
CP-violating portal to Dark Sectors

Nicola Valori

University of Valencia & IFIC

DISCRETE 2024
Ljubljana 04/12/2024

Based on **JHEP11(2024)049** with **M. Ardu, M.H. Rahat, O.Vives**

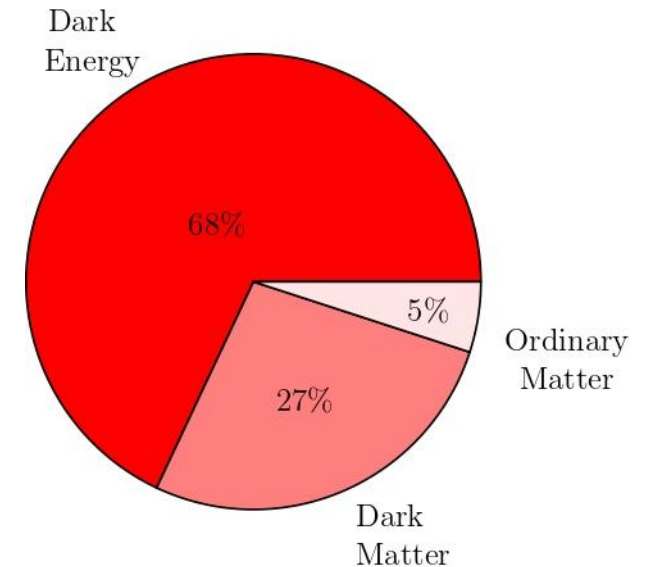


Motivation

Motivation

Particle Dark Matter:

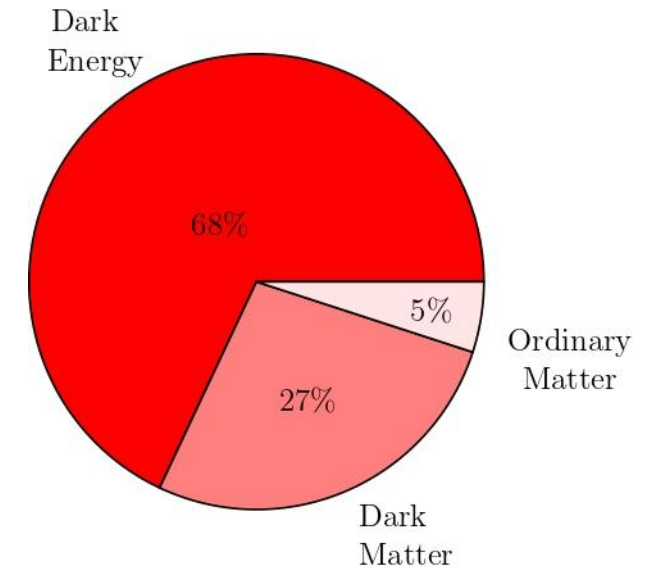
- **Dark Matter** is almost $\frac{1}{4}$ of the whole energy budget.
- **Dark Matter Production** (usually) requires interaction with SM.
- If $DM \in DS$: **Portals** between the visible and dark sector.



Motivation

Particle Dark Matter:

- **Dark Matter** is almost $\frac{1}{4}$ of the whole energy budget.
- **Dark Matter Production** (usually) requires interaction with SM.
- If $DM \in DS$: **Portals** between the visible and dark sector.



CP-violation:

- The SM of particle physics allows for **CP-violation** (CKM matrix)
- CP-violation in the SM is not enough to explain matter-antimatter asymmetry
- CP-violation in Hidden sectors or **Portals** ?

Portals

Portals

A **portal** is a Lagrangian term with fields from different sectors

Portals

A **portal** is a Lagrangian term with fields from different sectors

Some of the most commonly used in literature are:

Portals

A **portal** is a Lagrangian term with fields from different sectors

Some of the most commonly used in literature are:

Abelian Kinetic Mixing:

- Additional **U(1) abelian** dark gauge group
- **Kinetic Mixing** at dim 4: $\frac{\epsilon}{2} B^{\mu\nu} X_{\mu\nu}$
- ϵ naturally $O(1)$ but experiments yields $\epsilon \ll 1$



Portals

A **portal** is a Lagrangian term with fields from different sectors

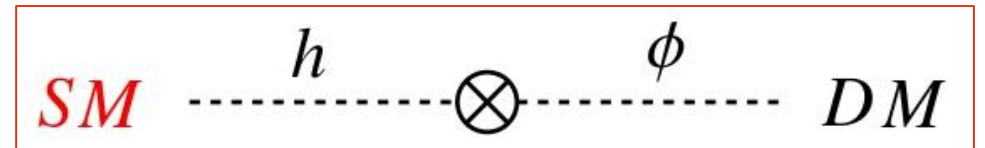
Some of the most commonly used in literature are:

Abelian Kinetic Mixing:

- Additional **U(1) abelian** dark gauge group
- **Kinetic Mixing** at dim 4: $\frac{\epsilon}{2} B^{\mu\nu} X_{\mu\nu}$
- ϵ naturally $O(1)$ but experiments yields $\epsilon \ll 1$

Scalar Mixing:

- Additional **Dark Scalar** neutral under SM
- **Interaction** at renormalizable level: $k |H|^2 |S|^2$
- $\langle S \rangle \neq 0$ and mixing.



