

Tracing two-neutron halos in N=28 isotones: A three-body adventure

Jagjit Singh

jagjit.singh@manchester.ac.uk

IOP Joint APP, HEP, and NP Conference, Liverpool.

Parallel Session: H, 10.04.2024 (11:00-11:15)

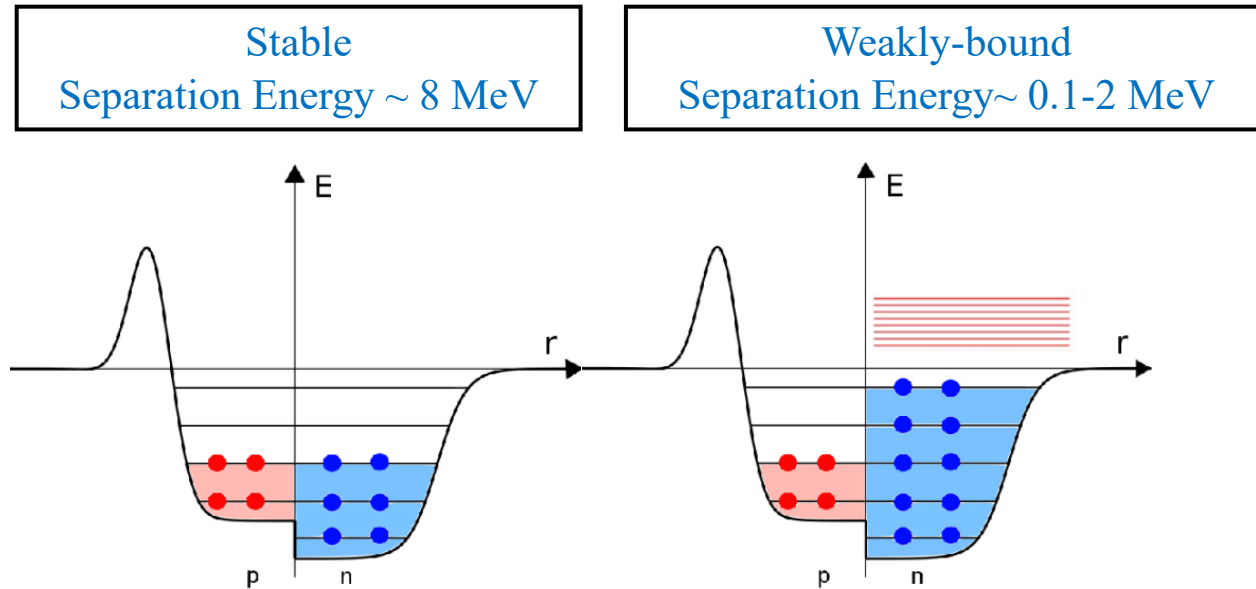
Based on:

Prediction of two-neutron halos in the N=28 isotones ^{40}Mg and ^{39}Na

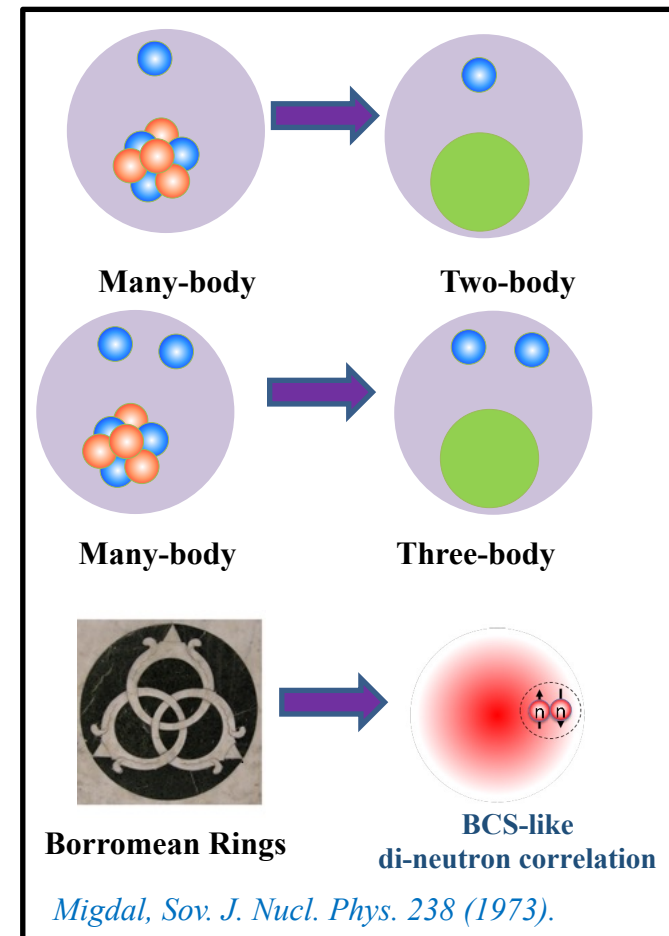
Jagjit Singh (UoM), J. Casal (Seville), W. Horiuchi (OMU, Osaka), N. R. Walet (UoM), W. Satula (Warsaw)

[arXiv:2401.05160 \(2024\).](https://arxiv.org/abs/2401.05160)

Introduction: Weakly-bound neutron-rich systems



Picture: L. Moschini Master Thesis (ArXiv:1410.7167)



