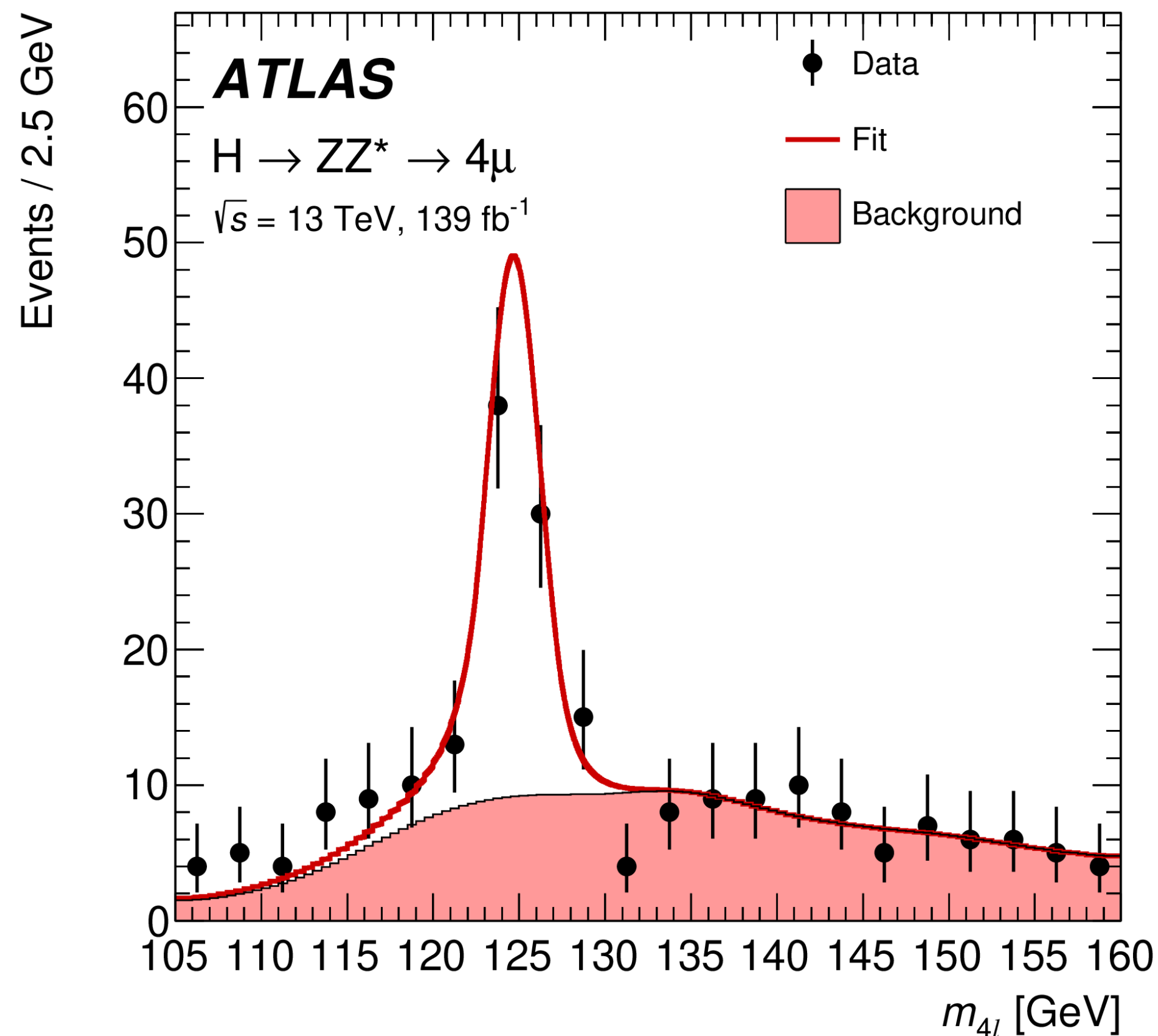


Indirect measurements of the Higgs Boson natural width with the ATLAS detector

Will Leight, for the ATLAS Experiment
HIGGS2024, Uppsala, Sweden

Off-shell Production and the Higgs Boson Width

[Phys. Lett. B 843 \(2023\) 137880](#)

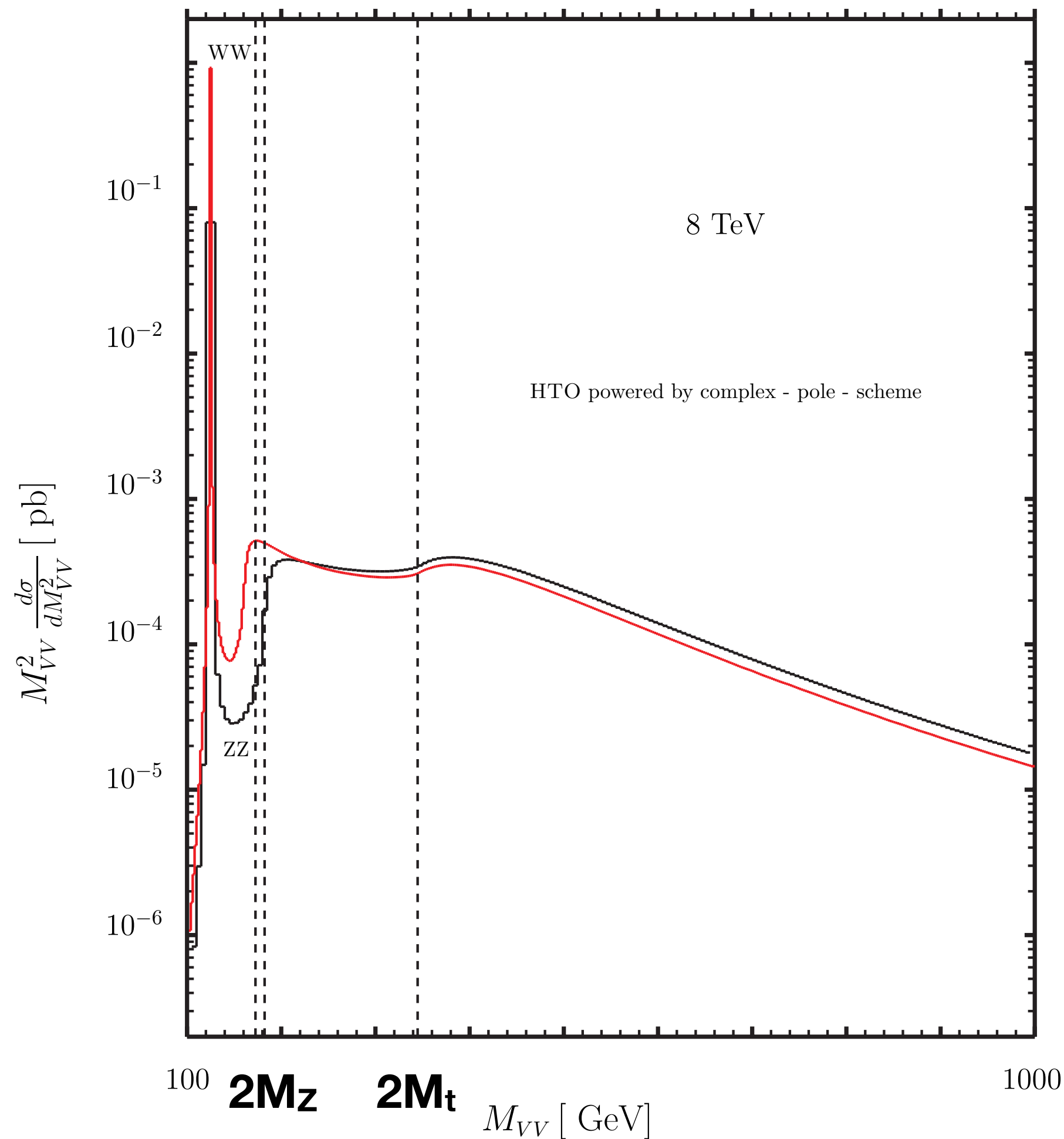


- Higgs peak in the $H \rightarrow ZZ \rightarrow 4\mu$ channel
- Best achievable resolution is still $>1 \text{ GeV}$
- It is challenging to measure Γ_H of 4.1 MeV at ATLAS from the on-shell lineshape
- But it can be measured using off-shell production, which does not depend on Γ_H

Off-shell Production in $H \rightarrow ZZ$

Kauer & Passarino, *JHEP* 2012, 116

$2 M_W$

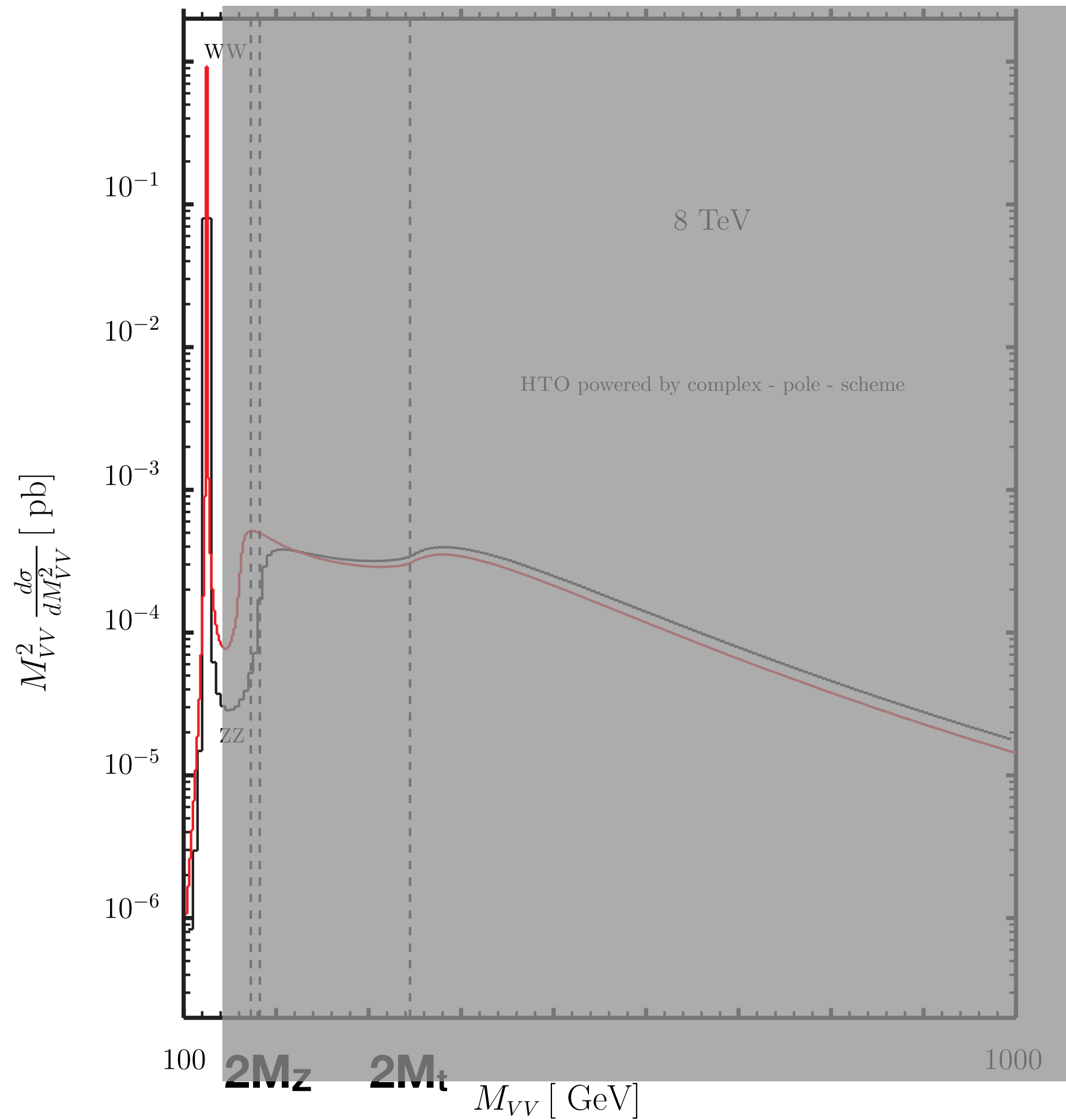


- The $H \rightarrow ZZ$ channel is a good candidate for measuring the width using off-shell production
 - Decays in this channel are enhanced when both Z bosons are on-shell

Off-shell Production in $H \rightarrow ZZ$

Kauer & Passarino, JHEP 2012, 116

$2 M_W$



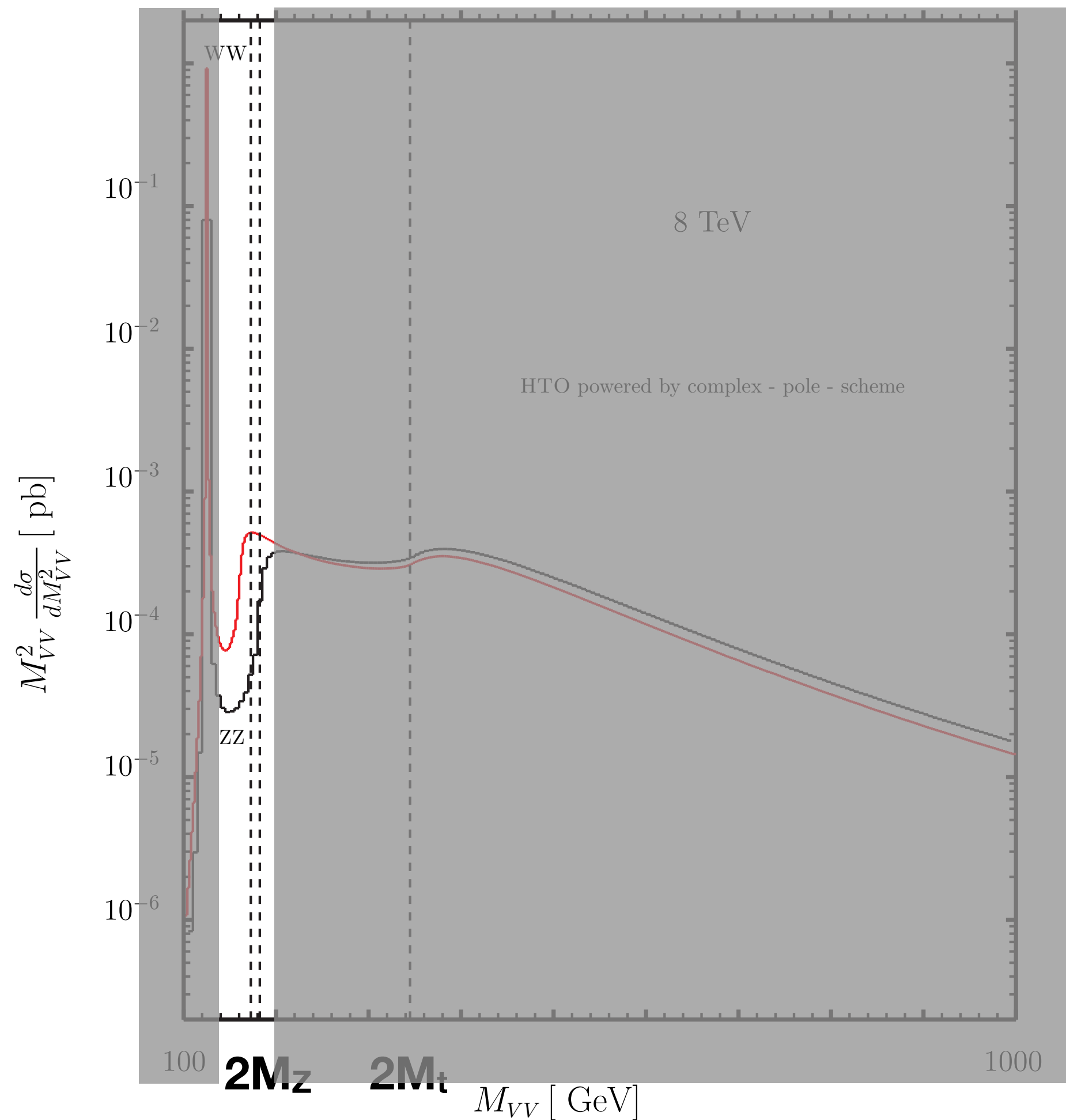
- On-shell, lineshape is Breit-Wigner, hence $\sigma_{\text{on-shell}} \sim 1/\Gamma_H$

$$\frac{1}{(q^2 - M_H^2)^2 + \Gamma_H^2 M_H^2}$$

Off-shell Production in $H \rightarrow ZZ$

Kauer & Passarino, *JHEP* 2012, 116

$2 M_W$

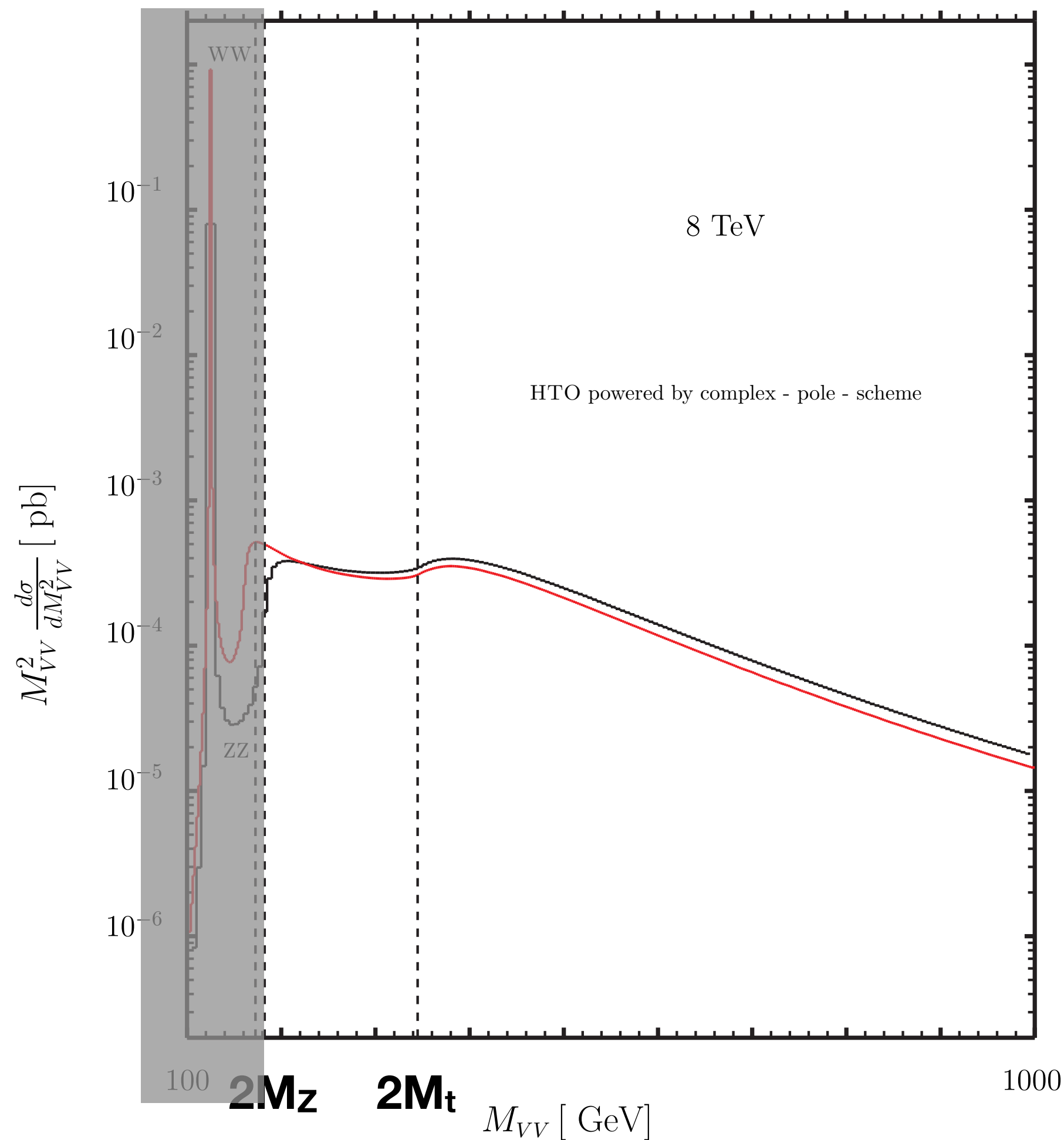


- But at higher masses, this is no longer the case
- In the ZZ channel, there is an increase in the cross-section as m_{ZZ} approaches $2m_Z$
 - There is a dropoff due to the q^2 dependence of the propagator
 - But the increase in decay phase space is much larger

