

ENERGY CORRELATION OF BOTTOM QUARKS FROM DECAYS OF TOP QUARKS IN ELECTRON-POSITRON ANNIHILATION AT HIGH ENERGY

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Motivation

- The study of top-quark properties is a significant area of interest in particle physics, particularly in our quest to uncover effects that go beyond the Standard Model (SM). One of the key properties we focus on is polarization. Polarization allows us to gain more detailed insights into the top quark and potential influences of what we call new physics.
- It is believed that more precise measurements of top-quark properties at e^+e^- colliders will be sensitive to physics beyond the SM.
- In this context our work is focused on the future electron-positron collider, The Compact Linear Collider (CLIC). During its first cycle of experiments, CLIC is expected to achieve a center-of-mass energy of 380 GeV.

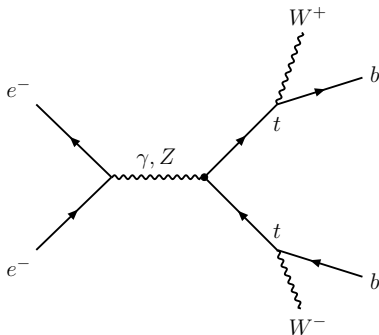


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Let us consider the process of electron-positron annihilation into a pair of quarks

$$e^+ e^- \rightarrow t \bar{t} \rightarrow b \bar{b} W^+ W^-$$

At the tree level, this process is described by two Feynman diagrams involving the exchange of a photon and Z boson:



The vertex of the interaction between γtt and Ztt

We assume that the Lagrangian describing the interaction of the photon and Z boson with quarks consists of Standard model (SM) contributions and terms beyond the SM (BSM).

The corresponding γtt and Ztt vertices are written as

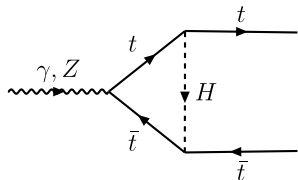
$$\Gamma_{t\bar{t}V}^{\mu} = -ie\bar{t} \left(\underbrace{\gamma^{\mu}(v_t^V - a_t^V \gamma_5)}_{SM} + \underbrace{\frac{1}{2m_t} \sigma^{\mu\nu} q_{\nu} (\overbrace{i\kappa^V(s)}^{CP \text{ even}} - \overbrace{\tilde{\kappa}^V(s) \gamma_5}^{CP \text{ odd}})}_{BSM} \right) t$$

where $V = \gamma, Z$.

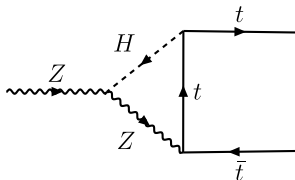
The terms proportional to the coupling κ^V are CP -even, while the terms $\tilde{\kappa}^V$ are responsible for the CP violation.



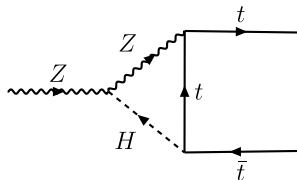
One-loop corrections to the vertices γtt and Ztt



(a)



(b)



(c)

To generate CP violation effects, the interaction of the Higgs boson with the top quarks is chosen in the form

$$\mathcal{L}_{htt} = -\frac{m_t}{v} h \bar{t} (\alpha + i \beta \gamma_5) t$$

which includes scalar and pseudoscalar parts. The values $\alpha = 1$, $\beta = 0$ correspond to the Lagrangian in the SM.



Using the Cutkosky rules (cut of the Feynman diagram) for the diagram with vertex γtt (“a”)

$$\text{Im } \tilde{\kappa}(s) = \alpha\beta \frac{Q_t m_t^4}{8\pi E_t p_t v^2} \left(1 - \frac{m_h^2}{4p_t^2} \log \frac{m_h^2 + 4p_t^2}{m_h^2} \right) \theta(s - 4m_t^2).$$

For vertex Ztt (“a”) similarly

$$\text{Im } \tilde{\kappa}_Z(s)_a = \frac{v_t}{Q_t} \text{Im } \tilde{\kappa}(s).$$