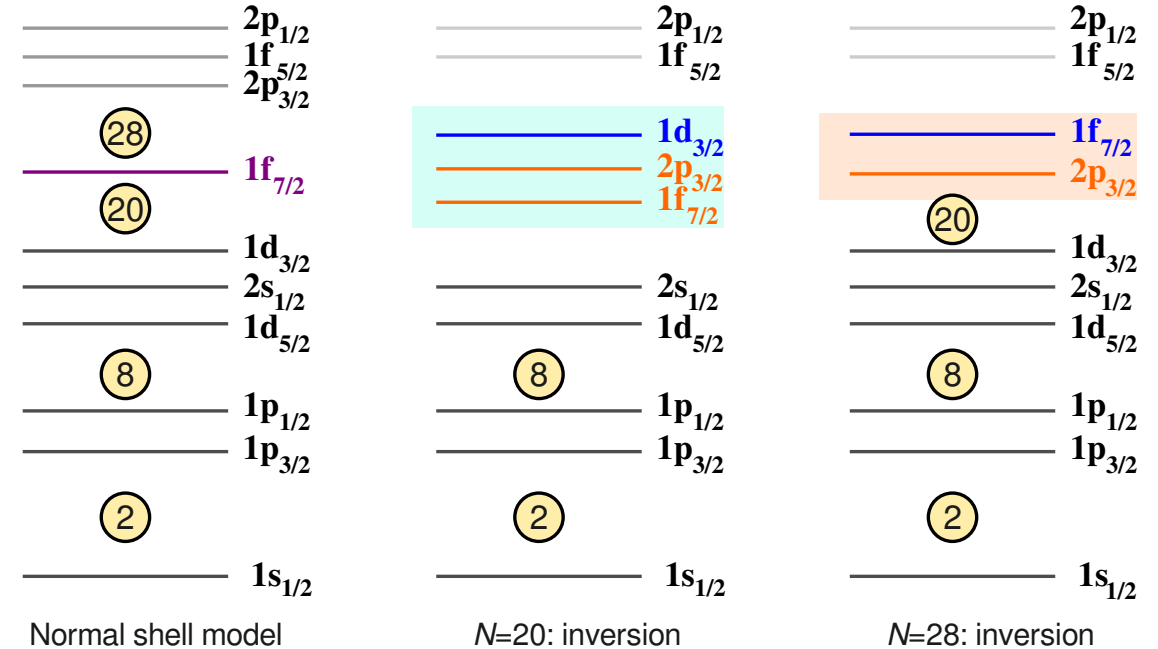
- Valence nucleon(s) - **weakly bound**. ^{11}Be ($^{10}\text{Be}+n$) $S_n=0.503$ MeV, ^6He ($^4\text{He}+n+n$) $S_{2n}=0.973$ MeV. *P. G. Hansen and B. Jonson, Europhys. Lett.* 4, 409 (1987).
- Role and **treatment of continuum** is more significant. *JS, PhD thesis, University of Padova (2016)*.
- Few-body approaches are convenient tool for such systems. *JS et al., PRC 101, 024310 (2020)*.
- Exotic features such as **halo/Borromean formation** *I. Tanihata et al., Phys. Rev. Lett.*, 55 (1985) 2676.
- Strong short-range cluster like correlation **di-neutron** - which plays important role in their binding mechanisms leads to weak binding of the system and low-density at the surface. *A.B.Migdal Sov. J. Nucl. Phys.* 16, 238 1973. *M.Matsuo PRC 73, 044309 (2006)*. *A.Gezerlis, PRC 81, 025803 (2010)*.

Halo formation on lower Z-side of the magic numbers N=20 and 28

³³ Si	³⁴ Si	³⁵ Si	³⁶ Si	³⁷ Si	³⁸ Si	³⁹ Si	⁴⁰ Si	⁴¹ Si	⁴² Si	⁴³ Si	⁴⁴ Si	⁴⁵ Si
³² Al	³³ Al	³⁴ Al	³⁵ Al	³⁶ Al	³⁷ Al	³⁸ Al	³⁹ Al	⁴⁰ Al	⁴¹ Al	⁴² Al	⁴³ Al	
³¹ Mg	³² Mg	³³ Mg	³⁴ Mg	³⁵ Mg	³⁶ Mg	³⁷ Mg	³⁸ Mg	³⁹ Mg	⁴⁰ Mg	⁴¹ Mg		
³⁰ Na	³¹ Na	³² Na	³³ Na	³⁴ Na	³⁵ Na	³⁶ Na	³⁷ Na	³⁸ Na	³⁹ Na			
²⁹ Ne	³⁰ Ne	³¹ Ne	³² Ne	³³ Ne	³⁴ Ne							
²⁸ F	²⁹ F	³⁰ F	³¹ F									

N=20

N=28



- Recent observation of the disappearance of the $N = 20$ shell gap at the low-Z side of the $N = 20$ chain, led to the identification of the ^{29}F system as the heaviest known two-neutron Borromean-halo nucleus. *S. Bagchi et al., PRL 124, 222504 (2020). JS et al., PRC 101, 024310 (2020). LF, JC, WH, JS et al., Comm. Physics 3, 132 (2020). JC, JS, et al., PRC 102, 064227 (2020).*
- Motivated by this observation, it is interesting to explore the low-Z side of the $N = 28$ shell closure for potential two-neutron Borromean halos in the Na and Mg isotopes).

- The $1f_{7/2}$ orbit is bordered by two magic numbers, i.e., 20 and 28. **20** has a HO origin, whereas the **28** has a SO origin.
- Inversion occur, when energy gap, associated with filling of shell closures disappears.

Neutron-rich F, Na, and Mg isotopes

^{28}F (Z=9, N=19)

Extending the southern shore of the island of inversion to ^{28}F

A. Revel et al., PRL 124, 152502 (2020).

Unbound
 $^{27}\text{F}+n$

-- ground state resonance of ^{28}F at
0.199(6) MeV (l=1~79%)

Inversion !!

-- 1st excited state resonance around
0.966 MeV (l=2~72%)
PRL 124, 152502 (2020).

^{29}F (Z=9, N=20)

JS, JC, WH, LF and AV, PRC 101, 024310 (2020).

LF, JC, WH, JS and AV, Nature: Communication Physics 3, 132 (2020).

JC, JS, LF, WH and AV, PRC 102, 064227 (2020).

S. Bagchi et al., PRL 124, 222504 (2020).

Weakly
Bound
 $^{27}\text{F}+n+n$

$s_{2n} = 1.443$ (436) MeV

$= 1.440$ (650) MeV

PRL 109, 202503 (2012).

CPCC 41, 030003 (2017).



Tools : Three-body structure model and Glauber reaction theory.

^{40}Mg (Z=12, N=28)

H. L. Crawford et al. Phys. Rev. Lett. 122, 052501 (2019).

T. Baumann et al., Nature volume 449, 1022 (2007).

A. O. Macchiavelli, et al., Eur. Phys. J. A 58, 66 (2022).

Weakly
Bound
 $^{38}\text{Mg}+n+n$

$s_{2n} = 0.670 \pm 0.710$ MeV

Chinese Physics C 45 (3), 030003 (2021).

No experimental data on ^{39}Mg spectrum

Theoretical predictions within GSM

^{39}Na (Z=11, N=28)

Discovery of ^{39}Na

D. S. Ahn et al., PRL 129, 212502 (2022).

KY Zhang et al., PRC 107, L04130303 (2023).

Weakly
Bound
 $^{37}\text{Na}+n+n$

No data on ^{38}Na spectrum
and s_{2n} of ^{39}Na .

Three-Body Structure Model- Hyperspherical framework

Three-body Hamiltonian is given by

$$H = T + \sum_{i=1}^2 V_{core+n_i} + V_{nn} + V_{3b}$$

$$V_{core+n_i} = \left(-V_0 + V_{ls} \vec{l} \cdot \vec{s} \frac{1}{r} \frac{d}{dr} \right) \frac{1}{1 + \exp\left(\frac{r-R}{a}\right)}, \quad V_{3b}(\rho) = v_{3b} e^{-(\rho/\rho_0)^2},$$

PHYSICAL REVIEW C **102**, 064627 (2020)

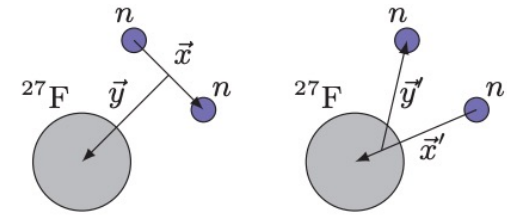


FIG. 1. Jacobi- T (left) and $-Y$ (right) coordinates for the ^{29}F nucleus described as $^{27}\text{F} + n + n$.

nn interaction- Gogny-Pires-Tourelil (GPT) interaction including central, spin-orbit and tensor terms.

