



Two-way communication between High-Energy Physics and Quantum Technologies

Enrique Rico Ortega Thursday, 28/03/2024

Theory Colloquia

A Two-Way Communication Between High-Energy Physics and Quantum Technologies

by Enrique Rico Ortega

■ Thursday Mar 28, 2024, 11:00 AM → 12:00 PM Europe/Zurich

💡 4/3-006 - TH Conference Room (CERN)





EHU QC EHU Quantum Center

A fruitful dialogue (two-way communication)







Quantum Information Science and Technology



A fruitful dialogue (two-way communication)



Simulating lattice gauge theories within quantum technologies



Collaborators: M. Dalmonte, S. Montangero, U.-J. Wiese, P. Zoller...

Eur. Phys. J. D (2020) 74: 165 https://doi.org/10.1140/epjd/e2020-100571-8

Colloquium

THE EUROPEAN PHYSICAL JOURNAL D

Simulating lattice gauge theories within quantum technologies

Mari Carmen Bañuls^{1,2}, Rainer Blatt^{3,4}, Jacopo Catani^{5,6,7}, Alessio Celi^{3,8}, Juan Ignacio Cirac^{1,2}, Marcello Dalmonte^{9,10}, Leonardo Fallani^{5,6,7}, Karl Jansen¹¹, Maciej Lewenstein^{8,12,13}, Simone Montangero^{14,15,a}, Christine A. Muschik³, Benni Reznik¹⁶, Enrique Rico^{17,18}, Luca Tagliacozzo¹⁹, Karel Van Acoleyen²⁰, Frank Verstraete^{20,21}, Uwe-Jens Wiese²², Matthew Wingate²³, Jakub Zakrzewski^{24,25}, and Peter Zoller³



A fruitful dialogue (two-way communication)



Quantum Simulation for High Energy Physics

C.W. Bauer, Z. Davoudi, A.B. Balantekin, T. Bhattacharya, M. Carena, W.A. de Jong, P. Draper, A. El-Khadra, N. Gemelke, M. Hanada, D. Kharzeev, H. Lamm, Y.-Y. Li, J. Liu, M. Lukin, Y. Meurice, C. Monroe, B. Nachman, G. Pagano, J. Preskill, E. Rinaldi, A. Roggero, D.I. Santiago, M.J. Savage, I. Siddiqi, G. Siopsis, D. Van Zanten, N. Wiebe, Y. Yamauchi, K. Yeter-Aydeniz, S. Zorzetti

PRX Quantum 4, 027001, (2023)

Lattice gauge theories simulations in the quantum information era

M. Dalmonte, S. Montangero Contemporary Physics 57, 388 (2016)

Quantum Simulations of Lattice Gauge Theories using Ultracold Atoms in Optical Lattices

E. Zohar, J.I. Cirac, B. Reznik Rep. Prog. Phys. 79, 014401 (2016)

Towards Quantum Simulating QCD

U.-J. Wiese Nucl.Phys. A931, 246-256 (2014)



A fruitful dialogue (two-way communication)



Quantum Computing for High-Energy Physics State of the Art and Challenges Summary of the QC4HEP Working Group

Alberto Di Meglio,^{1,*} Karl Jansen,^{2,3,†} Ivano Tavernelli,^{4,‡} Constantia Alexandrou,^{5,3} Srinivasan Arunachalam,⁶ Christian W. Bauer,⁷ Kerstin Borras,^{8,9} Stefano Carrazza,^{10,1} Arianna Crippa,^{2,11} Vincent Croft,¹² Roland de Putter,⁶ Andrea Delgado,¹³ Vedran Dunjko,¹² Daniel J. Egger,⁴ Elias Fernández-Combarro,¹⁴ Elina Fuchs,^{1,15,16} Lena Funcke,¹⁷ Daniel González-Cuadra,^{18,19} Michele Grossi,¹ Jad C. Halimeh,^{20,21} Zoë Holmes,²² Stefan Kühn,² Denis Lacroix,²³ Randy Lewis,²⁴ Donatella Lucchesi,^{25,26,1} Miriam Lucio Martinez,^{27,28} Federico Meloni,⁸ Antonio Mezzacapo,⁶ Simone Montangero,^{25,26} Lento Nagano,²⁹ Voica Radescu,³⁰ Enrique Rico Ortega,^{31,32,33,34} Alessandro Roggero,^{35,36} Julian Schuhmacher,⁴ Joao Seixas,^{37,38,39} Pietro Silvi,^{25,26} Panagiotis Spentzouris,⁴⁰ Francesco Tacchino,⁴ Kristan Temme,⁶ Koji Terashi,²⁹ Jordi Tura,^{12,41} Cenk Tüysüz,^{2,11} Sofia Vallecorsa,¹ Uwe-Jens Wiese,⁴² Shinjae Yoo,⁴³ and Jinglei Zhang^{44,45}

arXiv:2307.03236 (2023)







