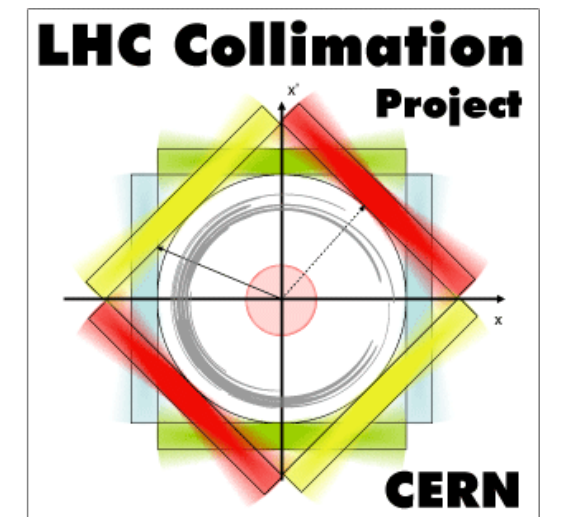
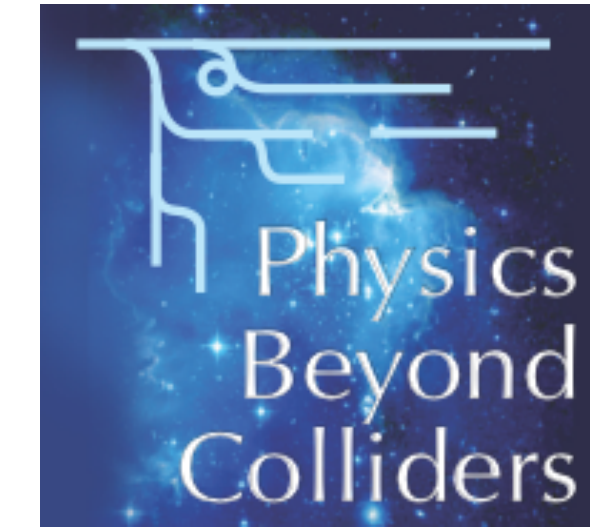




université
PARIS-SACLAY

IJCLab
Irène Joliot-Curie
Laboratoire de Physique
des 2 Infinis



Fixed-target experiments at LHC with bent crystals. MDM measurement of short lived particles.

Alex Fomin

contributed by:

S. Barsuk (IJCLab), L. Burmistrov (IJCLab), G. Calderini (LPNHE), S.P. Fomin (KIPT, KhNU), F. Galluccio (CERN), C. Hadjidakis (IJCLab),
I.V. Kirillin (KIPT, KhNU), A.Yu. Korchin (KIPT, KhNU), V.A. Kovalchuk (KIPT, KhNU), E. Kou (IJCLab), M. Liul (LAL, KhNU),
L. Massacrier (IJCLab), D. Mirarchi (CERN), E. Niel (IJCLab), M. Patecki (CERN), S. Redaelli (CERN), P. Robbe (IJCLab),
W. Scandale (CERN), N.F. Shul'ga (KIPT, KhNU), A. Stocchi (IJCLab)



Introduction

- Electromagnetic moments of short lived particles
- Spin precession in a bent crystal

Optimal crystal orientation for MDM and EDM measurement

- Spin precession in a bent crystal
- Initial polarisation of baryons [1,2]
- Quantitative analysis [1]

Performance assessment of layouts in IR2, IR3 and IR8 of LHC

- Double and single crystal layouts at LHC [3,4,5]
- Precision of measurement [1]
- Possible improvements [1,3]
- Direct measurement of τ -lepton MDM and EDM [6]

[1] [A.S. Fomin et al. Eur. Phys. J. C \(2020\) 80:358](#)

[2] [A.S. Fomin, JHEP 08 \(2017\) 120](#)

[3] [D. Mirarchi et al. Eur. Phys. J. C 80 \(2020\) 10, 929](#)

[4] [CERN Yellow Reports: Monographs, 4/2020](#)

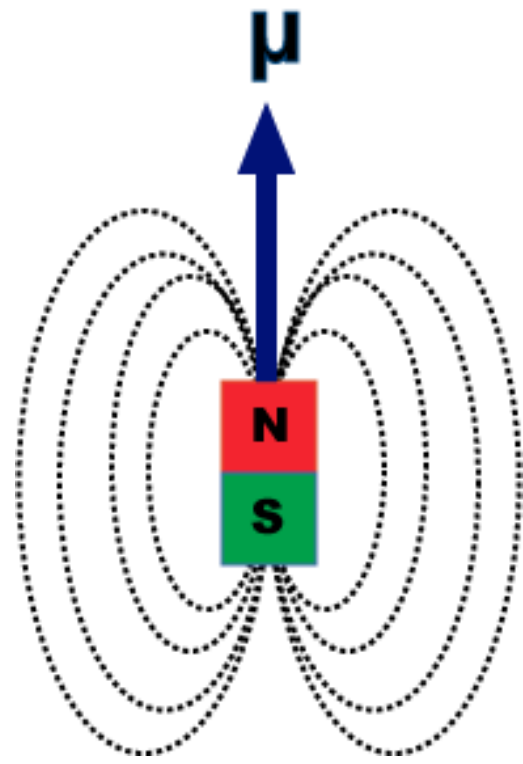
[5] [M. Patecki et al. Eur. Phys. J. C \(2023\) 83, 1053](#)

[6] [A.S. Fomin et al. JHEP 1903 \(2019\) 156](#)



MDM and EDM of short lived particles

Magnetic Dipole Moment:



$$\vec{\mu} = \frac{g}{2} \frac{e}{m} \vec{S}, \quad \vec{S} = \frac{\hbar}{2} \vec{\sigma}$$

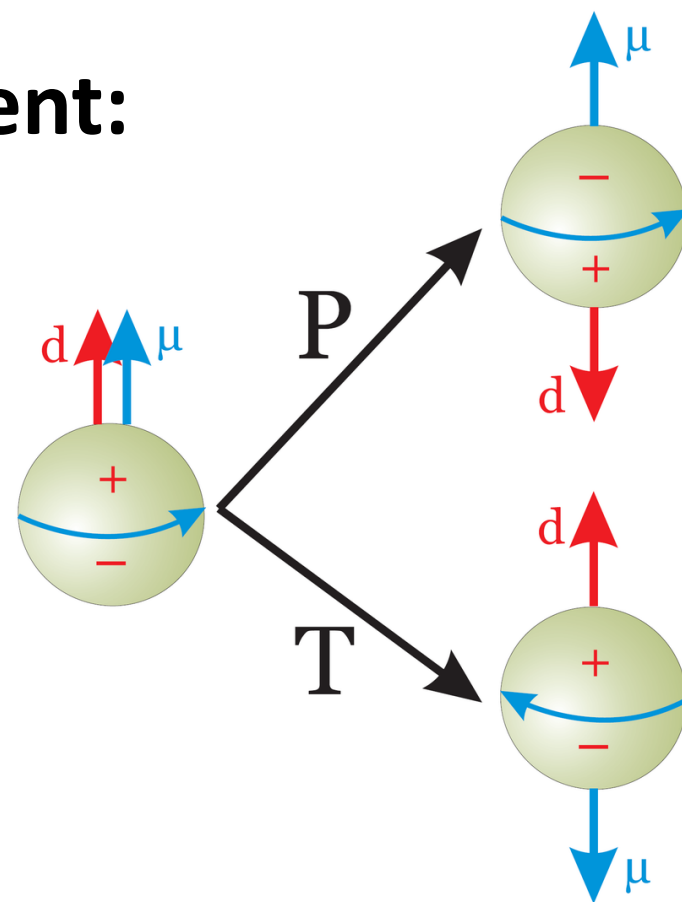
$|g| = 2 \rightarrow$ a point-like Dirac particle

$|g| \approx 2 \rightarrow$ a radiative corrections

$|g| \neq 2 \rightarrow$ a composite structure or NP

| Particle | $c\tau$ | g-factor | Comments |
|---------------|--------------------------|---|--|
| e^- | | $-2.002\,319\,304\,361\,82\,(52)$ | exp. most accurate determinations of α (Harvard 2008) |
| μ^- | 659 m | $-2.002\,331\,8361\,(10)$ $-2.002\,331\,84110\,(43_{\text{stat}}, 19_{\text{syst}})$ | theor. SM prediction (diff. from lattice gauge theory) exp. running (3/6 year of data taking at Fermilab 2021) |
| τ^- | 87 μm | $-2.002\,354\,42\,(10)$ $-2.036\,(34)$ $-2.002\,(6)$ no direct measurement | theor. SM prediction exp. $\sigma(e^+e^- \rightarrow e^+e^- \tau^+\tau^-)$ (LEP2: DELPHI 2004) exp. assuming EDM $\tau = 0$ (from LEP and SLD 2000) exp. Feasibility studies at LHC |
| p n | ∞ $\sim\infty$ | $+5.585\,694\,702\,(17)$ $-3.826\,085\,45\,(90)$ | exp. |
| Σ^+ | 2.4 cm | $+6.233\,(25)$ $+6.1\,(1.2)_{\text{stat}}\,(1.0)_{\text{syst}}$ | exp. world-average value exp. using Bent Crystals (at Fermilab 1990) |
| Λ_c^+ | 60 μm | $+1.90\,(15)$ not measured | theor. assuming $g_c \approx 2$ exp. Feasibility studies at LHC |

Electric Dipole Moment:



$$\vec{\delta} = \frac{f}{2} \frac{e}{m} \vec{S}, \quad \vec{S} = \frac{\hbar}{2} \vec{\sigma}$$

A nonzero value is forbidden by both:
T invariance and P invariance.

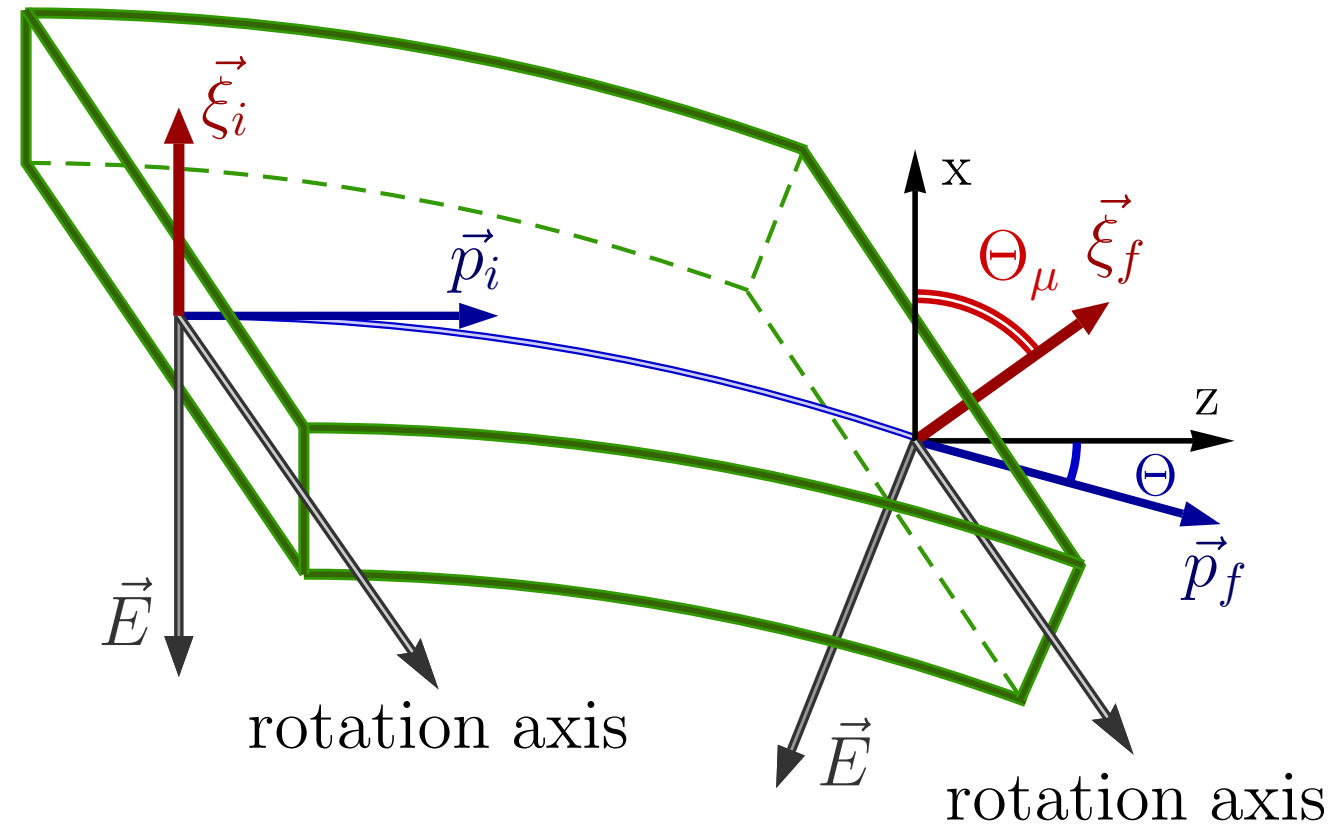
| Particle | $ \delta , e \text{ cm } 10^{-25}$ |
|---------------|------------------------------------|
| p | < 2.1 |
| n | < 0.18 |
| Σ^+ | not measured |
| Λ_c^+ | not measured |



Spin precession in a bent crystal

V.G. Baryshevsky, Sov. Tech. Phys. Lett. 5 (1979) 73.

V.L. Lyuboshits, Sov. J. Nucl. Phys. 31 (1980) 509 [[inSPIRE](#)].



$$\Theta_\mu \equiv \angle(\xi_i, \xi_f) = (1 + \gamma a) \Theta$$

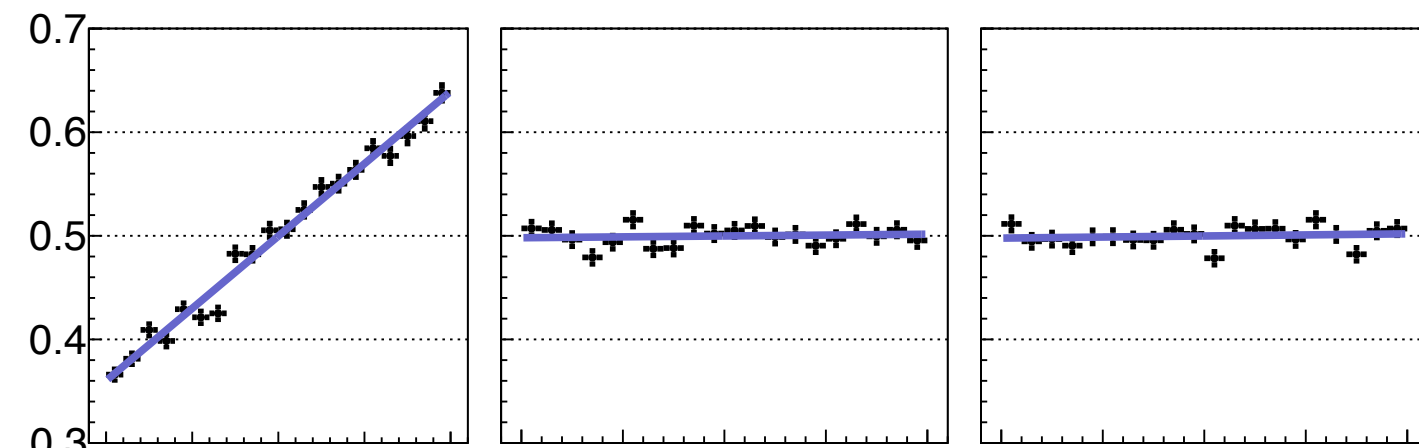
$$a = \frac{g - 2}{2}, \quad \Theta = \frac{L}{R}$$