Non Abelian Kinetic Mixing

- Introduction of a $SU(N)$ **Non Abelian Dark Sector** \supset $\left\{ \begin{array}{l} \Sigma_a : \text{Scalar fields in the adjoint of } SU(N) \\ X_a^\mu : N^2 - 1 \text{ gauge bosons} \end{array} \right.$

Non Abelian Kinetic Mixing

- Introduction of a SU(N) **Non Abelian Dark Sector** \supset $\left\{ \begin{array}{l} \Sigma_a : \text{Scalar fields in the adjoint of SU(N)} \\ X_a^\mu : N^2 - 1 \text{ gauge bosons} \end{array} \right.$

- Dim. 5 operators: **EFT description** $\overset{\text{CP-even}}{-\frac{C}{\Lambda} \text{Tr} [\Sigma X^{\mu\nu}] B_{\mu\nu}}$

Non Abelian Kinetic Mixing

- Introduction of a SU(N) **Non Abelian Dark Sector** \supset $\left\{ \begin{array}{l} \Sigma_a : \text{Scalar fields in the adjoint of SU(N)} \\ X_a^\mu : N^2 - 1 \text{ gauge bosons} \end{array} \right.$

- Dim. 5 operators: **EFT description**
$$-\frac{C}{\Lambda} \text{Tr} [\Sigma X^{\mu\nu}] B_{\mu\nu} \quad \overset{\text{CP-even}}{\quad} \quad - \frac{\tilde{C}}{\Lambda} \text{Tr} [\Sigma X^{\mu\nu}] \tilde{B}_{\mu\nu} \quad \overset{\text{CP-odd}}{\quad}$$

Non Abelian Kinetic Mixing

- Introduction of a SU(N) **Non Abelian Dark Sector** \supset
 - Σ_a : Scalar fields in the adjoint of SU(N)
 - X_a^μ : $N^2 - 1$ gauge bosons

- Dim. 5 operators: **EFT description**

$$-\frac{C}{\Lambda} \text{Tr} [\Sigma X^{\mu\nu}] B_{\mu\nu} \quad \text{CP-even} \quad - \frac{\tilde{C}}{\Lambda} \text{Tr} [\Sigma X^{\mu\nu}] \tilde{B}_{\mu\nu} \quad \text{CP-odd}$$

- SSB of SU(N) $\rightarrow \Sigma_a = v_a + \phi_a$: **Scalar Mixing** and low energy operators:

$$-\frac{\epsilon_a}{2} X_a^{\mu\nu} B_{\mu\nu} - \frac{\tilde{\epsilon}}{2} \phi^a X_a^{\mu\nu} \tilde{B}_{\mu\nu}$$

Non Abelian Kinetic Mixing

- Introduction of a SU(N) **Non Abelian Dark Sector** \supset
 - Σ_a : Scalar fields in the adjoint of SU(N)
 - X_a^μ : $N^2 - 1$ gauge bosons

- Dim. 5 operators: **EFT description**

$$-\frac{C}{\Lambda} \text{Tr} [\Sigma X^{\mu\nu}] B_{\mu\nu} \quad \text{CP-even} \quad - \frac{\tilde{C}}{\Lambda} \text{Tr} [\Sigma X^{\mu\nu}] \tilde{B}_{\mu\nu} \quad \text{CP-odd}$$

- SSB of SU(N) $\rightarrow \Sigma_a = v_a + \phi_a$: **Scalar Mixing** and low energy operators:

$$-\frac{\epsilon_a}{2} X_a^{\mu\nu} B_{\mu\nu} - \frac{\tilde{\epsilon}}{2} \phi^a X_a^{\mu\nu} \tilde{B}_{\mu\nu}$$

- Kinetic Mixing parameters **naturally** small
- New source of **CP-violation**

EDM

EDM

- CPV int. of fermions with EM fields
- QFT description: $\mathcal{L} = -\frac{i}{2} d \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$

EDM

- CPV int. of fermions with EM fields
- QFT description: $\mathcal{L} = -\frac{i}{2} d \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$
- No exp. evidence of EDMs

EDM

- CPV int. of fermions with EM fields
- QFT description: $\mathcal{L} = -\frac{i}{2} d \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$
- No exp. evidence of EDMs
- Flavor blind model means $d_i = \frac{m_i}{m_k} d_k$

EDM

- CPV int. of fermions with EM fields
- QFT description: $\mathcal{L} = -\frac{i}{2} d \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$
- No exp. evidence of EDMs
- Flavor blind model means $d_i = \frac{m_i}{m_k} d_k$

Electron EDM:

- eEDMs is the most sensitive to CPV

EDM

- CPV int. of fermions with EM fields
- QFT description: $\mathcal{L} = -\frac{i}{2} d \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$
- No exp. evidence of EDMs
- Flavor blind model means $d_i = \frac{m_i}{m_k} d_k$

Electron EDM:

- eEDMs is the most sensitive to CPV
- CPV in the SM predicts: $d_e^{eq} = 10^{-35} e \text{ cm}$
[Ema et al. (2022)]

EDM

- CPV int. of fermions with EM fields
- QFT description: $\mathcal{L} = -\frac{i}{2} d \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$
- No exp. evidence of EDMs
- Flavor blind model means $d_i = \frac{m_i}{m_k} d_k$

Electron EDM:

- eEDMs is the most sensitive to CPV
- CPV in the SM predicts: $d_e^{eq} = 10^{-35} e \text{ cm}$
[Ema et al. (2022)]
- Possible exp. deviations hint at New Physics

EDM

- CPV int. of fermions with EM fields
- QFT description: $\mathcal{L} = -\frac{i}{2} d \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$
- No exp. evidence of EDMs
- Flavor blind model means $d_i = \frac{m_i}{m_k} d_k$

