Unexpected large isospin violation in high-energy collisions

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+NA61/SHINE feedback

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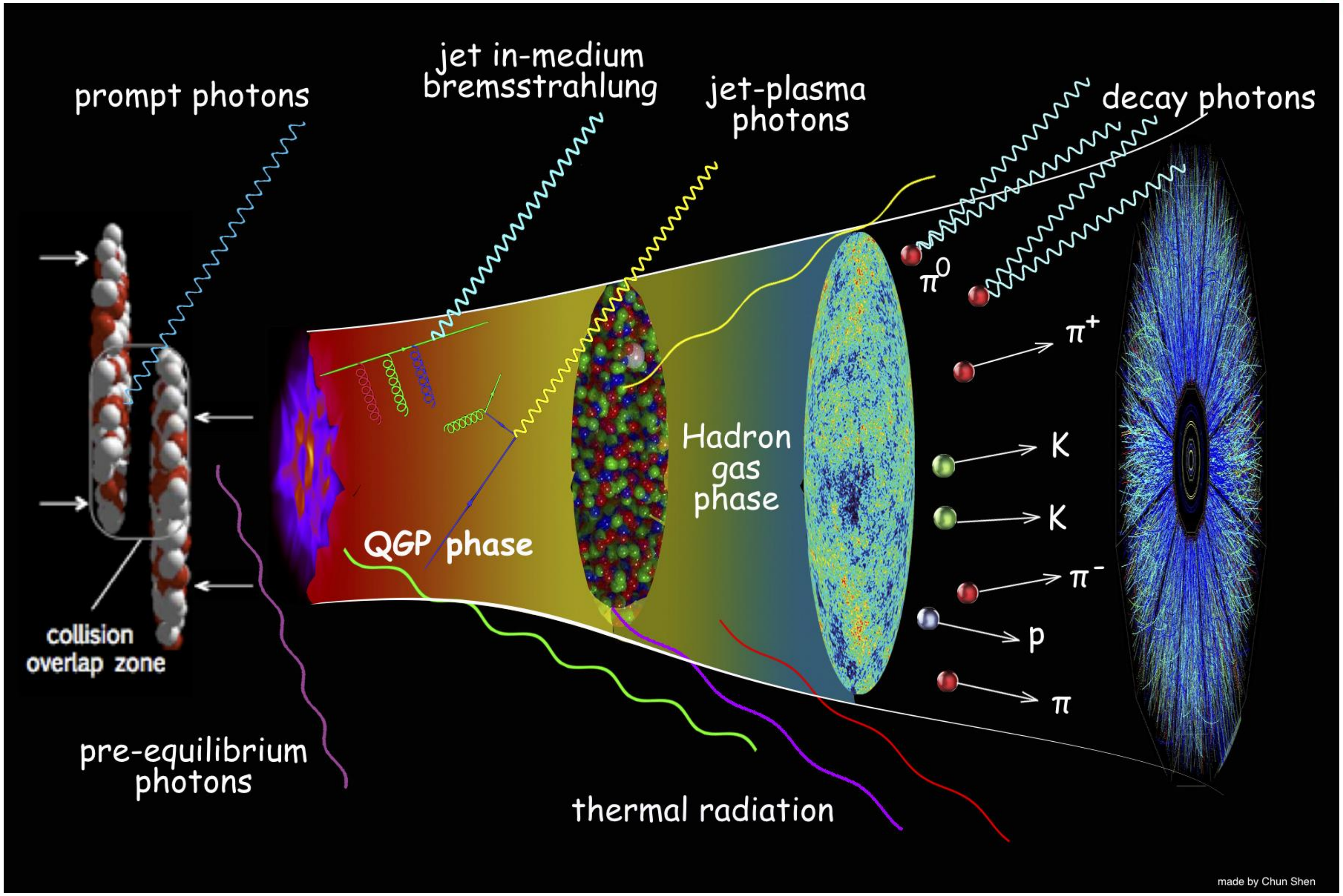
Present, Past, Future
28-29/5/2024
Kiev (+ zoom) Ukraine

Outline

1. Heavy-Ion collisions: brief recall
2. Isospin: brief recall
3. Kaon productions
4. Theory vs experiment
5. Conclusions

Heavy-ion collisions



C. Shen, U. Heinz, Nucl. Phys. News 25 (2015) 2, 6-11

At the freeze-out, the emission of hadrons is well described by e.g. thermal models.

- Here, we concentrate on kaon production, especially on an unexpected large violation of isospin in charged to neutral kaon ratio

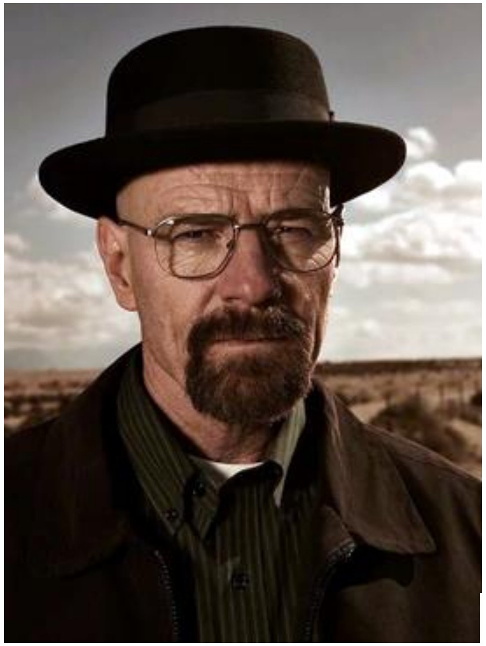
- Brylinski et al., Large isospin symmetry breaking in kaon production at high energies," [arXiv:2312.07176 [nucl-th]].

- Adhikary et al. [NA61/SHINE], Excess of Charged Over Neutral K Meson Production in High-Energy Collisions of Atomic Nuclei, [arXiv:2312.06572 [nucl-ex]]

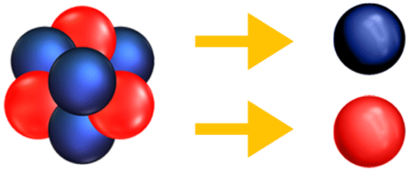
- …as well as to a compilation of other experiments

