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Composite Higgs

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HIGGS 2024

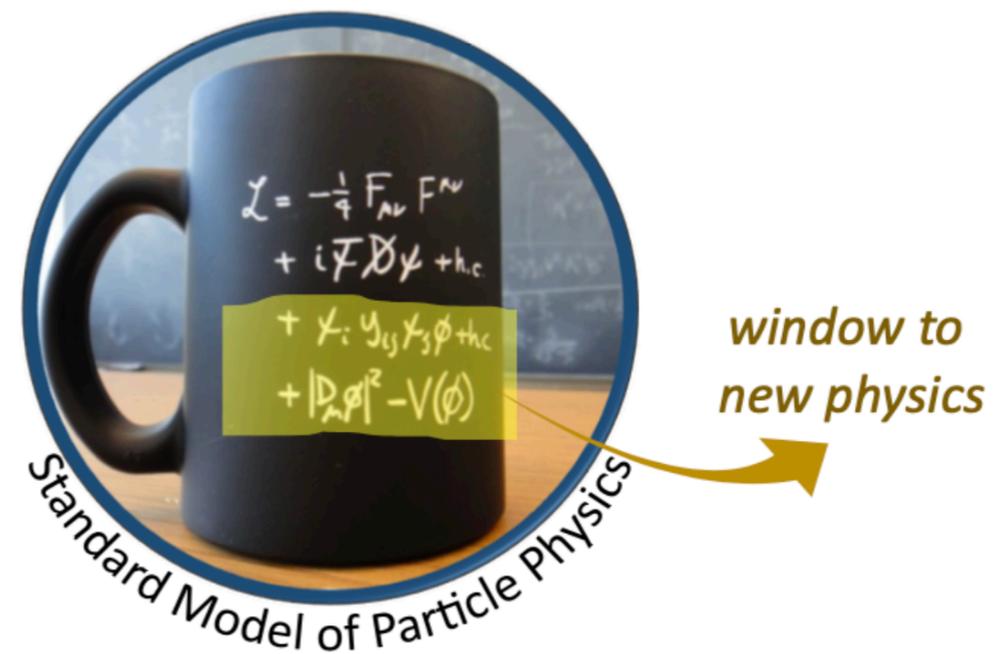
Uppsala, Sweden — 4-8 November 2024

Open fundamental questions

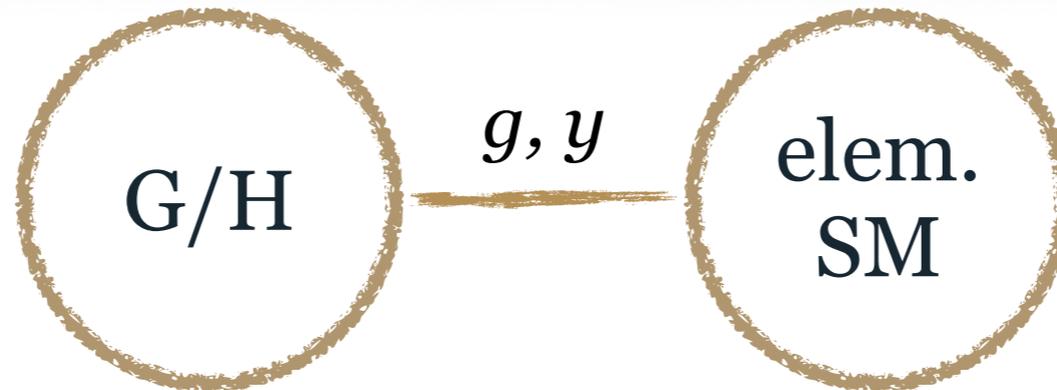
- ☑ Is it the SM Higgs?
- ☑ Is it small mass "natural"?
- ☑ Is it an elementary or composite particle?
- ☑ Is it unique?
- ☑ Is it the first supersymmetric particle ever observed?
- ☑ Is it the only responsible for the masses of all the elementary particles?
- ☑ Is it a portal to a hidden world?



The SM is a **"partial" description of the Nature**, it could be **part of a more general theory** which will manifest itself at energies higher than the ones explored till now



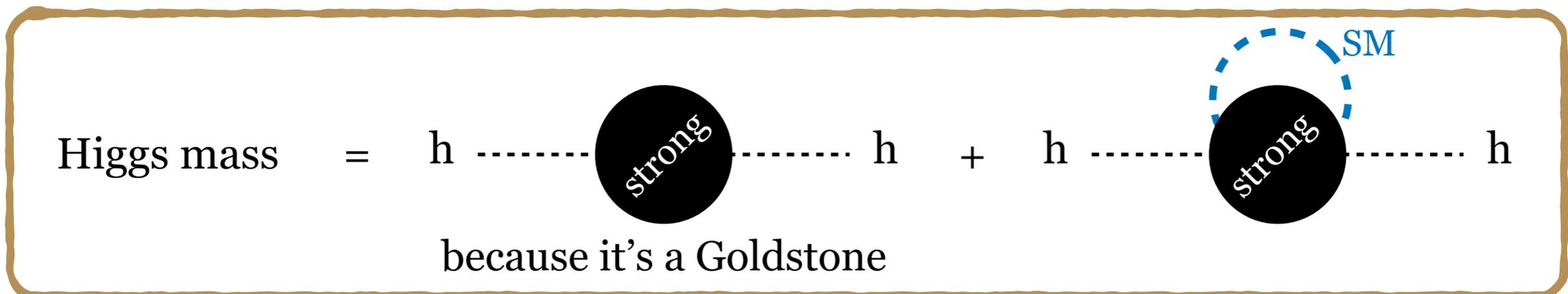
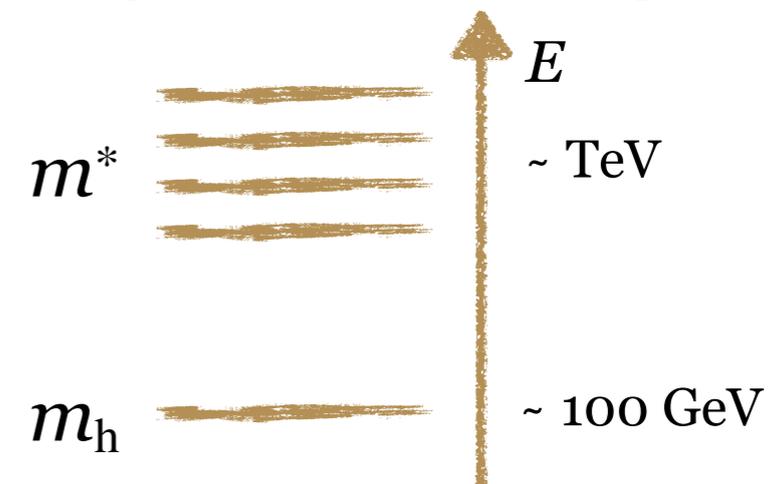
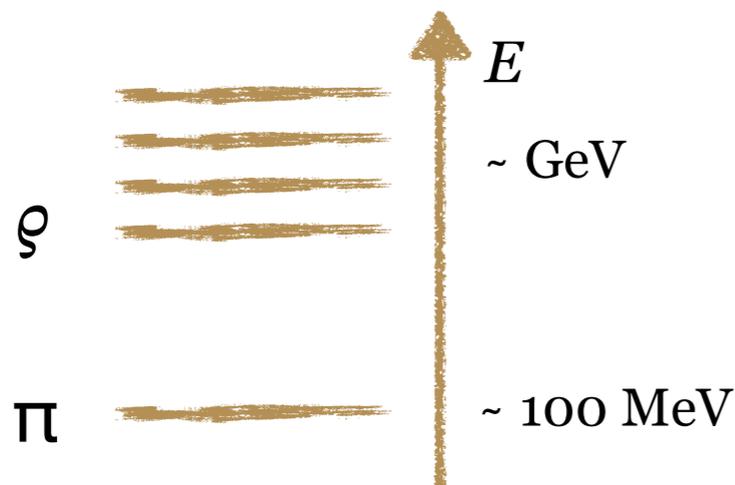
Higgs as a Composite pseudo Nambu Goldstone Boson



The basic idea:
 Higgs as Goldstone boson of G/H in a strong sector (Georgi, Kaplan '80s)

Inspired by QCD where one observes that the (pseudo) scalars are the lightest states

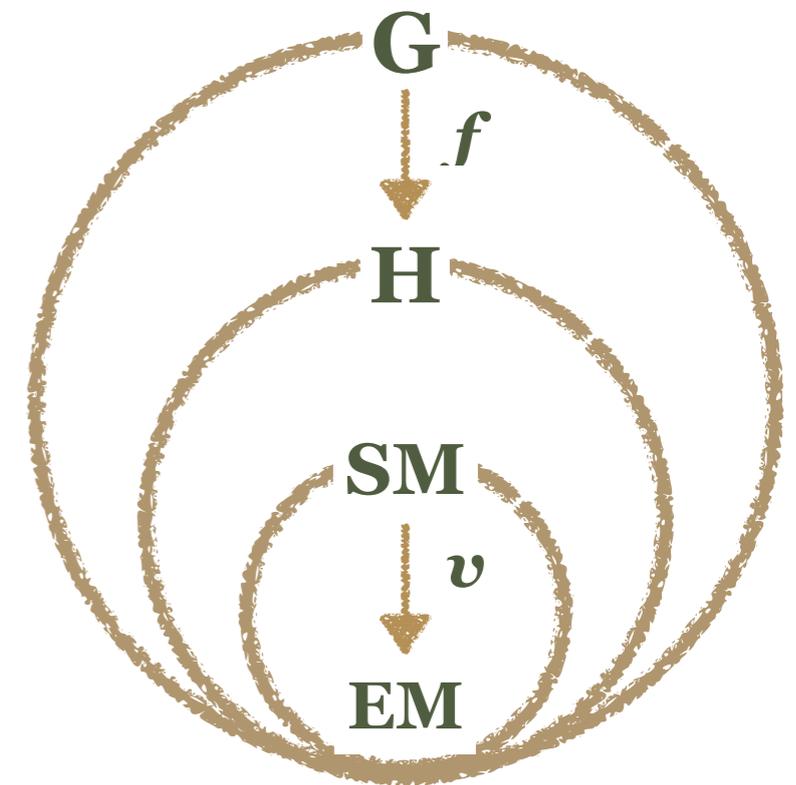
the Higgs could be a kind of pion arising from a new strong sector



Mass protected by the global QCD symmetry

Basic rules for explicit composite pNGB models

- ☑ Need to choose the correct $G \rightarrow H$ (spontaneous) breaking at f ($\sim \text{TeV}$) to have the required NGBs (≥ 4)
- ☑ Need to break H (explicitly, so pNGBs) via g_0 (gauge) and Y (Yukawa) couplings to generate the one-loop effective potential for EWSB
- ☑ Need to include new composite resonances from the confining strong dynamics



Composite Higgs Models

Elementary sector

$$A_\mu, \psi \in SU(2) \times U(1)$$

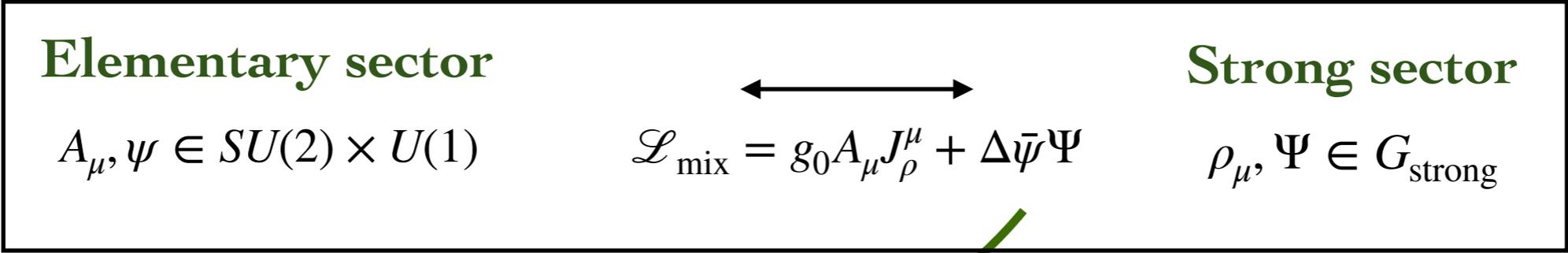


$$\mathcal{L}_{\text{mix}} = g_0 A_\mu J_\rho^\mu + \Delta \bar{\psi} \Psi$$

Strong sector

$$\rho_\mu, \Psi \in G_{\text{strong}}$$

Composite Higgs Models

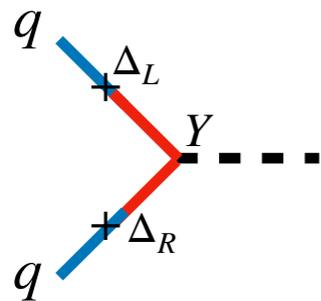


partial compositeness

Linear elementary-composite fermion mixings

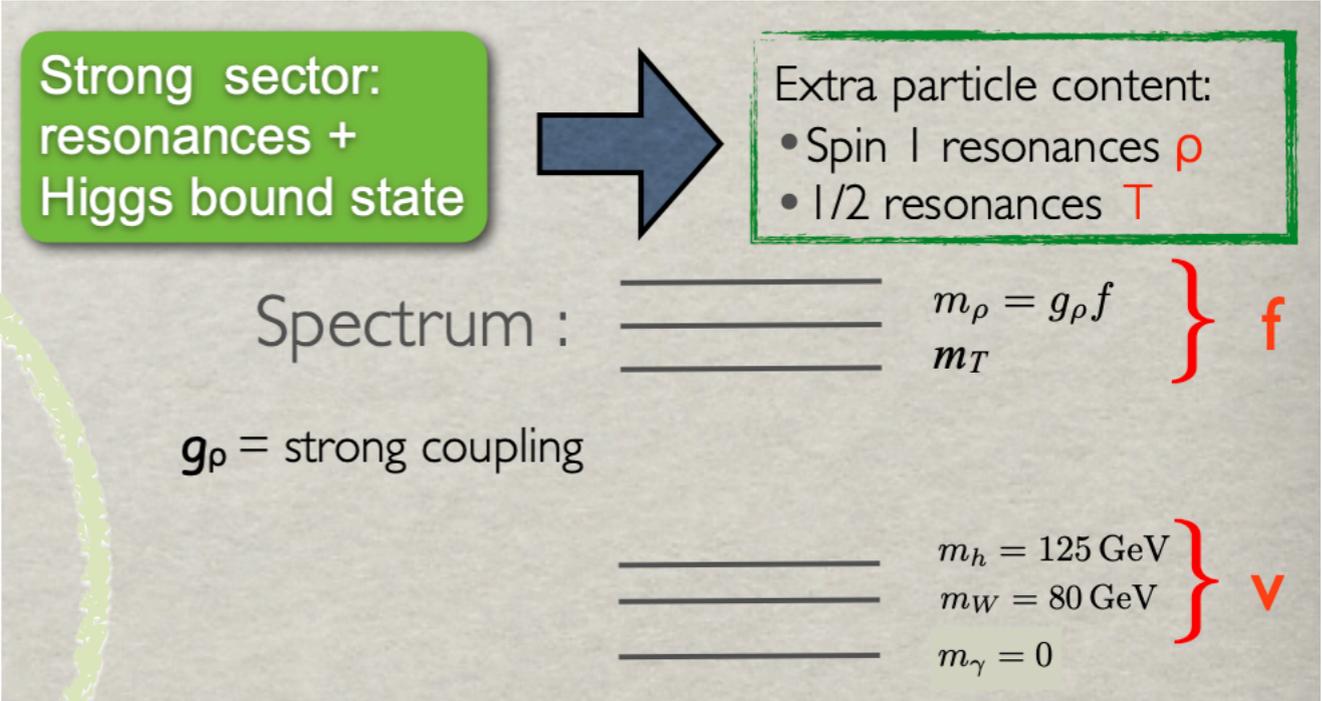
→ for the 3rd generation quarks

$$\Delta_R \bar{q}_R \mathcal{O}_L + \Delta_L \bar{q}_L \mathcal{O}_R + Y \bar{\mathcal{O}}_L H \mathcal{O}_L$$



top Yukawa couplings

$$m_t \sim \frac{v}{\sqrt{2}} \frac{\Delta_{t_L} \Delta_{t_R} Y_T}{m_T m_{\tilde{T}} f}$$



SM hierarchies are generated by the mixings:
light quarks mostly elementary, top mostly composite

Extended Composite Higgs Models

Models with a larger Higgs structure with respect to the SM have been largely discussed
Supersymmetry requires two Higgs doublets with specific Yukawa and potential terms
 2HDMs offer a rich phenomenology in EW and flavour physics

Look for a pNGB realisation of extended Higgs scenarios

The structure of the Higgs sector is determined by the **coset G/H**

G	H	PGB
SO(5)	SO(4)	4=(2,2)
SO(6)	SO(5)	5=(2,2)+(1,1)
SO(6)	SO(4)×SO(2)	8=(2,2)+(2,2)
SO(7)	SO(6)	6=(2,2)+(1,1)+(1,1)
	G ₂	7=(1,3)+(2,2)
...

