

# Rare charm decays at LHCb

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**FSP LHCb**  
Erforschung von  
Universum und Materie

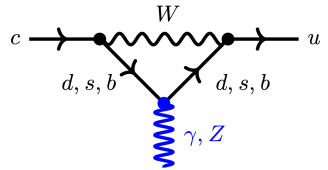
# Introduction: why do we study charm rare decays?

The  $c \rightarrow ul^+\ell^-$  transitions are **heavily suppressed**: may allow to **observe of New Physics**.

- ▶ Tree level **flavor-changing neutral-current** (FCNC) is not in the Standard Model (SM).
- ▶ Amplitude  $A(c \rightarrow u)$  depends on CKM elements ( $\lambda_i = V_{ci}^* V_{ui}$ ) and loop functions ( $f_i$ ):

$$A(c \rightarrow u) = \sum_{i=d,s,b} \lambda_i f_i = \lambda_s \left( (f_s - f_d) + \frac{\lambda_b}{\lambda_s} (f_b - f_d) \right)$$

**CKM suppression** and **Glashow-Iliopoulos-Maiani (GIM)** mechanism extremely suppress the branching fractions and CP asymmetries.



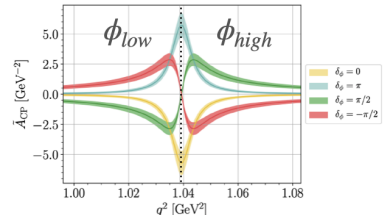
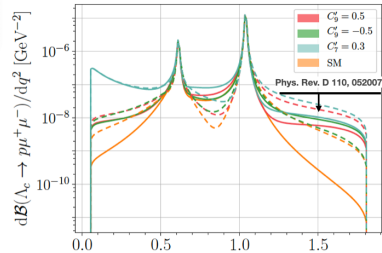
# Introduction: complementary approaches to probe New Physics

We have different strategies to probe New Physics:

- ▶ measure the branching fraction of the decays;
- ▶ NP sensitive in the **non-resonance region**;
- ▶ **CP and angular asymmetries (null tests)**;
- ▶ NP sensitive in the **resonance region**;
- ▶ **Lepton flavour universality tests**;

$$R_{P_1, P_2}^D = \frac{\int_{q_{min}^2}^{q_{max}^2} \frac{dB(D \rightarrow P_1 P_2 \mu^+ \mu^-)}{dq^2}}{\int_{q_{min}^2}^{q_{max}^2} \frac{dB(D \rightarrow P_1 P_2 e^+ e^-)}{dq^2}}$$

LHCb is providing **major contributions** in the field.  
Today we discuss the **most recent results** at LHCb.



Hiller et al, JHEP09(2021)208

# Search for $D^0 \rightarrow h^+ h^- e^+ e^-$ decays

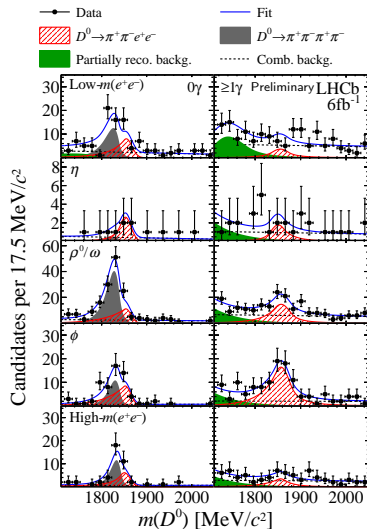
(*LHCb-PAPER-2024-047, in preparation*)



$$D^0 \rightarrow \pi^+ \pi^- e^+ e^-$$

- ▶ LHCb observed for the first time the  $D^0 \rightarrow \pi^+ \pi^- e^+ e^-$  in  $\rho/\omega$  and  $\phi$  dilepton mass regions.
- ▶ The current **best upper limits** on  $\mathcal{B}$  in other regions have been established.
- ▶ The  $\mathcal{B}$  in the electron mode is **compatible with the one of the muon mode**.

