

Normalizing Flows for importance sampling in lattice computations

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Talk based on PRD 108, 114501 (2023) and prior works PRL 126, 032001 (2021), PRE 101, 023304 (2020)

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Generative AI for Physics : Past, Present and Future

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F. Noé, et al., Science, eaaw1147 (2019)



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Generative AI for Physics : Past, Present and Future



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Scalar Field Theories (ϕ^4)

M. S. Albergo, et al., Phys. Rev. D 100, 034515 (2019) K. A. Nicoli, et al., Phys. Rev. Lett. 126, 032001 (2021) P. deHaan, et al., arXiv: 2110.02673 @ ML4Pys workshop (2021) L. Vaitl et al., arXiv: 2206.09016 @ ICML (2022) A. Matthews et al., arXiv:2201.13117 @ ICML (2022) M. Caselle, et al., J. High Energ. Phys. 2022, 15 (2022) M. Gerdes, et al., SciPost Phys. 15, 238 (2023) A. Singha, et al., Phys. Rev. D 107, 014512 (2023) Lattice Gauge Theories (U(1), SU(N)) G. Kanwar, et al., Phys. Rev. Lett. 125, 121601 (2020) S. Bacchio, et al., Phys. Rev. D 107, L051504 (2023) R. Abbott et al., Phys. Rev. D 106, 074506 (2022) M. S. Albergo, et al., Phys. Rev. D 106, 014514 (2022) R. Abbott, et al., arXiv:2305.02402 (2023) J. Finkenrath, arXiv: 2201.02216 (2022) Sampling Multimodal Densities in QFT D.C. Hackett et al., arXiv:2107.00734 (2021) K. A. Nicoli, et al., arXiv: 2111.11303 @ LATTICE21 (2021) K. A. Nicoli, et al., Phys. Rev. D 108, 114501 (2023) B. Maté et al., TMLR 2835-8856 (2023)

V. Kanaujia et al., arXiv:2401.15948 (2024)



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(Image credits: Lattice QCD © Derek Leinweber/CSSM/University of Adelaide)

Scaling to Larger Lattices

L. Del Debbio, et al., Phys. Rev. D 104, 094507 (2021) R. Abbott et al., Eur. Phys. J. A 59, 257 (2023) A. Faraz et al., arXiv:2308.08615 (2022) B. Maté et al., arXiv: 2401.00828 (2024) R. Abbott et al., arXiv: 2401.10874v1 (2024) J. Finkenrath, arXiv: 2402.12176 (2024)

Autoregressive Models in Stat. Mech.

D. Wu et al., Phys. Rev. Lett. 122 (8), 080602 (2019) K. A. Nicoli, et al., Phys. Rev. E 101 (2), 023304 (2019) P. Bialas et al., Computer Physics Communications 281 (2022) P. Bialas et al., Phys. Rev. E 107 (1), 015303 (2023)

Interesting Readings

F. Noé, et al., Science, eaaw1147 (2019)

S. Chen, et al., Phys. Rev. D 107, 056001 (2022)

B. Maté et al., arXiv: 2210.13772 (2022)

L. Vaitl et al., MLST 3 (4), 045006 (2022)

Caselle, M., et al., J. High Energ. Phys. 2024, 48 (2024)

Cranmer K. et al., Nature Reviews Physics 5 (9), 526-535

And much more...



$$\langle \mathcal{O} \rangle_p = \int D[\phi] \mathcal{O}(\phi) p(\phi) \approx \frac{1}{N} \sum_{i=1}^N \mathcal{O}(\phi_i)$$

The field configuration $\phi(x)$ is a **random variable** sampled with **MCMC** to estimate computed over a Boltzmann-like density:



HMC and its flaws



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MCMC: sequentially proposes new sample and guarantees to eventually converge to a target density.

However MCMC algorithms come at a cost:

Sequential \Rightarrow MCMC chains can't be parallelized.

 \mathbf{P} Critical slowing down \Rightarrow Phase transitions.

 \P Long-range autocorrelations \Rightarrow large statistical errors.

Figure The partition function Z is unknown.

No direct estimation of thermodynamic observables.



(Adapted from:

What about flow-based sampling?



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We use a parametric function f_{θ} (a **<u>diffeomorphism</u>**) to transform Gaussian samples $z \sim q_0$ into physical configurations $\phi \sim q_{\theta}$

$$f_{\theta}: z \in \mathcal{Z} \sim q_z \to x = f_{\theta}(z) \in \mathcal{X} \sim q_{\theta}.$$

The parametric function needs to fulfill certain criteria:

- **Bijective** transformation $\phi = f_{\theta}(z)$
- Invertible and differentiable
- Tractable Jacobian



What's flow-based sampling?



