

Long-distance contributions

from lattice QCD & QED

- Isospin-breaking and electromagnetic corrections to hadronic interactions
- Non-local matrix elements: $K \rightarrow \pi \ell^+ \ell^-$ decays
- Multi-hadrons interactions
- (Not $g - 2$)

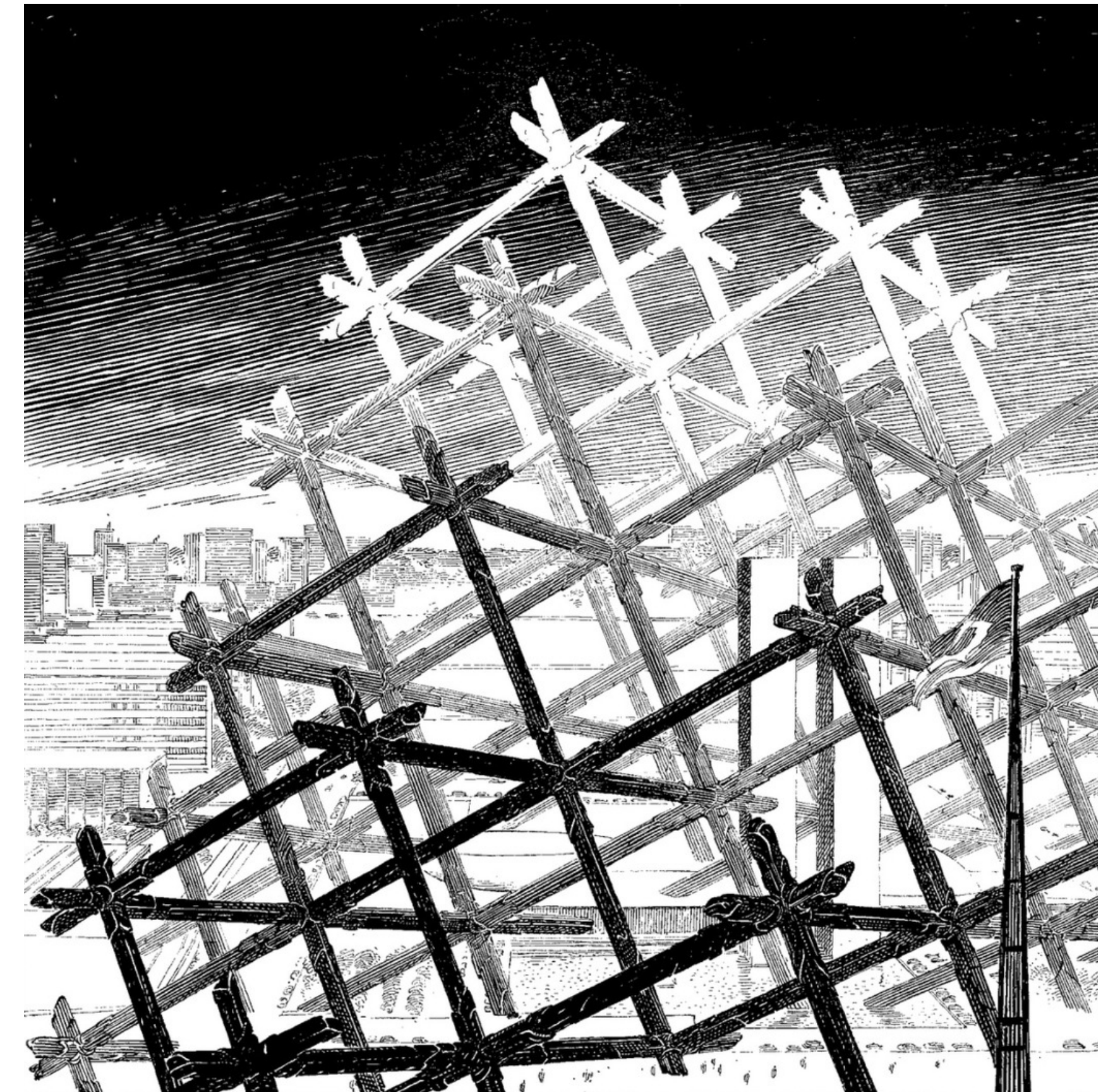
Lattice field theory

- Strong interactions described by **Quantum Chromodynamics (QCD)**
- In a discrete and Euclidean space-time, **QCD becomes equivalent to a statistical system**

📖 KG Wilson, Phys Rev D 10(8) (1974)

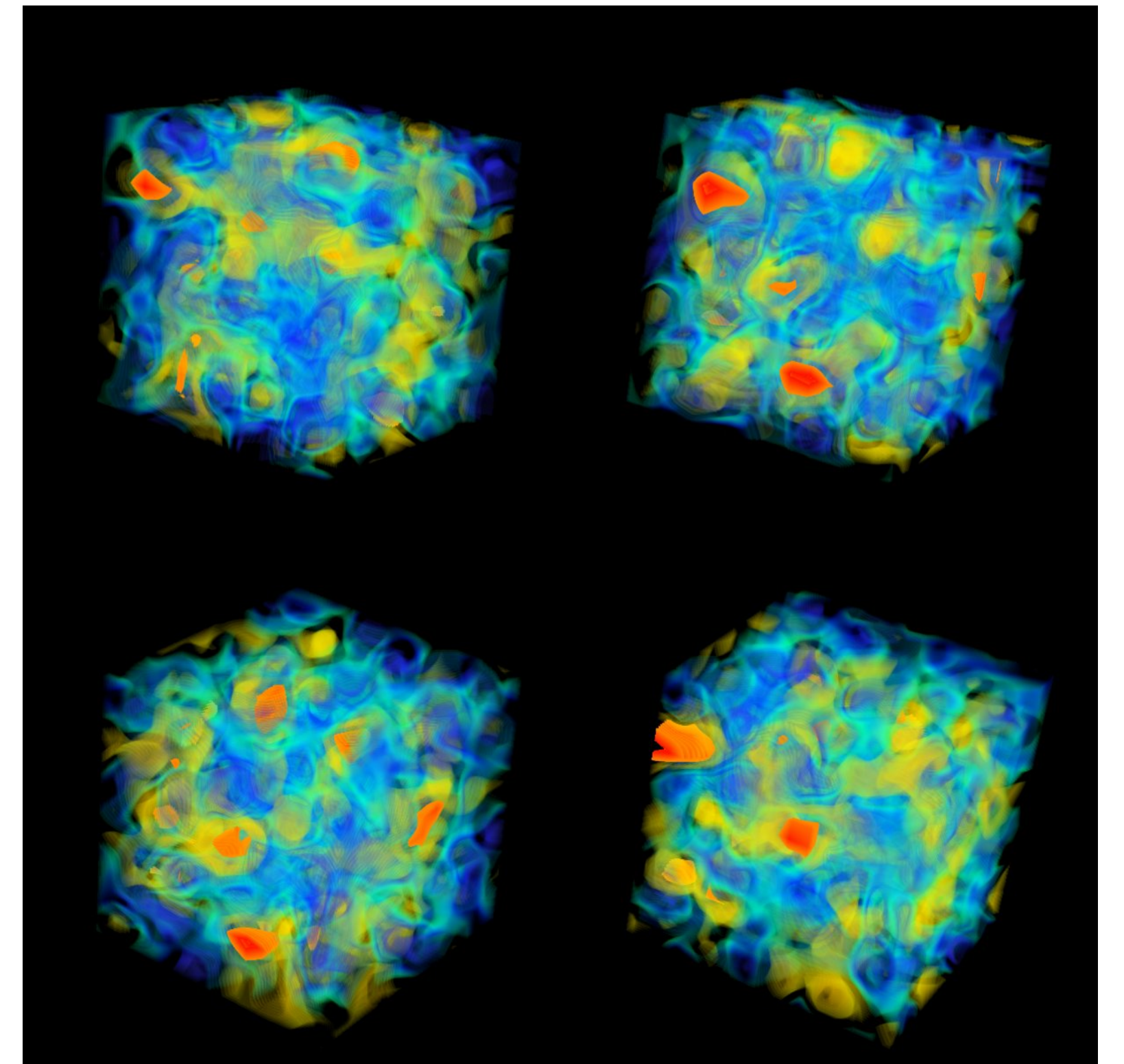
- Physical observables can be evaluated **through Monte-Carlo simulations**

$$\langle O \rangle = \frac{1}{\mathcal{Z}} \int DU O[U] \det(M[U]) e^{-S[U]}$$



Lattice field theory

- Lattice simulations have **billions of degrees of freedom**
- **They can potentially describe any strongly bound quantum field theory from first principles**
- Predictive capacity is directly bounded by
 - **available supercomputing power**
 - **algorithmic research progress**
 - **understanding of Euclidean field theory**



Isospin-breaking corrections to hadronic interactions



Isospin-breaking (IB) corrections

Beyond isospin symmetry

- Isospin symmetry assumed in most lattice calculations
- Violations generally expected to be $\mathcal{O}(1\%)$ of hadronic observables
- This is **highly relevant for** searches for new physics through precision measurements ($g - 2$ & weak decays)
- **Main challenge** for lattice QCD: **adding QED**

IB corrections to weak decays

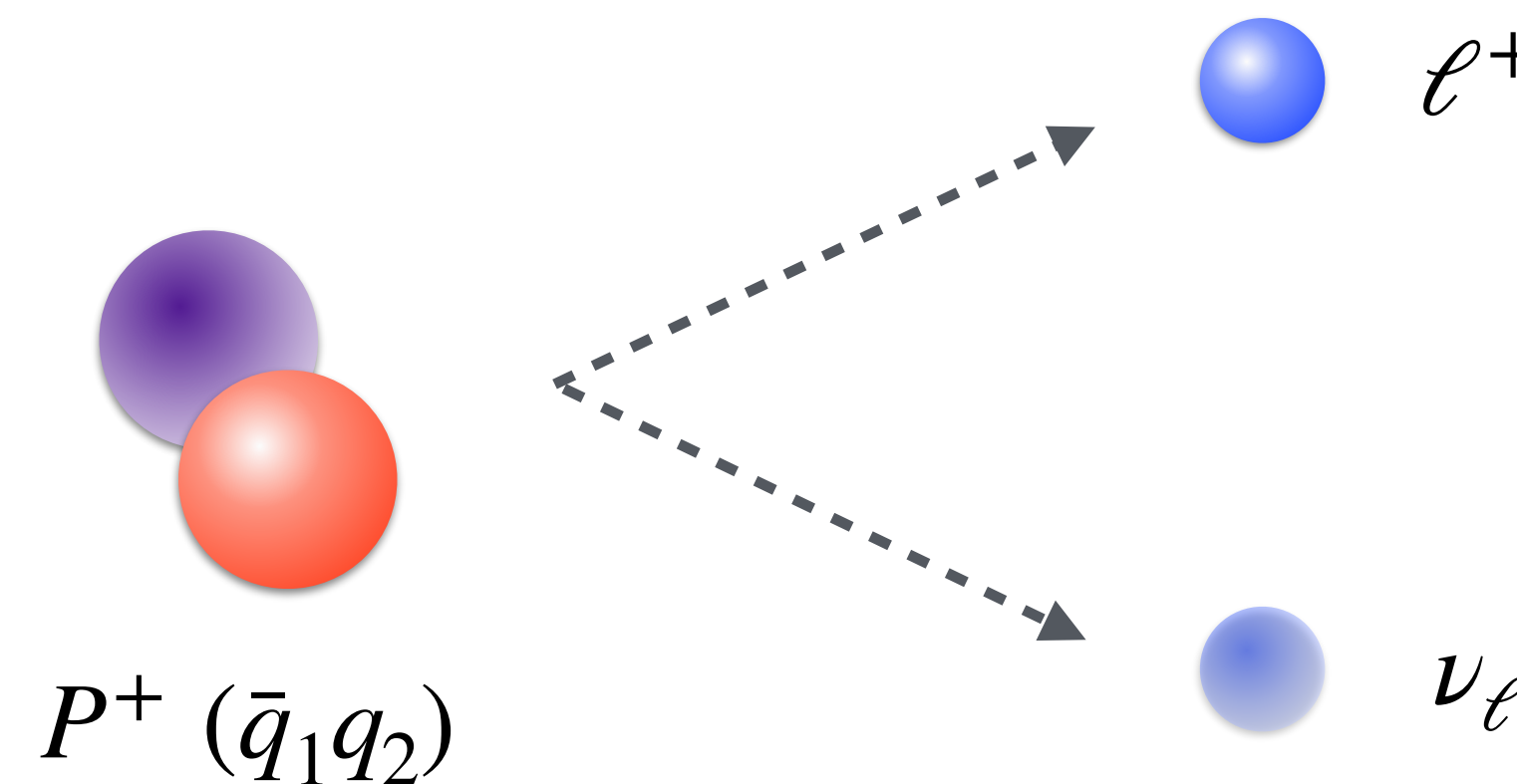
CKM matrix elements from leptonic decays

- Leptonic meson decay:

quark pair W -boson annihilation

- Rate proportional to $|V_{q_1q_2}|^2$
- Allows to determine CKM matrix element from experimental rate...