D. Gogny, P. Pires, and R. D. Tourrelil, PLB 32, 591 (1970).

The three-body force is modelled as a simple Gaussian potential, where $\rho = \sqrt{x^2 + y^2}$, is **hyper-radius** and $\rho_0 = 6$ fm and the strength v_{3b} is adjusted to recover s_{2n} .

We use the analytical transformed harmonic oscillator (THO) basis.

J. Casal, M. Rodrguez-Gallardo, and J. M. Arias, Phys. Rev. C 88, 014327 (2013).

The diagonalization of the three-body Hamiltonian requires the computation of the corresponding kinetic energy & potential matrix elements.

Reaction cross-section within Glauber model

The reaction cross section for a projectile-target collision integrating the reaction probability with respect to the impact parameter b

Y. Ogawa et al., Prog. Theor. Phys. Suppl. 142, 157 (2001a)

$$\sigma_R = \int db (1 - |e^{i\chi(b)}|^2),$$

Phase shift function

$$e^{i\chi(b)} = \langle \Phi_0^P \Phi_0^T | \prod_{A_P} \prod_{A_T} [1 - \Gamma_{NN}(s_i^P - s_j^T + b)] | \Phi_0^P \Phi_0^T \rangle$$

Profile function

$$\Gamma_{NN}(b) = \frac{1 - i\alpha}{4\pi\beta} \sigma_{NN}^{\text{tot}} \exp\left(-\frac{b^2}{2\beta}\right)$$

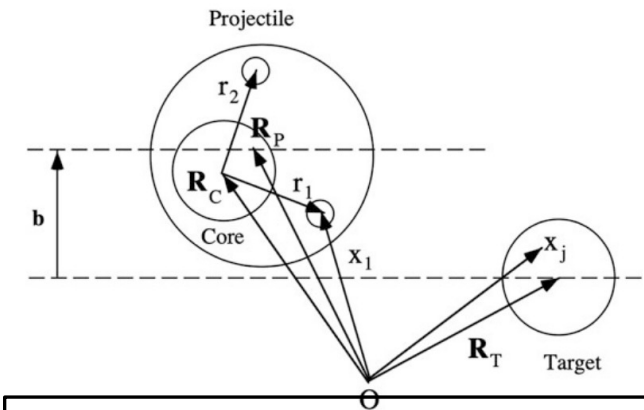
α, β : determined so as to reproduce the NN scattering

B. Abu-Ibrahim et al., PRC 77, 034607 (2008). W. Horiuchi et al., PRC 75, 044607 (2007).

Phase-shift function: Many-body operator

Approximate using a cumulant expansion: Nucleon-Target profile function

→ Input: **Nuclear density and Profile function no adjustable parameters**



- $\mathbf{r}_i = (\mathbf{s}_i, \mathbf{z}_i)$ is the coordinate of the two-halo neutron(s) with index $i = 1, 2$.
- \mathbf{R}_C , \mathbf{R}_P , and \mathbf{R}_T are the position vectors to the COM of the core nucleus, projectile nucleus, and target nucleus, respectively. $\mathbf{x}_i - \mathbf{R}_C = \mathbf{r}_i$.
- \mathbf{s}_i and \mathbf{s}_j are the two-dimensional vectors of the coordinates for the P and T, measured from their COM which lies on a plane perpendicular to the incident momentum of the projectile

Two-body (core+n) models for ^{39}Mg and ^{38}Na

Jagjit Singh et al., arXiv:2401.05160 [nucl-th] (2024).

[K. Fosse, et. al., PHYSICAL REVIEW C 94, 054302 \(2016\).](#)

Unbound ground state of ^{39}Mg is predicted to be either a $J^\pi = 7/2^-$ or $3/2^-$ state.

A narrow $J^\pi = 7/2^-$ or $3/2^-$ ground-state candidate exhibits a resonant structure at 129 KeV.

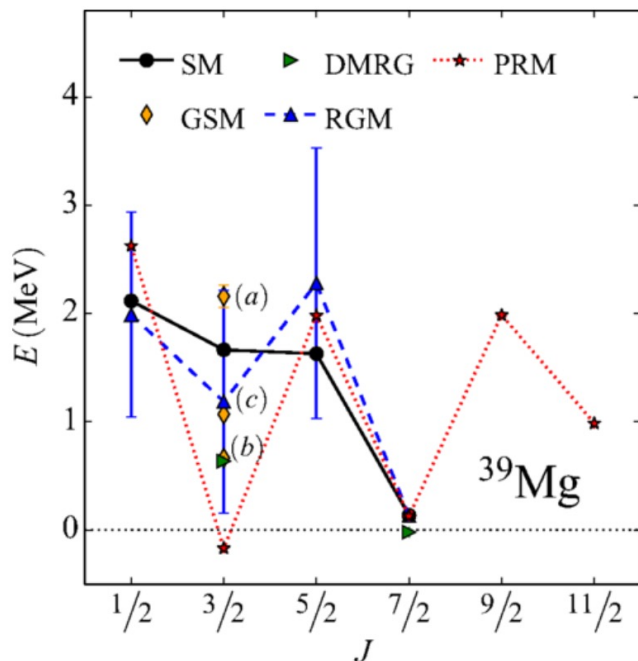


Table 1: Parameter sets for the $^{38}\text{Mg} + n$ (upper-panel) and $^{37}\text{Na} + n$ (lower-panel) Woods-Saxon interactions, Eq. (6). Here a is diffuseness, $V_0^{(l)}$ is the potential depth and E_R is the position of the resonances. Note that $r_0 = 1.25$ fm, and $V_{ls} = 16.842$ MeV (for $^{38}\text{Mg} + n$) and 16.324 MeV (for $^{37}\text{Na} + n$) are fixed.

