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Status of Triplet Extended Higgs Sector Models in Light of NLO Unitarity and Latest LHC Data

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Based on:

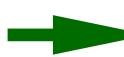
D. Chowdhury, P. Mondal, S.S.

arXiv: 2404.18996 [hep-ph]

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Beyond the Standard Model

Standard Model (SM) of particle physics is a framework



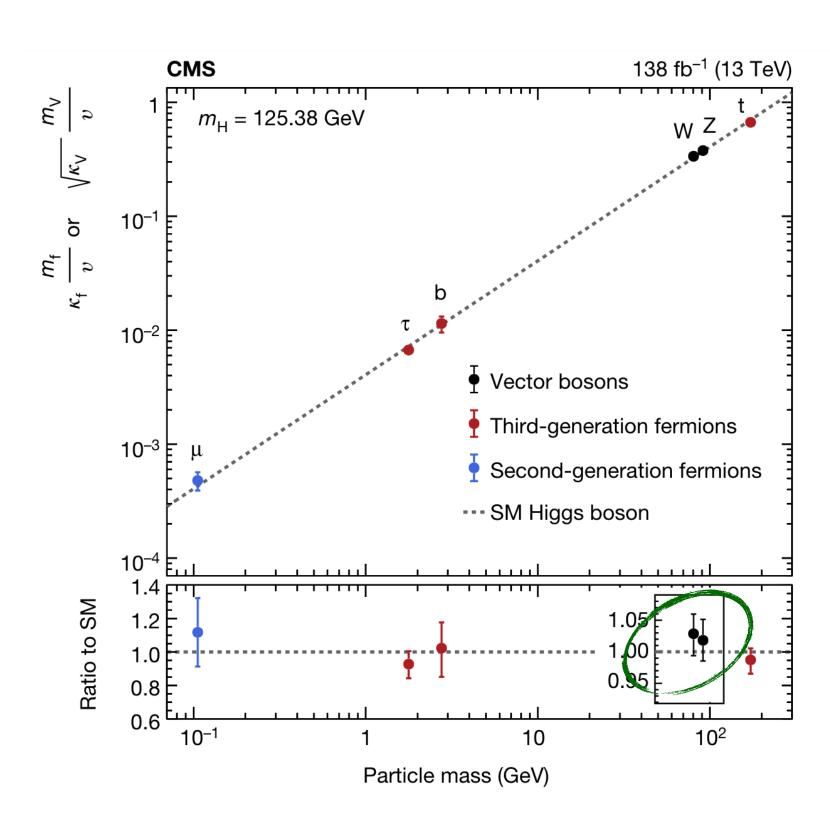
good agreement with collider data

Open questions:

Neutrino masses, Dark Matter, Baryon asymmetry of the Universe (BAU), Electroweak Phase Transitions (EWPT)

- SM might be a simplified version of a more complicated model
- Do the LHC data preclude the existence of additional multiplets in the scalar sector of the SM?

Search for BSM @ forefront of particle physics research



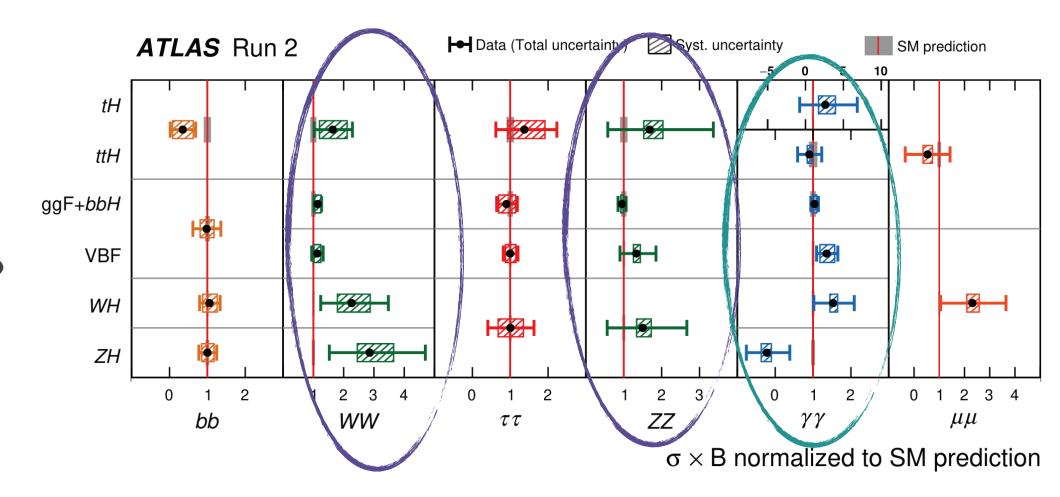
Triplet Extended Higgs Sector

ullet Both ATLAS and CMS suggest an enhanced rate of WW and $Z\!Z$ relative to the SM

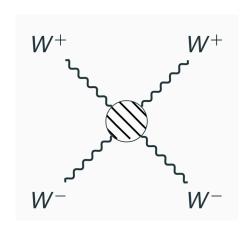
The questions we ask is:

If WW and ZZ rates are enhanced, how far beyond the SM one must to describe them?

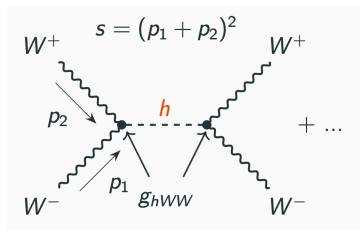
Additional charged Higgs?



Recap: WW scatterings are weakly interacting or strongly interacting?



Higgs portal restore unitarity



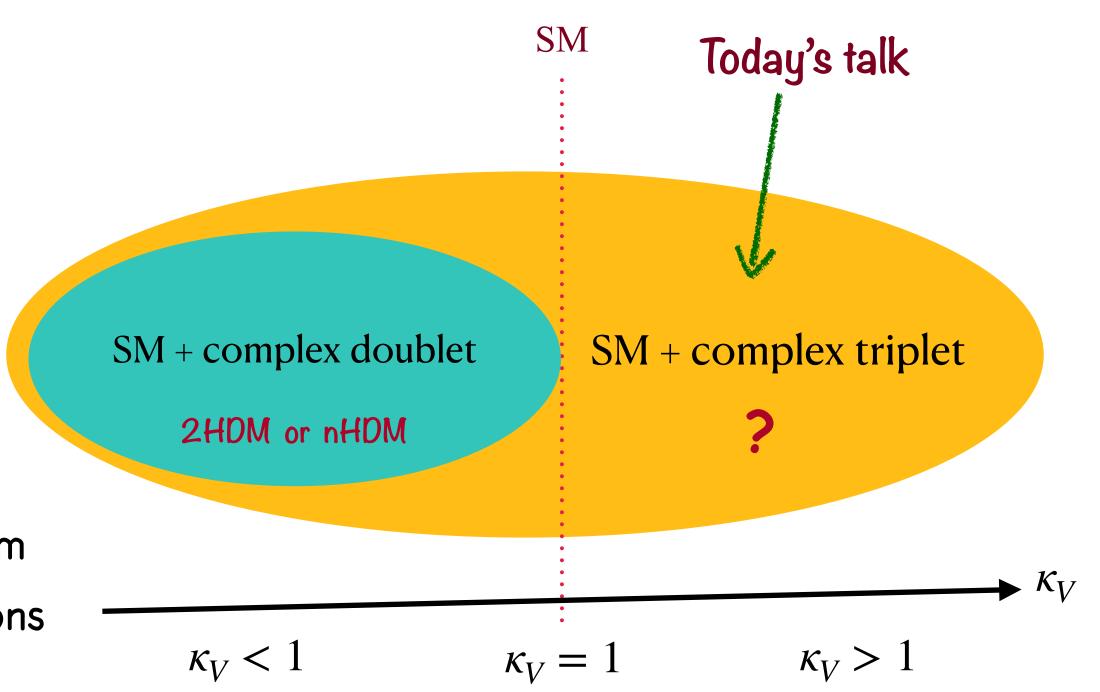
 $g_{hWW} > (g_{hWW})_{SM}$ $g_{hWW} = (g_{hWW})_{SM}$ $g_{hWW} = (g_{hWW})_{SM}$

