



Electromagnetic moments of ground and excited states calculated in nearly spherical and well-deformed odd nuclei

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Danneaux, A. Nagpal, A.E. Stuchbery, and H. Wibowo

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Outline

1. Methodology
 - a) Polarization
 - b) Self-consistency
 - c) Symmetry restoration
2. Odd near neighbours of doubly magic nuclei
3. Excited quasiparticle states in odd-N open-shell isotopes from gadolinium to osmium
4. Conclusions



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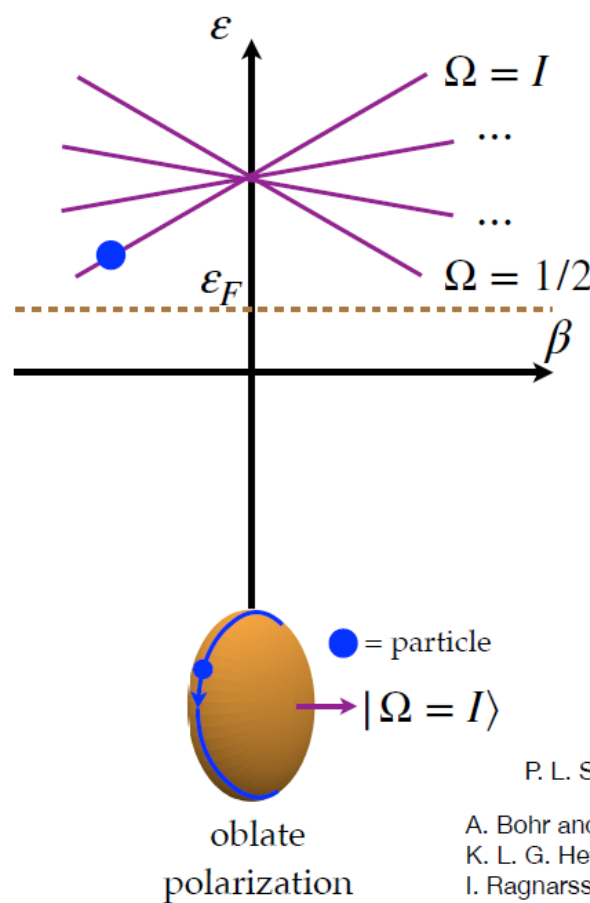
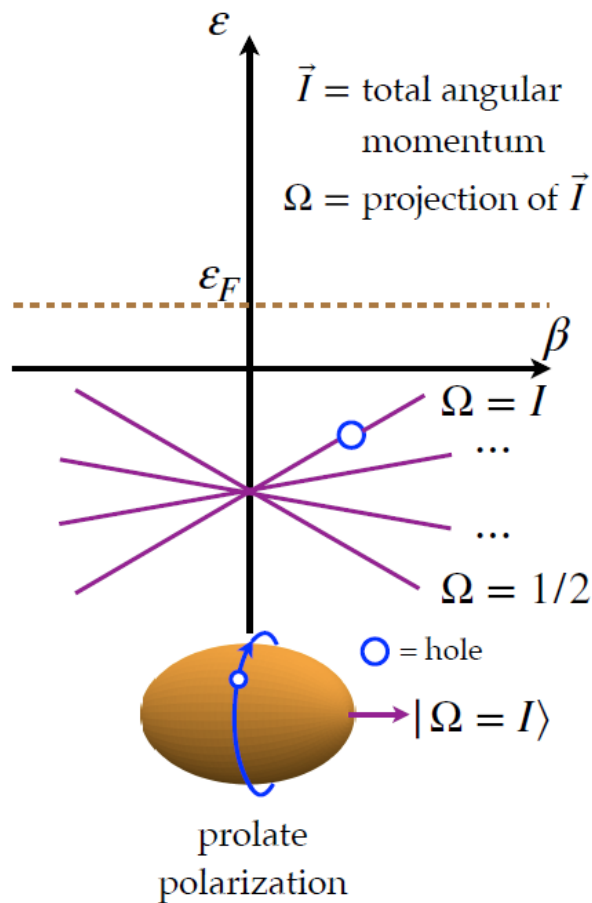


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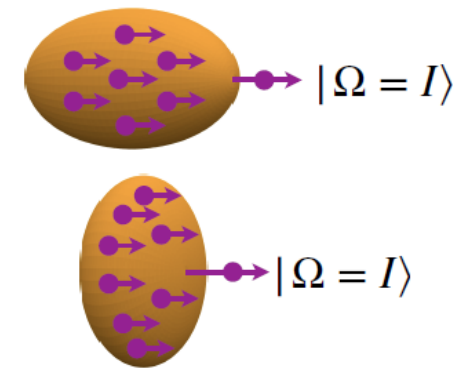


2

Shape and spin core polarizations



Spin polarization



Landau parameter g'_0 ($g'_0 = 1.7$)

$$g'_0 = N_0 (2C_1^S + 2C_1^T (3\pi^2 \rho_0 / 2)^{2/3})$$

$$\frac{1}{N_0} \approx 150 \frac{m}{m^*} \text{ MeV} \cdot \text{fm}^3$$

P. L. Sassarini et al., *J. Phys. G: Nucl. Part. Phys.* **49**, 11LT01 (2022)

A. Bohr and B. R. Mottelson, *Nuclear Structure* Vol. 1

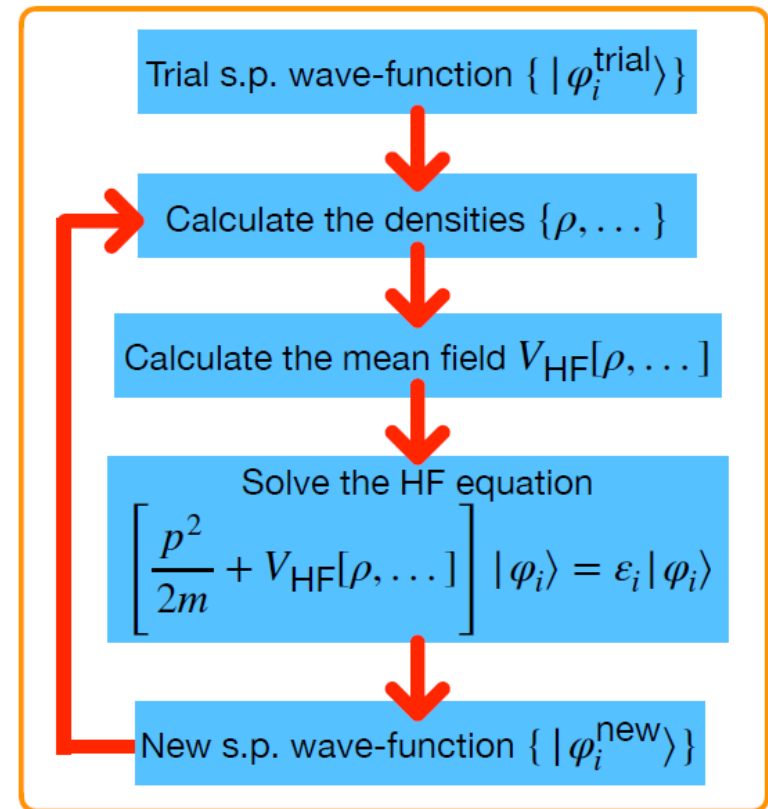
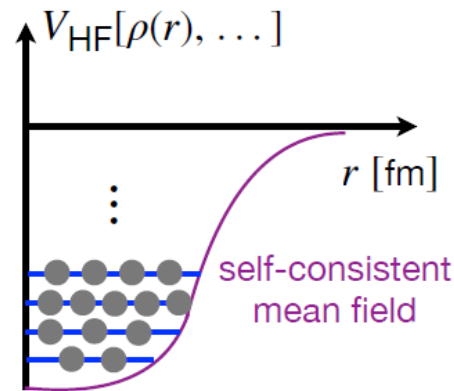
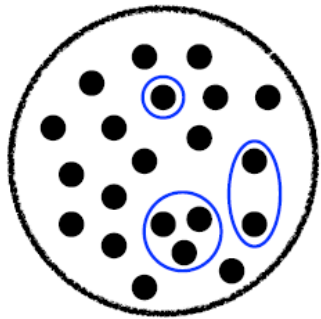
K. L. G. Heyde, *The Nuclear Shell Model*

I. Ragnarsson and S. G. Nilsson, *Shapes and Shells in Nuclear Structure*



1

Nuclear density functional theory



Energy density functional
 $\mathcal{E}[\rho(\mathbf{r}), \mathbf{s}(\mathbf{r}), \boldsymbol{\tau}(\mathbf{r}), \mathbf{T}(\mathbf{r}), \mathbf{j}(\mathbf{r}), \vec{\mathbf{J}}(\mathbf{r})]$