$^{38}\text{Mg} + n$					
Set	Scenario	a (fm)	lj	$V_0^{(l)}$ (MeV)	E_R (MeV)
1	Normal	0.70	$f_{7/2}$	38.225	0.129
			$p_{3/2}$	38.225	0.349
2	Degenerate	0.75	$f_{7/2}$	38.400	0.129
			$p_{3/2}$	38.400	0.135
3	Inverted	0.75	$f_{7/2}$	37.880	0.349
			$p_{3/2}$	38.425	0.130
$^{37}\text{Na} + n$					
1	Degenerate	0.70	$f_{7/2}$	38.225	0.539
			$p_{3/2}$	38.225	0.599
2	Inverted	0.75	$f_{7/2}$	38.400	0.522
			$p_{3/2}$	38.400	0.271
3	Inverted	0.75	$f_{7/2}$	37.880	0.734
			$p_{3/2}$	38.425	0.265

Configuration mixing and matter radii for ^{40}Mg and ^{39}Na

^{40}Mg (s_{2n}) = 0.670 (0.710) MeV

M. Wang, *et. al.*, Chinese Physics C **45** (3), 030003 (2021).

Matter radius of core ^{38}Mg =3.60 fm

S. Watanabe, *et. al.*, PRC **89**, 044610 (2014).

^{39}Na (s_{2n}) = unbound

M. Wang, *et. al.*, Chinese Physics C **45** (3), 030003 (2021).

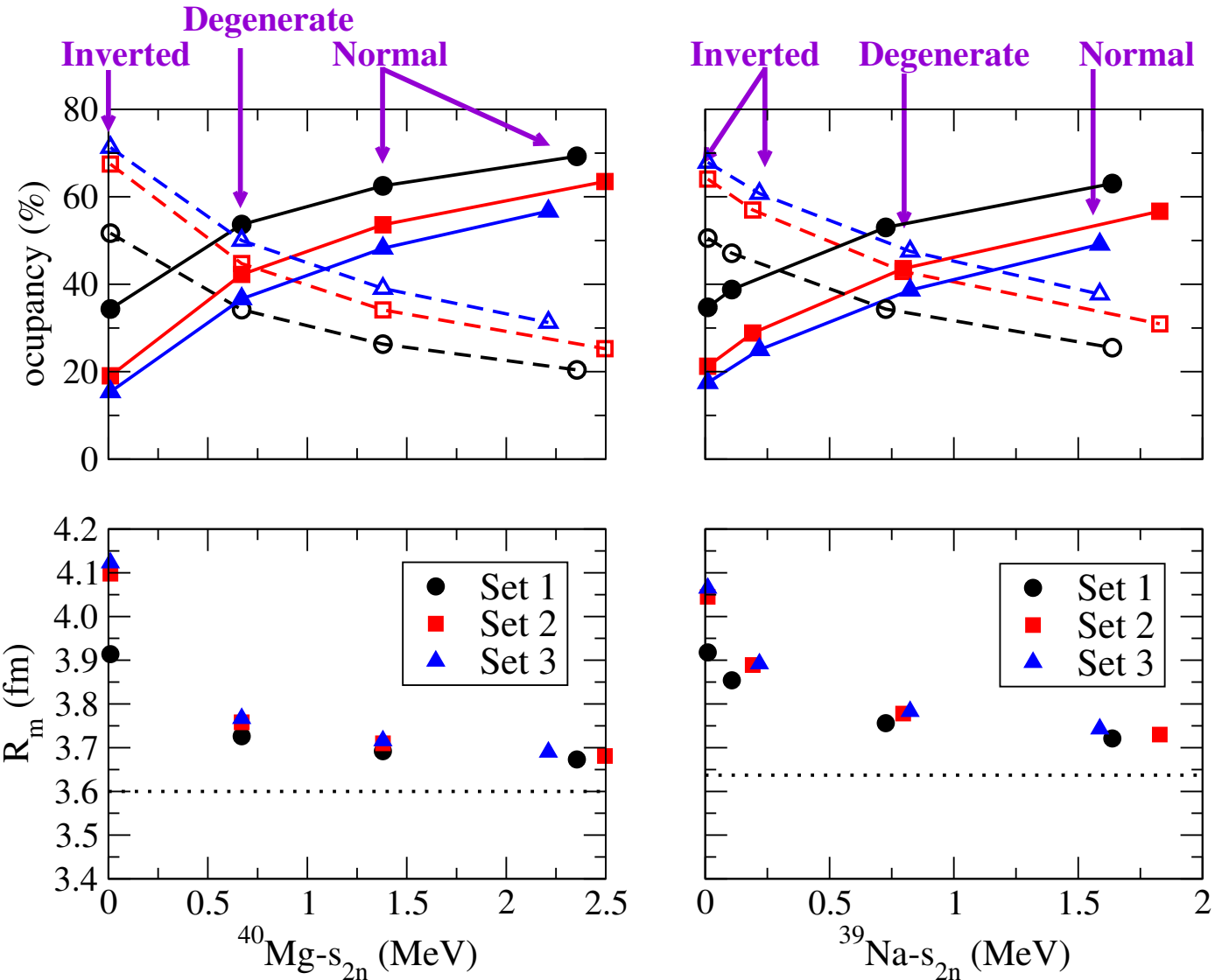
^{39}Na (s_{2n}) = bound (Contradicts AME)

