

Suppressed electric quadrupole collectivity in ^{32}Si

Jacob Heery
Research Fellow
University of Surrey

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10/04/2024

Nuclear deformation

Oblate



$$dy = dx > dz$$

Spherical



$$dx = dy = dz$$

Prolate

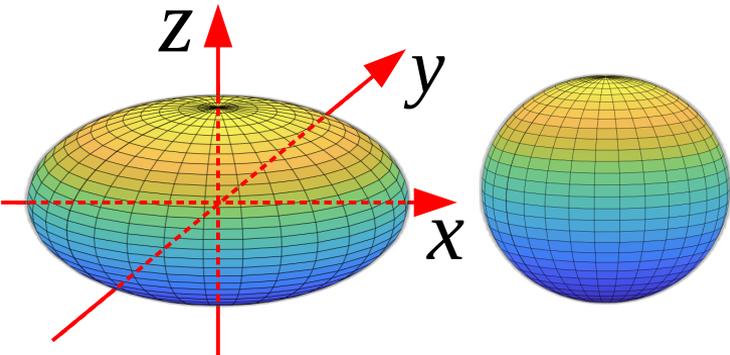
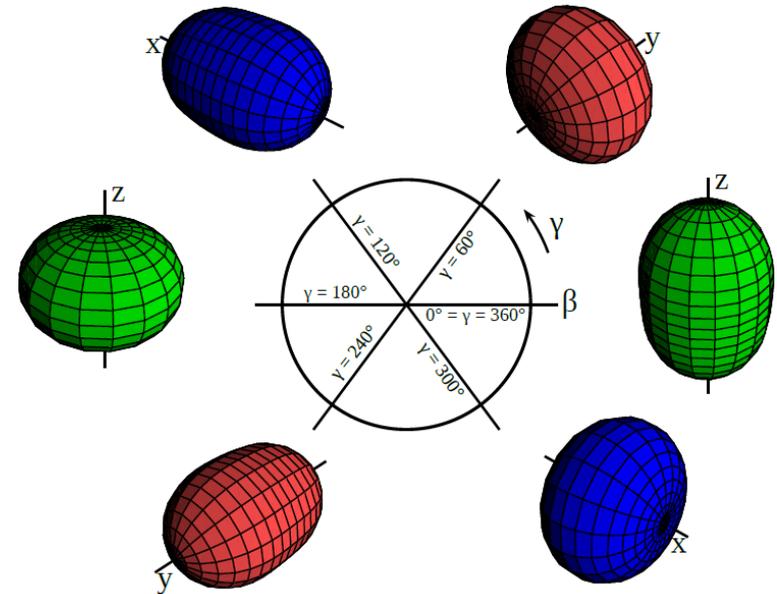


$$dx = dy < dz$$

Triaxial shapes – $dx \neq dy \neq dz$

Shape defined by two parameters:

- β – Magnitude
- γ – Form (prolate/oblate)



Oblate



$$dy = dx > dz$$

Spherical



$$dx = dy = dz$$

Prolate



$$dx = dy < dz$$

Triaxial shapes – $dx \neq dy \neq dz$

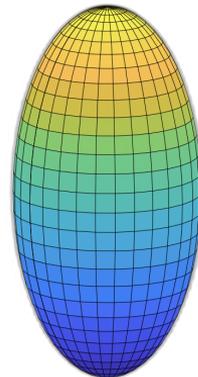
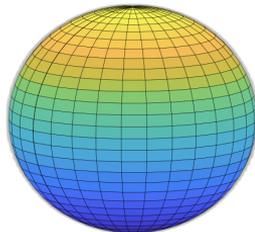
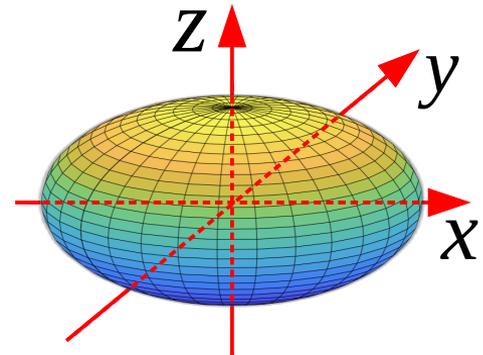
Shape defined by two parameters:

- β – Magnitude
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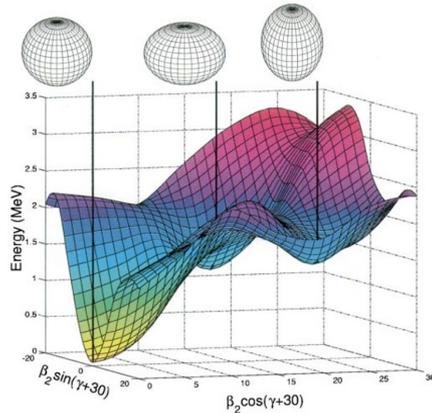
Two experimental observables that allow us to calculate β and γ :

- Reduced transition probability, $B(E2)$
- Spectroscopic quadrupole moment, Q_s

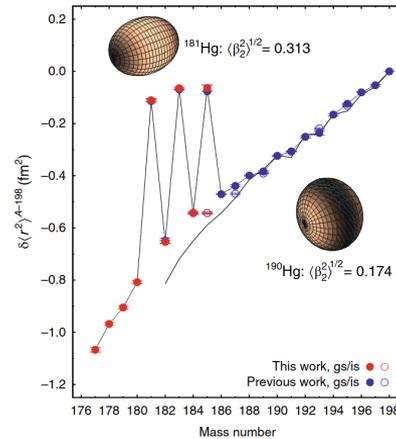
(Usually) measure these observables for the first 2^+ state (which can be constructed using the quadrupole operator)