Additional doubly charged Higgs boson plays an important role to restore unitarity at high energies

Triplet Extended Higgs Sector

- ullet Model with triplet extended Higgs sector can explain an enhanced rate of WW and ZZ
 - Doubly charged Higgs boson comes from $SU(2)_L$ triplet through EWSB mechanism

Triplet or higher multiplets leave an imprint on the EWSB mechanism that can be detected via SM-like Higgs couplings to the vector bosons



- ✓ It can also explain neutrino oscillation , EWBG , Dark-Matter puzzle
- It offers a much richer prospect for collider experiments

Can be probed in particle colliders and in cosmological observatories

Higgs Triplet Models with Custodial Symmetry

• Custodial Symmetry : $M_W^2 = M_Z^2 \cos^2 \theta_W$

Rho-parameter,
$$\rho = \frac{M_W^2}{M_Z^2 \cos^2 \theta_W}$$

At tree-level,
$$\rho=1$$

Triplet extended Higgs sector with Custodial Symmetry:

SM doublet ϕ (T=1/2, Y=1/2) + Real triplet ξ (T=1, Y=0) + Complex triplet χ (T=1, Y=1)

$$\langle \phi \rangle = v_{\phi}, \, \langle \xi \rangle = v_{\xi}, \, \langle \chi \rangle = v_{\chi}$$
 \longrightarrow $\rho = 1$ in tree-level if $(v_{\chi} = v_{\xi})$

* Georgi-Machacek (GM) model:

Equality of triplet VEVs is preserved by the Higgs potential [Georgi, Malinteractions among the Higgs fields maintained $SU(2)_L \times SU(2)_R$ symmetry

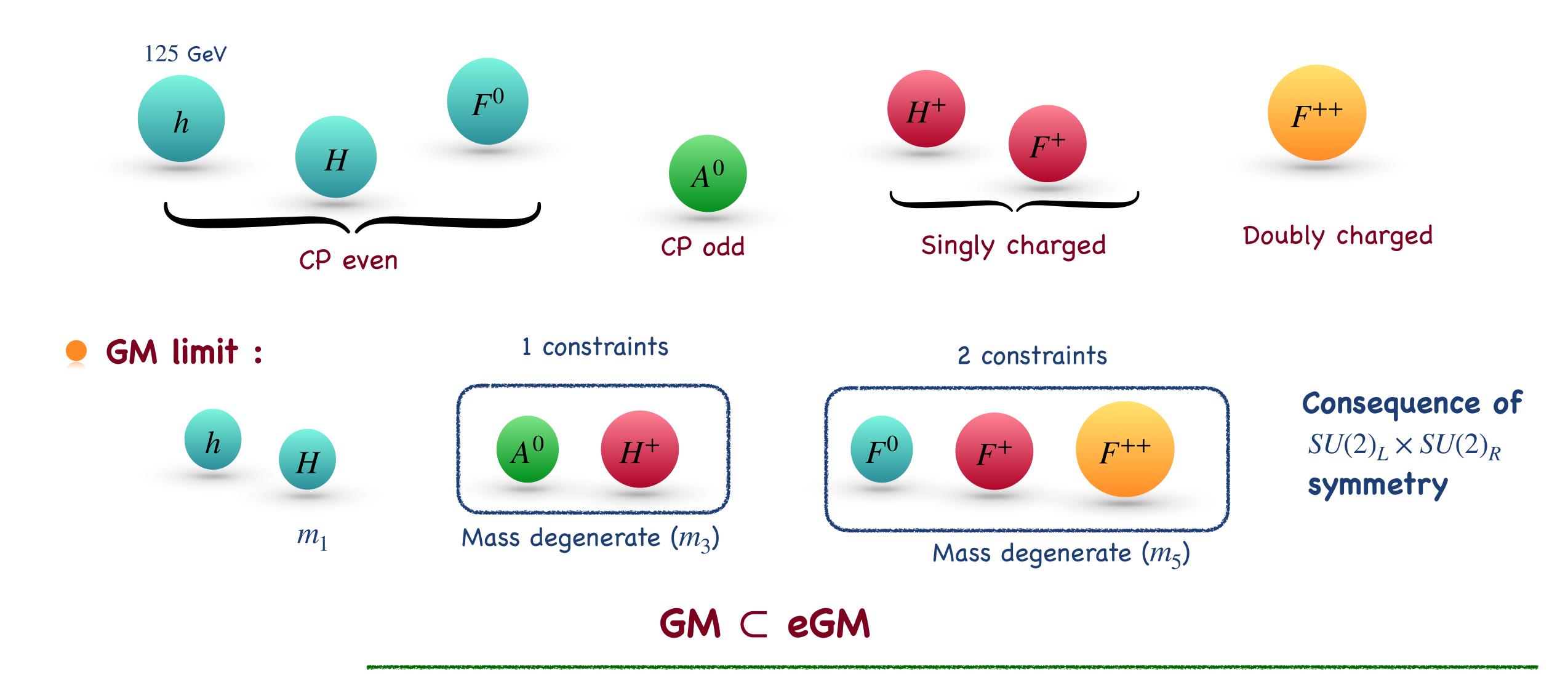
[Georgi, Machacek '85; Chanowitz, Golden '85]

* extended Georgi-Machacek (eGM) model:

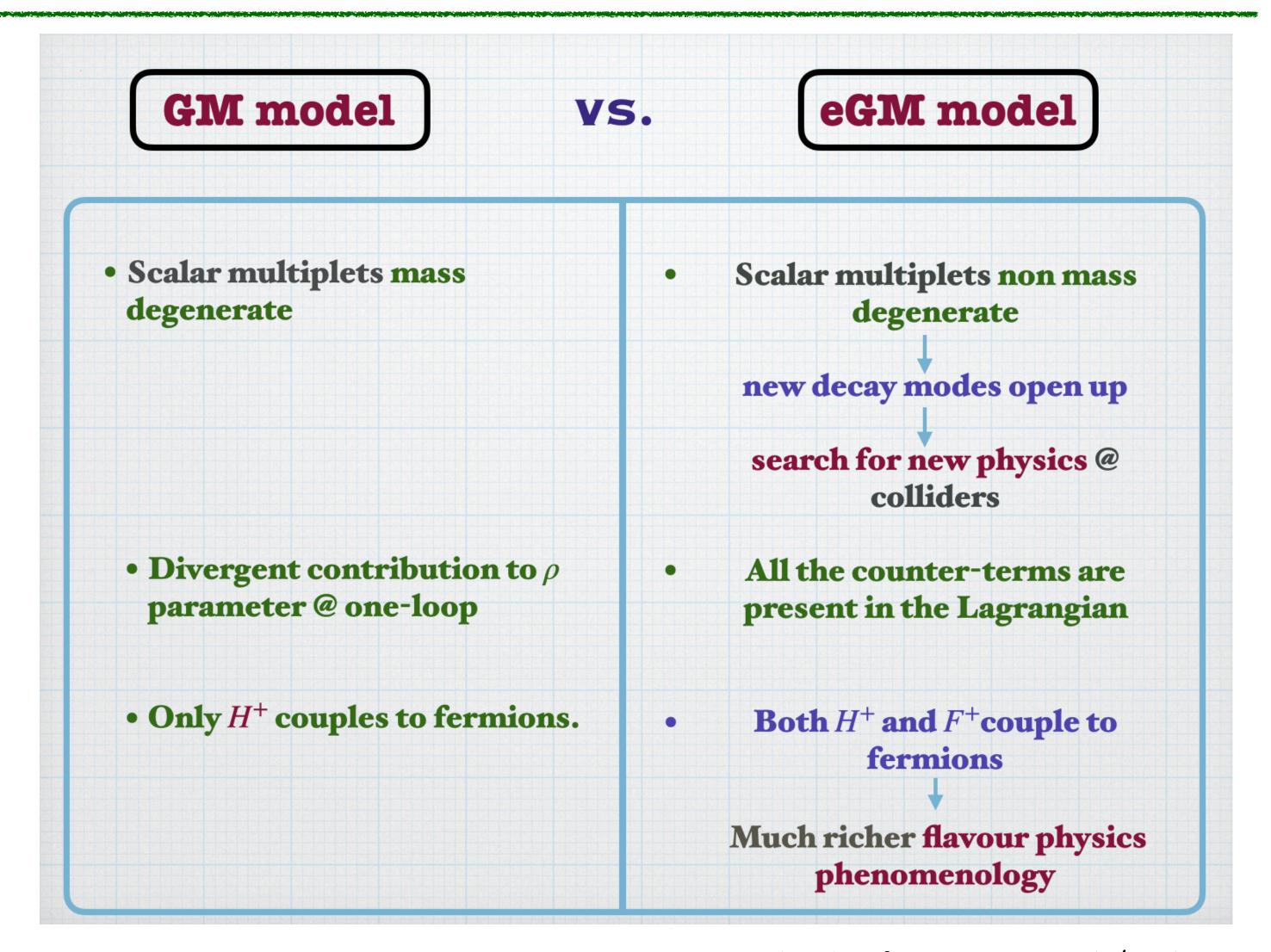
Equality of triplet VEVs is obtained by tuning the potential parameters @ tree-level [Kundu, Mondal, Pal '21] Higgs interactions does not maintained $SU(2)_L \times SU(2)_R$ symmetry

