

# Were there any anomalies in the gluon jets in ALEPH?

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Youngkwon Jo, Inkyu Park (University of Seoul),  
& Yi Chen, Yen-Jie LEE (MIT)



# About this talk

- **Title**
  - Were there any anomalies in the gluon jets in ALEPH?
- **Abstract**
  - According to the **Abelian decomposition of QCD**, there is a theoretical prediction that **gluons can be classified into two types**, each **exhibiting distinct experimental signatures**. The optimal setting for experimental verification of this theory is a clean environment such as the LEP, rather than the LHC. We have investigated whether there were any anomalies observed already in the gluon jets recorded in the ALEPH experiment and **revisited the analyses with the archived ALEPH data**. In this presentation, we will show our latest updates on our study on the gluon jet properties in ALEPH.
- **Authors**
  - **Youngkwon Jo, Inkyu Park (University of Seoul)**
  - **Yi Chen, Yen-Jie Lee (MIT)**



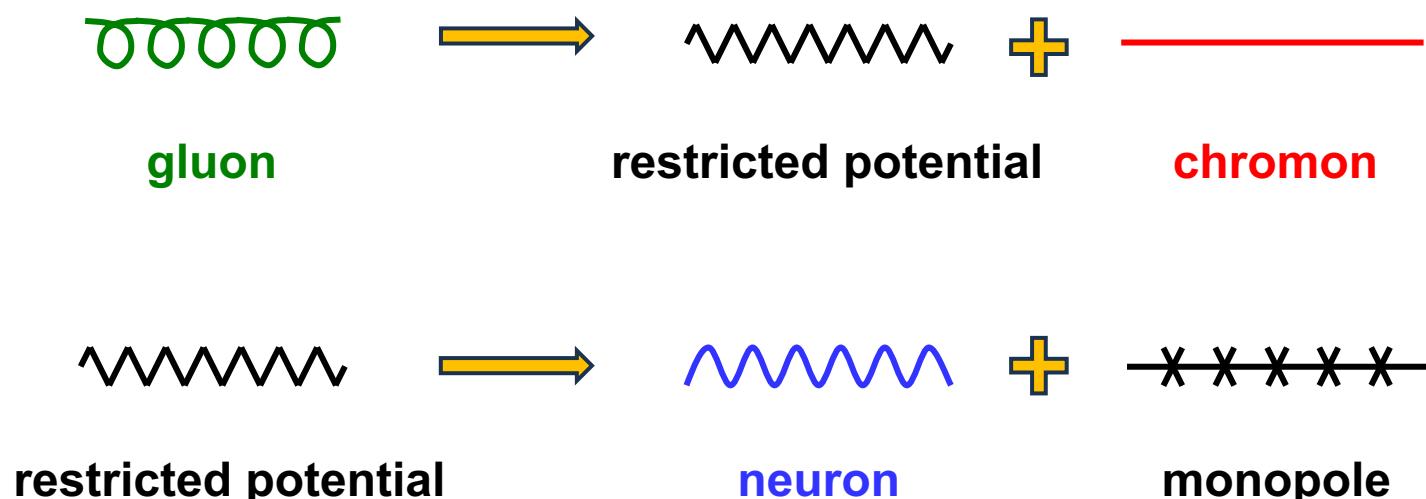
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- 3 Resurrection of the ALEPH data
- 4 Event selections & observables
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- 6 Event display
- 7 Remarks & Discussions

Motivation: Are there two types of gluons?

# Theoretical motivation

- “**Experimental verification of Two types of Gluon Jets in QCD**”
  - **Y. M. Cho, Pengming Zhang, and Li-Ping Zou**
  - ✓ **Phys. Rev. D 107, 054024 – Published 17 March 2023**



PHYSICAL REVIEW D 107, 054024 (2023)

## Experimental verification of two types of gluon jets in QCD

Y. M. Cho<sup>1,2,\*</sup>, Pengming Zhang,<sup>3,†</sup> and Li-Ping Zou<sup>3,‡</sup>

<sup>1</sup>School of Physics and Astronomy, Seoul National University, Seoul 08826, Korea

<sup>2</sup>Center for Quantum Spacetime, Sogang University, Seoul 04107, Korea

<sup>3</sup>School of Physics and Astronomy, Sun Yat-Sen University, Zhuhai 519082, China

<sup>4</sup>Sino-French Institute of Nuclear Engineering and Technology,

Sun Yat-Sen University, Zhuhai 519082, China

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The Abelian decomposition of QCD tells that there are two types of gluons, the color neutral neurons and colored chromons. We propose to confirm the Abelian decomposition testing the existence of two types of gluons jets experimentally. We predict that one quarter of the gluon jet is made of the neurons which has the color factor 3/4 and the sharpest jet radius and smallest charged particle multiplicity, while the three quarters of the gluon jet are made of the chromons with the color factor 9/4, which have the broadest jet radius (broader than the quark jet). Moreover, we argue that the neuron jet has a distinct color flow which forms an ideal color dipole, while the quark and chromon jets have distorted dipole pattern. To test the plausibility of this proposal, we suggest to analyze the gluon distribution against the jet shape (the sphericity) and/or particle multiplicity from the existing gluon jet events and look for two distinct peaks in the distribution.

DOI: 10.1103/PhysRevD.107.054024

## I. INTRODUCTION

A common misunderstanding on QCD is that the non-Abelian color gauge symmetry is so tight that it defines the theory almost uniquely and thus, does not allow any simplification. This is not true. The Abelian decomposition of QCD tells that we can construct the restricted QCD (RCD), which inherits the full non-Abelian color gauge symmetry with the restricted potential obtained by the Abelian projection. This tells that QCD has a nontrivial core, RCD, which describes the Abelian subdynamics of QCD but has the full color gauge symmetry. Moreover, it tells that QCD can be viewed as RCD which has the gauge covariant valence gluons as the colored source [1,2]. This is because the Abelian decomposition decomposes the color gauge potential to the restricted potential made of the color neutral gluon potential, the topological monopole potential, and the gauge covariant valence potential, which describes the colored gluon gauge independently.

There are ample motivations for the Abelian decomposition. Consider the proton made of three quarks. Obviously, we need the gluons to bind the quarks in the proton. However, the quark model tells that the proton has no valence gluon. If so, what is the binding gluon that bind the quarks in proton, and how do we distinguish it from the valence gluon?

Moreover, the simple group theory tells that the color gauge group has the Abelian subgroup generated by the diagonal generators and that the gauge potential, which corresponds to these generators, must be color neutral, while the potential which corresponds to the off diagonal generators must carry the color. This strongly implies that there are two types of gluon, the color neutral ones and colored ones. And they should behave differently, because they have different color charges. If so, how can we distinguish them?

Another motivation is the color confinement in QCD. Two popular proposals for the confinement are the monopole condensation [2,3] and the Abelian dominance [4,5]. To prove the monopole condensation, we first have to separate the monopole potential gauge independently. Similarly, to prove the Abelian dominance, we have to know what is the Abelian part and how to separate it.

The Abelian decomposition tells how to do this. It decomposes the non-Abelian gauge potential to two parts, the restricted Abelian part, which has the full non-Abelian gauge symmetry and the gauge covariant valence part,

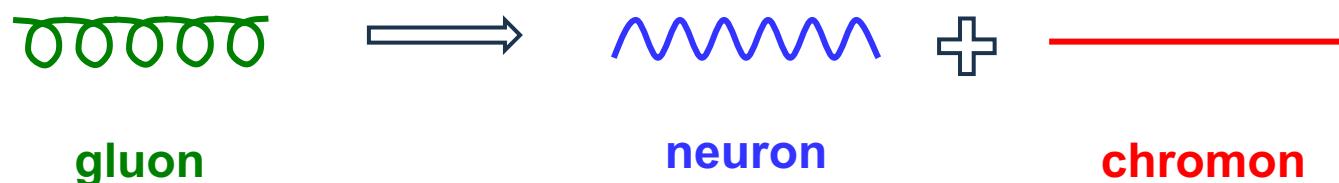
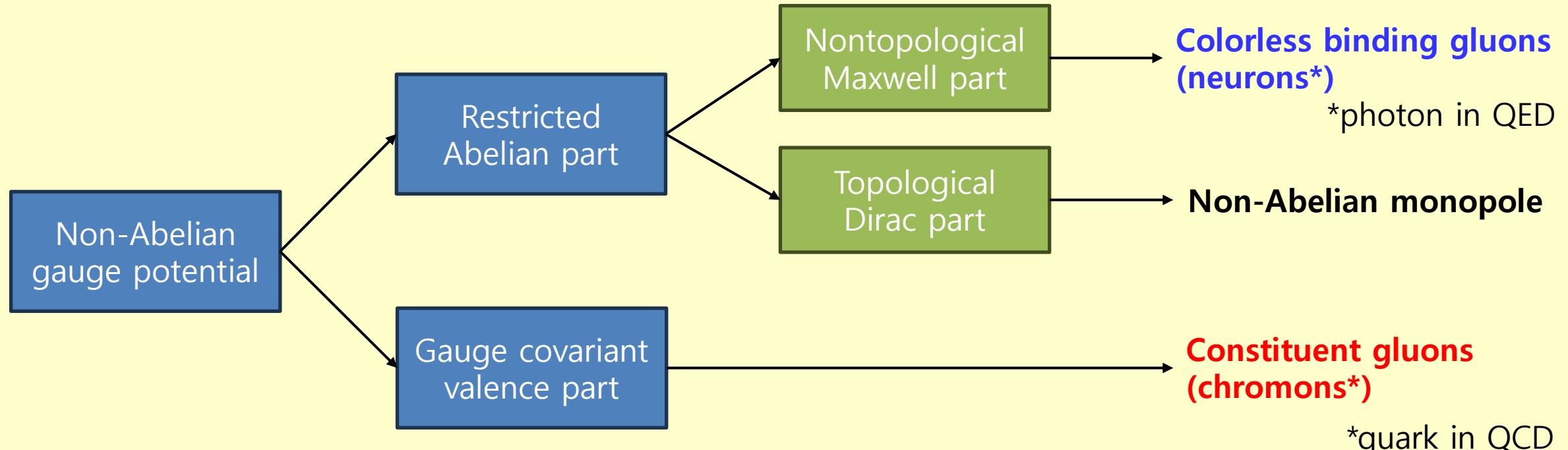
\*ymcho0416@gmail.com

†zhangpm5@mail.sysu.edu.cn

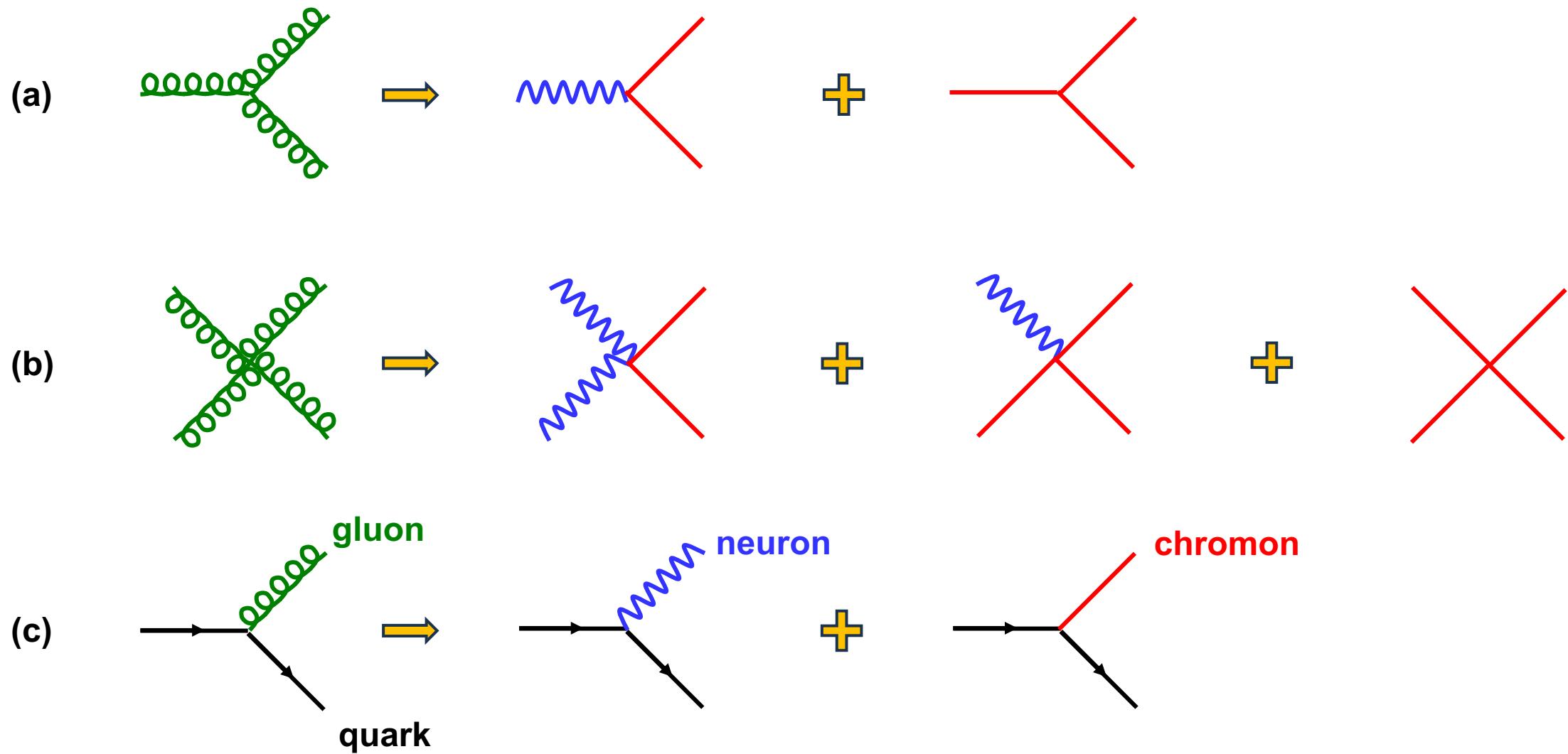
‡zoulp5@mail.sysu.edu.cn

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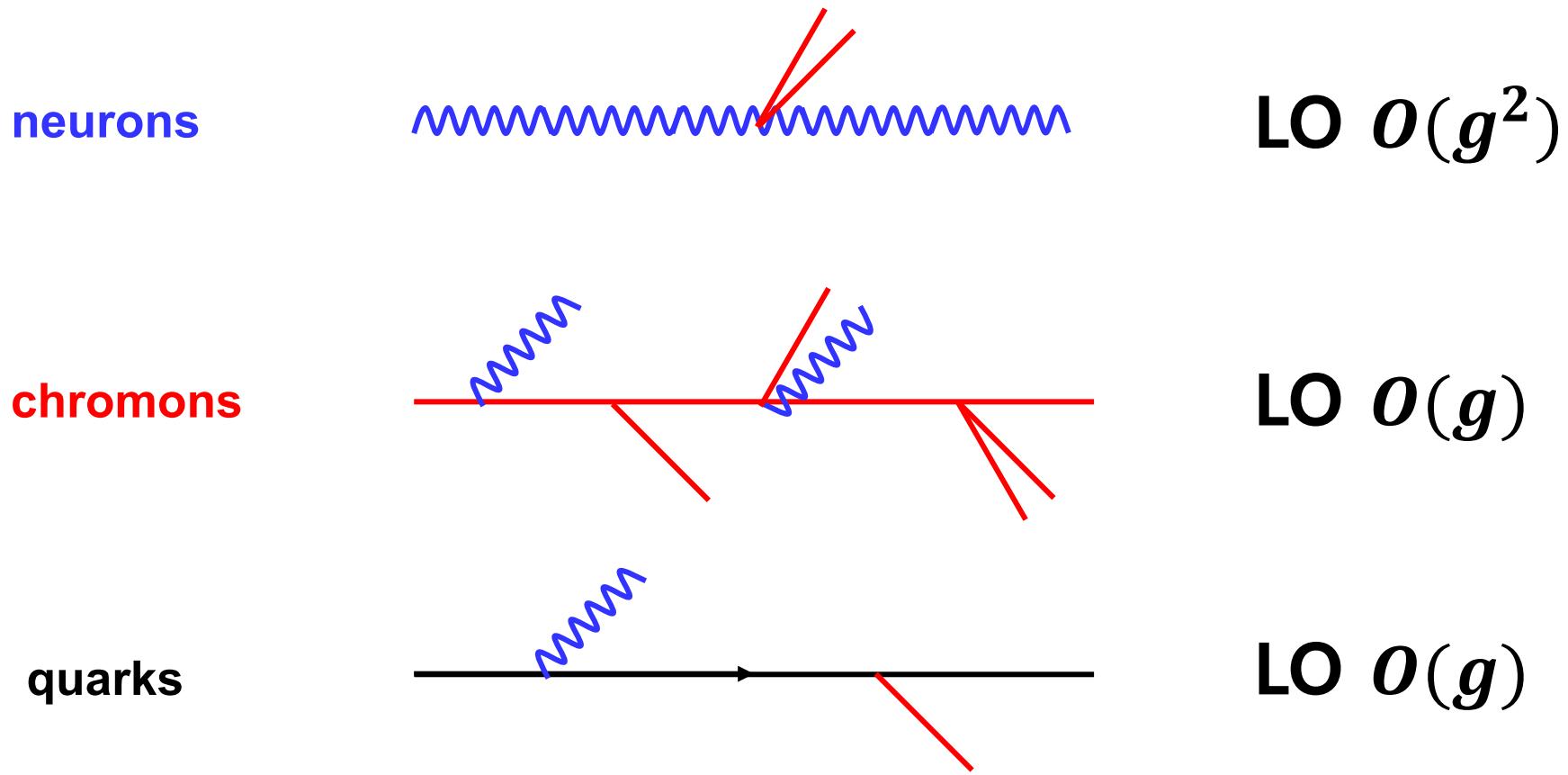
# Abelian decomposition of QCD