Three (take-home) messages



Tensor network algorithms: dynamical string breaking and hadronization





Three (take-home) messages



Tensor network algorithms: dynamical string breaking and hadronization



Hybrid algorithms (classical/quantum): Real time evolution and light front parton correlators

$$f_{f/P}(\xi) = \sum_{S} \int \frac{dy^{-}}{2\pi} e^{-i\xi p^{+}y^{-}} \langle PS | \left[\bar{\psi}\mathcal{U}\right] \left(y^{-}\right) \frac{\gamma^{+}}{2} \left[\mathcal{U}^{\dagger}\psi\right](0) | PS \rangle$$

We show how to quantum simulate non-local Wilson loops in space and real-time



Three (take-home) messages



Tensor network algorithms: dynamical string breaking and hadronization



The role of the anomaly on decoherence robust qubits



Hybrid algorithms (classical/quantum): Real time evolution and light front parton correlators

$$f_{f/P}(\xi) = \sum_{S} \int \frac{dy^{-}}{2\pi} e^{-i\xi p^{+}y^{-}} \langle PS | \left[\bar{\psi}\mathcal{U}\right](y^{-})\frac{\gamma^{+}}{2} \left[\mathcal{U}^{\dagger}\psi\right](0) | PS \rangle$$

We show how to quantum simulate non-local Wilson loops in space and real-time









• Quantum matter as the basic building block







• Quantum matter as the basic building block

• Gauge symmetry as a fundamental principle and at the origin of every force







• Quantum matter as the basic building block

• Gauge symmetry as a fundamental principle and at the origin of every force

• Renormalisation group as a tool to study Nature at different scales





Memory

Quantum matter as the basic building block



R.P. Feynman, Int. J. Theor. Phys. (1982)



Preparation of a general quantum state

quantum correlations = entanglement





Quantum matter as the basic building block



R.P. Feynman, Int. J. Theor. Phys. (1982)



Preparation of a general quantum state

 $|\psi\rangle = c_1 |\uparrow\uparrow\uparrow\cdots\uparrow\rangle + c_2 |\uparrow\uparrow\cdots\downarrow\rangle + \cdots + c_{2^N} |\downarrow\downarrow\cdots\downarrow\rangle$

quantum correlations = entanglement

S. Lloyd, Science (1996)

Evolution of a general quantum state

 $|\psi(t)\rangle = U(t)|\psi\rangle$





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Tensor network algorithms: an overview



Variational (non-perturbative) method for Hamiltonian systems

Extremely useful in 1D systems (MPS) Proposals and extensions in higher dimensions (TNS)

Related works at: ICFO, Barcelona (Lewenstein's group); MPQ, Munich (Cirac's group); Gent - Vienna (Verstraete's group)



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Tensor network algorithms: an overview



Variational (non-perturbative) method for Hamiltonian systems Extremely useful in 1D systems (MPS) Proposals and extensions in higher dimensions (TNS)



Ground states Low-energy excitations Thermal states Time evolution Proposal for fermionic systems

Related works at: ICFO, Barcelona (Lewenstein's group); MPQ, Munich (Cirac's group); Gent - Vienna (Verstraete's group)





Tensor network algorithms: an overview

A class of tailored variational ansatz states on a lattice many-body quantum system





 $\dim(MPS) = N d D^2$

Tensor network algorithms: an overview

A class of tailored variational ansatz states on a lattice many-body quantum system

$\Psi\,$ is obtained contracting smaller tensors over auxiliary indexes

$$|\Psi_{\rm MPS}\rangle = \sum_{\{s_i\},\{\alpha_i\}} A_{\alpha_1}^{(s_1)} A_{\alpha_1,\alpha_2}^{(s_2)} \cdots A_{\alpha_{N-1}}^{(s_N)} |s_1, s_2, \cdots, s_N\rangle$$







Tensor network algorithms: an overview

A class of tailored variational ansatz states on a lattice many-body quantum system