Electron EDM:

- eEDMs is the most sensitive to CPV
- CPV in the SM predicts: $d_e^{eq} = 10^{-35} e \text{ cm}$
[Ema et al. (2022)]
- Possible exp. deviations hint at New Physics

Upper bound on $|d_e|$ (e · cm)

JILAeEDM	4.1×10^{-30}
ACMEIII	$\sim 1 \times 10^{-30}$
YBF	$\sim 1 \times 10^{-31}$
BaF(EDM ³)	$\sim 1 \times 10^{-33}$

[Roussy et al. (2023)]

[Hiramoto et al. (2023)]

[Fitch et al. (2021)]

[Vutha et al. (2018)]

EDM

- CPV int. of fermions with EM fields
- QFT description: $\mathcal{L} = -\frac{i}{2} d \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$
- No exp. evidence of EDMs
- Flavor blind model means $d_i = \frac{m_i}{m_k} d_k$

Electron EDM:

- eEDMs is the most sensitive to CPV
- CPV in the SM predicts: $d_e^{eq} = 10^{-35} e \text{ cm}$
[Ema et al. (2022)]
- Possible exp. deviations hint at New Physics

Upper bound on $|d_e|$ (e · cm)

JILAeEDM	4.1×10^{-30}
ACMEIII	$\sim 1 \times 10^{-30}$
YBF	$\sim 1 \times 10^{-31}$
BaF(EDM ³)	$\sim 1 \times 10^{-33}$

[Roussy et al. (2023)]

[Hiramoto et al. (2023)]

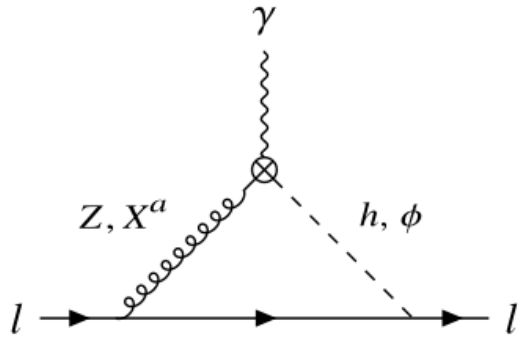
[Fitch et al. (2021)]

[Vutha et al. (2018)]

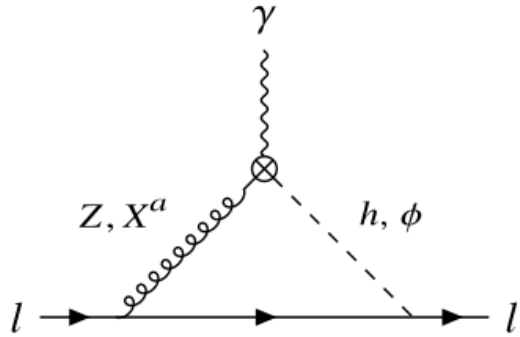
Expect significant improvements of the current JILAeEDM sensitivity in the coming years!

eEDM: prediction

eEDM: prediction



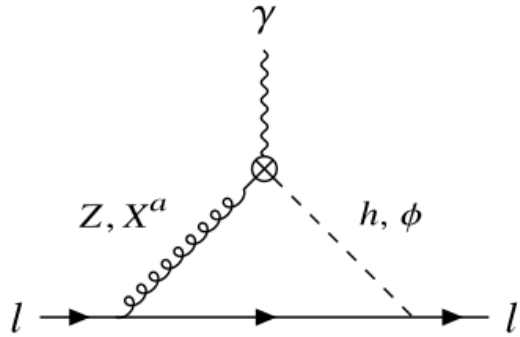
eEDM: prediction



Assumptions:

- One Dark boson and one Dark scalar mix with SM
- Dark Sector at the same energy scale: $v_D \sim m_\phi \sim M_X$

eEDM: prediction

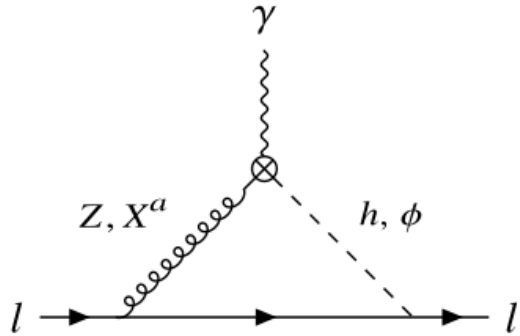


Assumptions:

- One Dark boson and one Dark scalar mix with SM
- Dark Sector at the same energy scale: $v_D \sim m_\phi \sim M_X$

$$d_l = \frac{Y_l}{8\pi^2 v_D} \epsilon^2 \tan \chi \beta c_\theta^2 e f(M_X, m_\phi, m_h)$$