...Similar interactions were considered in 2HDM with softly broken Z_2 symmetry

Extended Georgi-Machacek Model



Salient features



... slide taken from Poulami Mondal's talk @ Higgs Hunting 2024

Theoretical Constraints

Positivity of the Higgs potential

$$V^{(4)} = \lambda_{\phi} (\phi^{\dagger} \phi)^{2} + \lambda_{\xi} (\xi^{\dagger} \xi)^{2} + \lambda_{\chi} (\chi^{\dagger} \chi)^{2} + \tilde{\lambda}_{\chi} \left| \tilde{\chi}^{\dagger} \chi \right|^{2} + \lambda_{\phi \xi} (\phi^{\dagger} \phi) (\xi^{\dagger} \xi)$$

$$+ \lambda_{\phi \chi} (\phi^{\dagger} \phi) (\chi^{\dagger} \chi) + \lambda_{\chi \xi} (\chi^{\dagger} \chi) (\xi^{\dagger} \xi) + \kappa_{1} \left| \xi^{\dagger} \chi \right|^{2} + \kappa_{2} (\phi^{\dagger} \tau_{a} \phi) (\chi^{\dagger} t_{a} \chi) + \kappa_{3} \left[(\phi^{T} \epsilon \tau_{a} \phi) (\chi^{\dagger} t_{a} \xi) + \text{h.c.} \right] > 0$$

Yukawa and quartic couplings of the theory need to be in perturbative regime

$$y_i < \sqrt{4\pi}$$
 and $\lambda_i < 4\pi$

Quartic couplings should satisfy the unitarity conditions @ one-loop

$$\left| a_{\ell}^{2 \to 2} - \frac{1}{2}i \right|^2 + \sum_{k > 2} \left| a_{\ell}^{2 \to k} \right|^2 = \frac{1}{4}.$$

NLO corrections to the LO amplitudes should be smaller in magnitude

$$|a_{\ell}^{NLO}| < |a_{\ell}^{LO}|$$

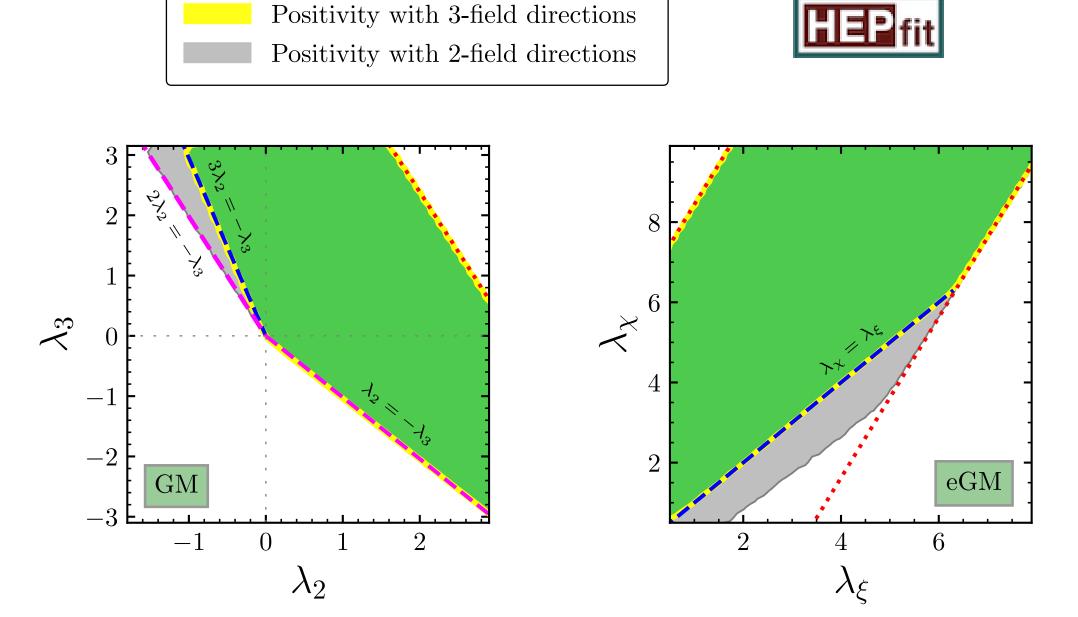
Positivity of the Higgs Potential

Ensure that boundedness of the potential in any directions of field space

Numerically, 3-field direction BFB conditions (neither necessary nor sufficient) are a very good approximation of all 13-field direction BFB conditions

[Moultaka, Peyranère '21, D. Chowdhury, P. Mondal, S.S. 2404.18996]

These 3-field direction BFB conditions are faster numerically



Positivity with 13-field directions

NLO unitarity

For a given $2 \to 2$ process, the unitarity bounds : $|a_{\ell} - i\frac{1}{2}| \le \frac{1}{2}$

- * LO unitarity: $a_{\ell}^{LO} \in \mathbb{R}$, $|\operatorname{Re}(a_{\ell}^{LO})| \leq \frac{1}{2}$
- * NLO unitarity: $a_{\ell}^{NLO} \notin \mathbb{R}$, $|a_{\ell}^{NLO} i\frac{1}{2}| \leq \frac{1}{2}$

Prior to the Higgs discovery: LO unitarity: $\lambda \leq \frac{8\pi}{3}$

Weakly interacting theories:

$$a_\ell^{LO} > a_\ell^{NLO}$$
 @ 1-loop

These are used to put bound on the potential parameters or exotic Higgs masses in a weakly interacting theory

[Lee, Quigg, Thacker '77]

NLO unitarity: $\lambda \leq 2-2.5$ [Dawson, Eillenbrock '89; Durand, Johnson, Lopez'92]