Coupling constants

T-even : $C_t^\rho, C_t^{\Delta\rho}, C_t^\tau, C_t^J, C_t^{\nabla J}$

T-odd : $C_t^s, C_t^{\Delta s}, C_t^T, C_t^j, C_t^{\nabla j}$

Parametrization: UNEDF1

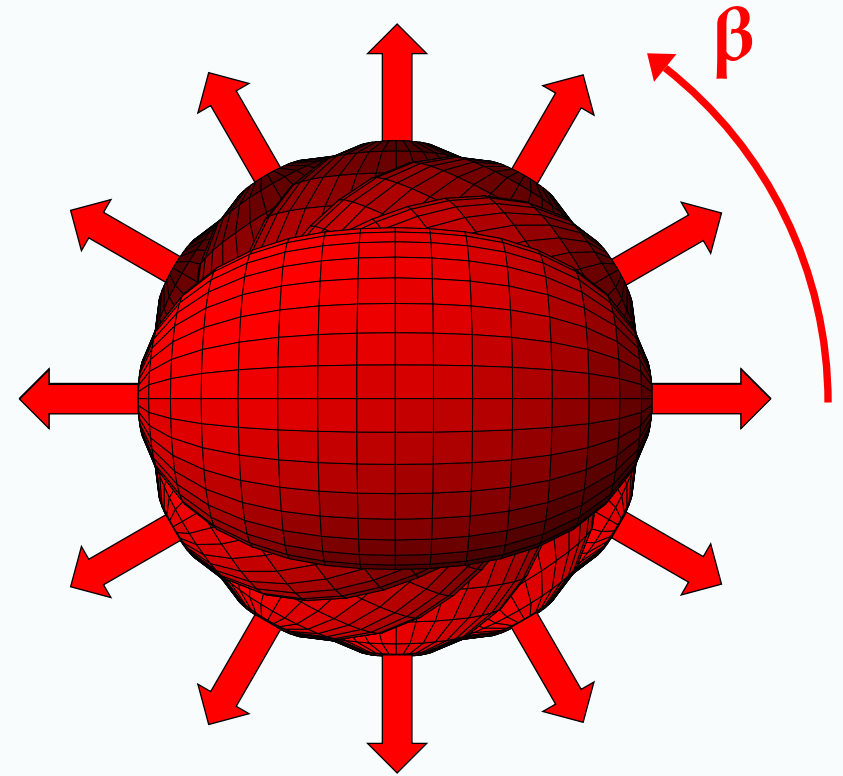
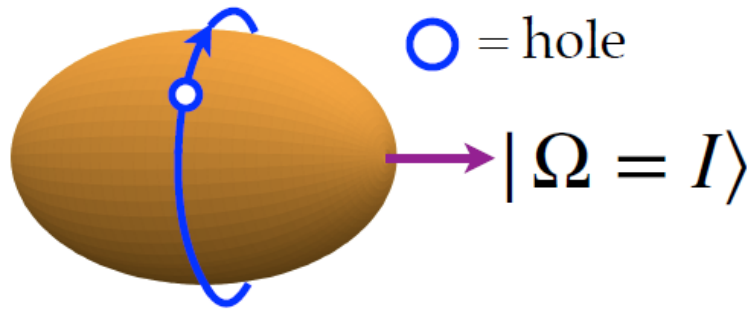
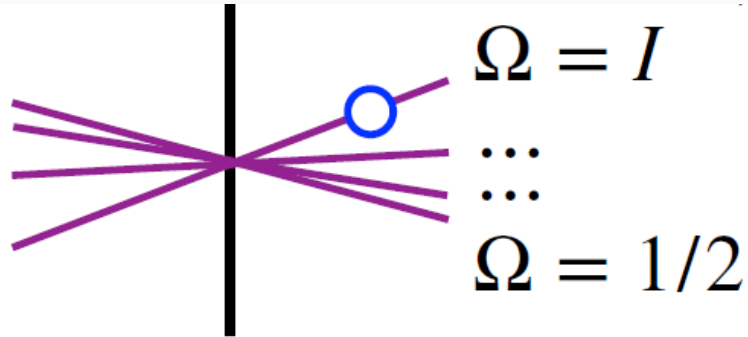
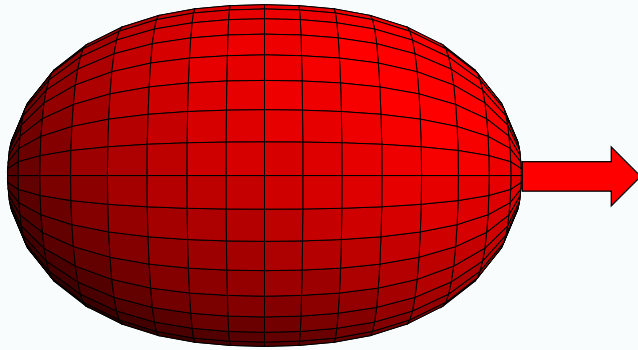
Hartree-Fock (HF) equation

M. Kortelainen et al., Phys. Rev. C 85, 024304 (2012)



Time-odd spin alignment & symmetry restoration

**“Intrinsic”
Symmetry broken**



**“Laboratory”
Symmetry restored**

$$|IM\rangle = \mathcal{N}_I \int_{\beta=0}^{\pi} d\beta d_{M\Omega}^I(\beta) |\Omega, \beta\rangle$$

J. A. Sheikh et al., J. Phys. G48, 123001 (2021)



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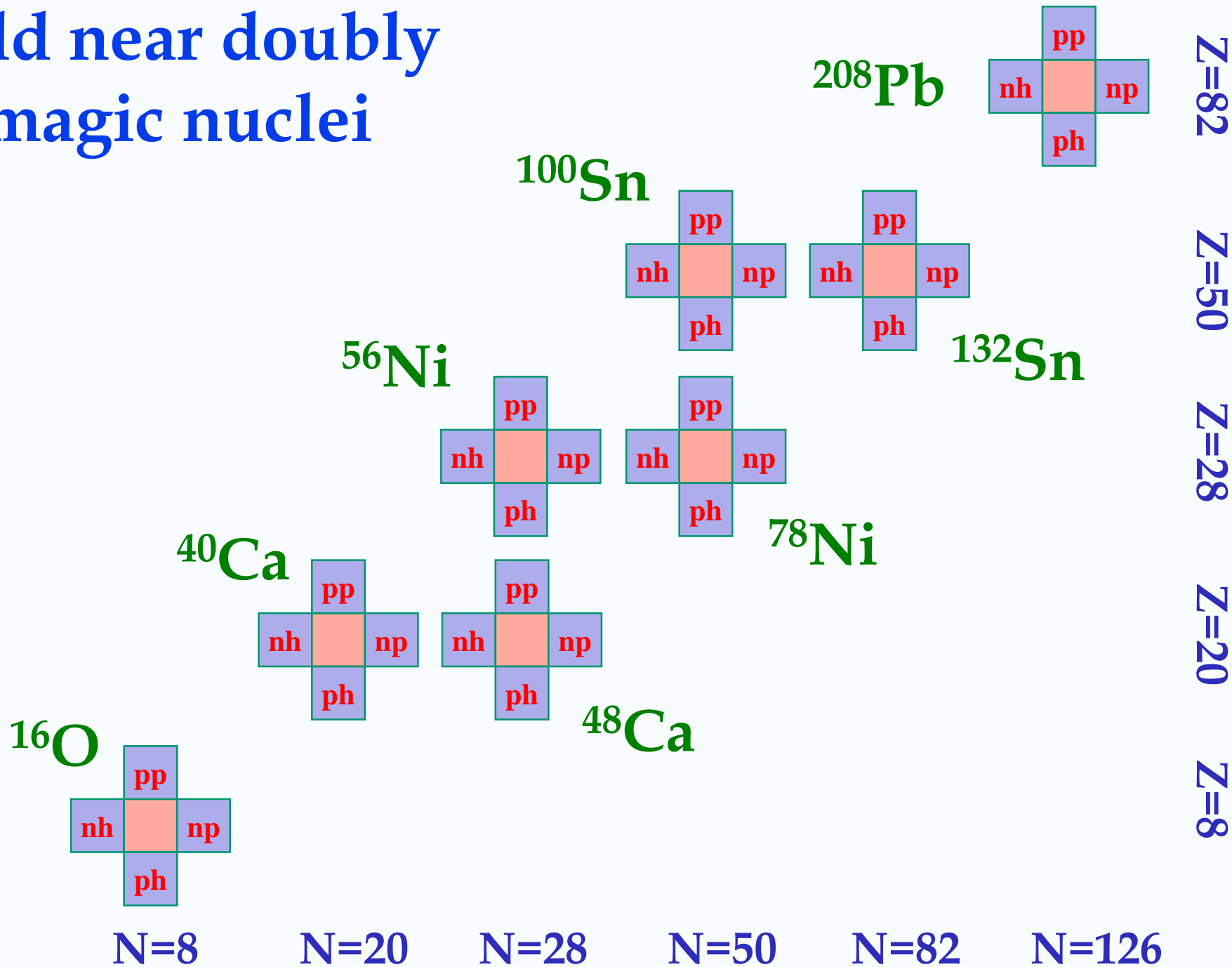
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Odd near doubly magic nuclei



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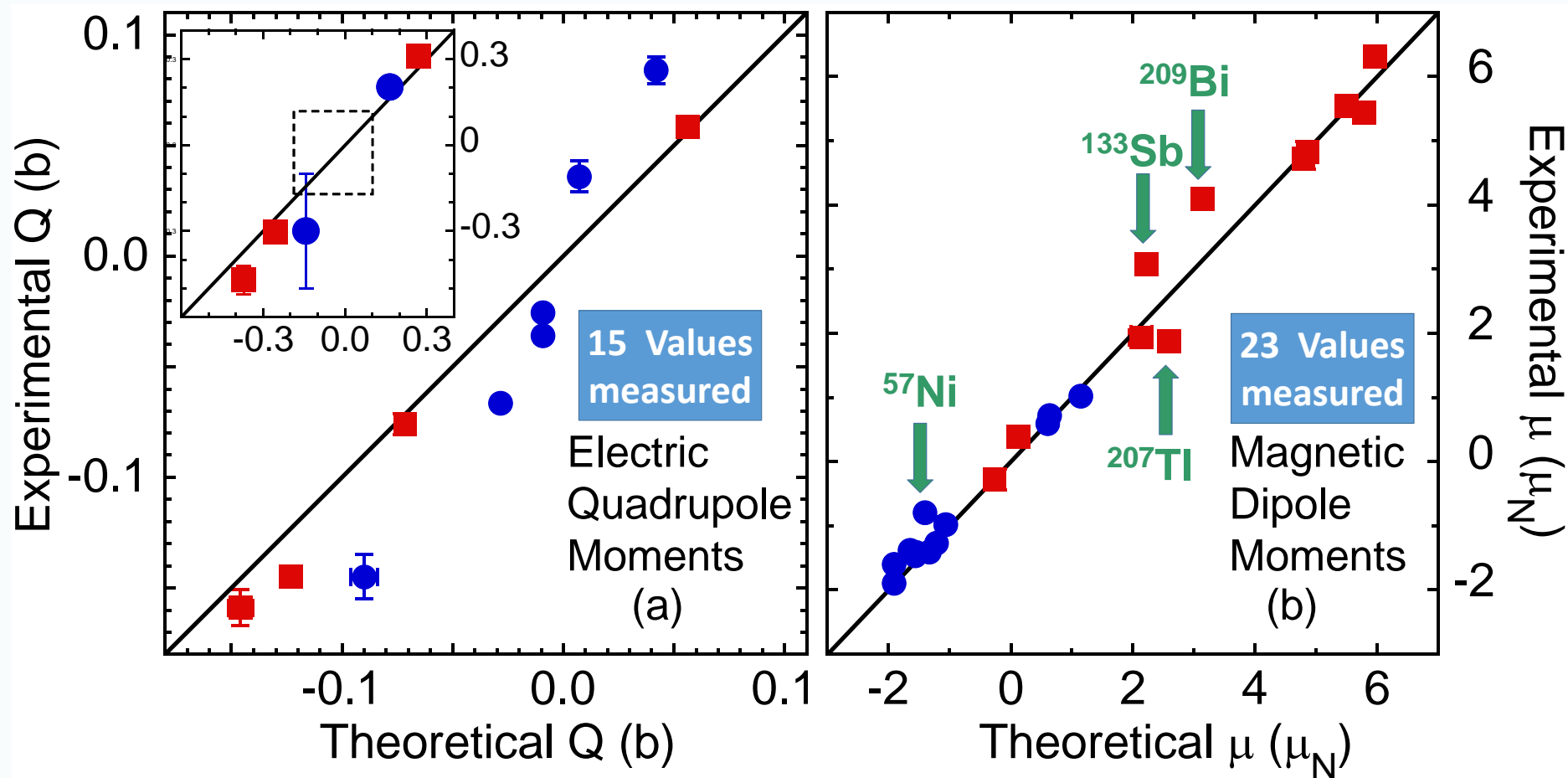


