

$H \rightarrow \ell^+ \ell^- Z$ at NLO in the SMEFT

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See also related talk on $e^+ e^- \rightarrow ZH$ by Konstantin Asteriadis

Standard Model Effective Field Theory

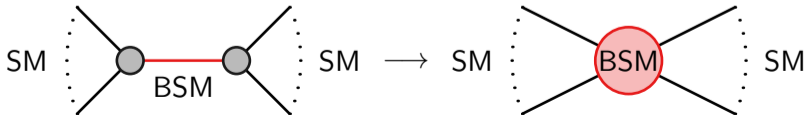
For heavy BSM physics at scale Λ , effects in low energy observables can be computed using effective field theory techniques:

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_i \frac{C_i}{\Lambda^2} \mathcal{O}_i + \dots$$

Systematically improvable, both in loops and powers of Λ

Comprehensive framework for constraining heavy BSM models with LHC and low-energy data

For precise constraints, need to go to higher loop orders



$$H \rightarrow 4\ell$$

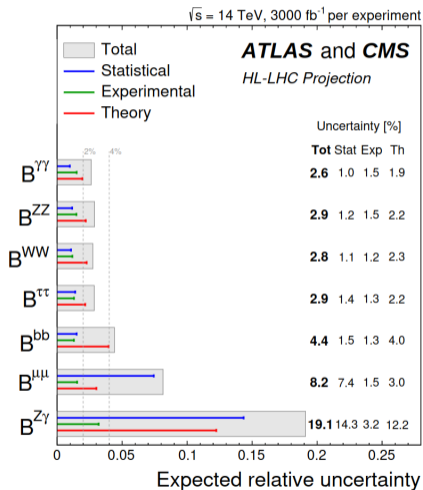
$H \rightarrow 4\ell$ is one of the best measured decay modes!

SM known to NLO EW [hep-ph/0604011], [1912.02010]

Most Higgs decays are known at one-loop in the SMEFT, but $H \rightarrow 4f$ still missing

As precision improves, Higgs decays become sensitive to loop effects

Necessary ingredient for NLO accurate fits



HL-LHC working group [1902.00134]

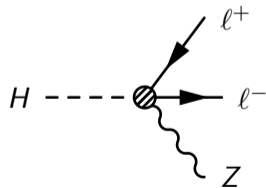
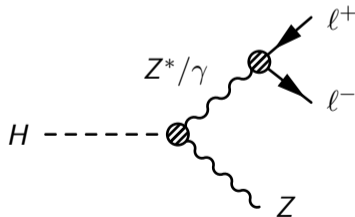
$H \rightarrow \ell^+ \ell^- Z$ in the SMEFT

We have the complete one-loop calculation of $H \rightarrow \ell^+ \ell^- Z$ in the SMEFT at dimension-6 with fully general flavour structure

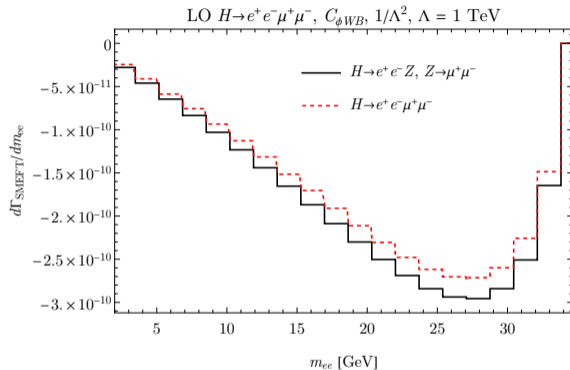
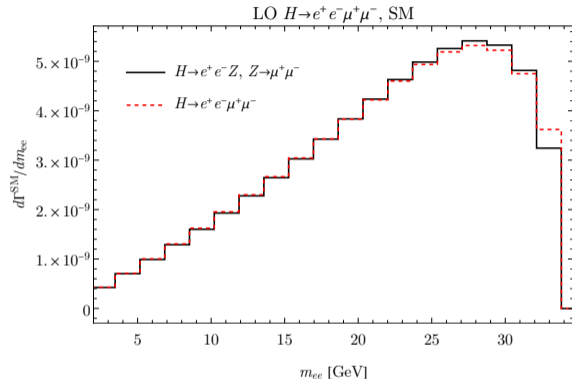
With the narrow width approximation, dominant contributions to $H \rightarrow 4\ell$