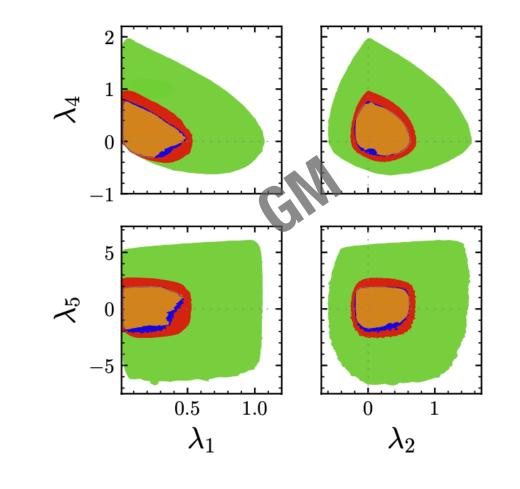
No revised limit @2-loop

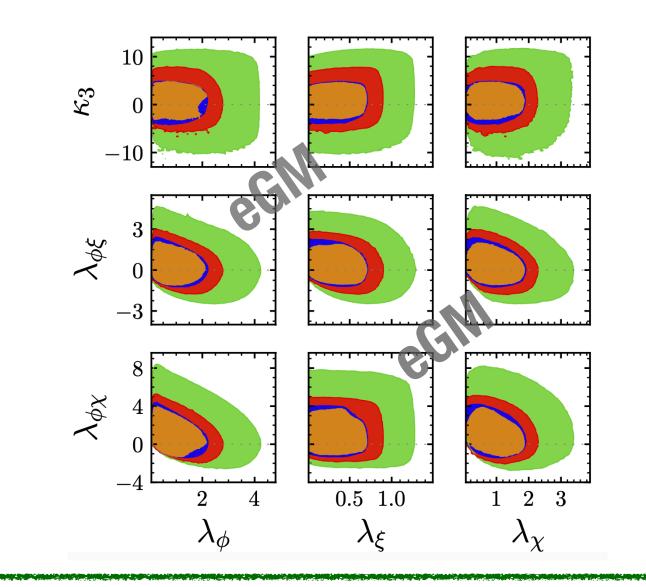
[Durand, Maher, Riesselmann, 92]

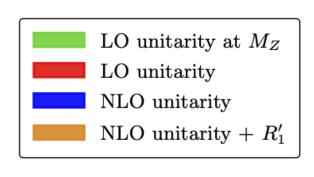
Weakly interacting SM Higgs scenario

* GM and eGM model:

NLO unitarity significantly refine the parameter space









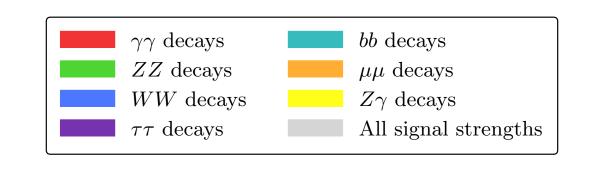
D. Chowdhury, P. Mondal, S.S. 2404.18996

Higgs Signal Strengths

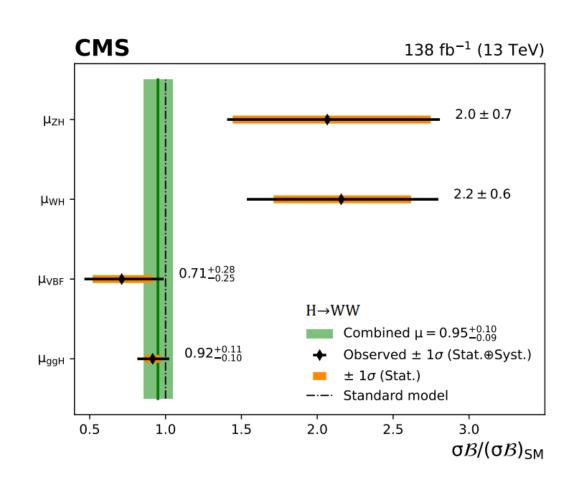
$$\mu_i^f = \frac{\sigma B \ (i \to h \to f)}{\sigma B_{SM} \ (i \to h \to f)}$$

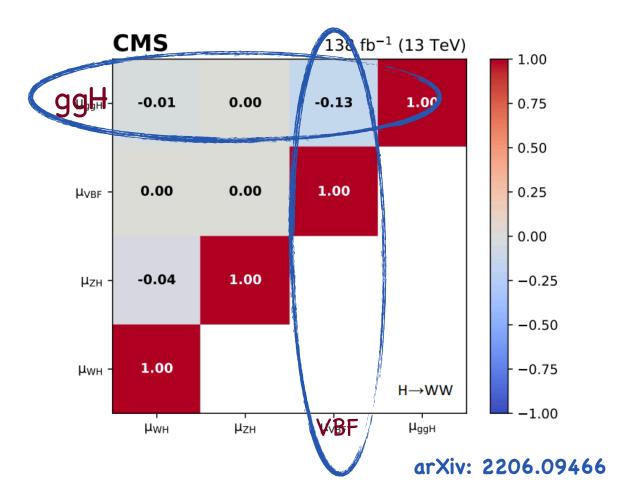
 $i \in \{ggF, bbh, VBF, Wh, Zh, tth, th\}$

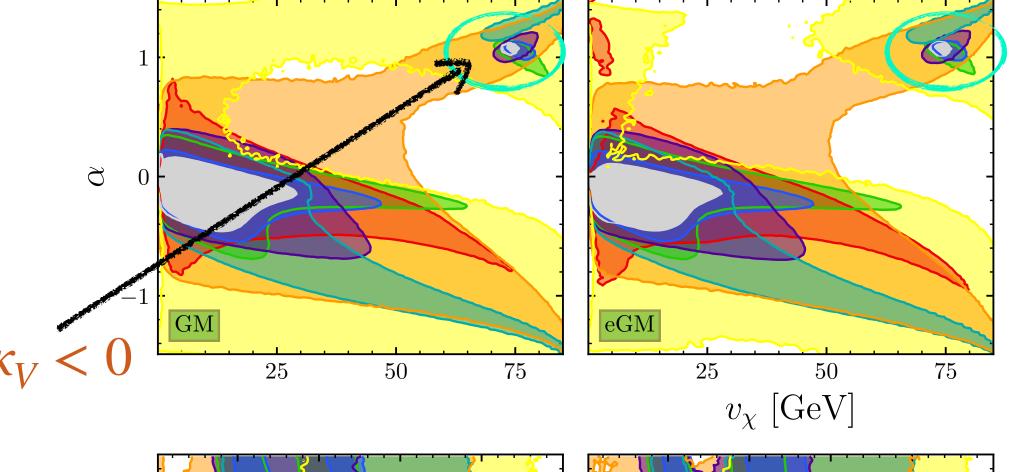
 $f \in \{ZZ, WW, \gamma\gamma, Z\gamma, \mu\mu, bb, \tau\tau\}$