eEDM: prediction

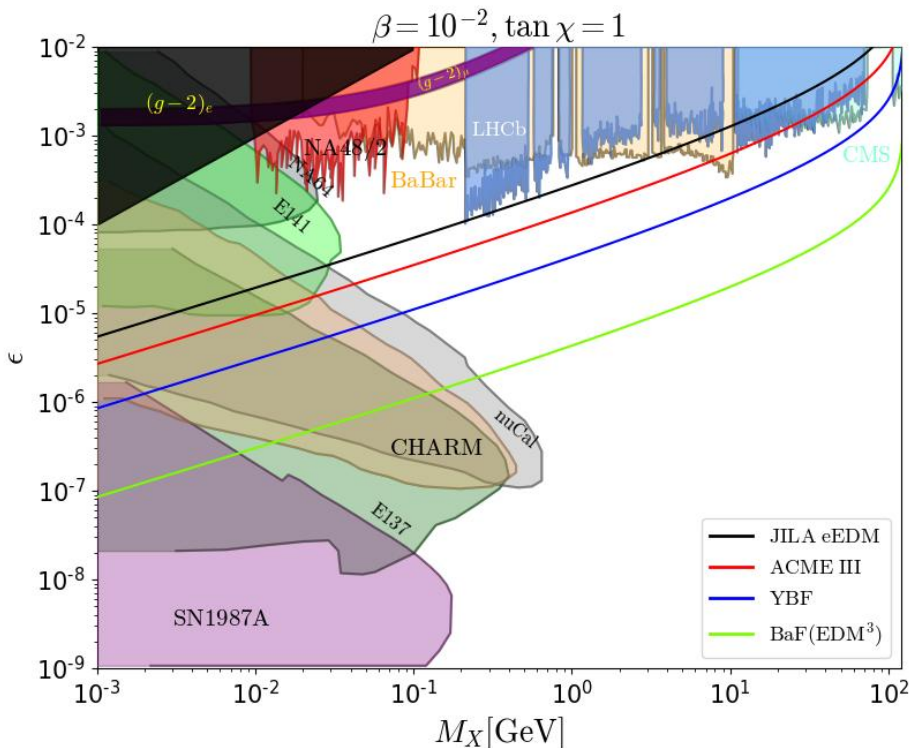


Assumptions:

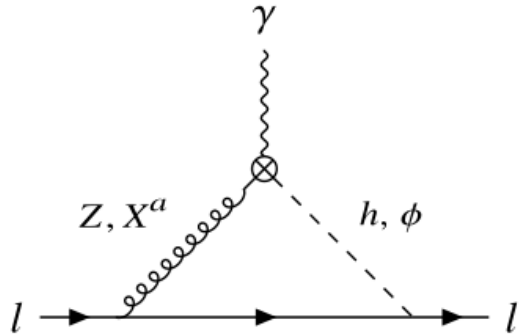
- One Dark boson and one Dark scalar mix with SM
- Dark Sector at the same energy scale: $v_D \sim m_\phi \sim M_X$

$$d_l = \frac{Y_l}{8\pi^2 v_D} \epsilon^2 \tan \chi \beta c_\theta^2 e f(M_X, m_\phi, m_h)$$

- Scalar mixing parameter $\beta \lesssim 10^{-2}$
[T.Ferber et al. (2024)]
- Constraints on ϵ from colliders and beam dump exp.



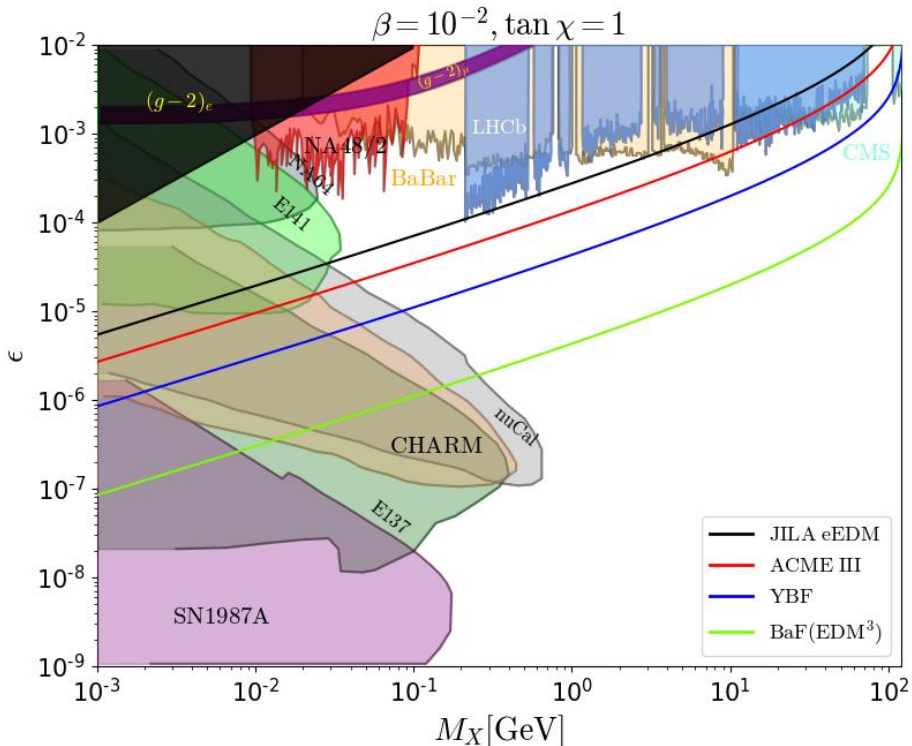
eEDM: prediction



Assumptions:

- One Dark boson and one Dark scalar mix with SM
- Dark Sector at the same energy scale: $v_D \sim m_\phi \sim M_X$

$$d_l = \frac{Y_l}{8\pi^2 v_D} \epsilon^2 \tan \chi \beta c_\theta^2 e f(M_X, m_\phi, m_h)$$



- Scalar mixing parameter $\beta \lesssim 10^{-2}$
[T.Ferber et al. (2024)]
- Constraints on ϵ from colliders and beam dump exp.
- Sizeable eEDM for $\epsilon \sim 10^{-5} \div 10^{-3}$

