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Indirect probes of the Higgs sector

(In)direct probes: theorist's view



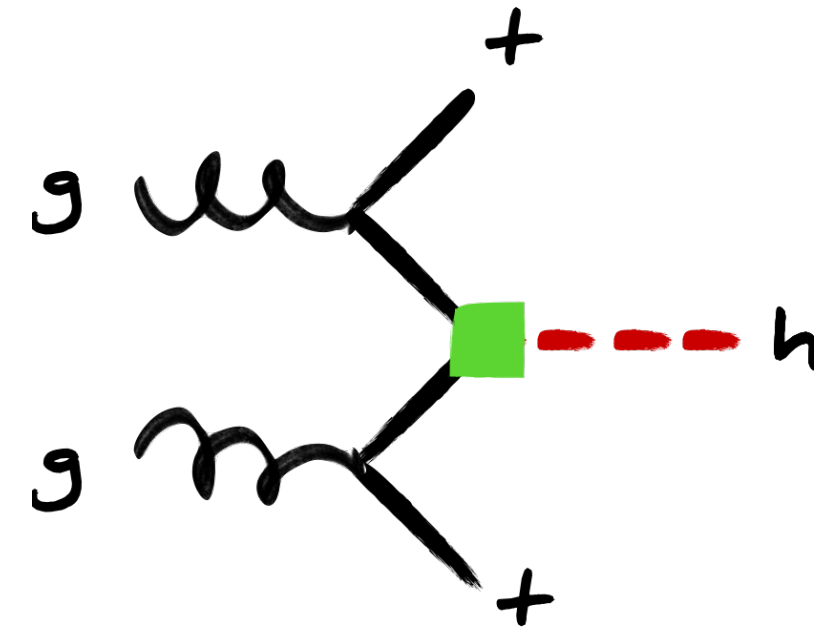
Direct probes: tests of tree-level interactions

Indirect probes: tests of loop modifications of interactions

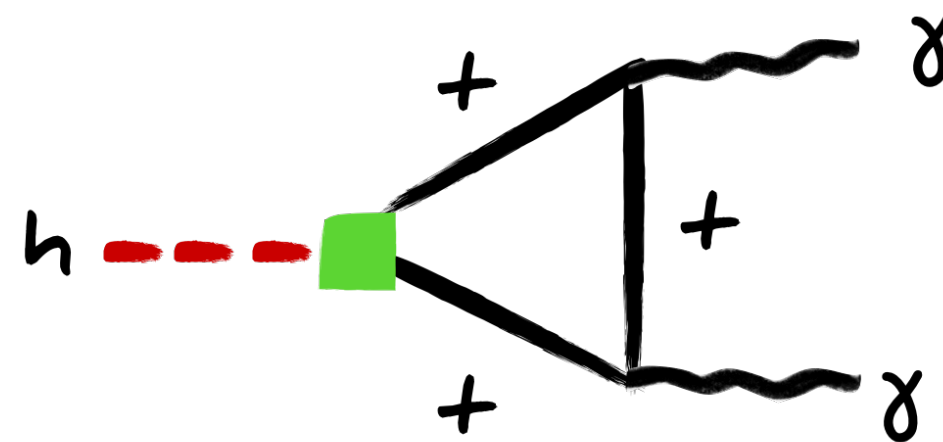
(In)direct probes: example top Yukawa (y_t)



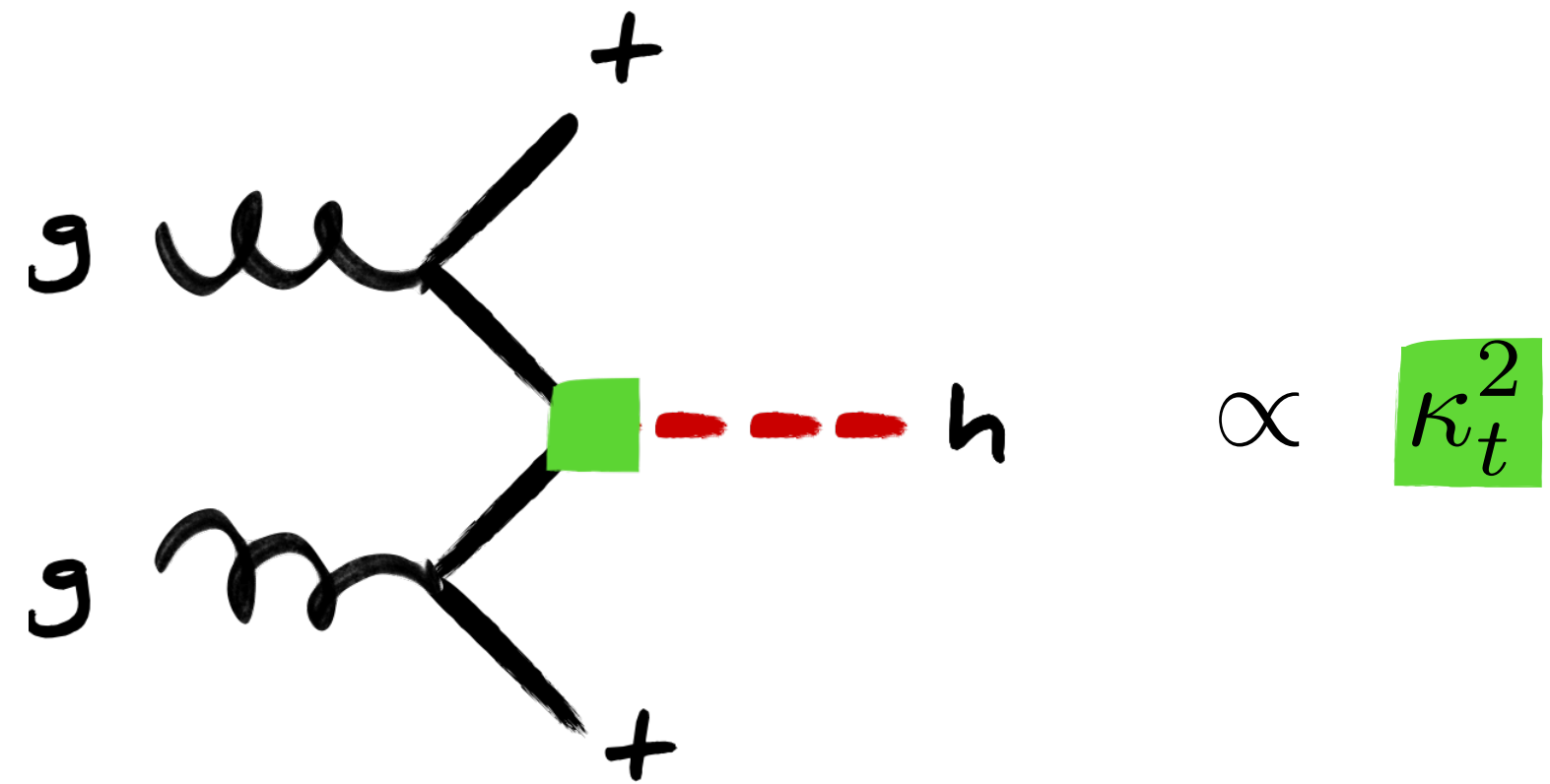
Direct probe of y_t :



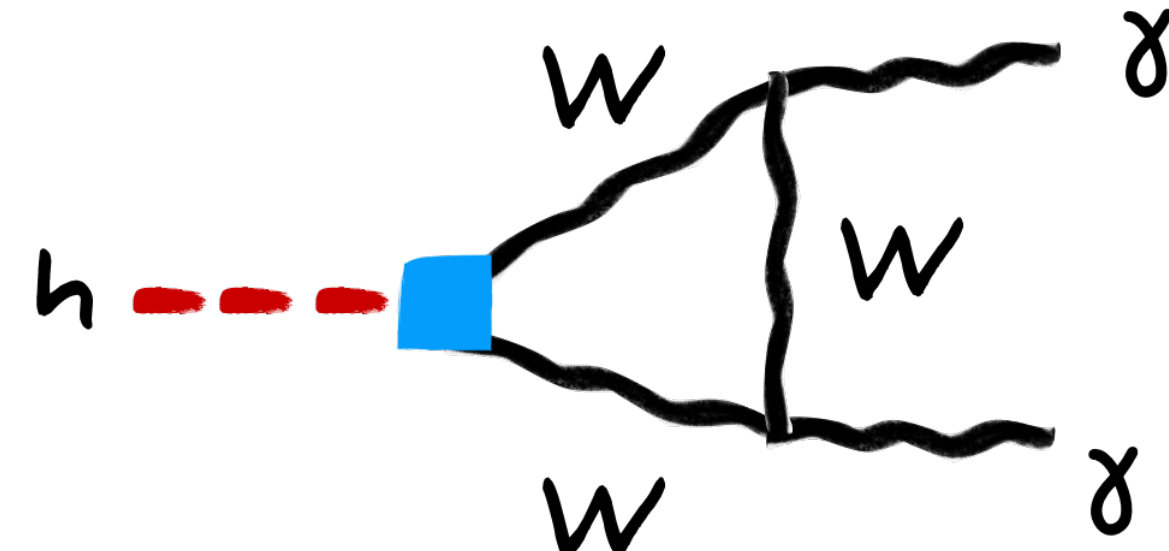
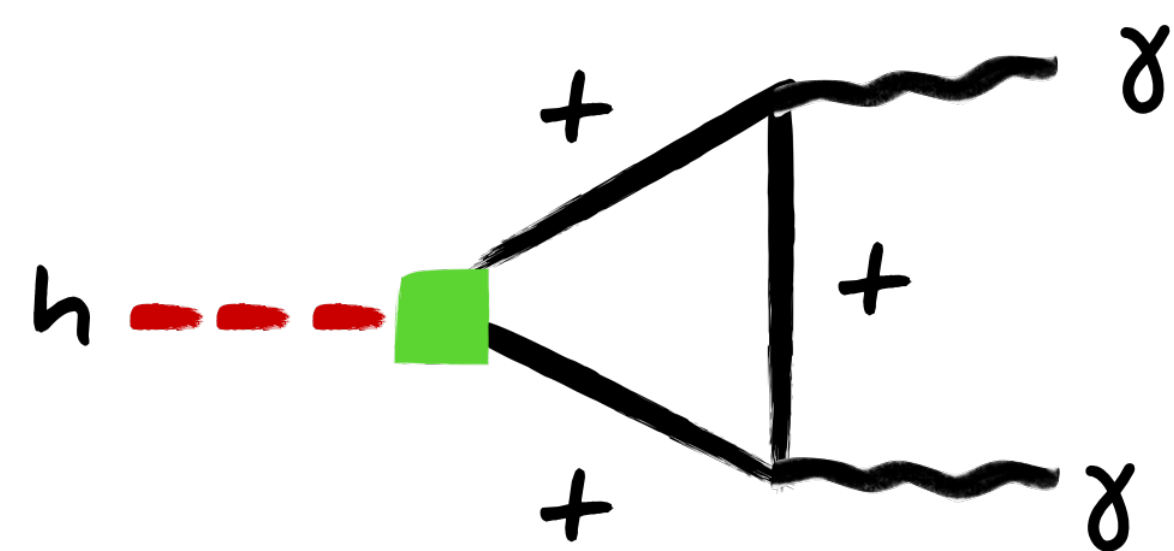
Indirect probe of y_t :



(In)direct probes: synergy & complementarity



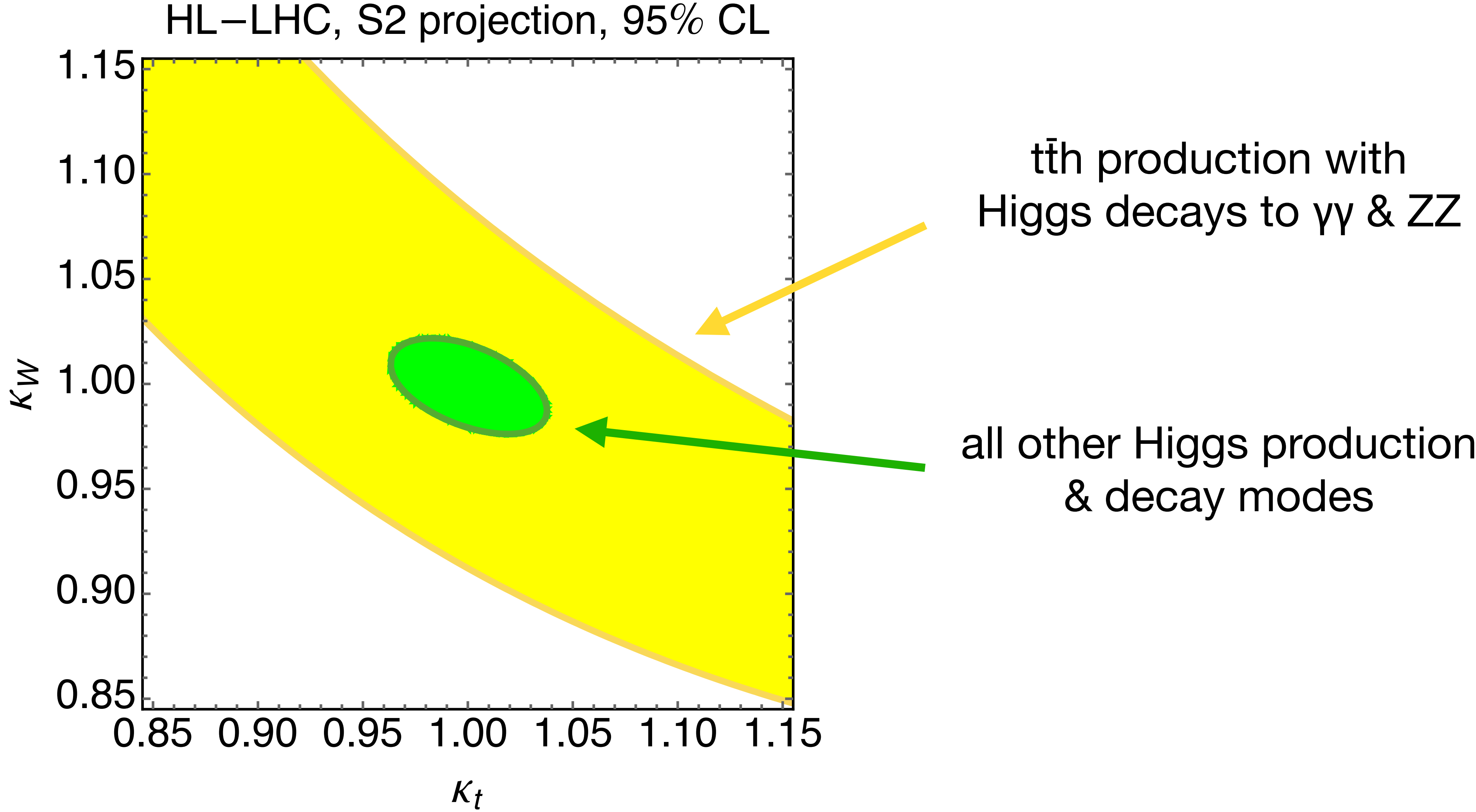
$$\propto \kappa_t^2$$



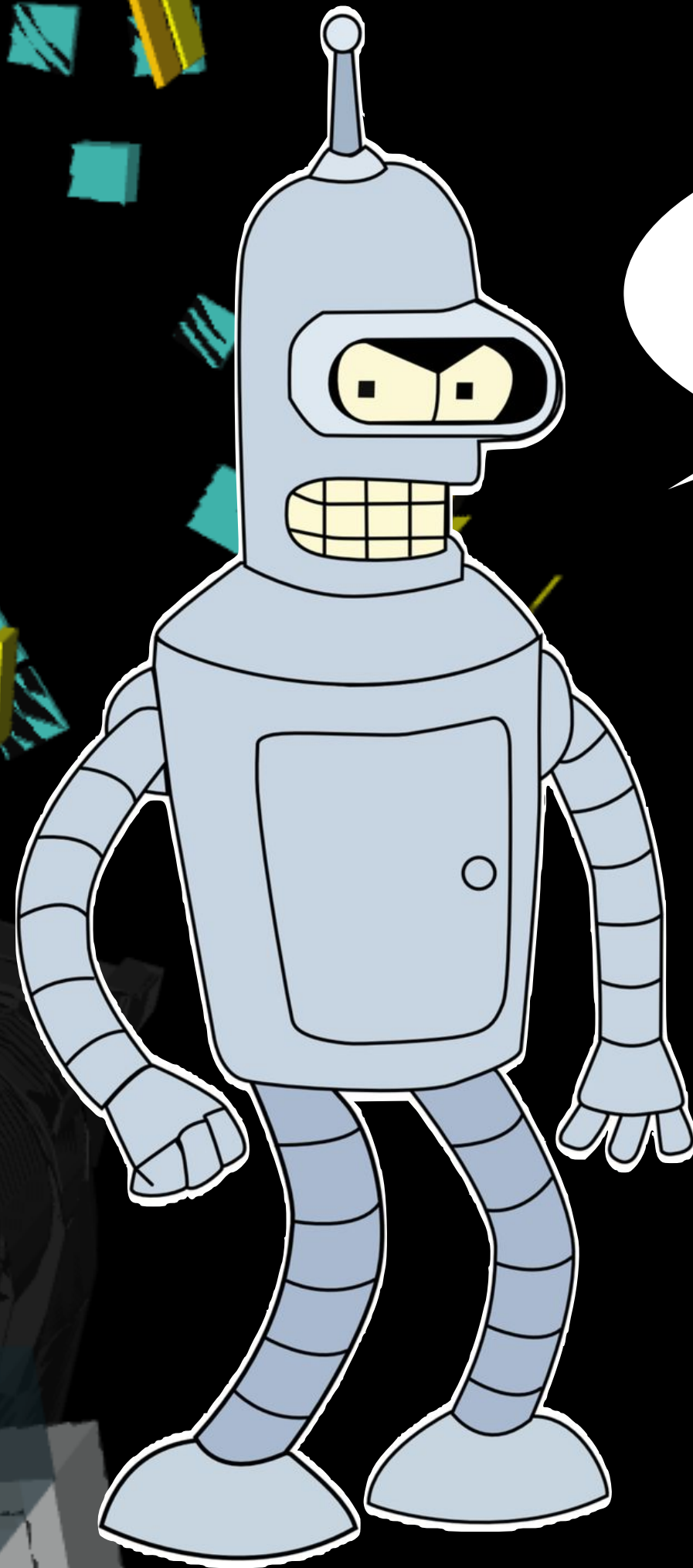
$$\propto 0.07 \kappa_t^2 + 1.59 \kappa_W^2 - 0.66 \kappa_W \kappa_t$$

Under a given new-physics hypothesis, say for example κ framework, it is always possible to combine direct & indirect probes to strengthen constraints

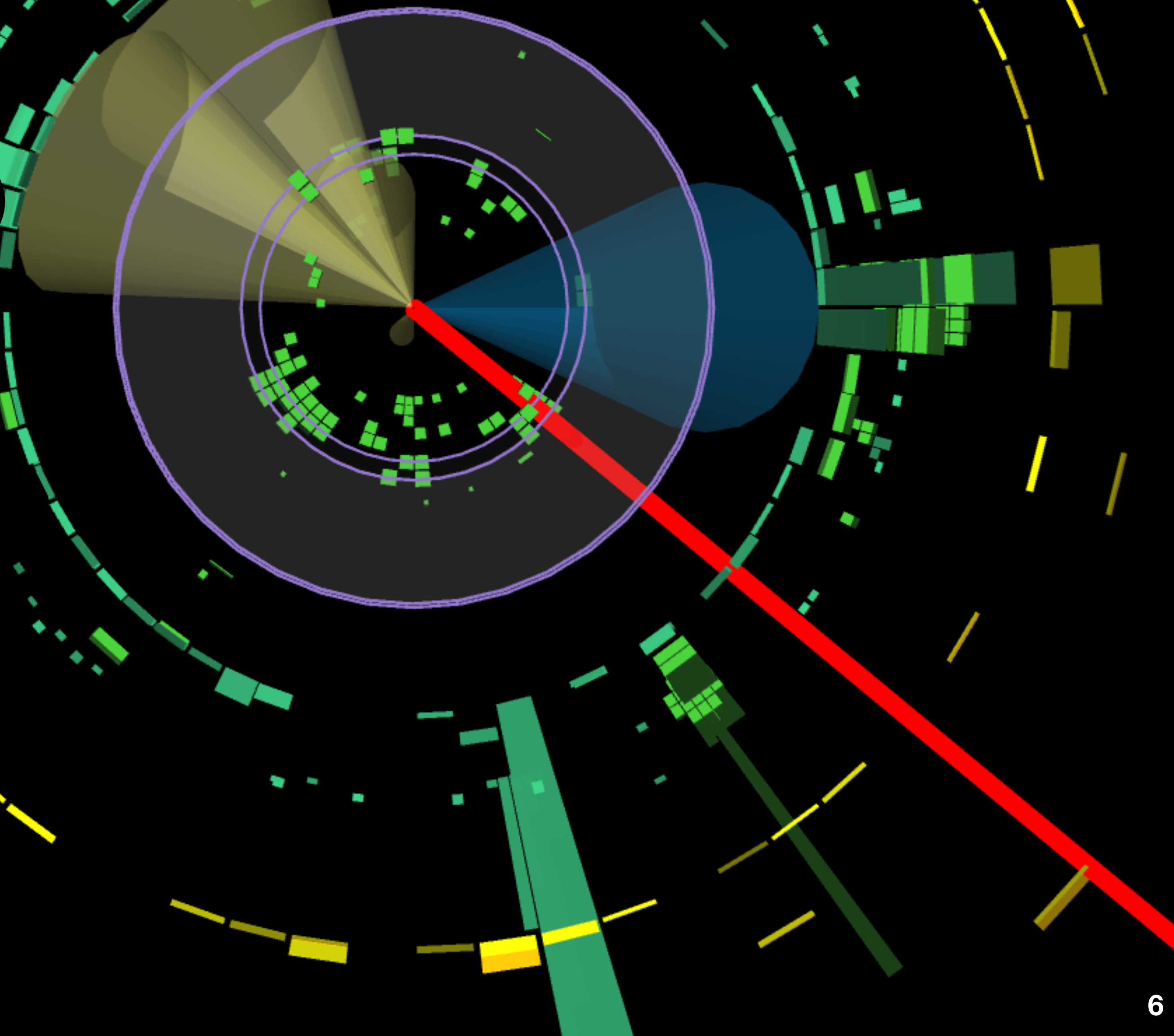
(In)direct probes: synergy & complementarity



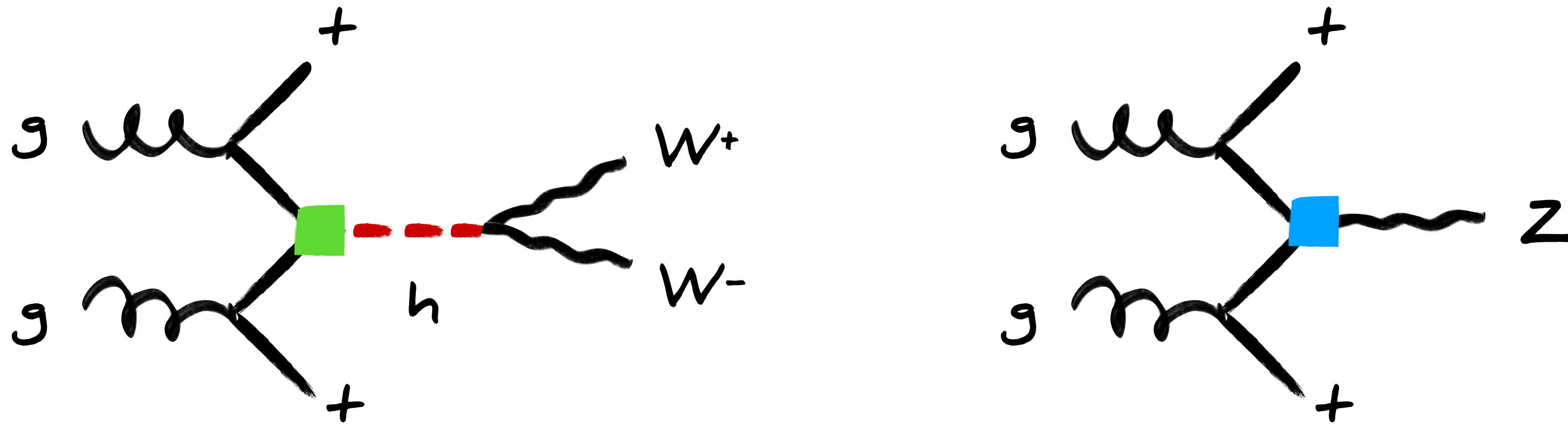
(In)direct probes: experimentalist's view



What the f@*#
is going on?

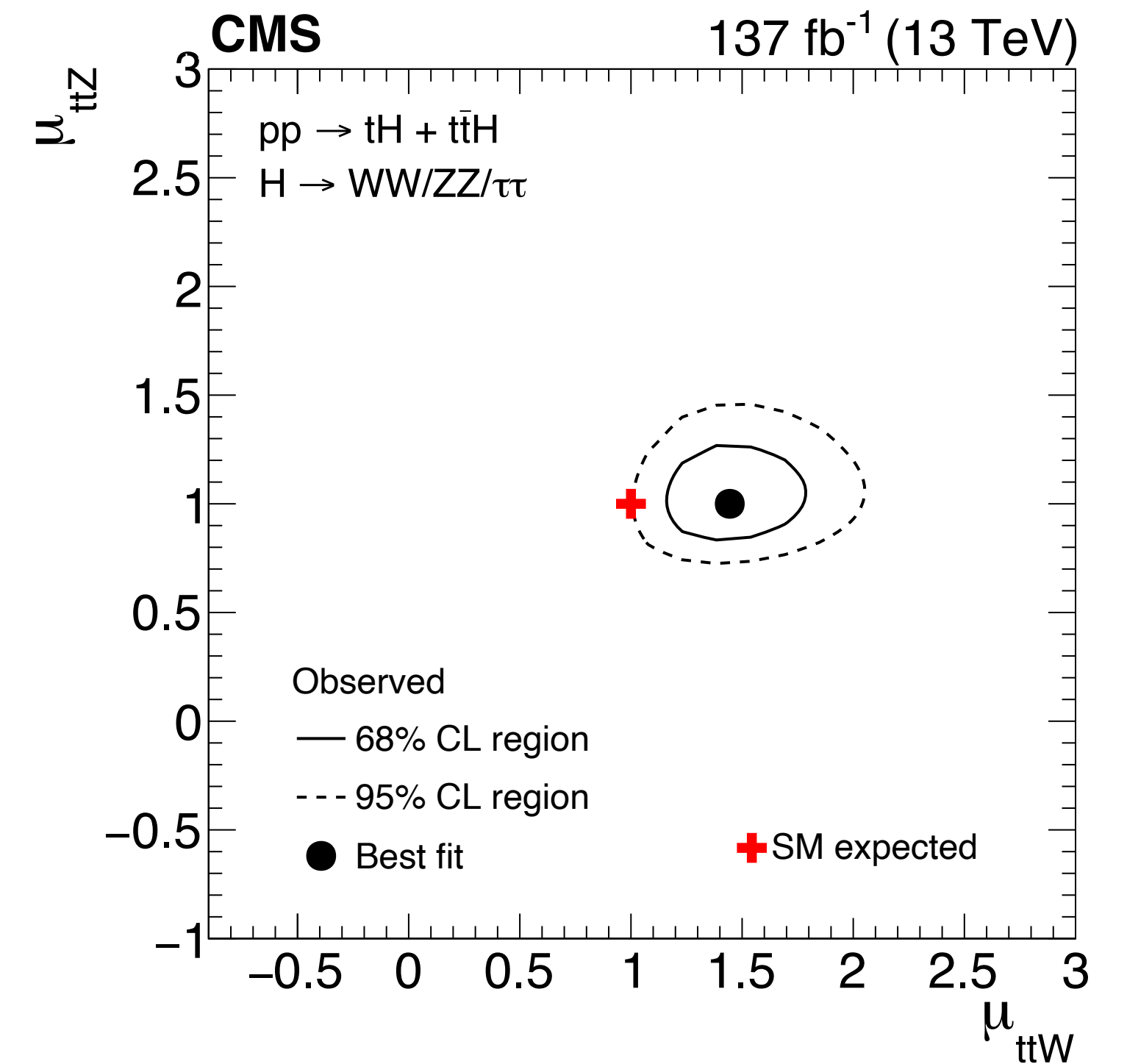
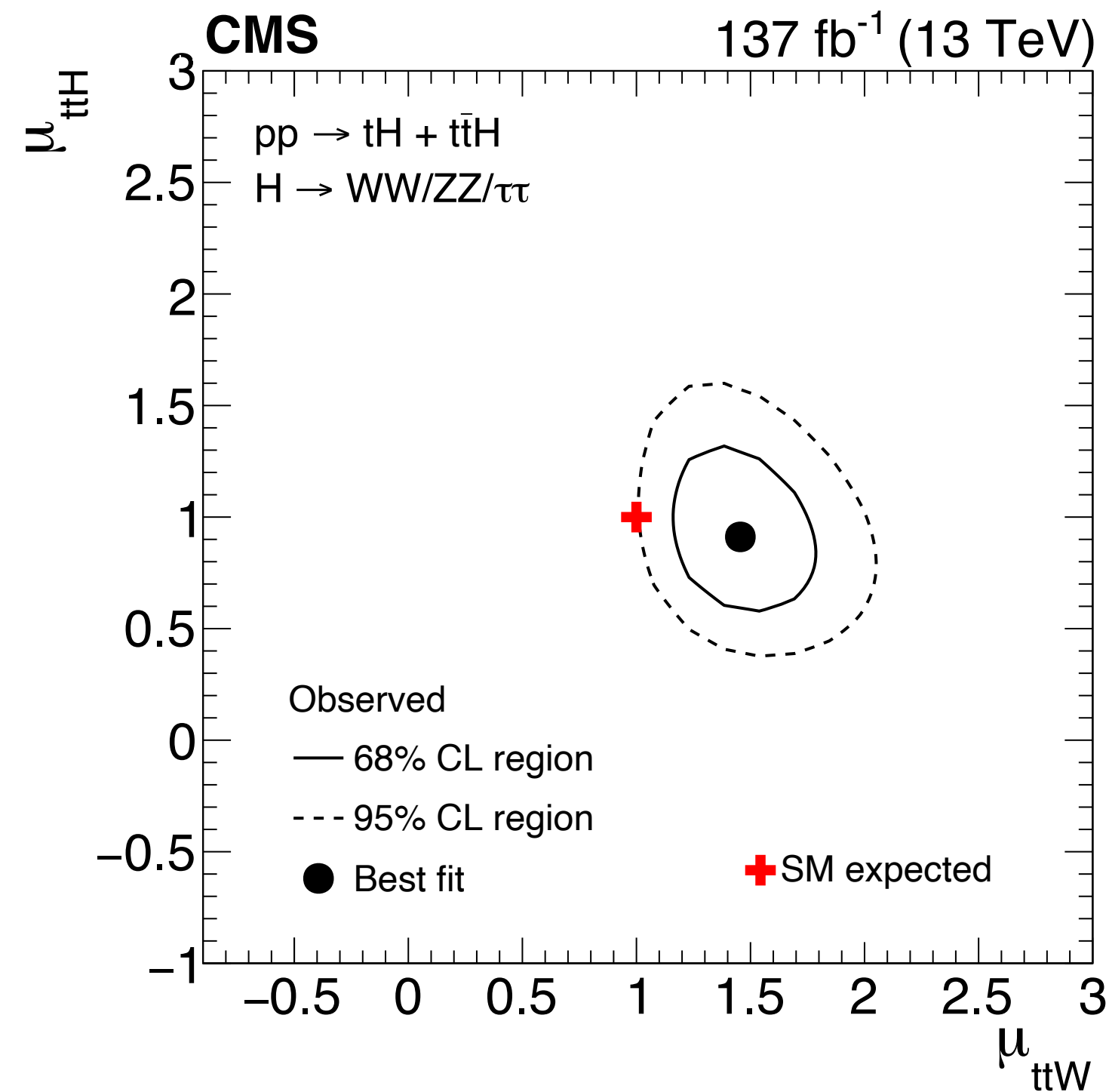
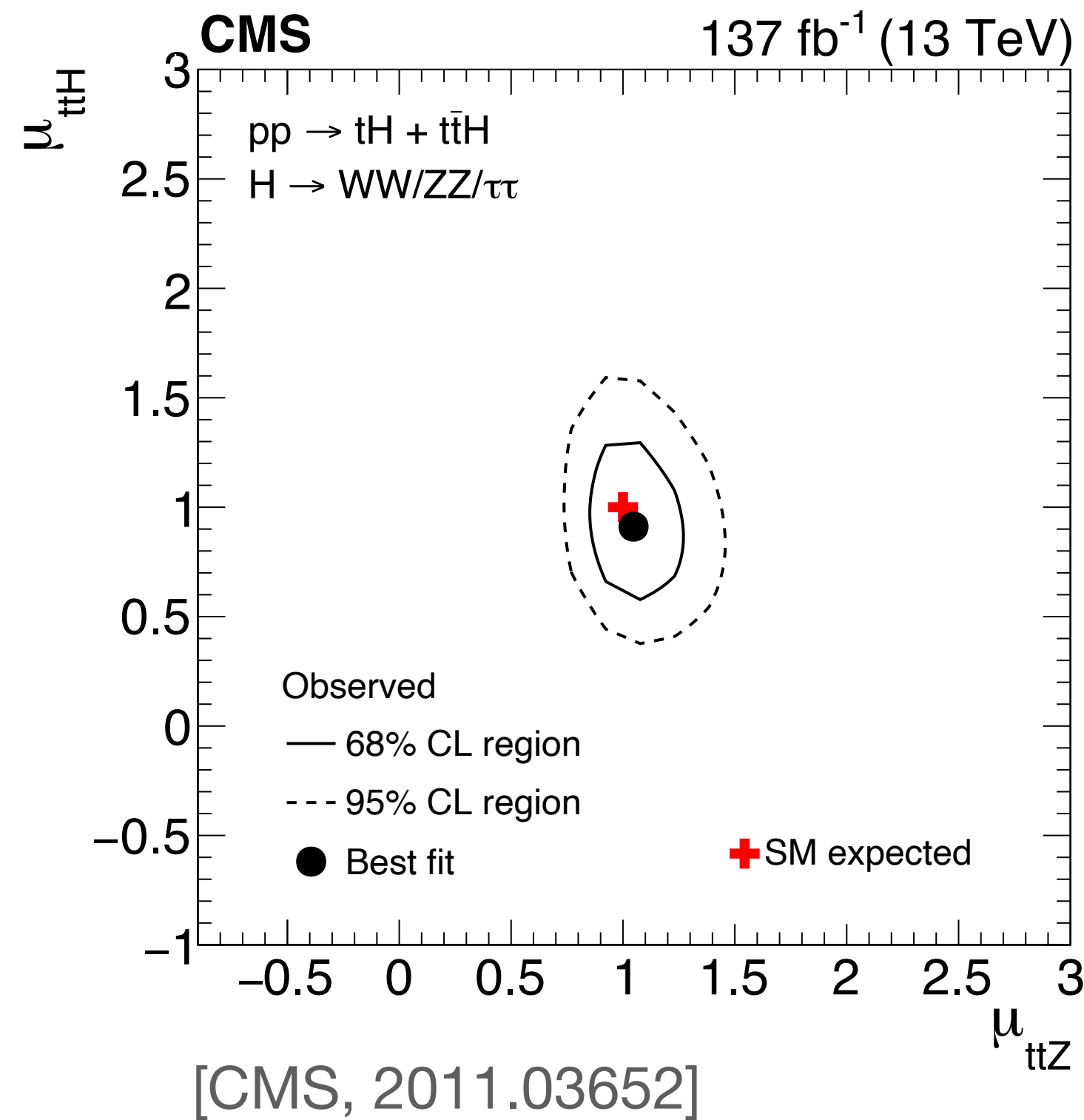


(In)direct probes: experimentalist's view



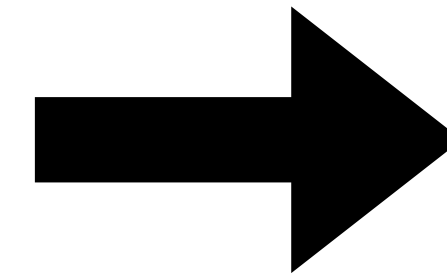
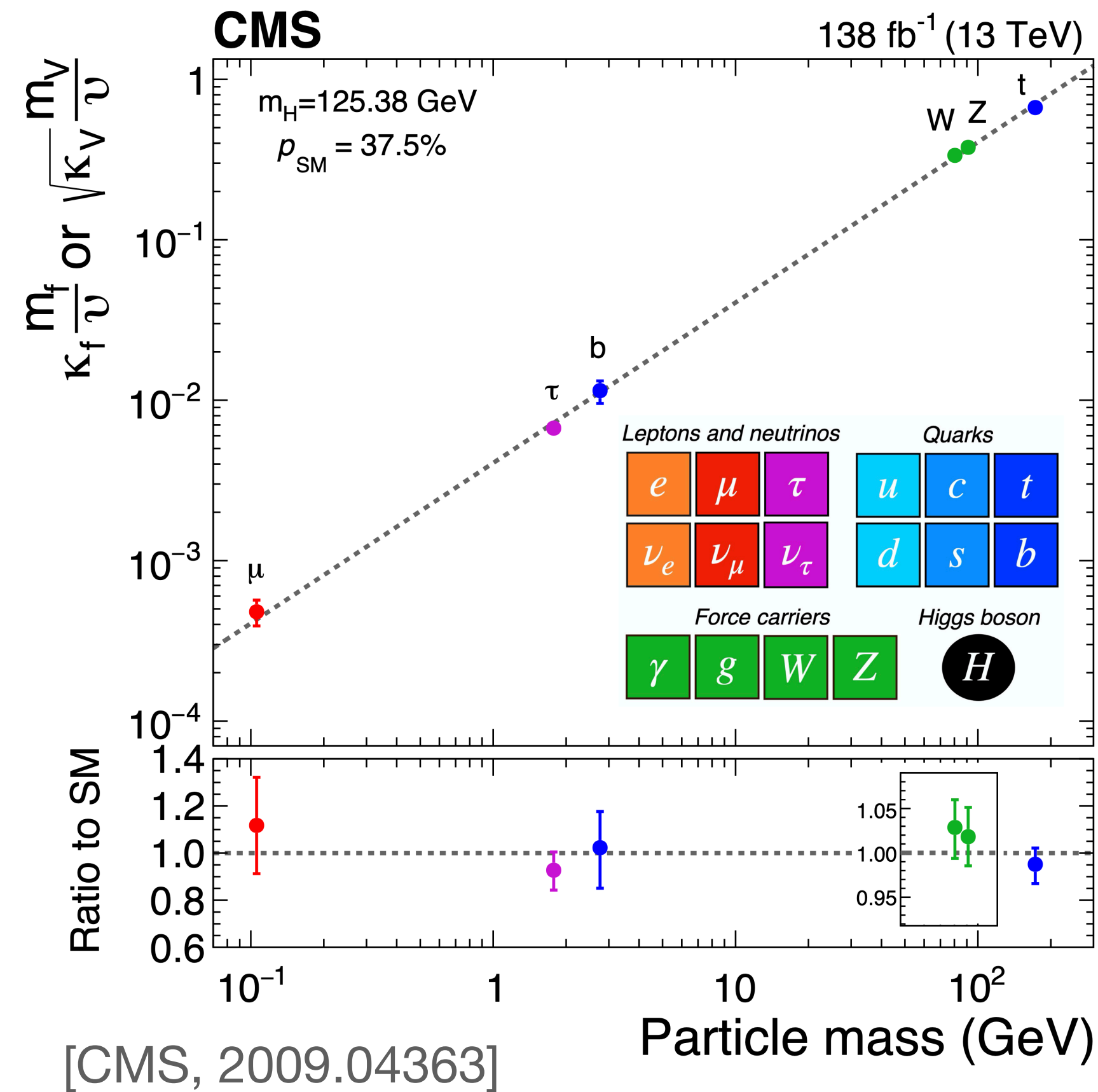
Important constraint on $t\bar{t}h$ production provided by multilepton search, resulting from Higgs decays to W^+W^- , ZZ & $\tau^+\tau^-$. Background from $t\bar{t}Z$ production, which depends on Z-boson coupling to top quarks, which is not very well known

(In)direct probes: experimentalist's view



Extraction of y_t in $t\bar{t}h$ production direct only under model assumptions. Consider $t\bar{t}h$, $t\bar{t}Z$, $t\bar{t}W$, etc. together to extract bounds on y_t , Z-boson coupling to top-quarks, etc.