Knowing the expressions for the imaginary parts of $\tilde{\kappa}(s)$ and $\tilde{\kappa}_Z(s)$ the real parts of the form factors can be calculated using the dispersion relations:

$$\begin{aligned} \text{Re } \tilde{\kappa}(s) &= \frac{1}{\pi} \text{PV} \int_{s_0}^{\infty} \frac{\text{Im } \tilde{\kappa}(s')}{s' - s} ds', \\ \text{Re } \tilde{\kappa}_Z(s) &= \frac{1}{\pi} \text{PV} \int_{s_1}^{\infty} \frac{\text{Im } \tilde{\kappa}_Z(s')}{s' - s} ds' \end{aligned}$$

PV denotes the integral in the sense of the principal value.



Constraints on anomalous Htt couplings are obtained by CMS analysis through the combination of the Hgg results with measurements of the $t\bar{t}H$ and tH processes in the $H \rightarrow 4l$ and $H \rightarrow \gamma\gamma$ channels.

Channels	Coupling	Observed	Expected
tH & $\bar{t}H$ & ggH	α	$1.05^{+0.25}_{-0.20}$	$1.00^{+0.34}_{-0.26}$
	β	$-0.01^{+0.69}_{-0.67}$	$0.00^{+0.71}_{-0.71}$

Table: Summary of constraints on the Htt , Hgg , and HVV coupling parameters in the Higgs basis of SMEFT. CMS Analysis: HIG-19-009



Differential cross section of the process

The evaluation of the cross section of two-step processes is based on the formalism that the b -quark and \bar{b} -quark are produced by decays of top-quarks on the mass surface.

$$d\sigma_{e^+e^- \rightarrow b\bar{b}W^+W^-} = 4 \int \frac{d\sigma_{e^+e^- \rightarrow t\bar{t}}(n^\mu, n'^\mu)}{d\Omega_t} \frac{d\Gamma_{t \rightarrow bW^+}^0}{\Gamma_t} \frac{d\Gamma_{\bar{t} \rightarrow \bar{b}W^-}^0}{\Gamma_t} d\Omega_t,$$
$$\frac{d^2\sigma_{e^+e^- \rightarrow b\bar{b}W^+W^-}}{dE_b dE_{\bar{b}}} = 4 \left(\frac{m_t}{4\pi p_t p_b^0} \right)^2 \int \frac{d\sigma_{e^+e^- \rightarrow t\bar{t}}(n^\mu, n'^\mu)}{d\Omega_t} d\Omega_t d\phi_b d\phi_{\bar{b}}$$

The 4-polarization vectors of the top-quark and antitop-quark n^μ and n'^μ are defined in the following way

$$n^\mu = \alpha_b \left(-\frac{p_t^\mu}{m_t} + \frac{m_t p_b^\mu}{p_t \cdot p_b} \right), \quad n'^\mu = \alpha_{\bar{b}} \left(-\frac{p_t'^\mu}{m_t} + \frac{m_t p_b'^\mu}{p_t' \cdot p_{\bar{b}}} \right),$$

$\alpha_b = -\alpha_{\bar{b}}$ is the asymmetry parameters in the decays of the top-quark.



Energy distribution

To investigate the effects of CP violating, let us analyze the cross section of the combined process.

$$W(E_b, E_{\bar{b}}) = \frac{1}{\sigma_0} \frac{d^2\sigma_{e^+e^- \rightarrow b\bar{b}W^+W^-}}{dE_b dE_{\bar{b}}}$$

where $\sigma_0 \equiv \sigma_{e^+e^- \rightarrow t\bar{t}}$ is the total unpolarized cross section of the process $e^+e^- \rightarrow t\bar{t}$. For convenience, let us re-designate the energies $\varepsilon \equiv E_b$ and $\bar{\varepsilon} \equiv E_{\bar{b}}$.

The overall structure of the energy distribution looks like this

$$W(\varepsilon, \bar{\varepsilon}) = S(\varepsilon, \bar{\varepsilon}) + (\varepsilon - \bar{\varepsilon}) T(s),$$

$$T(s) = [a(s) \operatorname{Im} \tilde{\kappa}(s) + b(s) \operatorname{Im} \tilde{\kappa}_z(s) + c(s) \operatorname{Re} \tilde{\kappa}(s) + d(s) \operatorname{Re} \tilde{\kappa}_z(s)],$$

where $S(\varepsilon, \bar{\varepsilon})$ is the symmetric term of the energies ε and $\bar{\varepsilon}$.



