

Third-family lepton-quark fusion

Resonant U_1 leptoquark production @ LHC

in collaboration with Peter Krack (NIKHEF) & Nudžeim Selimović (INFN Padova)

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Arman Korajac (INFN Pisa)

OUTLINE

- **UV model for U_1**
- **NLO QCD + QED results**
- **Simulation details & exclusion limits**
- **Conclusions**

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- **Conclusions**

4321 for U1

- Passing step in many UV flavor theories
- Decoupling the 3rd gen and the light fermions in the UV, effectively starting out with a $U(2)^5$ flavor symmetry
- Motivated by solving the $R(D^*)$ anomaly

$$\begin{array}{c}
 [SU(4) \times SU(2)_L \times SU(2)_R]^{[3]} \times [SU(3) \times SU(2)_L \times U(1)_{B-L} \times U(1)_R]^{[12]} \\
 \downarrow \langle \phi \rangle \\
 [SU(4) \times SU(2)_L \times SU(2)_R]^{[3]} \times [SU(3) \times SU(2)_L \times U(1)_Y]^{[12]} \\
 \begin{array}{c} \langle \Sigma_L \rangle \downarrow \langle \Sigma_R \rangle \end{array} \\
 SU(4)^{[3]} \times SU(3)^{[12]} \times SU(2)_L \times U(1)_X \\
 \begin{array}{c} \langle \Omega_3 \rangle \downarrow \langle \Omega_1 \rangle \end{array} \\
 SU(3)_c \times SU(2)_L \times U(1)_Y ,
 \end{array}$$

[Courtesy of N. Selimović]

[L. Di Luzio, A. Greljo and M. Nardecchia, arXiv: 1708.08450]

[L. Di Luzio, J. Fuentes-Martin, A. Greljo, M. Nardecchia, S. Renner, arXiv: 1808.00942]

[M. Bordone, C. Cornella, J. Fuentes-Martin, G. Isidori, arXiv: 1712.01368, 1805.09328]

[H. Georgi, Y. Nakai, arXiv: 1606.05865]

[Fuentes-Martin et al, arXiv: 1910.13474, 2006.16250, 2009.11296]

[D. Guadagnoli, M. Reboud, P. Stangl, arXiv: 2005.10117] ...

4321 for U1

- Using the minimal set-up, enough for gauge boson dynamics:

$$\mathcal{L} \supset -\frac{1}{4}H_{\mu\nu}^a H_{\mu\nu}^a + \sum_{i=1,3} (D_\mu \Omega_i)^\dagger (D_\mu \Omega_i) + \sum_{f=L,R^\pm} i\bar{\psi}_f \not{D}\psi_f + V(\Omega_i)$$

$$- Y_{\text{light,u}}^{ij} \bar{q}_L^i \tilde{H} u_R^j - Y_{\text{light,d}}^{ij} \bar{q}_L^i H d_R^j - y_{LR}^- \bar{\psi}_L H \psi_R^- - y_{LR}^+ \bar{\psi}_L \tilde{H} \psi_R^+ + \text{h.c.}$$

Field	$SU(4)_3$	$SU(3)_{12}$	$SU(2)_L$	$U(1)_X$
ψ_L	4	1	2	0
ψ_R^+	4	1	1	1/2
ψ_R^-	4	1	1	-1/2
q_L^i	1	3	2	1/6
ℓ_L^i	1	1	2	-1/2
u_R^i	1	3	1	2/3
d_R^i	1	3	1	-1/3
e_R^i	1	1	1	-1
H	1	1	2	1/2
Ω_3	$\bar{\mathbf{4}}$	3	1	1/6
Ω_1	$\bar{\mathbf{4}}$	1	1	-1/2

$i = 1, 2$

$$\psi_L \equiv (q_L^3 \ell_L^3)^\top$$

$$\psi_R^- \equiv (d_R^3 e_R^3)^\top \quad \psi_R^+ \equiv (u_R^3 \nu_R^3)^\top$$

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SSB : $\mathcal{G}_{4321} \rightarrow \mathcal{G}_{\text{SM}}$

$$G' \sim (\mathbf{8}, \mathbf{1}, 0), \quad U_1 \sim (\mathbf{3}, \mathbf{1}, 2/3), \quad Z' \sim (\mathbf{1}, \mathbf{1}, 0)$$

SM fields

$$G_\mu^a = s_3 H_\mu^a + c_3 C_\mu^a, \quad B_\mu = s_1 H_\mu^{15} + c_1 X_\mu$$

Massive gauge bosons

$$G_\mu^{\prime a} = c_3 H_\mu^a - s_3 C_\mu^a, \quad Z'_\mu = c_1 H_\mu^{15} - s_1 X_\mu$$

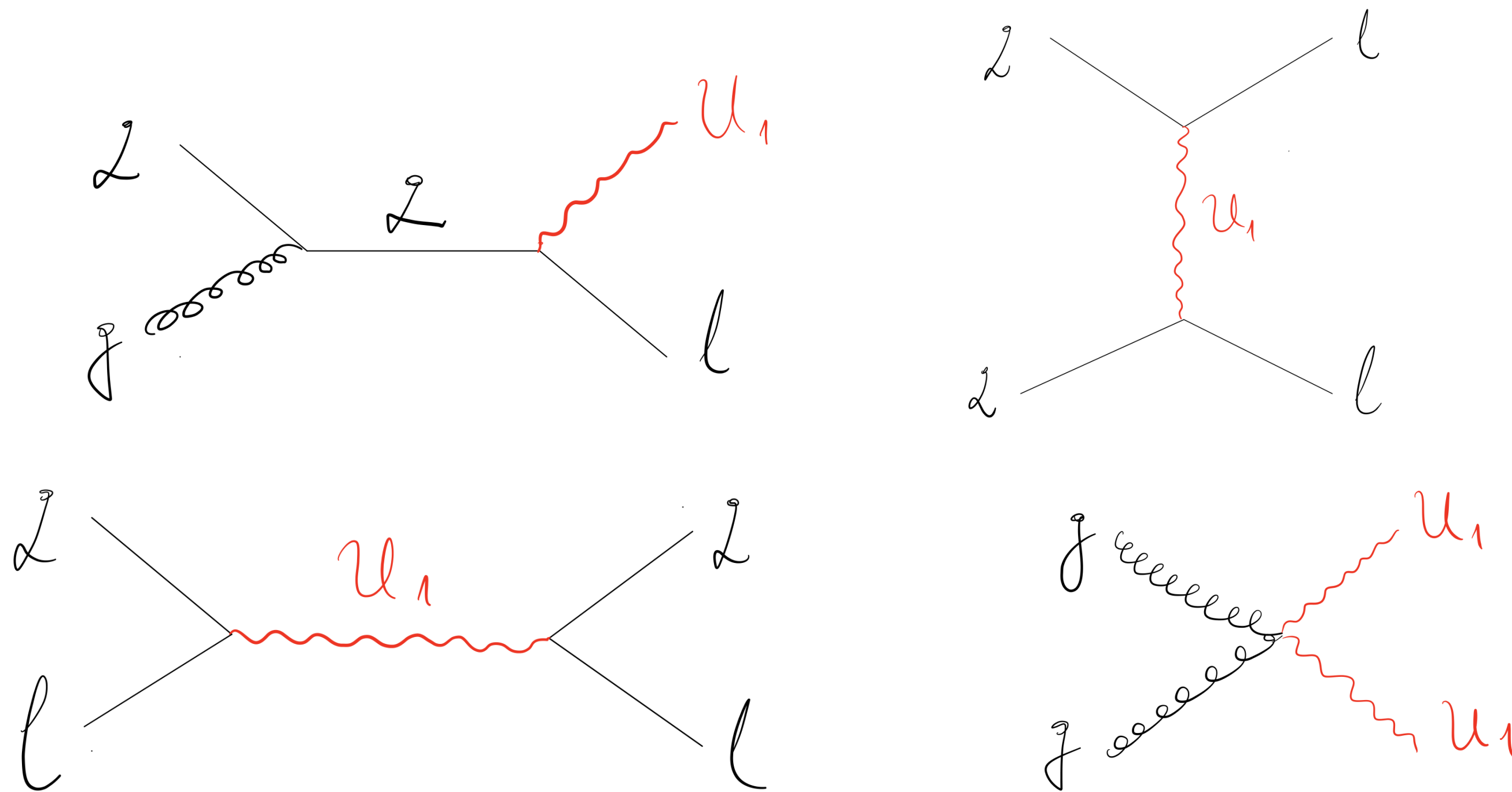
$$U_\mu^{1,2,3} = \frac{1}{\sqrt{2}} (H_\mu^{9,11,13} - i H_\mu^{10,12,14})$$