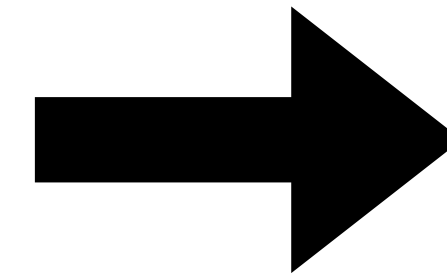
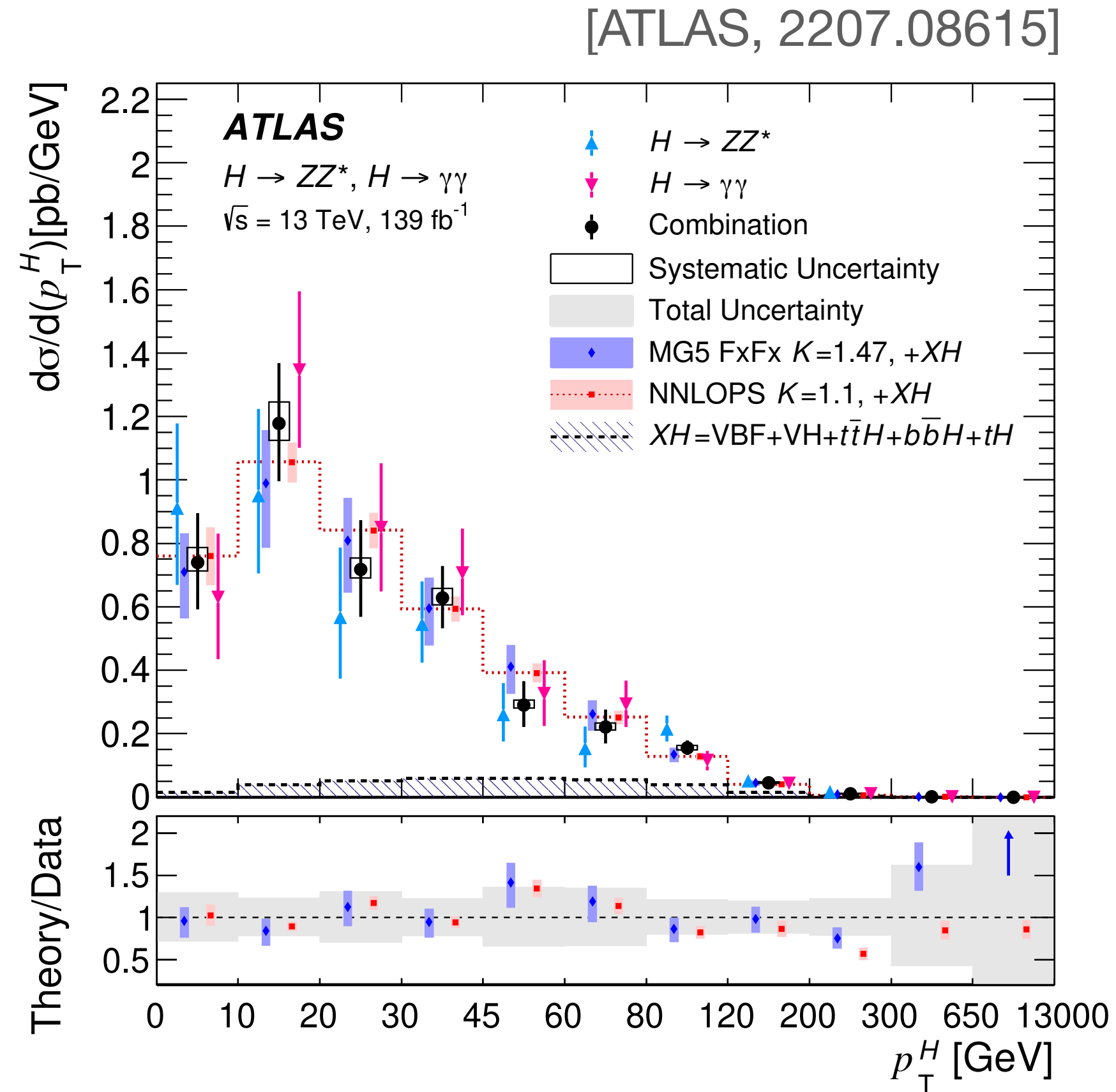
When are indirect probes helpful?



Yukawa couplings of
light fermions, Higgs
self-couplings, ...

If present LHC data is only able to put very loose bounds on a given coupling

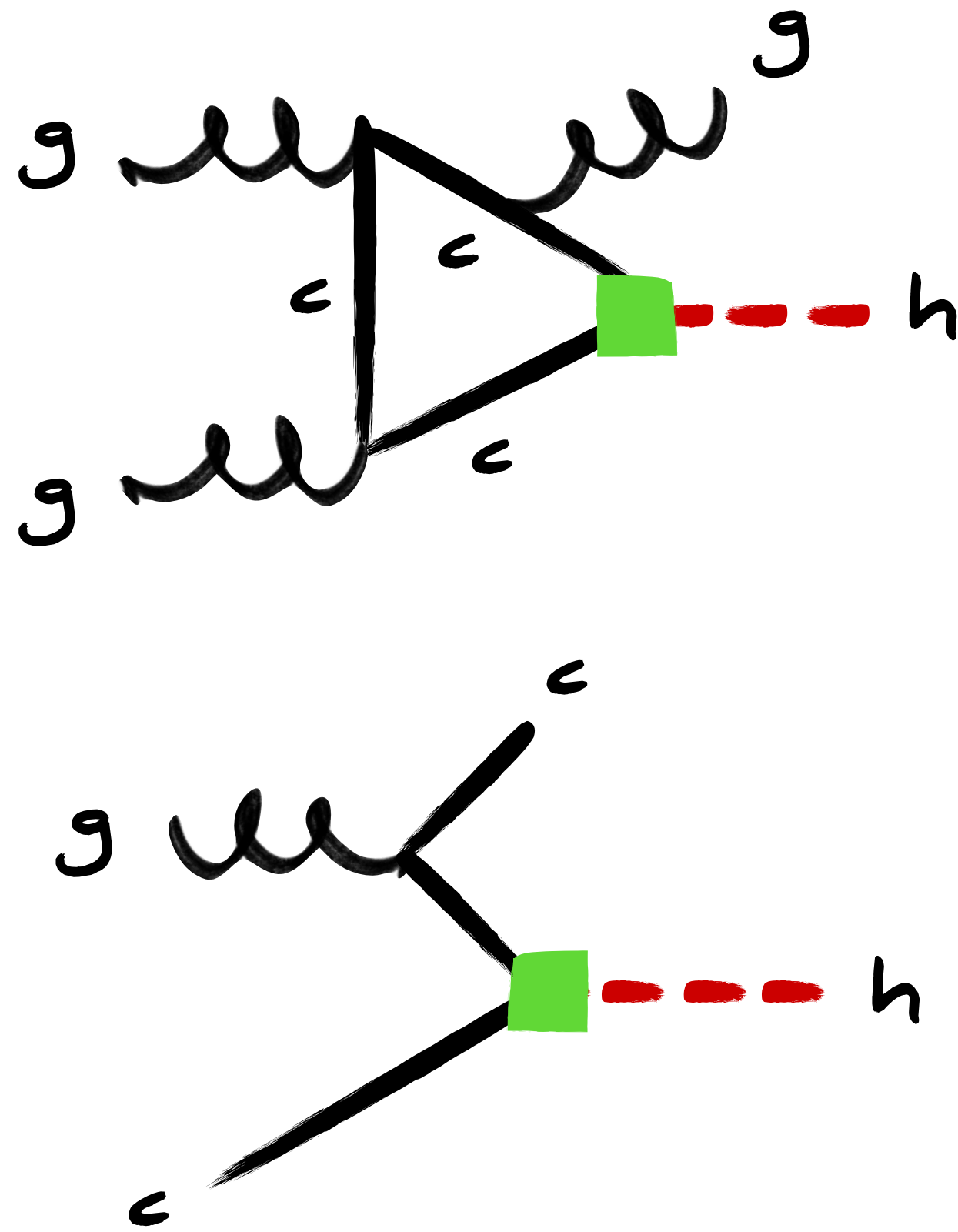
When are indirect probes helpful?



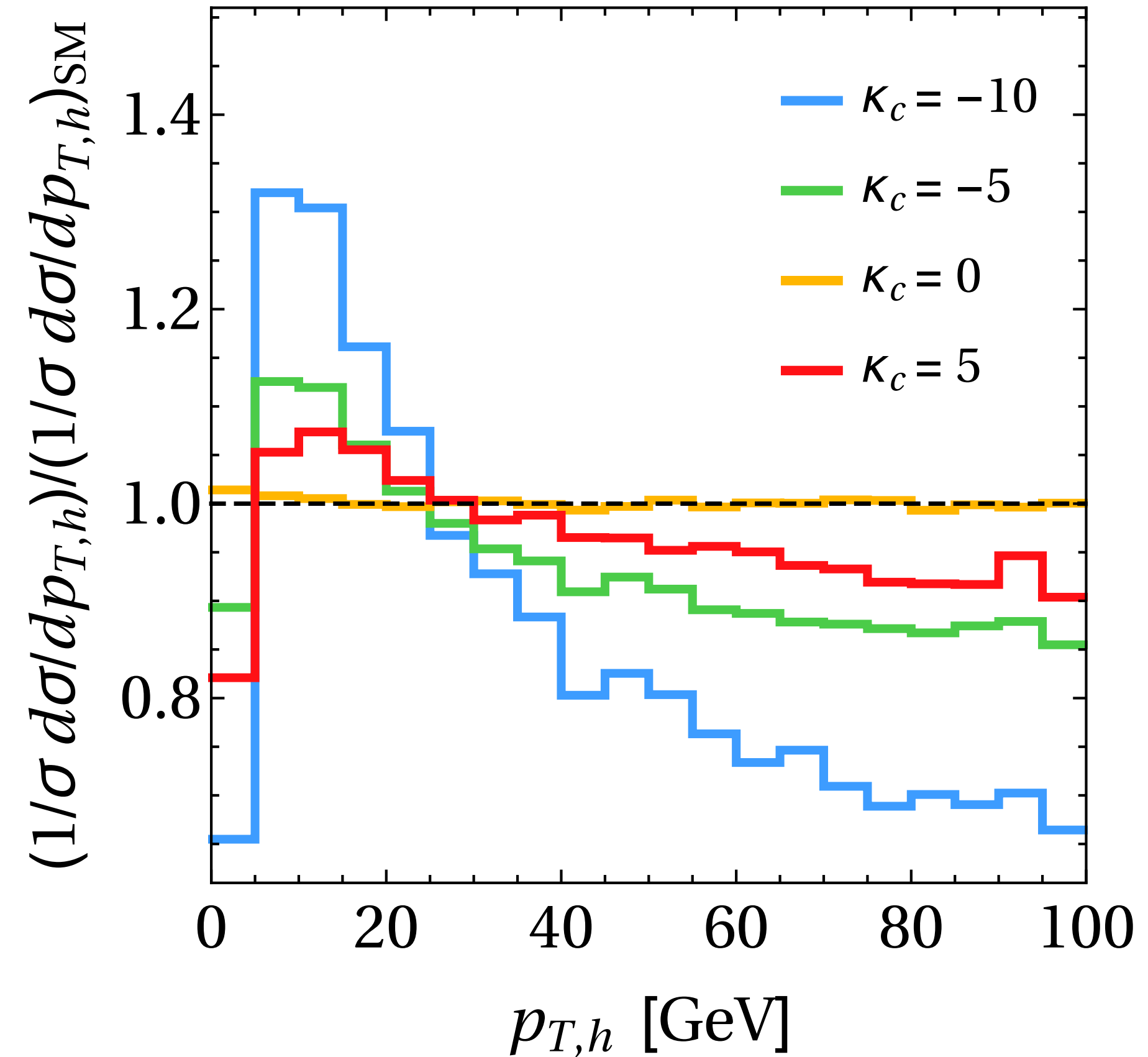
transverse momentum (p_{\perp})
spectrum of Higgs, Higgs
off-shell measurements, ...

If LHC able to measure a theoretically clean, differential Higgs observable precisely

Charm Yukawa coupling (y_c) from Higgs p_T



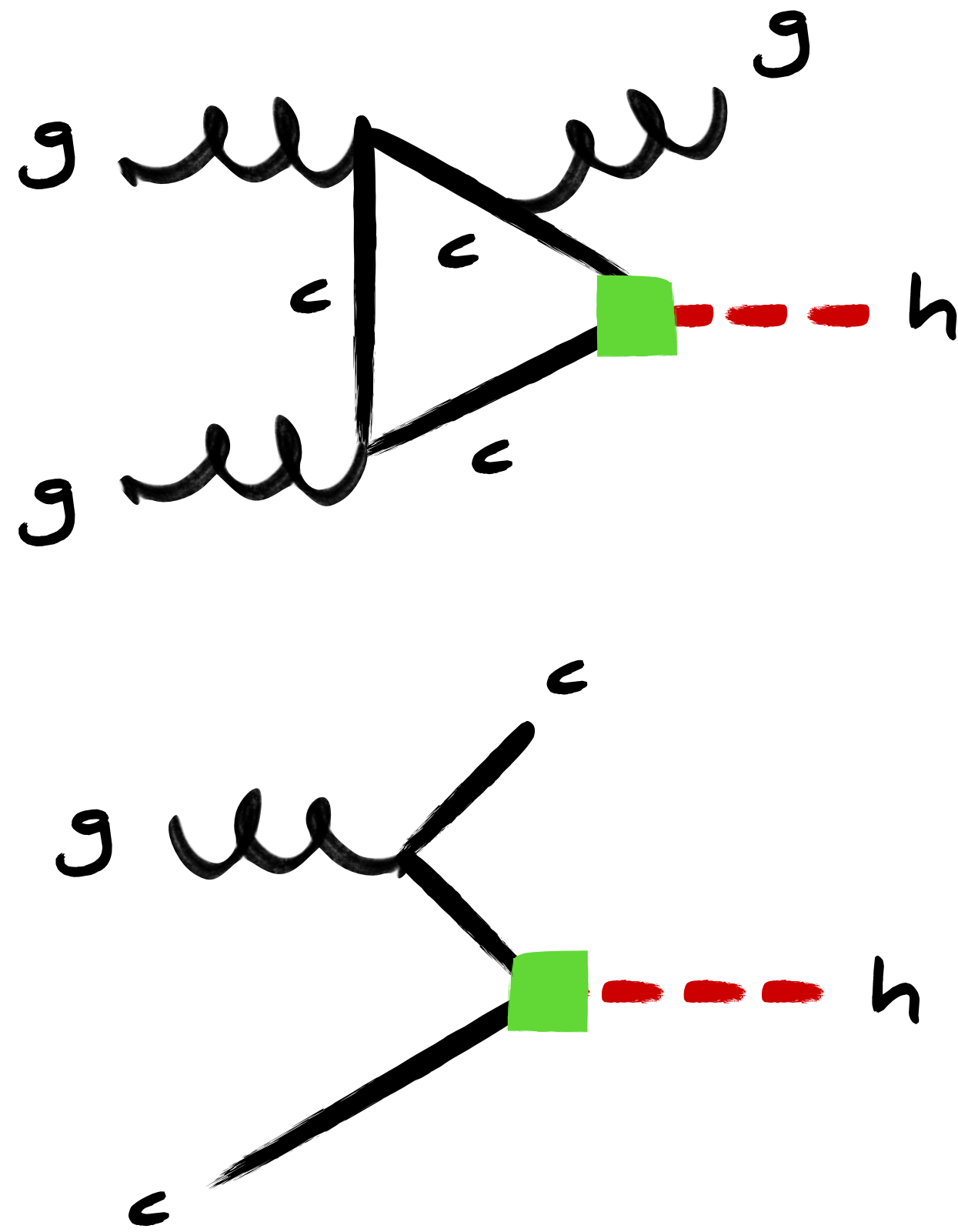
[Bishara, UH, Monni & Re, 1606.09253]



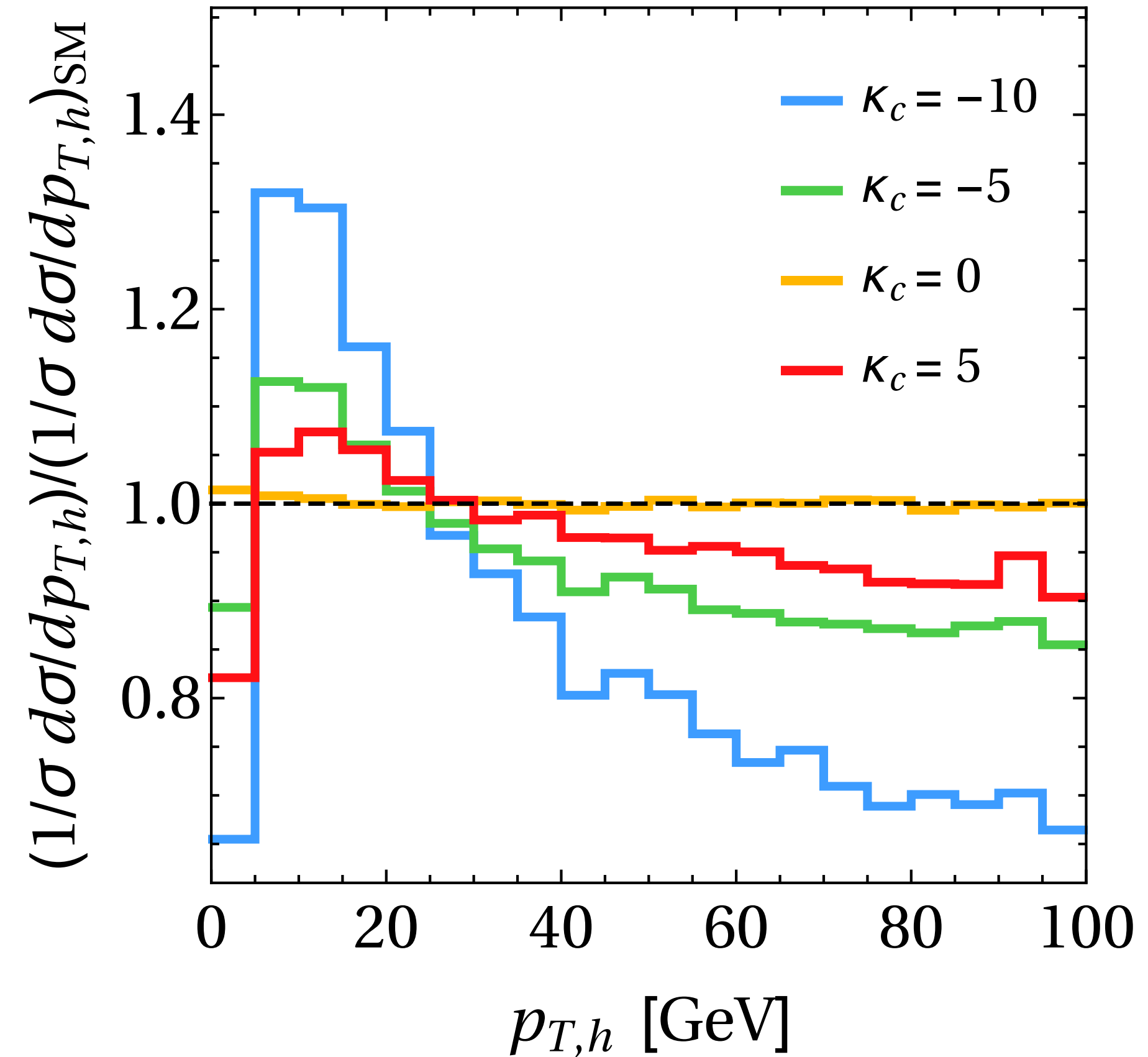
Emissions from charm loops or charm-initiated production distort $p_{T,h}$ spectrum

[see also Soreq, Zhu & Zupan, 1606.09621]

Charm Yukawa coupling (y_c) from Higgs p_T



[Bishara, UH, Monni & Re, 1606.09253]

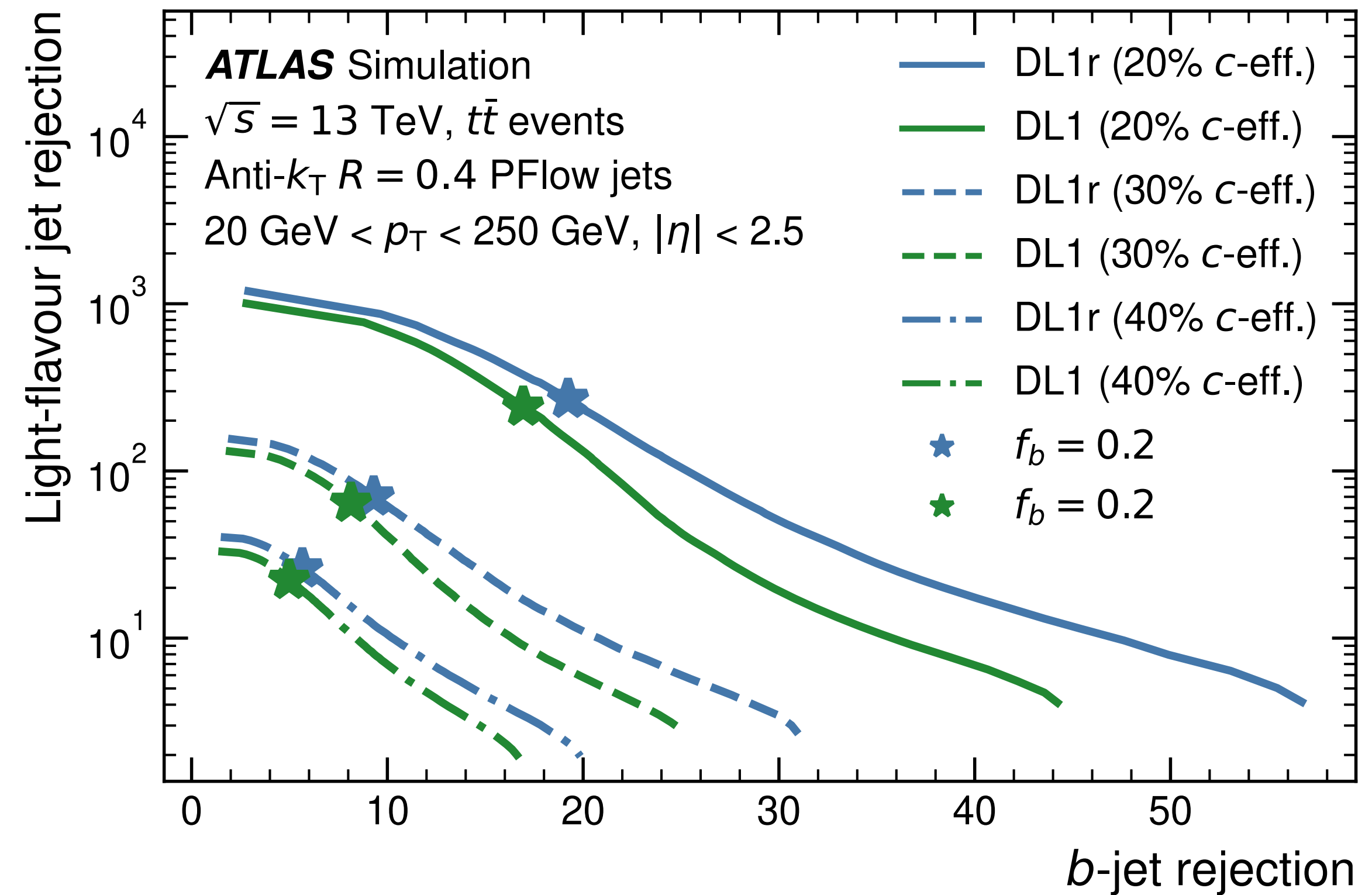
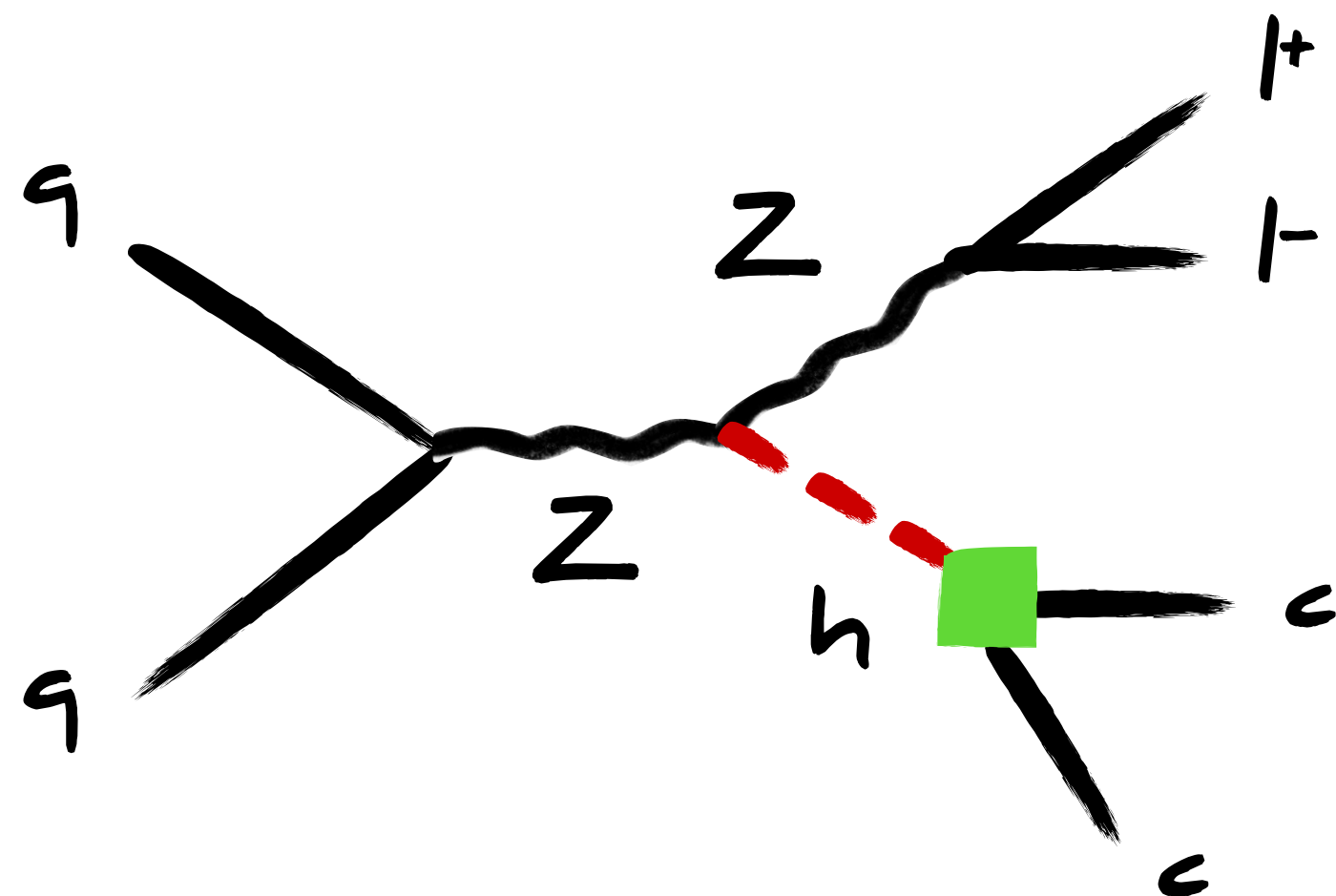


Same idea applies in light-quark case, but detectable effects only for very large κ_q

[see also Soreq, Zhu & Zupan, 1606.09621]

Constraints on y_c from Vh production

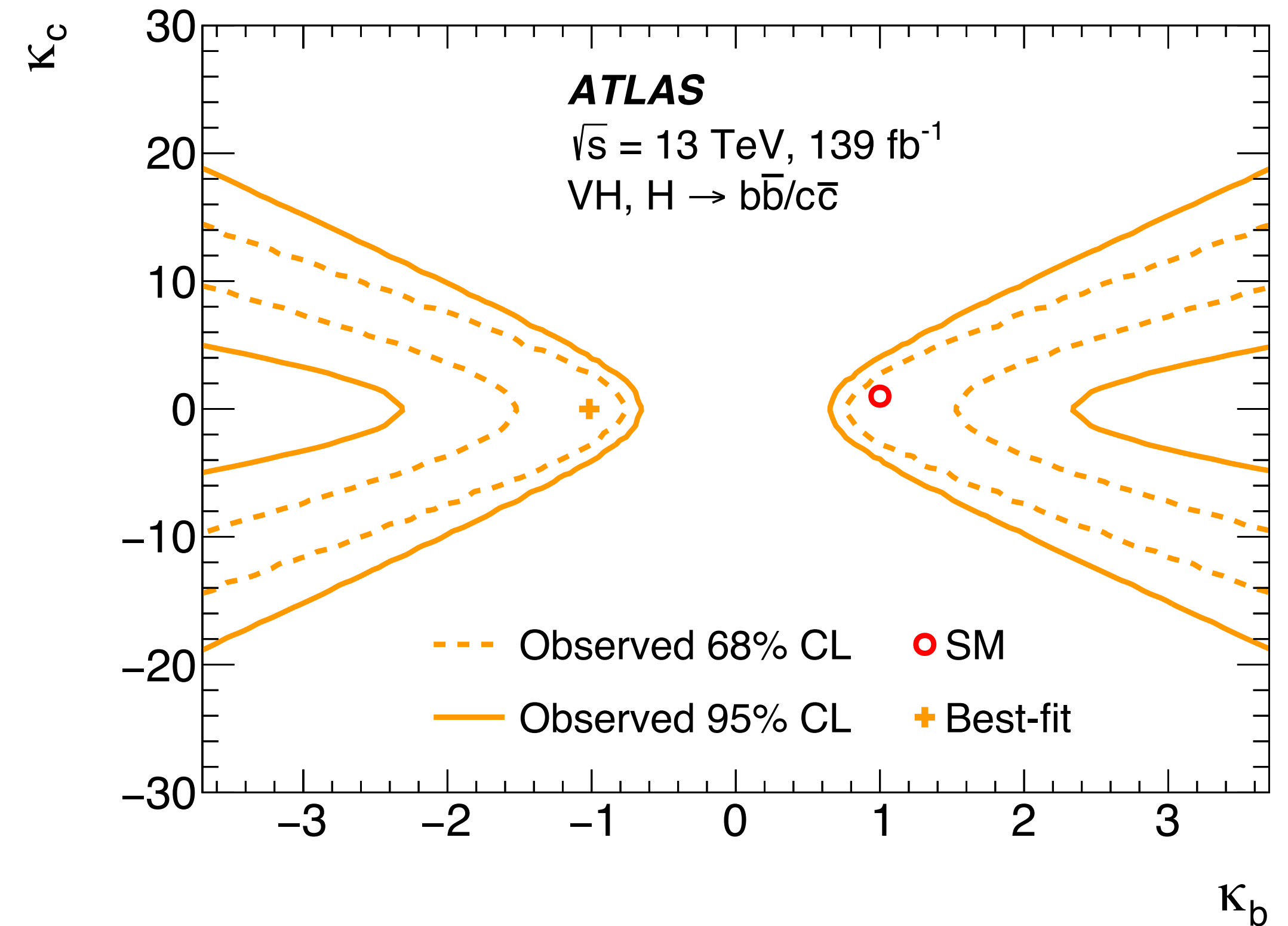
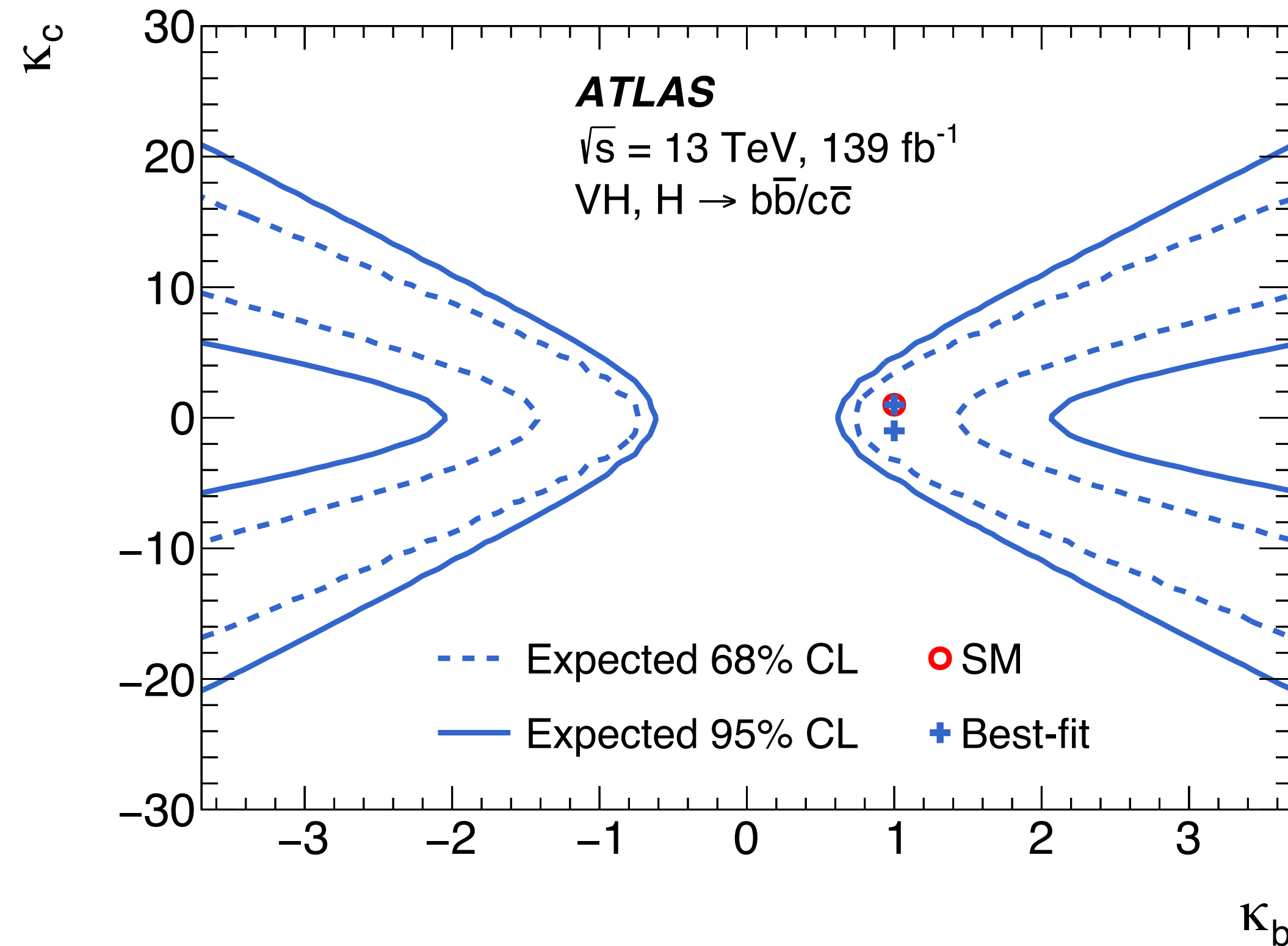
[ATLAS, 2211.16345]



Charm-tagging enables to probe y_c directly @ LHC in associated Vh production

Direct constraints on κ_b & κ_c @ LHC

[ATLAS, 2201.11428]