γ, g, a – Lorentz factor, g-factor, anomalous MDM of Λ_c

Θ, L, R – deflecting angle, length, curvature radius of the crystal

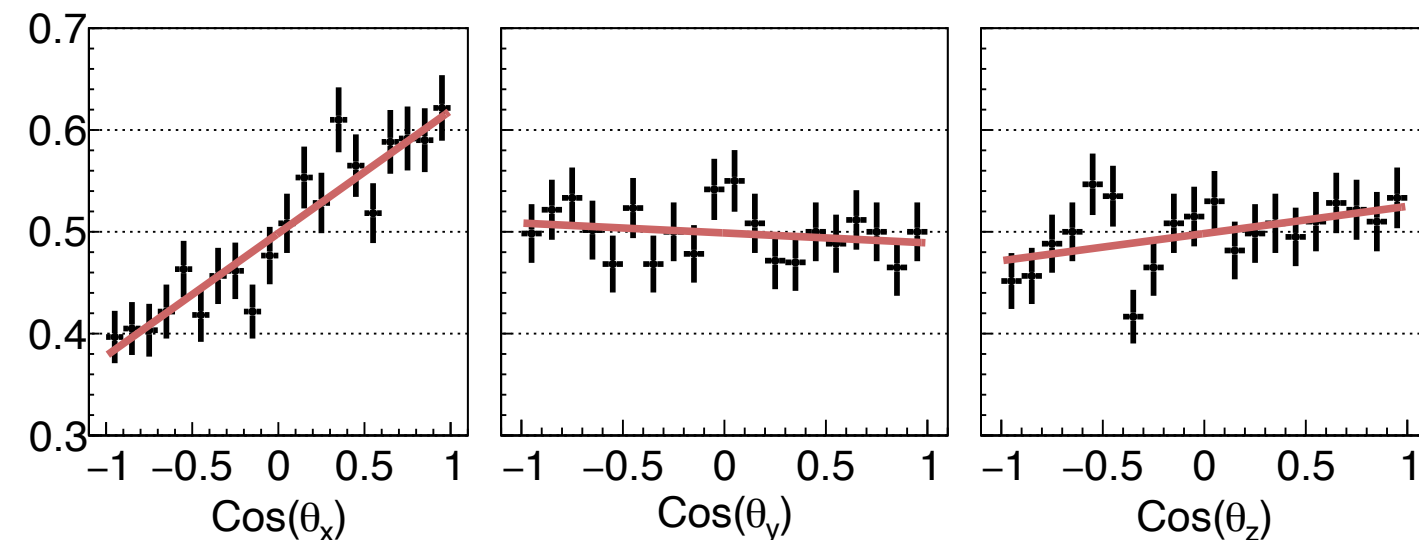
Initial polarisation

$$\vec{\xi}_i = \xi (1, 0, 0)$$

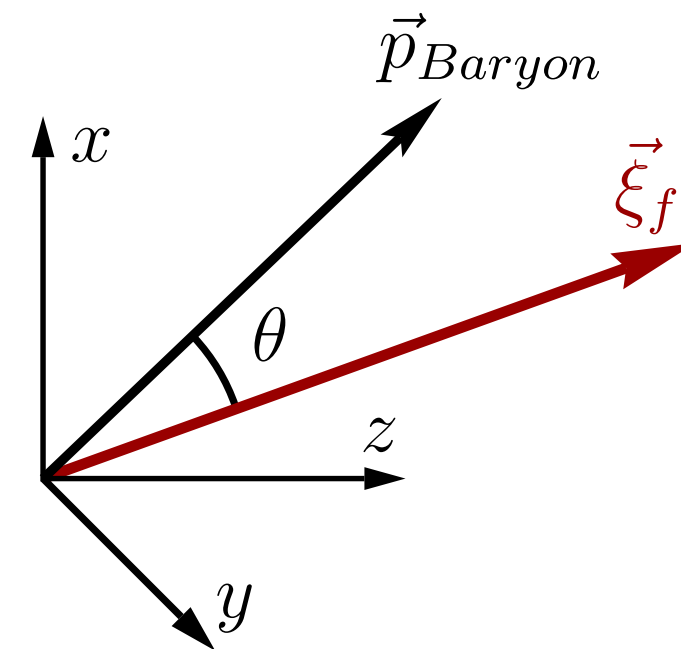


Final polarisation

$$\vec{\xi}_f = \xi (\cos \Theta_\mu, 0, \sin \Theta_\mu)$$



$\Lambda_c^+ \rightarrow \text{Meson} + \text{Baryon}$



$$\frac{dN}{d \cos \theta_z} = \frac{1}{2} \left(1 + \alpha \xi_{fz} \cos \theta_z \right)$$

$$b \equiv \alpha \xi \Theta_\mu \quad \Delta b = \sqrt{\frac{3}{N}}$$

$$\Delta g = \frac{2}{\alpha \langle \xi \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

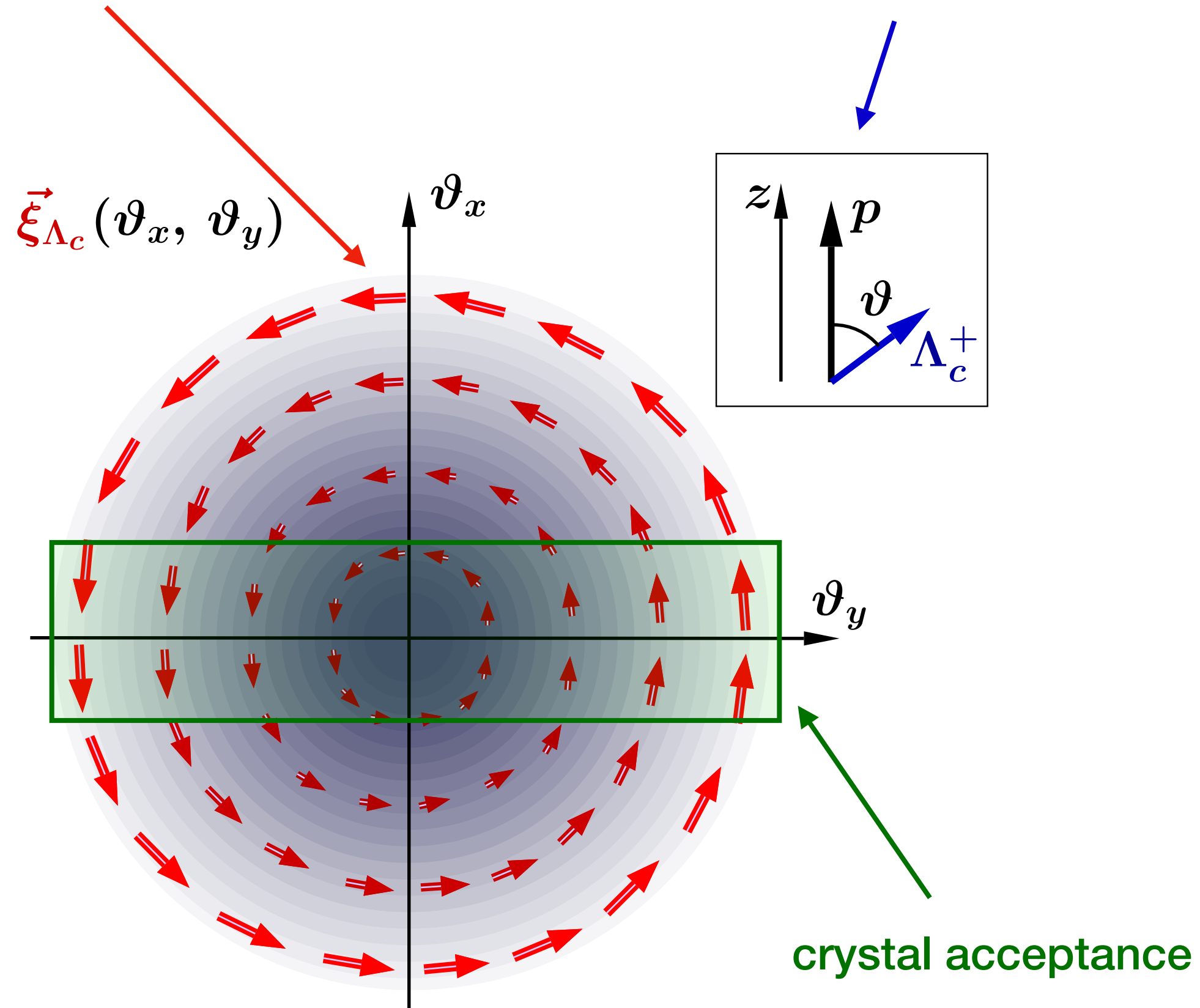


Optimal crystal orientation for MDM and EDM measurements

A. Fomin et al. Eur. Phys. J. C (2020) 80:358 [1909.04654]

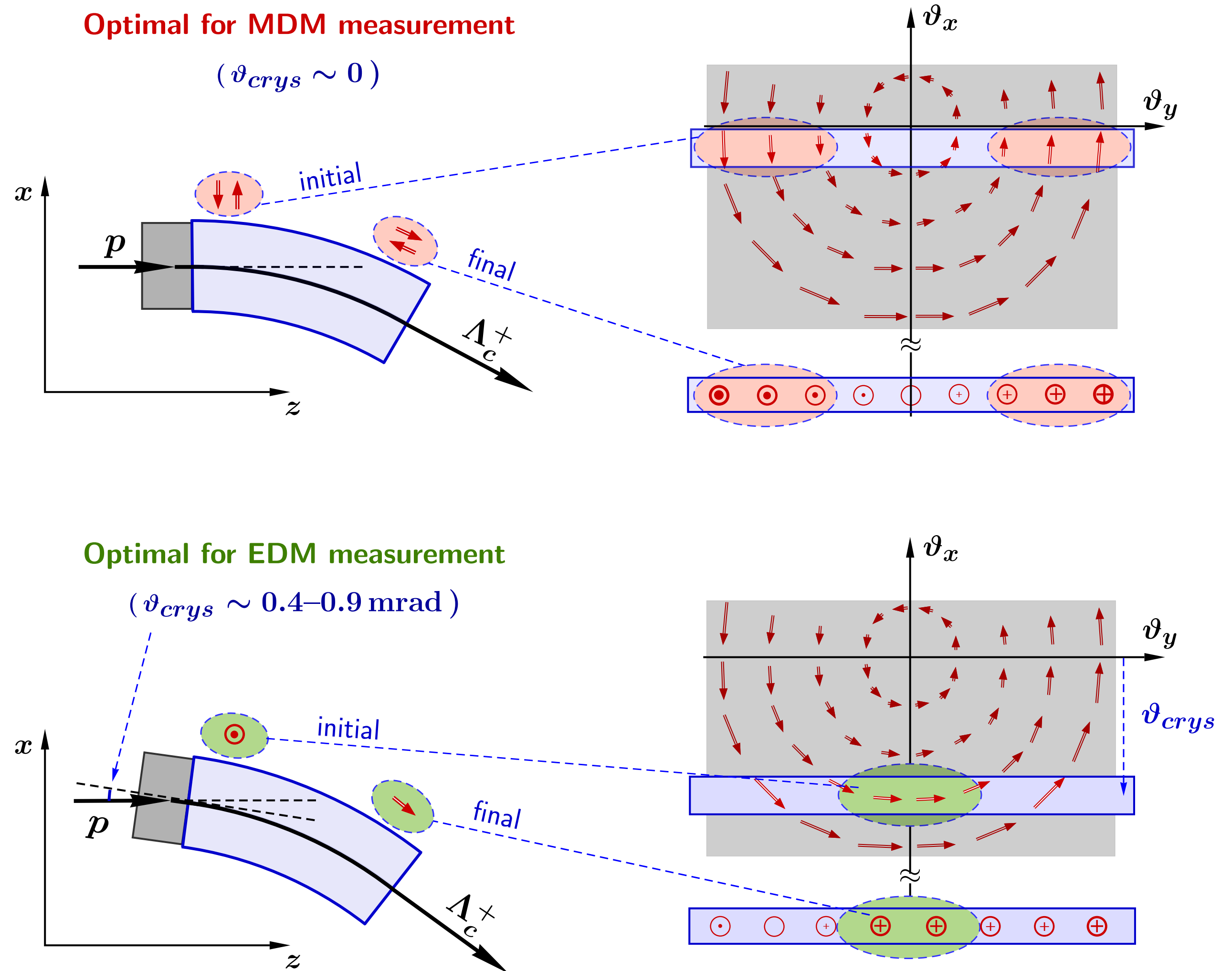
Production of Λ_c^+ in a fixed target $p + p \rightarrow \Lambda_c^+ + X$

Due to the space-inversion symmetry of the strong interaction Λ_c^+ **polarisation** is perpendicular to the **reaction plane**



