

Indications for New Higgs Bosons in **Associated Di-Photon Production**

HIGGS 2024 6th November 2024

Swiss National Science Foundation



Universität Zürich

Sumit Banik







Motivation Hints for new Higgs Bosons

O Minimality of the scalar sector of the SM not guaranteed theoretically

O ATLAS recently performed Model-Independent analysis of $\gamma\gamma + X$ for SM Higgs

O Analysis involves 22 signal regions

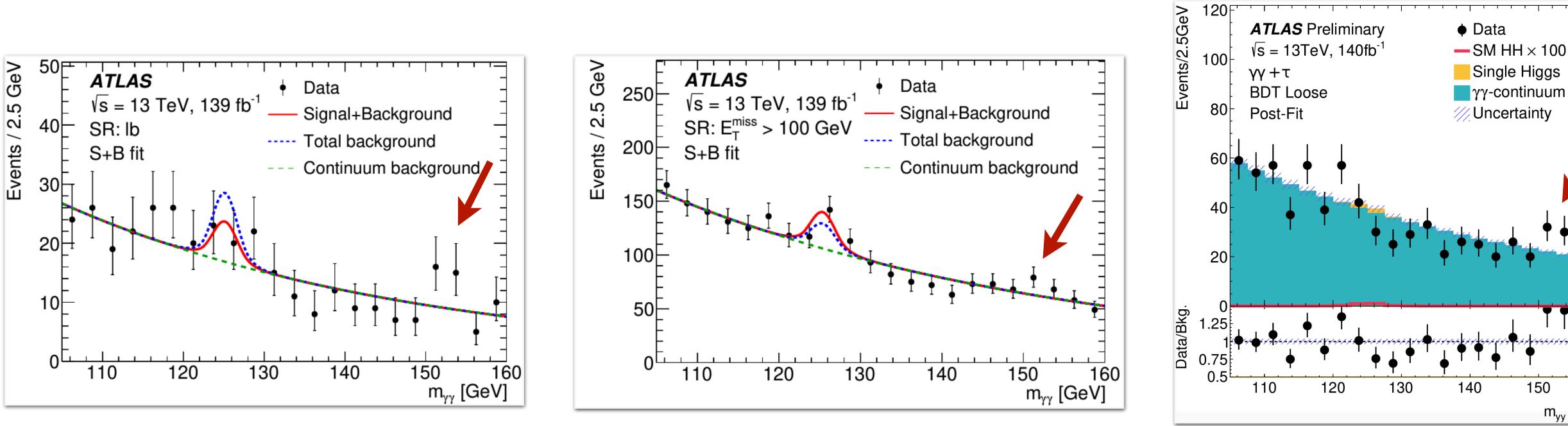
Heavy Flavor	Jets	Lepton	E_T^{miss}	Тор	H_T	Photo
$\geq 3b, \geq 4b$	$\geq 4j, \geq 6j,$ etc.	$ \begin{array}{rcl} 1\ell, \ 2\ell, \ \geq \ 3\ell, \\ 1\tau, \ 2\tau \end{array} $	$\begin{array}{l} E_T^{\rm miss} > 100 {\rm GeV}, \\ E_T^{\rm miss} > 200 {\rm GeV} \end{array}$	$\ell b, t_{\rm lep}, t_{\rm had}$	$H_T > 500 \text{ GeV}, \ H_T > 1000 \text{ GeV}$	$m_{\gamma\gamma}^{12}, r$

Full Run 2 Data



Motivation Hints for new Higgs Bosons

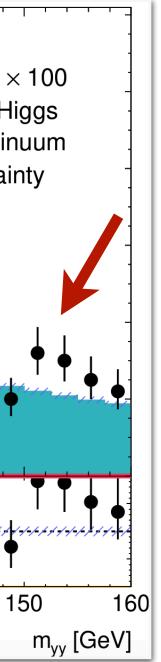
O Excesses Most Pronounced: $\gamma\gamma + \ell b$, $\gamma\gamma + MET$, $\gamma\gamma + 1\tau$, $\gamma\gamma + 4j$, $\gamma\gamma + 1\ell$



[ATLAS: CERN-EP-2022-232]

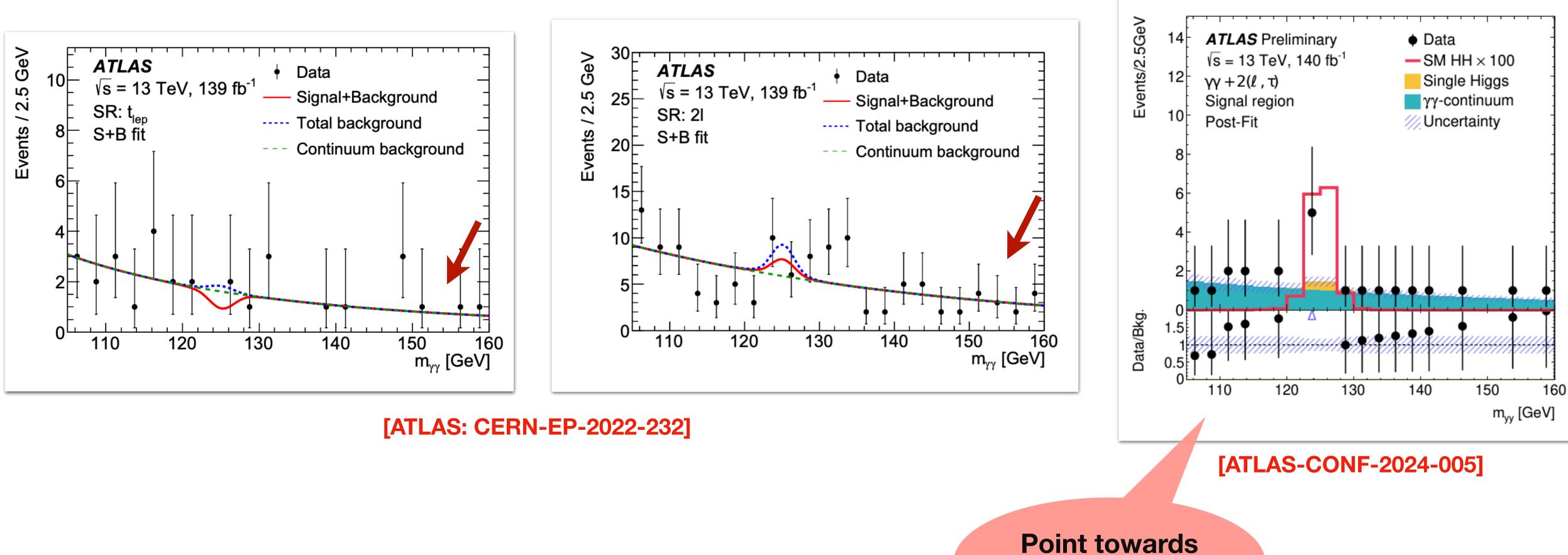
O Possible new Higgs Boson?

[ATLAS-CONF-2024-005]



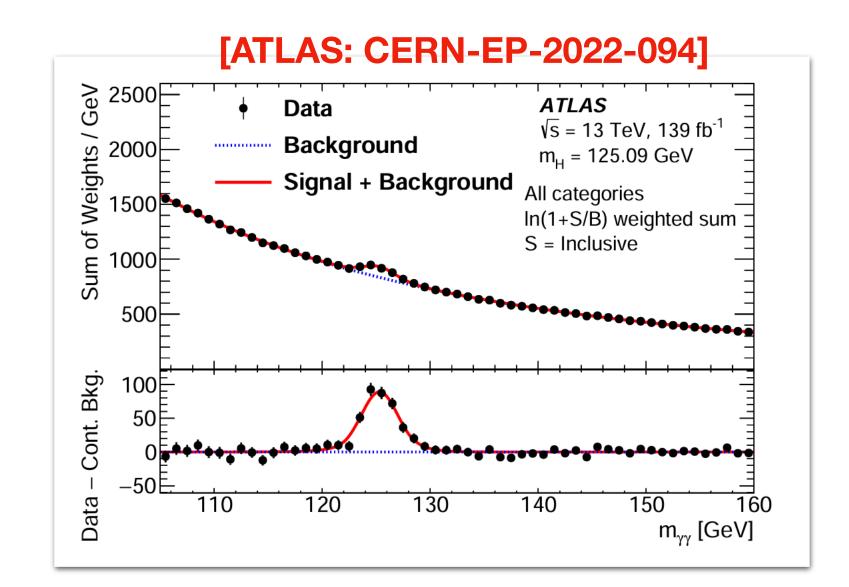
Motivation **Hints for new Higgs Bosons**

o No Excesses at 152 GeV in SRs: $\gamma\gamma + t_{lep}$, $\gamma\gamma + 2\ell$, $\gamma\gamma + 2\tau$,