→ Minimal = One doublet

→ Doublet + Singlet

→ Two Doublets

New players in the game

Composite 2-Higgs Doublet Model (C2HDM)

J.Mrazek et al. '11; De Curtis, Moretti, Yagyu, Yildirim '16, De Curtis, Delle Rose, Moretti, Yagyu '18

- ☑ EWWSB is driven by 2 Higgs doublets as pNGBs of $SO(6)/SO(4) \times SO(2)$. The unbroken group contains the custodial $SO(4)$
- ☑ Alignment conditions on the strong Yukawa couplings must be imposed to suppress FCNCs (composite version of an Aligned 2HDM Pich, Tuzón, '09)

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- ☑ Alignment conditions on the strong Yukawa couplings must be imposed to suppress FCNCs (composite version of an Aligned 2HDM Pich, Tuzón, '09)
- ☑ The SM fields are linearly coupled to operators of the strong sector and explicitly break its symmetry
A potential for the Higgses is radiatively generated, couplings and masses determined by the strong sector
- ☑ Fermion sector: embed the 3rd generation quarks into $SO(6)$ reps. + linear couplings $\Delta_{L,R}$ between composite and elementary fermions (partial compositeness)

$$\begin{aligned} \mathcal{L}_{\text{mix}} + \mathcal{L}_{\text{strong}} &= \Delta_L^I \bar{q}_L^{\mathbf{6}} \Psi_R^I + \Delta_R^I \bar{t}_R^{\mathbf{6}} \Psi_L^I + h.c. \\ &+ \bar{\Psi}^I i \not{D} \Psi^I - \bar{\Psi}_L^I M_{\Psi}^{IJ} \Psi_R^J - \bar{\Psi}_L^I (Y_1^{IJ} \Sigma + Y_2^{IJ} \Sigma^2) \Psi_R^J \end{aligned} \quad \Sigma = U_1 \Sigma_2 U_1^T$$

- ☑ Two heavy fermions' sextuplets ψ^J needed for an UV finite effective potential $I, J=1,2$

scale of compositeness, f ; linear mixings, $\Delta_L^{1,2}, \Delta_R^{1,2}$; Yukawas, $Y_{1,2}^{IJ}$; heavy fermion mass parameters, M_{Ψ}^{IJ} , $I, J = 1, 2$ (partial compositeness for the top)

2-Higgs Doublets as pNGBs

Aligned 2HDM realised in a composite scenario

- Same physical Higgs states as in the elementary 2HDM: h, H, A, H^\pm (h =SM-like Higgs)

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- CP-even states: h, H

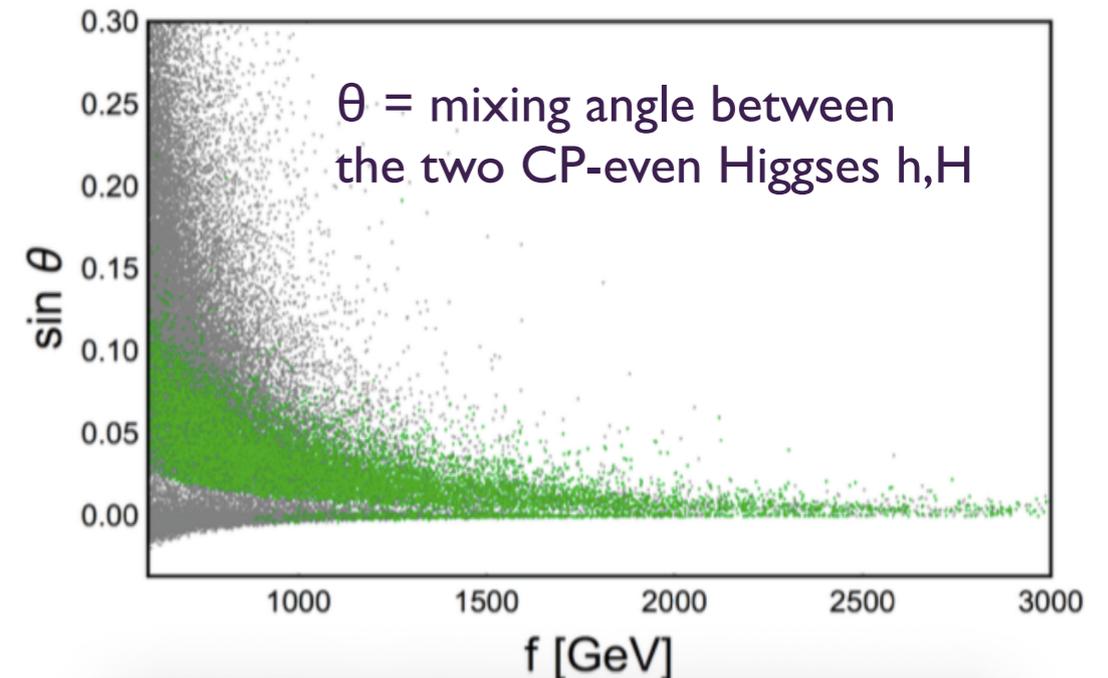
$$m_h \sim v \quad m_H \sim f + O(v) \quad \xi = v^2/f^2$$

θ is predicted to be small: $O(\xi)$ for large f

- CP-odd states: A, H^\pm

$$m_A \sim m_{H^\pm} \sim f + O(v)$$

$f \rightarrow \infty$ SM limit
 H, A, H^\pm decouple
 $h \rightarrow h^{\text{SM}}$



green points satisfy the bounds from direct and indirect Higgs searches

tested against HiggsBounds and HiggsSignals

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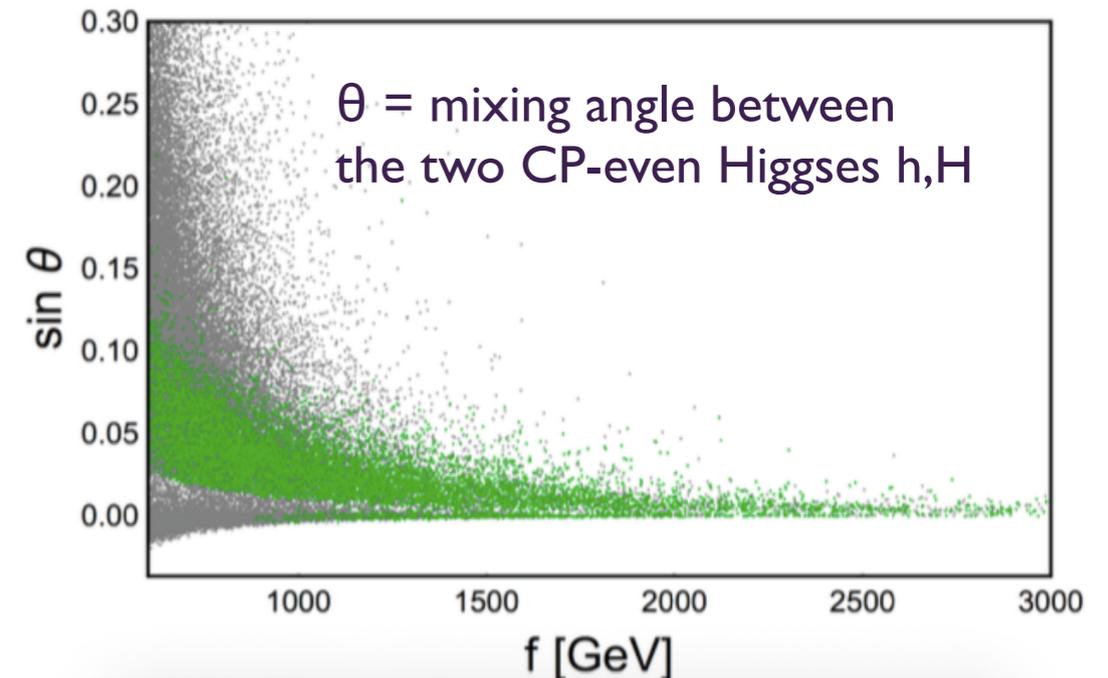
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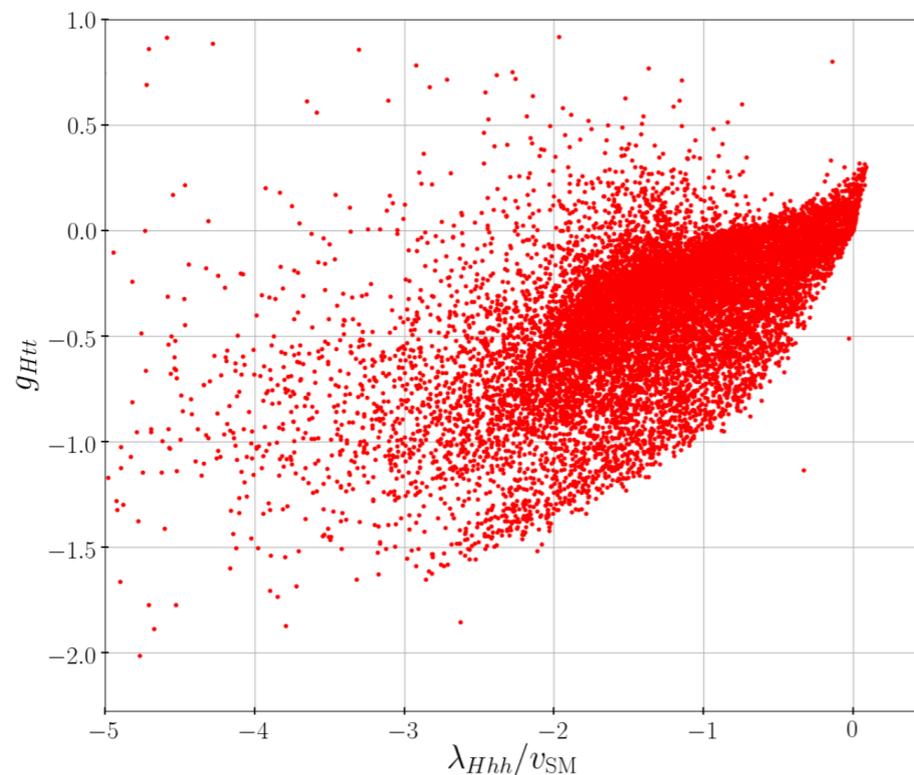
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in the C2HDM the Higgs sector parameters are correlated and carry the imprint of compositeness

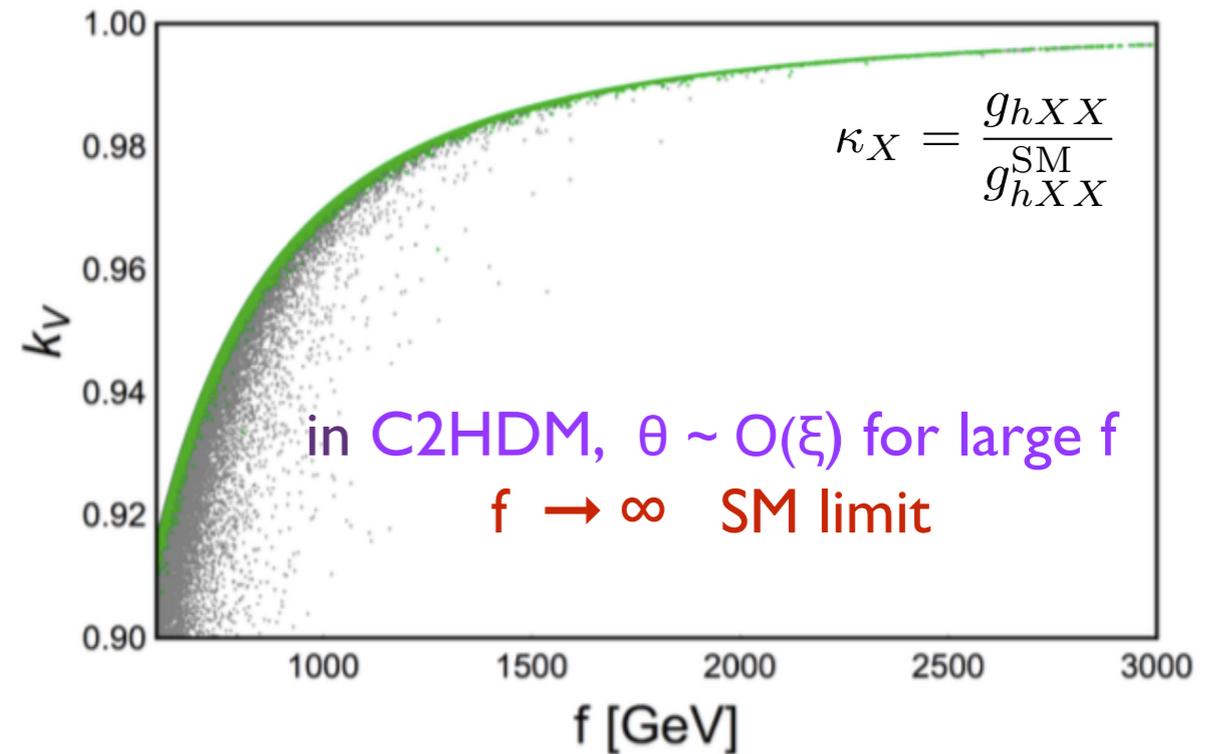
→ Ex: H_{tt} and H_{hh}

C2HDM - facing the data

- **h couplings to SM particles:**

dictated by symmetries (as in QCD chiral Lagrangian) Ex: **corrections of order ξ** to the **hVV** couplings. Also modified by the mixing angle θ