D. S. Ahn *et al.*, PRL **129**, 212502 (2022).

Matter radius of core ^{37}Na =3.64 fm

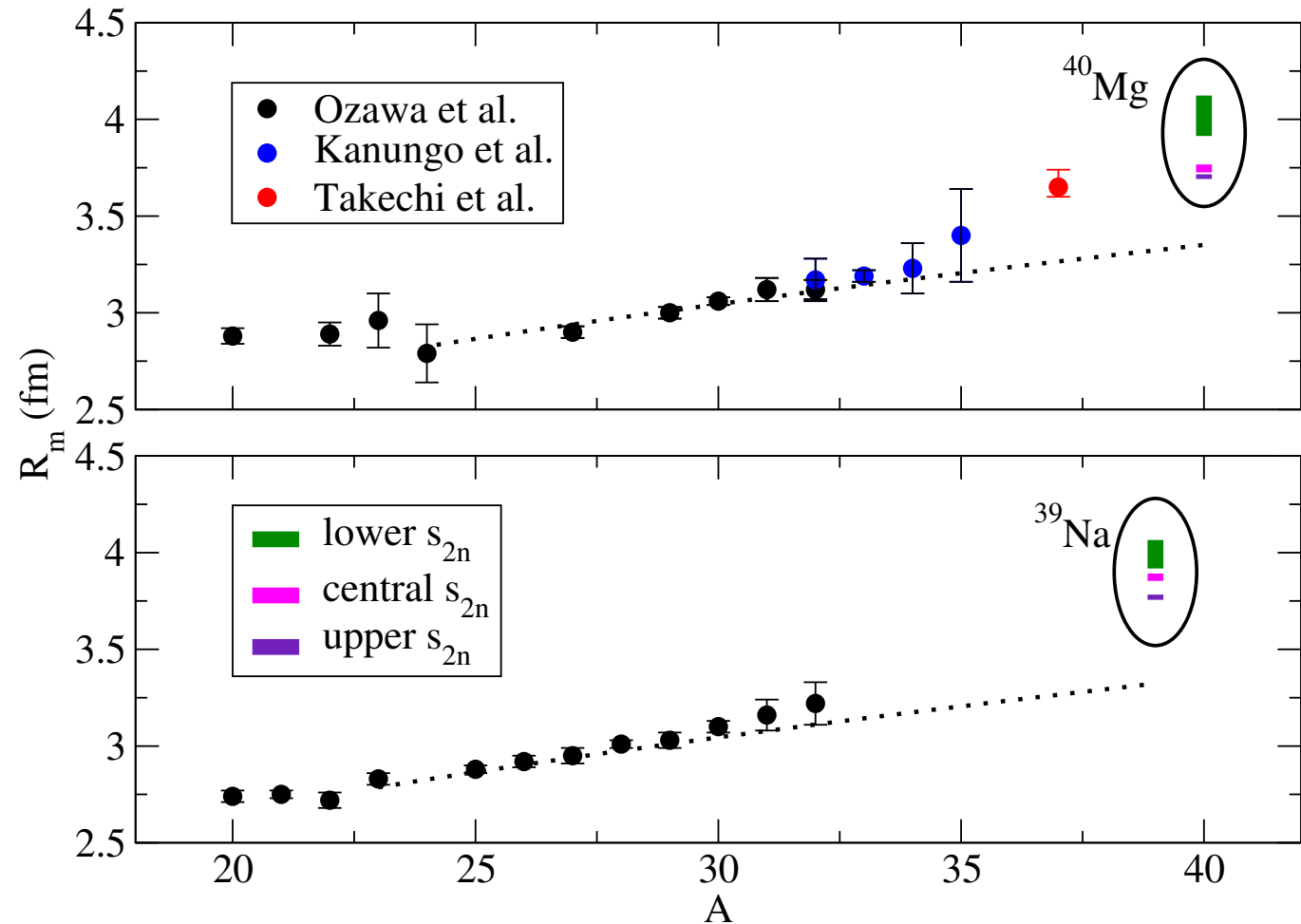
L. Geng, *et. al.*, NPA **730**, 80 (2004).

The larger change in R_m *w.r.t* core involves wave function which contains significant $(p_{3/2})^2$ (dotted lines) weight, pointing toward the necessity of intruder configurations to sustain halo formation. $(f_{7/2})^2$ (solid lines)



Matter radii for ^{40}Mg and ^{39}Na

- Dotted black lines in the figure correspond to the weighted fit of the experimental data points with the standard $R_0 A^{1/3}$ formula.
- The radii of ^{40}Mg and ^{39}Na are higher than the standard fitted value.
- This observation implies a likely two-neutron halo structure in the ground state of ^{40}Mg and ^{39}Na , and the corresponding melting of the traditional $N = 28$ shell gap is due to the intrusion of the $p_{3/2}$ orbital.



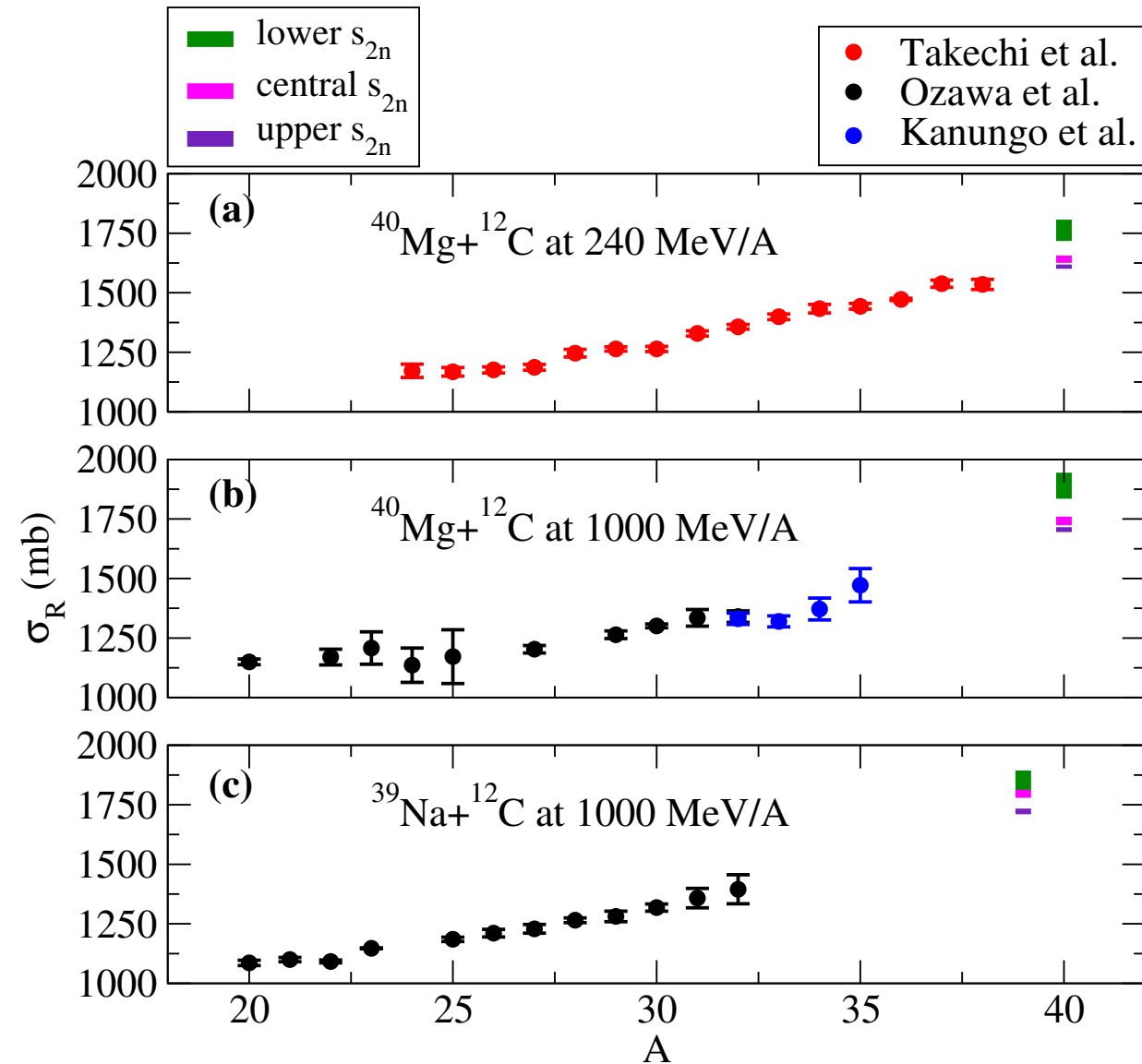
- Ozawa *et al.*, Nuclear Physics A **691** (3), 599 (2001).
- Kanungo *et al.*, Phys. Rev. C **83**, 021302 (2011).
- Takechi *et al.*, Phys. Rev. C **90**, 061305(R) (2014).
- Jagjit Singh *et al.*, arXiv:2401.05160 [nucl- th] (2024).