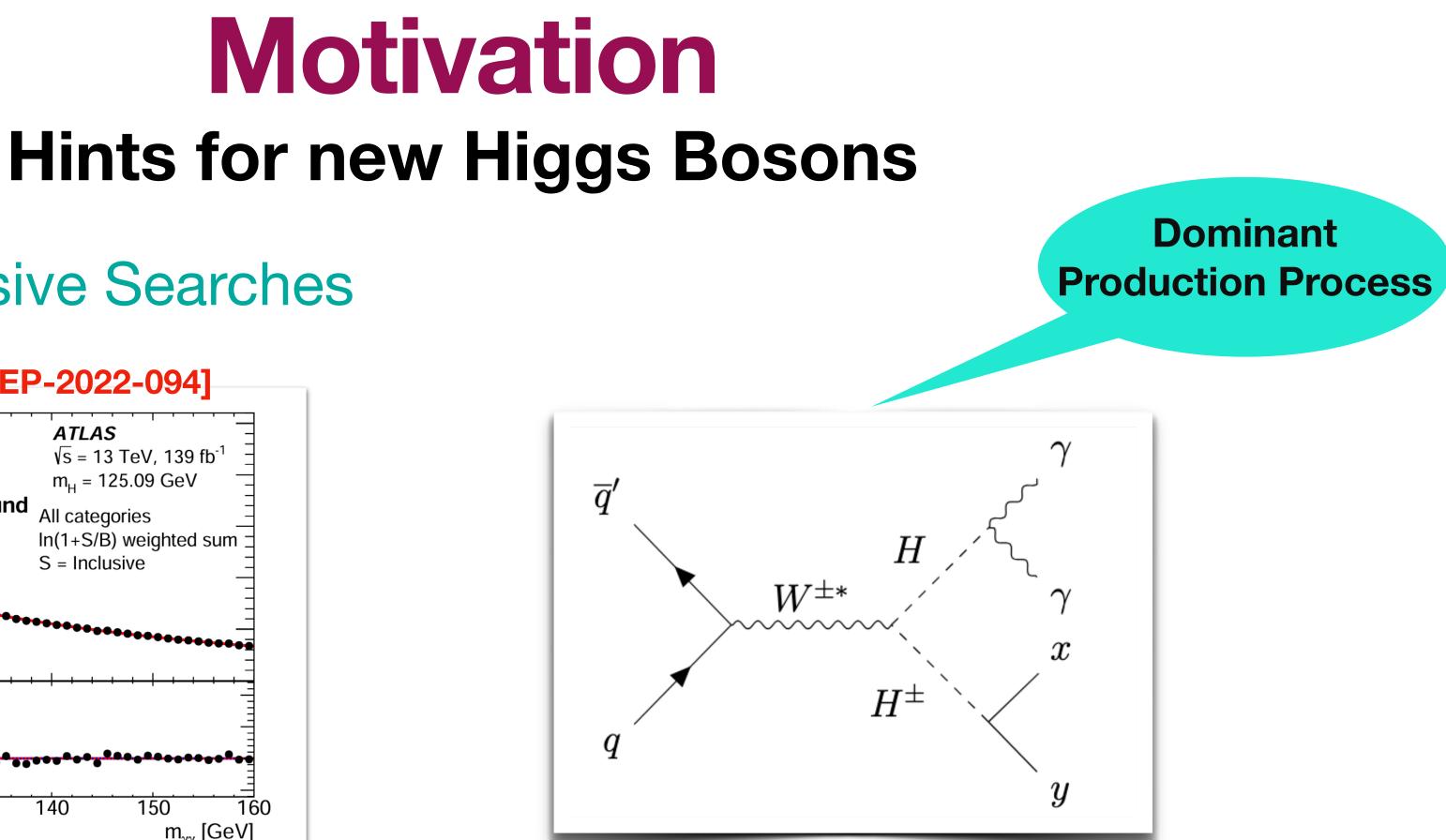
asssociated H^{\pm}

O No excess in Inclusive Searches



O Hints towards DY production of new Higgs at LHC

• Properties of H^{\pm} unknown



Simplified Model Model Description

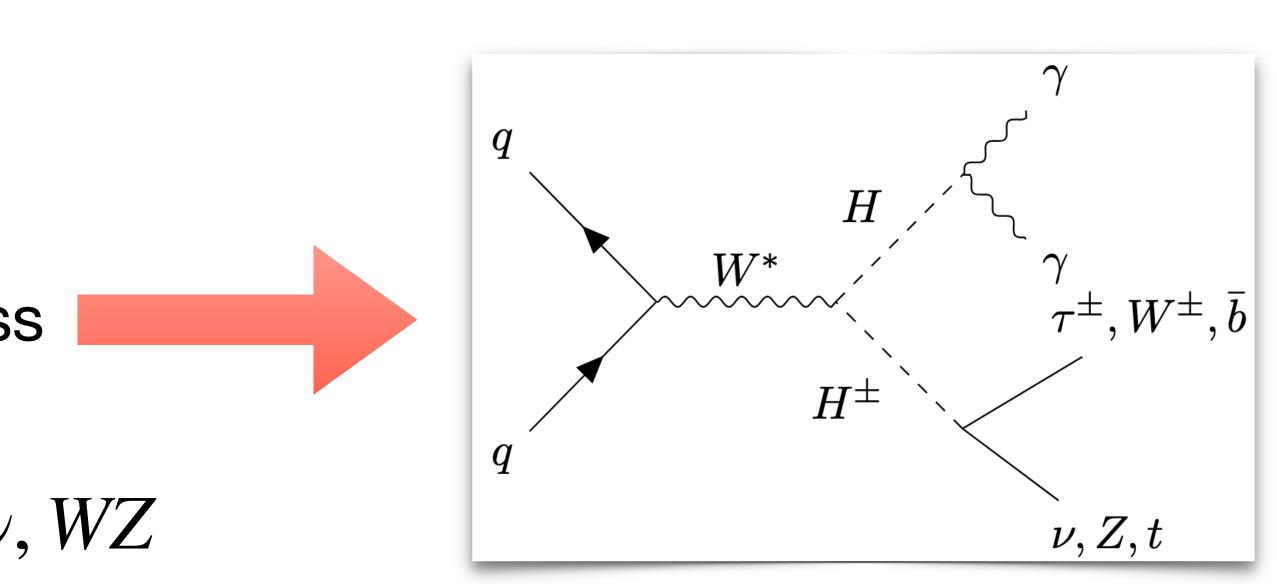
O Two New Particles: H, H^{\pm}

o H produced only via DY process

O Dominant decays of H^{\pm} : $tb, \tau\nu, WZ$

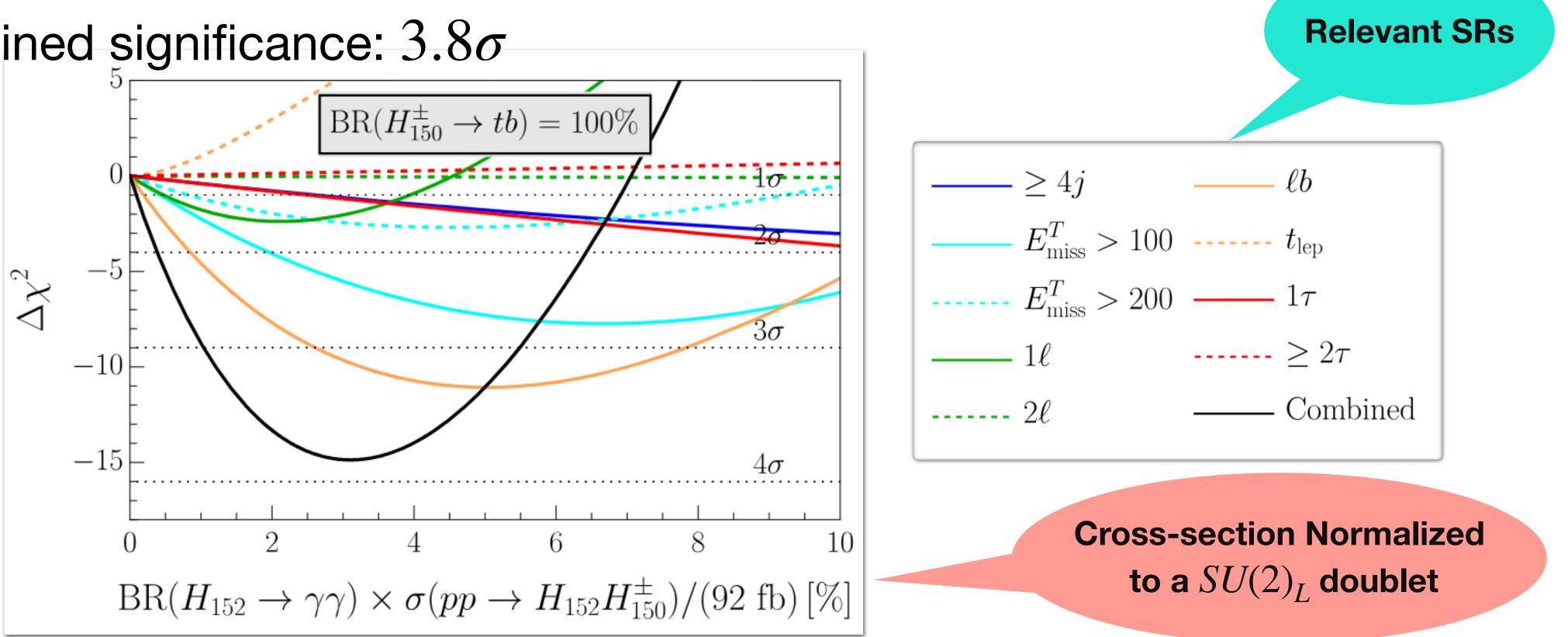
• Simulation Setup: MadGraph + Pythia + Delphes

O Log-Likelihood Fit performed using Poisson Statistics



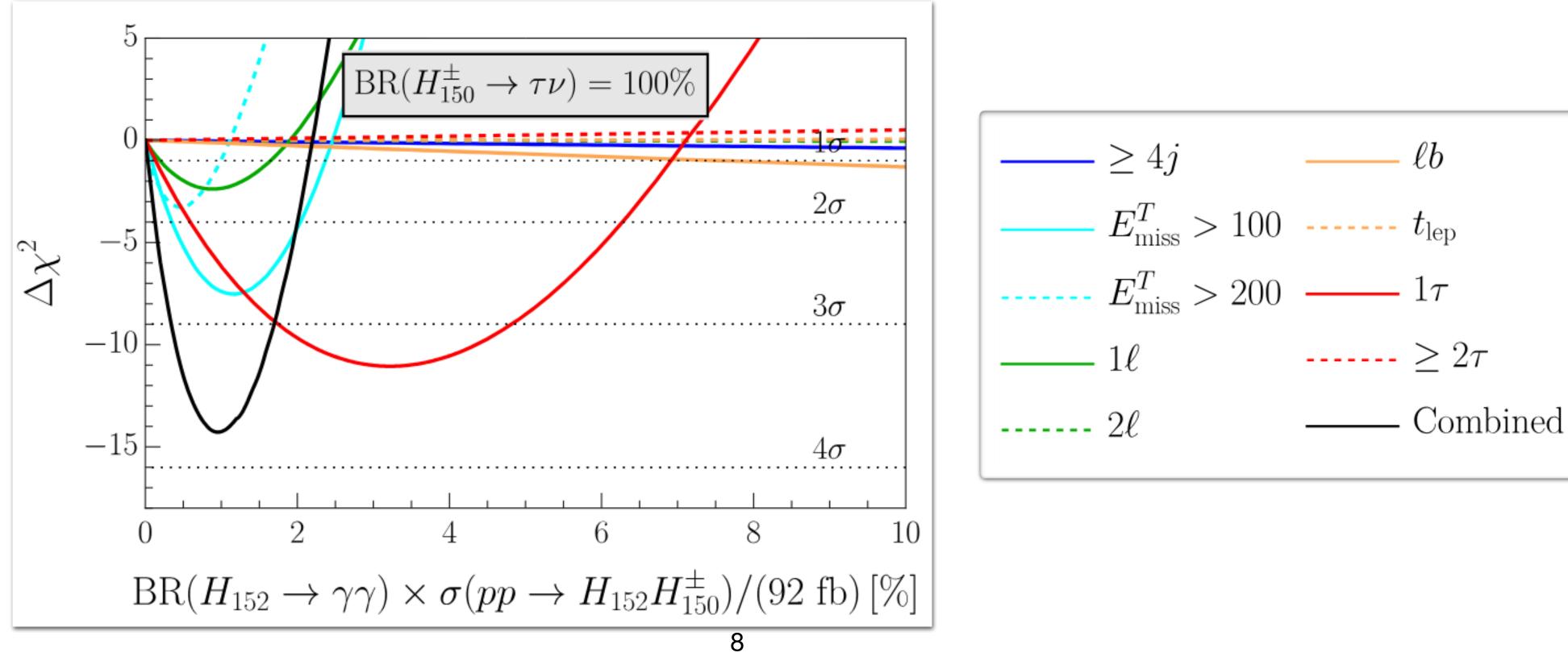
Simplified Model Charged Higgs Decay

- $O BR(H^{\pm} \rightarrow tb \rightarrow bbW) = 100\%$
- o Dominant Effect: $\gamma\gamma + \ell b, \gamma\gamma + MET, \gamma\gamma + 1\ell, \gamma\gamma + t_{ep}$
- O Combined significance: 3.8σ



Simplified Model Charged Higgs Decay $\circ BR(H^{\pm} \to \tau \nu) = 100\%$

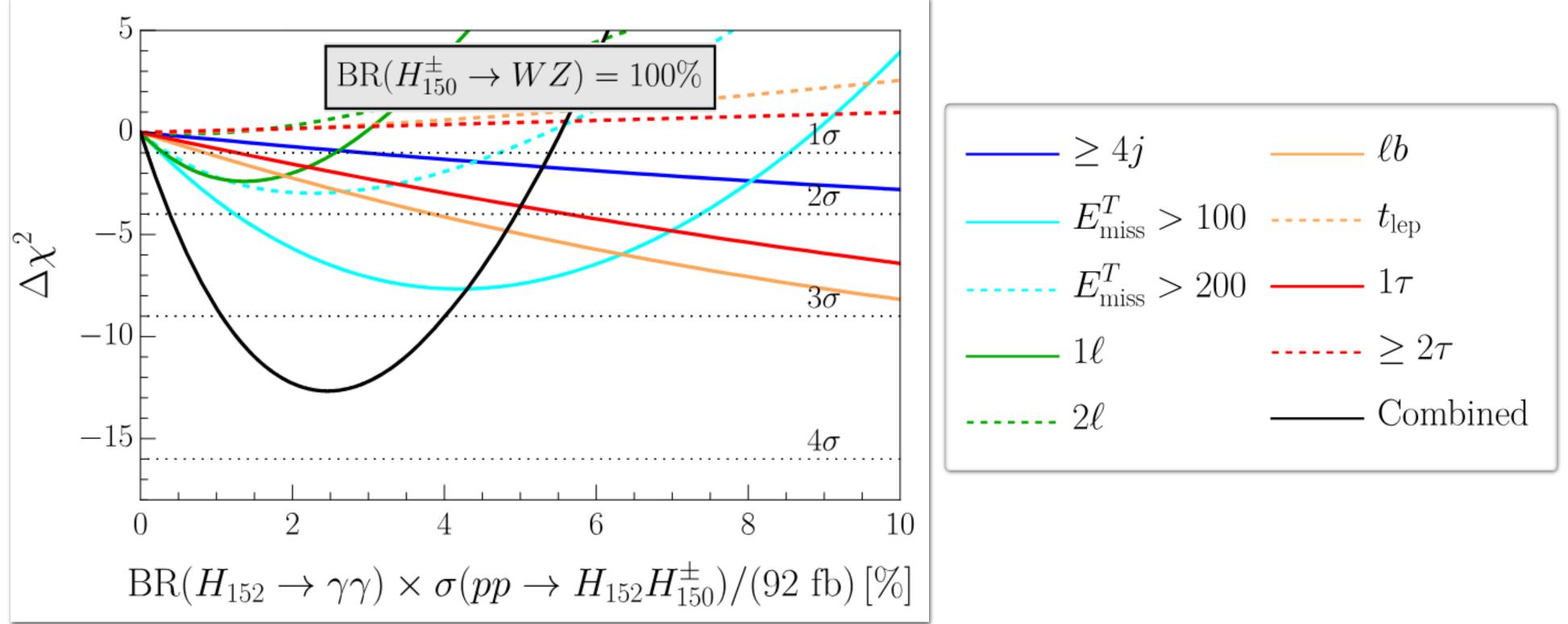
- O Dominant Effect: $\gamma\gamma + MET, \gamma\gamma + 1\tau, \gamma\gamma + 1\ell$
- o Combined significance: 3.8σ



Simplified Model Charged Higgs Decay

$\circ BR(H^{\pm} \rightarrow WZ) = 100\%$

- O Dominant Effect: $\gamma\gamma + MET$, $\gamma\gamma + 1\ell$, $\gamma\gamma + 2\ell$, $\gamma\gamma + 2\tau$
- O Combined significance: 3.5σ



