

LIGHT NEW PHYSICS & FUNDAMENTAL CONSTANTS

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Challenging the Standard Model

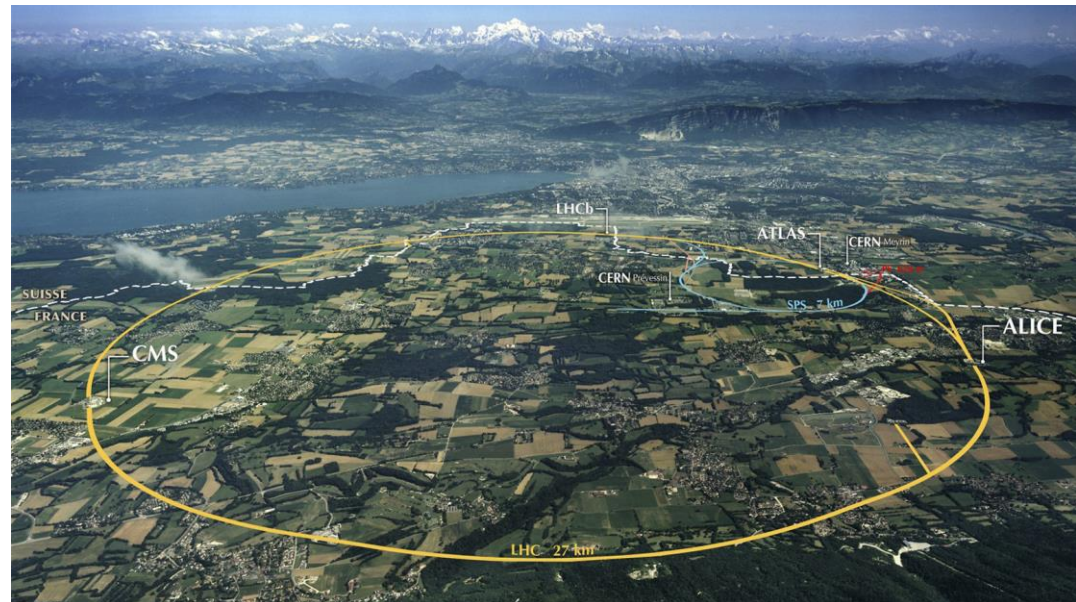
$$E \sim \hbar c/L$$



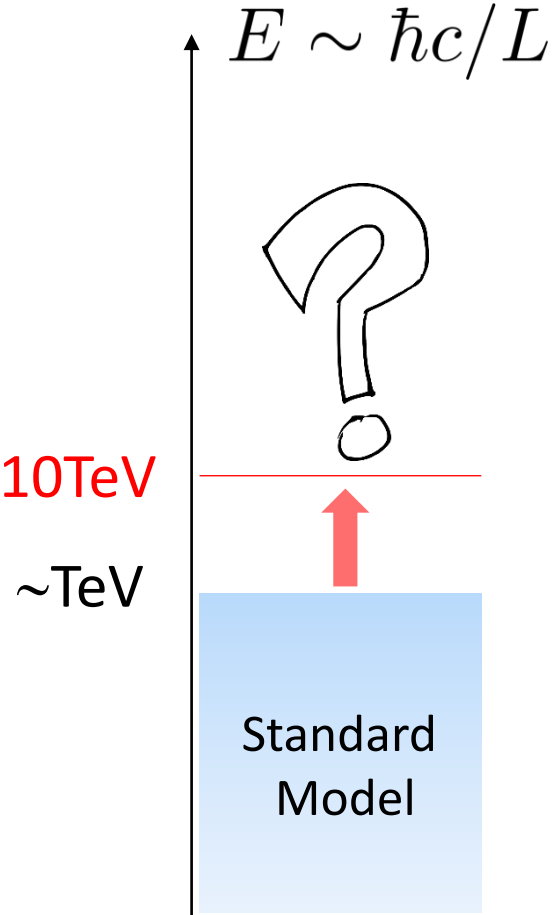
~TeV

Standard
Model

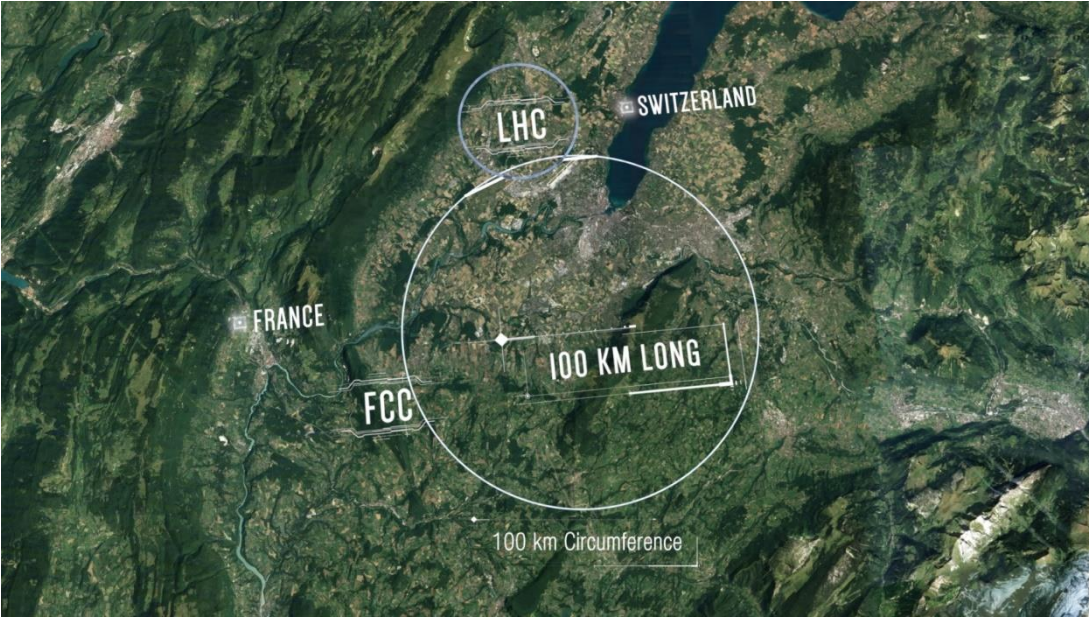
The Standard Model was tested experimentally
up to TeV energies $\sim 10^{-18} \text{m}$
at the **Large Hadron Collider**
(and indirectly at LEP and flavor factories)



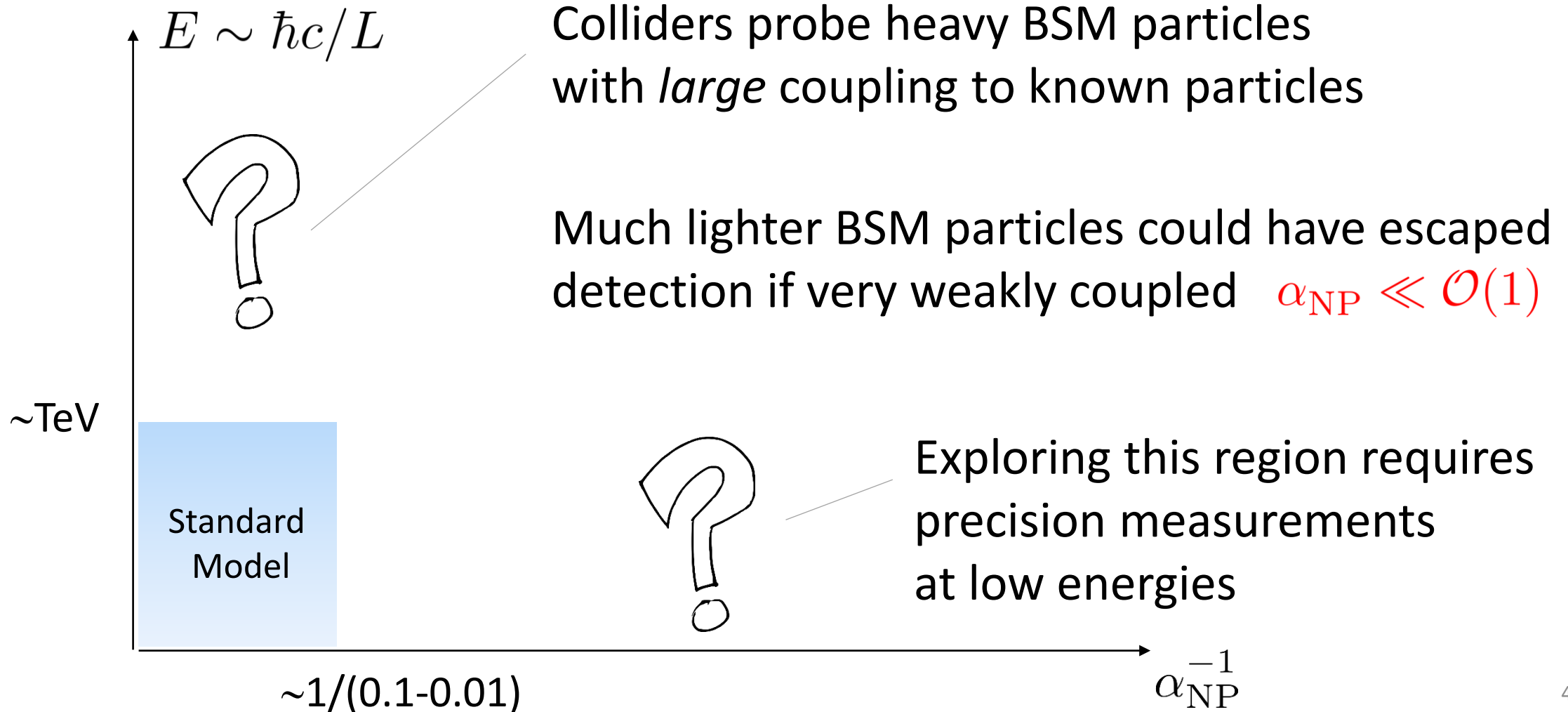
Challenging the Standard Model



Next generation of high-energy colliders plan to extend the reach by a factor ~ 10



Challenging the Standard Model



Light New Physics

Sub-GeV physics beyond the SM is well-motivated theoretically :

- moduli fields in String theory compactifications
- approximate new global symmetries, *e.g.* QCD axion, ALPs, dilaton
- light (mediators for) non-thermal DM, *e.g.* DM produced by freeze-in
- cosmological solutions to the Higgs hierarchy problem, *e.g.* relaxion

The Precision Frontier

Atomic Spectrometry

Accuracy of spectroscopic measurements is very (very!) high:

18 digits in optical clock frequency comparison, e.g. ytterbium vs. strontium

BACON coll. (Nature **591**, 2021)

$$\nu_{\text{Yb}}/\nu_{\text{Sr}} = 1.207\ 507\ 039\ 343\ 337\ 8482(82)$$

Can it be used to probe BSM physics?