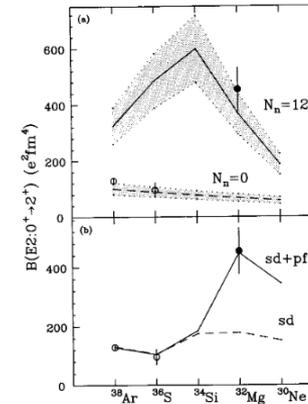
Shape coexistence



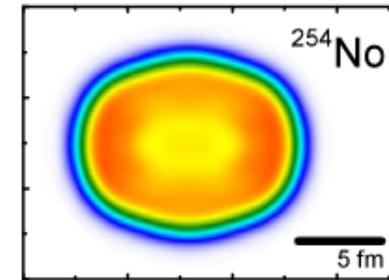
Phase transitions



Shell evolution



Deformation in superheavy nuclei

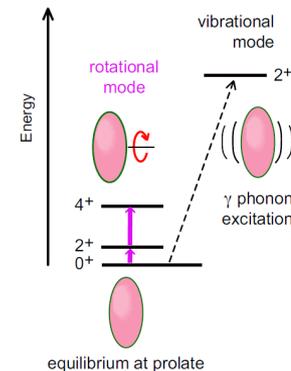


Exotic nuclear shapes – e.g. octupole deformation

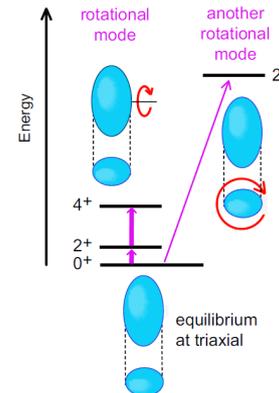


Different modes of excitation

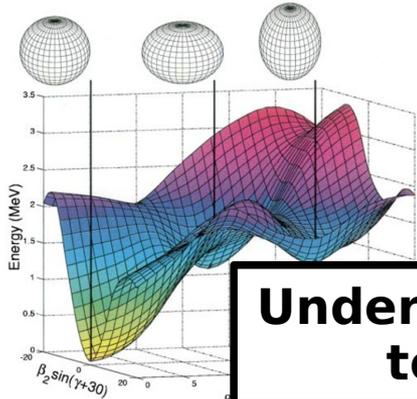
a. conventional picture (prolate)



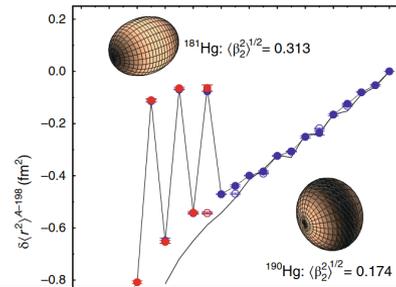
b. present picture (triaxial)



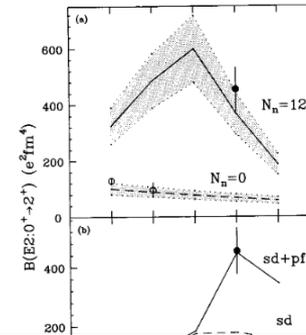
Shape coexistence



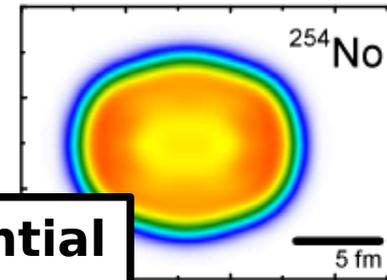
Phase transitions



Shell evolution



Deformation in superheavy nuclei

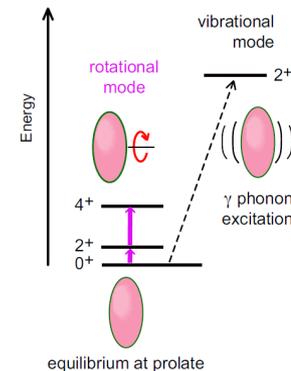


Understanding self-organisation of nuclei is essential to solving the quantum many-body problem

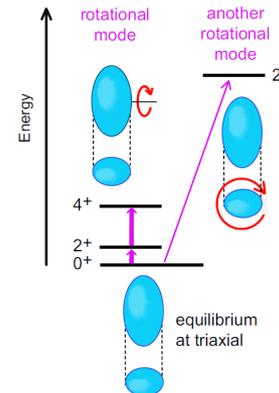
Exotic nuclear shapes – e.g. octupole deformation



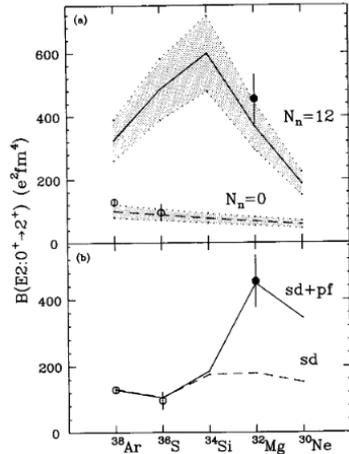
a. conventional picture (prolate)



b. present picture (triaxial)

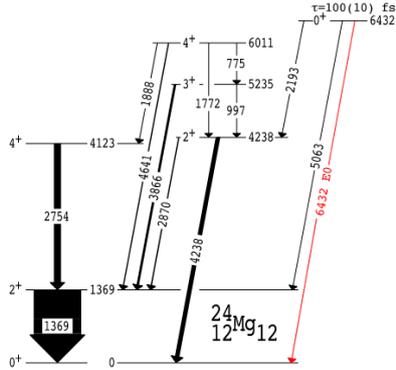


Shell evolution



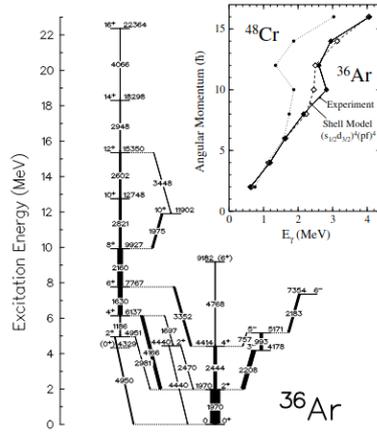
T. Motobayashi, et al., PLB **346** (1995)

Shape coexistence



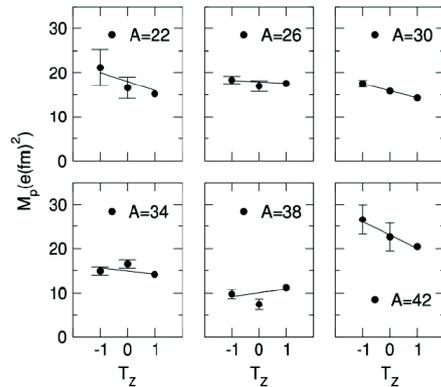
J.T.H. Dowie, et al., PLB, **811** 135855 (2020)

Superdeformation

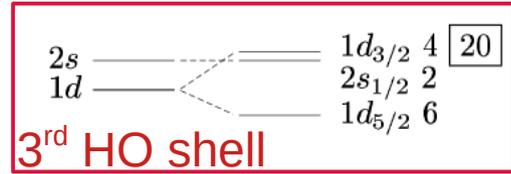
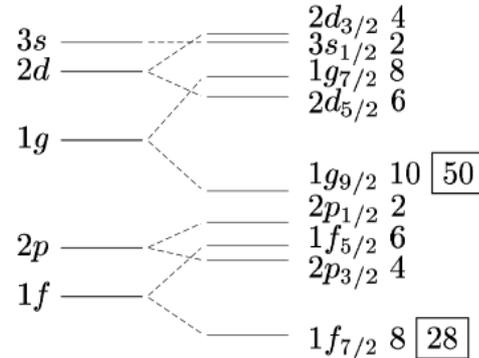


C. E. Svensson, et al., PRL **85** 13 (2000)

Isospin breaking (N=Z)