Reaction cross-sections for ^{40}Mg and ^{39}Na : within Glauber reaction theory

- Experimentally, a very obvious way to determine whether a nucleus is a halo nucleus, is to look for an enhanced reaction cross section. Thus, we examine the total reaction cross section by employing the conventional Glauber theory. *B. Abu-Ibrahim et al., PRC 77, 034607 (2008). W. Horiuchi et al., PRC 75, 044607 (2007).*
- Using this prescription, we predict the σ_R for ^{40}Mg and ^{39}Na at different incident energies. The predicted values of σ_R for ^{40}Mg and ^{39}Na show significant enhancement with respect to the observed σ_R in the lower-A isotopes for both choices of energy.

Thus, our results provide a clear signal of the 2n-halo structure formation in ^{40}Mg and ^{39}Na and hence melting of the $N = 28$ shell closure.

- Ozawa et al., Nuclear Physics A 691 (3), 599 (2001).*
- Kanungo et al., Phys. Rev. C 83, 021302 (2011).*
- Takechi et al., Phys. Rev. C 90, 061305(R) (2014).*
- Jagjit Singh et al., arXiv:2401.05160 [nucl-th] (2024).*

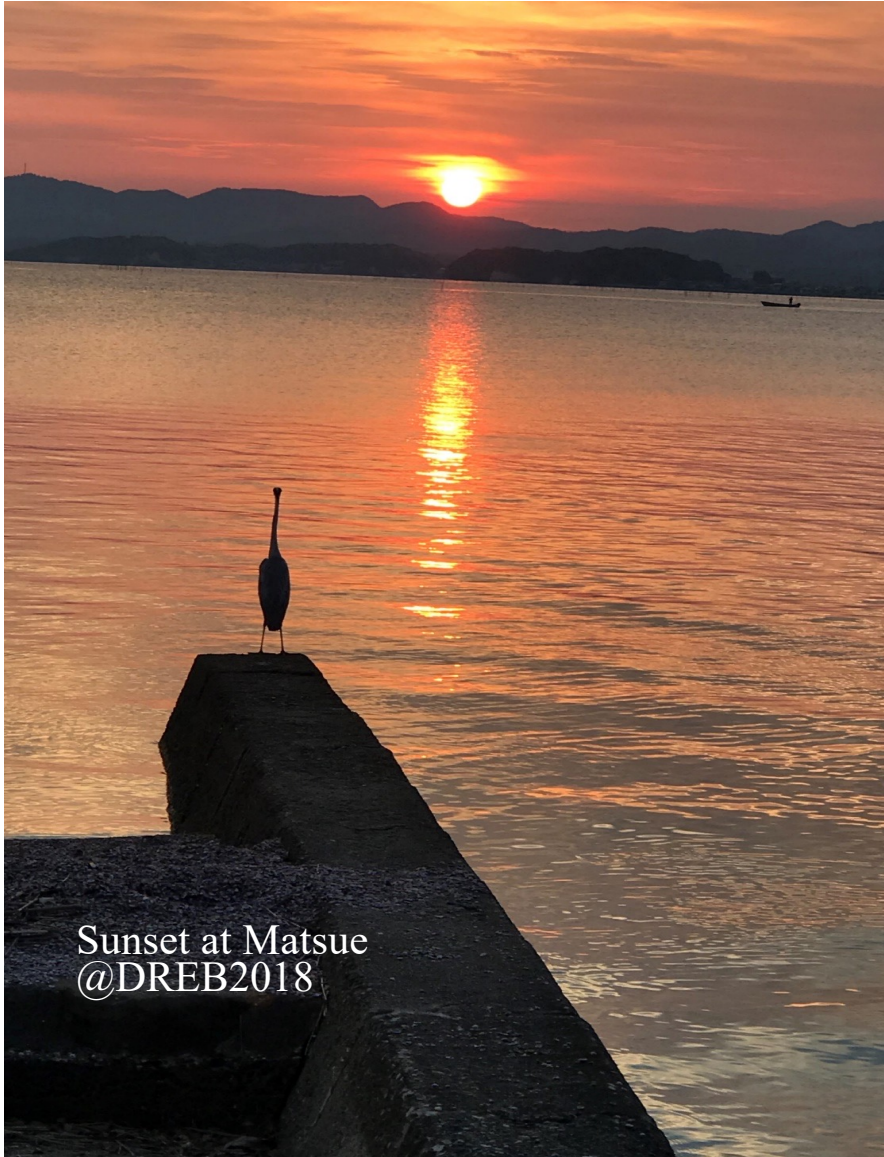


Summary:

- We started with studying melting of $N=20$ (*JS, JC, WH, LF, and AV, PRC 101, 024310 (2020)*) for ^{29}F . Our predictions got boost with new measurements (*S. Bagchi et al., PRL 124, 222504 (2020)* and *A. Revel et al., PRL 124, 152502 (2020)*). We updated our calculations with precise calculations along with detailed analysis of electric-dipole response and reaction calculations (*LF, JC, WH, JS, and AV, Nature: Commun. Phys. 3, 132 (2020)* and *JC, JS, LF, WH, and AV, PRC 102, 064627 (2020)*).
- Motivated by melting $N=20$ ends up in formation of Borromean in ^{29}F , by using same prescription we reported first three body calculations for ^{39}Na and ^{40}Mg lying on low- Z side of $N=28$. (*JS, JC, WH, NW, and WS arXiv:2401.05160 [nucl- th] (2024)*),
- Our results calls for new precise mass measurements for s_{2n} of three-body systems and the low-lying continuum spectrum of subsystems to better constrain the theoretical models.
- The disappearance of the conventional $N=28$ shell gap and emergence of the halo leads to significant occupancy of intruder $p_{3/2}$ orbit in the ground state of ^{39}Na and ^{40}Mg . Nevertheless, it is imperative to verify this conclusion through experimental measurements of interaction cross sections.

Future Perspectives:

- It is interesting to see how our predictions/results alter with inclusion of core deformation effects.



Sunset at Matsue
@DREB2018

Thanks for kind attention.