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$$|\Psi_{\rm MPS}\rangle = \sum_{\{s_i\},\{\alpha_i\}} A_{\alpha_1}^{(s_1)} A_{\alpha_1,\alpha_2}^{(s_2)} \cdots A_{\alpha_{N-1}}^{(s_N)} |s_1, s_2, \cdots, s_N\rangle$$



$\dim(MPS) = N d D^2$

quantum correlations = entanglement =

$$\log\left(D\right) \propto \frac{c+\bar{c}}{3}\log\left(N\right)$$

C. Holzhey, F. Larsen, F. Wilczek, Nucl. Phys. B (1994) G. Vidal, J.I. Latorre, E. Rico, A. Kitaev, Phys. Rev. Lett. (2003) B.-Q. Jin, V.E. Korepin, J. Stat. Phys. (2004) P. Calabrese, J.J. Cardy, Stat. Mech. (2004)



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Tensor network algorithms: an overview



Matrix Product States: M.P.S.

Well-suited to described translational invariant systems

Encoded the entropic boundary law (VBS picture)

Optimal to minimize the energy (DMRG)



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Tensor network algorithms: an overview



Matrix Product States: M.P.S.

Well-suited to described translational invariant systems

Encoded the entropic boundary law (VBS picture)

Optimal to minimize the energy (DMRG)

Simple way to obtain any expectation value (Transfer matrix)





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Tensor network algorithms: an overview

Exact description of the gauge invariant subspace with tensor network states

$$|\text{phys}\rangle = \sum_{s_1, \dots, s_x, \dots} a(s_1, \dots, s_x, \dots) \text{Tr} \left[A^{(s_1)} \dots A^{(s_x)} \dots \right] |s_1, \dots, s_x, \cdot\rangle$$



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CP symmetry breaking in the Schwinger model in a background electric field

Imaginary time evolution (Phase diagram)

Finite density phase diagram of a SU(2) gauge invariant Fermi-Hubbard model



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Real time evolution (Quench experiment)

Entanglement characterisation of dynamical string breaking



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Imaginary time evolution (Phase diagram)

CP symmetry breaking in the Schwinger model in a background electric field

Finite density phase diagram of a SU(2) gauge invariant Fermi-Hubbard model

Real time evolution (Quench experiment)

Embedding and model building (Chiral symmetry breaking) Entanglement characterisation of dynamical string breaking

"Nuclear Physics" in a SO(3) gauge invariant model



PHYSICAL REVIEW X **6,** 011023 (2016)

Real-Time Dynamics in U(1) Lattice Gauge Theories with Tensor Networks

T. Pichler,¹ M. Dalmonte,^{2,3} E. Rico,^{4,5,6} P. Zoller,^{2,3} and S. Montangero¹

Hamiltonian: Staggered fermions in 1D coupled to a U(1) gauge field

$$H = \frac{g^2}{2} \sum_{x} \left[E_{x,x+1} \right]^2 - J \sum_{x} \left[\psi_x^{\dagger} U_{x,x+1} \psi_{x+1} + \text{h.c.} \right] + \mu \sum_{x} \left(-1 \right)^x \psi_x^{\dagger} \psi_x$$

Electric term

Matter-gauge coupling

Staggered mass



PHYSICAL REVIEW X 6, 011023 (2016)

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Electric term

Matter-gauge coupling

Staggered mass



Staggered fermions in 1D at half-filling:

- massless case: linear dispersion relation and two chiral modes
- mass gap: particle-hole excitation proportional to the staggered mass



Confinement and string breaking: QED in (1+1)-d (Schwinger model)

Spin-1 representation

 $|1\rangle$

 $|0\rangle \bigcirc$



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Confinement and string breaking: QED in (1+1)-d (Schwinger model)

Spin-1 representation



Vacuum (reference) state

$$H = \frac{g^2}{2} \sum_{\langle x, y \rangle} \left(S_{x, y}^{(3)} \right)^2 + m \sum_{x} (-1)^x \psi_x^{\dagger} \psi_x$$



Confinement and string breaking: QED in (1+1)-d (Schwinger model)

Spin-1 representation

 $|0\rangle \bigcirc |1\rangle \bigcirc |-1\rangle |0\rangle |+1\rangle$

Vacuum (reference) state

$$H = \frac{g^2}{2} \sum_{\langle x, y \rangle} \left(S_{x,y}^{(3)} \right)^2 + m \sum_{x} (-1)^x \psi_x^{\dagger} \psi_x$$