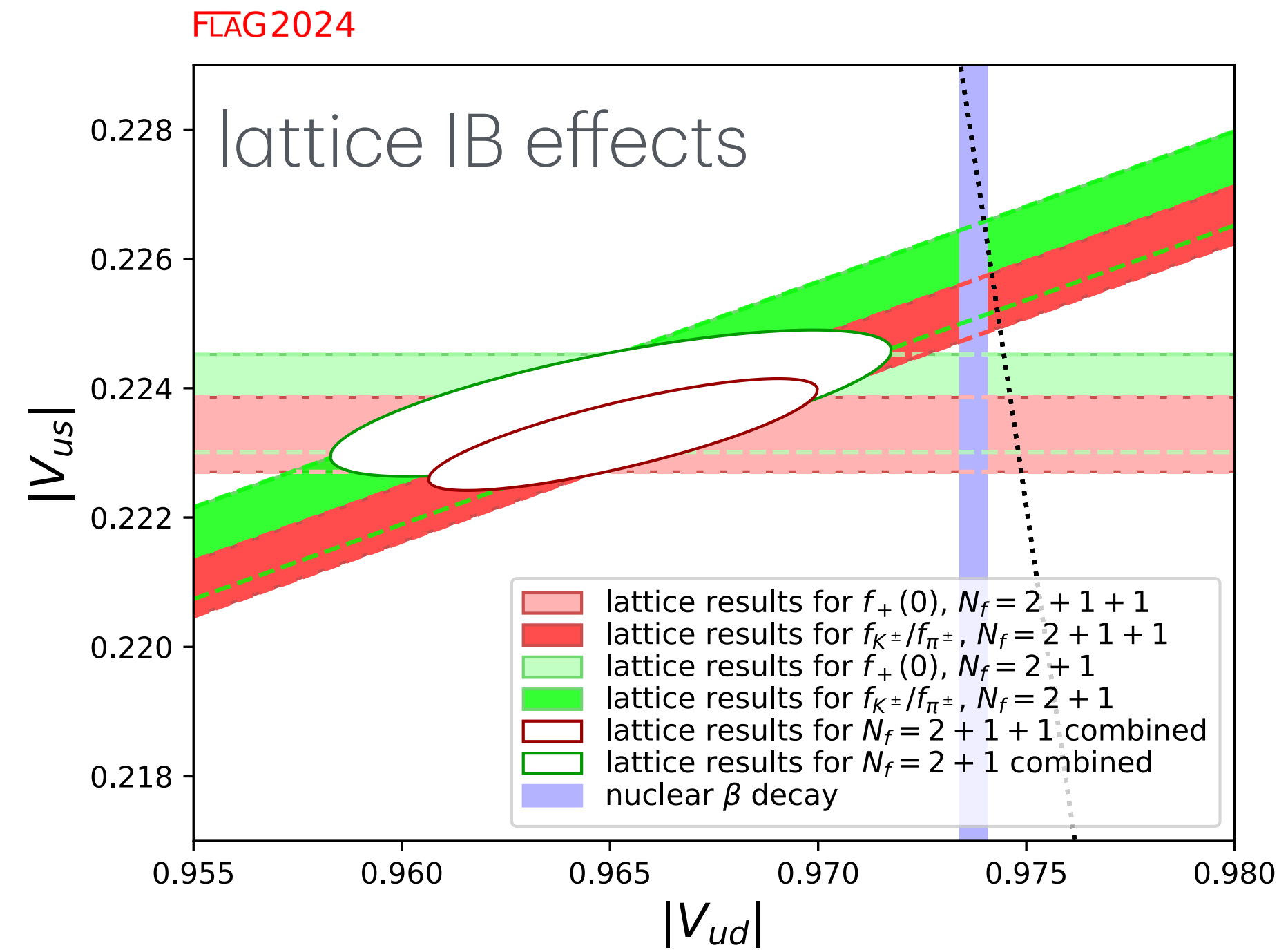
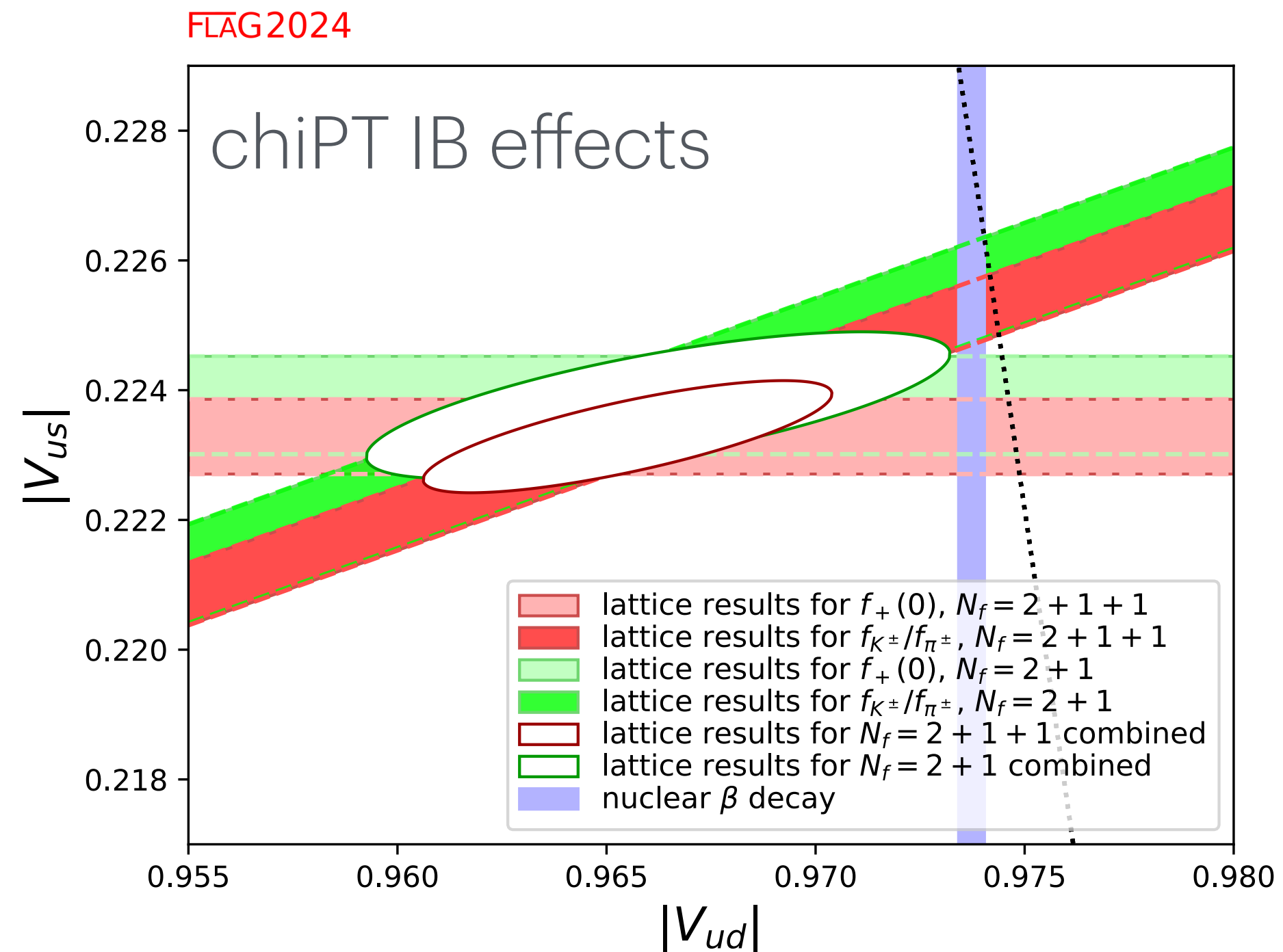
- ...but **needs a high-precision description of the hadronic dynamics**



$$\Gamma(P^+ \rightarrow \ell^+ \nu_\ell [\gamma]) = \frac{G_F^2}{8\pi} f_P^2 m_\ell^2 M_P \left(1 - \frac{m_\ell^2}{M_P^2}\right)^2 |V_{q_1q_2}|^2 (1 + \delta R_P)$$

IB corrections to weak decays

CKM first row



FLAG Review 2024

- Neutron lifetime under scrutiny, but **IB corrections to leptonic decays also relevant**

IB corrections to weak decays

CKM charm and bottom coefficients

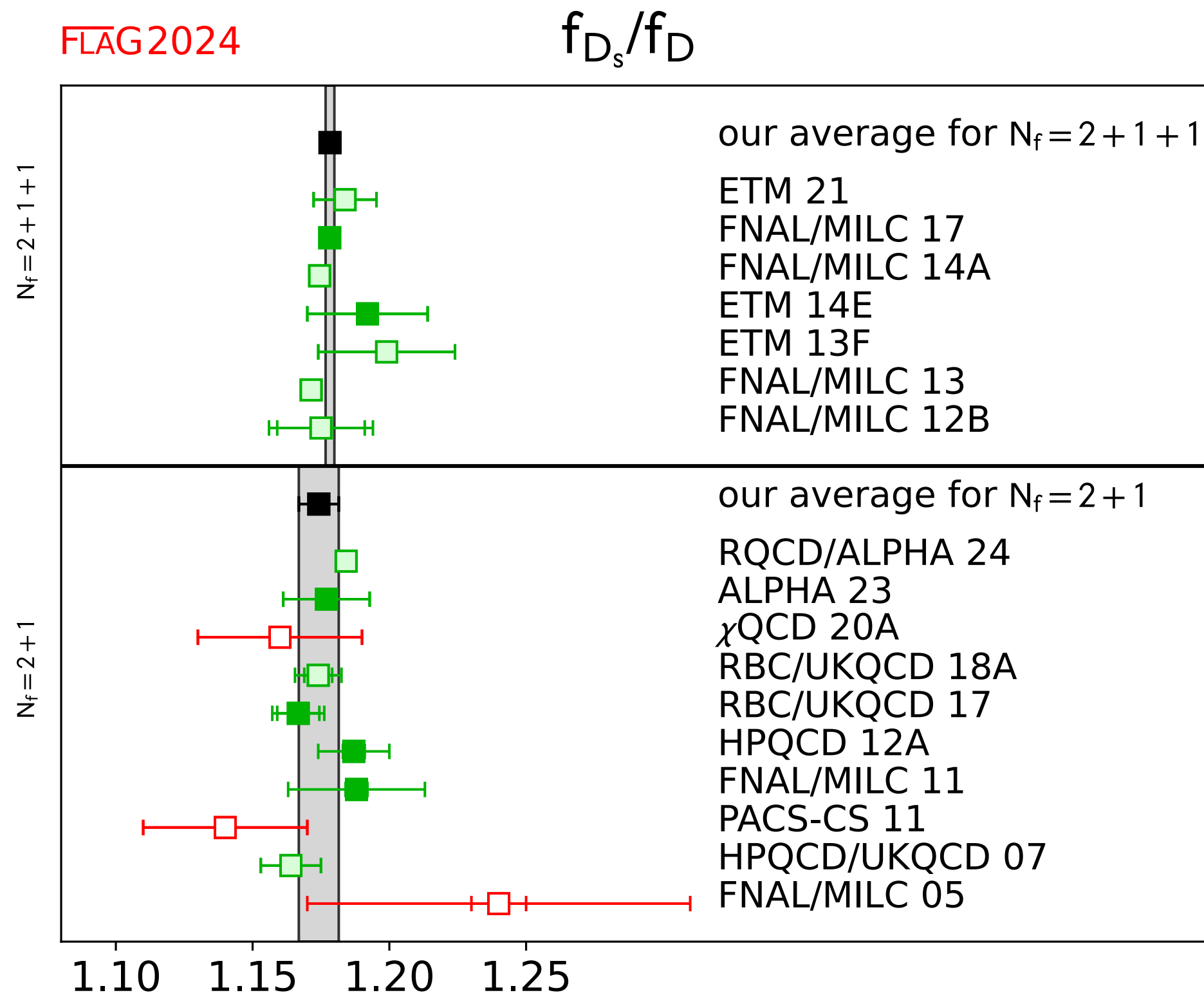


Table 30. Lattice inputs for decay constants $f_{B(s)}$ and bag parameters $B_{B(s)}$ in the SM. The current average of $f_{B(s)}$ for $N_f=2+1$ and $2+1+1$ are obtained from Refs. [150,213–216] and Refs. [212,217], respectively. The average of $B_{B(s)}$ is obtained from Refs. [148,150,151]. $f_{B(s)}\sqrt{B_{B(s)}}$ is in units of MeV.

N_f	Input	f_B [MeV]	f_{B_s} [MeV]	f_{B_s}/f_B
2+1+1	Current	188(3)	227(4)	1.203(0.007)
	5 yr w/o EM	188(1.5)	227(2.0)	1.203(0.0035)
	5 yr with EM	188(2.4)	227(3.0)	1.203(0.013)
	10 yr w/o EM	188(0.60)	227(0.80)	1.203(0.0014)
	10 yr with EM	188(2.0)	227(2.4)	1.203(0.012)
	our average			1.203(0.007)
2+1	Current	192.0(4.3)	228.4(3.7)	1.201(0.016)
	5 yr w/o EM	192.0(2.2)	228.4(1.9)	1.201(0.0080)
	5 yr with EM	192.0(2.9)	228.4(2.9)	1.201(0.014)
	10 yr w/o EM	192.0(0.86)	228.4(0.74)	1.201(0.0032)
	10 yr with EM	192.0(2.1)	228.4(2.4)	1.201(0.012)
our average			1.201(0.007)	
N_f	Input	$f_B\sqrt{B_B}$	$f_{B_s}\sqrt{B_{B_s}}$	ξ
2+1	Current	225(9)	274(8)	1.206(0.017)
	5 yr w/o EM	225(4.5)	274(4.0)	1.206(0.0085)
	5 yr with EM	225(5.0)	274(4.8)	1.206(0.015)
	10 yr w/o EM	225(1.8)	274(1.6)	1.206(0.0034)
	10 yr with EM	225(2.9)	274(3.2)	1.206(0.013)
	our average			1.206(0.007)
N_f	Input	B_B	B_{B_s}	B_{B_s}/B_B
2+1	Current	1.30(0.09)	1.35(0.06)	1.032(0.036)
	5 yr w/o EM	1.30(0.045)	1.35(0.030)	1.032(0.018)
	5 yr with EM	1.30(0.047)	1.35(0.033)	1.032(0.021)
	10 yr w/o EM	1.30(0.018)	1.35(0.012)	1.032(0.0072)
	10 yr with EM	1.30(0.022)	1.35(0.018)	1.032(0.013)
	our average			1.032(0.007)