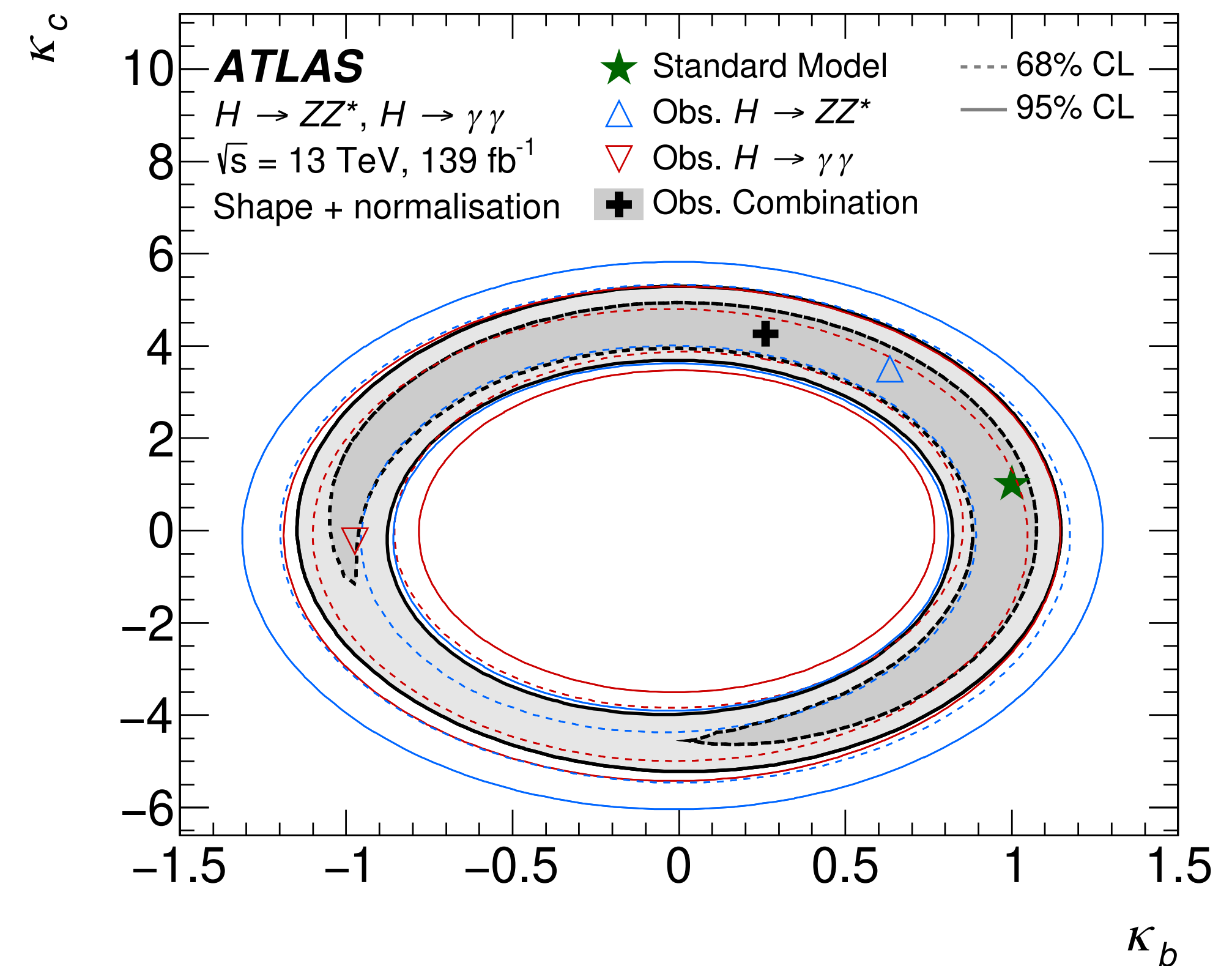
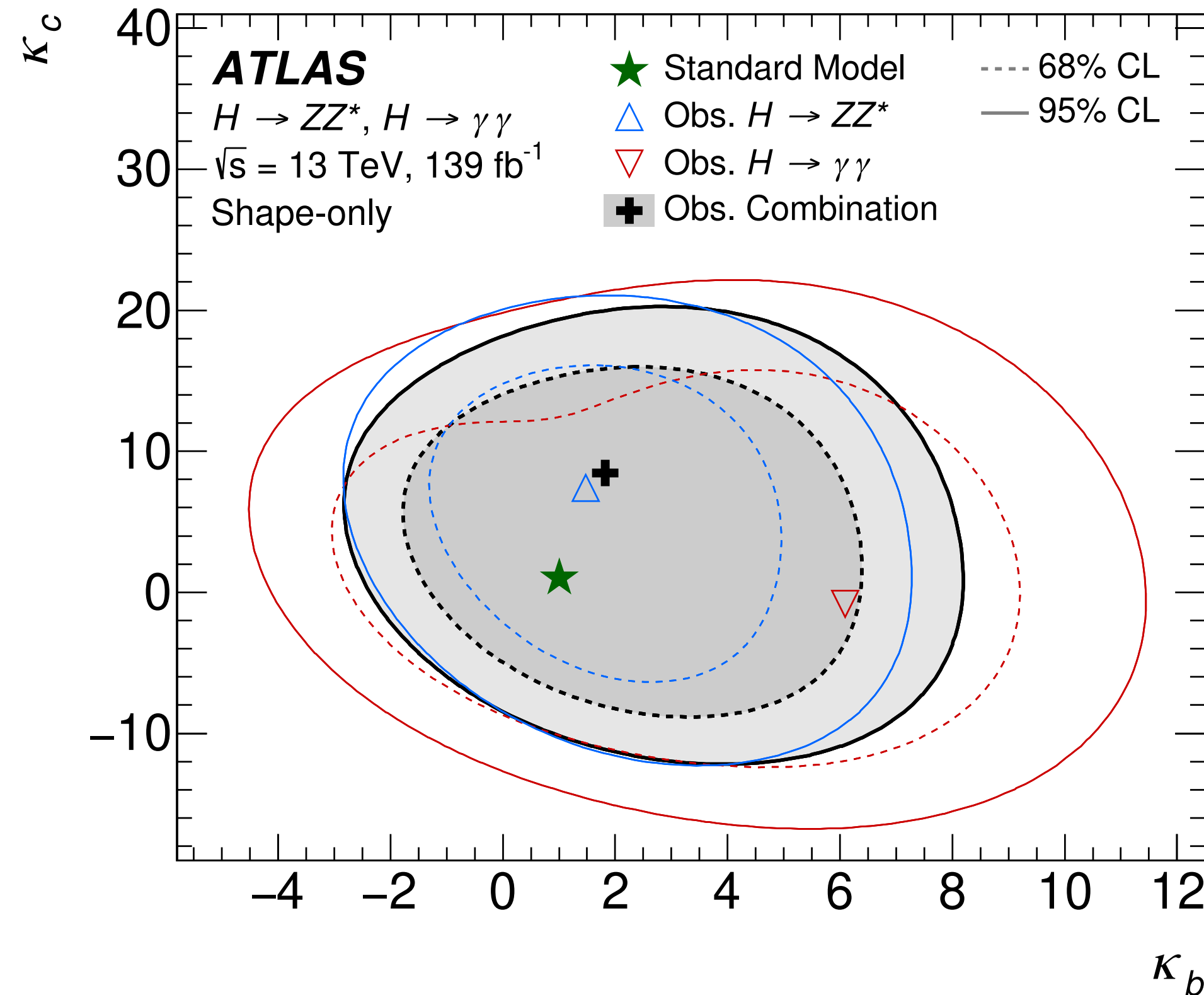


Vh production constrains κ_b & κ_c to two hyperbola-like strips in 2D plane

[see also CMS, 2205.05550; ATLAS, 2410.19611; talks by Missio, Örddek & Wuchterl]

Indirect constraints on κ_b & κ_c @ LHC

[ATLAS, 2207.08615]

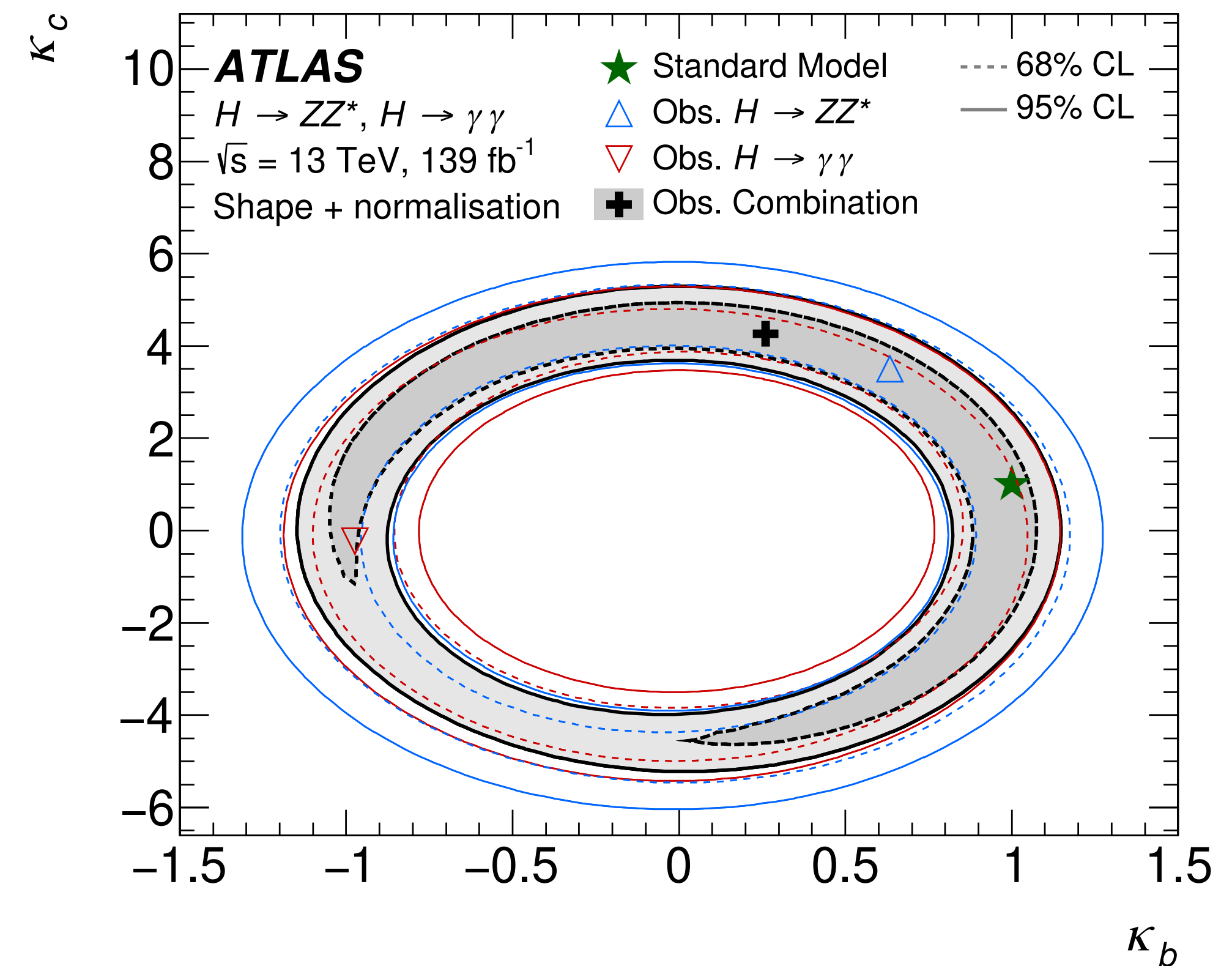
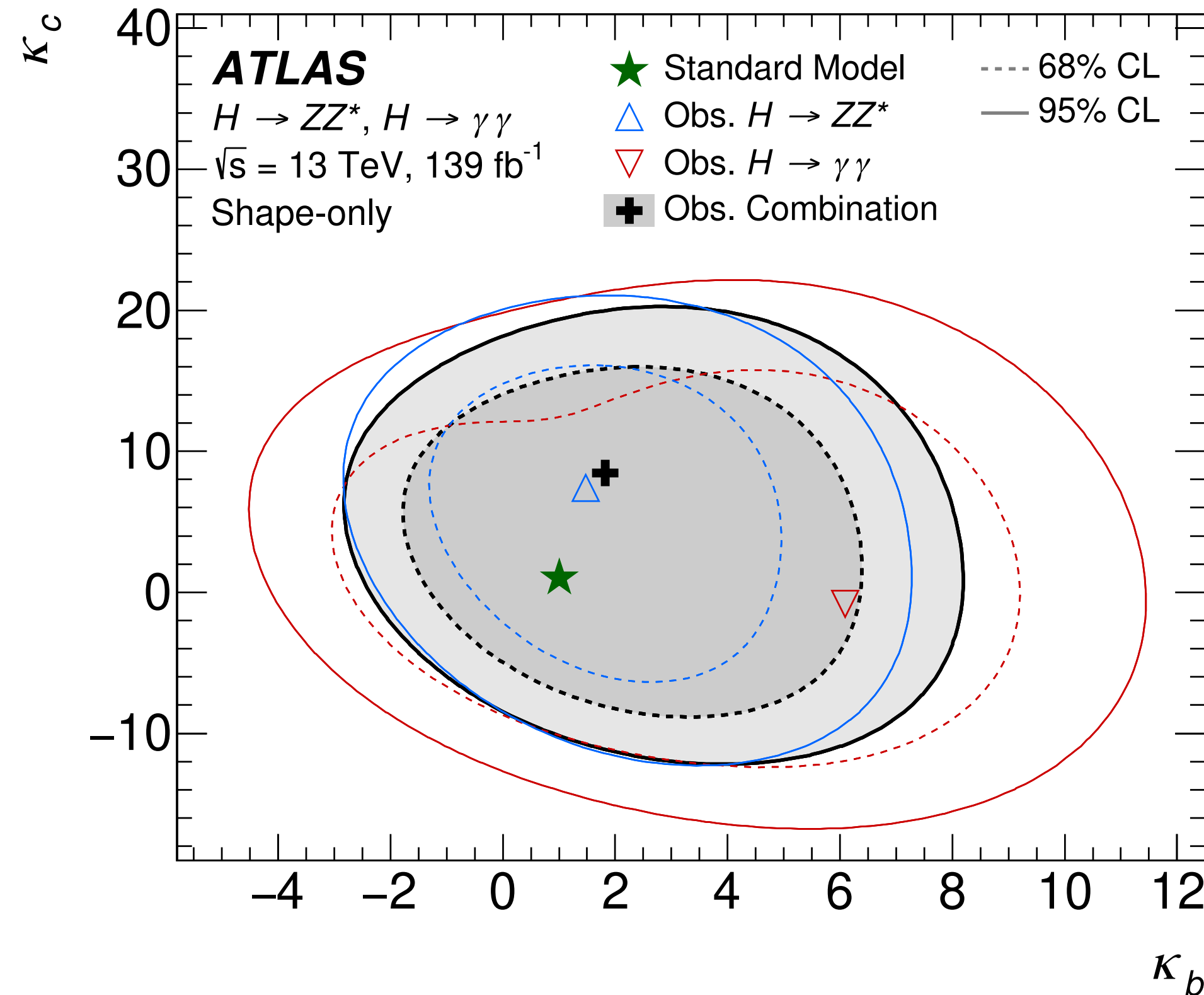


At present, shape changes in $p_{T,h}$ spectrum lead only to weak oval exclusions

[see also CMS, 2305.07532; CMS-PAS-HIG-23-013; talks by Galli, Winterbottom & Wuchterl]

Indirect constraints on κ_b & κ_c @ LHC

[ATLAS, 2207.08615]

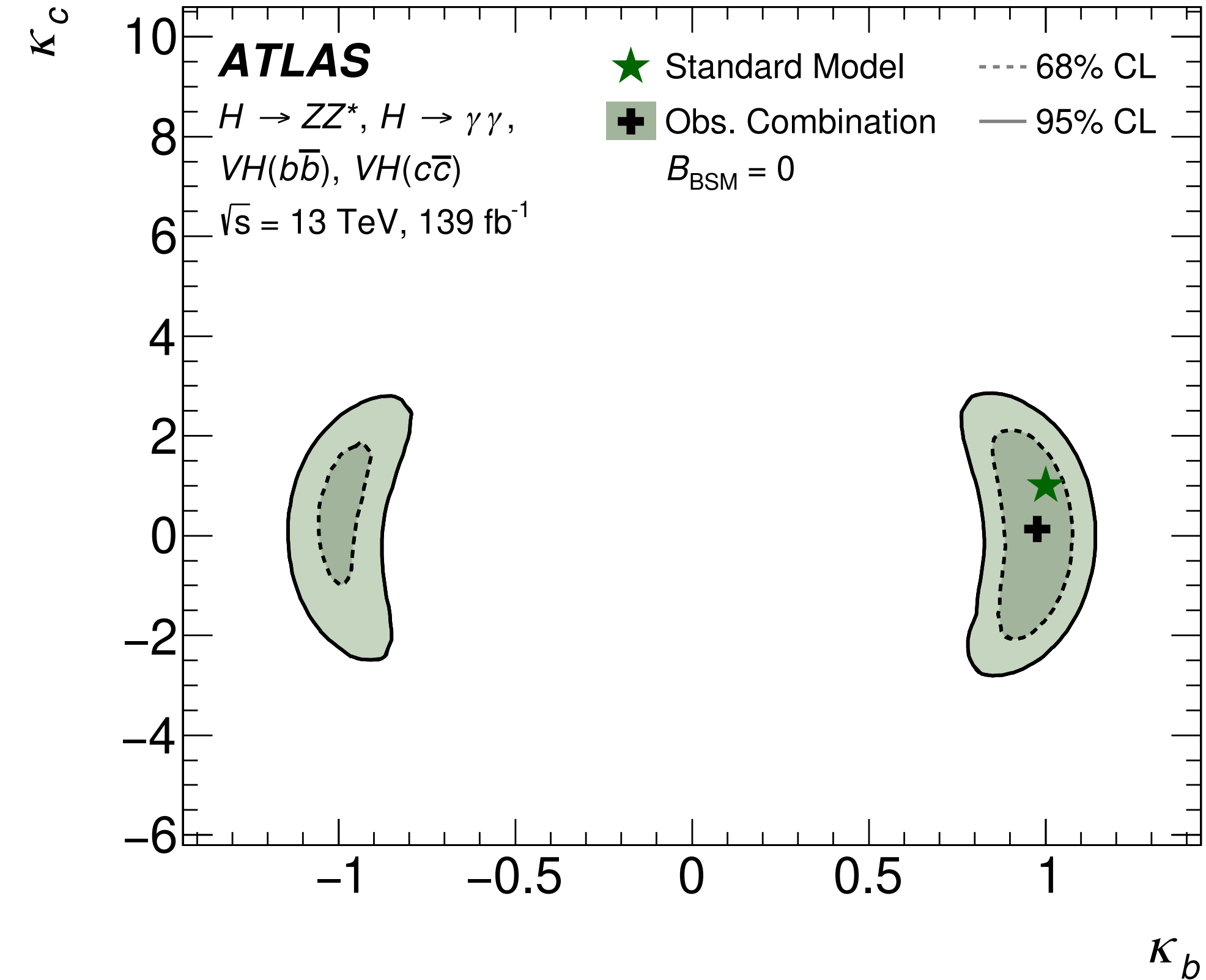
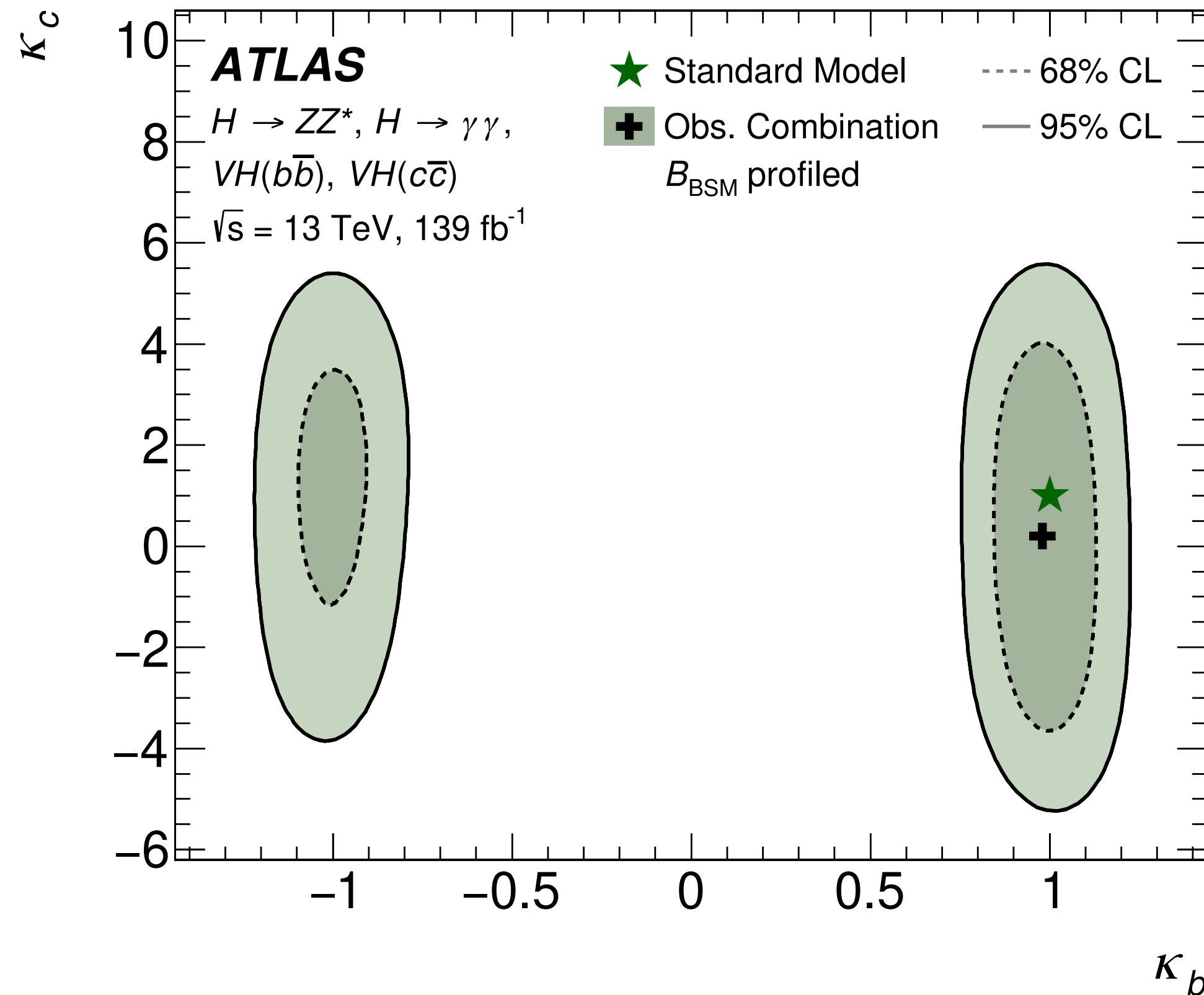


Adding information on normalisation, constrains κ_b & κ_c to elliptic band in 2D plane

[see also CMS, 2305.07532; CMS-PAS-HIG-23-013; talks by Galli, Winterbottom & Wuchterl]

Combined constraints on κ_b & κ_c @ LHC

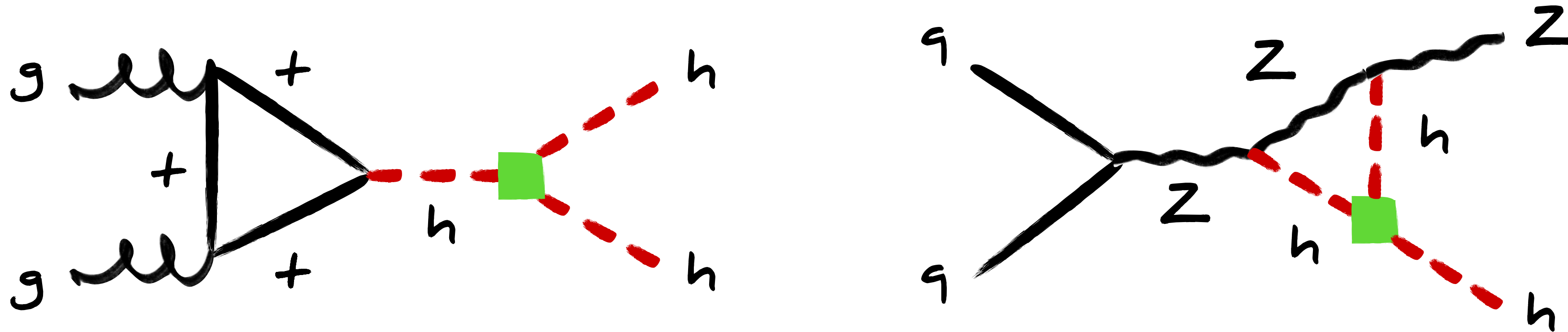
[ATLAS, 2207.08615]



Combination of bounds leads to two islands of solution centred around $(1,1)$ & $(-1,-1)$.

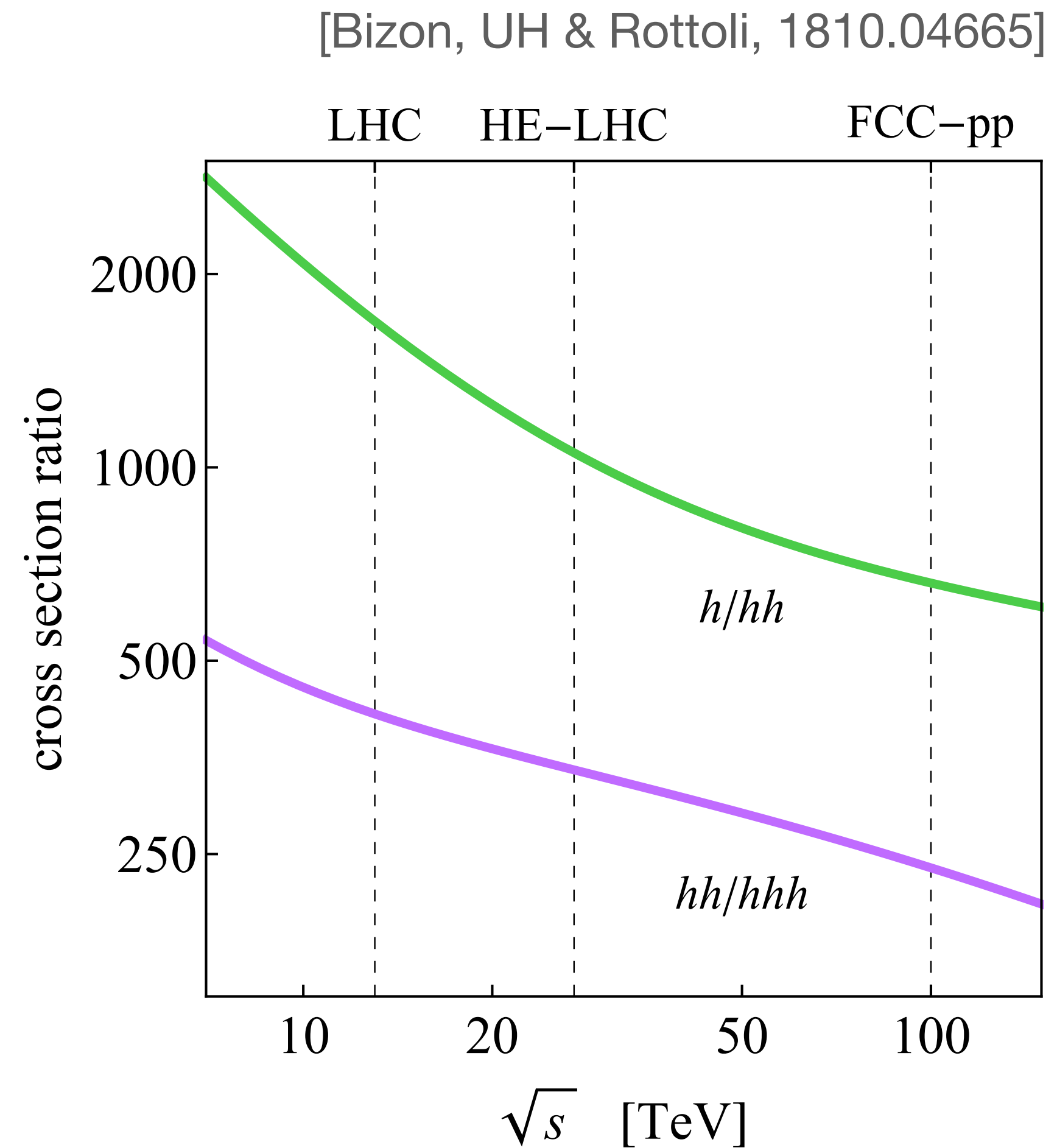
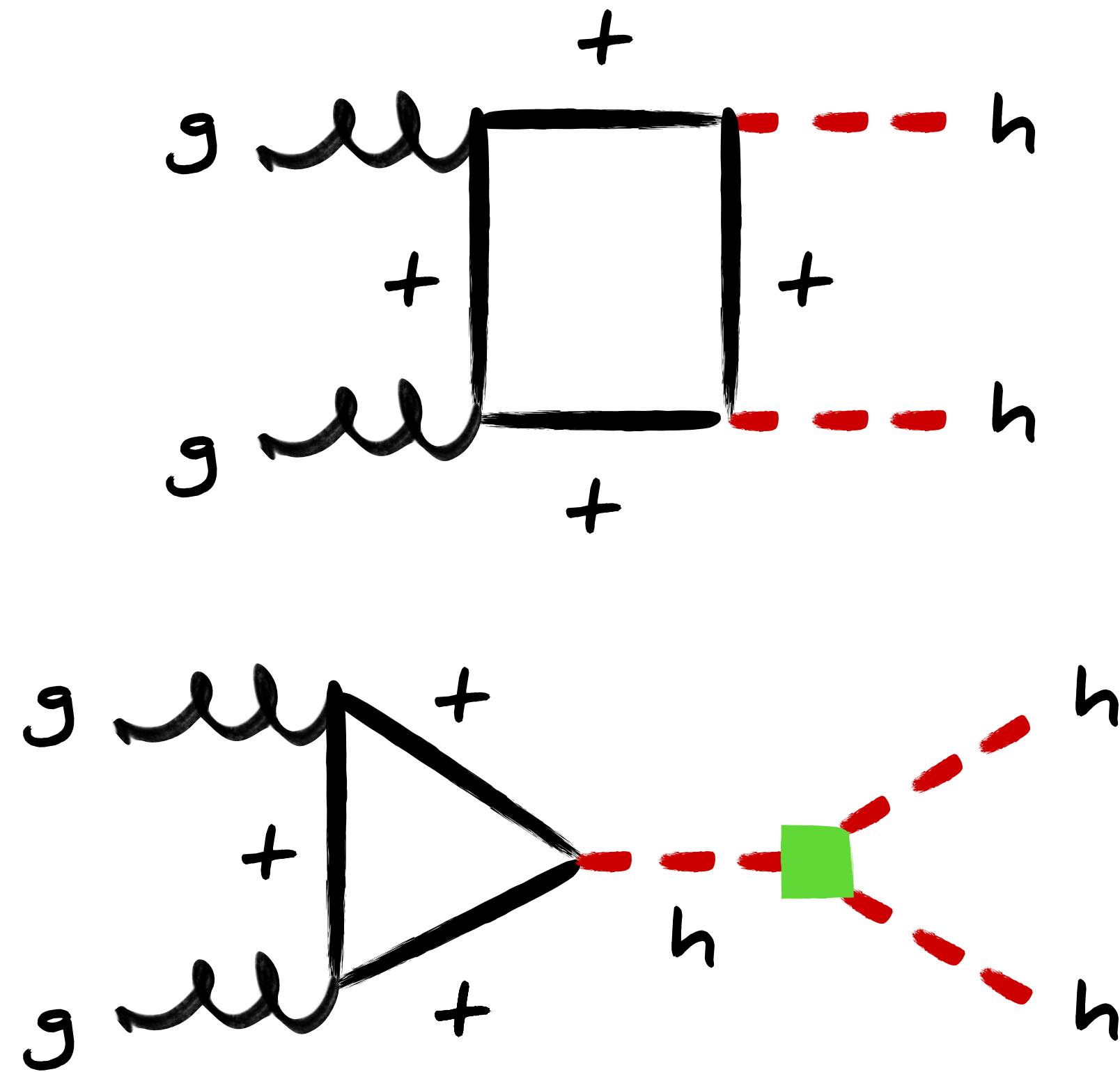
Limits depend on assumption about size of new-physics effects in Higgs width

(In)direct probes of κ_λ



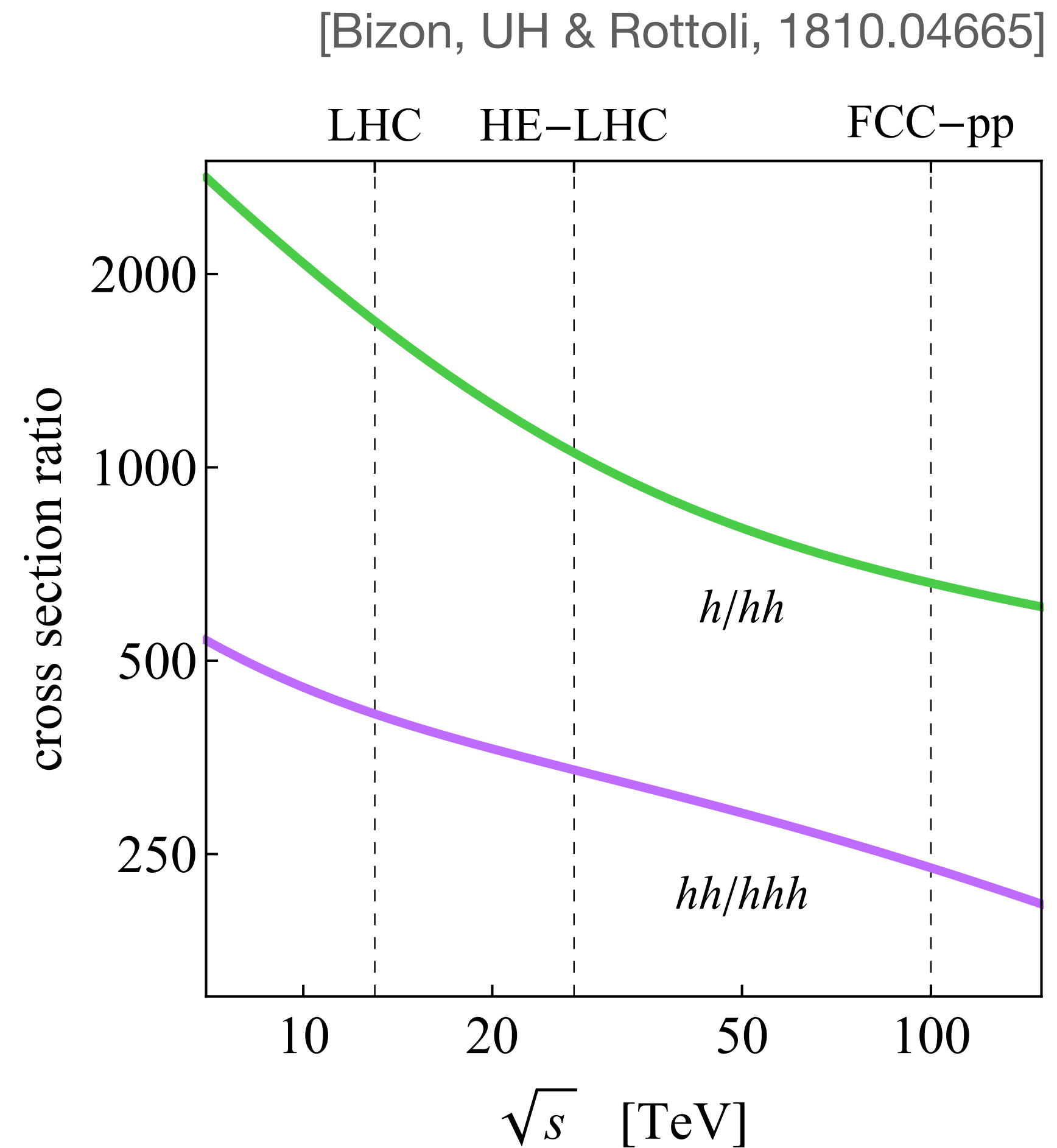
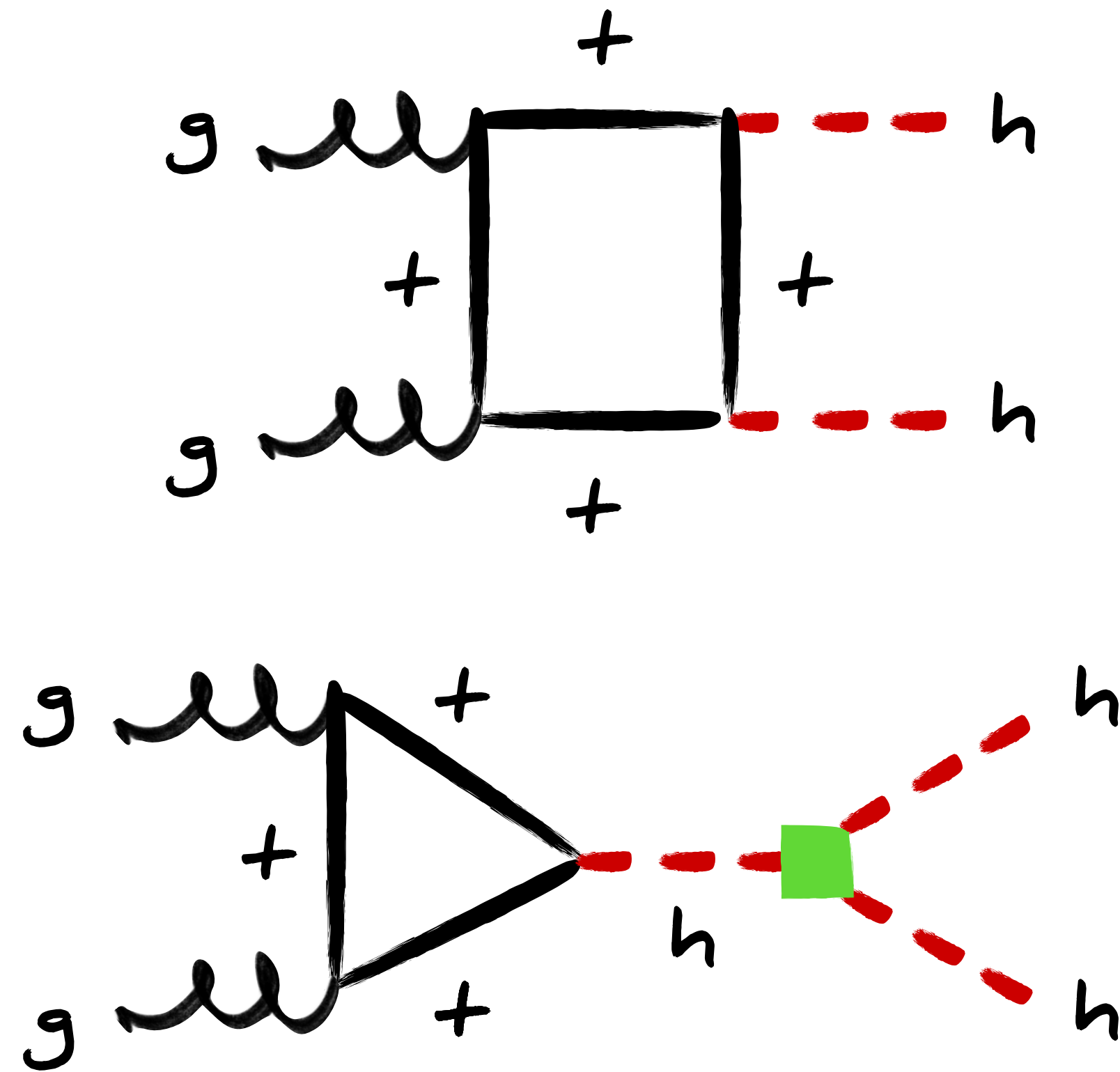
Direct probe of trilinear Higgs self-coupling (κ_λ) provided by hh production, while all single- h production & decay channels @ LHC become indirectly sensitive to κ_λ at 1-loop (Vh , VBF , $t\bar{t}h$, $h \rightarrow b\bar{b}$, etc.) or 2-loop level (ggF , $h \rightarrow \gamma\gamma$, etc.)