Creating a quark - antiquark pair:

$$\psi_{2x}^{\dagger} S_{2x,2x+1}^{+} \psi_{2x+1}$$



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Confinement and string breaking: QED in (1+1)-d (Schwinger model)





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Confinement and string breaking: QED in (1+1)-d (Schwinger model)





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Confinement and string breaking: QED in (1+1)-d (Schwinger model)

Confinement

$$E_{\text{string}} - E_0 = \frac{g^2}{2}(L-1)$$





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Confinement and string breaking: QED in (1+1)-d (Schwinger model)





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Confinement and string breaking: QED in (1+1)-d (Schwinger model)

String breaking and hadronization



$$E_{\rm meson} - E_0 = g^2 + 2m$$

$$L_c = 3 + \frac{4m}{g^2}$$


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Meson scattering

We prepare two mesons in a dynamical state giving them momentum towards the center



Electric field of two mesons during the scattering evolution



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Entanglement entropy during the scattering



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Tensor network algorithms and machine learning

Exploring the Phase Diagram of the quantum one-dimensional ANNNI model

M. Cea,^{1,2,*} M. Grossi,^{3,†} S. Monaco,^{4,5,‡} E. Rico,^{6,7,8,9,§} L. Tagliacozzo,^{10,¶} and S. Vallecorsa^{3,**}

arXiv:2402.11022 (2024)

O(400) sites simulation

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Tensor network algorithms and machine learning

Exploring the Phase Diagram of the quantum one-dimensional ANNNI model

M. Cea,^{1, 2}, M. Grossi,³, S. Monaco,^{4, 5}, E. Rico,^{6, 7, 8, 9}, L. Tagliacozzo,¹⁰, and S. Vallecorsa³, **

arXiv:2402.11022 (2024)

O(400) sites simulation

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Quantum Many-Body phase diagram characterization using Fidelity-based Kernels

Francesco Di Marcantonio^{2,3}

Nicola Mariella¹, Enrique Rico², Sofia Vallecorsa³, Sergiy Zhuk¹

in preparation





Core objectives for Quantum Simulation:





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Developing Tensor Network Algorithms for Quantum Many-body Problems: advancing tensor network algorithms for real-time dynamics and phase studies in high-energy process simulations, thereby facilitating in-depth fermionic system analyses and real-time simulations.





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Developing Quantum Machine Learning Models for Trigger Applications: Collaborating with CERN, to apply quantum kernel methods to enhance phase estimation, classification, and anomaly detection in many-body physics. This will contribute to a deeper theoretical understanding of these methods from a quantum information perspective.





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Prototyping TPUs and GPUs Implementation of Tensor Network Simulations: investigating how hardware accelerators like graphical processing units (GPUs) can scale tensor network algorithms to unprecedented bond dimensions, speed, and accuracy, thereby revolutionising quantum simulations.



Goal: Simulate the physics of a quantum system of interest by another quantum device that is easier to control and to measure

Quantum Simulation, Rev. Mod. Phys (2014)



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Quantum simulator



Quantum Simulation, Rev. Mod. Phys (2014)



Quantum simulation approaches



Analog simulation: Single purposed simulator

$$|\psi(0)\rangle \equiv e^{-iHt} \equiv |\psi(t)\rangle$$

Engineer the interactions to emulate the Hamiltonian of the model





Quantum simulation approaches



Analog simulation: Single purposed simulator

$$|\psi(0)\rangle \equiv e^{-iHt} \equiv |\psi(t)\rangle$$

Engineer the interactions to emulate the Hamiltonian of the model

Digital simulation: Universal simulator



Decompose dynamics into sequence of quantum gates







Feynman's universal quantum simulator:

controlled quantum device which efficiently reproduces the dynamics of any other many-particle quantum system.





Feynman's universal quantum simulator:

controlled quantum device which efficiently reproduces the dynamics of any other many-particle quantum system.

How?... cold atoms, ions, photons, superconducting circuit, etc.