SU(2) and thermal DM

SU(2) and thermal DM

SU(2) → U(1) :

SU(2) and thermal DM

SU(2) → U(1) :

SSB of SU(2) via $\langle \Sigma \rangle \neq 0$ → unbroken U(1) → mCPs only small fraction of DM
[McDermott et al., 2011]

SU(2) and thermal DM

SU(2) → U(1) :

SSB of SU(2) via $\langle \Sigma \rangle \neq 0$ → unbroken U(1) → mCPs only small fraction of DM
[McDermott et al., 2011]

SU(2) → ∅ :

SU(2) and thermal DM

SU(2) → U(1) :

SSB of SU(2) via $\langle \Sigma \rangle \neq 0$ → unbroken U(1) → mCPs only small fraction of DM
[McDermott et al., 2011]

SU(2) → ∅ :

INDIRECT DETECTION:

WIMP DM mass ≥ 30 GeV

[Planck,2018]

SU(2) and thermal DM

SU(2) → U(1) :

SSB of SU(2) via $\langle \Sigma \rangle \neq 0$ \longrightarrow unbroken U(1) \longrightarrow mCPs only small fraction of DM
[McDermott et al., 2011]

SU(2) → \emptyset :

INDIRECT DETECTION:

WIMP DM mass ≥ 30 GeV

[Planck,2018]

DIRECT DETECTION:

Severe constrains ϵ for DM $>$ few GeV

SU(2) and thermal DM

SU(2) → U(1) :

SSB of SU(2) via $\langle \Sigma \rangle \neq 0$ \longrightarrow unbroken U(1) \longrightarrow mCPs only small fraction of DM
[McDermott et al., 2011]

SU(2) → ∅ :

INDIRECT DETECTION:

WIMP DM mass ≥ 30 GeV

[Planck,2018]

DIRECT DETECTION:

Severe constrains ϵ for DM $>$ few GeV

Standard WIMP scenario highly disfavored !

SU(2) and thermal DM

SU(2) → U(1) :

SSB of SU(2) via $\langle \Sigma \rangle \neq 0$ \longrightarrow unbroken U(1) \longrightarrow mCPs only small fraction of DM
[McDermott et al., 2011]

SU(2) → \emptyset :

INDIRECT DETECTION:

WIMP DM mass ≥ 30 GeV

[Planck,2018]

DIRECT DETECTION:

Severe constrains ϵ for DM $>$ few GeV

Standard WIMP scenario highly disfavored !

Inelastic Dark Matter:

SU(2) and thermal DM

SU(2) → U(1) :

SSB of SU(2) via $\langle \Sigma \rangle \neq 0$ → unbroken U(1) → mCPs only small fraction of DM
[McDermott et al., 2011]

SU(2) → ∅ :

INDIRECT DETECTION:

WIMP DM mass ≥ 30 GeV

[Planck,2018]

DIRECT DETECTION:

Severe constrains ϵ for DM $>$ few GeV

Standard WIMP scenario highly disfavored !

Inelastic Dark Matter:

- Fermionic DS with at least 2 states (χ_H, χ_S)
- Mass splitting between χ_H and χ_S (DM)
- $\chi_S \chi_S \rightarrow$ SM forbidden (or highly suppressed)

SU(2) and thermal DM

SU(2) → U(1) :

SSB of SU(2) via $\langle \Sigma \rangle \neq 0$ → unbroken U(1) → mCPs only small fraction of DM
[McDermott et al., 2011]

SU(2) → ∅ :

INDIRECT DETECTION:

WIMP DM mass ≥ 30 GeV

[Planck, 2018]

DIRECT DETECTION:

Severe constrains ϵ for DM $>$ few GeV

Standard WIMP scenario highly disfavored !

Inelastic Dark Matter:

- Fermionic DS with at least 2 states (χ_H, χ_S)
- Mass splitting between χ_H and χ_S (DM) →
- $\chi_S \chi_S \rightarrow$ SM forbidden (or highly suppressed)

No DM annihilation → No I-D bounds

$\Delta m \gtrsim 1$ MeV → Negligible D-D bounds

Inelastic Dark Matter SU(2) model

Inelastic Dark Matter SU(2) model

- 3 gauge fields X_i^μ
- 2 scalar fields in the adj. Σ_2^a, Σ_3^a
- 2 SU(2) doublets $\left\{ \begin{array}{l} \bullet \chi_L = (\chi_L^1, \chi_L^2) \\ \bullet \psi_R = (\psi_R^1, \psi_R^2) \end{array} \right.$

Inelastic Dark Matter SU(2) model

- 3 gauge fields X_i^μ
 - 2 scalar fields in the adj. Σ_2^a, Σ_3^a
 - 2 SU(2) doublets $\left\{ \begin{array}{l} \bullet \chi_L = (\chi_L^1, \chi_L^2) \\ \bullet \psi_R = (\psi_R^1, \psi_R^2) \end{array} \right.$
-
- SSB and mass basis: Ψ_S, Ψ_H Dirac fields
 - X_2^μ, X_3^μ mix with SM gauge bosons
 - Off-diagonal current: $g_D X_2^\mu \bar{\Psi}_H \gamma_\mu \Psi_S$

Inelastic Dark Matter SU(2) model

- 3 gauge fields X_i^μ
 - 2 scalar fields in the adj. Σ_2^a, Σ_3^a
 - 2 SU(2) doublets $\left\{ \begin{array}{l} \bullet \chi_L = (\chi_L^1, \chi_L^2) \\ \bullet \psi_R = (\psi_R^1, \psi_R^2) \end{array} \right.$
-
- SSB and mass basis: Ψ_S, Ψ_H Dirac fields
 - X_2^μ, X_3^μ mix with SM gauge bosons
 - Off-diagonal current: $g_D X_2^\mu \bar{\Psi}_H \gamma_\mu \Psi_S$