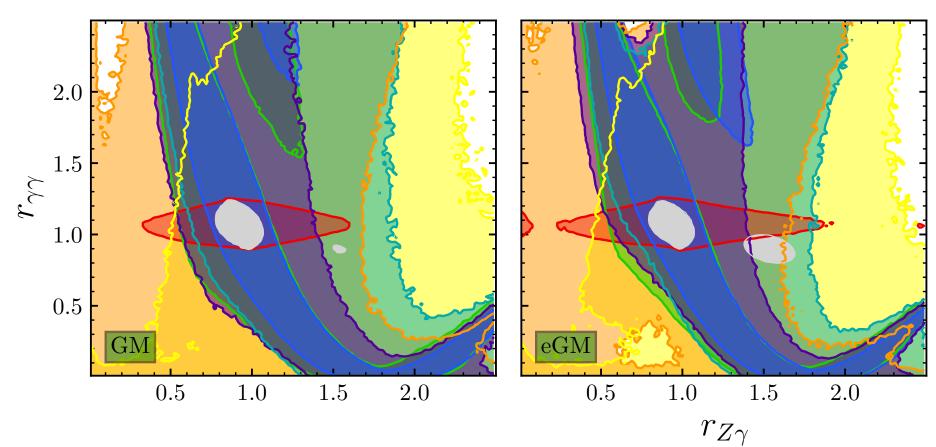




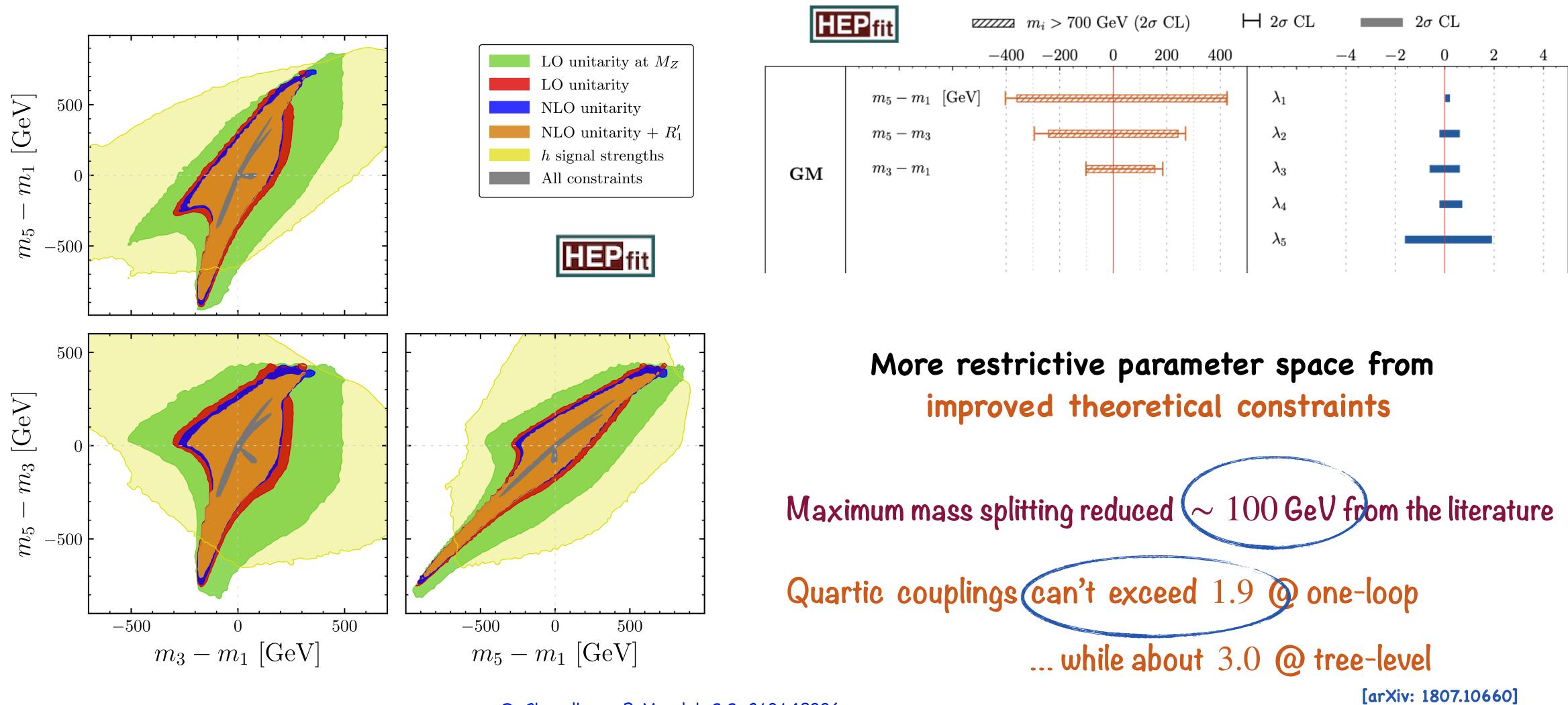
Latest Run 2 LHC data put a stringent bound on triplet VEV , $v_\chi < 32$ GeV

Strongly disfavour



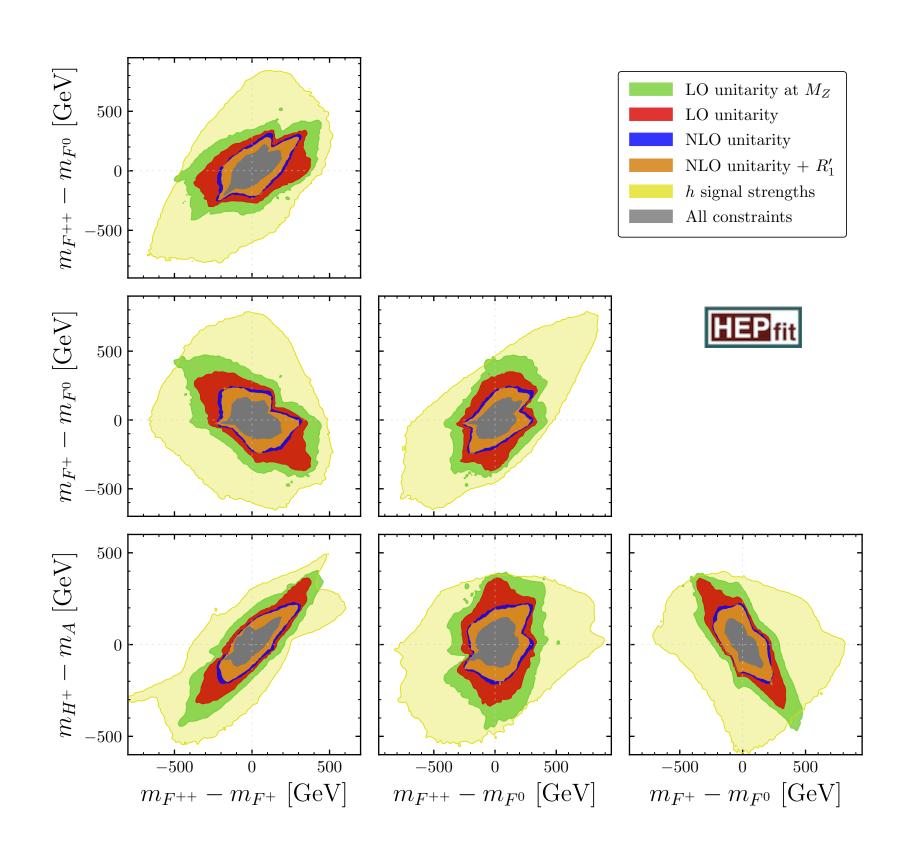


Status of GM Model (combined fit)

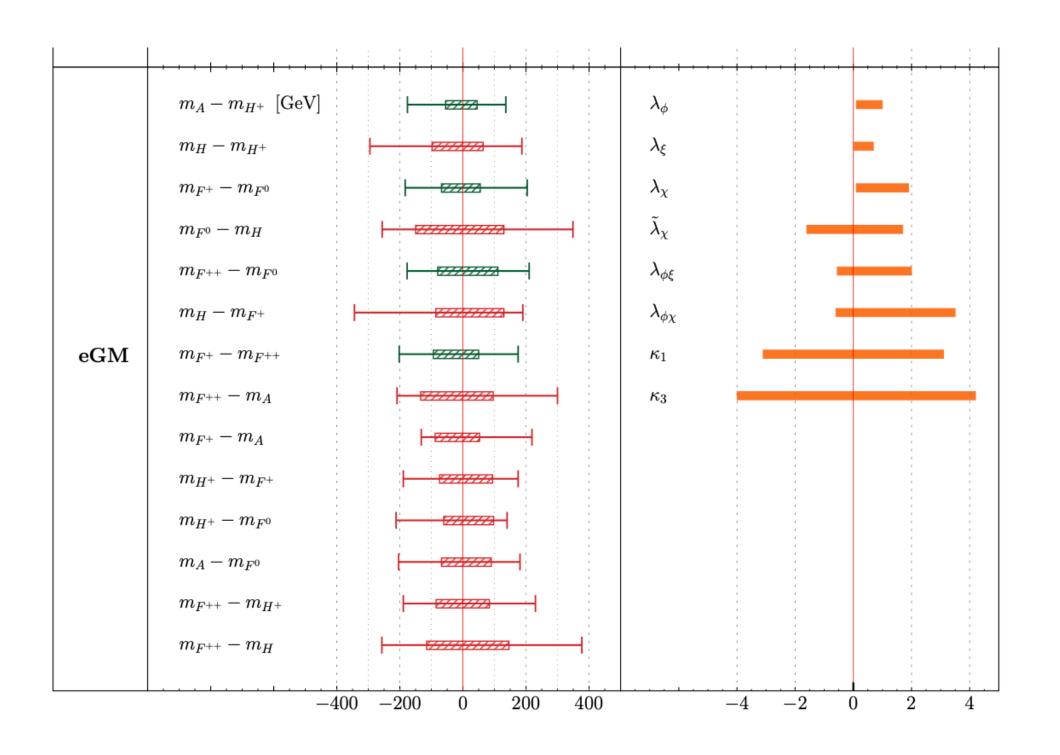


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Status of eGM Model (combined fit)