Off-shell Production in $H \rightarrow ZZ$

Kauer & Passarino, *JHEP* 2012, 116

$2 M_W$

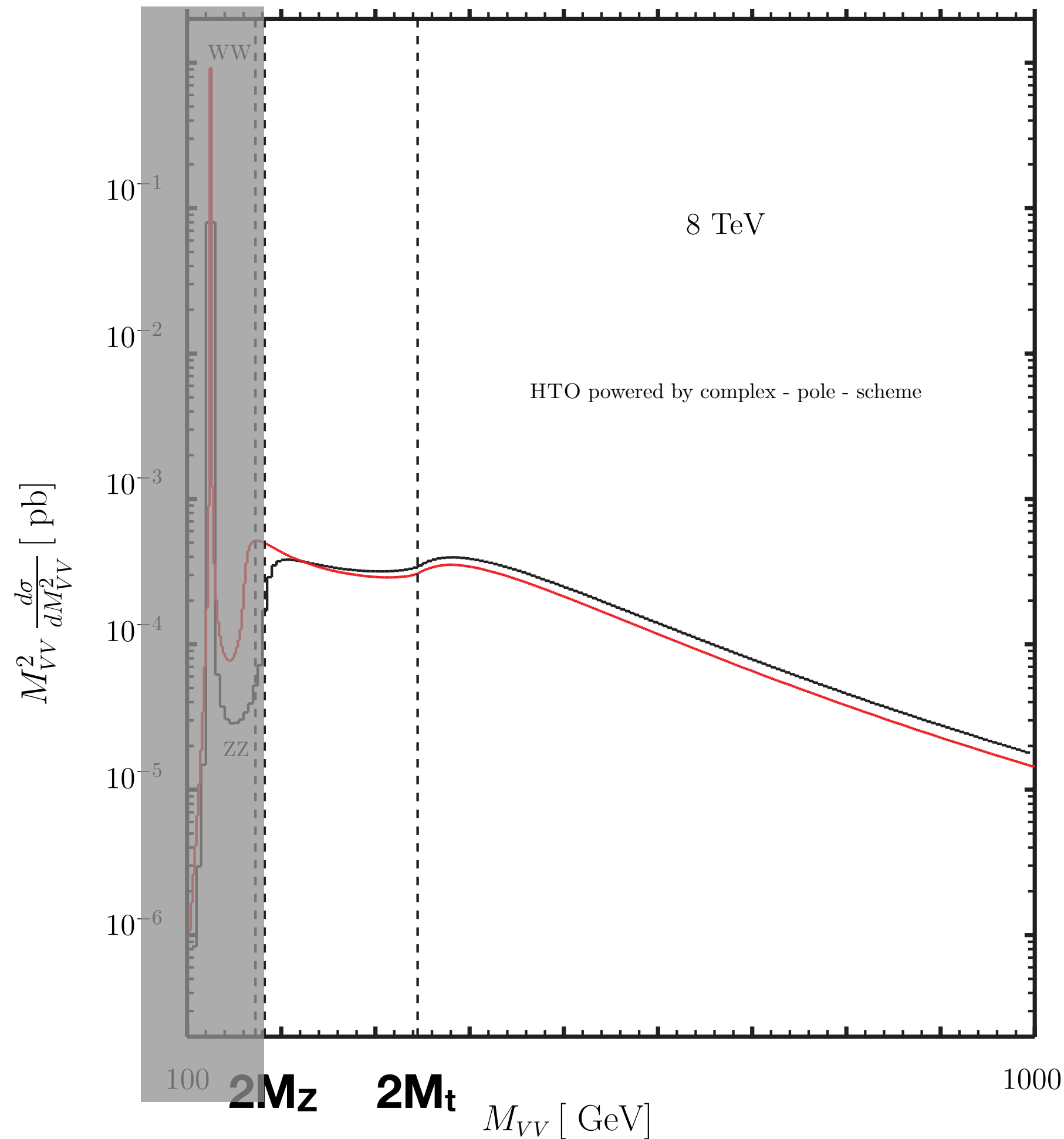


- Above $2m_Z$, the production cross-section is no longer dependent on the width
- Therefore Γ_H can be obtained from a ratio of on-shell and off-shell production

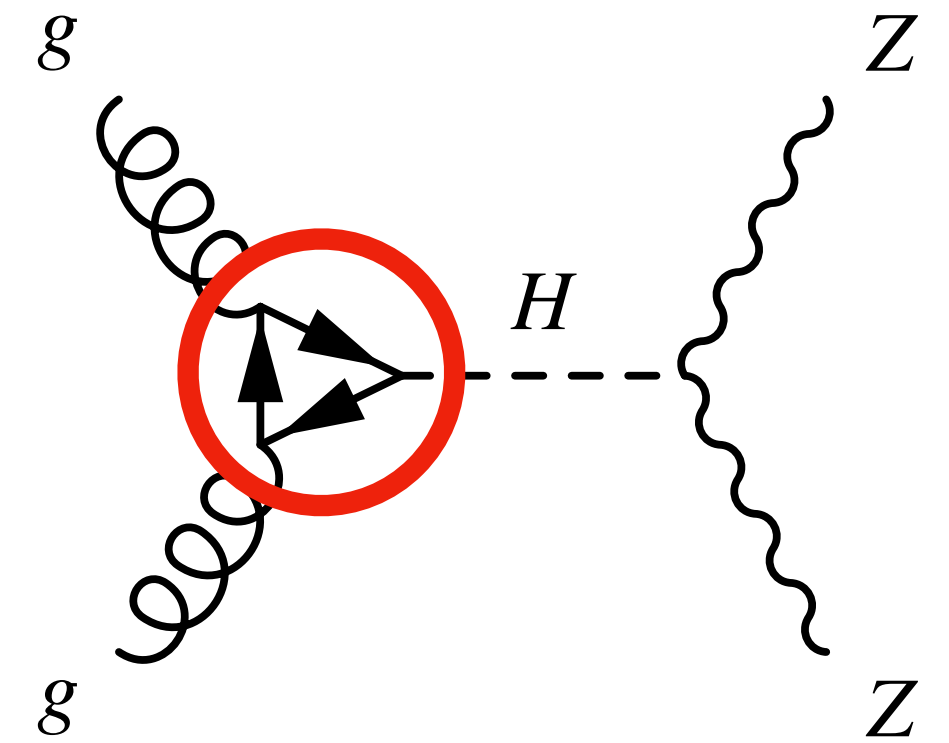
Off-shell Production in $H \rightarrow ZZ$

Kauer & Passarino, *JHEP* 2012, 116

$2 M_W$

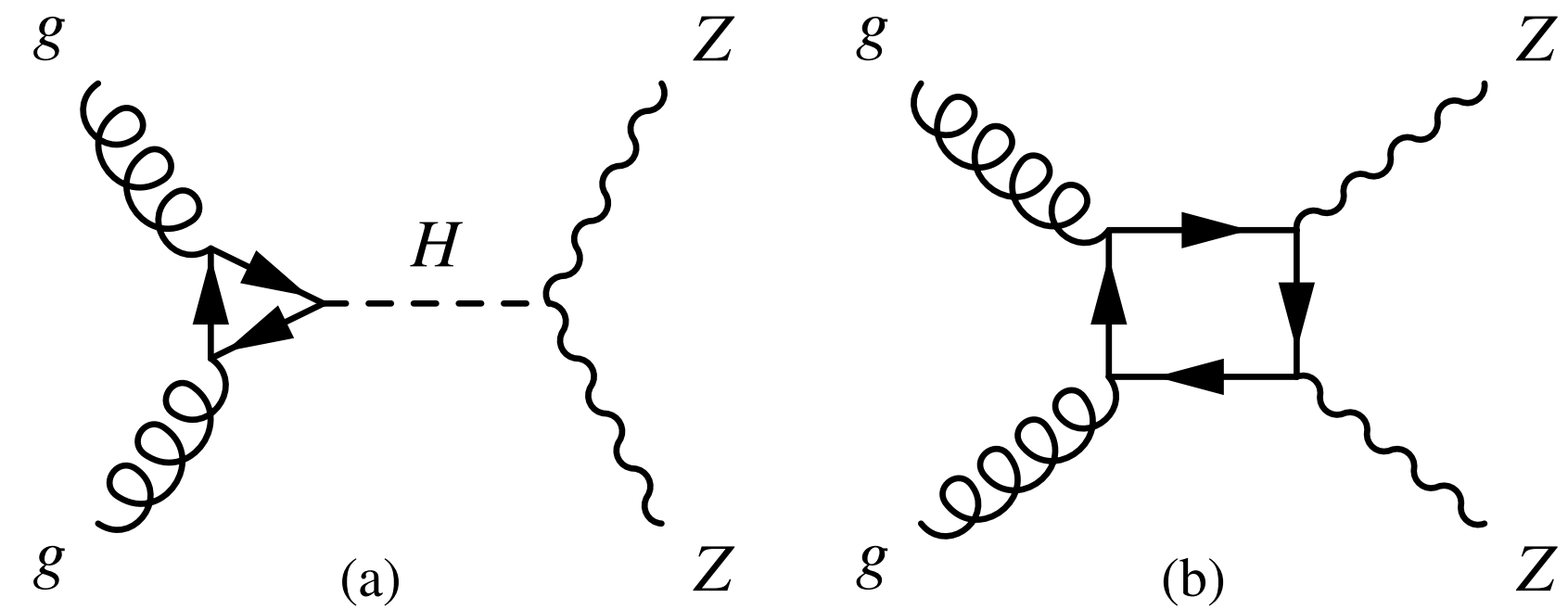
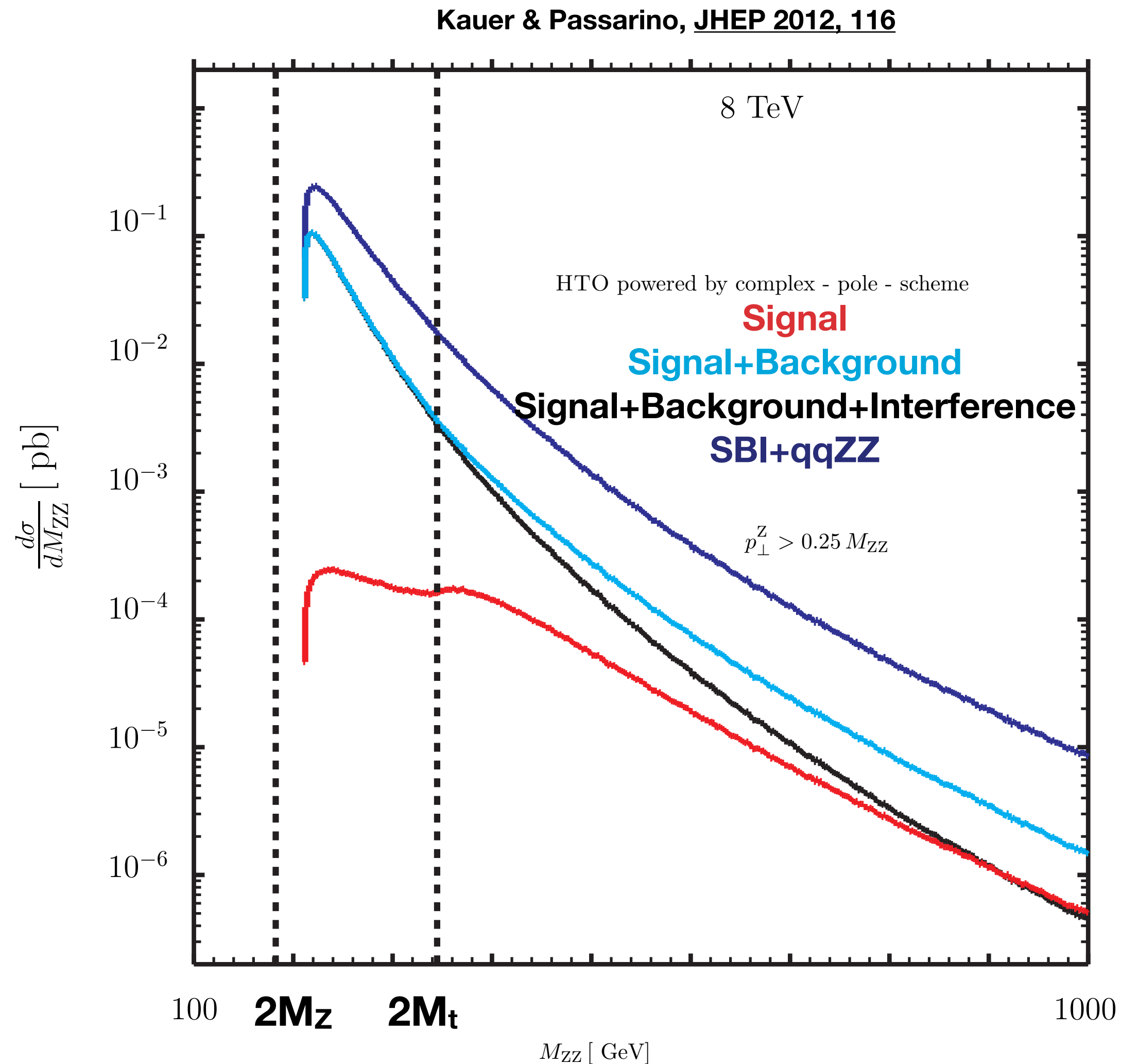


- Requires the assumption that **all couplings follow SM predictions**
- In particular, the effective coupling g_{ggH}
- \rightarrow **No new particles enter the quark loop, so on-shell and off-shell production are the same**