DM vs eEDM

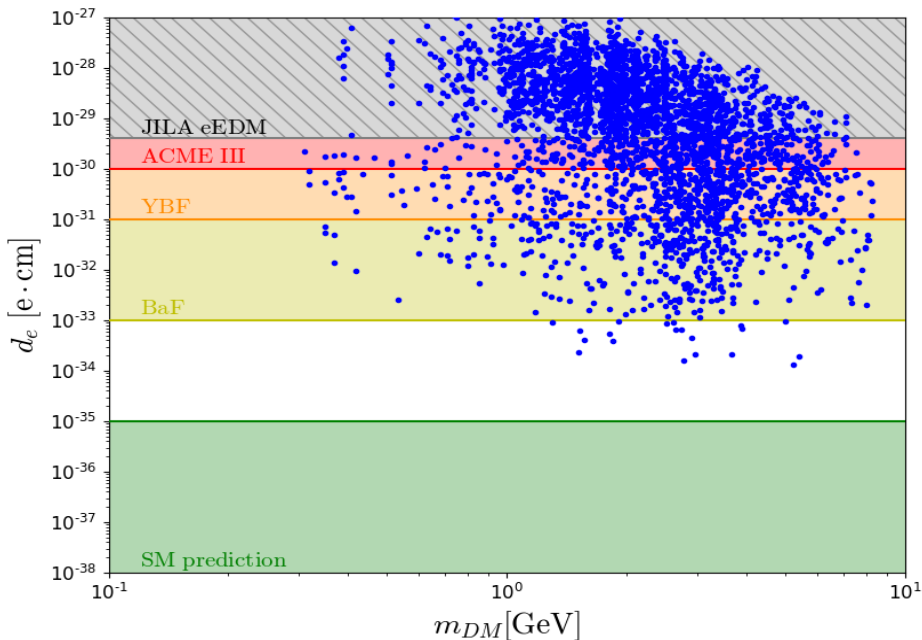
- Freeze out via coannihilation $\Psi_S \Psi_H \rightarrow \text{SM}$
- $m_{\Psi_S} \sim m_{\Psi_H} < M_X \sim 1-10 \text{ GeV}$
- $\Omega_\Psi h^2 = 0.12$ for $\epsilon \sim 10^{-5} \div 10^{-3}$

Inelastic Dark Matter SU(2) model

- 3 gauge fields X_i^μ
 - 2 scalar fields in the adj. Σ_2^a, Σ_3^a
 - 2 SU(2) doublets
 - $\chi_L = (\chi_L^1, \chi_L^2)$
 - $\psi_R = (\psi_R^1, \psi_R^2)$
-
- SSB and mass basis: Ψ_S, Ψ_H Dirac fields
 - X_2^μ, X_3^μ mix with SM gauge bosons
 - Off-diagonal current: $g_D X_2^\mu \bar{\Psi}_H \gamma_\mu \Psi_S$

DM vs eEDM

- Freeze out via coannihilation $\Psi_S \Psi_H \rightarrow \text{SM}$
- $m_{\Psi_S} \sim m_{\Psi_H} < M_X \sim 1-10 \text{ GeV}$
- $\Omega_\Psi h^2 = 0.12$ for $\epsilon \sim 10^{-5} \div 10^{-3}$

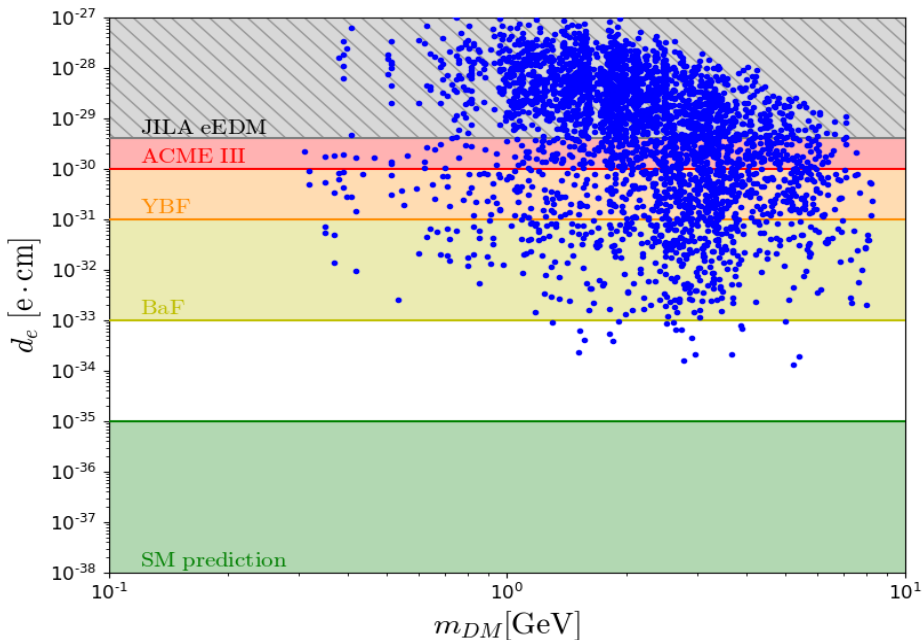


Inelastic Dark Matter SU(2) model

- 3 gauge fields X_i^μ
 - 2 scalar fields in the adj. Σ_2^a, Σ_3^a
 - 2 SU(2) doublets
 - $\chi_L = (\chi_L^1, \chi_L^2)$
 - $\psi_R = (\psi_R^1, \psi_R^2)$
-
- SSB and mass basis: Ψ_S, Ψ_H Dirac fields
 - X_2^μ, X_3^μ mix with SM gauge bosons
 - Off-diagonal current: $g_D X_2^\mu \bar{\Psi}_H \gamma_\mu \Psi_S$