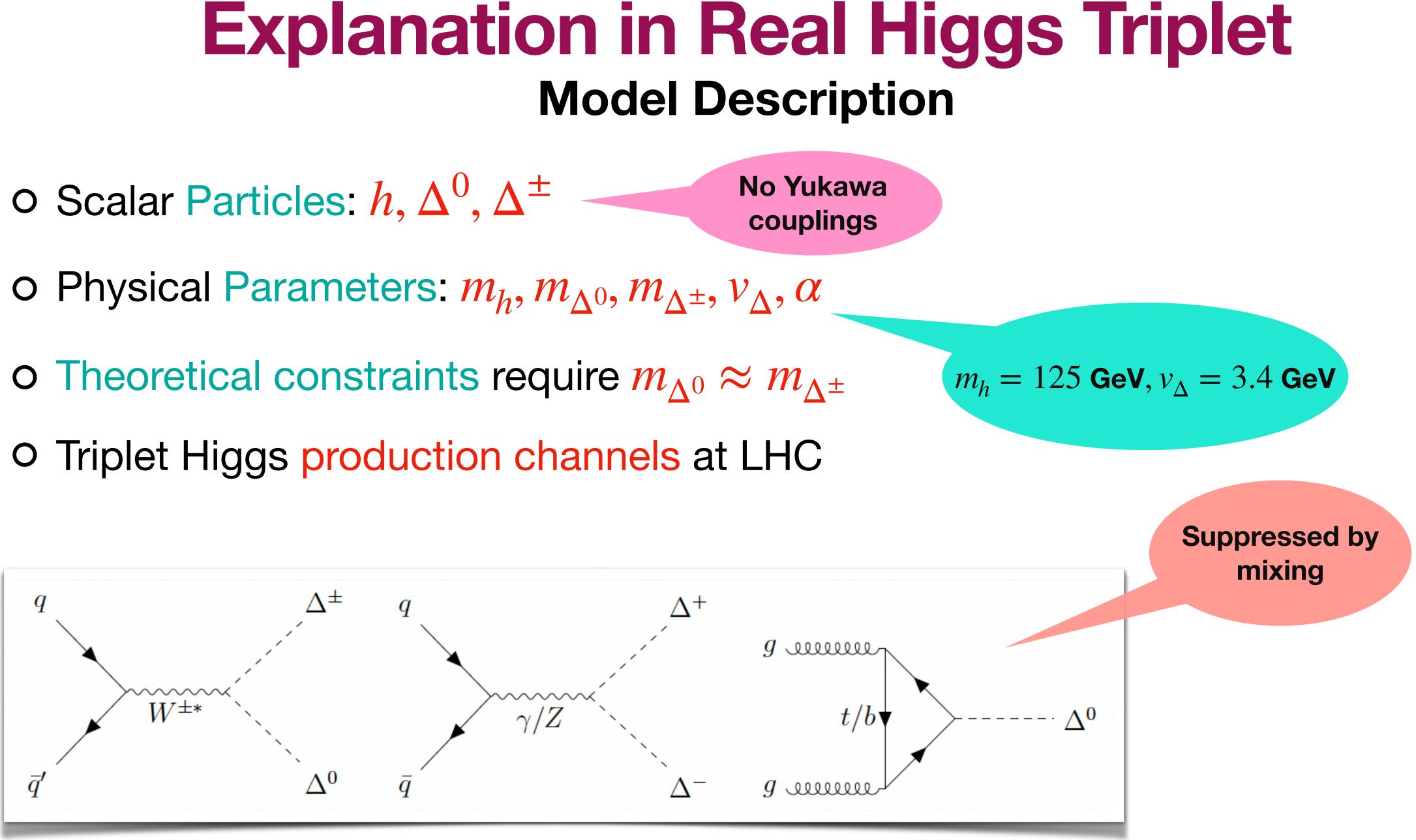
9

Dominant in Triplet Model



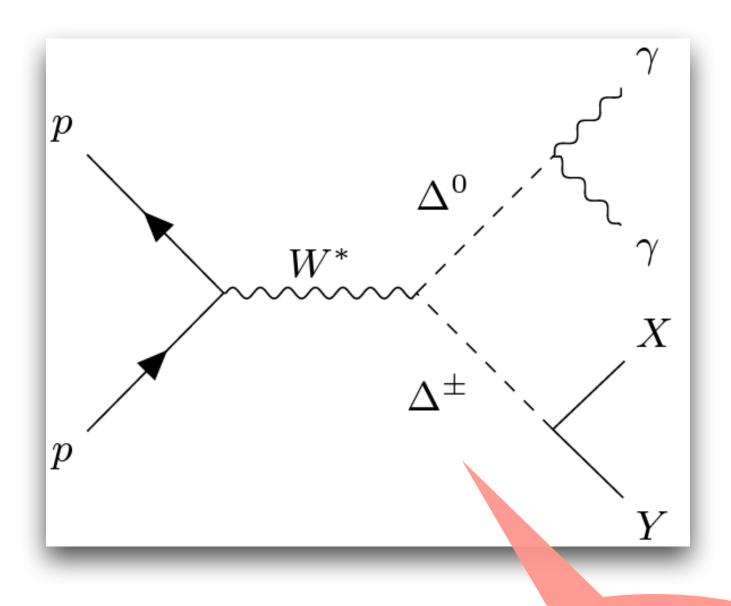
Model Building Key Points

- O Small total production cross-section
- O Dominant DY production cross-section
- o Large BR($H^{\pm} \rightarrow tb$) and BR($H^{\pm} \rightarrow \tau\nu$)
- O Small BR($H^{\pm} \rightarrow WZ$) to avoid multiple leptons
- Sizable BR($H \rightarrow \gamma \gamma$)

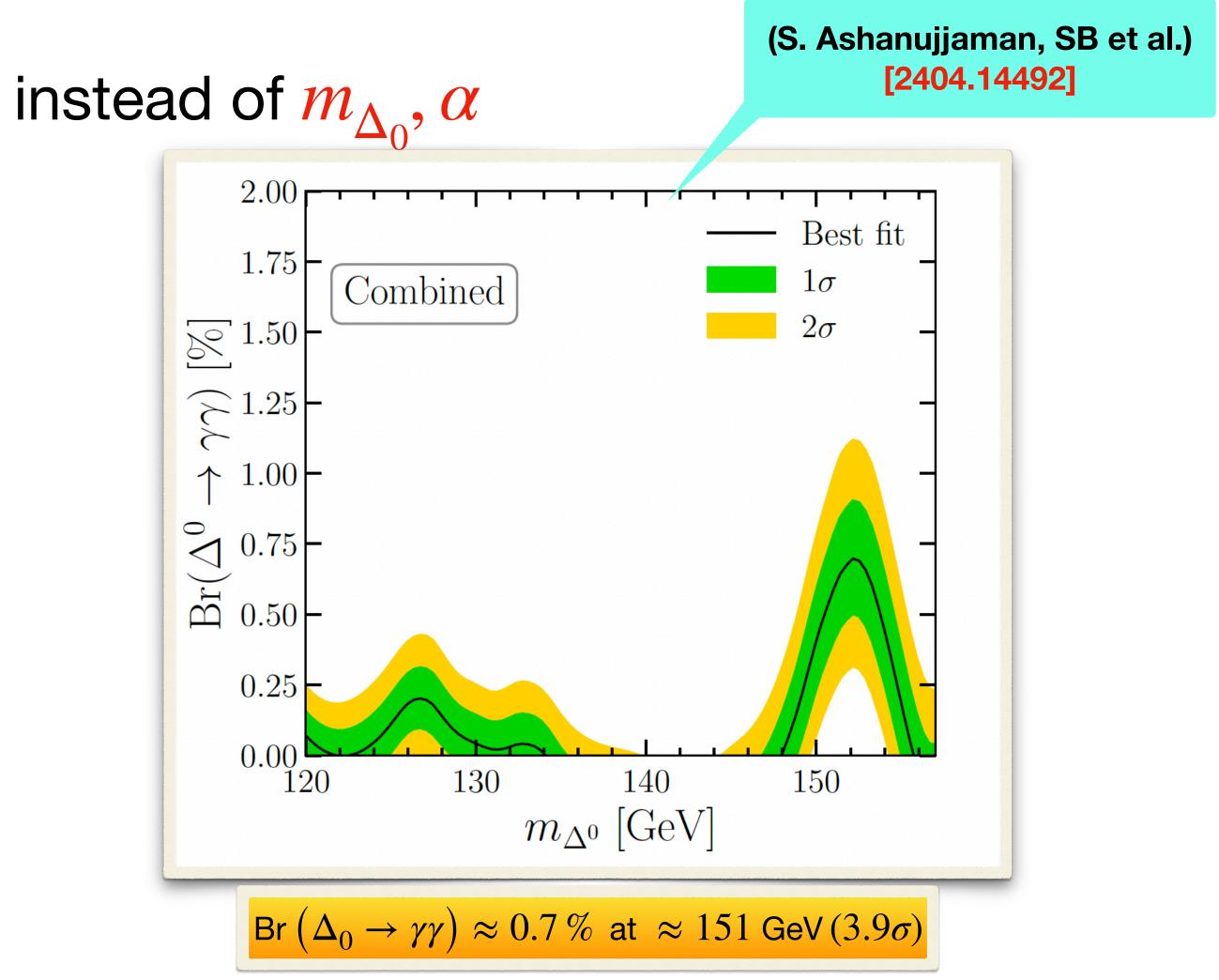


Explanation in Real Higgs Triplet Explanation of $\gamma\gamma + X$ Excesses

- O Model generated using SARAH
- Free Variables: m_{Δ_0} , Br $(\Delta_0 \rightarrow \gamma \gamma)$ instead of m_{Δ_0} , α

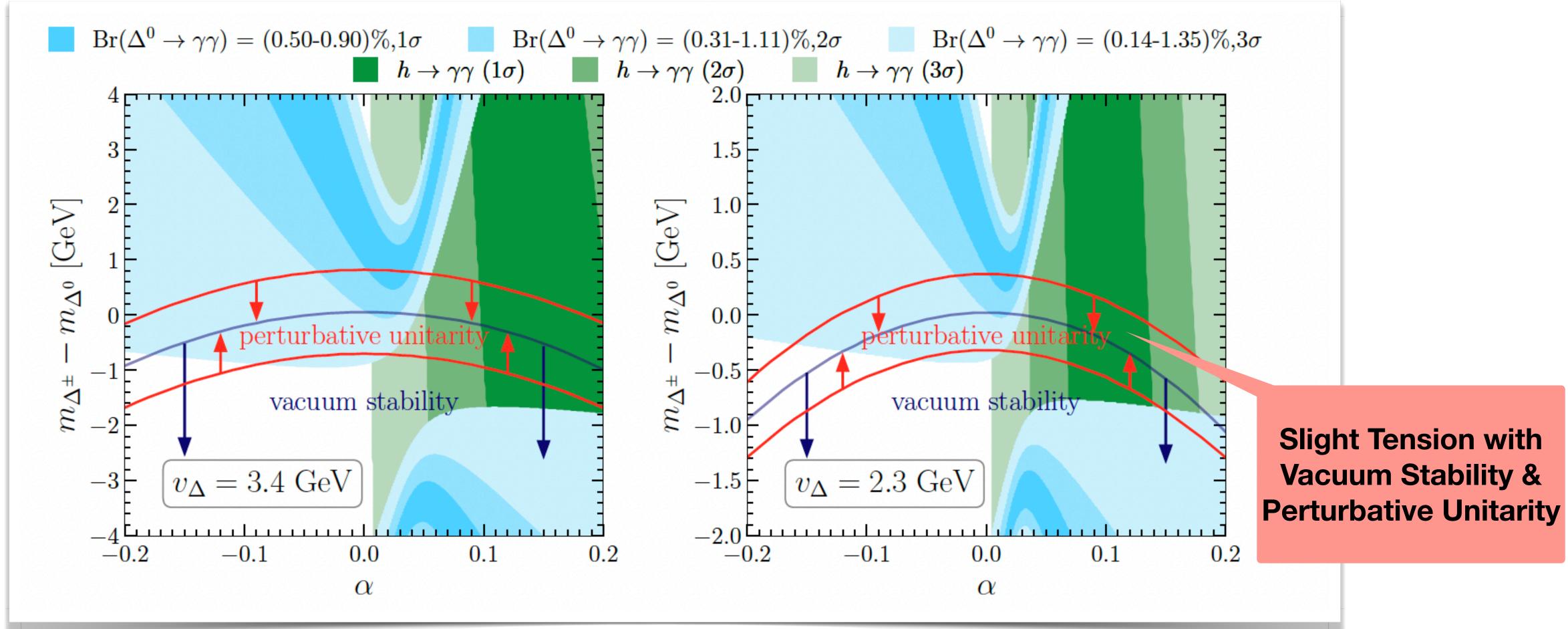


Process Simulated



Explanation in Real Higgs Triplet Explanation of $\gamma\gamma + X$ **Excesses**

O BR($\Delta^0 \rightarrow \gamma \gamma$) compatible with SM Higgs signal strength to $\gamma \gamma$





Explanation in 2HDM Description • Two $SU(2)_L$ doublets: ϕ_1 and ϕ_2 • Scalar potential $V(\phi_1, \phi_2) = m_{11}\phi_1^{\dagger}\phi_1 + m_{22}\phi_2^{\dagger}\phi_2 - m_{12}^2(\phi_1^{\dagger}\phi_2 + \text{h.c.})$

 $+\lambda_2(\phi_2^{\dagger}\phi_2)^2 + \lambda_3(\phi_1^{\dagger}\phi_1)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_1^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_2) + \lambda_3(\phi_2^{\dagger}\phi_2)(\phi_2^{\dagger}\phi_$

O Scalar Particles: h, H, A, H^{\pm}

o Free Parameters: m_h, m_H, m_A, m_H

$$b_2 - m_{12}^2 (\phi_1^{\dagger} \phi_2 + \text{h.c.}) + \lambda_1 (\phi_1^{\dagger} \phi_1)^2 + \lambda_4 (\phi_1^{\dagger} \phi_2) (\phi_2^{\dagger} \phi_1) + \lambda_5 ((\phi_1^{\dagger} \phi_2)^2 + \text{h.c})$$

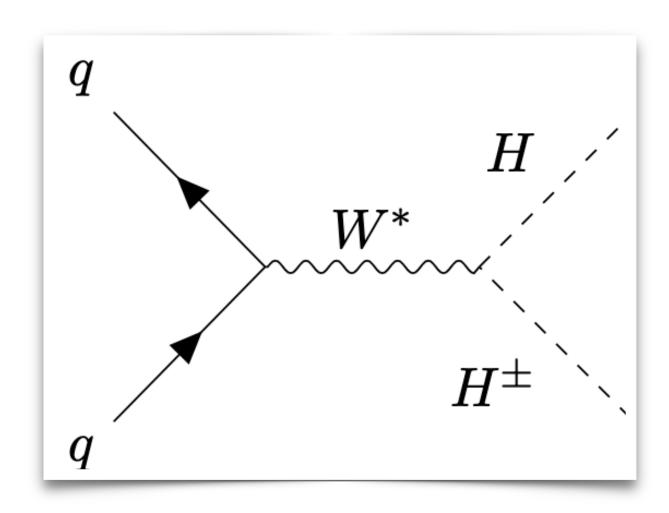
$$_{I^{\pm}}, m_{12}^2, \tan\beta = v_2/v_1, \alpha$$