4321 for U1

$$\mathcal{L}_{U_1} \supset \frac{g_4}{\sqrt{2}} U_\mu (\beta_L^{33} \bar{q}_L^3 \gamma^\mu \tau_L + \beta_R^{33} \bar{b}_R \gamma^\mu \tau_R + \text{h.c.})$$

$$- ig_s U_\mu^\dagger T^a U_\nu G^{a,\mu\nu} - \frac{2}{3} ie U_\mu^\dagger U_\nu F^{\mu\nu}$$

Can be extended to couplings with light-gen fermions; additional fields required



- Signature can be seen through t-channel production, LQ + lepton or double LQ production
- **Goal: Examine the resonant s-channel production for U1**
- Already done, *however*:
 - no lepton PDFs at that moment
 - no experimental results at that time
 - @ LO.

[U. Haisch, G. Polesello, arXiv:2012.11474]

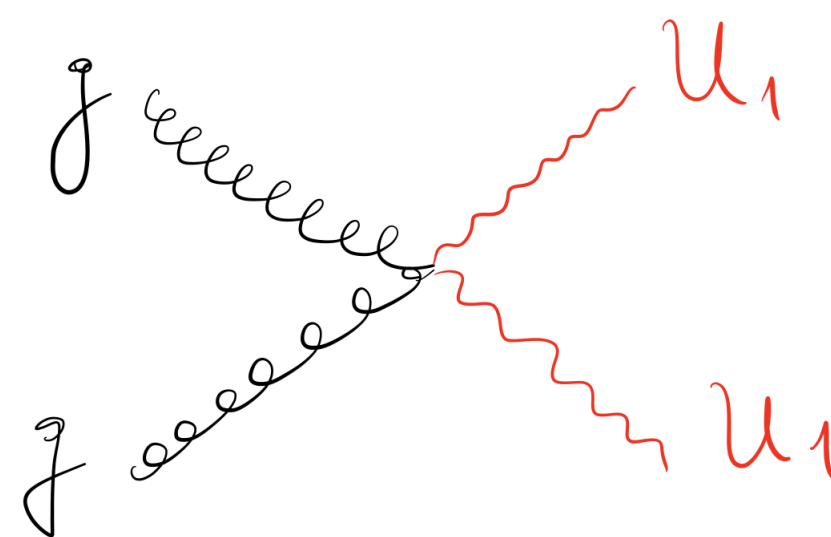
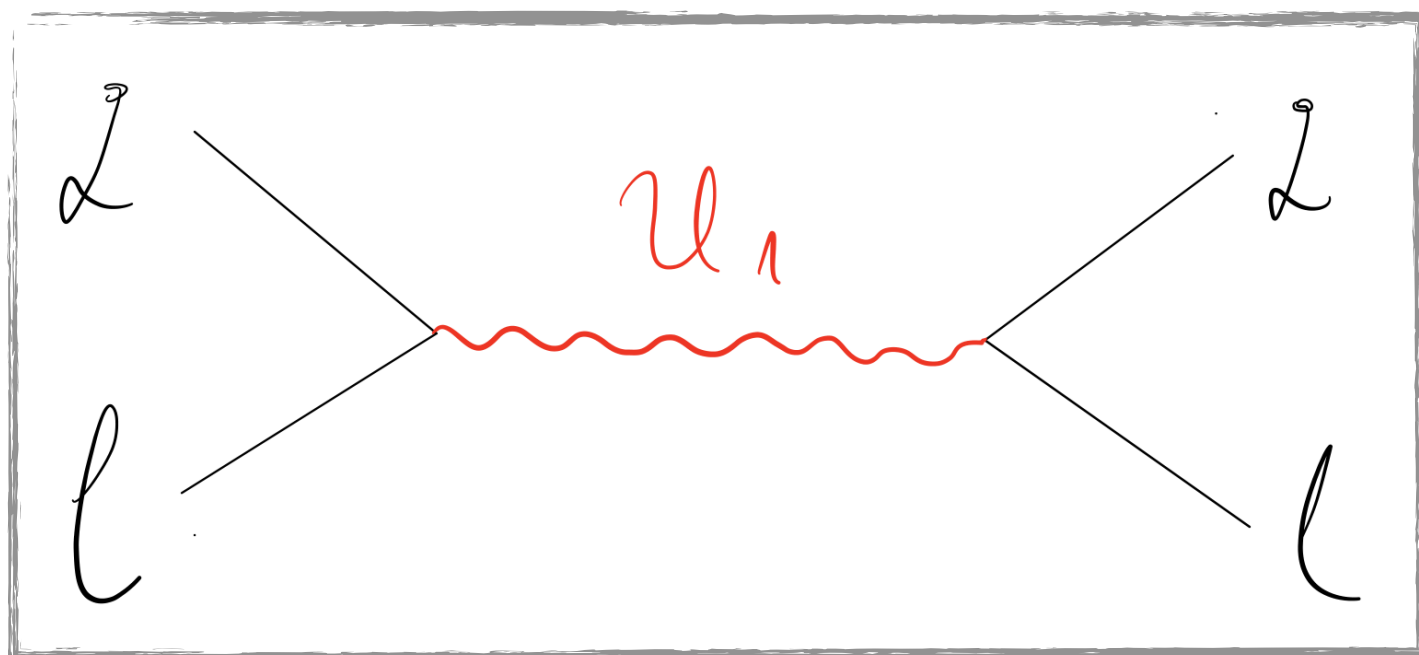
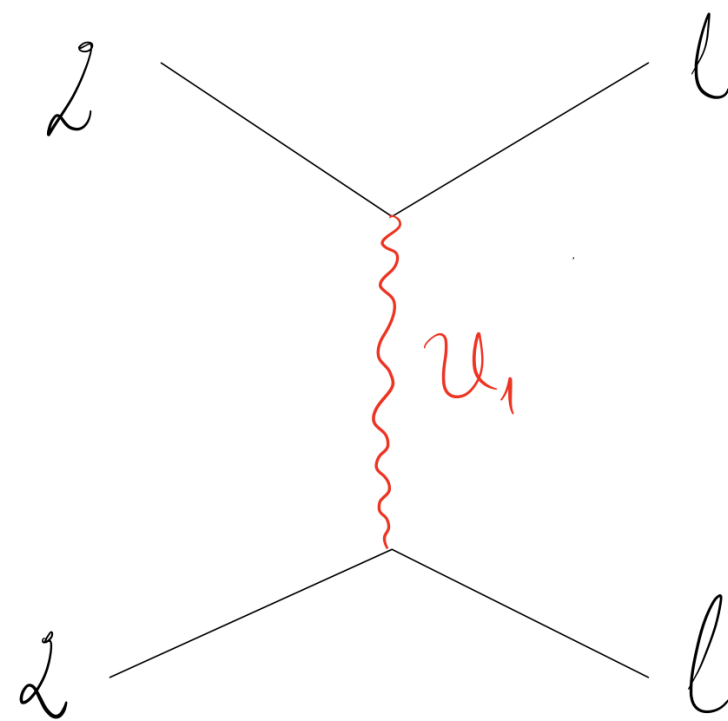
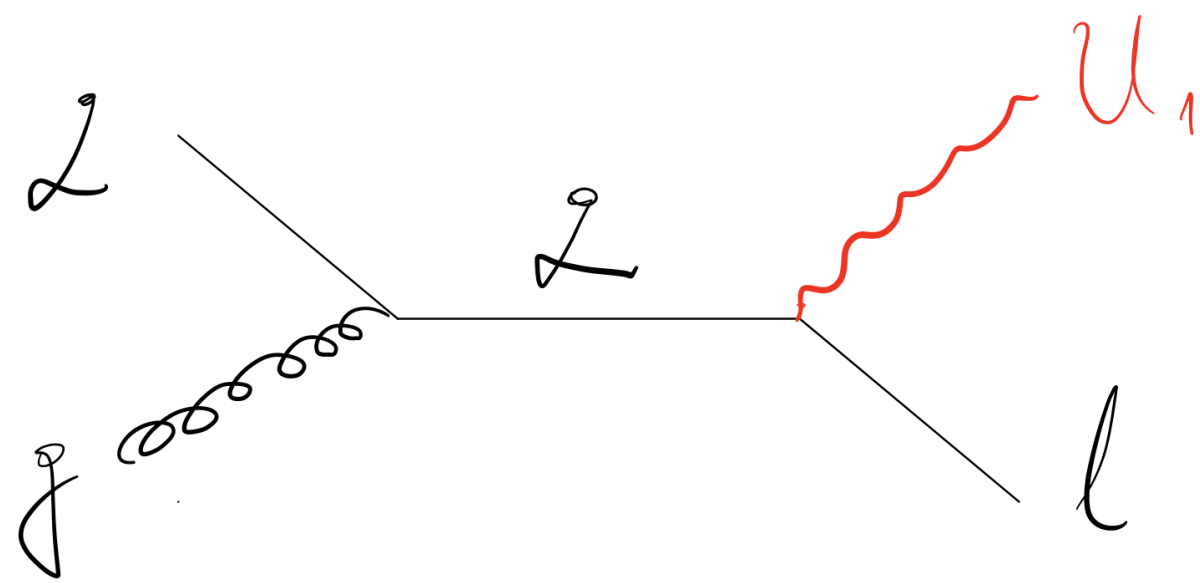
[Buonocore et al, arXiv: 2005.06475]