IB corrections to weak decays

CKM charm and bottom coefficients

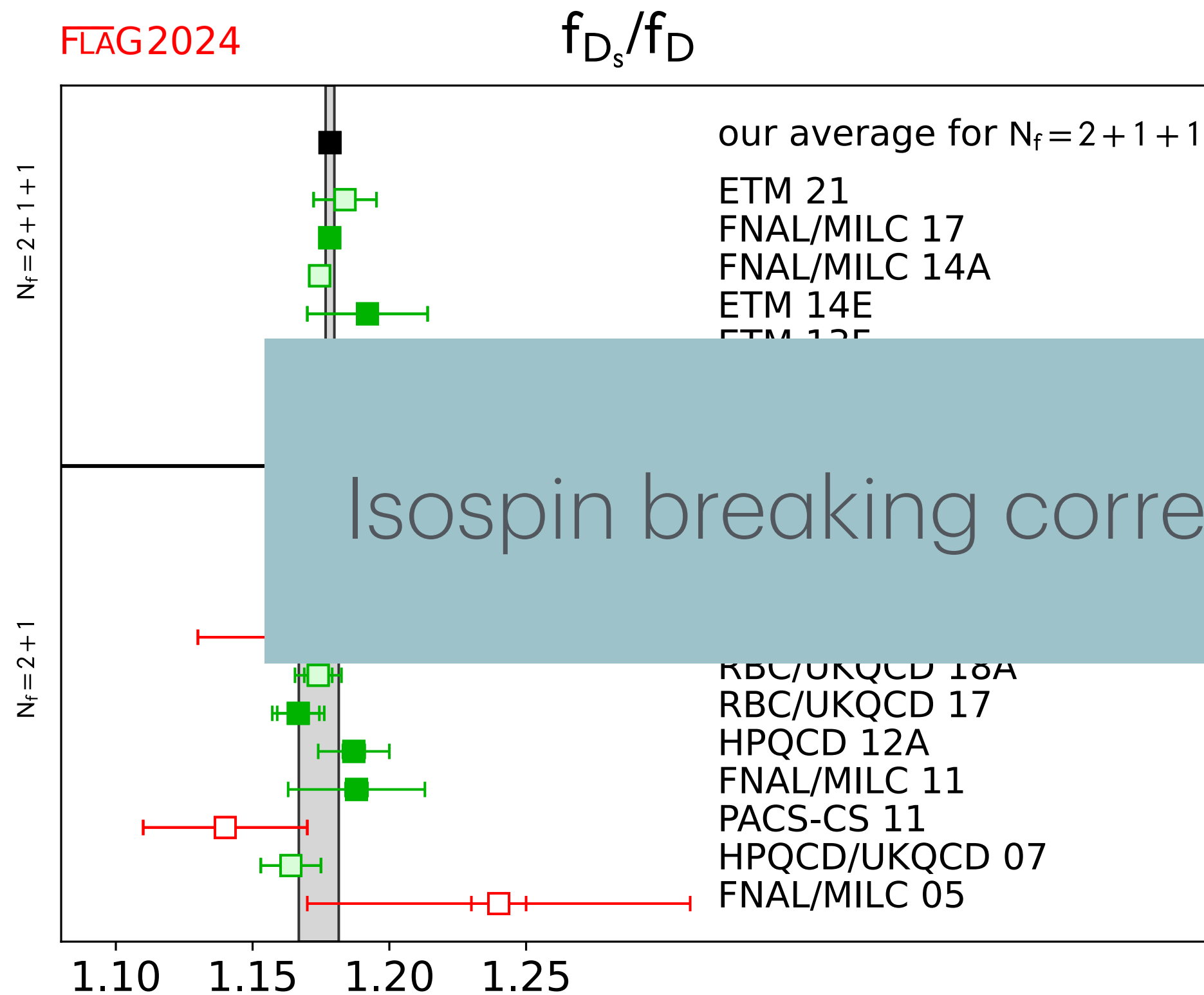


Table 30. Lattice inputs for decay constants $f_{B(s)}$ and bag parameters $B_{B(s)}$ in the SM. The current average of $f_{B(s)}$ for $N_f=2+1$ and $2+1+1$ are obtained from Refs. [150,213–216] and Refs. [212,217], respectively. The average of $B_{B(s)}$ is obtained from Refs. [148,150,151]. $f_{B(s)}\sqrt{B_{B(s)}}$ is in units of MeV.

N_f	Input	f_B [MeV]	f_{B_s} [MeV]	f_{B_s}/f_B
	Current	188(3)	227(4)	1.203(0.007)
2+1+1	5 yr w/o EM	188(1.5)	227(2.0)	1.203(0.0035)
	5 yr with EM	188(2.4)	227(3.0)	1.203(0.013)
	10 yr w/o EM	188(0.60)	227(0.80)	1.203(0.0014)
	10 yr with EM	188(0.8)	227(0.4)	1.203(0.012)
	Current	1.30(0.09)	1.35(0.06)	1.032(0.036)
2+1	5 yr w/o EM	225(4.5)	274(4.0)	1.206(0.0085)
	5 yr with EM	225(5.0)	274(4.8)	1.206(0.015)
	10 yr w/o EM	225(1.8)	274(1.6)	1.206(0.0034)
	10 yr with EM	225(2.9)	274(3.2)	1.206(0.013)
	Current	1.30(0.045)	1.35(0.030)	1.032(0.018)
2+1	5 yr with EM	1.30(0.047)	1.35(0.033)	1.032(0.021)
	10 yr w/o EM	1.30(0.018)	1.35(0.012)	1.032(0.0072)
	10 yr with EM	1.30(0.022)	1.35(0.018)	1.032(0.013)

Isospin breaking corrections critically needed in all cases

IB corrections to weak decays

First physical K & π leptonic decay calculation

- First calculation at the physical point of **IB**

corrections to K & π leptonic decay rate ratio

 P Boyle, **AP**, *et al.* JHEP 02 (2023)

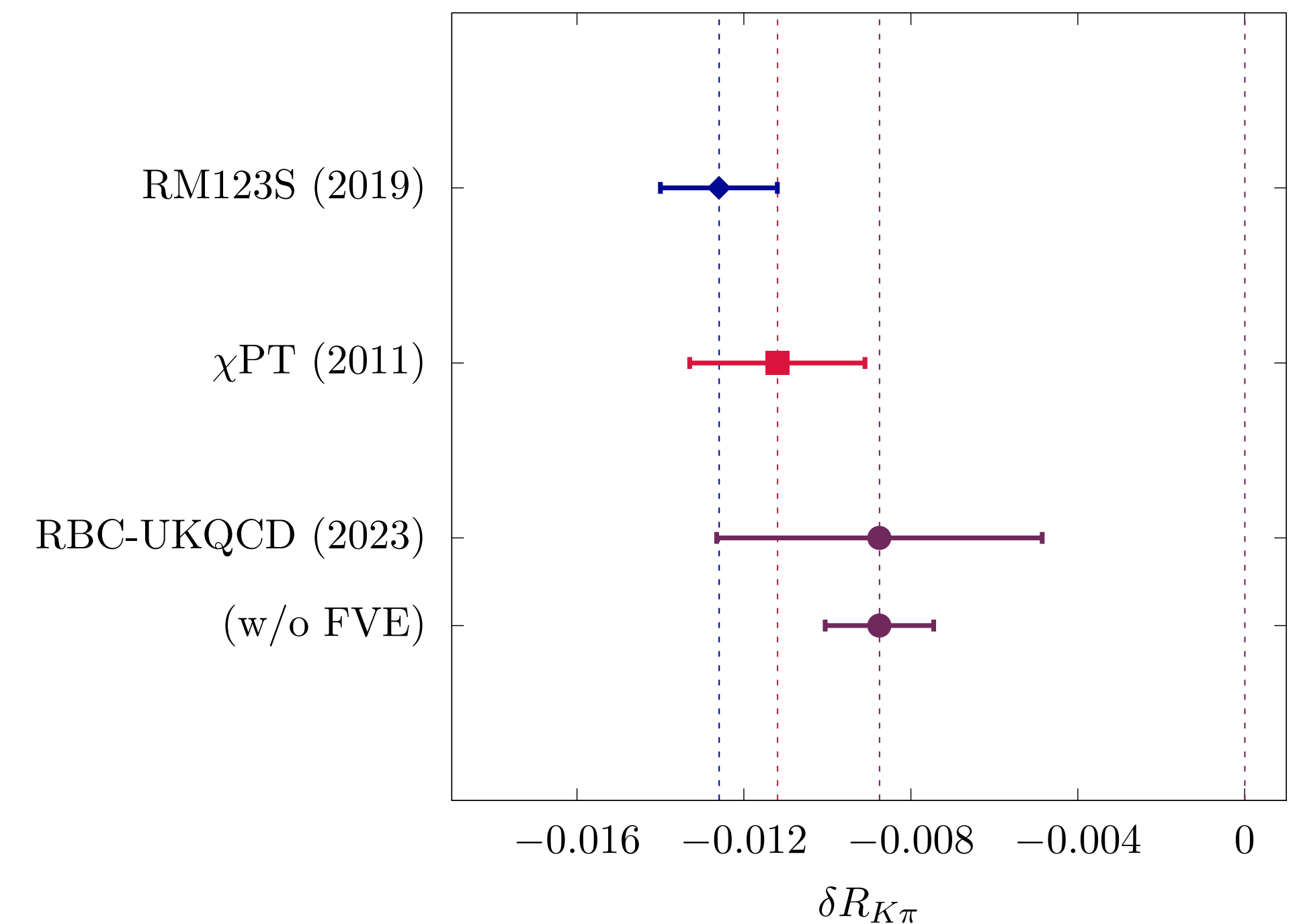
$$\delta R_{K\pi} = -0.0086(3)_{\text{stat.}} \left(\begin{matrix} +11 \\ -4 \end{matrix} \right)_{\text{fit}} (5)_{\text{disc.}} (5)_{\text{quench.}} (39)_{\text{vol.}}$$

- Largely based on the RM123S formalism

 N Carrasco, *et al.* PRD 91(7) (2015)

- Still **uncontrolled systematics**

FV effects, QED quenching, continuum limit

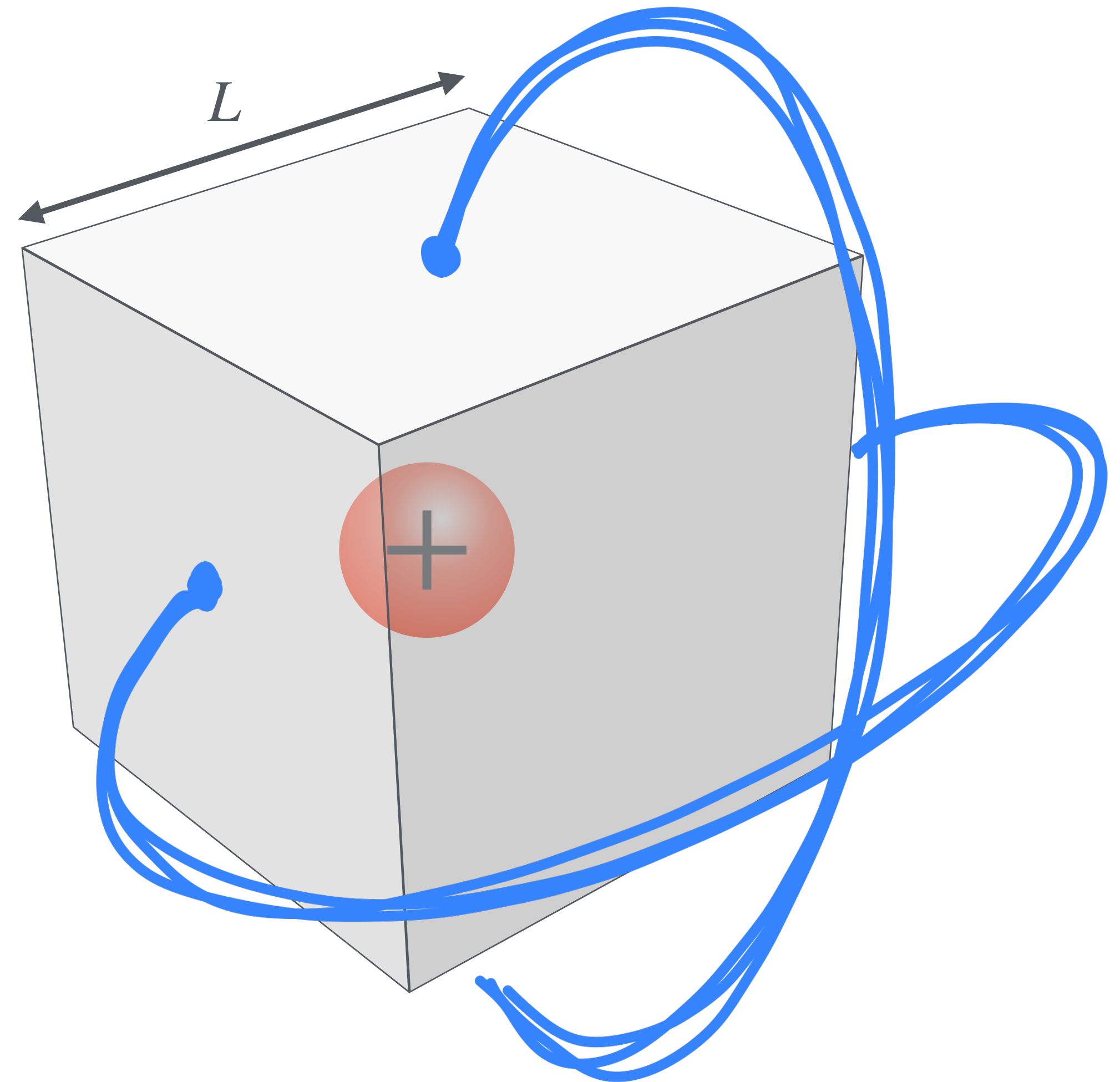


Finite-volume QED

Zero-mode singularities

- Periodic boundary conditions
 \implies **EM field feedback loop**
- **Large finite-volume (FV) effects expected**
- In reality, it is worse than that:
 Feedback loop diverges

(think about a lattice of Coulomb potentials)