Heisenberg (1932): the nucleon

Uniwersytet
Jana Kochanowskiego w Kielcach

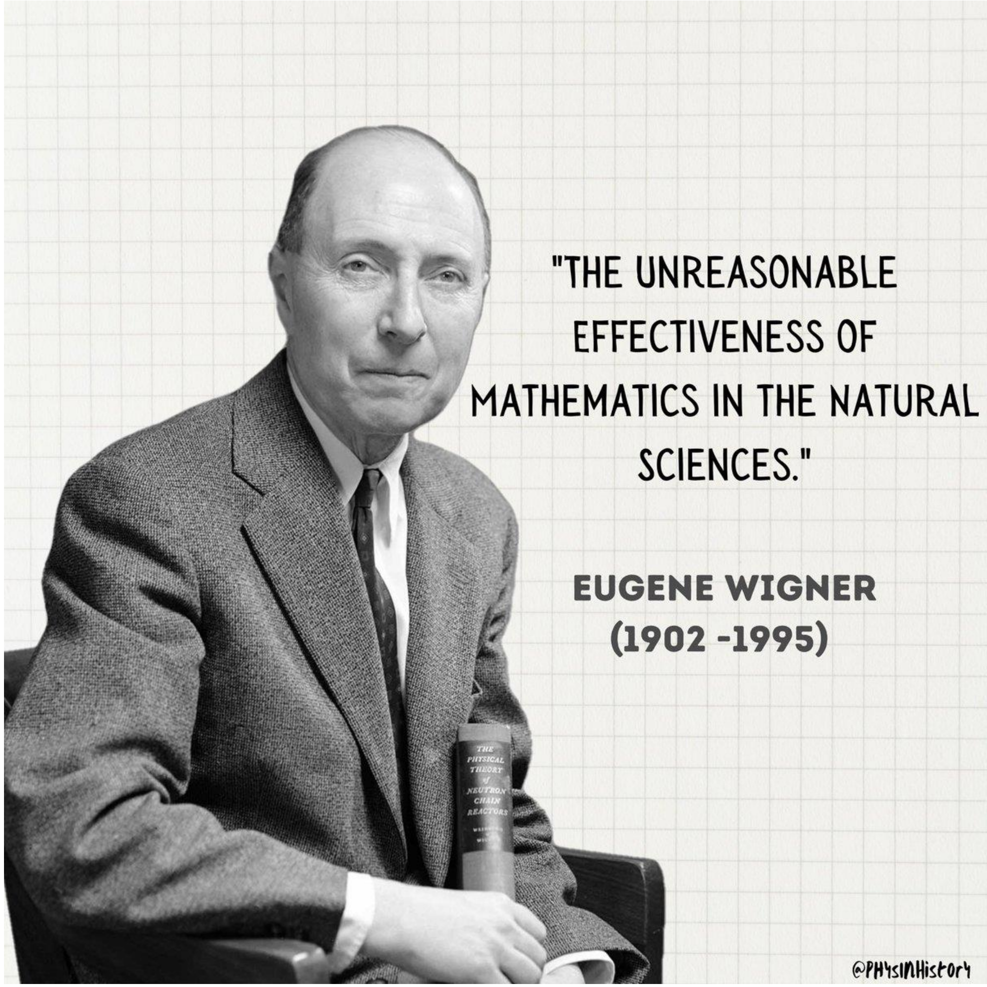


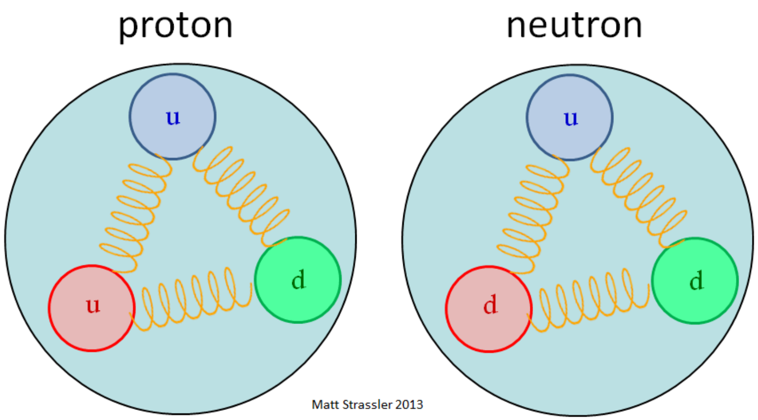
A nucleon is either a proton or a neutron as a component of an atomic nucleus

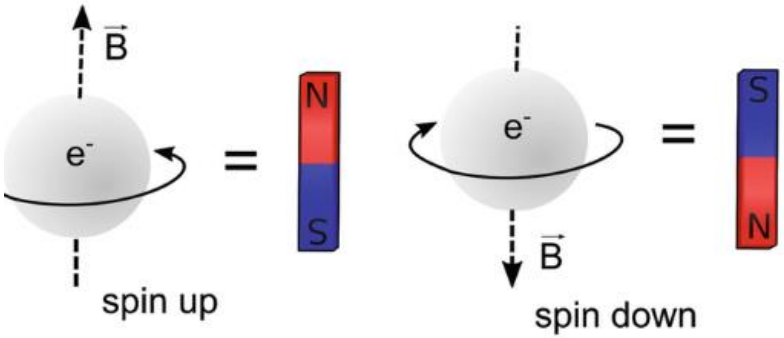


Proton and neutron merge into the nucleon
Masses very similar.

Wigner (1932): isotopic spin, thus isospin







Nucleon doublet: I=1/2

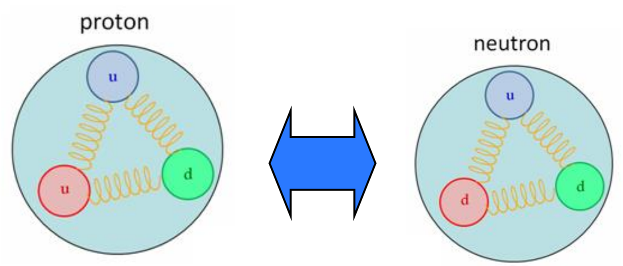
$$\begin{pmatrix} p \\ n \end{pmatrix} \to \hat{O} \begin{pmatrix} p \\ n \end{pmatrix}$$

$\hat{O}$ is a 2×2 unitary matrix. $\hat{O} = e^{i\theta_i \sigma_i / 2}$

A specific isospin transformation is the so-called charge transformation:

$$\hat{C} = e^{i\pi\sigma_2/2} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

Then under $\hat{C}$: $p \Longleftrightarrow n$



Kaons form isospin doublets, just as the nucleon

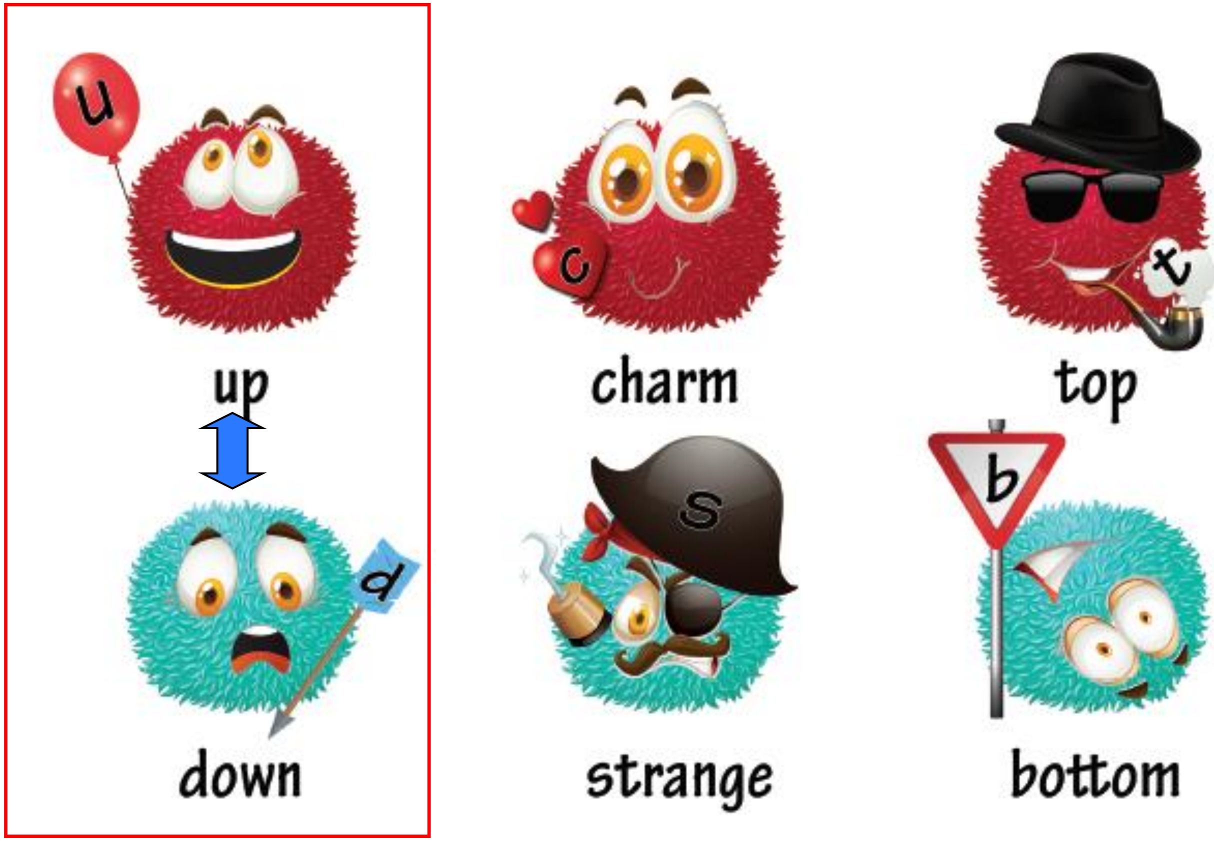
$$\begin{pmatrix} p \\ n \end{pmatrix} \begin{pmatrix} K^+ \\ K^0 \end{pmatrix} \begin{pmatrix} -\bar{K}^0 \\ K^- \end{pmatrix} \cdots$$

under $\hat{C}$:

$$\begin{aligned} p &\Longleftrightarrow n \\ K^+ &\Longleftrightarrow K^0 \\ \bar{K}^0 &\Longleftrightarrow K^- \end{aligned}$$