Atomic Spectrometry

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18 digits in optical clock frequency comparison, e.g. ytterbium vs. strontium

BACON coll. (Nature **591**, 2021)

$$\nu_{\text{Yb}}/\nu_{\text{Sr}} = 1.207\ 507\ 039\ 343\ 337\ 8482(82)$$

New bosons coupled to electrons and/or nuclei in a *parity-preserving* way will induce an exotic atomic force

$$V(r) = \begin{array}{c} \begin{array}{c} \text{Diagram 1: } \gamma \\ \text{Two fermion lines connected by a wavy line labeled } \gamma \end{array} \\ -\alpha/r \end{array} + \begin{array}{c} \begin{array}{c} \text{Diagram 2: } \phi \\ \text{Two fermion lines connected by a wavy line labeled } \phi \end{array} \\ -(-1)^s \alpha_{\text{NP}} e^{-m_\phi r} / r \end{array}$$

Extracting New Physics

Nuclear finite size is a major obstacle for $m_\phi \gtrsim \Lambda_{\text{QCD}} \simeq 200 \text{ MeV}$

$$\nu_{\text{Yb}}/\nu_{\text{Sr}} = 1.207\ 507\ 039\ 348\ 33$$

**MBPT theory
uncertainty**

$$\sim \mathcal{O}(1\%)$$

finite nuclear size

$$\sim \alpha m_e / \Lambda_{\text{QCD}}$$

The range of the NP interaction must extend beyond the nucleus.

(NP interactions that breaks P or CP are well-known exceptions. Those are better probed in APV and EDM experiments)

Many-body effects need to be controlled for multi-electron atoms.

A Possible Strategy

Focus on few-body systems where precision QED calculation is possible, *e.g.* Hydrogen, μH , Helium or positronium (e^+e^-), muonium (μ^+e^-)

Hydrogen and muonic hydrogen spectral lines are best measured. They have the highest sensitivity to NP.

However, these lines are used to determine the values of fundamental (= not predicted by the SM) constants, *assuming the SM holds!*

FCs and NP parameters must be determined together.

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and J. Zupan (U. Cincinnati)

The BSM CODATA

The Hydrogen Frontier

Measurements of atomic lines
in **hydrogen** are very precise:

$$\nu_{1S-2S} = 2\,466\,061\,413\,187\,035(10) \text{ Hz}$$
$$u_\nu = 4.2 \times 10^{-15} \quad \text{Parthey et al. (2011)}$$

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QED prediction is even better:

$$\nu_{1S-2S} = \frac{3}{4} \frac{R_\infty c}{(1+m_e/m_p)} \left[1 + \delta_{1S-2S}^{\text{QED}}(\alpha) + \delta_{1S-2S}^{\text{FNS}}(r_p) \right]$$

$R_\infty \equiv \alpha^2 m_e c / 2h$

TH uncertainty $\sim 2 \text{ Hz}$

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$R_\infty \equiv \alpha^2 m_e c / 2h$

TH uncertainty ~ 2 Hz

limited by proton radius

Direct comparison fixes the Rydberg constant:

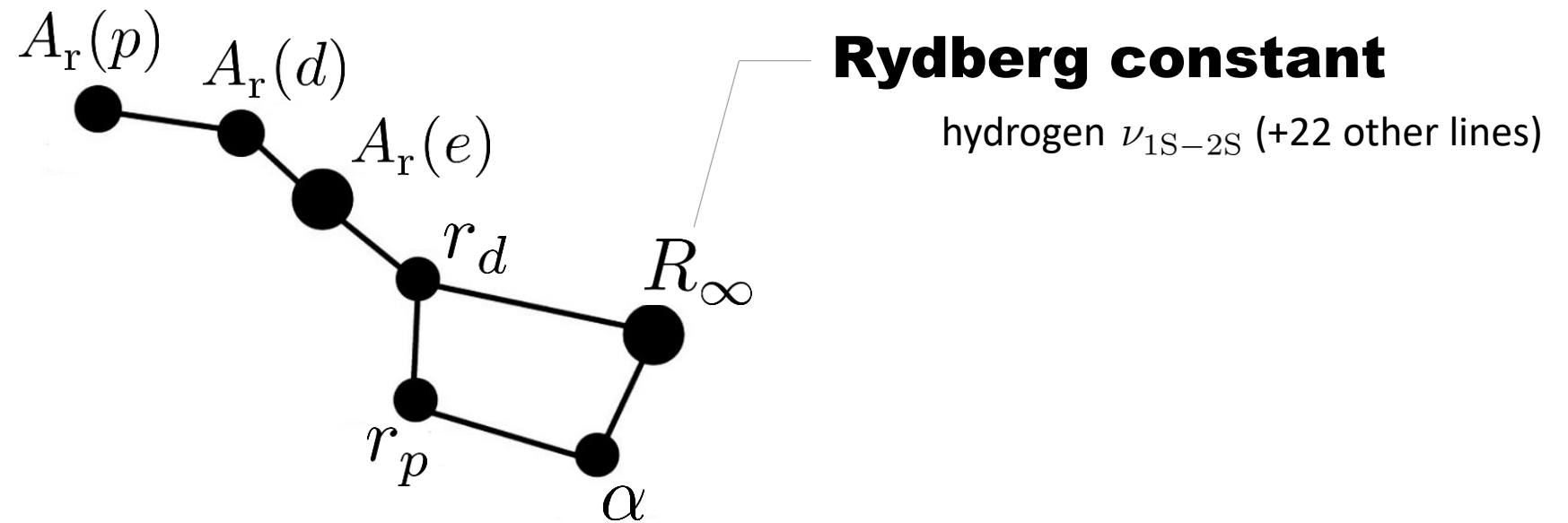
$$u_{R_\infty} = 1.9 \times 10^{-12}$$

Tiesinga et al.
[CODATA 2018]

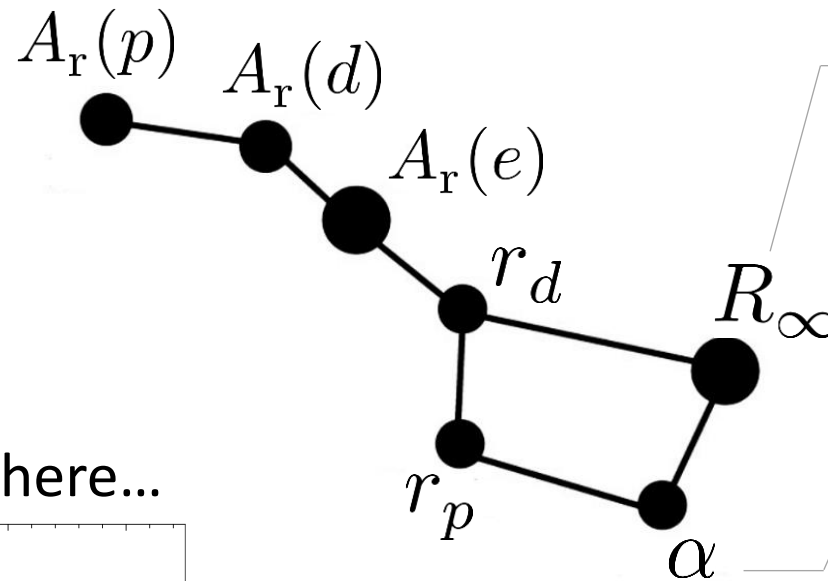
$$R_\infty c = 3.289\,841\,960\,2508(64) \times 10^{15} \text{ Hz}$$

most precisely known
fundamental constant
in physics!