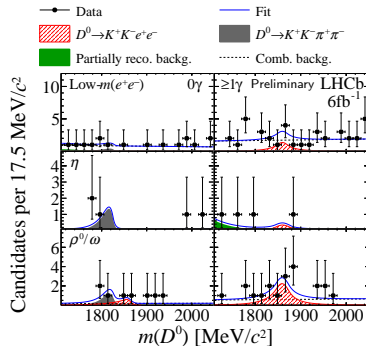
$m(e^+e^-)$ region	$D^0 \rightarrow \pi^+ \pi^- e^+ e^-$ [MeV/c <sup>2</sup> ]	$\mathcal{B}$ [10 <sup>-7</sup> ]
Low mass	211–525	< 4.8 (5.4)
$\eta$	525–565	< 2.3 (2.7)
$\rho^0/\omega$	565–950	$4.5 \pm 1.0 \pm 0.7 \pm 0.6$
$\phi$	950–1100	$3.8 \pm 0.7 \pm 0.4 \pm 0.5$
High mass	> 1100	< 2.0 (2.2)
Total	–	$13.3 \pm 1.1 \pm 1.7 \pm 1.8$



$$D^0 \rightarrow K^+ K^- e^+ e^-$$

- ▶ No evidences with the current precision.
- ▶ The current **best upper limits** on the  $\mathcal{B}$  have been established.

$D^0 \rightarrow K^+ K^- e^+ e^-$		
$m(e^+ e^-)$ region	[MeV/c <sup>2</sup> ]	$\mathcal{B}$ [10 <sup>-7</sup> ]
Low mass	211–525	< 1.0 (1.1)
$\eta$	525–565	< 0.4 (0.5)
$\rho^0/\omega$	> 565	< 2.2 (2.5)

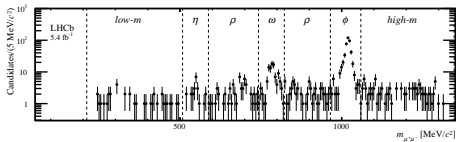


# Search for $\Lambda_c^+ \rightarrow p\mu^+\mu^-$ decays

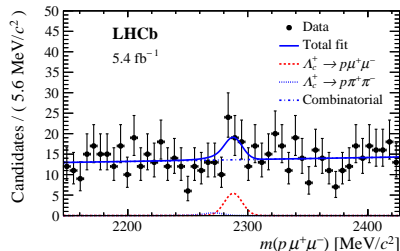
(*Phys. Rev. D* 110, 052007)

Study the **resonant and non resonant dimuon regions** with LHCb RUN2 ( $5.4 \text{ fb}^{-1}$ ).

- ▶ Measure the  $\mathcal{B}$ :
  - ▶  $\rho$ ,  $\omega$  and  $\eta$  resonant regions;
- ▶ **Search in the non resonant region;**
- ▶ uses the  $\phi$  resonant region as **normalization.**



**No signal evidence in the non-resonant region:**



- ▶ The **best upper limits** on the  $\mathcal{B}$  has been established:

$$\mathcal{B} \left( \Lambda_c^+ \rightarrow p\mu^+\mu^- \right) < 2.9(3.2) \times 10^{-8} \text{ at } 90\%(95\%) \text{ CL.}$$

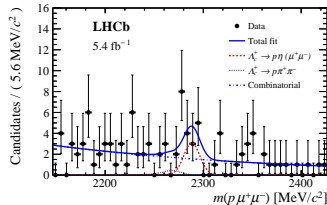
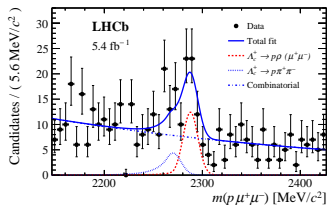
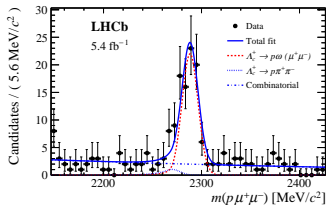