At LO ~ 10 operators contribute, with new kinematic dependence

- $\mathcal{O}_{\phi B} = \phi^\dagger \phi B_{\mu\nu} B^{\mu\nu}$
- $\mathcal{O}_{\phi \square} = (\phi^\dagger \phi) \square (\phi^\dagger \phi)$
- $\mathcal{O}_{\phi D} = |\phi^\dagger D^\mu \phi|^2$
- $\mathcal{O}_{\phi e} = (\phi^\dagger i \overleftrightarrow{D}_\mu \phi) (\bar{e}_R \gamma^\mu e_R)$
- ...



How good is the narrow width approximation?



Examples for the SM and an example EFT coefficient ($C_{\phi_{WB}}$)

Other operators similar – quite good agreement for $H \rightarrow e^+ e^- \mu^+ \mu^-$

We don't have full NLO $H \rightarrow 4\ell$, but we do in the NWA!

Going to NLO: real emission

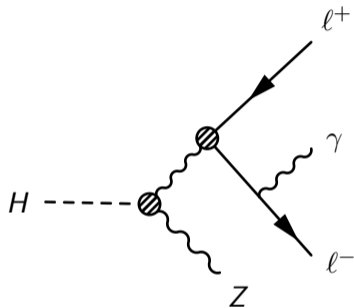
At NLO, virtual photon contributions have soft and collinear divergences that cancel against $H \rightarrow \ell^+ \ell^- Z \gamma$ contributions

Treat with standard dipole subtraction techniques

Requires 4-body phase space!

For non-inclusive observables, logarithms $\sim \log Q^2/m_\ell^2$ appear

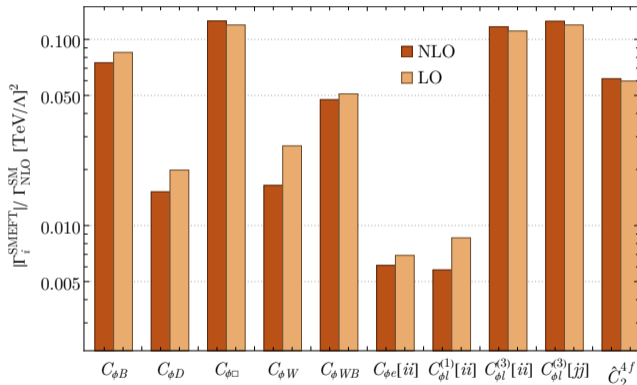
$\rightarrow e^+e^-$ and $\mu^+\mu^-$ modes differ at NLO after experimental cuts



NLO corrections are large

Up to a 40% correction for some operators at NLO

(Results shown with $\Lambda = 1$ TeV, $C_i = 1$)



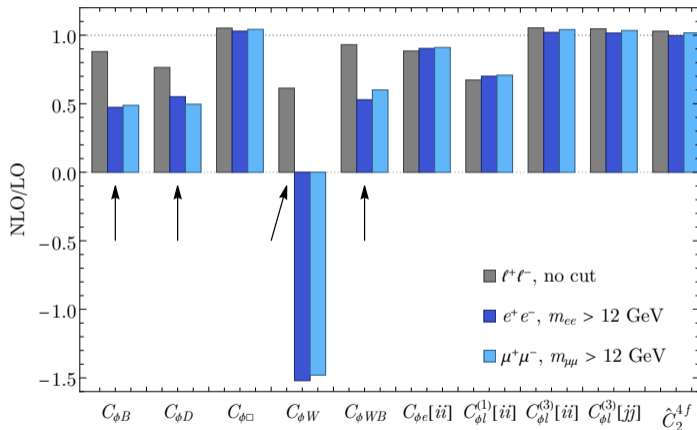
Realistic experimental cuts

Becomes even larger ($> 50\%$) with realistic $m_{\ell\ell} > 12$ GeV cut

Significant differences for $C_{\phi B}$, $C_{\phi D}$, $C_{\phi W}$, $C_{\phi WB}$