A Constellation of Constants



A Constellation of Constants



Rydberg constant

hydrogen ν_{1S-2S} (+22 other lines)

fine structure constant

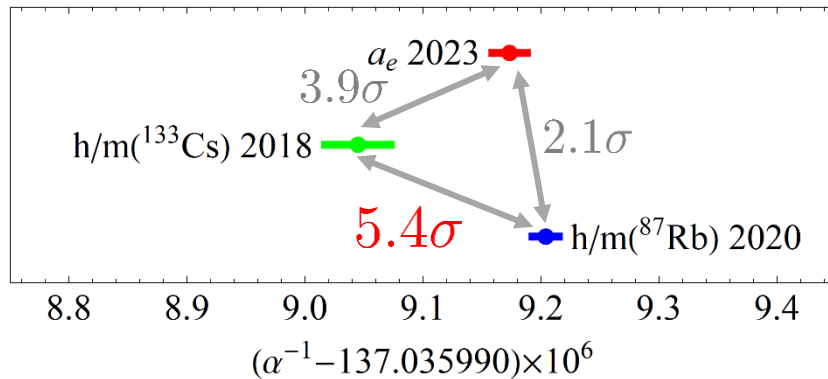
electron $g - 2$ [Fan et al. \[2023\]](#)
or atomic recoil

$$\alpha^2 = 2R_\infty/c \times \frac{m}{m_e} \times \frac{h}{m}$$

⁸⁷Rb [Morel et al. \[2020\]](#)

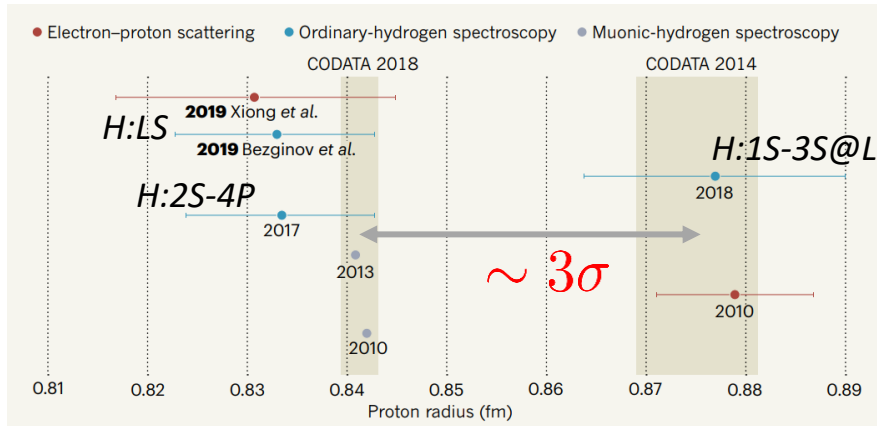
¹³³Cs [Parker et al. \[2018\]](#)

There are **tensions** there...



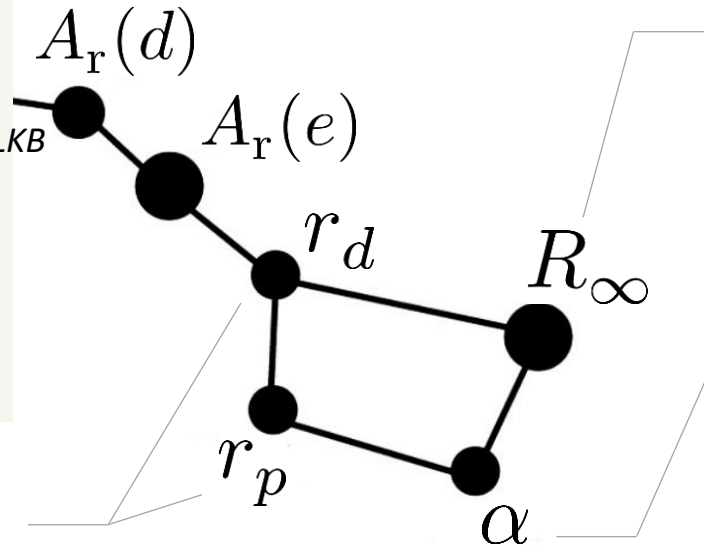
A Constellation of Constants

still a **proton size puzzle**...



**proton|deuteron
(charge) radius**

muonic hydrogen|deuterium Lamb shifts
 or ordinary hydrogen|deuterium lines
 or e-proton|e-deuteron scattering data



Rydberg constant

hydrogen ν_{1S-2S} (+22 other lines)

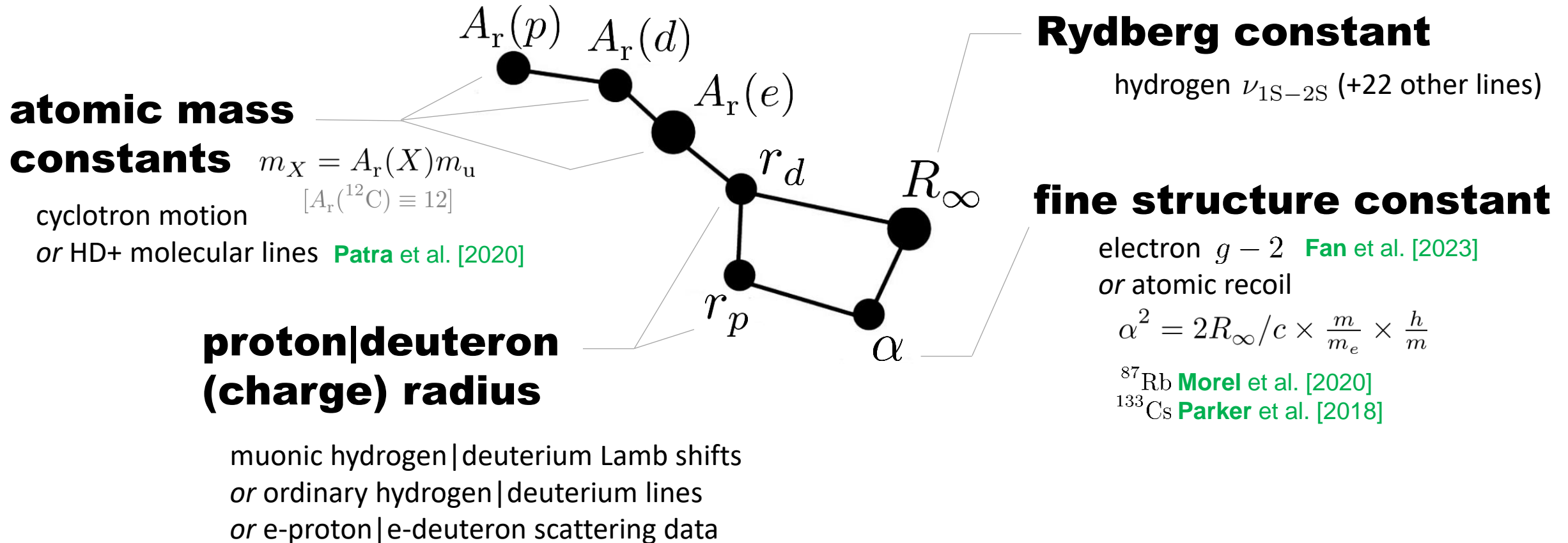
fine structure constant

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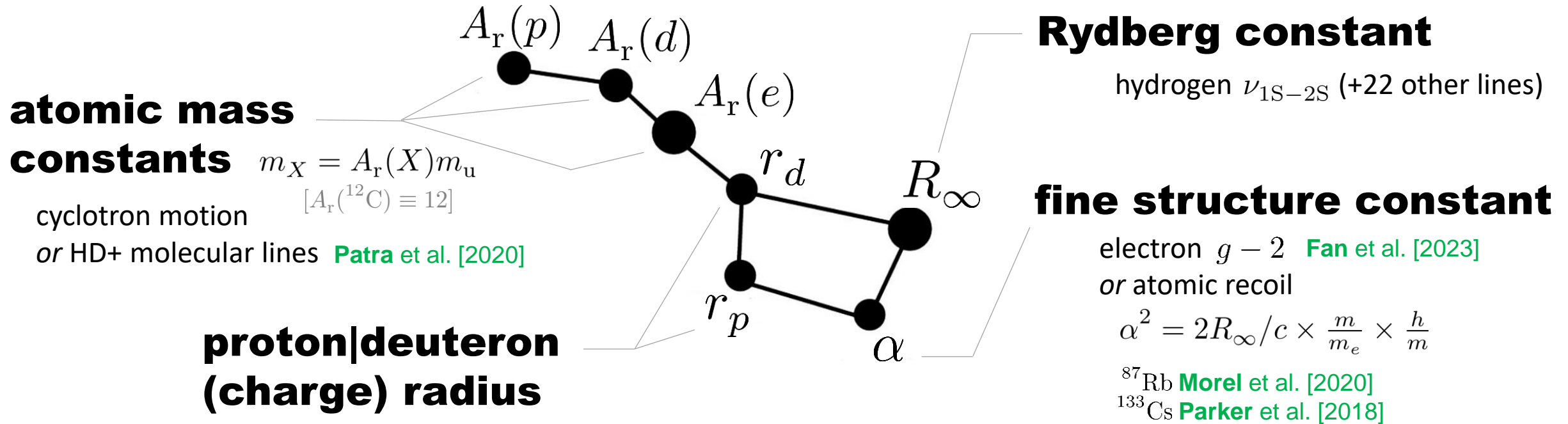
$$\alpha^2 = 2R_\infty/c \times \frac{m}{m_e} \times \frac{h}{m}$$

^{87}Rb **Morel et al. [2020]**
 ^{133}Cs **Parker et al. [2018]**

A Constellation of Constants



A Constellation of Constants



to be determined together in a global fit

→ **CODATA** recommended values

<https://pml.nist.gov/cuu/Constants/>

CODATA 2018 values

TABLE XXXI. The CODATA recommended values of the fundamental constants of physics and chemistry based on the 2018 adjustment.