Quarks and QCD

Quarks and QCD, isospin:



In terms of quarks: $\begin{pmatrix} u \\ d \end{pmatrix} \to \hat{O} \begin{pmatrix} u \\ d \end{pmatrix}$

Then under $\hat{C}$: $u \Longleftrightarrow d$

Isospin is an approximate symmetry of QCD

- Mesonic multiplets (nucleon doublet, pion triplet, kaon doublets).

- Reactions: if an initial state has a certain (I, I_z), then the final state is also such. Indeed, pion-pion, pion-nucleon and nucleon-nulceon scattering conserve isospin (to a good level of accuracy).

Example: ($I=I_z=1$)

$$p + p \to \Lambda + K^+ + p$$

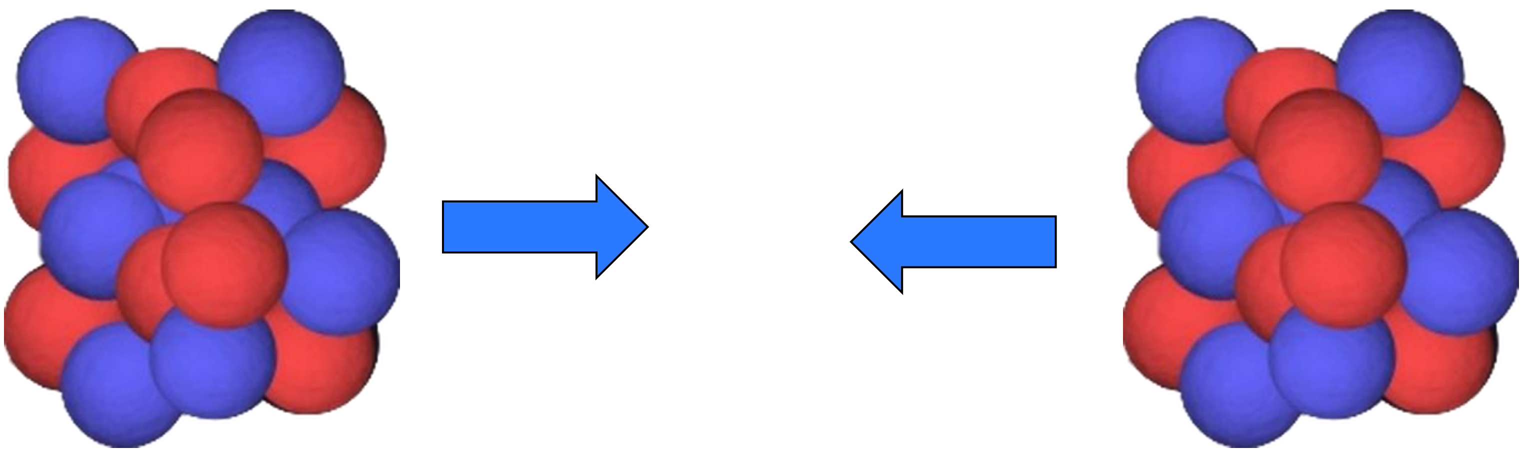
- Isospin symmetry is good, but not exact. Masses of u and d not equal (explicit symmetry breaking).

- Isopsin transformations are a subset of flavor transformations.

Nucleus-nucleus collion with equal numbers of protons and neutrons

$$Z = N = A/2 \qquad Q/B = 1/2$$

$$|A + A\rangle$$



Oxygen-16 **Oxygen-16**

$I_z = 0$ (typically also I =0 for each nucleus, thus total isospin also vanishing)

Toward the general initial state

- For total initial I = 0 it is easy to show that $\langle K^+ \rangle = \langle K^0 \rangle$
- The result can be easily extended to any **fixed** total initial isospin $I=I_0$.
- It can be even generalized to initial states that are not isospin eigenstates, provided that an appropriate average is performed.

Expected kaon multiplicities

Charge symmetry means that strong interactions are invariant under the inversion of the third component of the isospin of hadron of the initial and final states.

Let us consider an ensemble of initial states being invariant under the charge transformation - probabilities of having initial states related by this transformation are equal. This is the case of nucleus-nucleus collisions where each nucleus has an equal number of protons and neutrons (thus, $I_z = 0$). Then, the invariance under C-transformation holds also for the final state ensemble:

$$\langle K^+ \rangle = \langle K^0 \rangle$$

$$\langle K^- \rangle = \langle \bar{K}^0 \rangle$$

Neutral kaons and the ratio R_K

$$\begin{pmatrix} |K_S^0\rangle \\ |K_L^0\rangle \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} |K^0\rangle \\ |\bar{K}^0\rangle \end{pmatrix}$$

$$\langle K_S^0 \rangle = \frac{1}{2}\langle K^0 \rangle + \frac{1}{2}\langle \bar{K}^0 \rangle = \langle K_L^0 \rangle$$