# Decomposition of Feynman diagrams



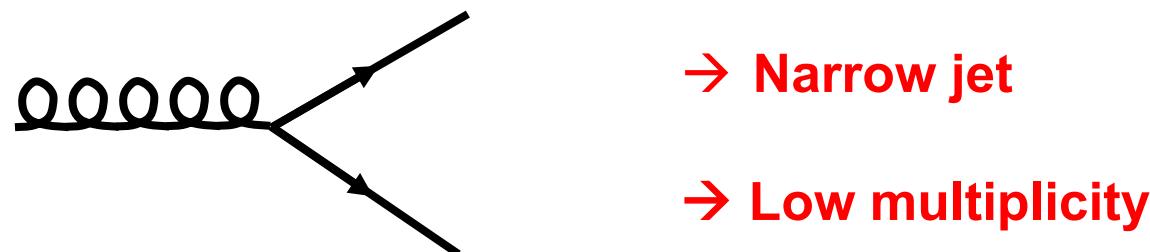
# Neuron jet and Chromon jet



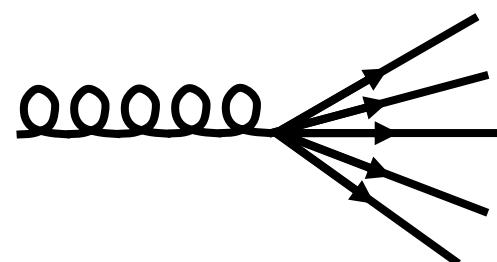
# Experimental signatures of gluon jets

- if there exist two distinct types of gluons, then

- Color neutral gluons (neurons) → EM jet like shape



- Colored gluons (chromons) → Typical hadronic jet shape

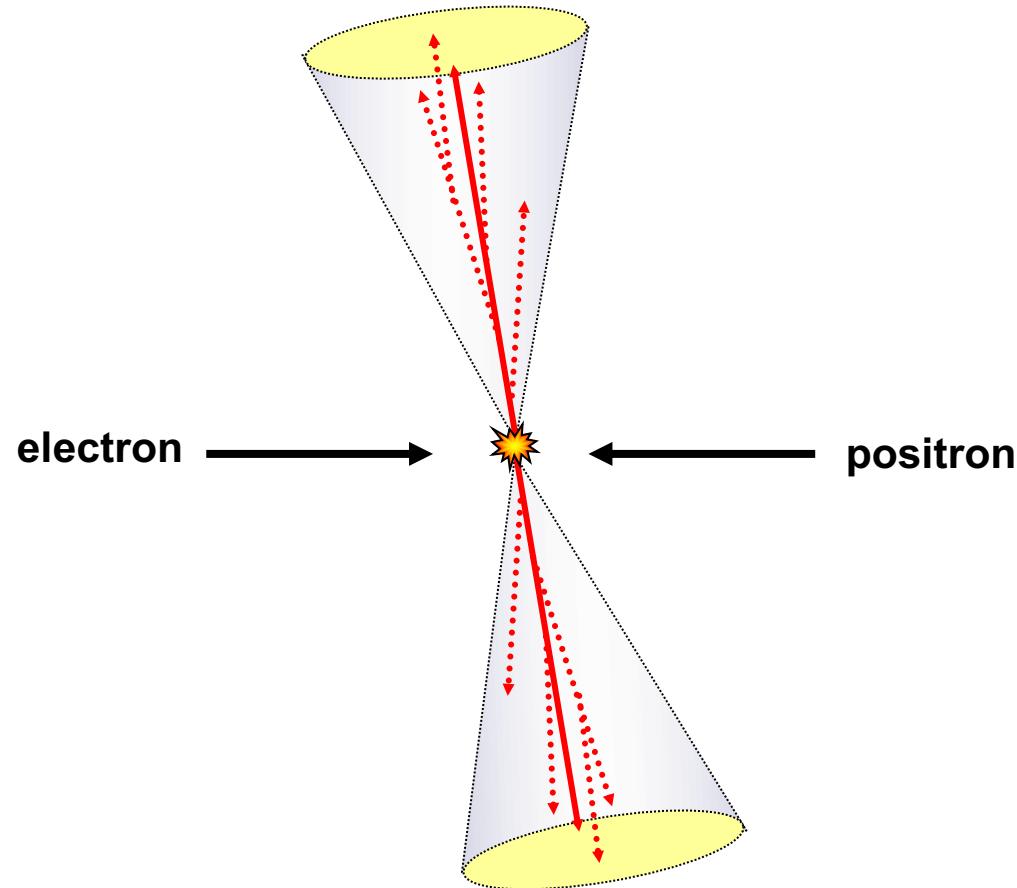


# then, we might observed...

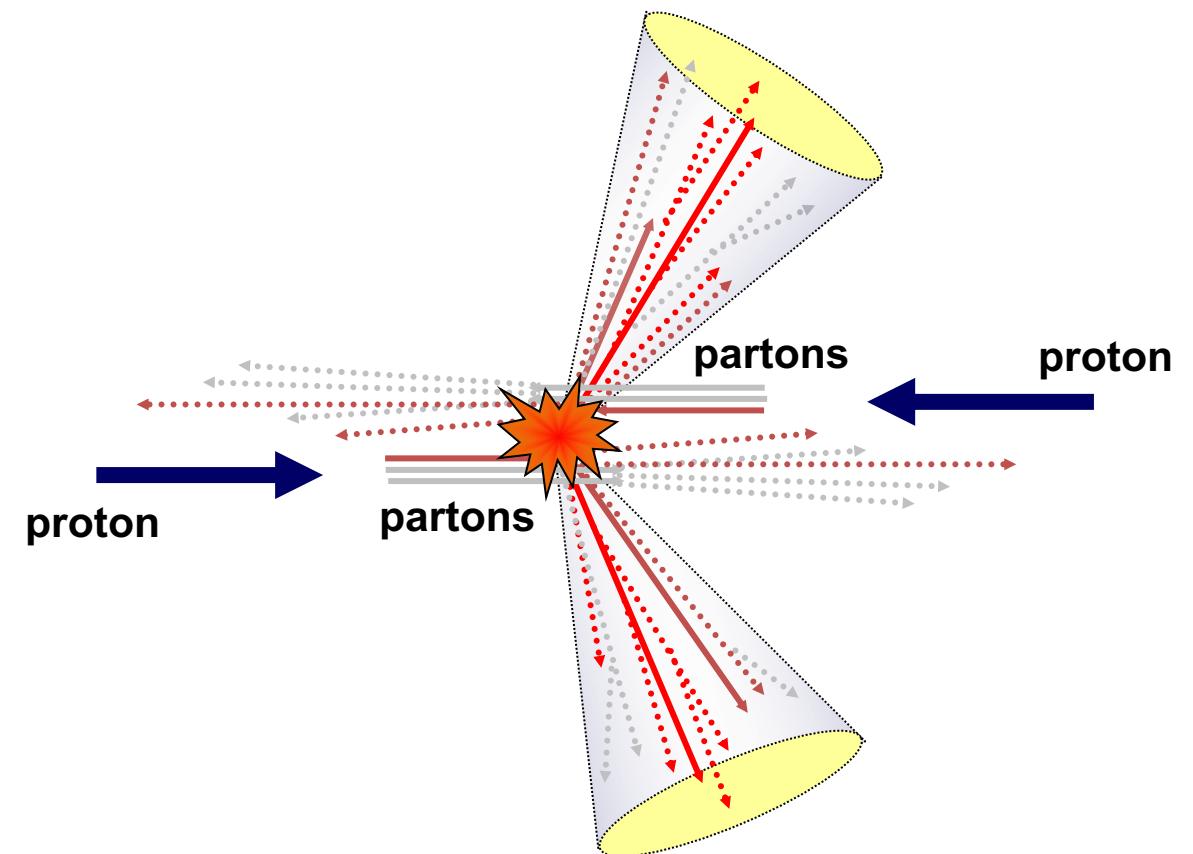
- due to the typical jet reconstruction techniques,
  - deficit of color neutral gluons
    - ✓ might not be reconstructed as jets
  - excess of jets with small number of tracks
    - ✓ check this easily by seeing # of tracks in jets
  - two distinct jet shapes
    - ✓ narrow jet (neurons) .vs. broad jet (chromons)
  - anomalies in energy shares of tracks in a jet
    - ✓ high share (neurons) .vs. low share (chromons)

Recall: QCD Jet study @ LEP

# LEP vs LHC for QCD study



No initial state colors, back-to-back jets



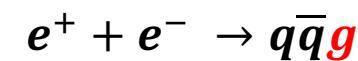
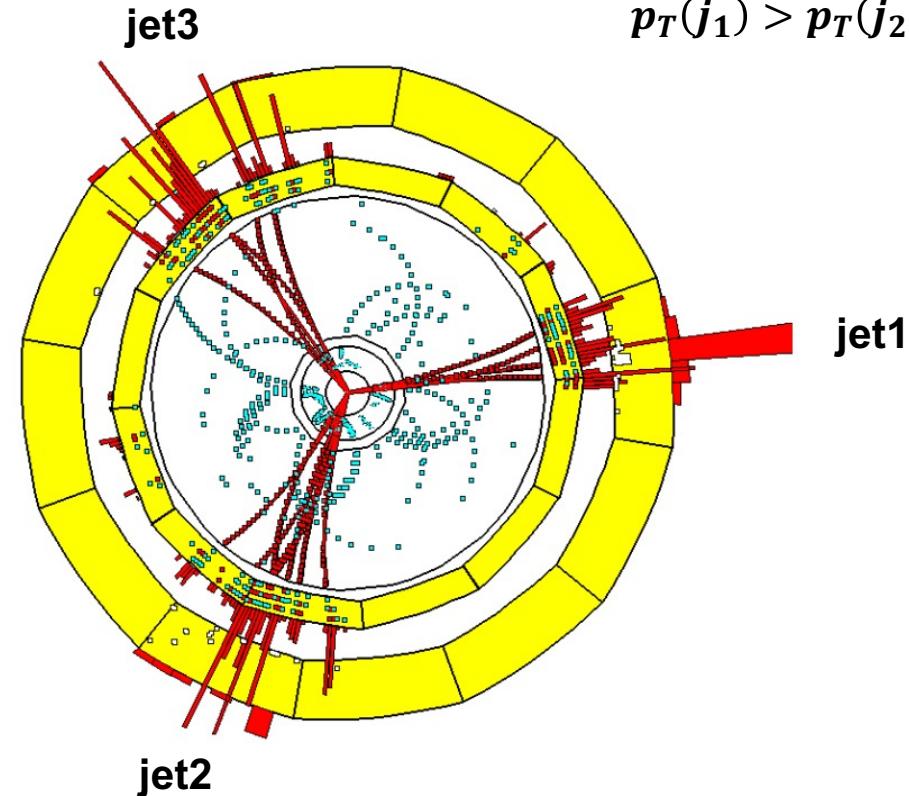
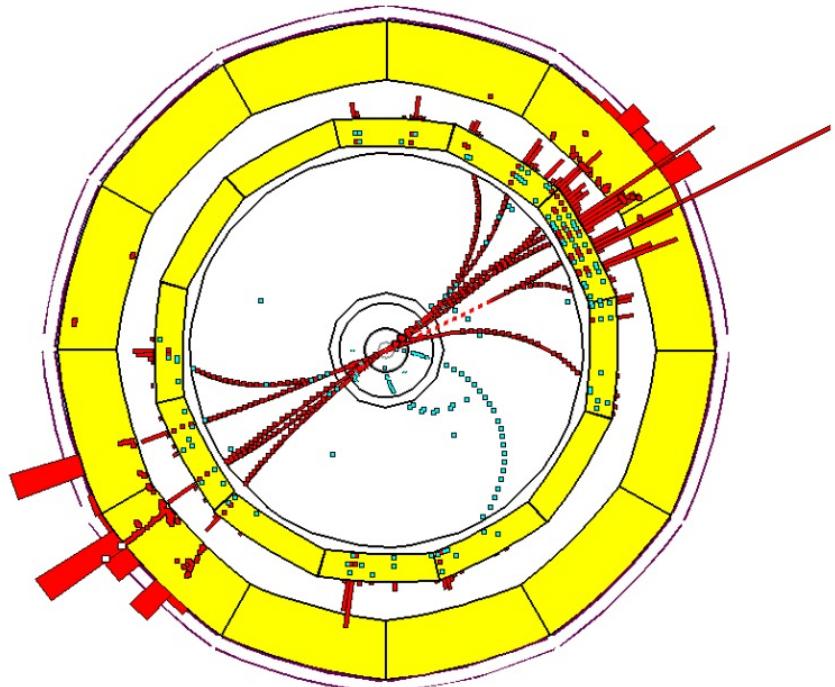
Initial state colors, boosted jets

# QCD jet study @ LEP1

- LEP1 : Z0 peak energy
  - $e^+ + e^- \rightarrow (\gamma, Z^0) \rightarrow q\bar{q}$
  - ✓  $\sqrt{s} = 91 \text{ GeV}$ 
    - Z mass
  - Pros and Cons for QCD study
    - ✓ Pros: No initial state color → clean jets
      - up to 90% gluon jet purity!!
    - ✓ Cons: difficult to access LEP data (already 30 years have passed)
      - there were ALEPH, DELPHI, L3, OPAL, however, can we access to the data??

“But! You’ve got a friend”

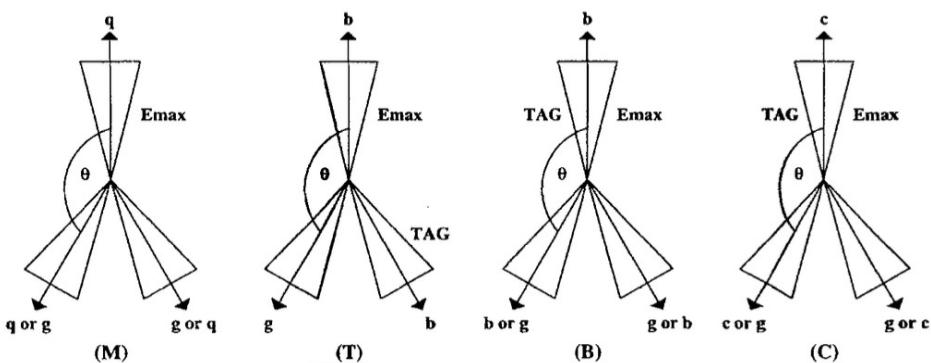
# ALEPH 2-jet vs 3-jet events



$$\begin{aligned} E(j_1) &> E(j_2) > E(j_3) \\ p_T(j_1) &> p_T(j_2) > p_T(j_3) \end{aligned}$$

# ALEPH study on gluon jets

- “Quark and Gluon jet properties in symmetric three-jet events”
- Phys. Lett. B 384, 353-364 (1996)
- ✓ after the ALEPH LEP1 runs



Sample	Number of jets	<i>uds</i>	<i>c</i>	<i>b</i>	<i>g</i>
M (J2 and J3)	45280	31.4	8.8	11.3	48.5
T (J2 or J3)	2071	6.1	1.7	2.2	90.0
B (J2 and J3)	872	2.5	3.0	44.5	50.0
C (J2 and J3)	40	0	35	15	50

19 September 1996  
PHYSICS LETTERS B  
Physics Letters B 384 (1996) 353-364

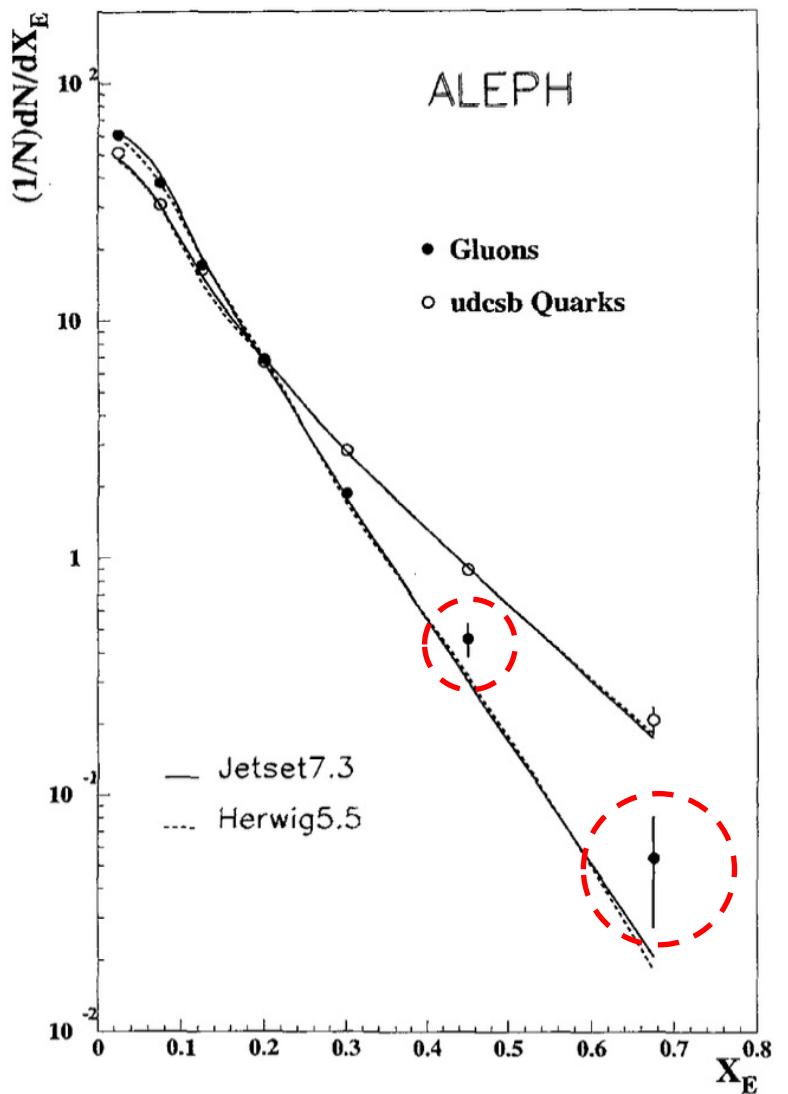
Quark and gluon jet properties in symmetric three-jet events  
ALEPH Collaboration

D. Buskulic<sup>a</sup>, D. Casper<sup>a</sup>, I. De Bonis<sup>a</sup>, D. Decamp<sup>a</sup>, P. Ghez<sup>a</sup>, C. Goy<sup>a</sup>, J.-P. Lees<sup>a</sup>, A. Lucotte<sup>a</sup>, M.-N. Minard<sup>a</sup>, P. Odier<sup>a</sup>, B. Pietrzyk<sup>a</sup>, M. Chmeissani<sup>b</sup>, J.M. Crespo<sup>b</sup>, I. Efthymopoulos<sup>b</sup>, E. Fernandez<sup>b</sup>, M. Fernandez-Bosman<sup>b</sup>, Ll. Garrido<sup>b,15</sup>, A. Juste<sup>b</sup>, M. Martinez<sup>b</sup>, S. Orteu<sup>b</sup>, A. Pacheco<sup>b</sup>, C. Padilla<sup>b</sup>, F. Palla<sup>b</sup>, A. Pascual<sup>b</sup>, J.A. Perlas<sup>b</sup>, I. Riu<sup>b</sup>, F. Sanchez<sup>b</sup>, F. Teubert<sup>b</sup>, A. Colaleo<sup>c</sup>, D. Creanza<sup>c</sup>, M. de Palma<sup>c</sup>, A. Farilla<sup>c</sup>, G. Gelao<sup>c</sup>, M. Girone<sup>c</sup>, G. Jaselli<sup>c</sup>, G. Maggi<sup>c,3</sup>, M. Maggi<sup>c</sup>, N. Marinelli<sup>c</sup>, S. Natali<sup>c</sup>, S. Nuzzo<sup>c</sup>, A. Ranieri<sup>c</sup>, G. Raso<sup>c</sup>, F. Romano<sup>c</sup>, F. Ruggieri<sup>c</sup>, G. Selvaggi<sup>c</sup>, L. Silvestris<sup>c</sup>, P. Tempesta<sup>c</sup>, G. Zito<sup>c</sup>, X. Huang<sup>d</sup>, J. Lin<sup>d</sup>, Q. Ouyang<sup>d</sup>, T. Wang<sup>d</sup>, Y. Xie<sup>d</sup>, R. Xu<sup>d</sup>, S. Xue<sup>d</sup>, J. Zhang<sup>d</sup>, L. Zhang<sup>d</sup>, W. Zhao<sup>d</sup>, R. Alemany<sup>e</sup>, A.O. Bazarko<sup>e</sup>, G. Bonvicini<sup>e,23</sup>, M. Cattaneo<sup>e</sup>, P. Comas<sup>e</sup>, P. Coyle<sup>e</sup>, H. Drevermann<sup>e</sup>, R.W. Forty<sup>e</sup>, M. Frank<sup>e</sup>, R. Hagelberg<sup>e</sup>, J. Harvey<sup>e</sup>, R. Jacobsen<sup>e,24</sup>, P. Janot<sup>e</sup>, B. Jost<sup>e</sup>, E. Kneringer<sup>e</sup>, J. Knobloch<sup>e</sup>, I. Lehraus<sup>e</sup>, E.B. Martin<sup>e</sup>, P. Mato<sup>e</sup>, A. Minten<sup>e</sup>, R. Miquel<sup>e</sup>, L.I.M. Mir<sup>e,2</sup>, L. Moneta<sup>e</sup>, T. Oest<sup>e</sup>, P. Palazzi<sup>e</sup>, J.R. Pater<sup>e,27</sup>, J.-F. Pusztaszeri<sup>e</sup>, F. Ranjard<sup>e</sup>, P. Rensing<sup>e</sup>, L. Rolandi<sup>e</sup>, D. Schlatter<sup>e</sup>, M. Schmelling<sup>e</sup>, O. Schneider<sup>e</sup>, W. Tejessy<sup>e</sup>, I.R. Tomalin<sup>e</sup>, A. Venturi<sup>f</sup>, H. Wachsmuth<sup>e</sup>, T. Wildish<sup>e</sup>, W. Witzeling<sup>e</sup>, J. Wotschack<sup>e</sup>, Z. Ajaltouni<sup>f</sup>, A. Barres<sup>f</sup>, C. Boyer<sup>f</sup>, A. Falvard<sup>f</sup>, P. Gay<sup>f</sup>, C. Guicheneuy<sup>f</sup>, P. Henrard<sup>f</sup>, J. Jousset<sup>f</sup>, B. Michel<sup>f</sup>, S. Monteil<sup>f</sup>, J.-C. Montret<sup>f</sup>, D. Pallin<sup>f</sup>, P. Perret<sup>f</sup>, F. Podlyski<sup>f</sup>, J. Proriol<sup>f</sup>, J.-M. Rossignol<sup>f</sup>, T. Fearnley<sup>f</sup>, J.B. Hansen<sup>g</sup>, J.D. Hansen<sup>g</sup>, J.R. Hansen<sup>g</sup>, P.H. Hansen<sup>g</sup>, B.S. Nilsson<sup>g</sup>, A. Wääänänen<sup>g</sup>, A. Kyriakis<sup>h</sup>, C. Markou<sup>h</sup>, E. Simopoulou<sup>h</sup>, I. Siotis<sup>h</sup>, A. Vayaki<sup>h</sup>, K. Zachariadou<sup>h</sup>, A. Blondel<sup>i,21</sup>, G. Bonneauud<sup>i</sup>, J.C. Brient<sup>i</sup>, P. Bourdon<sup>i</sup>, A. Rouge<sup>i</sup>, M. Rumpf<sup>i</sup>, R. Tanaka<sup>i</sup>, A. Valassi<sup>i,6</sup>, M. Verderi<sup>i</sup>, H. Videau<sup>i,21</sup>, D.J. Candlin<sup>j</sup>, M.I. Parsons<sup>j</sup>, E. Focardi<sup>k</sup>, G. Parrini<sup>k</sup>, M. Corden<sup>l</sup>, M. Delfino<sup>i,12</sup>, C. Georgopoulos<sup>l</sup>, D.E. Jaffe<sup>l</sup>, A. Antonelli<sup>m</sup>, G. Bencivenni<sup>m</sup>, G. Bologna<sup>m,4</sup>, F. Bossi<sup>m</sup>, P. Campana<sup>m</sup>, G. Capon<sup>m</sup>, V. Chiarella<sup>m</sup>, G. Felici<sup>m</sup>, P. Laurelli<sup>m</sup>, G. Mannocchi<sup>m,5</sup>, F. Murta<sup>m</sup>, G.P. Murta<sup>m</sup>, L. Passalacqua<sup>m</sup>, M. Pepe-Altarelli<sup>m</sup>, L. Curtis<sup>n</sup>, S.J. Dorris<sup>n</sup>, A.W. Halley<sup>n</sup>, I.G. Knowles<sup>n</sup>, J.G. Lynch<sup>n</sup>, V. O'Shea<sup>n</sup>, C. Raine<sup>n</sup>, P. Reeves<sup>n</sup>, J.M. Scarr<sup>n</sup>, K. Smith<sup>n</sup>, A.S. Thompson<sup>n</sup>, F. Thomson<sup>n</sup>, S. Thorn<sup>n</sup>, R.M. Turnbull<sup>n</sup>, U. Becker<sup>o</sup>, O. Braun<sup>o</sup>, C. Geweniger<sup>o</sup>, G. Graefe<sup>o</sup>, P. Hanke<sup>o</sup>, V. Hepp<sup>o</sup>, E.E. Kluge<sup>o</sup>, A. Putzer<sup>o</sup>, B. Rensch<sup>o</sup>, M. Schmidt<sup>o</sup>, J. Sommer<sup>o</sup>, H. Stenzel<sup>o</sup>, K. Tittel<sup>o</sup>, S. Werner<sup>o</sup>, M. Wunsch<sup>o</sup>, D. Abbaneo<sup>p</sup>, R. Beuselinck<sup>p</sup>, D.M. Binnie<sup>p</sup>, W. Cameron<sup>p</sup>, D.J. Colling<sup>p</sup>,