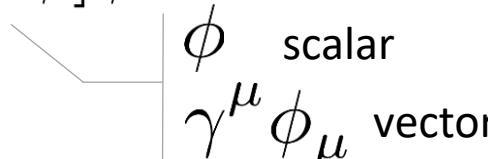
Quantity	Symbol	Numerical value	Unit	Relative std. uncert. u_r
Rydberg frequency $\alpha^2 m_e c^2 / 2h = E_h / 2h$	cR_∞	$3.289\,841\,960\,2508(64) \times 10^{15}$	Hz	1.9×10^{-12}
deuteron mass	m_d	$3.343\,583\,7724(10) \times 10^{-27}$	kg	3.0×10^{-10}
		$2.013\,553\,212\,745(40)$	u	2.0×10^{-11}
electron mass	m_e	$9.109\,383\,7015(28) \times 10^{-31}$	kg	3.0×10^{-10}
		$5.485\,799\,090\,65(16) \times 10^{-4}$	u	2.9×10^{-11}
proton mass	m_p	$1.672\,621\,923\,69(51) \times 10^{-27}$	kg	3.1×10^{-10}
		$1.007\,276\,466\,621(53)$	u	5.3×10^{-11}
fine-structure constant $e^2 / 4\pi\epsilon_0 \hbar c$ inverse fine-structure constant	α	$7.297\,352\,5693(11) \times 10^{-3}$		1.5×10^{-10}
	α^{-1}	$137.035\,999\,084(21)$		1.5×10^{-10}
deuteron rms charge radius	r_d	$2.12799(74) \times 10^{-15}$	m	3.5×10^{-4}
proton rms charge radius	r_p	$8.414(19) \times 10^{-16}$	m	2.2×10^{-3}

This accuracy relies on *assuming the SM*. Is this robust to **BSM**?

BSM CODATA

New particles below $\sim \text{GeV}$ will affect the observables used to determine the fundamental constants:

$$\mathcal{L}_{\text{int}} = \sum_{\psi = e, \mu, p, n} g_{\psi} \bar{\psi} [\Gamma \cdot \phi] \psi$$



ϕ scalar
 $\gamma^{\mu} \phi_{\mu}$ vector

BSM CODATA

New particles below $\sim \text{GeV}$ will affect the observables used to determine the fundamental constants:

$$\mathcal{L}_{\text{int}} = \sum_{\psi = e, \mu, p, n} g_{\psi} \bar{\psi} [\Gamma \cdot \phi] \psi$$

ϕ scalar
 $\gamma^{\mu} \phi_{\mu}$ vector



$$V_{\text{NP}}^{ij} = (-1)^{s+1} \alpha_{\phi} q_i q_j \frac{e^{-m_{\phi} r}}{r}$$

spin s

$\equiv \frac{|g_e g_p|}{4\pi} \geq 0$

$q_i \equiv \frac{g_i}{\sqrt{|g_e g_p|}}$

Yukawa potentials

contributing to spectral lines

$$\text{hydrogen} \sim \alpha_{\phi} q_e q_p = \pm \alpha_{\phi}$$

$$\text{deuterium} \sim \alpha_{\phi} [1 + q_e q_n]$$

$$\mu\text{H}/\mu\text{D} \sim \alpha_{\phi} q_{\mu} q_{p/n}$$

$$\text{HD}^+ \sim \alpha_{\phi} [1, q_e q_n, q_p q_n]$$

BSM CODATA

New particles below $\sim \text{GeV}$ will affect the observables used to determine the fundamental constants:

$$\mathcal{L}_{\text{int}} = \sum_{\psi = e, \mu, p, n} g_{\psi} \bar{\psi} [\Gamma \cdot \phi] \psi$$

ϕ scalar
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$$V_{\text{NP}}^{ij} = (-1)^{\text{spin} + 1} \alpha_{\phi} q_i q_j \frac{e^{-m_{\phi} r}}{r}$$

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hydrogen $\sim \alpha_{\phi} q_e q_p = \pm \alpha_{\phi}$

deuterium $\sim \alpha_{\phi} [1 + q_e q_n]$

$\mu\text{H}/\mu\text{D} \sim \alpha_{\phi} q_{\mu} q_p / n$

$\text{HD}^+ \sim \alpha_{\phi} [1, q_e q_n, q_p q_n]$

One-loop correction to a_e

$$\sim \alpha_{\phi} q_e^2 / 4\pi$$

BSM CODATA

New particles below $\sim \text{GeV}$ will affect the observables used to determine the fundamental constants:

$$\mathcal{L}_{\text{int}} = \sum_{\psi = e, \mu, p, n} g_{\psi} \bar{\psi} [\Gamma \cdot \phi] \psi$$

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$$V_{\text{NP}}^{ij} = (-1)^{\text{spin} + 1} \alpha_{\phi} q_i q_j \frac{e^{-m_{\phi} r}}{r} \equiv \frac{|g_e g_p|}{4\pi} \geq 0$$

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Yukawa potentials

contributing to spectral lines

One-loop correction to a_e

$$\sim \alpha_{\phi} q_e^2 / 4\pi$$

Theoretical prediction for an observable \mathcal{O} :

$$\mathcal{O}_{\text{th}} = \mathcal{O}_{\text{SM}}(g_{\text{SM}}) + \mathcal{O}_{\text{NP}}(g_{\text{SM}}, \alpha_{\phi}, m_{\phi}) + \delta \mathcal{O}_{\text{th}}$$

$R_{\infty}, \alpha, r_p, \dots$

evaluated @LO in α_{ϕ}

TH uncert. 24

Datasets

CODATA18 ← used for validation

Hydrogen/Deuterium

Label	Input datum	Value (kHz)
A1	$\nu_{\text{H}}(2S_{1/2} - 4S_{1/2}) - \frac{1}{4}\nu_{\text{H}}(1S_{1/2} - 2S_{1/2})$	4 797 338(10)
A2	$\nu_{\text{H}}(2S_{1/2} - 4D_{5/2}) - \frac{1}{4}\nu_{\text{H}}(1S_{1/2} - 2S_{1/2})$	6 490 144(24)
A3	$\nu_{\text{D}}(2S_{1/2} - 4S_{1/2}) - \frac{1}{4}\nu_{\text{D}}(1S_{1/2} - 2S_{1/2})$	4 801 693(20)
A4	$\nu_{\text{D}}(2S_{1/2} - 4D_{5/2}) - \frac{1}{4}\nu_{\text{D}}(1S_{1/2} - 2S_{1/2})$	6 494 841(41)
A5	$\nu_{\text{D}}(1S_{1/2} - 2S_{1/2}) - \nu_{\text{H}}(1S_{1/2} - 2S_{1/2})$	670 994 334.606(15)
A6	$\nu_{\text{H}}(1S_{1/2} - 2S_{1/2})$	2 466 061 413 187.035(10)
A7	$\nu_{\text{H}}(1S_{1/2} - 2S_{1/2})$	2 466 061 413 187.018(11)
A8	$\nu_{\text{H}}(1S_{1/2} - 3S_{1/2})$	2 922 743 278 659(17)
A9	$\nu_{\text{H}}(2S_{1/2} - 4P)$	616 520 931 626.8(2.3)
A10	$\nu_{\text{H}}(2S_{1/2} - 8S_{1/2})$	770 649 350 012.0(8.6)
A11	$\nu_{\text{H}}(2S_{1/2} - 8D_{3/2})$	770 649 504 450.0(8.3)
A12	$\nu_{\text{H}}(2S_{1/2} - 8D_{5/2})$	770 649 561 584.2(6.4)
A13	$\nu_{\text{D}}(2S_{1/2} - 8S_{1/2})$	770 859 041 245.7(6.9)
A14	$\nu_{\text{D}}(2S_{1/2} - 8D_{3/2})$	770 859 195 701.8(6.3)
A15	$\nu_{\text{D}}(2S_{1/2} - 8D_{5/2})$	770 859 252 849.5(5.9)
A16	$\nu_{\text{H}}(2S_{1/2} - 12D_{3/2})$	799 191 710 472.7(9.4)
A17	$\nu_{\text{H}}(2S_{1/2} - 12D_{5/2})$	799 191 727 403.7(7.0)
A18	$\nu_{\text{D}}(2S_{1/2} - 12D_{3/2})$	799 409 168 038.0(8.6)
A19	$\nu_{\text{D}}(2S_{1/2} - 12D_{5/2})$	799 409 184 966.8(6.8)
A20	$\nu_{\text{H}}(2S_{1/2} - 6S_{1/2}) - \frac{1}{4}\nu_{\text{H}}(1S_{1/2} - 3S_{1/2})$	4 197 604(21)
A21	$\nu_{\text{H}}(2S_{1/2} - 6D_{5/2}) - \frac{1}{4}\nu_{\text{H}}(1S_{1/2} - 4S_{1/2})$	4 699 099(10)
A22	$\nu_{\text{H}}(1S_{1/2} - 3S_{1/2})$	2 922 743 278 678(13)
A23	$\nu_{\text{H}}(1S_{1/2} - 3S_{1/2})$	2 922 743 278 671.5(2.6)
A24	$\nu_{\text{H}}(2S_{1/2} - 4P_{1/2}) - \frac{1}{4}\nu_{\text{H}}(1S_{1/2} - 2S_{1/2})$	4 664 269(15)
A25	$\nu_{\text{H}}(2S_{1/2} - 4P_{3/2}) - \frac{1}{4}\nu_{\text{H}}(1S_{1/2} - 2S_{1/2})$	6 035 373(10)
A26	$\nu_{\text{H}}(2S_{1/2} - 2P_{3/2})$	9 911 200(12)
A27	$\nu_{\text{H}}(2P_{1/2} - 2S_{1/2})$	1 057 862(20)
A28	$\nu_{\text{H}}(2P_{1/2} - 2S_{1/2})$	1 057 845.0(9.0)
A29	$\nu_{\text{H}}(2P_{1/2} - 2S_{1/2})$	1 057 829.8(3.2)

g_e -2, masses...