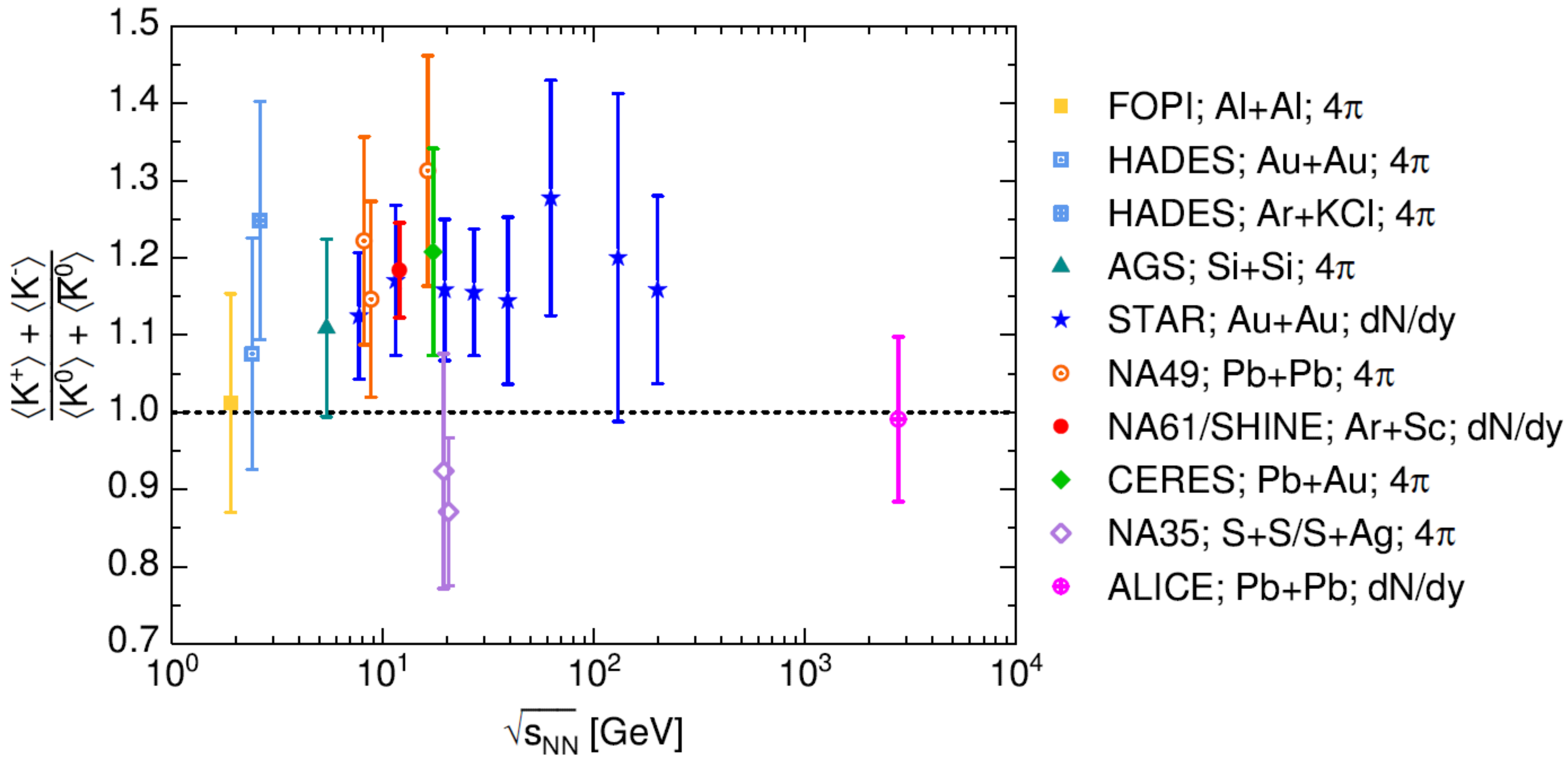
$$\langle K^+ \rangle + \langle K^- \rangle = 2\langle K_S^0 \rangle$$

$Q/B = 1/2$

+ isospin exact…

$$R_K \equiv \frac{\langle K^+ \rangle + \langle K^- \rangle}{\langle K^0 \rangle + \langle \bar{K}^0 \rangle} = \frac{\langle K^+ \rangle + \langle K^- \rangle}{2\langle K_S^0 \rangle} = 1$$

Experimental results (NA61/SHINE plus others)



Latest NA61/SHINE result: $R_K = 1.184 \pm 0.061$

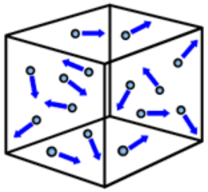
Note, however, most experiments have Q/B < 0.5

Theoretical approaches

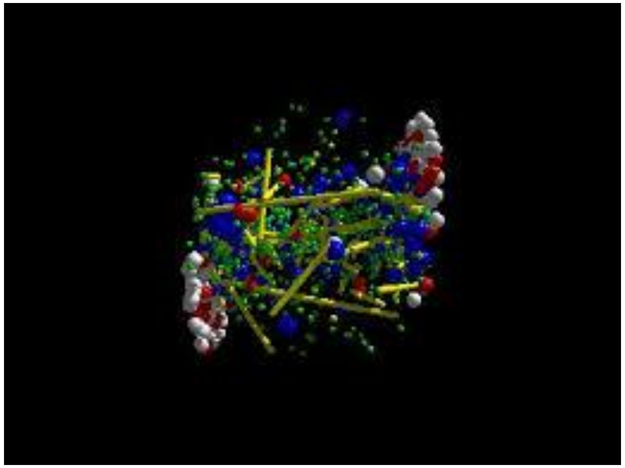
- HRG (hadron resonance gas approach)

$$\ln Z = \sum_k \ln Z_k^{\mathrm{stable}} + \sum_k \ln Z_k^{\mathrm{res}}$$

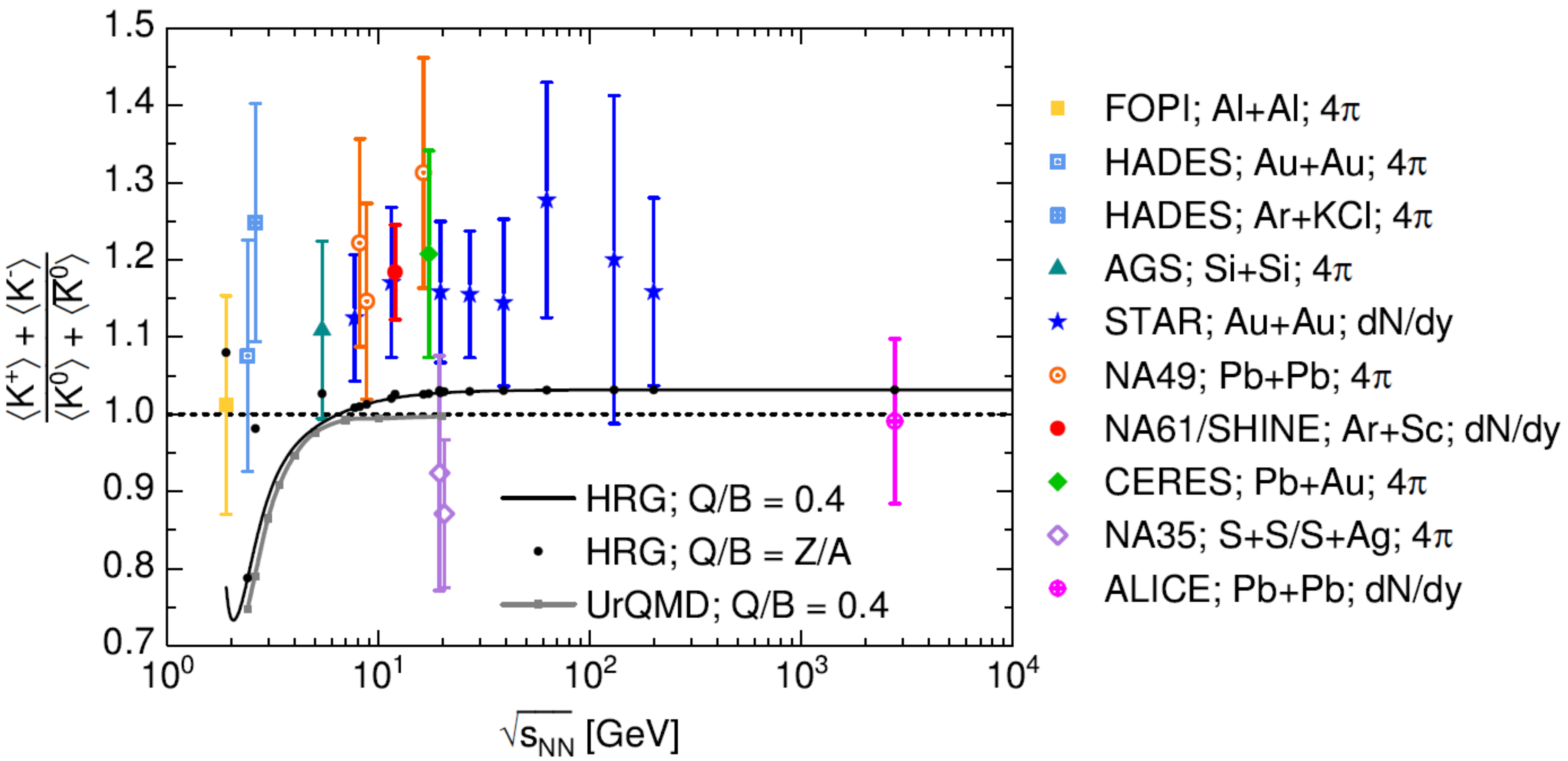
$$\ln Z_k^{\mathrm{stable,}} \quad = f_k V \int \frac{d^3p}{(2\pi)^3} \ln \left[1 \pm e^{-E_p/T}\right]^{\pm 1}$$