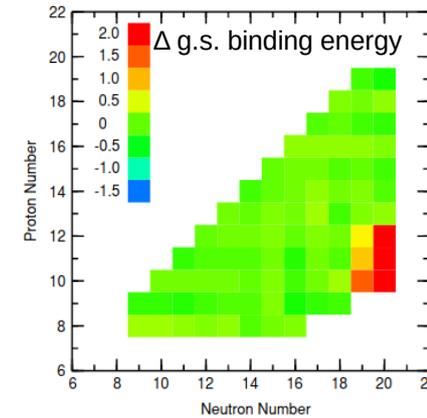
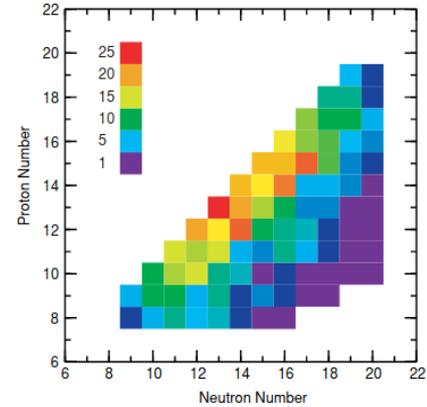
F. M. Prados Estévez, et al., PRC **75** 014309 (2007)



3rd HO shell

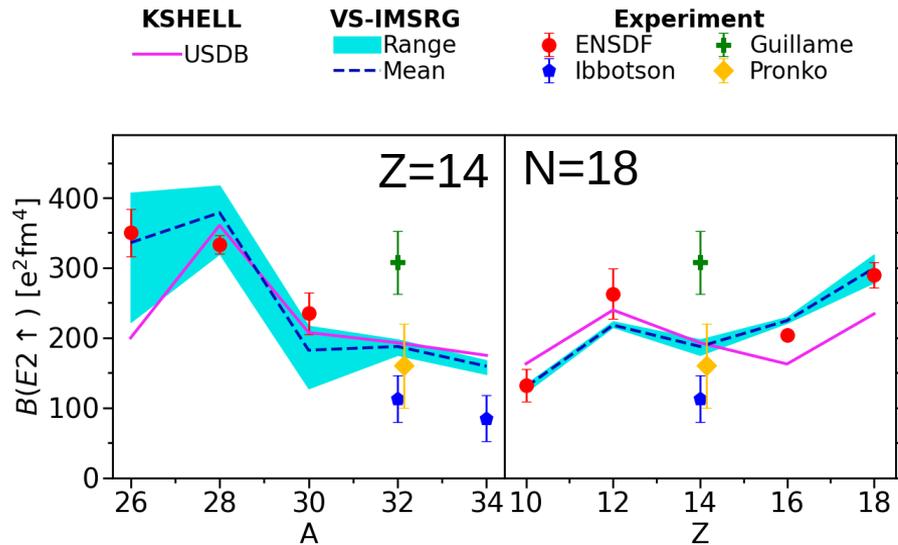


USDB



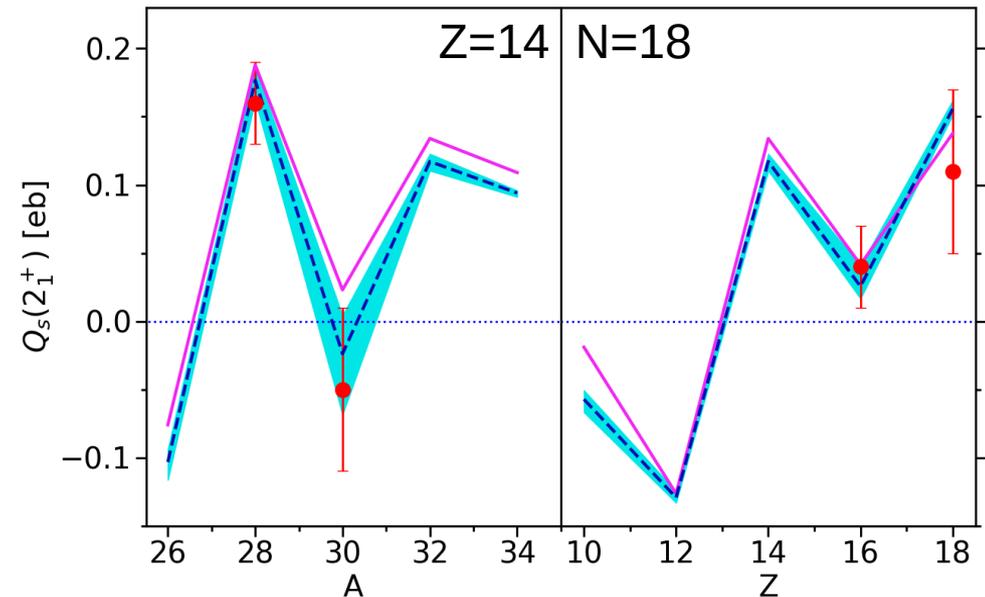
B. A. Brown and W. A. Richter, PRC **74** 034315 (2006)

Disagreement between previous B(E2) values



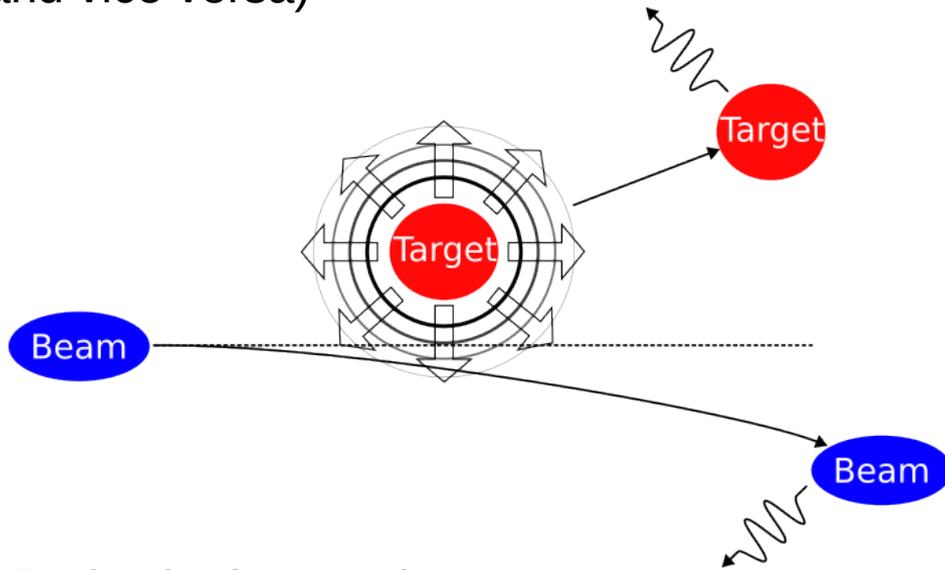
USDB shell model calculations do relatively good job in the region