Finite-volume QED

Zero-mode singularities, quantum field theory

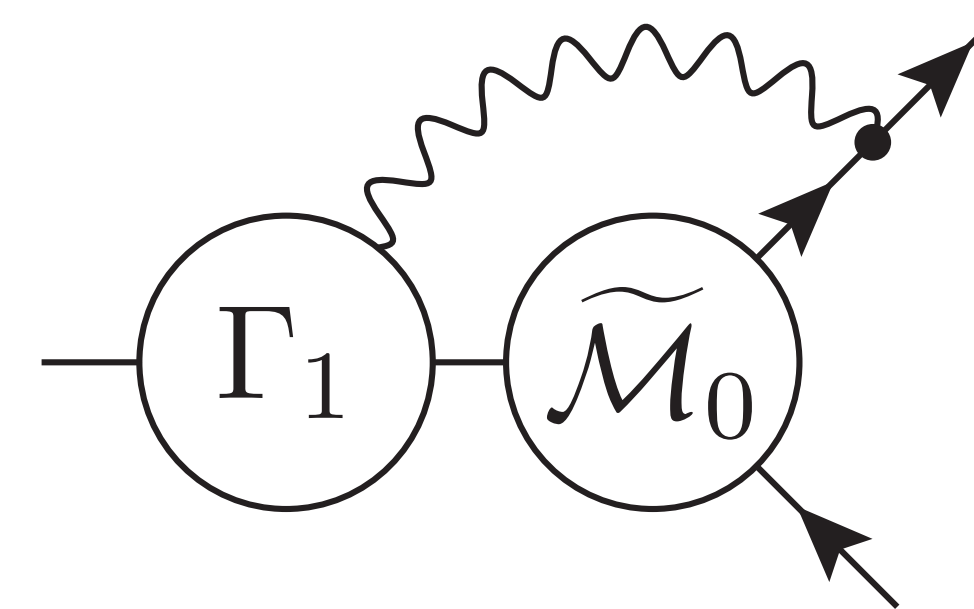
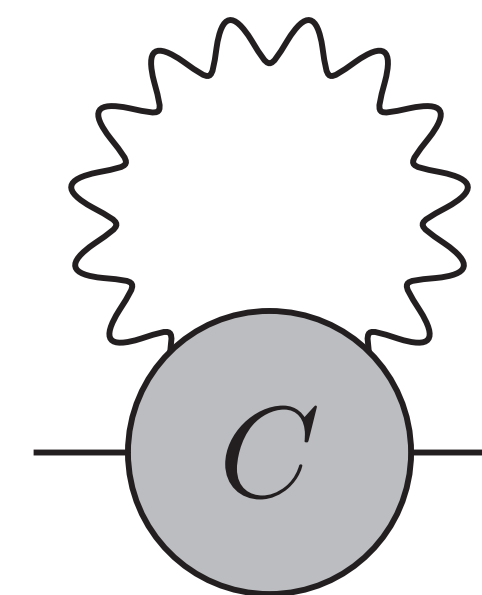
- One-loop QED amplitude

$$\int \frac{d^3\mathbf{k}}{(2\pi)^3} \frac{f(\mathbf{k})}{\mathbf{k}^2} \mapsto \frac{1}{L^3} \sum_{\mathbf{k}} \frac{f(\mathbf{k})}{\mathbf{k}^2}, \quad \text{with } \mathbf{k} = \frac{2\pi}{L}\mathbf{n}$$

maybe divergent
IR divergences

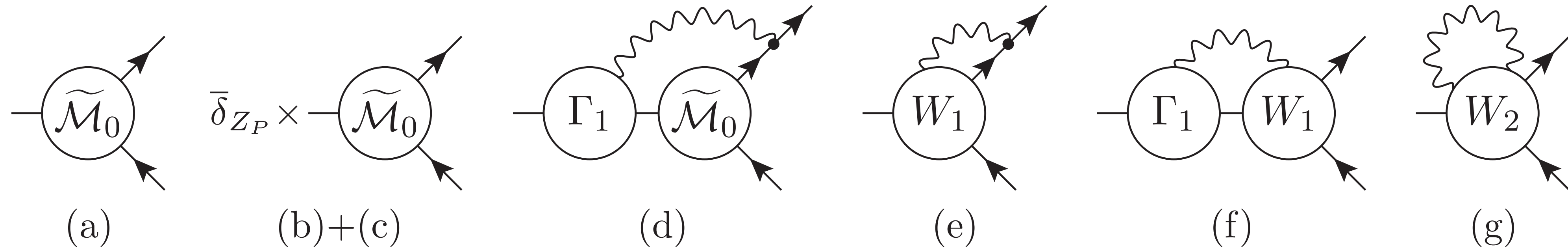
undefined because
of $f(\mathbf{0})/0$ term

- Regularisation or change of BC required
- QED_L : remove all 3D zero-modes $\mathbf{k} = \mathbf{0}$.
Non-local modification of QED



Finite-volume QED

Leptonic decays



- Known, universal $\log(ML)$, $1/L$ finite-size effects

 B Lucini, *et al.* PRD 95(3) (2017)

- Known structure-dependent $1/L^2$ finite-size effects

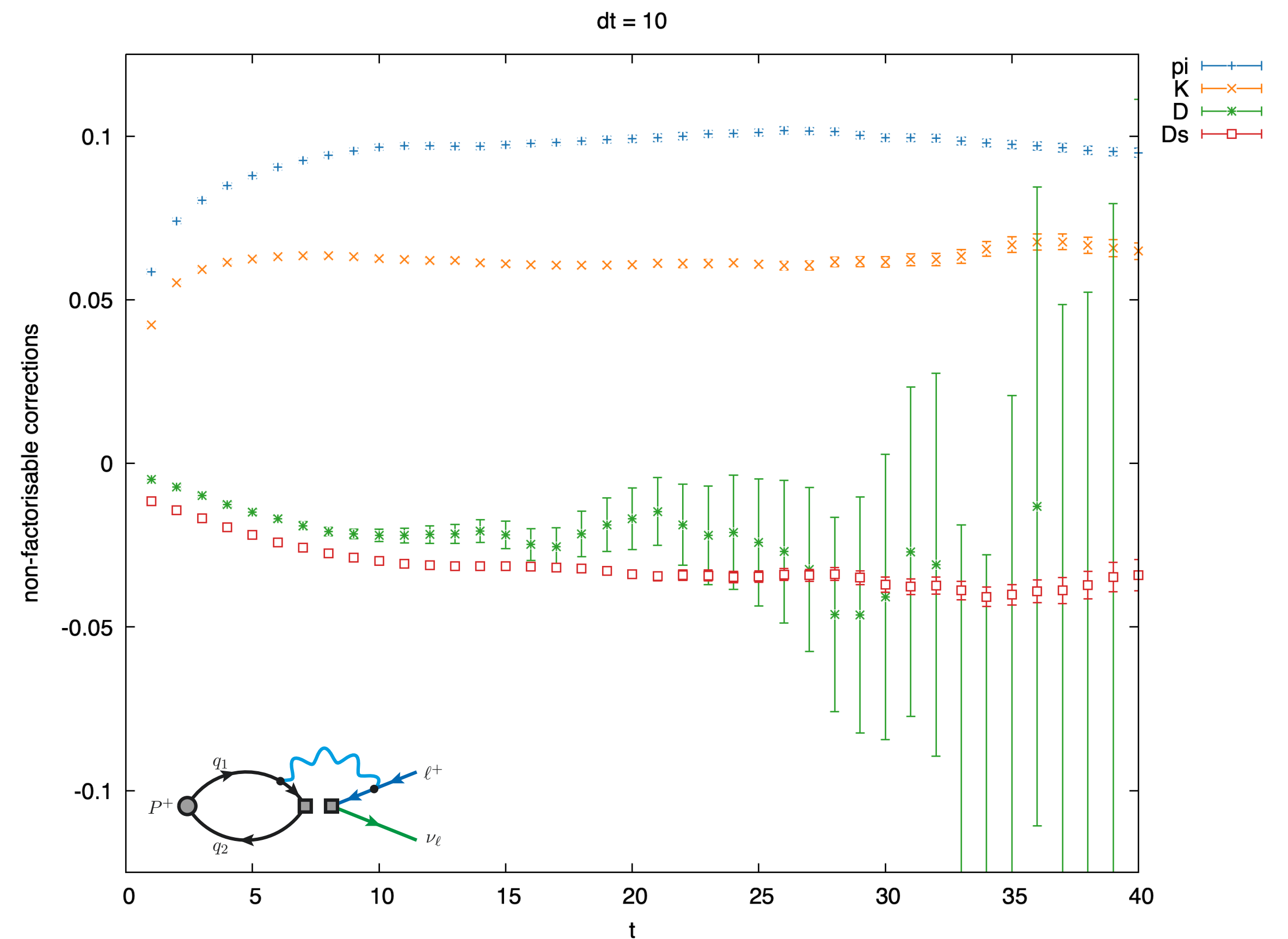
 M Di Carlo, *et al.* PRD 105(1) (2022)

- Unknown structure-dependent $1/L^3$, **potentially large**:
source of large uncertainty in  P Boyle, **AP**, *et al.* JHEP 02 (2023)

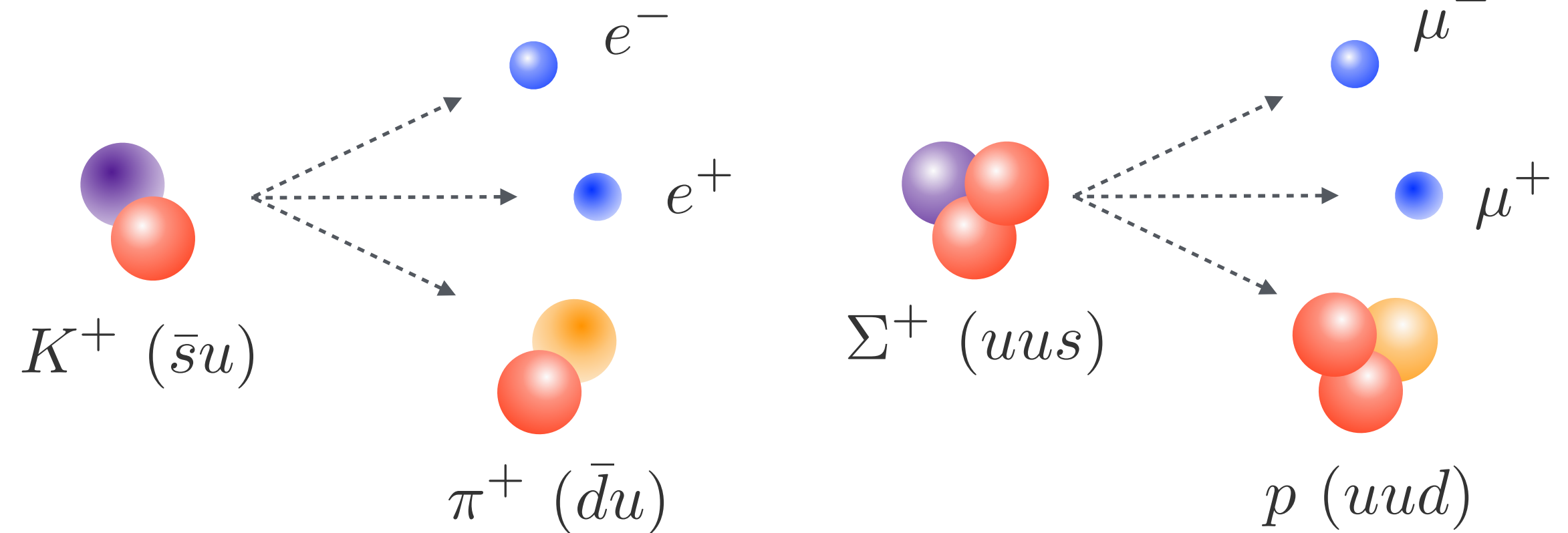
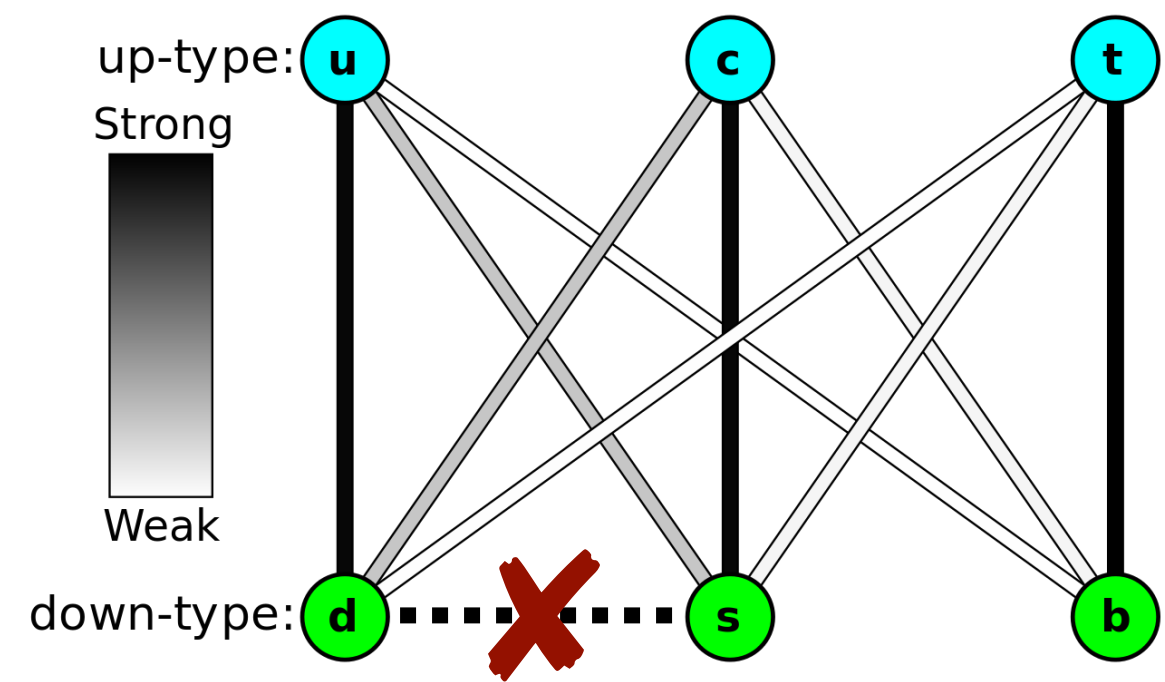
IB corrections to weak decays