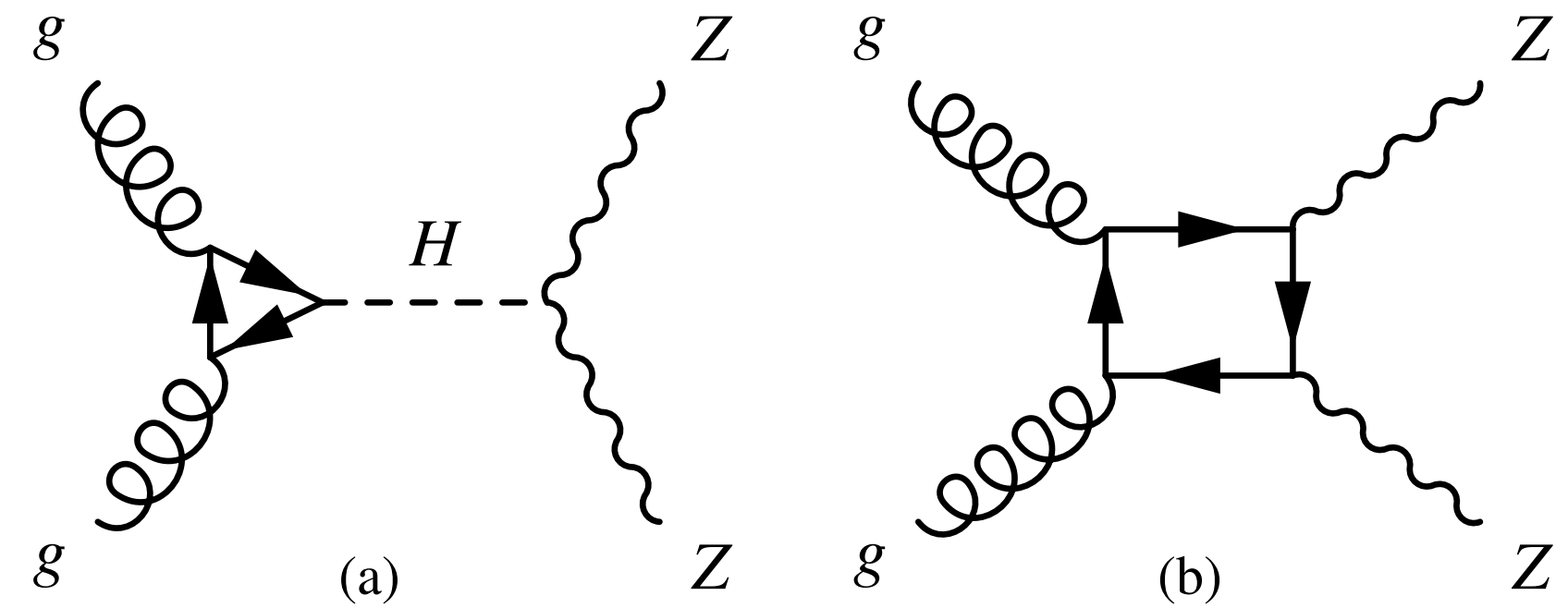
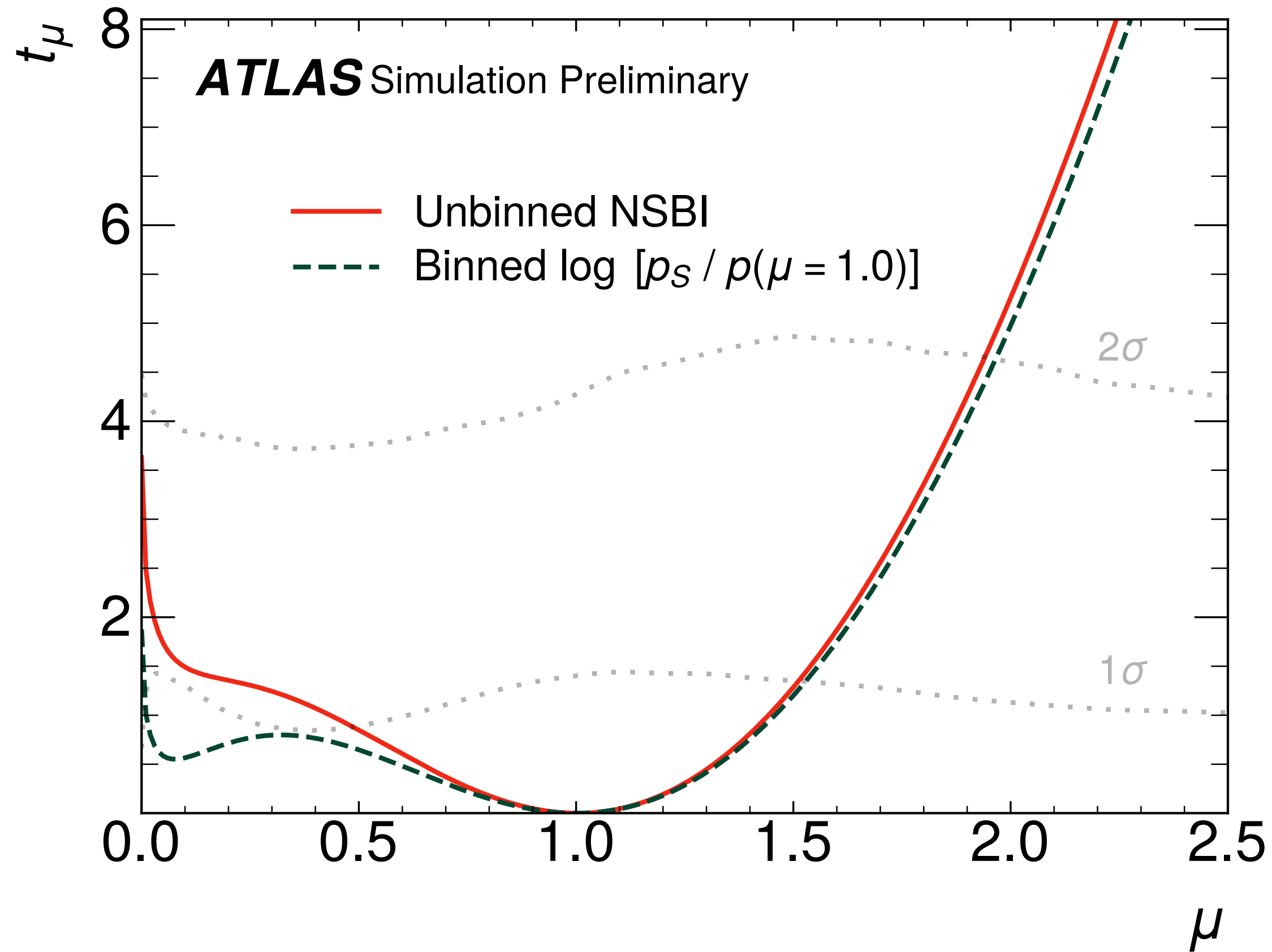
- Above $2m_Z$, the production cross-section is no longer dependent on the width
- Therefore Γ_H can be obtained from a ratio of on-shell and off-shell production

Signal-Background Interference in $H \rightarrow ZZ$



- The signal $gg \rightarrow H^* \rightarrow ZZ$ and background $gg \rightarrow ZZ$ process interfere
 - We measure a deficit wrt the background, not a signal
- Signal scales with μ_{offshell} but the interference goes as $\sqrt{\mu_{\text{offshell}}}$

Signal-Background Interference in $H \rightarrow ZZ$

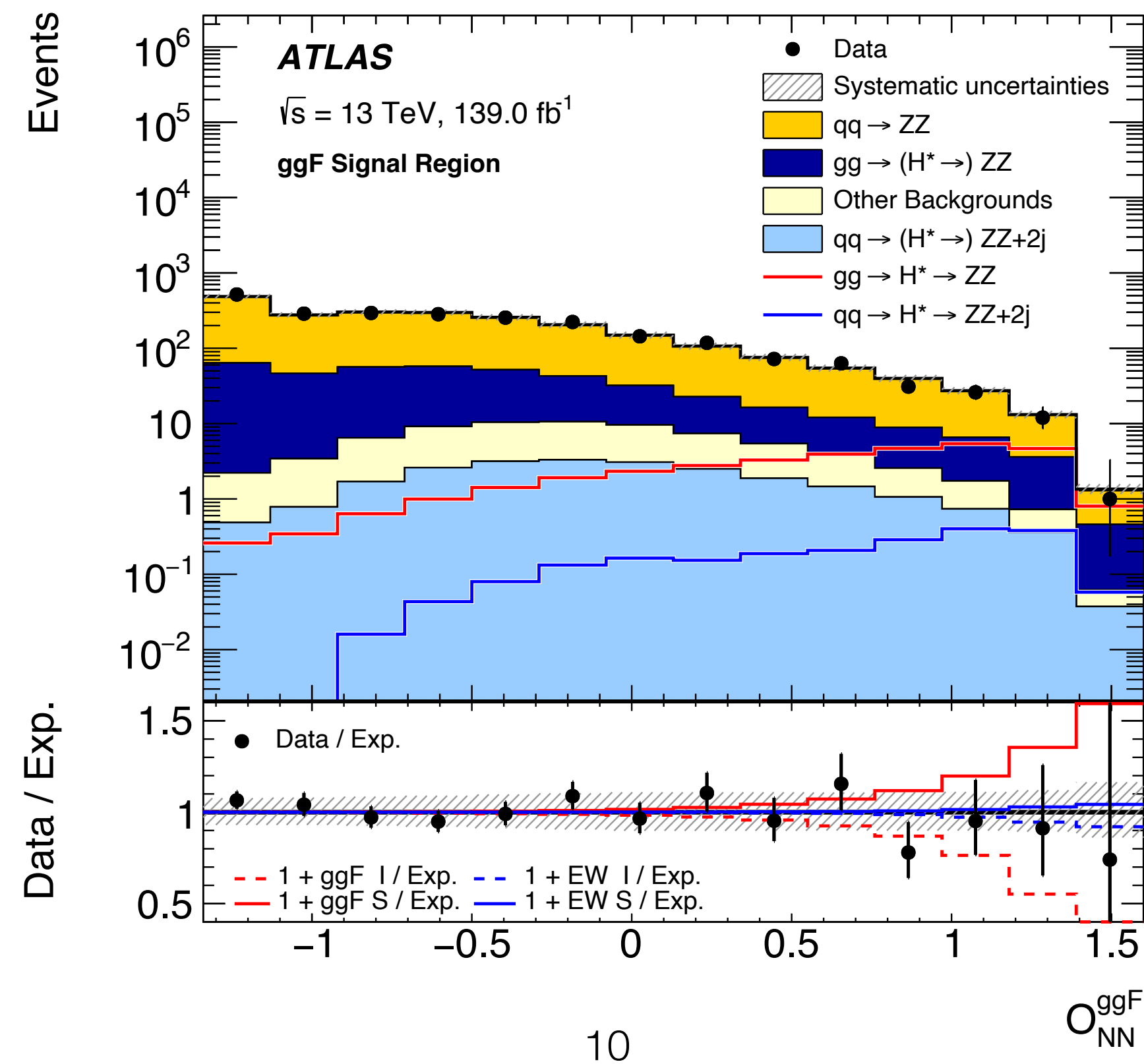


- $\sqrt{\mu_{\text{off-shell}}}$ dependence means that asymptotic approximation does not hold
 - Introduces double minimum
 - Requires cutoff at $\mu_{\text{off-shell}}=0$
 - \rightarrow Confidence intervals have to be derived using the Neyman construction

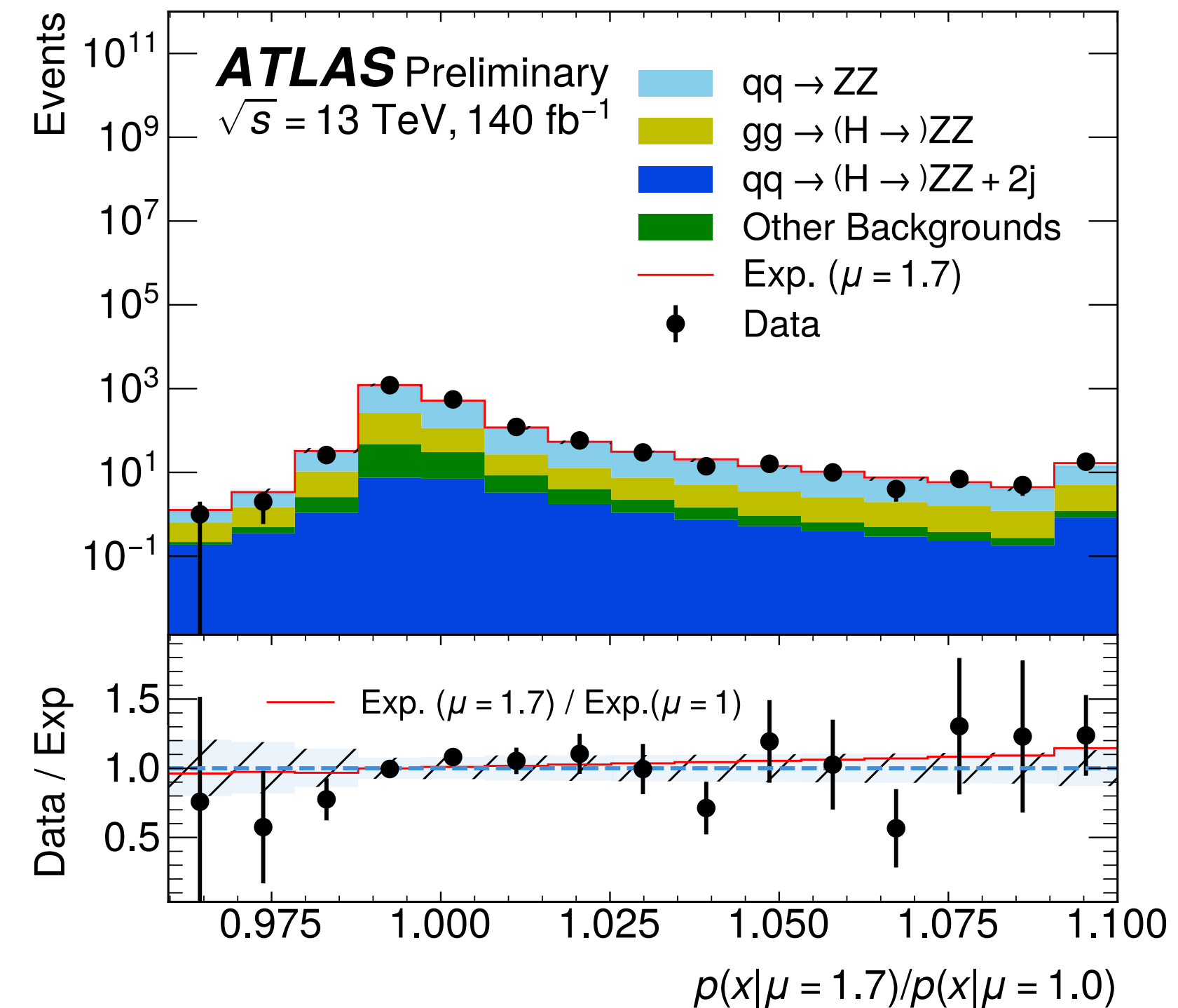
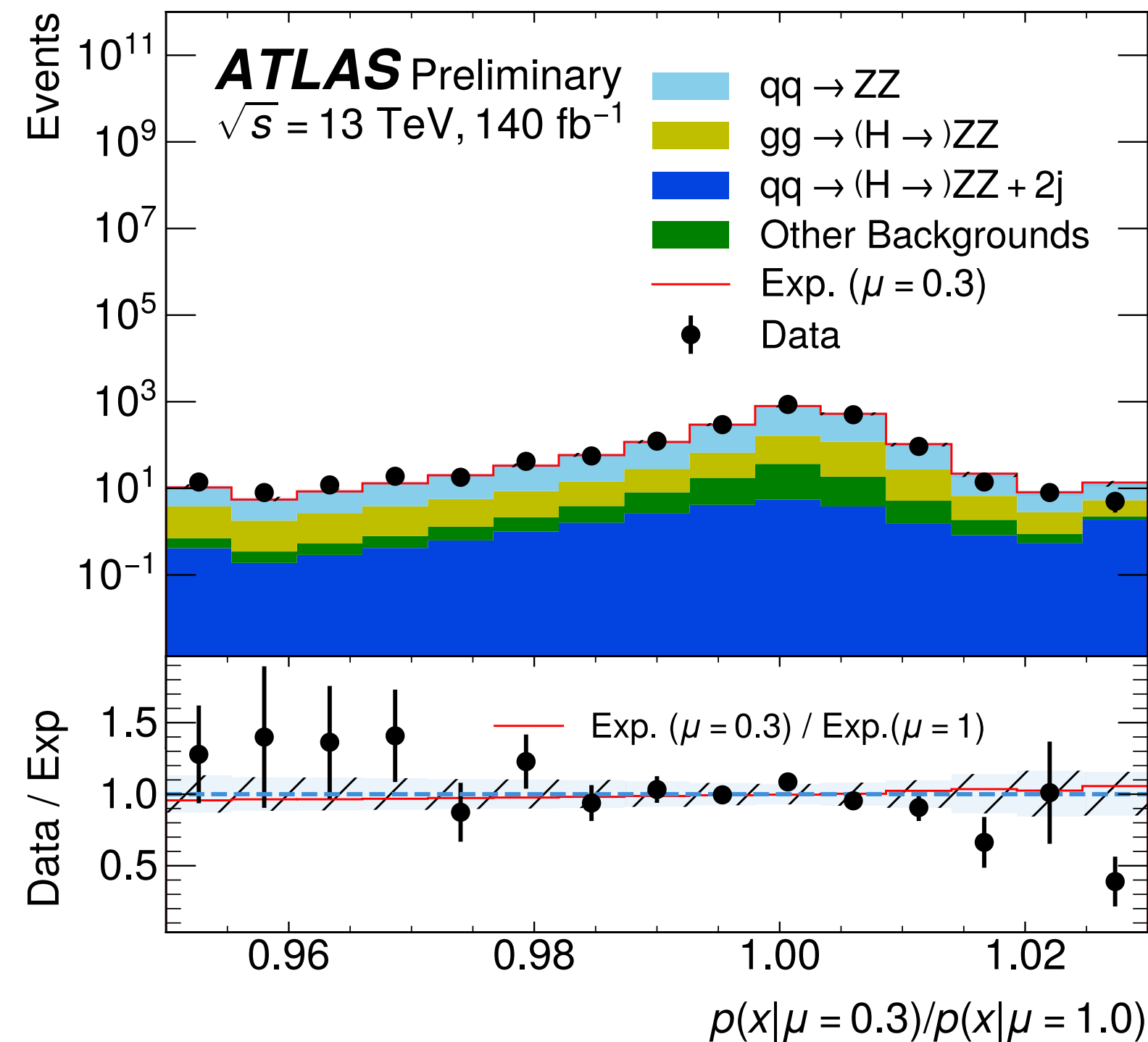
See talk by J. Sandesara

Neural Simulation Based Inference

- Standard approach uses histograms of kinematic observables to approximate density ratio
 - As done in the previous ATLAS result, [Phys. Lett. B 846 \(2023\) 138223](#)
 - NN trained to distinguish between signal and background

