) 2831

Igarashi, Itoh, Sonoda, PTP Suppl. 181 (2010) 1

Pagani, PRD 94 (2016) 045001





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Pagani, PRD 94 (2016) 045001







QCD phase structure

Locating the QCD phase boundary and the critical end point

Three remarks on Functional Approaches for QCD

• off-shell representation of thermodynamic observables

e.g. $\operatorname{Tr} \langle q(x) \bar{q}(x) \rangle$





pressure, trace anomaly, fluctuations, volume flucs., ...

'... and now for something completely different ...'

e.g. hadron resonances



Three remarks on Functional Approaches for QCD

• off-shell representation of thermodynamic observables

e.g. $\operatorname{Tr} \langle q(x) \bar{q}(x) \rangle$





pressure, trace anomaly, fluctuations, volume flucs., ...

gauge fixing = parameterisation

Consequences

I: simple correlations

'... and now for something completely different ...'

e.g. hadron resonances

$\langle q(x_1)\cdots \bar{q}(x_{2n})A_{\mu}(y_1)\cdots A_{\mu}(y_m)h(z_1)\cdots h(z_l)\rangle$

II: Difficult access to some observables

'No free lunch theorem'



Three remarks on Functional Approaches for QCD

• off-shell representation of thermodynamic observables

e.g. $\operatorname{Tr} \langle q(x) \bar{q}(x) \rangle$





pressure, trace anomaly, fluctuations, volume flucs., ...

gauge fixing = parameterisation

 $\langle q(x_1)\cdots \bar{q}(x_{2n})A_{\mu}(y_1)\cdots A_{\mu}(y_m)h(z_1)\cdots h(z_l)\rangle$

Consequences

I: simple correlations

Your mean field is not my mean field'

 $\frac{\delta S_{\rm cl}[\phi]}{\delta \phi}\Big|_{\phi=\bar{\phi}} =$ = 0

'... and now for something completely different ...'

e.g. hadron resonances

II: Difficult access to some observables

'No free lunch theorem'



full quantum equation of motion



Correlation functions at finite density from functional QCD To QCD or not to QCD....a minimal point of view

$$\partial_t \left\langle q(x)\bar{q}(y) \right\rangle^{-1}(\mu_B) = \operatorname{Loop}\left[\left\langle q(x)\bar{q}(y) \right\rangle(\mu_B), \ \left\langle q(x)A_\mu(y)\bar{q}(z) \right\rangle(\mu_B), \cdots; \ \mu_B \right]$$



'... and now for something completely different ...'

• Self-consistent truncations to functional relations define analytic functions in $\mu_{ m B}$, eg:





Dalian, Beijing, Darmstadt, Heidelberg, Gießen

Braun, Chen, Fu, Gao, Geissel, Huang, Lu, Ihssen, Pawlowski, Rennecke, Sattler, Schallmo, Stoll, Tan, Töpfel, Turnwald, Wessely, Wen, Wink, Yin, Zheng, Zorbach







- **fRG:** Fu, JMP, Rennecke, PRD 101 (2020) 054032
- **DSE:** Gao, JMP, PLB 820 (2021) 136584
 - Gunkel, Fischer, PRD 104 (2021) 054022





Collect all possible information/structure for physics understanding & extrapolations



- **fRG:** Fu, JMP, Rennecke, PRD 101 (2020) 054032
- **DSE:** Gao, JMP, PLB 820 (2021) 136584
 - Gunkel, Fischer, PRD 104 (2021) 054022









- **fRG:** Fu, JMP, Rennecke, PRD 101 (2020) 054032
- **DSE:** Gao, JMP, PLB 820 (2021) 136584
 - Gunkel, Fischer, PRD 104 (2021) 054022









- fRG: Fu, JMP, Rennecke, PRD 101 (2020) 054032
- **DSE:** Gao, JMP, PLB 820 (2021) 136584
 - Gunkel, Fischer, PRD 104 (2021) 054022



Extrapolations for Pheno

Requires a discussion of the

explicit & implicit assumptions







- fRG: Fu, JMP, Rennecke, PRD 101 (2020) 054032
- **DSE:** Gao, JMP, PLB 820 (2021) 136584

Gunkel, Fischer, PRD 104 (2021) 054022



Extrapolations for Pheno

Requires a discussion of the

explicit & implicit assumptions

Lattice extrapolations

Iow energy effective theories: QM, NJL, PQM, PNJL, ..., Holography



Functional QCD: systematic error estimates & the LEGO[®] principle



Example: 4-quark scattering vertex







Ihssen, JMP, Sattler, Wink, in preparation







Braun, Leonhardt, Pospiech, PRD 101 (2020) 036004

Dominance of scalar-pseudoscalar fluctuations

Pions & sigma mode







Braun, Leonhardt, Pospiech, PRD 101 (2020) 036004

Full chiral dynamics

Fu, JMP, Rennecke, PRD 101 (2020) 054032 (fRG)