No previous measurement of spectroscopic quadrupole moment, Q_s



Coulomb excitation

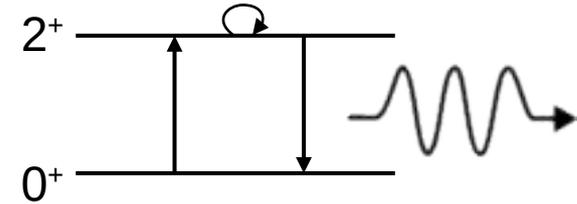
Beam nucleus is present in the electric field of target
(and vice versa)



Rutherford scattering

$$\frac{d\sigma}{d\Omega} = \left(\frac{zZe^2}{4\pi\epsilon_0}\right)^2 \left(\frac{1}{4T_a}\right)^2 \frac{1}{\sin^4 \frac{\theta}{2}}$$

Both nuclei are inelastically excited
through the Coulomb potential



We carefully choose the
energy and scattering angle to
suppress nuclear excitation

Cline criterion – 5 fm between
nuclear surfaces

To first order: $P \propto |\langle 0^+ || E2 || 2^+ \rangle|^2$

Or...

The B(E2) value

Coulomb excitation

We can also access Q_s from angular distribution for state population

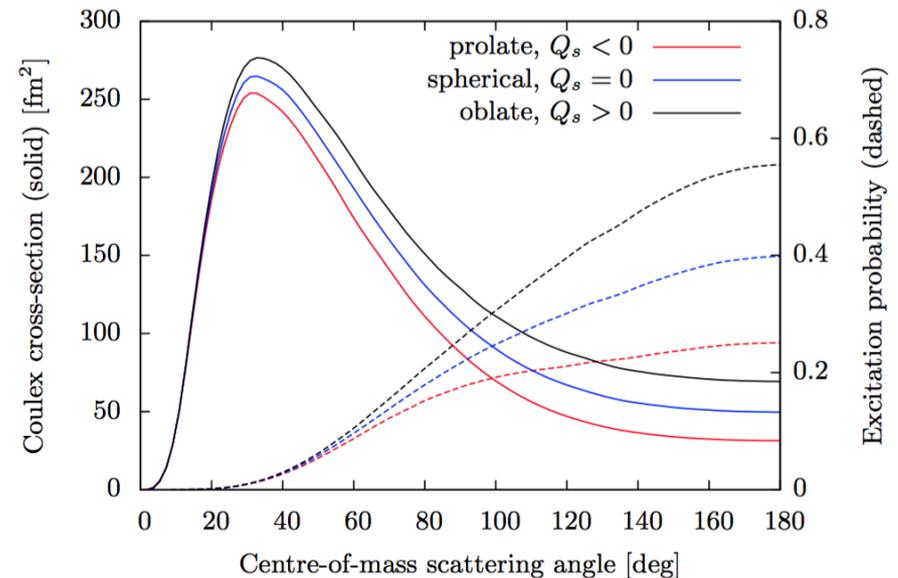
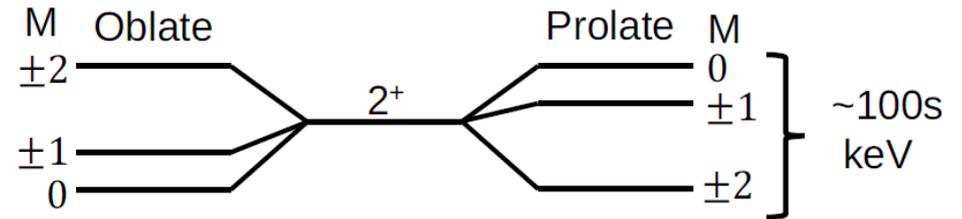
The reorientation effect:

Nuclei reorient in electric field gradient to minimise their energy

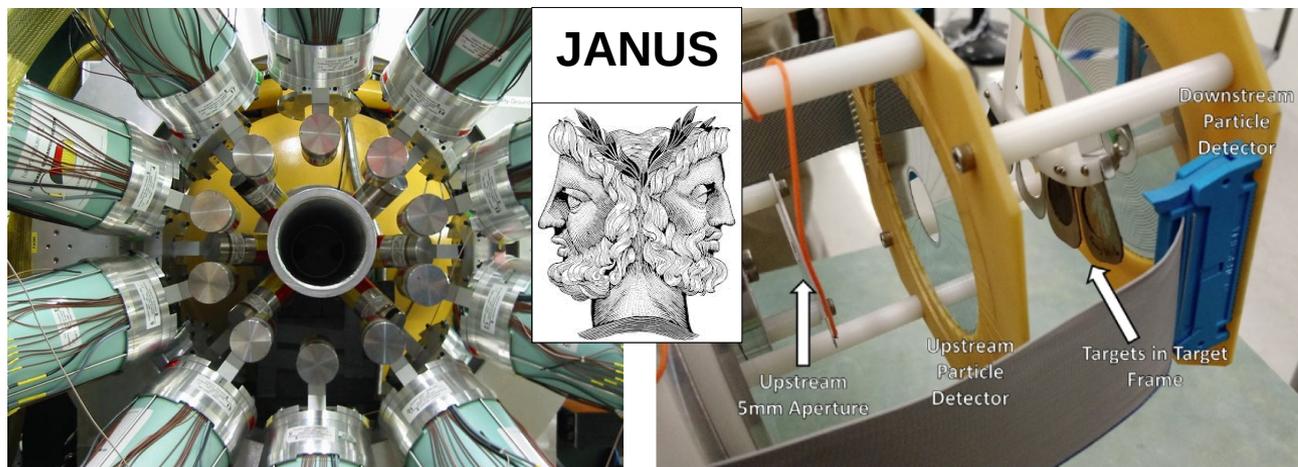
Typical Coulomb excitation reaction: $dV/dr = 10^{30}$ V/cm

$$E(t) \propto eQ_s Z/r^3(t)$$

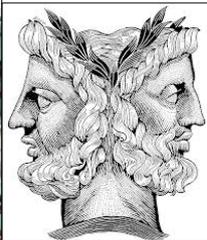
Breaking of m-state degeneracy depends on Q_s