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Quadrupole & dipole moments

P. Sassarini *et al.*, J. Phys. G49 (2022) 11LT01



- Proton-odd (squares) & neutron-odd (circles) nuclei
- Average of UNEDF1, SLy4, SkO', D1S, N3LO functionals
- RMS deviations much smaller than the residuals



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Effective spin g-factor? Who ordered that?

P. Sassarini *et al.*, J. Phys. G49 (2022) 11LT01

$$g_s^{(i)} = g_p(g_n) = 5.59(-3.83) \times g^{\text{eff}} ???$$

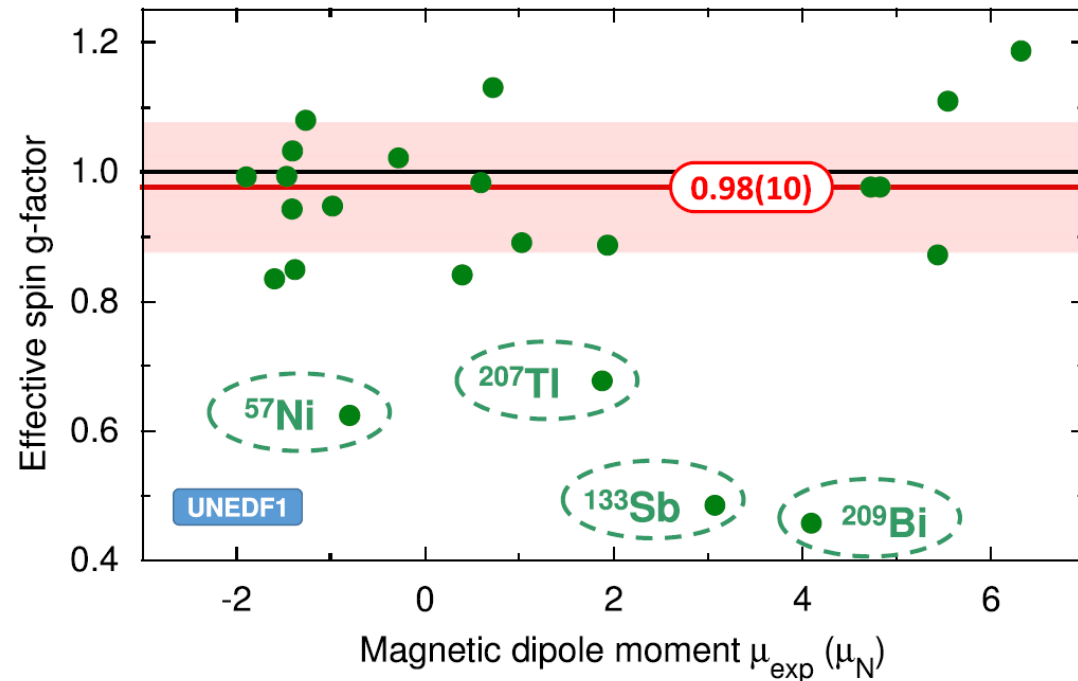
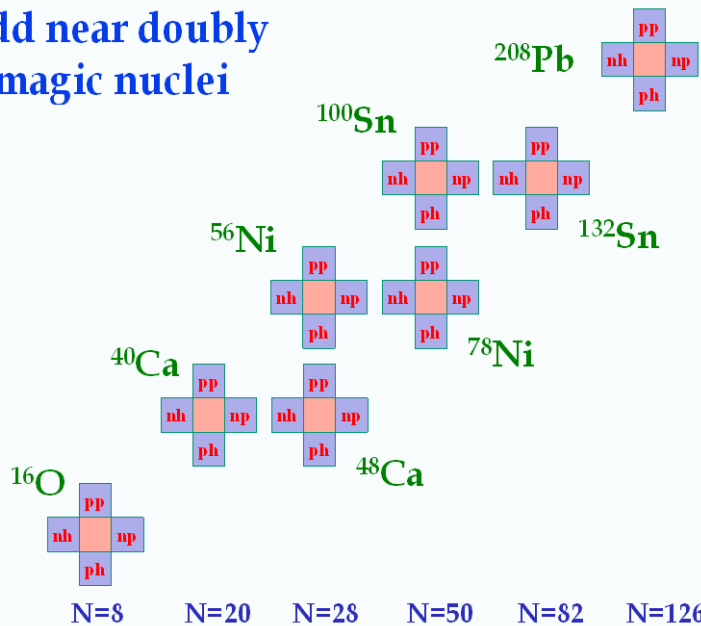
$$g_\ell^{(i)} = 1(0)$$

Landau parameter g'_0 ($g'_0 = 1.7$)

$$g'_0 = N_0 (2C_1^s + 2C_1^T (3\pi^2 \rho_0 / 2)^{2/3})$$

$$\frac{1}{N_0} \approx 150 \frac{m}{m^*} \text{ MeV} \cdot \text{fm}^3$$

Odd near doubly
magic nuclei



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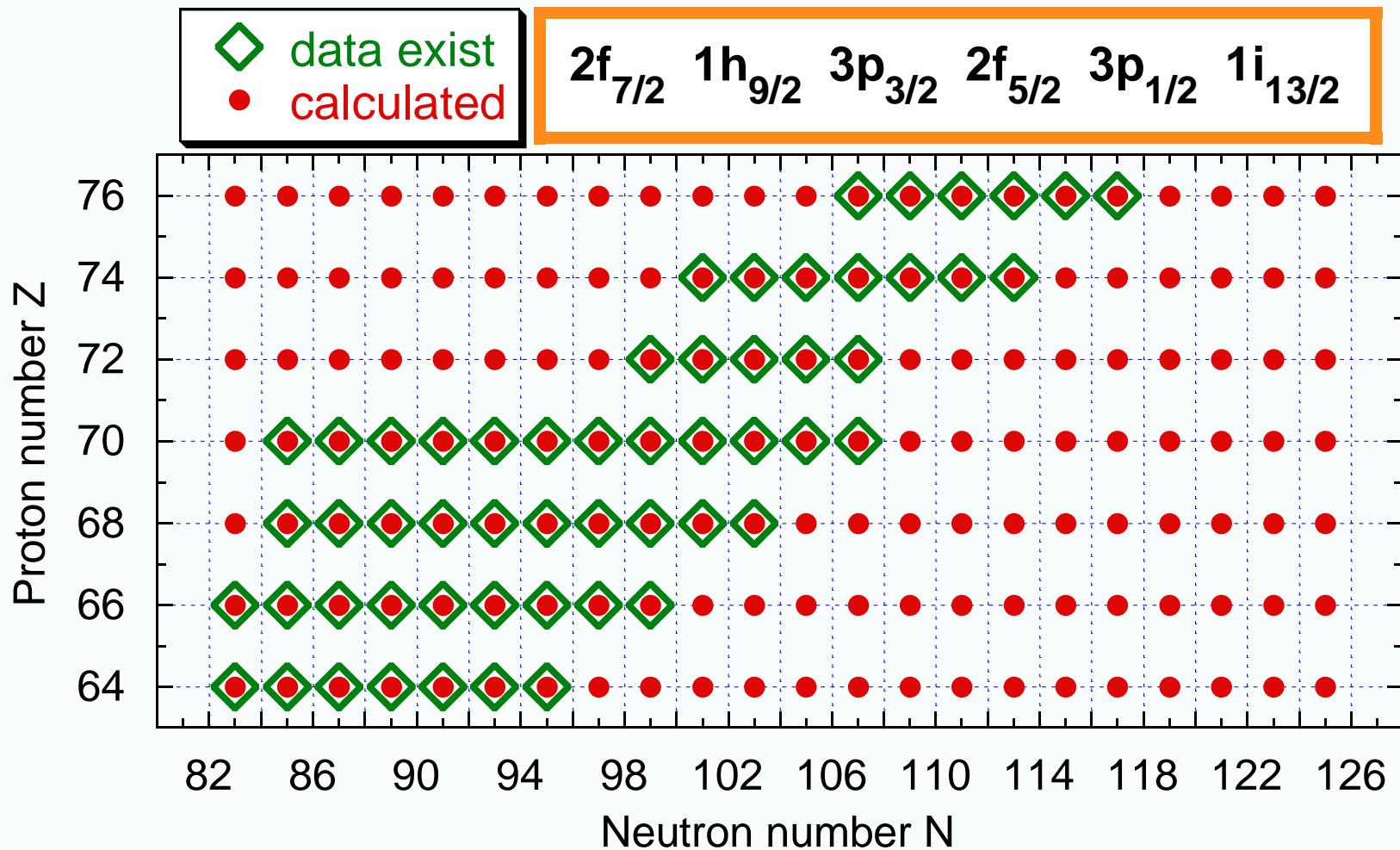
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The first systematic nuclear-DFT analysis of the electromagnetic moments in excited quasiparticle states