- UrQMD (Hadron-String transport model, fully integrated Monte Carlo simulation of nucleus-nucleus simulations)



Exp vs theory (HRG+UrQMD)



Almost all experimental dots are above the corresponding theoretical ones

Considerations

- HRG and UrQMD agree with each other
- Q/B <1/2 actually favors neutral kaons
- charged kaons are lighter than neutral ones: this favors charged kaons

Considerations/2

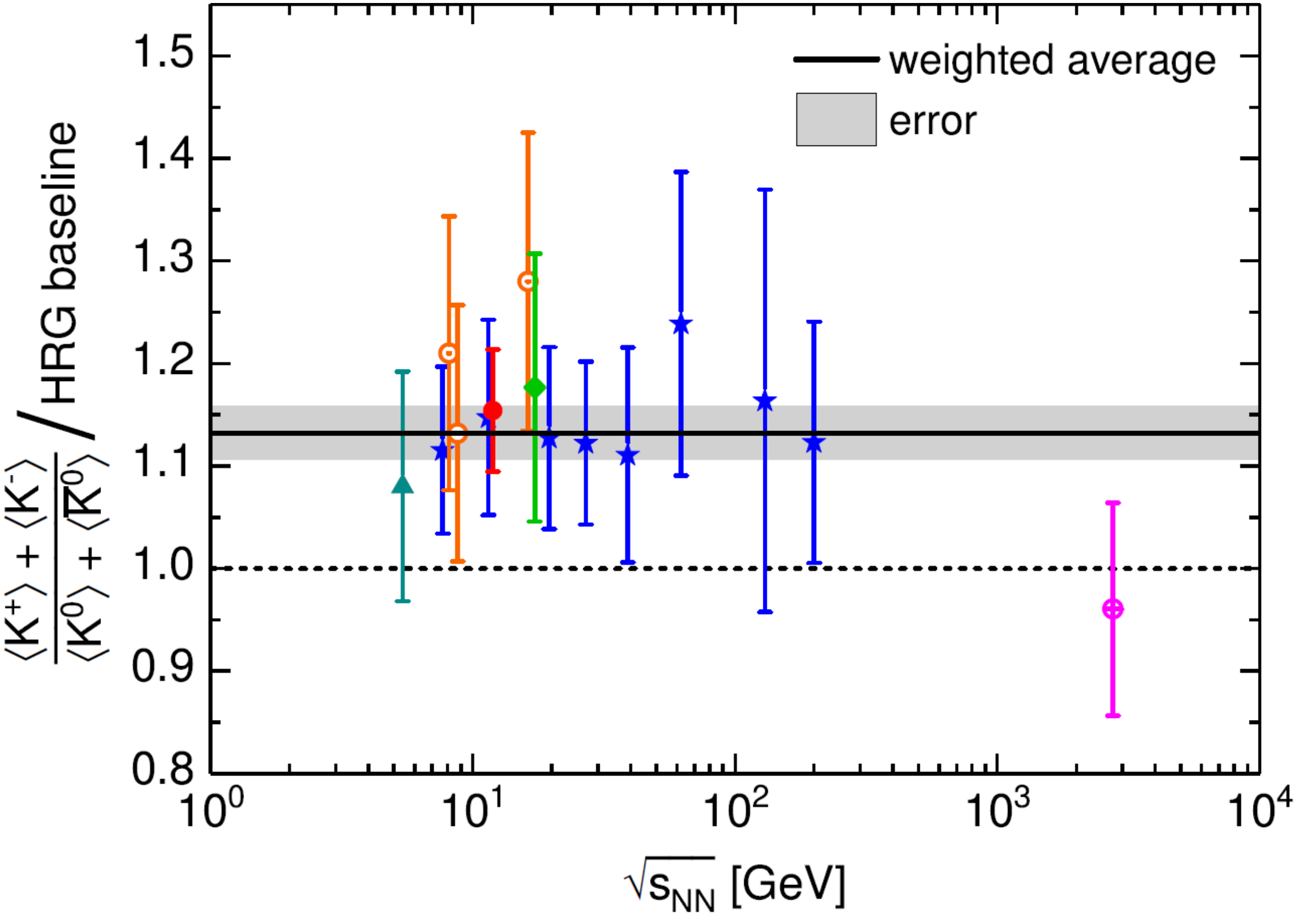
- Non-QCD effects: weak processes are negligible

- Non-QCD effects: electromagnetic processes are small, of the order of α^2

- Decays of ϕ(1020) meson as well as other asymmetries generate quite small effects

Experiment vs theory: ratio

1.132 ± 0.026 $\chi^2_{\min}/\mathrm{dof} = 0.38$



The exp/th missmatch is 5.06σ,
and increases to 8.25σ for the PDG-liked scaled errors.

Summary and conclusions

- Theory (HRG,UrQMD) cannot explain experiment
- Scattering of nuclei with Z=N=A/2 highly desired…
- Easier but equally good? Average over: $\pi^- + C \text{ and } \pi^+ + C$

NA61/SHINE
PRD 107 (2003) 062004

- Study other isospin multiplets
- Non-perturbative effects? Chiral anomaly (Pisarski&Wilczek,…)

Thanks!

Quarks and QCD, flavor symmetry:

Flavor transformation is a rotation in the (u,d,s) space.
Isospin is a subgroup of flavor.

Example of isospin breaking/1

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

CERN-EP/84-27

March 8th, 1984

THE ISOSPIN-VIOLATING DECAY $\eta' \to 3\pi^0$

IHEP[1]-IISN[2]-LAPP[3] Collaboration

$$BR(\eta' \to 3\pi^0) = 5.2 \left(1 - \frac{m_u}{m_d}\right)^2 10^{-3}$$

Example of isospin breaking/2

$\phi(1020)$

$I^G(J^{PC}) = 0^-(1^{--})$

ϕ(1020) MASS

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
1019.461±0.016 OUR AVERAGE				

ϕ(1020) DECAY MODES

	Mode	Fraction (Γ_i/Γ)	Scale factor/ Confidence level
Γ_1	K^+K^-	(49.1 ±0.5) %	S=1.3
Γ_2	$K^0_L K^0_S$	(33.9 ±0.4) %	S=1.2

Example of isospin breaking/3

Citation: R.L. Workman *et al.* (Particle Data Group), Prog.Theor.Exp.Phys. **2022**, 083C01 (2022) and 2023 update

$D^*(2007)^0$

$I(J^P) = \frac{1}{2}(1^-)$
I, J, P need confirmation.

J consistent with 1, value 0 ruled out (NGUYEN 77).

$D^*(2007)^0$ DECAY MODES

$\overline{D}^*(2007)^0$ modes are charge conjugates of modes below.