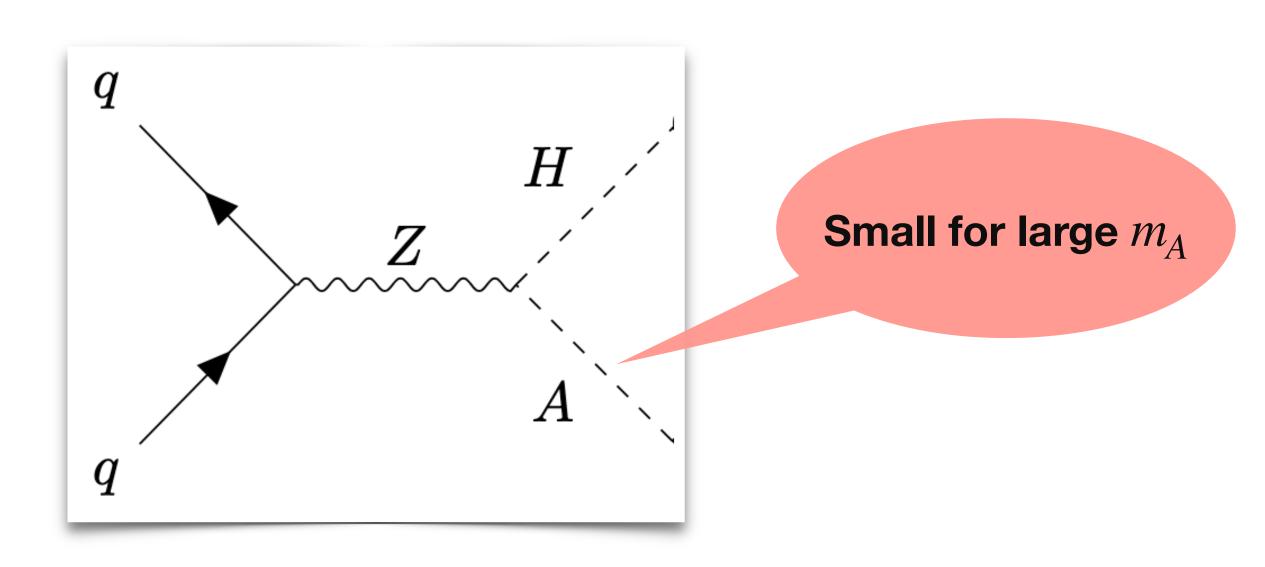
o Yukawa Sector

$Y = -\bar{Q}_{L}\phi_{2}d_{R} - \bar{Q}_{L}\phi_{2}^{c}u_{R} - \bar{L}_{L}\phi_{2}e_{R}$

- O Suppressed gluon-fusion, VBF, VH cross-section of H for large tan β
- O Dominant production channels for H



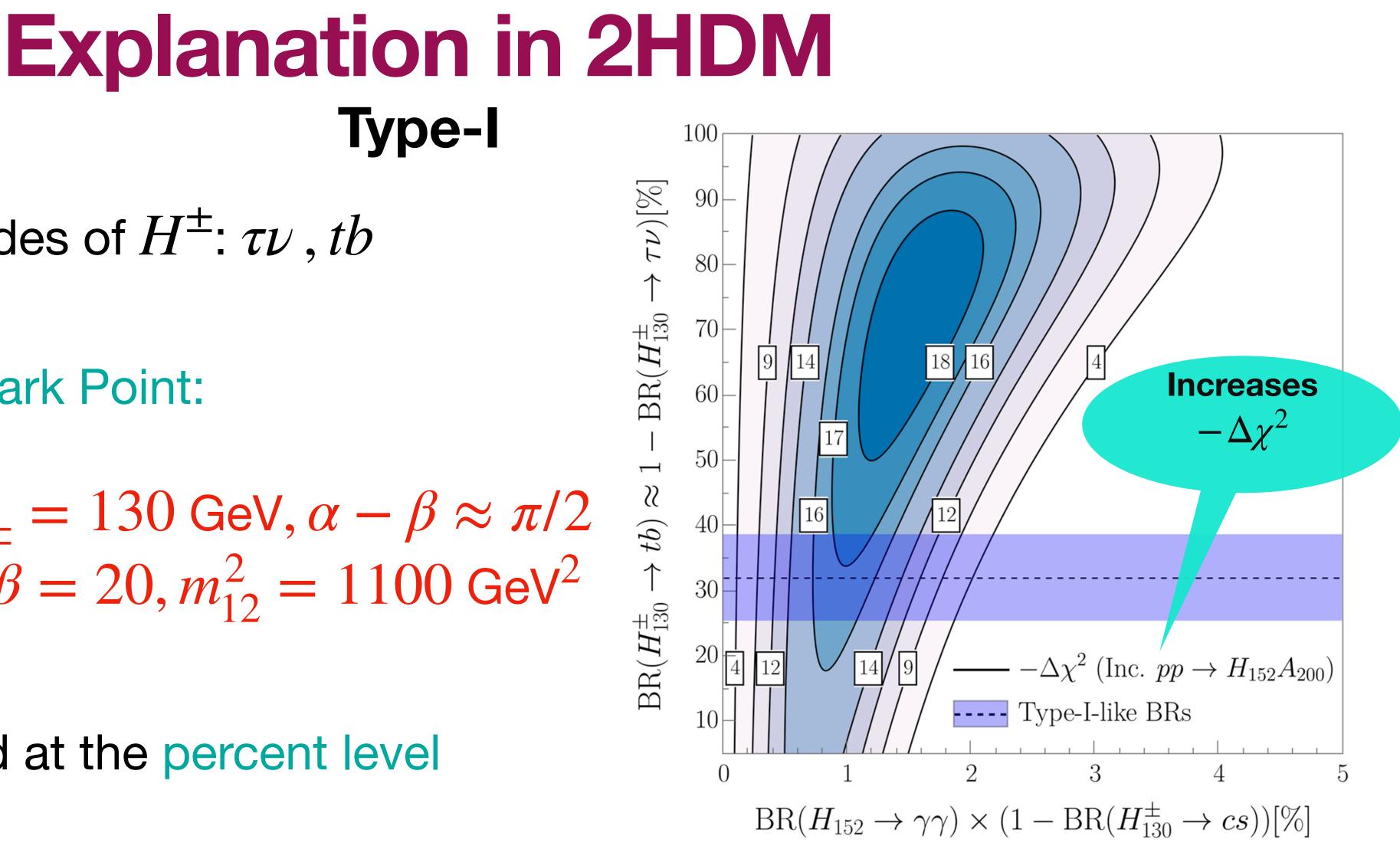
Explanation in 2HDM Type-I



- O Dominant decay modes of H^{\pm} : $\tau\nu$, tb
- **Considered Benchmark Point:** Ο

 $m_H = 152 \text{ GeV}, m_{H\pm} = 130 \text{ GeV}, \alpha - \beta \approx \pi/2$ $m_A = 200 \text{ GeV}, \tan \beta = 20, m_{12}^2 = 1100 \text{ GeV}^2$

- o $Br(H \rightarrow \gamma \gamma)$ required at the percent level
- O Possible with Effective Operator: $F_{\mu\nu}F^{\mu\nu}\phi_1^{\dagger}\phi_2$ or in general 2HDM





- O Model-Independent analysis by ATLAS of $\gamma\gamma + X$ in 22 SRs
- O Excesses observed in some SRs
- O Hints for associated production of Neutral Higgs Boson
- ^O Explanation possible in Δ SM and 2HDM
- O Large Br($H \rightarrow \gamma \gamma$) possible in general 2HDM

Thank you for your attention!







- O No direct coupling of Δ with fermions
- Scalar potential

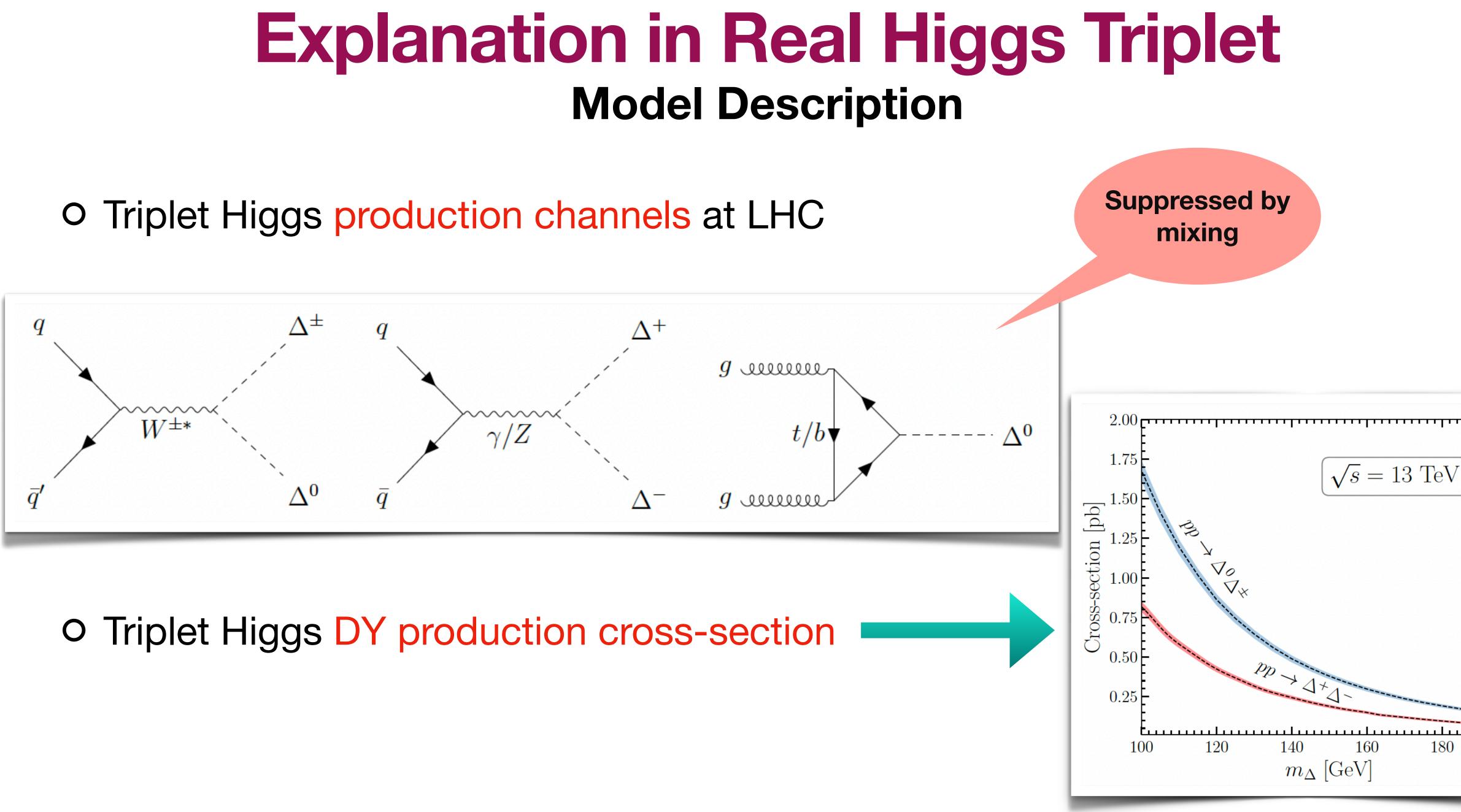
$$V = -\mu_{\phi}^2 \phi^{\dagger} \phi + \frac{\lambda_{\phi}}{4} (\phi^{\dagger} \phi)^2 - \mu_{\Delta}^2 \mathsf{T}$$

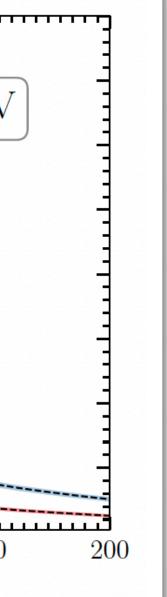
$$+A\phi^{\dagger}\Delta\phi + \lambda_{\phi\Delta}\phi^{\dagger}\phi \operatorname{Tr}(\Delta^{\dagger}\Delta)$$

where ϕ is the SM doublet.