Label	Input datum	Value	Rel. uncert.
D1	$a_e \equiv \frac{1}{2}(g - 2)_e$	$1.159\,652\,180\,73(28) \times 10^{-3}$	2.4×10^{-10}
D2	δ_e	$0.000(18) \times 10^{-12}$	1.5×10^{-11}
D3	$h/m_{\text{Rb}}(^{87}\text{Rb})$	$4.591\,359\,272\,9(57) \times 10^{-9} \text{ m}^2\text{s}^{-1}$	1.2×10^{-8}
D4	$h/m_{\text{Cs}}(^{133}\text{Cs})$	$3.002\,369\,472\,1(12) \times 10^{-9} \text{ m}^2\text{s}^{-1}$	4.0×10^{-9}
D5	$A_r(^{87}\text{Rb})$	86.909 180 531 2(65)	
D6	$A_r(^{133}\text{Cs})$	132.905 451 961 0(86)	
D7	$\omega_s/\omega_c(^{12}\text{C}^{5+})$	4376.210 500 87(12)	
D8	$\Delta E_{\text{B}}(^{12}\text{C}^{5+})/hc$	$43.563\,233(25) \times 10^7 \text{ m}^{-1}$	
D9	δ_{C}	$0.0(2.5) \times 10^{-11}$	
D10	$\omega_s/\omega_c(^{28}\text{Si}^{13+})$	3912.866 064 84(19)	
D11	$A_r(^{28}\text{Si})$	27.976 926 534 99(52)	
D12	$\Delta E_{\text{B}}(^{28}\text{Si}^{13+})/hc$	$420.6467(85) \times 10^7 \text{ m}^{-1}$	
D13	δ_{Si}	$0.0(1.7) \times 10^{-9}$	
D14	$\omega_c(\text{d})/\omega_c(^{12}\text{C}^{6+})$	0.992 996 654 743(20)	
D15	$\omega_c(^{12}\text{C}^{6+})/\omega_c(p)$	0.503 776 367 662(17)	
D19	$A_r(^1\text{H})$	1.007 825 032 241(94)	
D21	$\Delta E_{\text{B}}(^1\text{H}^+)/hc$	$1.096\,787\,717\,430\,7(10) \times 10^7 \text{ m}^{-1}$	
D23	$\Delta E_{\text{B}}(^{12}\text{C}^{6+})/hc$	$83.083\,850(25) \times 10^7 \text{ m}^{-1}$	

8.5×10^{-12}
 4.9×10^{-6}
 2.2×10^{-6}
 4.4×10^{-12}

$\mu\text{H}/\mu\text{D}$

Label	Input datum	Value	Rel. uncert.
C1	$E_{\text{LS}}(\mu\text{H})$	202.3706(23) meV	1.1×10^{-5}
C2	$E_{\text{LS}}(\mu\text{D})$	202.8785(34) meV	1.7×10^{-5}
C7	$\delta E_{\text{LS}}(\mu\text{H})$	0.0000(129) meV	6.4×10^{-5}
C8	$\delta E_{\text{LS}}(\mu\text{D})$	0.0000(210) meV	1.0×10^{-4}
C9	r_p	0.880(20) fm	2.3×10^{-2}
C10	r_d	2.111(19) fm	9.0×10^{-3}

DATA22

including post-CODATA18 improvements from

Hydrogen, HD+, pbar-He μHe lines, g_e -2 and atomic masses

Label	Input datum	Value	Rel. uncert.	Reference
A30	$\nu_{\text{H}}(1S_{1/2} - 3S_{1/2})$	2 922 743 278 665.79(72) kHz	2.5×10^{-13}	Grinin <i>et al.</i> [21]
A31	$\nu_{\text{H}}(2S_{1/2} - 8D_{5/2})$	770 649 561 570.9(2.0) kHz	2.6×10^{-12}	Brandt <i>et al.</i> [20]
D1	$a_e \equiv \frac{1}{2}(g - 2)_e$	$1.159\,652\,180\,59(13) \times 10^{-3}$	1.1×10^{-10}	Fan <i>et al.</i> [70]
D3	$h/m_{\text{Rb}}(^{87}\text{Rb})$	$4.591\,359\,258\,90(65) \times 10^{-9} \text{ m}^2\text{s}^{-1}$	1.4×10^{-10}	Morel <i>et al.</i> [69]
D5	$A_r(^{87}\text{Rb})$	86.909 180 529(6)	6.9×10^{-11}	AME 2020 [73]
D6	$A_r(^{133}\text{Cs})$	132.905 451 958(8)	6.0×10^{-11}	AME 2020 [73]
D9	δ_{C}	$0.0(9.4) \times 10^{-12}$	4.9×10^{-12}	Czarnecki <i>et al.</i> [71]
D13	δ_{Si}	$0.0(5.8) \times 10^{-10}$	2.8×10^{-10}	Czarnecki <i>et al.</i> [71]
D11	$A_r(^{28}\text{Si})$	27.976 926 534 42(55)	2.0×10^{-11}	AME 2020 [73]
D14	$A_r(^2\text{H})$	2.014 101 777 844(15)	7.4×10^{-12}	AME 2020 [73]
D15	$\Delta E_{\text{B}}(^2\text{H}^+)/hc$	$1.097\,086\,145\,529\,9(10) \times 10^7 \text{ m}^{-1}$	9.1×10^{-13}	NIST ASD 2021 [62]
D19	$A_r(^1\text{H})$	1.007 825 031 898(14)	1.4×10^{-11}	AME 2020 [73]
D23	$\Delta E_{\text{B}}(^{12}\text{C}^{6+})/hc$	---	---	---
E1	$\nu_{\text{HD}^+}((0, 0) - (0, 1))$	1 314 925 752.910(17) kHz	1.3×10^{-11}	Alighanbari <i>et al.</i> [33]
E2	$\nu_{\text{HD}^+}((0, 0) - (1, 1))$	58 605 052 164.24(86) kHz	1.5×10^{-11}	Kortunov <i>et al.</i> [35]
E3	$\nu_{\text{HD}^+}((0, 3) - (9, 3))$	415 264 925 501.8(1.3) kHz	3.1×10^{-12}	Patra <i>et al.</i> [34] + Germann <i>et al.</i> [14]
G1	$\nu_{\text{p}^4\text{He}}((32, 31) - (31, 30))$	1 132 609 226.7(4.0) MHz	3.5×10^{-9}	Hori <i>et al.</i> [37]
G2	$\nu_{\text{p}^4\text{He}}((33, 32) - (31, 30))$	2 145 054 858(7) MHz	3.4×10^{-9}	Hori <i>et al.</i> [36]
G3	$\nu_{\text{p}^3\text{He}}((32, 31) - (31, 30))$	1 043 128 581(6) MHz	6.2×10^{-9}	Hori <i>et al.</i> [37]
G4	$\nu_{\text{p}^3\text{He}}((35, 33) - (33, 31))$	1 553 643 100(10) MHz	6.7×10^{-9}	Hori <i>et al.</i> [36]
I1	$E_{\text{LS}}(\mu^4\text{He})$	1378.521(48) meV	3.5×10^{-5}	Krauth <i>et al.</i> [78]
I2	$E_{\text{LS}}(\mu^3\text{He})$	1258.586(49) meV	3.9×10^{-5}	Krauth [79]

New Physics Models

Dark photon

$$\mathcal{L}_{\text{int}} = -\frac{\epsilon}{2} F^{\mu\nu} F'_{\mu\nu}$$

$$\alpha_\phi = \alpha\epsilon^2$$

$$q_\ell = -q_p = -1$$

$$q_n = 0$$

$$\mathbf{U(1)}_{\mathbf{B-L}} \quad \alpha_\phi = g_{\mathbf{B-L}}^2/4\pi$$

$$q_\ell = -q_p = -1$$

$$q_n = 1 \leftarrow \text{highlights deuterium}$$

Higgs portal

$$\alpha_\phi = \sin^2\theta^2 m_e \kappa_p m_p / (4\pi v^2)$$

$$\kappa_p = 0.306(14), \kappa_n = 0.308(14) \leftarrow \text{from nucleon form-factors}$$