Optical lattices



Superconducting circuits

... and several others as quantum dots, NMR, NV centers



Trapped ions



Quantum photonics

J.I. Cirac, P. Zoller I. Bloch, J. Dalibard, S. Nascimbène R. Blatt, C.F. Roos, A. Aspuru-Guzik, P. Walther A.A. Hock, H.E. Türeci, J. Koch Nature Physics Insight -Quantum Simulation (2012)







• Implementing the gauge invariant dynamics





First experimental realisations in trapped ions platform





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Real-time dynamics of lattice gauge theories with a few-qubit quantum computer

Esteban A. Martinez [™], Christine A. Muschik [™], Philipp Schindler, Daniel Nigg, Alexander Erhard, Markus Heyl, Philipp Hauke, Marcello Dalmonte, Thomas Monz, Peter Zoller & Rainer Blatt

Nature **534**, 516–519 (23 June 2016) | Download Citation *↓*

Self-verifying variational quantum simulation of lattice models

C. Kokail, C. Maier, R. van Bijnen, T. Brydges, M. K. Joshi, P. Jurcevic, C. A. Muschik, P. Silvi, R. Blatt, C. F. Roos & P. Zoller [™]

Nature **569**, 355–360 (2019) **Download Citation** *±*





QUANTUM SIMULATIONS OF THE SCHWINGER MODEL

variational parameters heta



IQOQI Innsbruck

R. Blatt and P. Zoller's groups



21 lattice sites!

Nature (2016), arXiv:1810.03421



A scalable realization of local U(1) gauge invariance in cold atomic mixtures

Science **367**, 1128–1130 (2020)

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Alexander Mil¹*, Torsten V. Zache², Apoorva Hegde¹, Andy Xia¹, Rohit P. Bhatt¹, Markus K. Oberthaler¹, Philipp Hauke^{1,2,3}, Jürgen Berges², Fred Jendrzejewski¹





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Observation of gauge invariance in a 71-site Bose– Hubbard quantum simulator

Bing Yang, Hui Sun, Robert Ott, Han-Yi Wang, Torsten V. Zache, Jad C. Halimeh, Zhen-Sheng Yuan

Philipp Hauke 🖂 & Jian-Wei Pan 🖂

Nature 587, 392–396 (2020) Cite this article





b Gauss's law G=0



c Hubbard simulator







Probing topological spin liquids on a programmable quantum simulator



SCIENCE • 2 Dec 2021 • Vol 374, Issue 6572 • pp. 1242-1247 • <u>DOI: 10.1126/science.abi8794</u>

Gauss law = Dimer Constraint = Rydberg blockade



219 atoms



Quantum Computing for High-Energy Physics State of the Art and Challenges Summary of the QC4HEP Working Group



arXiv:2307.03236 (2023)

Alberto Di Meglio,^{1,*} Karl Jansen,^{2,3,†} Ivano Tavernelli,^{4,‡} Constantia Alexandrou,^{5,3} Srinivasan Arunachalam,⁶ Christian W. Bauer,⁷ Kerstin Borras,^{8,9} Stefano Carrazza,^{10,1} Arianna Crippa,^{2,11} Vincent Croft,¹² Roland de Putter,⁶ Andrea Delgado,¹³ Vedran Dunjko,¹² Daniel J. Egger,⁴ Elias Fernández-Combarro,¹⁴ Elina Fuchs,^{1,15,16} Lena Funcke,¹⁷ Daniel González-Cuadra,^{18,19} Michele Grossi,¹ Jad C. Halimeh,^{20,21} Zoë Holmes,²² Stefan Kühn,² Denis Lacroix,²³ Randy Lewis,²⁴ Donatella Lucchesi,^{25,26,1}
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Development Roadmap

Executed by IBM ♥ On target ७







modern microscopes

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(semi-inclusive) deep-inelastic lepton scattering





modern microscopes

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(semi-inclusive) deep-inelastic lepton scattering highly virtual photons resolve inner (partonic) structure





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(semi-inclusive) deep-inelastic lepton scattering highly virtual photons resolve inner (partonic) structure





partonic cross section: calculable



non-perturbative parametrization of nucleon: PDFs, TMDs etc.





$$f_{f/P}(\xi) = \sum_{S} \int \frac{dy^{-}}{2\pi} e^{-i\xi p^{+}y^{-}} \langle PS | \left[\bar{\psi}\mathcal{U}\right] \left(y^{-}\right) \frac{\gamma^{+}}{2} \left[\mathcal{U}^{\dagger}\psi\right] (0) | PS \rangle$$



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Non-local (space-time) matrix elements require Wilson lines for gauge invariance We study the quantum simulation of Wilson loops in space and real-time