$$\mathbf{k_V} \simeq (1 - \xi/2) \cos\theta \quad V=W,Z \quad \xi = v^2/f^2$$



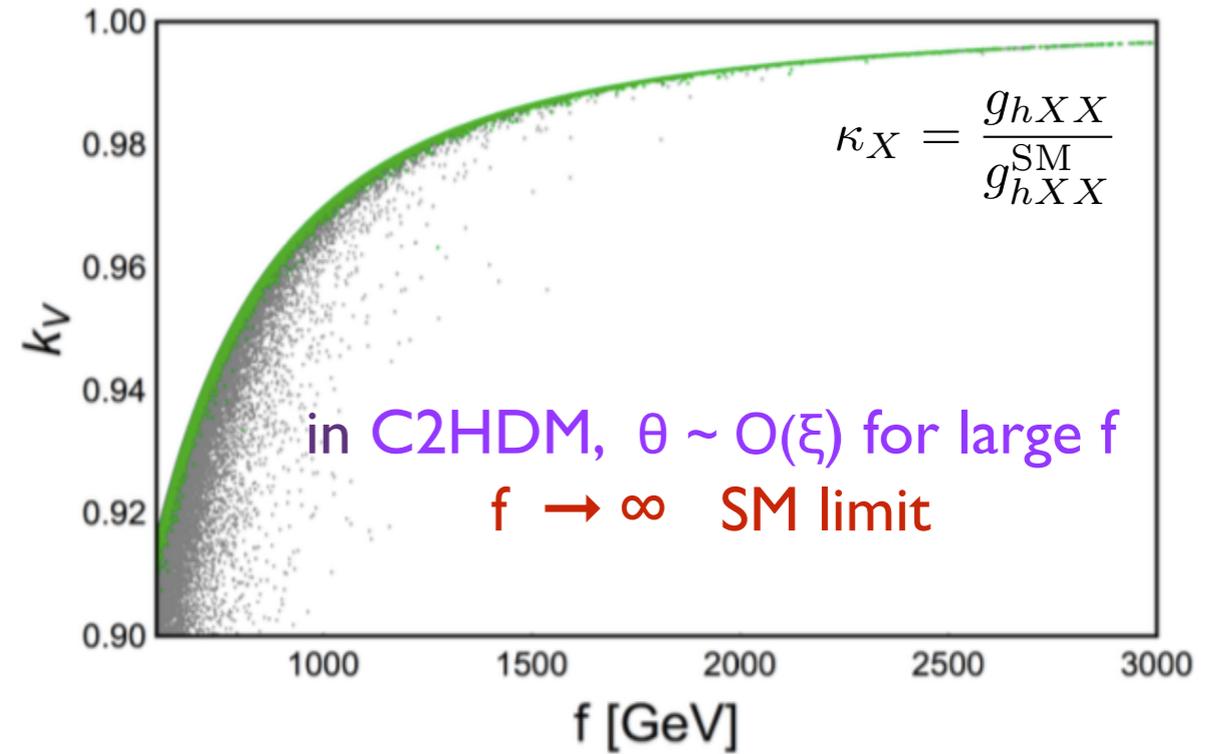
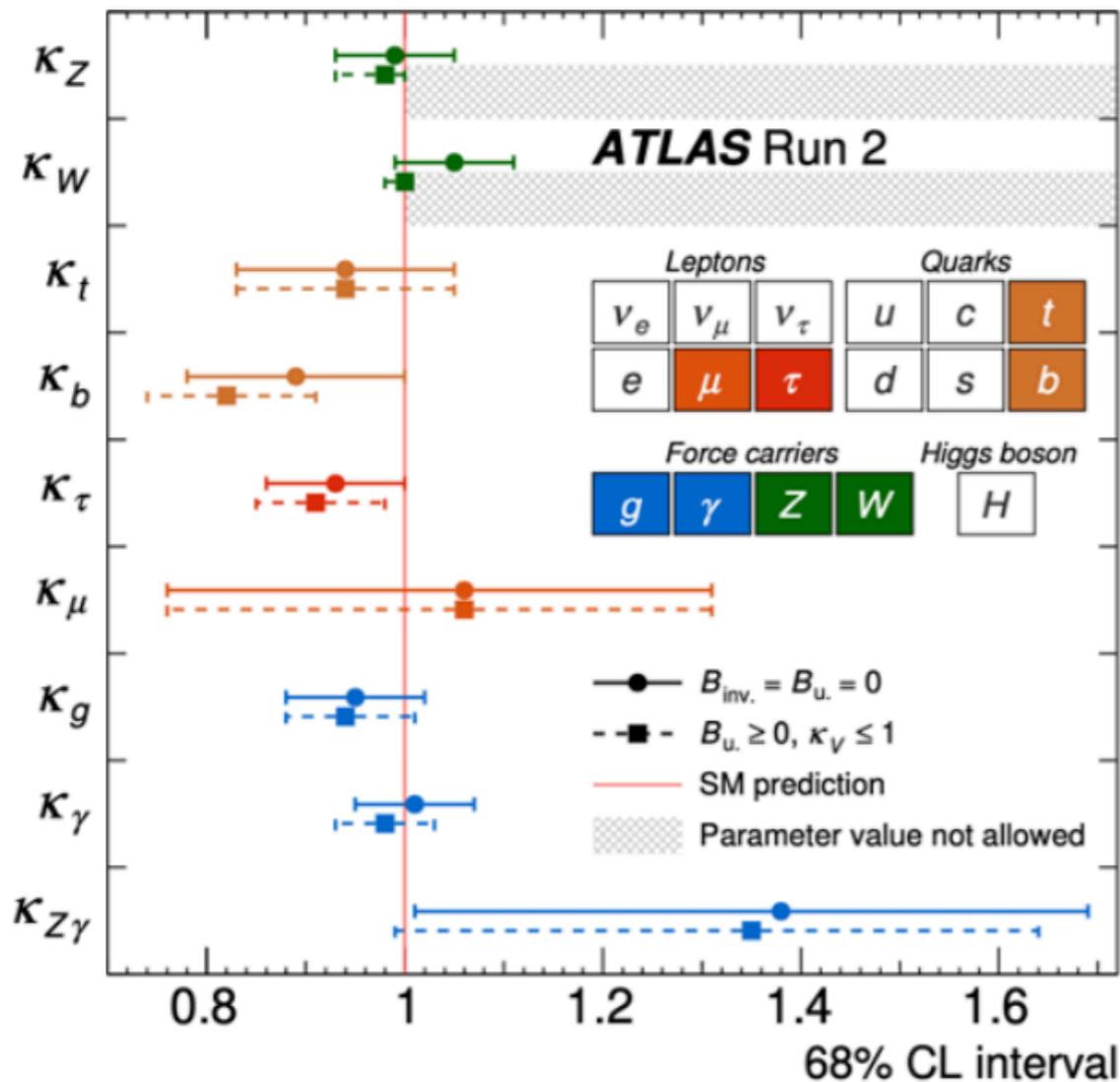
green points satisfy the present bounds

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green points satisfy the present bounds

NOW: the Higgs couplings are constrained at 10-20% level

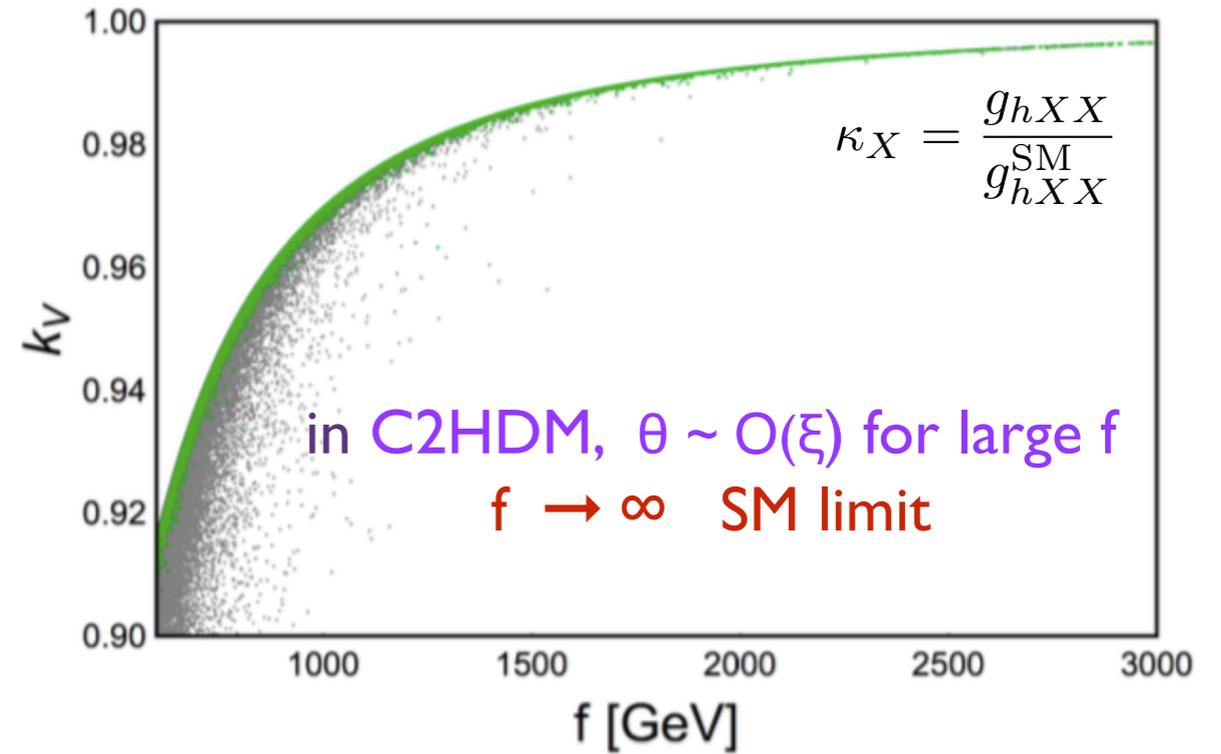
$$\xi \leq 0.1 \quad f \geq 750 \text{ GeV}$$

C2HDM - facing the data

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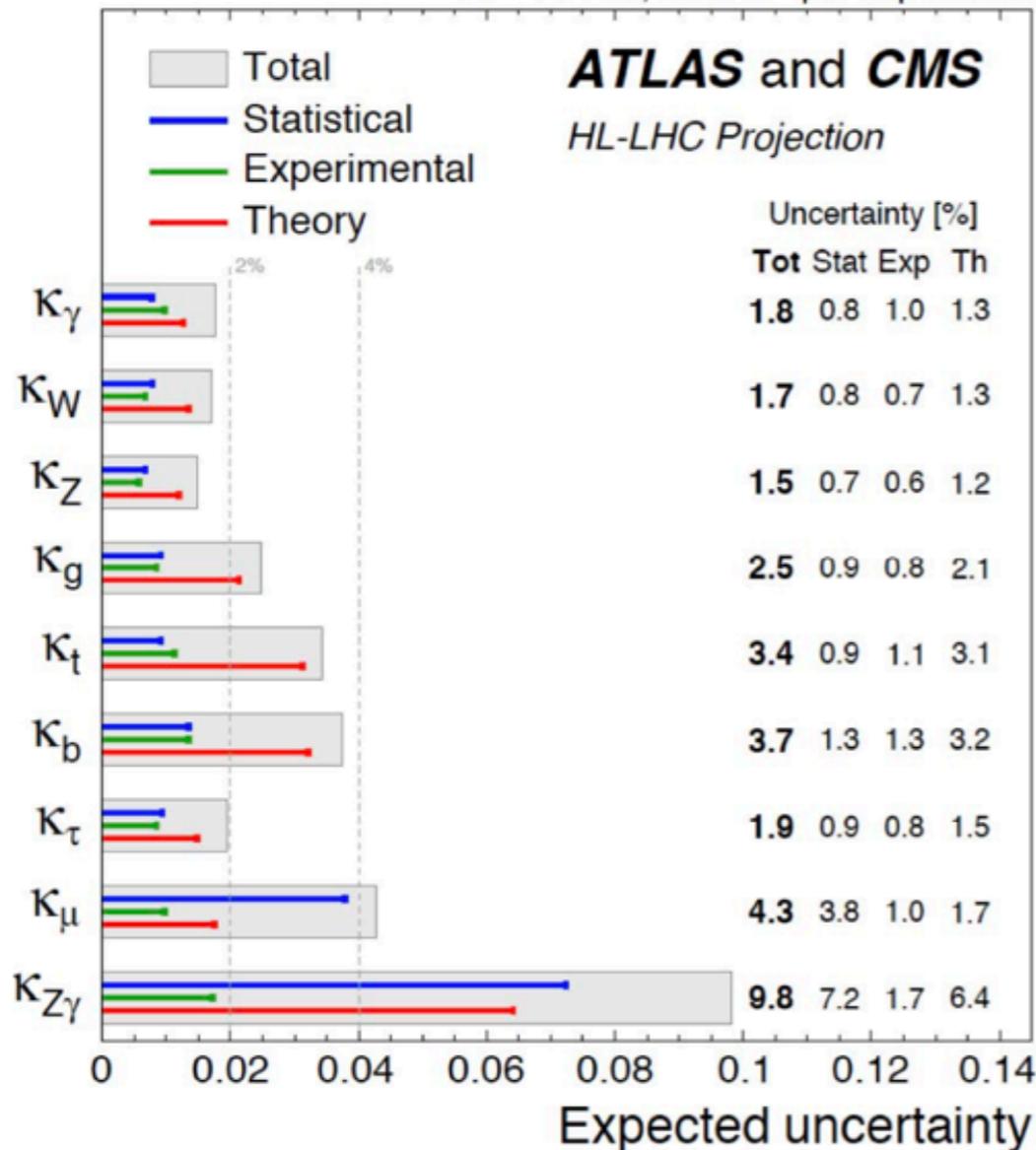
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green points satisfy the present bounds

[De Blas et al., 2020] $\sqrt{s} = 14 \text{ TeV}, 3000 \text{ fb}^{-1}$ per experiment



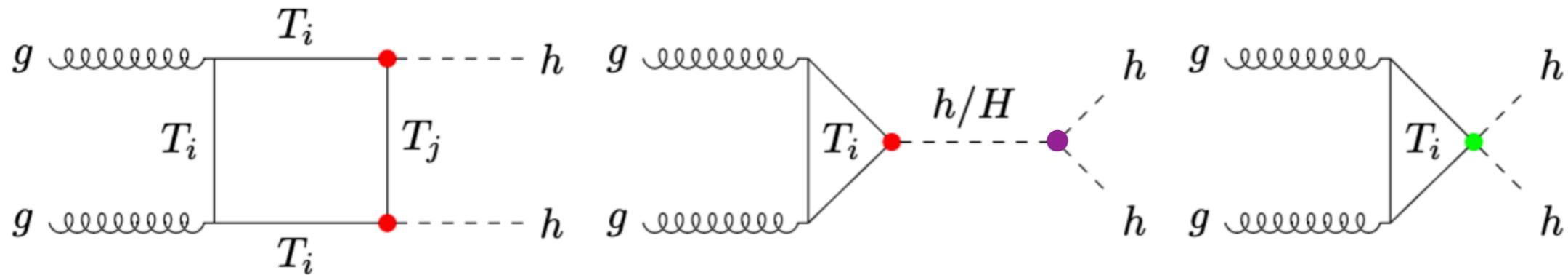
HL-LHC : the Higgs couplings will be constrained at 2-4% level

$$\xi \leq 0.04 \quad f \geq 1200 \text{ GeV}$$

CHMs
NOT
ruled out

Can di-Higgs production at LHC reveal the underlying EWSB?