M. Zielinska et al. Eur. J. Phys. A **52** 99 (2016)



JANUS



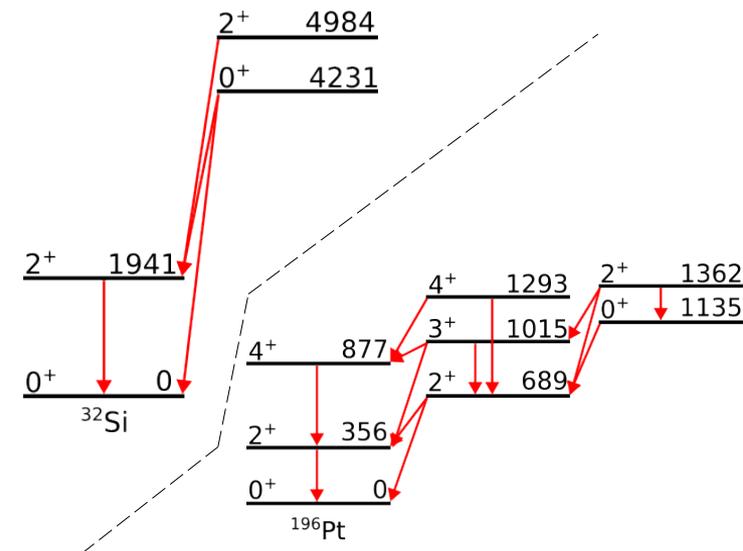
- “Safe” Coulomb excitation of ^{32}Si beam on a ^{196}Pt target.
- Beam energies: 3.57 and 3.48 MeV/u
- Beam intensity: 10^6 pps
- Target thickness: 1 mg/cm^2

SeGA

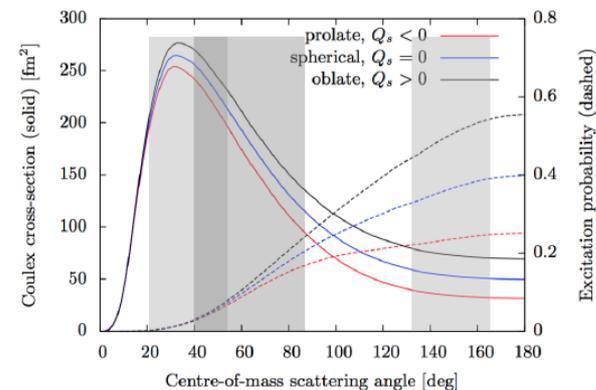
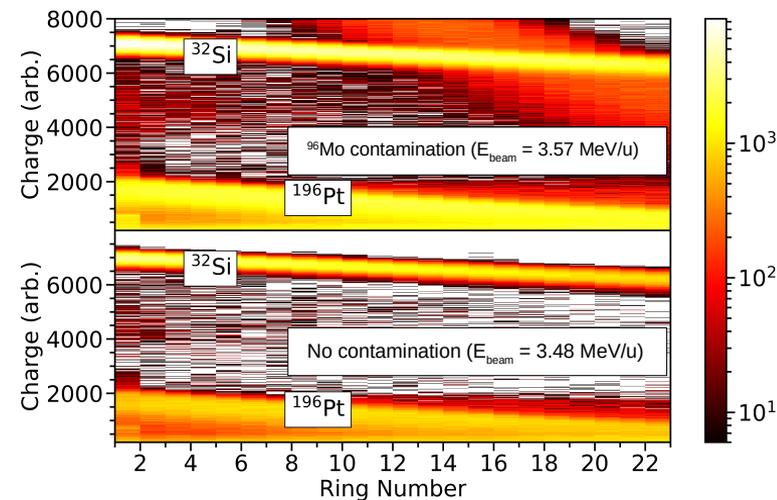
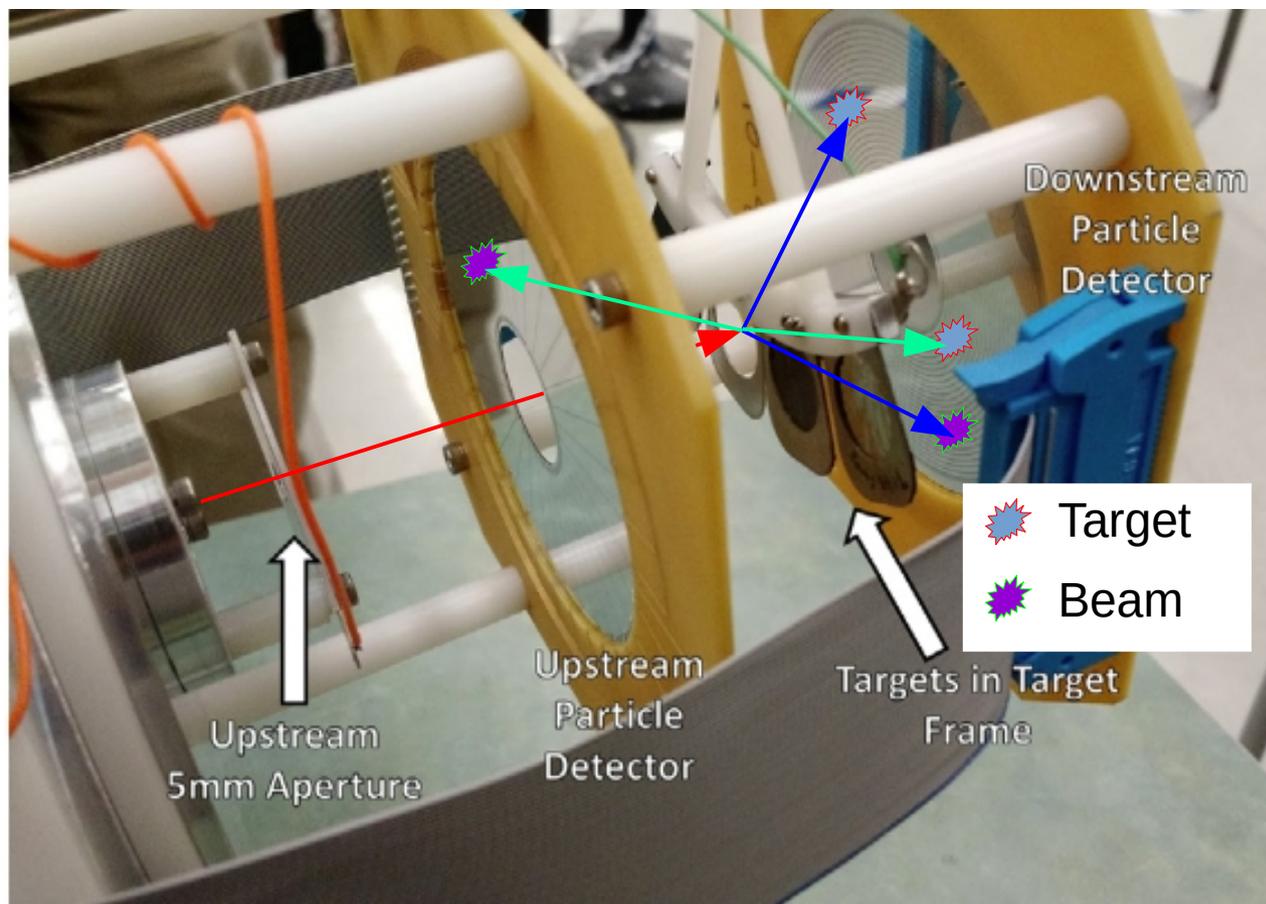
- Array of 16 32-fold segmented germanium crystals.
- Energy resolution of ~ 2.5 keV FWHM at 1332 keV.
- Efficiency of $\sim 6\%$ at 1 MeV.

