

vowing chapter will give a more thorough introabiish the basics for following chapters.





Peter Kicsiny | CERN Openlab Summer Student Lectures | 15th July 2024



What is this lecture about?

- Introduction to some of the main concepts in the field
- Walkthrough of the concepts using toy examples
- Some application highlights

What is not covered in this lecture?

- Introduction of numerical simulations / differential equations
- Introduction to machine learning

Literature

The lecture will cover:

- 1. Surrogate models
- 2. Physics informed neural networks
- 3. Differentiable physics

Content based on:

Physics-based Deep Learning by N. Thuerey, P . Holl, M. Mueller, P. Schnell, F. Trost, K. Um

+ supplemented by real physics examples









30 Apr 2022 18:50Z NESDIS/STAR GOES-East GEOCOLOR

























Goal: make classical simulations faster & more accurate



Goal: make classical simulations faster & more accurate



• Supervised learning



- Supervised learning
- Given arbitrary unknown function describing a physical system: $f^*(x) = y^*$



- Supervised learning
- Given arbitrary unknown function describing a physical system: $f^*(x) = y^*$
- Data: measured/simulated samples $[x_0, y_0^*], ... [x_n, y_n^*]$



- Supervised learning
- Given arbitrary unknown function describing a physical system: $f^*(x) = y^*$
- Data: measured/simulated samples $[x_0, y_0^*], ... [x_n, y_n^*]$
- Goal: approximate f^* with a neural network (NN) denoted by f, trained on this data



- Supervised learning
- Given arbitrary unknown function describing a physical system: $f^*(x) = y^*$
- Data: measured/simulated samples $[x_0, y_0^*], ... [x_n, y_n^*]$
- Goal: approximate f^* with a neural network (NN) denoted by f, trained on this data
- Evaluate loss function & optimize weights of NN via backpropagation



Surrogate models: example

Output Input 128 x 128 x 1 Freestream X Pressure 128 x 128 x 1 128 x 128 x 1 Freestream Y Velocity X Navier-Stokes numerical solver 128 x 128 x 1 Velocity Y Mask 128 x 128 x 1



$$\underset{\theta}{\operatorname{argmin}} \sum_{i} (f(x_i; \theta) - y_i^*)^2$$





- 0.6

- 0.4

- 0.2

0.0

1.0

 Reconstructing pressure seems to be the most challenging

- Supervised training setup is a good first approach in many situations
- Always start with a 1-sample overfitting test
- Check how many trainable parameters your network has
- Slowly increase the amount of training data (and potentially network parameters and depth)
- Adjust hyperparameters (especially the learning rate)
- For any structured data, like spatial functions, or data of any physical field, convolutional NNs are preferable to fully connected NNs

- Surrogate modeling of LHC beam lifetime
 - Machine learning for beam dynamics studies at the CERN Large Hadron Collider, P. Arpaia et al. [10.1016/j.nima.2020.164652]
 - Data: lifetime measurements from real-life LHC parameter scans





- Surrogate modeling of LHC dynamic aperture
 - Modeling Particle Stability Plots for Accelerator Optimization Using Adaptive Sampling, M. Schenk et al. [10.18429/JACoW-IPAC2021-TUPAB216]
 - Data: numerical simulations of different LHC configurations



Colors: how many turns the particles survive





Fast training & constant time evaluation compared to numerical solvers or measurements









Physical losses: PDEs

- Partial differential equation (PDE)
- Time evolution of a physical field e.g. velocity expressed with spatial derivatives

$$\mathbf{u}_t = \mathcal{F}(\mathbf{u}_x, \mathbf{u}_{xx}, \dots, \mathbf{u}_{xx\dots x}) \qquad \mathbf{u} = \mathbf{u}(\mathbf{x}, t) \qquad \mathbf{x}, \mathbf{u} \in \mathbb{R}^N$$
N dimensional space

- Initial condition: $\mathbf{u}(\mathbf{x}, t = 0) = \mathbf{u}^0(\mathbf{x})$
- Boundary condition: $\mathbf{u}(\mathbf{x} = \partial \mathbf{x}, t) = 0$
- With this (usually) a unique solution for **u** exists (=well-posed PDE)
- Space-time domain is discretized into a computational grid

 $t \in \mathbb{R}^+$

• Partial differential equation (PDE)

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2} \qquad \begin{array}{c} x, u \in \mathbb{R} \\ t \in \mathbb{R}^+ \end{array}$$