Future progress

- First calculation for D & D_s decays
- **Collaboration with KEK** on B & B_s decays
- Continuum limit for K & π decays
- Improving FV QED / removing $1/L^3$ terms
 - 📖 Z Davoudi, **AP**, et al. Phys Rev D 99(3) (2019)
 - 📖 M Di Carlo, PoS Lattice 2023 (2024)
 - 📖 M Di Carlo, **AP**, et al. in preparation (2025)



Rare $s \rightarrow d$ decays



- A promising avenue for new physics is to study **flavour-changing neutral current** decays
- Those are forbidden at leading order in the SM, **sensitivity to new physics is increased**
- They generally feature **long-distance multi-hadron corrections**

📖 NH Christ, **AP**, *et al.* Phys Rev D 92(9) (2015)

📖 PA Boyle, **AP**, *et al.* Phys Rev D Lett 107(1) (2023)

📖 F Erben, **AP**, *et al.* JHEP 04 (2024)

📖 PA Boyle, **AP**, *et al.* arXiv:2406.19193 (2024)

Amplitude parameterisation

- $K^c \rightarrow \pi^c \gamma^*$ amplitude ($c \in \{0, +\}$)


$$\begin{aligned}\mathcal{A}_\mu^c(q^2) &= \int d^4x \langle \pi^c(\mathbf{p}) | T[J_\mu(0)H_W(x)] | K^c(\mathbf{k}) \rangle \\ &= -i \frac{G_F}{(4\pi)^2} [q^2(k+p)_\mu - (M_K^2 - M_\pi^2)q_\mu] V_c(z)\end{aligned}$$

$$\begin{aligned}q &= p - k \\ z &= q^2/M_K^2\end{aligned}$$

- Low-energy parameterisation

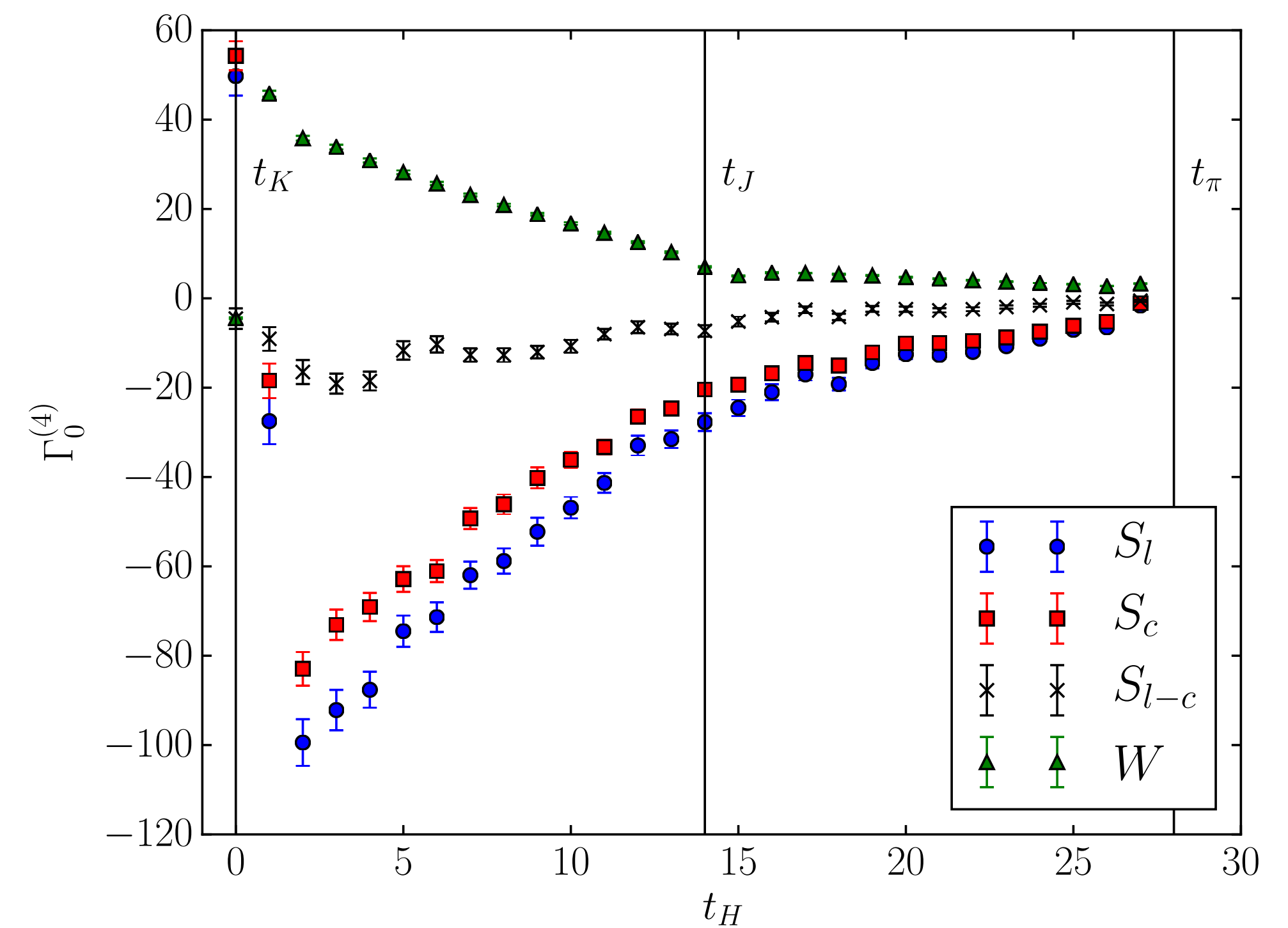
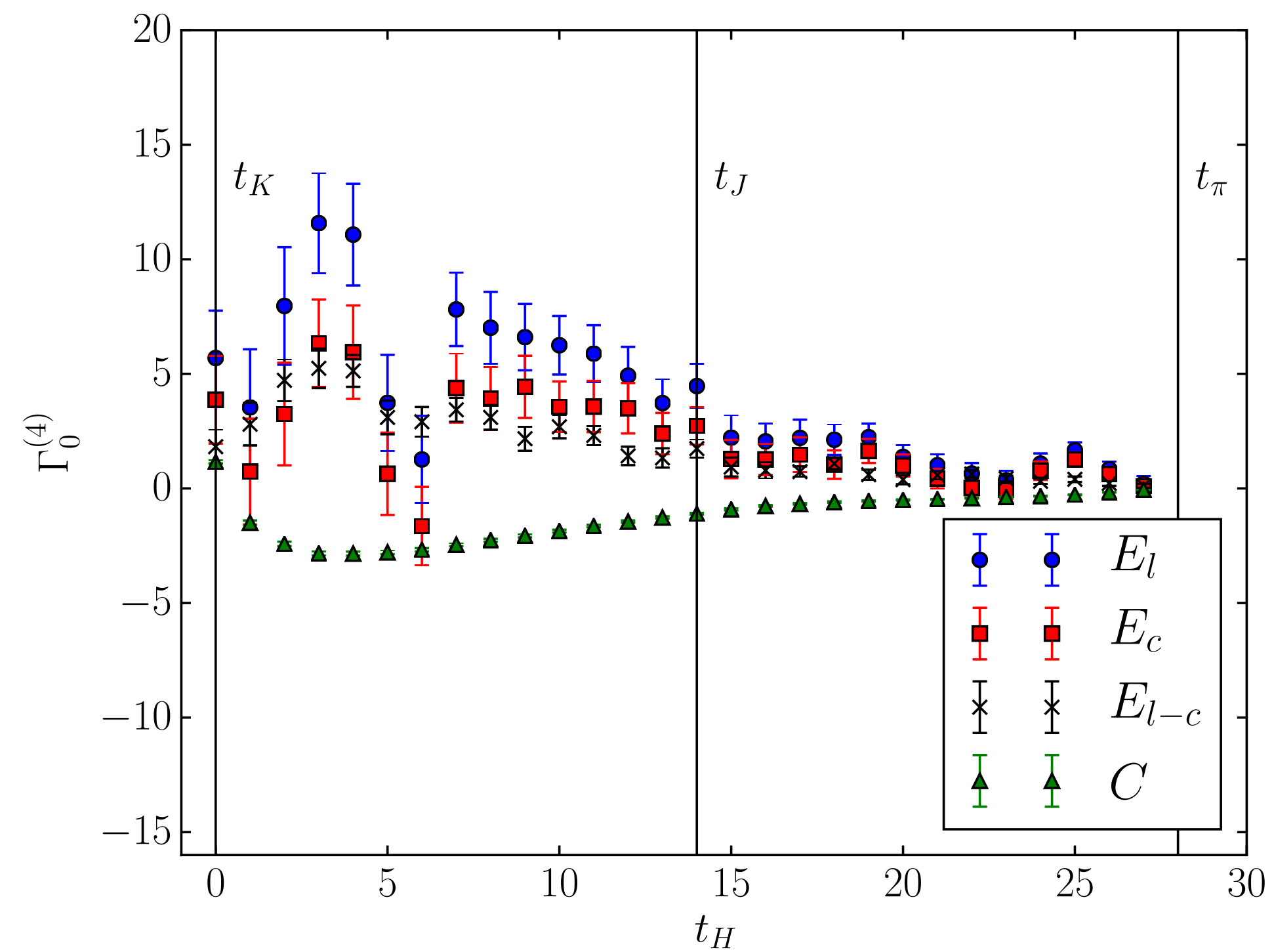
$$V_c(z) = a_c + b_c z + V_c^{\pi\pi}(z)$$

Analytical continuation issues

- **In Euclidean space-time, states below initial energy generate large contamination of Euclidean the time integral**
 - Happens potentially for $K \rightarrow \pi, \pi\pi, \pi\pi\pi \rightarrow \pi\gamma^*$
 - $\pi\pi$ forbidden in $K \rightarrow \pi\ell^+\ell^-$ (allowed in $K \rightarrow \pi\nu\bar{\nu}$)
 - Several subtraction strategies possible
-  NH Christ, AP, et al., PRD 92(9) 094512 (2015)

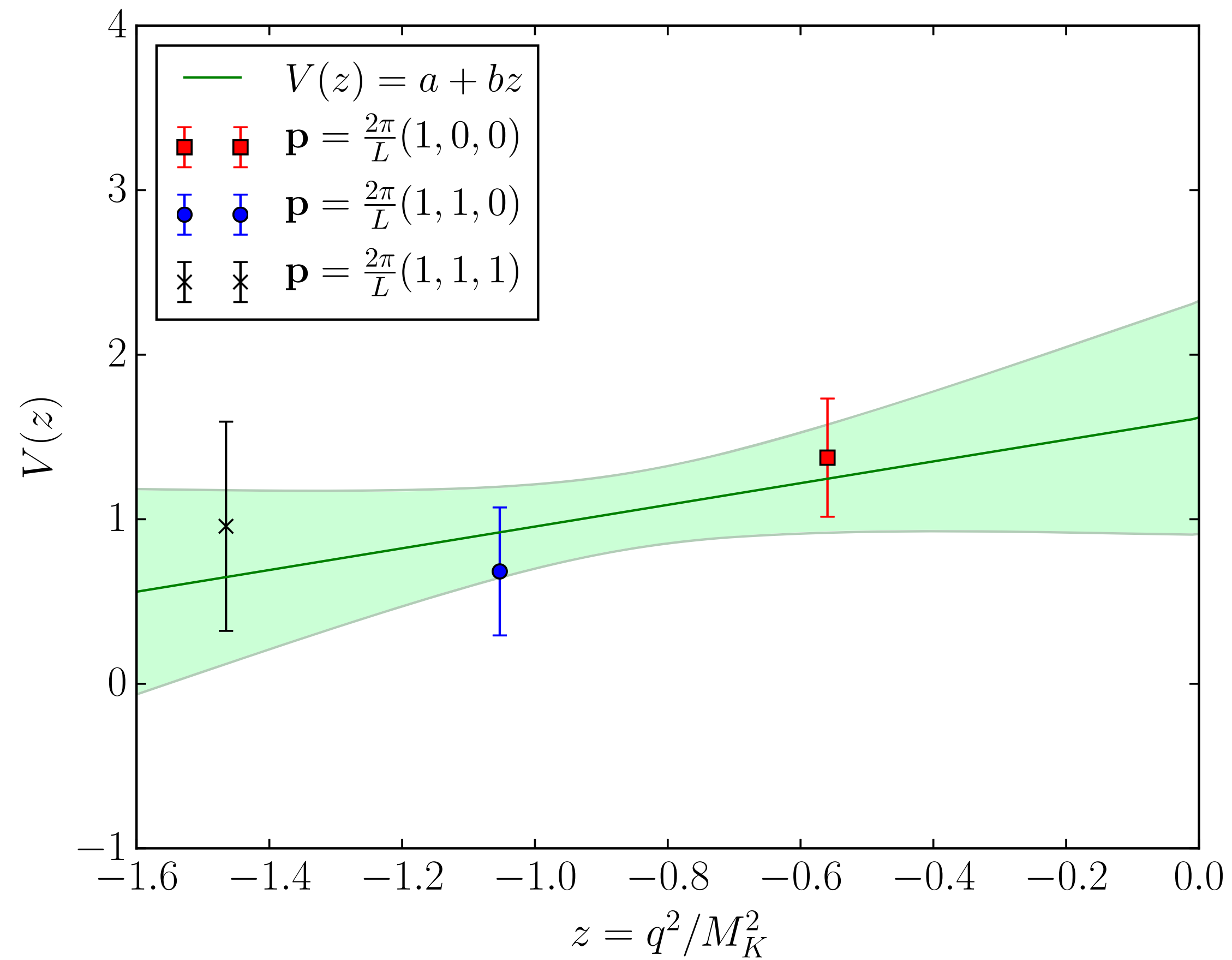
4-point correlators

Unphysical point  NH Christ, AP, et al., PRD 94(11) 114516 (2016)