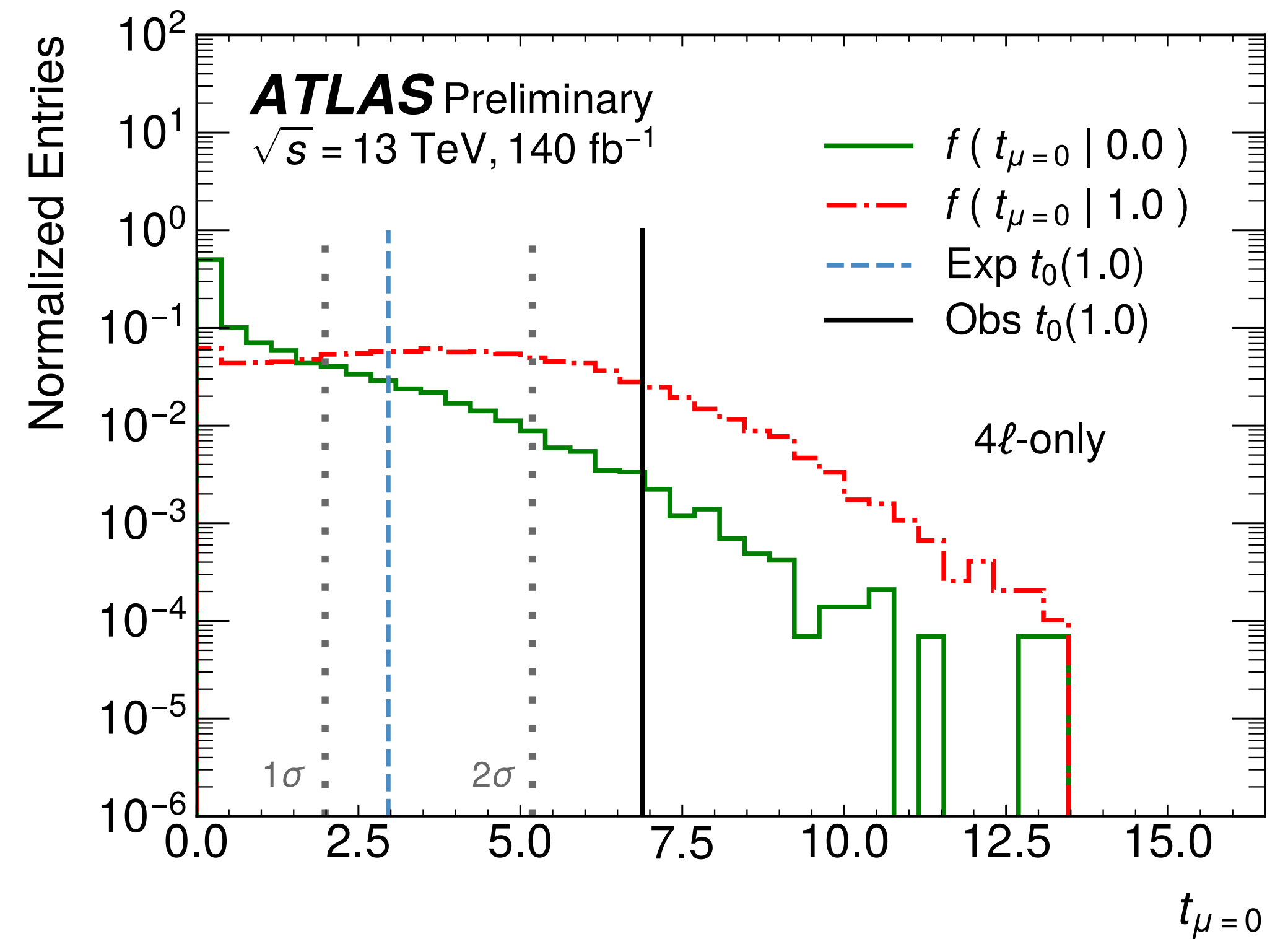
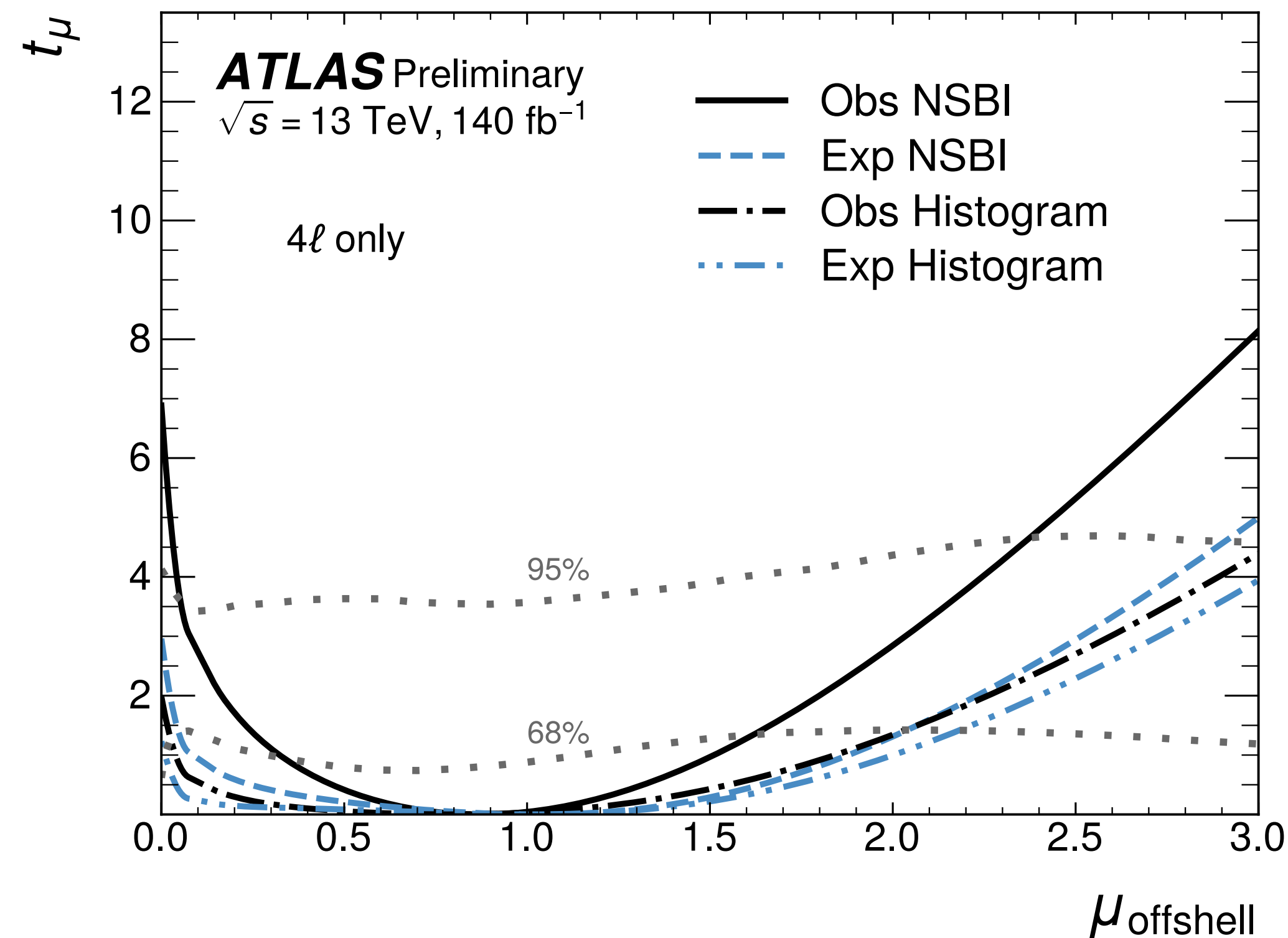


Neural Simulation Based-Inference



- Event-by-event approximation of probability density ratios using NSBI offers potential for improvement
 - Allows building an optimal observable for any value of $\mu_{\text{off-shell}}$

Results in the 4l Final State



- Improvement compared to histogram analysis
- Observe off-shell Higgs boson production with significance 2.3σ using only the $ZZ \rightarrow 4l$ channel

2l2v Final State

- The analysis in this final state is not changed, it remains a histogram analysis

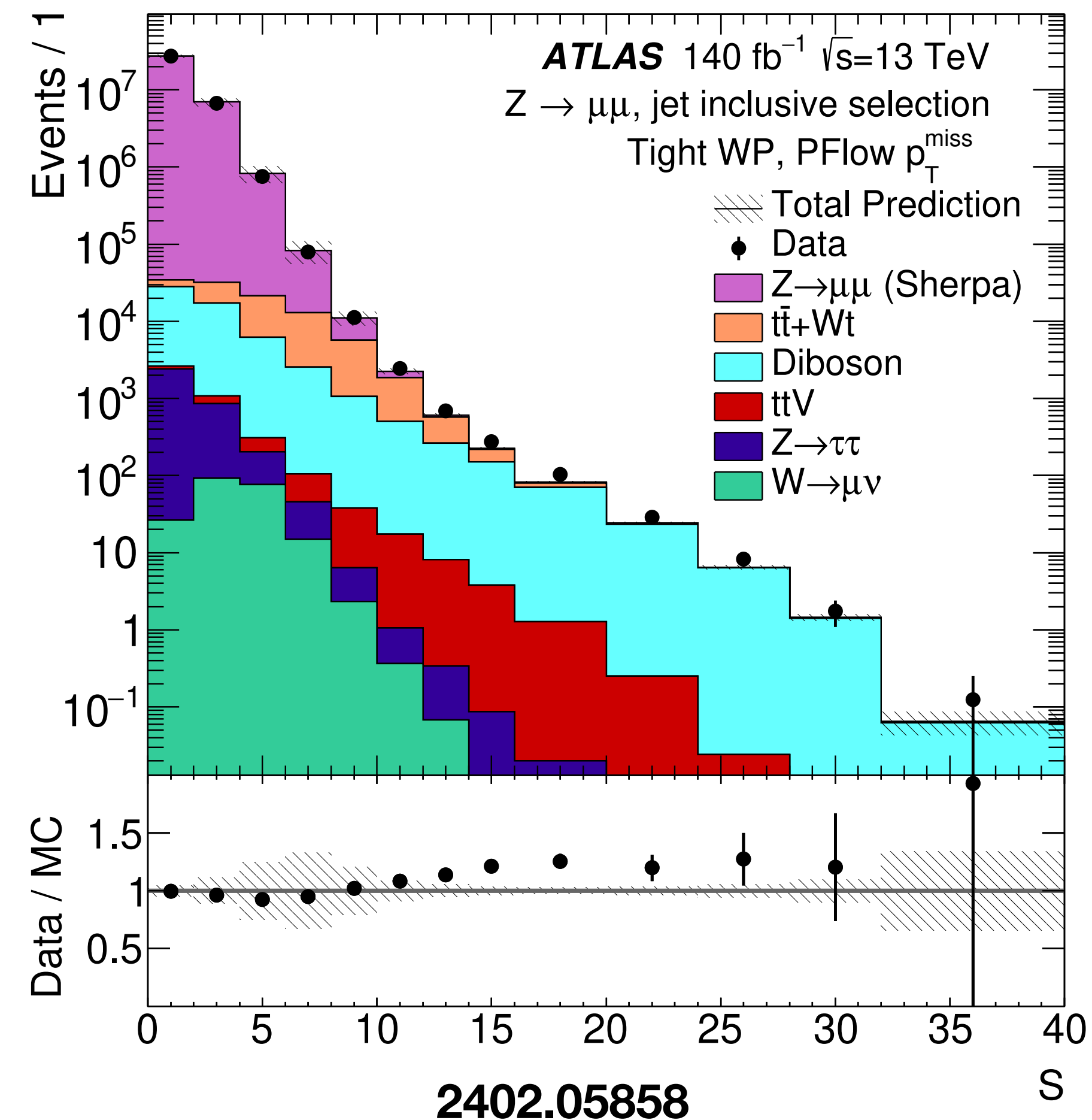
- Observable is transverse mass m_T^{ZZ}

$$m_T^{ZZ} \equiv \sqrt{\left[\sqrt{m_Z^2 + (p_T^{\ell\ell})^2} + \sqrt{m_Z^2 + (E_T^{\text{miss}})^2} \right]^2 - \left| \vec{p}_T^{\ell\ell} + \vec{E}_T^{\text{miss}} \right|^2}$$

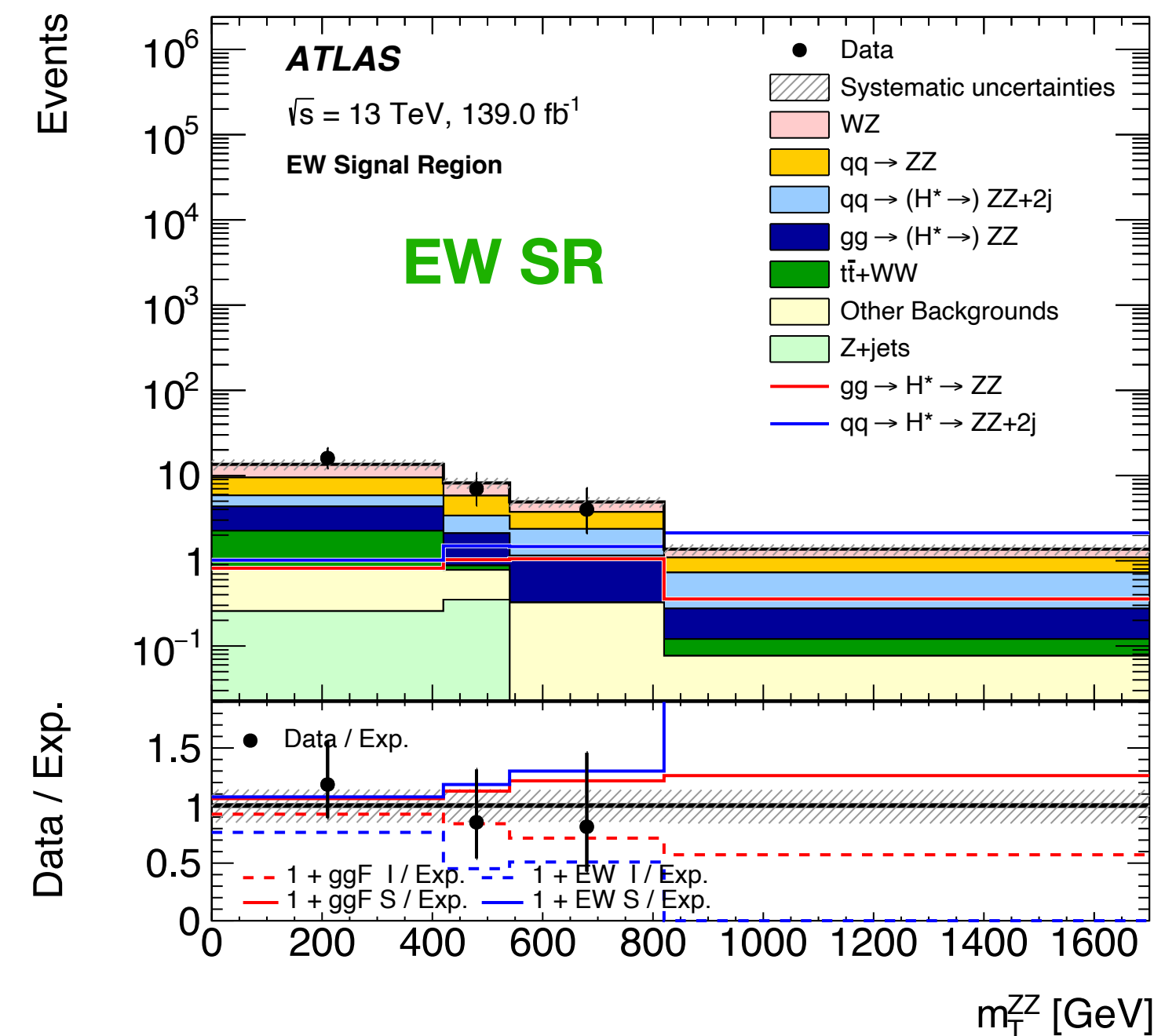
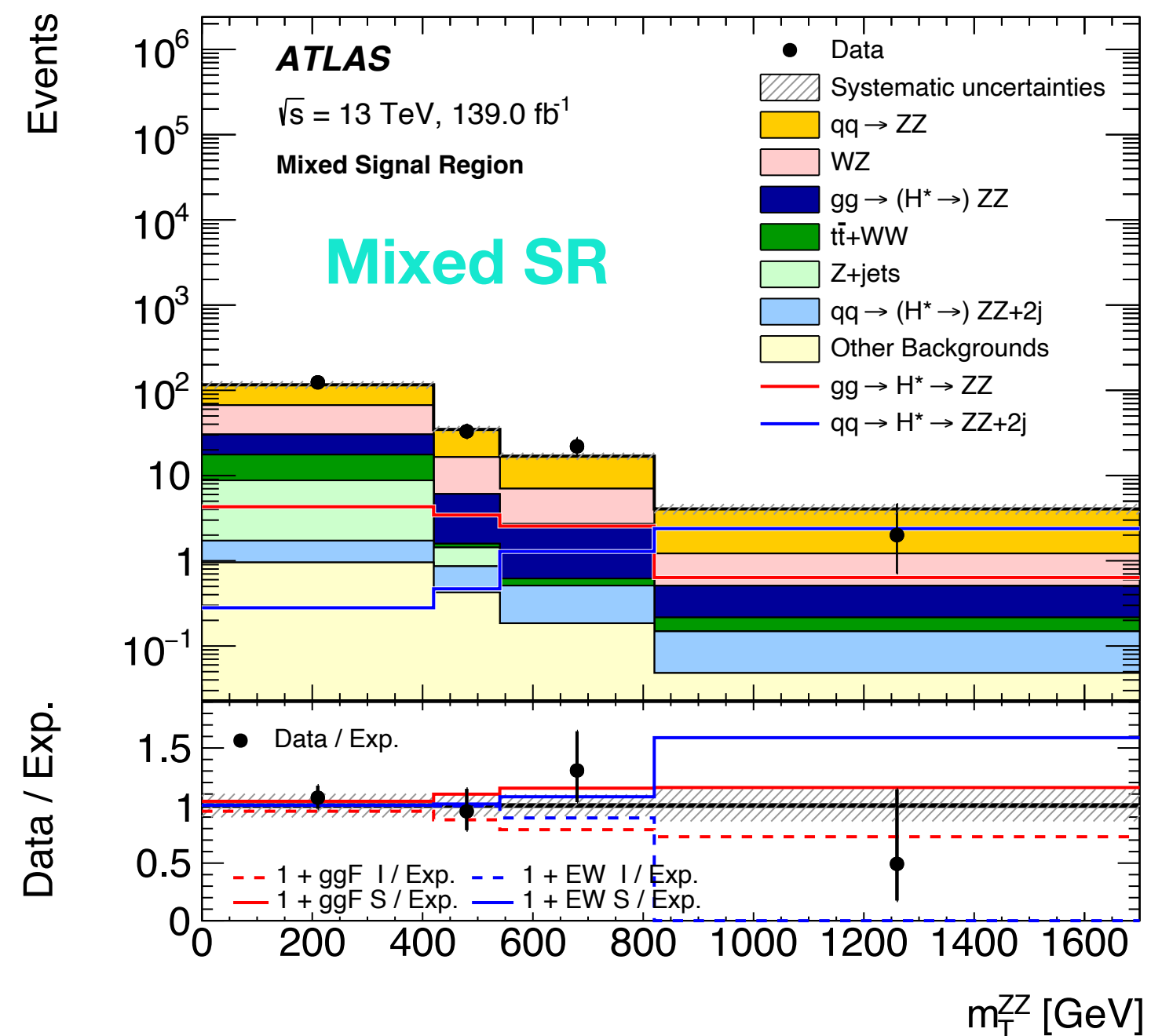
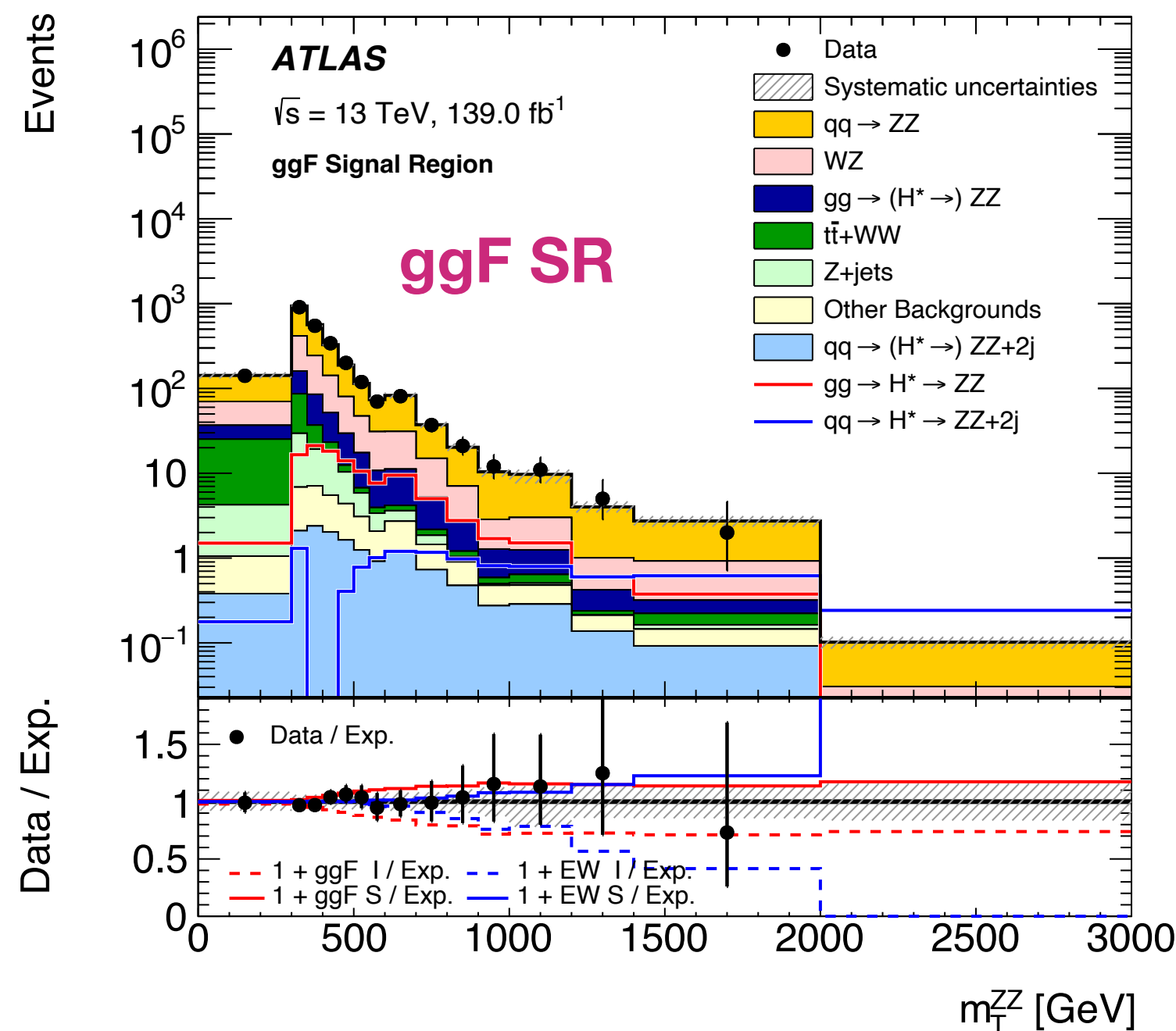
- The off-shell signal is more significant at higher m , m_T^{ZZ} is a useful proxy