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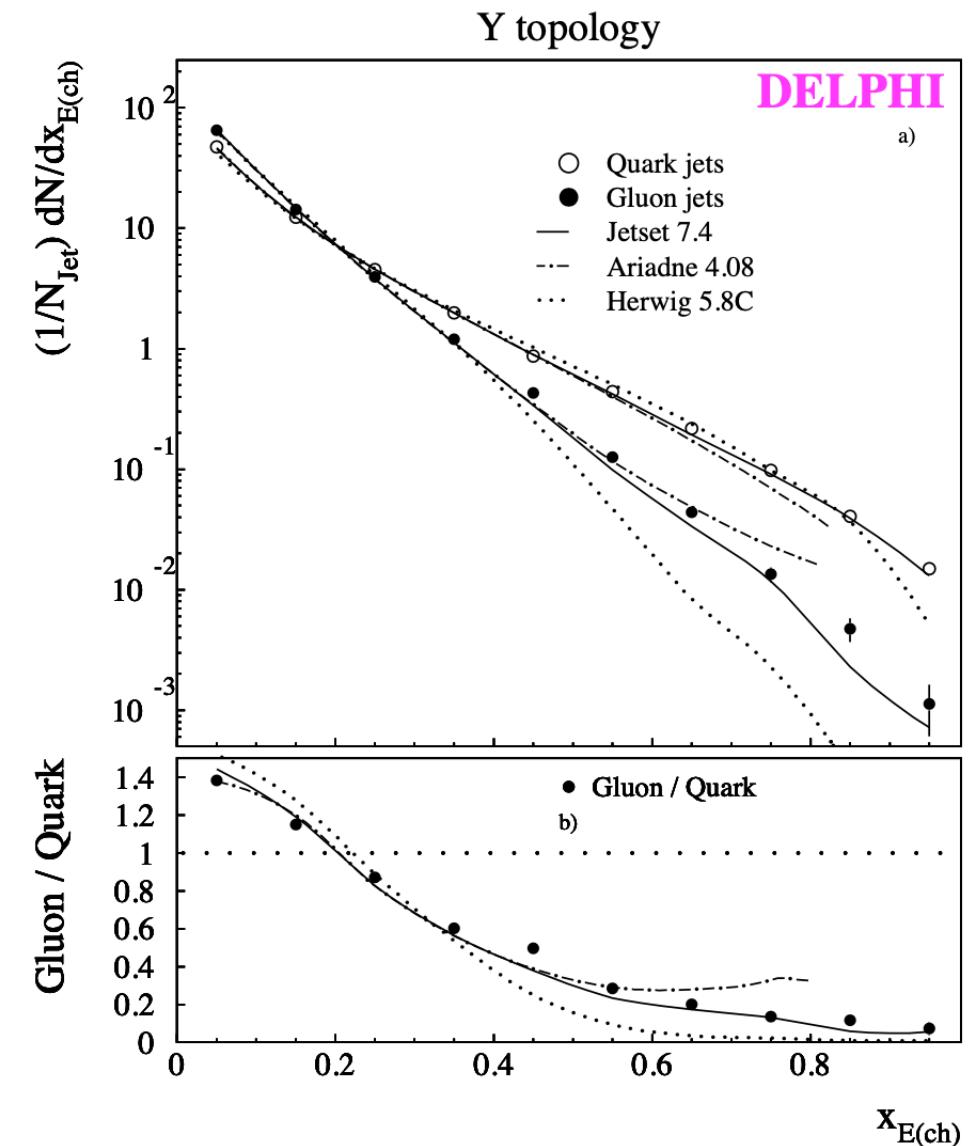
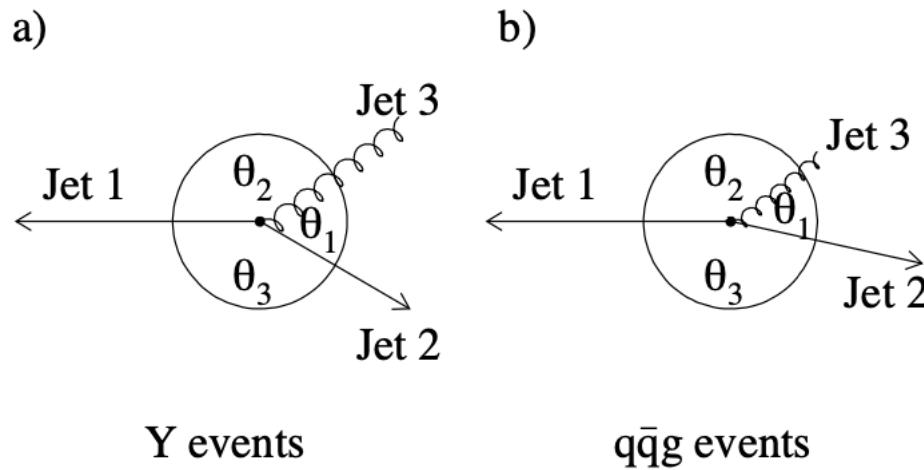
# Fragmentation function $X_E$ (ALEPH)

- Fragmentation function ( $X_E$ )
  - like Björken  $x$  (track's energy share)
    - ✓  $X_E = \frac{E_p}{E_j} \rightarrow \frac{1}{N} \frac{dN}{dX_E}$  as a function of  $X_E$
- ALEPH observations (PLB-384, 353-364)
  - Gluons are softer
    - ✓ broader
  - Quarks are harder
    - ✓ collimated
- Were there any ALEPH gluon anomalies?
  - Quarks
    - ✓ good agreement between Data and MC
  - Gluons
    - ✓ Neurons → excess in high  $X_E$
    - ✓ Chromons → excess in low  $X_E$



# Fragmentation function $X_E$ (DELPHI)

- “Measurement of the gluon fragmentation function and a comparison of the scaling violation in gluon and quark jets”
  - Eur. Phys. J. C. 13, 573-589 (2000)
  - ✓ LEP1 runs 1992-1995



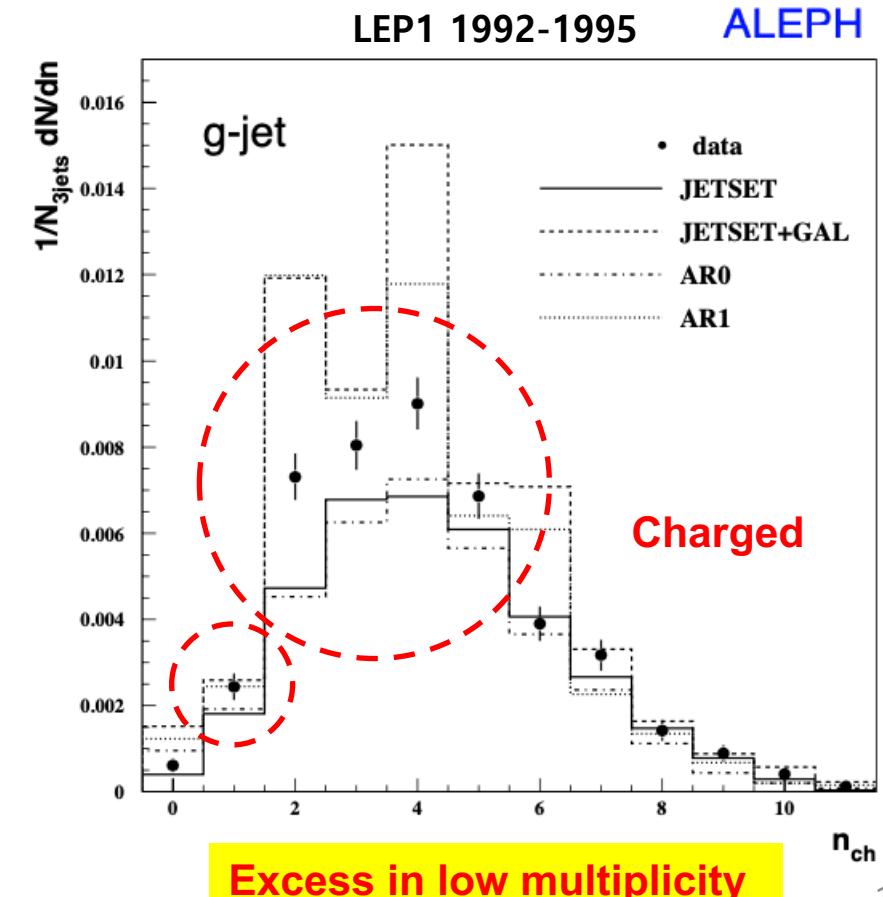
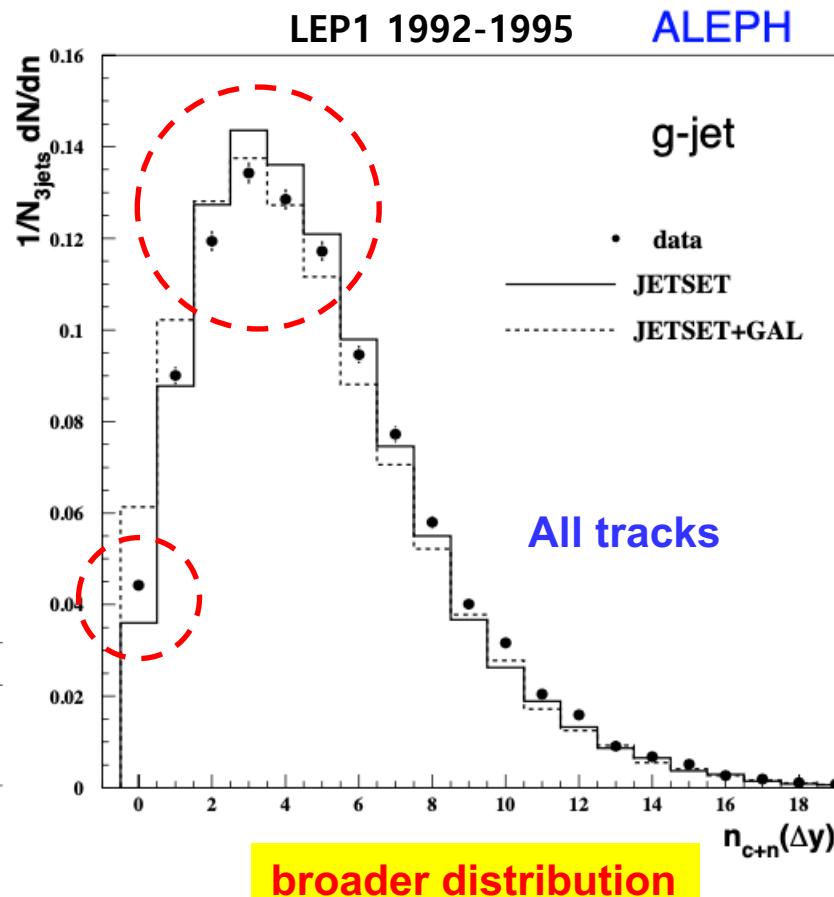
# Anomaly in charged multiplicity

- “Test of colour reconnection models using three-jet events in hadronic Z decays”
  - Eur. Phys. J. C. 48, 685-698 (2006)

QCD model	million events
JETSET	7.6
JETSET+GAL, $R_0 = 0.04$	0.5
JETSET+GAL, $R_0 = 0.10$	3.4
JETSET+BE <sub>32</sub>	0.5
ARIADNE AR0	3.4
ARIADNE AR1	3.4
HERWIG	0.5
HERWIG CR	0.5

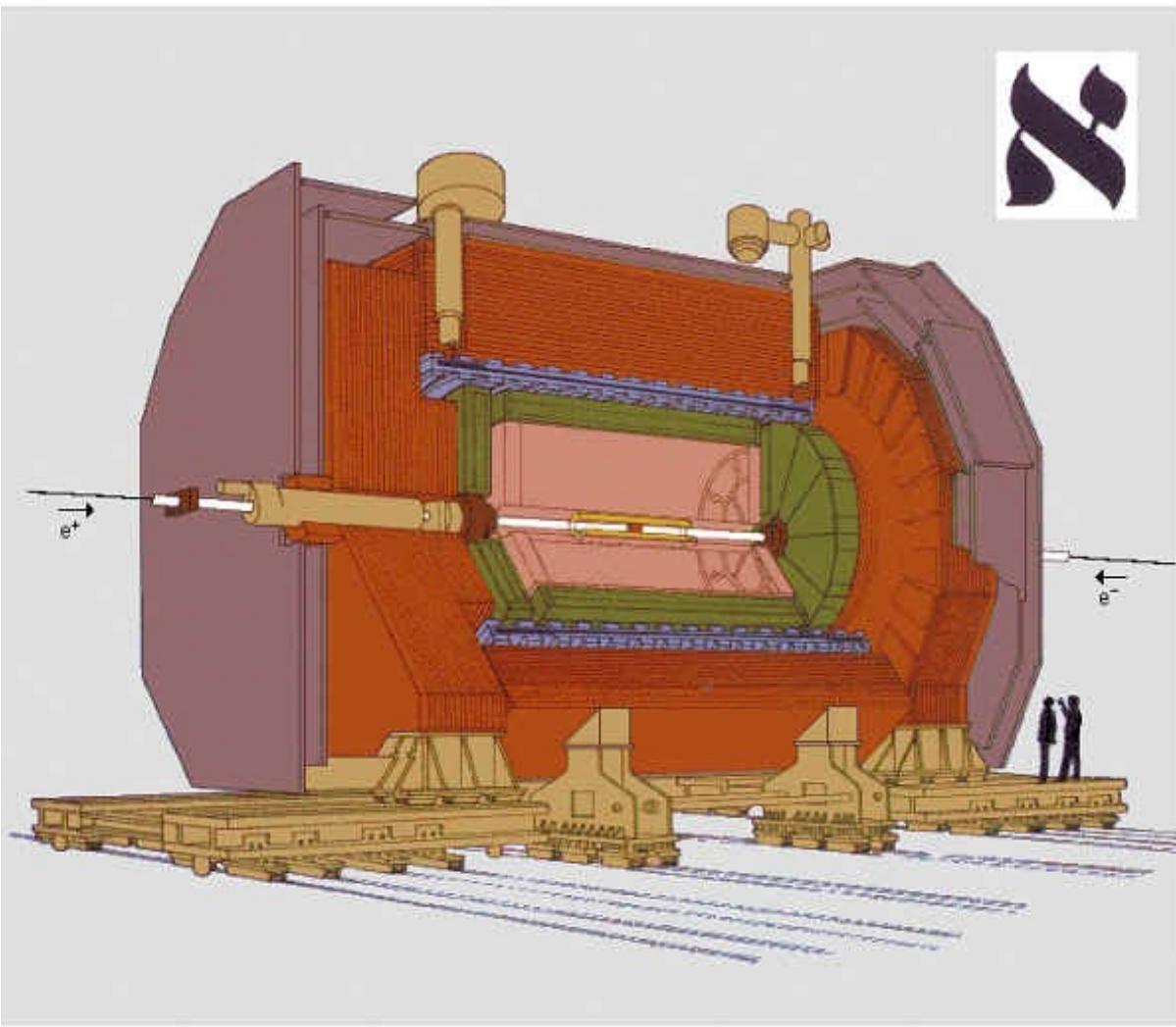
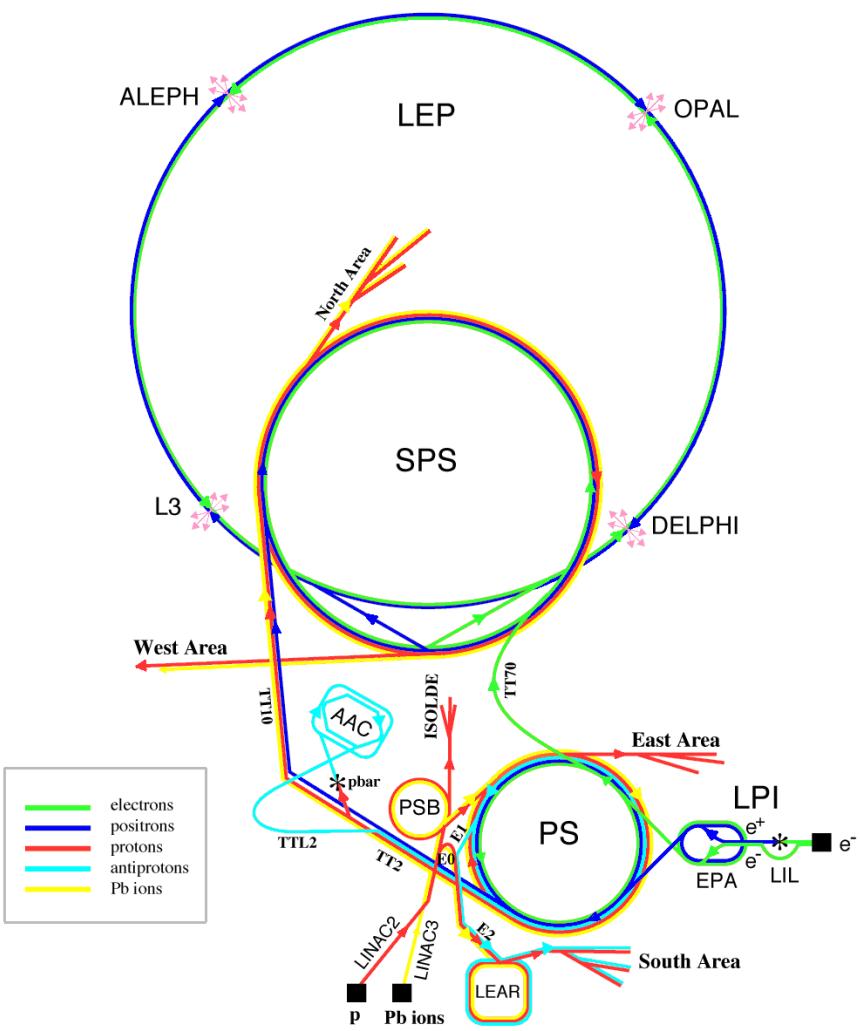
  

jet number	$\langle E_{\text{jet}} \rangle$ , GeV	spread, GeV	$\langle n_{\text{ch}} \rangle$	$P_g$
1	40.8	2.7	8.81	0.059
2	32.7	4.5	8.28	0.248
3	17.7	5.1	7.02	0.693