J. Dobaczewski *et al.*, to be published

Standard UNEDF1 nuclear functional used, no parameters (re)adjusted in this work

75 measured magnetic dipole moments (plus 3 rotational bands)

48 measured electric quadrupole moments (plus 3 rotational bands)



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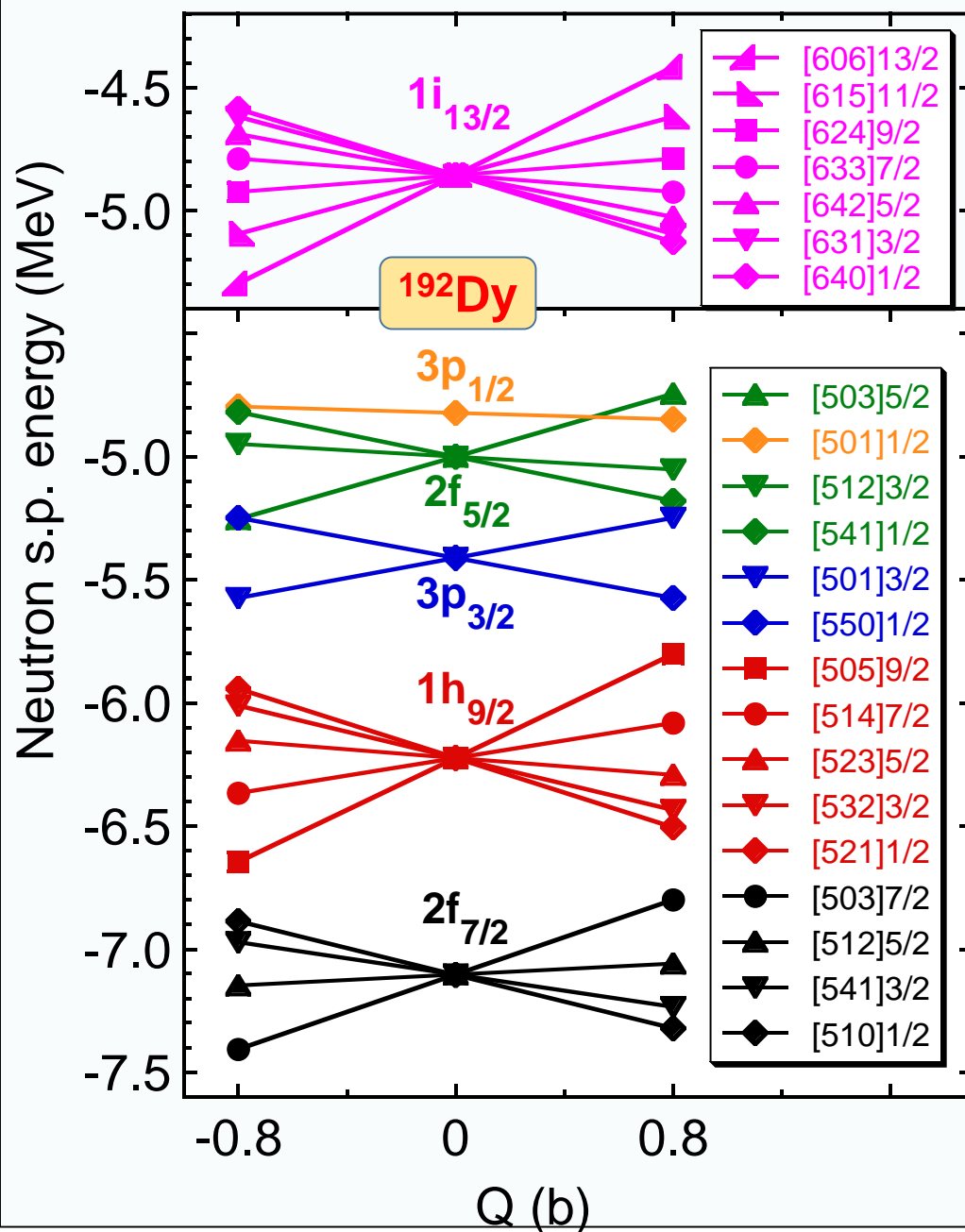
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How to calculate odd nuclei in nuclear DFT?



without pairing

A even, $p > A$, $h \leq A$

$$|\Psi\rangle_{\text{HF}}^{\text{even}} = a_A^+ \dots a_2^+ a_1^+ |0\rangle$$

$$|\Psi\rangle_{\text{HF}}^{\text{odd}} = \begin{cases} a_p^+ |\Psi\rangle_{\text{HF}}^{\text{even}} \\ a_h |\Psi\rangle_{\text{HF}}^{\text{even}} \end{cases}$$

with pairing

$$|\Psi\rangle_{\text{HFB}}^{\text{even}} = \prod_{\mu>0} (u_\mu + v_\mu a_\mu^+ a_\mu^+) |0\rangle$$

$$|\Psi\rangle_{\text{HFB}}^{\text{odd}} = \beta_\nu^+ |\Psi\rangle_{\text{HFB}}^{\text{even}}$$

$$= a_\nu^+ \prod_{\nu \neq \mu > 0} (u_\mu + v_\mu a_\mu^+ a_\mu^+) |0\rangle$$

tagging quasiparticle states

$$\max_\mu \{ \langle \varphi_\nu | \phi_\mu^{\text{upper}} \rangle, \langle \varphi_\nu | \phi_\mu^{\text{lower}} \rangle \}$$



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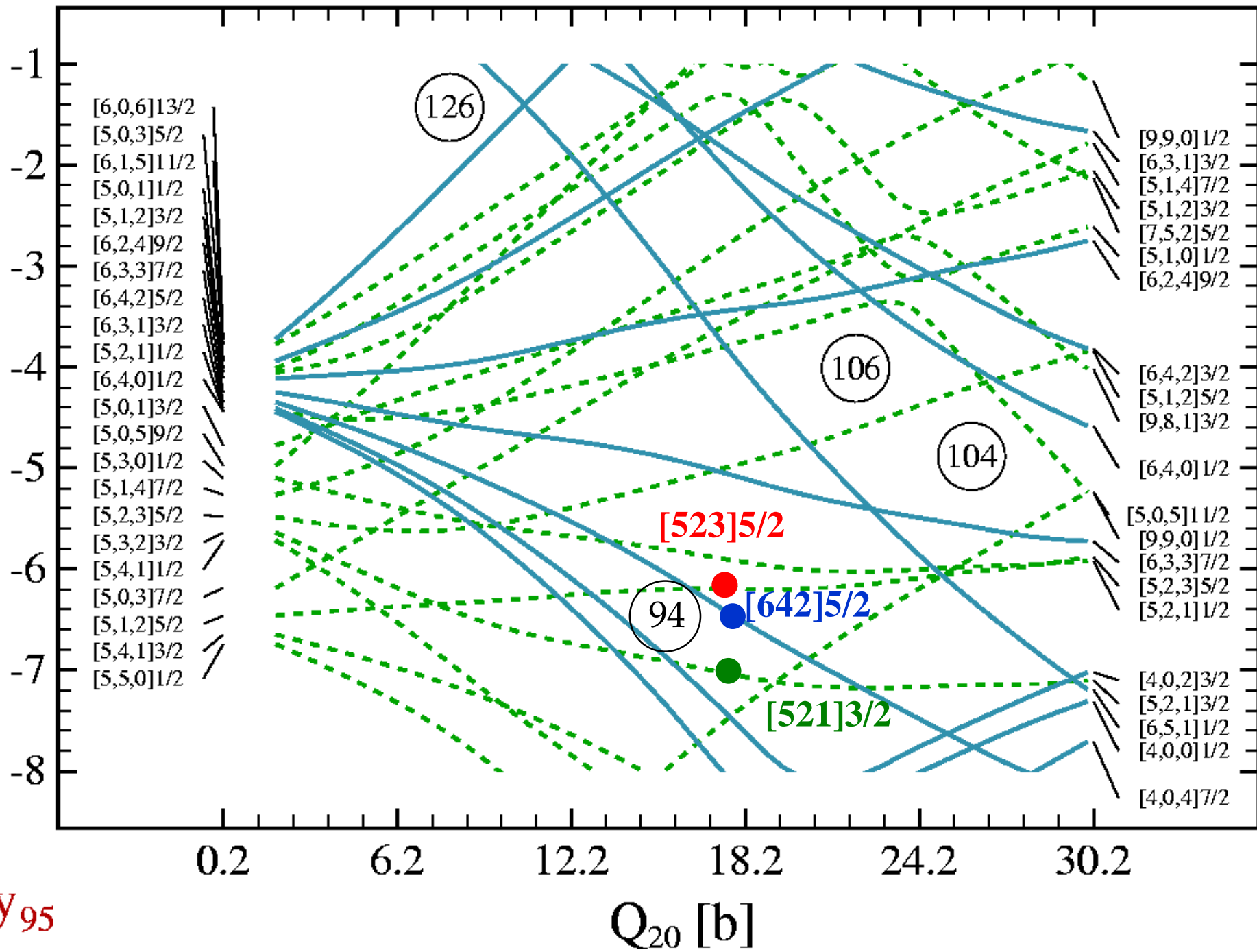
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Single-neutron Energies [MeV]