	Mode	Fraction (Γ_i/Γ)
Γ_1	$D^0\pi^0$	(64.7 ± 0.9) %
Γ_2	$D^0\gamma$	(35.3 ± 0.9) %
Γ_3	$D^0 e^+ e^-$	$(3.91 \pm 0.33) \times 10^{-3}$

Citation: R.L. Workman *et al.* (Particle Data Group), Prog.Theor.Exp.Phys. **2022**, 083C01 (2022) and 2023 update

$D^*(2010)^\pm$

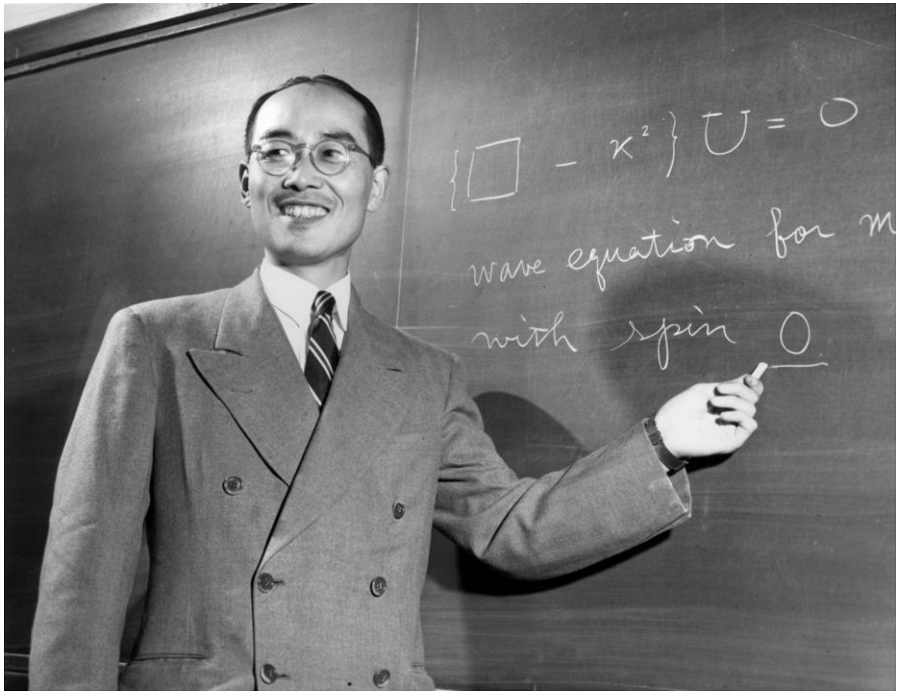
$I(J^P) = \frac{1}{2}(1^-)$
I, J, P need confirmation.

$D^*(2010)^\pm$ DECAY MODES

$D^*(2010)^-$ modes are charge conjugates of the modes below.

	Mode	Fraction (Γ_i/Γ)
Γ_1	$D^0\pi^+$	(67.7 ± 0.5) %
Γ_2	$D^+\pi^0$	(30.7 ± 0.5) %
Γ_3	$D^+\gamma$	(1.6 ± 0.4) %

Yukawa (1932) and Kemmer (1939): isospin triplet I=1



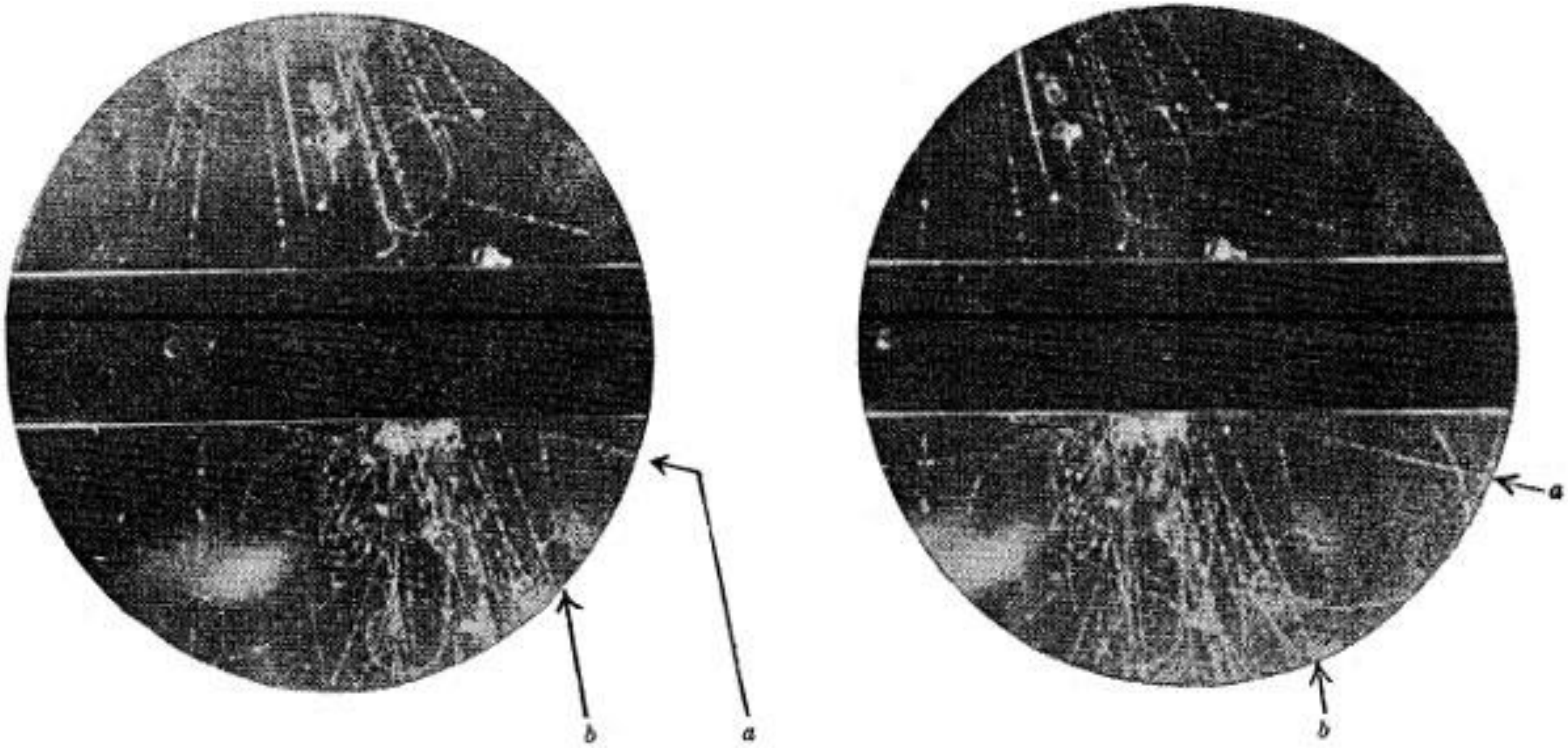
$$\begin{pmatrix} \pi^+ \\ \pi^0 \\ \pi^- \end{pmatrix}$$

under $\hat{C}$: $\pi^+ \Longleftrightarrow \pi^-$

Kaons

20 DECEMBER 1947

Clifford Butler and George Rochester discover the kaon; first strange particle



Historical recall: „Shmushkevich” rule

An initial ‘uniform’ ensemble of hadronic state (that is, one with an equal mean number of each member of any isospin multiplet, such as the scattering of two isosinglet nuclei) evolves into a uniform final-state ensemble.

Uniform stays uniform

Shmushkevich, I.: . Dokl. Akad. Nauk SSSR **103**, 235 (1955)

Dushin, N., Shmushkevich, I.: . Dokl. Akad. Nauk SSSR **106**, 801 (1956)