# Resurrection of the ALEPH data

# ALEPH @ LEP1



- Vertex Detector
- Inner Tracking Chamber
- Time Projection Chamber
- Electromagnetic Calorimeter
- Superconducting Magnet Coil
- Hadron Calorimeter
- Muon Chambers
- Luminosity Monitors

$$\text{TPC} : \sigma\left(\frac{1}{p_T}\right) \sim 0.06\%$$

$$\text{ECAL} : \frac{\sigma(E)}{E} \sim \frac{0.18}{\sqrt{E}} + 1\%$$

$$\text{HCAL} : \frac{\sigma(E)}{E} \sim \frac{0.85}{\sqrt{E}}$$

# MIT archived data

- Data Samples

- Archived  $e^+e^-$  data

- ✓  $e^+e^-$  collisions in LEP1 @  $\sqrt{s} = 91$  GeV
    - a total of 3.3M hadronic events

- MC Samples

- Archived  $e^+e^-$  simulation data

- ✓  $e^+e^-$  collisions in PYTHIA 6.1 @  $\sqrt{s} = 91$  GeV
    - a total of ~ 1M hadronic events

stable GEN particles from PYTHIA

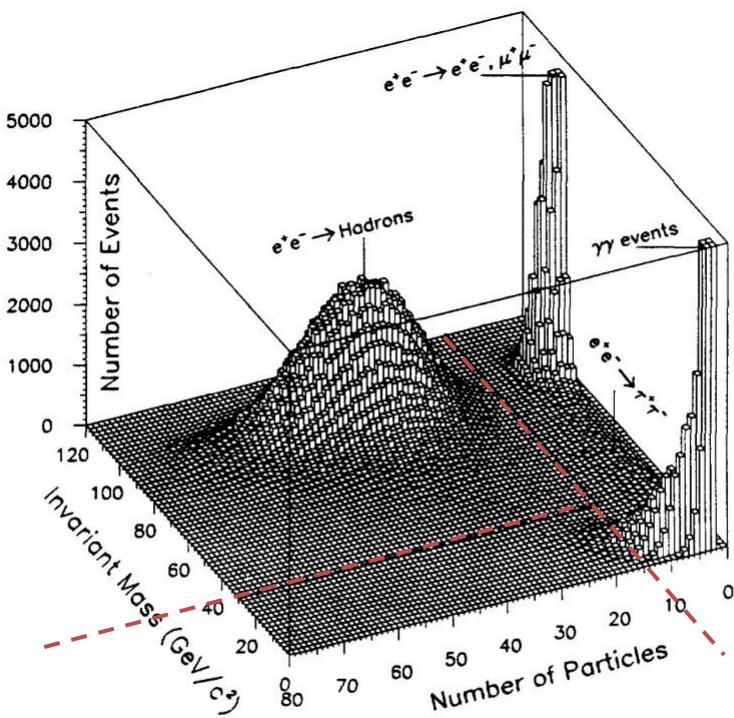
reconstructed particles after GALEPH (GEANT3+ALEPH) detector simulation

	Data	MC
1991	230,513	
1992	551,474	
1993	538,601	
1994	1,365,440	
1995	595,095	
Total	3,281,123	973,769

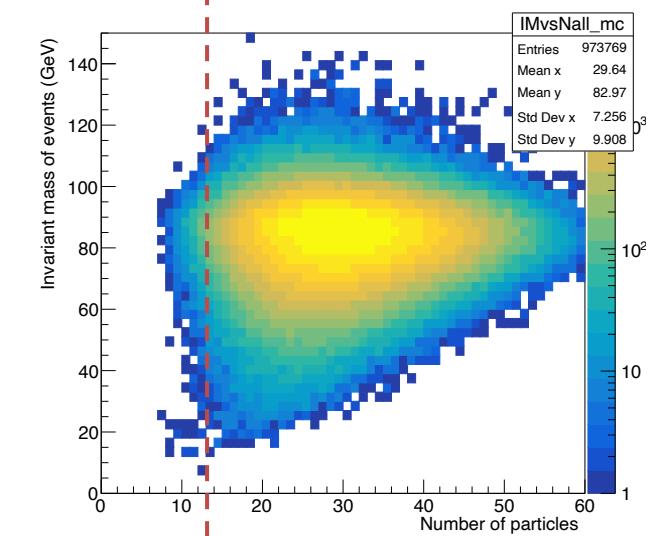
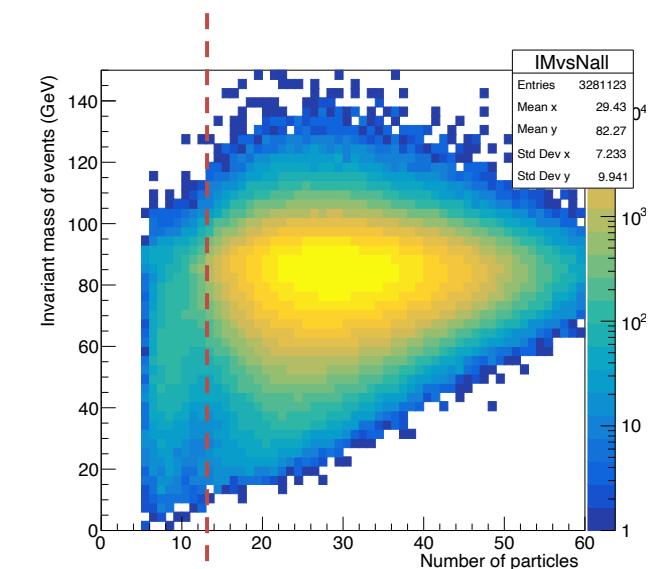
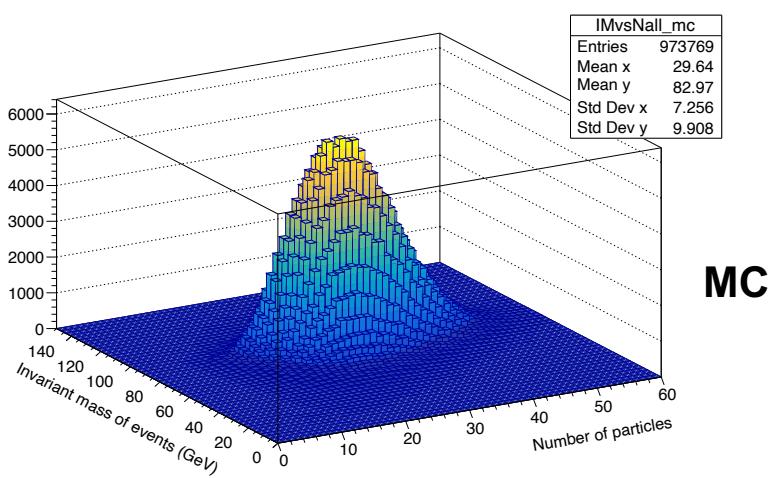
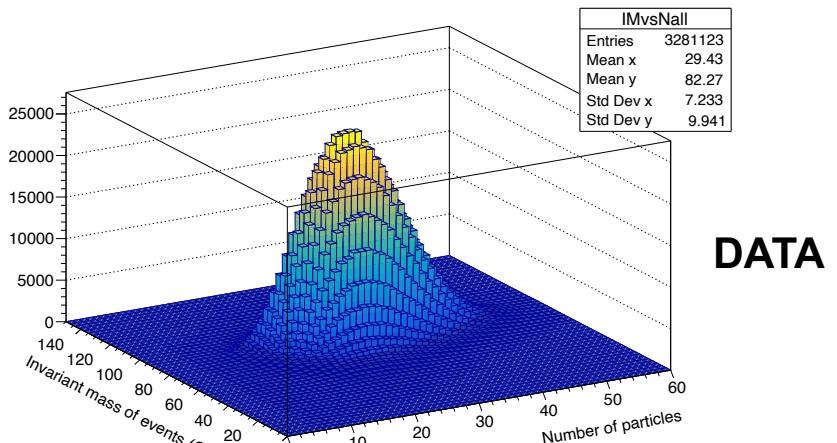


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# Hadronic event selection



$N_{trk} \geq 13 \text{ & } M_{inv} \geq 35$

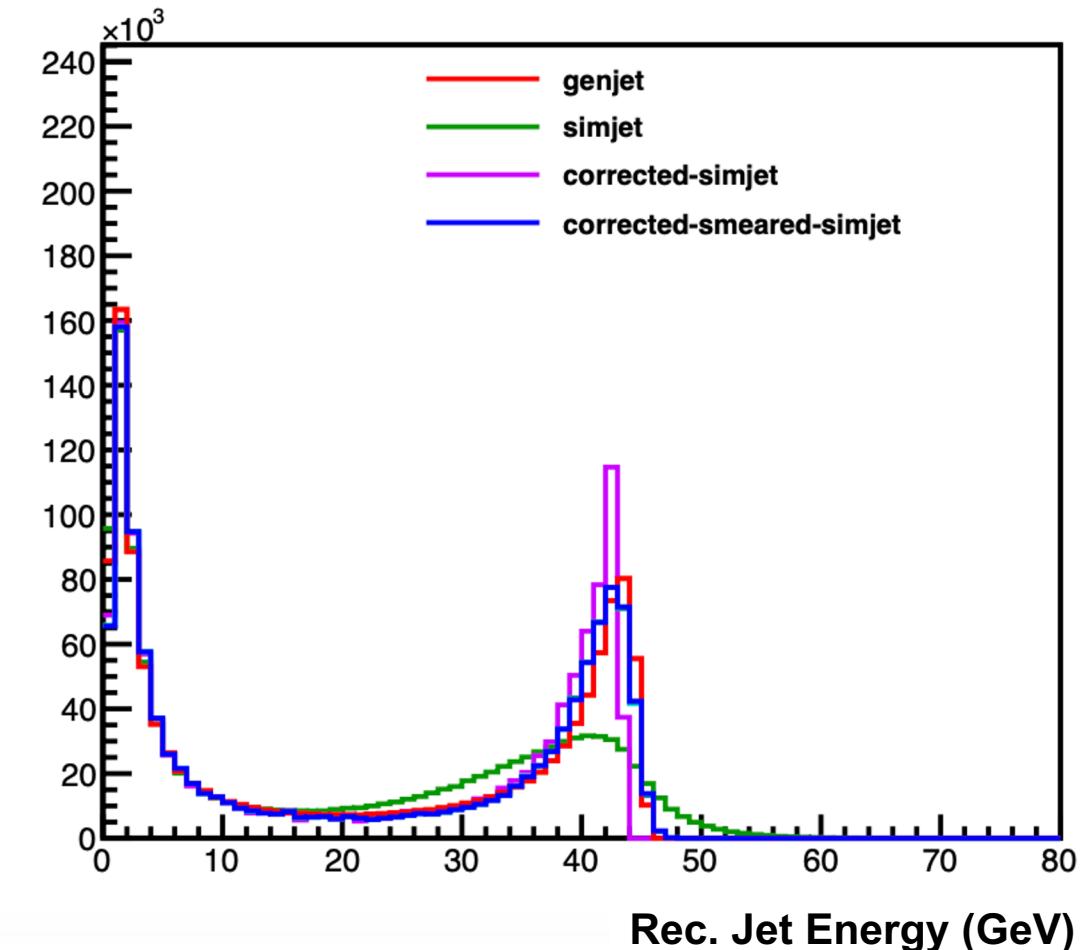
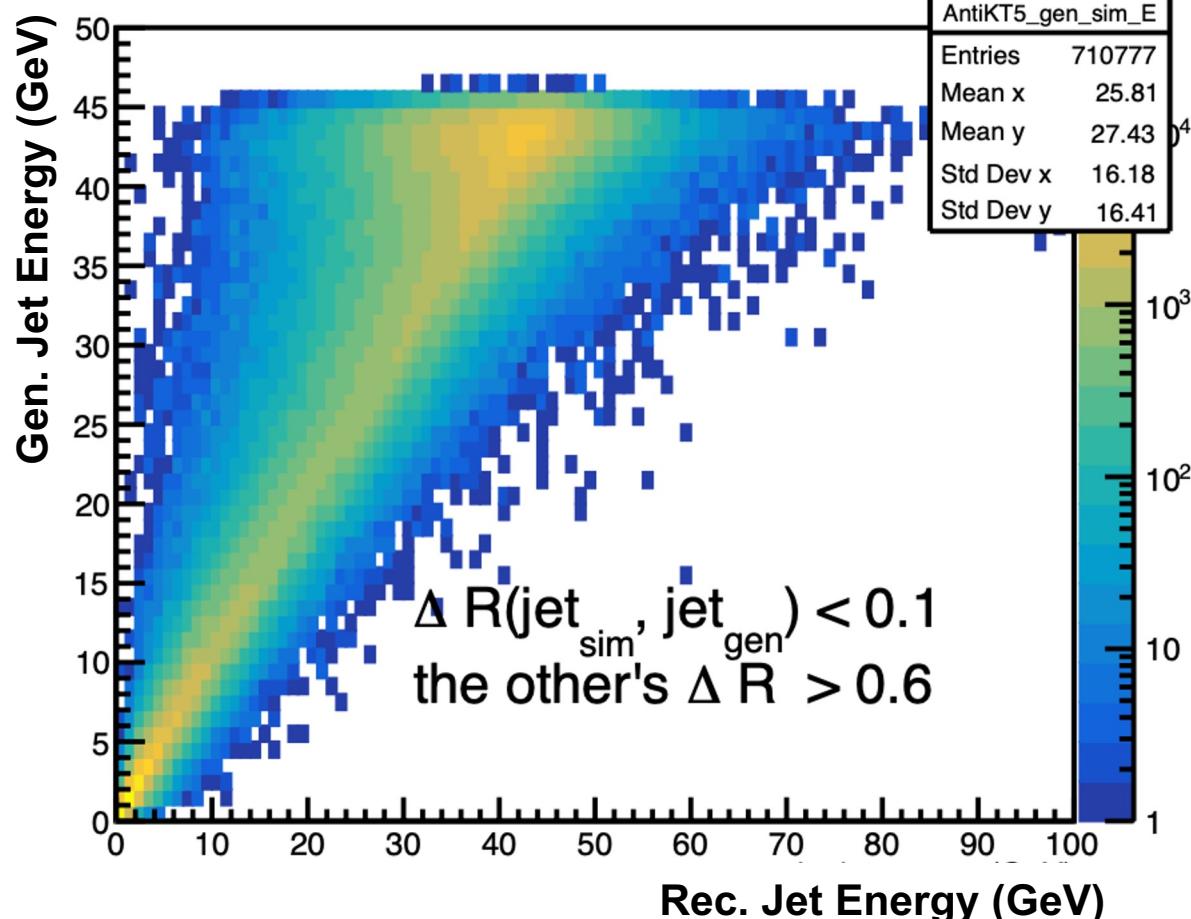


# Jet reconstruction

- Jet reconstruction
  - ALEPH used Durham ( $k_\perp$ ) algorithm
  - We use Fastjet (version 3.4.1.3)
    - ✓ Jet algorithm : "Anti-Kt algorithm"
      - using 4-vector (pt, eta, phi, mass) of particles with an opening angle  $R = 0.5$
- Jet energy resolution
  - ALEPH achieved a JER of  $\frac{\sigma_E}{E} \sim 0.6/\sqrt{E/\text{GeV}}$  with Particle Flow
    - ✓ TPC+VTX:  $\sigma\left(\frac{1}{p_T}\right) \sim 0.006 (\text{GeV}/c)^{-1}$
    - ✓ ECAL:  $\frac{\sigma_E}{E} \sim 0.18/\sqrt{E/\text{GeV}}$ , HCAL:  $\frac{\sigma_E}{E} \sim 0.85/\sqrt{E/\text{GeV}}$ 
      - e.g. 36 GeV jet  $\Rightarrow$  10% error

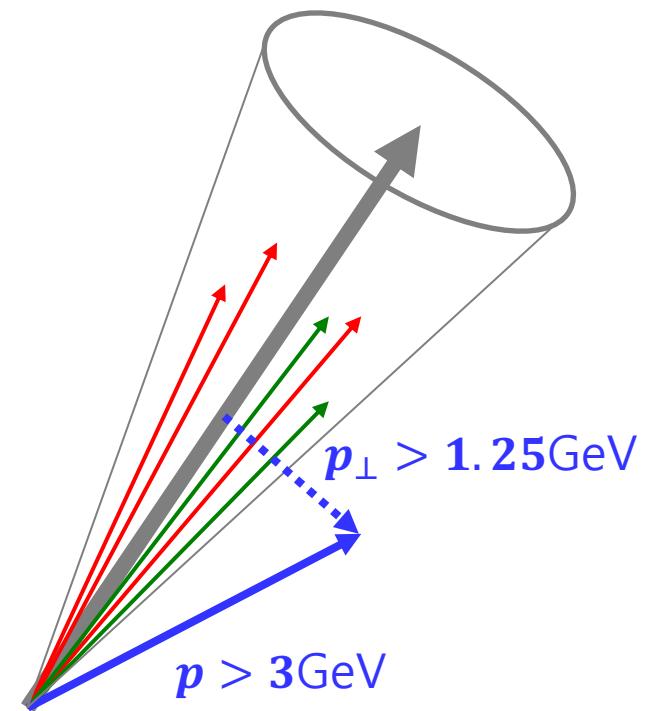
# Jet energy correction

Correction factors = Egen/Erec



# B tagging

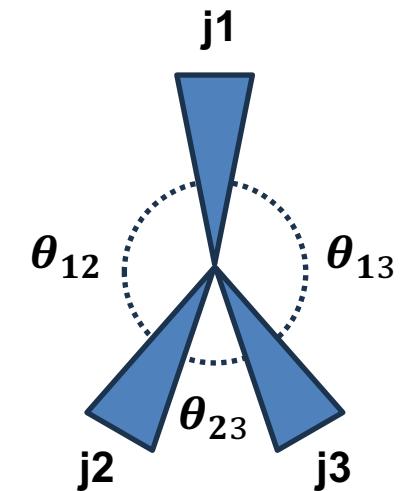
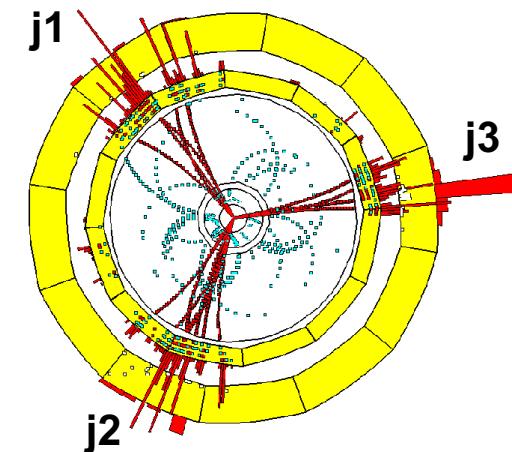
- B tagging
  - Lepton tagging method
    - ✓ lepton with  $p > 3\text{GeV}$  &  $p_{\perp} > 1.25\text{GeV}$ 
      - $p_{\perp}$  calculated from the jet reconstructed without lepton
  - Tagging efficiency
    - ✓ According to Jetset study
      - b :  $88.9 \pm 1.2\%$
      - c :  $6.1 \pm 1.1\%$
      - uds:  $\sim 5\%$