4321 for U1

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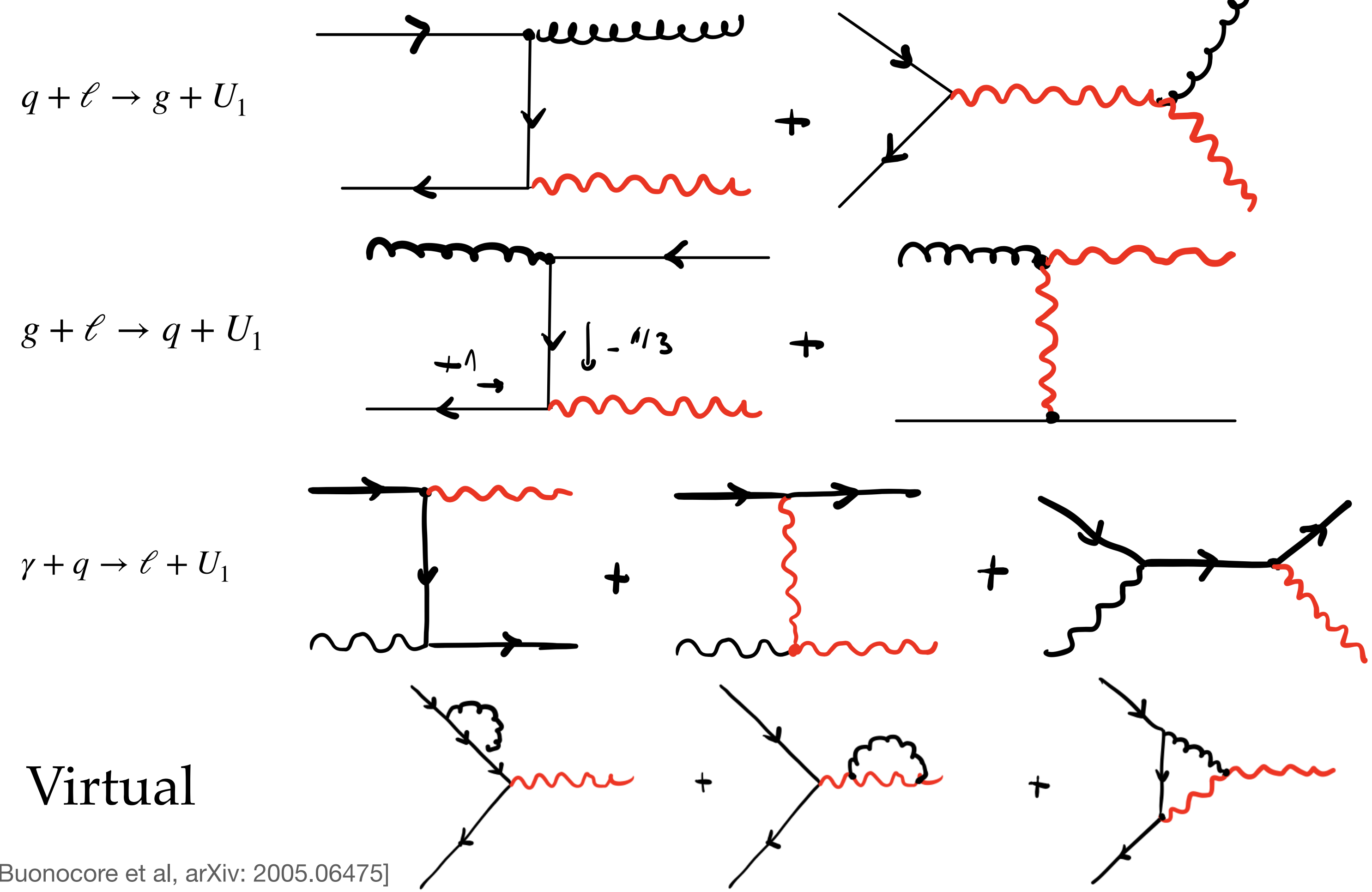
[Buonocore et al, arXiv: 2005.06475]