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The likelihood of q_{θ} can be computed exactly: $q_{\theta}(\phi) = q_0(f_{\theta}^{-1}(\phi)) \left| \det\left(\frac{\partial f_{\theta}}{\partial z}\right) \right|$ $z = f_{\theta}^{-1}(\phi)$





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Often the variational density q_{θ} is trained by minimizing the **<u>Reverse-KL</u>** divergence:

$$KL(q_{\theta} | | p) = \int D[\phi] q_{\theta}(\phi) \ln \frac{q_{\theta}(\phi)}{p(\phi)} \equiv \mathbb{E}_{q_{\theta}}\left[\ln \frac{q_{\theta}(\phi)}{p(\phi)}\right]$$



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since we know the target $p(\phi)$ is a Boltzmann distribution $p(\phi) = Z^{-1} \exp\{-S(\phi)\}$

$$KL(q_{\theta} | | p) = \mathbb{E}_{q_{\theta}} \left[\ln \frac{q_{\theta}(\phi)}{p(\phi)} \right] = \mathbb{E}_{q_{\theta}} \left[\ln q_{\theta}(\phi) + S(\phi) + \ln Z \right]$$



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$$\nabla_{\theta} KL(q_{\theta} | | p) = \mathbb{E}_{q_{\theta}} \left[\nabla_{\theta} \ln q_{\theta}(\phi) + \nabla_{\theta} S(\phi) + \ln Z \right]$$



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Training can be performed by <u>self-sampling</u> from the model we are training!



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$$p(\phi) \approx q_{\theta} \sim \phi_i$$
 with $p(\phi) = \frac{\exp\{-S(\phi)\}}{Z}$

$$Z = \int D[\phi] \exp\{-S(\phi)\} = \int D[\phi] q_{\theta}(\phi) V(\phi) \text{ where } V(\phi) = \frac{\exp(-S(\phi))}{q_{\theta}(\phi)}$$



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$$Z \stackrel{\text{MC}}{\approx} \hat{Z} = \frac{1}{N} \sum_{i=1}^{N} \frac{\tilde{w}(\phi_i)}{\phi_i \sim q_{\theta}}$$



$$p(\phi) \approx q_{\theta} \sim \phi_{i} \quad \text{with} \quad p(\phi) = \frac{\exp\{-S(\phi)\}}{Z}$$
$$= \left[D[\phi] \exp\{-S(\phi)\} = \left[D[\phi] q_{\theta}(\phi) \, \tilde{w}(\phi) \text{ where } \quad \tilde{w}(\phi) = \frac{\exp\{-S(\phi)\}}{\tilde{x}(\phi)} \right]$$

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$$Z \stackrel{\text{MC}}{\approx} \hat{Z} = \frac{1}{N} \sum_{i=1}^{N} \tilde{w}(\phi_{i}) \longrightarrow \hat{F} = -T \ln \hat{Z}$$
$$\stackrel{\text{Nicolit, Phys. Rev. Lett. (2021)}}{\stackrel{\text{Nicolit, Phys. Rev. Lett. (2021)}}}$$

Asymptotically Unbiased Estimators



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$$\langle \mathcal{O} \rangle_p = \langle w \mathcal{O} \rangle_{q_{\theta}} \stackrel{\text{\tiny MC}}{\approx} \frac{1}{N} \sum_{i=1}^{N} w(\phi_i) \mathcal{O}(\phi_i) \quad \phi_i \sim q_{\theta}$$

Asymptotically Unbiased Estimators



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Asymptotically Unbiased Estimators



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The 2D Ising Model



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The 2D Ising Model



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Real Scalar ϕ^4 -Theory in (1+1) D



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Real Scalar ϕ^4 -Theory in (1+1) D



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Nicoli+, Phys. Rev. D (2023)

Real Scalar ϕ^4 -Theory in (1+1) D



 $\phi_i \sim q_{ heta}$

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Nicoli+, Phys. Rev. D (2023)







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Reverse-KL Div.

$$KL_{R}(q_{\theta} | | p) = \int D[\phi] q_{\theta}(\phi) \ln \frac{q_{\theta}(\phi)}{p(\phi)}$$



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Reverse-KL Div.

$$KL_{R}(q_{\theta} | | p) = \int D[\phi] q_{\theta}(\phi) \ln \frac{q_{\theta}(\phi)}{p(\phi)}$$

Forward-KL Div.

$$KL_F(p \mid \mid q_{\theta}) = \int D[\phi] p(\phi) \ln \frac{p(\phi)}{q_{\theta}(\phi)}$$