# Event selections & observables

# 3-jet event selection

- S1 : Hadronic event selection
  - $N_{trk} \geq 13, N_{chg} \geq 5, \Sigma E_{chg} \geq 15 \text{ GeV}$
- S2 : 3-jet event selection
  - $N_{jets} = 3 \quad (E_j > 10 \text{ GeV})$
- S3 : 3-jet geometry selection
  - $\theta_{12}, \theta_{13}, \theta_{23}$
  - ✓  $\theta_{12}, \theta_{13} = 150^\circ \pm 7.5^\circ$  (Tight: ALEPH revisited)
    - $E_{quark} \sim E_{gluon} \sim 24 \text{ GeV}$
  - ✓  $0.5 < \theta_{23} < \min(\theta_{12}, \theta_{13})$  (Loose: This analysis)
- S4 : b-jet tagging
  - Only one of J2 and J3 tagged as b-jet



# Event selections

3-jet event selection	Data	MC
S0: 1991-1995	2900317	973769
S1: $N_{trk} \geq 13$ , $N_{chg} \geq 5$ , $\sum E_{chg} \geq 15 \text{ GeV}$	2876547	767769
S2: $N_{jets} = 3$ ( $E_j > 10 \text{ GeV}$ )	432565	119655
<b>Tight selection (ALEPH revisited)</b>		
S3: $\theta_{12}, \theta_{13} = 150^\circ \pm 7.5^\circ$	26649	8137
S4: one b-jet among J2 & J3	3064	907
<b>Loose selection (This analysis)</b>		
S3: $0.5 < \theta_{23} < \min(\theta_{12}, \theta_{13})$	341152	96406
S4: one b-jet among J2 & J3	40984	11243

# Analysis and results

# What observables should we measure?

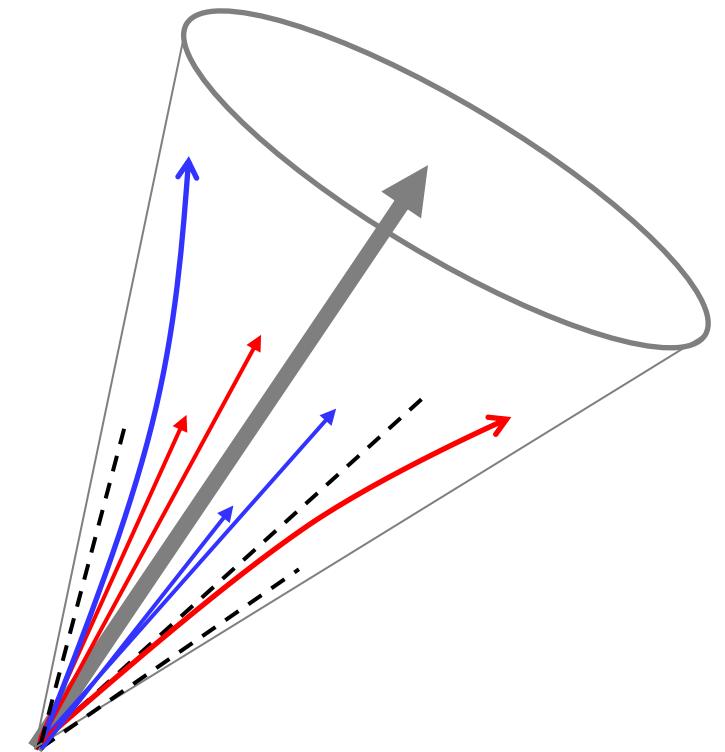
- Particle multiplicities
  - number of particles in a jet

✓  $n_{all}$ ,  $n_{chg}$ ,  $n_{neu}$   
– particles: charged + neutrals  
charged: Track objects ( $e^\pm, \mu^\pm, \pi^\pm$ )  
neutrals : ECAL/HCAL objects ( $\gamma, n$ )

- Fragmentation functions

- particle's energy fraction:  $X_E = \frac{E_p}{E_{jet}}$

✓  $X_E^{all}$ ,  $X_E^{chg}$ ,  $X_E^{neu}$   
– particles: charged + neutrals

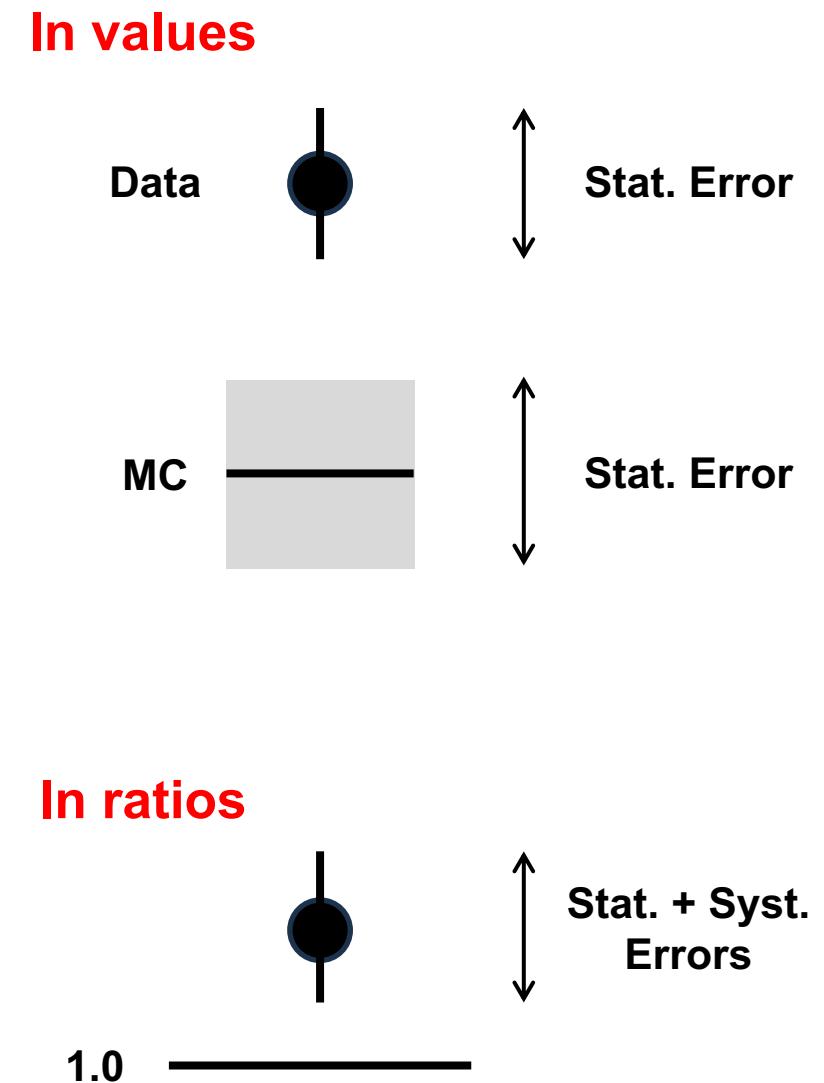


# Systematic error estimation

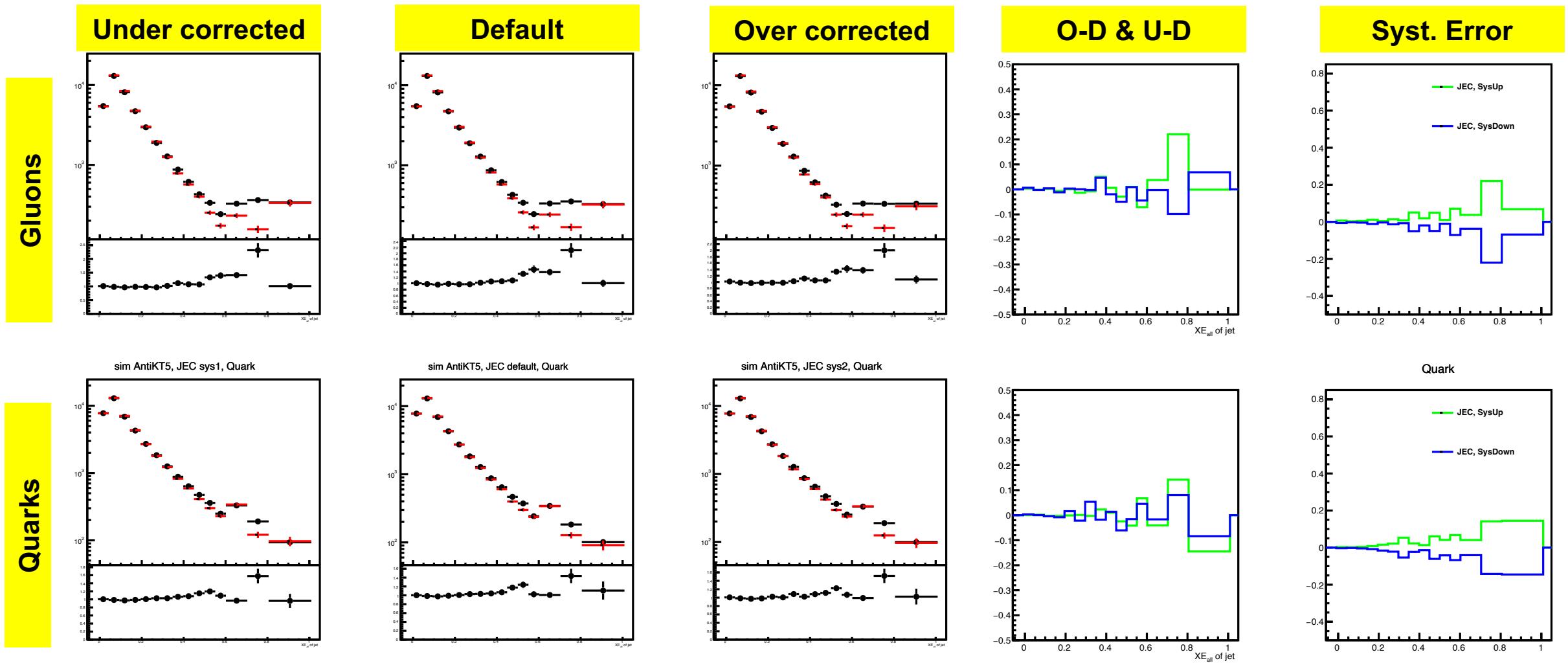
- Basics: Model comparison study
  - To see the data anomaly
  - To test the MC model
- Most of detector related systematic effects cancel out in the ratio plot

$$R = \frac{A}{B} \quad \Rightarrow \quad \frac{\sigma_R}{R} = \sqrt{\left(\frac{\sigma_A}{A}\right)^2 + \left(\frac{\sigma_B}{B}\right)^2}, \quad \sigma_R^{Stat} = R \sqrt{\left(\frac{\sigma_A}{A}\right)^2 + \left(\frac{\sigma_B}{B}\right)^2}$$

$$\sigma_R = \sqrt{(\sigma_R^{Stat})^2 + (\sigma_R^{Syst})^2}$$

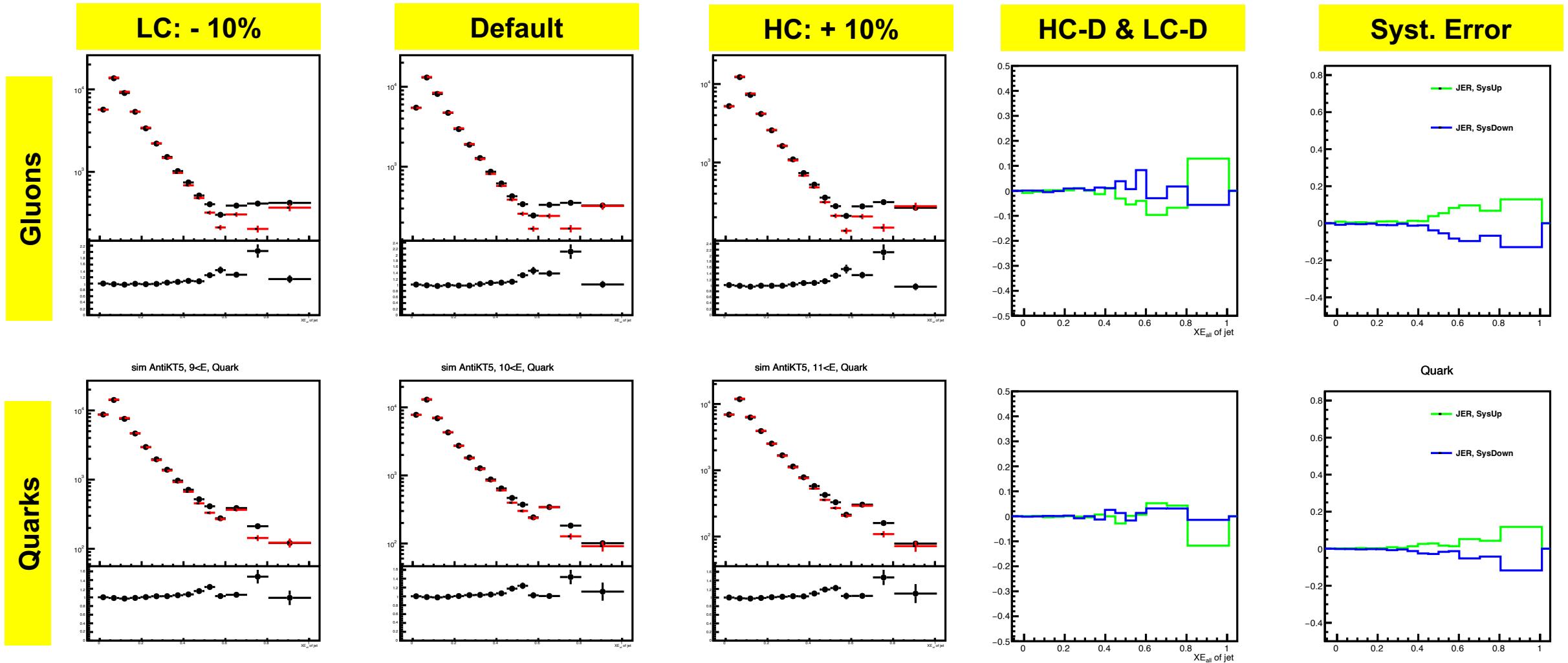


# Systematic error: Jet Energy Correction



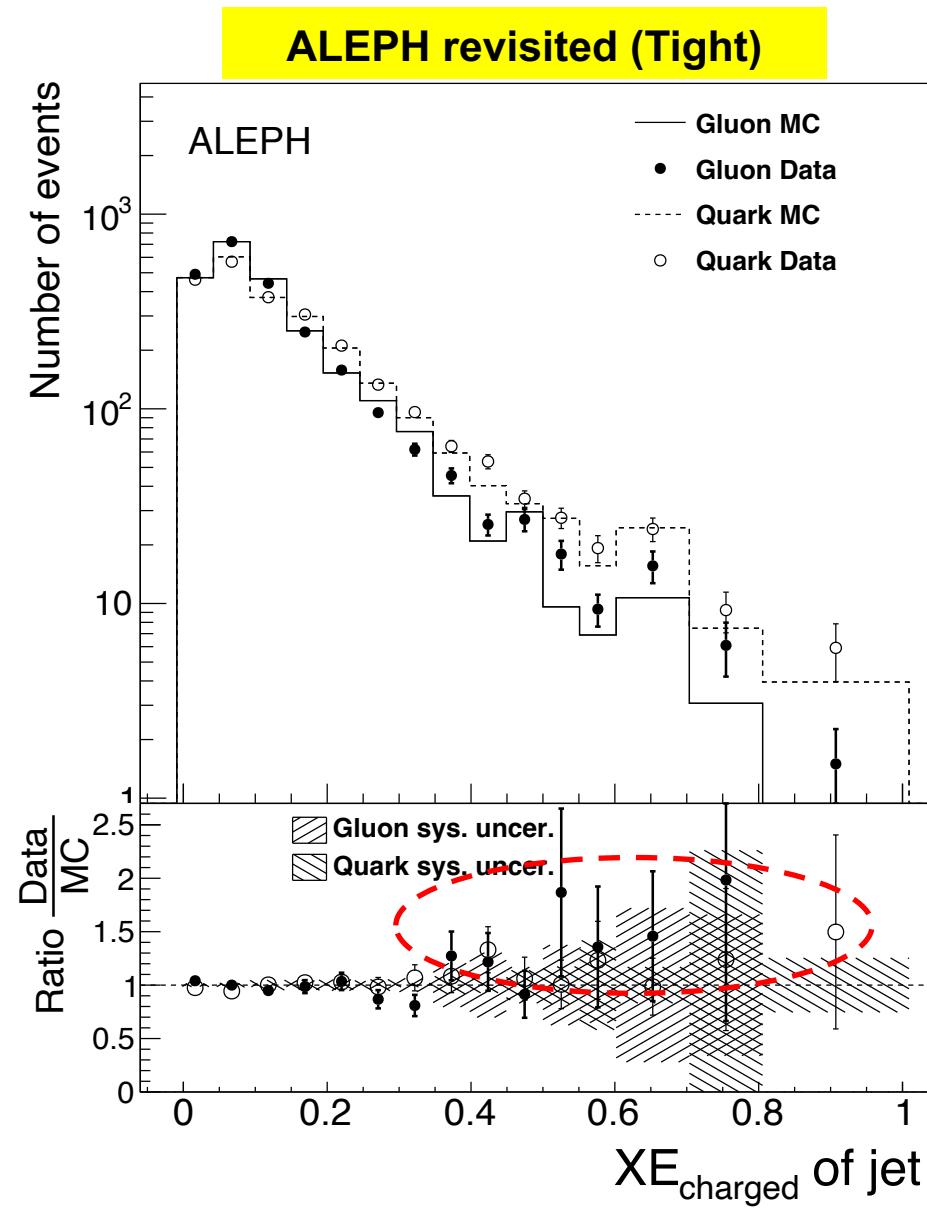
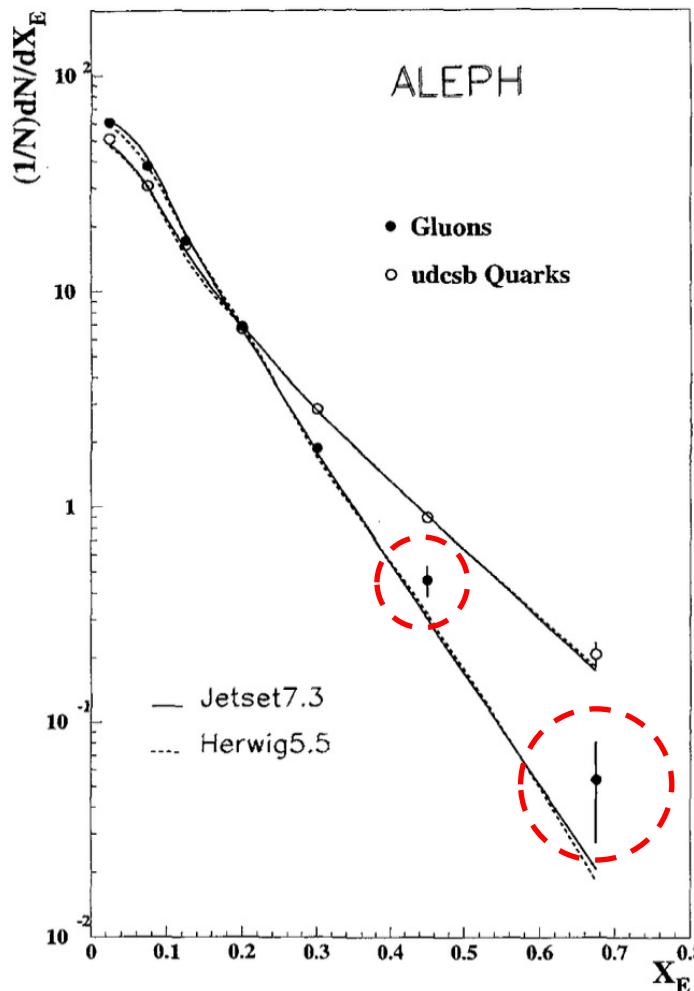
Based on statistical errors from JEC factor estimation (Profile histogram)

# Systematic error: Jet Energy Resolution



Based on 10% variations on Jet Energy Cut : Default  $\pm 1$  GeV

# Fragmentation function $X_E$ : charged only (Tight cut)



No meaningful excesses were found from our ALEPH reanalysis, within statistical and systematical errors.