Signals of New Physics in $gg \rightarrow hh$



$T_i = t, T$'s with $Q=2/3$

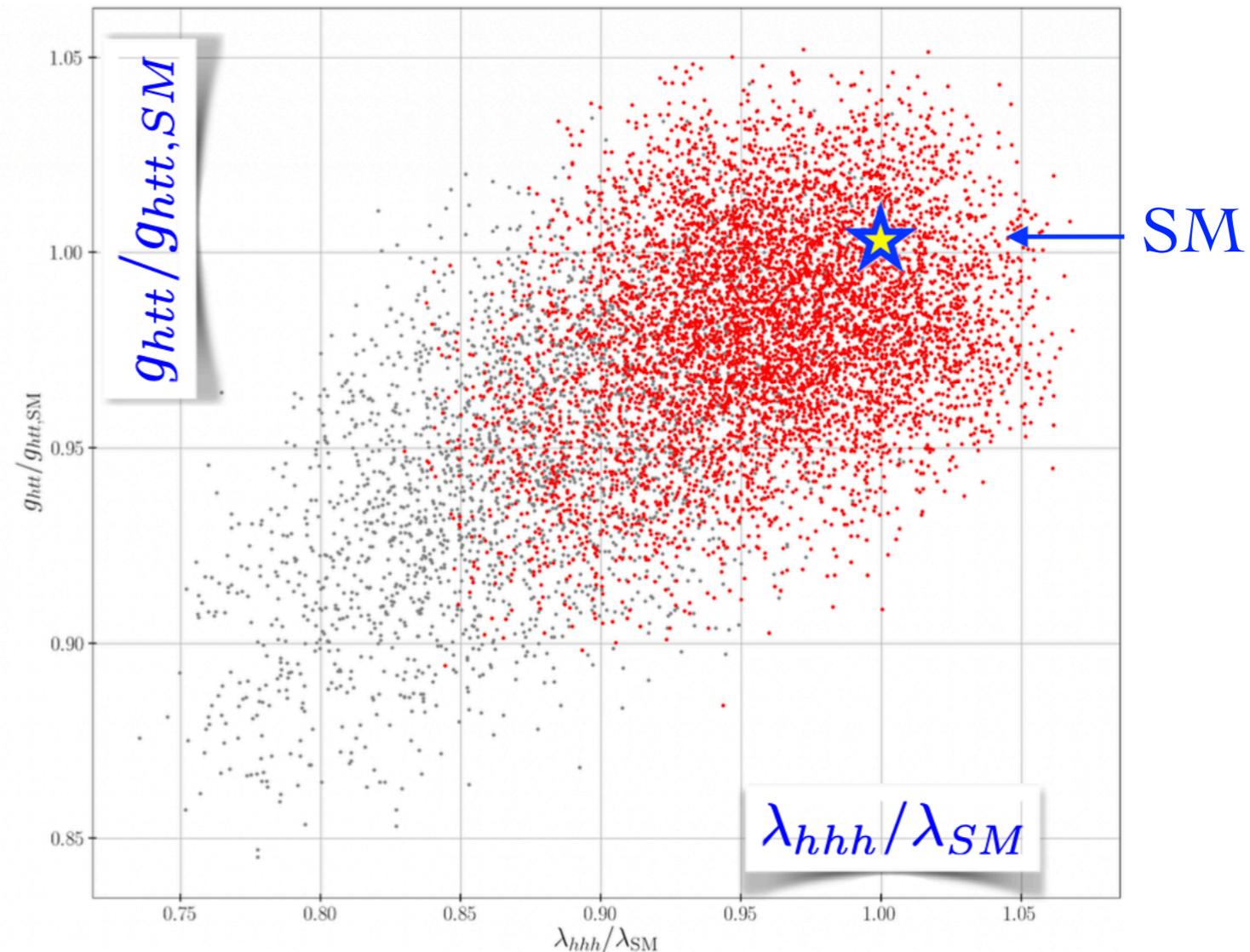
INGREDIENTS: modified h couplings, s -channel H exchange, new heavy tops in the loops, new quartic $hhTT$ (typical of pNGBs)

In C2HDM both resonant and non-resonant modes yield to a change in the integrated cross-section and to peculiar kinematic features in its differential distributions

h-top Yukawa and h-trilinear couplings in the C2HDM

scan over the model parameters $700 \leq f(\text{GeV}) \leq 3000$, $0 \leq \Delta, Y, M_\psi \leq 10f$
with the constraints to reconstruct v_{SM} , m_h , m_{top} exp. values, and $M_T \geq 1.3 \text{ TeV}$

the grey points are
excluded by the
present **direct** and
indirect Higgs
searches
(enforced with
HiggsBounds and
HiggsSignals)



deviations up to 10% in g_{htt} and 15% in λ_{hhh}

Numerical analysis

De Curtis, Delle Rose, Egle, Mühlleitner, Moretti, Sakurai, 2310.10471

The di-Higgs production cross sections through gluon fusion are computed by adapting the public code **HPAIR** (M. Spira), that has been **extended** to **include** the **C2HDM**

INCLUSIVE RESULTS

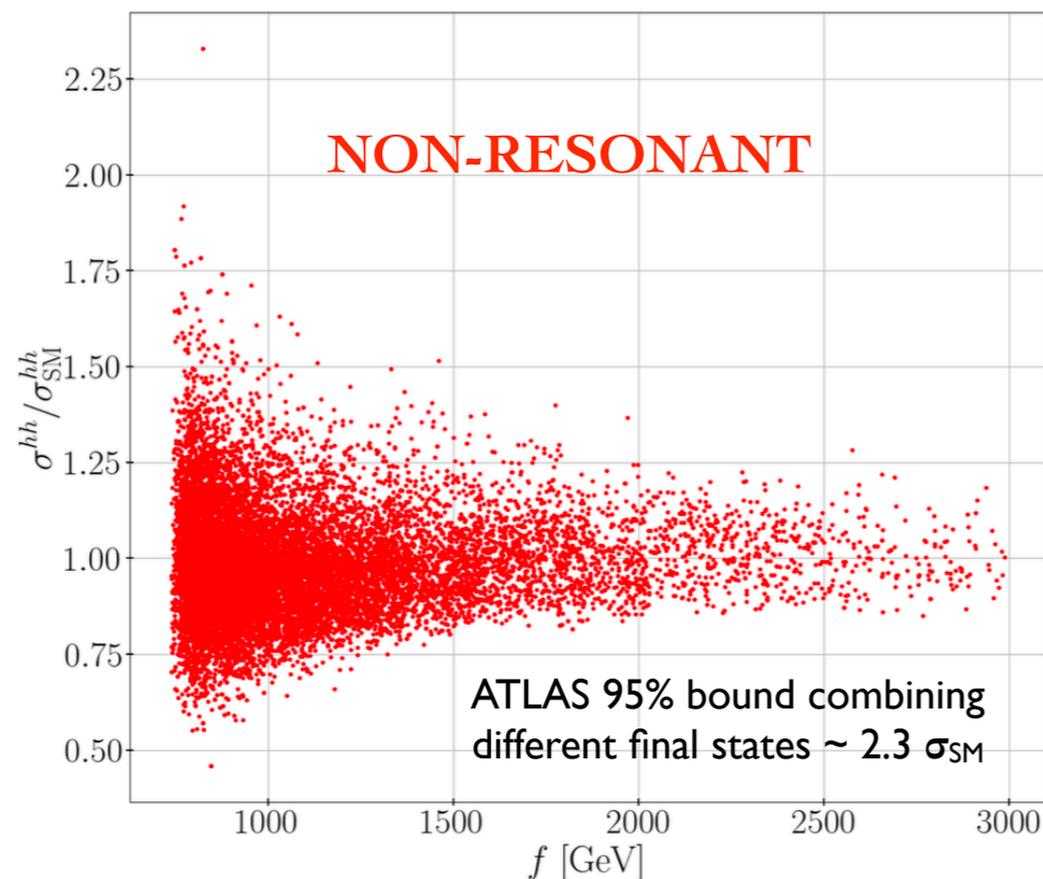
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INCLUSIVE RESULTS

NON-RESONANT: $M_H < 2 m_h$ + cases
with **suppressed resonant contribution**
(small H couplings, large m_H , large Γ_H ,
destructive interferences between diagrams)
 $\sigma(gg \rightarrow H) \times \text{BR}(H \rightarrow hh) / \sigma(gg \rightarrow hh) < 0.1$



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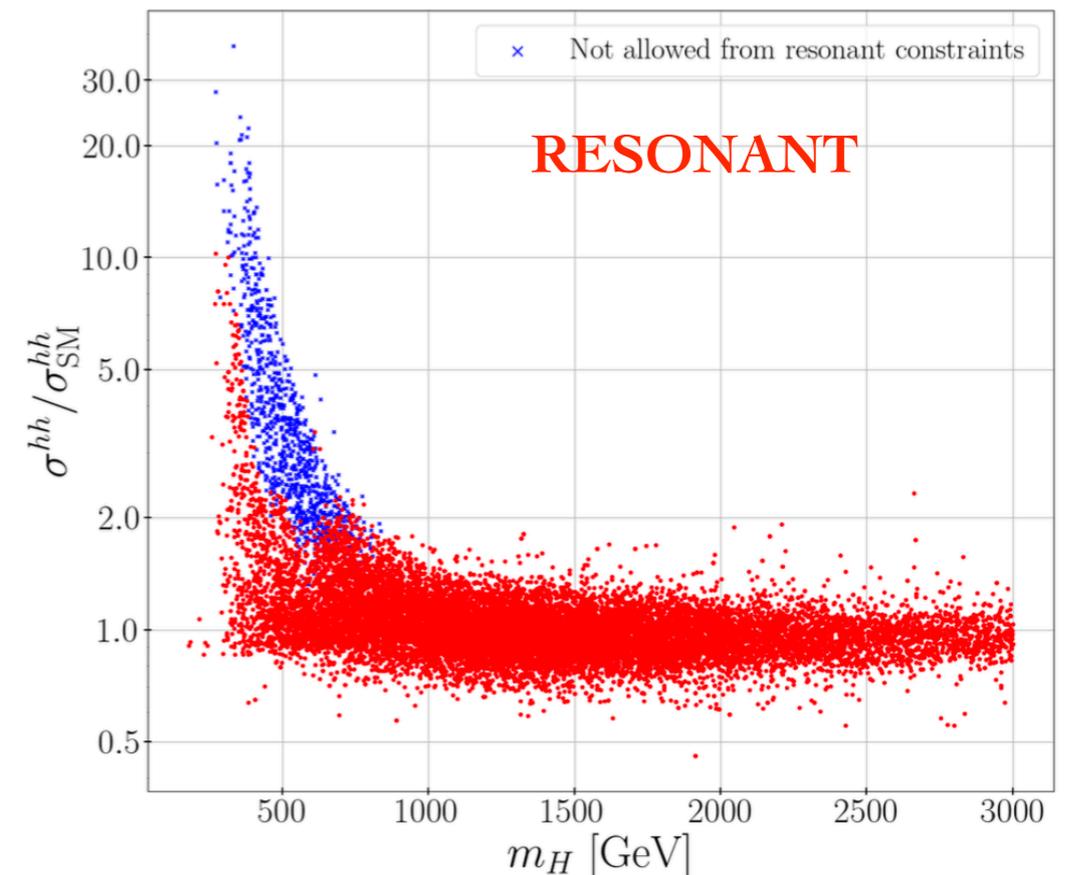
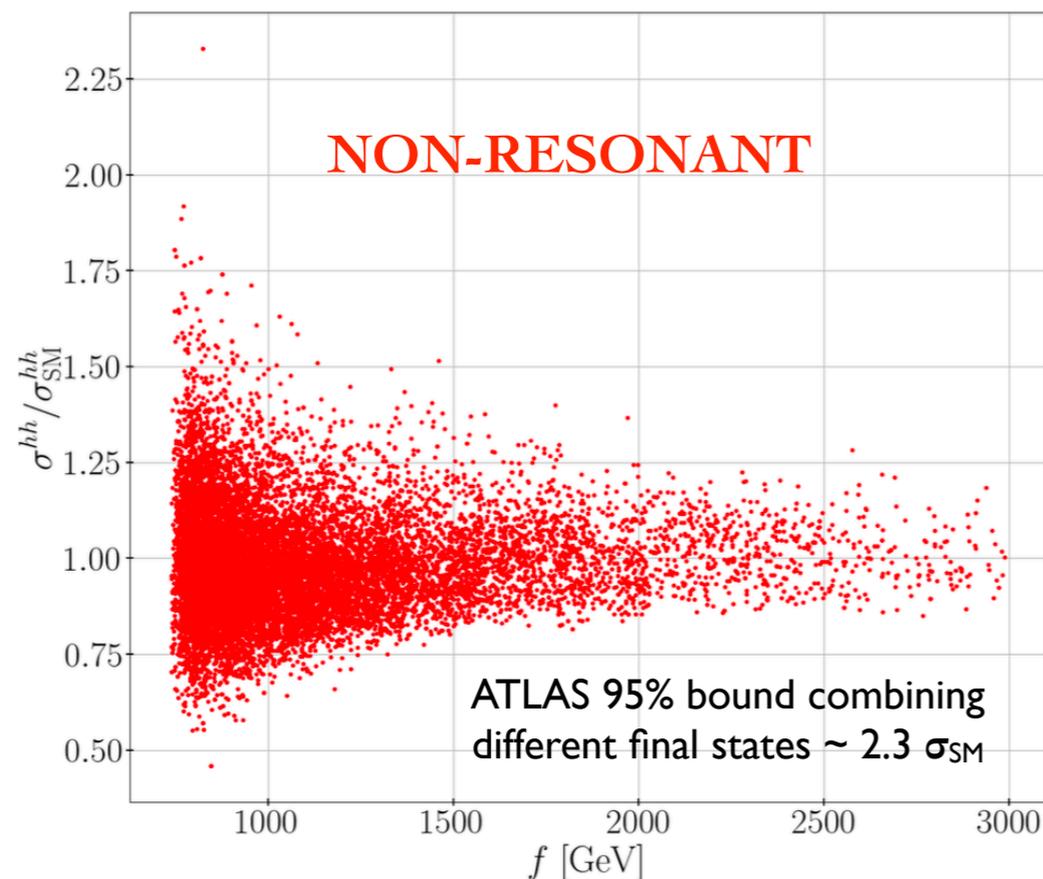
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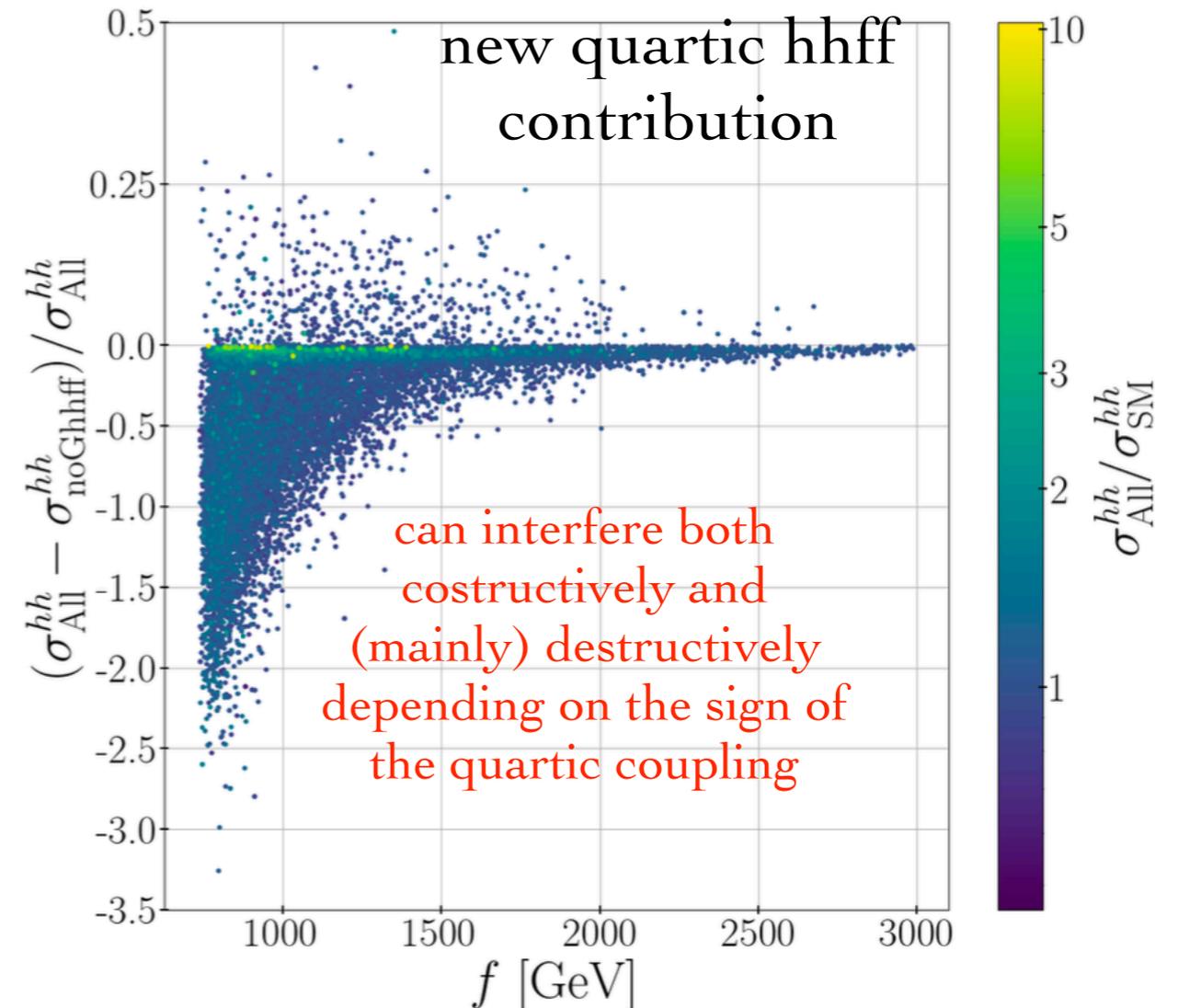
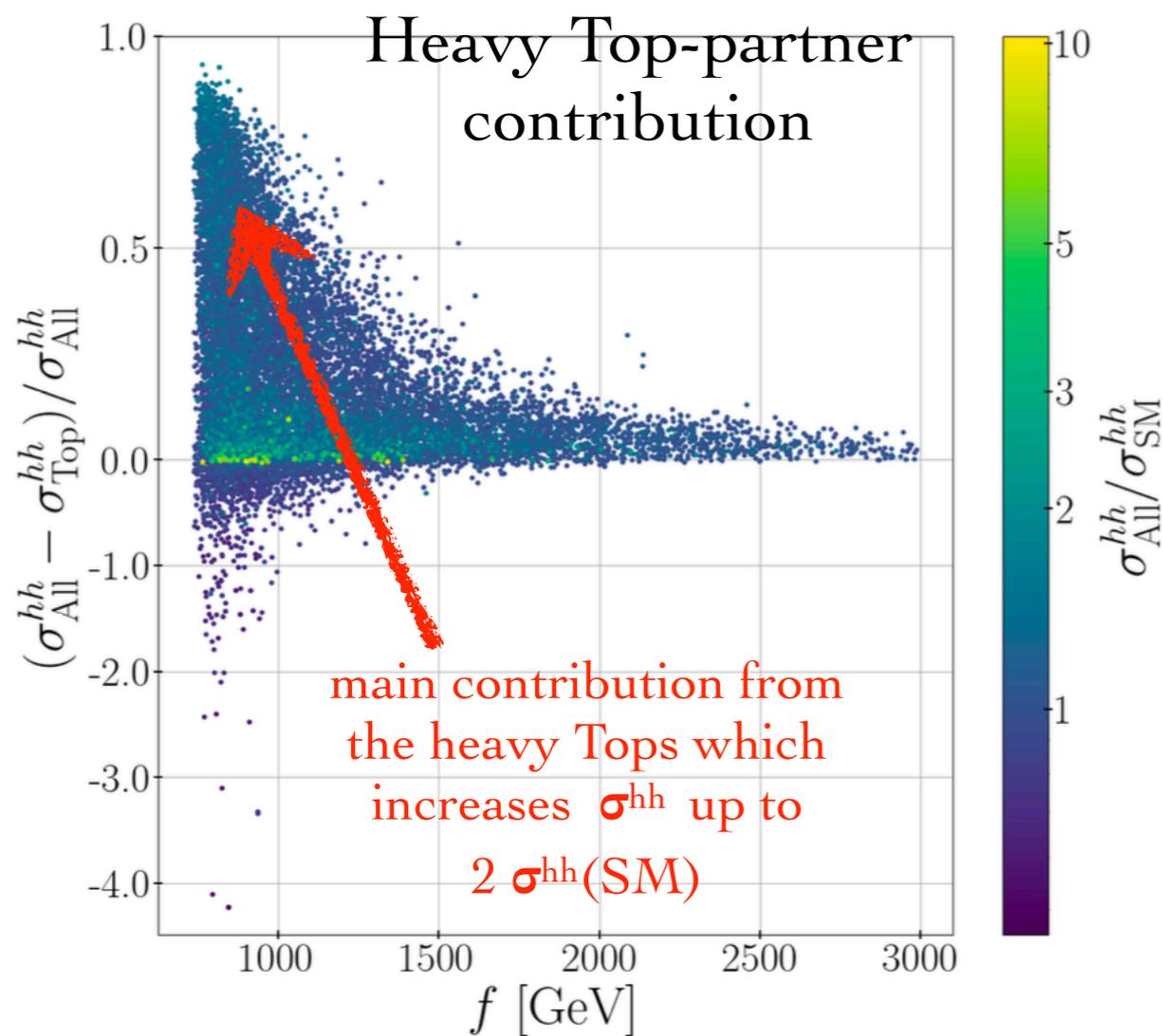
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RESONANT: $M_H > 2 m_h$