@ 95.4% CL limit on mass differences and quartic couplings



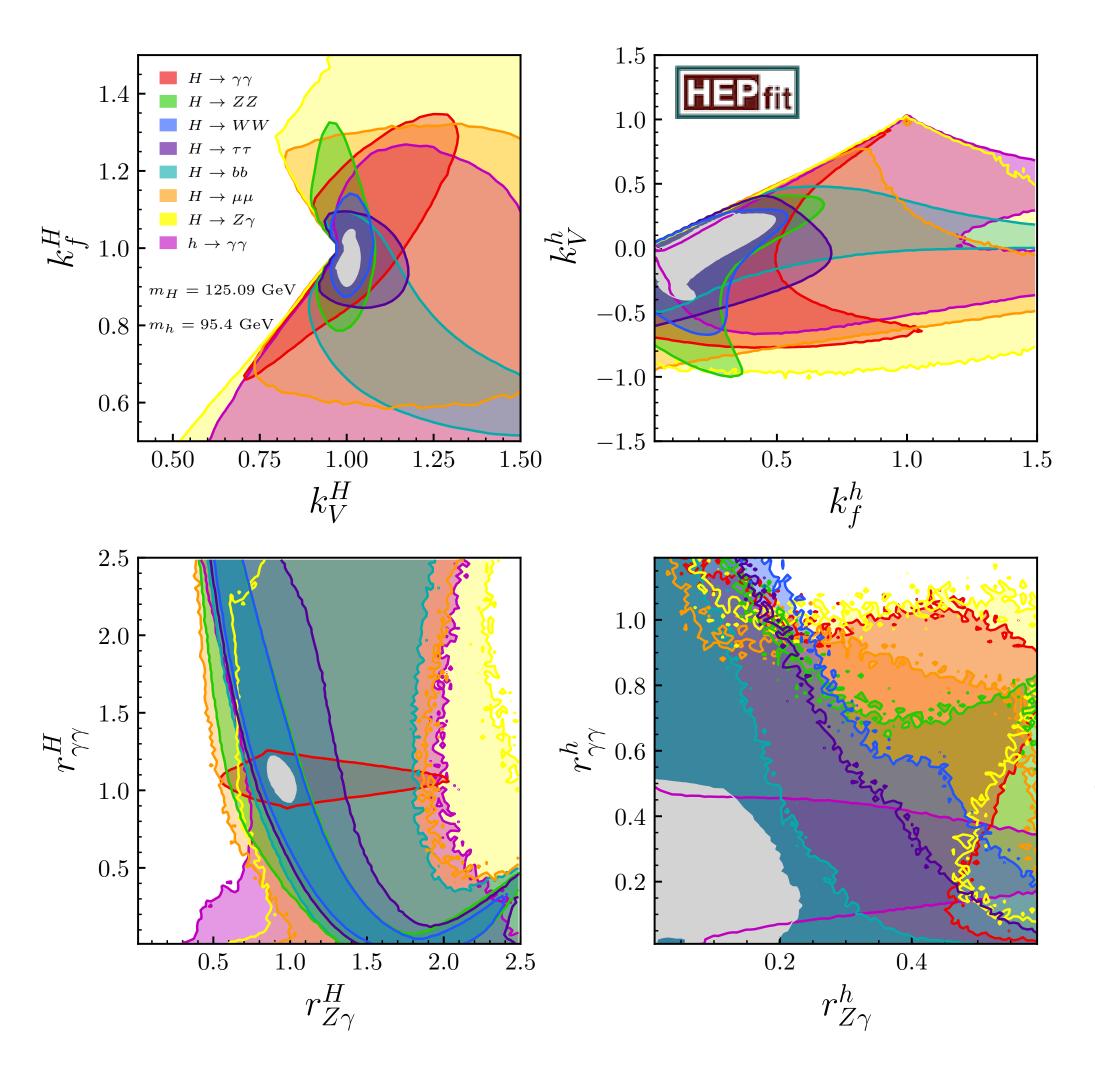
Maximum mass splitting within custodial multiplets



D. Chowdhury, P. Mondal, S.S. 2404.18996

Flavor or electroweak precision data could be used to constrain the model further. (Work in progress ...)

95 GeV Higgs Results



$$\kappa_f^h \kappa_V^h + \kappa_f^H \kappa_V^H = 1$$

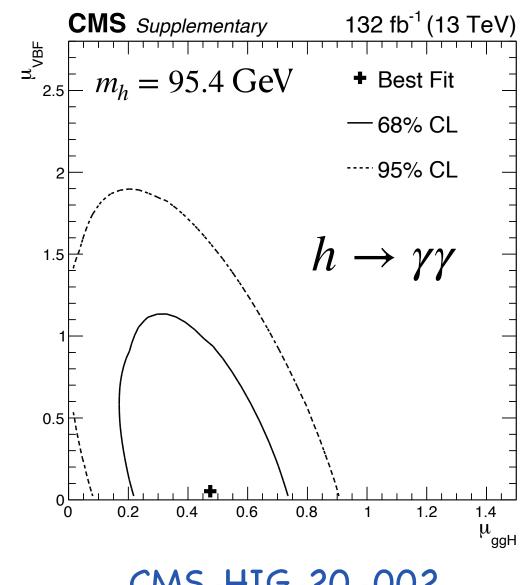
$$|\kappa_V^h| > 0.4$$
 $|\kappa_V^H| > 1.05$ excluded @ 95.4% CL

$$m_h = 95.4 \text{ GeV}$$

 $m_H = 125.09 \text{ GeV}$

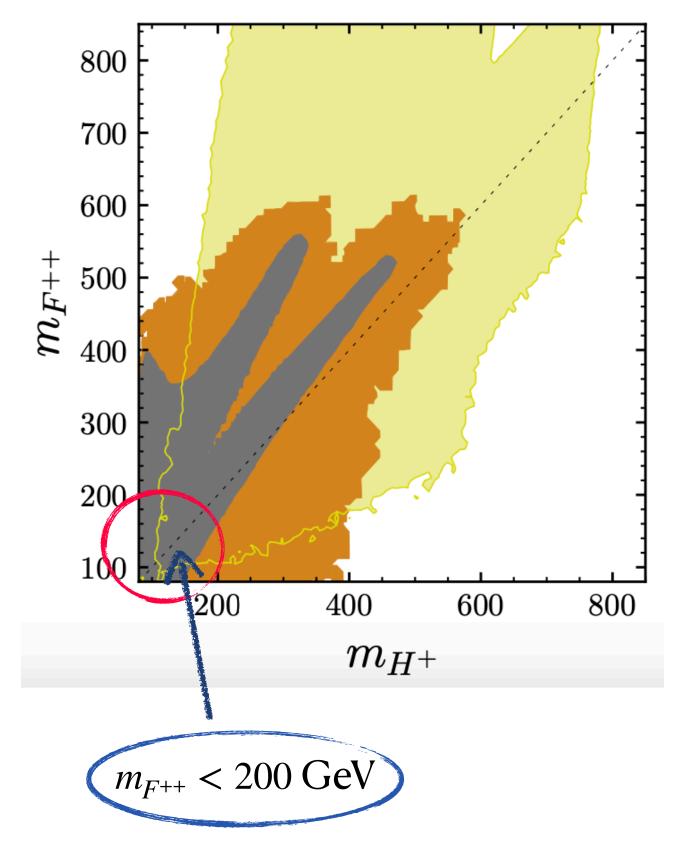
$$\mu_{\gamma\gamma}^{\text{ATLAS}} = \frac{\sigma^{\text{exp}} (pp \to \phi \to \gamma\gamma)}{\sigma^{\text{SM}} (pp \to h \to \gamma\gamma)} = 0.18^{+0.10}_{-0.10}$$