**3064 events only**

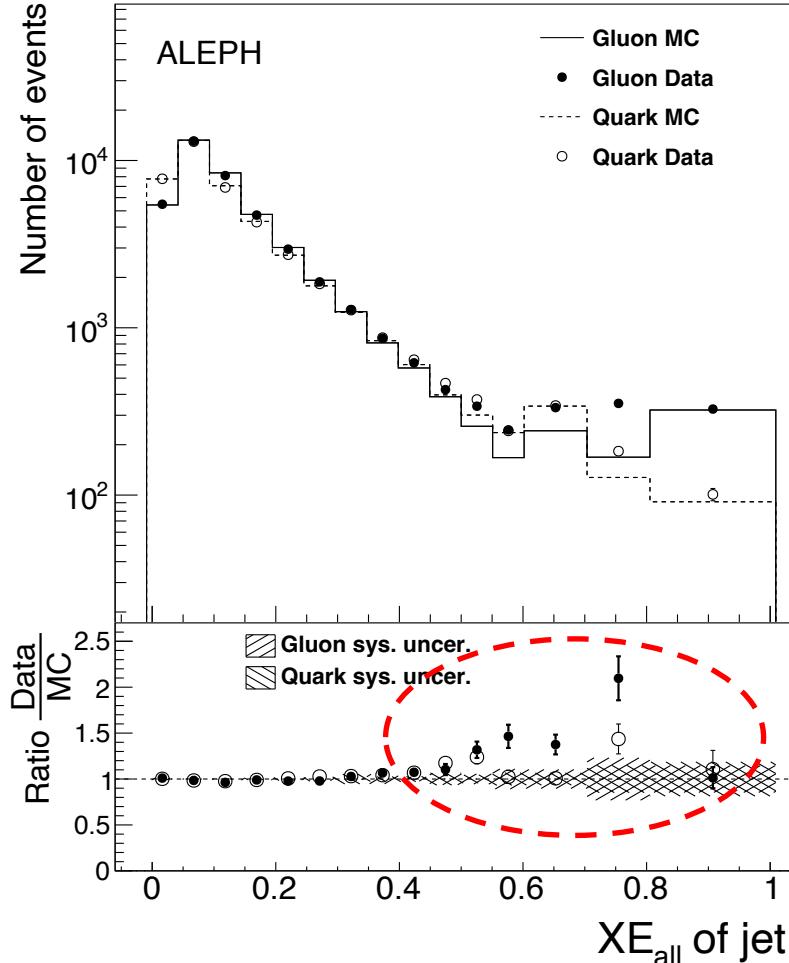
- more statistics needed
- control systematic errors

⇒ Use the loose cut

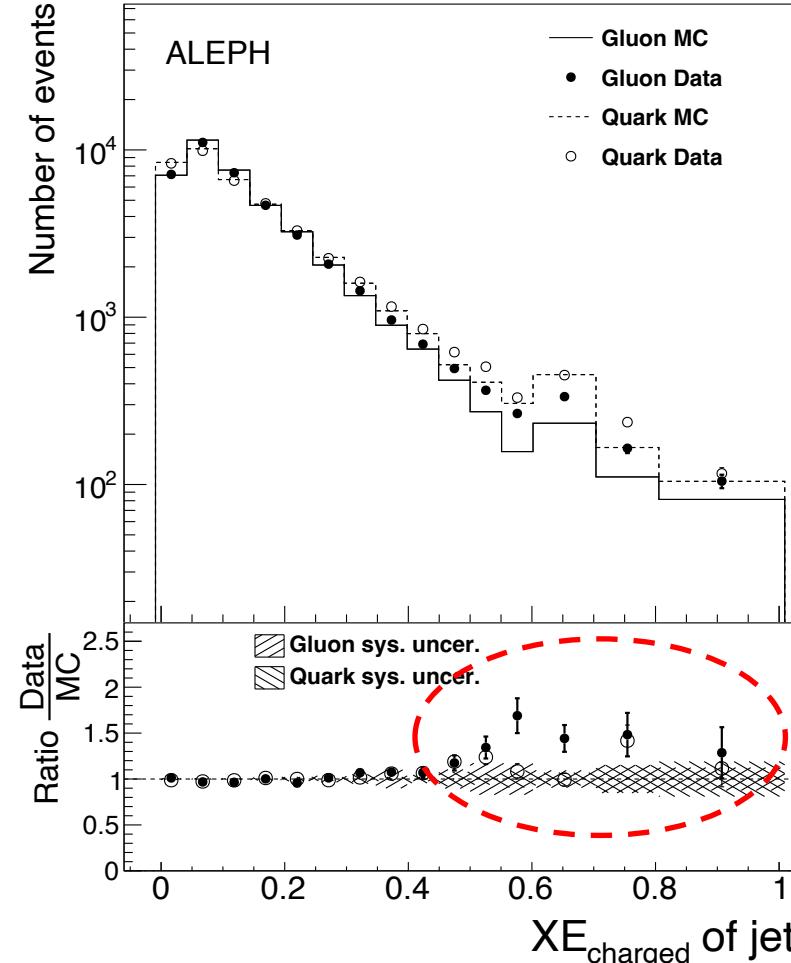
**40984 events are available**

# $X_E$ : charged vs neutral (Loose cut)

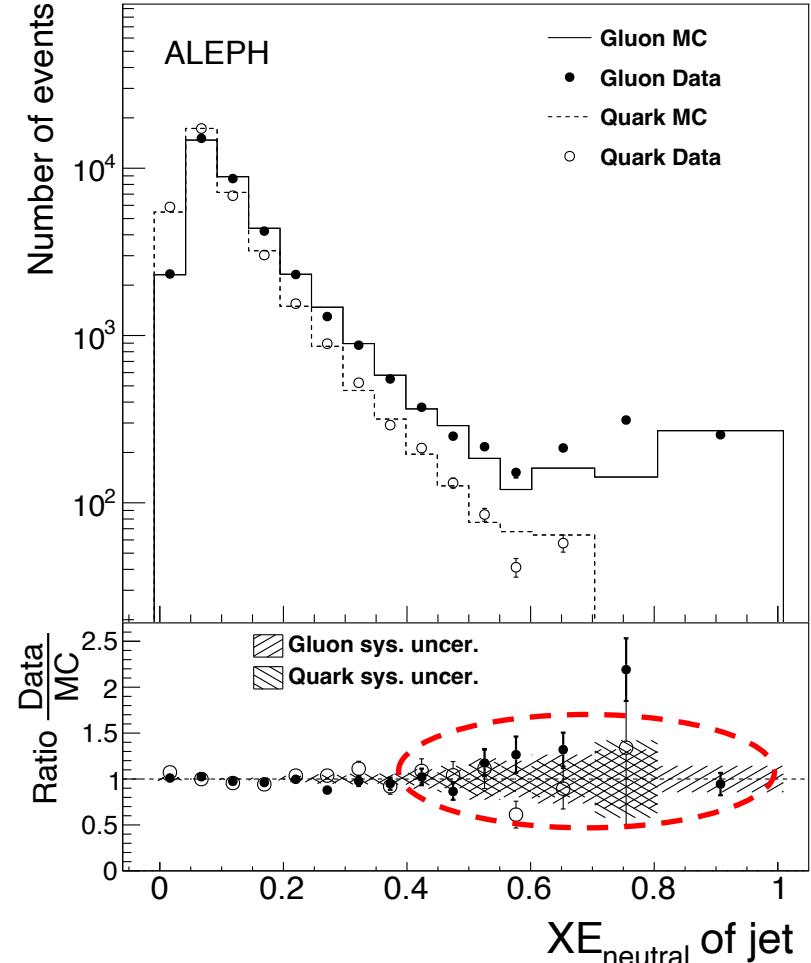
All particles



Charged particles

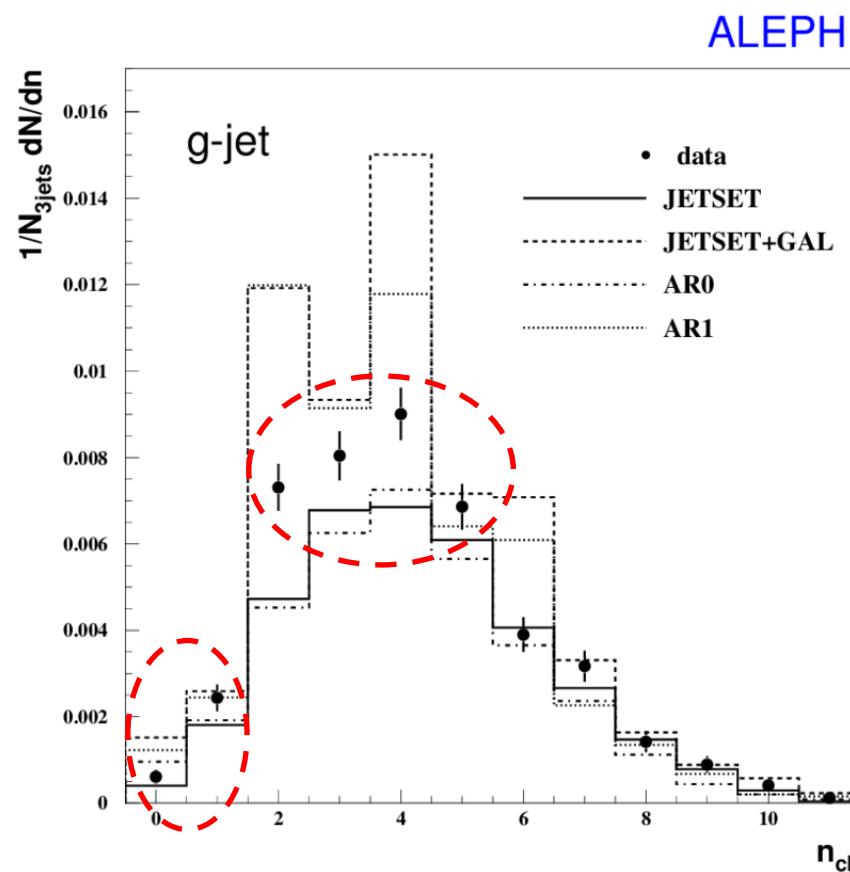


Neutral particles

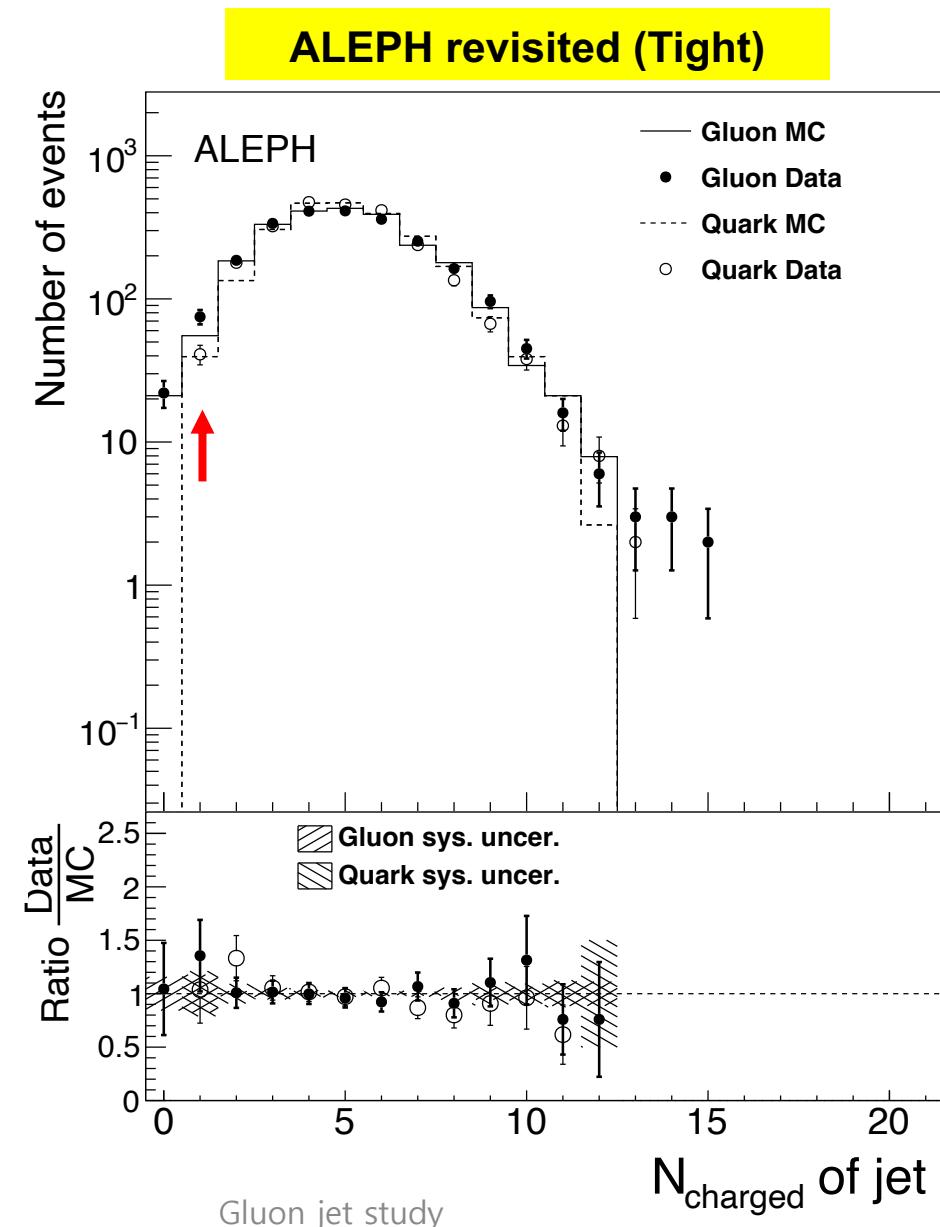


Particles with high energy share are more often in gluons in data than those of in MC.

# Charged particle multiplicity (Tight cut)

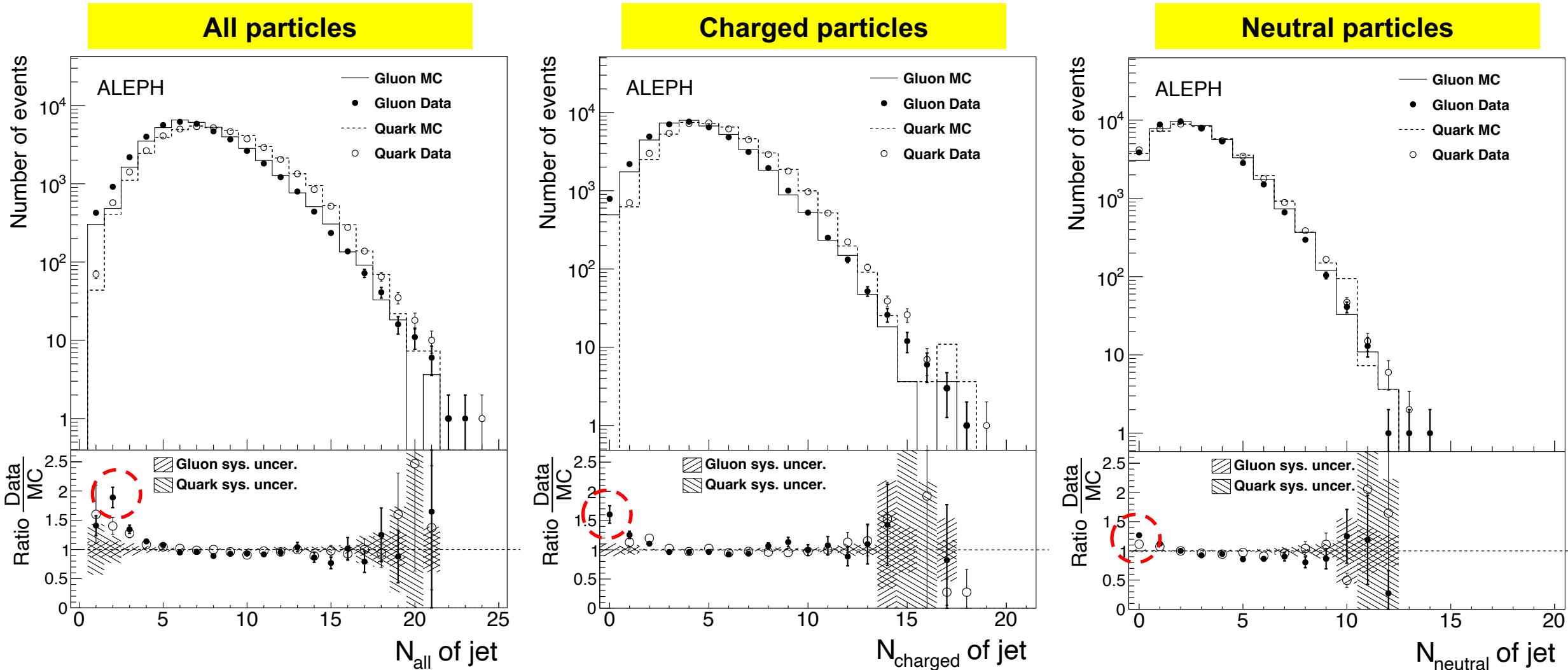


???