OUTLINE

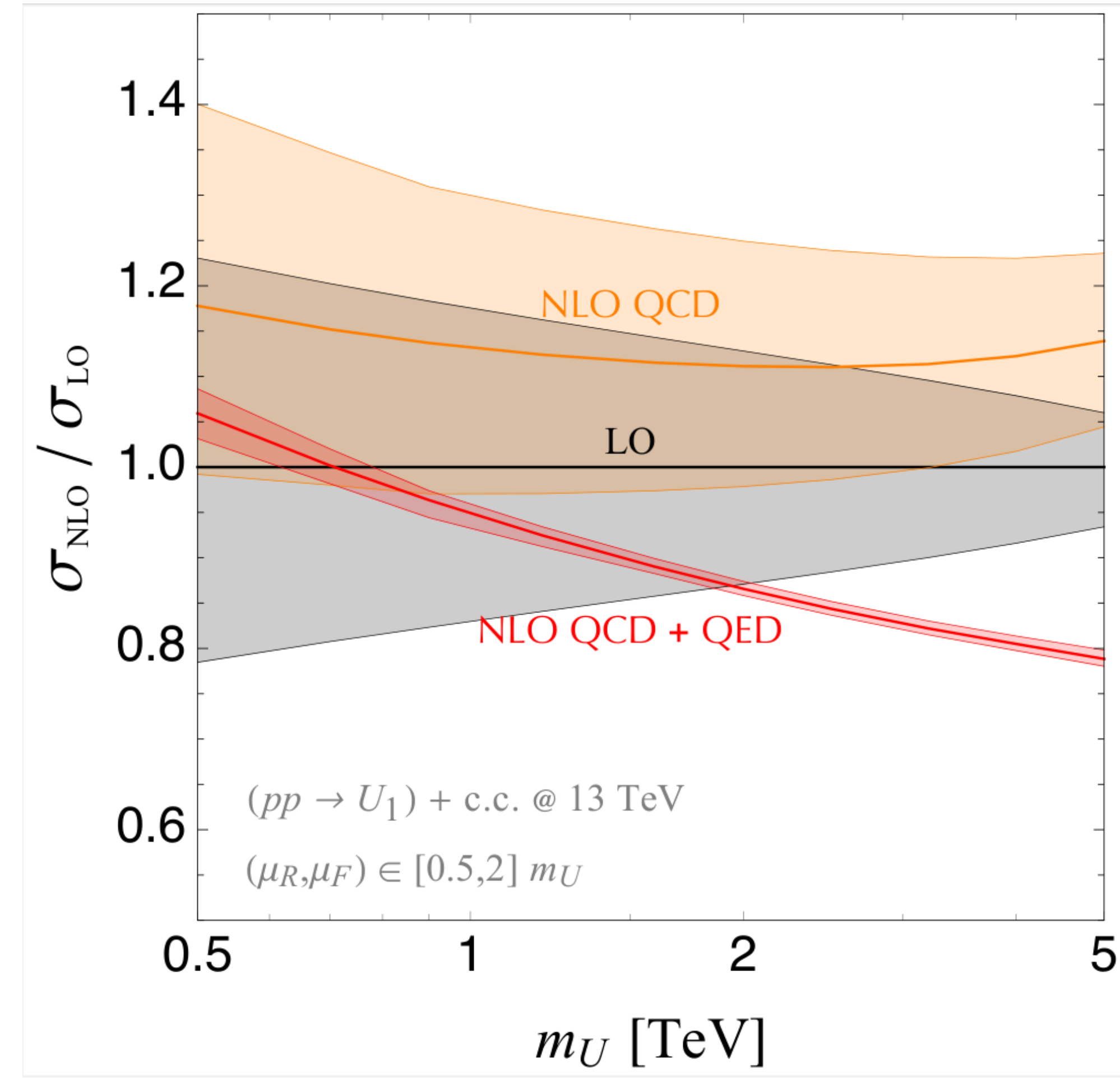
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NLO QCD + QED results

QED coupling size: $\alpha_{em} \sim \alpha_s^2$



K-factor



[Buonocore et al, arXiv: 2005.06475]
 [A. Greljo and N. Selimović, arXiv: 2012.02092]

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Simulation details & exclusion limits

- With the matrix elements squared and the finite renormalization piece, build a POWHEG-BOX event generator

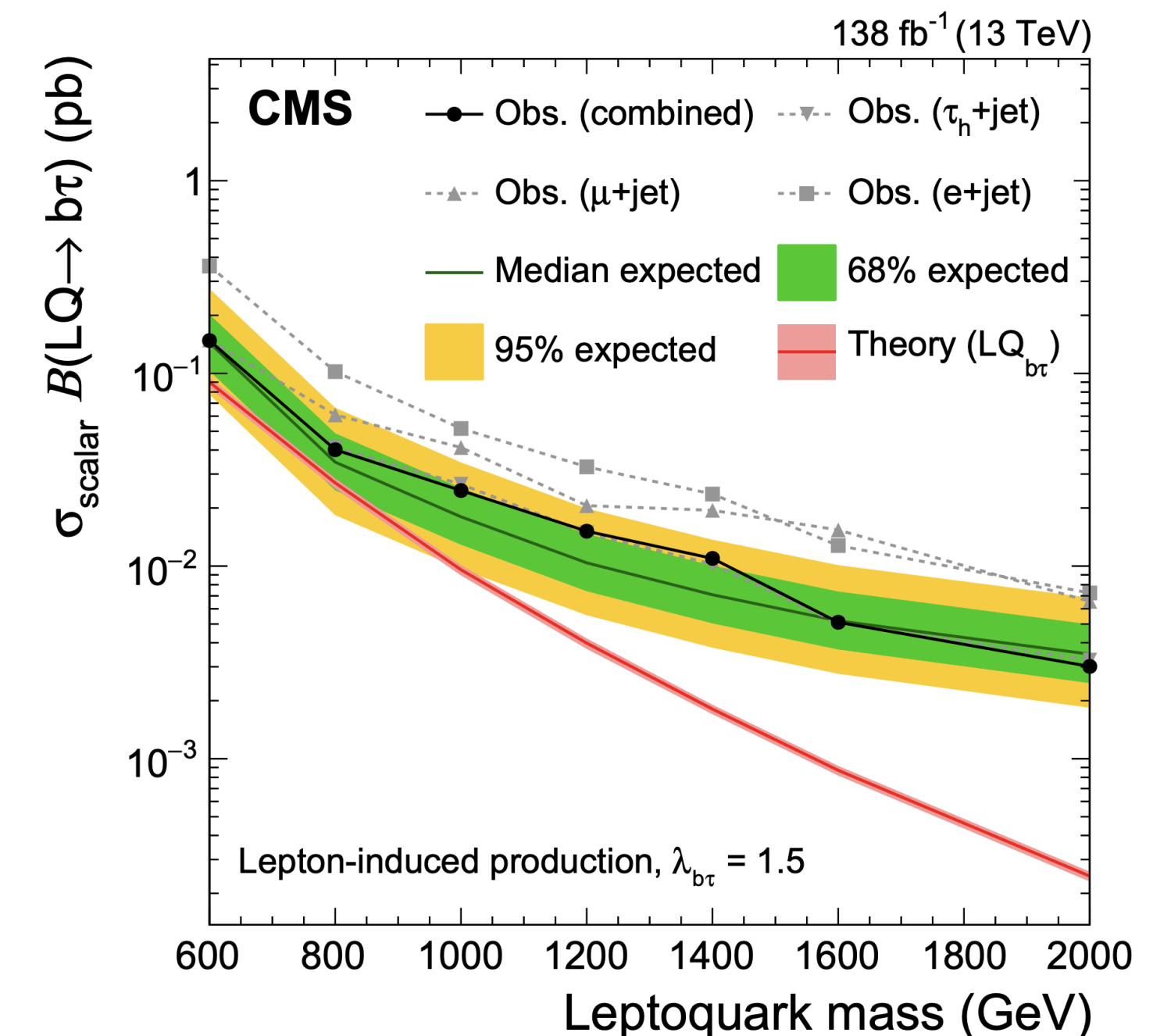
[Frixione et al, arXiv: 0707.3088]

[Buonocore et al, arXiv: 2209.02599]

- HERWIG has implemented lepton showering algorithm, contrary to Pythia - correction $\sim 30\%$ on the cross-section
- Using the NNPDF3.1 LUXlep PDF set - **with lepton PDFs included**
- Jet clustering - FastJet
- Detector effects, tagging and reconstruction - Delphes or own POWHEG algorithm
- <https://github.com/peterkrack/3rd-Lepton-Quark-Fusion>