Bambino 2

- 2 annular double-sided silicon detectors (S3).
- **Upstream** angular range,
 θ : $130.6^\circ - 159.9^\circ$
- **Downstream** angular range,
 θ : $20.1^\circ - 49.4^\circ$
- Angular resolution: 1.5° in θ and 3° in ϕ

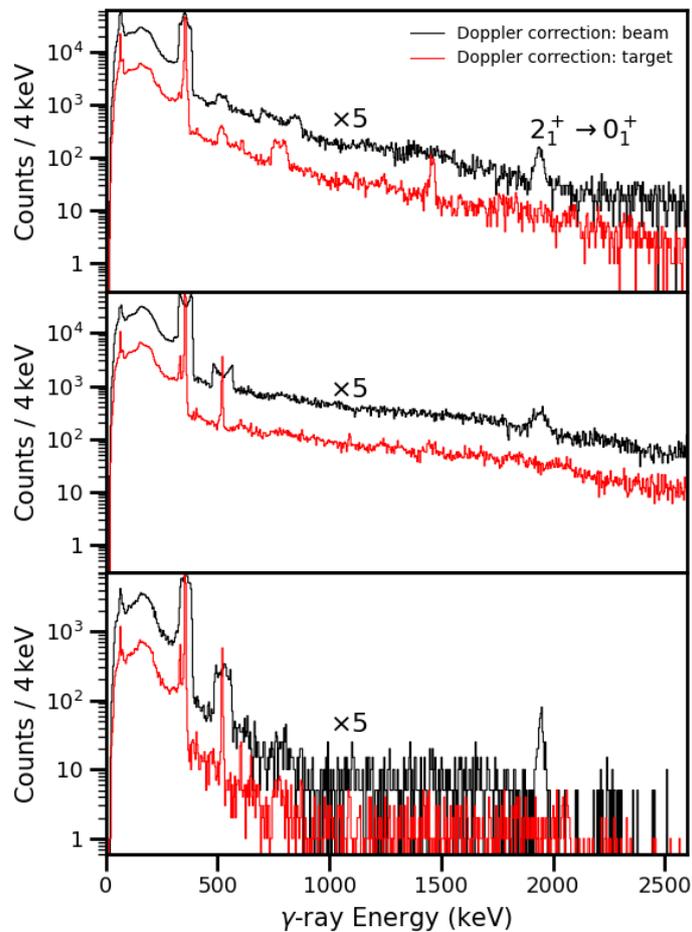


Experiment at NSCL



M. Zielinska et al. Eur. J. Phys. A **52** 99 (2016)

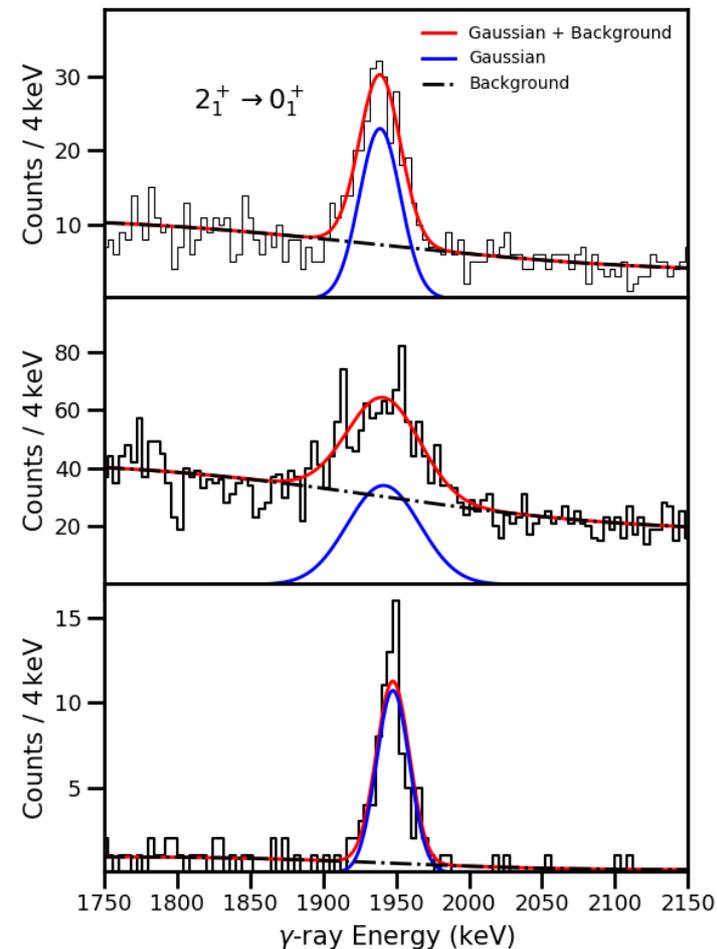
Experiment at NSCL



^{32}Si detected downstream

^{196}Pt detected downstream

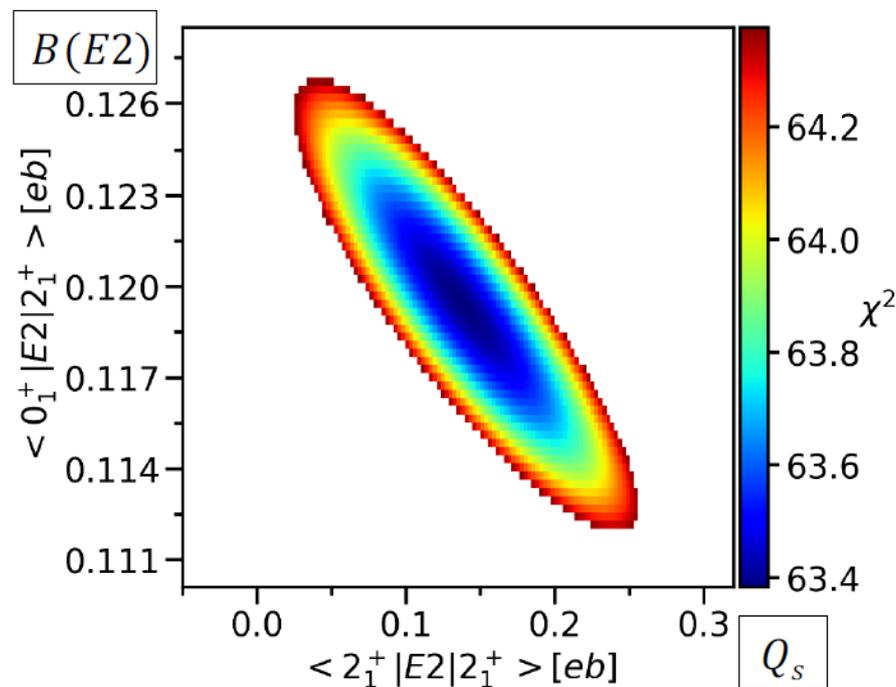
^{32}Si detected upstream



Yields evaluated using GOSIA.