(In)direct probes of κ_λ



Loop probes of κ_λ can only compete with hh production because of destructive interference between box & triangle contribution in $gg \rightarrow hh$ amplitude

(In)direct probes of κ_λ



Suppression of $pp \rightarrow hh$ rate more pronounced @ LHC than @ HE-LHC or FCC-pp energies, rendering indirect tests of Higgs self-couplings more promising @ LHC

Combined constraints on κ_λ @ LHC

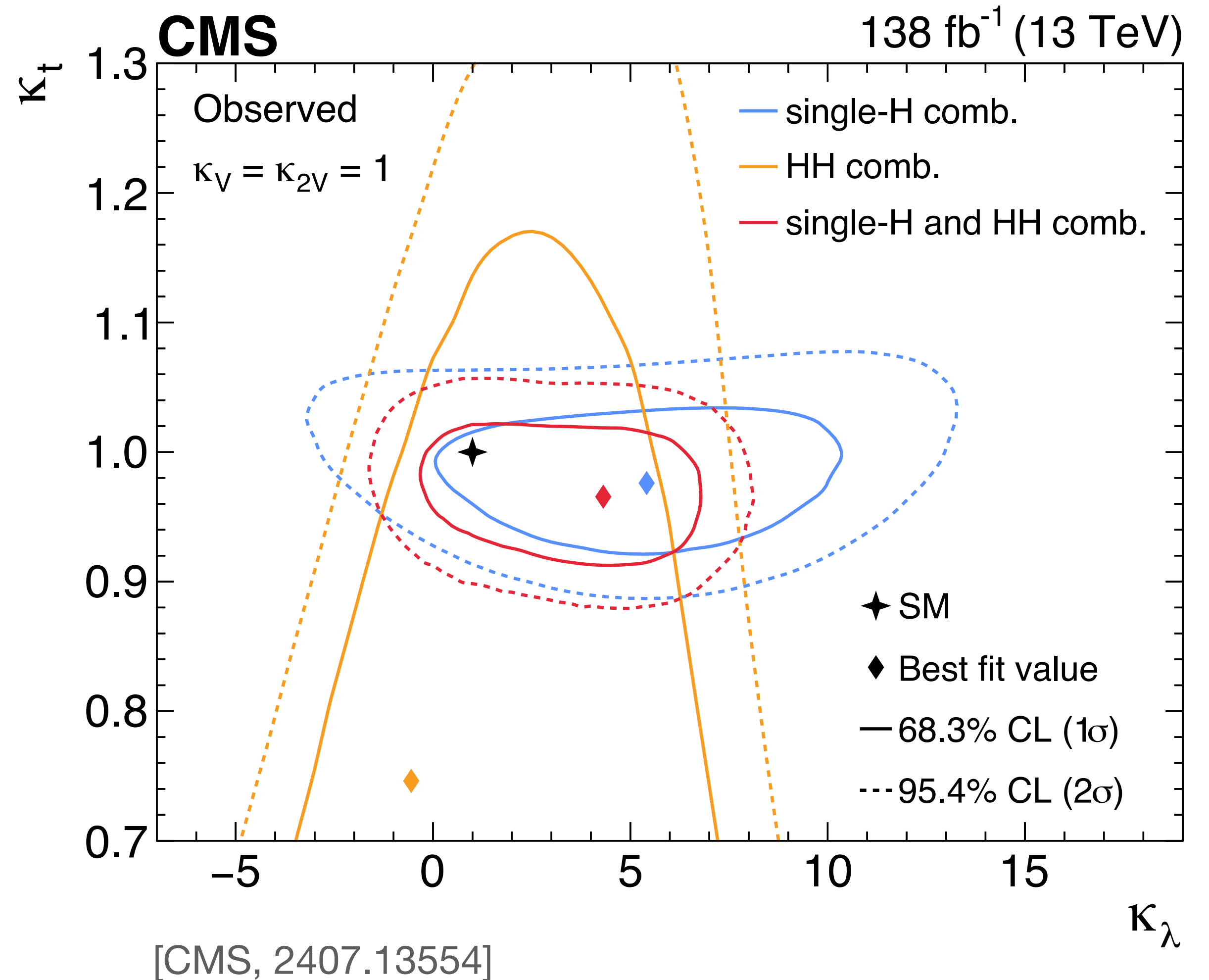
[ATLAS, 2211.01216]

Combination assumption	Obs. 95% CL	Exp. 95% CL	Obs. value $^{+1\sigma}_{-1\sigma}$
<i>HH</i> combination	$-0.6 < \kappa_\lambda < 6.6$	$-2.1 < \kappa_\lambda < 7.8$	$\kappa_\lambda = 3.1^{+1.9}_{-2.0}$
Single- <i>H</i> combination	$-4.0 < \kappa_\lambda < 10.3$	$-5.2 < \kappa_\lambda < 11.5$	$\kappa_\lambda = 2.5^{+4.6}_{-3.9}$
<i>HH+H</i> combination	$-0.4 < \kappa_\lambda < 6.3$	$-1.9 < \kappa_\lambda < 7.6$	$\kappa_\lambda = 3.0^{+1.8}_{-1.9}$

If only κ_λ is considered, constraints from hh production are notable better than bounds that arise from combination of single-h measurements

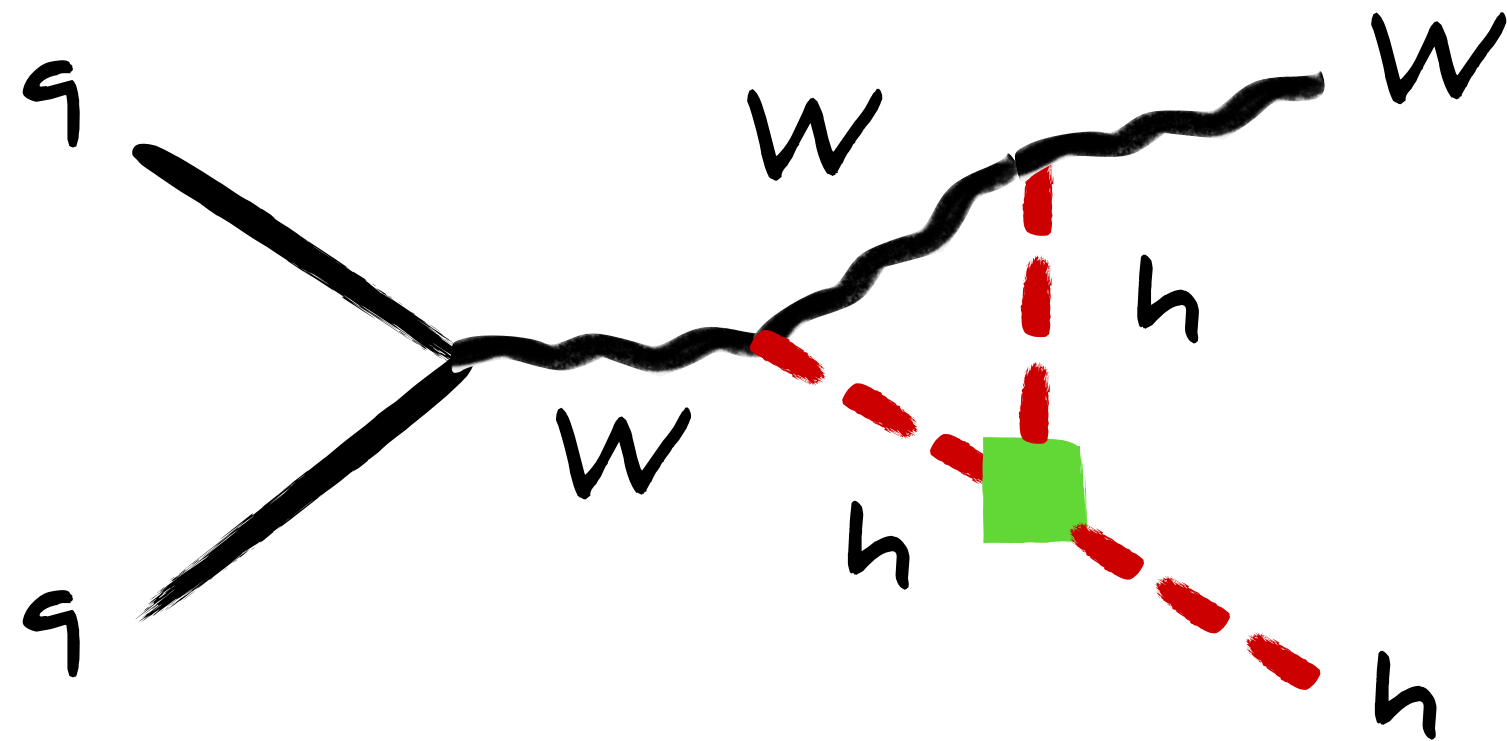
Combined constraints on κ_λ @ LHC

But, large degeneracy of hh production cross section to κ_λ & κ_t , limits κ_λ sensitivity of hh process in 2D plane. Instead, single-h combination provides stringent bounds on κ_t , which is utilized in hh+h combination

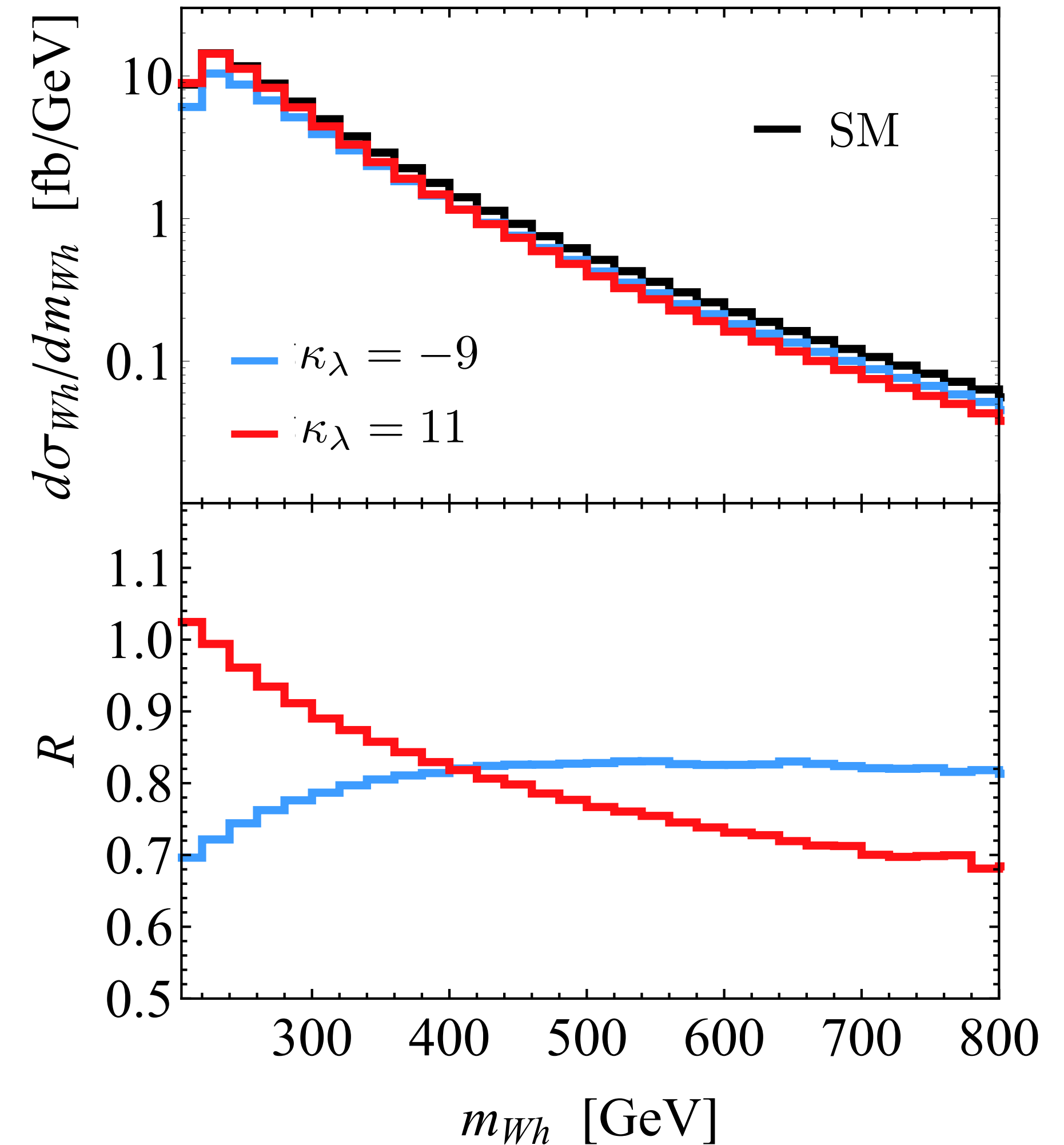


Differential single-h measurements & κ_λ

[Bizon et al., 1610.05771]

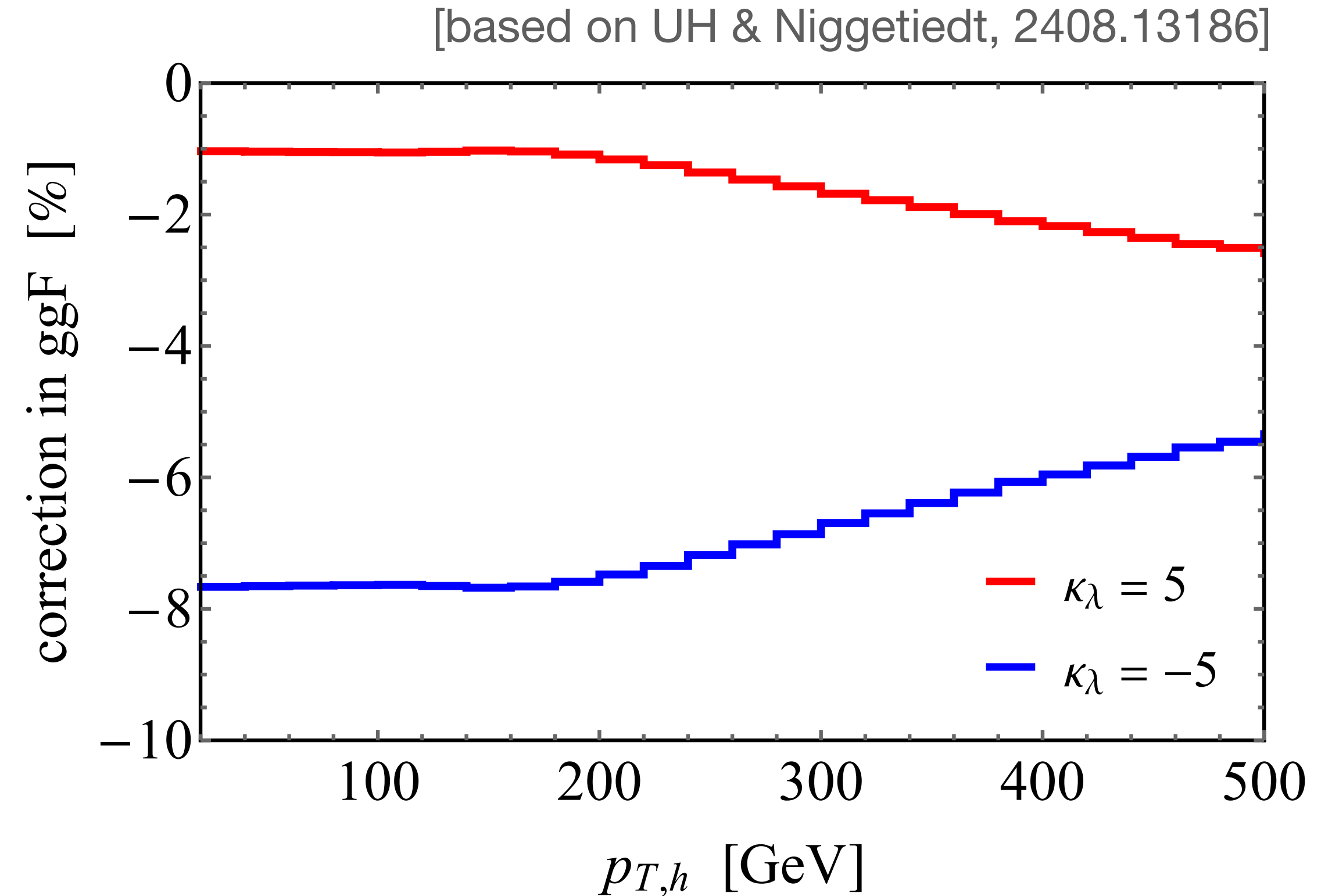
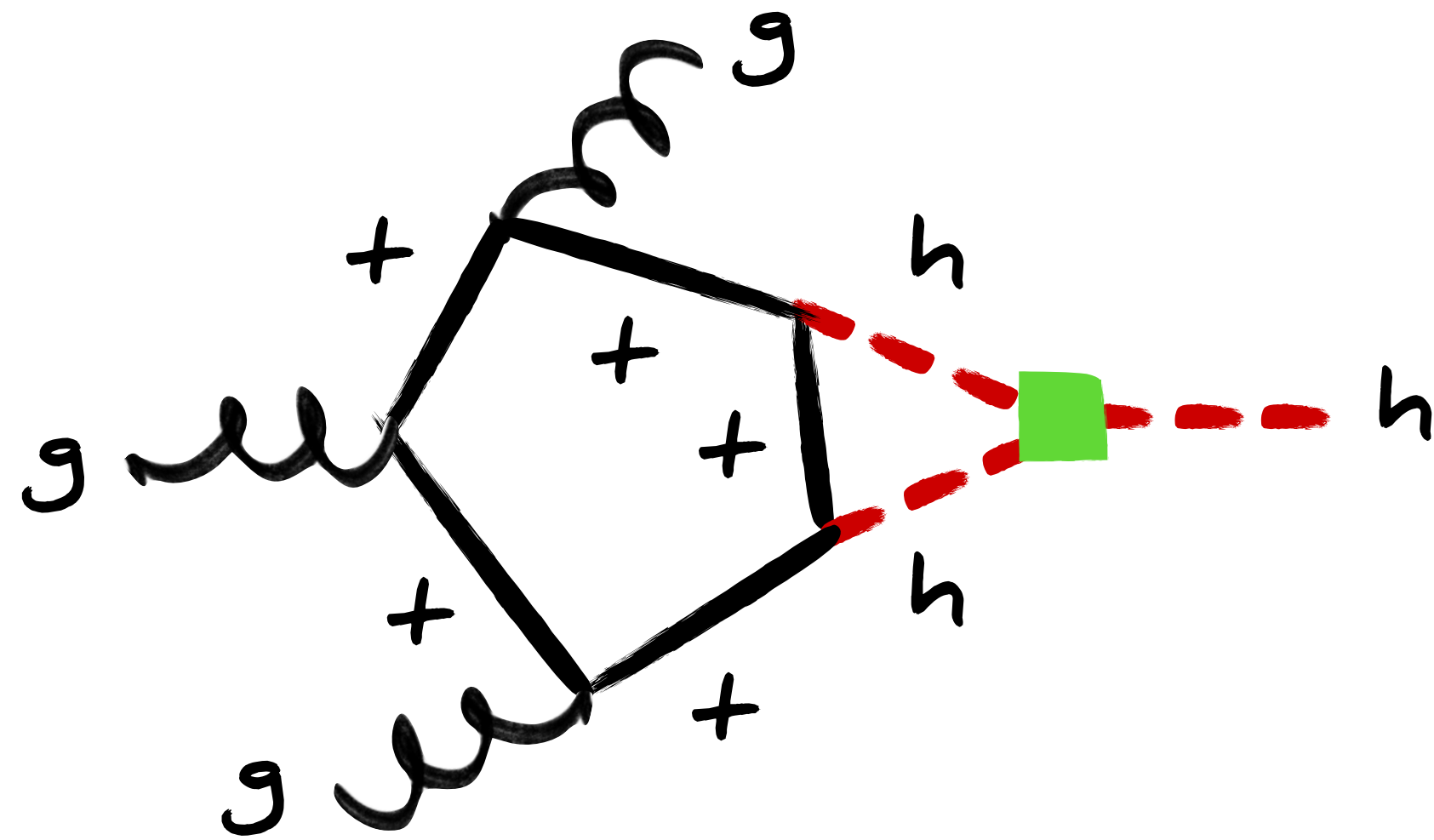


Non-trivial modifications of kinematic distributions due to heavy particles in loops of single-h probes



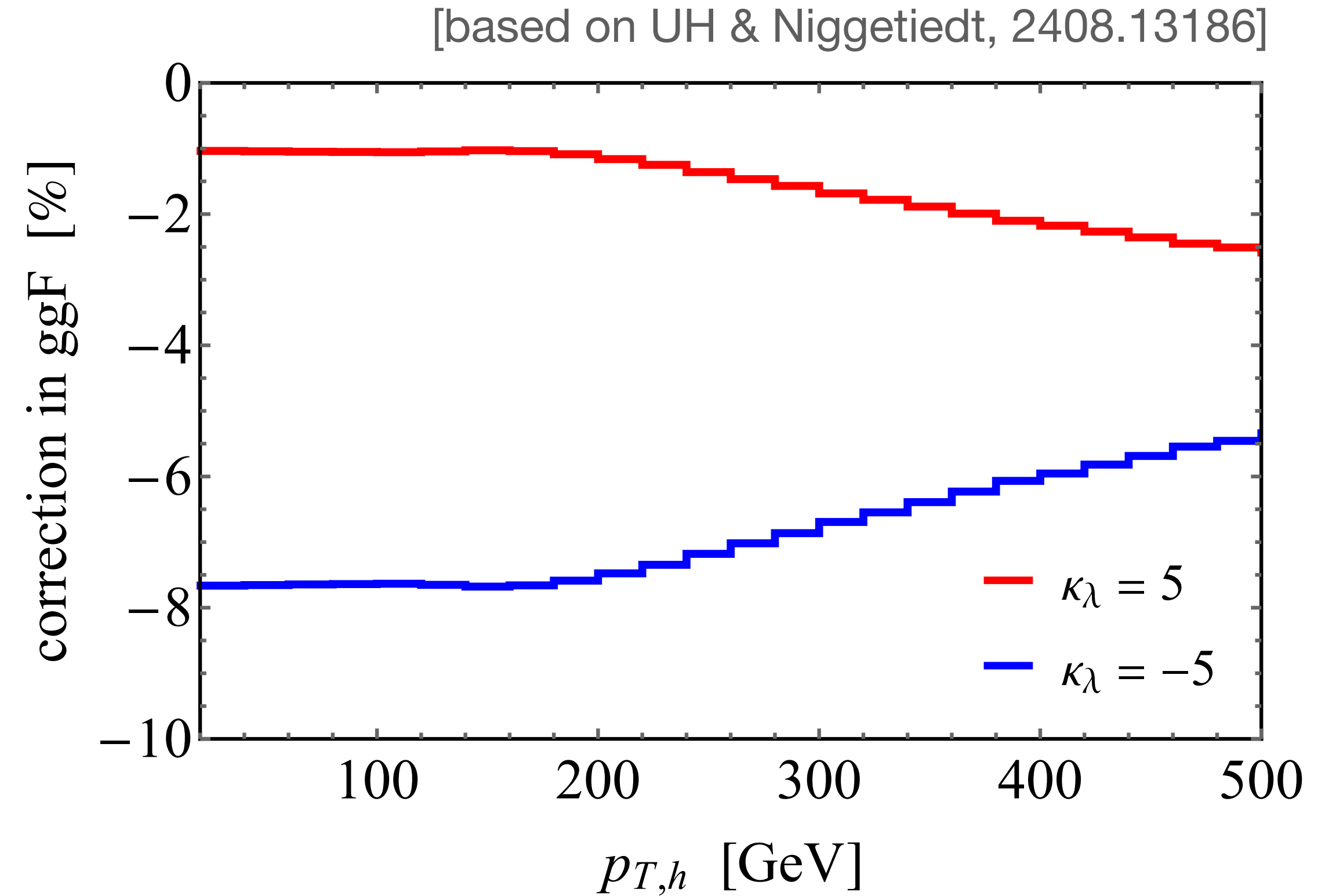
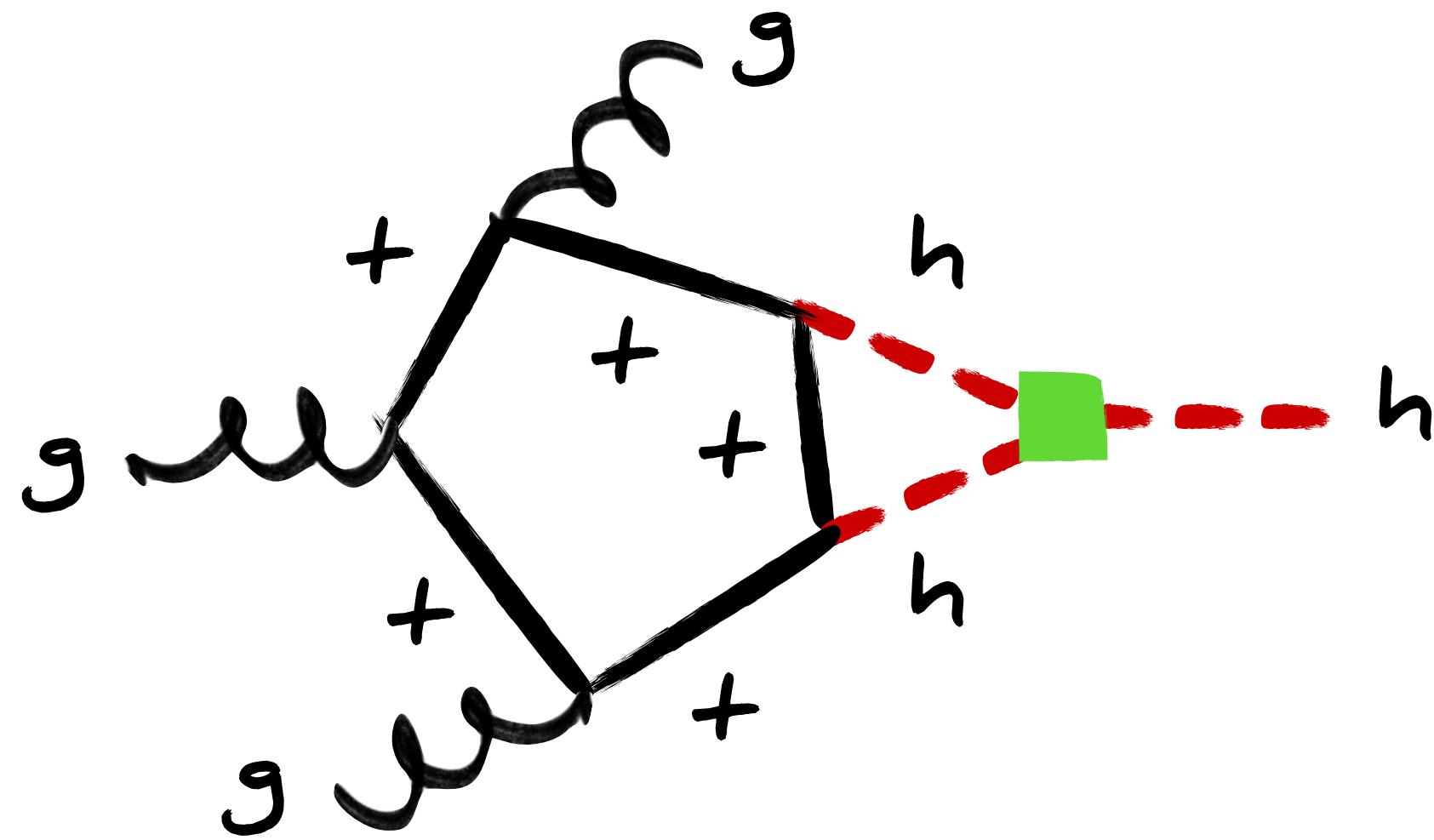
[see also Maltoni, Pagani, Shivaji & Zhao, 1709.08649]

Differential single-h measurements & κ_λ



Exact κ_λ dependence computed for differential single-Higgs predictions in all cases relevant for LHC, i.e. ggF, Vh, VBF, $t\bar{t}h$ & th production

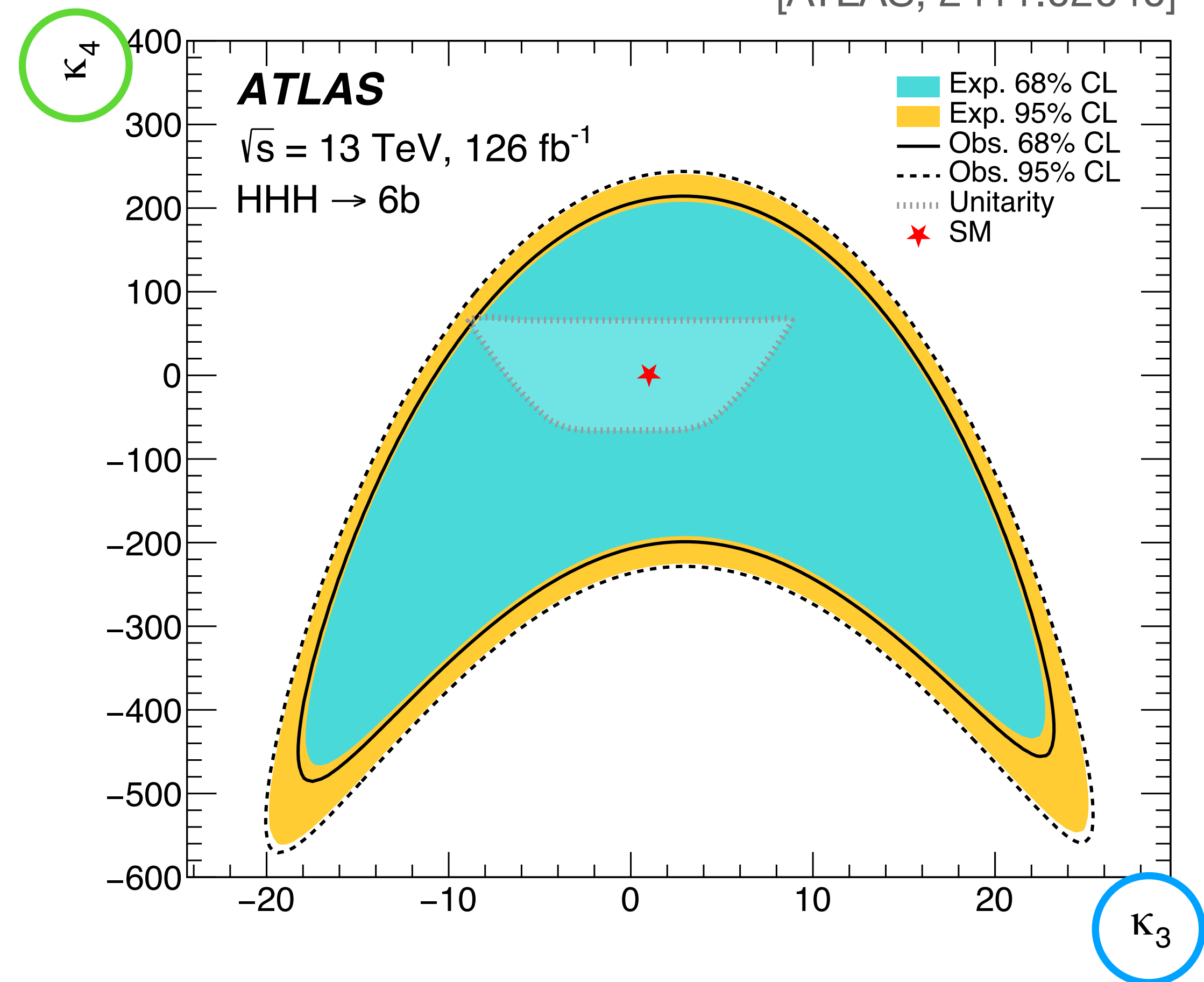
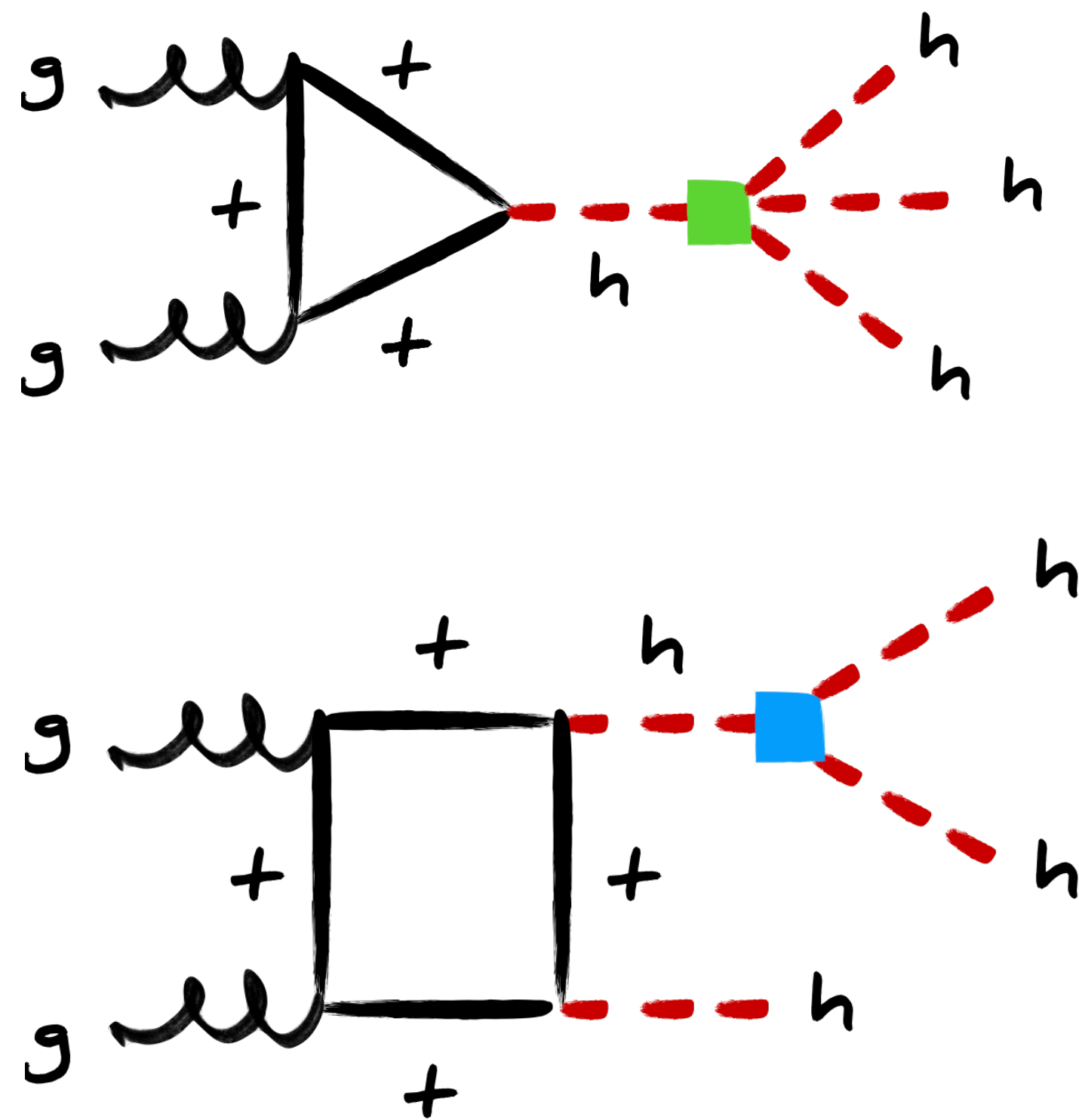
Differential single-h measurements & κ_λ



Dedicated studies needed to quantify precise impact of differential single-h measurements in global $hh+h$ analyses to constrain κ_λ

Nice, first LHC limit on 3h production!