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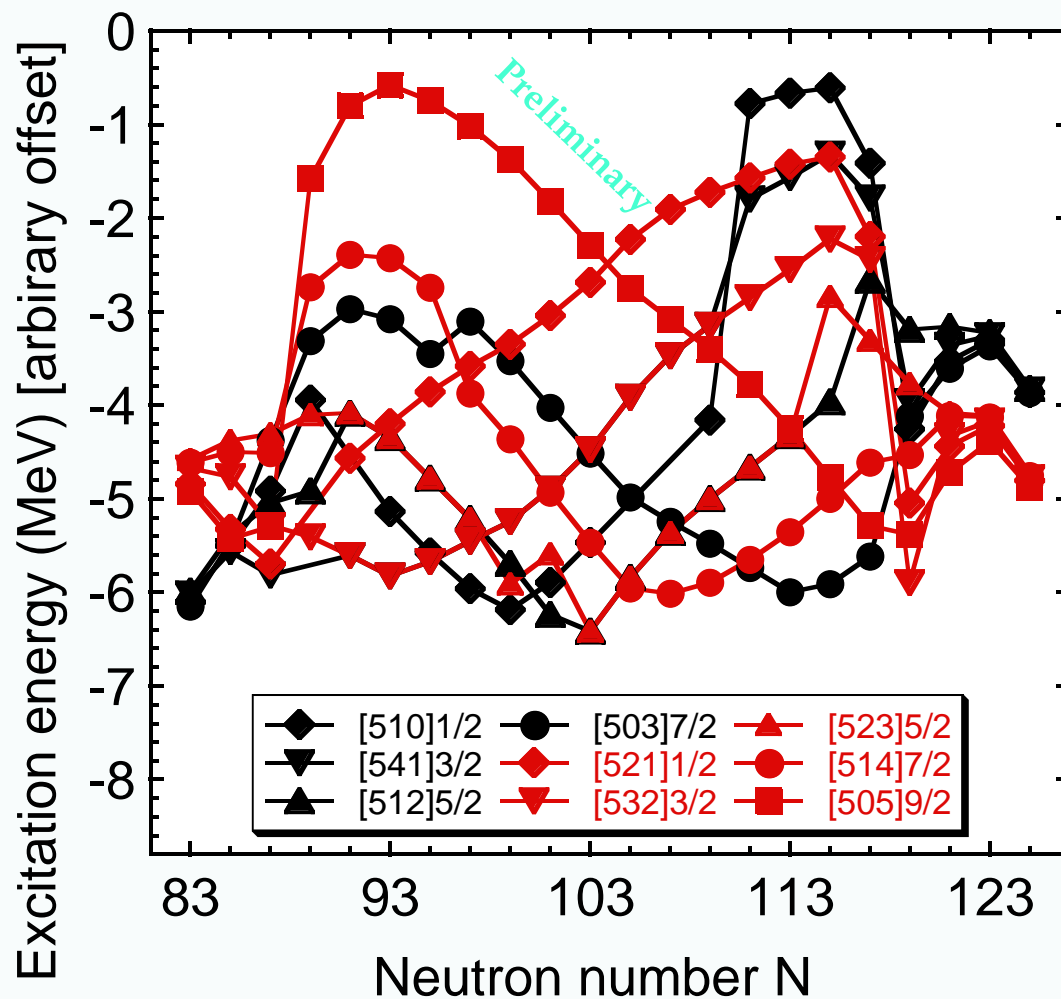
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Excitation energies of odd dysprosium isotopes



J. Dobaczewski *et al.*, to be published

66 band-head states were associated with the lowest calculated quasiparticle states of the given spin and parity. Among those were 27 calculated ground states and 21 were calculated low-lying excited states below 300 keV.



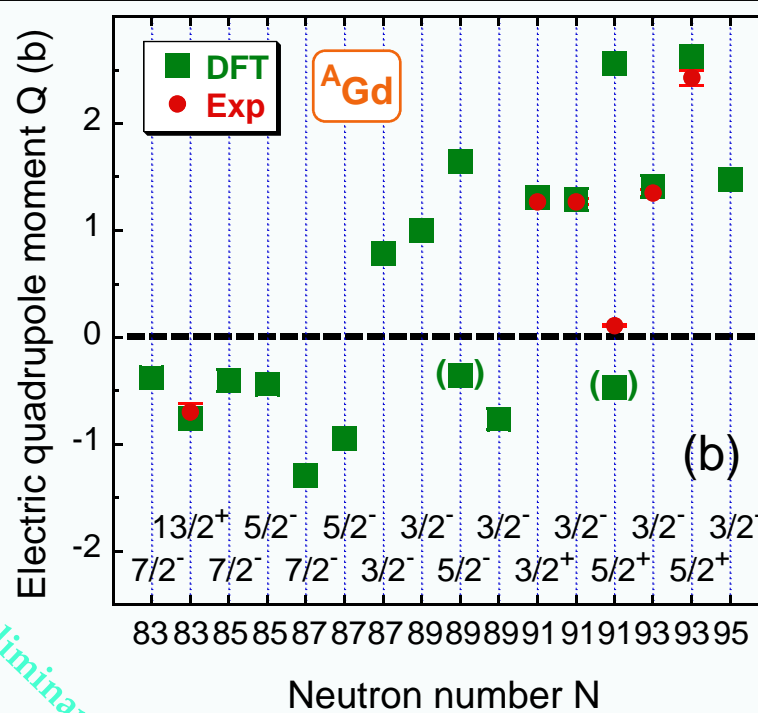
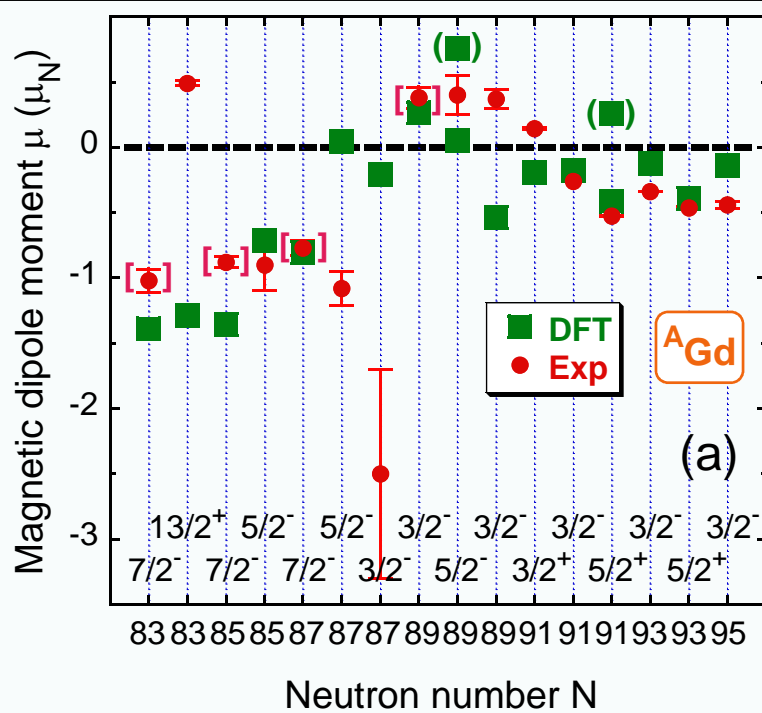
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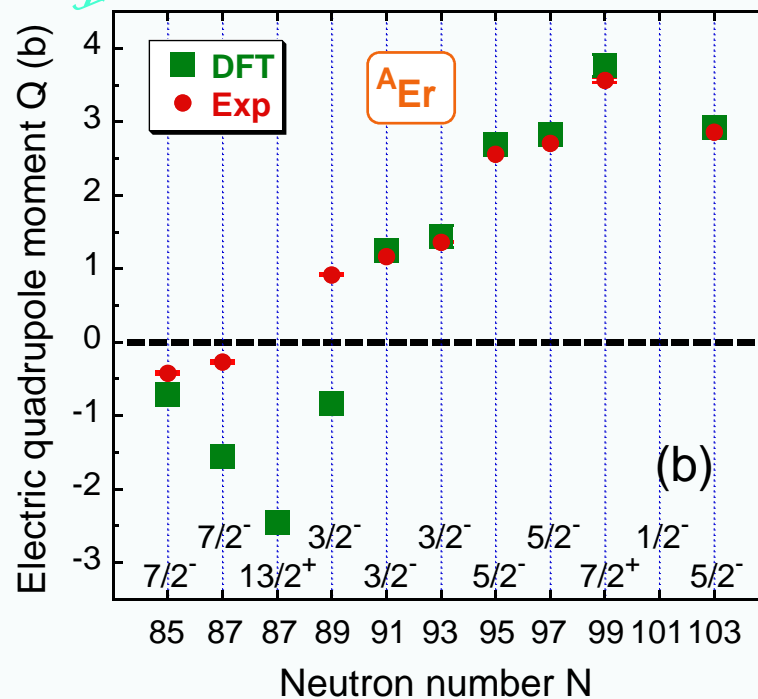
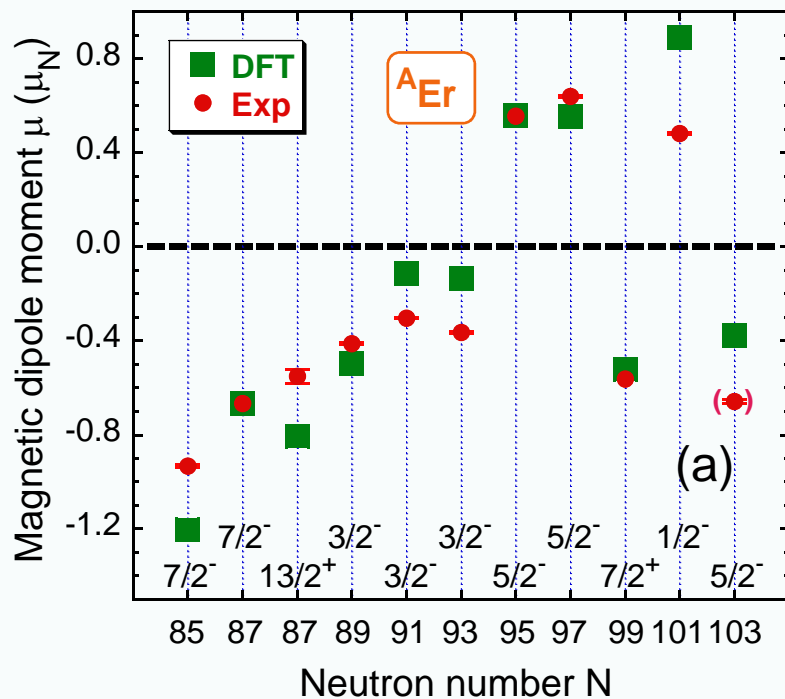