$$q_\ell = m_\ell / \sqrt{m_e \kappa_p m_p}$$

$$q_{p,n} = \kappa_{p,n} m_{p,n} / \sqrt{m_e \kappa_p m_p}$$

\leftarrow larger effects in muonic atoms and molecules

Hadrophilic scalar

$$\alpha_\phi = \sin^2\theta^2 m_e \kappa_p m_p / (4\pi v^2)$$

$$q_\ell = 0 \leftarrow \text{highlights molecules}$$

$$q_{p,n} = \kappa_{p,n} m_{p,n} / \sqrt{m_e \kappa_p m_p}$$

Up-Lepto-Darko-philic

(ULD) scalar

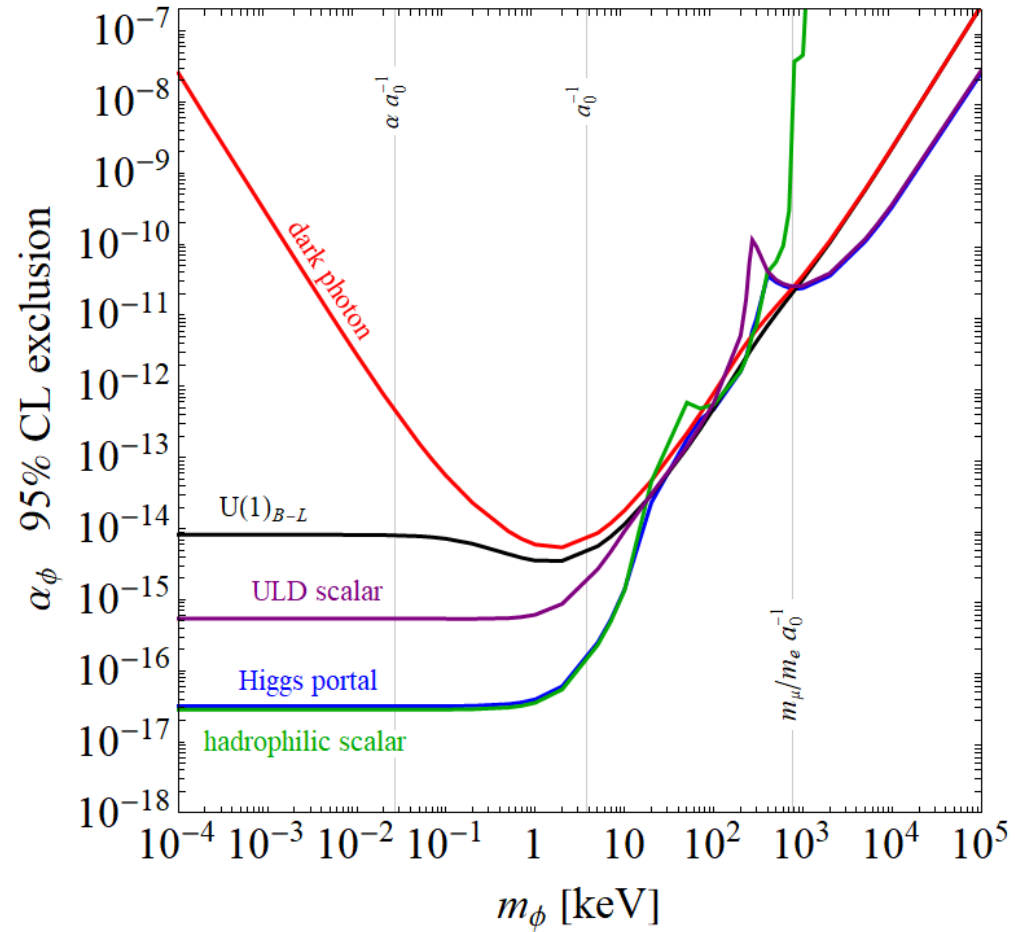
$$\alpha_\phi = k^2 m_e \kappa'_p m_p / (4\pi v^2)$$

$$q_\ell = m_\ell / \sqrt{m_e \kappa'_p m_p}, \quad q_{p,n} = \kappa'_{p,n} m_{p,n} / \sqrt{m_e \kappa'_p m_p}$$

$$\kappa'_p = 0.018(5), \kappa'_n = 0.016(5) \leftarrow \text{couples only to up-quark}$$