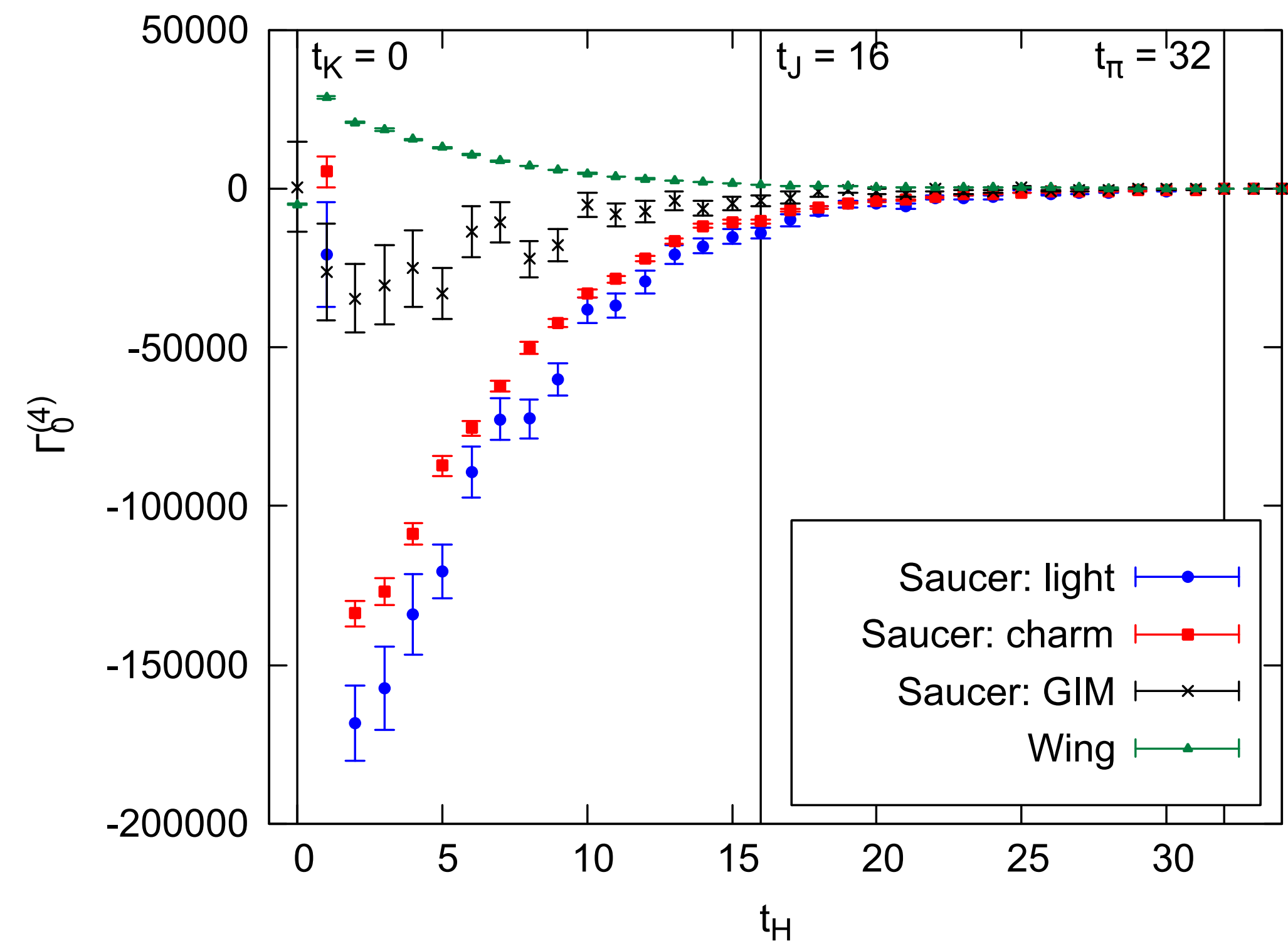
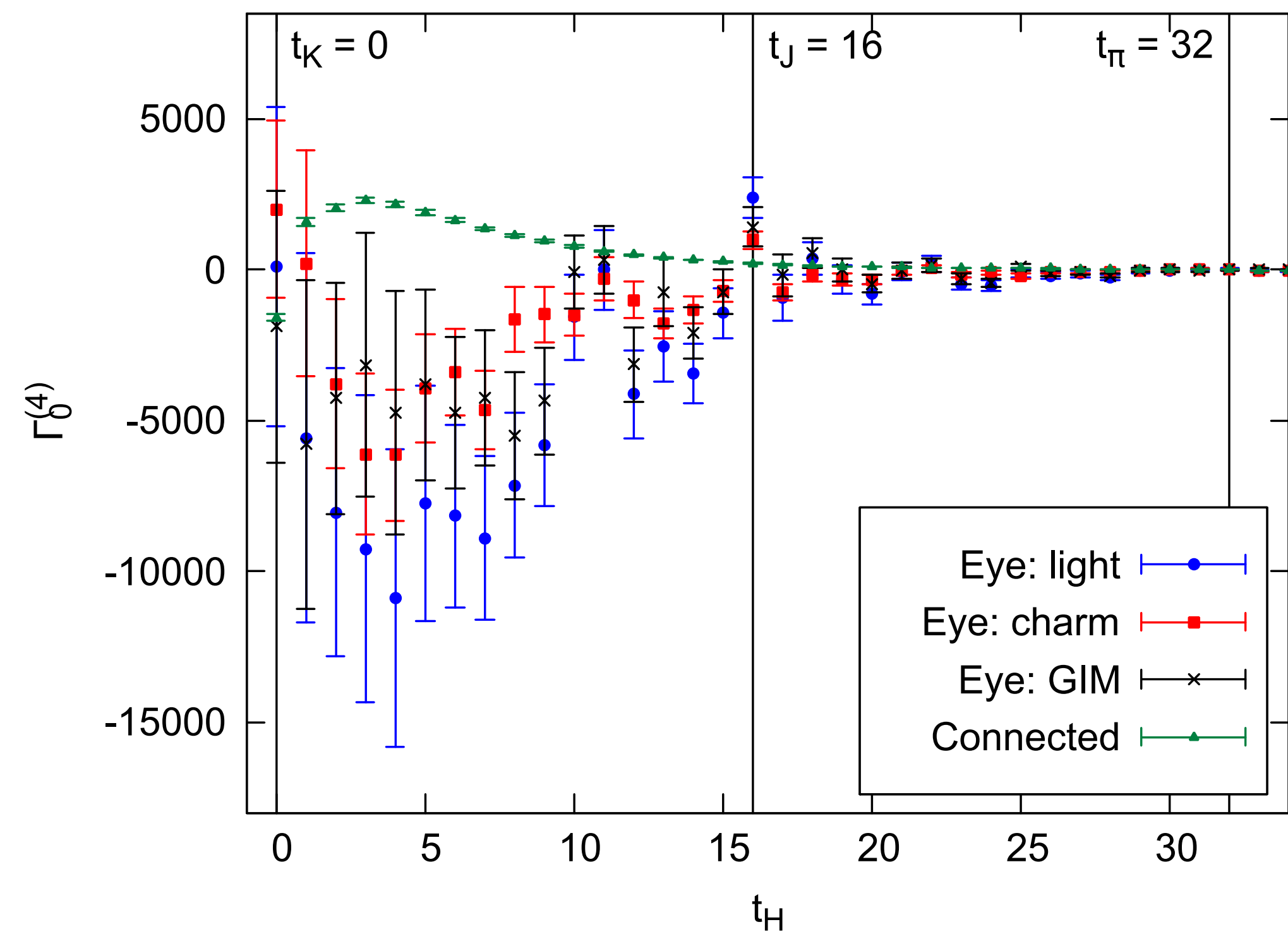
Amplitude result

Unphysical point  NH Christ, AP, et al., PRD 94(11) 114516 (2016)



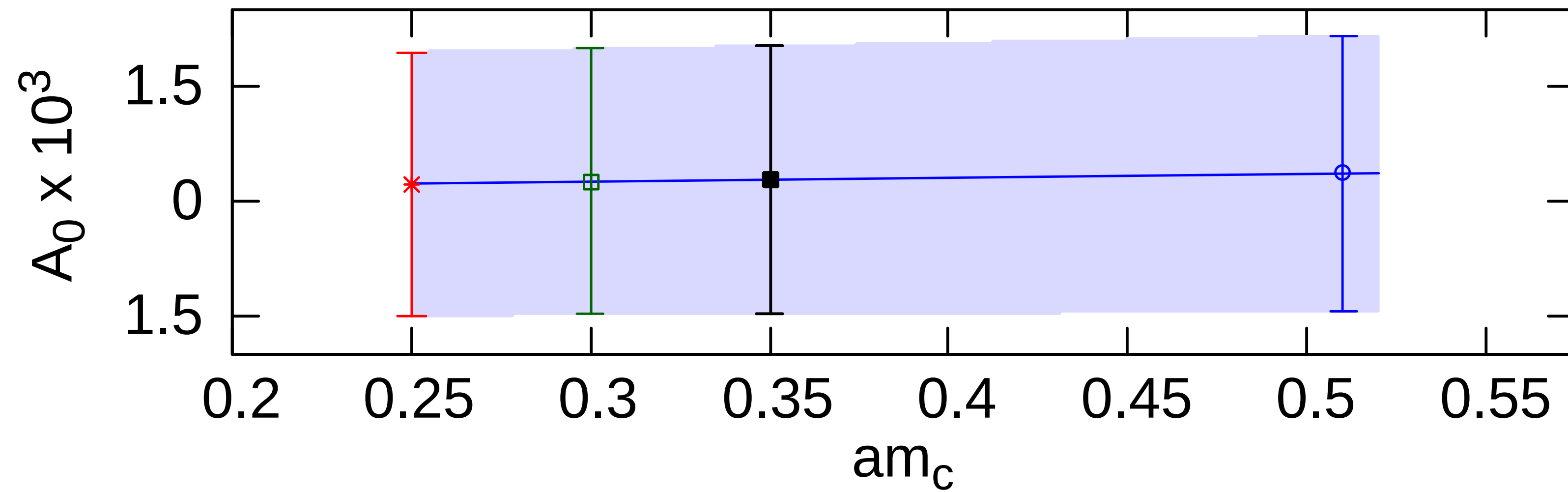
4-point correlators

Physical point  PA Boyle, AP, et al., PRD 107(1) L011503 (2022)



Amplitude result

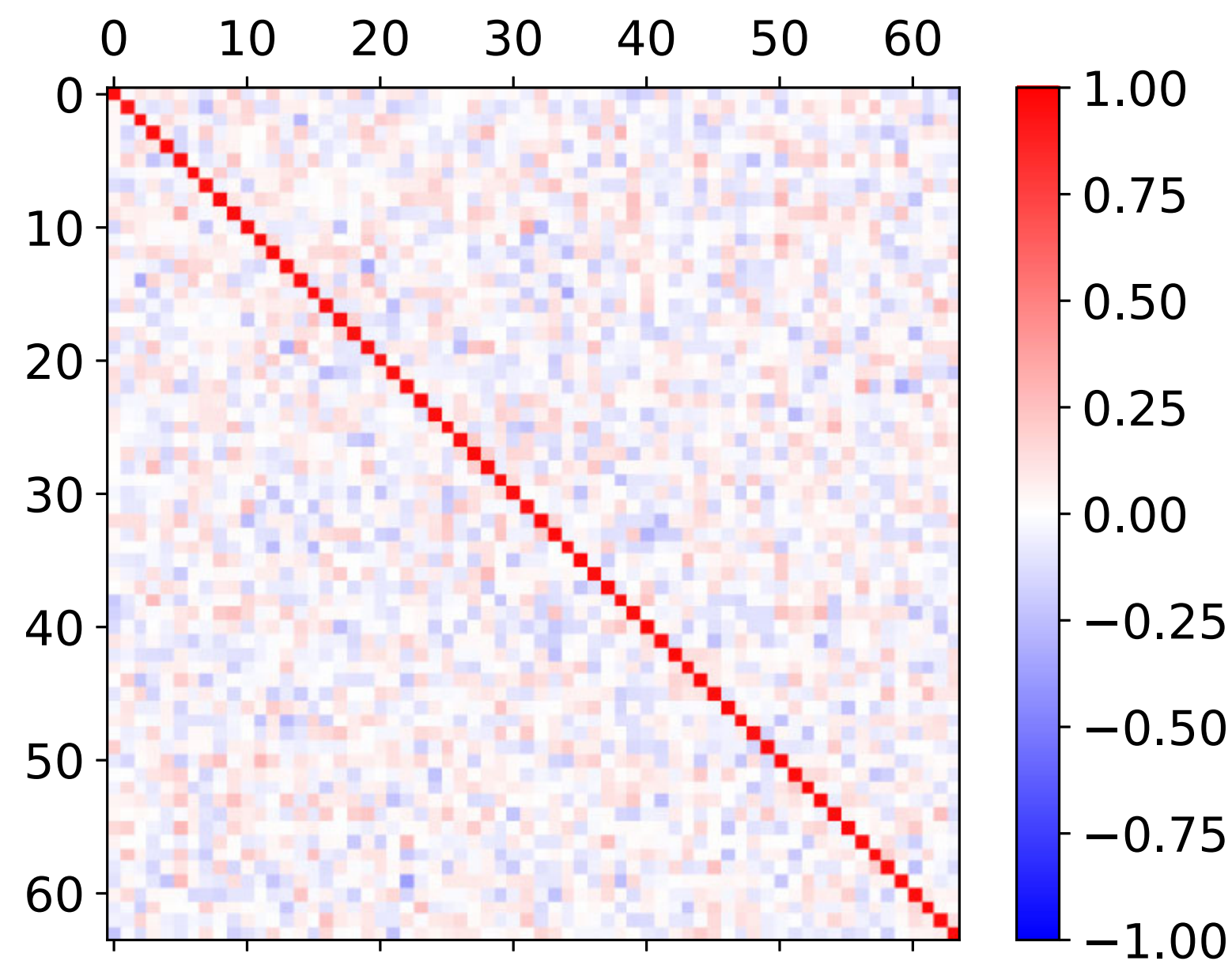
Physical point  PA Boyle, AP, et al., PRD 107(1) L011503 (2022)