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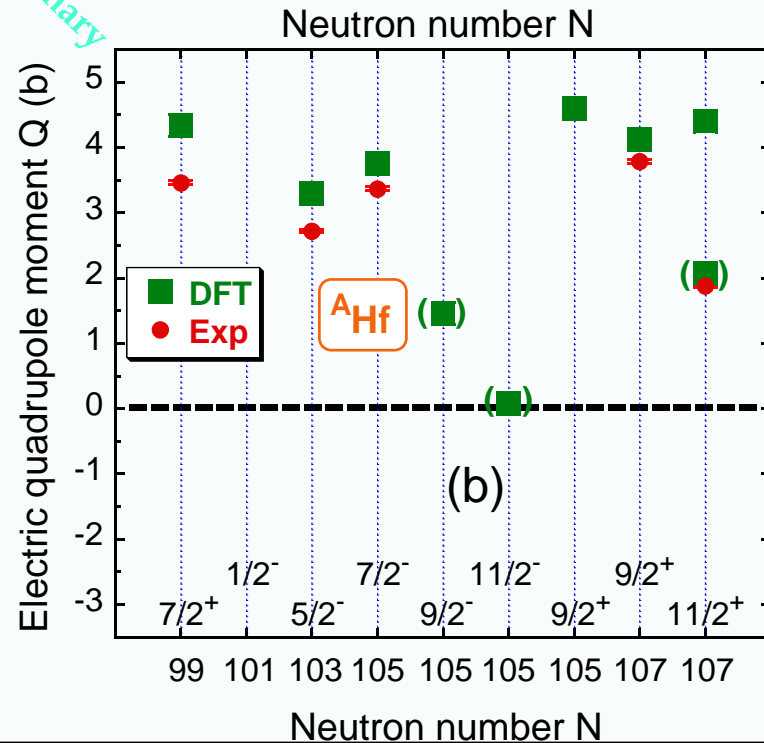
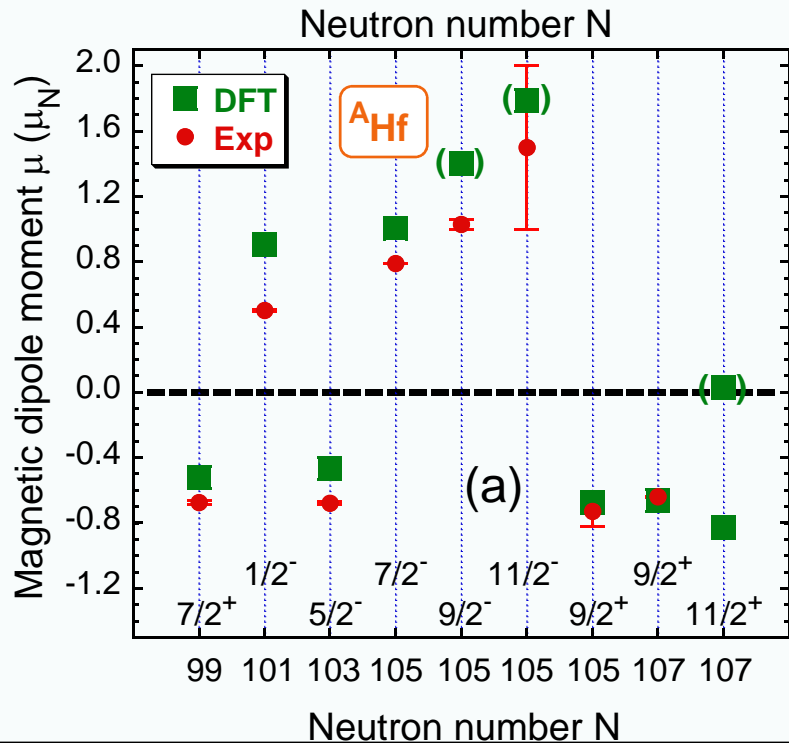
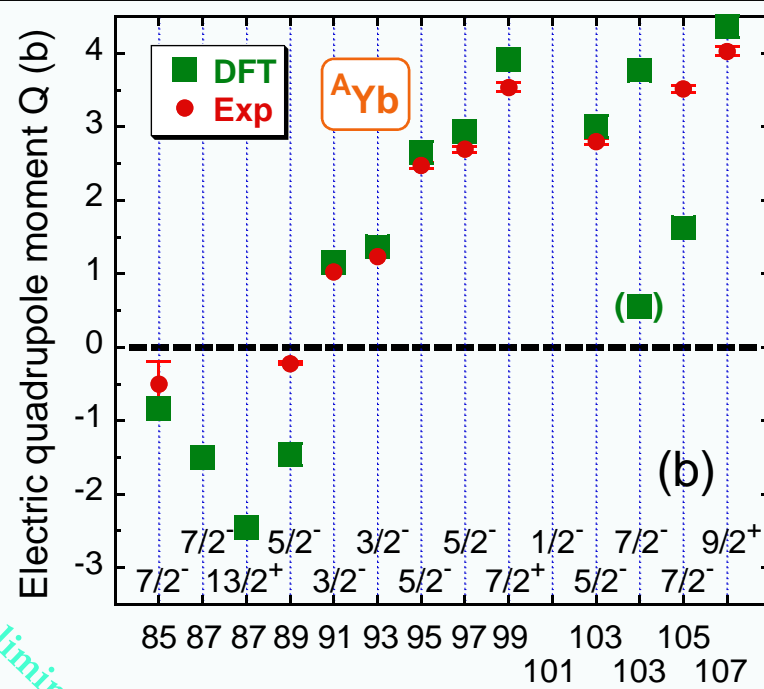
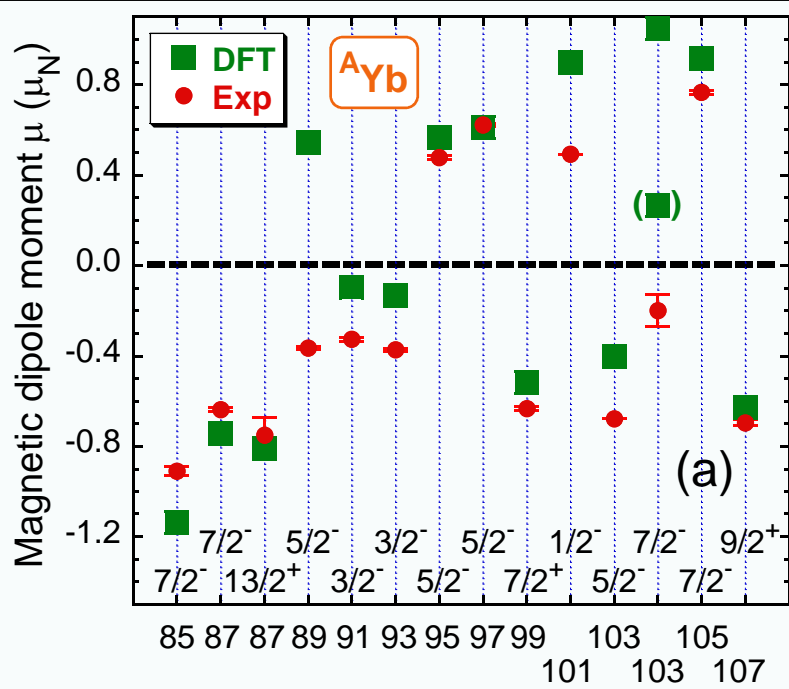
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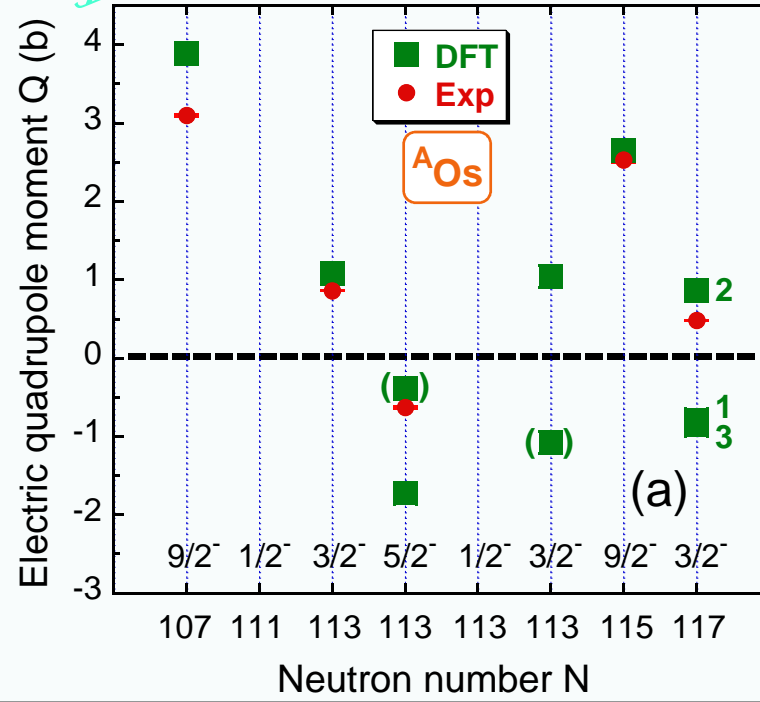
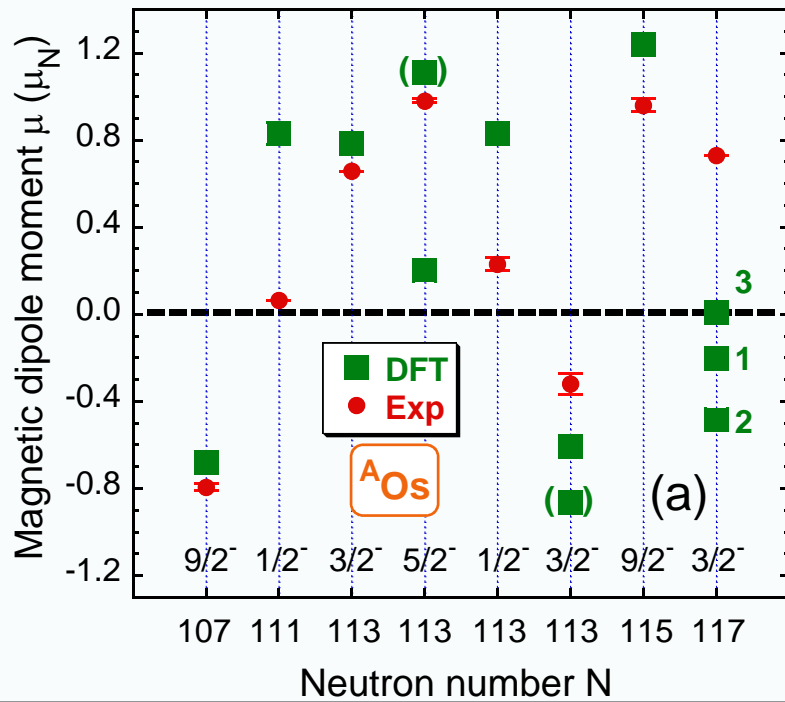
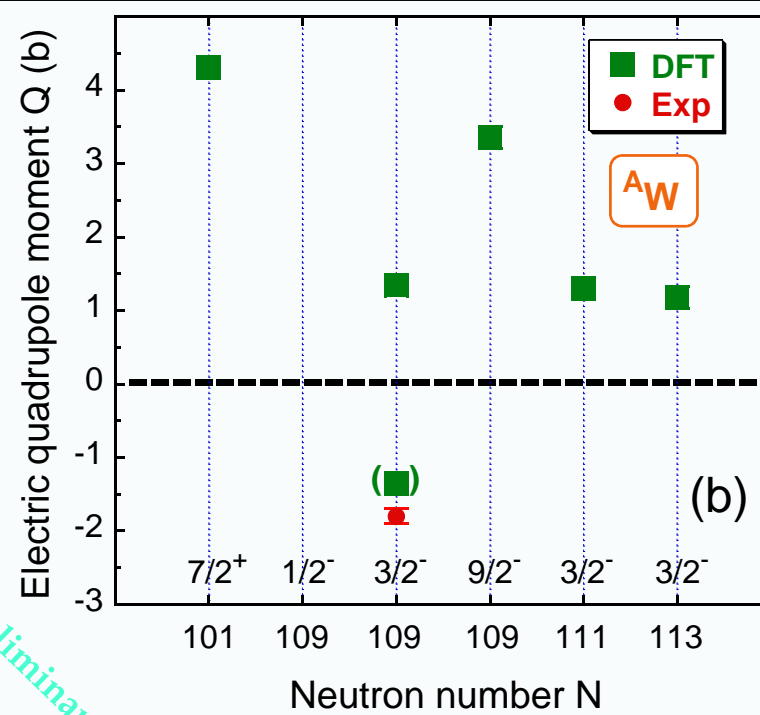
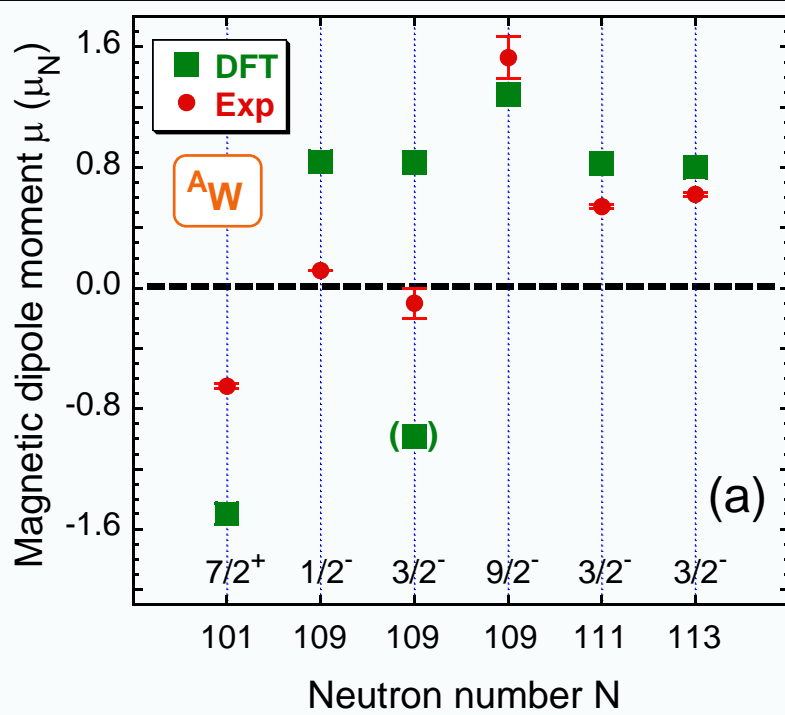