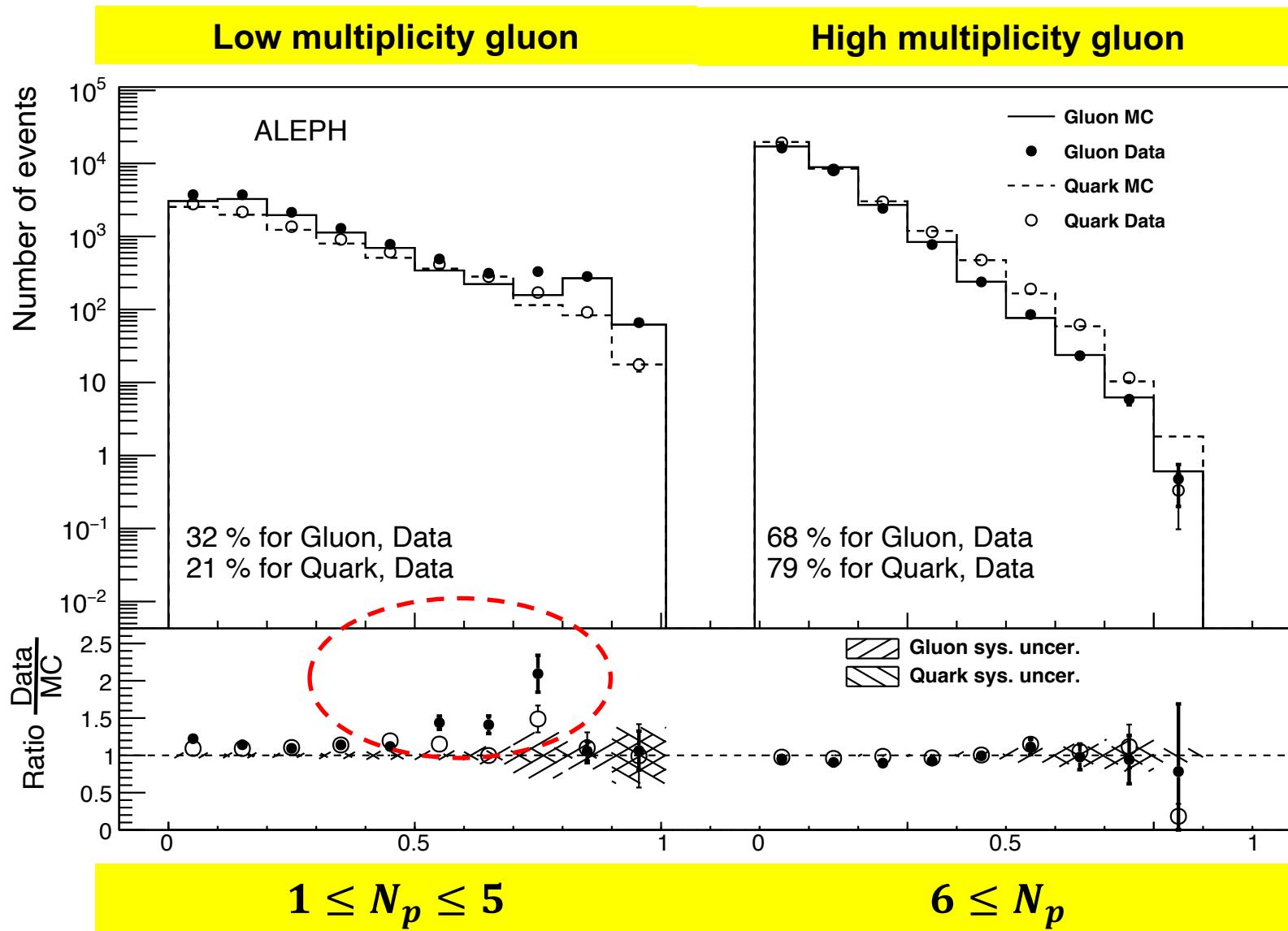
No meaningful discrepancies were found from our ALEPH reanalysis, within statistical and systematical errors.

# Gluon jet particle multiplicity anomaly? (Loose cut)



Gluons with two tracks! Gluons without charged tracks.

# Correlation between $X_E$ and (all)



# Event display

Year: 1991, EventNo: 806, RunNo: 12716

# chromon like

jets  
anti-b-tagged  
b-tagged

particles  
charged hadrons  
leptons  
photons  
neutral hadrons

jet1 (36.9, 4, 1, 2, 1)  
jet2 (27, 3, 1, 4, 0)  
jet3 (11.3, 3, 0, 4, 0)  
jet4 (1.4, 0, 0, 0, 1)

b-jet  
(37 GeV)

b-jet  
(27 GeV)

gluon-jet  
(11 GeV)  
 $n=7$

Year: 1993, EventNo: 5149, RunNo: 22501

# chromon like

jets  
anti-b-tagged  
b-tagged

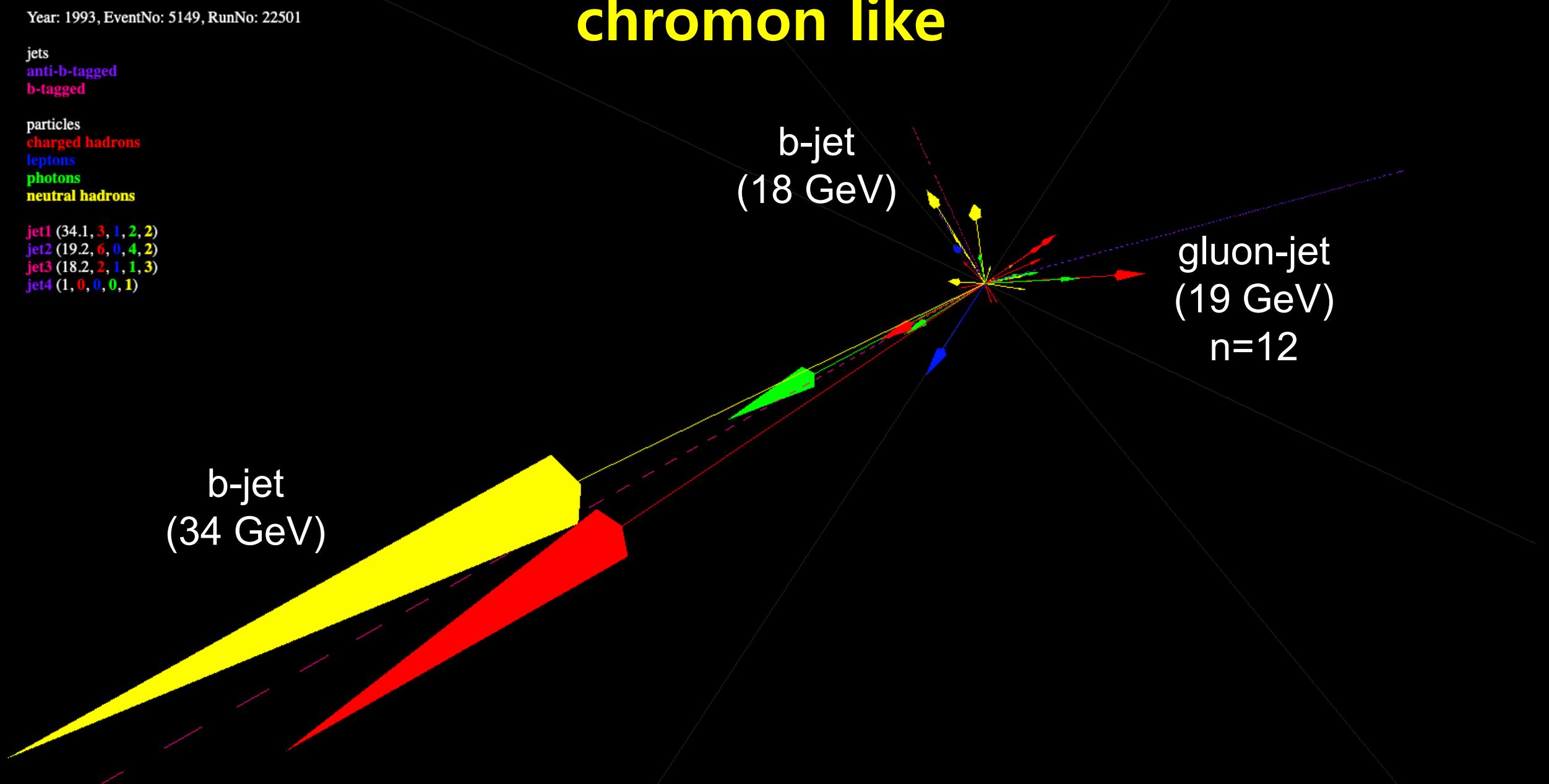
particles  
charged hadrons  
leptons  
photons  
neutral hadrons

jet1 (34.1, 3, 1, 2, 2)  
jet2 (19.2, 6, 0, 4, 2)  
jet3 (18.2, 2, 1, 1, 3)  
jet4 (1, 0, 0, 0, 1)

b-jet  
(18 GeV)

gluon-jet  
(19 GeV)  
n=12

b-jet  
(34 GeV)



Year: 1993, EventNo: 6991, RunNo: 22504

# chromon like

jets  
anti-b-tagged  
b-tagged

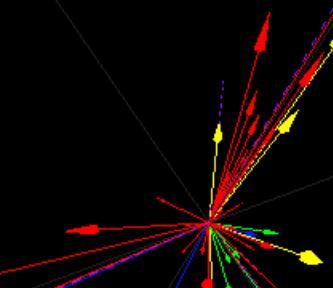
particles  
charged hadrons  
leptons  
photons  
neutral hadrons

jet1 (28.8, 5, 2, 0, 0)  
jet2 (26.3, 4, 2, 5, 1)  
jet3 (19.7, 8, 0, 0, 2)  
jet4 (1.9, 0, 0, 0, 1)

b-jet  
(29 GeV)

b-jet  
(26 GeV)

gluon-jet  
(20 GeV)  
 $n=10$



Year: 1994, EventNo: 7452, RunNo: 27200

# neuron like

jets  
anti-b-tagged  
b-tagged

particles  
charged hadrons  
leptons  
photons  
neutral hadrons

jet1 (41.9, 4, 1, 4, 0)  
jet2 (25.6, 5, 1, 1, 0)  
jet3 (20.6, 0, 0, 1, 0)  
jet4 (1.3, 1, 0, 0, 0)

b-jet  
(42 GeV)

gluon-jet  
(21 GeV)  
 $n=1$

b-jet  
(26 GeV)

Year: 1994, EventNo: 1288, RunNo: 27510

# neuron like

jets  
anti-b-tagged  
b-tagged

particles  
charged hadrons  
leptons  
photons  
neutral hadrons

jet1 (39.3, 3, 1, 7, 2)  
jet2 (23.9, 6, 2, 6, 1)  
jet3 (18.1, 1, 0, 0, 1)  
jet4 (0.9, 0, 0, 1, 0)

b-jet  
(39 GeV)

b-jet  
(24 GeV)

gluon-jet  
(18 GeV)  
n=2

Year: 1994, EventNo: 2251, RunNo: 28512

# neuron like

jets  
anti-b-tagged  
b-tagged

particles  
charged hadrons  
leptons  
photons  
neutral hadrons

jet1 (43.1, 5, 1, 5, 2)  
jet2 (28.3, 5, 1, 2, 1)  
jet3 (16.6, 1, 0, 1, 1)  
jet4 (0.7, 1, 0, 0, 0)  
jet5 (0.5, 1, 0, 0, 0)

b-jet  
(43 GeV)

gluon-jet  
(17 GeV)  
n=3

b-jet  
(28 GeV)

Year: 1994, EventNo: 9035, RunNo: 28848

# neuron like

jets  
anti-b-tagged  
b-tagged

particles  
charged hadrons  
leptons  
photons  
neutral hadrons

jet1 (30.3, 3, 1, 4, 1)  
jet2 (18.5, 5, 1, 3, 0)  
jet3 (12.3, 1, 0, 2, 0)  
jet4 (0.8, 2, 0, 0, 0)

b-jet  
(30 GeV)

b-jet  
(19 GeV)

gluon-jet  
(12 GeV)  
n=3

Year: 1994, EventNo: 5975, RunNo: 29653

# neuron like

jets  
anti-b-tagged  
b-tagged

particles  
charged hadrons  
leptons  
photons  
neutral hadrons

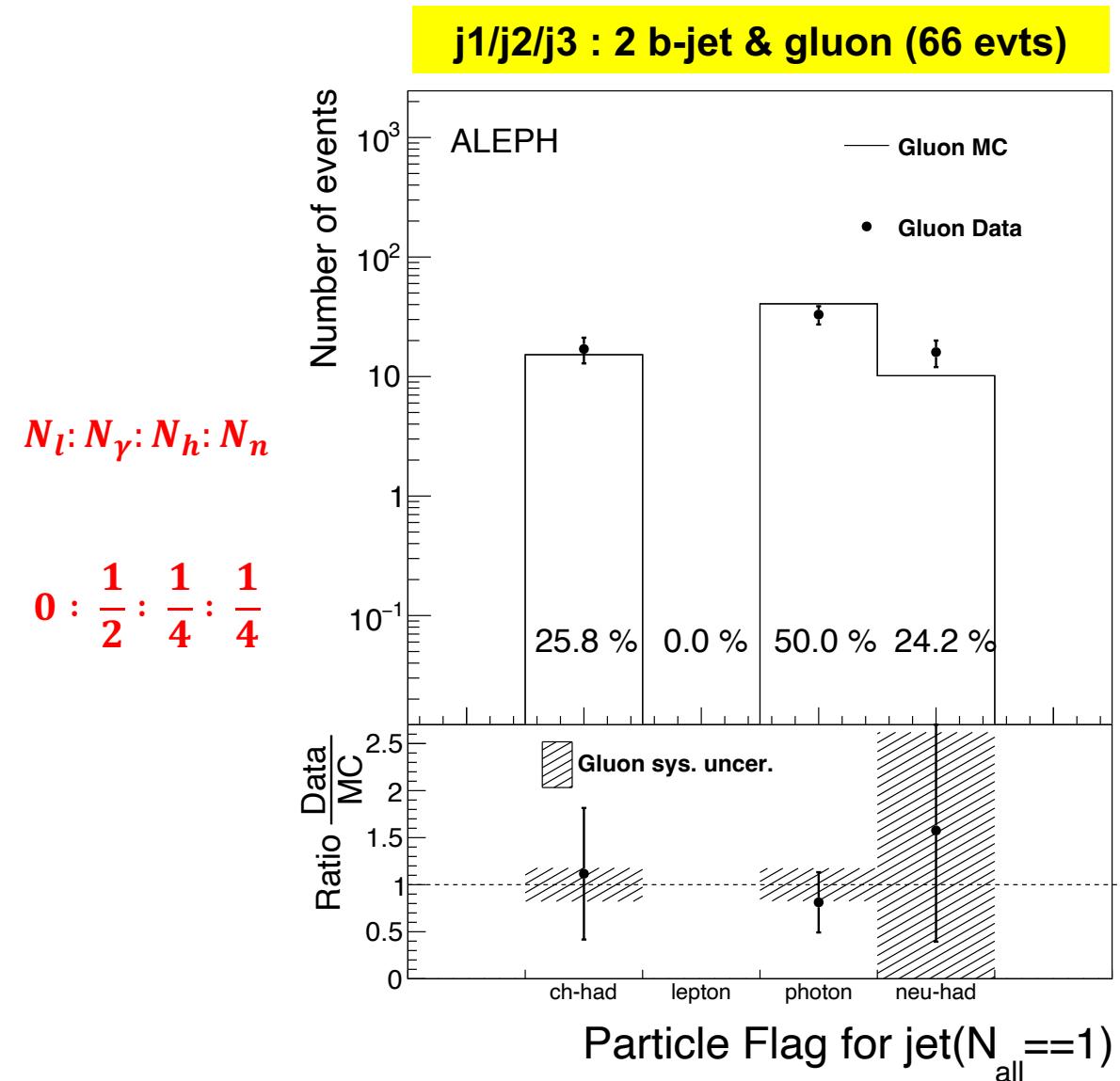
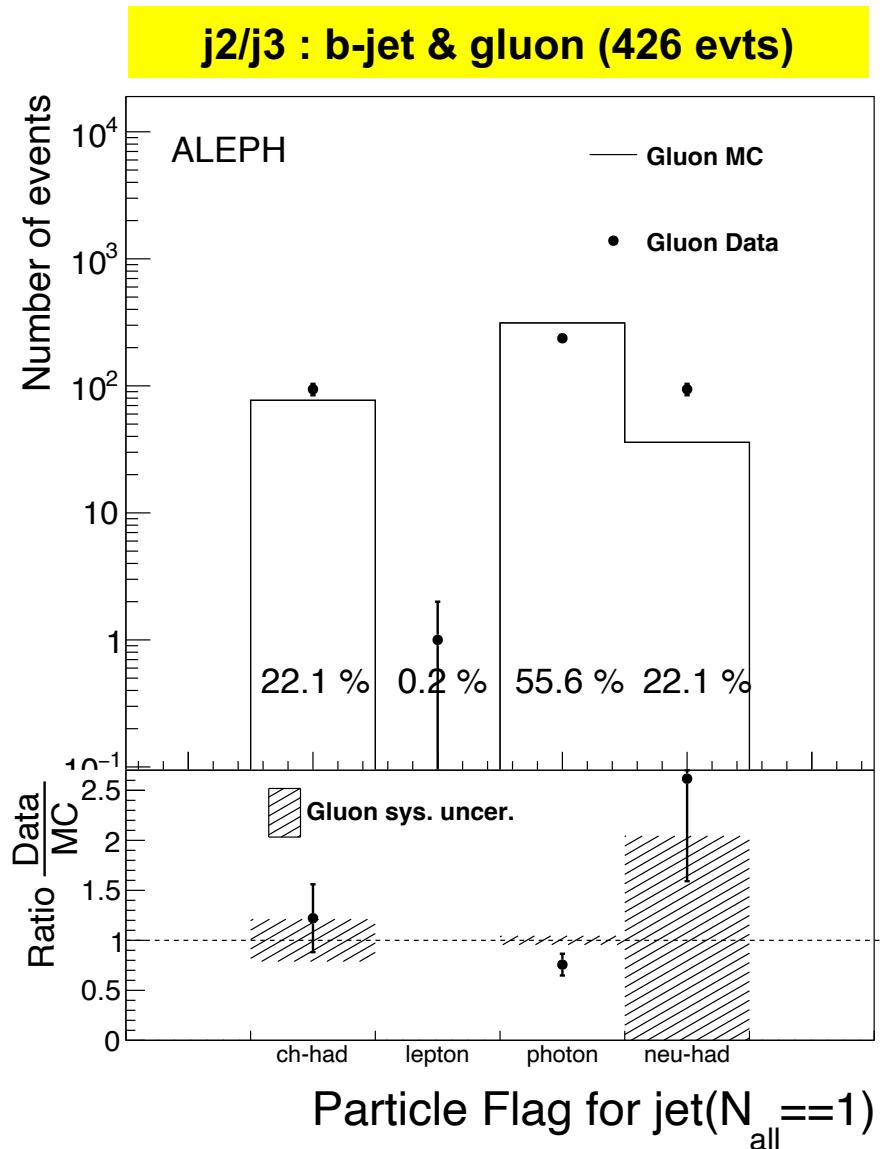
jet1 (38, 9, 1, 0, 0)  
jet2 (23.3, 1, 2, 5, 1)  
jet3 (15.3, 3, 0, 0, 0)  
jet4 (1.1, 1, 0, 0, 0)

b-jet  
(38 GeV)