compare with the exp. limits on resonant
di-Higgs production obtained in the
narrow width approximation (points with
 $\Gamma_H / M_H > 5\%$ are not excluded)

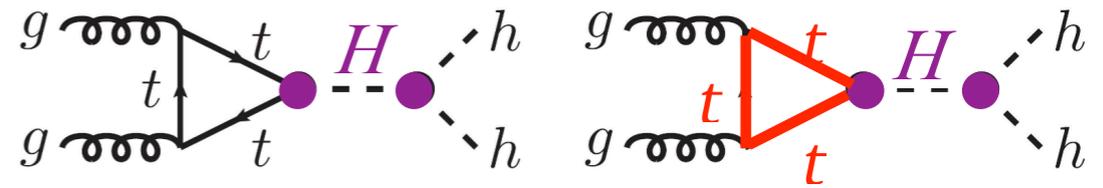


Impact of new C2HDM effects (not present in 2HDM)

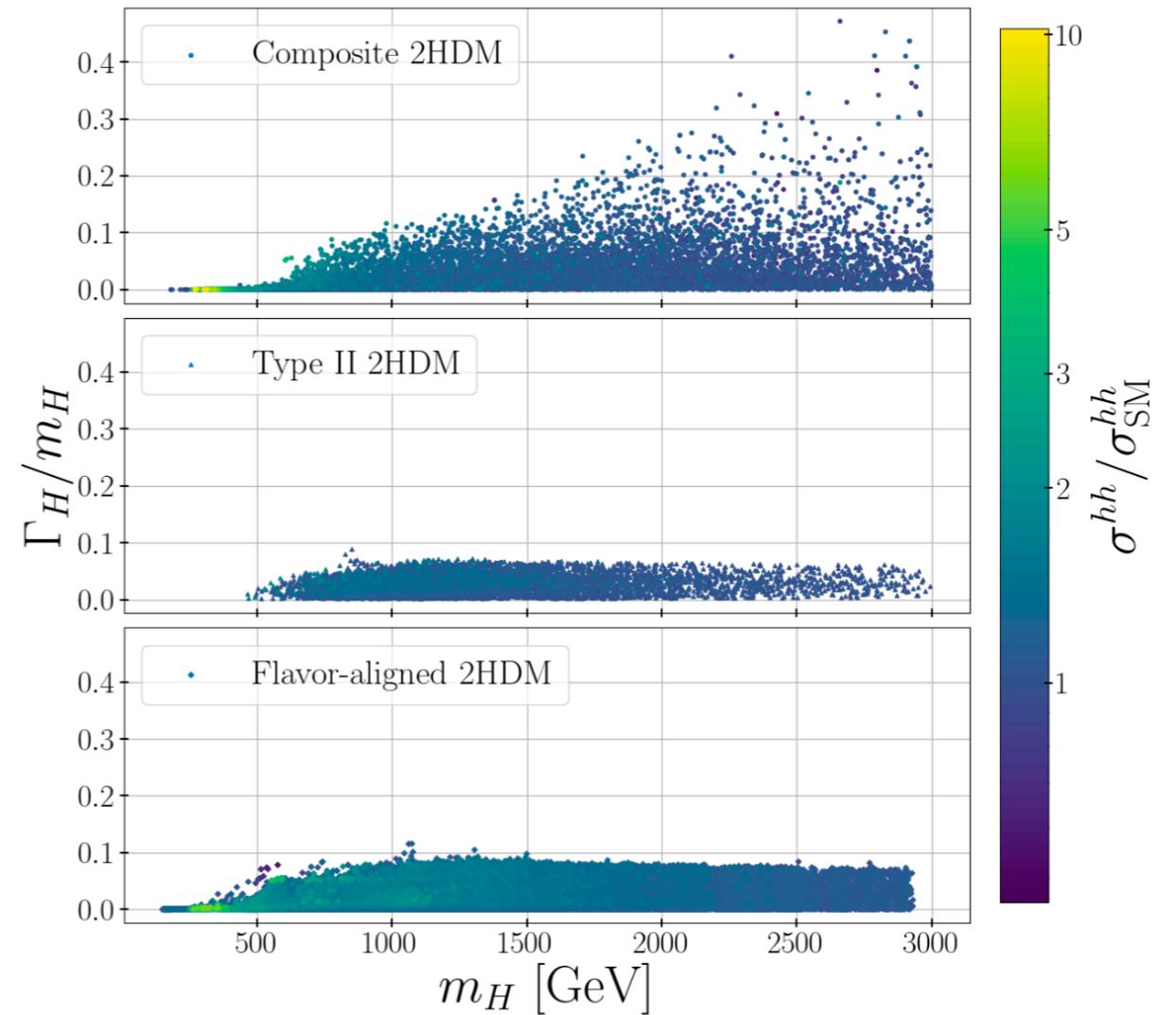
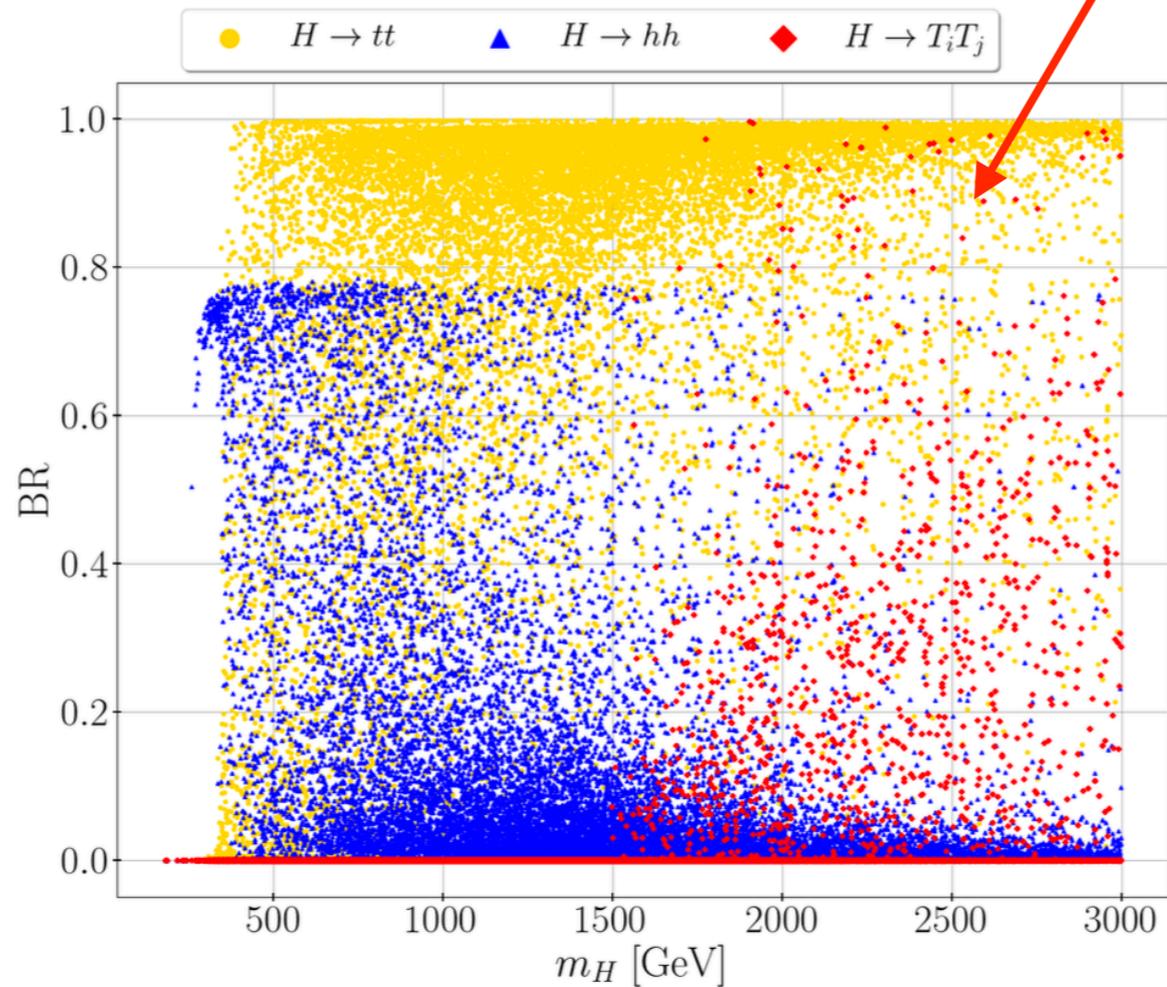


The largest cross-sections are the resonant ones (yellow and green BPs) are not affected by heavy Tops and new quartic terms