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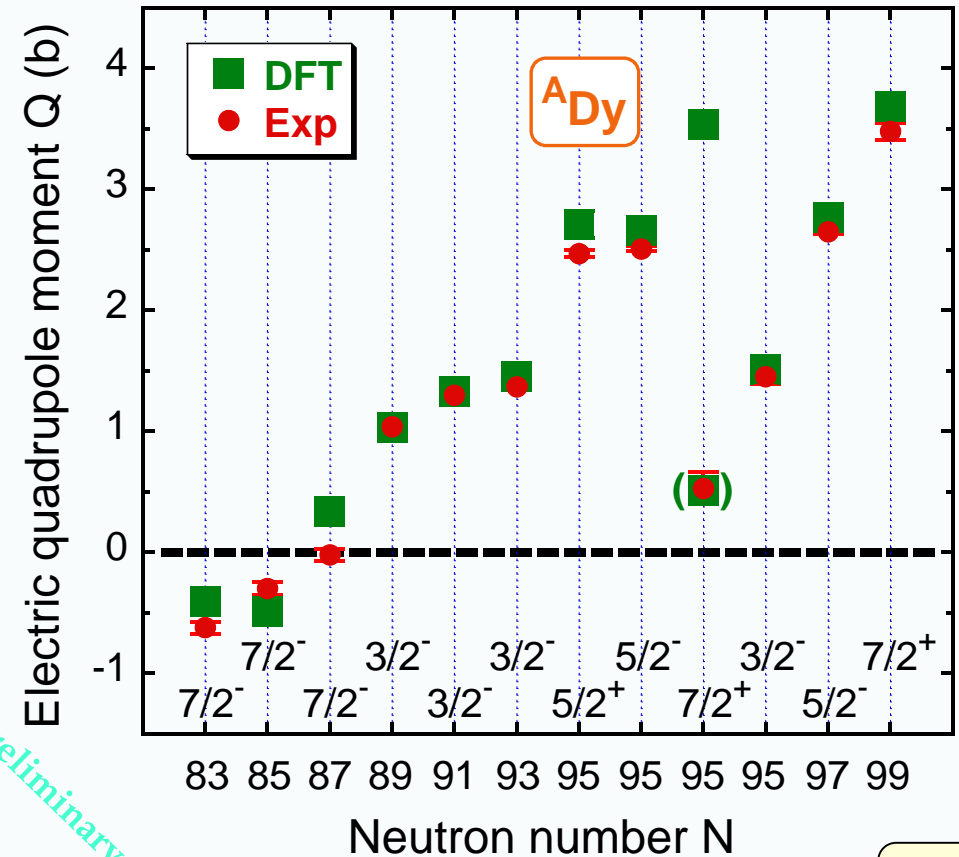
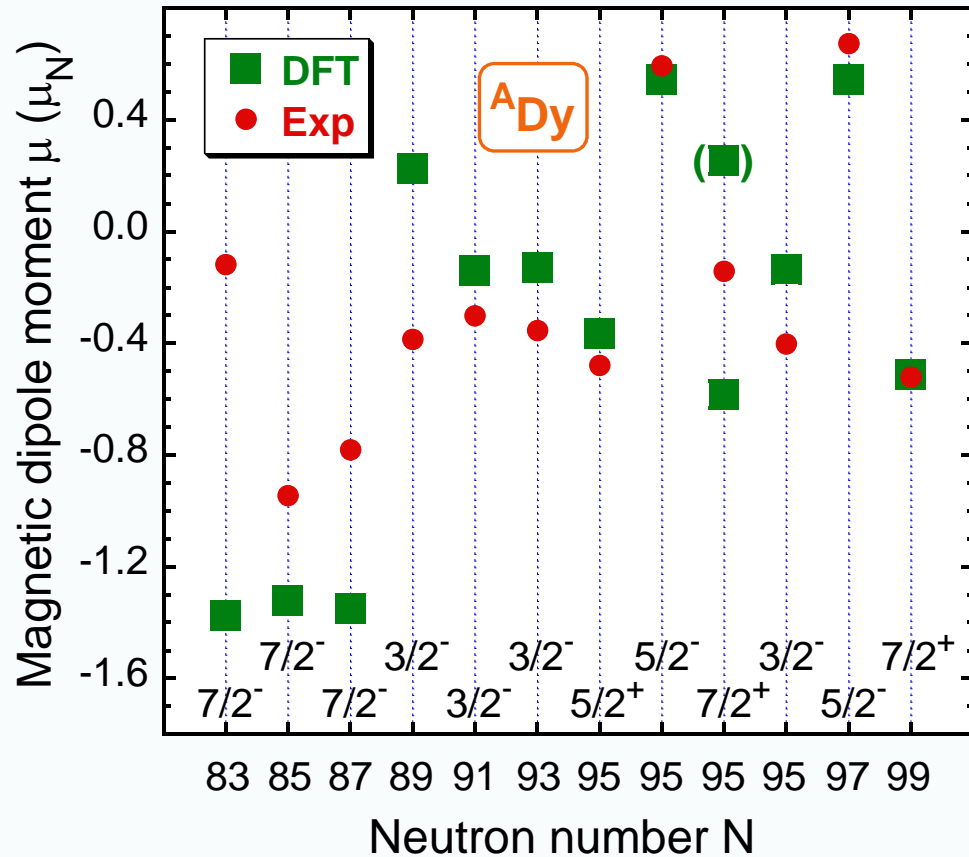