Explanation in Real Higgs Triplet Model Description o Extend SM with Y = 0, $SU(2)_L$ triplet: $\Delta = \frac{1}{2} \begin{pmatrix} v_\Delta + h_\Delta^0 & \sqrt{2}h_\Delta^+ \\ \sqrt{2}h_\Delta^- & -v_\Delta - h_\Delta^0 \end{pmatrix}$ Suppressed ggH $\operatorname{Fr}(\Delta^{\dagger}\Delta) + \frac{\lambda}{\Lambda} [\operatorname{Tr}(\Delta^{\dagger}\Delta)]^2$

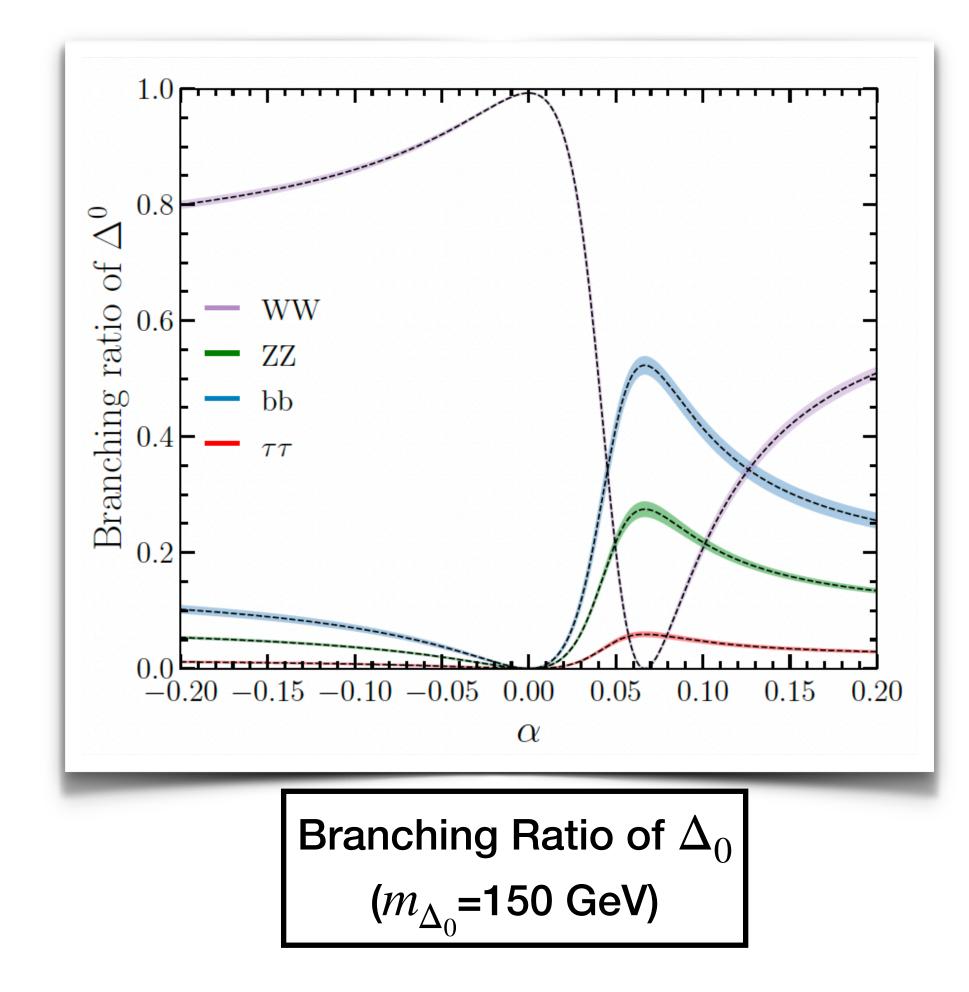
20

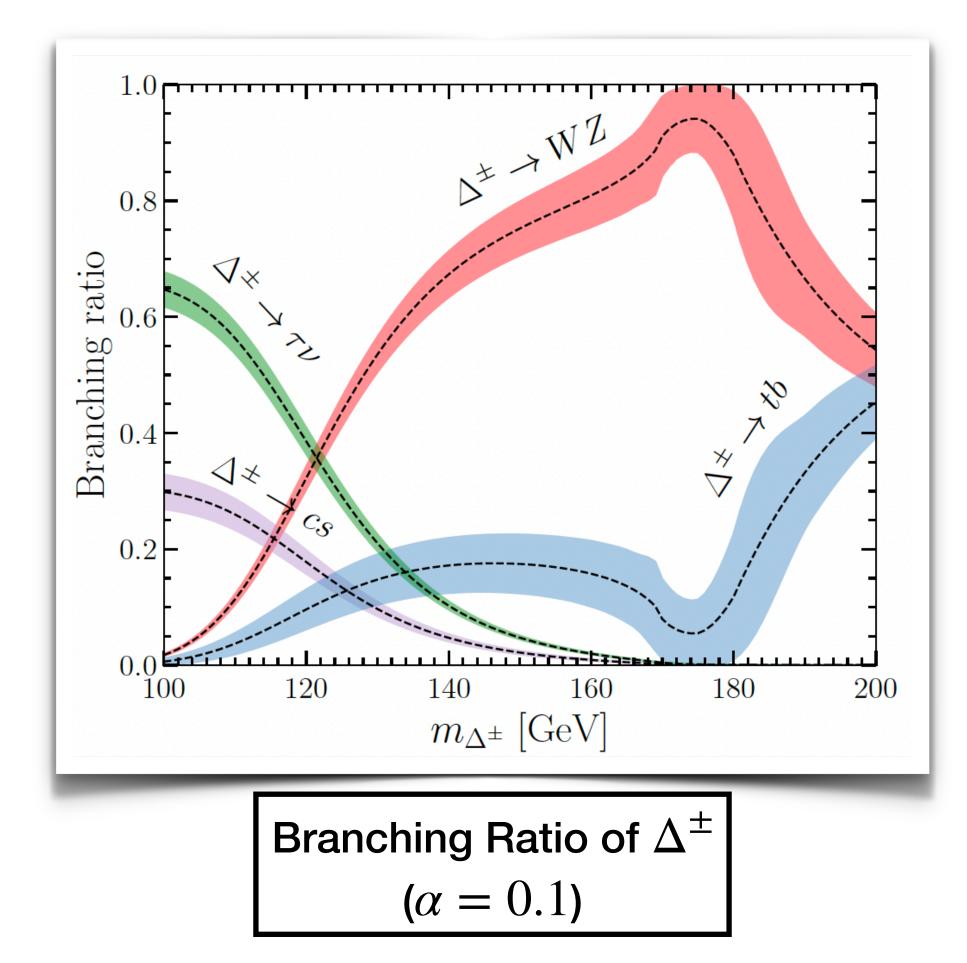




Explanation in Real Higgs Triplet Model Description

O Dominant Triplet Higgses decay channels





Real Higgs Triplet Basis Transformation

Physical to Lagrangian Basis

$$\begin{split} m_h^2 &= \frac{\lambda_\Phi v_\Phi^2}{2} + \tan \alpha \left(\lambda_{\Phi\Delta} v_\Delta - \frac{A}{2} \right) v_\Phi, \\ m_{\Delta^0}^2 &= \frac{\lambda_\Delta v_\Delta^2}{2} + \frac{A v_\Phi^2}{4 v_\Delta} - \tan \alpha \left(\lambda_{\Phi\Delta} v_\Delta - \frac{A}{2} \right) v_\Phi, \\ m_{\Delta^\pm}^2 &= A \frac{v_\Phi^2 + 4 v_\Delta^2}{4 v_\Delta}, \end{split}$$

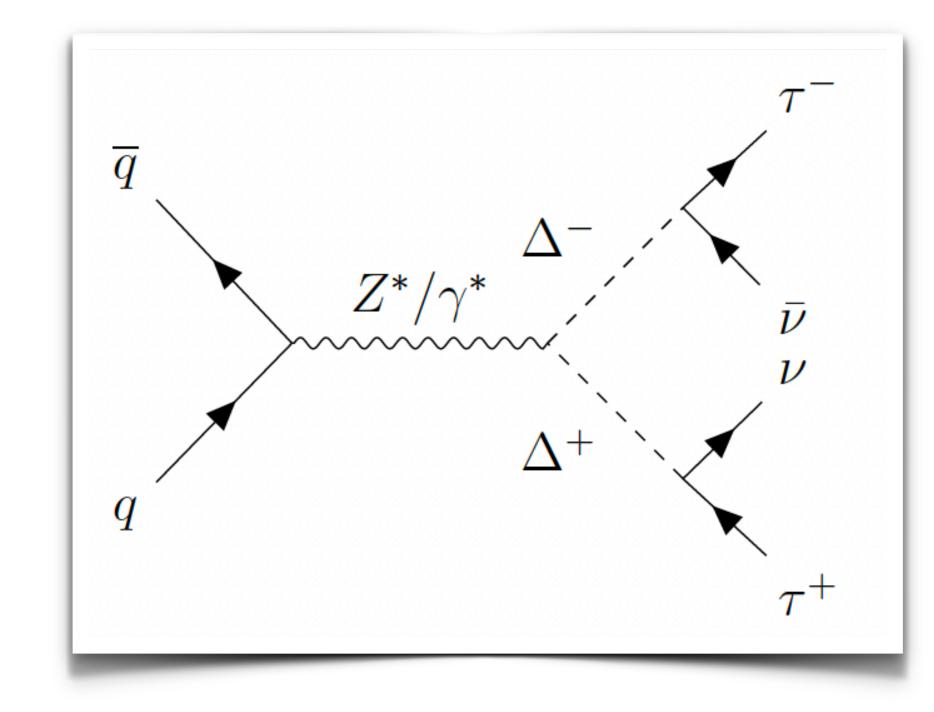
Lagrangian to Physical Basis

$$\begin{split} \lambda_{\Phi} &= \frac{2m_h^2}{v^2} ,\\ \lambda_{\Delta} &= \frac{2}{v_{\Delta}^2} \left[m_{\Delta^0}^2 - m_{\Delta^{\pm}}^2 \right] ,\\ \lambda_{\Phi\Delta} &= \frac{\alpha}{vv_{\Delta}} \left(m_{\Delta^0}^2 - m_{\Delta^{\pm}}^2 \right) + \frac{2}{v^2} m_{\Delta^{\pm}}^2 ,\\ A &= \frac{4v_{\Delta}}{v^2} m_{\Delta^{\pm}}^2 . \end{split}$$

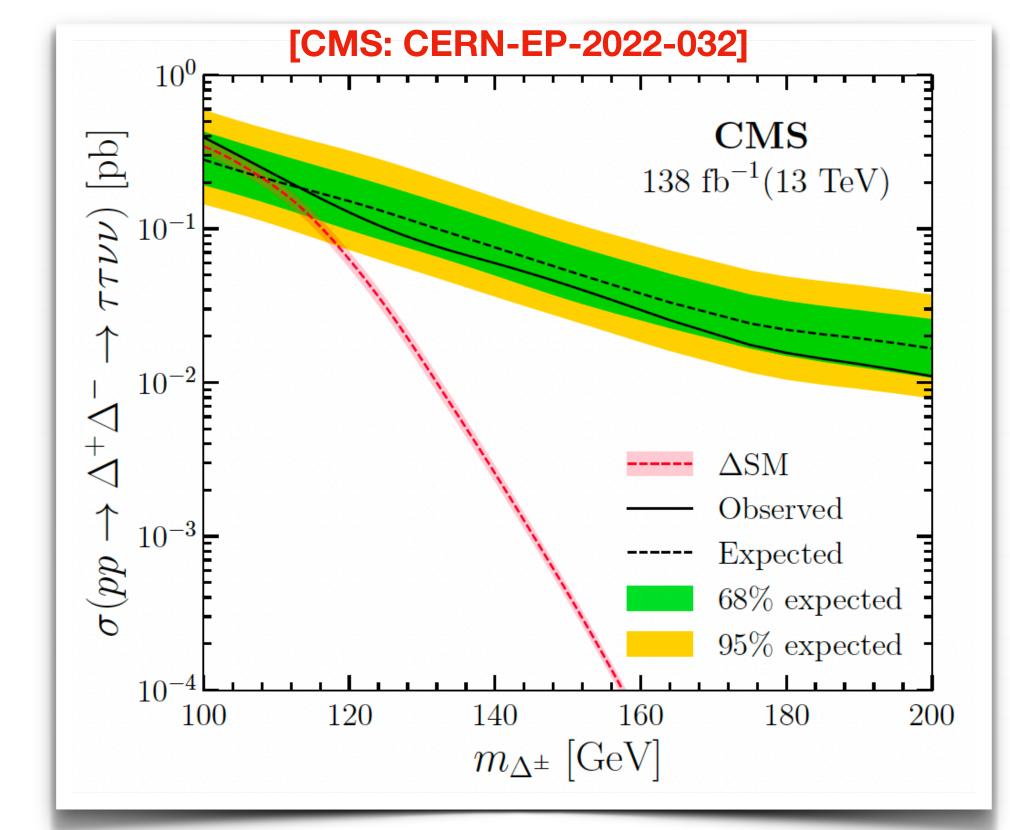


Explanation in Real Higgs Triplet Model Constraints

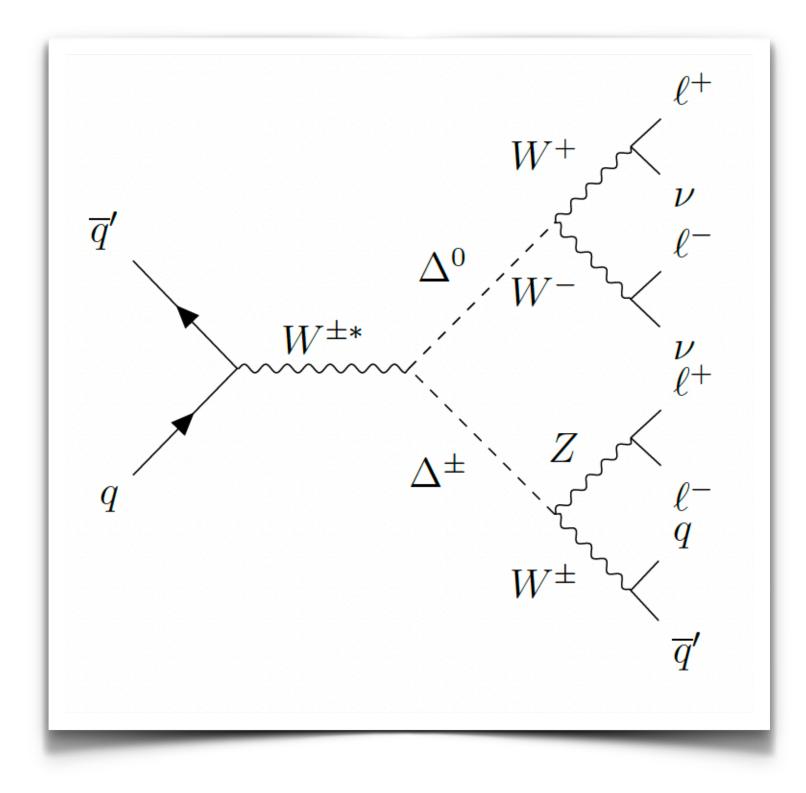
O Mass of Δ^{\pm} constrained from stau-like searches



o $m_{\Lambda^{\pm}} < 110$ GeV excluded at 95 % confidence level.