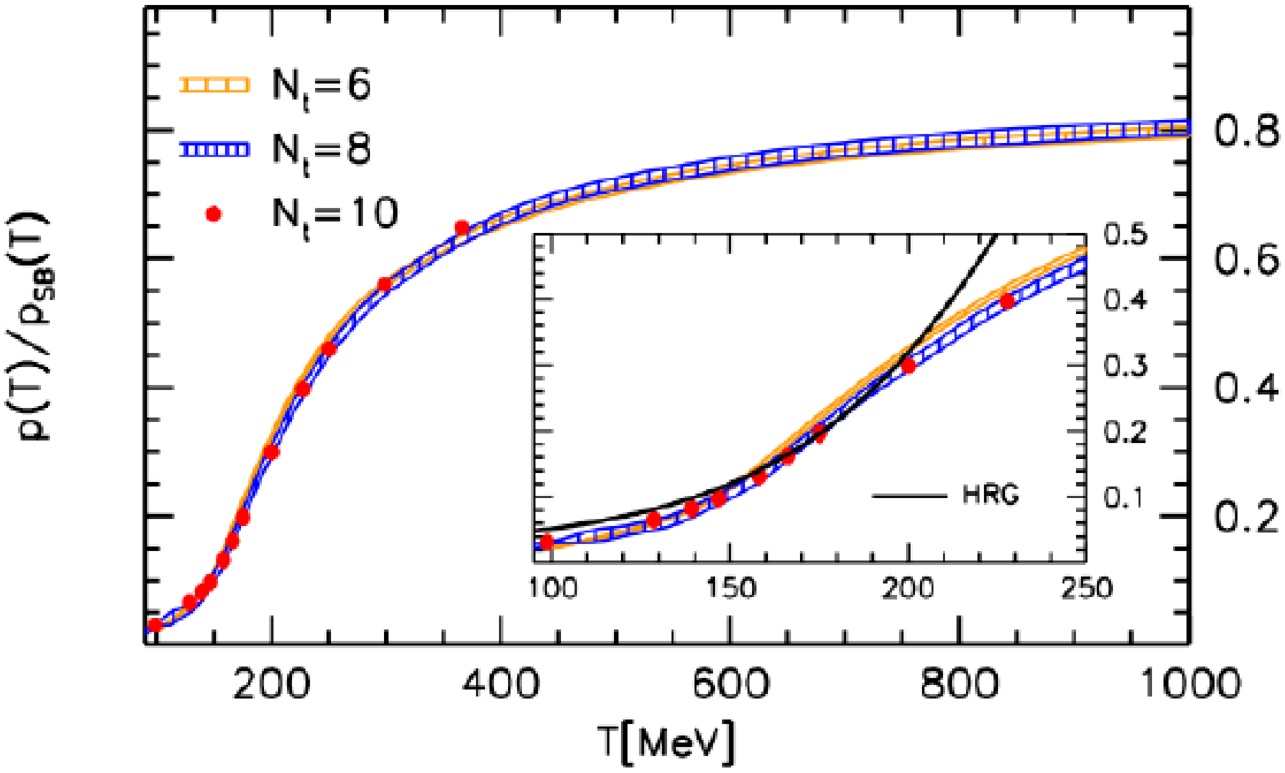
MacFarlane, A.J., Pinski, G., Sudarshan, G.: Shmushkevich's method for a charge independent theory. Phys. Rev. **140**, 1045 (1965) https://doi.org/10.1103/PhysRev.140.B1045

Wohl, C.G.: Isospin relations by counting. American Journal of Physics **50**(8), 748–753 (1982) https://doi.org/10.1119/1.12743

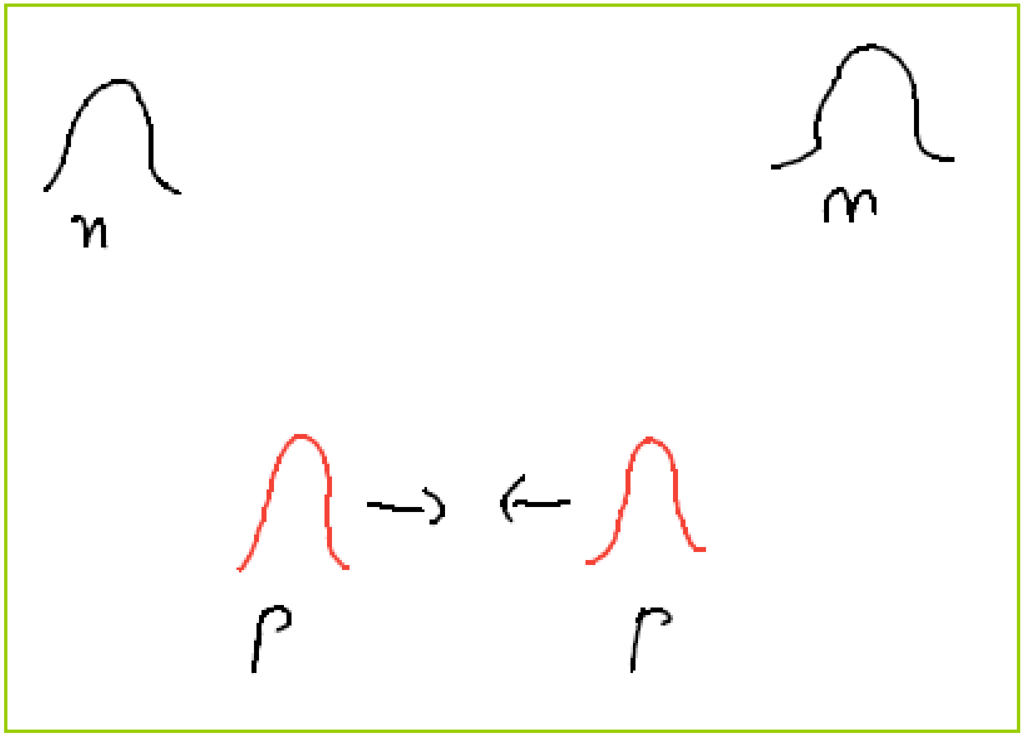
Pal, P.: An Introductory Course of Particle Physics -CRC Press, (2014)

Hadron resonance gas vs lattice results

- All baryons and mesons (m < 2.5 GeV) from PDG [Borsnayi et al.JHEP11(2010)077]



ppnn ⟼ ?



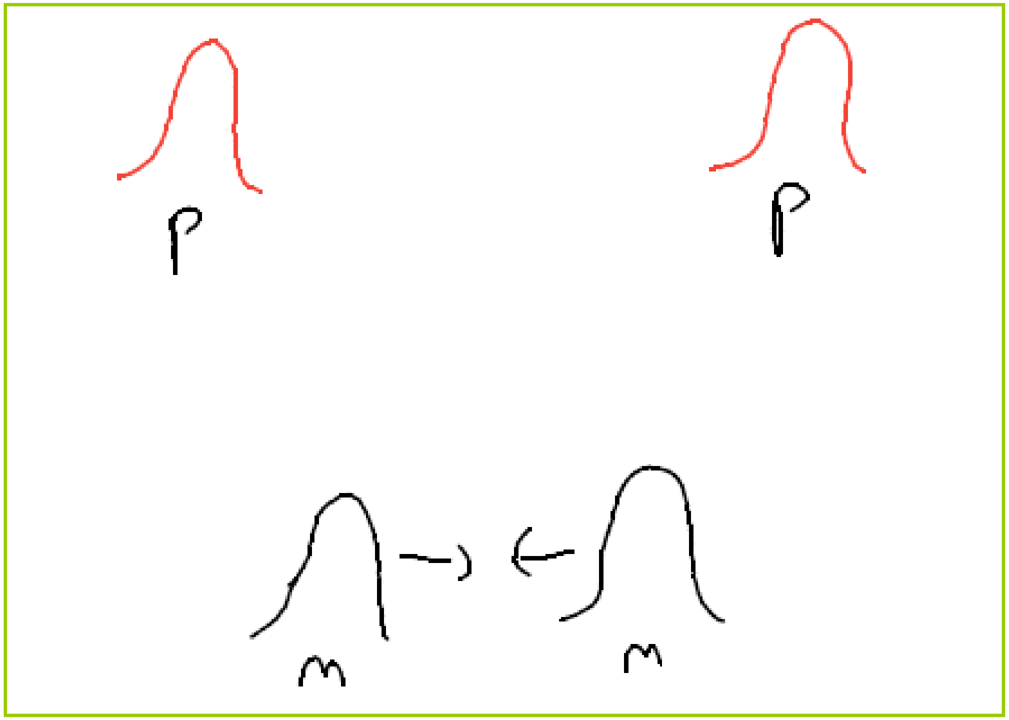
Just as pp!

More K^+ than K^0

Is then the previous argumentation wrong?

No.
One needs to average.