H contribution



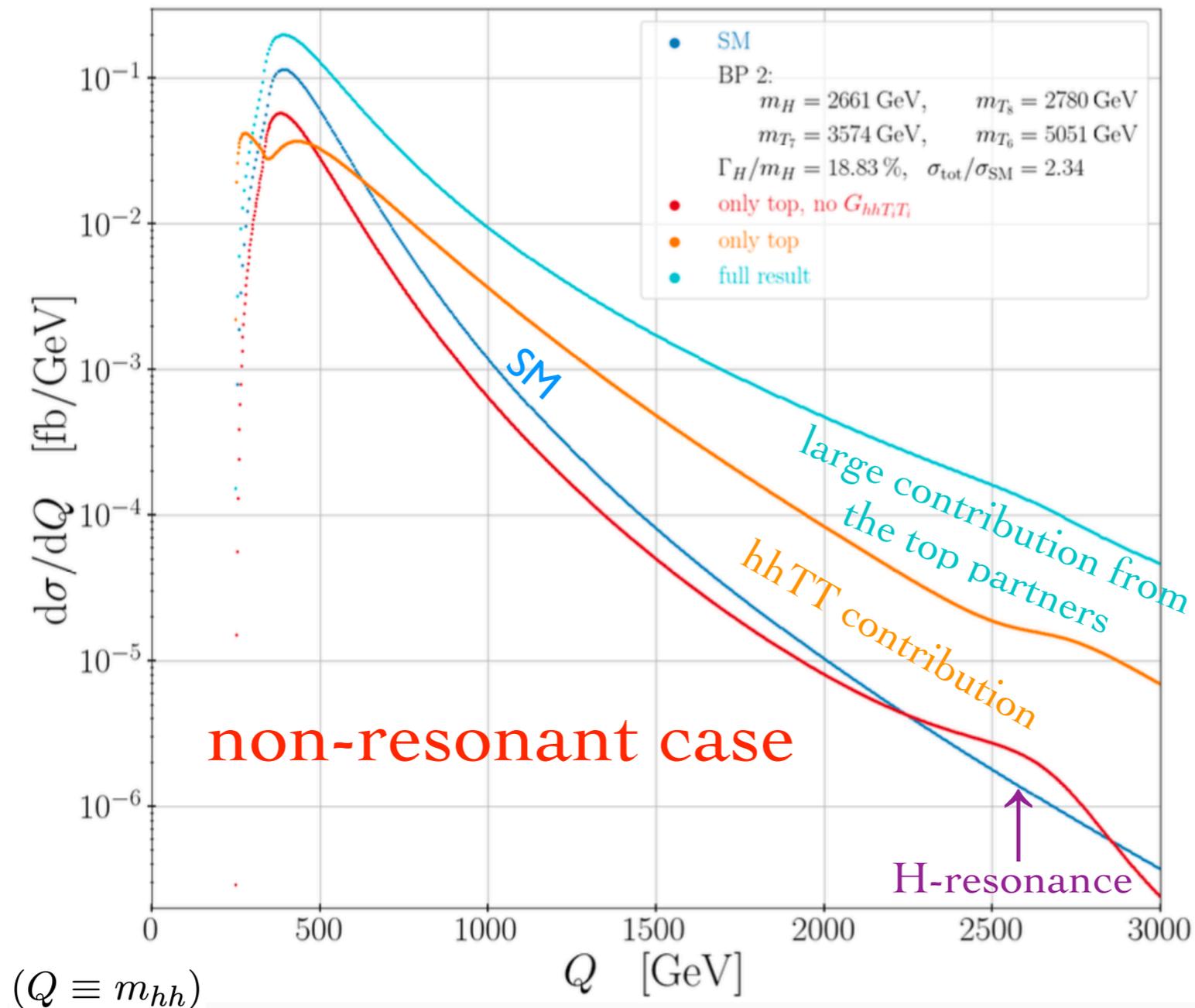
the **heavy Higgs H** can have a **sizeable BR in $T_9 T_{8,7}$**
 $T_9 = \text{top}$, $T_{8,7} = \text{lightest heavy tops}$



Peculiar feature of the C2HDM: Γ_H / M_H can be $\sim 10\text{-}20\%$

di-Higgs production in C2HDM - invariant mass distributions

1. modified hhh (κ_λ) and tth (κ_t) couplings \rightarrow small deviations
2. additional H contributions \rightarrow present in several BSM schemes (MSSM, 2HDM, ...)
3. additional heavy top contributions ($t' = T_i, i = 1, \dots, 8$) \rightarrow naturally present in CHMs
4. quartic $tthh$ coupling (since h is pNGB) \rightarrow naturally present in CHMs



$m_H \sim 2.6$ TeV

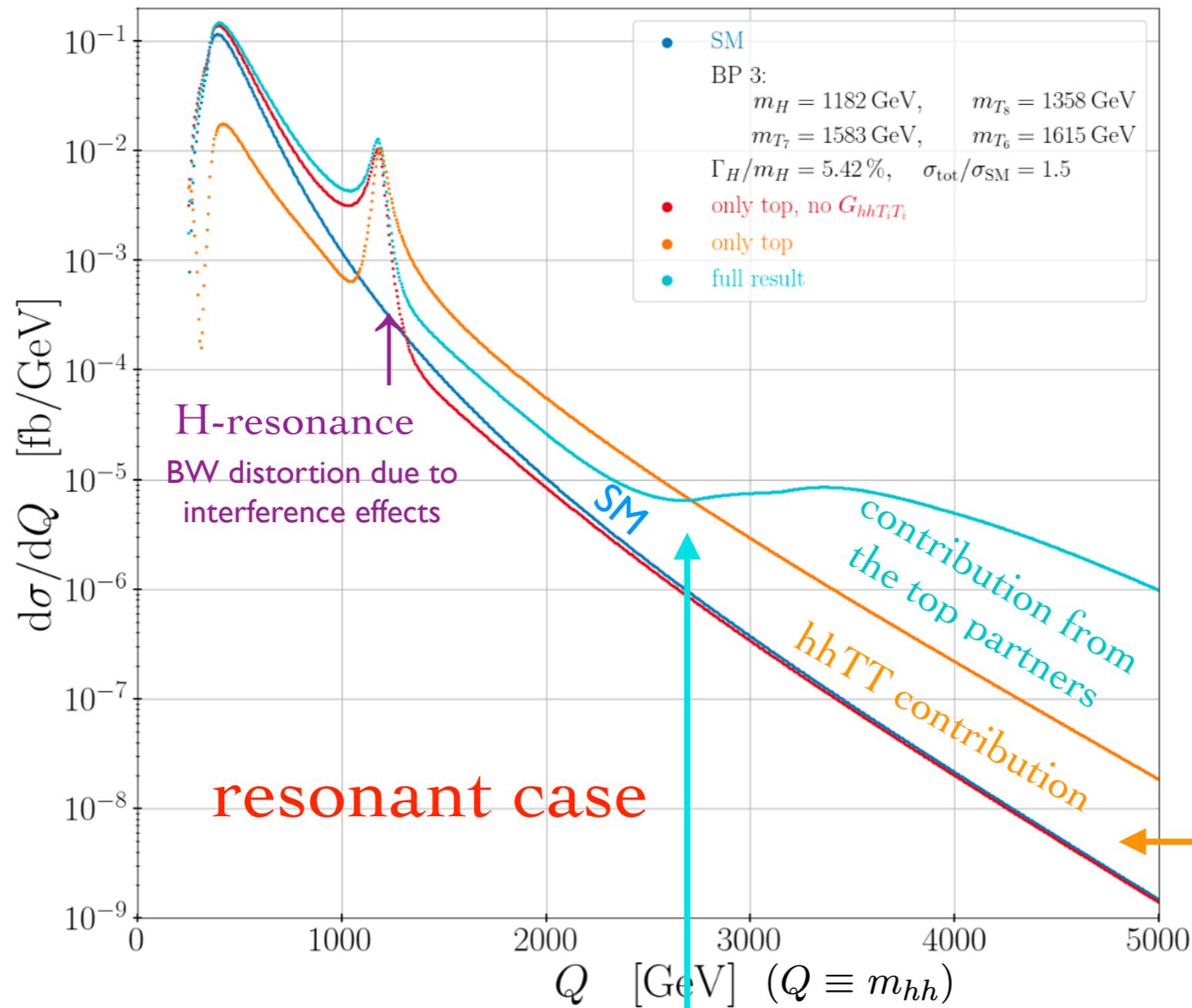
$\Gamma_H/m_H \sim 19\%$

Large Width

$m_T > 2.7$ TeV

$\sigma_{tot}/\sigma_{SM} = 2.34$

di-Higgs production in C2HDM - invariant mass distributions



$$\sigma_{\text{tot}}/\sigma_{\text{SM}} = 1.5$$

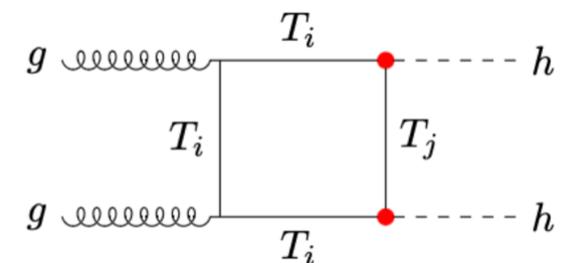
$$m_H \sim 1.2 \text{ TeV}$$

$$\Gamma_H/m_H \sim 5.4\%$$

$$m_T > 1.3 \text{ TeV}$$

destructive interference before the peak and constructive interference after the peak

start to see the threshold shape at $2MT$ induced by boxes



di-Higgs production in C2HDM - invariant mass distributions

The results of the present analysis are primarily of theoretical nature and serve to demonstrate that a computable framework exists within composite scenarios that can eventually be tested experimentally

$$\sigma_{\text{tot}}/\sigma_{\text{SM}} = 1.5$$

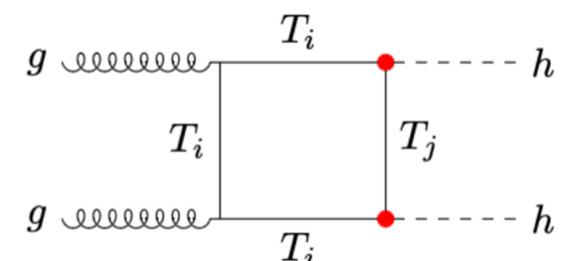
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destructive interference before the peak and constructive interference after the peak

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**New Physics
in the Higgs sector**

```
graph TD; A["New Physics in the Higgs sector"] --> B["First order EW phase transition"]; A --> C["deviations in the Higgs couplings"]; B --> D["Gravitational Wave signals"]; B --> E["EW Baryogenesis"]; C --> F["e.g. Signals in di-Higgs production"];
```

**First order EW
phase transition**

**deviations in the
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EW Baryogenesis

**New Physics
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Cosmology - Collider synergy

**Gravitational Wave
signals**

**e.g. Signals in di-Higgs
production**

*observables at
future interferometers*

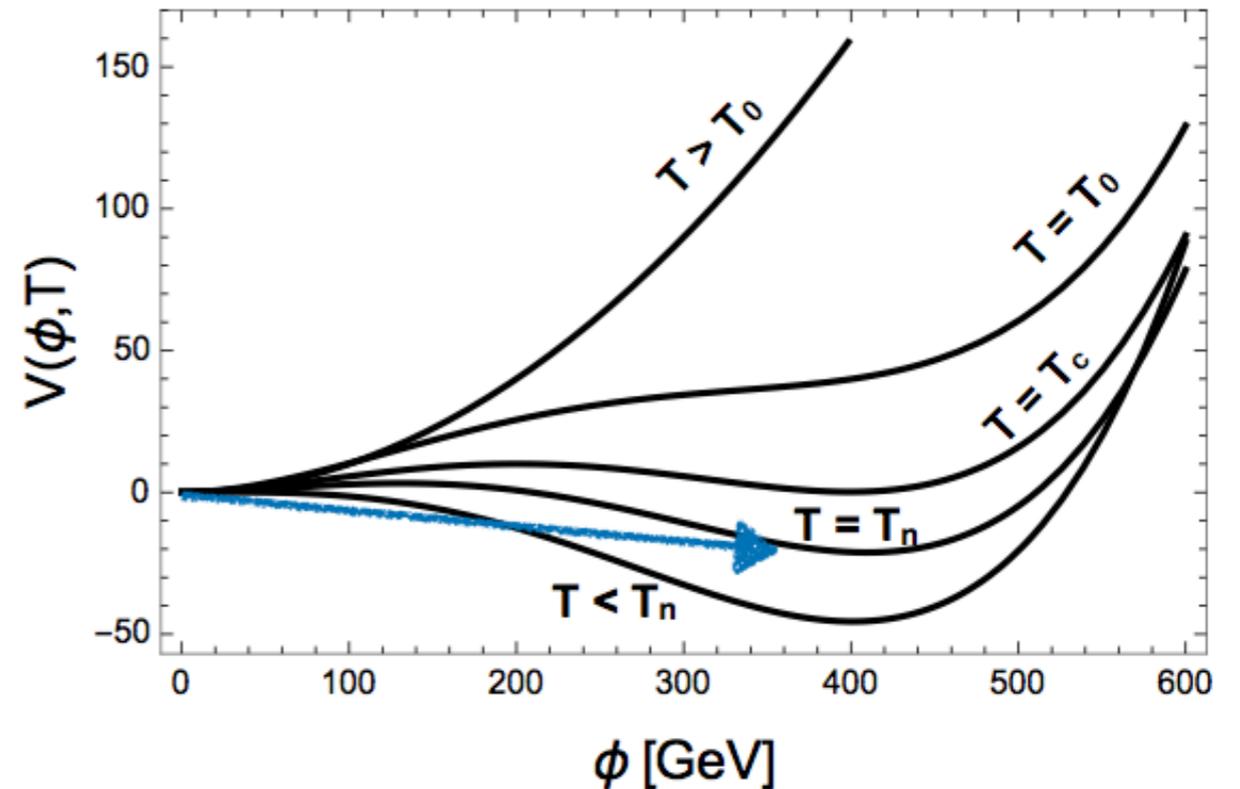
*observables at
present and future colliders*