Initial conditionBoundary condition
$$u(x,t=0) = -\sin(\pi x)$$
 $u(x=\partial x,t) = 0$

Physical losses: Burger's equation



Initial condition	Boundary condition
$u(x,t=0) = -\sin(\pi x)$	$u(x = \partial x, t) = 0$

Physical losses: Burger's equation



$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2} \qquad x, u \in \mathbb{R} + \frac{1}{2}$$

1.00

t=0.0

• Approximate velocity field **u** with a NN **f**

$$f(X_i, \theta) \approx u(X_i)$$
 $X_i = [x_i, t_i]$

Burger's equation

$$\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

• Approximate velocity field **u** with a NN **f**

$$f(X_i, \theta) \approx u(X_i)$$
 $X_i = [x_i, t_i]$

Burger's equation

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

$$0 = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} - \nu \frac{\partial^2 u}{\partial x^2}$$

• Approximate velocity field **u** with a NN **f**

$$f(X_i, \theta) \approx u(X_i)$$
 $X_i = [x_i, t_i]$

Burger's equation

$$\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

$$0 = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} - \nu \frac{\partial^2 u}{\partial x^2}$$

Residual loss

$$R = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} - \nu \frac{\partial^2 u}{\partial x^2}$$

• Approximate velocity field *u* with a NN *f*

$$f(X_i, \theta) \approx u(X_i)$$
 $X_i = [x_i, t_i]$

Physics informed neural network (PINN)

Burger's equation

Residual loss

$$\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

$$0 = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} - \nu \frac{\partial^2 u}{\partial x^2}$$

$$L(X_{i};\theta) = \sum_{i} \alpha_{0} \left(f(X_{i};\theta) - y_{i}^{*} \right)^{2} + \alpha_{1} R(X_{i};\theta) \qquad R = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} - \nu \frac{\partial^{2} u}{\partial x^{2}}$$
replace u by f

$$R = \frac{\partial f(X_{i};\theta)}{\partial t} + f(X_{i};\theta) \frac{\partial f(X_{i};\theta)}{\partial x} - \nu \frac{\partial^{2} f(X_{i};\theta)}{\partial x^{2}}$$

Physical losses: obtaining partial derivatives

$$R = \frac{\partial f(X_i;\theta)}{\partial t} + f(X_i;\theta)\frac{\partial f(X_i;\theta)}{\partial x} - \nu \frac{\partial^2 f(X_i;\theta)}{\partial x^2}$$

sample grid point

$$X_i = [x_i, t_i]$$

$$L(X_{i};\theta) = \sum_{i} \alpha_{0} (f(X_{i};\theta) - y_{i}^{*})^{2} + \alpha_{1}R(X_{i};\theta)$$

$$\underset{\theta}{\operatorname{argmin}} L(X_{i};\theta)$$

$$\underset{\theta}{\operatorname{argmin}} \sum_{i} \alpha_{0} (f(X_{i};\theta) - y_{i}^{*})^{2} + \alpha_{1}R(X_{i};\theta)$$
Physical losses: obtaining partial derivatives

$$R = \frac{\partial f(X_i;\theta)}{\partial t} + f(X_i;\theta) \frac{\partial f(X_i;\theta)}{\partial x} - \nu \frac{\partial^2 f(X_i;\theta)}{\partial x^2}$$

argmin
$$\sum_{\theta} \alpha_0 \left(f(X_i;\theta) - y_i^* \right)^2 + \alpha_1 R(X_i;\theta)$$

• Derivatives in *R* are obtained from the NN via standard backpropagation

Physical losses: obtaining partial derivatives

$$R = \frac{\partial f(X_i;\theta)}{\partial t} + f(X_i;\theta) \frac{\partial f(X_i;\theta)}{\partial x} - \nu \frac{\partial^2 f(X_i;\theta)}{\partial x^2}$$

argmin
$$\sum_{\theta} \alpha_0 \left(f(X_i;\theta) - y_i^* \right)^2 + \alpha_1 R(X_i;\theta)$$

• Derivatives in *R* are obtained from the NN via standard backpropagation

NN output (f) at sample point [x_i , t_i]

Physical losses: obtaining partial derivatives

$$R = \frac{\partial f(X_i;\theta)}{\partial t} + f(X_i;\theta) \frac{\partial f(X_i;\theta)}{\partial x} - \nu \frac{\partial^2 f(X_i;\theta)}{\partial x^2}$$

argmin
$$\sum_{\theta} \alpha_0 \left(f(X_i;\theta) - y_i^* \right)^2 + \alpha_1 R(X_i;\theta)$$