Quantum simulation of light-front parton correlators

M. G. Echevarria^(D),^{1,*} I. L. Egusquiza,^{2,†} E. Rico^(D),^{3,4,‡} and G. Schnell^(D),^{2,4,§}

arXiv:2011.01275 Phys. Rev. D 104, 014512 (2021)



Quantum simulation of light-front parton correlators



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Non-local (space-time) matrix elements require Wilson lines for gauge invariance We study the quantum simulation of Wilson loops in space and real-time

$$f_{f/P}(\xi) = \sum_{S} \int \frac{dy^{-}}{2\pi} e^{-i\xi p^{+}y^{-}} \langle PS | \left[\bar{\psi}\mathcal{U}\right](y^{-})\frac{\gamma^{+}}{2} \left[\mathcal{U}^{\dagger}\psi\right](0) | PS \rangle$$

Requirements for the quantum simulation of parton correlators:



Non-local (space-time) matrix elements require Wilson lines for gauge invariance We study the quantum simulation of Wilson loops in space and real-time

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Requirements for the quantum simulation of parton correlators:

- encode in quantum degrees of freedom both matter and gauge fields
- preparation of a reference state, e.g., vacuum, proton, glue-ball
- simulate gauge-invariant quantities, e.g., minimal gauge-matter coupling
- real-time evolution, since the Wilson line is non-local in time
- carry out measurements after the evolution, i.e., quantum interferometer


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Digital simulation: Universal simulator



Decompose dynamics into sequence of quantum gates

Stroboscopic simulation in an analog simulator



Discretisation of space-time in a Hamiltonian formulation

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Digital simulation: Universal simulator



Decompose dynamics into sequence of quantum gates

Stroboscopic simulation in an analog simulator

Note: in the Hamiltonian formulation the temporal gauge $A_0=0$ is chosen





Moving a single quark:

$$\begin{split} u_{12} &= \exp\left\{\frac{-i\pi}{2}\sum_{\alpha\beta}\left[\psi_{\alpha,1}^{\dagger}U_{\alpha\beta}(e)\psi_{\beta,2} + \mathrm{h.c.}\right]\right\}\\ &\to (-i)\left[\psi_{\alpha,1}^{\dagger}U_{\alpha\beta}(e)\psi_{\beta,2} + \mathrm{h.c.}\right], \end{split}$$

2



Moving a single quark:

$$u_{12} = \exp\left\{\frac{-i\pi}{2}\sum_{\alpha\beta} \left[\psi_{\alpha,1}^{\dagger}U_{\alpha\beta}(e)\psi_{\beta,2} + h.c.\right]\right\}$$
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Moving a single quark:

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Moving a single quark:

$$\begin{split} u_{12} &= \exp\left\{\frac{-i\pi}{2}\sum_{\alpha\beta}\left[\psi_{\alpha,1}^{\dagger}U_{\alpha\beta}(e)\psi_{\beta,2} + \mathrm{h.c.}\right]\right\}\\ &\to (-i)\left[\psi_{\alpha,1}^{\dagger}U_{\alpha\beta}(e)\psi_{\beta,2} + \mathrm{h.c.}\right], \end{split}$$

Starting from a "meson" state:

$$|\mathbf{m}\rangle \equiv \frac{1}{N^{1/2}} \sum_{\alpha=1}^{N} |\alpha(1), \bar{\alpha}(2)\rangle$$



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$$\begin{split} |\mathbf{m}\rangle &\equiv \frac{1}{N^{1/2}} \sum_{\alpha=1}^{N} |\alpha(1), \bar{\alpha}(2)\rangle \\ \mathscr{U}(A_1, B_L) &= \frac{1}{N^{1/2}} \sum_{\alpha\beta\cdots\mu\nu\omega\cdots\theta\phi} |\alpha(A_1)\rangle U_{\alpha\beta}(e_1)\cdots U_{\mu\nu}(e_{L/2-1}) U_{\omega\nu}^*(e_{L/2})\cdots U_{\phi\theta}^*(e_{L-1}) |\bar{\phi}(B_L)\rangle \\ &= \frac{1}{N^{1/2}} \sum_{\alpha\phi} |\alpha(A_1)\rangle \mathscr{U}_{\alpha\phi}(e_1, \cdots, e_{L-1}) |\bar{\phi}(B_L)\rangle \\ \end{split}$$
 we built a spatial Wilson line