Dominant chiral dynamics Gunkel, Fischer, PRD 104 (2021) 054022 (DSE)

Full quark-gluon dynamics

Gao, JMP, PLB 820 (2021) 136584 (DSE)

Dominance of scalar-pseudoscalar fluctuations

Pions & sigma mode











Pisarski, Rennecke, PRL 127 (2021) 152302

see talk of Fabian Rennecke

T=114 MeV & μ_B =630 MeV



Fu, JMP, Pisarski, Rennecke, Wen, Yin, in prep

T=160 MeV & μ_B =0 MeV

Pion spectral functions









Pisarski, Rennecke, PRL 127 (2021) 152302

see talk of Fabian Rennecke

T=114 MeV & μ_B =630 MeV





Fu, JMP, Pisarski, Rennecke, Wen, Yin, in prep

Moat regime is not captured quantitatively

Pion spectral functions









Emergent diquarks

Braun, Leonhardt, Pospiech, PRD 101 (2020) 036004

Regime of quantitative reliability of current best truncation

Estimate







Emergent diquarks

Braun, Leonhardt, Pospiech, PRD 101 (2020) 036004

Regime of quantitative reliability of current best truncation

Estimate

Emergent diquarks are not captured by extrapolations







Braun, Leonhardt, Pospiech, PRD 101 (2020) 036004



Regime of quantitative reliability of current best truncation

Estimate



Eichmann, Fischer, Welzbacher, PRD 93 (2016) 034013













Xiaofeng Luo

Most functional computations (LEFT or QCD) have not been set-up for CEP-predictions!

Location of CP : Theoretical Prediction

Preliminary collection from Lattice, DSE, FRG and PNJL (2004-2020)

Large uncertainties for the estimation of CP location.

The 10th RHIC BES theory and experiment online seminar, Oct. 6th, 2020

RHIC-BES Seminar Oct. 6th 2020, Xiaofeng Luo

9

Disclaimer

Lack of predictive power for CEP-predictions is no quality measure!









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Lack of predictive power for CEP-predictions is no quality measure!









Xiaofeng Luo

Remove CEP-predictions

(i) 'old' CEPs: lattice, Functional QCD approaches, LEFTS (updated computations available)

(ii) LEFTs & Functional Results (qualitative approximations) that miss lattice benchmarks at $\mu_{\rm B}$ =0

(iii) LEFTs with CEPs at large density (missing quark-gluon back reaction)

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The 10th RHIC BES theory and experiment online seminar, Oct. 6th, 2020 Xiaofeng Luo

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Location of CP : Theoretical Prediction



RHIC-BES Seminar Oct. 6th 2020, Xiaofeng Luo

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Location of CP : Theoretical Prediction



RHIC-BES Seminar Oct. 6th 2020, Xiaofeng Luo







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Location of CP : Theoretical Prediction



RHIC-BES Seminar Oct. 6th 2020, Xiaofeng Luo



Scenario I





Scenario II



Scenario III

Scaling and/or new phases







Soft modes in hot QCD matter: Braun, Chen, Fu, Gao, Huang, Ihssen, JMP, Rennecke, Sattler, Tan, Wen, Yin, arXiv:2310.19853







Scenario II







Scenario II



see talk of Wei-jie Fu



Fu, Luo, JMP, Rennecke, Wen, Yin, arXiv:2308.15508







The unreasonable effectiveness of low energy effective theories

or the LEGO[®] principle at work

The LEGO[®] principle at work





The unreasonable effectiveness of low energy effective theories



99



On the unreasonable effectiveness of low energy effective theories



Sequential decoupling of gluon, quark, sigma, pion fluctuations





Braun, Fister, Haas, JMP, Rennecke, PRD 94 (2016) 034016

Based on:

Rennecke, PRD 92 (2015) 076012

On the unreasonable effectiveness of low energy effective theories







Based on:

Braun, Fister, Haas, JMP, Rennecke, PRD 94 (2016) 034016

Rennecke, PRD 92 (2015) 076012


On the unreasonable effectiveness of low energy effective theories







Based on:

Braun, Fister, Haas, JMP, Rennecke, PRD 94 (2016) 034016

Rennecke, PRD 92 (2015) 076012



On the unreasonable effectiveness of low energy effective theories







Chiral condensates

101



renormalised condensate

lattice: S. Borsanyi, Z. Fodor, C. Hoelbling, S. D. Katz, S. Krieg, C. Ratti, and K. K. Szabo, JHEP 09, 073 (2010)