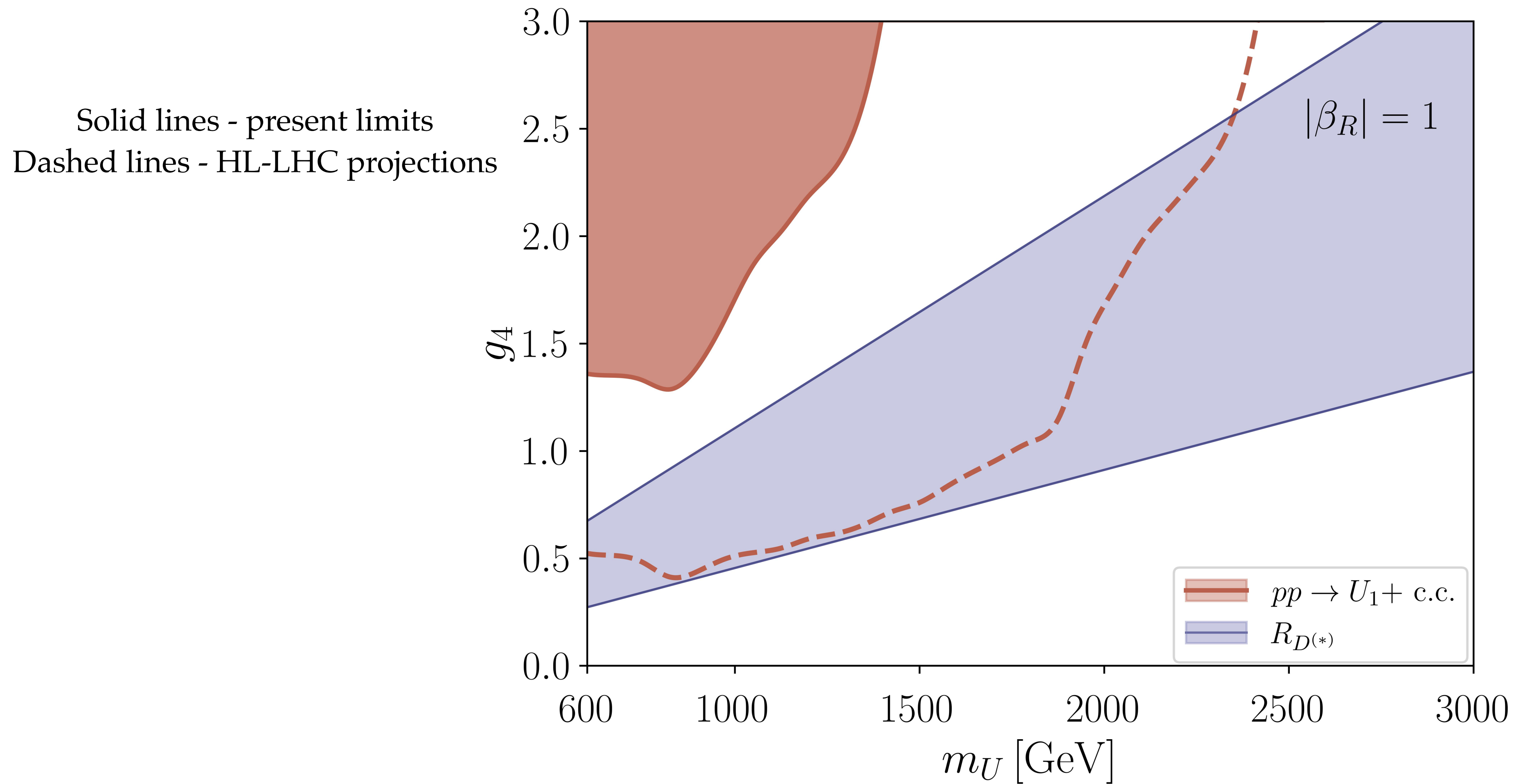
- **First experimental search on resonant scalar LQ production, targeting the $b\tau$ final state**

[CMS Collaboration, arXiv: 2308.06143]



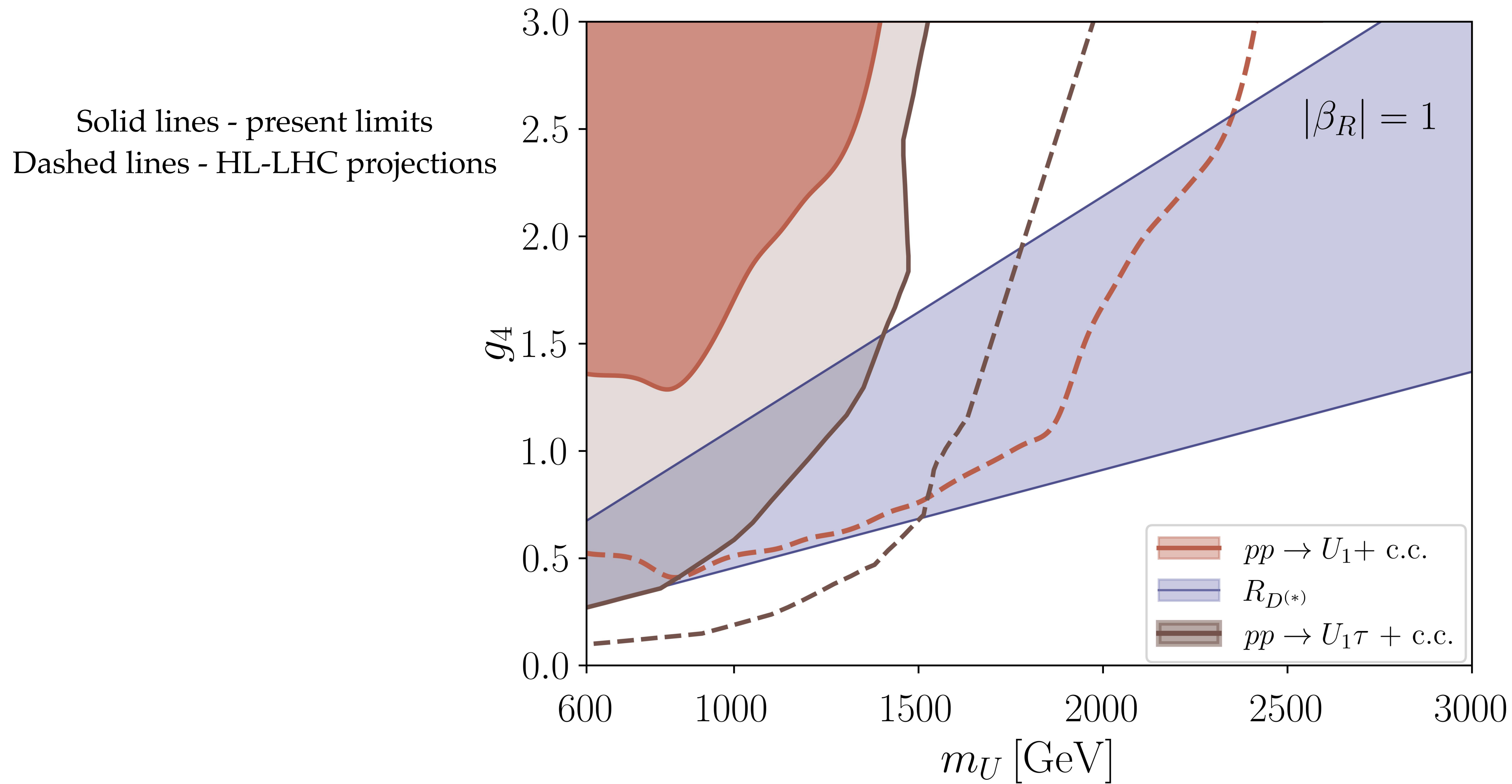
- Data for Brazilian bands (for different couplings) available on HEPData <http://dx.doi.org/10.17182/hepdata.141335>
- Input our signal, check when exclusion bounds are saturated

Simulation details & exclusion limits



[ATLAS, arXiv: 2305.15962]