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Definition. The effective support of the variational density q_{θ} <u>relative</u> to p

$$\widetilde{\operatorname{supp}}_{p,\epsilon}(q_{\theta}) = \{ \phi \in \operatorname{supp}(q_{\theta}); \ q_{\theta}(\phi) > \epsilon p(\phi) \}$$

for a given numerical threshold ϵ . The mode dropping set is then given by

$$\mathcal{S} := \operatorname{supp}(p) \setminus \widetilde{\operatorname{supp}}_{p,\epsilon}(q_{\theta})$$

Nicoli+, Phys. Rev. D (2023)



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$$\mathcal{S} := \operatorname{supp}(p) \setminus \widetilde{\operatorname{supp}}_{p,\epsilon}(q_{\theta})$$

if the flow is **effectively mode-dropping**, the importance-weighted estimator, with a finite number of samples *N*, will miss a contribution from the mass $\int_{\mathcal{S}} p(\phi) d\phi$ with approximately the probability $1 - \epsilon N \int_{\mathcal{S}} p(\phi) d\phi$.

Nicoli+, Phys. Rev. D (2023)



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Definition. We define the effective sampler distribution

$$\widetilde{q}_{\theta}(\phi) = \begin{cases} q_{\theta}(\phi)/\zeta & \text{if } \phi \in \widetilde{\operatorname{supp}}_{p,\epsilon}(q_{\theta}) \\ 0 & \text{otherwise,} \end{cases} \quad where \quad \zeta = \int_{\widetilde{\supp}_{p,\epsilon}} \mathcal{D}[\phi]q_{\theta}(\phi) \le 1$$

is the multiplicative renormalization factor necessary to guarantee the normalization of $\tilde{q}_{\theta}(\phi)$.



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It follows that the importance-weighted estimator misses the contribution from the mode-dropping set \mathcal{S}

$$\hat{\mathcal{O}} \equiv \frac{1}{N} \sum_{i=1}^{N} \frac{p(\phi_i)}{q_{\theta}(\phi_i)} \mathcal{O}(\phi_i) \approx \mathbb{E}_{\phi \sim \widetilde{q}_{\theta}} \left[\frac{p(\phi)}{q_{\theta}(\phi)} \mathcal{O}(\phi) \right] \equiv \bar{\mathcal{O}}$$



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the typical values of the estimator $\hat{O} \approx \bar{O}$ can be **significantly different** from the true expectation value!



The mode-dropping estimator



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When q_{θ} has **full effective support** on the domain of p

$$w^* = \mathbb{E}_{q_{\theta}}\left[\frac{p(\phi)}{q_{\theta}(\phi)}\right] = \int_{\mathrm{supp}(q_{\theta})} q_{\theta}(\phi) \frac{p(\phi)}{q_{\theta}(\phi)} \mathcal{D}[\phi] = \int_{\mathrm{supp}(q_{\theta})} p(\phi) \mathcal{D}[\phi] = 1.$$

The mode-dropping estimator



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however if q_{θ} is **effectively mode-dropping** this expectation value becomes

$$\bar{w} \equiv \frac{1}{Z} \mathbb{E}_{\phi \sim \tilde{q}_{\theta}} \left[\frac{e^{-S(\phi)}}{q_{\theta}(\phi)} \right] \in (0, 1]$$

The mode-dropping estimator



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for which we can get the corresponding Monte Carlo estimator, i.e., the mode-dropping estimator

$$\bar{w} \approx \frac{1}{\hat{Z}_p} \left(\frac{1}{N} \sum_{i=1}^N \frac{e^{-S(\phi_i)}}{q_\theta(\phi_i)} \right) = \left(\frac{1}{N} \sum_{j=1}^N \frac{q_\theta(\phi_j)}{e^{-S(\phi_j)}} \right) \left(\frac{1}{N} \sum_{i=1}^N \frac{e^{-S(\phi_i)}}{q_\theta(\phi_i)} \right) \equiv \hat{w}$$









- i) Asymptotically unbiased samplers can be constructed from trained DGMs (NIS or NMCMC).
- ii) **Direct** estimation of the **partition function** and **thermodynamic observables**.
- iii) Sampling from DGMs is **embarrassingly parallelizable** \neq MCMC (sequential).
- iv) Training with **forward-KL** leads to better models though requires training samples.
- v) Derivation of **mode-dropping estimator** to reliably assess the goodness of the model.
- vi) Mitigation of mode-dropping using different objectives (FWD-KL) or stochastic approaches (SNFs).



Thank you for your attention!