DSE: quark condensates

See also

Fischer, Luecker, PLB 718 (2013) 1036 Fischer, Luecker, Welzbacher, PRD 90 (2014) 034022 Isserstedt, Buballa, Fischer, Gunkel, PRD 100 (2019) 074011

$$\Delta_{l,R}(T,\mu_B) \simeq \Delta_l(T,\mu_B) - \Delta_l(0,0)$$

$$\Delta_q(T,\mu_B) = \frac{T}{\mathcal{V}} m_q^0 \int_x \langle \bar{q}(x)q(x) \rangle$$

fRG: Fu, JMP, Rennecke, PRD 101 (2020) 054032 **DSE:** Gao, JMP, PLB 820 (2021) 136584



Chiral condensates



renormalised condensate

lattice: S. Borsanyi, Z. Fodor, C. Hoelbling, S. D. Katz, S. Krieg, C. Ratti, and K. K. Szabo, JHEP 09, 073 (2010)



reduced condensate

$$\Delta_{l,R}(T,\mu_B) \simeq \Delta_l(T,\mu_B) - \Delta_l(0,0)$$

$$\Delta_q(T,\mu_B) = \frac{T}{\mathcal{V}} m_q^0 \int_x \langle \bar{q}(x)q(x) \rangle$$

$$\Delta_{l,s}(T,\mu_B) = \frac{\Delta_l(T,\mu_B) - \left(\frac{m_l^0}{m_s^0}\right)^2 \Delta_s(T,\mu_B)}{\Delta_l(0,0) - \left(\frac{m_l^0}{m_s^0}\right)^2 \Delta_s(0,0)}$$

fRG: Fu, JMP, Rennecke, PRD 101 (2020) 054032 **DSE:** Gao, JMP, PLB 820 (2021) 136584

101

300



QCD-assisted low energy effective theory



Direct QCD input





QCD-assisted low energy effective theory



Low energy quantum, thermal & density fluctuations via fRG (QCD-assisted PQM model)

Direct QCD input



Fu, Luo, JMP, Rennecke, Wen, Yin, arXiv:2308.15508



QCD-assisted low energy effective theory





Direct QCD input



Low energy quantum, thermal & density fluctuations via fRG (QCD-assisted PQM model)

Benchmarks with lattice and fQCD at vanishing density and fQCD at finite density



Fu, Luo, JMP, Rennecke, Wen, Yin, arXiv:2308.15508



QCD-assisted heavy ion physics: compilation of functional QCD results

Thermodynamics & spectral properties

Sneak preview on the QCD moat









Fu, JMP, Rennecke, PRD 101 (2020) 054032

iminary

Energy

Sneak preview on the QCD moat



Pion correlation function





Sneak preview on the QCD moat



0

0



Pion correlation function



105

Fu, JMP, Pisarski, Rennecke, Wen, Shi Yin, in preparation



EoS with the minimal DSE



Speed of sound





Pressure over energy



Yi Lu, Gao, Liu, JMP, 2310.18383 (accepted with PRD)



Chiral dynamics & soft modes

To be (critical) or not (to be)





Braun, Chen, Fu, Gao, Huang, Ihssen, JMP, Rennecke, Sattler, Tan, Wen, Yin, 2310.19853

Chiral transition temperature





Braun, Fu, JMP, Rennecke, Rosenblüh, Yin, PRD 102 (2020) 056010 Gao, JMP, PRD 105 (2022) 094020



To be (critical) or not (to be)



Order parameter potential & scaling



(Critical) exponent:
$$rac{1}{\delta}=rac{1}{n-1}$$



Braun, Chen, Fu, Gao, Huang, Ihssen, JMP, Rennecke, Sattler, Tan, Wen, Yin, 2310.19853

Chiral transition temperature

Braun, Fu, JMP, Rennecke, Rosenblüh, Yin, PRD 102 (2020) 056010 Gao, JMP, PRD 105 (2022) 094020







$$\Delta_l(m_\pi) \propto m_\pi^{2/\delta} \left[1 + a_m m_\pi^{2\theta_H} + \cdots \right]$$
 Braun, C



QM: Chen, Wen, WF, PRD 104 (2021) 054009

Chen, Fu, Gao, Huang, Ihssen, JMP, Rennecke, Sattler, Tan, Wen, Yin, 2310.19853



$$\Delta_l(m_\pi) \propto m_\pi^{2/\delta} ig[1 + a_m m_\pi^{2 heta_H} + \cdots ig]$$
Braun, C



'chiral scaling' Trivial $\Delta_l^{1+\delta}$ scaling

QM: Chen, Wen, WF, PRD 104 (2021) 054009

Chen, Fu, Gao, Huang, Ihssen, JMP, Rennecke, Sattler, Tan, Wen, Yin, 2310.19853











'chiral scaling'