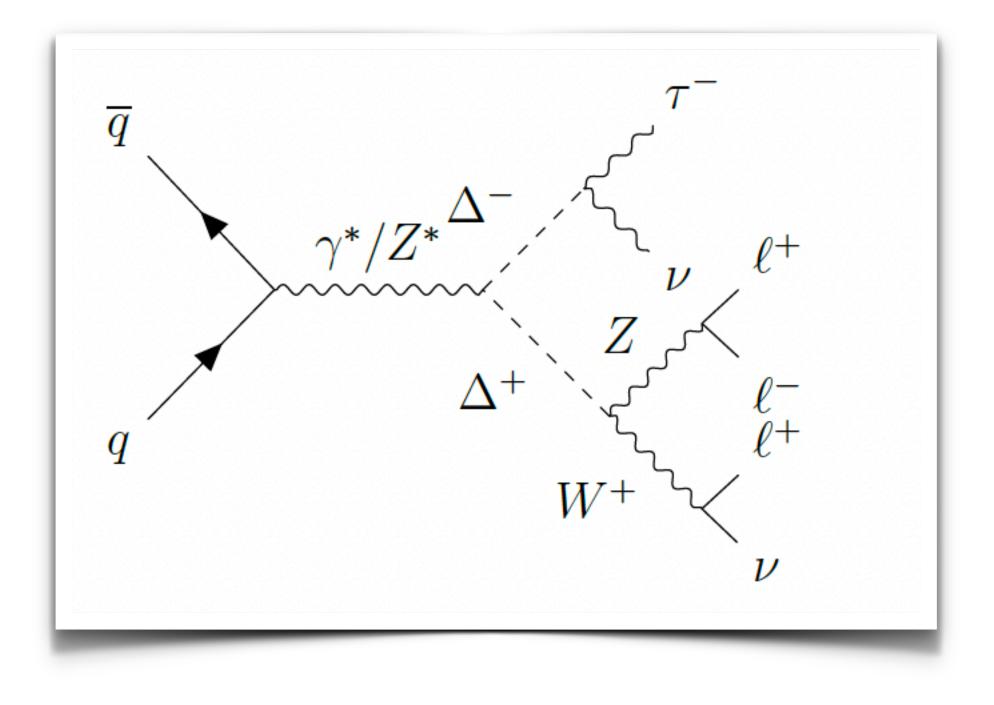


Explanation in Real Higgs Triplet Model Constraints



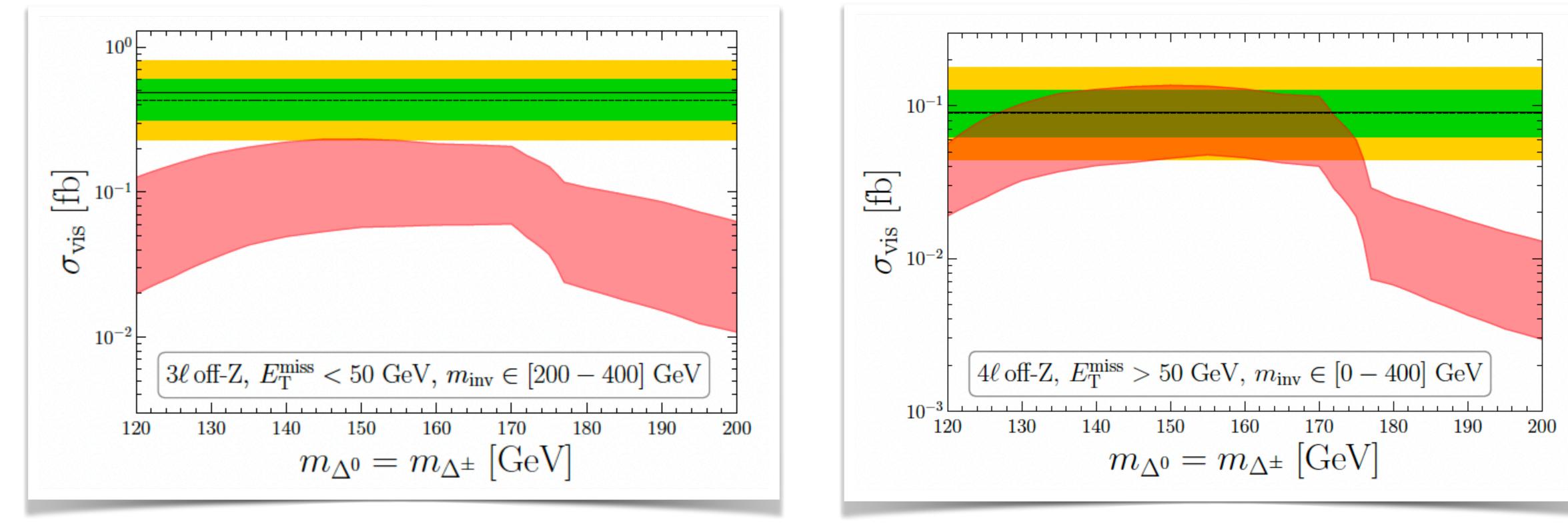
• ATLAS provides upper limit on visible cross-section for 22 SRs [CMS: CERN-EP-EP-2021-063]

O Triplet Higgs produces multiple lepton final states searched by ATLAS & CMS



Explanation in Real Higgs Triplet Model Constraints

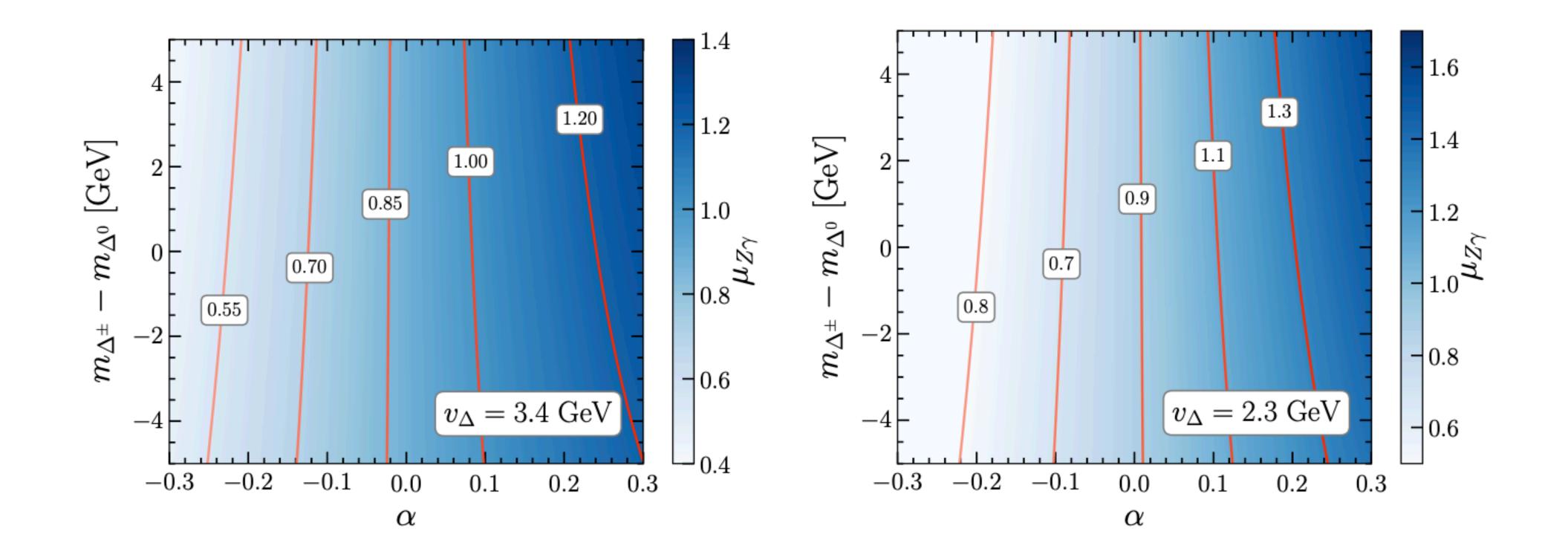
O Within 95 % CL upper limits of ATLAS



Simulated $pp \rightarrow \Delta^0 \Delta^{\pm}$ and $pp \rightarrow \Delta^{\mp} \Delta^{\pm}$ Ο

O Upper band obtained for $Br(\Delta^0 \rightarrow WW) = 100\%$

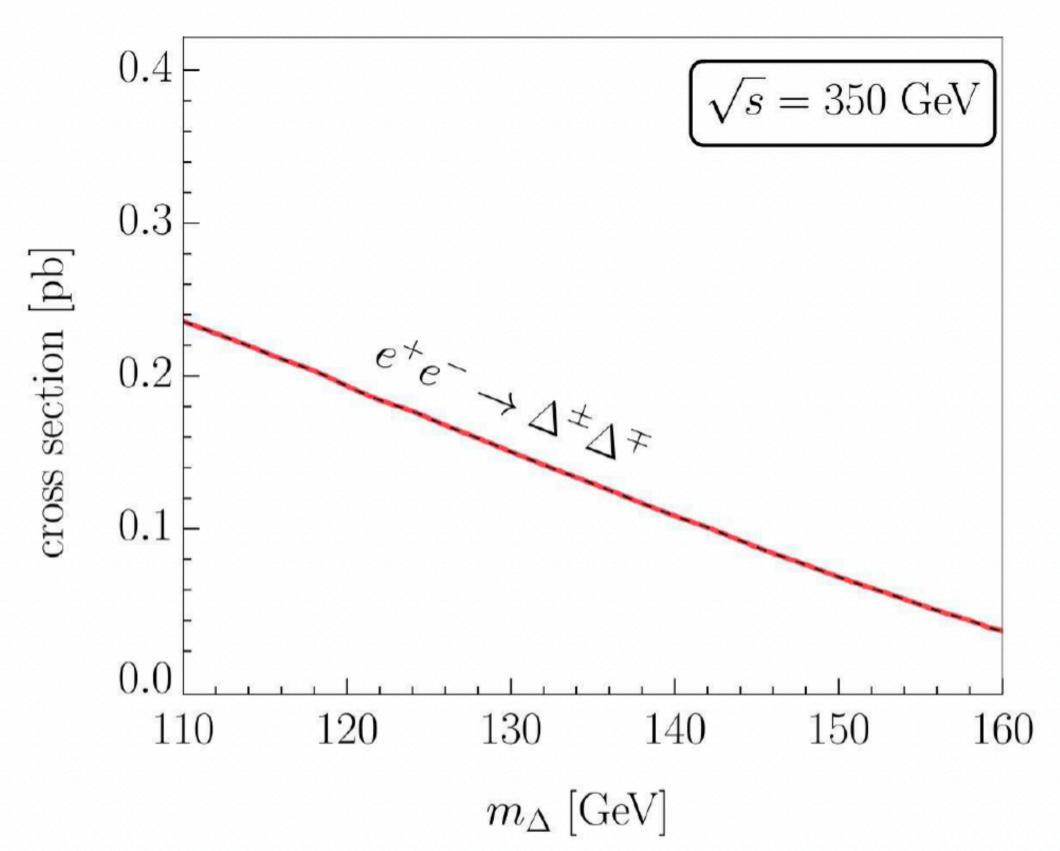


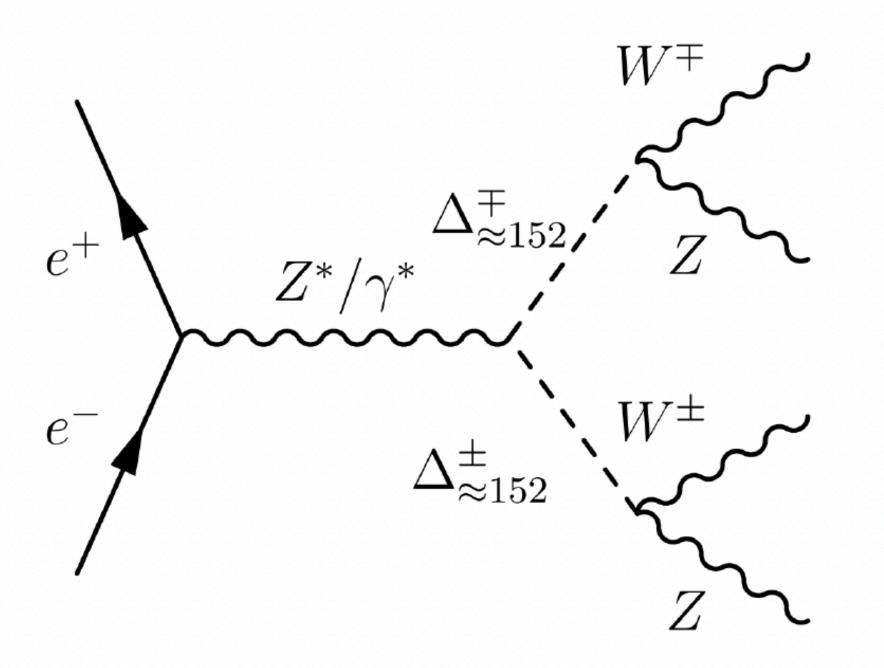




Real Higgs Triplet Prospects FCC-ee

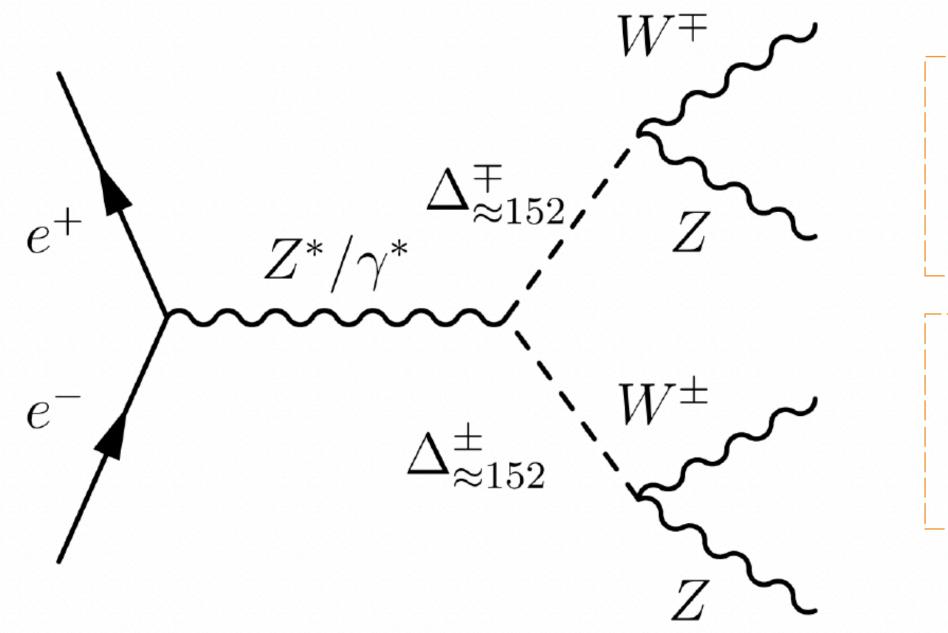
- Only Z^*/γ^* s-channel
- Suppressed $\Delta^0 \Delta^0$ production for a real triplet
- Pair production of the charged components