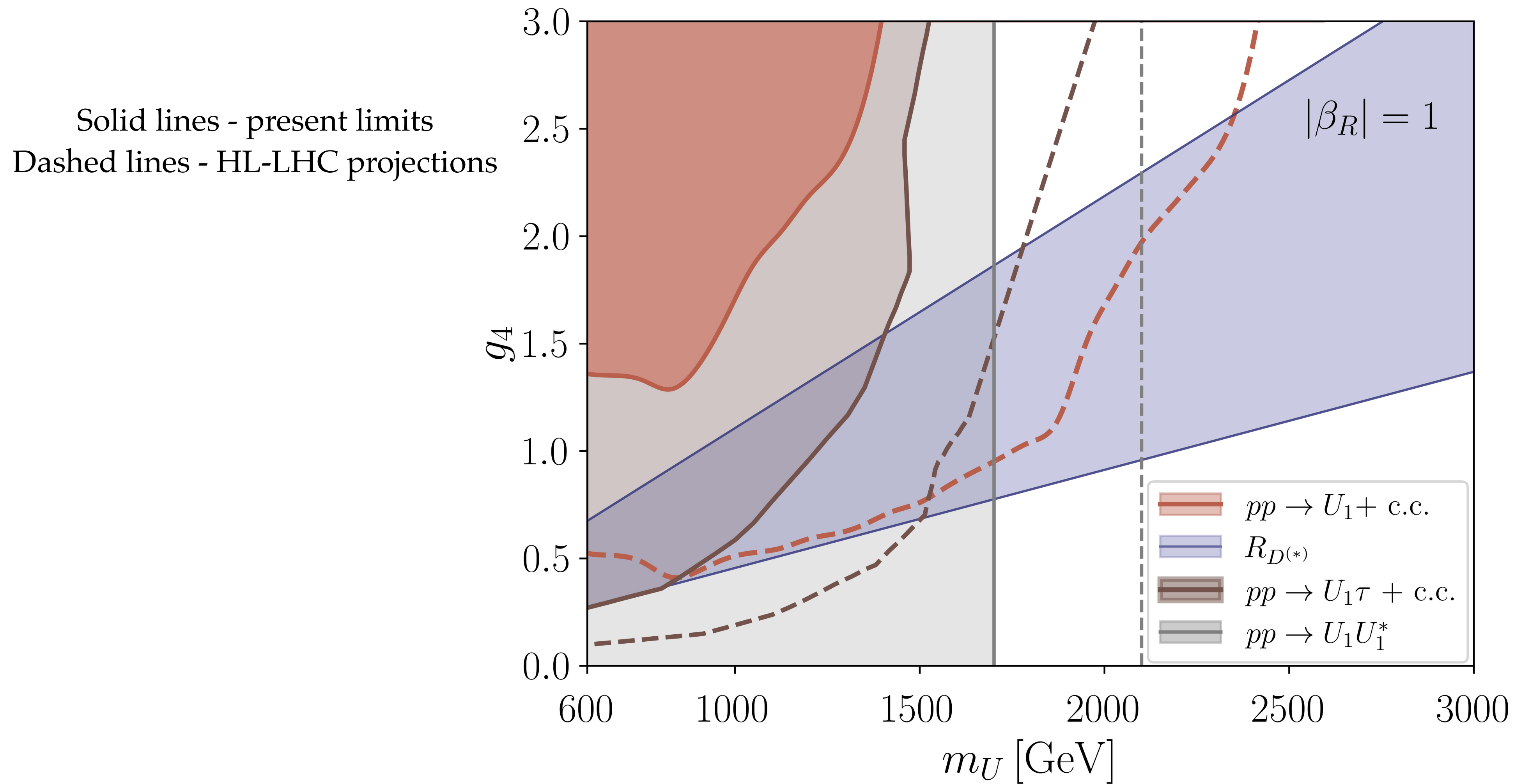
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[ATLAS, arXiv: 2305.15962]

[CMS, arXiv: 2012.04178]

Simulation details & exclusion limits

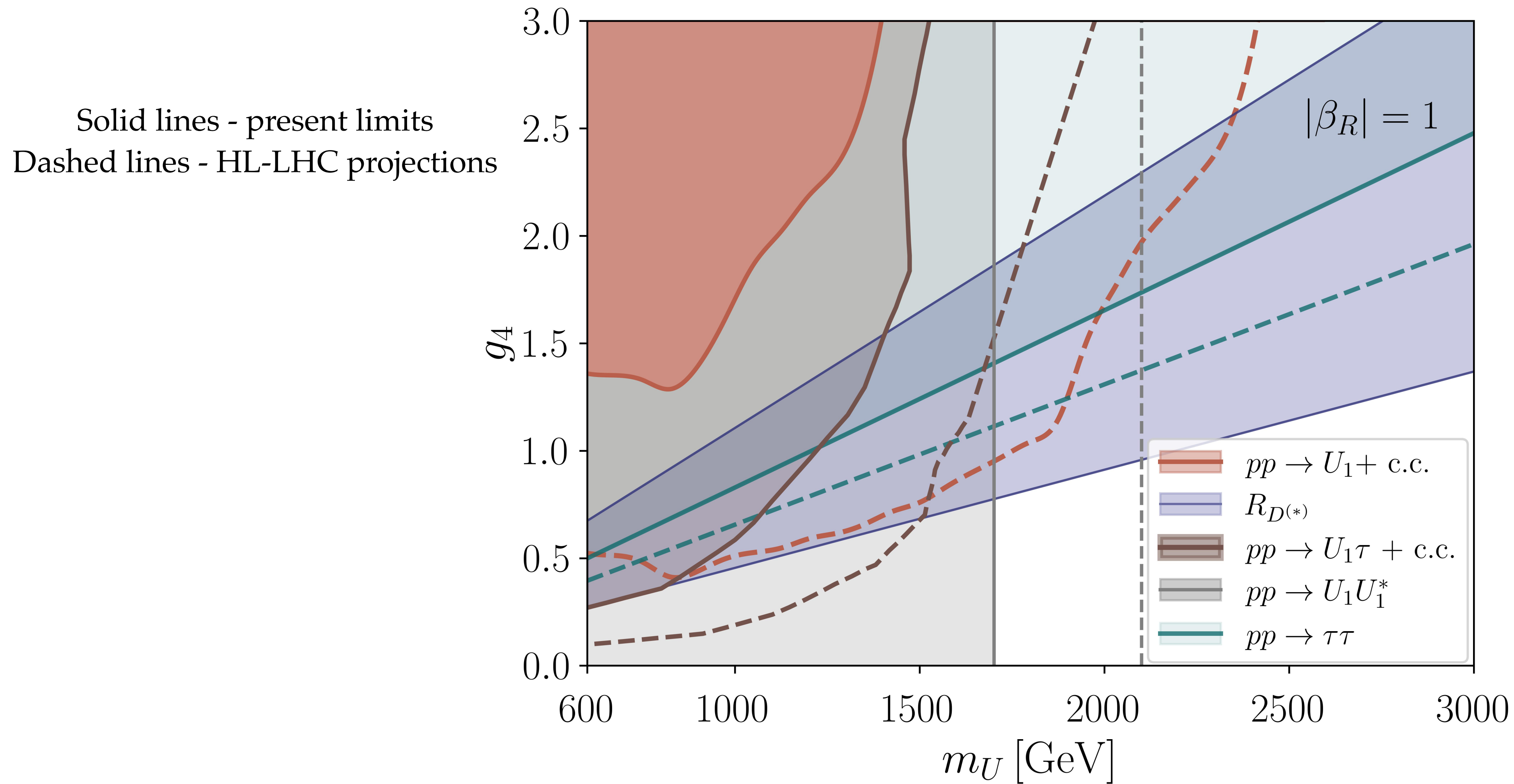


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Simulation details & exclusion limits