Time-evolution by a single time step

Universidad

del País Vasco

$$|\psi(0)\rangle \equiv e^{-iH\tau} \equiv |\psi(\tau)\rangle$$

$$W(\tau,\lambda) = W_{C_1}W_{\tau_1}W_{C_2}W_{\tau_2}\cdots W_{C_k}W_{\tau_k}\cdots$$





Decompose dynamics induced by systems' Hamiltonian into sequence of quantum gates

Digital simulation can simulate any model but requires many gate operations

Stroboscopic simulation in an analog simulator

 $H = H_{\rm el} + H_{\rm mag}$

 $e^{-iH} \simeq \left[e^{-iH_{\rm el}/2n_T} e^{-i\lambda H_{\rm mag}/n_T} e^{-iH_{\rm el}/2n_T} \right]^{n_T}$ Trotter-Suzuki approximation

Efficient for local interactions

S. Lloyd, Science (1996)



Proof of principle: Z₂ pure gauge model



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within a few Trotter steps a fidelity closed to one is achieved

operator norm:

 $\left| \operatorname{Tr} \left[\mathscr{W}^{\dagger} \mathscr{W}_{\mathbf{n}_{\mathrm{T}}} \right] \right|$

ground state fidelity:

 $\langle g.s. | \mathcal{W}^{\dagger} \mathcal{W}_{n_{T}} | g.s. \rangle$



CERN's new Next Generation Triggers Project (NGT), supported by a grant from the E&W Schmidt Fund for Strategic Innovation.



Core objectives for Quantum Simulation:

Quantum Simulations on Classical Hardware: quantum algorithms for state preparation and real-time dynamics simulations, including quantum simulation of parton distribution functions from first principles. Initially, I plan to utilise classical hardware to achieve a hybrid simulation, blending classical and quantum approaches seamlessly.



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Developing and Deploying Quantum Circuit Simulations for Large Systems: Engaged in spearheading the 100x100 IBM challenge within the QC4HEP collaboration, I aim to leverage near-term quantum hardware for high-energy physics applications by developing and deploying quantum circuit simulations for large systems comprising O(100) qubits.



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Quantum Simulations on Quantum Hardware: exploring Rydberg quantum platforms and strengthening collaborations with experimentalists in superconducting circuits to advance the quantum simulation of high-energy problems, extending existing partnerships from projects such as QuantERA.







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Letter

Role of anomalous symmetry in $0-\pi$ qubits

I. L. Egusquiza,^{1,2,*} A. Iñiguez,^{3,†} E. Rico,^{2,4,5,‡} and A. Villarino^{4,§} ¹Department of Physics, University of the Basque Country UPV/EHU, Apartado 644, 48080 Bilbao, Spain ²EHU Quantum Center, University of the Basque Country, UPV/EHU, Barrio Sarriena s/n, 48940 Leioa, Biscay, Spain ³Department of Mathematics, University of the Basque Country UPV/EHU, Apartado 644, 48080 Bilbao, Spain ⁴Department of Physical Chemistry, University of the Basque Country UPV/EHU, Apartado 644, 48080 Bilbao, Spain ⁵IKERBASQUE, Basque Foundation for Science, Plaza Euskadi 5, 48009 Bilbao, Spain



The role of the anomaly on decoherence robust qubits





Sources of decoherence: decay, dephasing, spin-flip



 $\hat{H} \propto \epsilon \mathbb{I}$



The role of the anomaly on decoherence robust qubits





Qubit alive thanks to the anomaly

In quantum physics an anomaly or quantum anomaly appears when the symmetry of a classical theory is not equally represented by the quantum theory.



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Qubit alive thanks to the anomaly

In quantum physics an anomaly or quantum anomaly appears when the symmetry of a classical theory is not equally represented by the quantum theory.

The degeneracies in the protected regime of the $0 - \pi$ qubit are a remnant of the anomalous symmetry. Degeneracies independent of energy parameters





Classical group symmetry of the ring









Classical group symmetry of the ring





Quantum particle on a ring

$$\hat{H} = E_c \left(\hat{n} - n_g\right)^2$$

Any rotation is a symmetry of the quantum Hamiltonian:

 $\hat{U}_{\alpha} = e^{i\hat{n}\alpha}$ $SO(2) \sim U(1)$





Classical group symmetry of the ring





Quantum particle on a ring

Any rotation is a symmetry of the quantum Hamiltonian: About the reflexion symmetry...