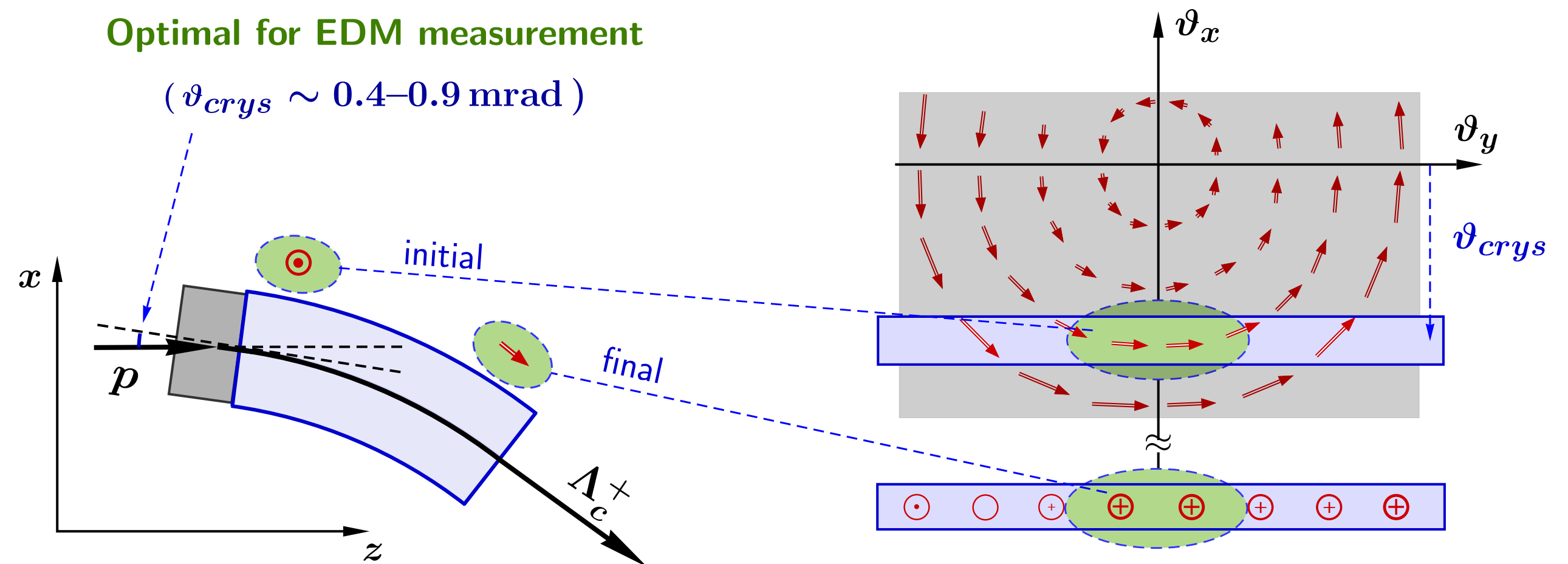
Optimal for MDM measurement

$(\vartheta_{crys} \sim 0)$



Optimal for EDM measurement

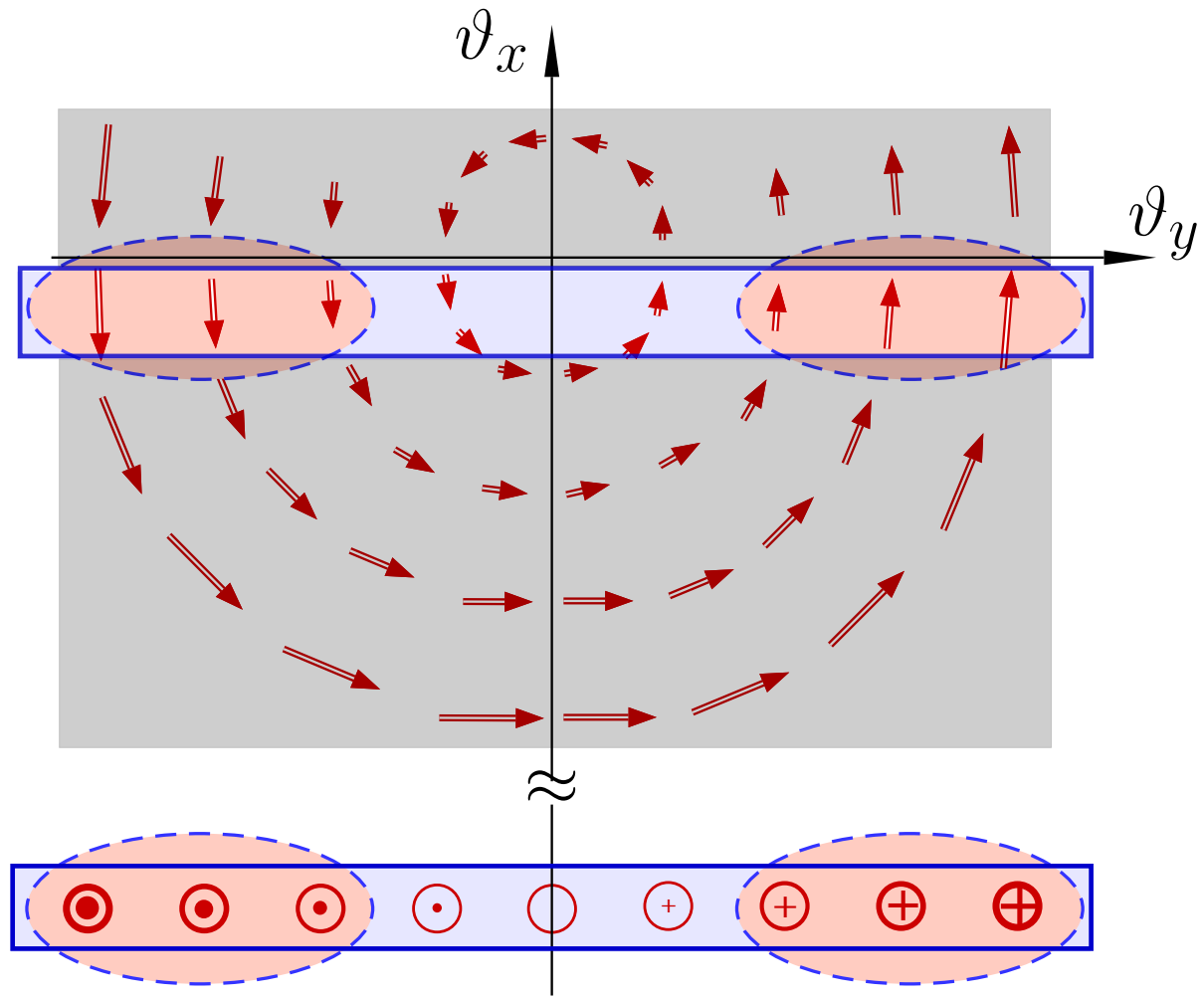
$(\vartheta_{crys} \sim 0.4-0.9 \text{ mrad})$



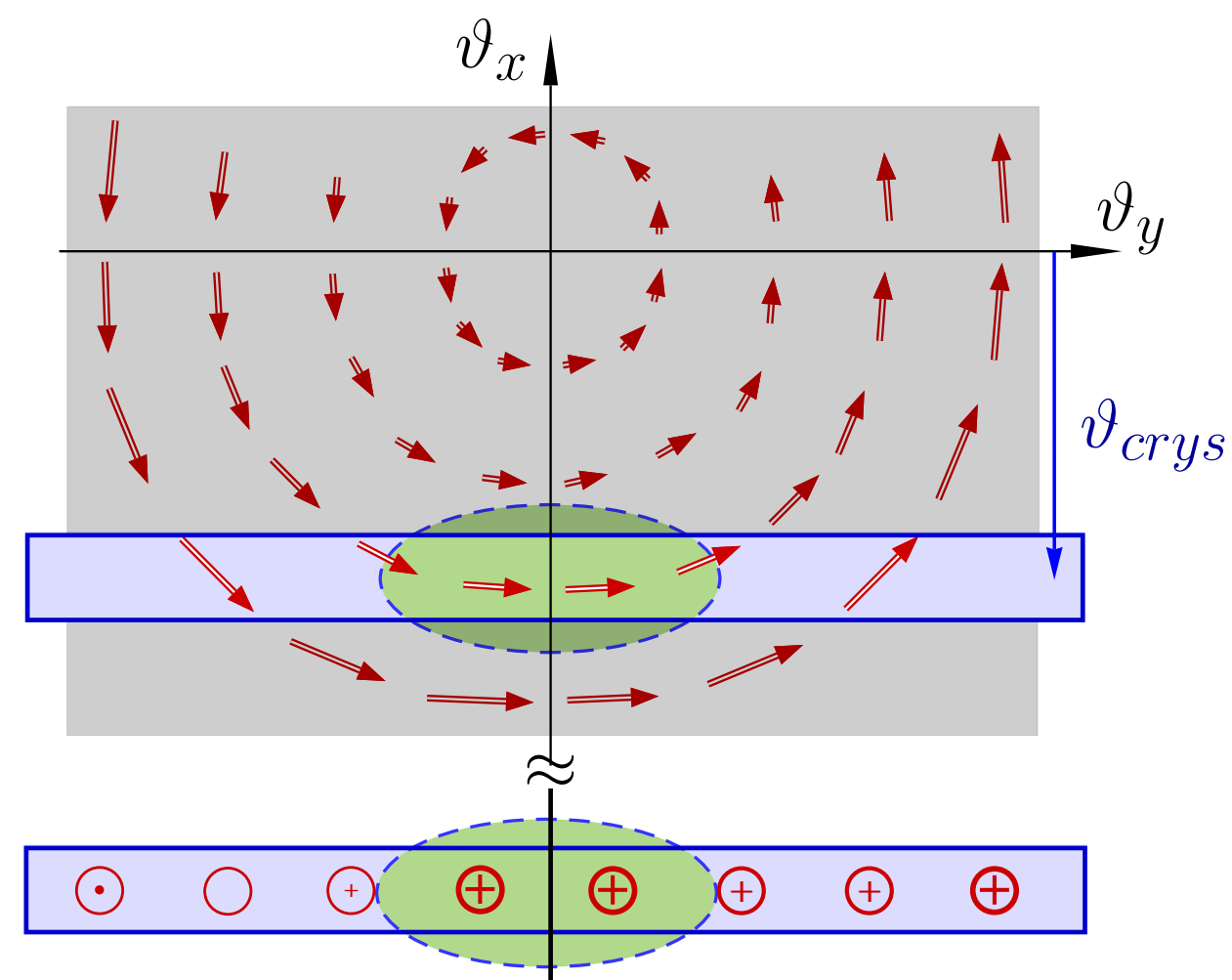


Optimal crystal orientation for EDM measurement: Quantitative analysis

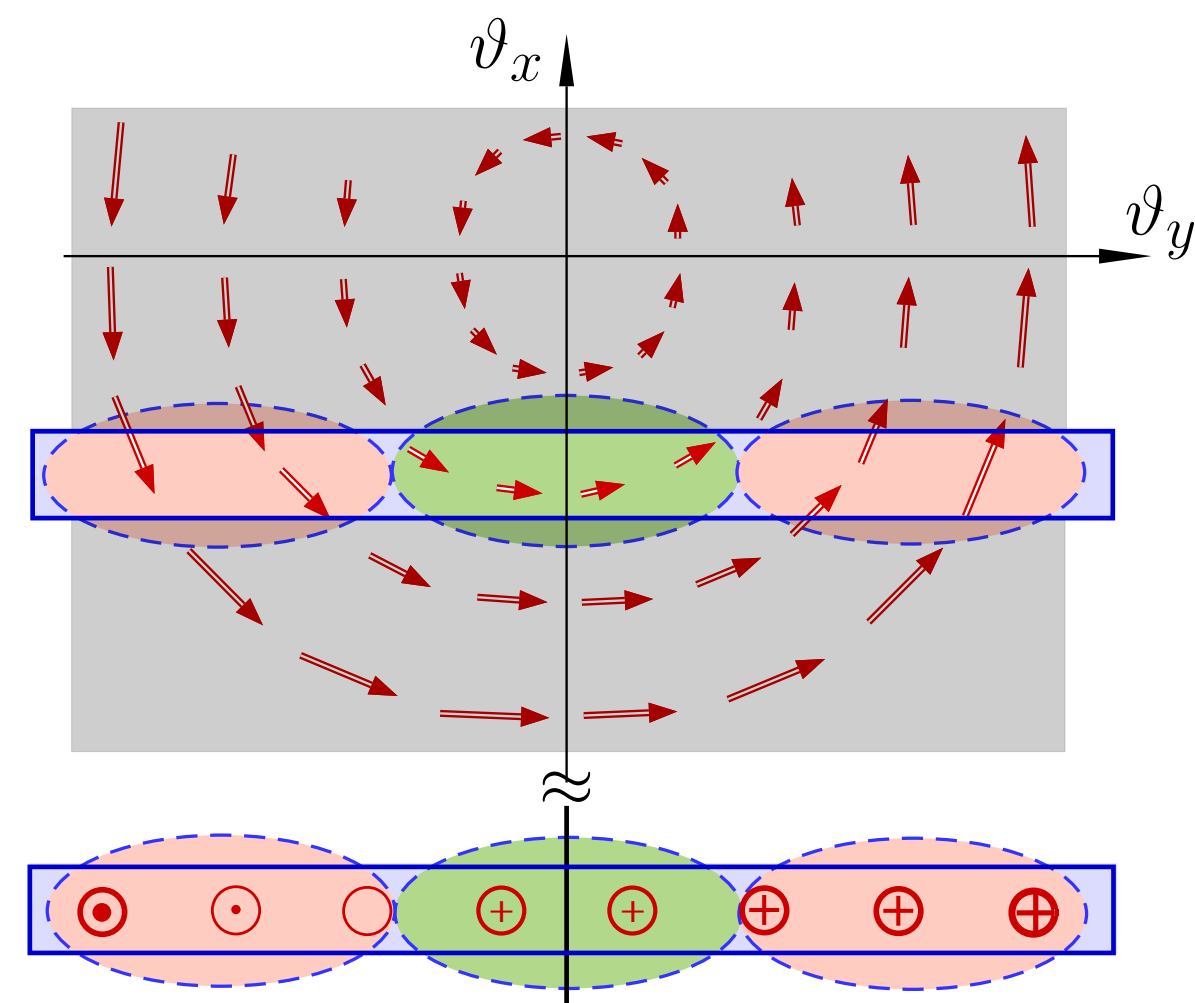
Optimal for MDM measurement



Optimal for EDM measurement



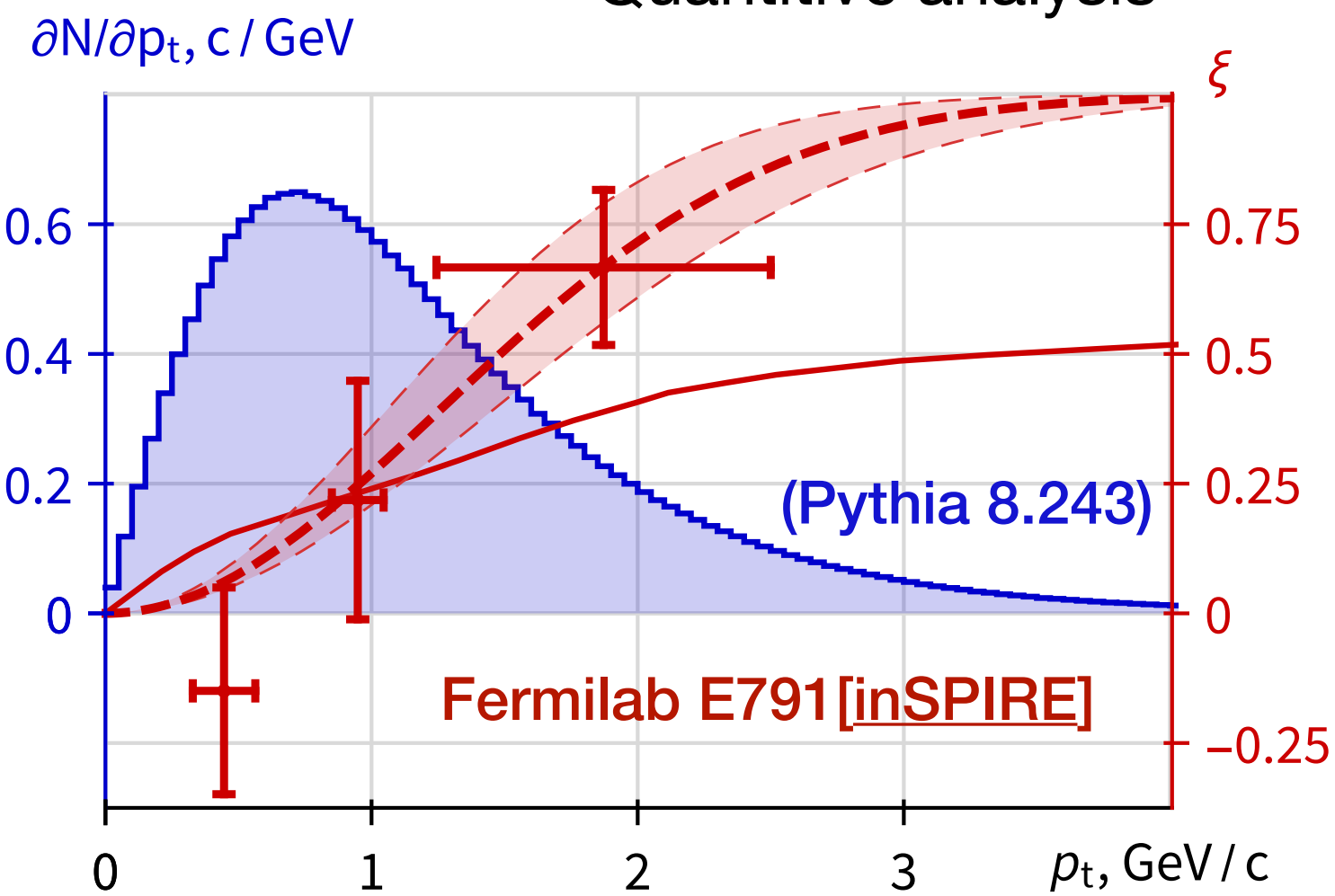
Simultaneous measurement



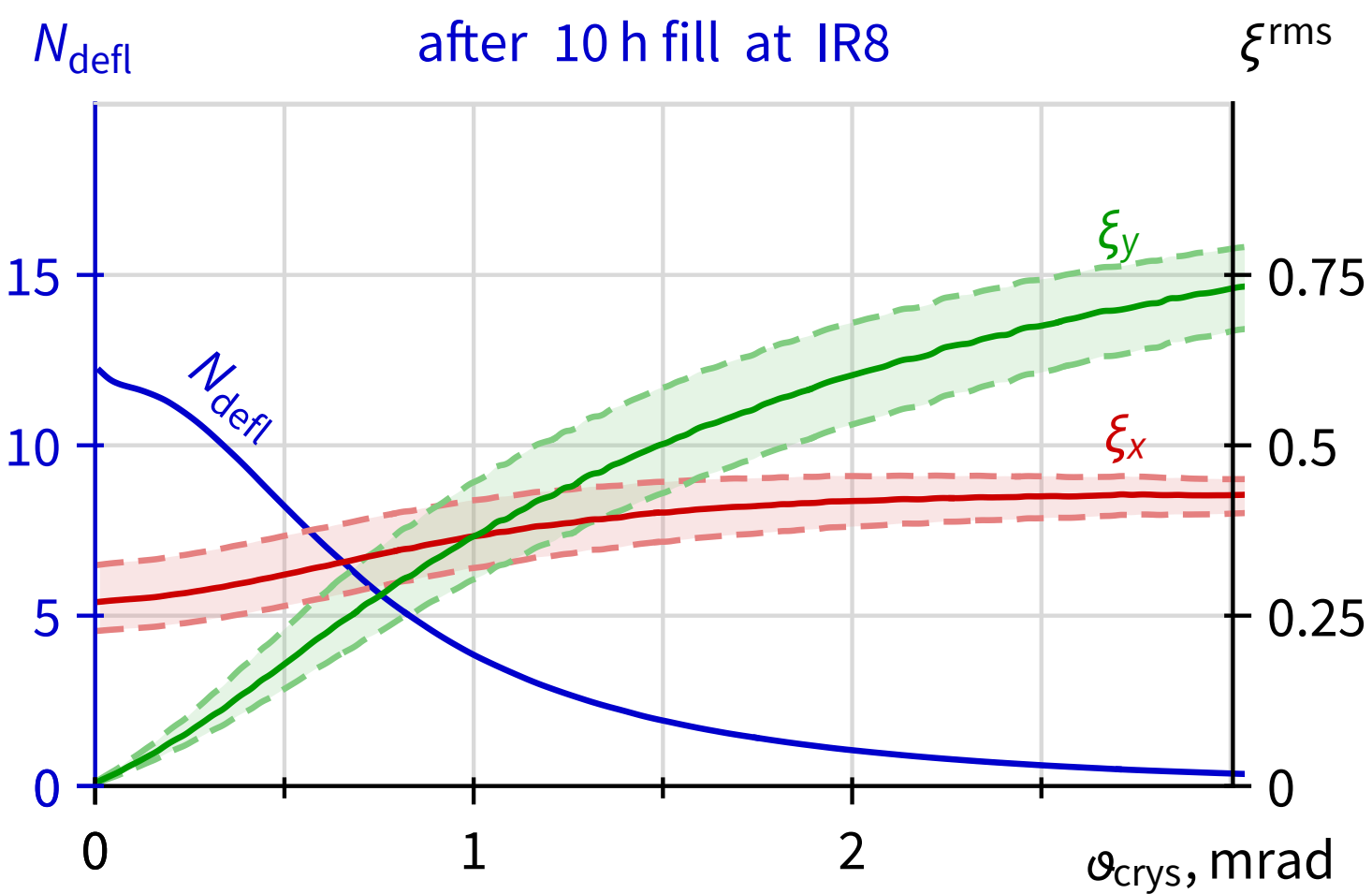
$$\Delta g = \frac{2}{\alpha \langle \xi_x \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

$$\Delta f = \frac{2}{\alpha \langle \xi_y \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