EW Baryogenesis

Strong EW Phase Transition can trigger Baryogenesis

Thermal History

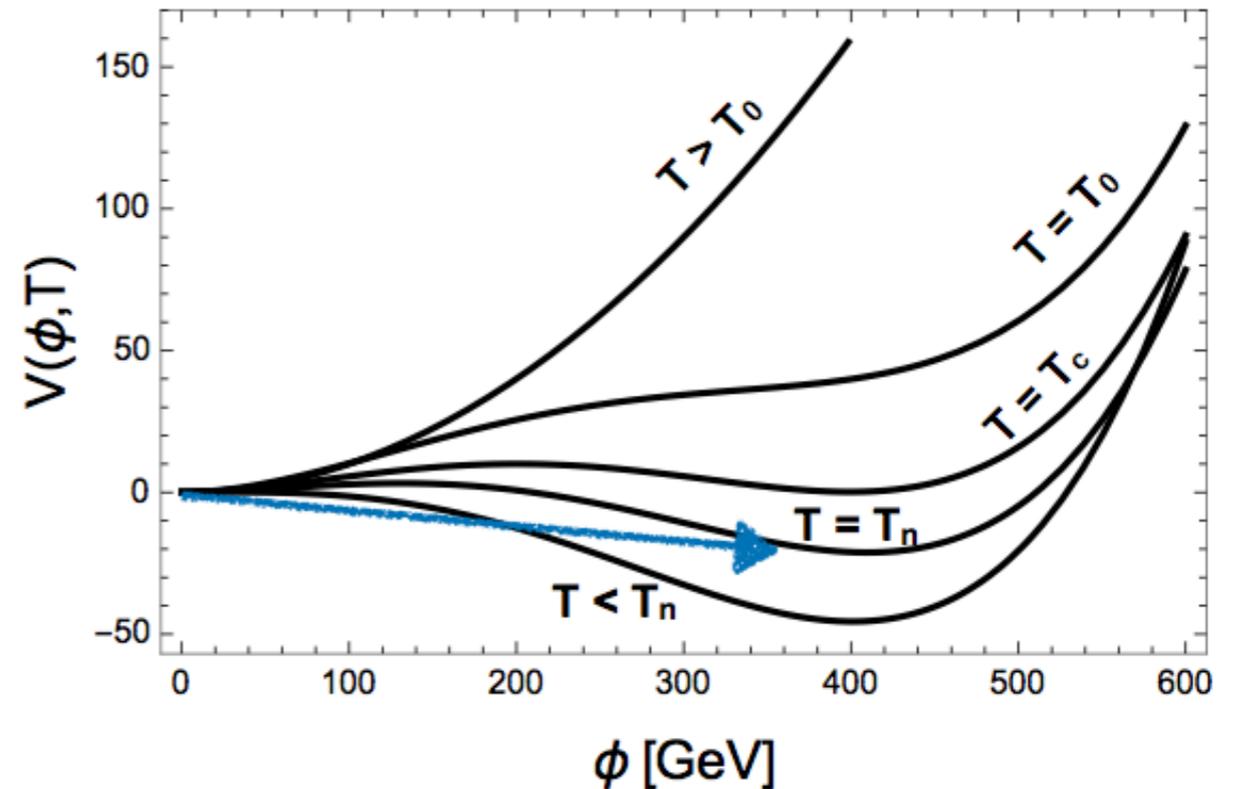
- ☑ The EW symmetry is restored at $T > T_0$
below T_0 a new (local) minimum appears
- ☑ At a critical T_c the two minima are degenerate
and separated by a barrier (two phases
coexist)
- ☑ The transition starts at the bubble nucleation
temperature $T_n < T_c$



Strong EW Phase Transition can trigger Baryogenesis

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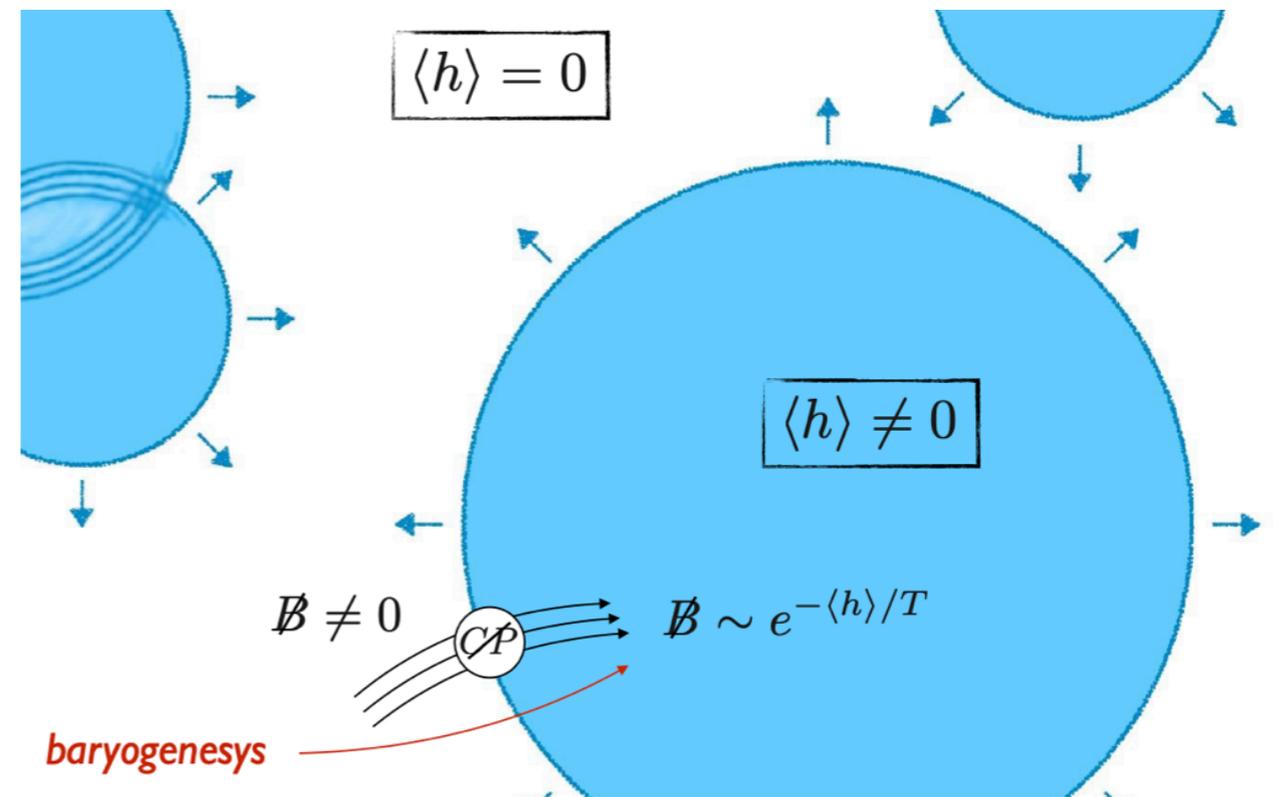
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Sakharov Conditions for Baryogenesis

- ☑ Baryon number violation
- ☑ C and CP violation
- ☑ **Out of equilibrium dynamics:**
(strong) 1st order phase transition

In the SM phase transition is a smooth
crossover, also not enough CP violation
from CKM → **NP needed !!**



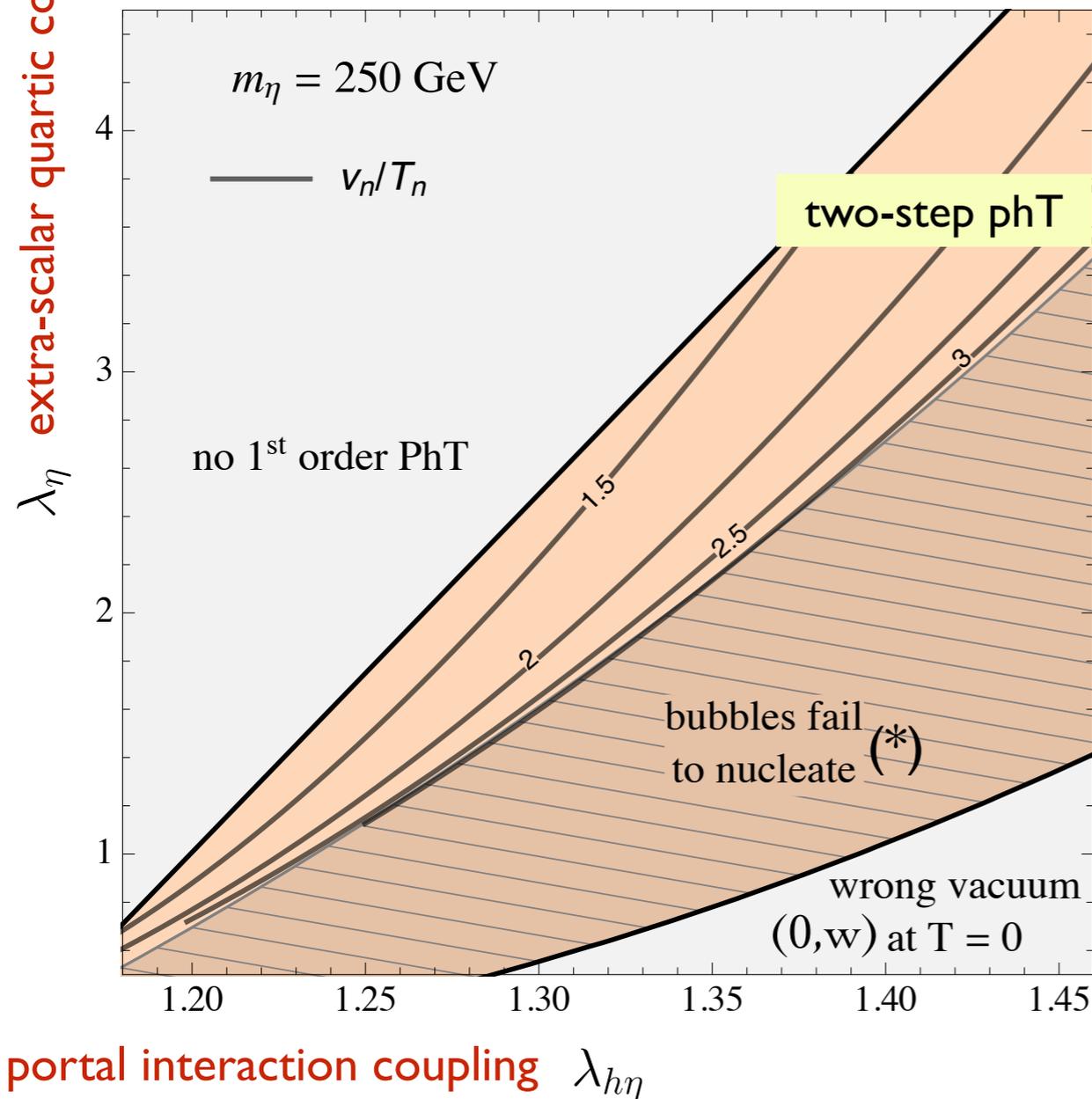
Composite Dynamics in the Early Universe

Properties of the EWPhT

H+ η pNGBs of SO(6) \rightarrow SO(5)

(De Curtis, Delle Rose, Panico, 2019)

extra-scalar quartic coupling λ_η



The EWPhT starts at $T_n < T_c$ determined by requiring:
Rate of nucleation of bubbles / Hubble volume ~ 1

The computation of T_n requires solving (numerically) a two-field bounce equation (use CosmoTransition package)

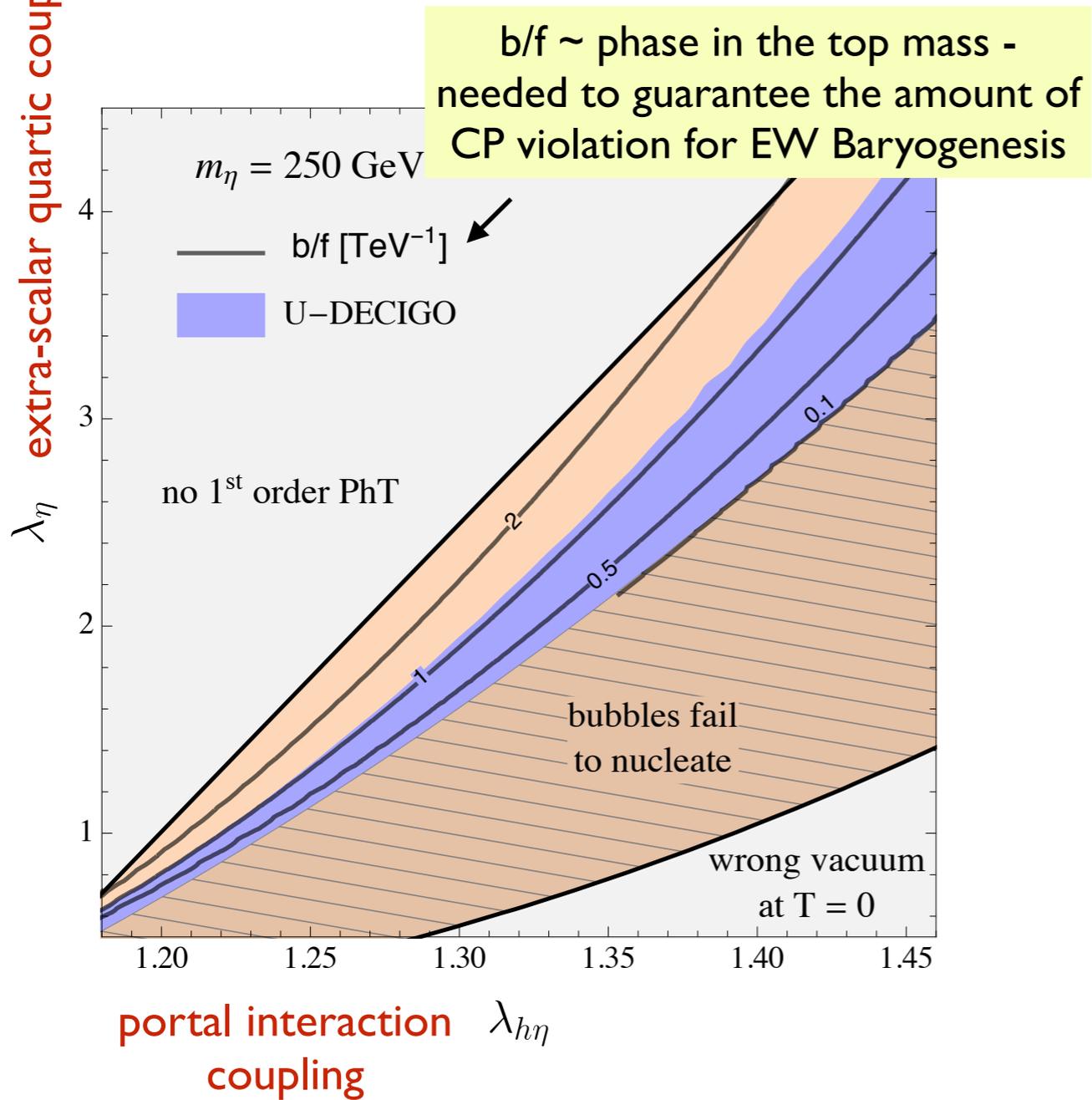
— Strength of the phase transition
 v_n/T_n ($v_n = \langle h \rangle|_{T_n}$)
a crucial parameter for EWBG

T_n is one of the parameter characterising the amplitude and the frequency peak of the GW spectrum