+ dominant ϕ decay to invisible states (see later) ²⁶

New Physics Bounds

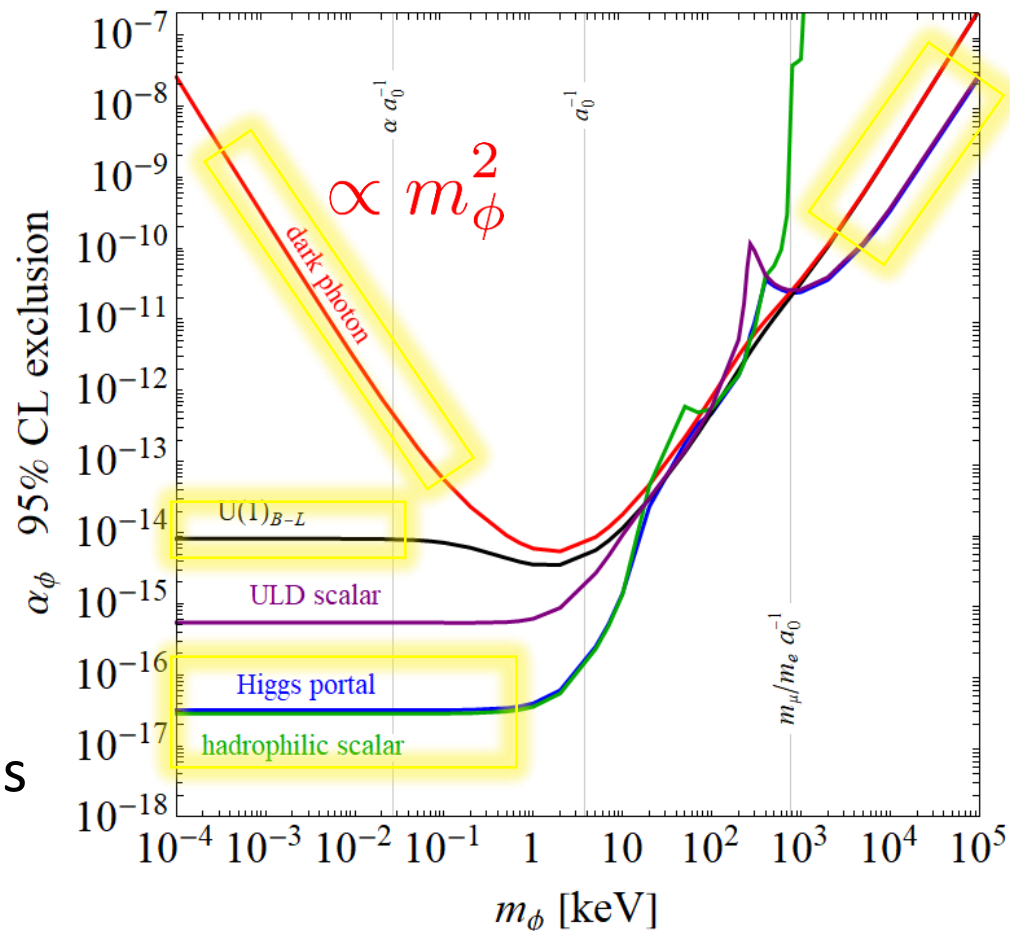


New Physics Bounds

$\propto m_\phi^0$ thanks to
Deuterium data

stronger sensitivity
from internuclear forces
in **molecules** in models
where

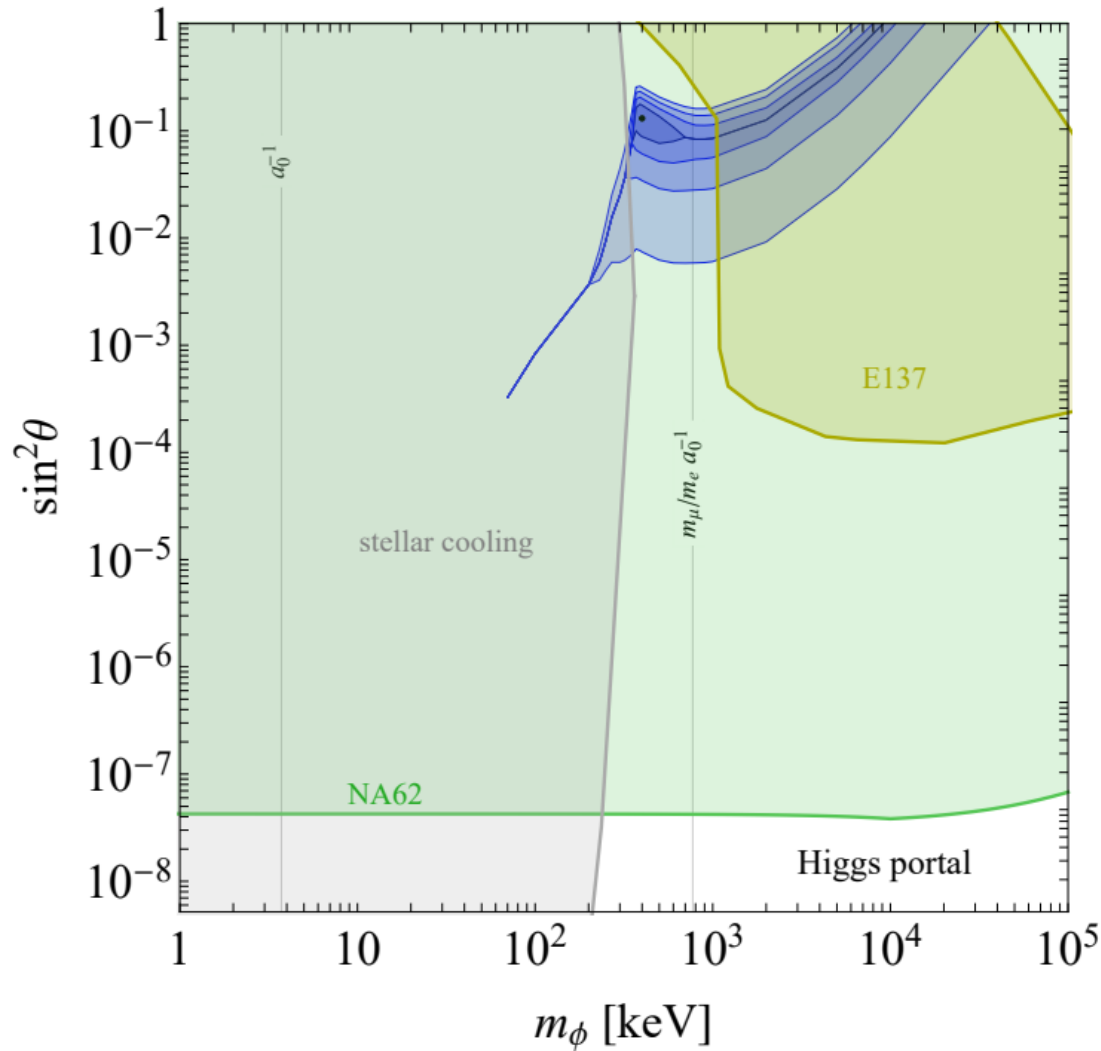
$$q_N/q_e \sim m_N/m_e \sim 10^3$$



stronger sensitivity
from **muonic** atoms
in models where

$$q_\mu/q_e \sim m_\mu/m_e \sim 200$$

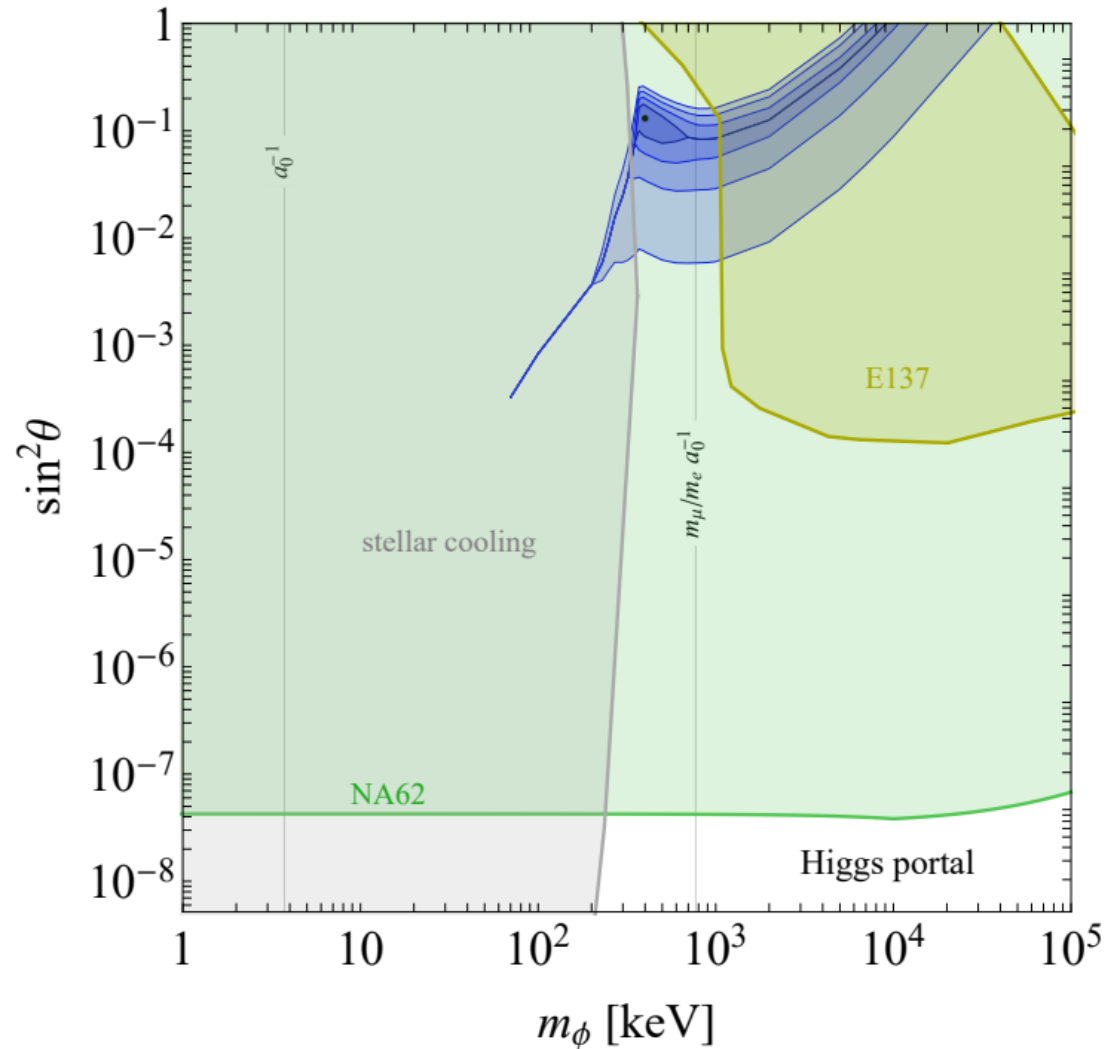
Non-zero Higgs Portal?



Best-fit point $\left| \begin{array}{l} \sin \theta \simeq 0.35 \\ m_\phi \simeq 400 \text{ keV} \end{array} \right.$

is largely **excluded** by
 $K^+ \rightarrow \pi^+ X_{\text{inv}}$ searches

Non-zero Higgs Portal?



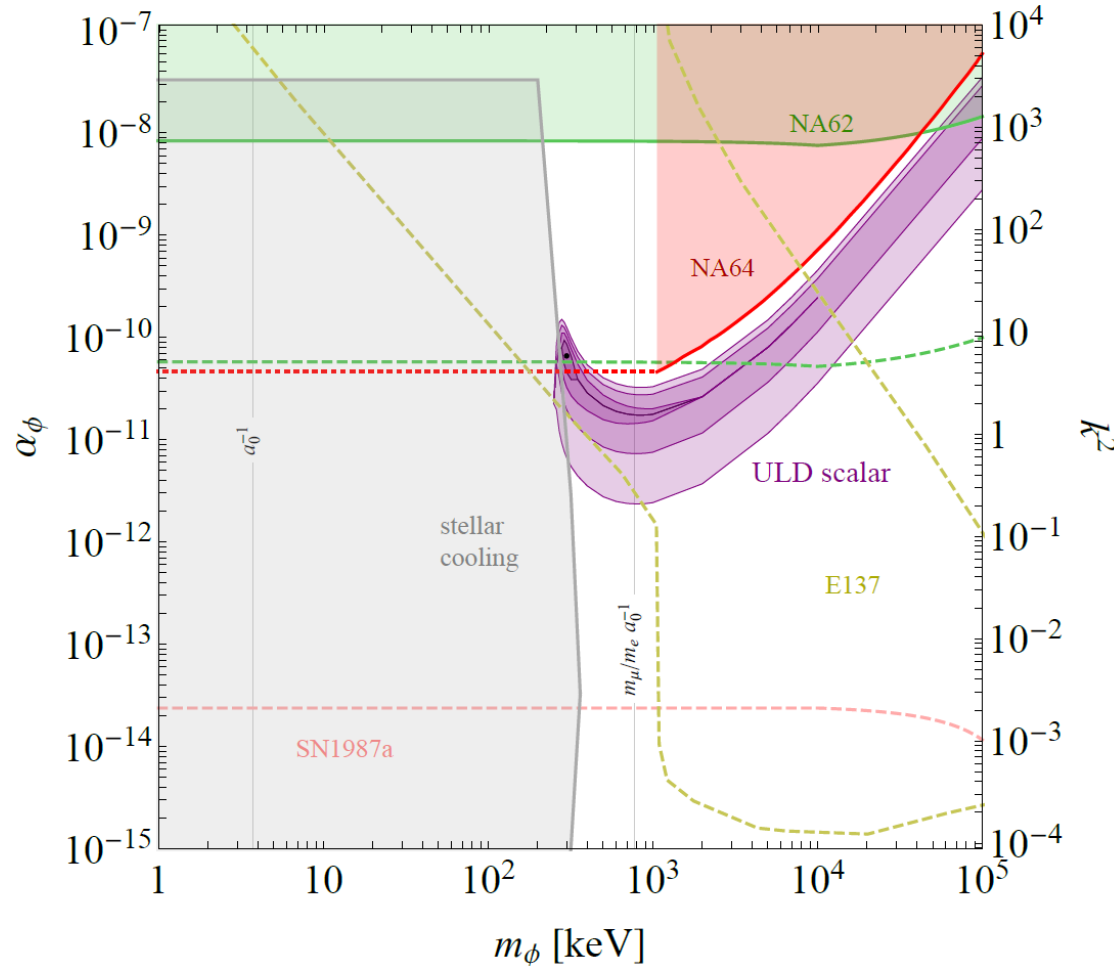
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The **NA62** bound is driven by coupling to heavy quarks through one-loop penguins

The **E137** beam-dump bound relies on scalars dominantly decaying to $\phi \rightarrow e^+ e^-$

Evidence for ULD scalar



Best-fit point $\left\{ \begin{array}{l} \alpha_\phi \simeq 6.7 \times 10^{-11} \\ m_\phi \simeq 300 \text{ keV} \end{array} \right.$

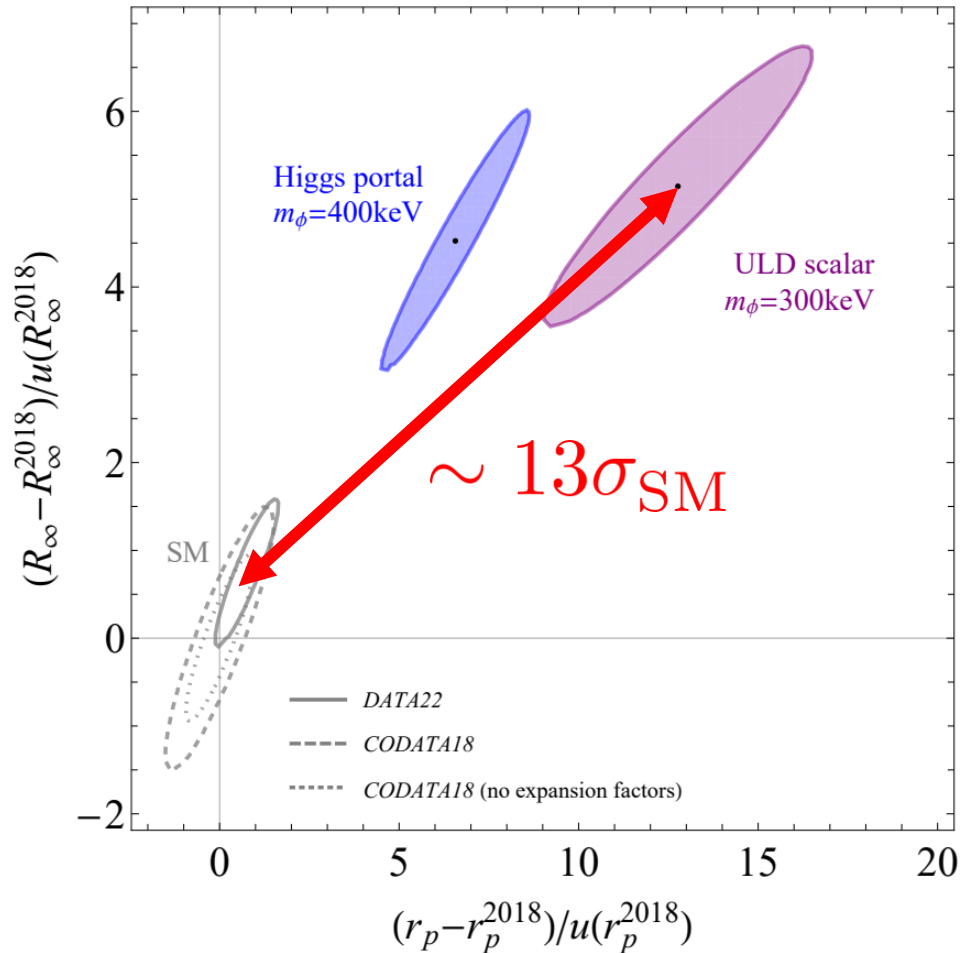
evades the **NA62** bound by coupling only to up quarks