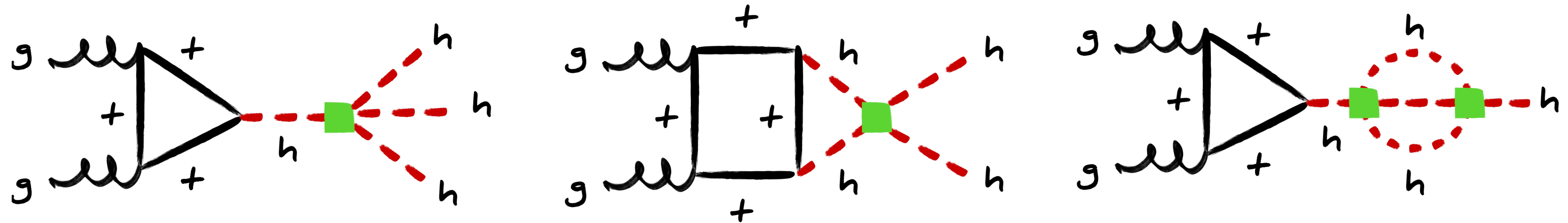
[ATLAS, 2411.02040]



ATLAS puts first constraint on 3h production & interprets its result in κ_3 - κ_4 plane

[see also talk by Balunas & Chen]

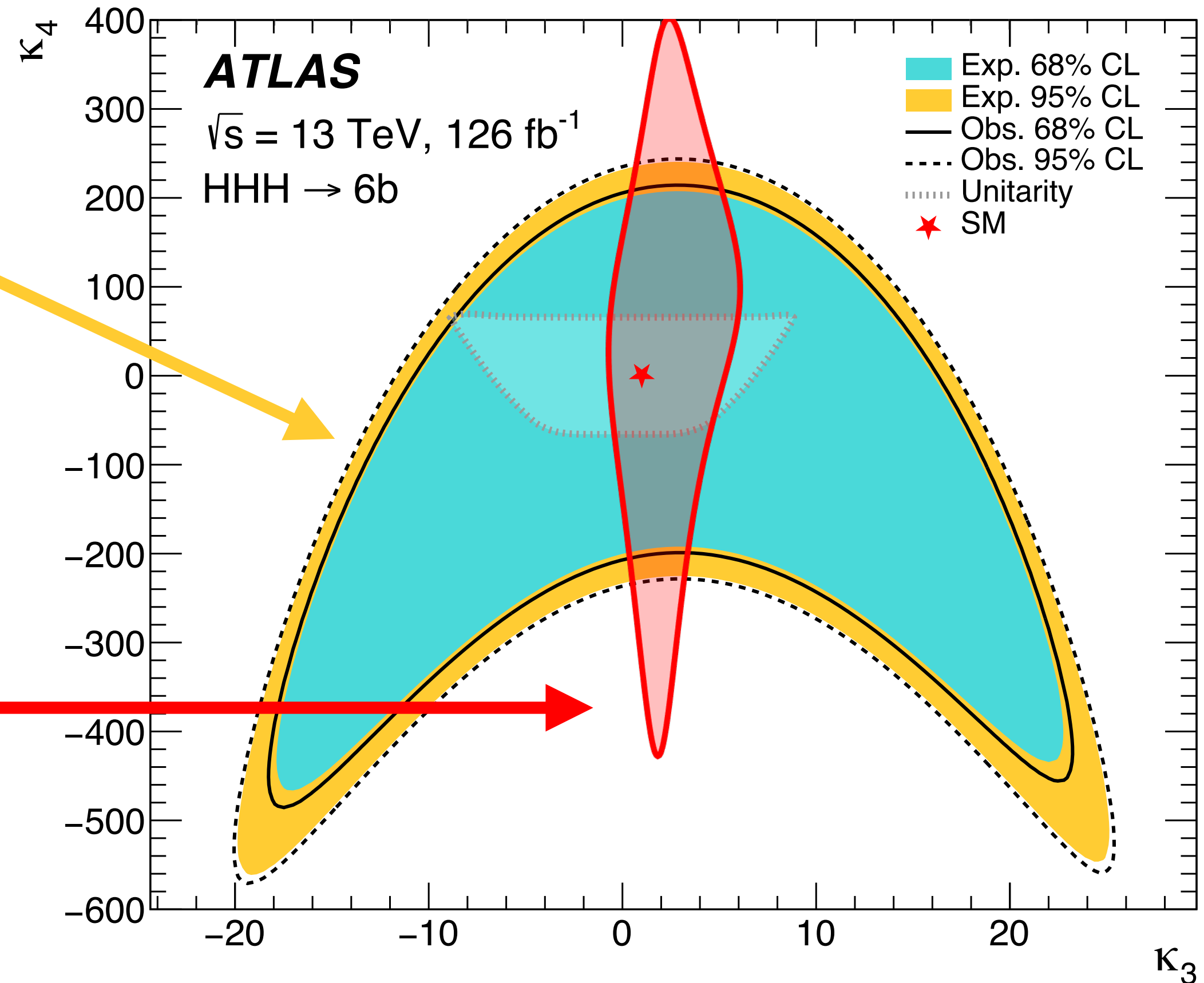
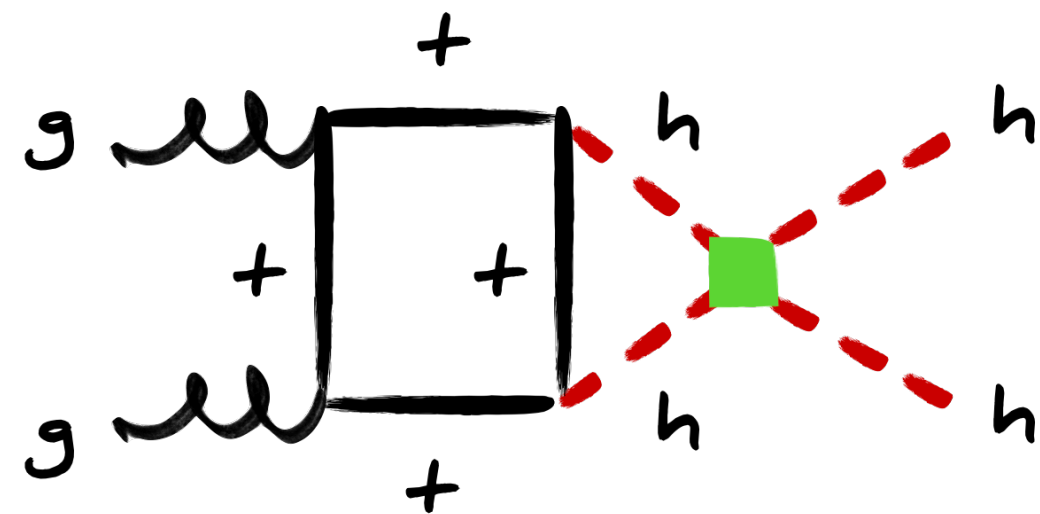
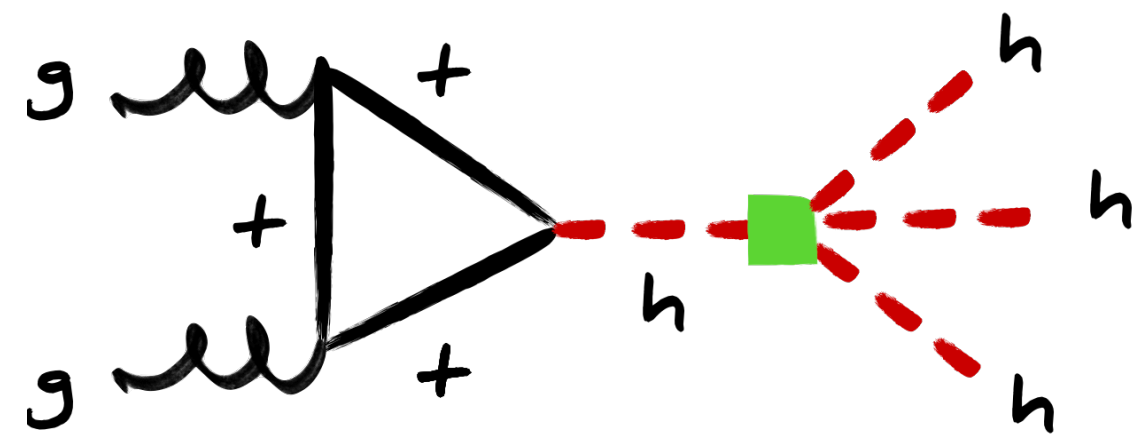
(In)direct probes of κ_4



Direct probe of quartic Higgs self-coupling (κ_4) provided @ 1-loop by 3h production, while indirect sensitivity in hh & h production through 2-loop & 3-loop corrections

Higgs self-couplings after LHC Run 2

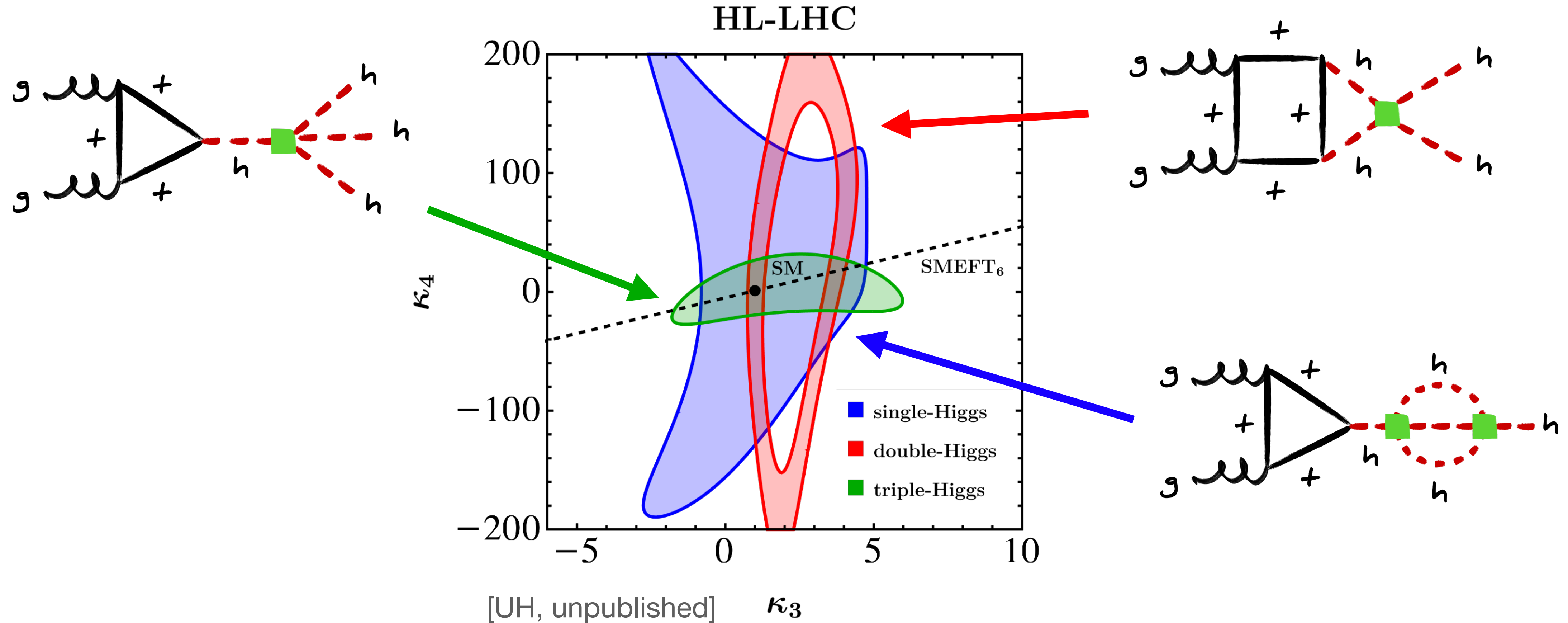
[ATLAS, 2411.02040]



Bounds on Higgs self-couplings from hh & 3h production orthogonal in κ_3 - κ_4 plane

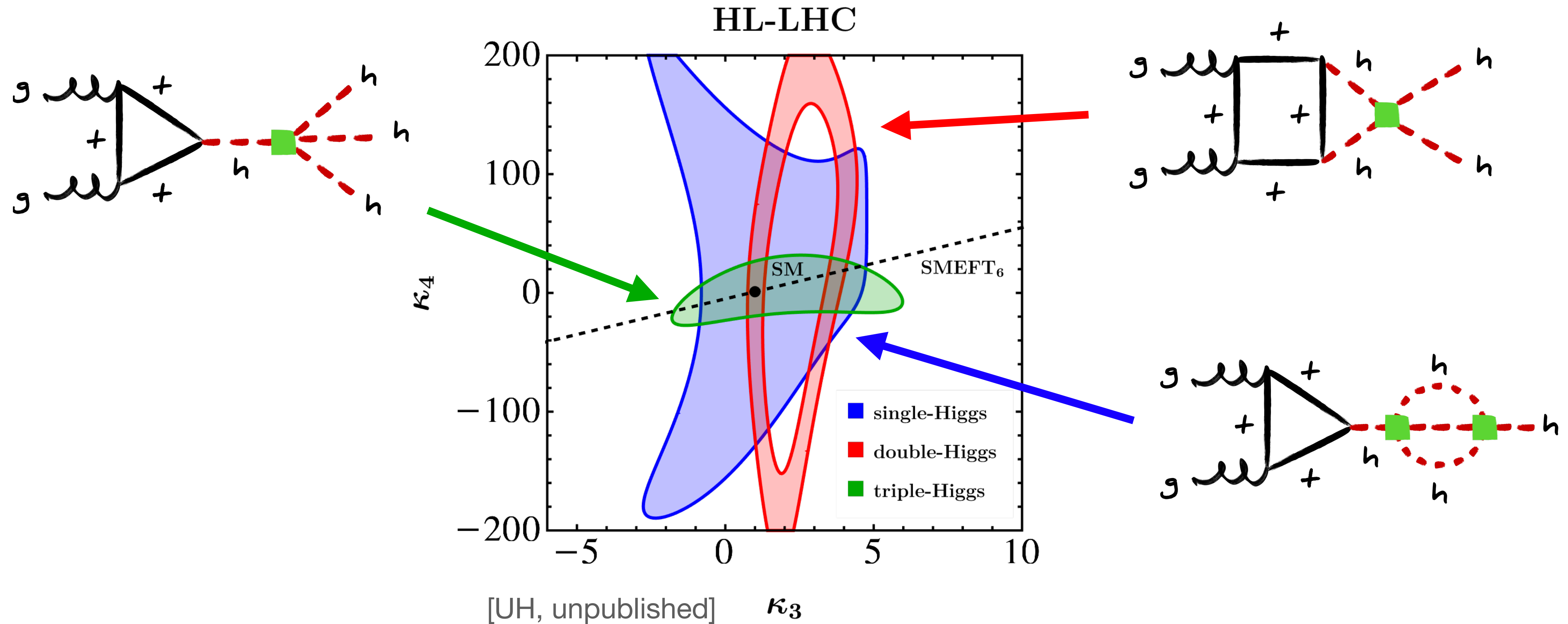
[pp→hh bound obtained using results from Bizon et al., 2402.03463]

Higgs self-couplings in HL-LHC era



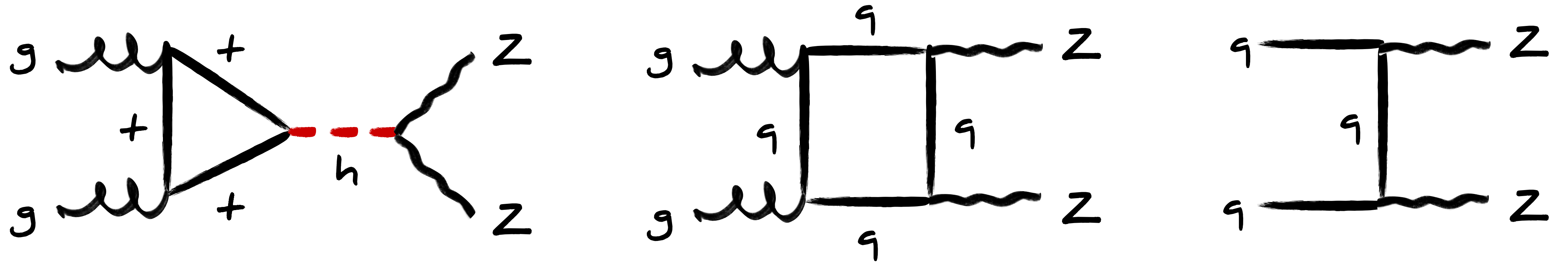
Hypothetical HL-LHC bound of $O(10)$ on $3h$ signal strength will set best bound on κ_4

Higgs self-couplings in HL-LHC era



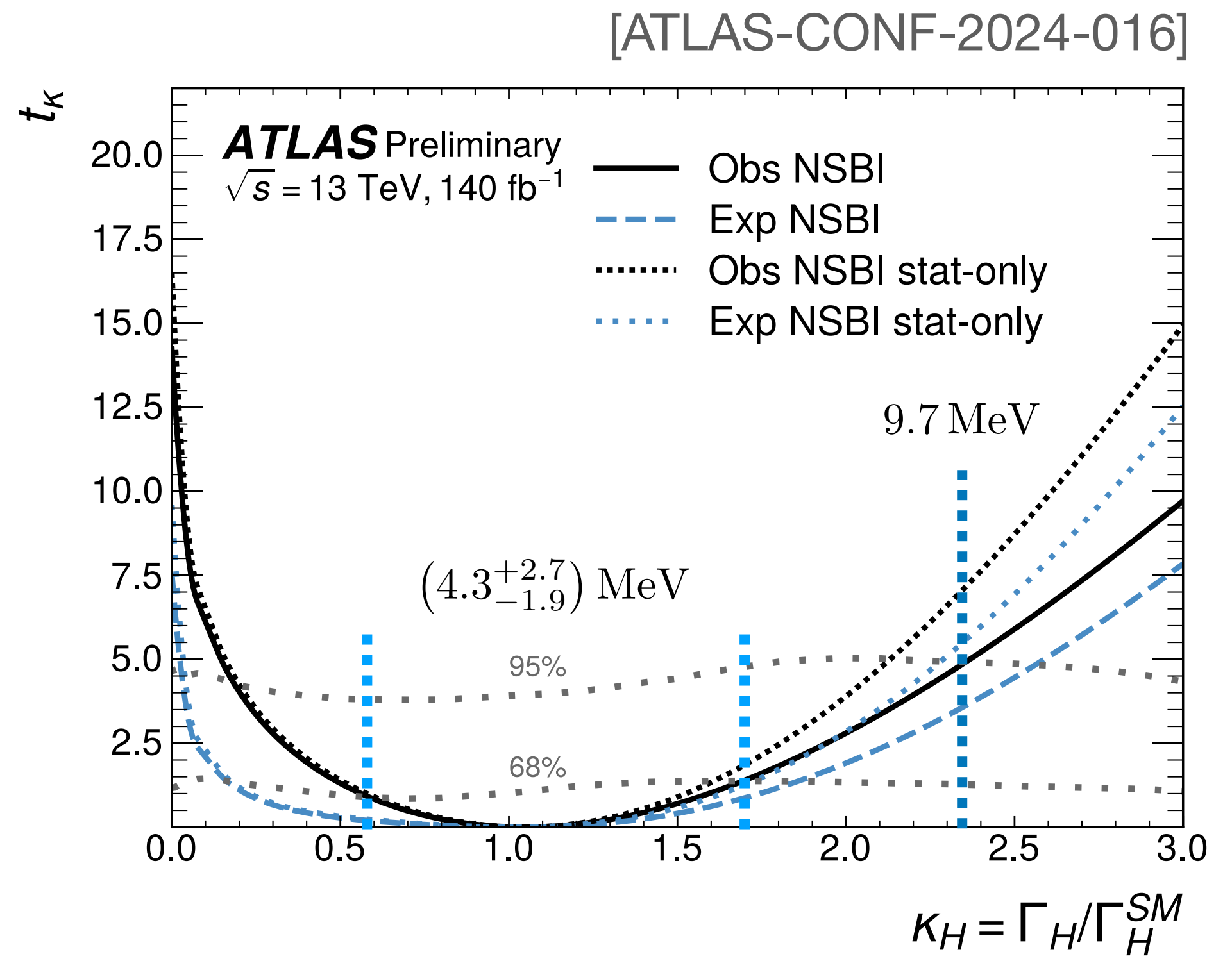
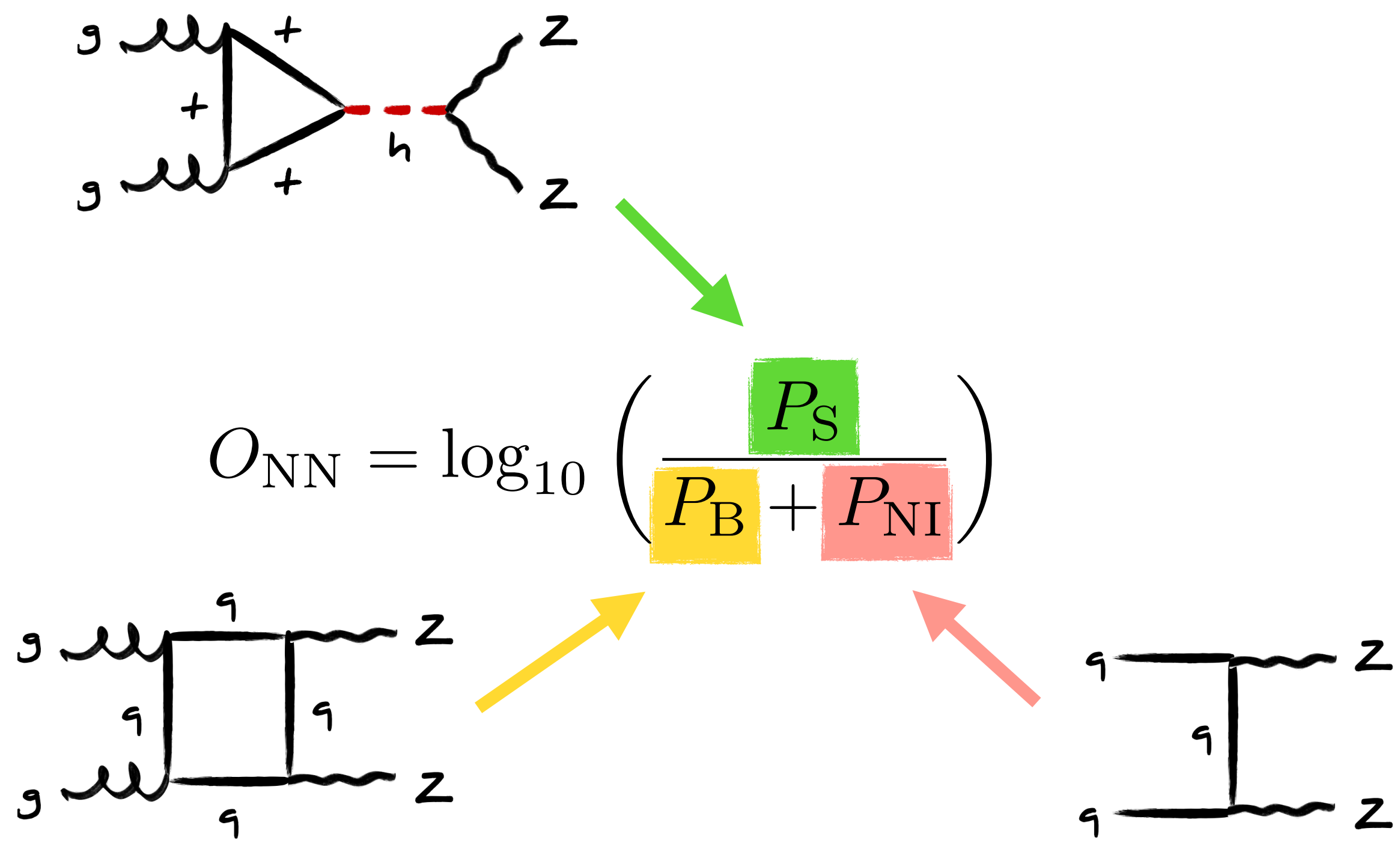
Flat direction in κ_3 of 3h constraint, partly resolved by indirect hh & h probes

Higgs off-shell measurements



For $m_{4l} > m_h$, $pp \rightarrow ZZ \rightarrow 4l$ distributions have an enhanced sensitivity to $gg \rightarrow h \rightarrow ZZ$ process & its interference with $gg \rightarrow ZZ$ channel. Assuming that on-shell Higgs signal strengths are SM-like, possible to set bounds on total Higgs width @ LHC

Higgs width measurements @ LHC

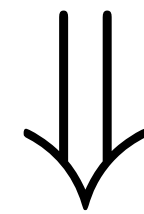


To enhance sensitivity to Higgs contribution with respect to m_{4l} , latest LHC searches for $ZZ \rightarrow 4l$ use machine learning approaches to matrix-element method

[see also CMS, 2409.13663; talks by Gargiulo, Leight, Sandesara & Winterbottom]

Trilinear Higgs self-coupling from $pp \rightarrow 4l$

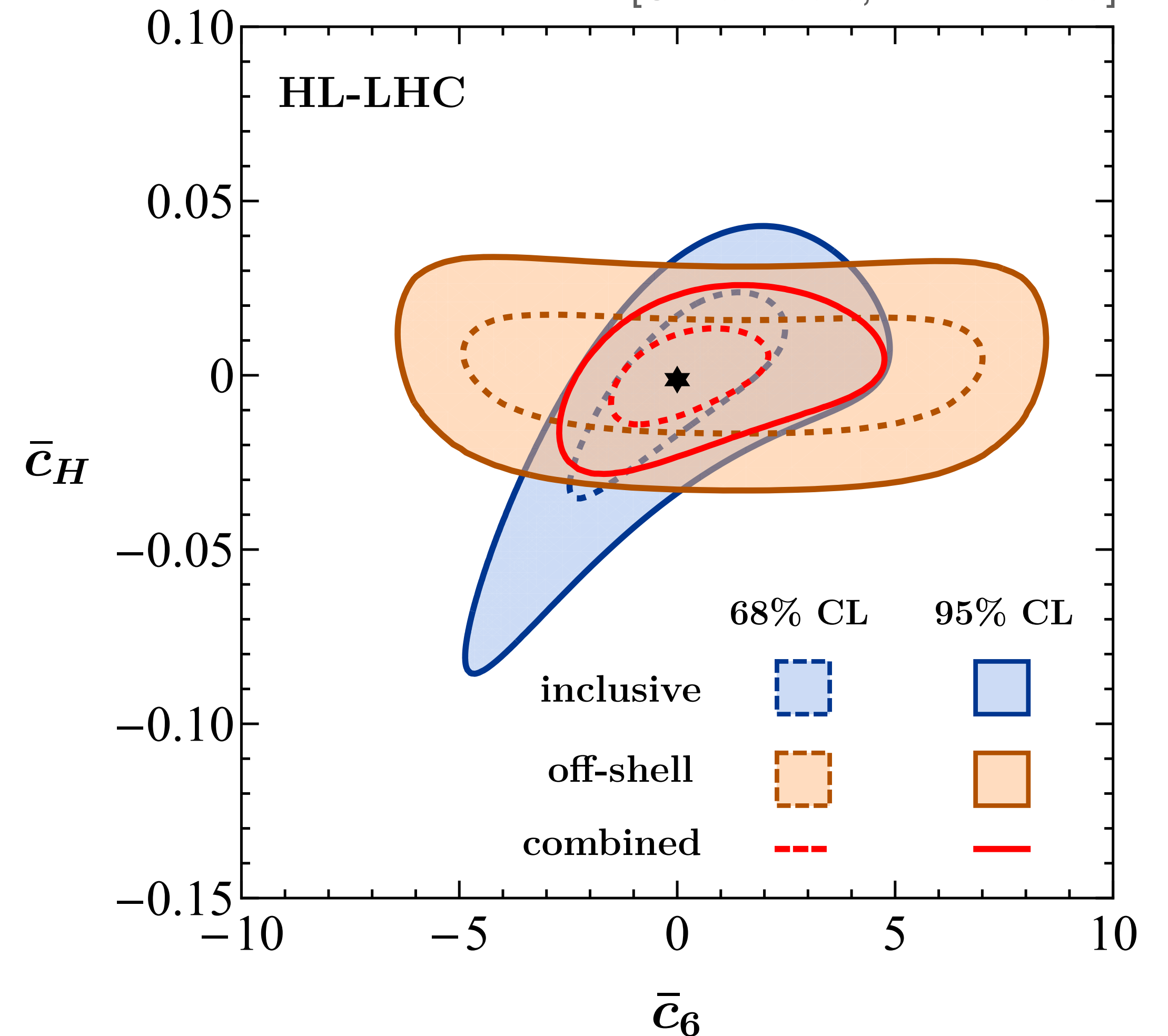
$$\mathcal{L}_{\text{SMEFT}} \supset -\frac{\lambda \bar{c}_6}{v^2} |H|^6 + \frac{\bar{c}_H}{2v^2} (\partial_\mu |H|^2)^2$$



$$\mathcal{L} \supset -\lambda \kappa_\lambda v h^3, \quad \kappa_\lambda = 1 + \bar{c}_6 - \frac{3}{2} \bar{c}_H$$

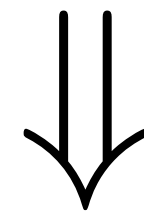
In SM effective field theory (SMEFT), κ_λ receives contributions from more than a single Wilson coefficient

[UH & Koole, 2111.12589]



Trilinear Higgs self-coupling from $pp \rightarrow 4l$

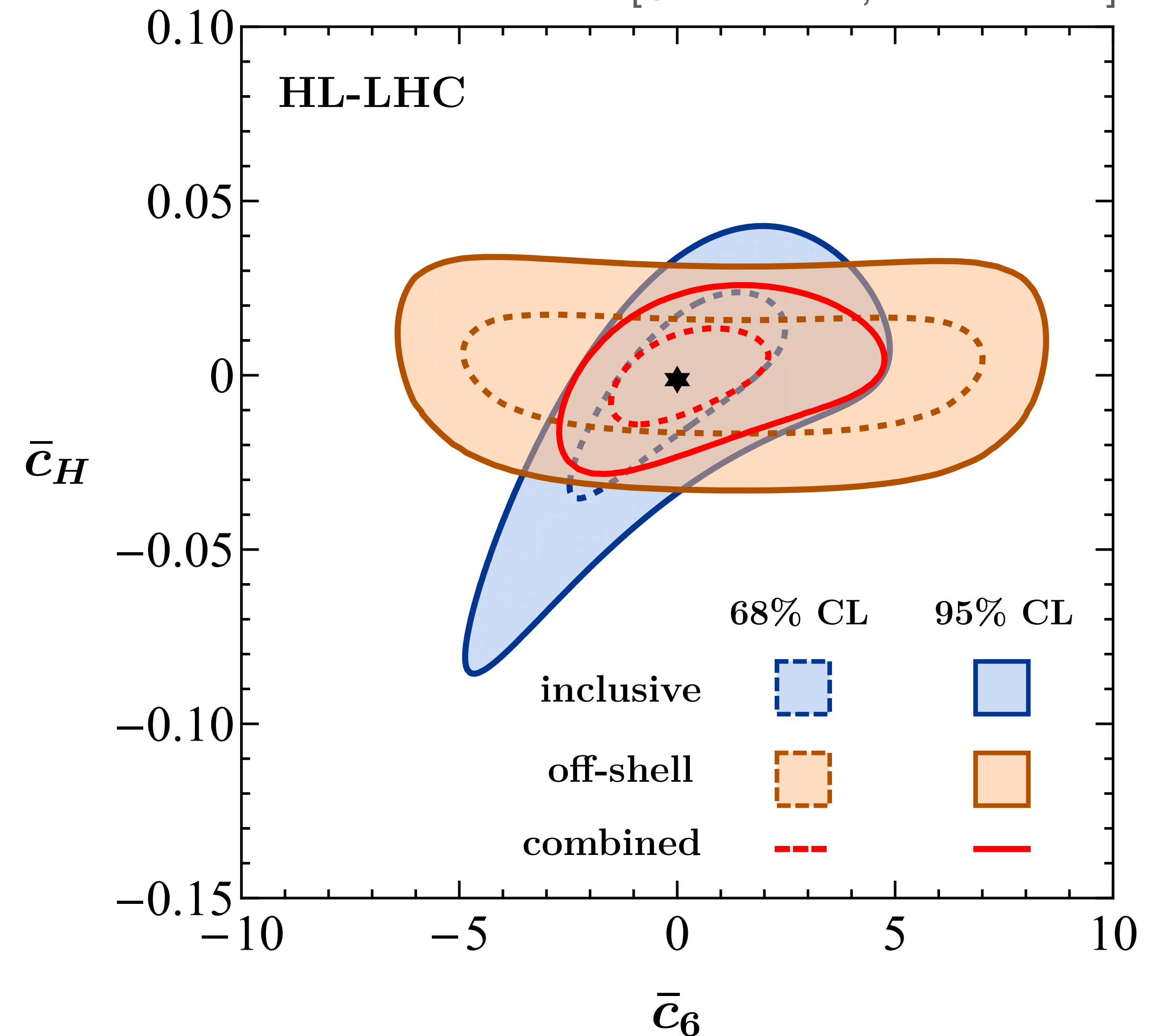
$$\mathcal{L}_{\text{SMEFT}} \supset -\frac{\lambda \bar{c}_6}{v^2} |H|^6 + \frac{\bar{c}_H}{2v^2} (\partial_\mu |H|^2)^2$$



$$\mathcal{L} \supset -\lambda \kappa_\lambda v h^3, \quad \kappa_\lambda = 1 + \bar{c}_6 - \frac{3}{2} \bar{c}_H$$

On-shell & off-shell measurements
depend differently on O_6 & O_H operators.
Combination improves 2D constraints

[UH & Koole, 2111.12589]



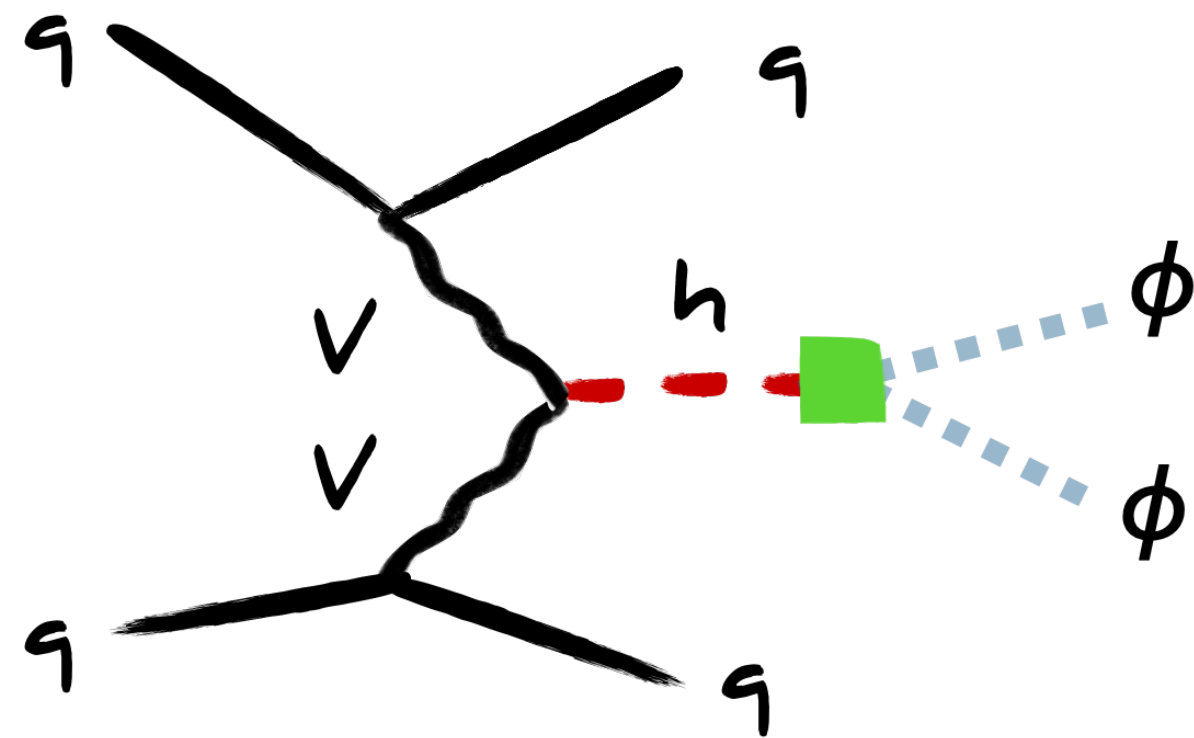
Higgs portal interactions

$$\mathcal{L} \supset c_\phi |H|^2 |\phi|^2 + \frac{c_{\partial\phi}}{f^2} (\partial_\mu |H|^2) (\partial^\mu |\phi|^2) + \frac{c_\psi}{f} |H|^2 \bar{\psi}\psi + c_V |H|^2 V_\mu V^\mu$$

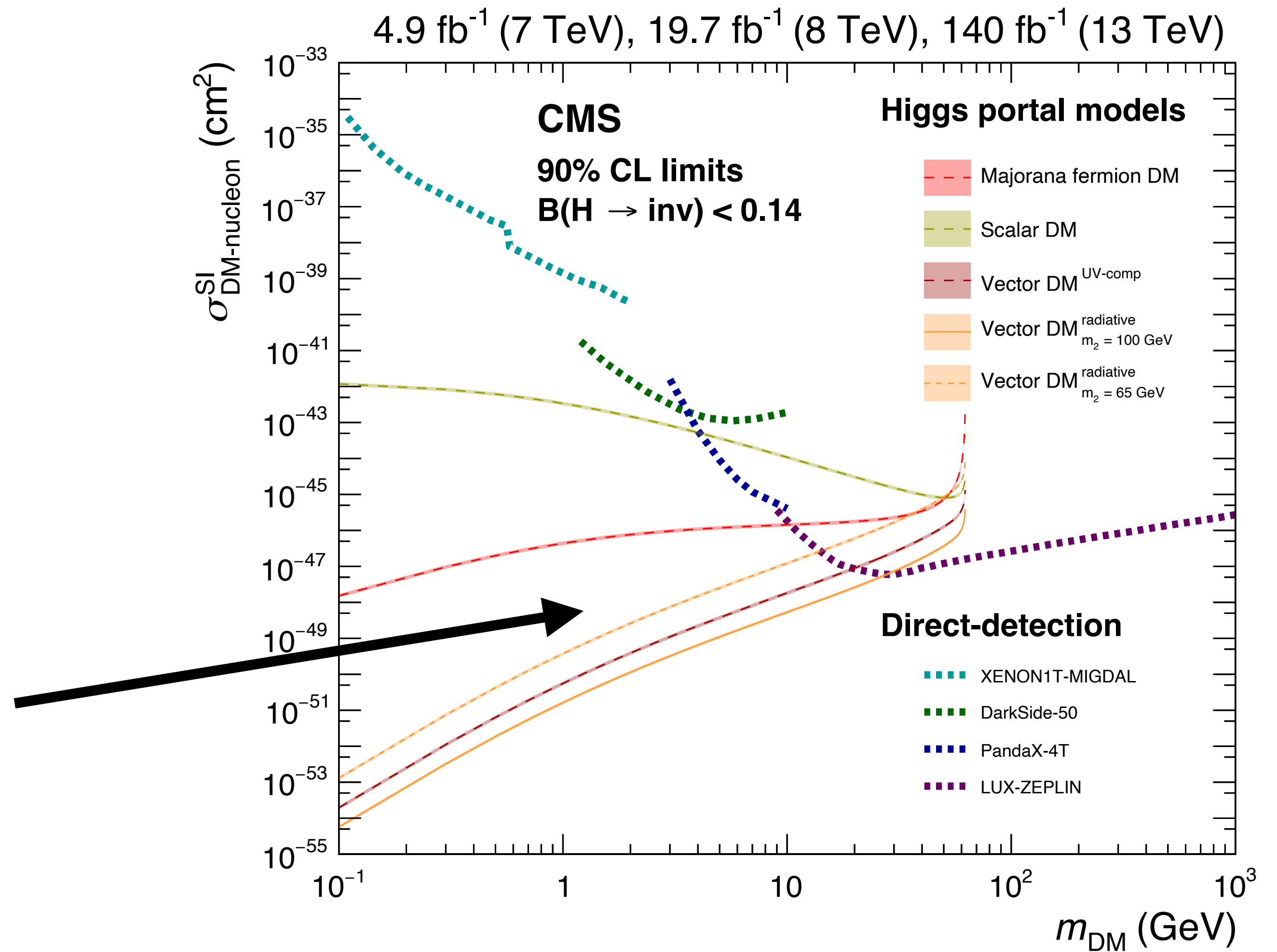
Diagram illustrating Higgs portal interactions. The Lagrangian \mathcal{L} is shown with four terms. The first term, $c_\phi |H|^2 |\phi|^2$, is labeled "scalar" with two green arrows pointing to the $|\phi|^2$ and $|\phi|^2$ terms. The second term, $\frac{c_{\partial\phi}}{f^2} (\partial_\mu |H|^2) (\partial^\mu |\phi|^2)$, is also labeled "scalar" with two green arrows pointing to the $|\phi|^2$ and $|\phi|^2$ terms. The third term, $\frac{c_\psi}{f} |H|^2 \bar{\psi}\psi$, is labeled "fermion" with a yellow arrow pointing to the $\bar{\psi}\psi$ term. The fourth term, $c_V |H|^2 V_\mu V^\mu$, is labeled "vector" with a red arrow pointing to the $V_\mu V^\mu$ term.