Acknowledgements:

J. Casal (Seville)

W. Horiuchi (OMU, Osaka)

N. R. Walet (Manchester)

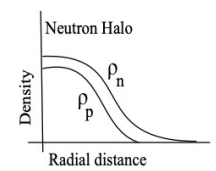
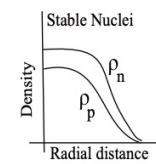
W. Satula (Warsaw)

L. Fortunato (Padova)

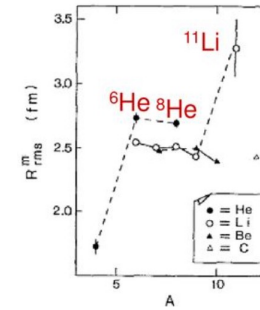
A. Vitturi (Padova)

Key features of halo nuclei

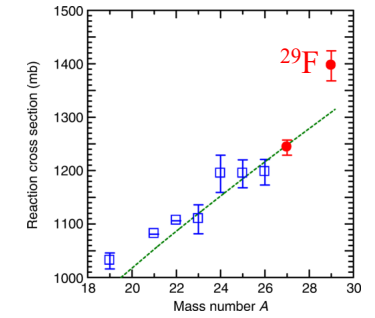
- Halo nuclei (manifestation of nuclear clustering)** exhibit: a *diffuse density distribution*, low one or more valence neutron(s) separation energies (s_n/s_{2n}), *abnormally large matter radius* and *large interaction cross sections*. P. G. Hansen and B. Jonson, *Europhys. Lett.* **4**, 409 (1987).
- Borromean systems (2n-halo):** Corresponding *core+n* subsystems being unbound, the strong correlations between the valence neutrons are key in binding two-neutron halos M. Zhukov *et al.*, *Phys. Rep.* **231**, 151 (1993). Y. Kikuchi *et al.*, *PTEP* **2016**, 103D03 (2016). Hagino & Sagawa, *PRC* **72**, 044321 (2005). Observed examples: ${}^6\text{He}$, ${}^{11}\text{Li}$, and ${}^{14}\text{Be}$, ${}^{17}\text{B}$, ${}^{19}\text{B}$, ${}^{22}\text{C}$ and ${}^{29}\text{F}$.
- Low-lying spectra** of the unbound *core+n* subsystems play an important role in shaping the properties of Borromean nuclei. One of the other salient features of halo nuclei is an *enhancement of the low-lying E1 (electric dipole) strength* into the continuum T. Aumann, *EPJA* **55**, 234 (2019). Experimentally this has been observed via *invariant mass spectroscopy in Coulomb dissociation (CD) experiments*.
- Theoretically three-body (core+n+n) models** have been found to describe reasonably well these features in *Borromean nuclei*.



R. Kanungo, *Handbook of Nuclear Physics* (2023).



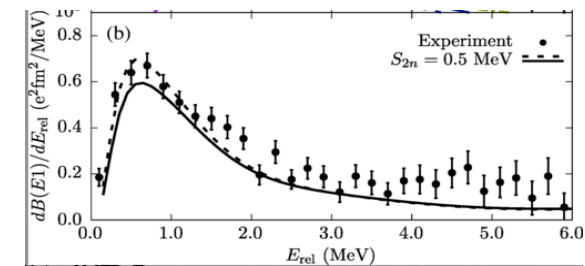
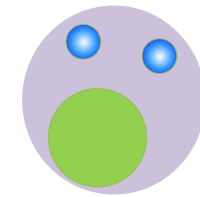
I. Tanihata *et al.*, *PRL* **55**, 2676 (1985).



S. Bagchi *et al.*, *PRL* **124**, 222504 (2020).



Borromean Rings



K. Cook *et al.*, *PRL* **124**, 212503 (2020).

Key inputs:

- Information on core+n low-lying spectrum
- Two-neutron separation energy

JS *et al.*, *PRC* **101**, 024310 (2020).

JS *et al.*, *arXiv:2401.05160 [nucl-th]* (2024).

Success story of ^{29}F

New Experiment confirms ground state resonance of ^{28}F at $0.199(6)$ MeV ($l=1\sim 79\%$) and 1st excited state resonance around 0.966 MeV ($l=2\sim 72\%$). **Inversion !!**

A. Revel et al., PRL 124, 152502 (2020).

Total reaction Cross section within Glauber Model

The calculated total reaction cross section using the standard Glauber theory are **1370 mb** if we assume $s_{2n} = 1.44$ MeV, and **1390 mb** if we take the lower limit ($S_{2n} \approx 1$ MeV), which are in good agreement with the observed interaction cross section **1396 ± 28 mb**.

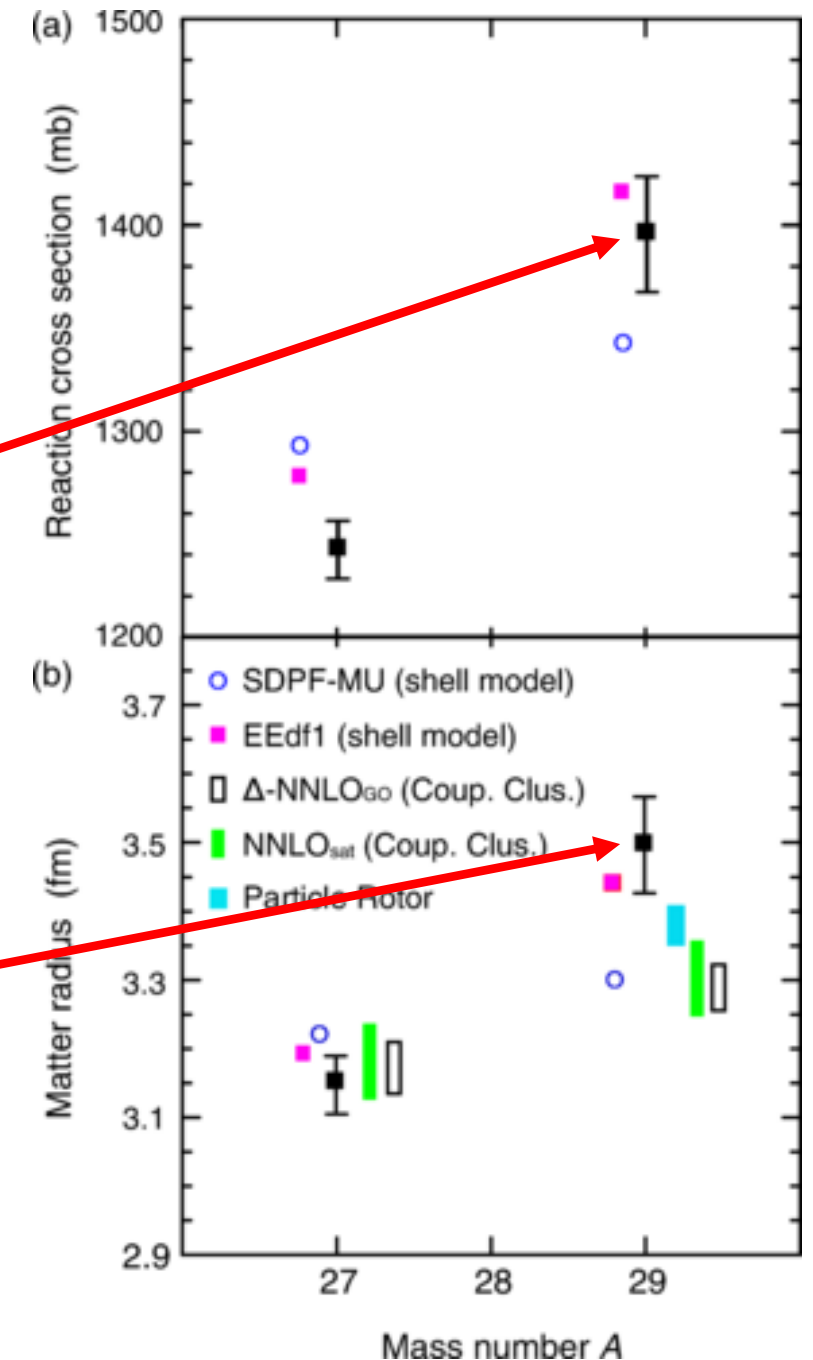
JC, JS, et al., PRC 102, 064627 (2020).