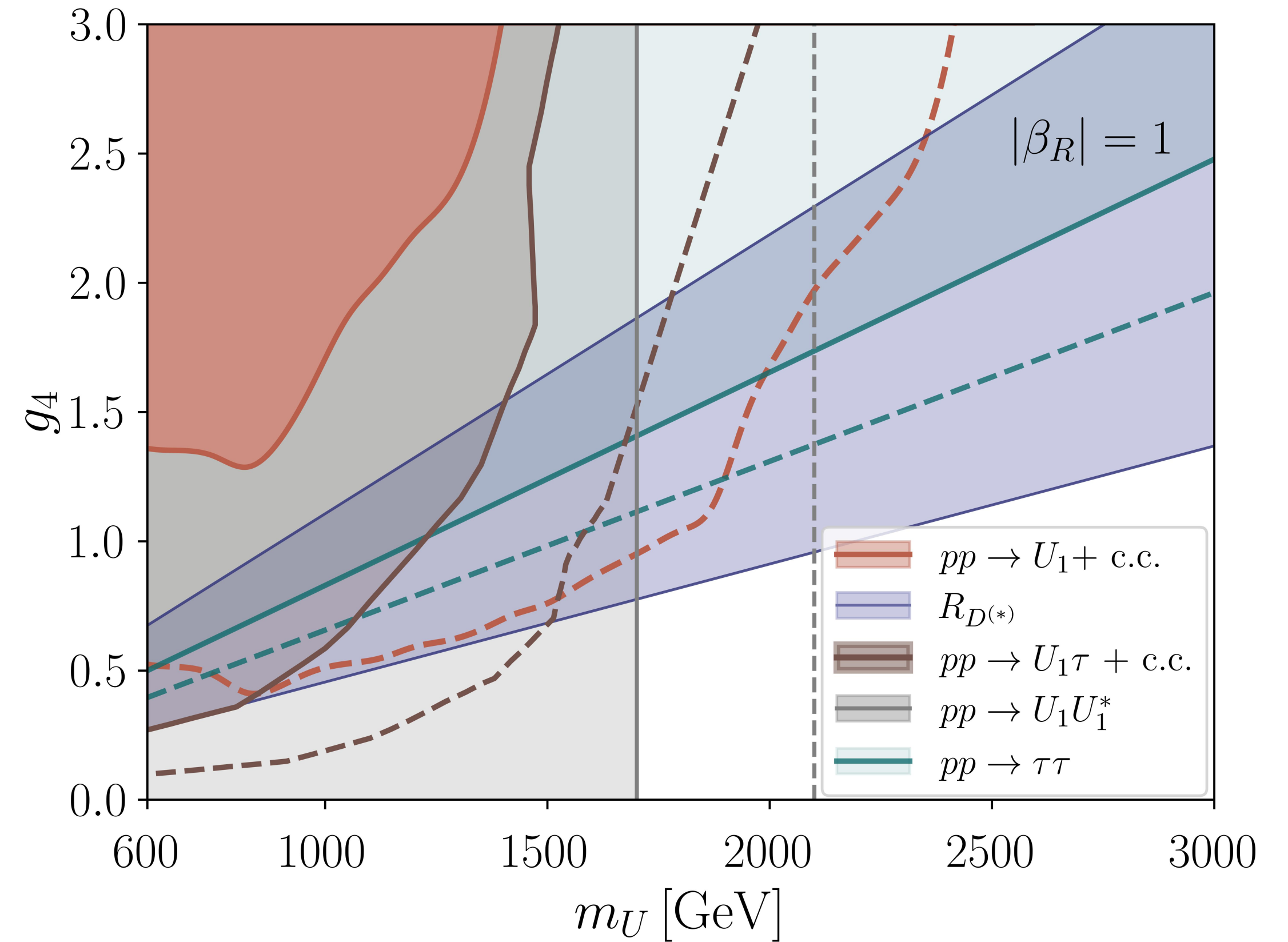
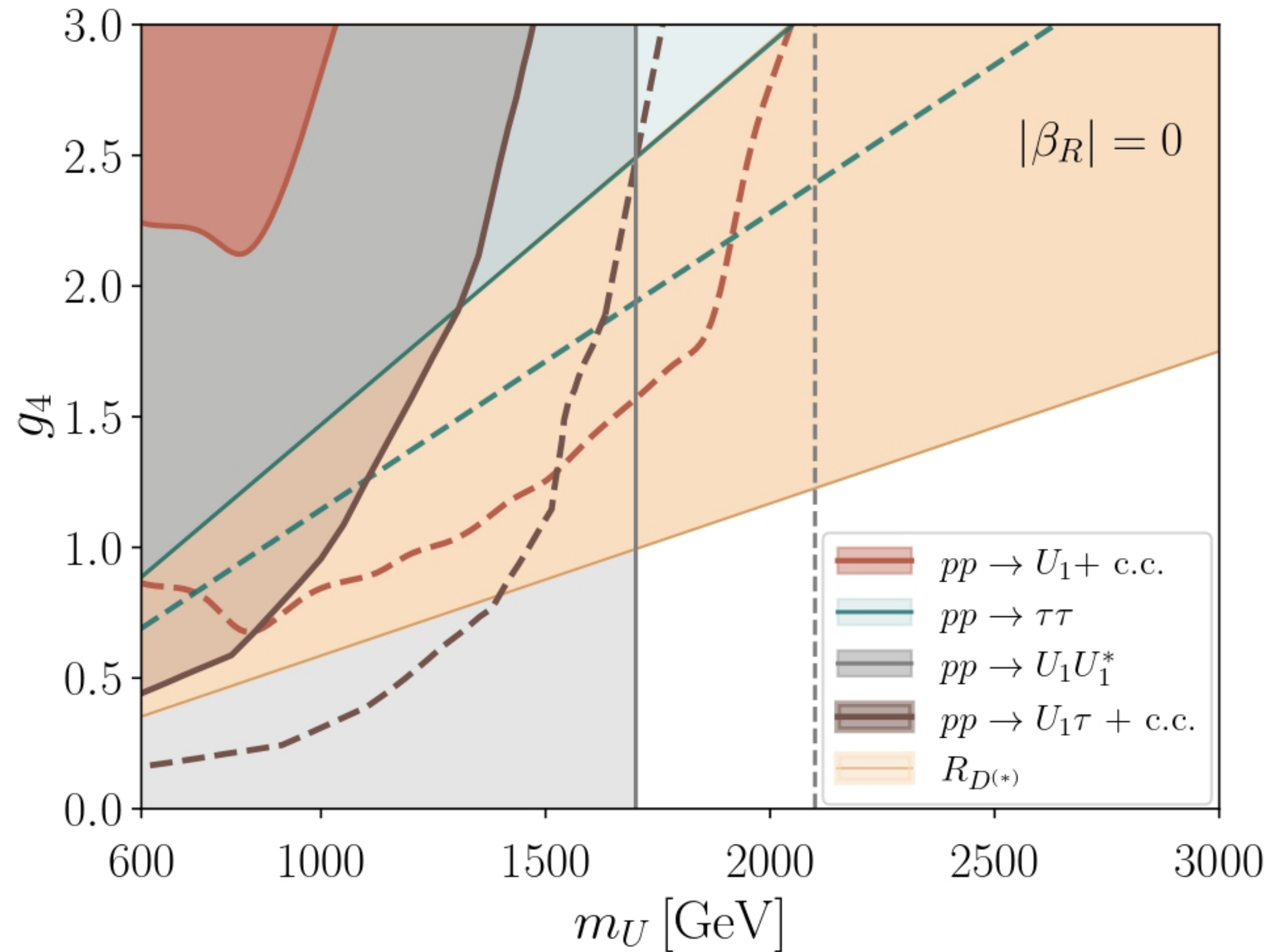


[ATLAS, arXiv: 2305.15962]

[CMS, arXiv: 2012.04178]

[Baker et al, arXiv: 2012.11474]

Simulation details & exclusion limits



- Current bounds not competitive with other production mechanisms, mainly due to the lepton PDF suppression
- BUT: The enhancement of the COM energy very beneficial for the resonant production
- A discovery channel (opposed to t-channel limits)