- Much larger background than 4l, mainly from Z+jets events

- These are suppressed by cutting on the E_T^{miss} significance

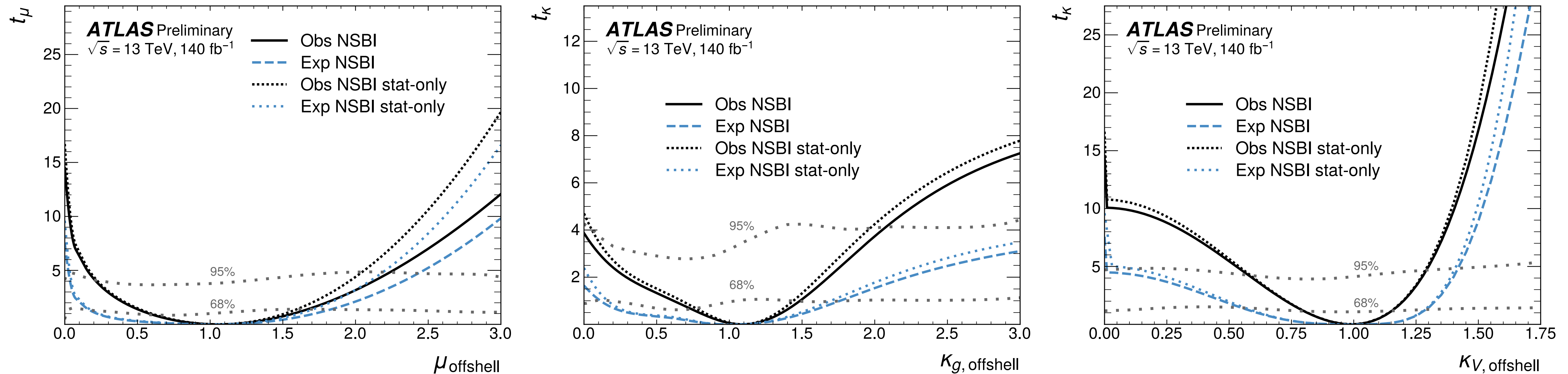


2l2v Final State



- SR further divided by N_{jets} , 2-jet SR targets **EW (VBF+VH)** production, others **ggF**
- **qqZZ** production is the dominant background
 - Constrained using 4l CRs, 1 per jet bin
 - Separate CRs constructed to constrain **WZ**, **Z+jets**, and **non-resonant** backgrounds

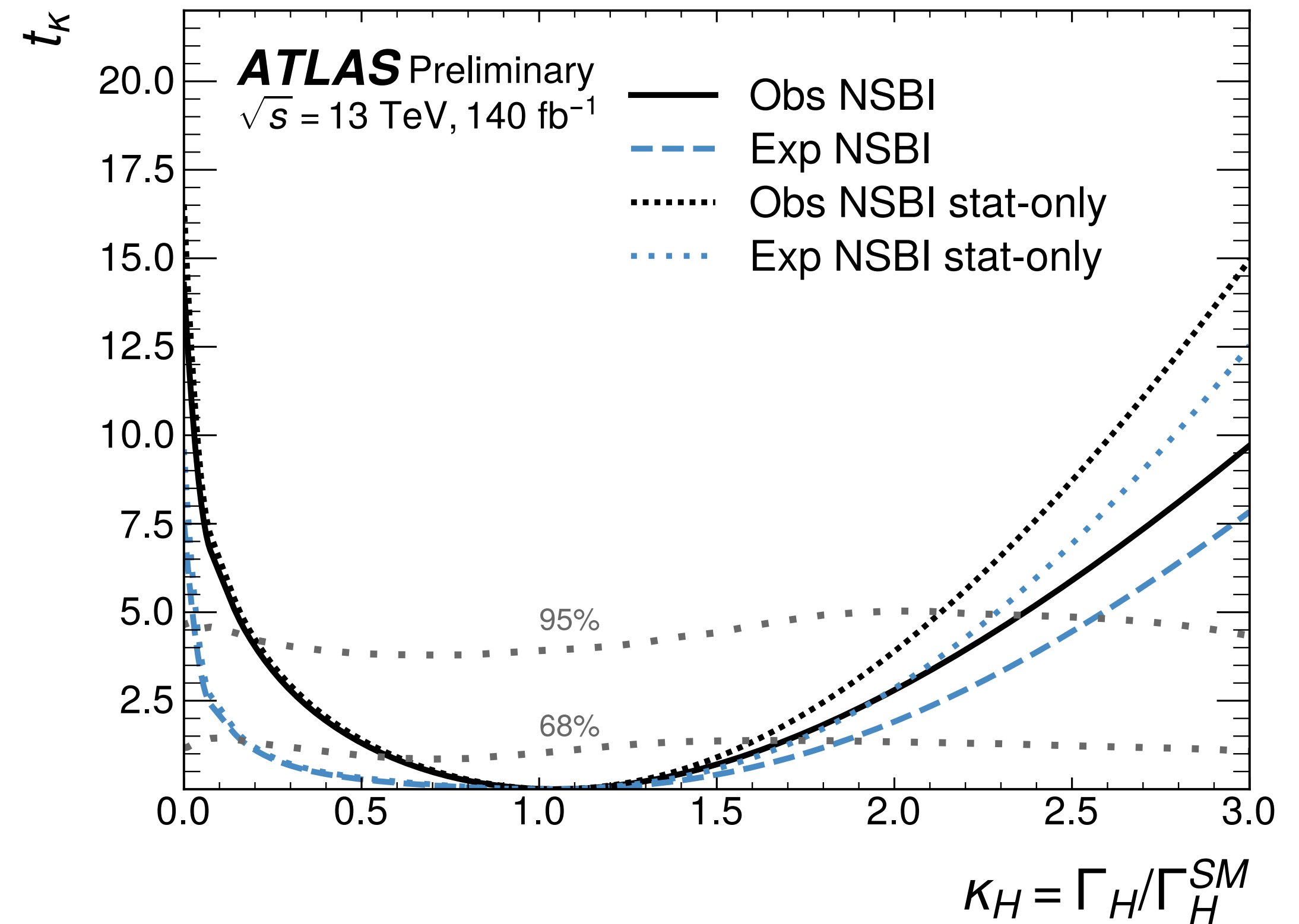
Combined HZZ Off-Shell Analysis



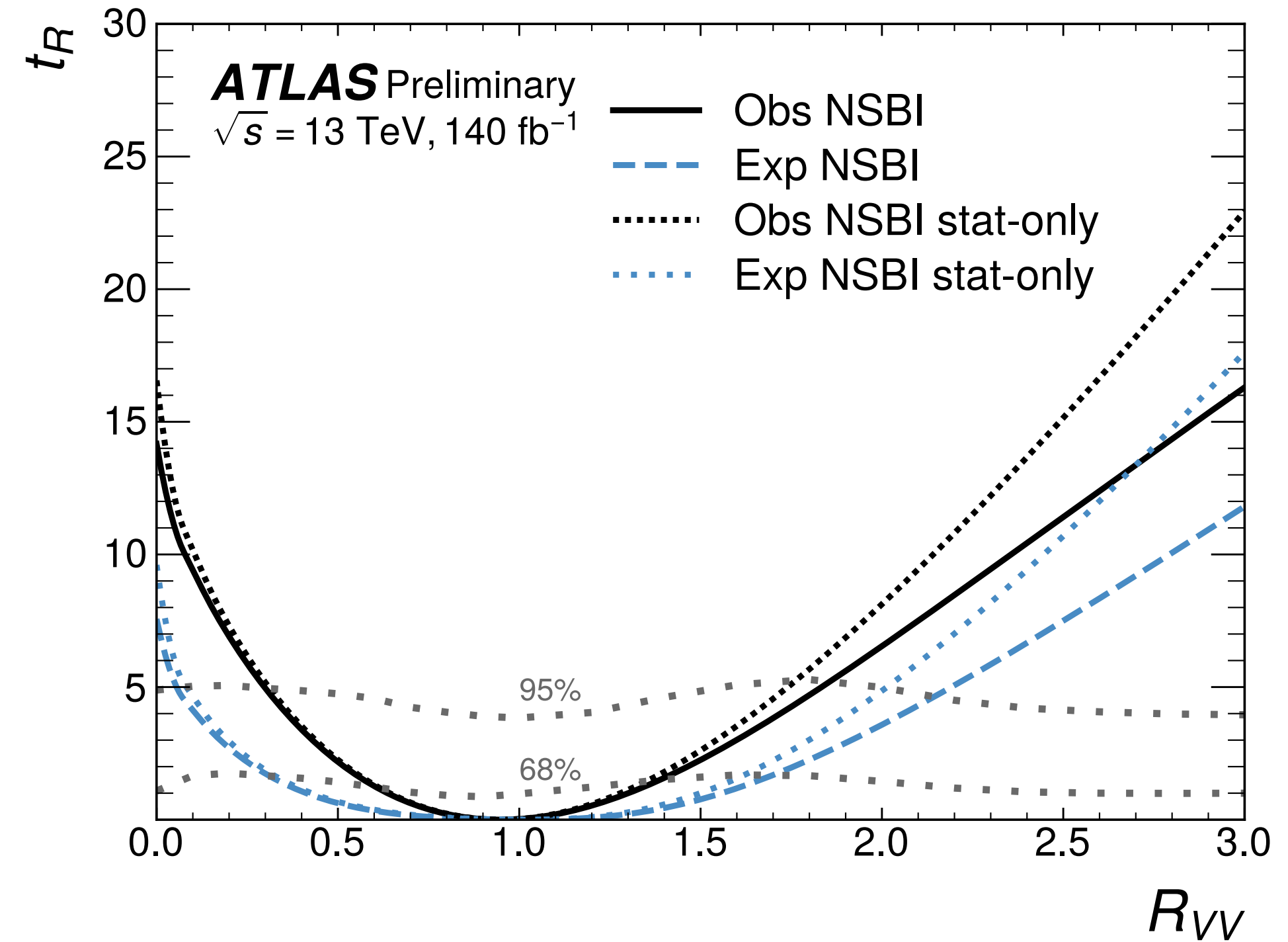
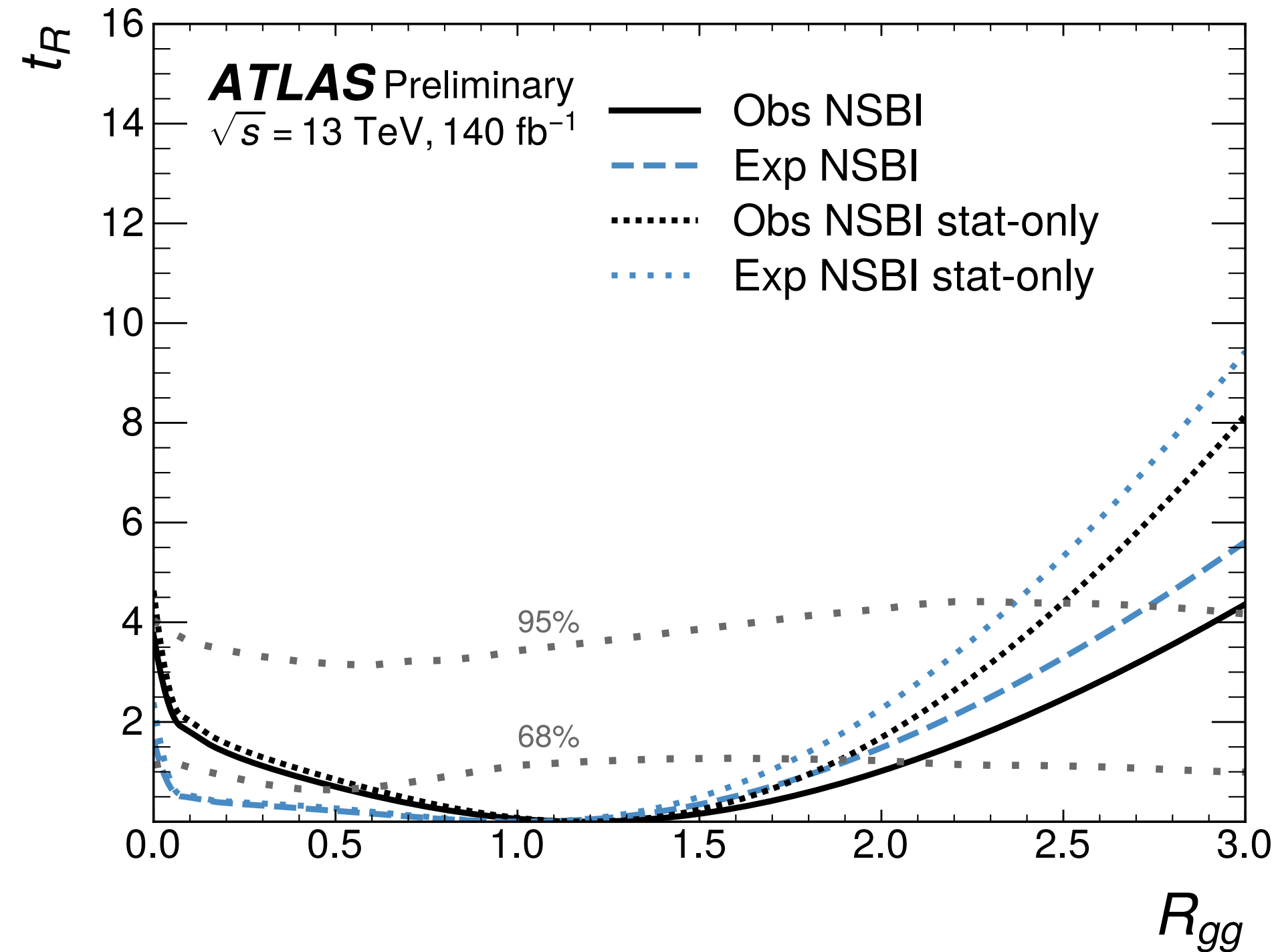
- Evidence for off-shell Higgs boson production with 3.7σ significance
- Can also constrain Higgs-gluon and Higgs-vector boson couplings
 - Without any assumption on the total Higgs width
 - Precision on $\kappa_{V,\text{off-shell}} = 0.99^{+0.016}_{-0.19}$ approaches on-shell measurement of 1.02 ± 0.06

Width Interpretation

- Combine with previous HZZ on-shell analysis to constrain Γ_H
- $\Gamma_H = 4.3^{+2.7}_{-1.9}$ MeV, with 95% CL limits of $[0.8, 9.8]$ MeV
 - Compare to previous result of $4.5^{+3.0}_{-2.5}$ and $[0.1, 10.2]$ MeV
 - The $2l2\nu$ channel provides \sim half of the sensitivity here