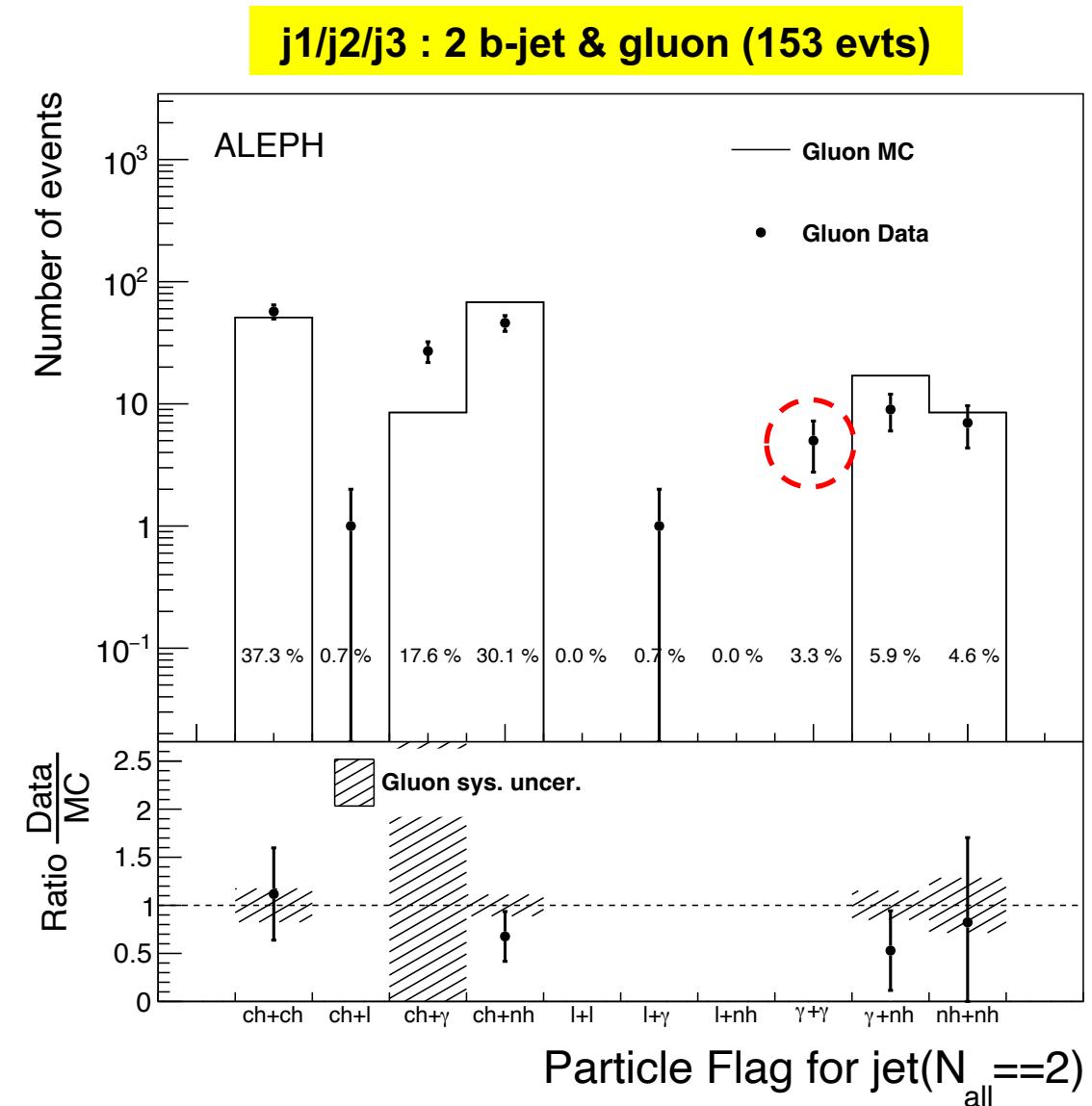
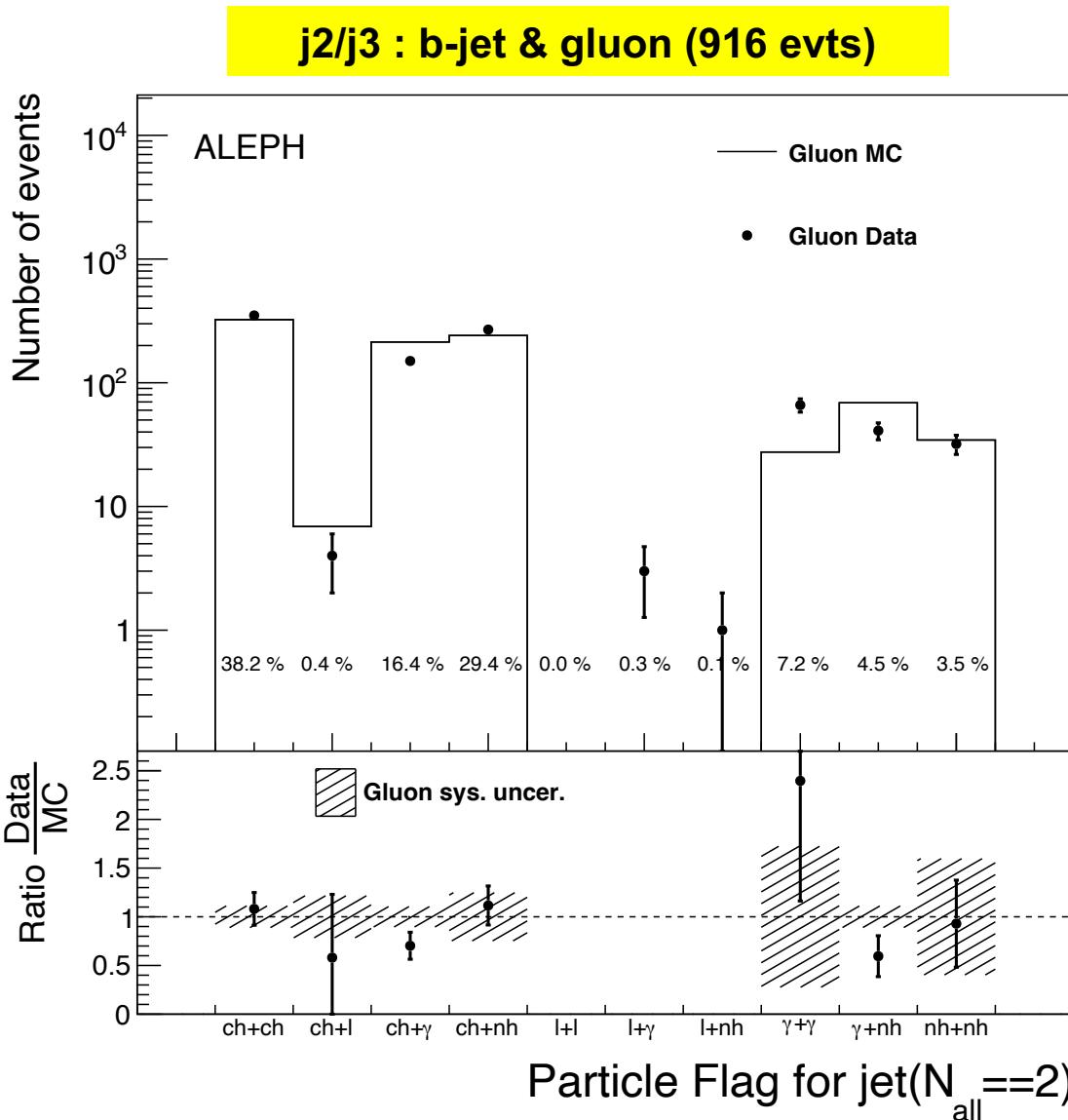
b-jet  
(23 GeV)

gluon-jet  
(15 GeV)  
 $n=3$

# Single track gluons? Who are they?



# 2-track gluons? Who are they?



# Remarks & Discussions

# Conclusions

- We followed ALPEH's previous QCD jet analysis.
  - Reconfirm the ALEPH's uncertain anomalies in  $N$  and  $X_E$ 
    - ✓ Excess of charged tracks with a large energy share
    - ✓ Correlation between two anomalies in  $N$  and  $X_E$ 
      - Excess of low multiplicity jet in ALEPH data
- We cannot claim the existence of two types of gluons; however, we see that Cho-gluons could explain such anomalies.
  - In-depth studies of the Gluon anomaly are needed
    - ✓ using other variables
    - ✓ using other elaborated jet algorithms that are sensitive to the jet shape.
  - MC generator level studies are also necessary
    - ✓ for unfolded analysis
    - ✓ with implementation of Neurons & Chromons

Thank you

# **backup slides**

# charged multiplicity

- ALEPH

- $\langle n_g \rangle = 9.90 \pm 0.10 \pm 0.27$
- $\langle n_{uds} \rangle = 7.90 \pm 0.44 \pm 0.26$
- $\langle n_b \rangle = 9.32 \pm 0.27 \pm 0.27$
- $\langle n_c \rangle = 8.37 \pm 1.64 \pm 0.28$

\* NFM=Natural Flavor Mix @ Z

$$R_{g/\text{NFM}} = 1.194 \pm 0.027(\text{stat}) \pm 0.019(\text{syst})$$

$$R_{g/uds} = 1.249 \pm 0.084(\text{stat}) \pm 0.022(\text{syst})$$

$$R_{g/b} = 1.060 \pm 0.041(\text{stat}) \pm 0.020(\text{syst})$$

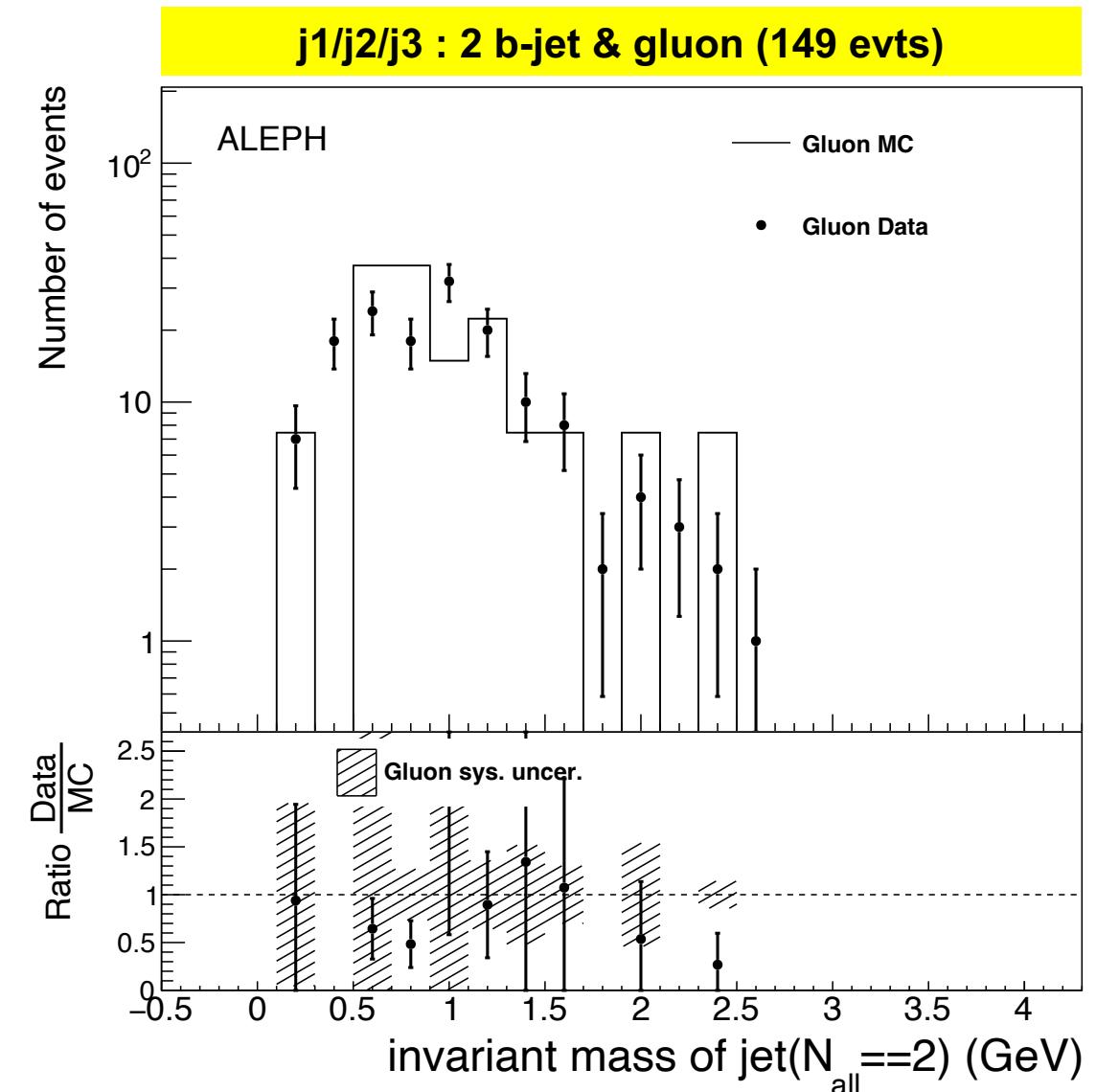
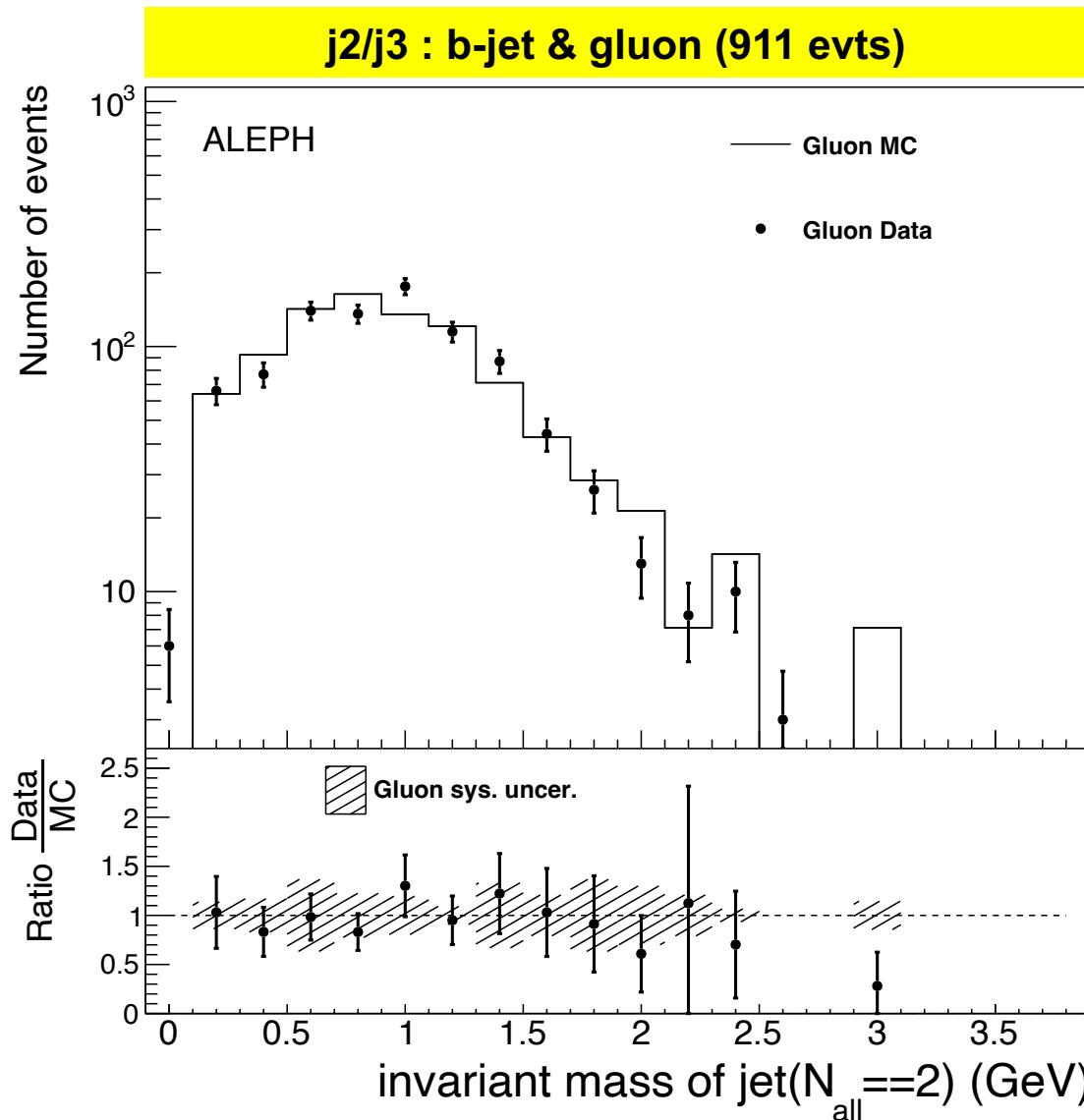
$$R_{g/c} = 1.183 \pm 0.221(\text{stat}) \pm 0.021(\text{syst})$$

- Model

- $R_{g/b} = 1.077 (\text{JETSET}) 1.003 (\text{HERWIG})$

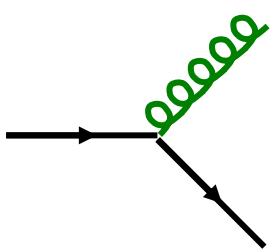
Physics Letters B 384 (1996) 353-364

# mass of 2-trk gluons (inclusive)

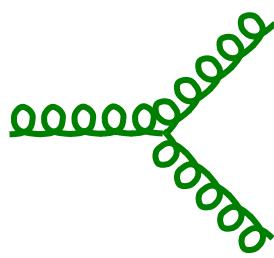


# Casimir operators

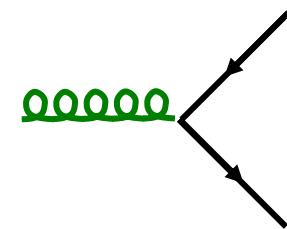
- In QCD, the strong force is described by the SU(3) gauge theory, where:
    - Quarks are in the fundamental representation (triplet) of SU(3).
    - Gluons are in the adjoint representation (octet) of SU(3).
  - $C_F$  : Casimir Operator for the Fundamental Representation
    - strength of the interaction between quarks
  - $C_A$  : Casimir Operator for the Adjoint Representation
    - characterizes the self-interaction of gluons
- ✓  $C_F = \frac{N_c^2 - 1}{2N_c} = \frac{4}{3}$  (for QCD)
- ✓  $C_A = N_c = 3$  (for QCD)



$$c_F = \frac{4}{3}$$



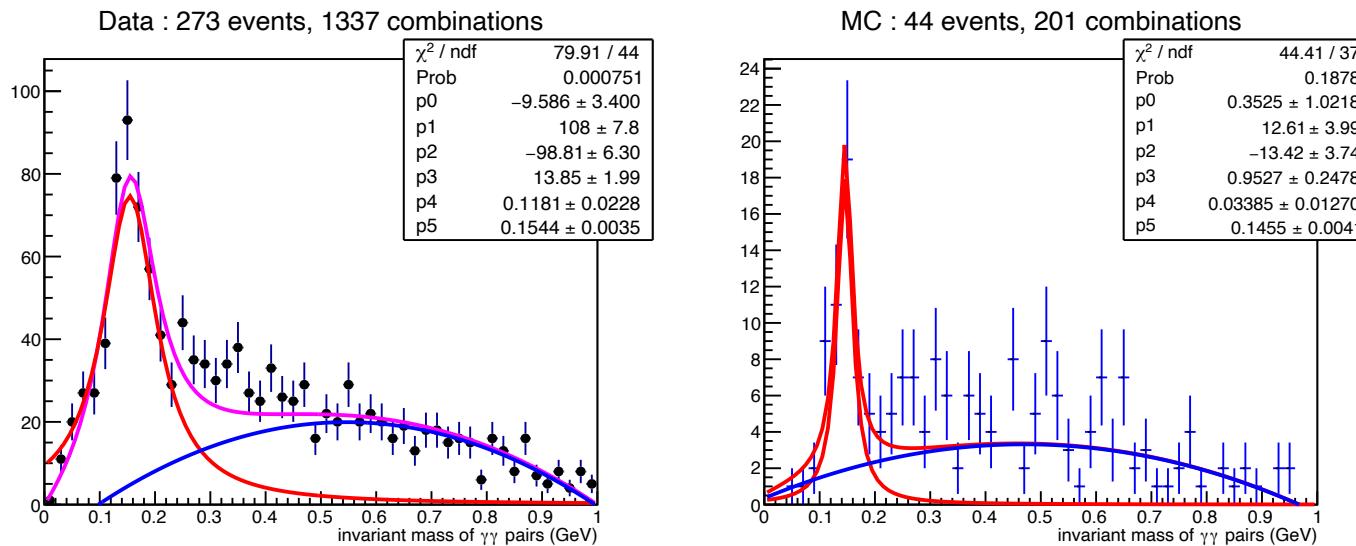
$$c_A = 3$$



$$T_F = \frac{1}{2}$$

# gluons without a charged track

- two photon mass



- such gluons have pi0, however no noticeable difference is found

# Our obsevations & implications

- Data excesses in gluon jets (compared to MC)
  - $X_E$ 
    - ✓ Gluon jets have more charged tracks with high energy share (high  $X_E$ )
  - $N_{trk}$ 
    - ✓ Gluon jets with small number of tracks (2 tracks) are more
- Implications
  - It is unclear whether this is simply MC's fault or if there is any gluon anomaly in the data.
  - ✓ MC tuning for fragmentations is required.
  - ✓ Can this be natually explained by Cho-gluon theory?