χ^2 minimization of matrix elements performed using the MIGRAD algorithm in the ROOT MINUIT library.

Simultaneous fitting of ^{196}Pt matrix elements accounts for systematic errors.

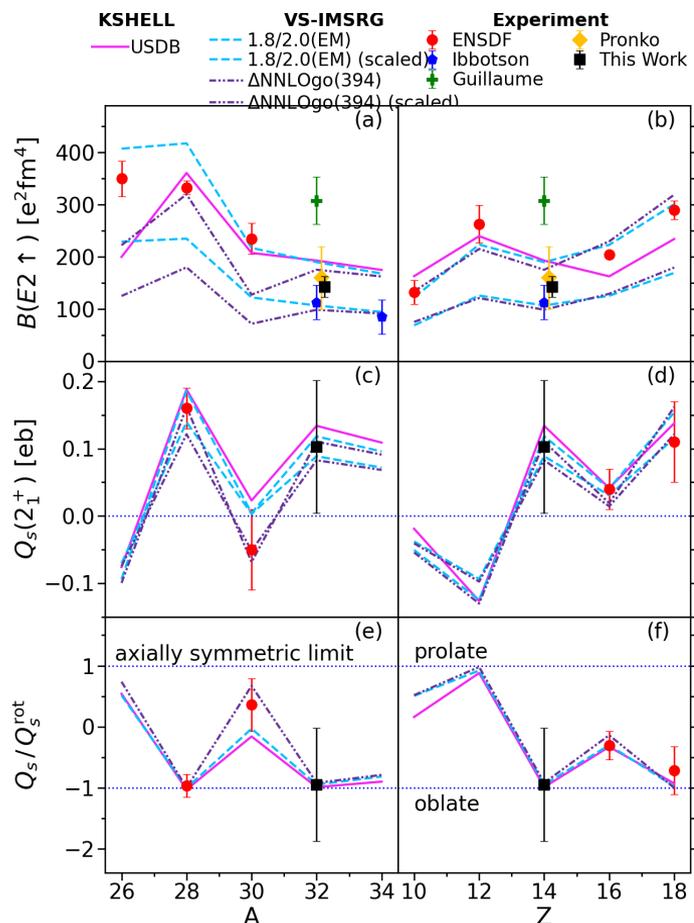


Final Result

$$B(E2; 0_1^+ \rightarrow 2_1^+) = 135(19) \text{ e}^2\text{fm}^4 = 22(3) \text{ W.u.}$$

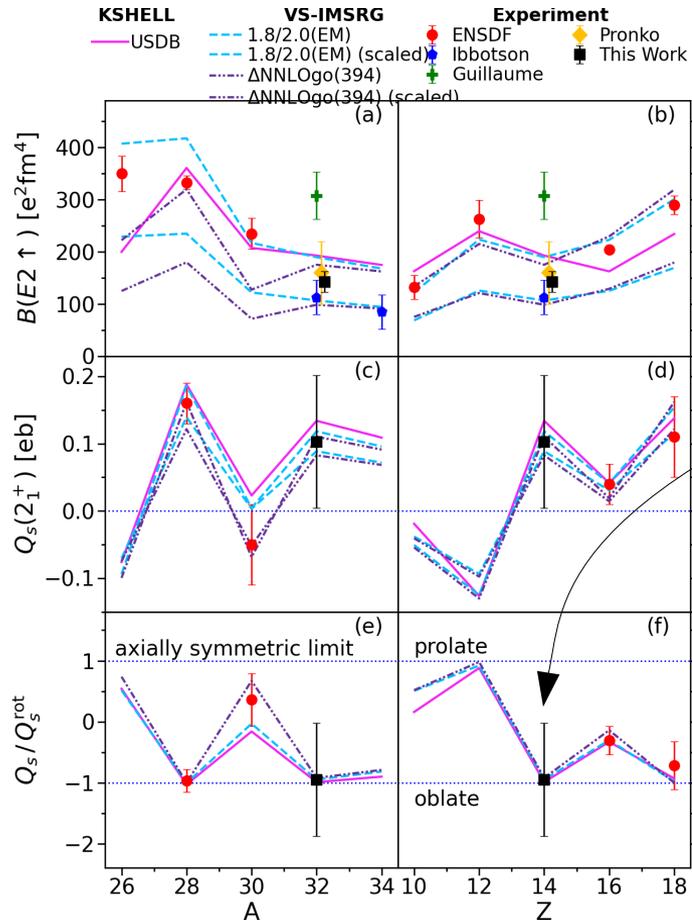
$$Q_s(2_1^+) = 0.14(8) \text{ eb}$$

How do the results compare...



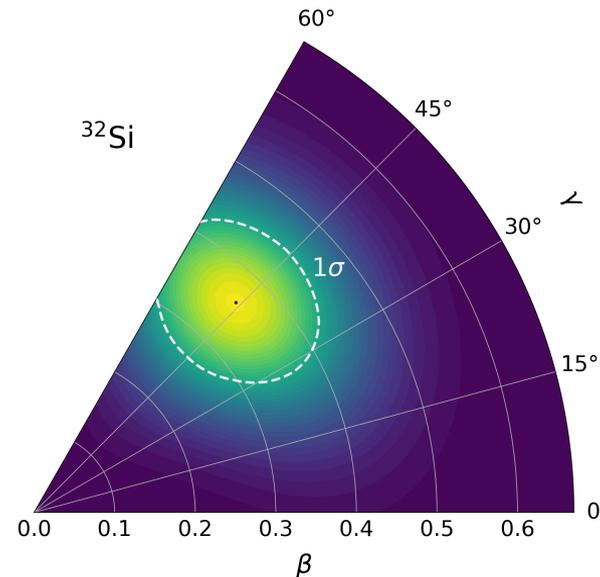
- Result for $Q_s(2_1^+)$ compares well to USDB (although note large errors)
- $B(E2)$ value is significantly overestimated by USDB for ^{32}Si and ^{34}Si
- USDB is reproducing shape well, but underpredicts the magnitude of deformation
- Ab-initio VS-IMSRG compares well... but note there is a truncation to the evolution of the electromagnetic operators at the two-body level [IMSRG(2) approximation]
- Comparison of data for several sd-shell nuclei shows calculations underpredict E2 matrix elements by $\approx 25\%$ [S. R. Stroberg *et al.*, PRC **105**, 034333 (2022)]
- Scaled results are similar to USDB

How do the results compare...