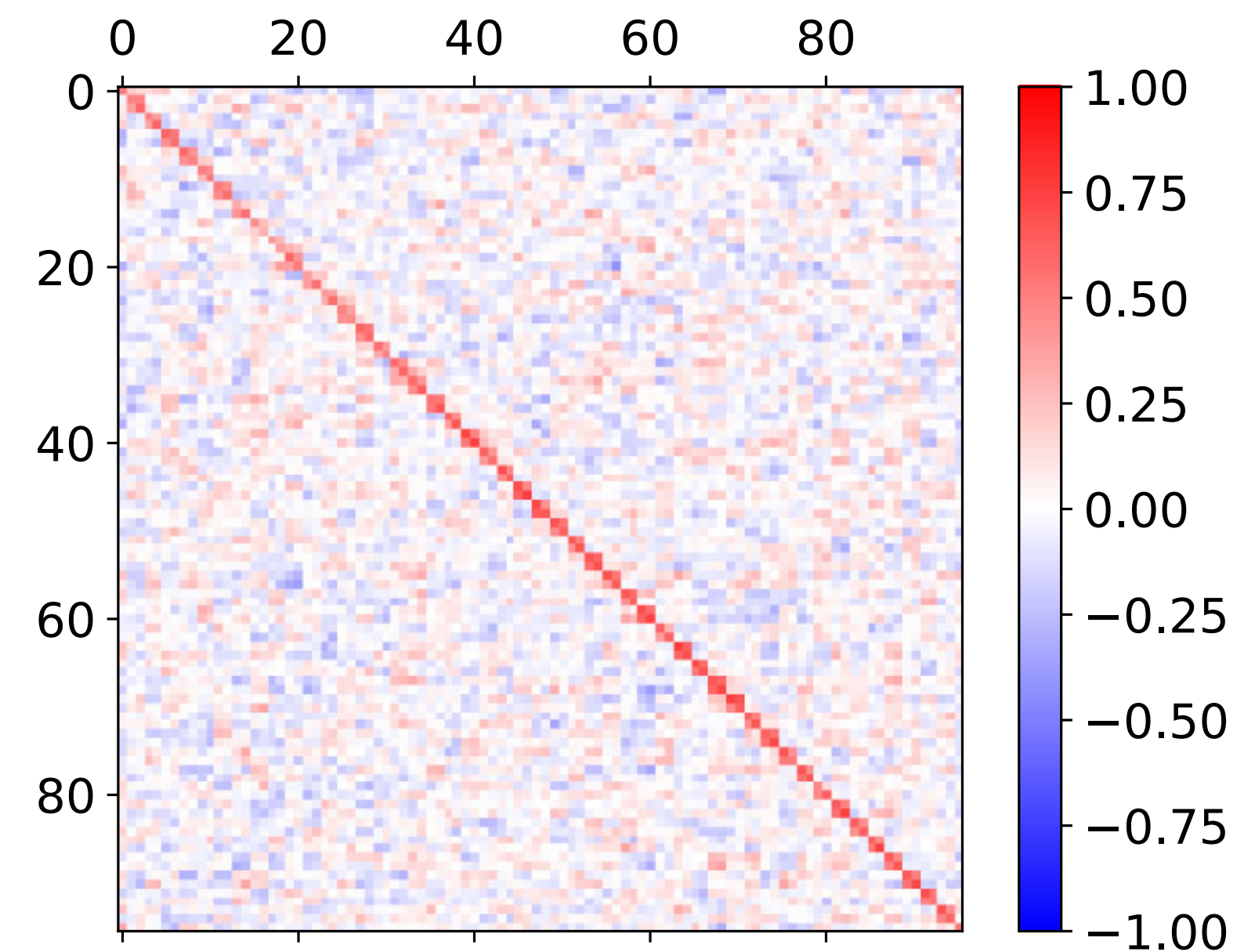
$$a_+ = -0.87(4.44)$$

A hint about the noise issue

- Correlations between up and charm loops is a **huge factor in GIM loops uncertainty**