The **E137** bound does not apply assuming invisible decay dominantes ($\phi \rightarrow \text{DMDM}?$)

In that case **NA64** is relevant
 $e^- Z \rightarrow e^- Z \phi$ Andreev et al. [2021]

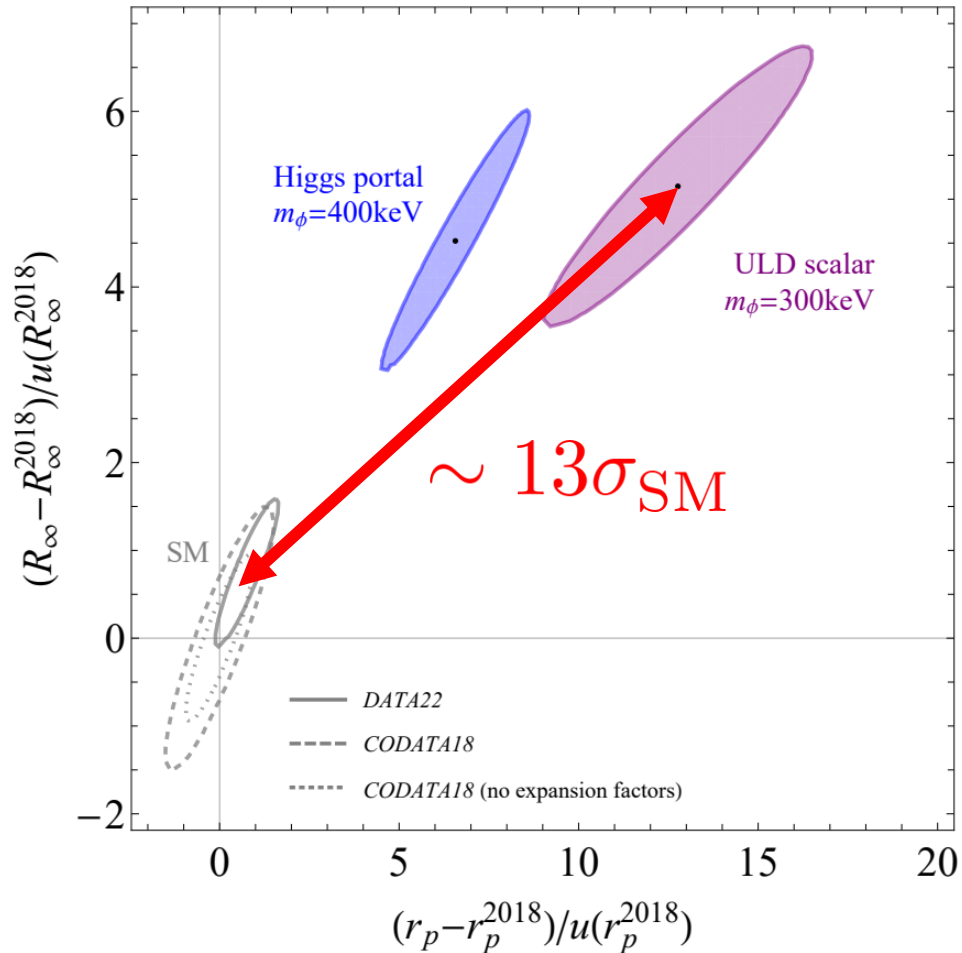
yielding a weaker bound but NP sensitivity not clear below MeV

Impact on fundamental constants

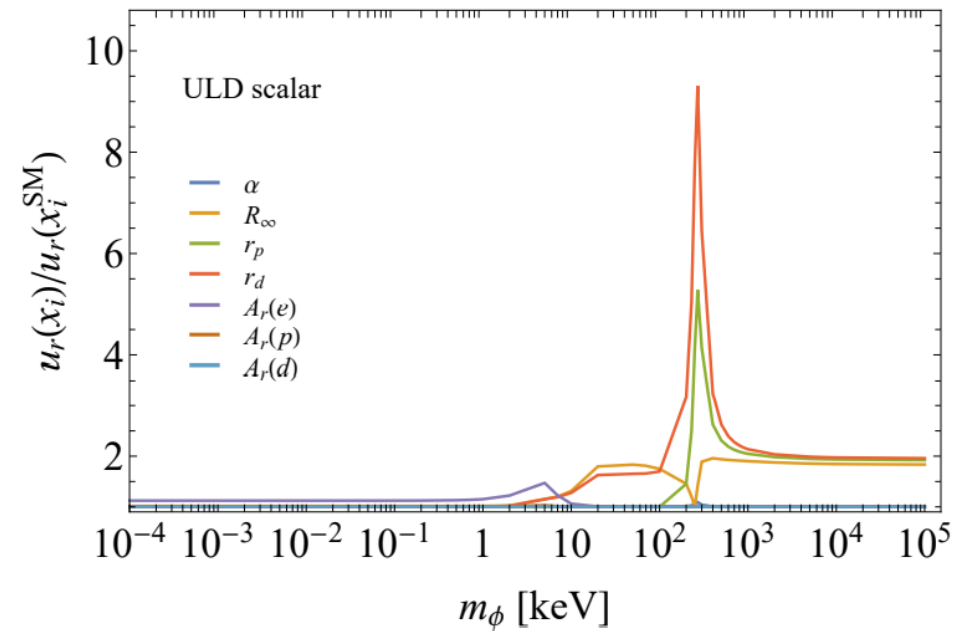


Fundamental Constants undergo **huge shifts** in the presence of NP

Impact on fundamental constants



Fundamental Constants undergo **huge shifts** in the presence of NP



and their uncertainty *significantly* inflates relative to the SM-only hypothesis

Conclusions

Conclusions

Physicists (almost) exclusively searched for BSM at TeV energies and above, thinking physics at lower energies is very well understood.

Light & weakly-coupled New Physics is well motivated theoretically.
Atomic spectroscopy can be repurposed to search for it.

Such low-energy probes complement the efforts (to be) invested at colliders

→ « multi-messenger » program for HEP searches

PHYS_TEV 2025 workshop in **Les Houches**

June 25 → July 4



backups

The Light Vector Case

Vectors with $m_\phi \ll \alpha m_e \simeq 4 \text{ keV}$ induce a long-range force

Then, effects are suppressed for couplings aligned with QED ($q_i \simeq Q_i$) because:

$$\mathcal{L}_{\text{QED}}(\alpha) + \mathcal{L}_{A'_\mu}(\alpha', m_{A'} \rightarrow 0) \rightarrow \mathcal{L}_{\text{QED}}(\alpha + \alpha')$$

massless dark photon is **unobservable!**

This behavior is only manifest for $\mathcal{O}_{\text{NP}}(\alpha')$ and $\mathcal{O}_{\text{SM}}(\alpha)$ calculated at the same order in couplings. Otherwise:

$$\mathcal{O} \rightarrow \mathcal{O}_{\text{SM}}^{\text{LO}}(\alpha + \alpha') + \mathcal{O}_{\text{SM}}^{\text{NLO}}(\alpha)$$

would fictitiously distinguish photon from darkphoton

The Light Vector Case

Vectors with $m_\phi \ll \alpha m_e \simeq 4 \text{ keV}$ induce a long-range force
Then, effects are suppressed for couplings aligned with QED ($q_i \simeq Q_i$)

Instead, we use a simple *prescription*:

$$V_{\text{NP}}^{ij} = \alpha_\phi \frac{Q_i Q_j}{r} + \tilde{V}_{\text{NP}}^{ij} \quad \text{with} \quad \tilde{V}_{\text{NP}}^{ij} \equiv \alpha_\phi (q_i q_j e^{-m_\phi r} - Q_i Q_j) / r$$

included to all orders
by shifting $\alpha \rightarrow \alpha + \alpha_\phi$
in \mathcal{O}_{SM}

deviations from either $m_\phi \neq 0$ or $q_i \neq Q_i$
can be treated as perturbations at **LO**

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 by shifting $\alpha \rightarrow \alpha + \alpha_\phi$
 in \mathcal{O}_{SM}

inverse Bohr radius

deviations from either $m_\phi \neq 0$ or $q_i \neq Q_i$
 can be treated as perturbations at **LO**

Hence: $\mathcal{O} = \mathcal{O}_{\text{SM}}(\alpha + \alpha_\phi) + \tilde{\mathcal{O}}_{\text{NP}}(\alpha + \alpha_\phi, \alpha_\phi, m_\phi) + \delta\mathcal{O}_{\text{th}}$

$\propto m_\phi^2$ or $\delta q_i Q_j + Q_i \delta q_j$