DM vs eEDM

- Freeze out via coannihilation $\Psi_S \Psi_H \rightarrow \text{SM}$
- $m_{\Psi_S} \sim m_{\Psi_H} < M_X \sim 1-10 \text{ GeV}$
- $\Omega_\Psi h^2 = 0.12$ for $\epsilon \sim 10^{-5} \div 10^{-3}$
- Future eEDM sensitivities can probe the model



Summary

- Non-abelian Dark sector allows for kinetic portals with small ϵ
- Non-abelian Dark sector allows for a CP-violating phase in portals
- Scalar and kinetic mixing + CP-violation signals can be traced in EDMs
- Model of iDM can be probed by future searches for a permanent eEDM!

Thank you for your attention!

BACK UP

UV completion

- **EFTs** call for UV completion
- Heavy vector-like fermion charged under $SU(N) \otimes U(1)_Y$
- Physical phase χ in Yukawa-like scalar couplings \mathcal{Y}

UV Lagrangian:

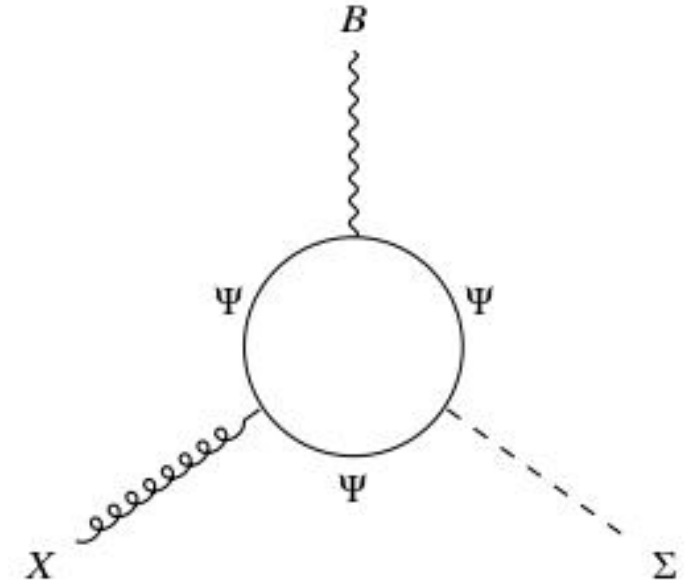
$$\mathcal{L}_\Theta \supset -g_1 Y \bar{\Theta}_i \gamma_\mu \Theta_i B^\mu - g_D \bar{\Theta}_i T^a \gamma_\mu \Theta_i X_a^\mu - M_{ii} \bar{\Theta}_{iR} \Theta_{iL} - \mathcal{Y}_{ij} \bar{\Theta}_{iR} \Sigma^a T^a \Theta_{jL} + \text{h.c.}$$

UV-EFT matching

$$\frac{C}{\Lambda} = \frac{g_D g_1 Y \text{Re}[\mathcal{Y}]}{12\pi^2 M};$$

$$\frac{\tilde{C}}{\Lambda} = \frac{g_D g_1 Y \text{Im}[\mathcal{Y}]}{16\pi^2 M}$$

$$\tilde{\epsilon} = \frac{3 \text{Im}[\mathcal{Y}]}{4 \text{Re}[\mathcal{Y}]} \epsilon \equiv (\tan \chi) \epsilon$$