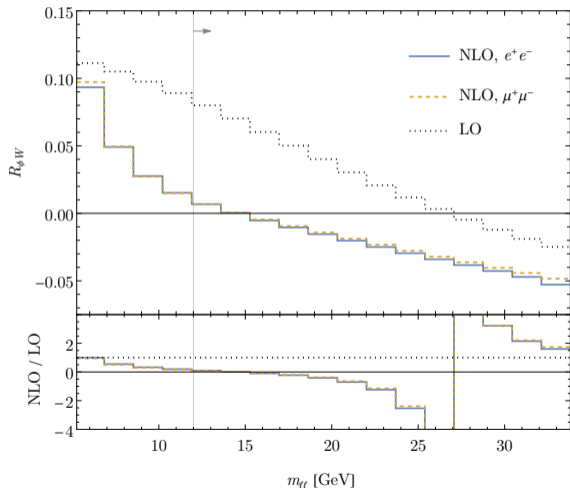
$C_{\phi W}$ switches sign!

Now mildly flavour dependent: up to $\sim 10\%$ differences

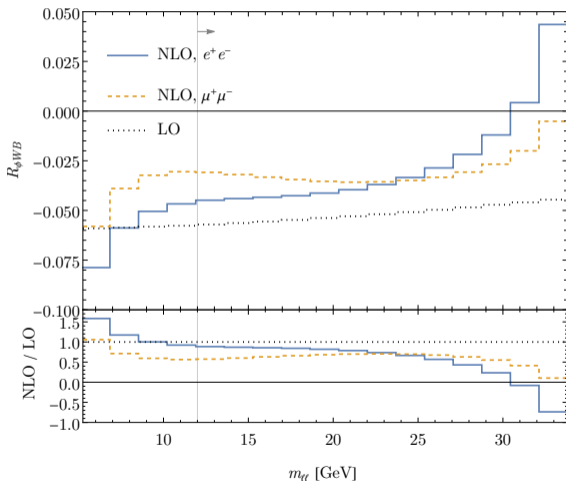


Differential distributions

$C_{\phi W}, 1/\Lambda^2, \Lambda = 1 \text{ TeV}$



$C_{\phi WB}, 1/\Lambda^2, \Lambda = 1 \text{ TeV}$



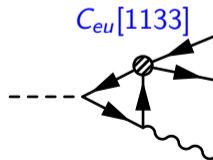
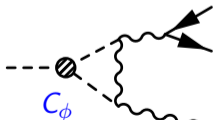
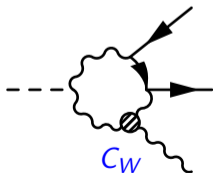
$$R_i = \text{SMEFT} / \text{SM}$$

New operators at NLO

About ~ 70 new operators first enter at NLO in ~ 25 combinations

Of these, most are quite small. Notable exceptions:

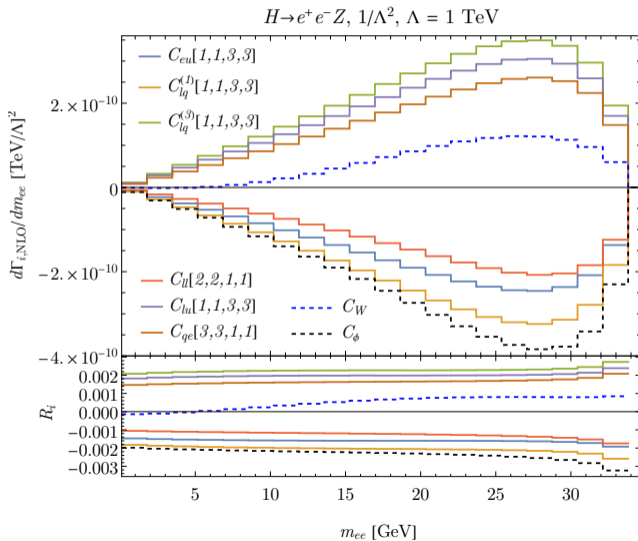
- Top-quark operators: $\mathcal{O}_{lq}^{(1)} = (\bar{L}\gamma_\mu L)(\bar{Q}\gamma^\mu Q)$, $\mathcal{O}_{\phi u} = (\phi^\dagger i \overleftrightarrow{D}_\mu \phi)(\bar{t}_R \gamma^\mu t_R)$, ...
- Higgs self-coupling: $\mathcal{O}_\phi = (\phi^\dagger \phi)^3$
- Anomalous triple gauge coupling: $\mathcal{O}_W = \epsilon_{IJK} W_\mu^{I\nu} W_\nu^{J\alpha} W_\alpha^{K\mu}$