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The role of the anomaly on decoherence robust qubits



 $0 - \pi$ Hamiltonian

$$H_{0-\pi} = 4E_{C_J}\hat{Q}_{\phi}^2 + E_L\hat{\phi}^2 + 4E_{C_s}\left(\hat{n}_{\theta} - n_g\right)^2 - 2E_J\cos\hat{\theta}\cos\left(\hat{\phi} - \frac{\varphi_{\text{ext}}}{2}\right)$$



The role of the anomaly on decoherence robust qubits



$0 - \pi$ Hamiltonian







E

EHU QC EHU Quantum Center

The role of the anomaly on decoherence robust qubits



 $0 - \pi$ Hamiltonian



high symmetry point

 $n_g = 1/2, \ \varphi_{\rm ext} = \pi$

$$\hat{V}_P = e^{-i\hat{\theta}}\hat{U}_P \qquad \begin{array}{c} n_\theta \to 1 - n_\theta \\ \theta \to -\theta \end{array}$$

$$\hat{U}_{\pi} = e^{i\hat{n}_{\theta}\pi}\hat{P}_{\phi} \qquad \begin{array}{c} \theta \to \theta + \pi \\ \phi \to -\phi \end{array}$$

 $\hat{V}_P \hat{U}_\pi = - \hat{U}_\pi \hat{V}_P$



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the whole spectrum is two-fold degenerate independent of any energy scales





A. Kitaev, Protected qubit based on a superconducting current mirror, arXiv:cond-mat/0609441.

P. Brooks, A. Kitaev, and J. Preskill, Protected gates for superconducting qubits, Phys. Rev. A 87, 052306 (2013).

B. Douçot and J. Vidal, Pairing of Cooper Pairs in a Fully Frustrated Josephson Junction Chain, Phys. Rev. Lett. 88, 227005 (2002).

L. B. Ioffe and M. V. Feigel'man, Possible realization of an ideal quantum computer in Josephson junction array, Phys. Rev. B 66, 224503 (2002).

J. M. Dempster, B. Fu, D. G. Ferguson, D. I. Schuster, and J. Koch, Understanding degenerate ground-states of a protected quantum circuit in the presence of disorder, Phys. Rev. B 90, 094518 (2014).

W. C. Smith, A. Kou, X. Xiao, U. Vool, and M. H. Devoret, Superconducting circuit protected by two-Cooper-pair tunneling, npj Quantum Inf. 6, 8 (2020).

P. Groszkowski, A. Di Paolo, A. L. Grimsmo, A. Blais, D. I. Schuster, A. A. Houck, and J. Koch, Coherence properties of the $0-\pi$ qubit, New J. Phys. **20**, 043053 (2018).

A. Di Paolo, A. L. Grimsmo, P. Groszkowski, J. Koch, and A. Blais, Control and coherence time enhancement of the $0-\pi$ qubit, New J. Phys. **21**, 043002 (2019).

A. Gyenis, P. S. Mundada, A. Di Paolo, T. M. Hazard, X. You, D. I. Schuster, J. Koch, A. Blais, and A. A. Houck, Experimental realization of a protected superconducting circuit derived from the $0-\pi$ qubit, PRX Quantum 2, 010339 (2021).



Three (take-home) messages



Tensor network algorithms: dynamical string breaking and hadronization





Three (take-home) messages



Tensor network algorithms: dynamical string breaking and hadronization



Hybrid algorithms (classical/quantum): Real time evolution and light front parton correlators

$$f_{f/P}(\xi) = \sum_{S} \int \frac{dy^{-}}{2\pi} e^{-i\xi p^{+}y^{-}} \langle PS | \left[\bar{\psi}\mathcal{U}\right] \left(y^{-}\right) \frac{\gamma^{+}}{2} \left[\mathcal{U}^{\dagger}\psi\right](0) | PS \rangle$$

We show how to quantum simulate non-local Wilson loops in space and real-time



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Nuclear Physics

Theory development:

Formulate QCD in the Hamiltonian language Optimal bases towards the continuum limit Importance of gauge invariance Systematics: finite volume, space-time discretisation



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Algorithmic developments:

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Resources for gauge theories in quantum algorithm Simulation of higher-dimensional non-abelian theories Computation of observables like scattering amplitudes ON Quantum state preparation Data analysis with Quantum Machine Learning





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Implementation, benchmark, and co-development:

Limits of actual hardware for gauge-theory quantum simulation Nature of noise in actual hardware and ways to mitigate it Co-design and co-develop quantum hardware for gauge theory simulations Digital, analog and hybrid ideas to simulate field theories.