Trivial $\Delta_I^{1+\delta}$ scaling

fQCD collaboration, in preparation QM: Chen, Wen, WF, PRD 104 (2021) 054009

Chiral dynamics & quasi-massless modes

'Non-critical chiral scaling' Far away from the critical regime for $m_{\pi} \gtrsim 1 \,\mathrm{MeV}$ $\Delta_l(T,H) \approx \Delta_{l,\chi}(0) \left(c_0 + c_{\frac{1}{5}} H^{\frac{1}{5}} + c_{\frac{1}{3}} H^{\frac{1}{3}} + c_1 H \right)$ Δ_l^6 $V_{\chi}(\Delta_l) \propto$ Δ_l^2 Δ_l^4



Full order parameter potential

Measure: correlation length





QCD: fQCD collaboration, in preparation



Dynamics and the size of the critical regime

Transport with fRG spectral functions & effective potential



Blum, Jiang, Nahrgang, JMP, Rennecke, Wink, NPA 982

QM: Roth, Schweitzer, Rieke, von Smekal, PRD 105 (2022)



Showcases in linear sigma models

Dynamical universality







Fluctuations of conserved charges: Ripples of the critical end point or the LEGO[®] principle at work

Baryon number fluctuations in the phase structure







Baryon number fluctuations in the phase structure







Variations of the CEP



Fu, Luo, JMP, Rennecke, Wen, Yin, arXiv:2308.15508



Canonical corrections via subensemble acceptance method

fixing the subensemble volume

subensemble volume

system volume

 $V_1 = \alpha V$



Vovchenko, Savchuk, Poberezhnyuk, Gorenstein, Koch, PLB 811, 135868 (2020)



Canonical corrections via subensemble acceptance method

fixing the subensemble volume

subensemble volume

system volume

 $V_1 = \alpha V$



qualitative adjustment

$$\alpha(\bar{s}) = a\left(1 - \sqrt{\bar{s}}\right)\theta\left(1 - \sqrt{\bar{s}}\right)\theta\left$$

$$a = 0.33 \qquad \qquad \sqrt{\bar{s}} = \frac{1}{1}$$

Vovchenko, Savchuk, Poberezhnyuk, Gorenstein, Koch, PLB 811, 135868 (2020)

 $(-\bar{s})$

 $/s_{\rm NN}$ $1.9\,\mathrm{GeV}$



fixing the subensemble volume

subensemble volume

system volume

 $V_1 = \alpha V$



qualitative adjustment

$$\alpha(\bar{s}) = a\left(1 - \sqrt{\bar{s}}\right)\theta\left(1 - \bar{s}\right)$$

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Canonical corrections via subensemble acceptance method

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baryon & proton number fluctuations



Reconstructing the CEP











Reconstructing the CEP









Unfolding the high density regime with new phases & physics

Great opportunity for a combined high precision analysis of high density QCD (Exp. data + lattice QCD + functional QCD)

Reconstructing the CEP







Summary

(I) Functional Renormalisation Group for QCD

- The renormalisation group is a one-loop exact functional approach
 - consistent RG-scaling
 - systematic expansion schemes & error control
 - compatibility with other functional approaches
- Resonances via dynamical hadronisation
 - hadrons as exchange fields of quarks scattering vertices
 - BSE wave function quark-hadron vertex
 - Stable dynamical emergence of low energy effective theories
- Quark-gluon-meson correlation functions
 - Self-consistent results: all correlation functions computed are fed back
 - **Dynamical chiral symmetry breaking & confinement**
 - Quantitative agreement with lattice results



(II) Functional QCD and the QCD phase structure

- QCD at finite temperature and density
 - all available benchmarks in the vacuum passed
 - confinement-deconfinement phase transition
 - compatibility with other functional approaches
- Locating the QCD phase structure and the critical end point
 - Quantitative predictions for μ_B/T
 - Estimate for the location of the CEF
 - Diquark domination for $\mu_B/T\gtrsim 8$
- Fluctuations of conserved charges: Ripples of the critical end point
 - Quantitative agreement of the fluctuations of conserved charges with lattice results Qualitative accounting for canonical effects with the sub-ensemble method
 - QCD-assisted low energy effective theory with the phase structure of QCD Remarkable compatibility with the new STAR data (baryon vs proton fluctuations,

 - finite volume effects, ...)

$$\lesssim 4$$
, estimates for $\mu_B/T \lesssim 800 \text{MeV}$
P: $(\mu_B, T)_{\text{CEP}} \sim (600 - 650, 105 - 115) \text{ MeV}$