- The  $\mathcal{B}$  in the **resonant regions** have been measured:

$$\Lambda_c^+ \rightarrow p\rho(\rightarrow \mu^+\mu^-) > 7\sigma$$

$$\Lambda_c^+ \rightarrow p\omega(\rightarrow \mu^+\mu^-) > 5.6\sigma$$

$$\Lambda_c^+ \rightarrow p\eta(\rightarrow \mu^+\mu^-) > 3.0\sigma$$



$$\mathcal{B}(\Lambda_c^+ \rightarrow p\rho) = (1.52 \pm 0.34(\text{stat}) \pm 0.14(\text{syst}) \pm 0.24(\text{ext})) \times 10^{-3}$$

$$\mathcal{B}(\Lambda_c^+ \rightarrow p\omega) = (9.82 \pm 1.23(\text{stat}) \pm 0.73(\text{syst}) \pm 2.79(\text{ext})) \times 10^{-4}$$

$$\mathcal{B}(\Lambda_c^+ \rightarrow p\eta) = (1.67 \pm 0.69(\text{stat}) \pm 0.23(\text{syst}) \pm 0.34(\text{ext})) \times 10^{-3}$$

# Search for resonance enhanced asymmetries on the $\Lambda_c^+ \rightarrow p\mu^+\mu^-$ channel

*(LHCb-PAPER-2024-051, in preparation)*



## CP asymmetry

We cannot access  $A_{CP}$  directly, we only can measure the so called raw-asymmetry ( $A_{CP}^{raw}$ ):

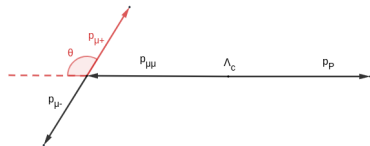
$$A_{CP}^{raw} = \frac{N(\Lambda_c^+) - N(\bar{\Lambda}_c^-)}{N(\Lambda_c^+) + N(\bar{\Lambda}_c^-)} = A_{CP} + A_P(\Lambda_c^+) + A_D(p) + (A_D(\mu^+\mu^-) = 0)$$

- ▶ a control sample of  $\Lambda_c^+ \rightarrow pK_s^0$  is used to estimate the nuisance asymmetries:

## FB asymmetry

The  $A_{FB}$  is defined with respect to  $\theta$ -angle:

$$A_{FB} = \frac{N(\cos \theta > 0) - N(\cos \theta < 0)}{N(\cos \theta > 0) + N(\cos \theta < 0)}$$



It is also possible to evaluate the relevant combinations:

$$\Sigma A_{FB} = \frac{1}{2} (A_{FB}^{\Lambda_c^+} + A_{FB}^{\bar{\Lambda}_c^-})$$

$$\Delta A_{FB} = \frac{1}{2} (A_{FB}^{\Lambda_c^+} - A_{FB}^{\bar{\Lambda}_c^-})$$

The detector and the offline selection may have different efficiencies in different regions of the  $(m_{\mu\mu}, \cos\theta)$ -phase space.

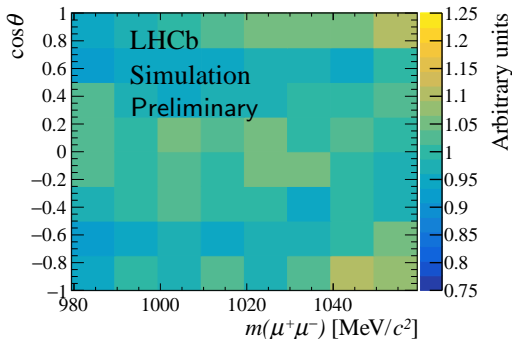
The distributions **at generation level** and **after the offline selection** are compared.