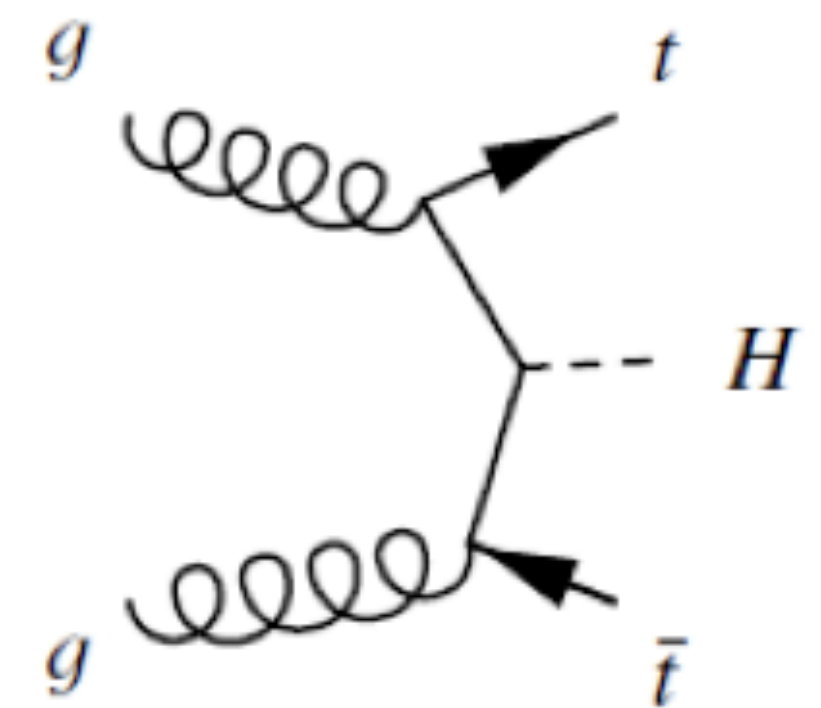
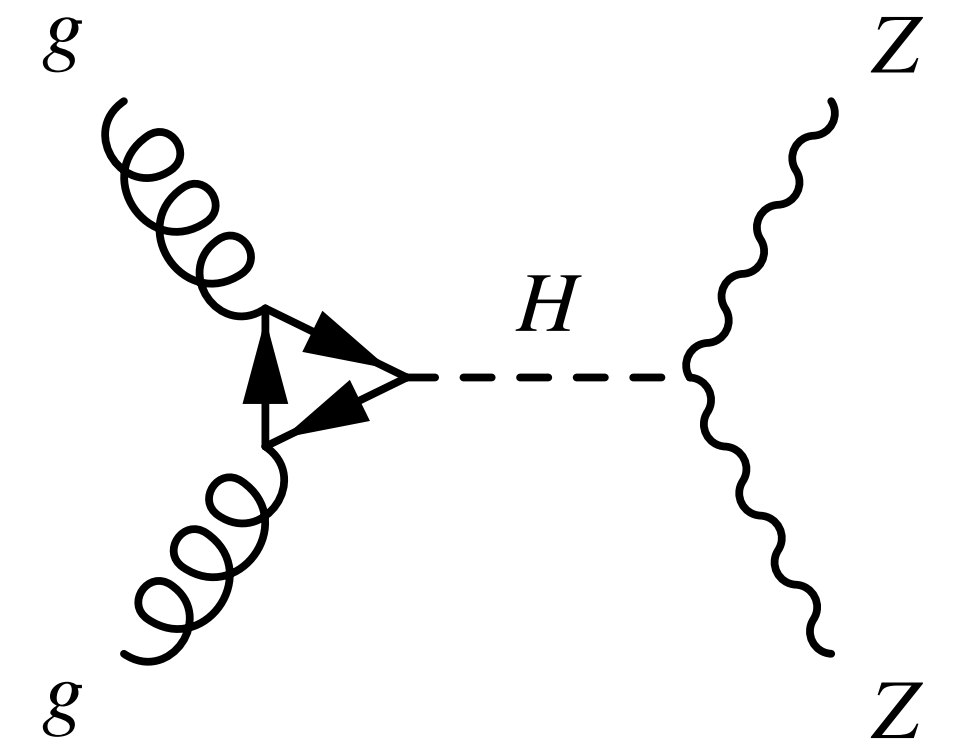
Couplings Interpretation



- Can also use the combination to constrain the ratio of on-shell and off-shell couplings
 - Separately for EW and ggF; $R_{VV}=1$ for R_{gg} determination and vice versa
- Assuming the SM width, verify that on-shell and off-shell couplings to vector bosons/ gluons agree: $R_{VV}=0.95^{+0.44}_{-0.34}$, $R_{gg}=1.19^{+0.89}_{-0.66}$

A Complementary Approach

- The analysis in the HZZ channel requires assuming no changes to the ggH or HVV couplings
 - Especially problematic for ggH production which requires assuming no new particles that could enter into the quark loop
- Important to try to constrain the width using other processes to avoid relying on the same assumptions
- ttH production is the largest cross-section Higgs production method that does not involve g or V couplings
 - Measurements of off-shell ttH could then be used to constrain Γ_H without overlapping assumptions
 - Use 4t production, recently observed by ATLAS, [Eur. Phys. J C 83 \(2023\) 496](#)



A Complementary Approach

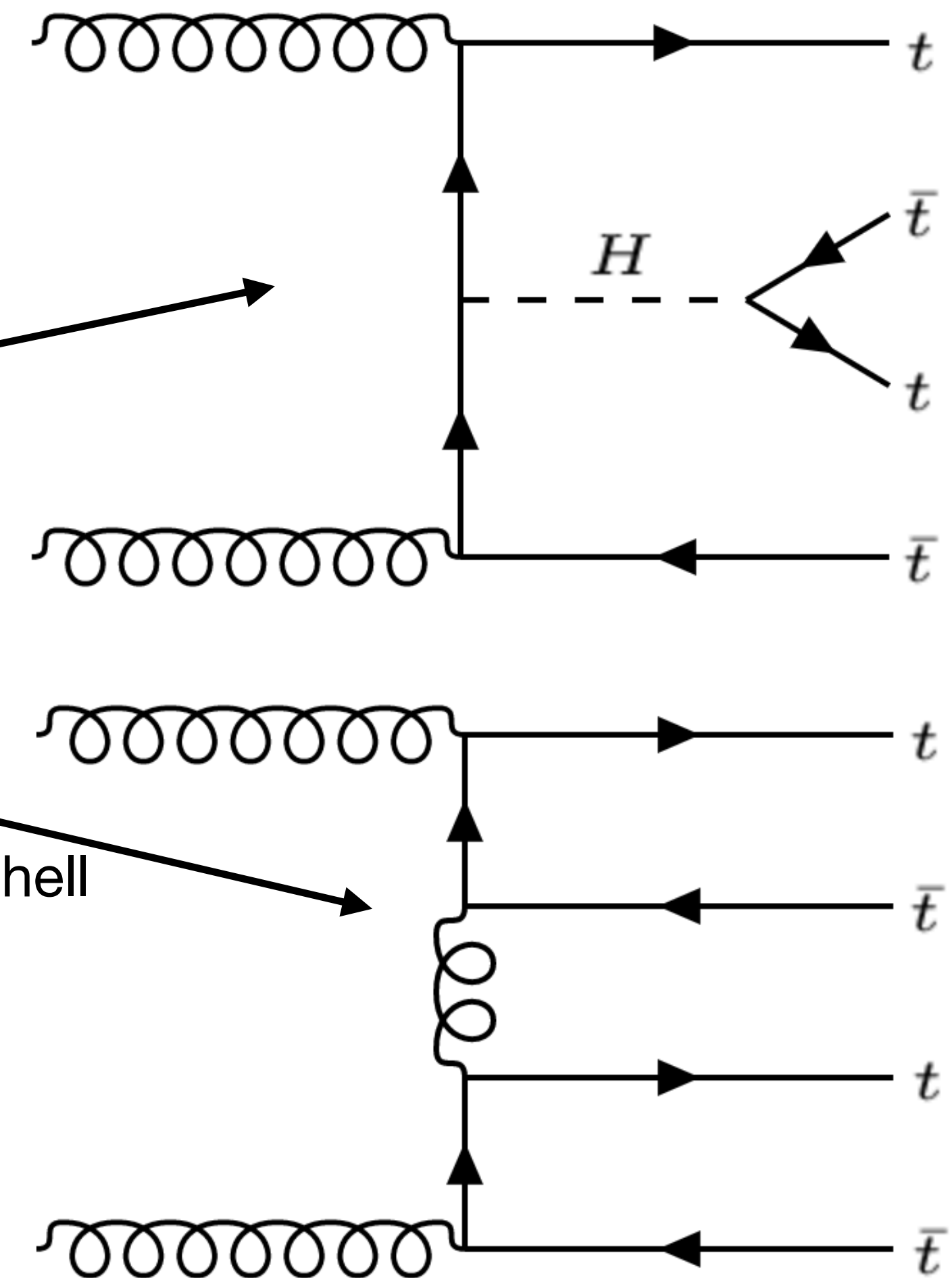
- 4top production is a combined signal + background + interference process, analogous to the combined HZZ off-shell process

- The mediator in 4t production can be a Higgs: signal
- Or $g/Z/\gamma^*$: background
- And the two interfere, of course

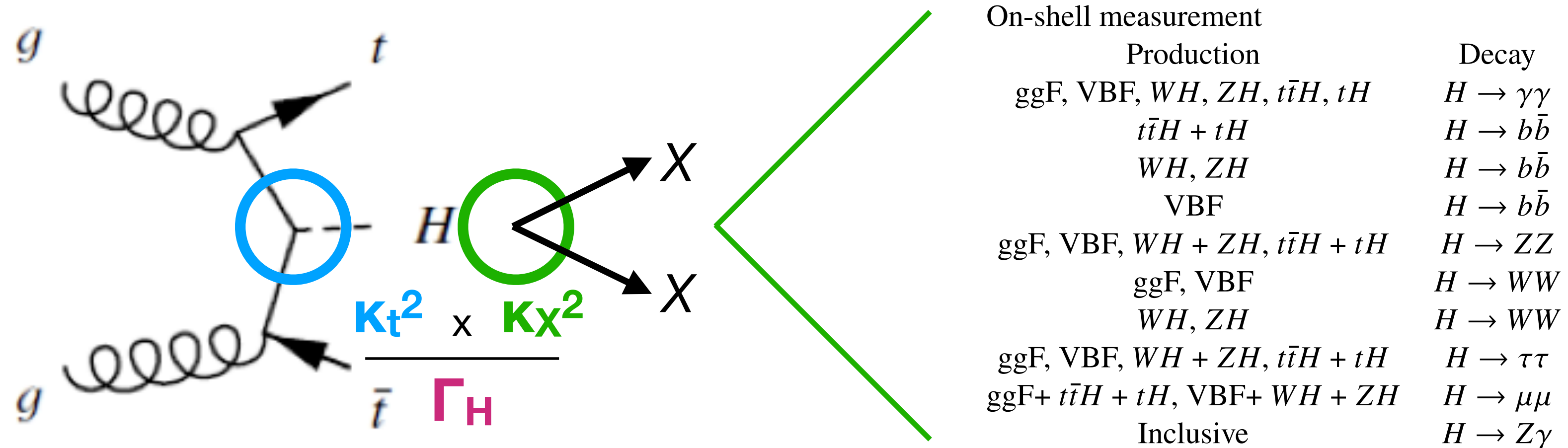
- Quartic dependence on κ_t instead of quadratic on $\mu_{\text{off-shell}}$

- Signal $\sim \kappa_t^4$, interference $\sim \kappa_t^2$, background is constant

- 2L and 3L final states used for the measurement



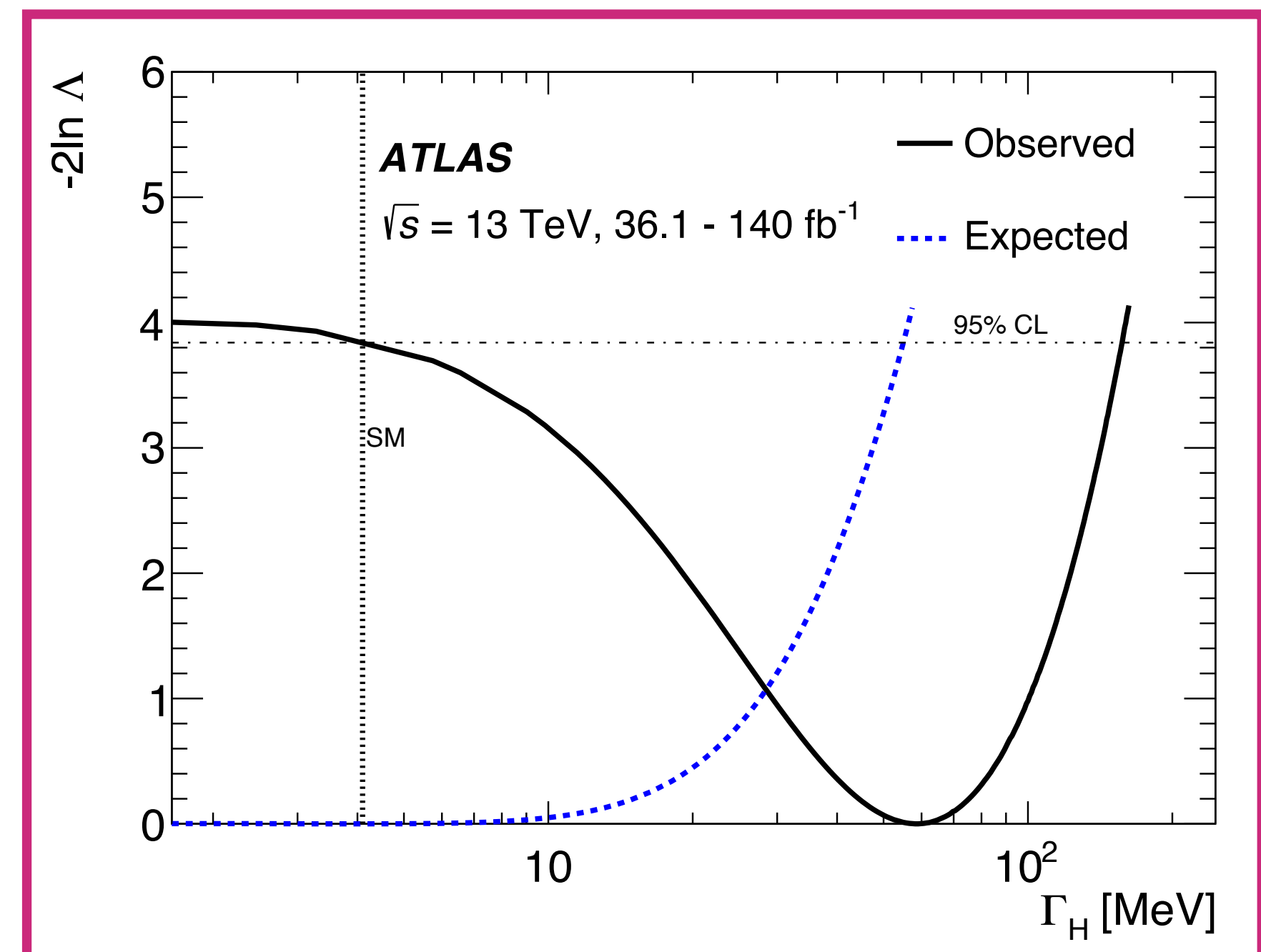
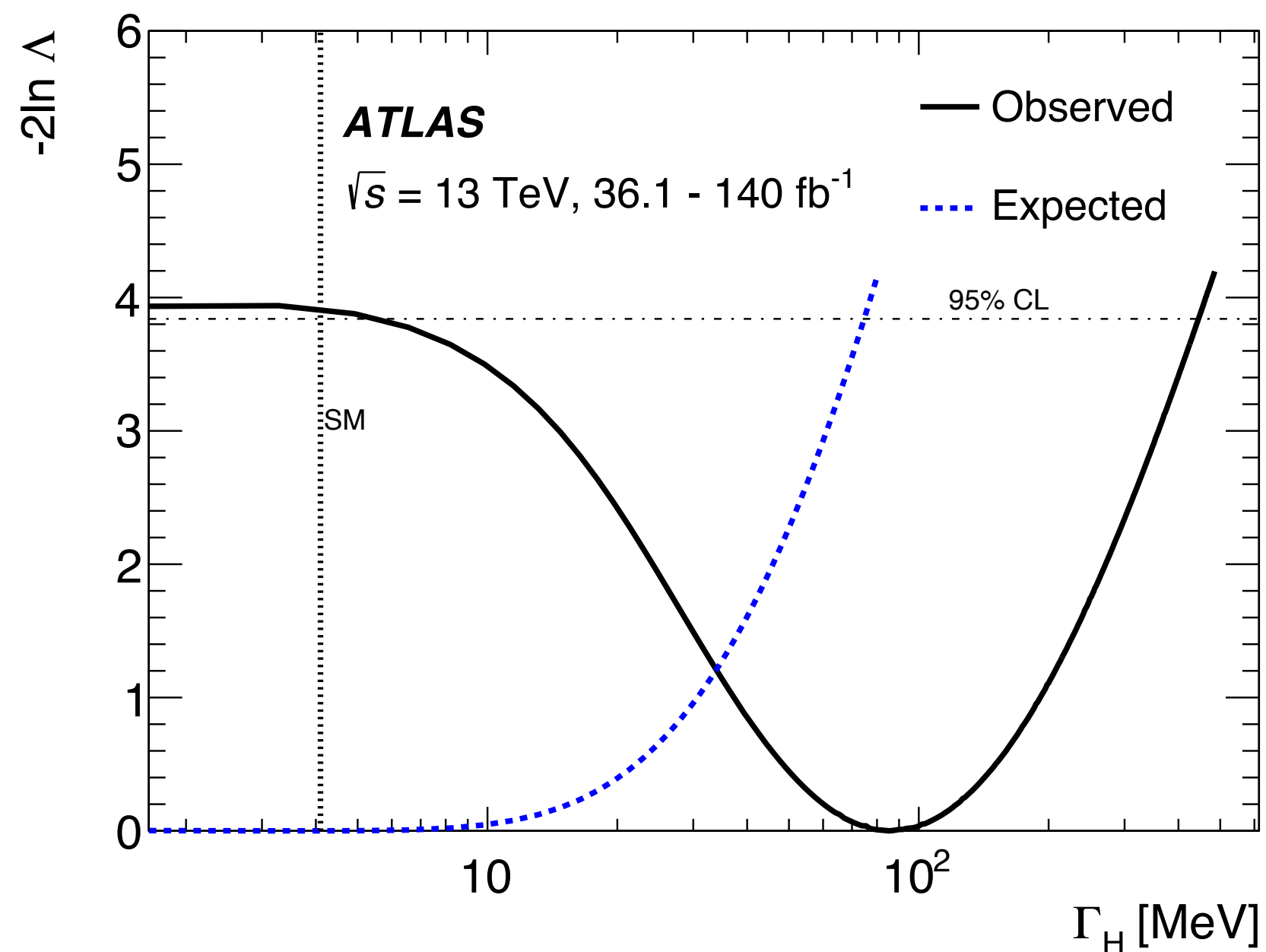
On-Shell Analysis