- Derivatives in *R* are obtained from the NN via standard backpropagation
- When R is minimized: *u* (NN output) approximately solves the PDE

$$f(X_i; \theta) \approx u(X_i)$$
 $X_i = [x_i, t_i]$



find *u* with a PINN

$$\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

1D computational "grid": $x \in [-1,1]$ 128 steps Forward simulation in time: $t \in [0,1]$ 33 steps



- Supervised loss: ground truth data y^{*}_i
 - Reference **u**

$$\underset{\theta}{\operatorname{argmin}} \sum_{i} \frac{\sum_{i} \left[\alpha_0 \left(f(X_i; \theta) - y_i^* \right)^2 \right] + \alpha_1 R(X_i; \theta)}{X_i = [x_i, t_i]}$$

1D computational "grid": $x \in [-1,1]$ 128 steps Forward simulation in time: $t \in [0,1]$ 33 steps



- Supervised loss: ground truth data y_i*
 - Reference **u**
 - E.g. we know u(x, t=0.5) (direct constraint)

$$\underset{\theta}{\operatorname{argmin}} \sum_{i} \frac{\sum_{i} \left[\alpha_0 \left(f(X_i; \theta) - y_i^* \right)^2 \right] + \alpha_1 R(X_i; \theta)}{X_i = [x_i, t_i]}$$

 $x \in [-1,1]$ 128 steps $t \in [0,1]$ 33 steps





- Reference **u**
- E.g. we know u(x, t=0.5) (direct constraint)
- We also know u(x=-1;1, t)=0 (boundary conditions of PDE)

$$\underset{\theta}{\operatorname{supervised loss term}} \sum_{i} \frac{\left[\alpha_0 \left(f(X_i; \theta) - y_i^*\right)^2\right] + \alpha_1 R(X_i; \theta)}{X_i = [x_i, t_i]}$$

a construction of the second second

 $x \in [-1,1]$ 128 steps $t \in [0,1]$ 33 steps



а





 $\underset{\theta}{\operatorname{argmin}} \sum_{i} \alpha_0 \left(f(X_i; \theta) - y_i^* \right)^2 + \alpha_1 R(X_i; \theta)$ $X_i = [x_i, t_i]$ Residual loss: no ground truth data, but we sample the the NN inside the domain

> $x\in |-1,1|$ 128 steps $t\in [0,1]$ 33 steps

residual loss term



 $R = \frac{\partial f(X_i;\theta)}{\partial t} + f(X_i;\theta)\frac{\partial f(X_i;\theta)}{\partial r} - \nu \frac{\partial^2 f(X_i;\theta)}{\partial r^2}$

- Residual loss: no ground truth data, but we sample the the NN inside the domain
 - i.e.: input random X_i to NN
 - NN estimates *u* as *f*(*X_i;θ*)

$$\underset{\theta}{\operatorname{argmin}} \sum_{i} \alpha_0 \left(f(X_i; \theta) - y_i^* \right)^2 + \alpha_1 R(X_i; \theta)$$
$$X_i = [x_i, t_i]$$

$$x \in [-1,1]$$
 128 steps $t \in [0,1]$ 33 steps

residual loss term

0



- NN in this example:
 - 8 fully connected layers
 - ~200 trainable parameters

$$\underset{\theta}{\operatorname{argmin}} \sum_{i} \alpha_0 \left(f(X_i; \theta) - y_i^* \right)^2 + \alpha_1 R(X_i; \theta)$$
$$X_i = [x_i, t_i]$$

 $x \in [-1,1]$ 128 steps $t \in [0,1]$ 33 steps



- NN in this example:
 - 8 fully connected layers
 - ~200 trainable parameters
 - Evaluate NN at constraint points , & x
 - Calculate loss & update weights
 - Repeat 10,000 iterations (~1 min)

 $\underset{\theta}{\operatorname{argmin}} \sum_{i} \alpha_0 \left(f(X_i; \theta) - y_i^* \right)^2 + \alpha_1 R(X_i; \theta)$ $X_i = [x_i, t_i]$ $\& \mathbf{x}$ $x \in [-1, 1] \text{ 128 steps}$

 $t\in [0,1]$ 33 steps



- After training: need to evaluate NN for all grid points!
 - Computationally expensive for large grids
- Why is this possible?
 - NN inherently supports calculation of derivatives