The **relative efficiencies** in the phase space are:

$$\epsilon(m_{\mu\mu}, \cos\theta) = \frac{f_{\text{selected}}(m_{\mu\mu}, \cos\theta)}{f_{\text{generated}}(m_{\mu\mu}, \cos\theta)}$$

and the **acceptance weights**:

$$\lambda_{\text{corr}} = \frac{1}{\epsilon}$$



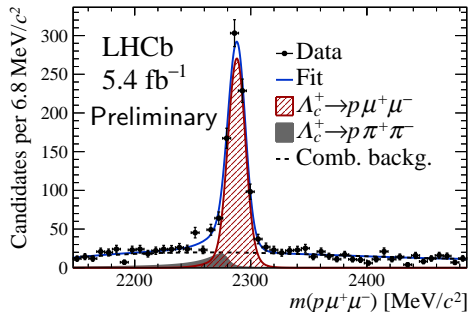
A total model to fit the **final selected mass distribution** has been studied:

$$f_{tot}(m; N_1, \dots, N_n) = \sum_{i=1}^n N_i f_i(m; \vec{x}_i)$$

The shapes of the different components ( $\vec{x}_i$ ) have been **constrained by previous fits** on simulation and data control samples.

The **yields from the total fit** are:

$m(\mu^+\mu^-)$	Efficiency-weighted yields		
	Signal	Misid. back.	Comb. back.
$\phi_{low}$	$346 \pm 22$	$57 \pm 21$	$437 \pm 26$
$\phi_{high}$	$435 \pm 22$	$35 \pm 17$	$390 \pm 25$



Finally the **values of the null tests**:

$m(\mu^+\mu^-)$	$A_{CP}[\%]$	$A_{FB}^{\Lambda_c^+}[\%]$	$A_{FB}^{\bar{\Lambda}_c^-}[\%]$	$\Sigma A_{FB}[\%]$	$\Delta A_{FB}[\%]$
$\phi_{\text{low}}$	$-0.8 \pm 6.2 \pm 0.6$	$11.7 \pm 8.5 \pm 1.1$	$2.2 \pm 8.7 \pm 1.4$	$6.9 \pm 6.1 \pm 1.0$	$4.8 \pm 6.1 \pm 0.8$
$\phi_{\text{high}}$	$-1.4 \pm 5.3 \pm 0.6$	$3.5 \pm 7.2 \pm 0.9$	$-0.3 \pm 7.4 \pm 1.1$	$1.6 \pm 5.2 \pm 0.8$	$1.9 \pm 5.2 \pm 0.6$

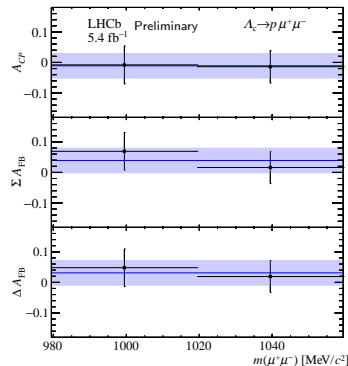
We also **combine the two dilepton mass bins**:

$$A_{CP} = -1.1 \pm 4.0 \pm 0.5\%$$

$$\Sigma A_{FB} = +3.9 \pm 4.0 \pm 0.6\%$$

$$\Delta A_{FB} = +3.1 \pm 4.0 \pm 0.4\%$$

- ▶ **No significant deviation** from the SM are observed.
- ▶ At this precision the  $\Lambda_c^+ \rightarrow p\mu^+\mu^-$  is a competitive channel for **setting constraints to BSM physics** (Hiller et al, arXiv 2410.00115).



# Conclusions and outlook

## Conclusions

- ▶ **Rare charm decays** are a powerful probe to investigate New Physics, complementary to the beauty-sector.
- ▶ LHCb is giving a **major contribution** in the field:
  - ▶  $D^0 \rightarrow h^+ h^- e^+ e^-$ : **first observation** in  $D^0 \rightarrow \pi^+ \pi^- e^+ e^-$  and **world's best upper limit** in  $D^0 \rightarrow K^+ K^- e^+ e^-$  decays;
  - ▶  $\Lambda_c^+ \rightarrow p \mu^+ \mu^-$ : **world's best upper limit** in non resonant region;
  - ▶  $\Lambda_c^+ \rightarrow p \mu^+ \mu^-$ : **first null test measurement** in the  $\phi$ -resonant region.

## Outlook

- ▶ The electron mode will allow **future LFU tests**.
- ▶  $\Lambda_c^+ \rightarrow p \mu^+ \mu^-$  is a powerful channel to **probe New Physics**.
- ▶ LHCb new data, with **improved trigger efficiency**, will provide **higher statistic dataset**.