- Only κ_t is present in the off-shell production, but **other couplings** are involved on-shell, depending on the decay
- Use other on-shell Higgs measurements to constrain them
 - Strength of constraint depends on whether ggH is treated as an effective coupling to gluons or as a quark loop
- Analyses not designed to avoid overlap, so some off-shell contamination in on-shell and vice versa
 - Drop ttH multilepton final state to try to minimize this

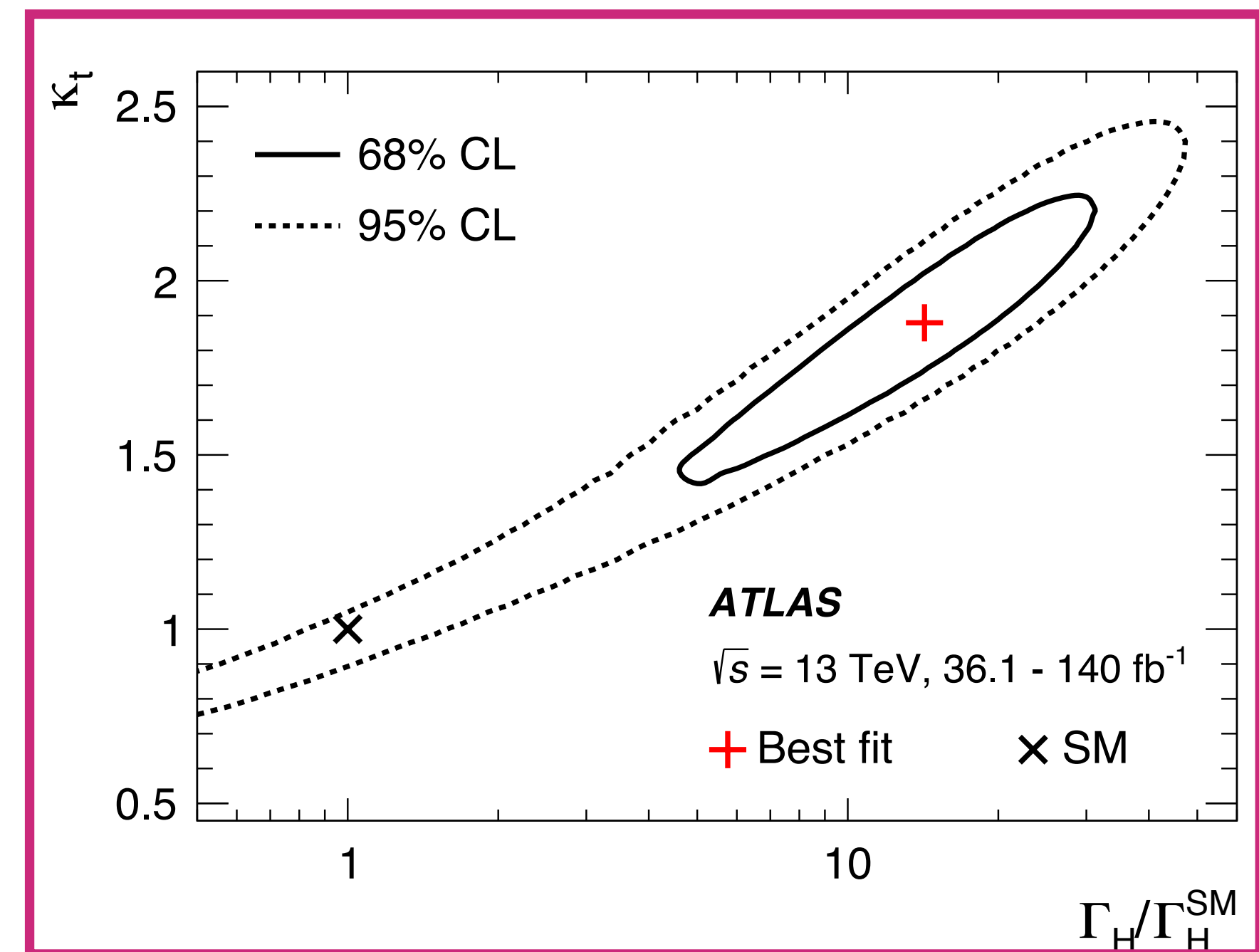
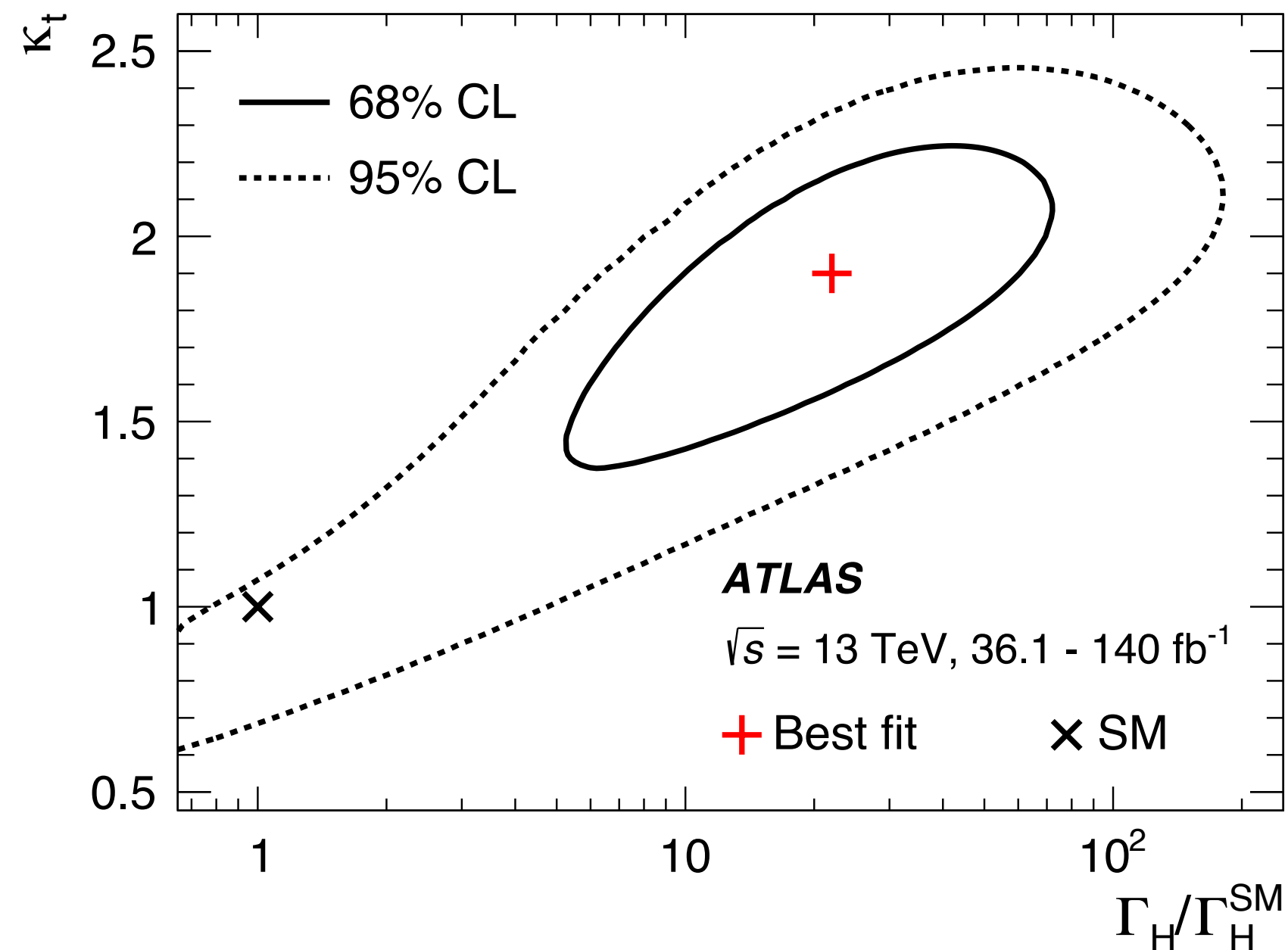
Width From ttH

- Observed 95% CL upper limit for Γ_H is 450 MeV, with $\Gamma_H = 86^{+110}_{-49}$ MeV
 - 2σ tension with SM due to measured $4t$ cross-section being larger than SM prediction
 - ▶ $\sigma_{4t} = 22.5^{+6.6}_{-5.5}$ pb, compared to SM 12 pb
 - Stronger limits from parameterizing κ_g as a function of κ_t
 - ▶ 95% CL upper limit of 160 MeV in this case



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Summary and Outlook

- Direct measurement of Γ_H is very difficult but indirect measurements using off-shell Higgs boson production have been quite successful
- The $H \rightarrow ZZ$ channel is good for this measurement due to the relatively large off-shell cross-section
- New statistical techniques allow us to make the best possible use of our limited dataset
- Using $4t$ provides a complementary measurement reliant on different assumptions
- Stay tuned, more results are coming!

Backup

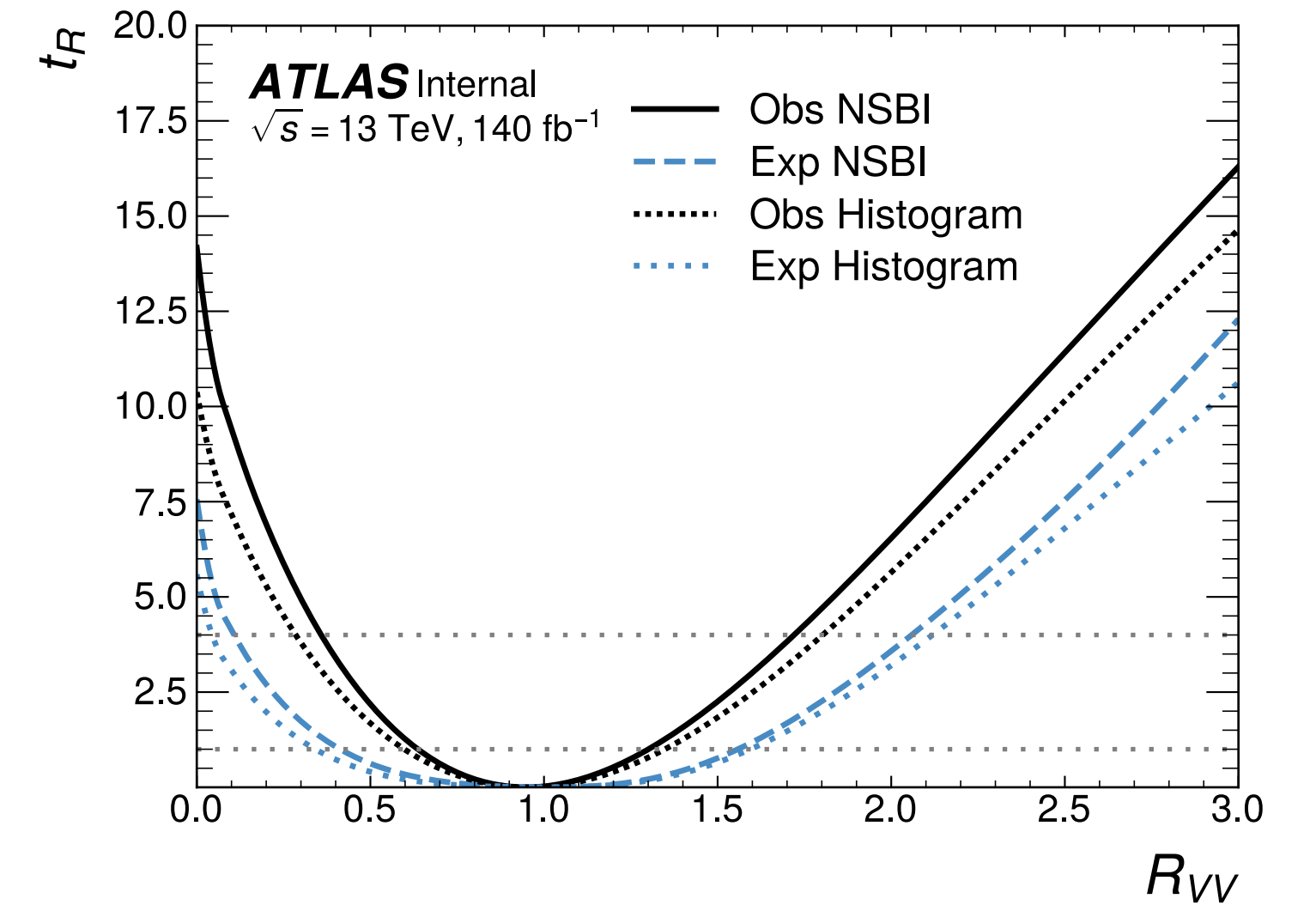
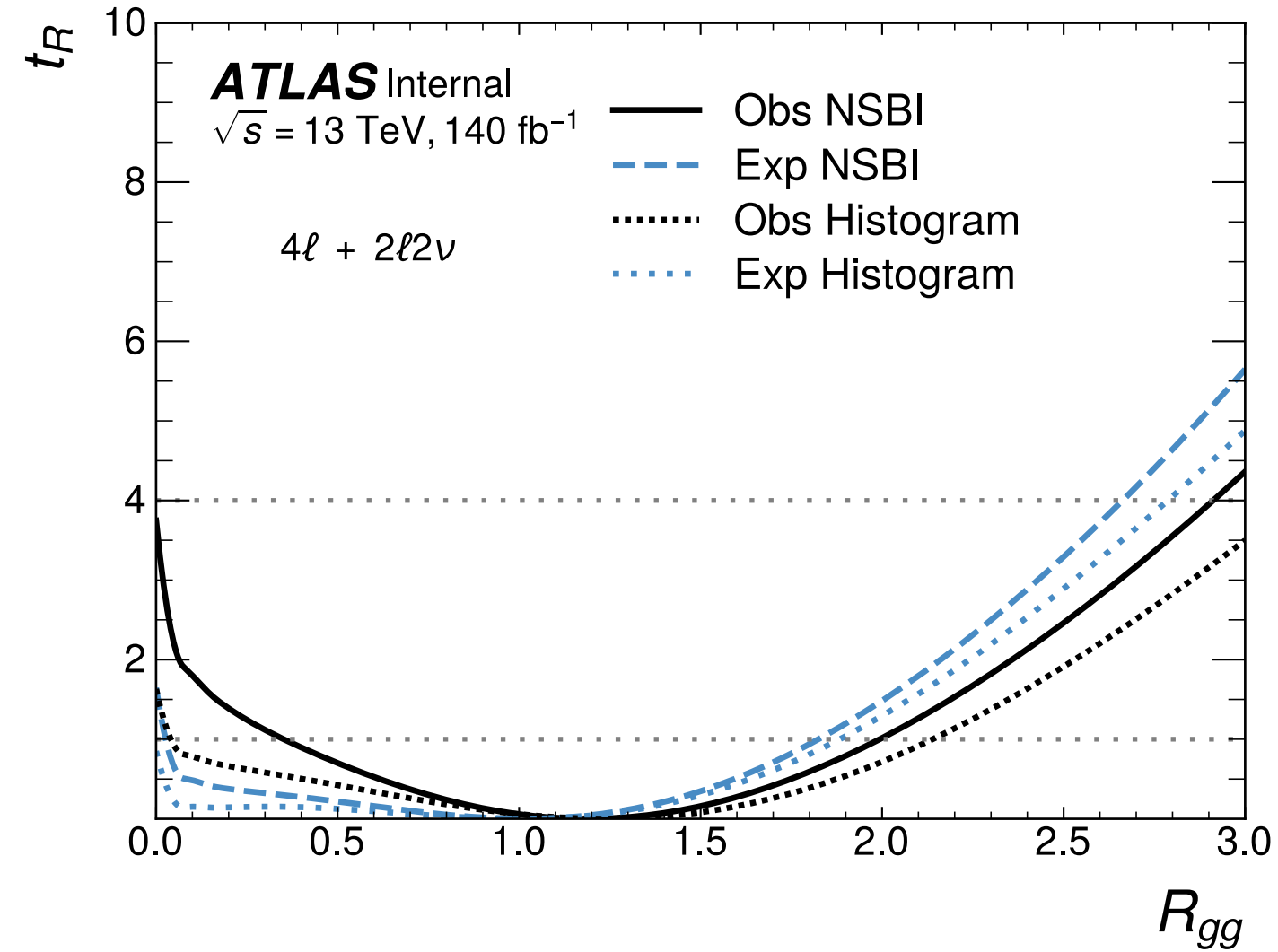
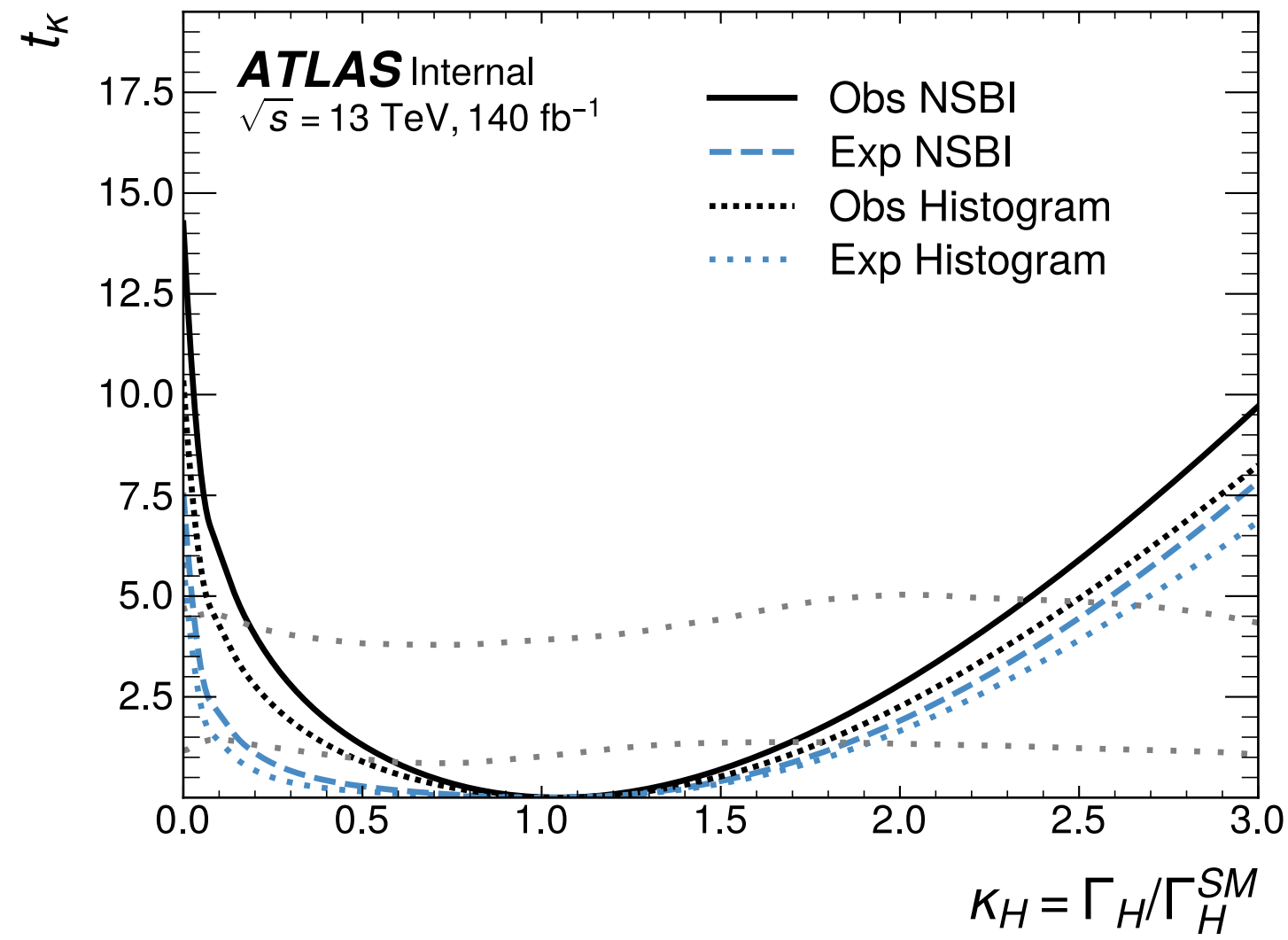
ttH+4t Systematic Uncertainty Impacts