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Conclusions

- Utilized the first resonant LQ search at LHC
- Phenomenology of lepton-initiated production mechanism is well-understood, simulationwise it is controllable
- Tools for simulations steadily improving, there are more products on the market (for hadronic and lepton colliders)

Resonant scalar LQ production @ LHC in POWHEG-BOX - Buonocore et al; arXiv: 2209.02599

LePDF - PDFs for lepton colliders (e^+e^- , $\mu^+\mu^-$) - Garosi et al; arXiv: 2303.16964

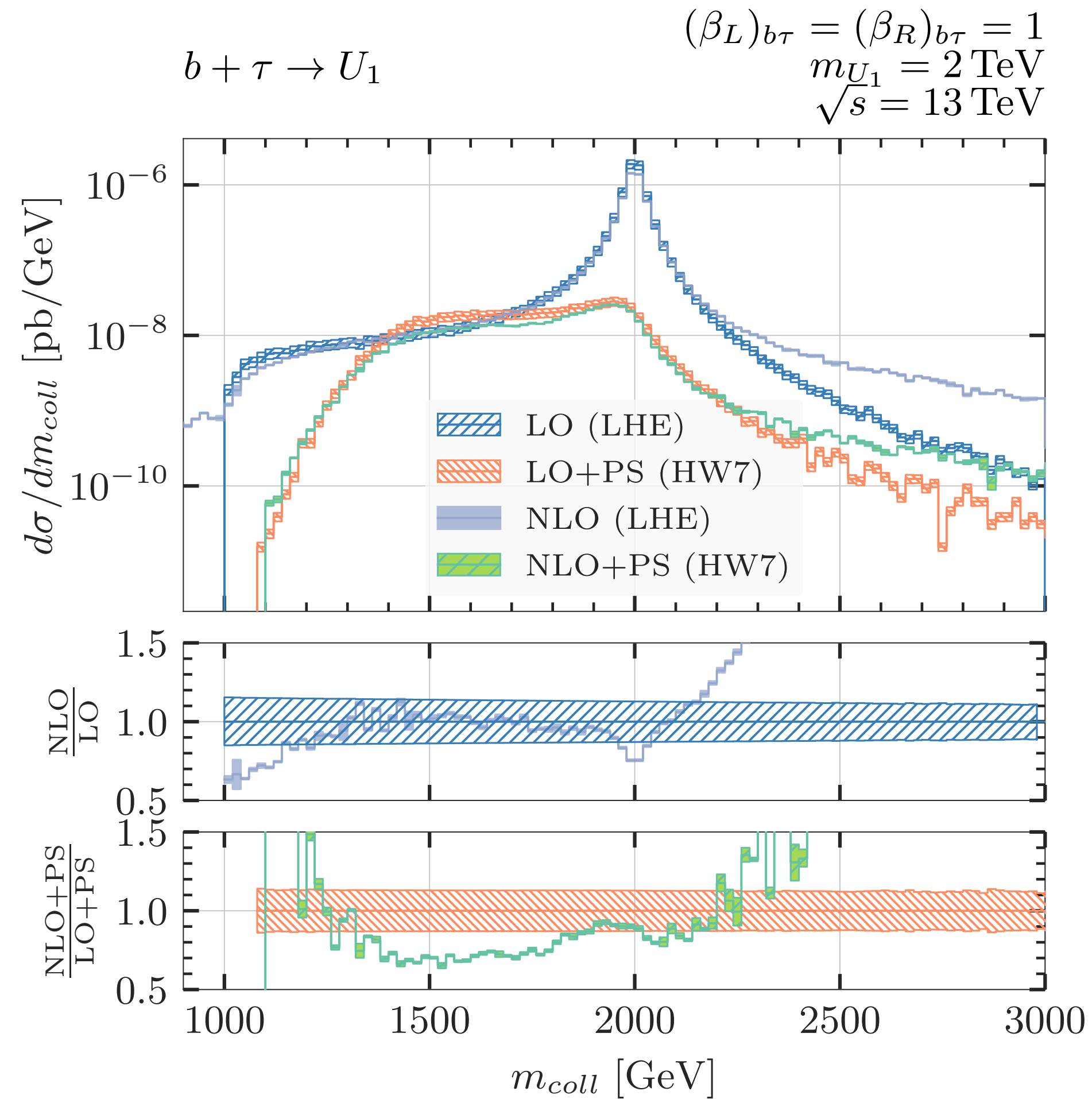
MadGraph/POWHEG/Herwig compatibility - arXiv: 2108.10261

simulate!
Shut up and ~~calculate~~

Thank you!

Backup

Simulation details - Reconstruction



- Taking into account the decay of τ in HERWIG
- Implementation of the CMS cuts in the post-processing:

Variable	$\tau_h + \text{jet}$	$e + \text{jet}$	$\mu + \text{jet}$
$p_T^\ell (\text{GeV})$	>200	>100	>100
$ \eta^\ell $	<2.1	<2.1	<2.1
$p_T^{\text{jet}} (\text{GeV})$	>300	>200	>200
$ \eta^{\text{jet}} $	<2.4	<2.4	<2.4
$p_T^{\text{miss}} (\text{GeV})$	>100	>150	>150
$p_T(\vec{\ell} + \vec{p}_T^{\text{miss}}) (\text{GeV})$	>100	>100	>100
$ \Delta\phi(\ell, \vec{p}_T^{\text{miss}}) (\text{radians})$	<0.3	<0.2	<0.2
$\Delta R(\ell, \text{jet})$	>0.5	>0.5	>0.5

[CMS Collaboration, arXiv: 2308.06143]

- Using the quantity $m_{\text{coll}} = m_{\text{vis}} \sqrt{1 + p_T^{\text{invis}}/p_T^{\text{vis}}}$ to reconstruct the peak

NLO QCD + QED results

- The QCD and QED corrections are of the same order:

[A. Greljo and N. Selimović, arXiv: 2012.02092]

$$f_q \sim f_g \sim \mathcal{O}\left(\sum_n (\alpha_s L)^n\right) \sim \mathcal{O}(1) \quad \alpha_s \sim 1/L \quad L = \log(\mu_F^2/\Lambda^2)$$

Λ - hadronic scale

- QED coupling size: $\alpha_{\text{em}} \sim \alpha_s^2$
 - γ - PDF, 1st order QED effect: $f_\gamma \sim \mathcal{O}(\alpha_{\text{em}} L) \sim \mathcal{O}(\alpha_s)$
 - ℓ - PDF, 2nd order QED effect: $f_\ell \sim \mathcal{O}((\alpha_{\text{em}} L)^2) \sim \mathcal{O}(\alpha_s^2)$

- Including real corrections:

- $q + \ell \rightarrow g + U_1$

$$\sigma_{q\ell}(s) \sim \int (f_q \otimes f_\ell) \hat{\sigma}_{q\ell} \sim 1 \times \alpha_s^2 \times \alpha_s \sim \alpha_s^3$$

- $g + \ell \rightarrow q + U_1$

- $\gamma + q \rightarrow \ell + U_1$

$$\sigma_{q\gamma}(s) \sim \int (f_q \otimes f_\gamma) \hat{\sigma}_{q\gamma} \sim 1 \times \alpha_s \times \alpha_{\text{em}} \sim \alpha_s^3$$

Finite-width effects

- The narrow-width approximation (NWA) can have an $\mathcal{O}(1)$ relative error compared to full Breit-Wigner result for $m_{LQ} \gtrsim 2 \text{ TeV}$ and $g_4 \gtrsim 1.5$!
- The Breit-Wigner prescription for the propagator can be turned on by setting the flag `BWgen = 1`.
- For larger values of coupling, the convolutional area between the PDFs and the Breit-Wigner is also larger, hence we are more sensitive to the low- x region of the PDFs (enhanced by orders of magnitude compared to the values of the PDF at the mass pole)