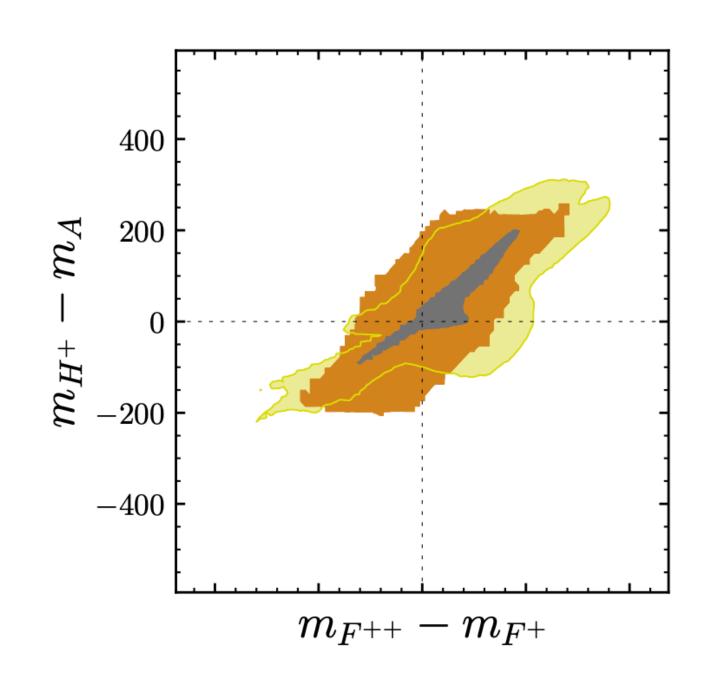
arXiv: 2306.03889

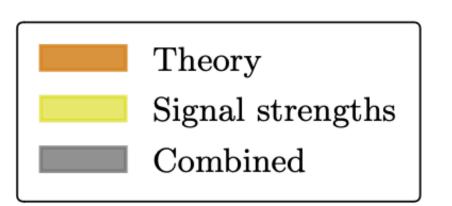


Constraints on additional Higgs bosons

Upper bound on the masses of additional heavy Higgs bosons is ~ 550 GeV @ 95.4% CL









allowed by NLO unitarity and Higgs data @ 95.4% CL

.....yet to be explored at the LHC or in future colliders

Summary

- * Triplet scalar extension of SM with custodial symmetry at tree-level without requiring global $SU(2)_L \times SU(2)_R$ symmetry in the potential gives extended Georgi-Machacek (eGM) model
- * Improved theoretical constraints (NLO unitarity with positivity) significantly refine the parameter space of the GM and eGM models
- * Triplet VEV gets more and more constrained from the LHC data
- * Regions where $|\kappa_V| > 1.05$ are disfavoured by the latest LHC data
- * Mass splittings within custodial multiplets introduce new decay modes in eGM model
- * In the presence of a 95 GeV Higgs, the upper bound on the masses of additional heavy Higgs bosons is \sim 550 GeV

Thank You

Backup slides

Collider Phenomenology

Yukawa sector :

Only the doublet couples to fermions

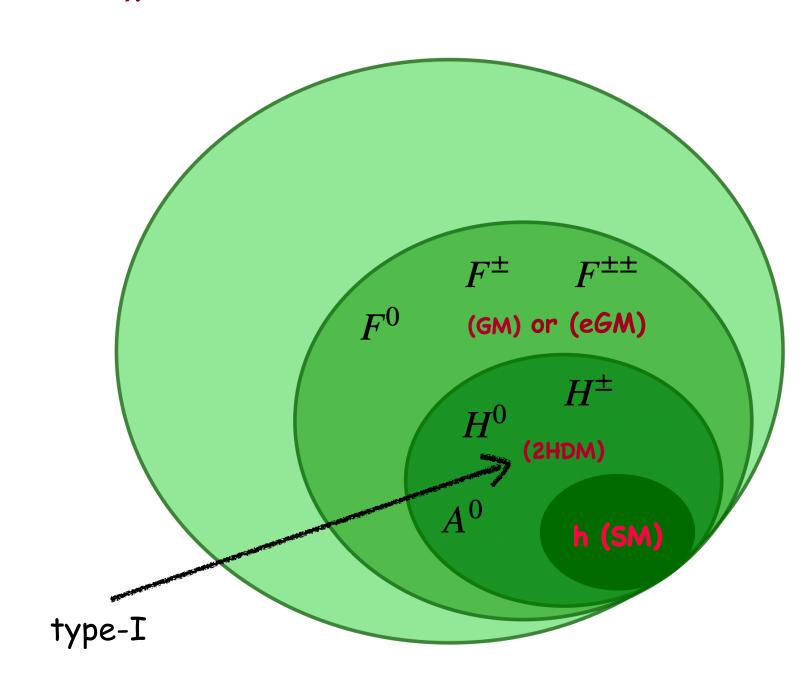
triplet VEV
$$(v_{\chi})$$
: $v_{\phi}^2 + 8v_{\chi}^2 = v^2$ and $\tan \beta = \frac{v_{\phi}}{2\sqrt{2}v_{\chi}}$



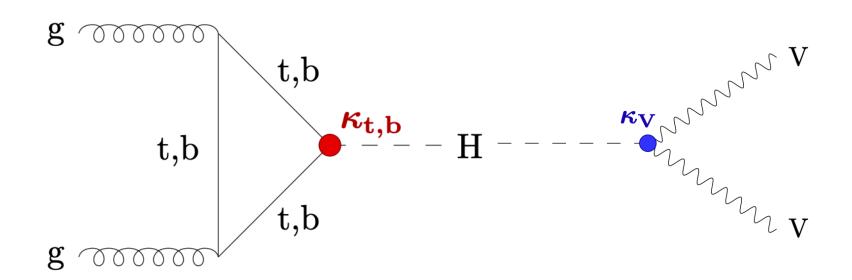
.....Similar phenomenology as in type-I 2HDM

Additional features:

The addition of a singly charged scalar (F^+) coupled to fermions, along with the presence of a doubly charged scalar, makes these models highly interesting for collider studies.

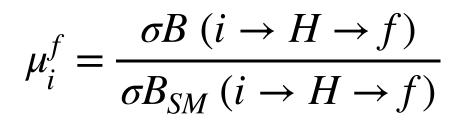


Higgs Signal Strengths: HEPfit Implementation



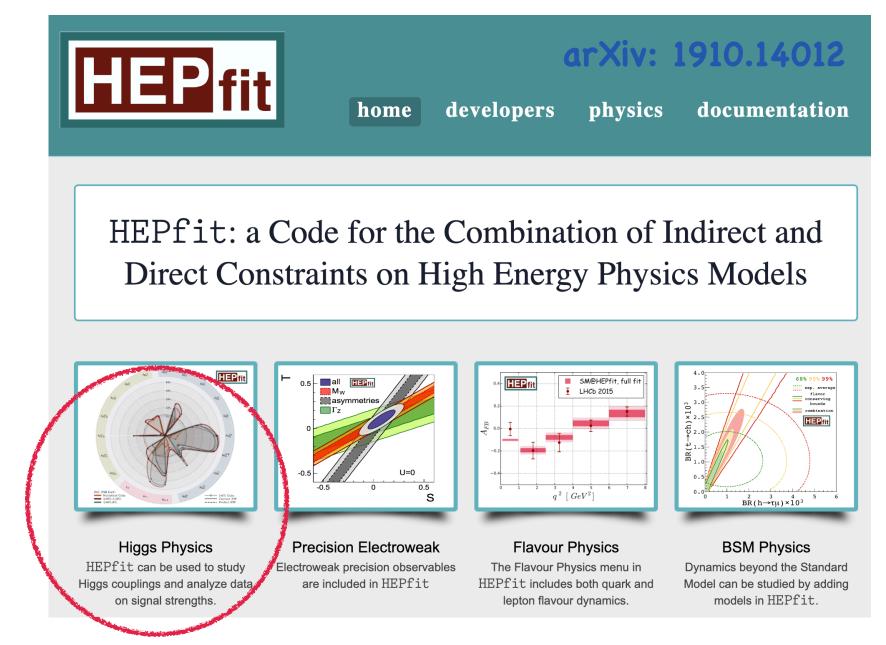
 $i \in \{ggF, bbh, VBF, Wh, Zh, tth, th\}$