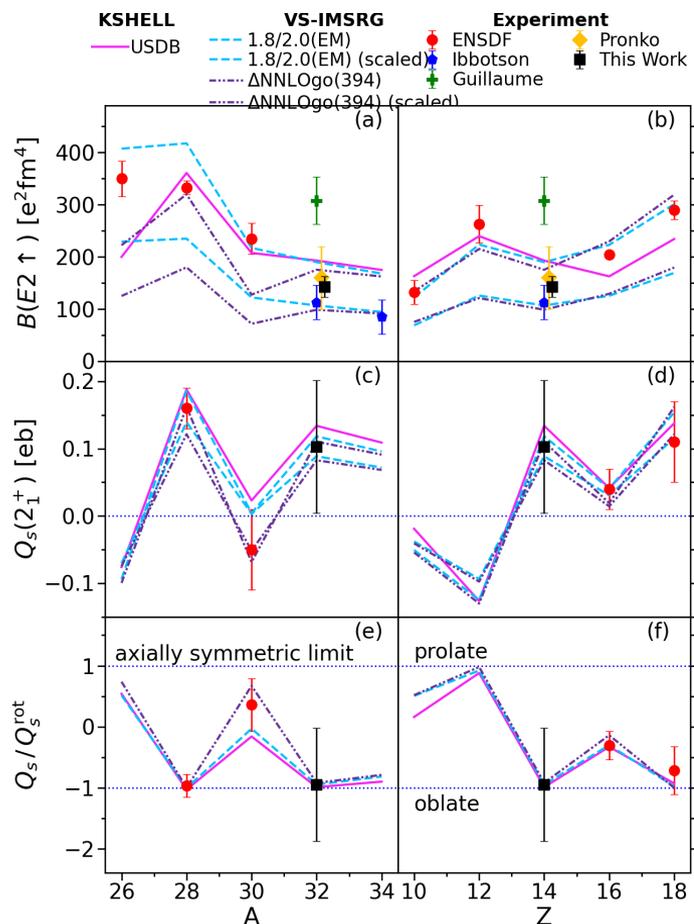


We can approximately calculate the γ parameter

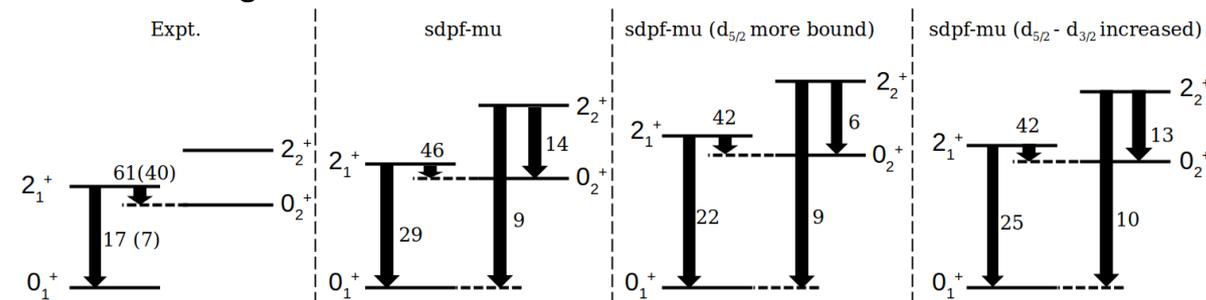
$$\cos(3\gamma) \approx \cos(3\gamma)_{2_1^+} = -\frac{Q_s(2_1^+)}{\frac{2}{7}\sqrt{\frac{16\pi}{5}} \cdot B(E2; 0_1^+ \rightarrow 2_1^+)}$$



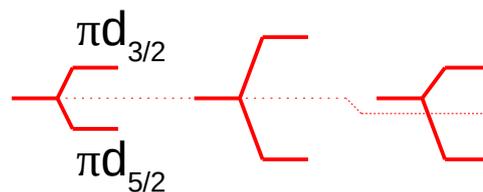
Full calculation using Kumar-Cline quadrupole invariants - exact within USDB model space



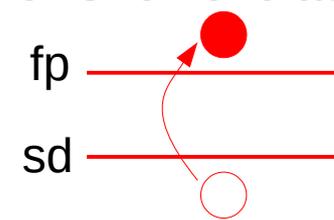
What do things look like in ^{34}Si ?



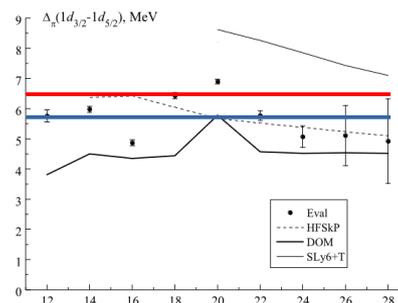
Changes to $\pi d_{5/2}$, $\pi d_{3/2}$ shells



Out-of-shell excitations



Some evidence for this at ^{32}Si and ^{34}Si
 [O.V. Bespalova, *et al.*, *Eur. Phys. J. A* **54**, 2 (2018).]



^{32}Si – Inhibited quadrupole deformation

- Nuclear deformation in ^{32}Si has been investigated through “safe” Coulomb excitation at NSCL, MI, USA
- $B(E2; 0_1^+ \rightarrow 2_1^+) = 135(19) \text{ e}^2\text{fm}^4$, $Q_s(2_1^+) = 0.14(8) \text{ eb}$
- Phenomenological and ab-initio calculations both reproduce oblate structure but overpredict the scale of deformation
- There is a reduced role of out-of-space excitations (core polarisation)

Collaborators:

J. Henderson¹, C. R. Hoffman², A. M. Hill³, B. Hu⁴, J. D. Holt^{4,5,6}, T. Beck^{3,6}, C. Cousins¹, P. Farris^{3,6}, A. Gade^{3,6}, S. A. Gillespie⁶, H. Iwasaki^{3,6}, S. Kisyov⁷, A. Kuchera⁸, B. Longfellow⁷, C. Müller-Gatermann², E. Rubino⁶, R. Russell¹, R. Salinas^{3,6}, A. Sanchez^{3,6}, D. Weisshaar⁶, C. Y. Wu⁷, J. Wu⁶

¹University of Surrey, Guildford GU2 7XH, UK

²Physics Division, Argonne National Laboratory, Lemont IL, USA

³Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan 48824 USA

⁴TRIUMF, Vancouver, BC V6T 2A3, Canada

⁵Department of Physics, McGill University, 3600 Rue University, Montréal, QC H3A 2T8, Canada

⁶National Superconducting Cyclotron Laboratory, Michigan State University, East Lansing Michigan 48824 USA

⁷Lawrence Livermore National Laboratory, Livermore, California 94550, USA

⁸Department of Physics, Davidson College, Davidson, North Carolina 28035, USA