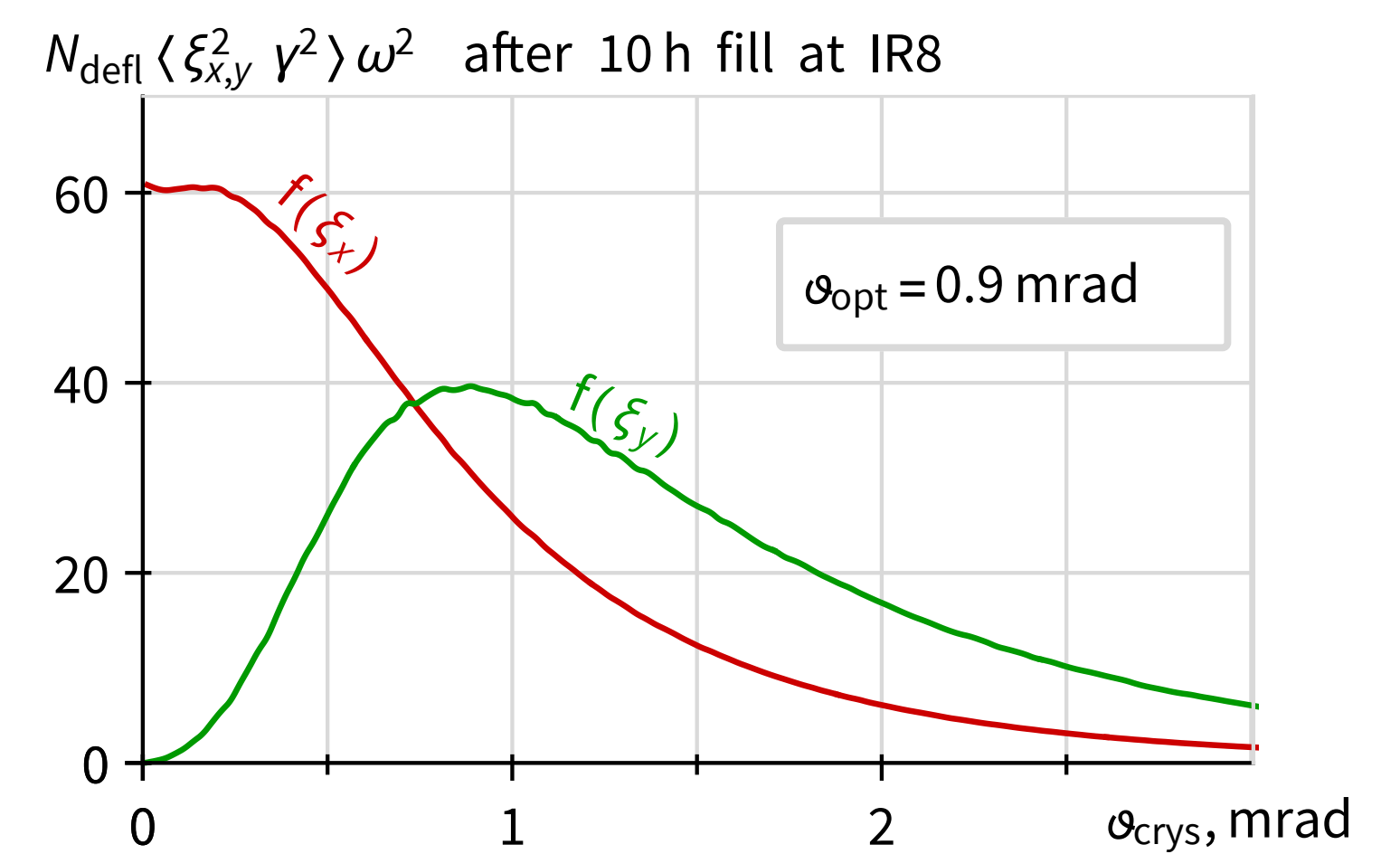
Quantitative analysis



Initial polarisation of deflected Λ_c^+



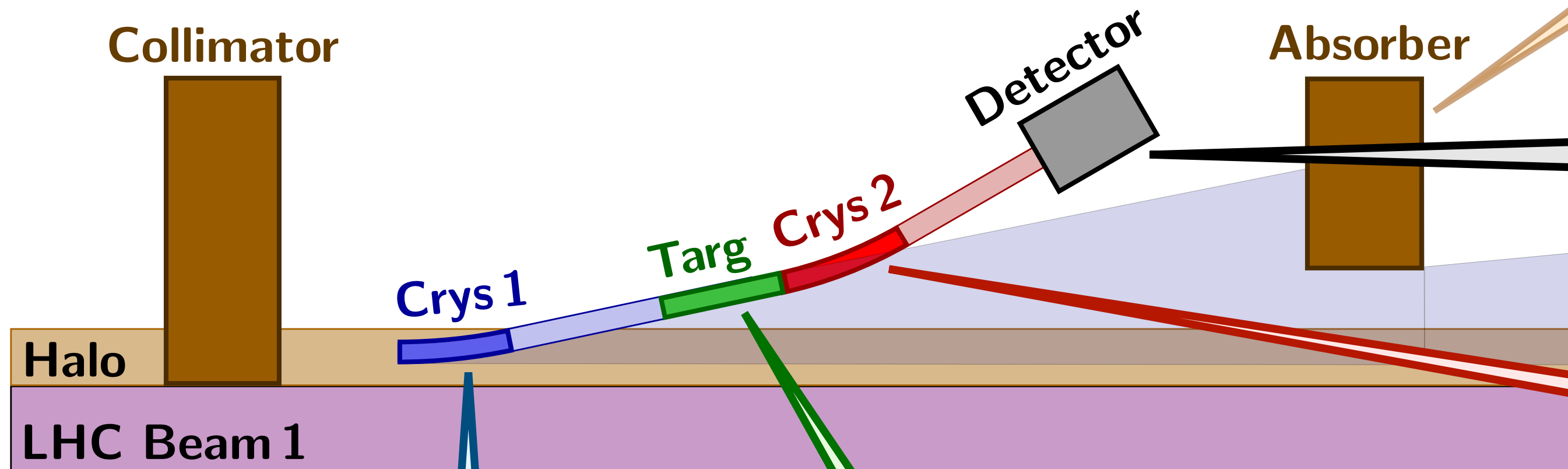
Measurement efficiency





MDM of charmed baryons: Fixed target at the LHC

- L. Burmistrov et al., CERN-SPSC-2016-030, CERN, Geneva Switzerland, June 2016 [[SPSC-EOI-012](#)].
- A. Stocchi, W. Scandale, [talks at Physics Beyond Collider Workshop](#), CERN, Geneva Switzerland, 6–7 September 2016.



Beam halo particles that do not interact with the Target+Crys2 assembly are intercepted by 4 double-sided LHC-type collimators

In the Detector the final polarisation of Λ_c is reconstructed from the distribution of decay products

The second Crystal deflects Λ_c with specific initial polarisation.

Λ_c spin precession in the electric field of crystal planes is proportional to MDM (or EDM)

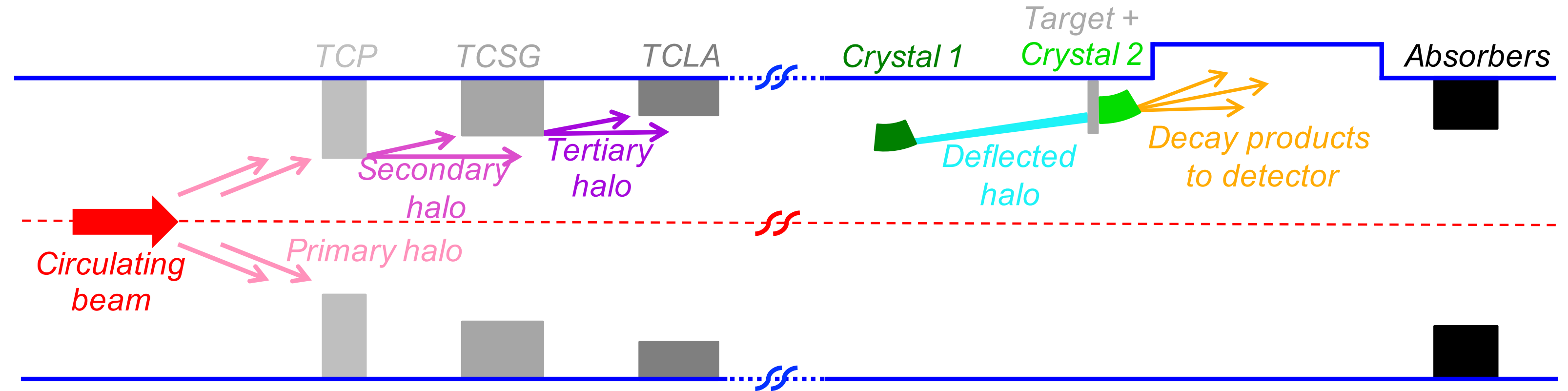
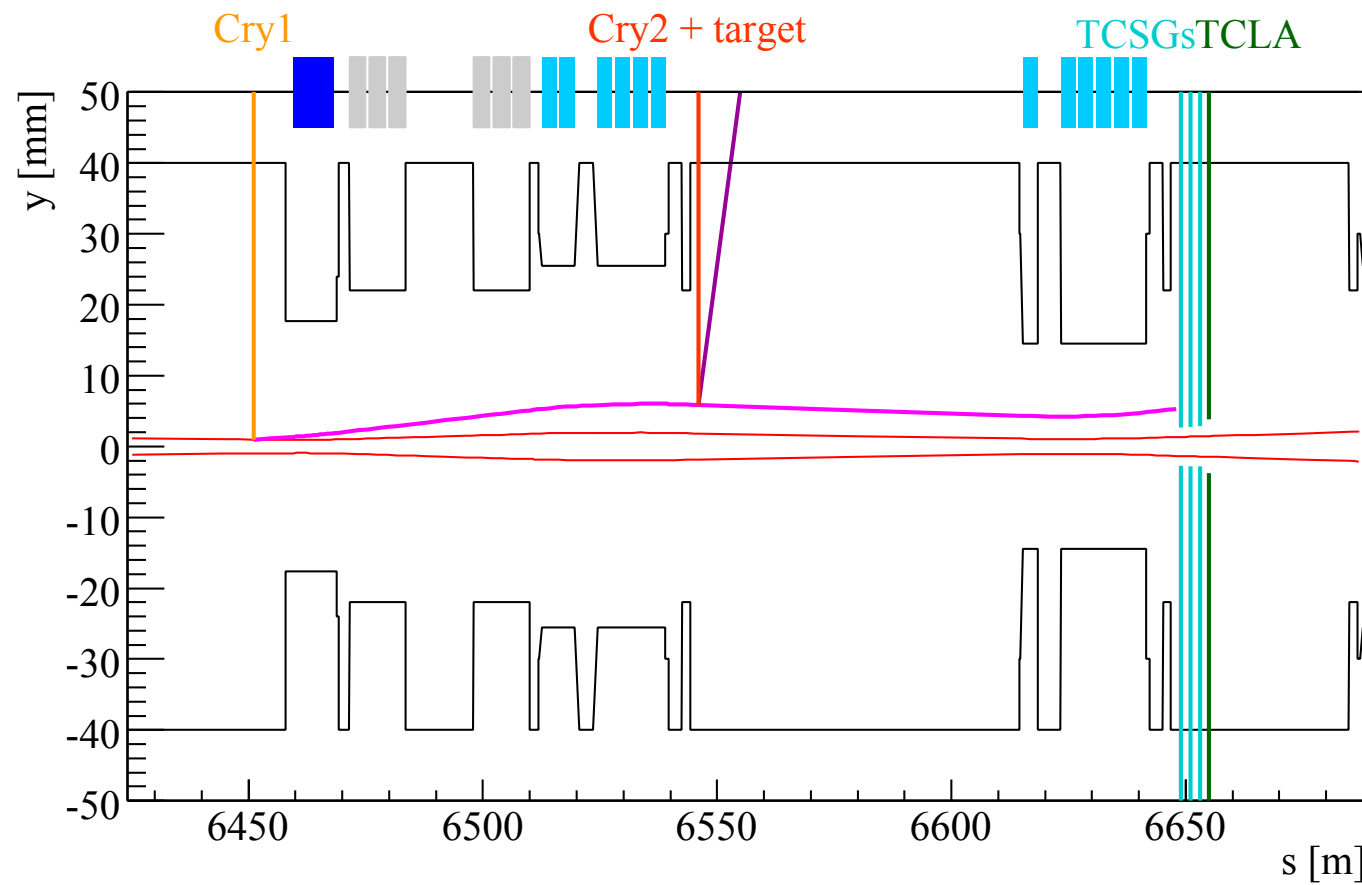
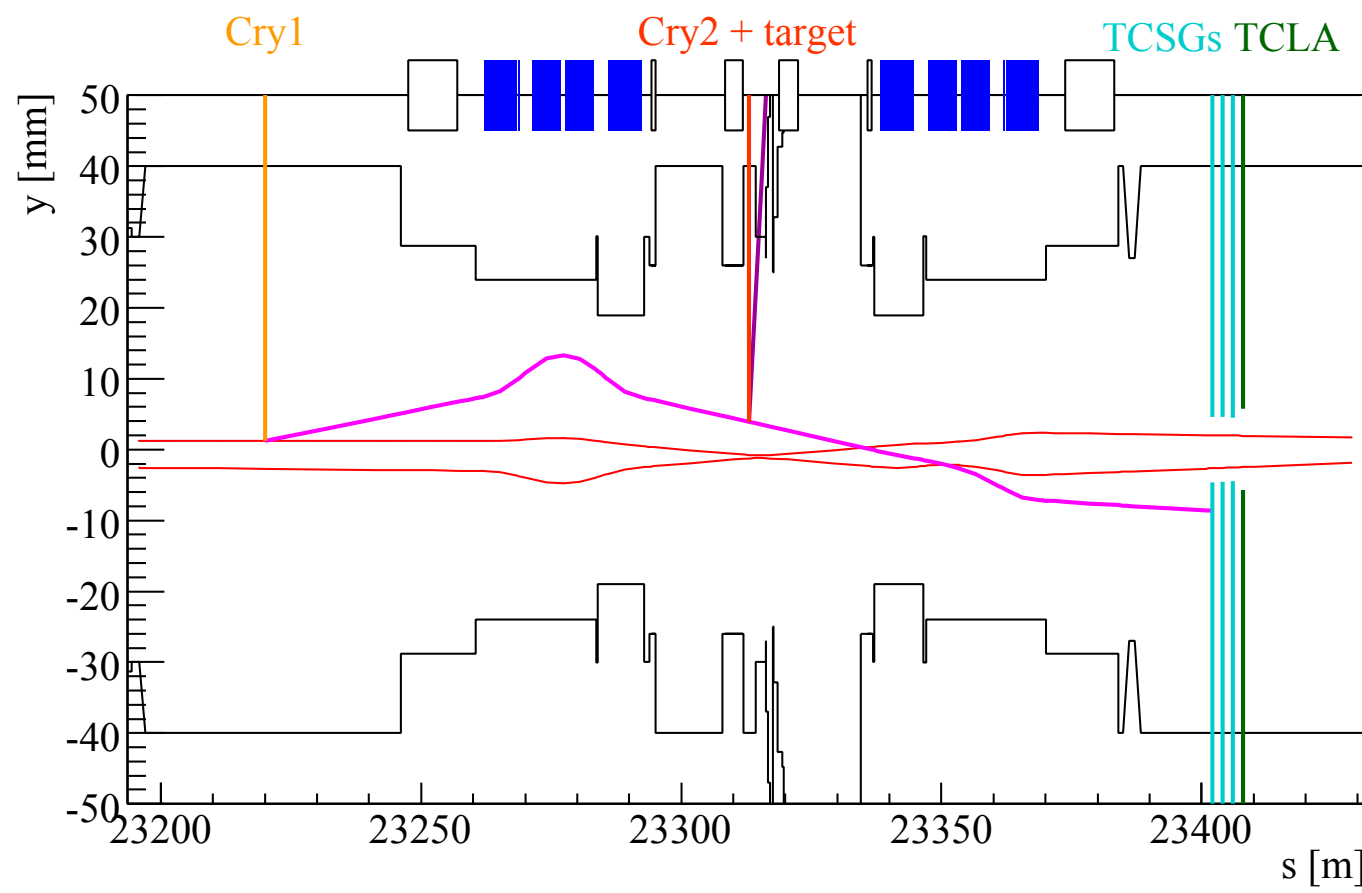
The first Crystal deflects protons from the LHC beam halo onto the Target