Physical losses: discussion



- Boundary conditions u=0 are fulfilled ۲
- Shock at center not well represented ٠
- More accurate representation requires significantly more iterations even for this simple case

0.75

1.00

• This is a conceptual starting point, but not very accurate



- Physical loss allows to encode (unique) solutions to PDEs with NNs, which allows to use NNs as universal function approximators
- Not really "machine learning": we reconstruct a single PDE solution in a known space-time region



Plasma dege Plasma edge Plasm

Reconstruction of turbulent fluctuations of plasma in a tokamak (A. Mathews et al., 2022, [doi.org/10.1063/5.0088216])



The learned two-dimensional n_e (a), T_e (b), and Cn_0 (c) for plasma discharge 1 120 711 021 along with the experimentally observed 587.6 nm photon emission (d) at t = 1.312815 s. The learned measurements are based on the collective predictions within the deep learning framework training against the neutral transport physics and $N_P = 1$ CR theory constraints. Multimedia view: https://doi.org/10.1063/5.0088216.1.

Reconstructing unknown parameters in Schrödinger equation (M. Raissi et al., 2017, [arxiv.org/abs/1708.00588])





Reconstruction of blood flow in a blood vessel (E. Hwuang, S. Wang et al. [doi.org/10.1038/s42254-021-00314-5])







Quite slow: need to evaluate NN at every grid point, i.e. "paint the image pixel by pixel"



Easy setup with simple NN

PDE derivatives in physical loss can be computed with backpropagation

Popular for inverse problems: given certain measurements or observations (=training data), find a PDE solution



Quite slow: need to evaluate NN at every grid point, i.e. "paint the image pixel by pixel"

Accuracy of computed derivatives relies on learned representation

Does not combine well with numerical solvers e.g. for refining solution





Differentiable physics = differentiable numerical simulations of physical systems

- Differentiable physics = differentiable numerical simulations of physical systems
- Equip classical numerical solvers (discretized PDE) with the ability to compute gradients with respect to their inputs

- Differentiable physics = differentiable numerical simulations of physical systems
- Equip classical numerical solvers (discretized PDE) with the ability to compute gradients with respect to their inputs
- This allows integration of numerical methods into the training process of an attached NN

• Linear advection equation with u = u(x) and d = d(x, t)

$$\frac{\partial d}{\partial t} + \mathbf{u} \cdot \nabla d = 0$$

• Linear advection equation with u = u(x) and d = d(x, t)

$$\frac{\partial d}{\partial t} + \mathbf{u} \cdot \nabla d = 0 \quad \Longrightarrow \quad \frac{d(x_i, t + \Delta t) - d(x_i, t)}{\Delta t} = -u(x_i) \frac{d(x_i + \Delta x, t) - d(x_i, t)}{\Delta x}$$

E.g. in 1D

Linear advection equation with u = u(x) and d = d(x, t)٠

Einear advection equation with
$$\mathbf{u} = \mathbf{u}(\mathbf{x})$$
 and $\mathbf{d} = \mathbf{d}(\mathbf{x}, t)$
 $\frac{\partial d}{\partial t} + \mathbf{u} \cdot \nabla d = 0 \implies d(x_i, t + \Delta t) = d(x_i, t) - \Delta t \left[u(x_i) \frac{d(x_{i+1}, t) - d(x_i, t)}{\Delta x} \right]$

Generally: •

 $d(t + \Delta t) = \mathcal{P}(d(t), u, t + \Delta t)$

• Linear advection equation with u = u(x) and d = d(x, t)

$$\frac{\partial d}{\partial t} + \mathbf{u} \cdot \nabla d = 0 \quad \Longrightarrow \quad d(x_i, t + \Delta t) = d(x_i, t) - \Delta t \left[u(x_i) \frac{d(x_{i+1}, t) - d(x_i, t)}{\Delta x} \right]$$

- Generally:
 - $d(t + \Delta t) = \mathcal{P}(d(t), u, t + \Delta t)$
- For N forward iterations:

$$d^e = \mathcal{P}(d^0, u, t + N\Delta t = t^e)$$



E.a. in 1D

• Linear advection equation with u = u(x) and d = d(x, t)

$$\frac{\partial d}{\partial t} + \mathbf{u} \cdot \nabla d = 0 \quad \Longrightarrow \quad d(x_i, t + \Delta t) = d(x_i, t) - \Delta t \left[u(x_i) \frac{d(x_{i+1}, t) - d(x_i, t)}{\Delta x} \right]$$

