Asymmetric dark matter semi-annihilation into long-lived particles

Based on 2412.01470

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DISCRETE 2024 in Ljubljana, Dec 2-6, 2024

Introduction - semi-annihilating dark matter



$$\frac{\mathrm{d}n_{\chi}}{\mathrm{d}t} + 3Hn_{\chi} = -\frac{1}{2}[n_{\chi}^2 - n_{\chi}n_{\chi}^{\mathrm{eq}}]\langle\sigma v\rangle_{\chi\chi\to\chi^{\dagger}\phi} \tag{1}$$

[D'Eramo, Thaler, '10]

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Asymmetric semi-annihilation in A. Ghosh, D Ghosh, Mukhopadhyay '20

$$\mathcal{L} \supset \frac{\mu}{3!} \chi^3 + \frac{\lambda}{3!} \phi \chi^3 + \text{H.c.} + \frac{\lambda_1}{4} |\chi|^4 + \frac{\lambda_2}{2} \phi^2 |\chi|^2 + \mu_1 \phi |\chi|^2 + \frac{\mu_2}{3!} \phi^3 + \frac{\lambda_3}{4!} \phi^4.$$
(2)

Dark matter asymmetry via semi-annihilations



[A. Ghosh, D Ghosh, Mukhopadhyay '20]

CP asymmetry and unitarity relations

$$SS^{\dagger} = S^{\dagger}S = 1 \quad \rightarrow \quad S^{\dagger} = 1 - iT^{\dagger} = (1 + iT)^{-1} \tag{3}$$

$$|T_{fi}|^{2} = -iT_{if}^{\dagger}iT_{fi} = -iT_{if}iT_{fi} + \sum_{n}iT_{in}iT_{nf}iT_{fi} - \dots$$
(4)

[Coster, Stapp '70; Bourjaily, Hannesdottir, et al. '21; Hannesdottir, Mizera '22; Blažek, Maták '21a]

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$$\Delta |T_{fi}|^{2} = \sum_{n} (iT_{in}iT_{nf}iT_{fi} - iT_{if}iT_{fn}iT_{ni})$$

$$-\sum_{n,m} (iT_{in}iT_{nm}iT_{mf}iT_{fi} - iT_{if}iT_{fm}iT_{mn}iT_{ni})$$

$$+ \dots$$
[Blažek, Maták '21a]

$$\sum_{f} \Delta |T_{fi}|^2 = 0 \tag{6}$$

[Dolgov '79; Kolb, Wolfram '80;

Additional ϕ_2 field enables non-vanishing asymmetry

$$\mathcal{L} \supset -\frac{\lambda_1}{6}\chi^3\phi_1 - \frac{\lambda_2}{6}\chi^3\phi_2 + \text{H.c.} - \lambda_{12}|\chi|^2\phi_1\phi_2 - \frac{\lambda_3}{2}|\chi|^2\phi_3^2.$$
(7)

- Strong annihilation into visible sector represented by ϕ_3 field
- Long-lived ϕ_1, ϕ_2 particles

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Asymmetric semi-annihilations from



Boltzmann equations for the asymmetry

$$\left(\frac{\mathrm{d}\Delta}{\mathrm{d}x}\right)_{\mathrm{source}} = -\frac{3}{8} \frac{s}{Hx} Y_0^{\mathrm{eq}} Y \left(\frac{Y_{\phi_1}}{Y_{\phi_1}^{\mathrm{eq}}} - \frac{Y_{\phi_2}}{Y_{\phi_2}^{\mathrm{eq}}}\right) \Delta \langle \sigma v \rangle_{\chi\chi \to \chi^{\dagger} \phi_1}$$

$$\left(\frac{\mathrm{d}\Delta}{\mathrm{d}x}\right)_{\mathrm{washout}} = -\frac{3}{2} \frac{s}{Hx} Y_0^{\mathrm{eq}} \Delta \sum_{i=1,2} \left(\frac{Y}{Y_0^{\mathrm{eq}}} + \frac{1}{2} \frac{Y_{\phi_i}}{Y_{\phi_i}^{\mathrm{eq}}}\right) \langle \sigma v \rangle_{\chi\chi \to \chi^{\dagger} \phi_i}$$

$$(9)$$

$$Y = Y_{\chi} + Y_{\bar{\chi}} = \frac{n_{\chi} + n_{\bar{\chi}}}{s} \quad \text{and similarly for } \phi_1 \text{ and } \phi_2 \tag{10}$$

$$\frac{Y_{\phi_1}}{Y_{\phi_1}^{\text{eq}}} = \frac{Y_{\phi_2}}{Y_{\phi_2}^{\text{eq}}} \implies \text{no asymmetry}$$
(11)



 $m = 200 \text{ GeV}, \ m_{\phi_2} = 400 \text{ GeV}, \ m_{\phi_1} = 0 \ (\text{left}), \ m_{\phi_1} = 260 \text{ GeV} \ (\text{right})$

Asymmetric semi-annihilations - fermionic model

Effective dimension-6 operators

$$\mathcal{L} \supset -\frac{\kappa_1}{6\Lambda^2} (\bar{\chi}^c P_R \chi) (\bar{\chi}^c P_L \psi_1) - \frac{\kappa_2}{6\Lambda^2} (\bar{\chi}^c P_R \chi) (\bar{\chi}^c P_L \psi_2) - \frac{\kappa_{12}}{2\Lambda^2} (\bar{\chi}^c P_L \psi_1)^{\dagger} (\bar{\chi}^c P_L \psi_2) + \text{H.c.}$$

Asymmetric semi-annihilations from



 $\sin \delta_1 = \operatorname{Im}[\kappa_1^* \kappa_{12}^* \kappa_2] / |\kappa_1^* \kappa_{12}^* \kappa_2|$

 $\sin \delta_2 = \operatorname{Im}[\kappa_1^* \kappa_{12} \kappa_2] / |\kappa_1^* \kappa_{12} \kappa_2|$

Asymmetric semi-annihilations - fermionic model



 $|\kappa_1| = |\kappa_2| = |\kappa_{12}| = 1$

Summary

- Even feeble asymmetric semi-annihilation can have substantial effect on dark matter relic density
- \bullet Semi-annihilations into long-lived particles $~\to~$ dark matter totally asymmetric after freeze-out

[Blažek, Maták, Zaujec, '24]

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[Blažek, Maták, Zaujec, '24]

Thank you!