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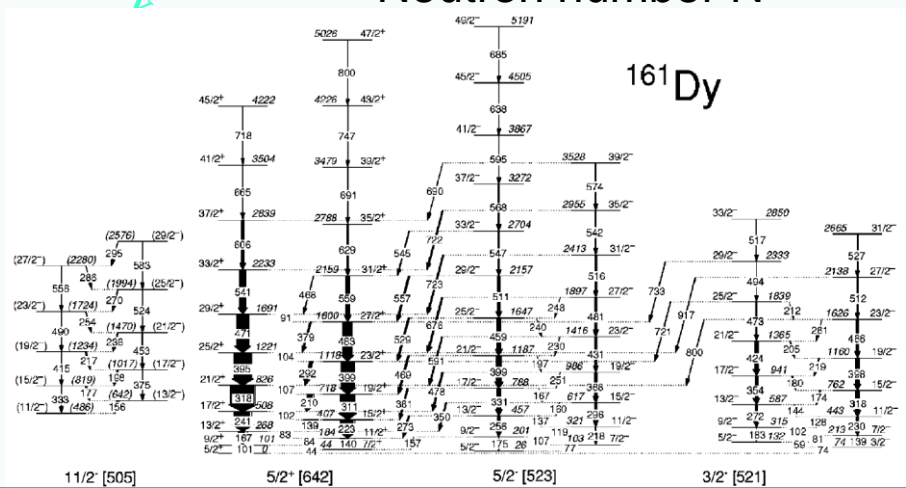


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$9/2+$	100.4	$7/2-$	103.0	$5/2-$	131.7
	*			$3/2-$	74.6
$7/2+$	43.8				*
	*	$5/2-$	25.7		
$5/2+$	0		*		
	*				

^{161}Dy

S. J. Margraf et al.,
Phys. Rev. C52, 2429 (1995)



A. Junglausa et al.,
Phys. Rev. C67, 034302 (2003)



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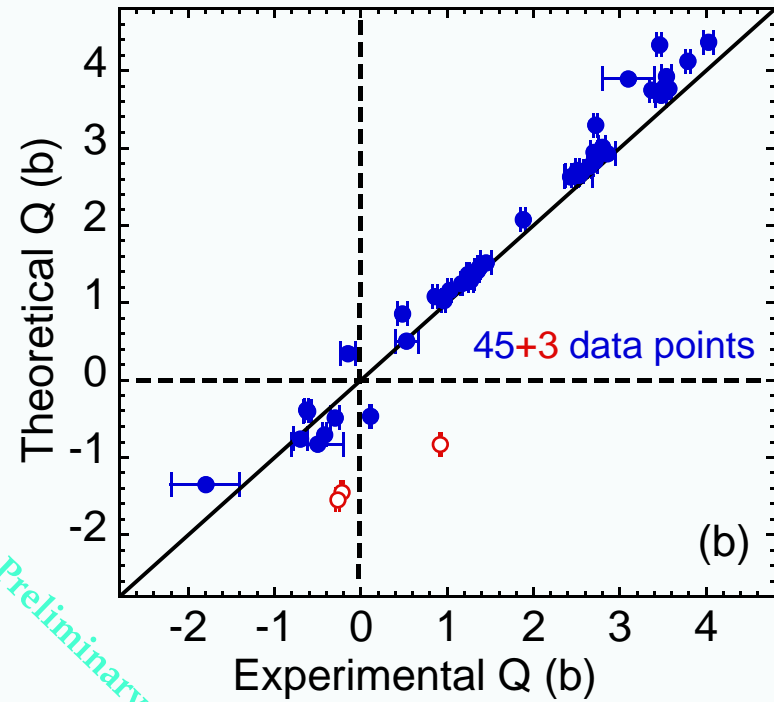
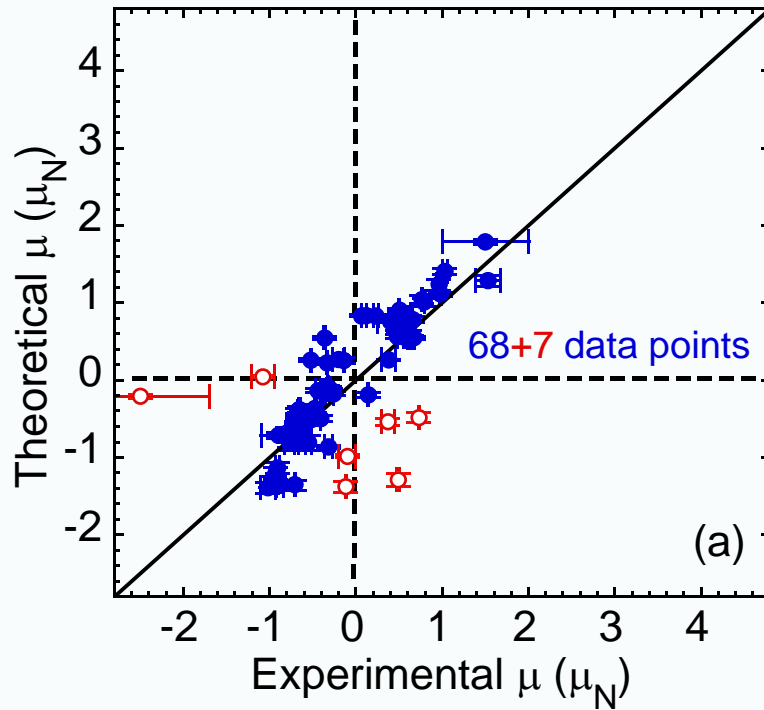
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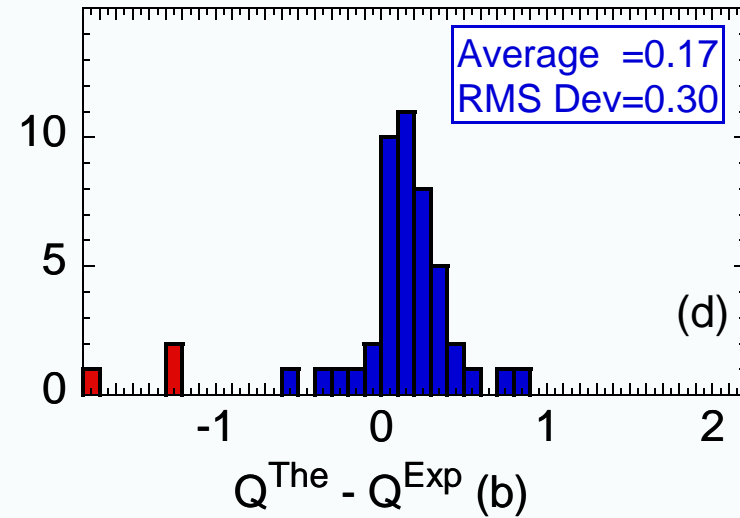
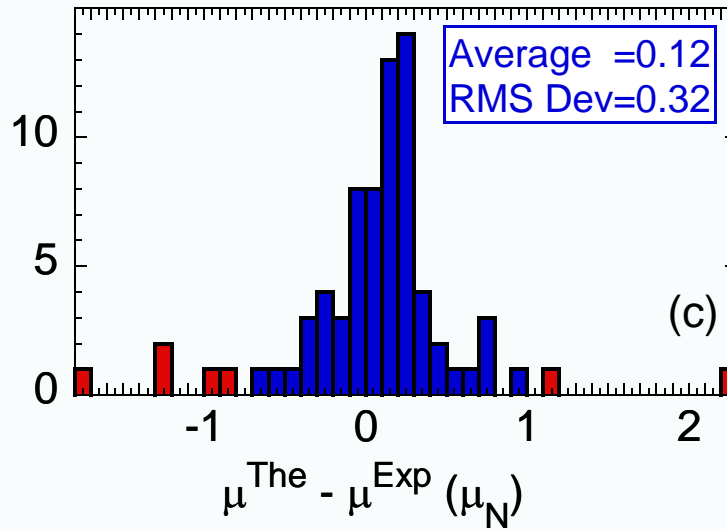
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Summary of results obtained in the Gd – Os isotopes



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Conclusions

1. Essential role of **simultaneously** taking into account:
 - a) **Polarization**
 - b) **Self-consistency**
 - c) **Symmetry restoration**
2. Isovector spin-spin interaction is essential in determining the **spin polarisation** and magnetic dipole moments.
3. A single parameter, the **isovector Landau parameter** of $g_0'=1.7(4)$, has been adjusted to data across the mass table.
4. Large single-particle phase space (well beyond the valence space) allows using the **bare effective charges** and **bare g-factors**. (No adjustable “effective” values are needed.)
5. The calculated **magnetic dipole moments μ** and **electric quadrupole moments Q** reproduce the known experimental data in odd-N open-shell isotopes of Gd-Os with the RMS deviations of $\Delta\mu=0.32 \mu_N$ and $\Delta Q=0.30 \text{ b}$, respectively.
6. The effects of the triaxiality, octupolarity, two-body currents, K-mixing, and configuration interaction (...) **remain to be studied.**



Thank you



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