• Generally:

$$d(t + \Delta t) = \mathcal{P}(d(t), u, t + \Delta t)$$

• For N forward iterations:

$$d^e = \mathcal{P}(d^0, u, t + N\Delta t = t^e)$$

Find velocity field *u* that brings a known initial density *d*⁰ = *d*(*t*⁰=0) into a known target density *d*^{target} = *d*(*t*^e=*t*+*N*Δ*t*)



 $F \alpha$ in 1D

Linear advection equation with u = u(x) and d = d(x, t)•

E.g. in 1D $\frac{\partial d}{\partial t} + \mathbf{u} \cdot \nabla d = 0 \quad \Longrightarrow \quad d(x_i, t + \Delta t) = d(x_i, t) - \Delta t \left[u(x_i) \frac{d(x_{i+1}, t) - d(x_i, t)}{\Delta x} \right]$

Generally:

$$d(t + \Delta t) = \mathcal{P}(d(t), u, t + \Delta t)$$

For N forward iterations:

$$d^e = \mathcal{P}(d^0, u, t + N\Delta t = t^e)$$

Optimization problem:

$$L = |d^e - d^{\text{target}}|^2$$

output of simulation known from observation

• Find velocity field *u* that brings a known initial density $d^0 = d(t^0=0)$ into a known target density $d^{target} = d(t^e = t + N\Delta t)$



• Linear advection equation with u = u(x) and d = d(x, t)

$$\frac{\partial d}{\partial t} + \mathbf{u} \cdot \nabla d = 0 \quad \Longrightarrow \quad d(x_i, t + \Delta t) = d(x_i, t) - \Delta t \left[u(x_i) \frac{d(x_{i+1}, t) - d(x_i, t)}{\Delta x} \right]$$

• Generally:

$$d(t + \Delta t) = \mathcal{P}(d(t), u, t + \Delta t)$$

• For N forward iterations:

$$d^e = \mathcal{P}(d^0, u, t + N\Delta t = t^e)$$

• Optimization problem:

$$L = |d^e - d^{\text{target}}|^2 = |\mathcal{P}(d^0, u, t^e) - d^{\text{target}}|^2$$

Find velocity field *u* that brings a known initial density *d*⁰ = *d*(*t*⁰=0) into a known target density *d*^{target} = *d*(*t*^e=*t*+*N*Δ*t*)



E.a. in 1D

• Linear advection equation with u = u(x) and d = d(x, t)

$$\frac{\partial d}{\partial t} + \mathbf{u} \cdot \nabla d = 0 \quad \Longrightarrow \quad d(x_i, t + \Delta t) = d(x_i, t) - \Delta t \left[u(x_i) \frac{d(x_{i+1}, t) - d(x_i, t)}{\Delta x} \right]$$

• Generally:

$$d(t + \Delta t) = \mathcal{P}(d(t), u, t + \Delta t)$$

• For N forward iterations:

$$d^e = \mathcal{P}(d^0, u, t + N\Delta t = t^e)$$

• Optimization problem:

$$L = |d^e - d^{\text{target}}|^2 = |\mathcal{P}(d^0, u, t^e) - d^{\text{target}}|^2$$

$$\underset{u}{\operatorname{argmin}} L(u) = \underset{u}{\operatorname{argmin}} |\mathcal{P}(d^{0}, u, t^{e}) - d^{\operatorname{target}}|^{2}$$

$$u \text{ unknown at start e.g. init with 0s}$$

Find velocity field *u* that brings a known initial density *d*⁰ = *d*(*t*⁰=0) into a known target density *d*^{target} = *d*(*t*^e=*t*+*N*Δ*t*)



E.a. in 1D
Differentiable physics: differentiable solver example

$$\underset{u}{\operatorname{argmin}} L(u) = \underset{u}{\operatorname{argmin}} |\mathcal{P}(d^{0}, u, t^{e}) - d^{\operatorname{target}}|^{2}$$

$$\underset{u}{\operatorname{learning rate}} |\mathcal{D}_{u_{i}} = \frac{\partial \mathcal{P}_{i}}{\partial u_{i}} \frac{\partial L}{\partial \mathcal{P}_{i}} \quad u_{i}^{\operatorname{new}} = u_{i} - \eta \Delta u_{i}$$