Differential distributions

New operators at NLO have SM-like distributions

Exception: C_W has novel momentum dependence



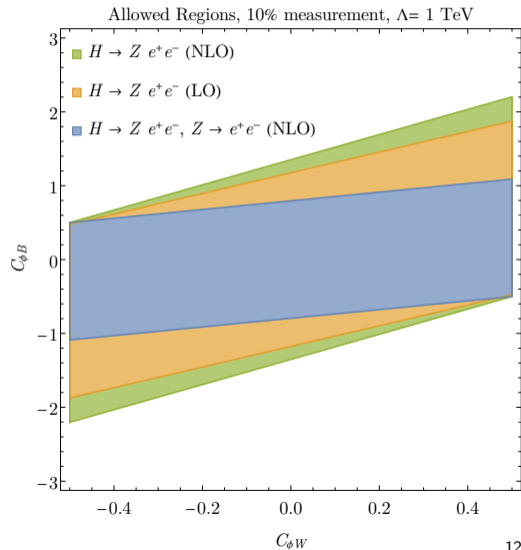
Estimating the impact on constraints

Toy example: consider a 10% measurement of $H \rightarrow 4\ell$

Combine with known $Z \rightarrow \ell^+\ell^-$ at NLO in the SMEFT [1909.02000]
→ Full $H \rightarrow 4\ell$ at NLO in NWA

Correlations change shape at NLO

Proper constraints require NLO accurate production modes as well in a general fit



Bounding C_ϕ

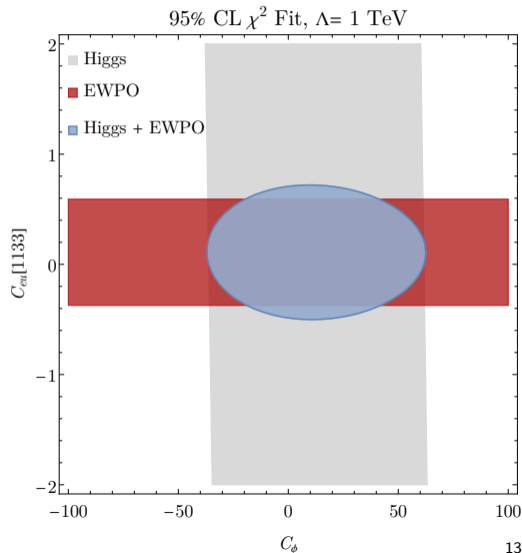
Higgs self-coupling constraints depend on other operators

An example: C_{eu} doesn't appear in the production mode

$$\mathcal{O}_{eu}[1133] = (\bar{e}_R \gamma_\mu e_R)(\bar{t}_R \gamma^\mu t_R)$$

Production mode C_ϕ dependence included

Would like to do this for all operators, but need to know production mode dependence



Conclusions

We have computed $H \rightarrow \ell^+ \ell^- Z$ to full NLO in the SMEFT

- One of the last remaining H decays to one-loop accuracy

Several operators appearing at LO experience large corrections at NLO

- In particular, $C_{\phi W}$, $C_{\phi B}$, $C_{\phi WB}$, $C_{\phi D}$, $C_{\phi I}^{(1)}$
- Experimental cut of $m_{\ell\ell} > 12$ GeV changes results significantly

Higgs self-coupling and anomalous top-quark couplings enter at NLO

- Extraction of Higgs trilinear depends on assumptions about other operators

Our calculation is one more step towards fully NLO accurate SMEFT fits

Thank you!