$|H|^2$ provides a simple portal to dark or hidden sectors. At dimension four one has couplings of $|H|^2$ to spin-0 & spin-1 fields, while interactions with spin-1/2 fields are of dimension five. Dimension-six derivative spin-0 coupling also interesting, since dark matter (DM) direct detection (DD) cross section momentum suppressed

Higgs portal searches @ LHC



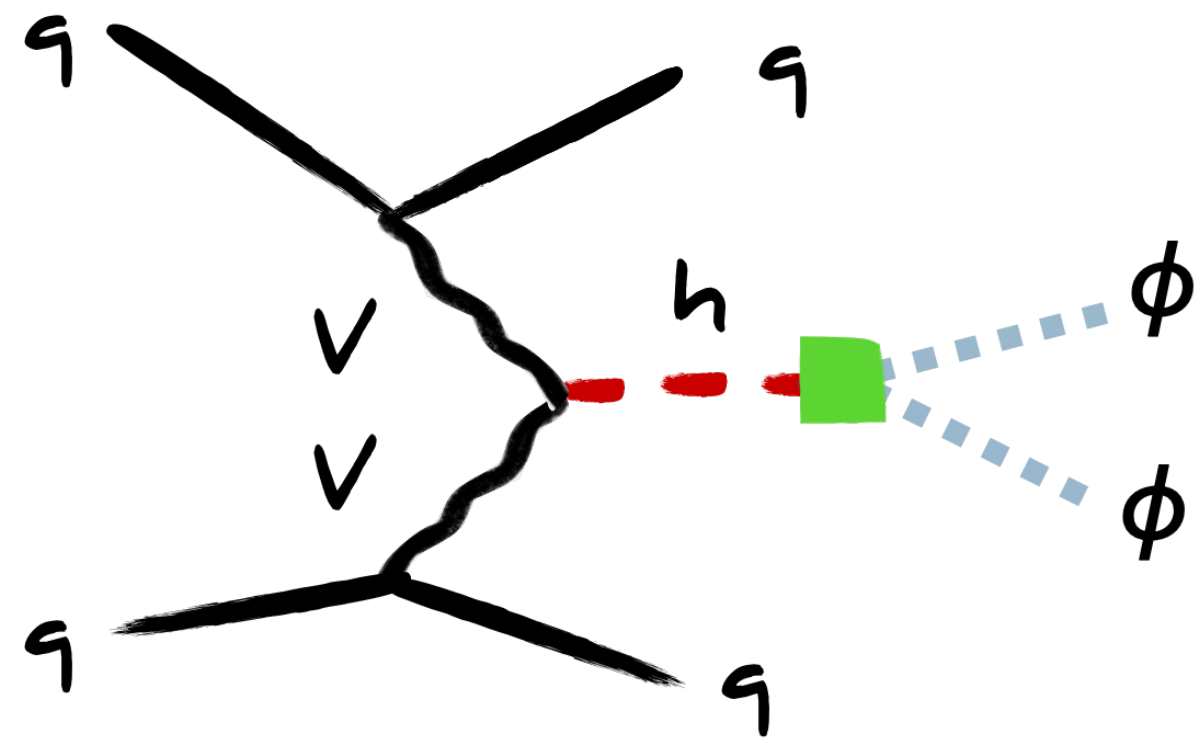
If DM states are kinematically accessible in Higgs decays, LHC searches for $h \rightarrow \text{invisible}$ superior to present DD limits



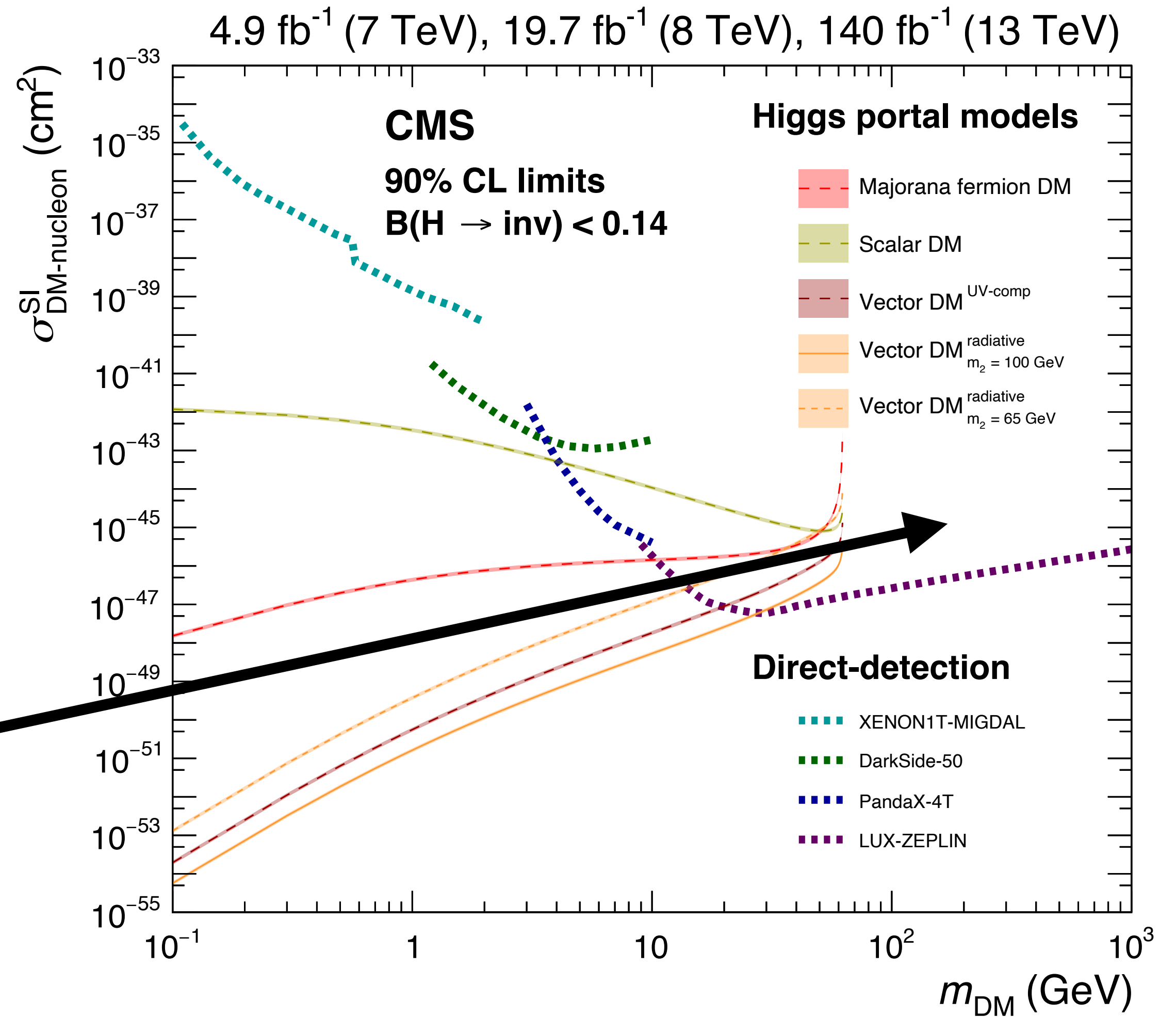
[CMS, 2303.01214]

[see also ATLAS, 2202.07953]

Higgs portal searches @ LHC



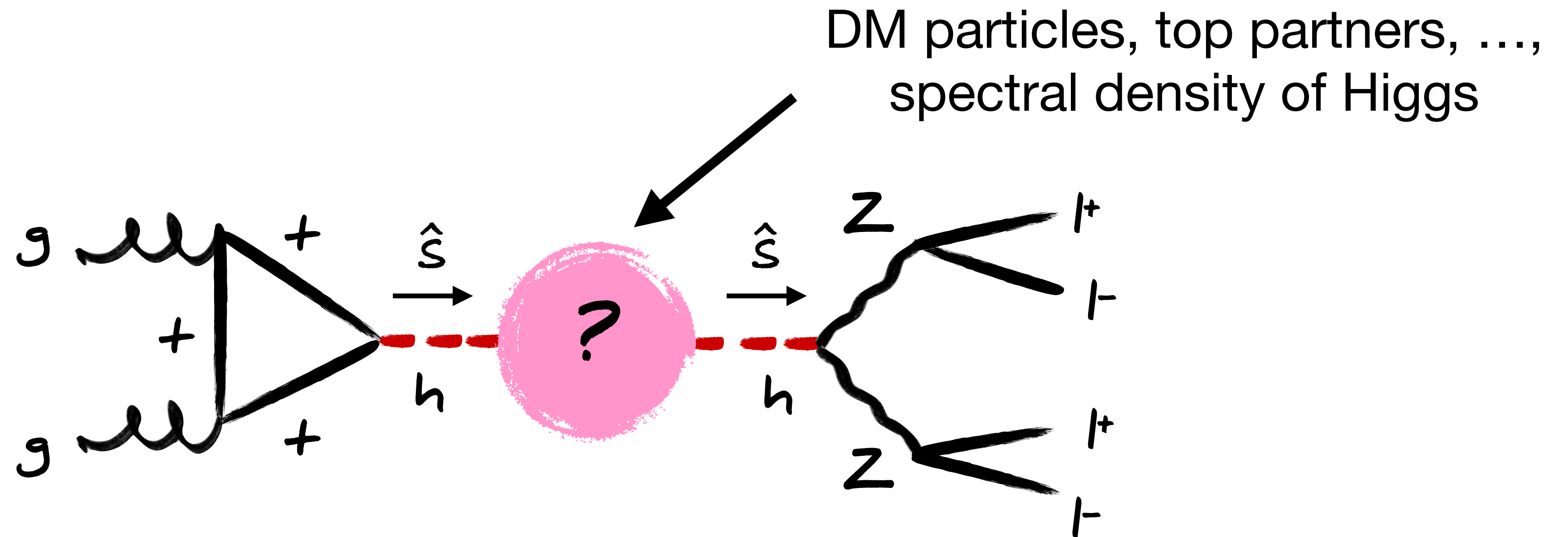
Can LHC say something about region where DM is inaccessible in $h \rightarrow$ invisible?



[CMS, 2303.01214]

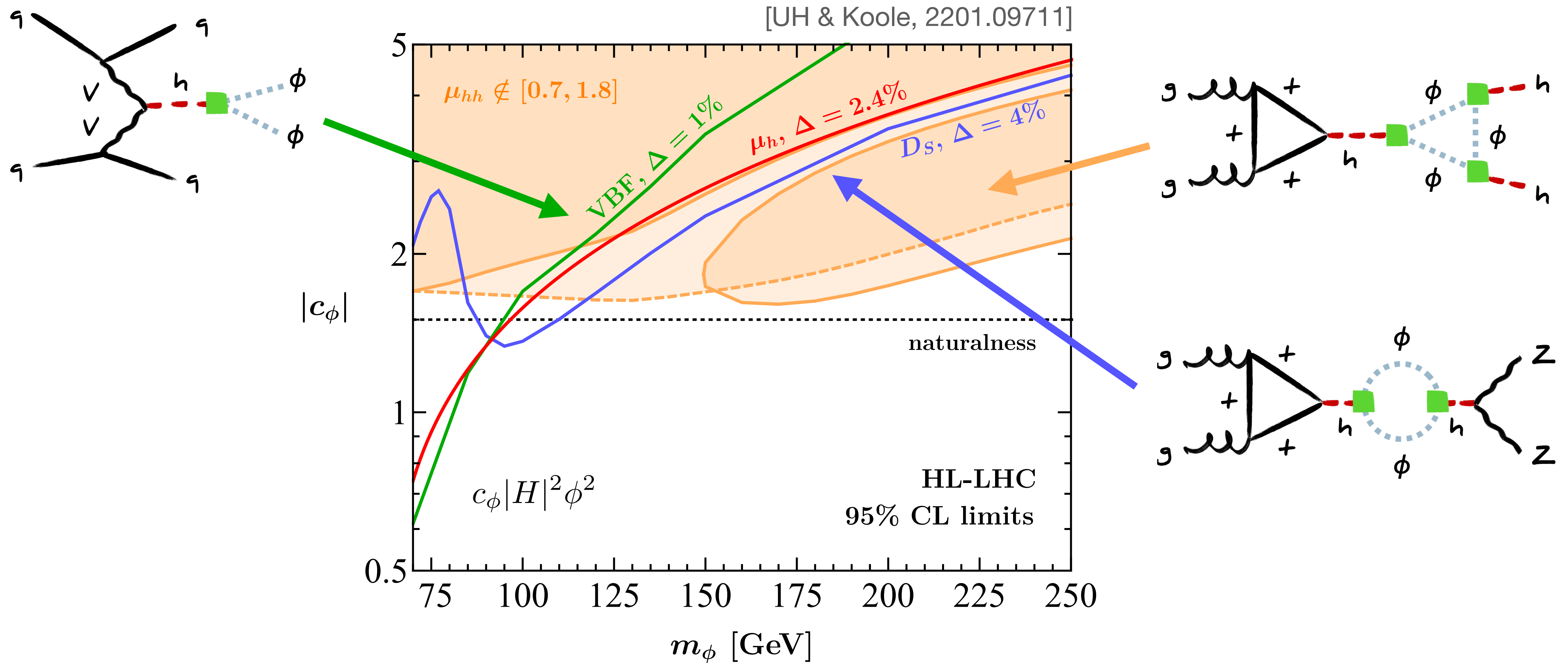
[see also ATLAS, 2202.07953]

Searches for Higgsphilics in $pp \rightarrow 4l$



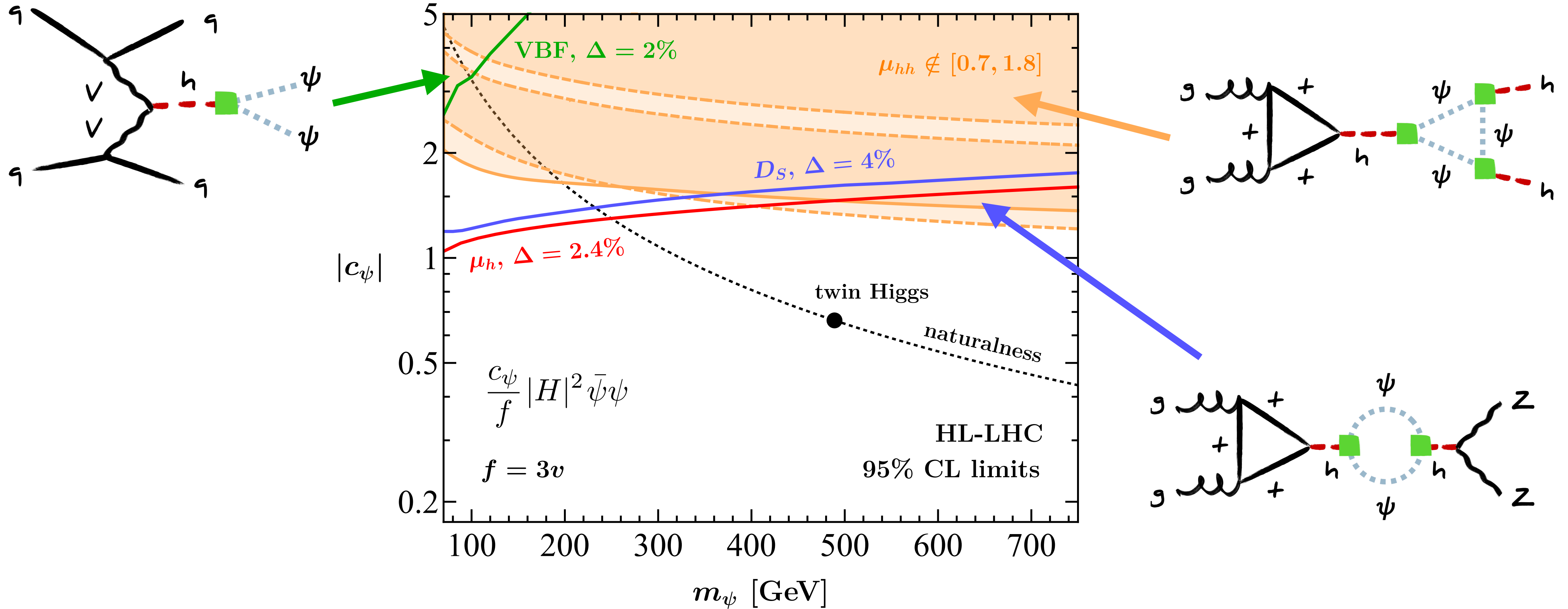
Off-shell Higgs measurements in $pp \rightarrow 4l$ allow to scan \hat{s} -dependence of Higgs propagator, which is sensitive to virtual exchange of light Higgsphilic states

HL-LHC bounds on marginal Higgs portal



[off-shell VBF bound taken from Ruhdorfer, Salvioni & Weiler, 1910.04170]

HL-LHC bounds on fermionic Higgs portal



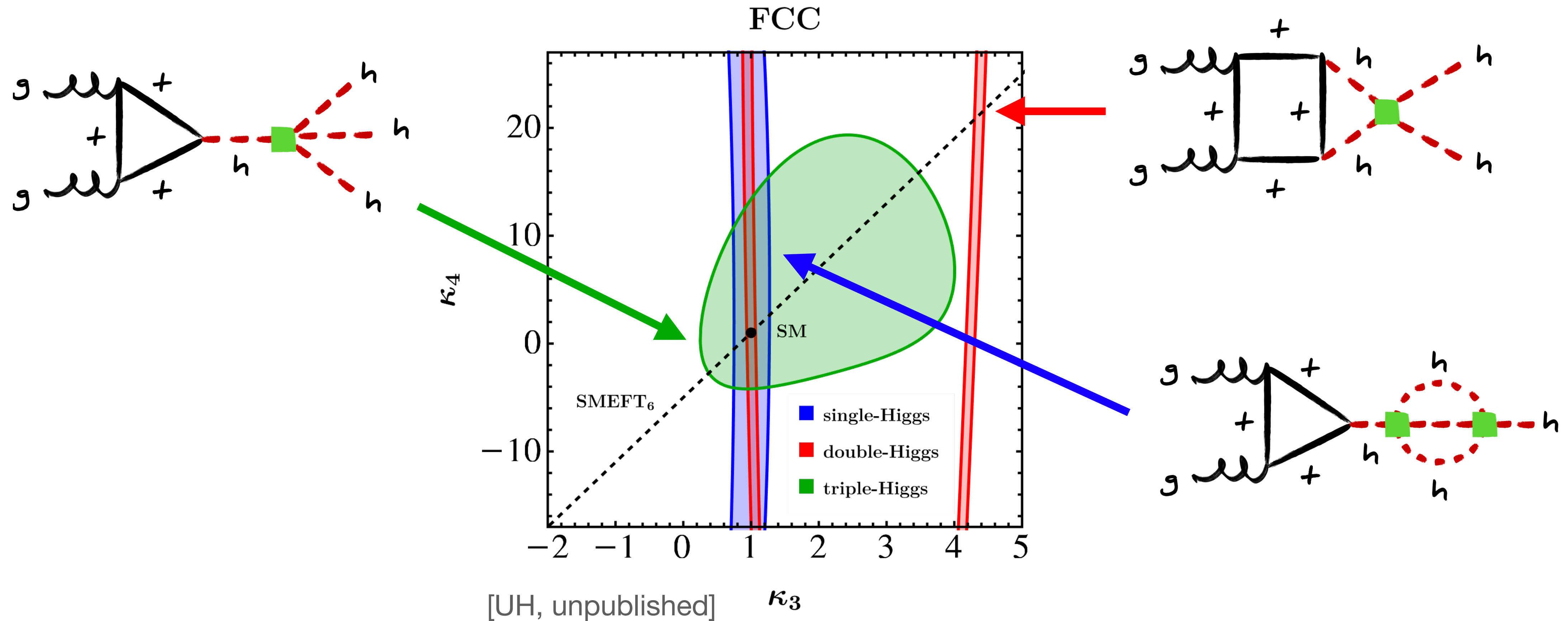
Conclusions

- Distinction between direct & indirect Higgs probes academic, because in practice both strategies test same physics in a given beyond SM model
- Indirect Higgs probes well-established @ LHC & employed by both ATLAS & CMS to extractions of charm Yukawa & trilinear Higgs self-coupling
- Differential & off-shell Higgs measurements @ HL-LHC will allow to tighten indirect constraints on κ 's, Wilson coefficients, light Higgsphilic physics, etc. Measurements have not been fully exploited experimentally @ LHC Run 2

Backup



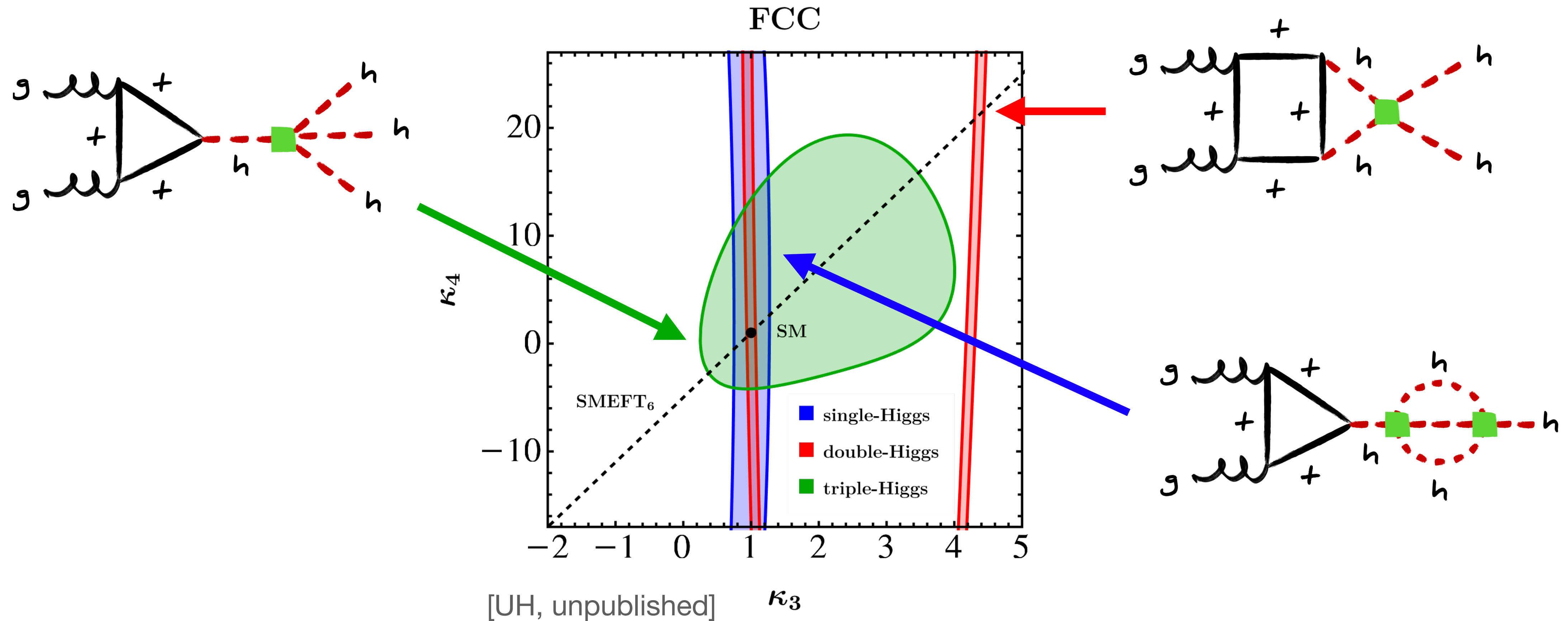
Higgs self-couplings in FCC era



Single-h bounds notable improved due to permille accuracy of Zh @ FCC-ee

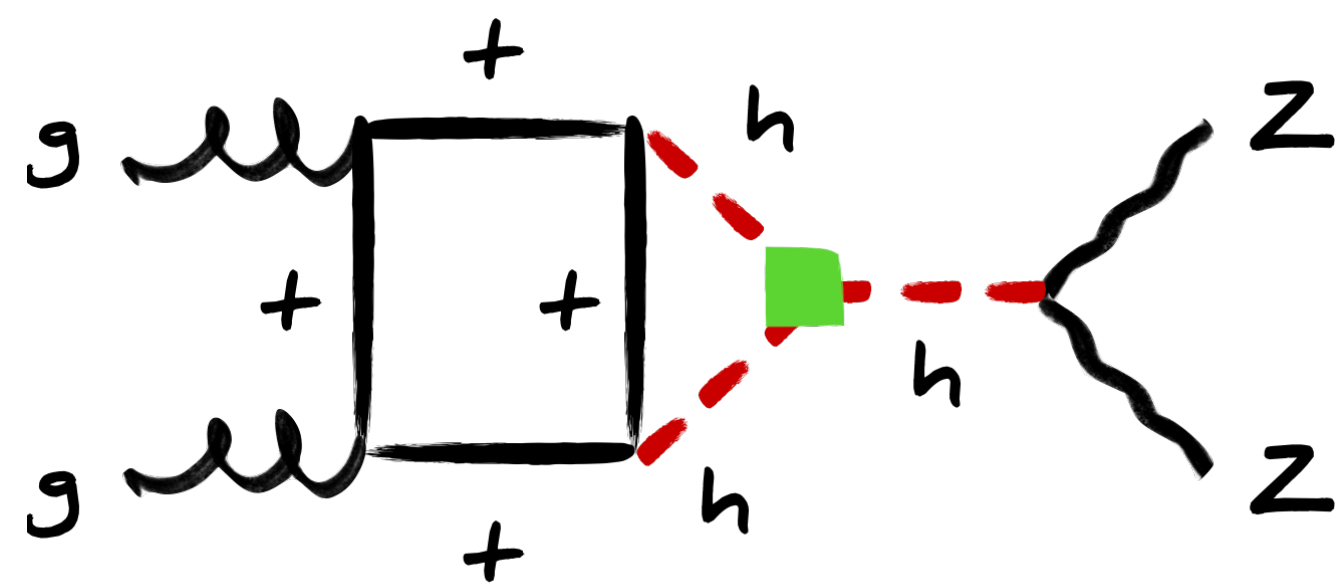
[see Bizon et al., 1810.04665, 2402.03463 for pp→hh & pp→3h FCC studies]

Higgs self-couplings in FCC era



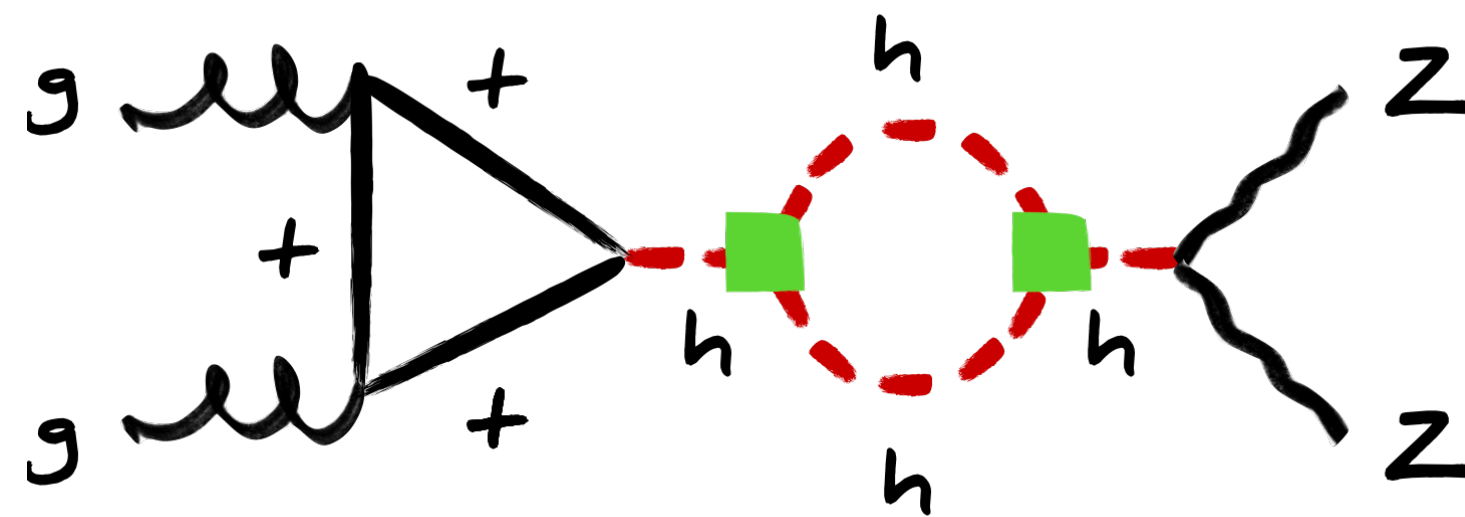
Indirect probes remove degeneracy of direct $pp \rightarrow 3h$ constraint in κ_3 - κ_4 plane

Anatomy of κ_λ corrections



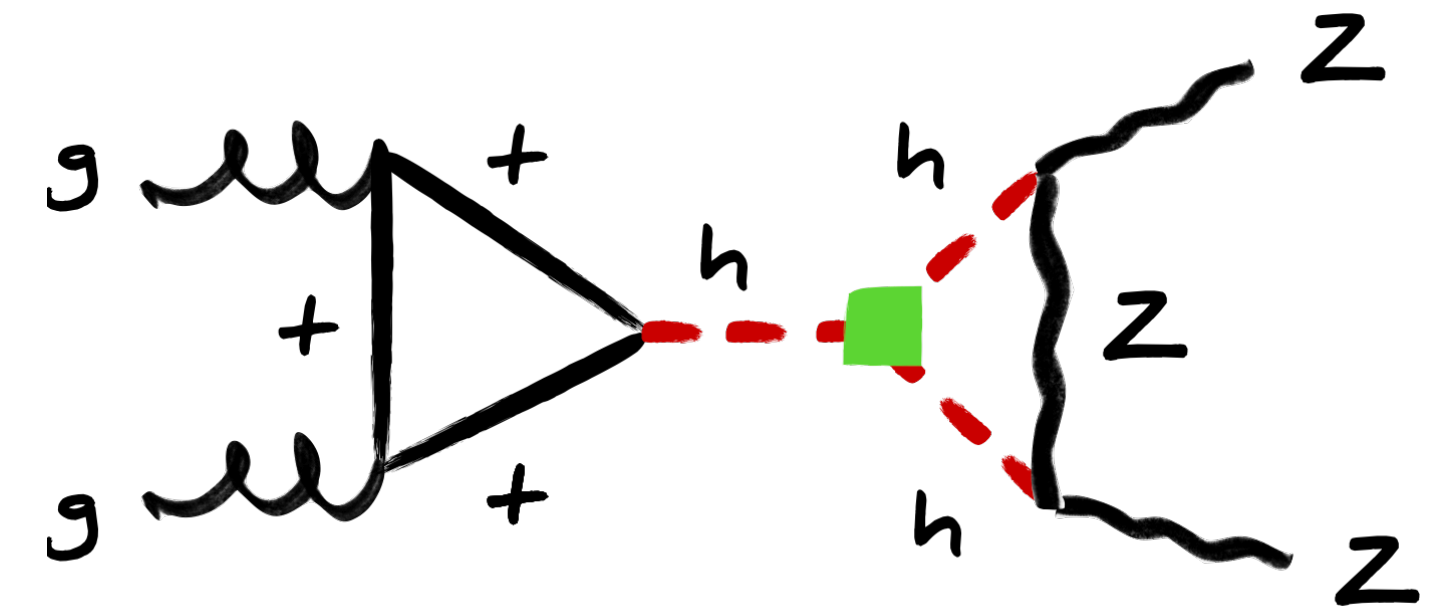
production side:

non-universal correction,
linear κ_λ dependence



Higgs propagator:

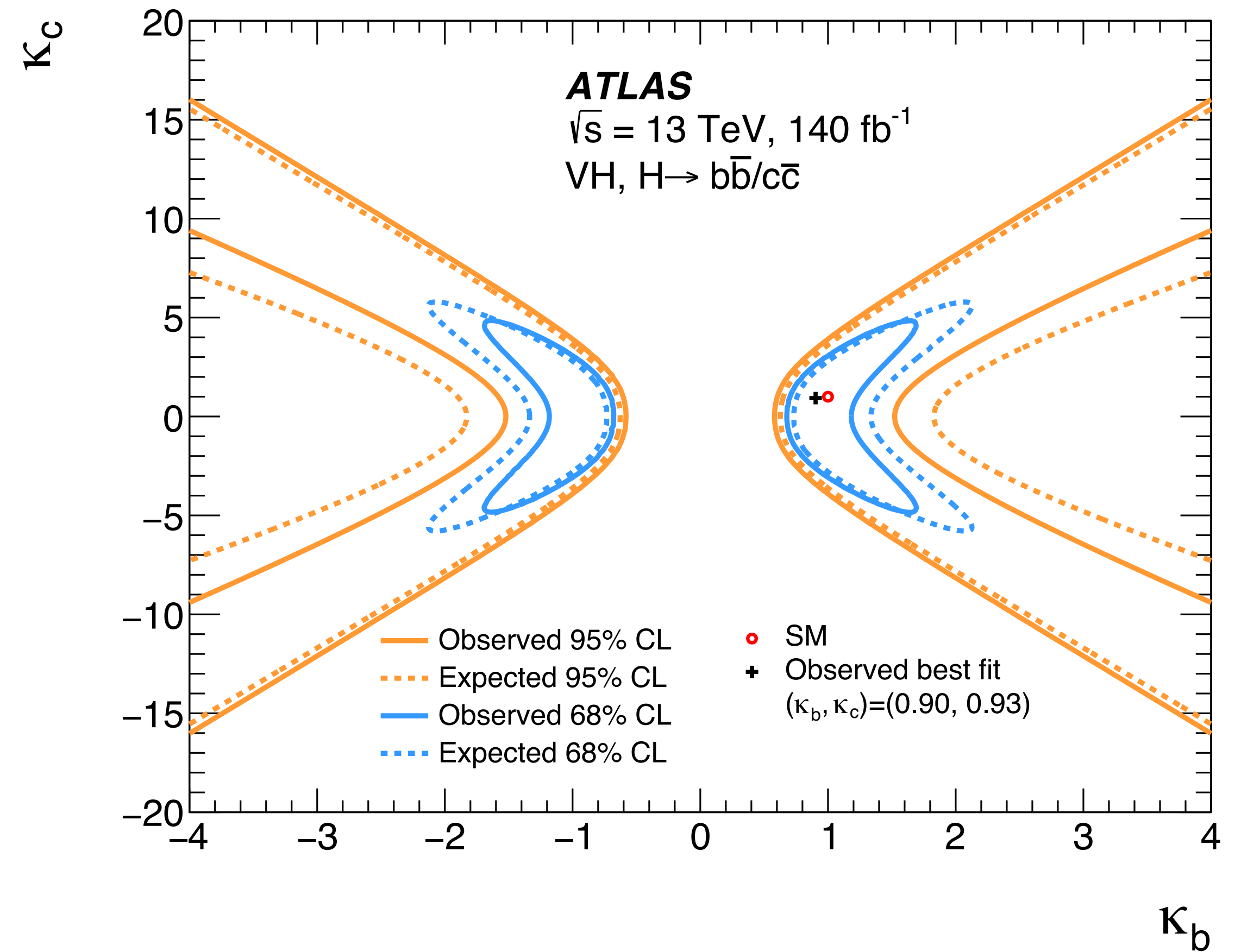
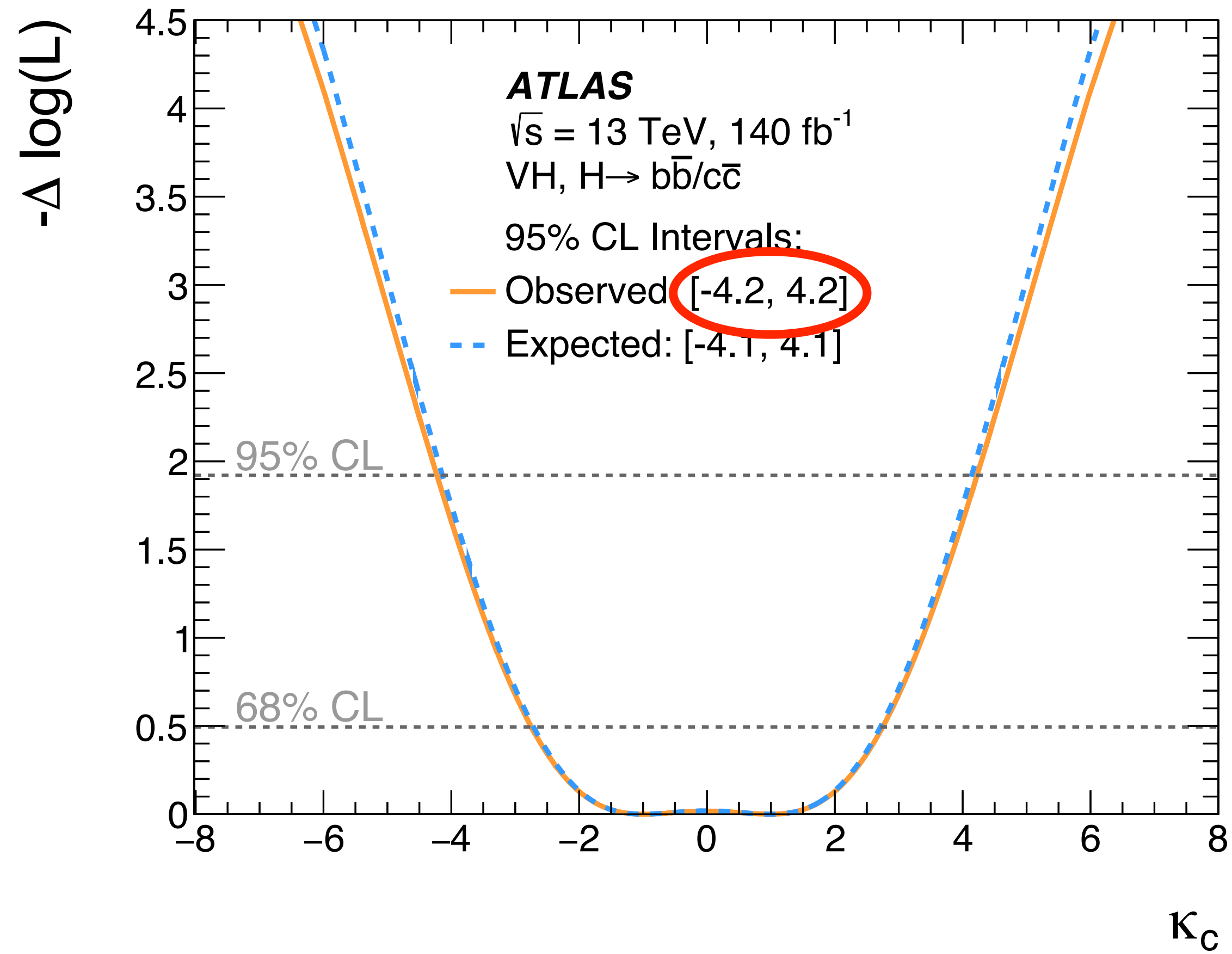
universal correction,
quadratic κ_λ dependence



decay side:

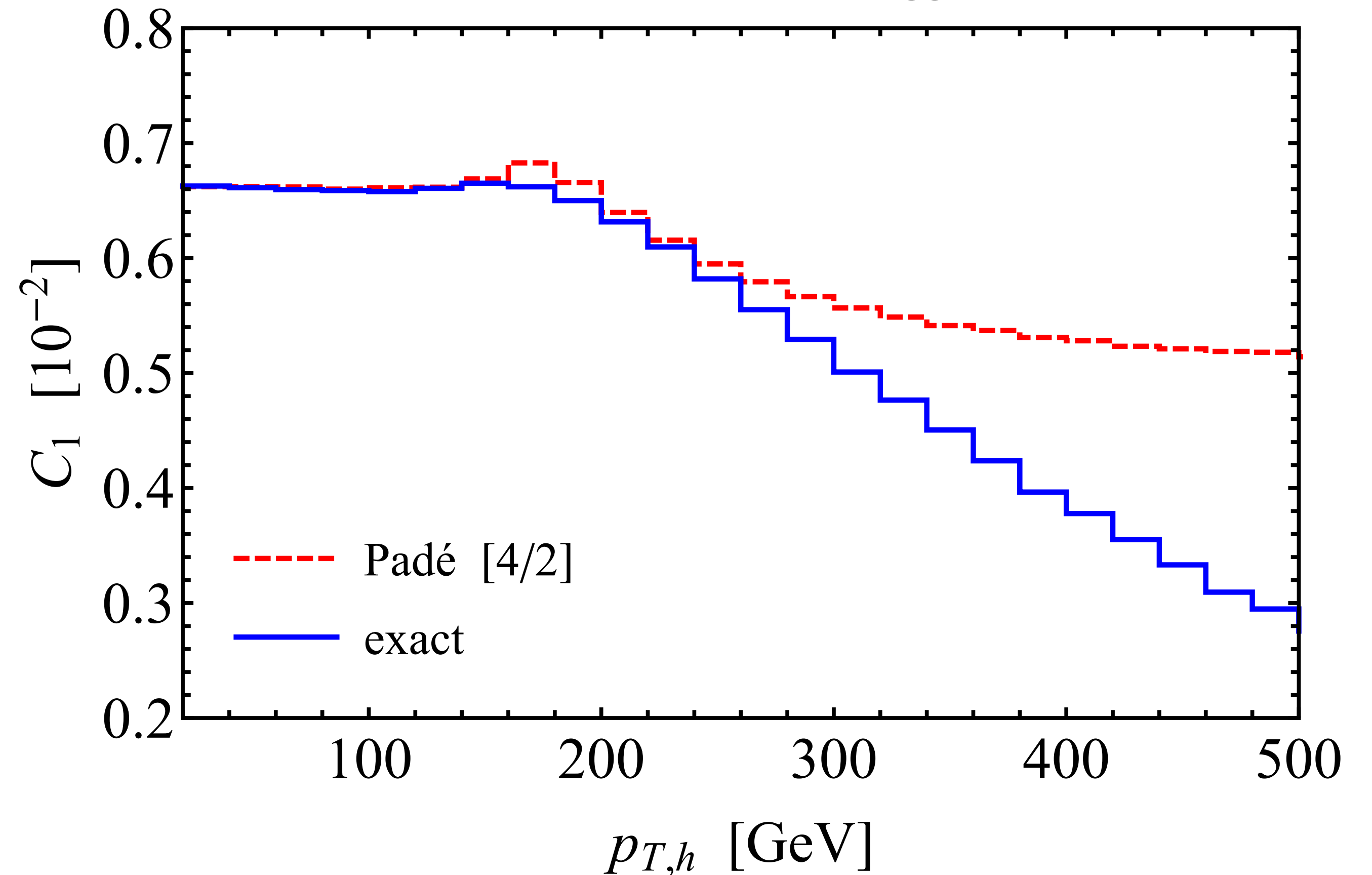
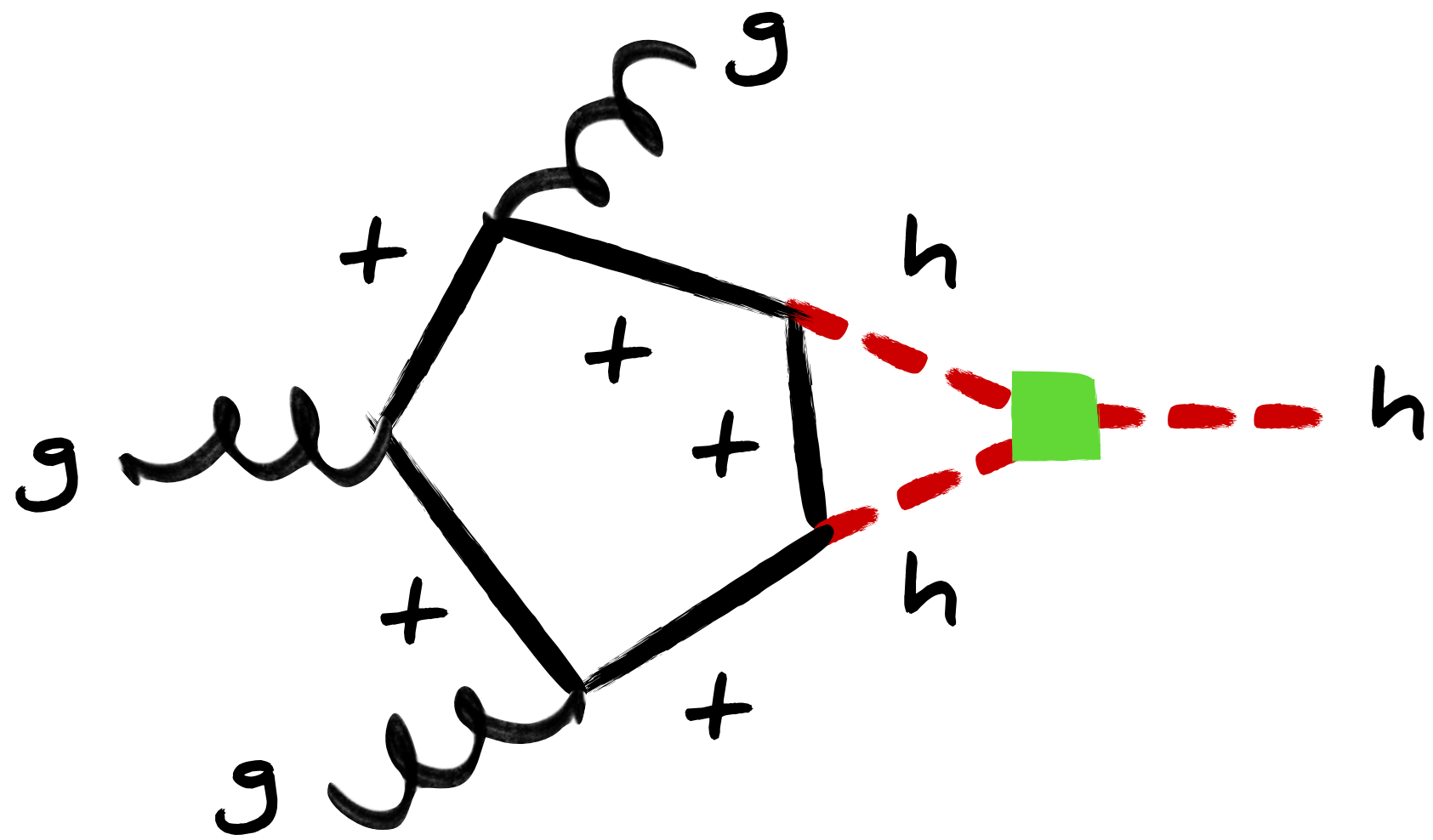
non-universal correction,
linear κ_λ dependence

Direct constraints on κ_b & κ_c @ LHC



Non-universal κ_λ effects in h +jet production

[UH & Niggetiedt, 2408.13186]



Only exact calculation of linear κ_λ terms allows to describe $p_{T,h}$ spectrum above m_t

[Padé approximation obtained in Gao, Shen, Wang, Yang & Zhou, 2302.04160]

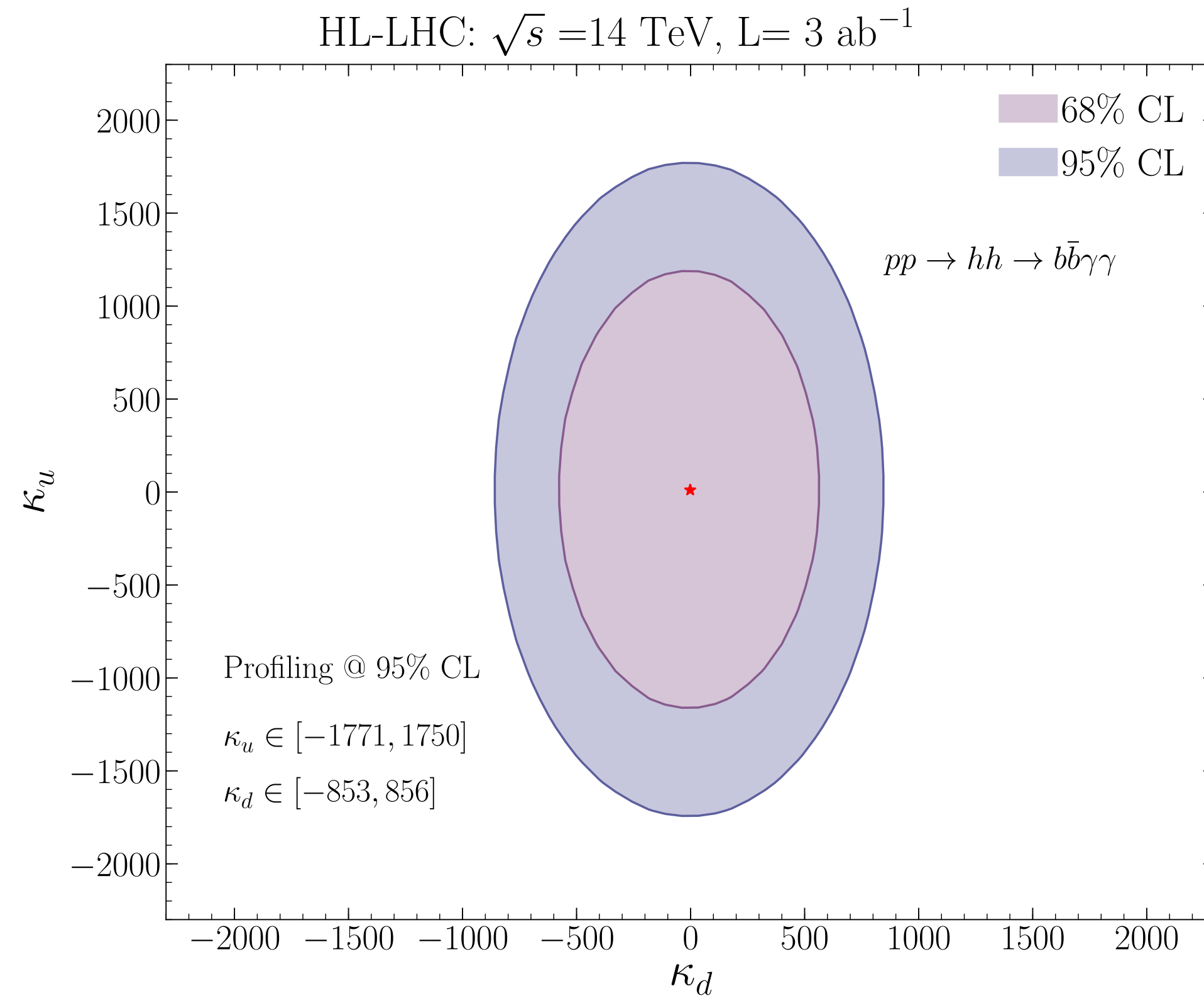
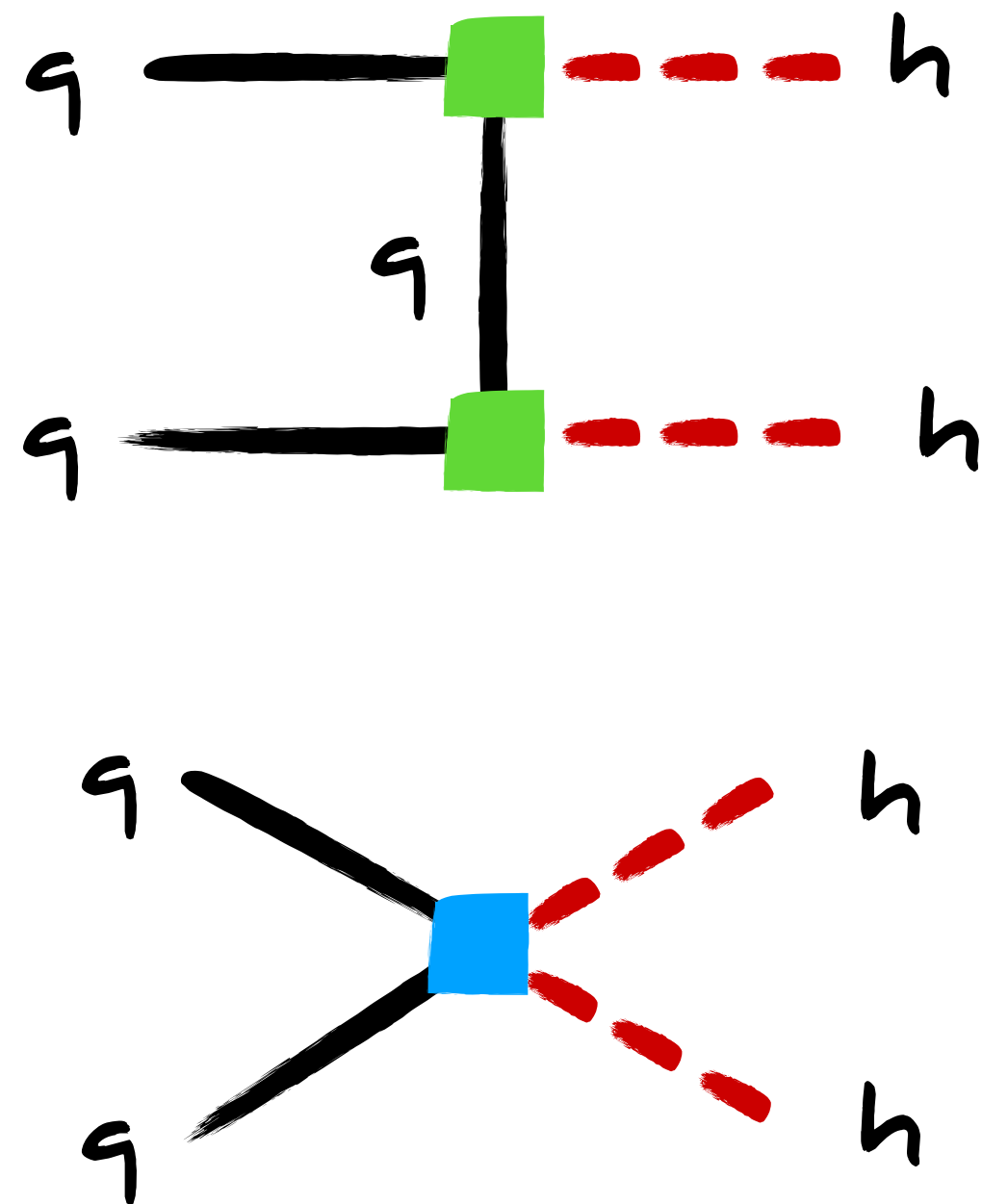
LHC constraints on light-quark Yukawas

[CMS-PAS-HIG-23-011]

Parameter	Scenario	Observed		Expected	
κ_u	float all	$(0.0 \pm 1.5) \times 10^3$	$[-2.4, 2.4] \times 10^3$	$(0.0 \pm 1.8) \times 10^3$	$[-2.6, 2.6] \times 10^3$
κ_u	fix others	$(0.0 \pm 1.4) \times 10^3$	$[-2.3, 2.3] \times 10^3$	$(0.0 \pm 1.6) \times 10^3$	$[-2.5, 2.5] \times 10^3$
κ_d	float all	$(0.0 \pm 7.1) \times 10^2$	$[-1.0, 1.0] \times 10^3$	$(0.0 \pm 7.4) \times 10^2$	$[-1.0, 1.0] \times 10^3$
κ_d	fix others	$(1.5^{+5.0}_{-8.0}) \times 10^2$	$[-9.7, 9.7] \times 10^2$	$(0.0 \pm 6.5) \times 10^2$	$[-9.7, 9.7] \times 10^2$
κ_s	float all	0^{+33}_{-34}	$[-46, 44]$	1^{+32}_{-31}	$[-44, 42]$
κ_s	fix others	11^{+19}_{-42}	$[-44, 42]$	1^{+26}_{-30}	$[-41, 40]$
κ_c	float all	$0.0^{+2.7}_{-3.0}$	$[-4.0, 3.4]$	$1.0^{+1.4}_{-3.8}$	$[-3.8, 3.2]$
κ_c	fix others	$1.4^{+1.2}_{-4.4}$	$[-4.0, 3.5]$	$1.0^{+1.3}_{-3.8}$	$[-3.8, 3.2]$
Γ_H^{BSM} (MeV)	float all	$0.0^{+0.9}_{-0.0}$	< 1.6	$0.0^{+0.7}_{-0.0}$	< 1.4

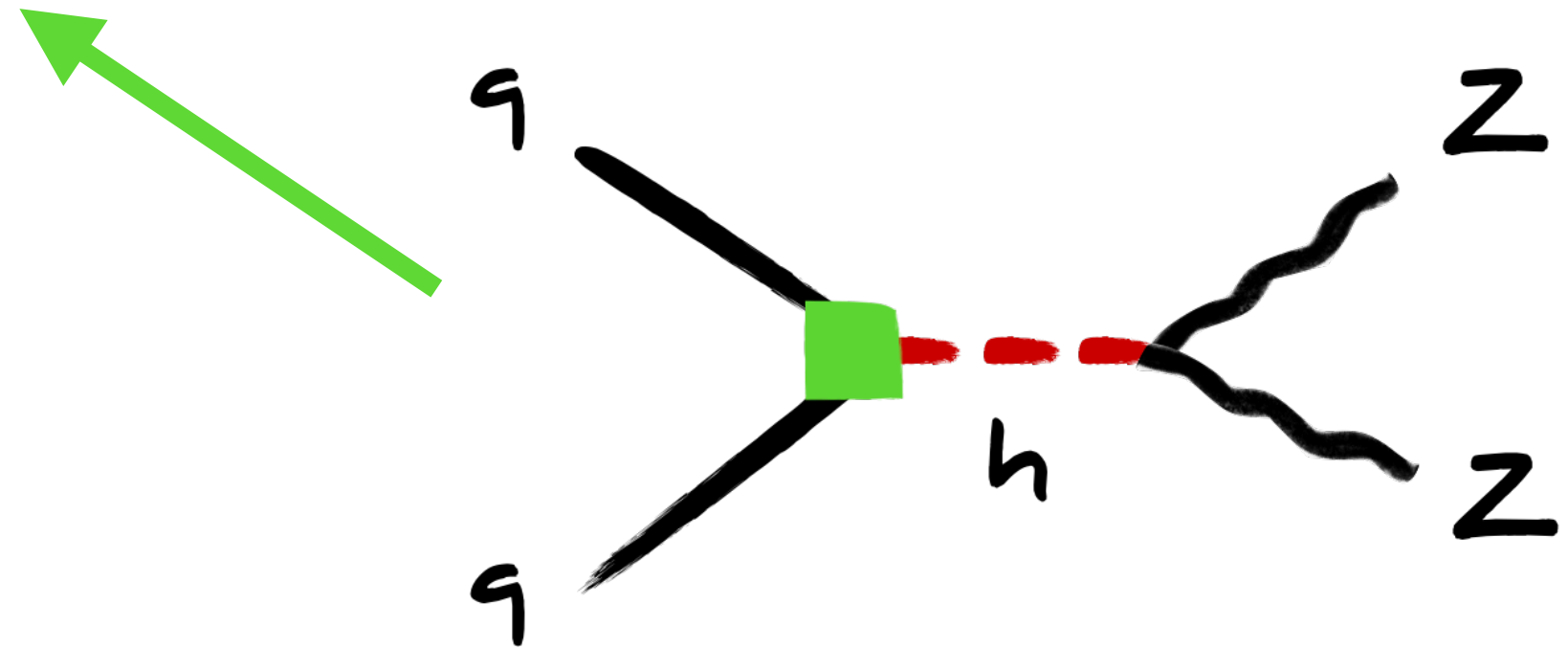
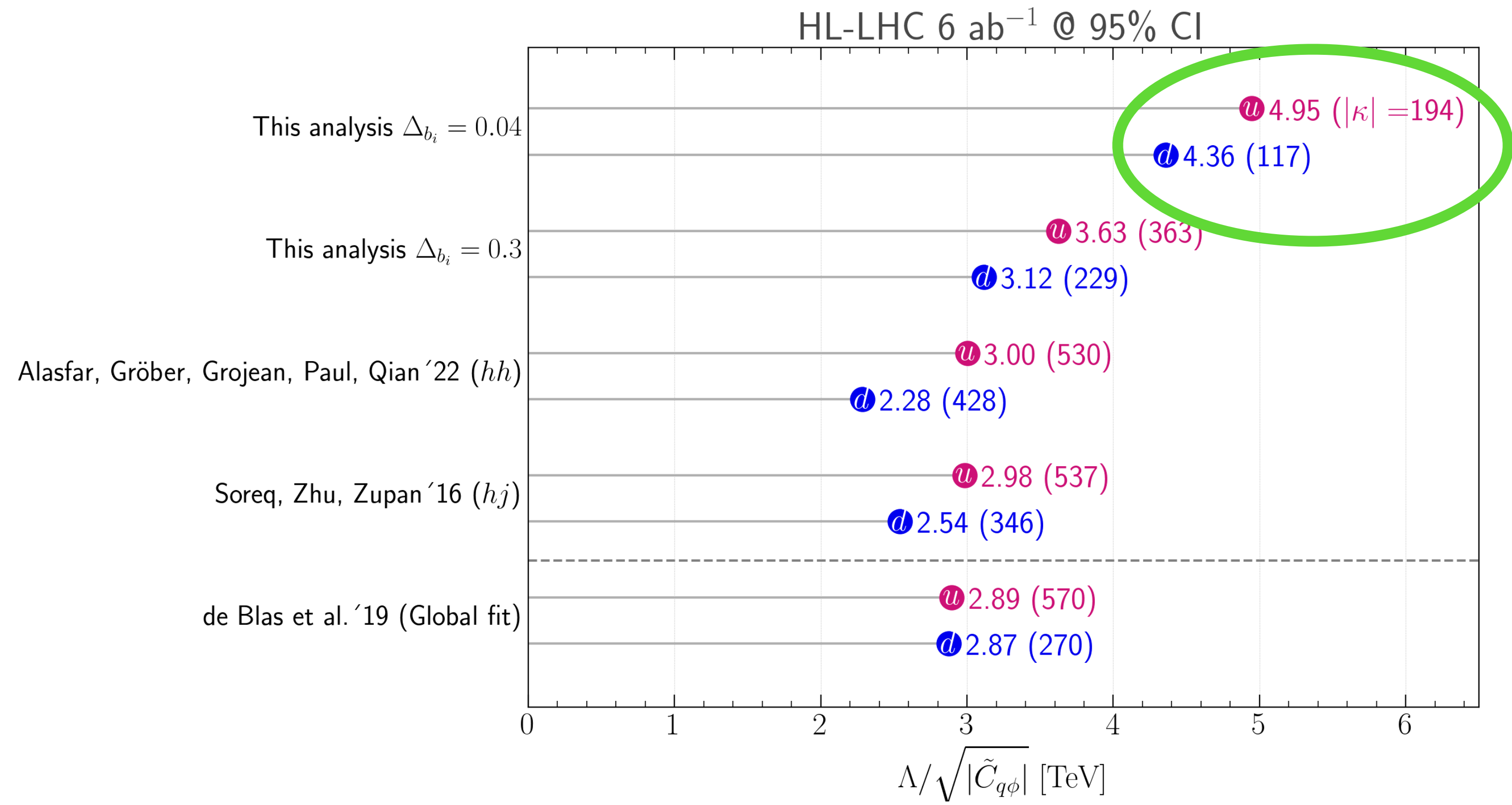
First indirect bounds on κ_q modifiers obtained from production rate of $h \rightarrow 4l$

hh constraints on light-quark Yukawas



$q\bar{q} \rightarrow hh$ process gives access to light-quark Yukawas. Dominant effect arises from possible Higgs non-linearities, which induce a 4-point interaction

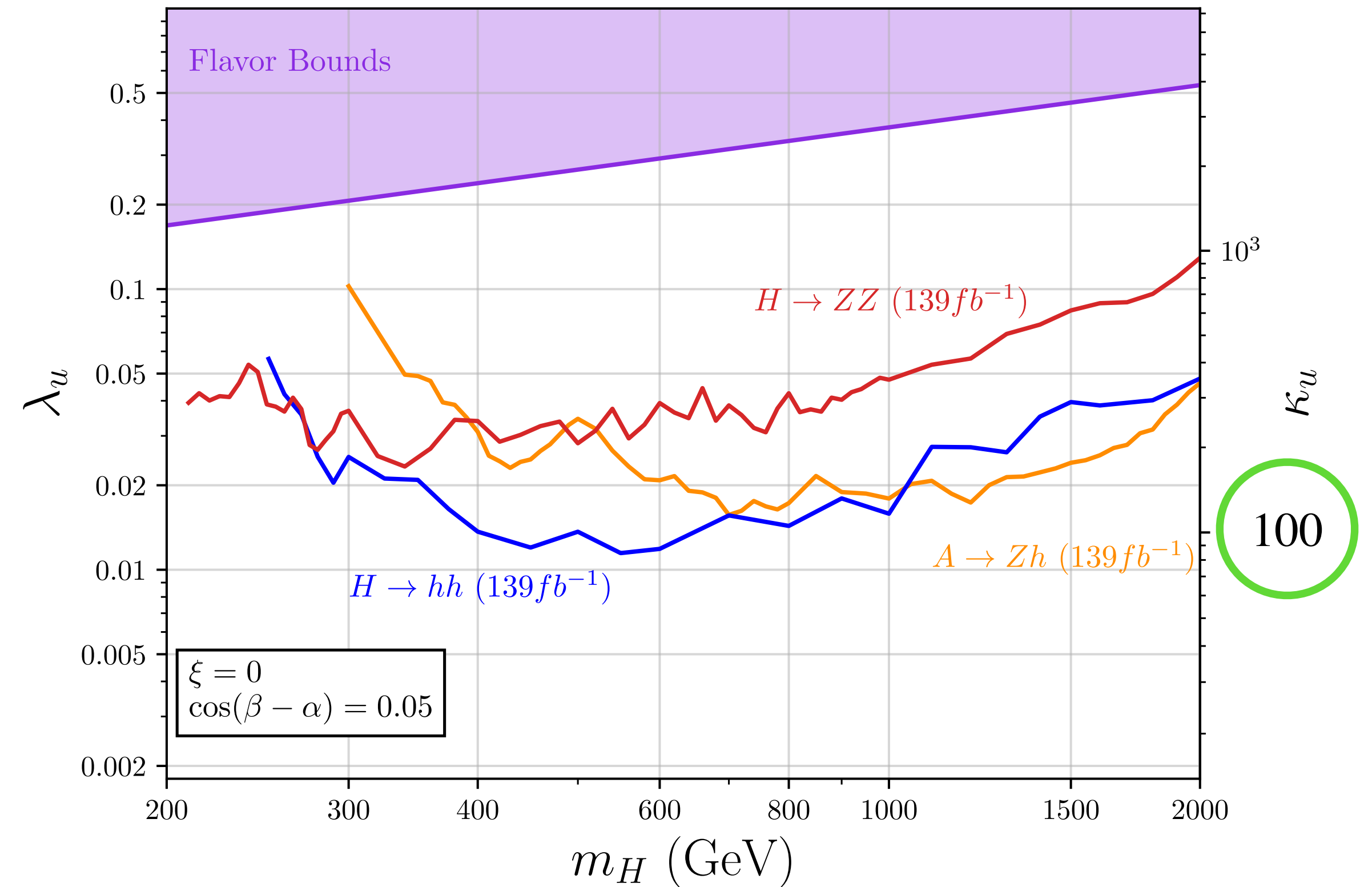
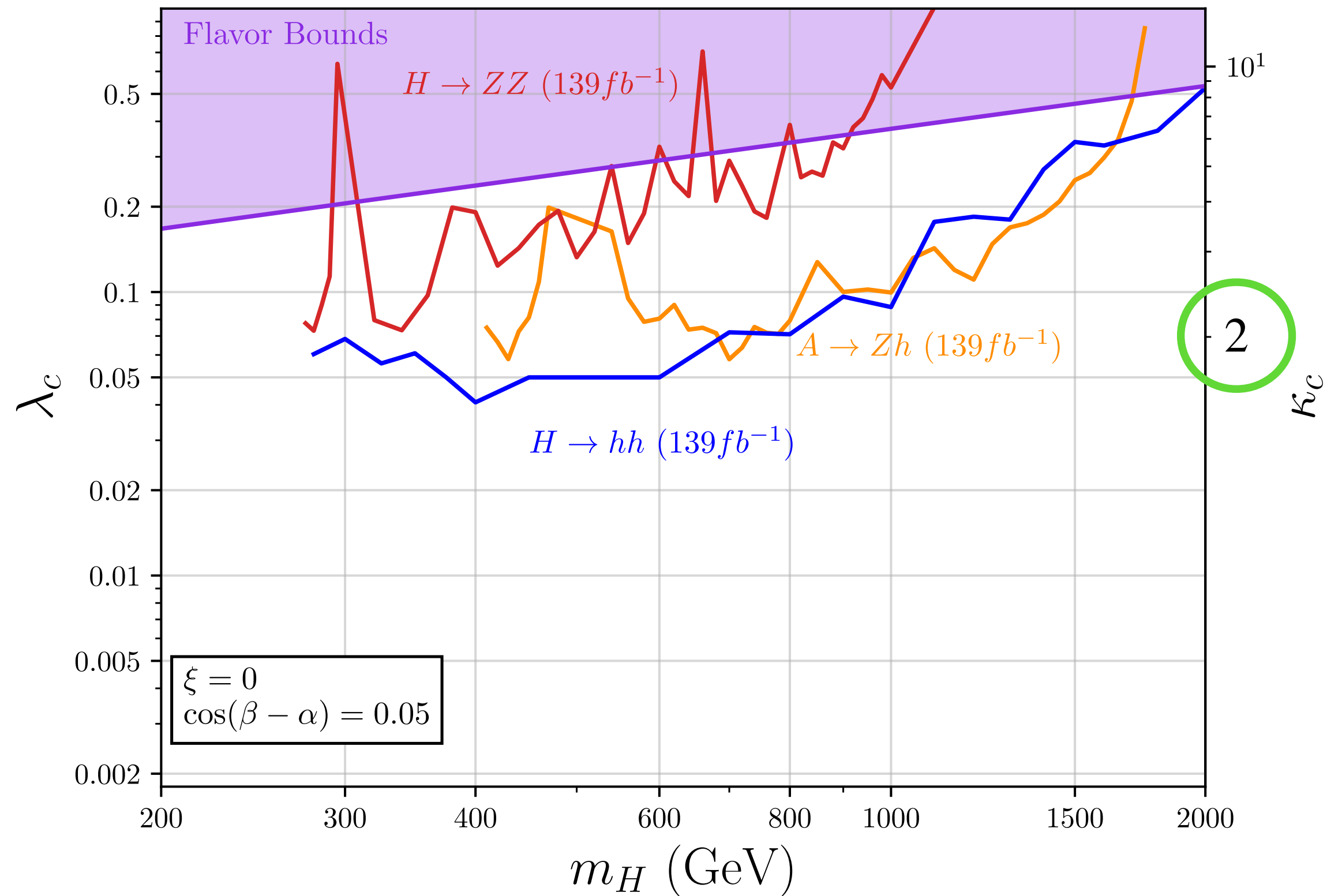
Light-quark Yukawas from $pp \rightarrow 4l$



Using D_S , $pp \rightarrow 4l$ found to be promising channel to extract information on κ_q 's

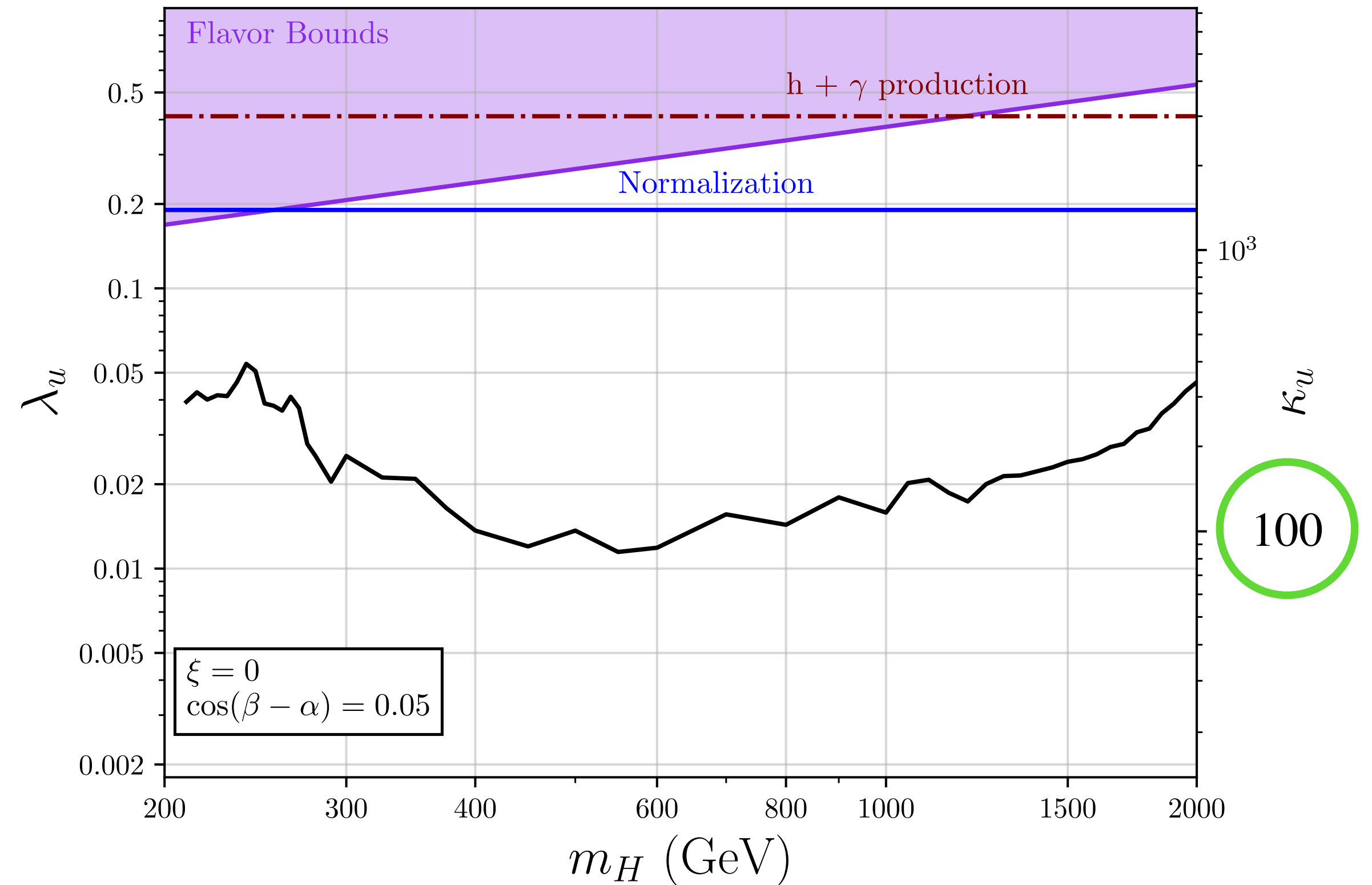
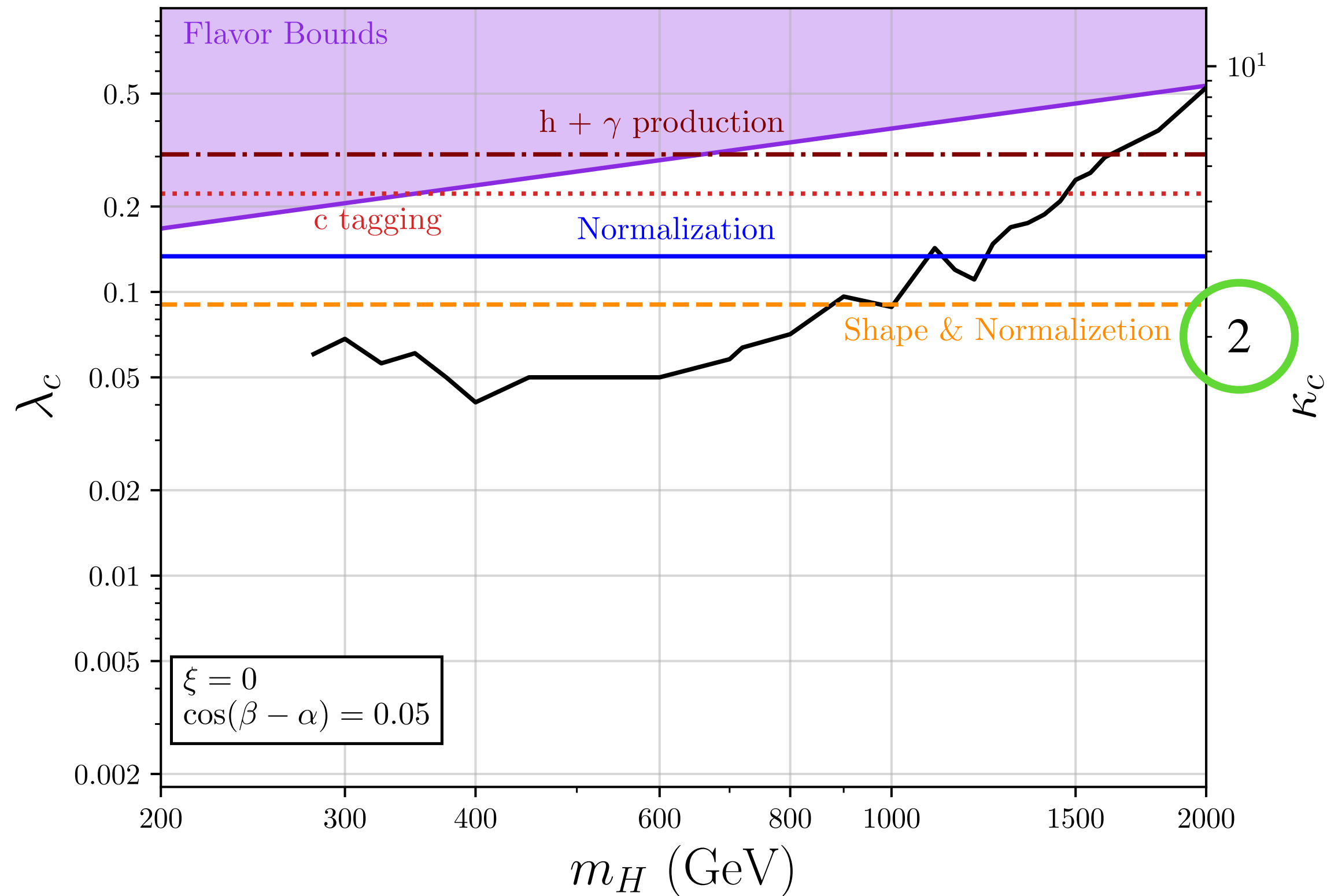
[Balzani, Gröber & Vitti, 2304.09772]

How charming/upsy can Higgs be?