$$\begin{aligned} \underset{u}{\operatorname{argmin}} & L(u) = \underset{u}{\operatorname{argmin}} |\mathcal{P}(d^{0}, u, t^{e}) - d^{\operatorname{target}}|^{2} \\ \underset{u}{\operatorname{learning rate}} \\ \Delta u_{i} = \frac{\partial L}{\partial u_{i}} = \frac{\partial \mathcal{P}_{i}}{\partial u_{i}} \underbrace{\frac{\partial L}{\partial \mathcal{P}_{i}}}_{i} \quad u_{i}^{\operatorname{new}} = u_{i} - \eta \Delta u_{i} \\ \underset{i}{\operatorname{ln case of a single forward step from t^{e-1} to t^{e:}}} \\ \underbrace{\frac{\partial L}{\partial \mathcal{P}_{i}} = \frac{\partial |\mathcal{P} - d^{\operatorname{target}}|^{2}}{\partial \mathcal{P}_{i}} = \frac{\partial L}{\partial d_{i}^{e}} = \frac{\partial |d^{e} - d^{\operatorname{target}}|^{2}}{\partial d_{i}^{e}} = 2\left(d_{i}^{e} - d_{i}^{\operatorname{target}}\right) \end{aligned}$$

$$\underset{u}{\operatorname{argmin}} L(u) = \underset{u}{\operatorname{argmin}} |\mathcal{P}(d^{0}, u, t^{e}) - d^{\operatorname{target}}|^{2}$$

$$\underset{u}{\operatorname{learning rate}} |\mathcal{D}_{u_{i}}| = \frac{\partial \mathcal{P}_{i}}{\partial u_{i}} \frac{\partial L}{\partial \mathcal{P}_{i}} \quad u_{i}^{\operatorname{new}} = u_{i} - \eta \Delta u_{i}$$

In case of a single forward step from t^{e-1} to t^e:

$$\frac{\partial L}{\partial \mathcal{P}_i} = \frac{\partial |\mathcal{P} - d^{\text{target}}|^2}{\partial \mathcal{P}_i} = \frac{\partial L}{\partial d_i^e} = \frac{\partial |d^e - d^{\text{target}}|^2}{\partial d_i^e} = 2\left(d_i^e - d_i^{\text{target}}\right)$$

$$\frac{\partial \mathcal{P}_i}{\partial u_i} = \frac{\partial d_i^e}{\partial u_i} = -\frac{\Delta t}{\Delta x} \left[d_{i+1}^{e-1} - d_i^{e-1} \right]$$

From:
$$d(x_i, t + \Delta t) = d(x_i, t) - \Delta t \left[u(x_i) \frac{d(x_{i+1}, t) - d(x_i, t)}{\Delta x} \right]$$
 with: $\begin{array}{c} t^e = t + \Delta t \\ t^{e-1} = t \end{array}$

$$\underset{u}{\operatorname{argmin}} L(u) = \underset{u}{\operatorname{argmin}} |\mathcal{P}(d^0, u, t^e) - d^{\operatorname{target}}|^2$$

$$\Delta u_i = \frac{\partial L}{\partial u_i} = \frac{\partial \mathcal{P}_i}{\partial u_i} \frac{\partial L}{\partial \mathcal{P}_i} \quad u_i^{\text{new}} = u_i - \eta \Delta u_i$$

In case of multiple forward steps from t⁰ to t^e: $\Delta u_i = \frac{\partial d^e}{\partial u_i} \frac{\partial L}{\partial d^e}$ $+ \frac{\partial d^{e-1}}{\partial u_i} \frac{\partial d^e}{\partial d^{e-1}} \frac{\partial L}{\partial d^e}$ $+ \cdots$ $+ \left(\frac{\partial d^0}{\partial u_i} \cdots \frac{\partial d^{e-1}}{\partial d^{e-2}} \frac{\partial d^e}{\partial d^{e-1}} \frac{\partial L}{\partial d^e}\right)$

Final density **d**^e depends on velocity **u**_i through all previous density states:



$$\underset{u}{\operatorname{argmin}} L(u) = \underset{u}{\operatorname{argmin}} |\mathcal{P}(d^0, u, t^e) - d^{\operatorname{target}}|^2$$

$$\Delta u_i = \frac{\partial L}{\partial u_i} = \frac{\partial \mathcal{P}_i}{\partial u_i} \frac{\partial L}{\partial \mathcal{P}_i} \quad u_i^{\text{new}} = u_i - \eta \Delta u_i$$

In case of multiple forward steps from t⁰ to t^e:

Final density **d**^e depends on velocity **u**_i through all previous density states:



$$\underset{u}{\operatorname{argmin}} L(u) = \underset{u}{\operatorname{argmin}} |\mathcal{P}(d^0, u, t^e) - d^{\operatorname{target}}|^2$$

$$\Delta u_i = \frac{\partial L}{\partial u_i} = \frac{\partial \mathcal{P}_i}{\partial u_i} \frac{\partial L}{\partial \mathcal{P}_i} \quad u_i^{\text{new}} = u_i - \eta \Delta u_i$$

In case of multiple forward steps from t⁰ to t^e:



Final density **d**^e depends on velocity **u**_i through all previous density states:





*potential contributions from cells i+1, i-1 etc...