In the Target protons are converted to polarised Λ_c

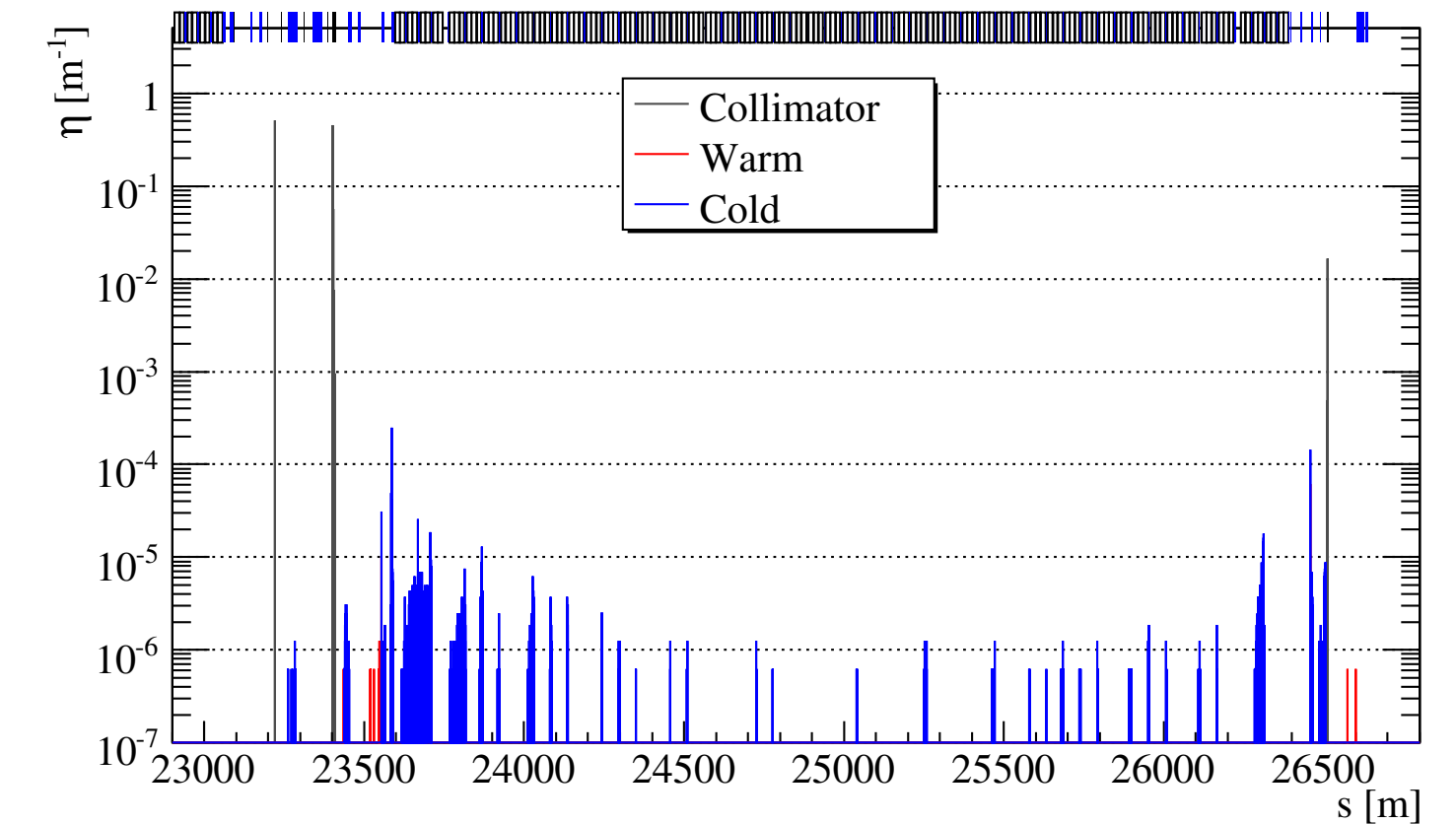


Performance assessment of layouts in IR3 and IR8 of LHC

D. Mirarchi et al. Eur. Phys. J. C 80 (2020) 10, 929



- impact on the machine
- optimisation of Crystal 1 and Absorbers positions
- running experiment in a parasitic mode
- layout in front of LHCb (IR8) 4.3×10^{10} POT/fill
- alternative layout at IR3 3.0×10^{10} POT/fill



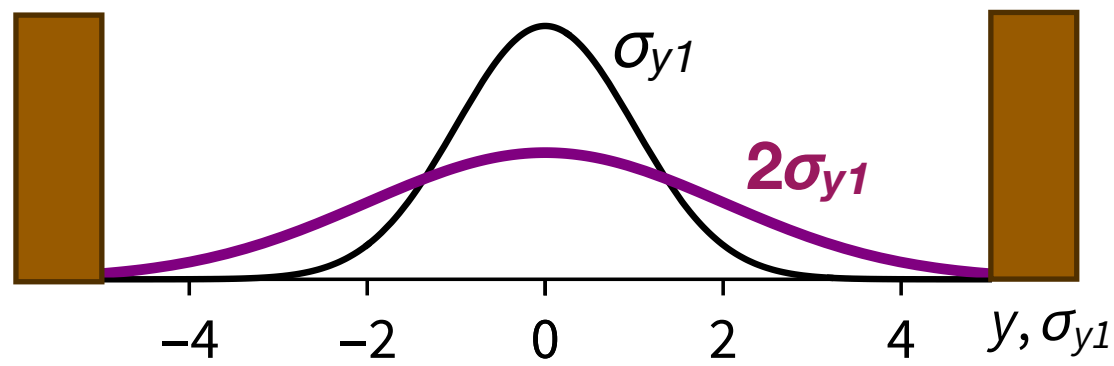


Increasing the PoT through active bunch excitation

General idea

- apply random noise excitations with the **transverse damper** when a selected trains of bunches pass by

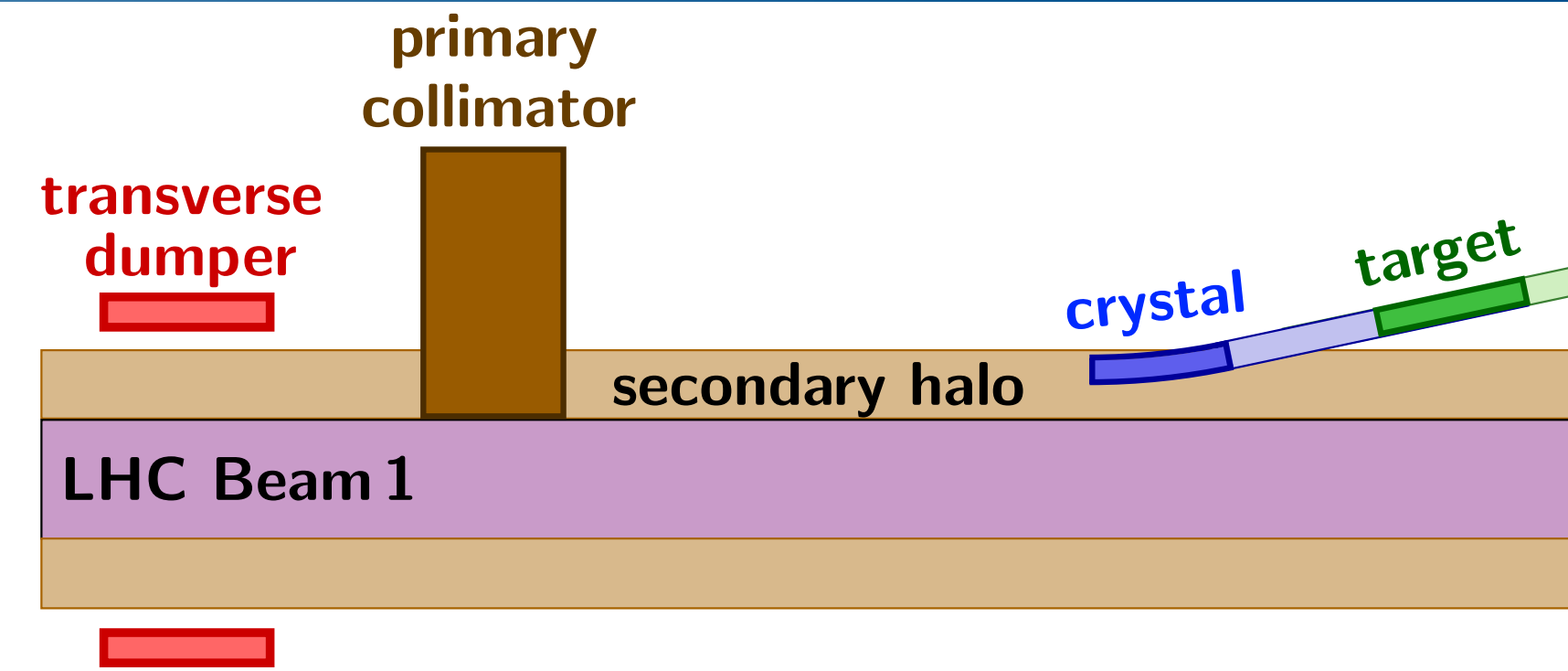
→ **emittance grow** of these bunches



→ increase the loss rate in **primary collimators** at IR7

→ enrich the **secondary halo** of the beam

→ increase the flux on the **crystal** → increase the deliverable rates of **PoT**

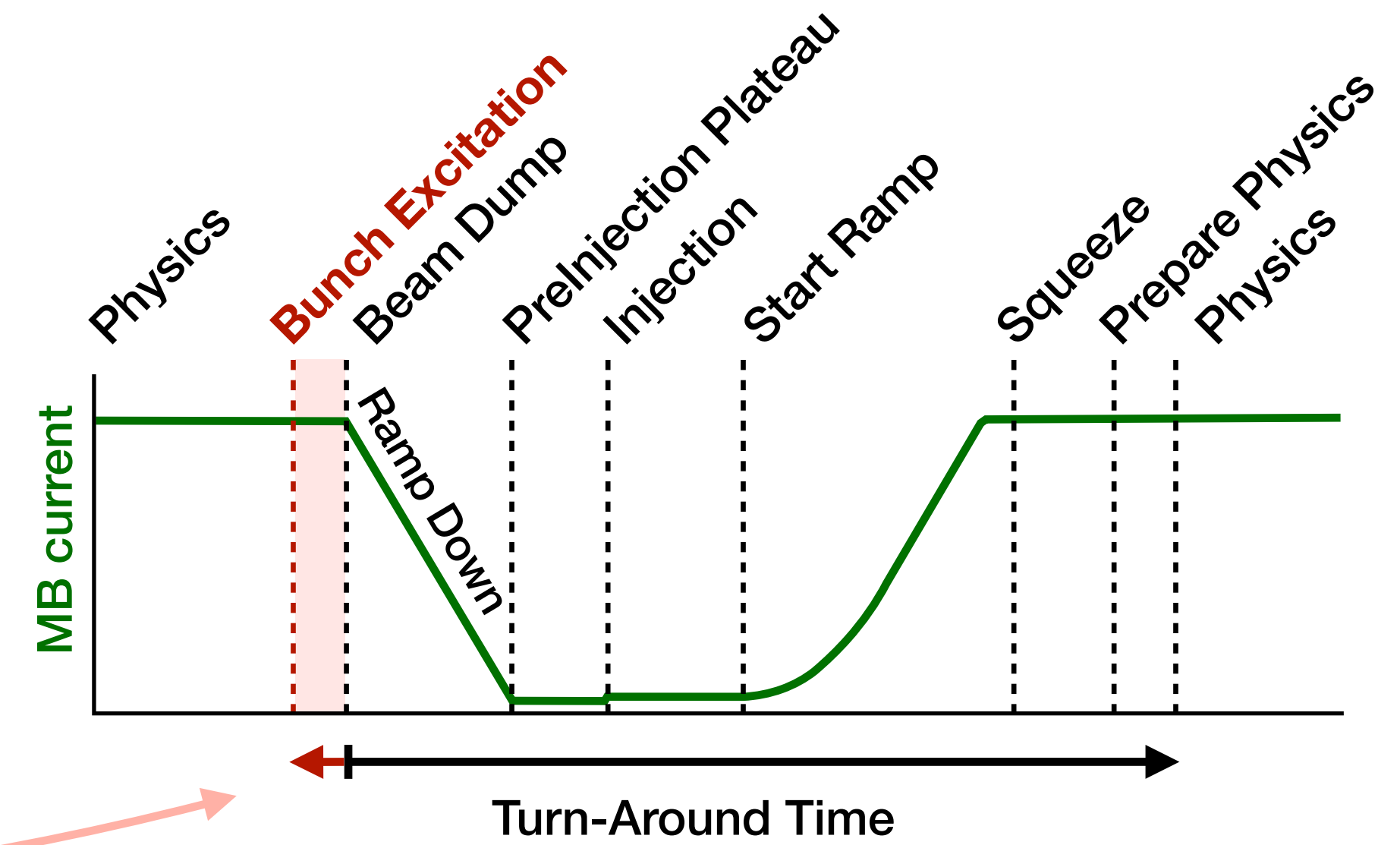
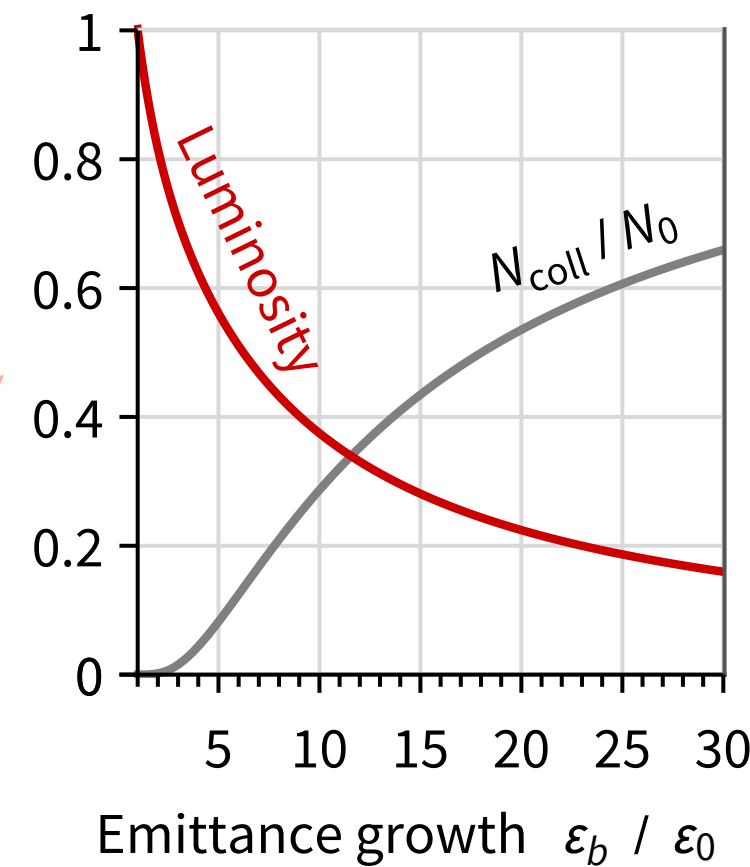