But ... $\hat{C}$ transform



This is the C-transformed version fo the previous reaction.

Here, the protons are spectactors and the neutrons interact.

Just as nn scattering!

More K^0 than K^+

Averaging leads to…

If both initial states are equally probable

$$\langle K^+ \rangle = \langle K^0 \rangle$$

holds!

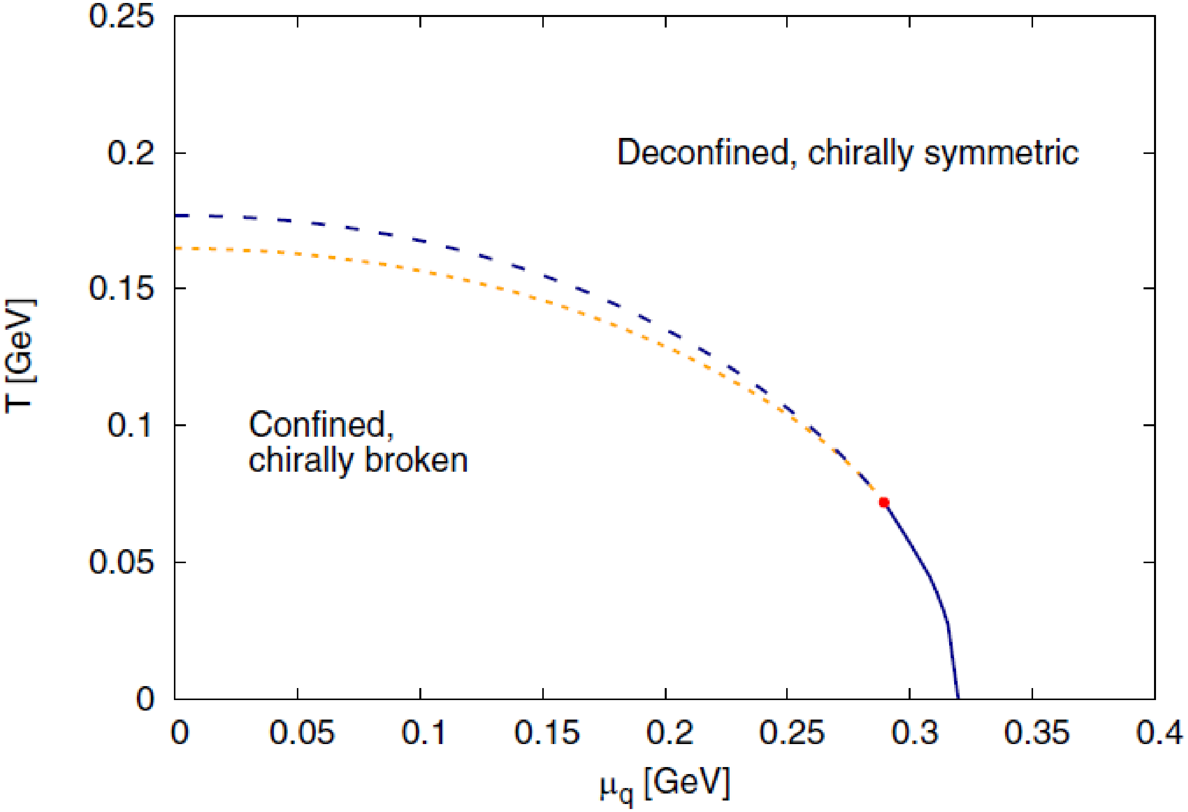
Formally:

$$\hat{\rho} = \sum_n p_n \left|\Psi_n\right\rangle \left\langle\Psi_n\right|$$

$$\hat{C}\hat{\rho}\hat{C}^\dagger = \hat{\rho}$$

This is a general result!

Phase diagram of the eLSM: Nc =3



Details in 2209.09568

Schematic phase diagram at large Nc

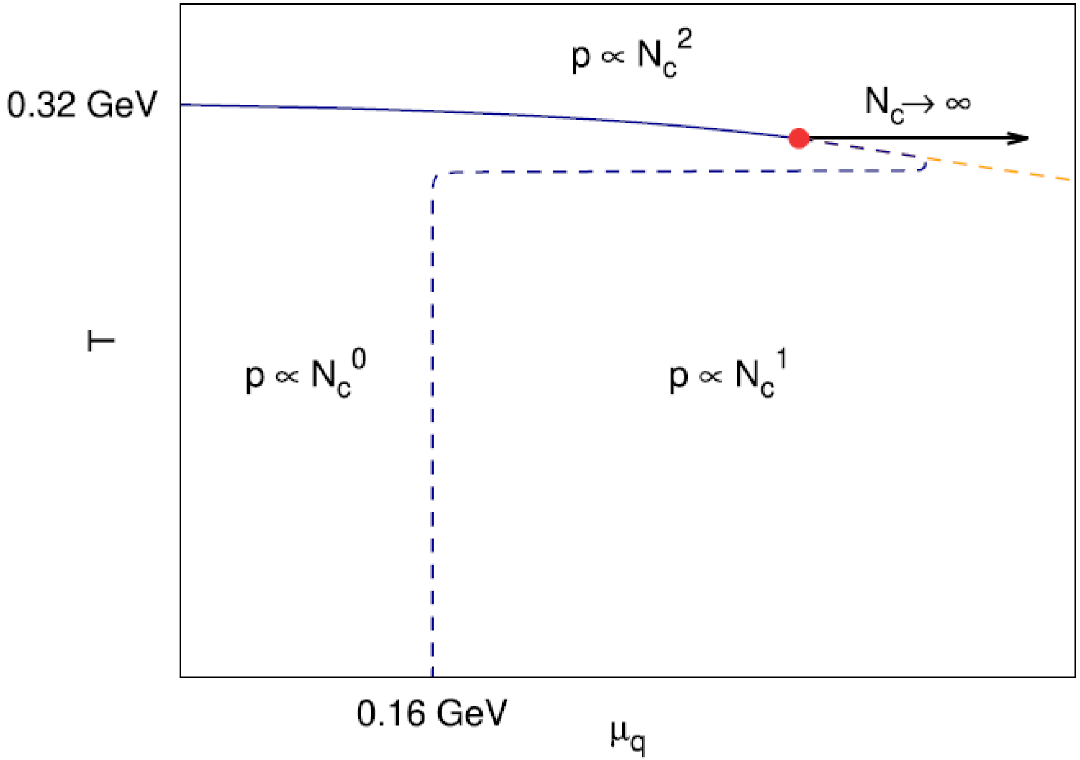


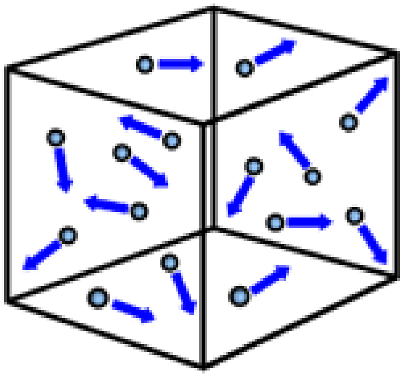
FIG. 13. The schematic phase diagram for large N_c and the N_c scaling of the pressure in the different phases.

Details in 2209.09568.

Then, for the QCD diagram: 3 is not a large number!!!!

Theoretical description of a thermal gas

$$\ln Z = \sum_k \ln Z_k^{\text{stable}} + \sum_k \ln Z_k^{\text{res}}$$



$$\ln Z_k^{\text{stable,}} \quad = f_k V \int \frac{d^3p}{(2\pi)^3} \ln\left[1 \pm e^{-E_p/T}\right]^{\pm 1}$$

$$E_p = \sqrt{\vec{p}^2 + M_k^2} \qquad\qquad P = \frac{T}{V} \ln Z$$

The stable part is “easy”: In first approximation, resonance as stable.
Then, correcting for width, itneraction,…