Differentiable physics: differentiable solver example

$$\begin{aligned} \underset{u}{\operatorname{argmin}} L(u) &= \underset{u}{\operatorname{argmin}} |\mathcal{P}(d^{0}, u, t^{e}) - d^{\operatorname{target}}|^{2} \\ \Delta u_{i} &= \frac{\partial L}{\partial u_{i}} \\ \\ \begin{array}{l} \text{In case of mi} \\ \Delta u_{i} &= \frac{\partial d^{e}}{\partial u_{i} \partial} \\ &+ \frac{\partial d}{\partial u_{i}} \\ &+ \frac{\partial d}{\partial u_{i}} \\ &+ \frac{\partial d}{\partial u_{i}} \\ &+ \frac{\partial d^{e}}{\partial u_{i}} \\ &+ \frac{\partial d^{e}}{\partial u_{i}} \\ \end{array} \\ \\ \begin{array}{l} \text{Protection of the set of$$

Differentiable physics: differentiable solver example

$$\operatorname{argmin}_{u} L(u) = \operatorname{argmin}_{u} |\mathcal{P}(d^{0}, u, t^{e}) - d^{\operatorname{target}}|^{2}$$

$$\Delta u_{i} = \frac{\partial L}{\partial u_{i}} = \operatorname{In \ case \ of \ m_{u}}_{u_{i}} = \operatorname{In \ case \ m_{u}}_{u_{i}$$

A typical PDE based numerical solver consists of arithmetic operations which are differentiable. The computation of gradients is typically not expensive and it can happen during forward simulation.

$$+ \left(\frac{\partial d^{\circ}}{\partial u_{i}} \cdots \frac{\partial d^{e-1}}{\partial d^{e-2}} \frac{\partial d^{e}}{\partial d^{e-1}} \frac{\partial L}{\partial d^{e}}\right) \left($$

$$\partial d_i^{e-1} \stackrel{-}{\longrightarrow} \Delta x^{a}$$

*potential contributions from cells i+1, i-1 etc...

- Classical gradient based optimization, no DL
- Start with u(x,t=0)=0, u^{target}(x,t=0.5) known





- Classical gradient based optimization, no DL
- Start with u(x,t=0)=0, u^{target}(x,t=0.5) known

- Burger's equation $\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$
- 1. Simulate \mathcal{P} (discretized PDE) from t=0 to t=1 in N=32 time steps



- Classical gradient based optimization, no DL
- Start with u(x,t=0)=0, u^{target}(x,t=0.5) known

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

- 1. Simulate \mathcal{P} (discretized PDE) from t=0 to t=1 in N=32 time steps
- 2. Backpropagate gradients from t=0.5, 16 steps back till first step at t=0



- Classical gradient based optimization, no DL
- Start with u(x,t=0)=0, u^{target}(x,t=0.5) known

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

- 1. Simulate \mathcal{P} (discretized PDE) from t=0 to t=1 in N=32 time steps
- 2. Backpropagate gradients from t=0.5, 16 steps back till first step at t=0
- 3. Update u(x,t=0) with gradients: $u_i^{\text{new}} = u_i \eta \Delta u_i$



- Classical gradient based optimization, no DL
- Start with u(x,t=0)=0, u^{target}(x,t=0.5) known
- 1. Simulate \mathcal{P} (discretized PDE) from t=0 to t=1 in N=32 time steps
- 2. Backpropagate gradients from t=0.5, 16 steps back till first step at t=0





Burger's equation

 $u \overline{\partial}_{\mathcal{X}}$

 ∂u

 ∂t





Physics informed NN (PINN)

1.00

0.75

0.50

0.25

- 0.00

-0.25

-0.50

-0.75

-1.00

- Recovers the overall shape of the solution
- Temporal constraints are at least partially fulfilled
- Poor reconstruction of amplitudes

Differentiable physics solver (DP)