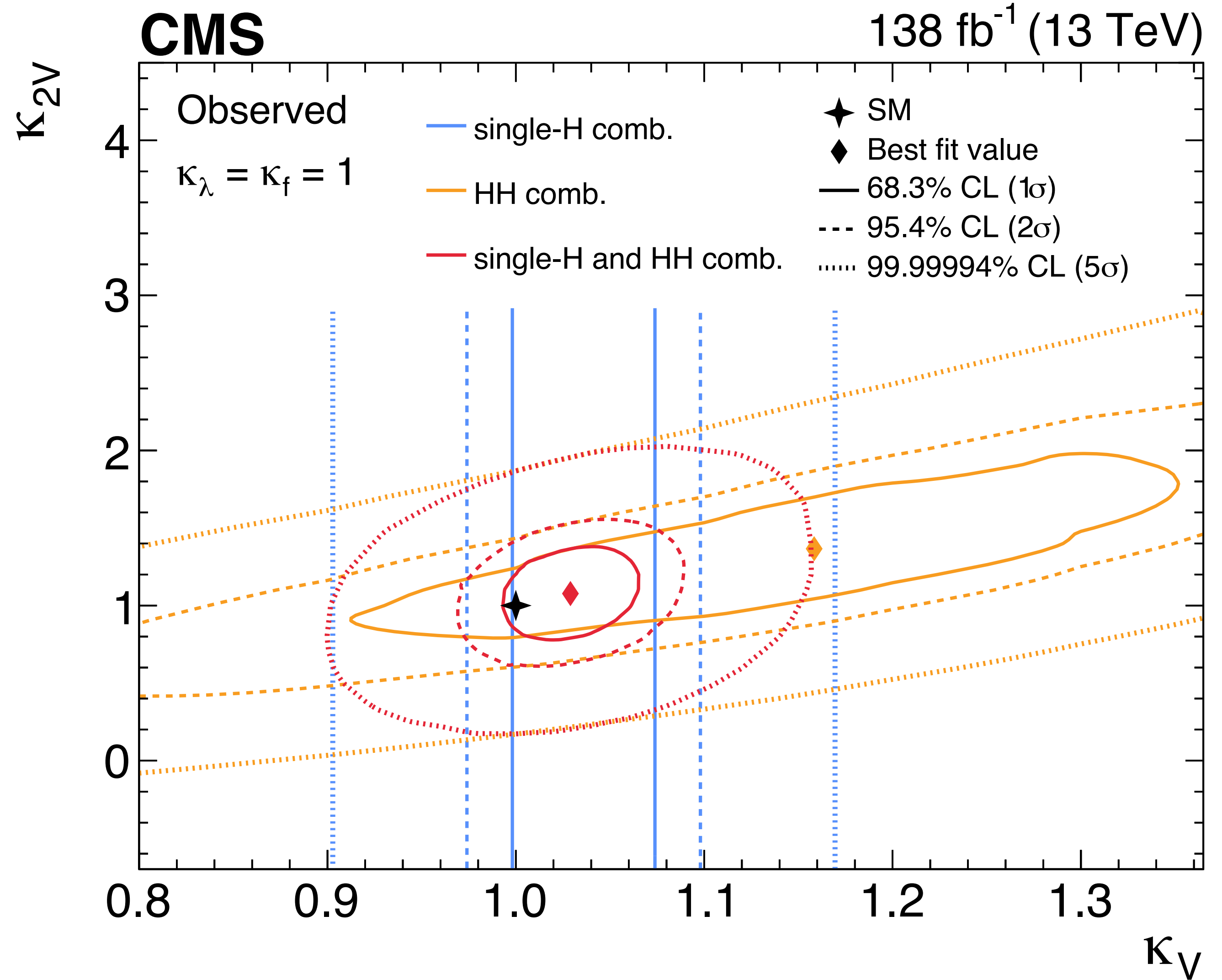
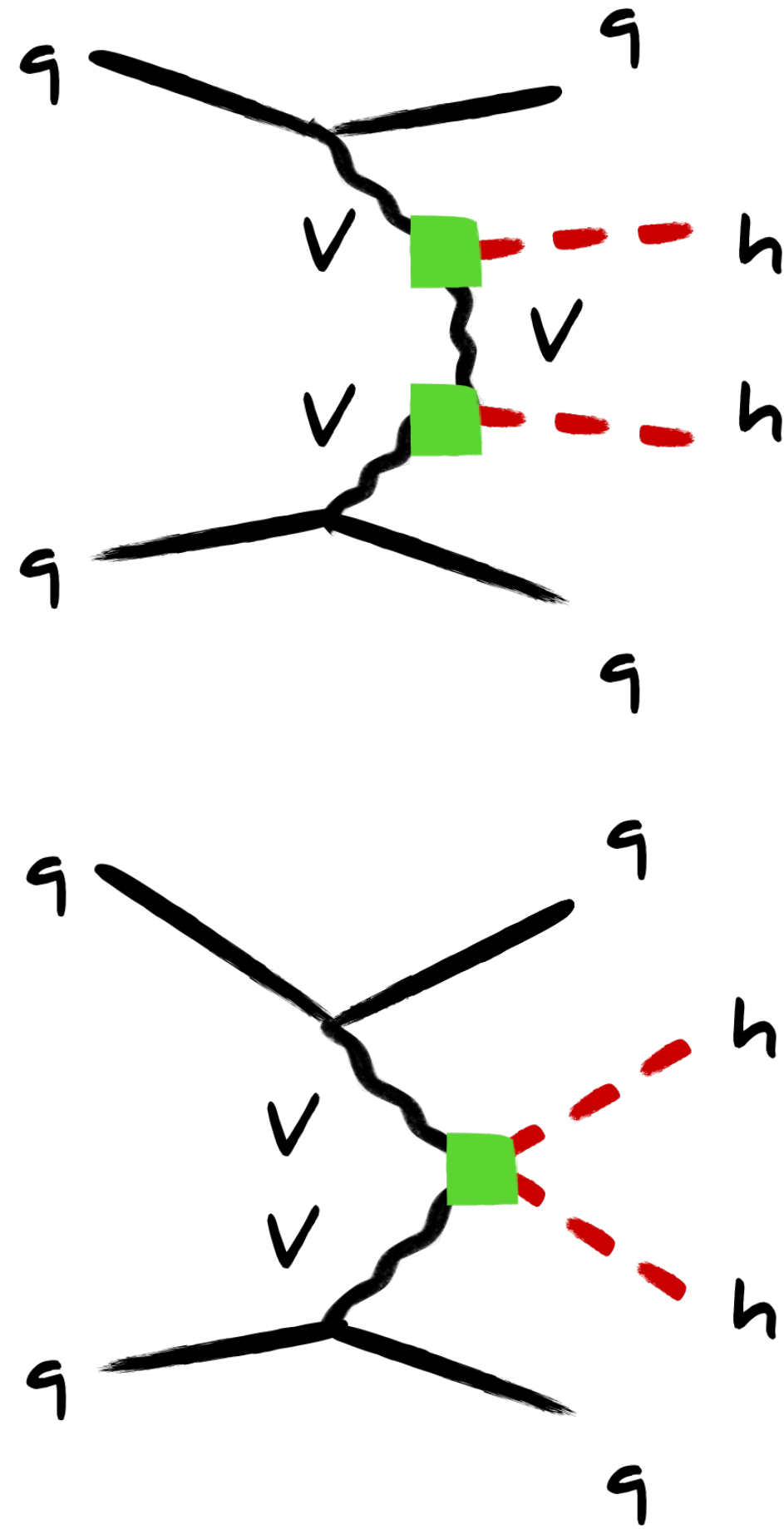
In 2HDM with non-alignment & down-type spontaneous flavor violation,
LHC constraints from heavy Higgs searches stronger than flavor physics

How charming/upsy can Higgs be?



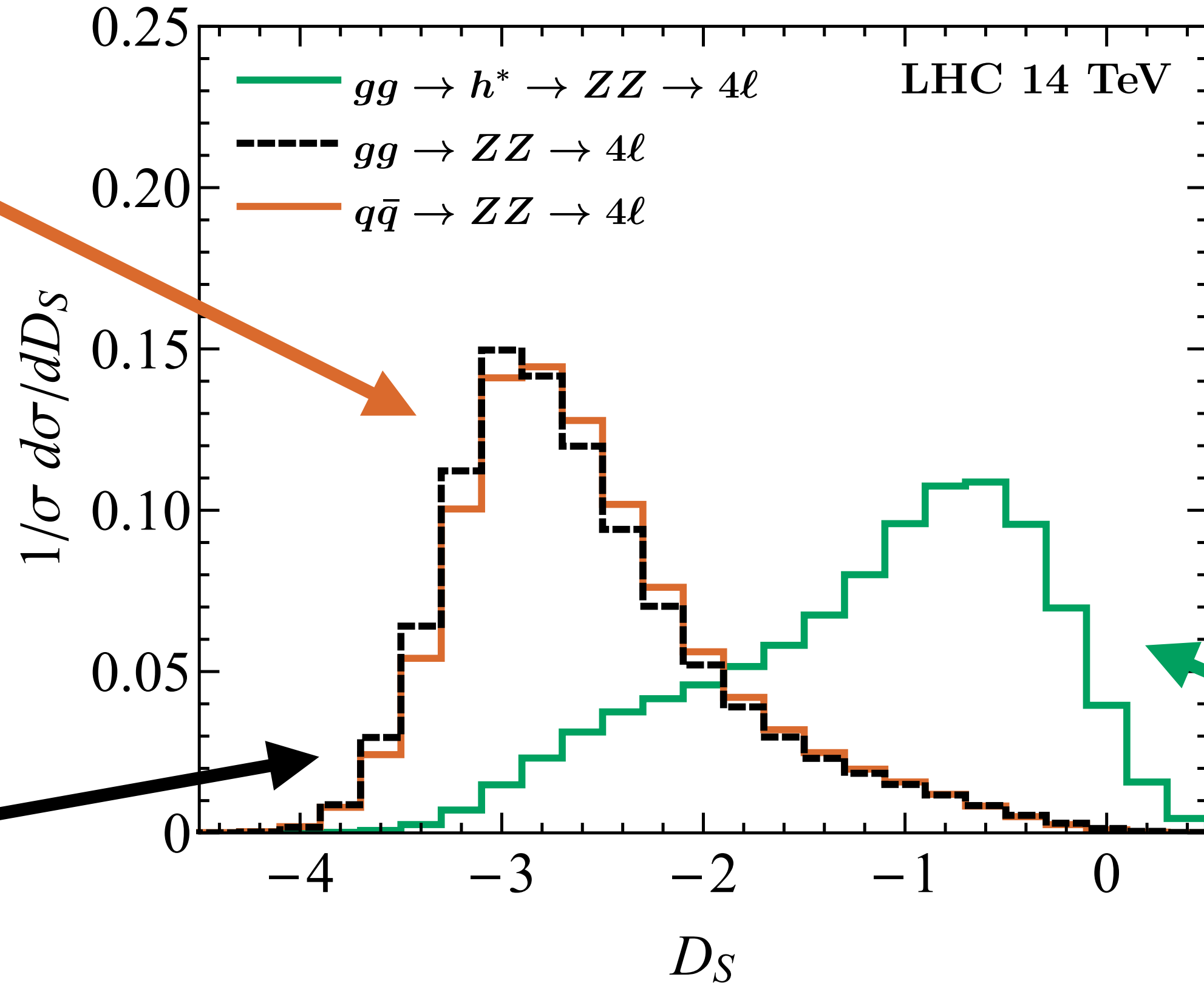
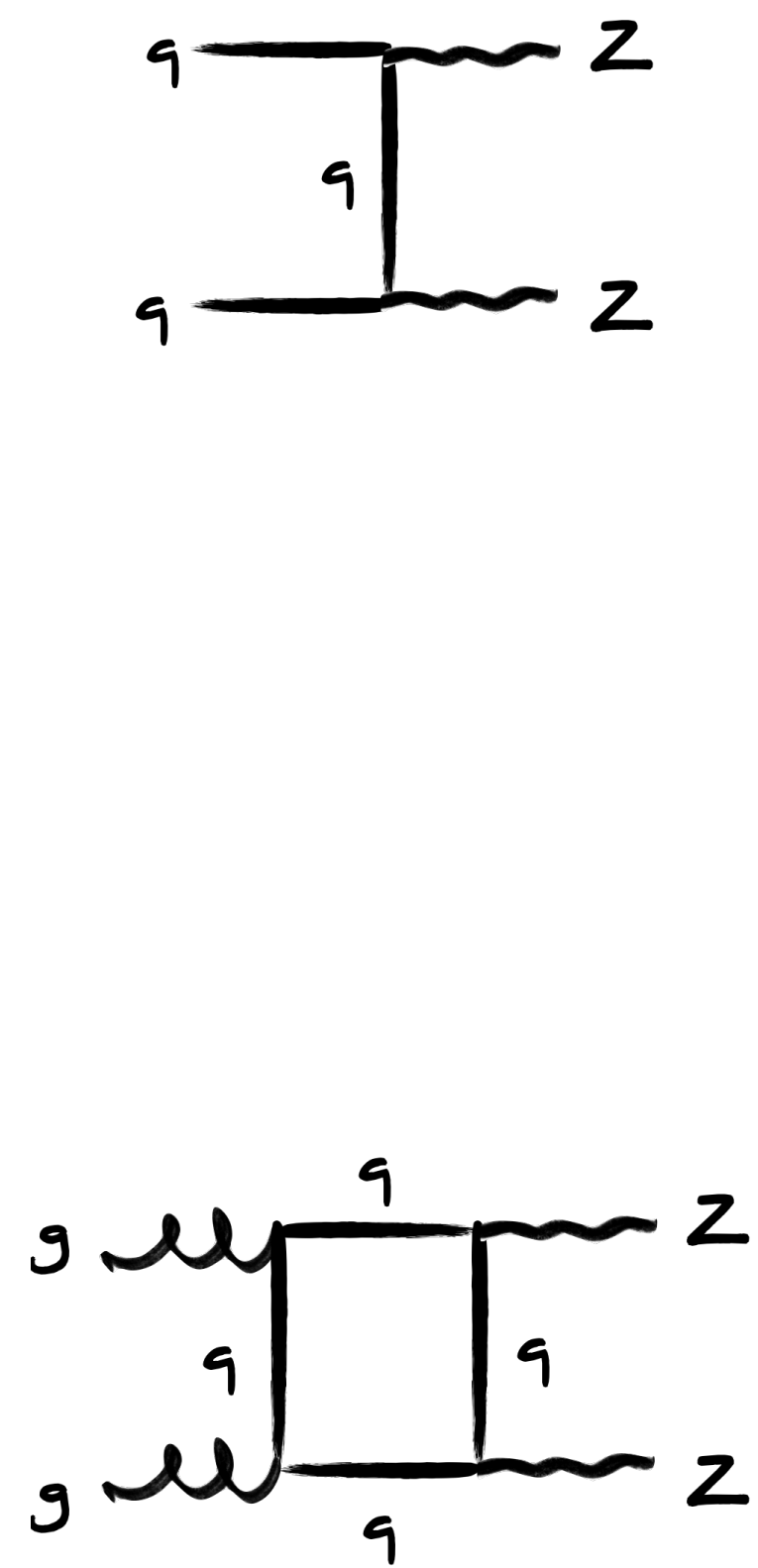
Direct resonant searches provide model-dependent bounds of $|\kappa_c| < O(2)$ & $|\kappa_u| < O(100)$, typically better what can be achieved by other means @ LHC

Synergy & complementarity: κ_V & κ_{2V}

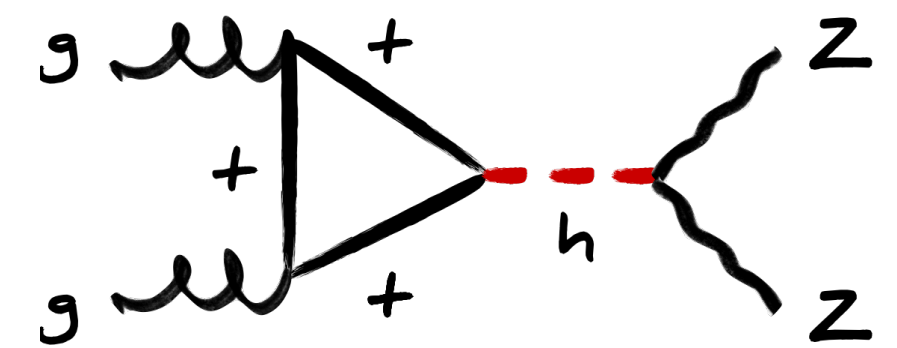


Matrix-element based discriminators

[UH & Koole, 2111.12589]



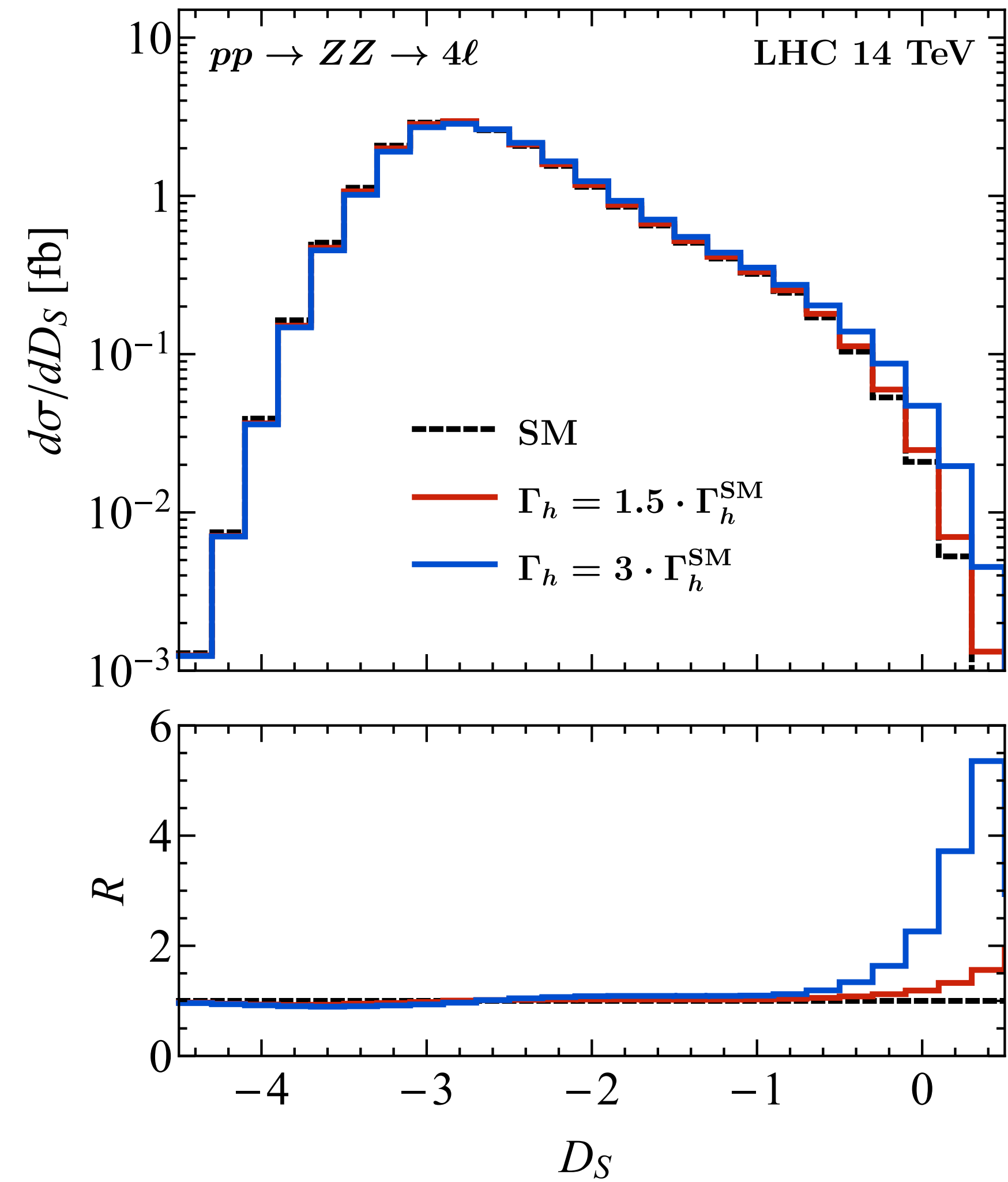
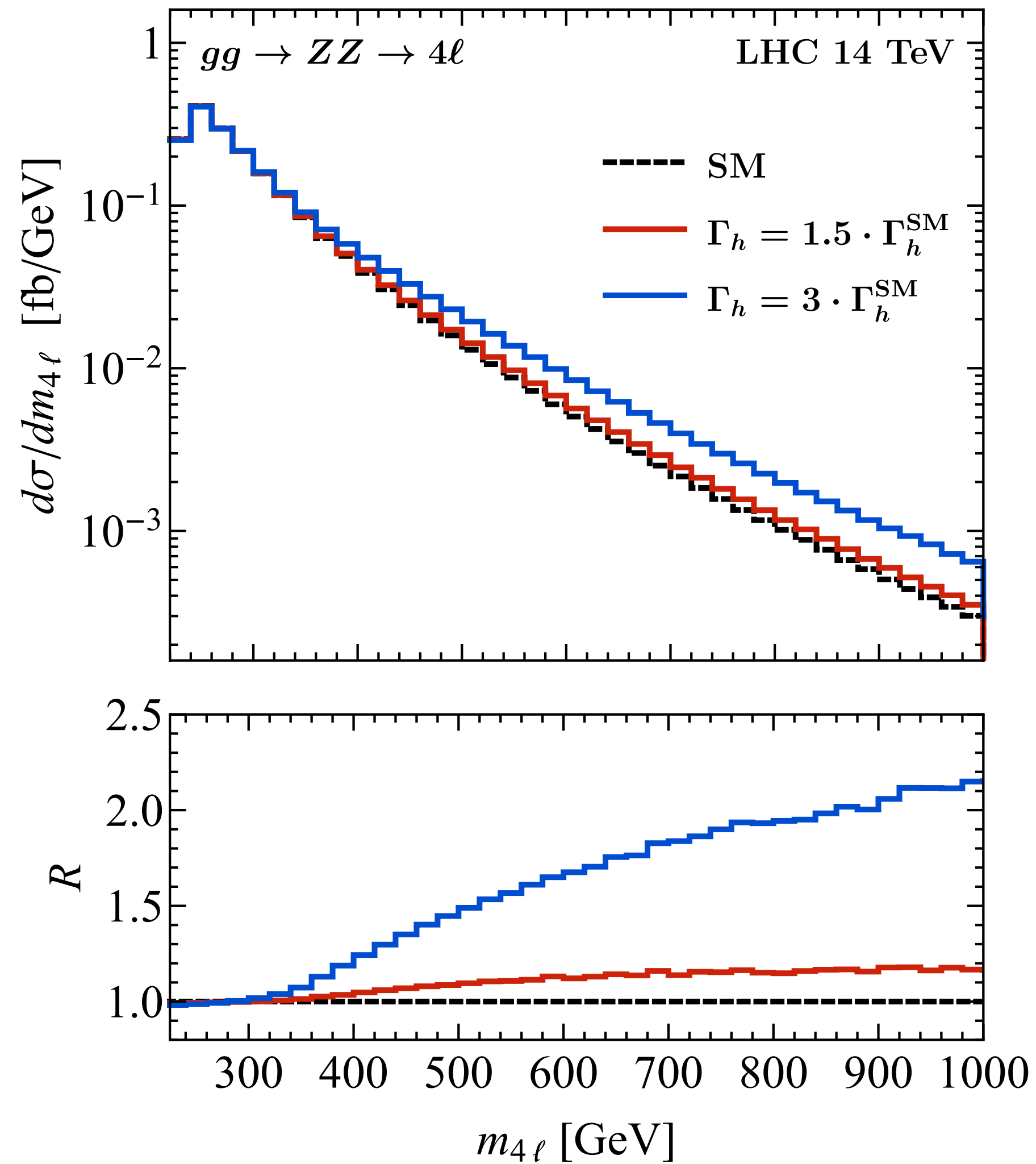
$$D_S = \log_{10} \left(\frac{P_h}{P_{gg} + c \cdot P_{q\bar{q}}} \right)$$



[D_S introduced in Campbell, Ellis & Williams, 1311.3589, 1312.1628]

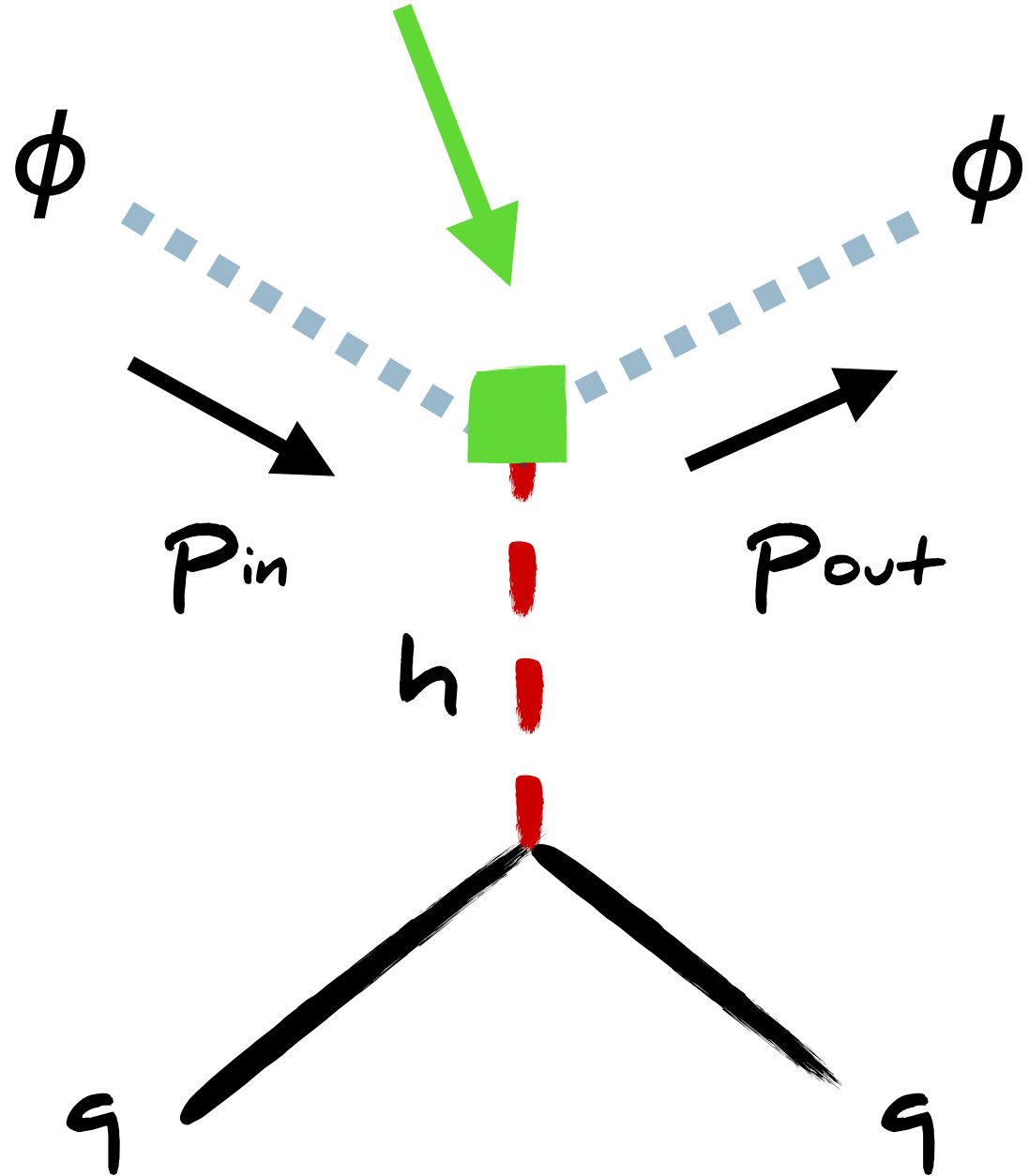
Higgs width measurements: $m_{4\ell}$ vs. D_S

[UH & Koole, 2111.12589]

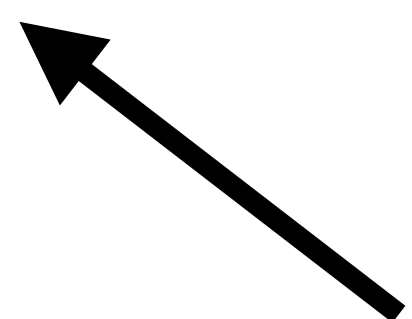


Derivative Higgs portal: DM-N scattering

$$\frac{c_d}{f^2} (\partial_\mu |H|^2) (\partial^\mu |\phi|^2)$$



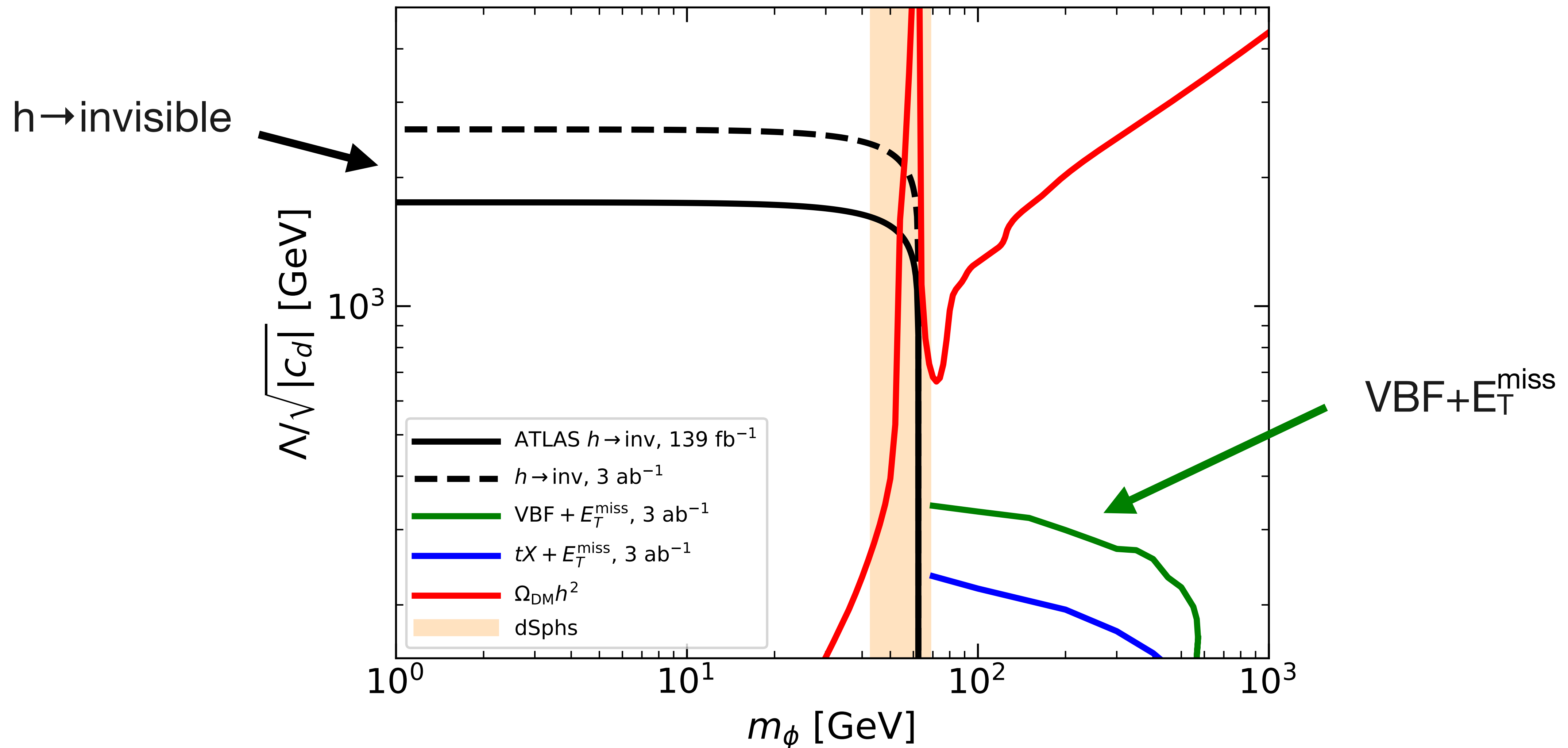
$$\propto \frac{(p_{in} - p_{out})^2}{f^2} \lesssim \frac{(100 \text{ MeV})^2}{f^2}$$



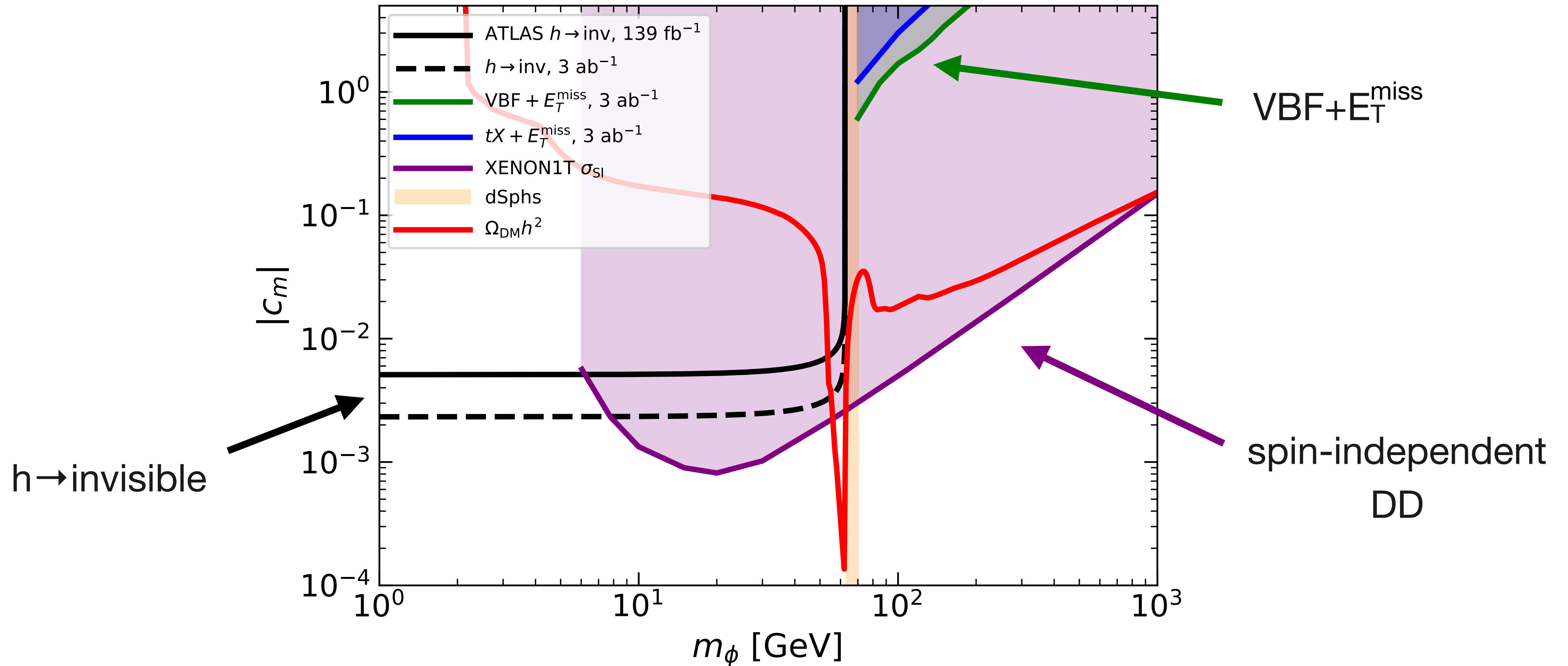
Due to momentum suppression, DD limits easily avoided for new-physics scales f of $O(1 \text{ TeV})$

[see for example Balkin, Ruhdorfer, Salvioni & Weiler, 1809.09106]

Global picture of derivative Higgs portal



Global picture of marginal Higgs portal



Future probes of neutrino floor