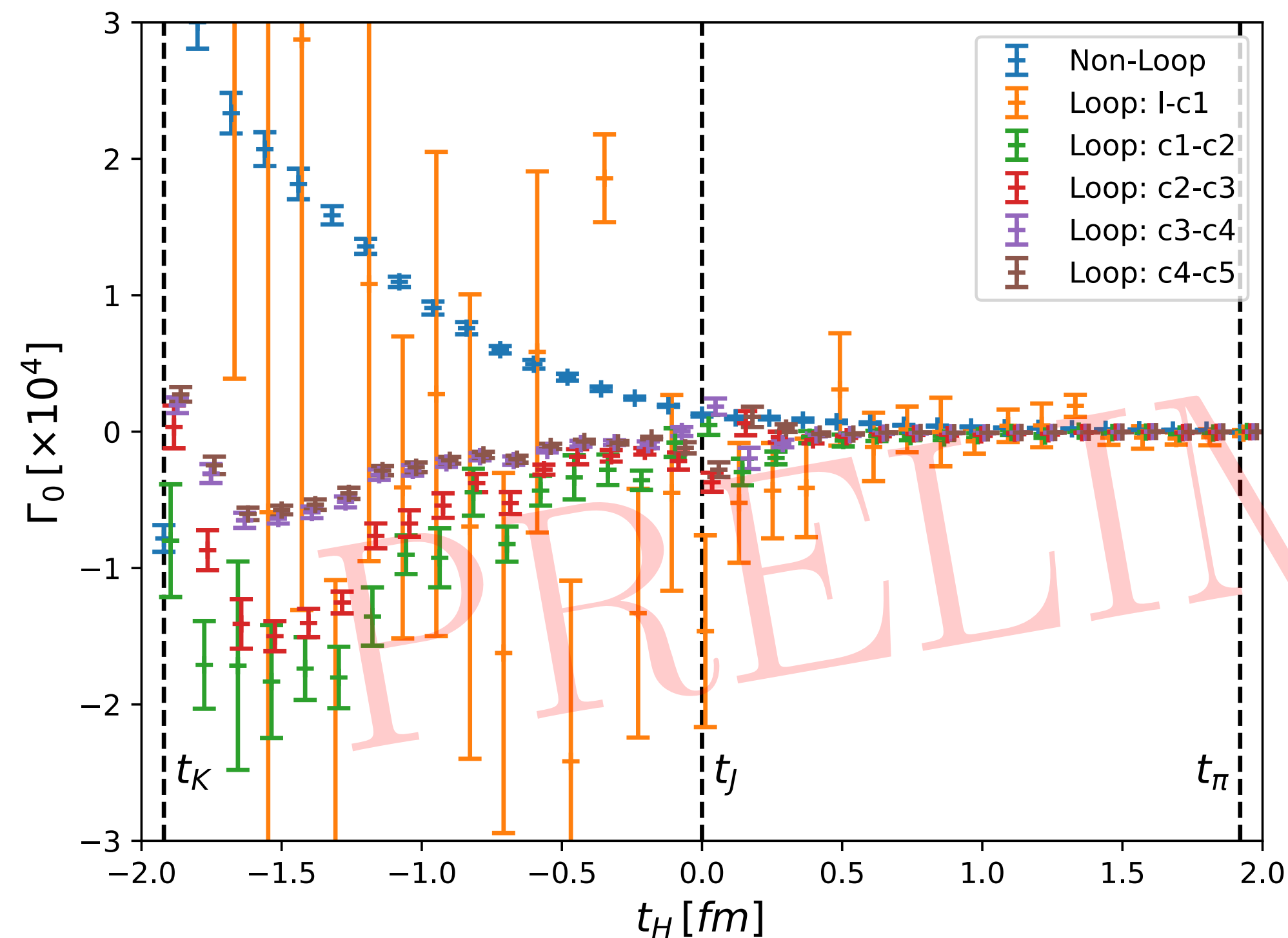
2016 unphysical data



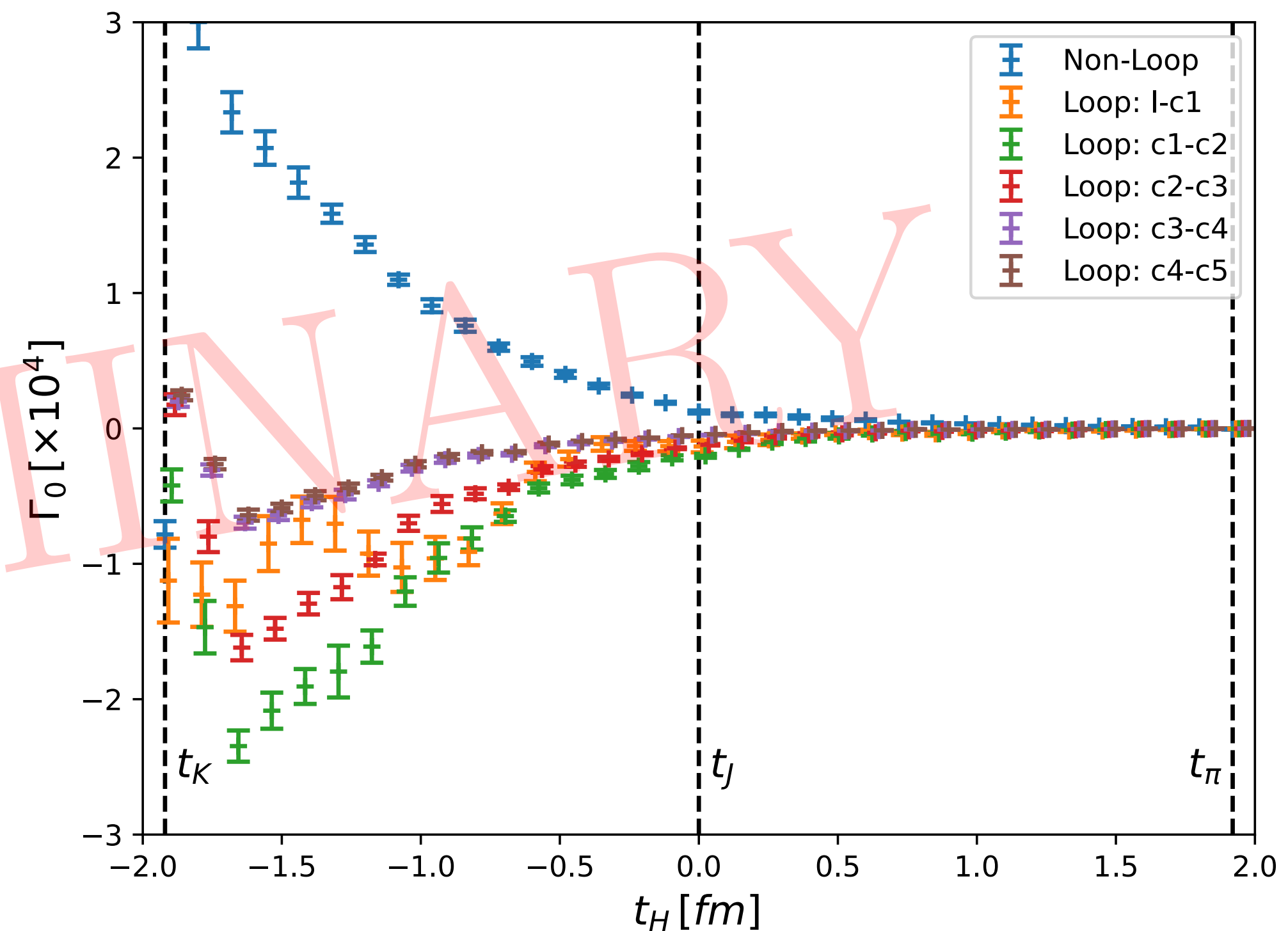
2022 physical point data

Improved estimators

4-point physical correlators / work with R Hill and R Hodsgon



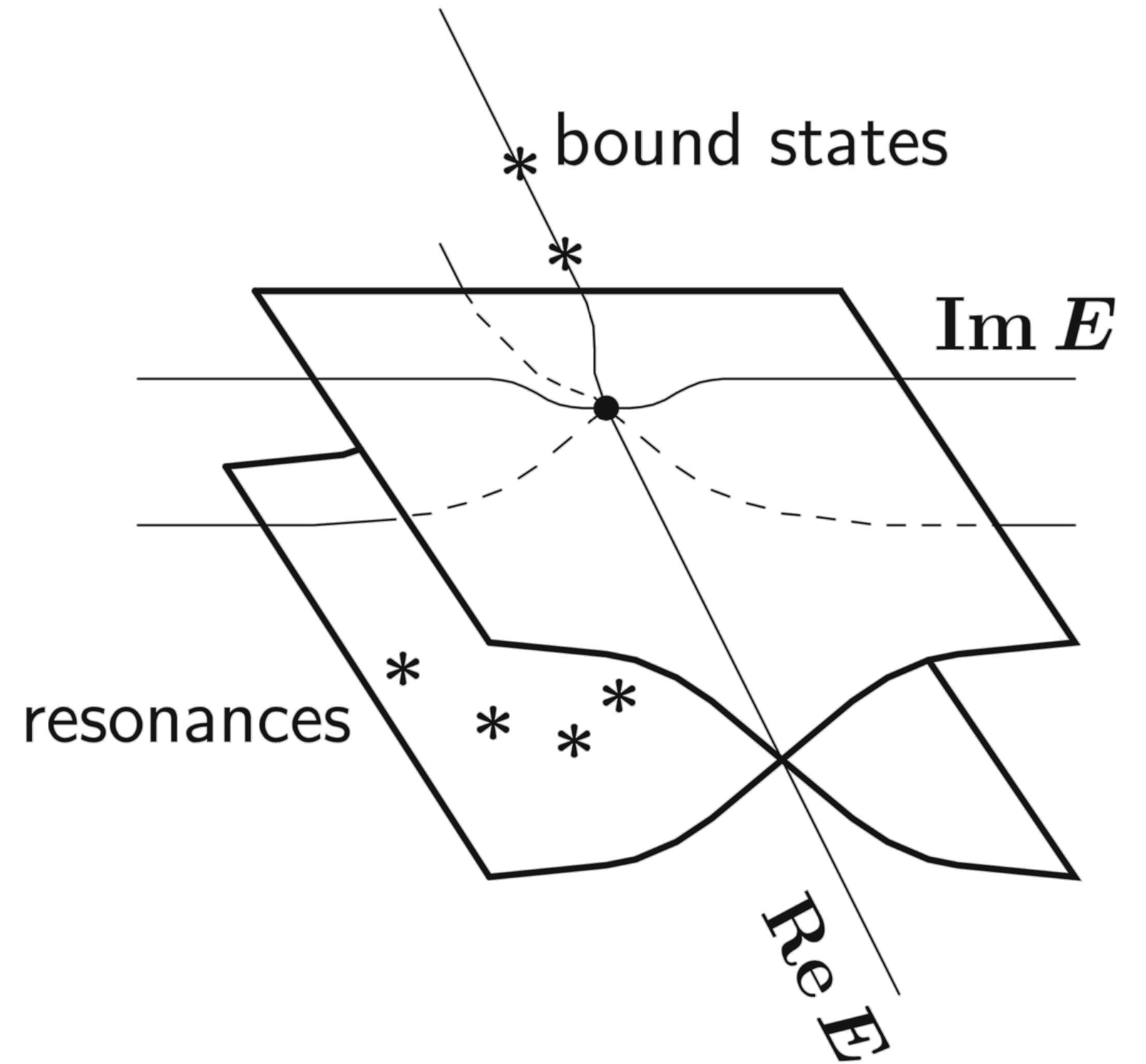
Standard estimator



Frequency-splitting estimator



 M Bordone, et al., to appear soon (Kaon@J-PARC 2024 Summary)

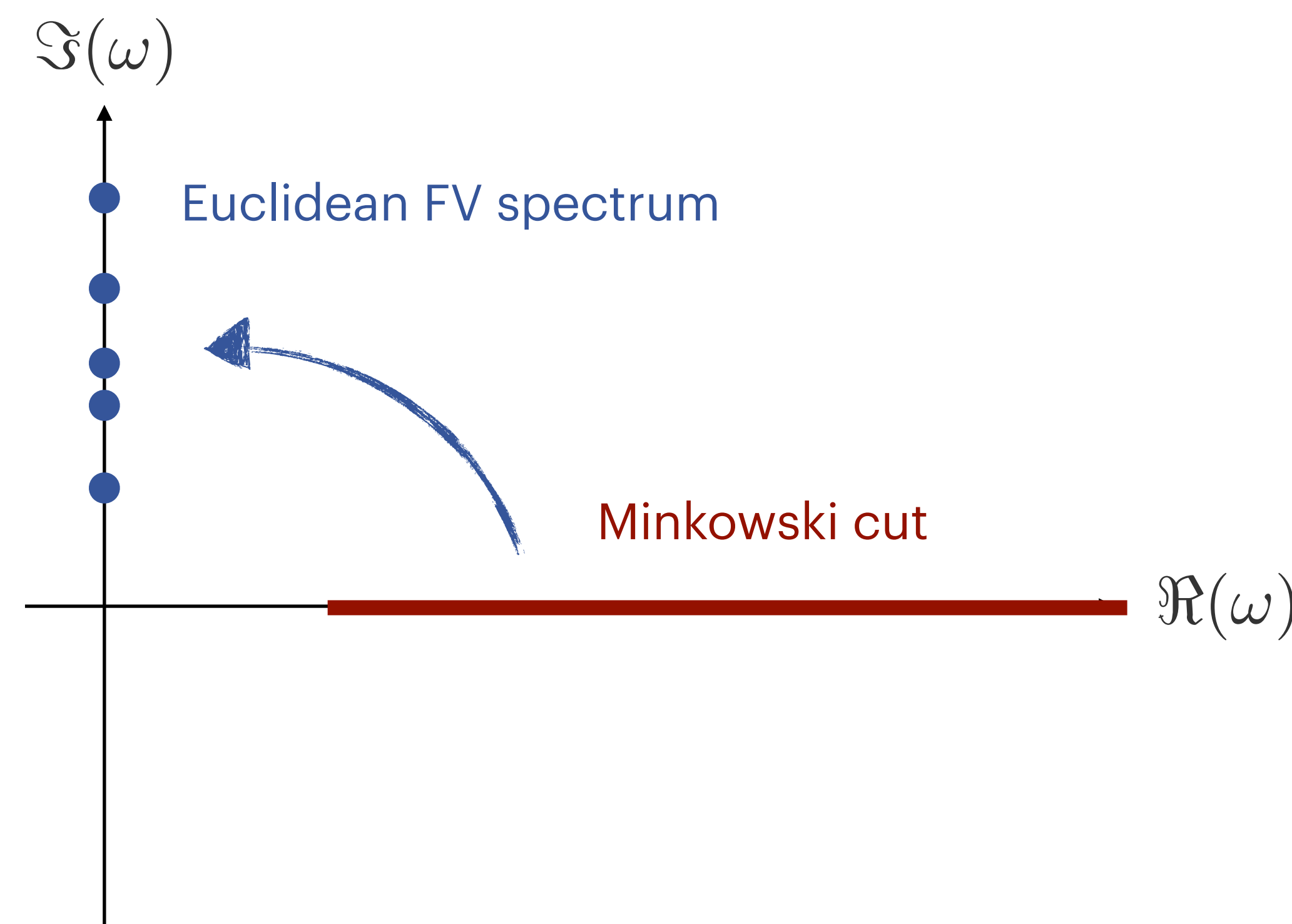
Multi-hadron interactions



Multi-hadron interactions

General issue for lattice simulations

- There is a theorem saying that hadronic scattering amplitudes cannot, in principle, be extracted from Euclidean field theory
 L Maiani & M Test, Phys Lett B 245 (1990)
- This was circumvented by noticing that splittings between **discrete energy levels in a finite volume encode information about scattering amplitudes**
- This is often referred as *Lüscher formalism*
 M Lüscher, Commun. Math. Phys. 105(2) (1986)



Hadronic resonances

Lattice determination

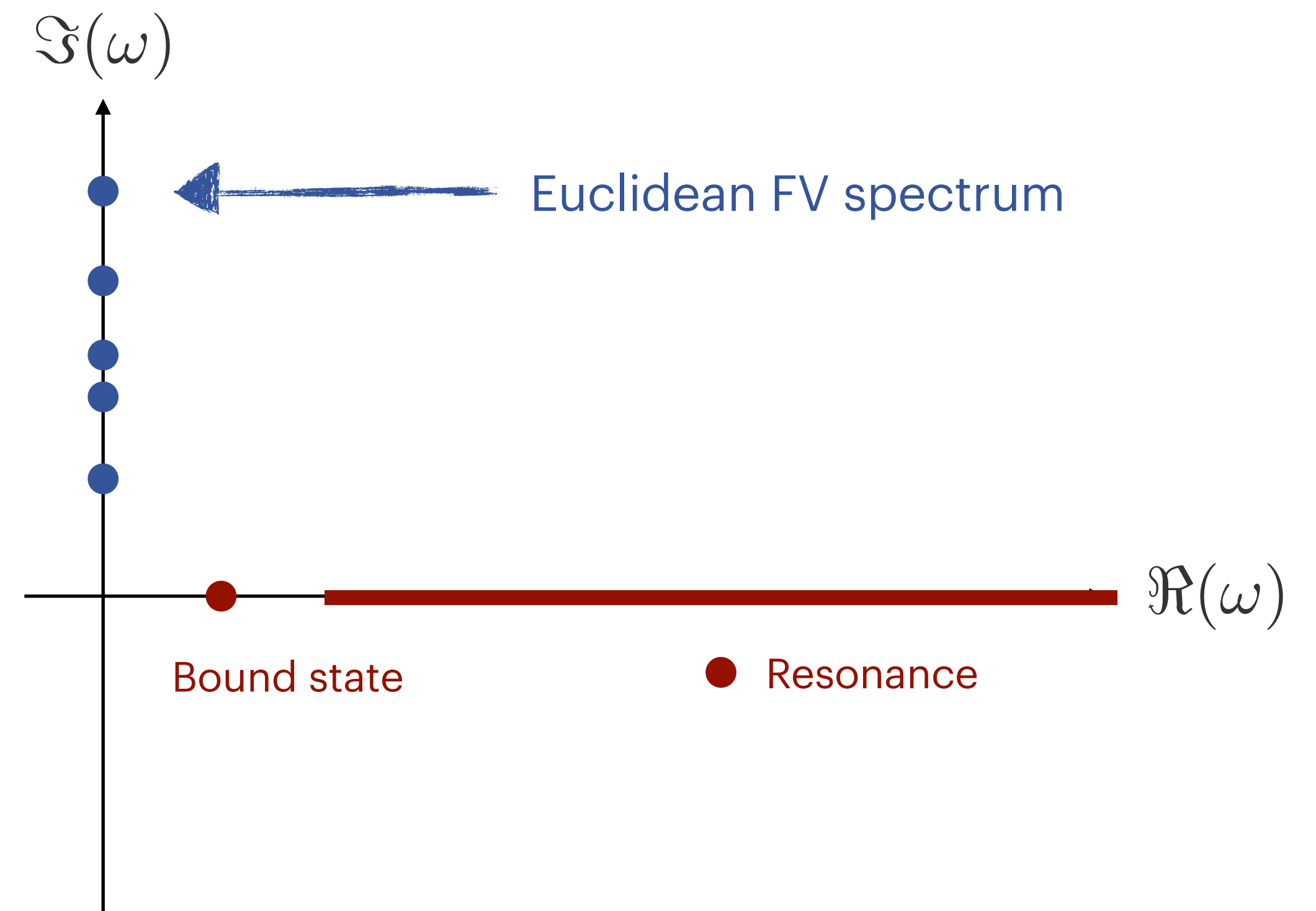
1. Determine **Euclidean FV energy levels**

2. Relate to phase-shift model using
Lüscher quantisation condition

 M Lüscher, Commun. Math. Phys. 105(2) (1986)

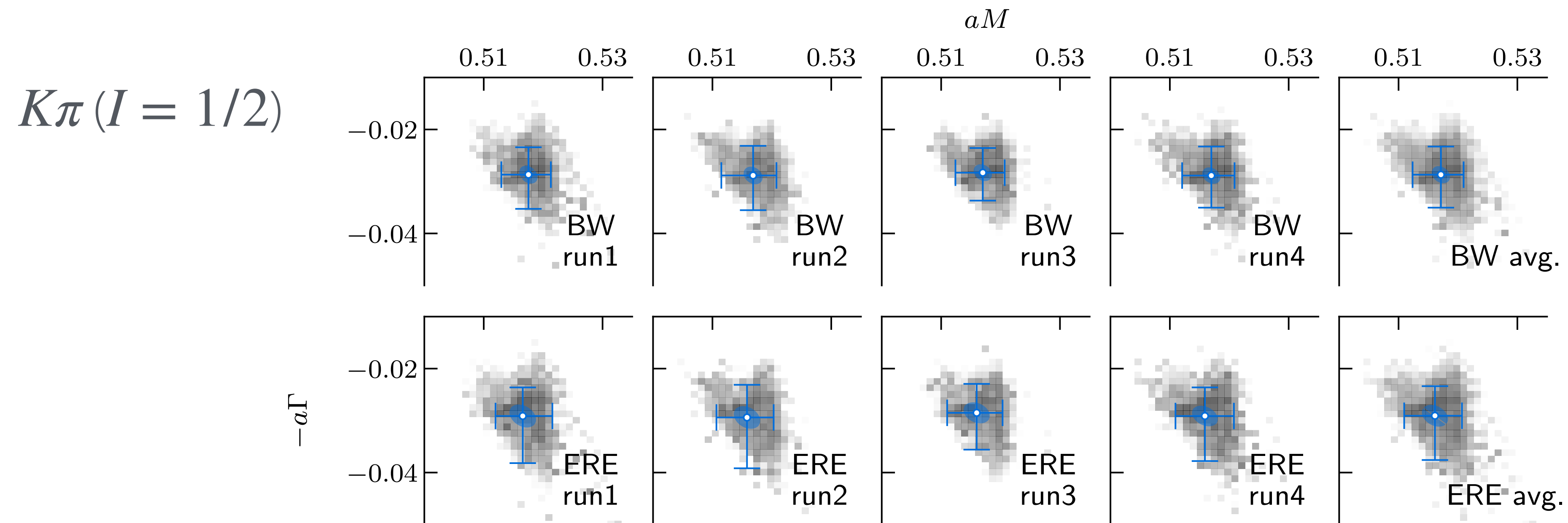
$$n\pi - \delta \left(\sqrt{\omega_n^2 - 4m^2} \right) = \phi \left(\frac{L}{2\pi} \sqrt{\omega_n^2 - 4m^2} \right)$$

3. **Solve for amplitude poles**



Lattice description of resonances

First physical point ρ & K^* simulation



- First determination of the ρ and K^* poles using **physical point 2+1 lattice simulations**
- First data-driven assessment of **analysis modelling and systematic errors**

PA Boyle, **AP**, et al. arXiv:2406.19193 (2024)

Conclusion

- Lattice QCD is entering the **era of physical, precise predictions** for
 - **Long-distance hadronic & electromagnetic** corrections to weak decays
 - Weak decays into **unstable states**
- More theoretical work on the way in key aspects e.g.
 - Treatment of **heavy quarks**
 - **Final state long-distance interactions**, particularly electromagnetic
- Hopefully crucial help for flavour physics measurements in future experiments

Thank you for your attention!