Real Higgs Triplet Prospects FCC-ee

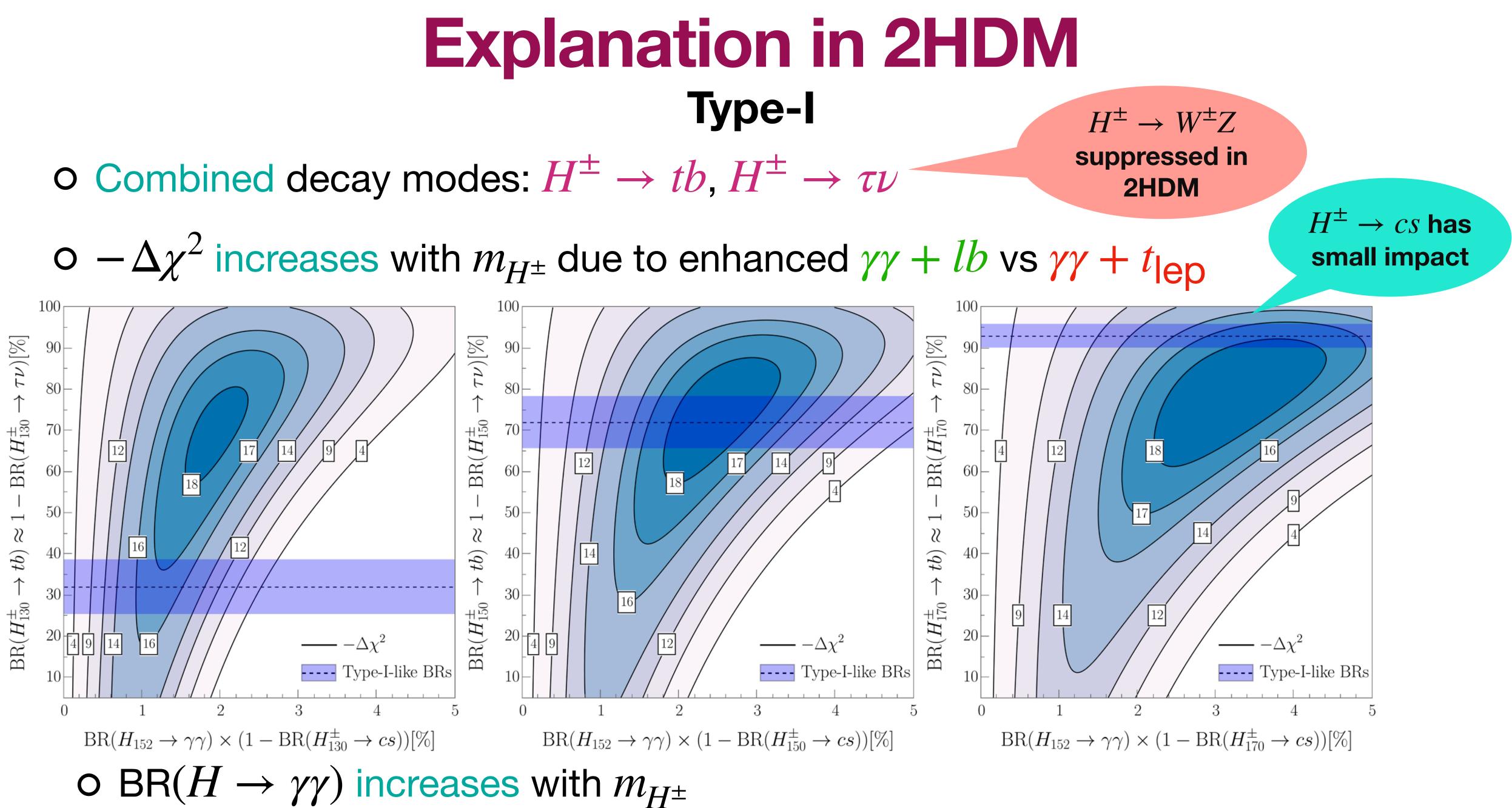
• The decay $\Delta^{\pm} \to W^{\pm}Z$ leads to a 6 ℓ (+ MET) signature



- Log-Likely-hood ratio yields
- $\sigma(e^+e^- \rightarrow \Delta^{\pm} \Delta^{\mp})$ could be measured at $\approx 9\sigma$

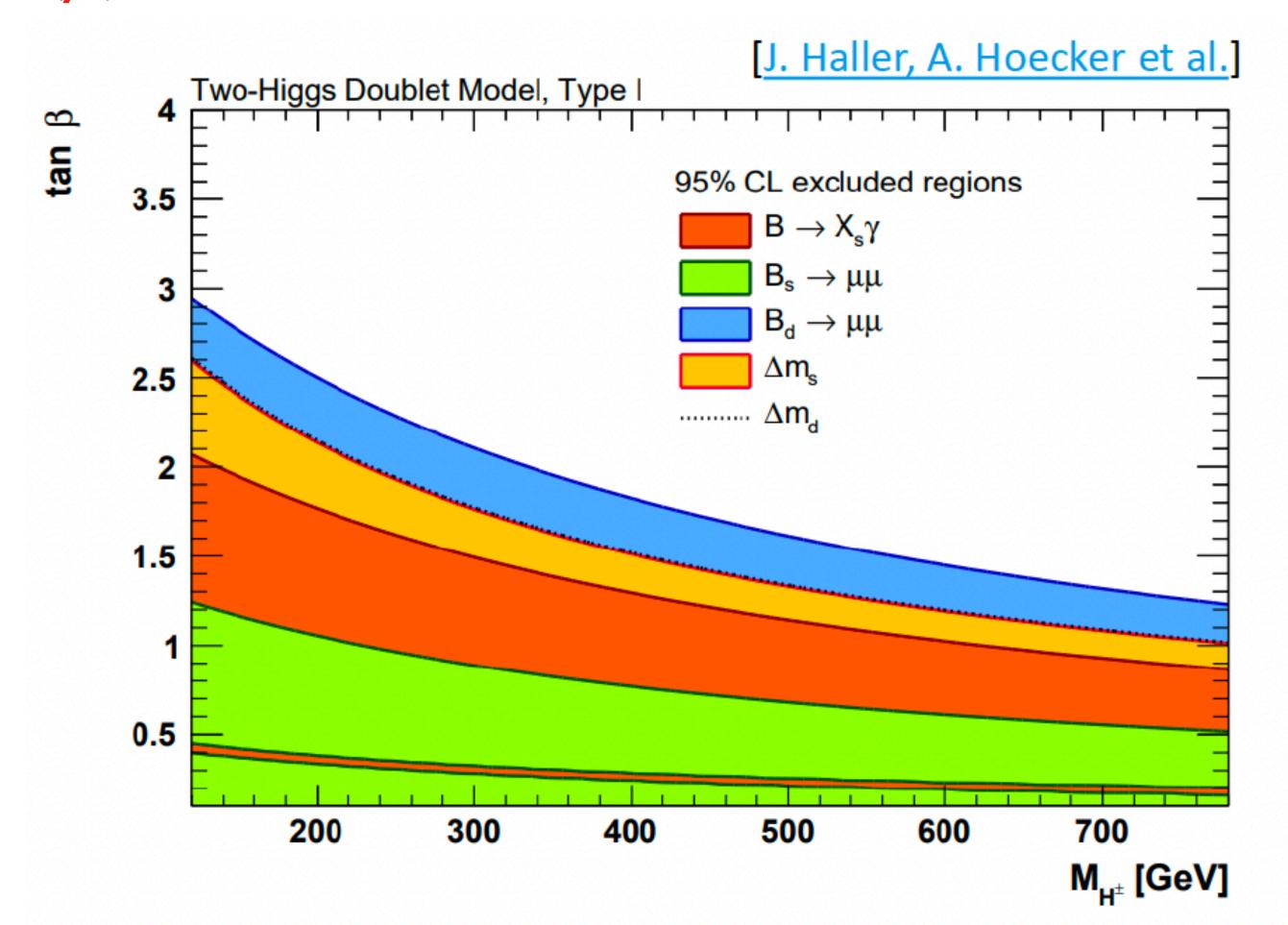
Events expected in the
$$\Delta$$
SM model
 $e^+e^- \rightarrow \Delta^{\pm} \Delta^{\mp} \rightarrow 6\ell + MET \approx 46$
Events expected in the SM model
 $e^+e^- \rightarrow 6\ell(+MET) \approx 1$

s
$$\chi^2 \approx 80$$



Explanation in 2HDM Explanation of $\gamma\gamma + X$ Excesses

• Bounds on $tan(\beta)$



Explanation in 2HDM FCNC & CP-Violation

- O General 2HDM may lead to FCNC at tree-level
- O Avoided in flavour aligned 2HDM

$$Y = -\bar{Q}_L Y_d (\phi_2 + \zeta_d \phi_1) d_R - \bar{Q}_L$$

O Complex parameters leads to CP-violation

Yukawa Sector: $\zeta_{\mu}, \zeta_{d}, \zeta_{l}$

We take them real

$_{L}Y_{u}(\phi_{2}^{c}+\zeta_{u}^{*}\phi_{1}^{c})u_{R}-\bar{L}_{L}Y_{l}(\phi_{2}+\zeta_{l}\phi_{1})e_{R}$

Higgs Sector: $\lambda_5, \lambda_6, \lambda_7$

Explanation in 2HDM EDM Constraints

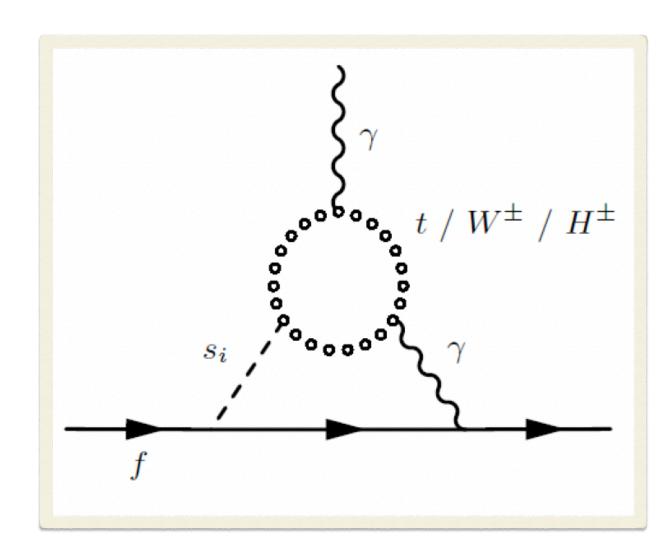
o $Im(\lambda_6)$ drives $Br(A \rightarrow \gamma \gamma)$

O Correlate with EDM constraints

O Transform Lagrangian to Higgs Basis

$$m_{11}^2, m_{22}^2, m_{12}^2$$
$$\lambda_1, \cdots, \lambda_7$$
$$\langle \phi_1 \rangle = v_1, \langle \phi_2 \rangle = v_2$$

O Used analytic expressions of [arXiv: 2009.01258]



 $\supset id_f \bar{u}\sigma^{\mu\nu}q_{\nu}\gamma_5 u$

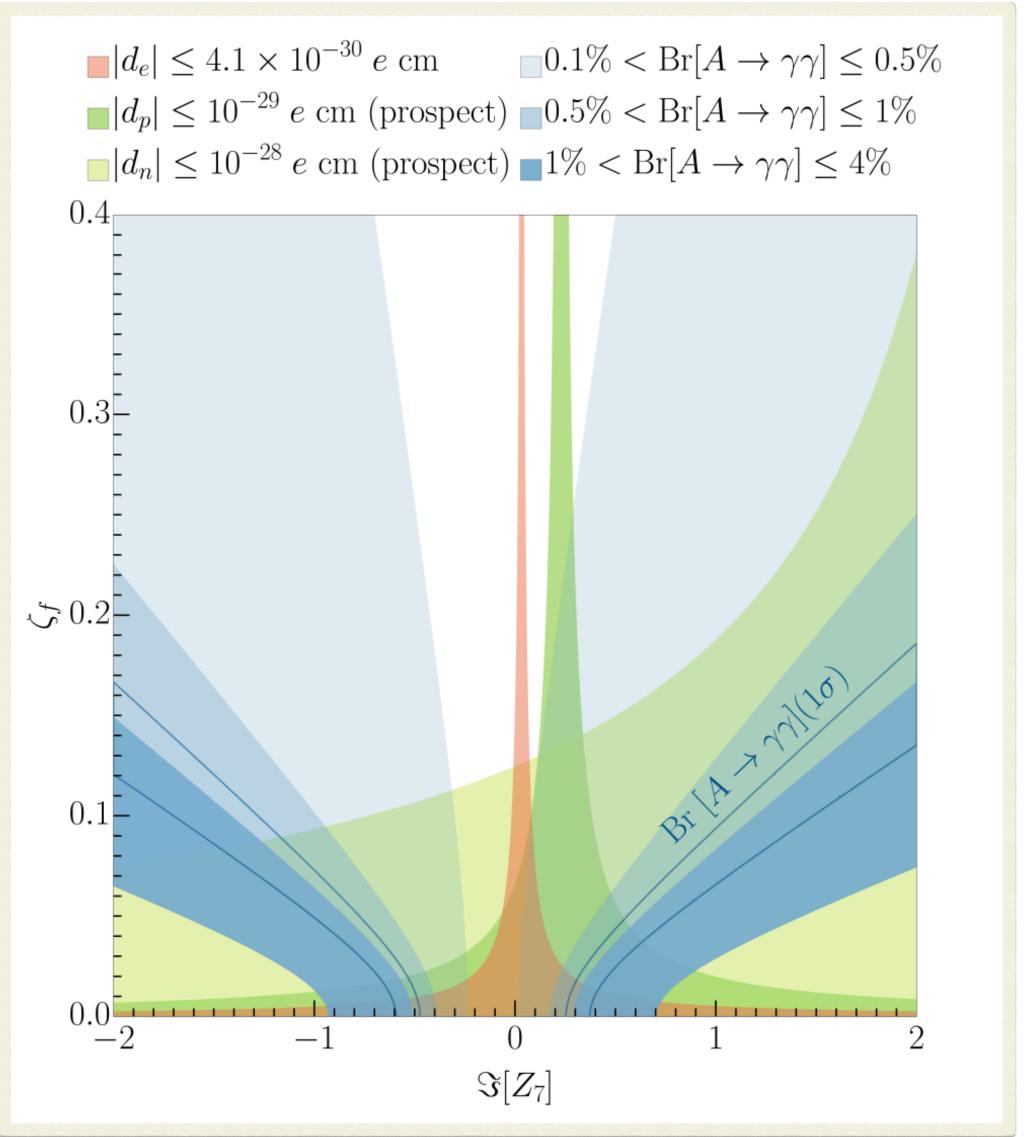


$$Y_1, Y_2, Y_3$$
$$Z_1, \cdots, Z_7$$
$$\langle \phi_1 \rangle = v, \langle \phi_2 \rangle = 0$$