Systematic uncertainty	Impact on 95% CL upper limit on Γ_H	
	Expected [%]	Observed [%]
Theory	37	33
$t\bar{t}t\bar{t}$ production	25	13
Higgs boson production/decay	5	6
Other processes	10	16
Experimental	2	2
Jet flavour tagging	2	1
Jet and missing transverse energy	< 1	< 1
Leptons and photons	< 1	< 1
All other systematic uncertainties	< 1	< 1

HZZ Systematic Uncertainty Impacts

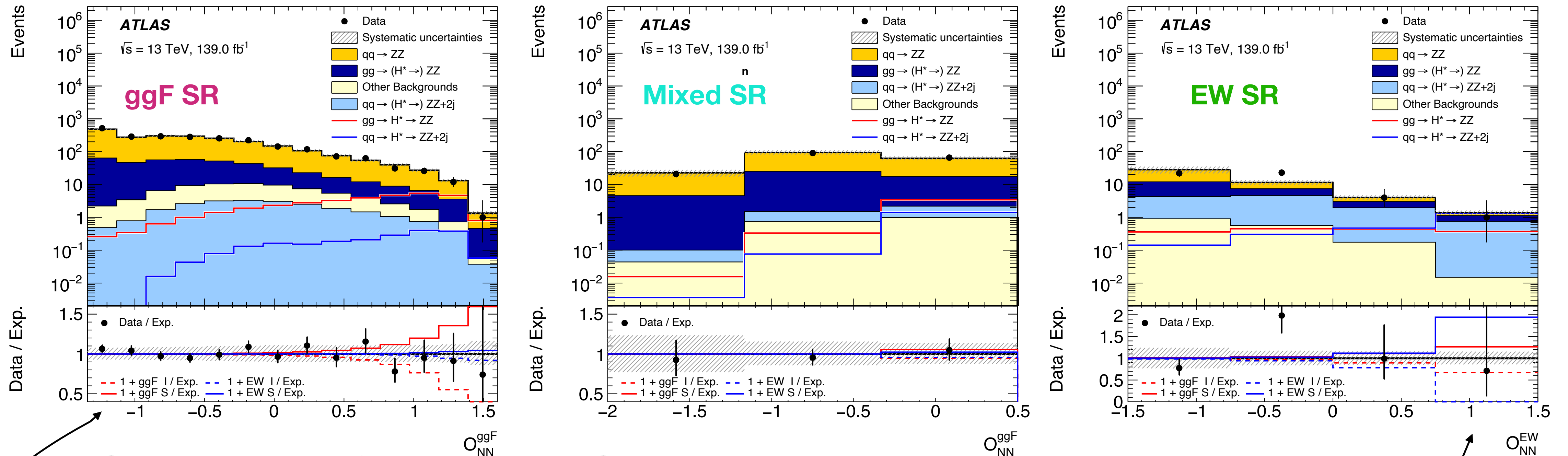
Uncertainty source	Absolute impact on $\mu_{\text{off-shell}}$	
	Nuisance Parameter	Global Observable
Electron uncertainties	(-5%, +6%)	(-5%, +6%)
Muon uncertainties	(-3%, +3%)	(-2%, +3%)
Jet uncertainties	(-10%, +10%)	(-9%, +11%)
Luminosity	(-1%, +1%)	(-1%, +1%)
Total experimental	-	(-11%, +12%)
$q\bar{q} \rightarrow ZZ$ modeling	(-6%, +7%)	(-6%, +7%)
$gg \rightarrow ZZ$ modeling	(-8%, +13%)	(-7%, +9%)
EW $q\bar{q} \rightarrow ZZ + 2j$ modeling	(-1%, +1%)	(-1%, +1%)
Total modeling	-	(-9%, +12%)
Systematic uncertainty	-	(-14%, +17%)
Statistical uncertainty	-	(-51%, +73%)
Total uncertainty	(-53%, +75%)	

Comparison With Histogram Analysis



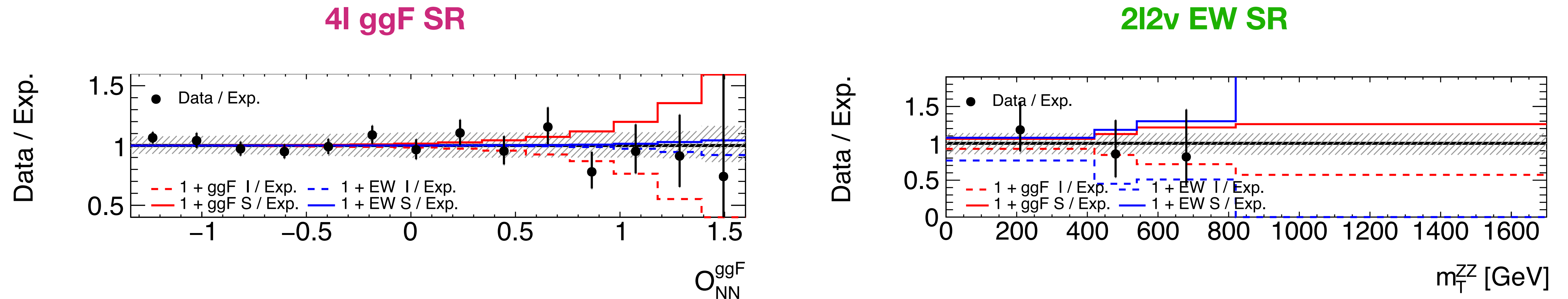
Parameter	Observed (expected)	
	NSBI analysis	histogram-based analysis
Γ_H [MeV] (4ℓ only)	$3.5^{+2.7}_{-2.2}$ ($4.1^{+3.6}_{-3.8}$)	$3.4^{+4.4}_{-3.4}$ ($4.1^{+4.2}_{-4.1}$)
Γ_H [MeV]	$4.3^{+2.7}_{-1.9}$ ($4.1^{+3.6}_{-3.8}$)	$4.3^{+3.0}_{-2.3}$ ($4.1^{+3.7}_{-3.7}$)
R_{gg}	$1.2^{+0.8}_{-0.8}$ ($1.0^{+0.8}_{-1.0}$)	$1.2^{+0.9}_{-1.1}$ ($1.0^{+0.8}_{-1.0}$)
R_{VV}	$0.95^{+0.35}_{-0.31}$ ($1.00^{+0.54}_{-0.60}$)	$0.95^{+0.40}_{-0.36}$ ($1.00^{+0.60}_{-0.66}$)

4l Final State



- Signal region is defined as $m_{4l} \geq 220$ GeV
 - Mixed (1-jet) and EW (2-jet) SRs are defined targeting EW (VBF+VH) production
 - Remaining events form the ggF SR
- qqZZ production is the dominant background
 - Constrained using data CRs, $180 < m_{4l} < 220$ GeV, $N_{\text{jets}} = 0, 1, \geq 2$
 - Other backgrounds are mainly ttV and VVV, fakes from Z+jet and ttbar are negligible
- NNs trained to separate ggF and EW signal from non-interfering qqZZ and interfering gg and qq backgrounds

Observables



- Observables do a good job of enhancing S/B at higher values
 - For both **ggF** and **EW** production
- Interference (dashed) goes opposite to signal (solid)

4l SR and CR Definitions

EW SR	Mixed SR	ggH SR	0-jet 4l CR	1-jet 4l CR	2-jet 4l CR
$N_{\text{jets}} \geq 2$	$N_{\text{jets}} = 1$	$N_{\text{jets}} = 0$ or ($N_{\text{jets}} = 1$ and $ \eta_j < 2.2$) or ($N_{\text{jets}} \geq 2$ and $\Delta\eta_{jj} < 4$)	$N_{\text{jets}} = 0$	$N_{\text{jets}} = 1$	$N_{\text{jets}} \geq 2$
$\Delta\eta_{jj} > 4$	$ \eta_j > 2.2$		$180 < m_{4l} < 220$ GeV	$180 < m_{4l} < 220$ GeV	$180 < m_{4l} < 220$ GeV

$m_{4l} > 220$ GeV

$m_{4l} > 220$ GeV

$m_{4l} > 220$ GeV

- Unprescaled single-lepton triggers
- 3 leading lepton p_T thresholds are 20, 15, and 10 GeV
- $\Delta R(l, l') > 0.1$
- $50 < m_{ll} < 115$ GeV
- Four leptons must pass a vertex fit

2l2v SR and CR Definitions

EW SR	Mixed SR	ggH SR	0-jet 3l CR	1-jet 3l CR	2-jet 3l CR	eμ CR	Zjets CR
$N_{\text{jets}} \geq 2$	$N_{\text{jets}} = 1$	$N_{\text{jets}} = 0$ OR	$N_{\text{jets}} = 0$	$N_{\text{jets}} = 1$	$N_{\text{jets}} \geq 2$	OFOS pair	
$\Delta\eta_{jj} > 4$	$ \eta_j > 2.2$	$(N_{\text{jets}} = 1, \eta_j < 2.2)$ OR $(N_{\text{jets}} \geq 2, \Delta\eta_{jj} < 4)$	3rd lepton with $p_{T} > 20$ GeV	3rd lepton with $p_{T} > 20$ GeV	3rd lepton with $p_{T} > 20$ GeV		
7 GeV l_3 veto	7 GeV l_3 veto	7 GeV l_3 veto	7 GeV l_4 veto	7 GeV l_4 veto	7 GeV l_4 veto	7 GeV l_3 veto	7 GeV l_3 veto
$E_{T}^{\text{miss}} > 120$ GeV	$E_{T}^{\text{miss}} > 120$ GeV	$E_{T}^{\text{miss}} > 120$ GeV	$m_t^W > 60$ GeV	$m_t^W > 60$ GeV	$m_t^W > 60$ GeV	$E_{T}^{\text{miss}} > 120$ GeV	$E_{T}^{\text{miss}} > 120$ GeV
$\Delta R_{ll} < 1.8$	$\Delta R_{ll} < 1.8$	$\Delta R_{ll} < 1.8$	$m_T(W) = \sqrt{2p_T^\ell E_T^{\text{miss}}(1 - \cos\Delta\phi)}$			$\Delta R_{ll} < 1.8$	$\Delta R_{ll} < 1.8$
$\Delta\phi(p_{T,ll}, E_T^{\text{miss}}) > 2.5$	$\Delta\phi(p_{T,ll}, E_T^{\text{miss}}) > 2.5$	$\Delta\phi(p_{T,ll}, E_T^{\text{miss}}) > 2.5$				$\Delta\phi(p_{T,ll}, E_T^{\text{miss}}) > 2.5$	$\Delta\phi(p_{T,ll}, E_T^{\text{miss}}) > 2.5$
$\Delta\phi(\text{jet}_{>100}, E_T^{\text{miss}}) > 0.4$	$\Delta\phi(\text{jet}_{>100}, E_T^{\text{miss}}) > 0.4$	$\Delta\phi(\text{jet}_{>100}, E_T^{\text{miss}}) > 0.4$				$\Delta\phi(\text{jet}_{>100}, E_T^{\text{miss}}) > 0.4$	
$E_{T}^{\text{miss}} \text{ signif} > 10$	$E_{T}^{\text{miss}} \text{ signif} > 10$	$E_{T}^{\text{miss}} \text{ signif} > 10$	$E_{T}^{\text{miss}} \text{ signif} > 3$	$E_{T}^{\text{miss}} \text{ signif} > 3$	$E_{T}^{\text{miss}} \text{ signif} > 3$	$E_{T}^{\text{miss}} \text{ signif} > 10$	$E_{T}^{\text{miss}} \text{ signif} < 9$
b-jet veto	b-jet veto	b-jet veto	b-jet veto	b-jet veto	b-jet veto	b-jet veto	b-jet veto

- Unprescaled single-lepton triggers

- Lepton $p_{T} > 20$ GeV

- $76 < m_{ll} < 106$ GeV

Leading Systematic Uncertainties

Process	Uncertainty	Final State	Value (%)
ggF Signal Region			
$qq \rightarrow ZZ$	QCD Scale	$2\ell 2\nu$	4–40
$qq \rightarrow ZZ + 2j$	QCD Scale	4ℓ	21–28
$qq \rightarrow ZZ + 2j$	QCD Scale	$2\ell 2\nu$	22–37
$qq \rightarrow ZZ + 2j$	Parton Shower	$2\ell 2\nu$	1–67
$gg \rightarrow H^* \rightarrow ZZ$	Parton Shower	4ℓ	27
$gg \rightarrow H^* \rightarrow ZZ$	Parton Shower	$2\ell 2\nu$	8–45
$gg \rightarrow ZZ$	Parton Shower	4ℓ	38
$gg \rightarrow ZZ$	Parton Shower	$2\ell 2\nu$	6–43
$WZ + 0j$	QCD Scale	$2\ell 2\nu$	1–54
1-jet Signal Region			
$gg \rightarrow H^* \rightarrow ZZ$	Parton Shower	4ℓ	27
$gg \rightarrow H^* \rightarrow ZZ$	QCD Scale	$2\ell 2\nu$	13–18
$gg \rightarrow ZZ$	Parton Shower	4ℓ	38
$gg \rightarrow ZZ$	QCD Scale	$2\ell 2\nu$	18–20
$qq \rightarrow ZZ$ (EW)	QCD Scale	$2\ell 2\nu$	7–18
2-jet Signal Region			
$qq \rightarrow ZZ + 2j$	QCD Scale	4ℓ	18–26
$qq \rightarrow ZZ + 2j$	QCD Scale	$2\ell 2\nu$	8–32
$gg \rightarrow H^* \rightarrow ZZ$	Parton Shower	4ℓ	27
$gg \rightarrow ZZ$	Parton Shower	4ℓ	38
$gg \rightarrow ZZ$	QCD Scale	$2\ell 2\nu$	18–20
$WZ + 2j$	QCD Scale	$2\ell 2\nu$	20–22
$qq \rightarrow ZZ$ Control Regions			
$qq \rightarrow ZZ + 2j$	QCD Scale	4ℓ	26
Three-lepton Control Regions			
$WZ + 2j$	QCD Scale	$2\ell 2\nu$	28

Systematic Uncertainty Fixed	$\mu_{\text{off-shell}}$ value at which $-2 \ln \lambda(\mu_{\text{off-shell}}) = 4$
Parton shower uncertainty for $gg \rightarrow ZZ$ (normalisation)	2.26
Parton shower uncertainty for $gg \rightarrow ZZ$ (shape)	2.29
NLO EW uncertainty for $qq \rightarrow ZZ$	2.27
NLO QCD uncertainty for $gg \rightarrow ZZ$	2.29
Parton shower uncertainty for $qq \rightarrow ZZ$ (shape)	2.29
Jet energy scale and resolution uncertainty	2.26
None	2.30

Yields

Process	ggF SR	Mixed SR	EW SR
$gg \rightarrow (H^* \rightarrow)ZZ$	341 ± 117	42.5 ± 14.9	11.8 ± 4.3
$gg \rightarrow H^* \rightarrow ZZ$	32.6 ± 9.07	3.68 ± 1.03	1.58 ± 0.47
$gg \rightarrow ZZ$	345 ± 119	43.0 ± 15.2	11.9 ± 4.4
$qq \rightarrow (H^* \rightarrow)ZZ + 2j$	23.2 ± 1.0	2.03 ± 0.16	9.89 ± 0.96
$qq \rightarrow ZZ$	1878 ± 151	135 ± 23	22.0 ± 8.3
Other backgrounds	50.6 ± 2.5	1.79 ± 0.16	1.65 ± 0.16
Total expected (SM)	2293 ± 209	181 ± 29	45.3 ± 10.0
Observed	2327	178	50

Process	ggF SR	Mixed SR	EW SR
$gg \rightarrow (H^* \rightarrow)ZZ$	210 ± 53	19.7 ± 4.9	4.29 ± 1.10
$gg \rightarrow H^* \rightarrow ZZ$	111 ± 26	10.9 ± 2.5	3.26 ± 0.82
$gg \rightarrow ZZ$	251 ± 66	23.4 ± 6.2	5.31 ± 1.46
$qq \rightarrow (H^* \rightarrow)ZZ + 2j$	14.0 ± 3.0	1.63 ± 0.17	4.46 ± 0.50
$qq \rightarrow ZZ$	1422 ± 112	80.4 ± 11.9	7.74 ± 2.99
WZ	678 ± 54	51.9 ± 6.9	7.89 ± 2.50
Z+jets	62.3 ± 24.3	7.51 ± 6.94	0.62 ± 0.54
Non-resonant- $l\bar{l}$	106 ± 39	9.17 ± 2.73	1.55 ± 0.42
Other backgrounds	22.6 ± 5.2	1.62 ± 0.25	1.40 ± 0.10
Total expected (SM)	2515 ± 165	172 ± 17	28.0 ± 4.1
Observed	2496	33	27