- Much closer to ground truth (GT) thanks to flow of gradients
 - Difficulty with sharper features: artifacts

Differentiable physics: coupling differentiable solver with NN



Differentiable physics: coupling differentiable solver with NN ⁸⁹

$$\Delta \mathbf{u} = \partial L / \partial \mathbf{u} = \frac{\partial \mathcal{P}}{\partial \mathbf{u}} \frac{\partial L}{\partial \mathcal{P}}$$



- Physical field **u** input to differentiable solver
- Find optimal **u**

Differentiable physics: coupling differentiable solver with NN ⁹⁰



- Physical field **u** input to differentiable solver
- Find optimal **u**

- NN approximates physical field u
- Physical field **u** input to differentiable solver
- Find optimal NN weights: gradients are guided by solver

Differentiable physics: coupling differentiable solver with NN⁹¹



• Gradients from differentiable solver allow to access previously "hidden" parts of the loss landscape



Uses existing numerical tools which can be coupled to the training of neural networks







Controlling fluid deformations & reduction of numerical errors, J. Tang et al. 2023 [link]





Differentiable physics: application examples



Super-resolution forecasting of precipitation rate, B. Teufel et al. 2023 [doi.org/10.1186/s40562-023-00272-z]

Differentiable physics: application examples

Reconstruction of 6D beam distribution in a particle accelerator, R. Roussel et al., 2022, [arXiv:2211.09077]







Physics based deep learning is an emerging topic with many exciting possibilities



- Physics based deep learning is an emerging topic with many exciting possibilities
- Al will not replace classical numerical simulations!





- Physics based deep learning is an emerging topic with many exciting possibilities
- Al will not replace classical numerical simulations!



- A. Adelmann et al., 2022, New directions for surrogate models and differentiable programming for High Energy Physics detector simulation [doi.org/10.48550/arXiv.2203.08806]
- T. Dorigo et al., 2023, Toward the end-to-end optimization of particle physics instruments with differentiable programming [cds.cern.ch/record/2807001]
- MODE Collaboration, 2021, Toward Machine Learning Optimization of Experimental Design [inspirehep.net/literature/1850892]
- R. Lehe et al., 2020, Machine learning and surrogate models for simulation-based optimization of accelerator design [link]

- K. Kashinath et al., 2021, Physics-informed machine learning: case studies for weather and climate modelling [doi.org/10.1098/rsta.2020.0093]
- S. Rasp et al., 2021, Data-Driven Medium-Range Weather Prediction With a Resnet Pretrained on Climate Simulations: A New Model for WeatherBench [doi.org/10.1029/2020MS002405]
- J. Pathak et al., 2018, Hybrid forecasting of chaotic processes: Using machine learning in conjunction with a knowledge-based model [doi.org/10.1063/1.5028373]

Further read: examples in computer graphics

- S. Zhao et al., 2020, Physics-Based Differentiable Rendering A Comprehensive Introduction [link]
- N. Thuerey et al., 2019, Simulation & Animation [link]
- Y. Wang et al., 2023, Amortizing Samples in Physics-Based Inverse Rendering Using ReSTIR [link]

- J. Degrave et al., 2019, A Differentiable Physics Engine for Deep Learning in Robotics [doi.org/10.3389/fnbot.2019.00006]
- F. de Avila Belbute-Pere et al., 2018, End-to-End Differentiable Physics for Learning and Control [link]
- S. Chen et al., 2022, Imitation Learning via Differentiable Physics [doi.org/10.48550/arXiv.2206.04873]
- J. Lv et al., 2022, SAM-RL: Sensing-Aware Model-Based Reinforcement Learning via Differentiable Physics-Based Simulation and Rendering [doi.org/10.48550/arXiv.2210.15185]

Further read: examples in predictive control/maintenance

- J. Morton et al., 2018, Deep Dynamical Modeling and Control of Unsteady Fluid Flows [doi.org/10.48550/arXiv.1805.07472]
- L. G. Huber et al., 2023, Physics-Informed Machine Learning for Predictive Maintenance: Applied Use-Cases
 [doi.org/10.1109/SDS57534.2023.00016]
- D. Di Lorenzo et al., 2023, Physics informed and data-based augmented learning in structural health diagnosis [doi.org/10.1016/j.cma.2023.116186]
- V. Jadhav et al., 2022, Physics Informed Neural Network for Health Monitoring of an Air Preheater [doi.org/10.36001/phme.2022.v7i1.3343]