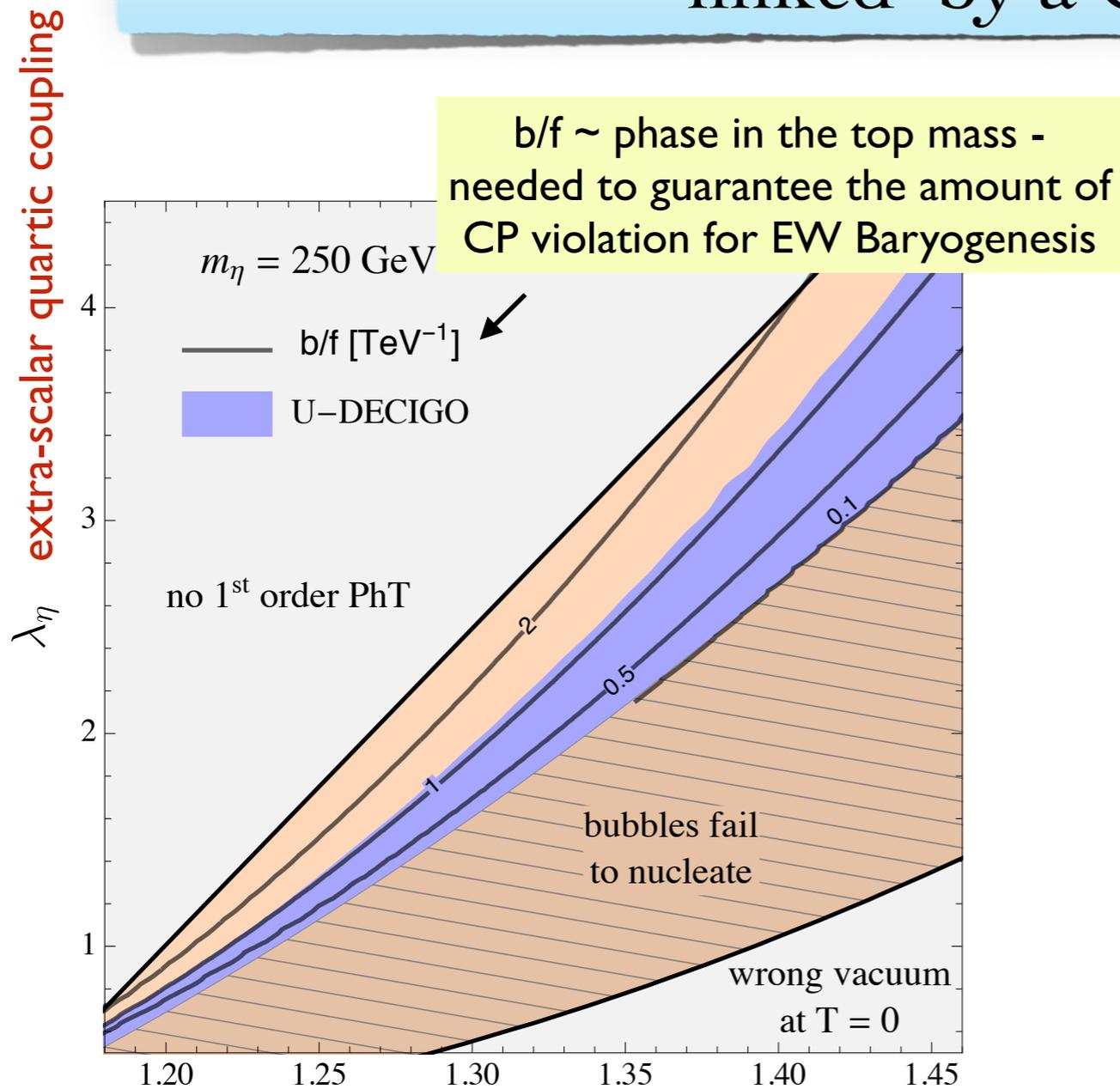
(*) the rate of bubble formation does not balance the Hubble expansion (ex. $\lambda_{h\eta}$ too large produces a high barrier) and no EWBSB occurs

Strong EWPhT, EWBG and GW spectrum linked by a CHM scenario

λ_η extra-scalar quartic coupling

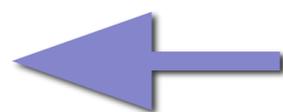


Strong EWPhT, EWBG and GW spectrum linked by a CHM scenario



portal interaction coupling $\lambda_{h\eta}$

same region where the EWBG could be achievable

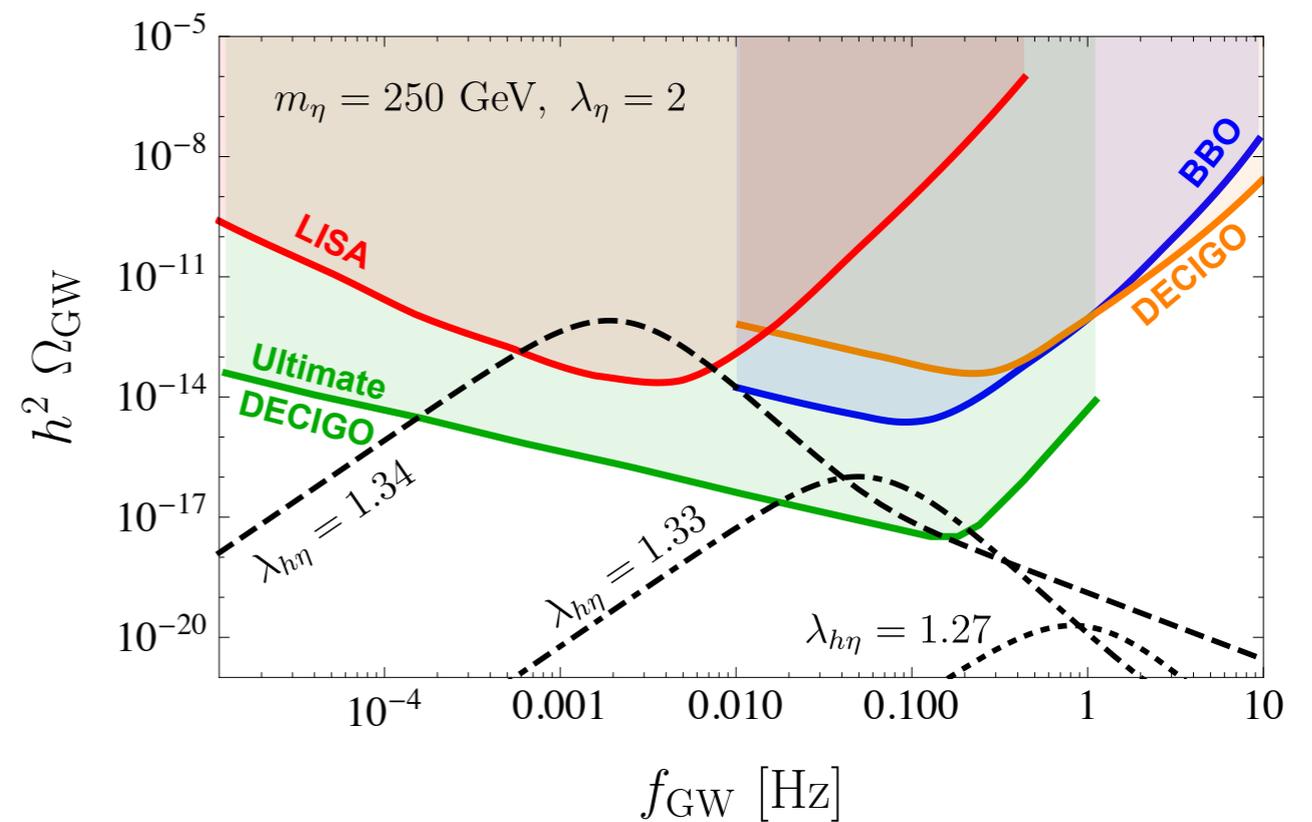


The bubbles expand, collide incoherently ...

Stochastic Background of GW's :
(bubble collisions, sound waves in the plasma, magnetohydrodynamic turbulence effects)

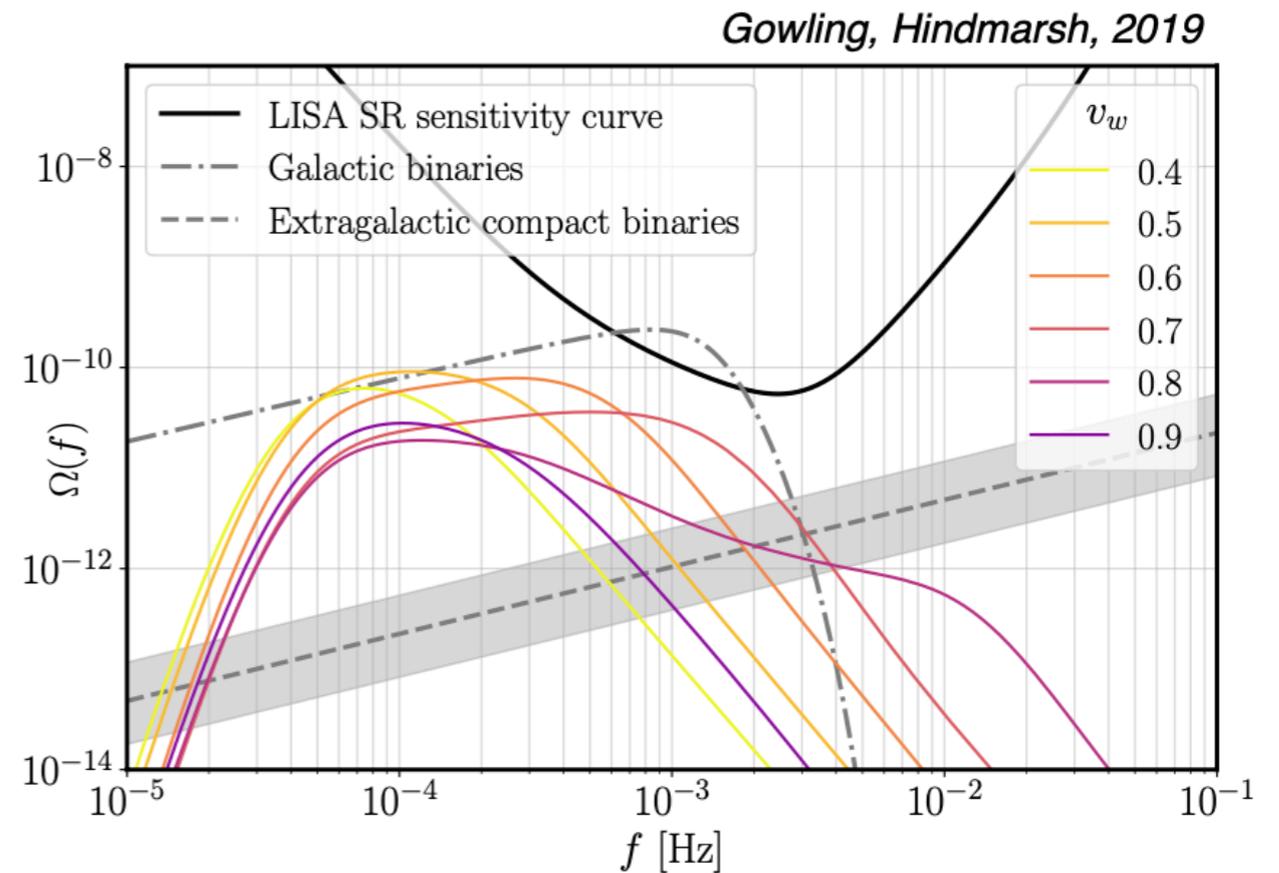
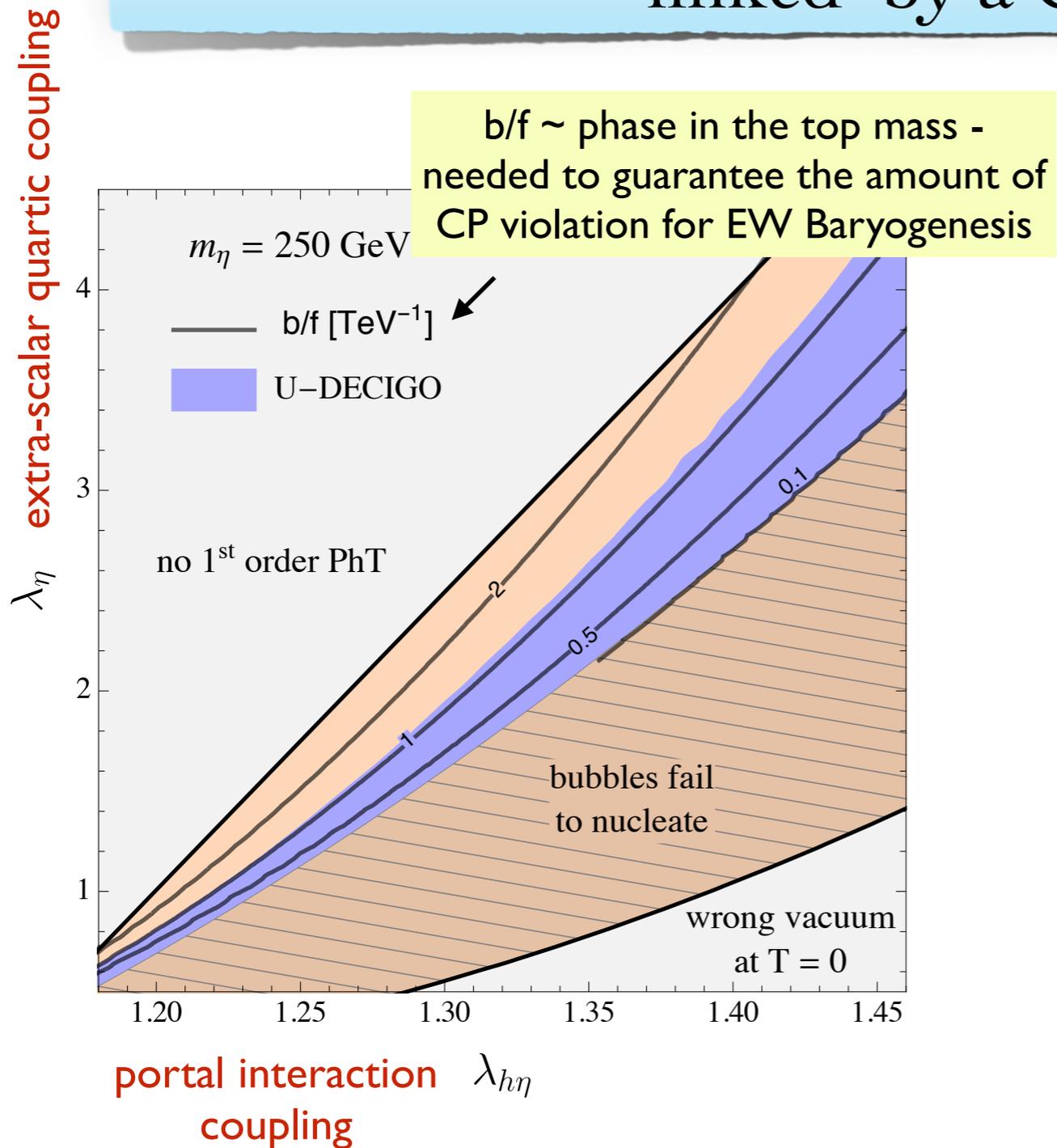
(Grojean, Servant '06, Caprini, Durrer, Servant '08, '09)

Gravitational Wave Spectrum



peak frequencies **within the sensitivity reach of future experiments** for a significant part of the parameter space

Strong EWPhT, EWBG and GW spectrum linked by a CHM scenario



the **wall speed** has a strong effect on the shape of the power spectrum

Can be determined by **solving the Boltzmann equation** which describes the plasma dynamics and its interactions with the bubble wall

De Curtis, Delle Rose, Guiggiani, Mayor, Panico JHEP 03(2022),163; JHEP 05(2023),194; JHEP 05(2024) 009

Conclusions

- ☑ **di-Higgs production** is a target process for the LHC, within the SM it is the experimental signature of the Higgs self-interaction, but also a **probe for BSM scenarios**
- ☑ We analysed $gg \rightarrow hh$ within the C2HDM with an approach which enables to **disentangle** the different NP ingredients: coupling modifications, **new resonance exchange**, **heavy fermions in the loops**, and the **extra quartic couplings**
- ☑ Sizeable effects both in the **integrated cross-section** and in the **differential distributions** open the prospect of **using di-Higgs production at the LHC as a probe for NP** with the possibility to **disentangle among different BSM schemes**
- ☑ **New Physics in the Higgs sector** can provide 1st order EWPhT, thus signals of gravitational waves and EW baryogenesis, along with modifications to the Higgs couplings and signatures at colliders