Asymmetry

We introduce a dimensionless quantity, the asymmetry of the energy distribution of b and \bar{b} quarks:

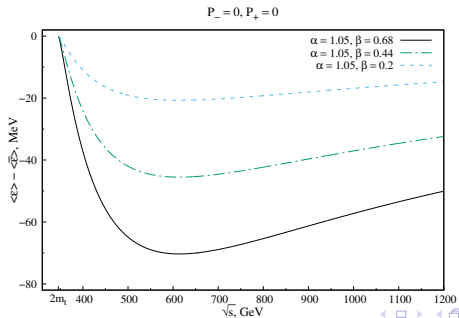
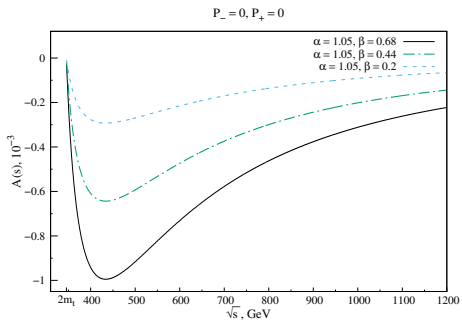
$$A(s) \equiv \int_{E_-}^{E_+} \int_{E_-}^{E_+} W(\varepsilon, \bar{\varepsilon}) [\theta(\varepsilon - \bar{\varepsilon}) - \theta(\bar{\varepsilon} - \varepsilon)] d\varepsilon d\bar{\varepsilon} = \frac{1}{3} (E_+ - E_-)^3 T(s)$$

$A(s)$ corresponds to the experimentally measured quantity in terms of the number of events $(N_{E_b > E_{\bar{b}}} - N_{E_{\bar{b}} > E_b}) / N_{tot}$.

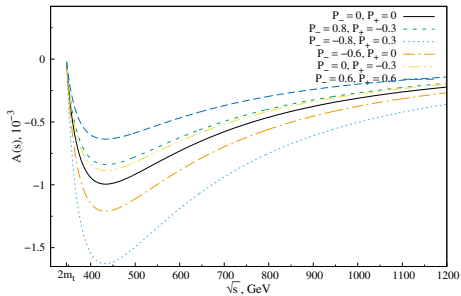
Another observable quantity carrying information about the form factors of $\tilde{\kappa}$ and $\tilde{\kappa}_Z$ is the difference in the mean energies of b and \bar{b} quarks:

$$\langle \varepsilon \rangle - \langle \bar{\varepsilon} \rangle \equiv \int_{E_-}^{E_+} \int_{E_-}^{E_+} W(\varepsilon, \bar{\varepsilon}) (\varepsilon - \bar{\varepsilon}) d\varepsilon d\bar{\varepsilon} = \frac{1}{6} (E_+ - E_-)^4 T(s).$$

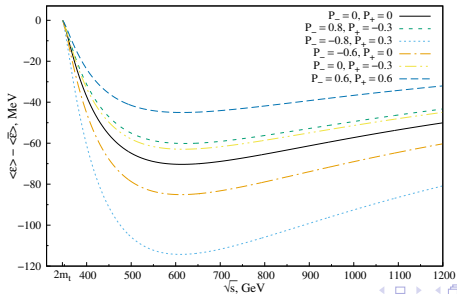




$\alpha = 1.05, \beta = 0.68$



$\alpha = 1.05, \beta = 0.68$



Conclusions

- The process of electron-positron annihilation into a pair of decaying top quarks in the conditions of the future e^+e^- collider is considered.
- Observables corresponding to the effects of CP invariance violation are proposed.
- The dependence of the observables on the β parameter is studied.
- I.V. Truten, A.Yu. Korchin. Energy correlation of bottom quarks from decays of top quarks in electron-positron annihilation. J. Phys. G 49 (2022) 4, 045003. arXiv:2109.10693



Thanks for your attention!