A model for Inelastic Dark Matter

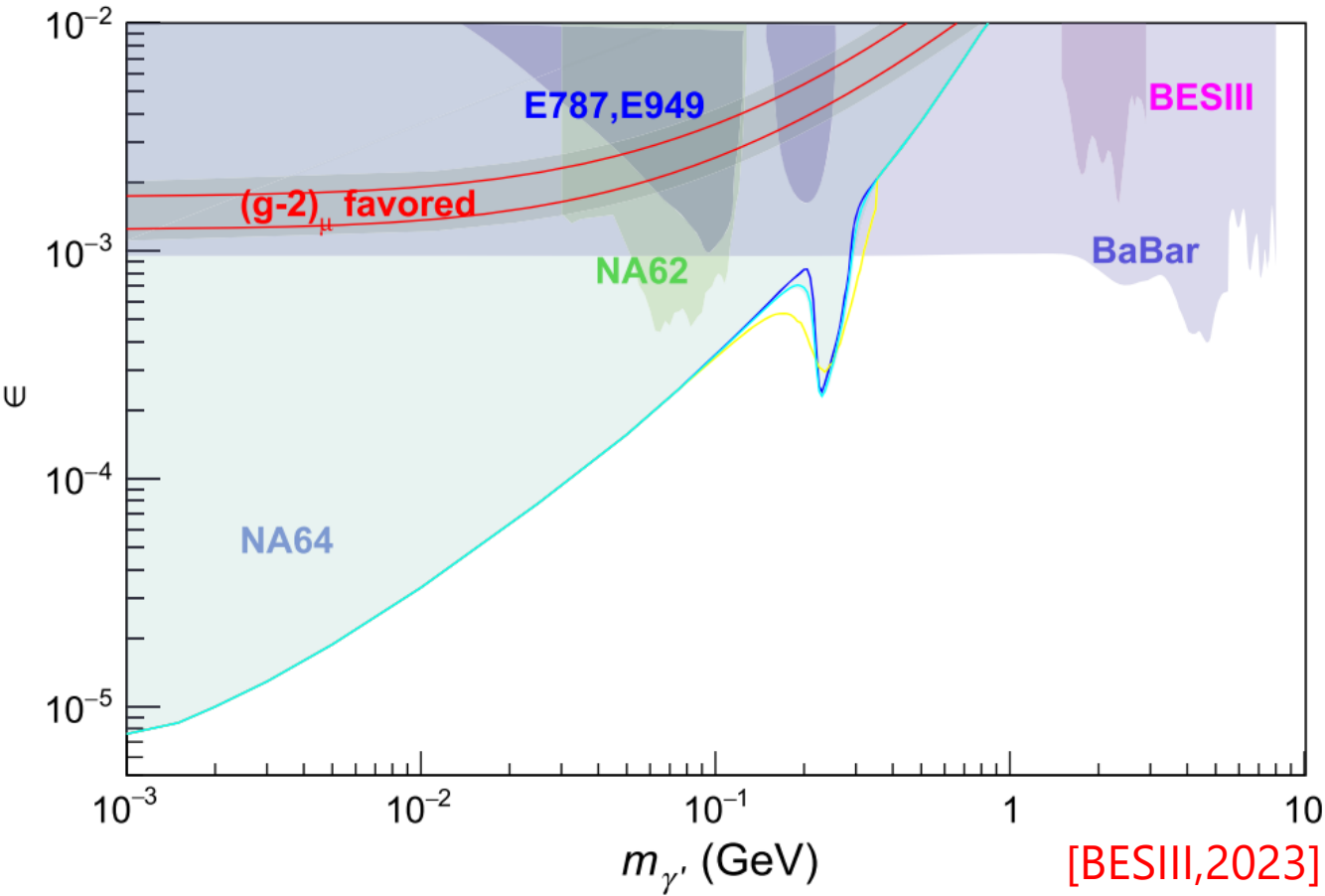
- SU(2) Dark group with matter content:
 - 3 gauge fields X_i^μ
 - 2 scalar fields in the adj. Σ_2^a, Σ_3^a
 - 2 Majorana SU(2) doublet
 - $\chi_L = (\chi_L^1, \chi_L^2)$
 - $\psi_R = (\psi_R^1, \psi_R^2)$

• Mass term:
$$\mathcal{L} \supset -m_D \overline{\chi_L} \psi_R - \sum_{i=1,2} Y_{D,i} \overline{\chi_L} \Sigma_i \psi_R - \sum_{i=1,2} Y_{L,i} \overline{\chi_L^c} i \sigma_2 \Sigma_i \chi_L - \sum_{i=1,2} Y_{R,i} \overline{\psi_R^c} i \sigma_2 \Sigma_i \psi_R + \text{h.c.}$$

$$- \frac{g_D}{2} \overline{\chi_L} \gamma_\mu \sigma^a X_a^\mu \chi_L - \frac{g_D}{2} \overline{\psi_R} \gamma_\mu \sigma^a X_a^\mu \psi_R. \quad (\text{C.2})$$

- SU(2) fully broken by: $\langle \Sigma_2 \rangle = (0, v_2, 0); \langle \Sigma_3 \rangle = (0, 0, v_3)$
- Dirac masses: $M_1 = m_D + vY_1 - vY_2; M_2 = m_D + vY_1 + vY_2$
- Off-diagonal currents with X_2 and X_1 and inelastic dark matter scenario
- X_3 diagonal current suppressed by either small eps or large M_{X_3} .

Laboratory bounds



$$\sum_f \Gamma(\Psi_H \rightarrow \Psi_S f \bar{f}) = \sum_f \frac{4\epsilon^2 \alpha \alpha_D \delta_\Psi^5}{15\pi M_X^4}$$

- Small mass splitting: favoured for DM a and long life time
- $\delta_\Psi > 2m_e$ to avoid too much long lifetime
- X to invisible searches: model independent constraints

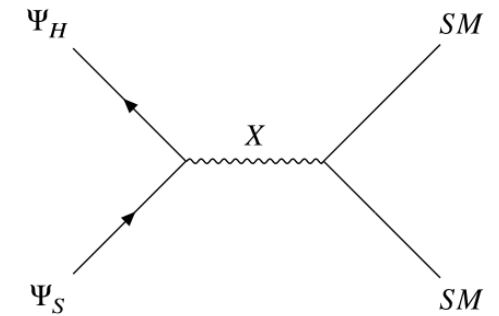
Inelastic DM set up

Scatter plot parameters:

Parameter	Lower limit	Upper limit
$\beta_{1,2}$	10^{-4}	10^{-2}
ϵ	10^{-6}	[102–106]
g_D	10^{-2}	1
$\tan \chi$	10^{-2}	1
v_D	1 GeV	20 GeV
m_ϕ	0.1 GeV	50 GeV

$$\langle \sigma_{eff} v \rangle \sim \frac{\langle \sigma v \rangle}{2} \sim 1.7 \times 10^{-9} \text{GeV}^{-2}$$

Total cross section $\chi\chi \rightarrow SM$



- Random values of parameters of the model and M_χ chosen to fulfill the relation above.
- eEDM computed with the formula slides 5