Transverse damper is currently used in LHC

- corrects the trajectory of the bunches during the fill
- increase the loss rate during alignment of collimators

Drawback — reduction of total luminosity

- excitation during the fill → higher emittance
- excitation before dumping → longer turn-around time





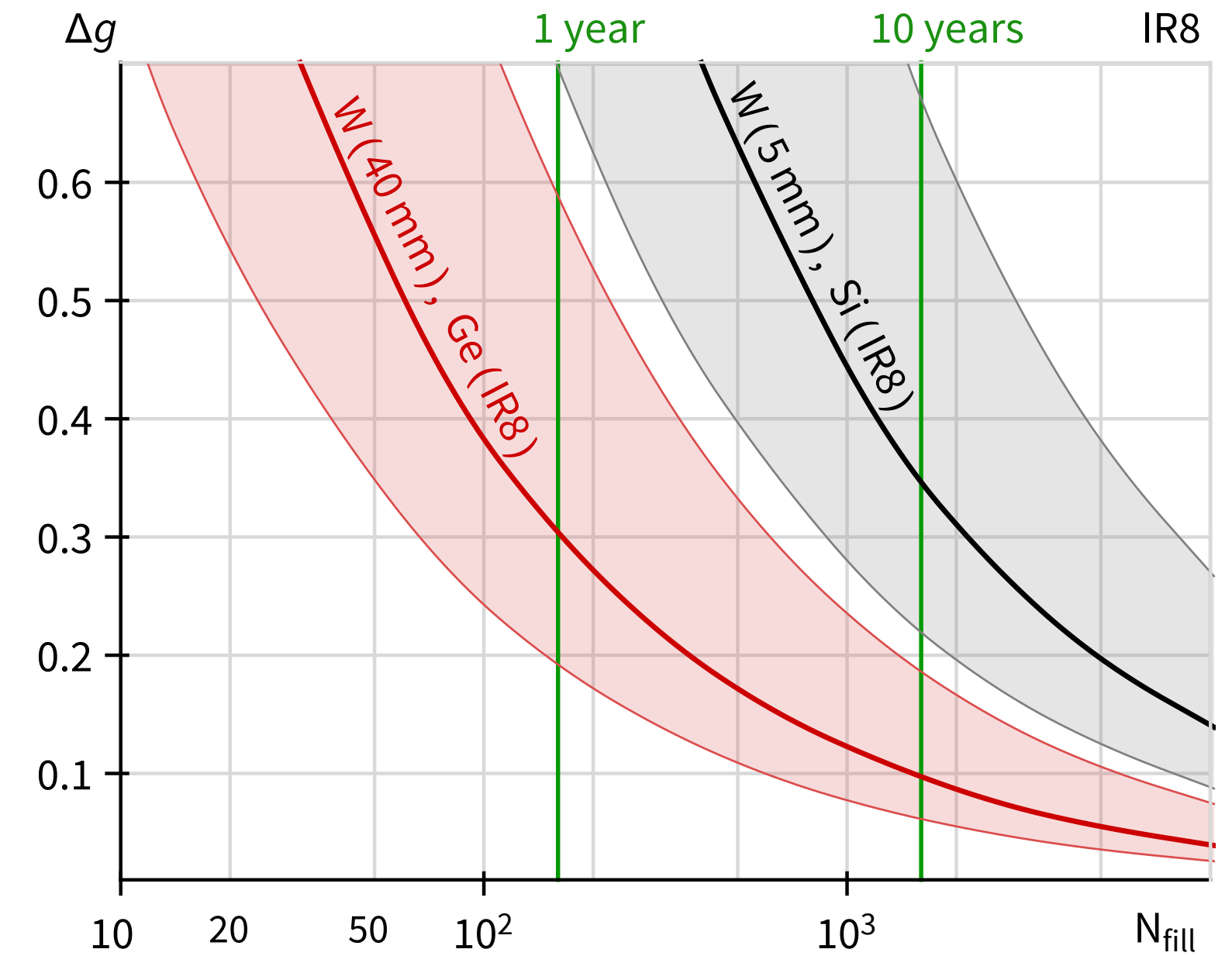
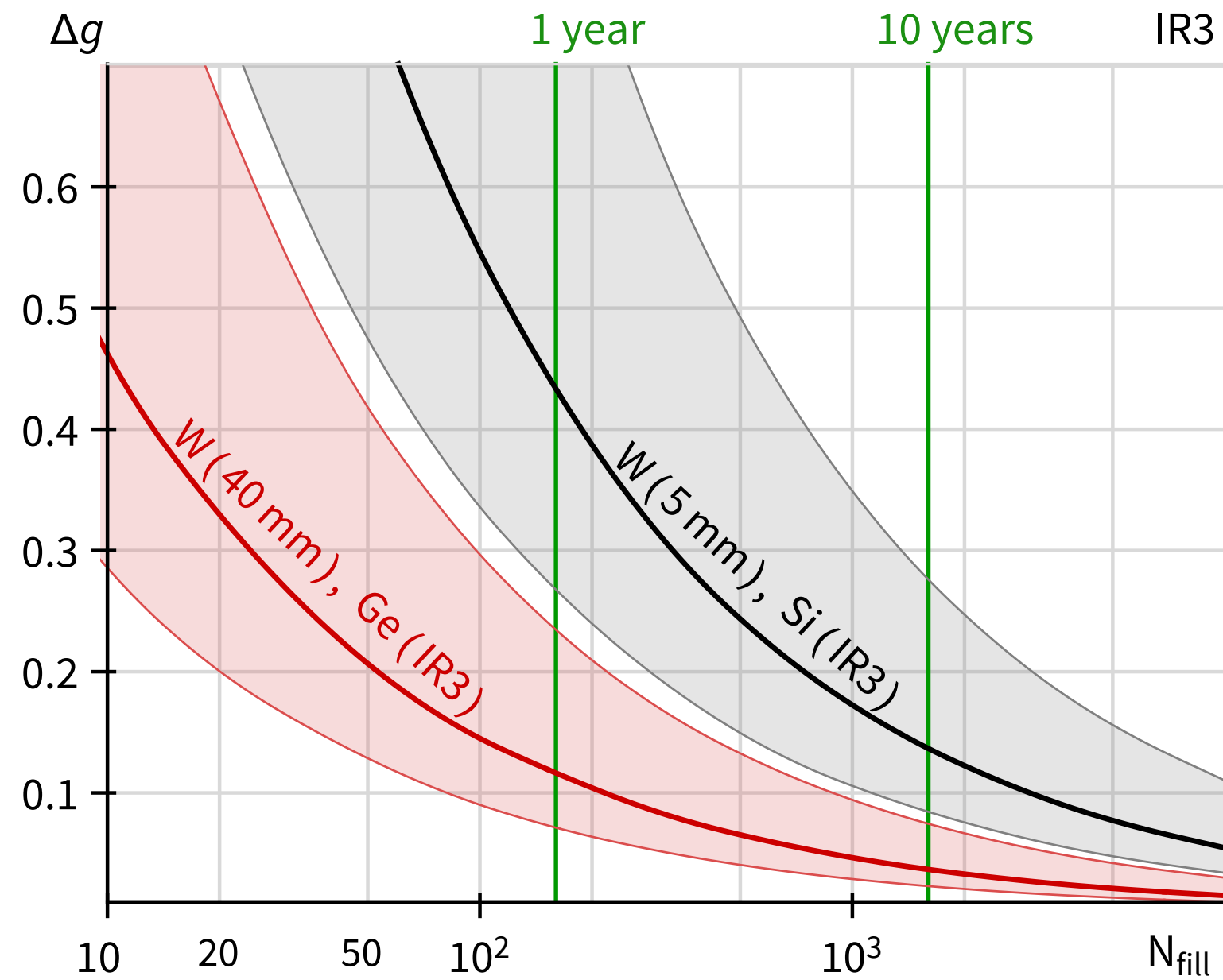
Performance assessment of layouts in IR3 and IR8: possible improvements

[A. Fomin et al. EPJ C80 \(2020\) 358](#)

- Thicker target 5 mm → 40 mm: ionisation energy losses and multiple scattering can be neglected, **showers production - to be checked**

- Proton rate, $3-4.3 \times 10^{10}$ per 10h fill

[D. Mirarchi et al. EPJ C80 \(2020\) 10, 929](#)



Possible improvements:

| | 1 → 2 | t1 / t2 |
|-----------------|----------------------------------|---------|
| Target | 5 mm → 40 mm | 4 |
| Crystal | silicon → germanium | 4 |
| Detector | LHCb (IR8) → dedicated at IR3 | 5 |
| Beam excitation | additional to parasitic scenario | 4 |

- **10 years** at LHCb, $\sim 7 \times 10^{13}$ POT, 5mm, Si → $\Delta g \sim \mathbf{0.35}$
- **1 year** at IR3, $\sim 0.5 \times 10^{13}$ POT, 40mm, Ge → $\Delta g \sim \mathbf{0.12}$
- **10 years** at LHCb, $\sim 7 \times 10^{13}$ POT, 40mm, Ge → $\Delta d \sim \mathbf{2.6 \cdot 10^{-16} e cm}$ (optimal orientation → data taking time reduced by **~170**)
- big uncertainty ($\times 10$) due to α parameter
→ significant improvement in PhD thesis of Elisabeth Neil



MDM of τ -lepton. Double-crystal setup.

A.S. Fomin et al. JHEP 1903 (2019) 156

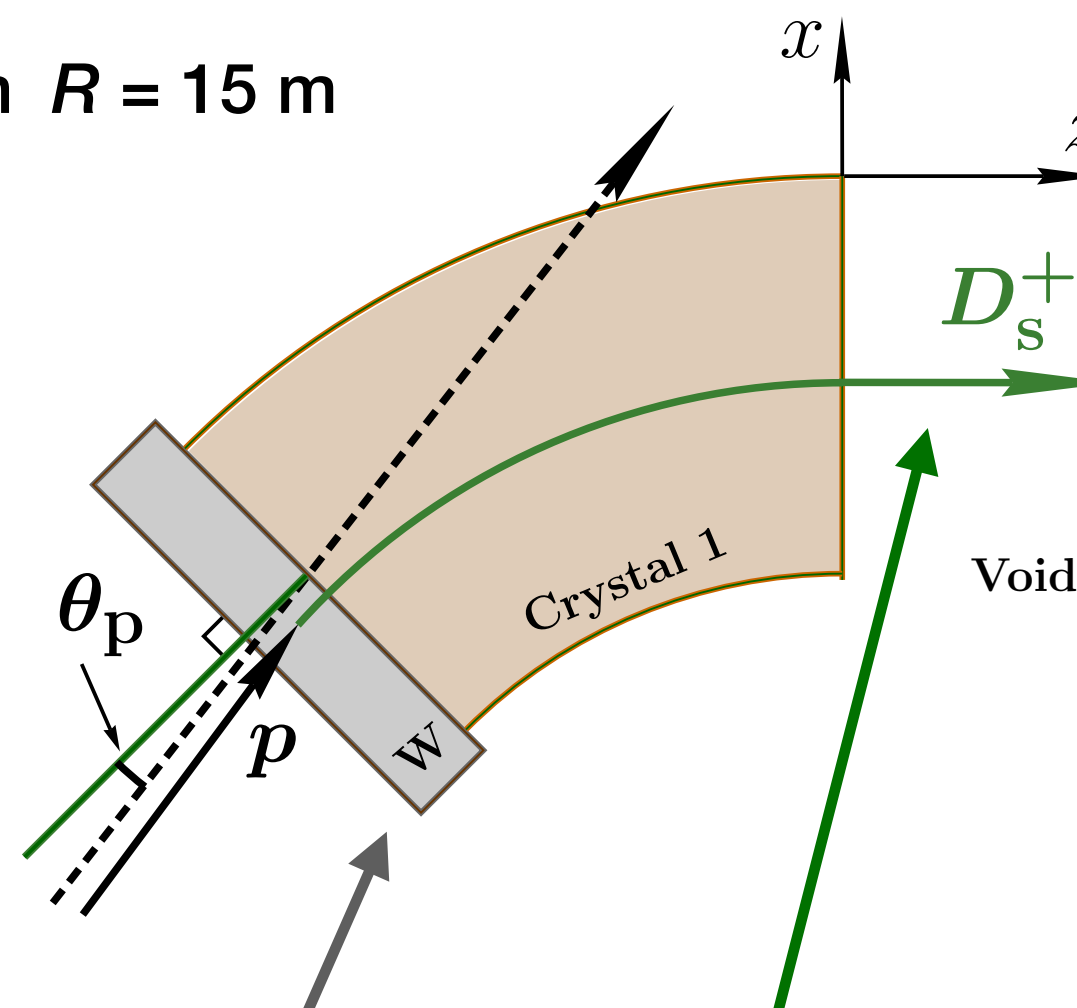
Crystal 1:

Ge: $L = 3$ cm $R = 10$ m

Si: $L = 4.5$ cm $R = 15$ m

$\theta_D = 3$ mrad

$\theta_p = 0.1$ mrad



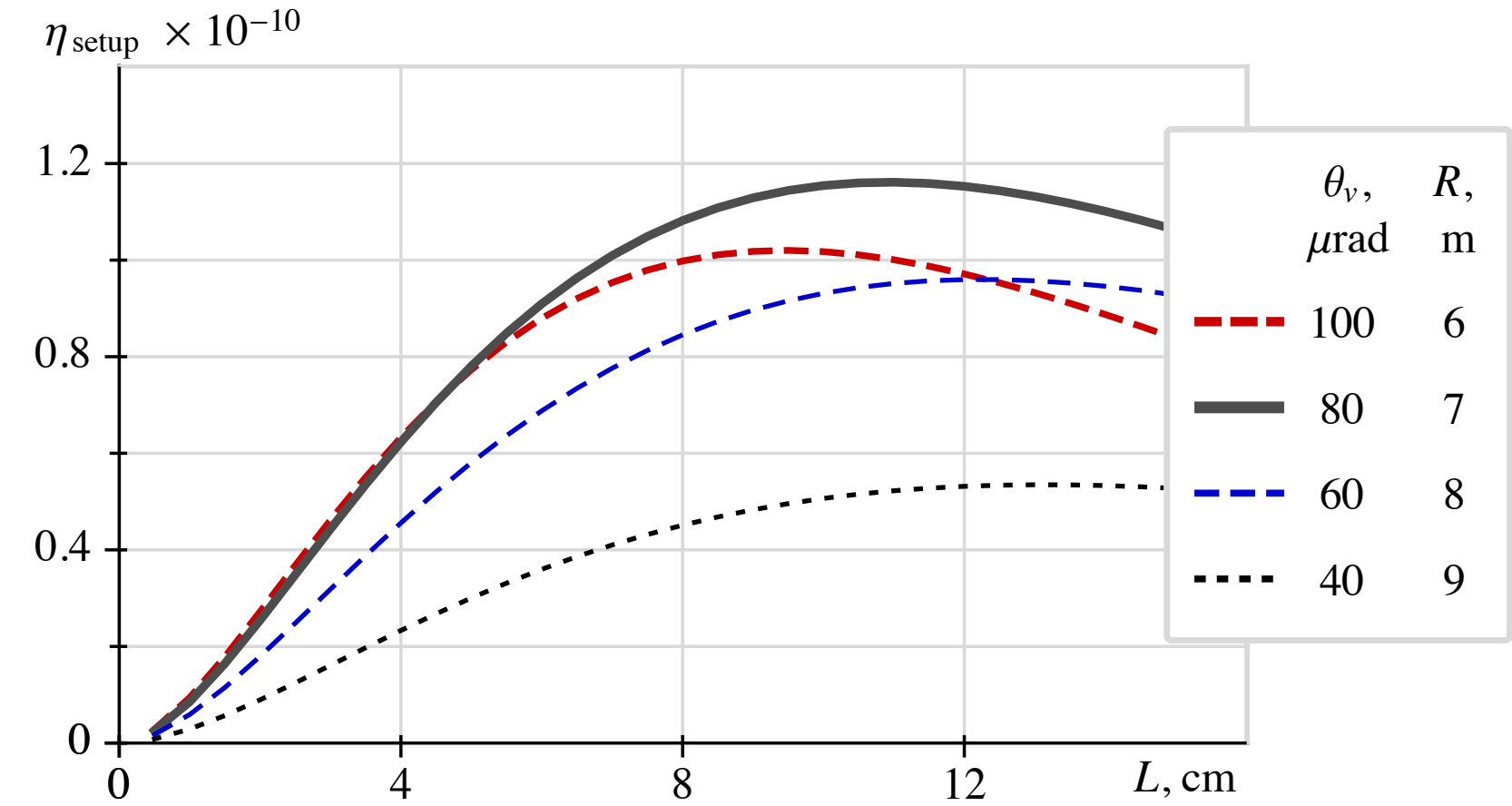
Crystal 2:

Ge: $L = 10$ cm $R = 7$ m

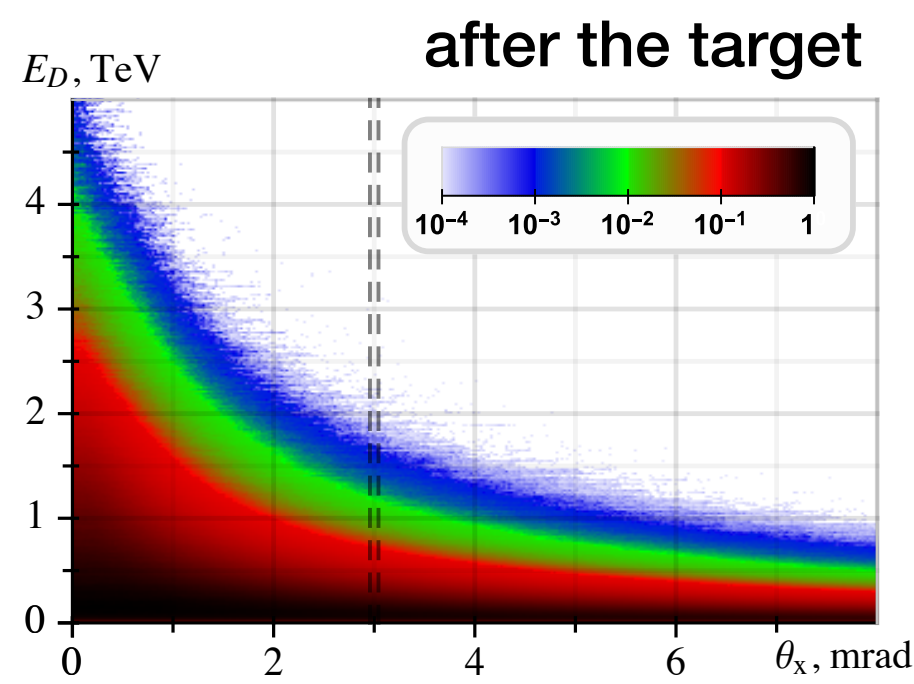
$\theta_\tau = 14$ mrad

$\theta_v = 0.08$ mrad

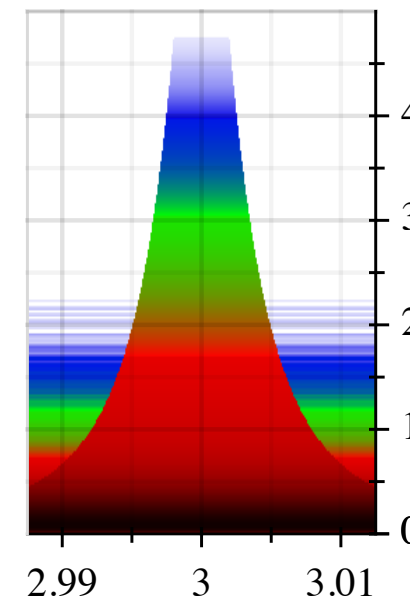
$L_v = 10$ cm



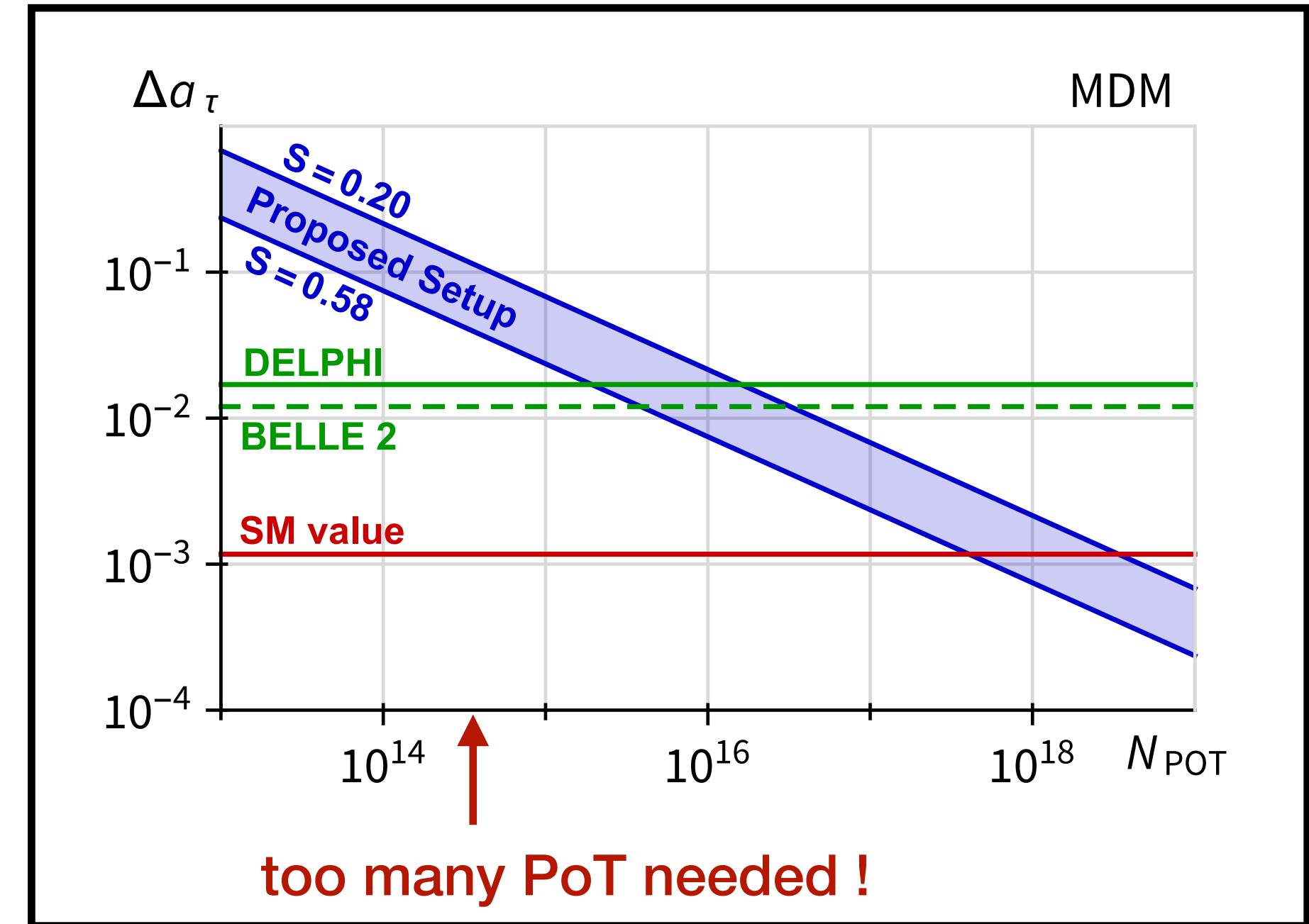
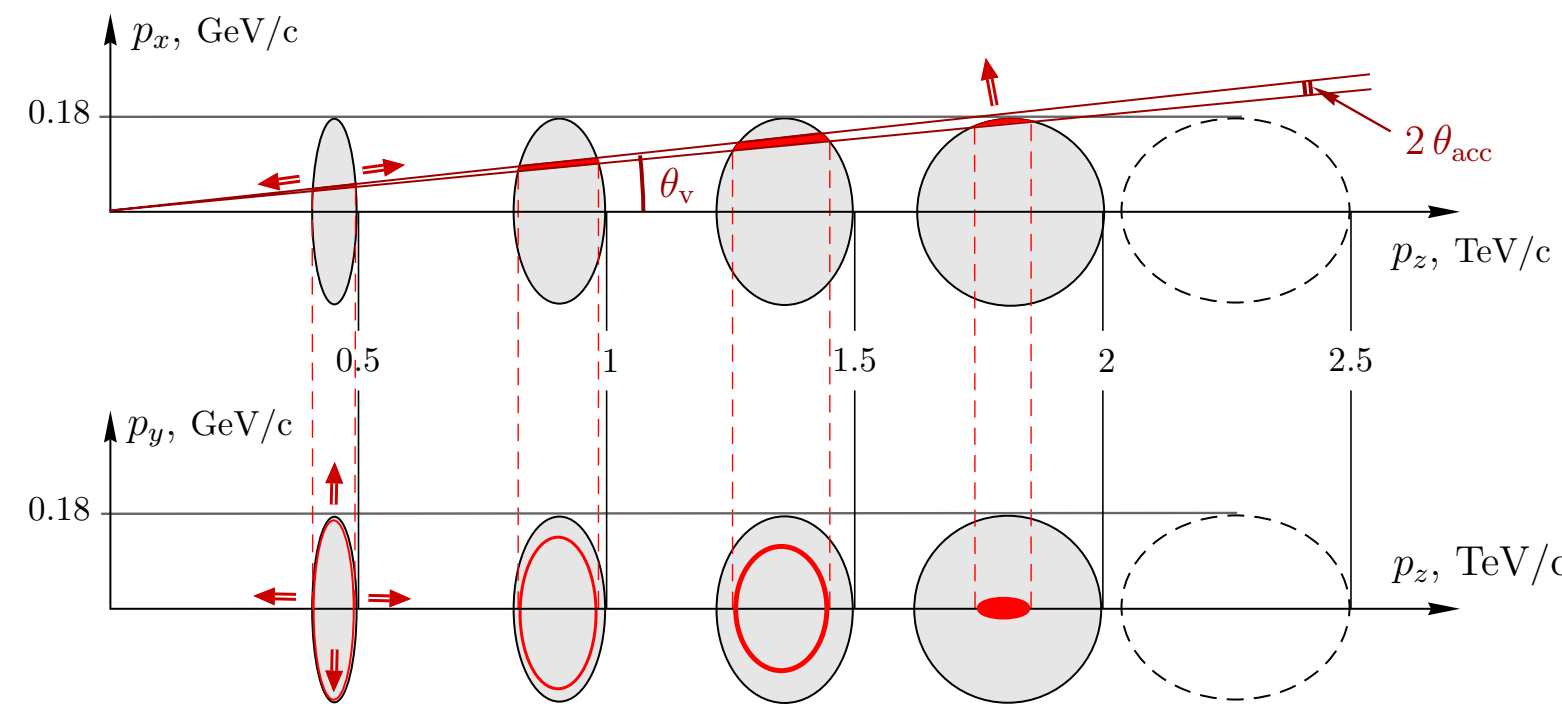
Sp.-ang. distribution of D_s



after Crystal 1



Angular collimation of τ by Crystal 2





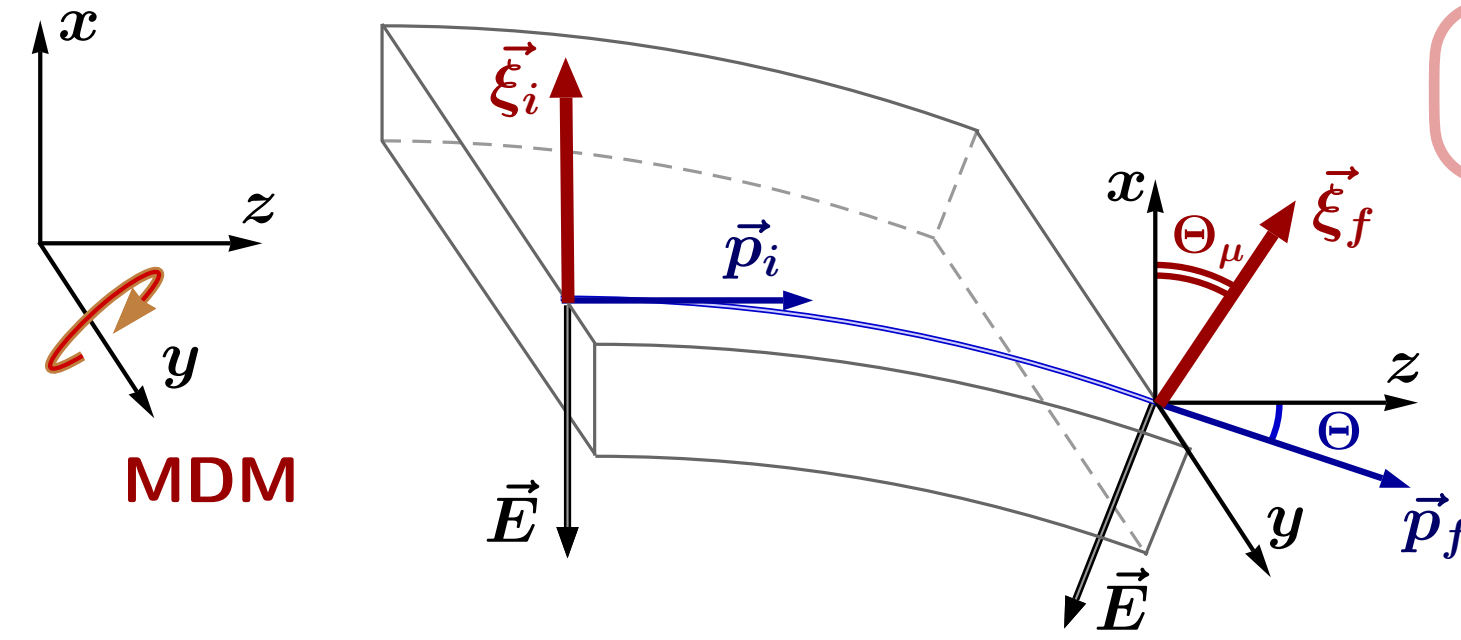
thank you



Optimal crystal orientation for MDM and EDM measurements

V.G. Baryshevsky,
Sov. Tech. Phys. Lett. 5 (1979) 73.