$s_{2n} (^{29}\text{F}) = 1.443 (436)$ MeV *PRL. 109, 202503 (2012).*

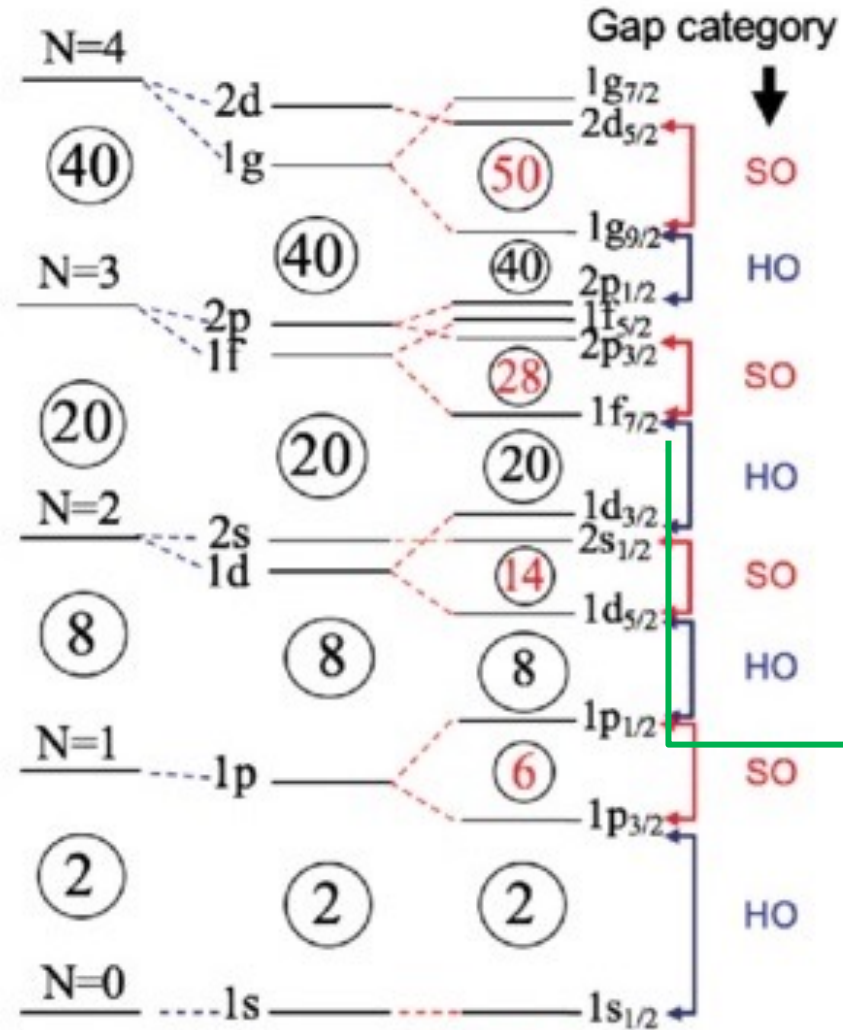
Matter radius of ^{29}F

The relative increase of matter radii with respect to ^{27}F core lies in the range **0.20-0.25** fm in the different choices of s_{2n} whereas experimental value is **0.35 (0.08) fm**.

JS, JC, WH, LF and AV, PRC 101, 024310 (2020).

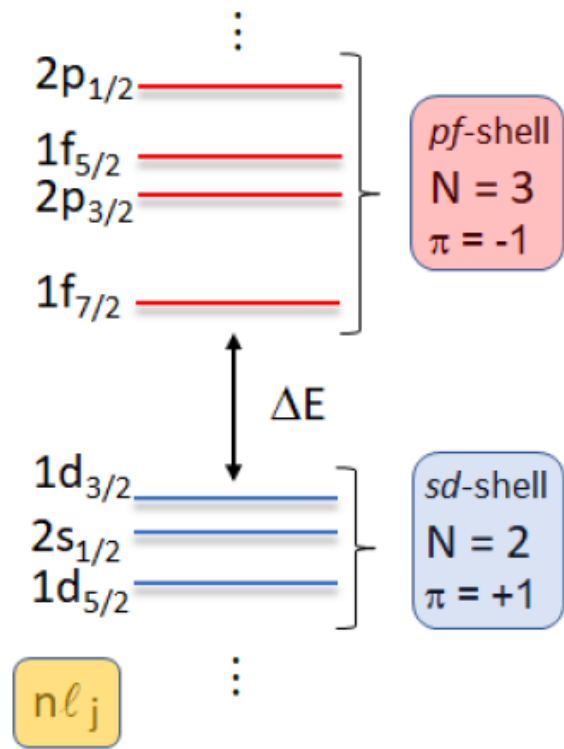


Background: Inversion mechanism

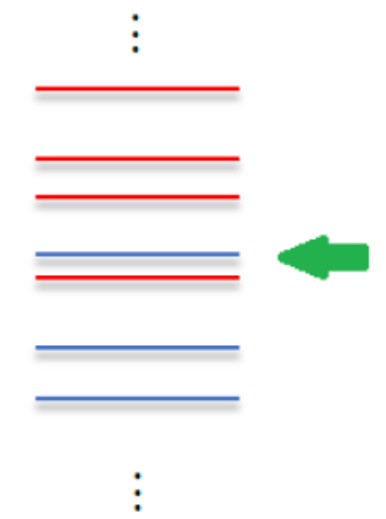


$$U(r) = \text{H.O.} + l^2 + \vec{l} \cdot \vec{s}$$

Standard ordering



Inversion occurs!



Inversion occur, when ΔE , associated with filling of **20 neutrons**, disappears and one level (or more) of the ***pf-shell* ($N = 3$)** gets lower than one (or more) of the levels of the ***sd-shell* ($N = 2$)**.

The $1f_{7/2}$ orbit is bordered by two magic numbers, i.e., 20 and 28. **20** has a HO origin, whereas the **28** has a SO origin.