$$f \in \{ZZ, WW, \gamma\gamma, Z\gamma, \mu\mu, bb, \tau\tau\}$$

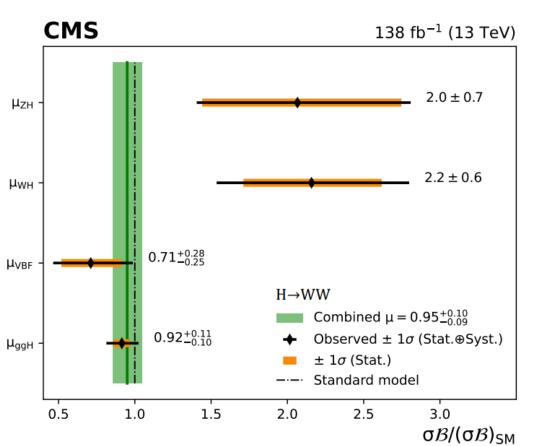


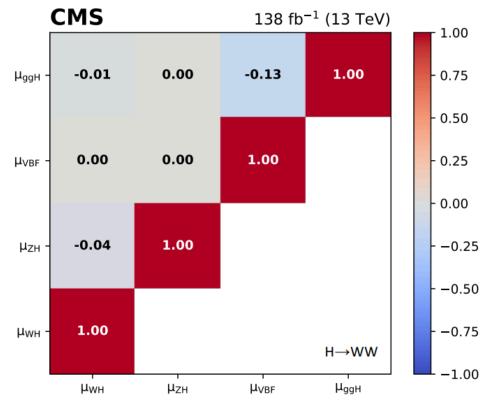
- \bullet Make all possible observables μ_i^f for different production and decay modes
- ullet Fit to the ALTAS and CMS data on (correlated) observables μ_i^f for a BSM model
- Present the results on the (new) observables from the combined fit

$$\kappa_V = c_{\alpha} c_{\beta} - \sqrt{\frac{8}{3}} s_{\alpha} s_{\beta} \,, \quad \text{and} \quad \kappa_f = \frac{c_{\alpha}}{c_{\beta}} \,,$$



http://hepfit.romal.infn.it





arXiv: 2206.09466

Higgs Signal Strengths: LHC data

ATLAS Run 2

Signal strength	Value		Co	rrelatio	n mat	rix		$\frac{\mathcal{L}}{[\mathbf{fb}^{-1}]}$	Source
	1041010		0.10					[ID]	
$\mu_{ m ggF,bbh}^{\gamma\gamma}$	1.04 ± 0.10	1	-0.13	0	0	0	0		
$\mu_{ ext{VBF}}^{\gamma\gamma}$	1.20 ± 0.26	-0.13	1	0	0	0	0		
$\mu_{\mathrm{Wh}}^{\gamma\gamma}$	1.5 ± 0.55	0	0	1	-0.37	0	-0.11	100	[18]
$\mu_{\mathrm{Zh}}^{\gamma\gamma}$	-0.2 ± 0.55	0	0	-0.37	1	0	0	139	[10]
$\mu_{ m tth}^{\gamma\gamma}$	0.89 ± 0.31	0	0	0	0	1	-0.44		
$\mu_{ m th}^{\gamma\gamma}$	3 ± 3.5	0	0	-0.11	0	-0.44	1		
$\mu_{ m ggF}^{ZZ}$	0.95 ± 0.1	1	-0.22	-0.27	0				
$\mu^{ZZ}_{ ext{VBF}}$	1.19 ± 0.45	-0.22	1	0	0				[4]
$\mu_{ m Vh}^{ZZ}$	1.43 ± 1.0	-0.27	0	1	-0.18			139	[4]
$\mu^{ZZ}_{ m tth}$	1.69 ± 1.45	0	0	-0.18	1				
$\mu^{ZZ}_{ m incl.}$	1.0 ± 0.1							139	[4]
$\mu_{ m ggF,bbh}^{WW}$	1.15 ± 0.135								
μ_{VBF}^{WW}	0.93 ± 0.21							139	[17]
$\mu_{ m ggF,bbh,VBF}^{WW}$	1.09 ± 0.11								
$\mu_{ ext{VBF}}^{ au au}$	0.90 ± 0.18	1	-0.24	0	0				
$\mu_{ m ggF,bbh}^{ au au}$	0.96 ± 0.31	-0.24	1	-0.29	0			400	[10]
$\mu_{\mathrm{Vh}}^{ au au}$	0.98 ± 0.60	0	-0.29	1	0			139	[13]
$\mu^{ au au}_{ m tth,th}$	1.06 ± 1.18	0	0	0	1				
$\mu_{ ext{VBF}}^{bb}$	0.95 ± 0.37							126	[9]
$\mu_{ m Wh}^{bb}$	0.95 ± 0.26							139	[<mark>6</mark>]
$\mu_{ m Zh}^{bb}$	1.08 ± 0.24							139	[<mark>6</mark>]
$\mu_{ m Vh}^{bb}$	1.02 ± 0.17							139	[<mark>6</mark>]
$\mu_{ m tth,th}^{bb}$	0.35 ± 0.35							139	[12]
$\mu_{ m pp}^{\mu\mu}$	1.2 ± 0.6							139	[7]
$\mu_{ m pp}^{Z\gamma}$	2.0 ± 0.95							139	[5]

CMS Run 2

Signal	Value	Correlation matrix	\mathcal{L}	Source
strength			$[\mathbf{f}\mathbf{b}^{-1}]$	
$\mu_{ m ggh,bbh}^{\gamma\gamma}$	1.07 ± 0.11			
$\mu_{ ext{VBF}}^{\gamma\gamma}$	1.04 ± 0.32			[4.4]
$\mu_{ m Vh}^{\gamma\gamma}$	1.34 ± 0.34		137	[11]
$\mu_{ m tth,th}^{\gamma\gamma}$	1.35 ± 0.31			
$\mu^{ZZ}_{ m ggh,bbh,tth,th}$	0.95 ± 0.13	13 1 -0.11		[10]
$\mu^{ZZ}_{ ext{VBF,Vh}}$	0.82 ± 0.34	-0.11 1	137	[10]
$\mu_{ m ggh}^{WW}$	0.92 ± 0.11	1 -0.13 0 0		
$\mu_{ ext{VBF}}^{WW}$	0.71 ± 0.26	-0.13 1 0 0		[16]
$\mu_{ m Zh}^{WW}$	2.0 ± 0.7	0 0 1 0	138	[16]
$\mu_{ m Wh}^{WW}$	2.2 ± 0.6	0 0 0 1		
$\mu_{ m incl.}^{ au au}$	0.93 ± 0.12			
$\mu_{ m ggh}^{ au au}$	0.97 ± 0.19			[15]
$\mu_{ m qqh}^{ au au}$	0.68 ± 0.23		138	[10]
$\mu_{ m Vh}^{ au au}$	1.80 ± 0.44			
$\mu_{ m qqh}^{bb}$	1.59 ± 0.60	1 -0.75		[10]
$\mu_{ m ggh}^{bb} = -2.7 \pm 3.$		-0.75 1	90.8	[19]
$\mu^{\mu\mu}_{ m ggh,tth}$	0.66 ± 0.67	1 -0.24	197	[0]
$\mu_{ ext{VBF,Vh}}^{\mu\mu}$	1.85 ± 0.86	-0.24 1	137	[8]
$\mu_{ m pp}^{Z\gamma}$	2.4 ± 0.9		138	[14]