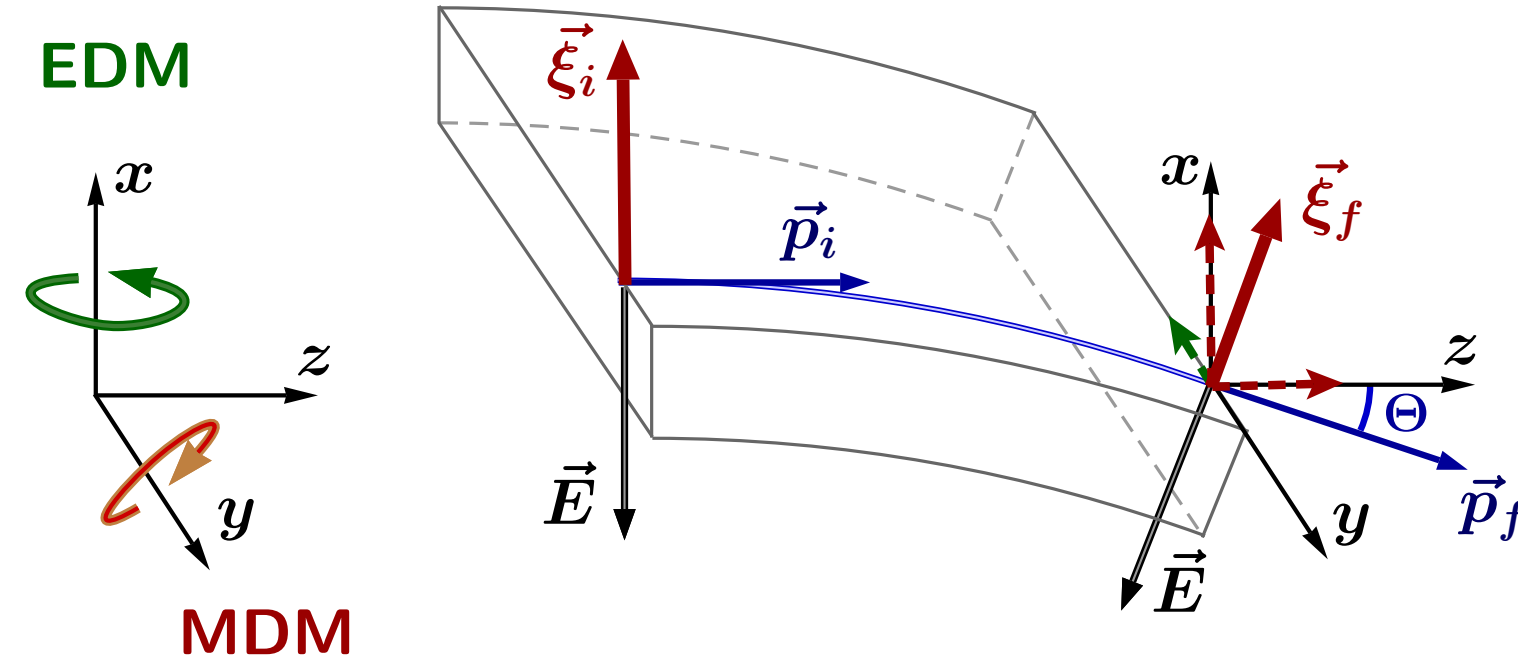
V.L. Lyuboshits,
Sov. J. Nucl. Phys. 31 (1980) 509
[\[inSPIRE\]](#).



$$\Theta_\mu \equiv \angle(\xi_i \xi_f) = (1 + \gamma a) \Theta$$

$$\Delta g = \frac{2}{\alpha \langle \xi_x \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

F. J. Botella et al.,
EPJ C77 (2017) 181 [\[inSPIRE\]](#)

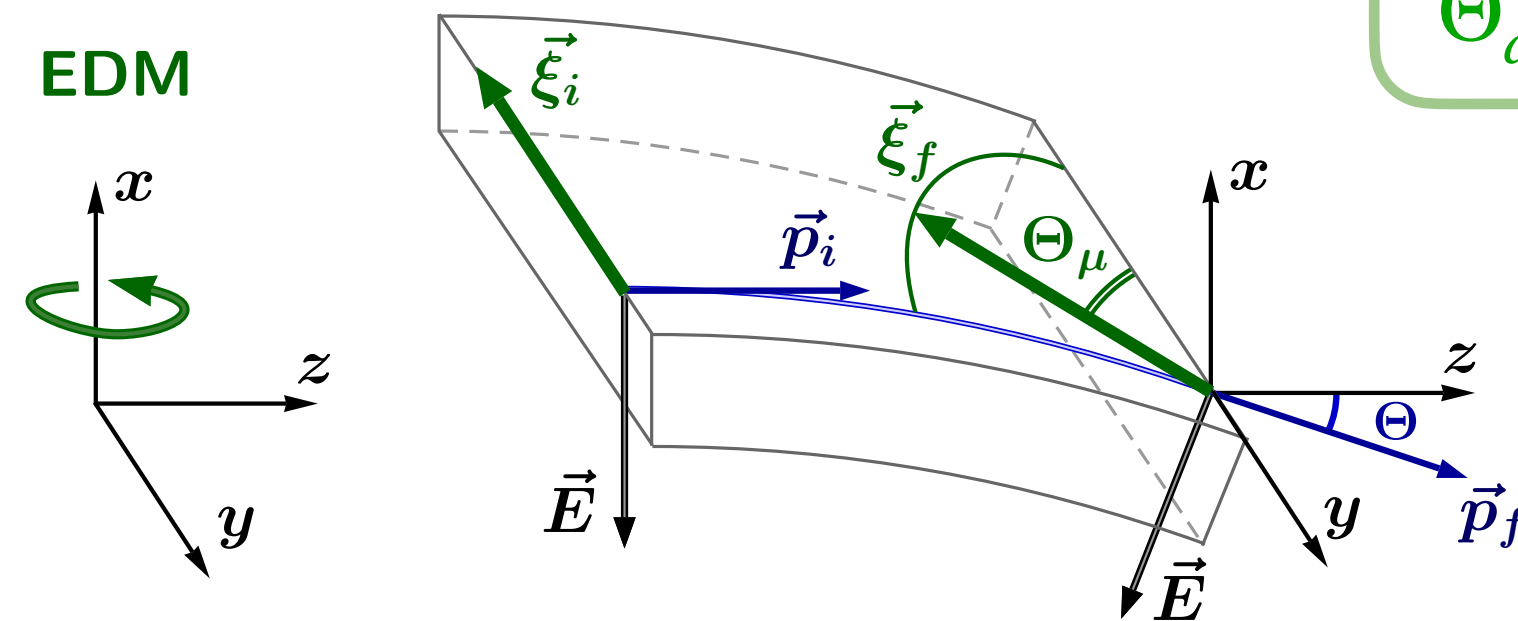


$$\frac{\Delta f}{\Delta g} = \frac{2 \gamma a}{\Theta (1 + \gamma a)^2}$$

effect is suppressed by a small bending angle

V.G. Baryshevsky,
EPJ C79 (2019) 350 [\[inSPIRE\]](#)

A.S. Fomin et al.,
EPJ C80 (2020) 358 [\[inSPIRE\]](#)



$$\Theta_d \equiv \angle(\xi_i \xi_f) = (1 + \gamma f) \Theta$$

$$\Delta f = \frac{2}{\alpha \langle \xi_y \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$



Performance assessment of layouts in IR3 and IR8: precision of measurement

The **gain factor** in PoT from the beam excitation with respect to the “Baseline” (no excitation):

$$\eta_{exc} = \frac{N_{PoT}^{(ex)}}{N_{PoT}^{(BL)}}$$

too high rates of PoT per bunch crossing (~49 @LHCb)

| Beam excitation | N_{coll} | N_{PoT} (IR3) | N_{PoT} (IR8) | η_{exc} | 200 kW in coll. system | | 3.5 p / bunch X @ LHCb | | |
|---|----------------------|----------------------|----------------------|--------------|-------------------------|----------------|------------------------|----------------|----------------|
| | | | | | <i>after collisions</i> | | after collis. | during collis. | |
| | | | | | t, h | $\Delta L / L$ | t, h | $\Delta L / L$ | $\Delta L / L$ |
| No excitation (“Baseline”) | 1.7×10^{15} | 4.7×10^{12} | 6.8×10^{12} | 1 | 1908 | 0 | 1908 | 0 | 0 |
| All fills (3–2556b) (no limit on Δt) | 4.7×10^{15} | 2.6×10^{13} | 3.8×10^{13} | 5.6 | 8 | > 0.4 % | 167 | 8.7 % | 2.9 % |
| Selected fills (2556b), $\Delta t < 1$ hour | 2.5×10^{15} | 1.4×10^{13} | 2.0×10^{13} | 3.0 | - | - | 54 | 2.8 % | 1.2 % |

The most efficient scenario:

The gradual excitation of bunches (3.5 p/bunch X) at the end of the fills with 2556 bunches, with duration under 1 hour:

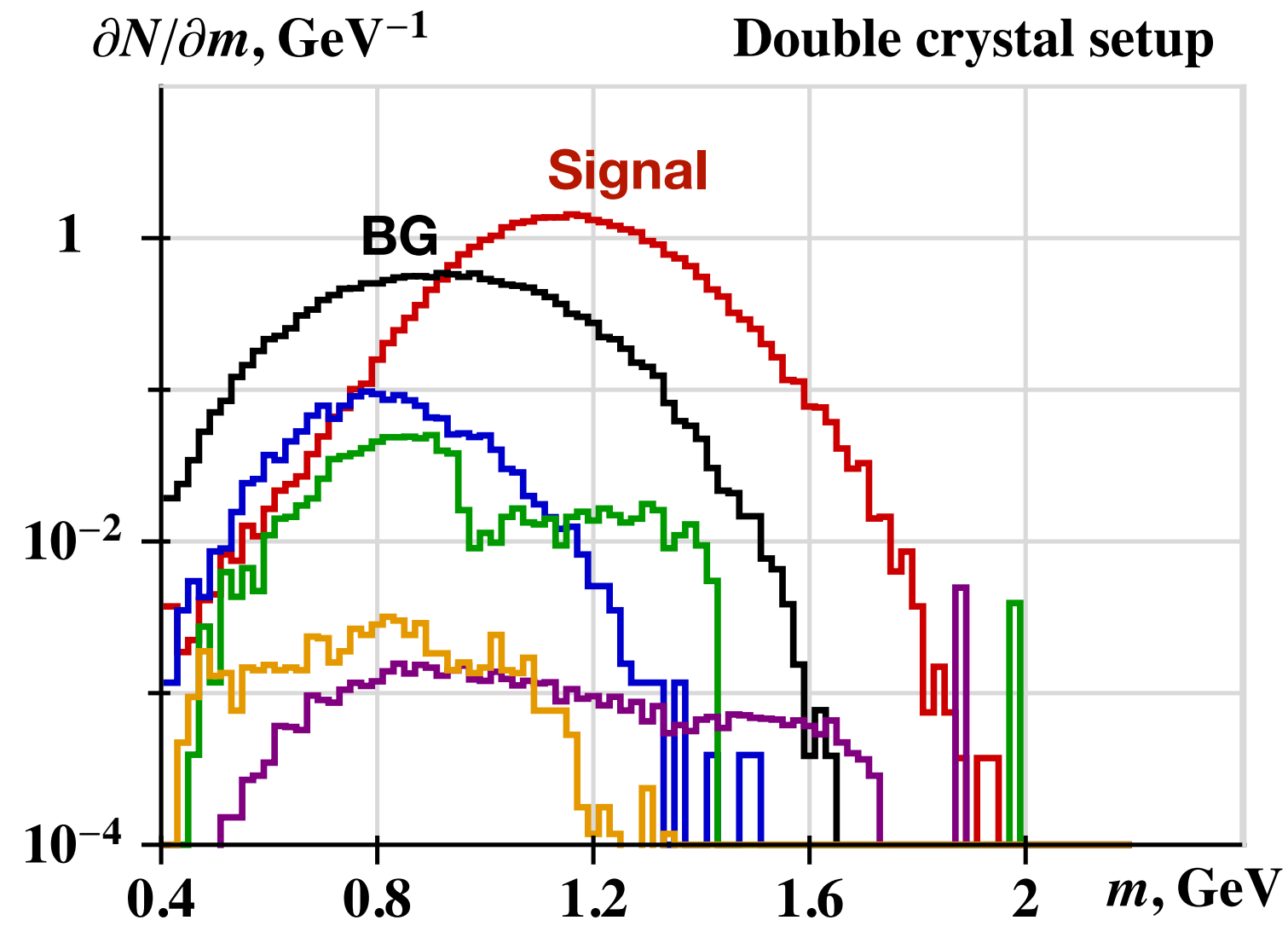
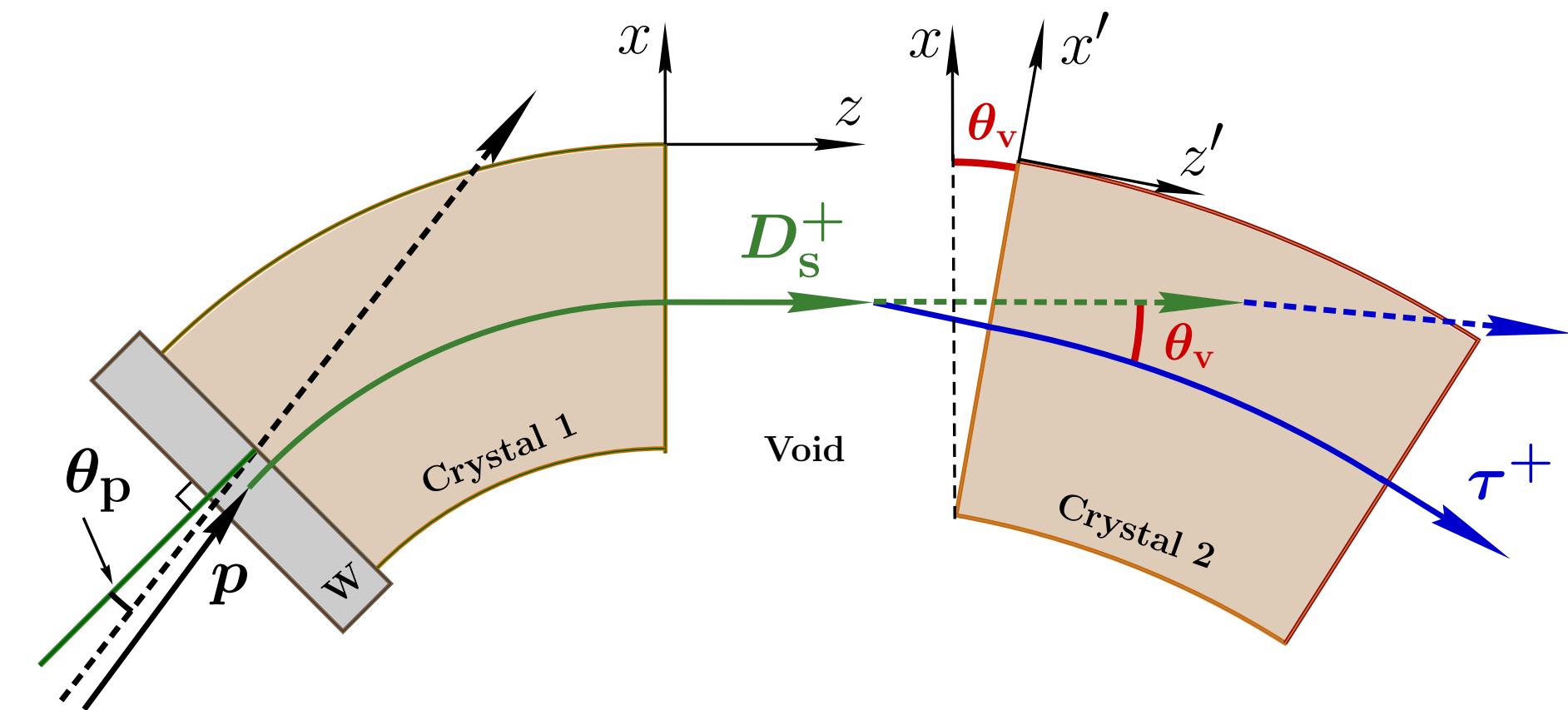
⇒ delivers **3 times more PoT** w.r.t. “Baseline”

⇒ reduces the total luminosity at ATLAS and CMS by **~1.2%**



MDM of τ -lepton. Background from other decay channels

A.S. Fomin et al. JHEP 1903 (2019) 156



- $D_s^+ \rightarrow \tau^+ \nu_\tau, \tau^+ \rightarrow \pi^+ \pi^+ \pi^- \bar{\nu}_\tau$
- $D_s^+ \rightarrow \tau^+ \nu_\tau, \tau^+ \rightarrow \pi^+ \pi^+ \pi^- \pi^0 \bar{\nu}_\tau$
- $D_s^+ \rightarrow \tau^+ \nu_\tau, \tau^+ \rightarrow \pi^+ \pi^+ \pi^- \pi^0 \pi^0 \bar{\nu}_\tau$
- $D_s^+, D^+, \Lambda_c^+ \rightarrow \pi^+ \pi^+ \pi^- X^0$
- $B^+ \rightarrow D_s^+ X, D_s^+ \rightarrow \pi^+ \pi^+ \pi^- X^0$
- $B^+ \rightarrow D^+ X, D^+ \rightarrow \pi^+ \pi^+ \pi^- X^0$
- $B^+ \rightarrow \Lambda_c^+ X, \Lambda_c^+ \rightarrow \pi^+ \pi^+ \pi^- X^0$