Explanation in 2HDM EDM Constraints

- O eEDM gives stringent bounds: $10^{-30}e \, cm^{-1}$ [arXiv:2212.11841]
- Projection for nEDM and pEDM considered
 - nEDM $\leq 10^{-28} e \, cm^{-1}$; pEDM $\leq 10^{-29} e \, cm^{-1}$ [EPJ Web Conf. 262 (2022) 01015] [arXiv:2007.10332]
- Benchmark Point:
 - $m_H = 200 \text{ GeV}, m_{H^+} = 130 \text{ GeV}, m_A = 152 \text{ GeV}$ $Z_2 = -Z_3 = 0.2, \operatorname{Re}(Z_7) = 0.1, \theta_{12} = 0.001$ $\theta_{13} = \theta_{23} = 0.01, \zeta_l = \zeta_u = \zeta_d = \zeta_f$

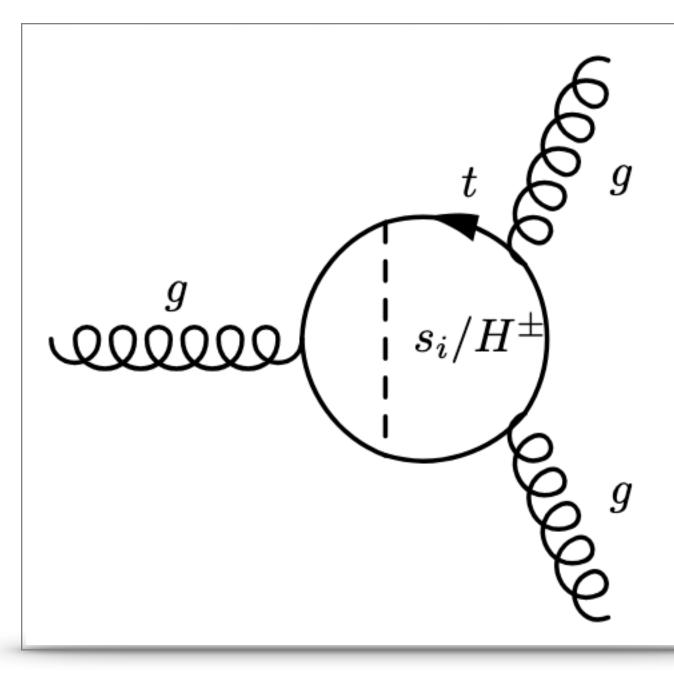


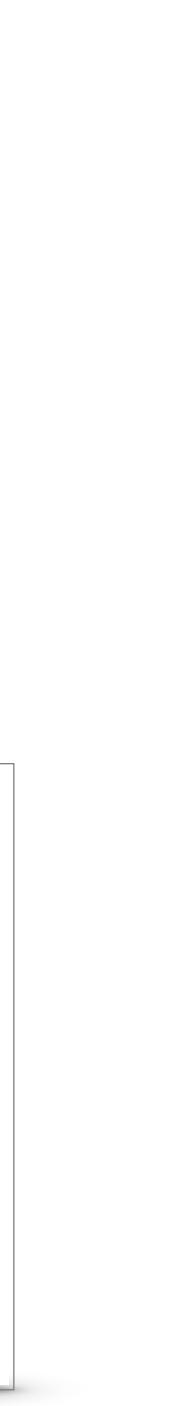
Explanation in 2HDM nEDM

o nEDM is expressed as

- $d_n = -(0.20 \pm 0.01)d_u + (0.78 \pm 0.03)d_d (0.55 \pm 0.28)e\tilde{d}_u$ $-(1.1 \pm 0.55) e \tilde{d}_d + (50 \pm 40) \text{ MeV } e \tilde{d}_G$
- d_a is the quark EDM and \tilde{d}_a is the chromo EDM

 $\circ d_G$ contribution from Weinberg operator

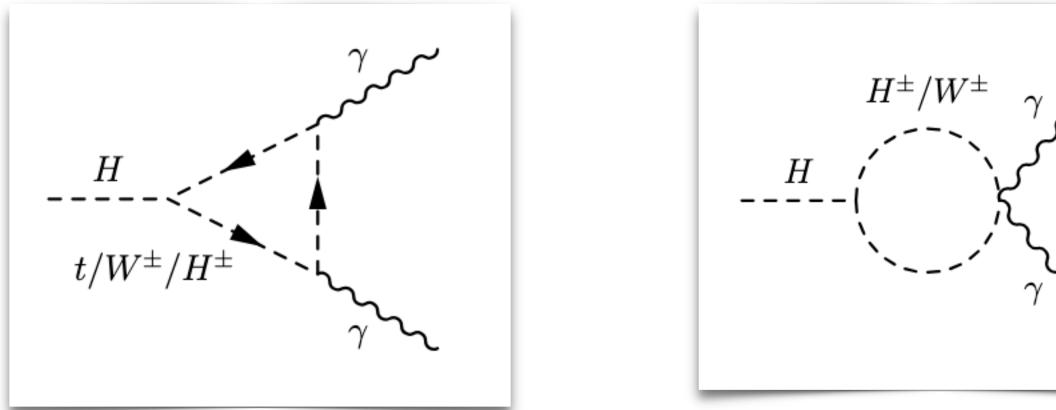




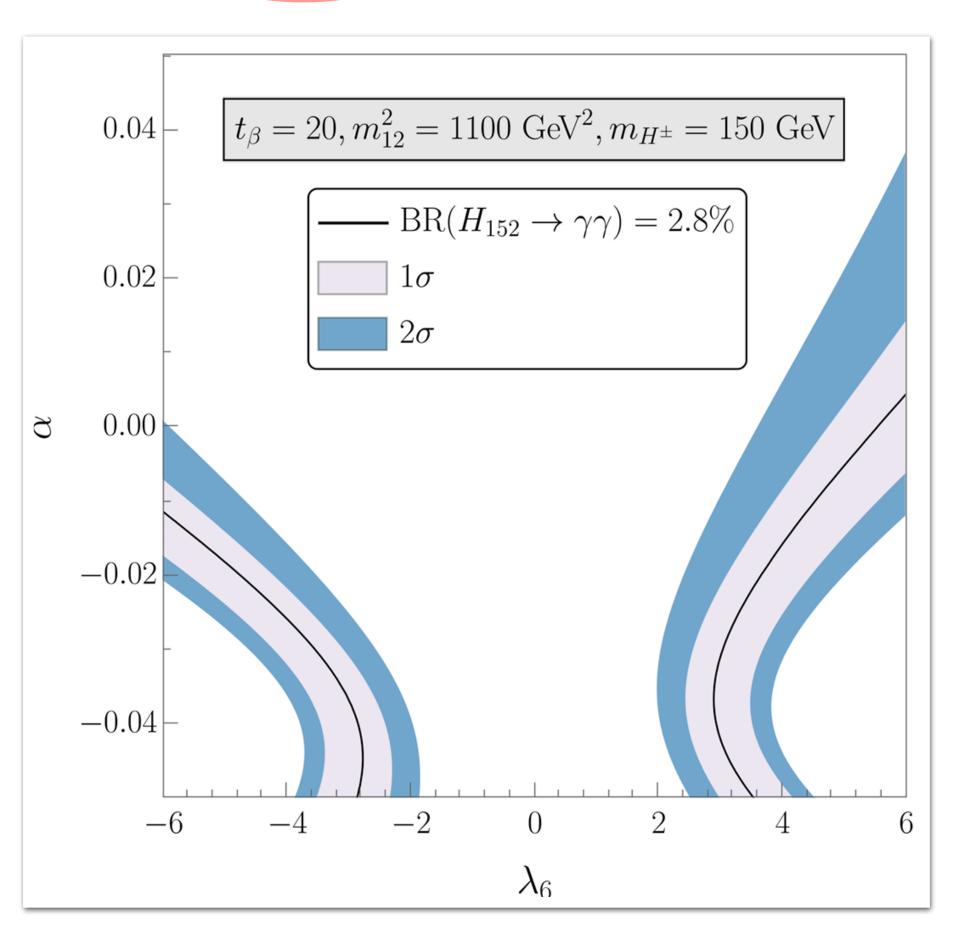
General 2HDM Large $H \rightarrow \gamma \gamma$

O Large $Br(H \rightarrow \gamma \gamma)$ possible in general 2HDM

- $\mathscr{L} \in -\lambda_6 H_1^{\dagger} H_1 H_2^{\dagger} H_1 + \text{h.c.},$
- Modifies the $HH^{\pm}H^{\mp}$ vertex
- O Enhanced $Br(H \rightarrow \gamma \gamma)$ via H^{\pm} loop

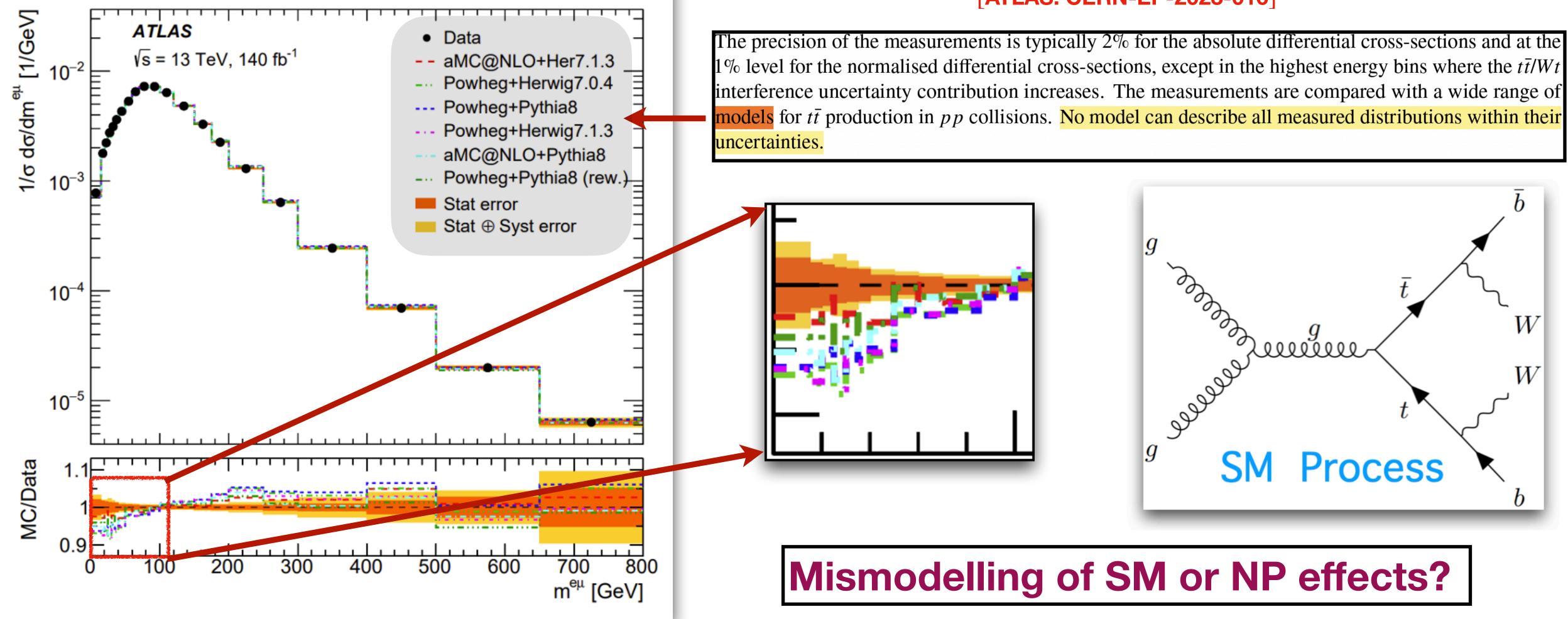


 Z_2 symmetry broken



Deviations in *tt* **Differential cross-section** No match with SM

[ATLAS: CERN-EP-2023-016]



[ATLAS: CERN-EP-2023-016]

- Simplified model with three Higgs bosons
- ^O Preferred over SM by atleast 5.8 σ
- Compatible with 95 GeV and 152 GeV Excesses

