

STRING PERCOLATION AND THE NEW LHC DATA

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Dept Particle Physics

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Low x 2011 Santiago de Compostela

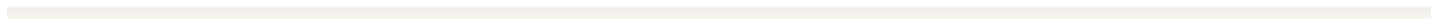
Collaboration with I.Bautista, J.Dias de Deus,
J.G. Milhano

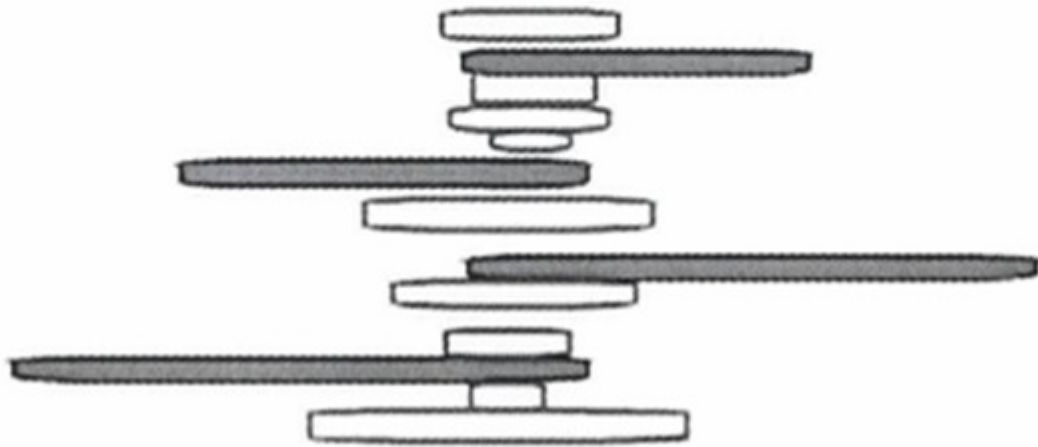
Outline

- Brief description of string percolation
 - dn/dy in pp and Pb-Pb
 - Rapidity long range correlations, Ridge in pp and AA
 - Elliptic flow in percolation, comparison at RHIC and LHC
- Shear viscosity/entropy in percolation



b

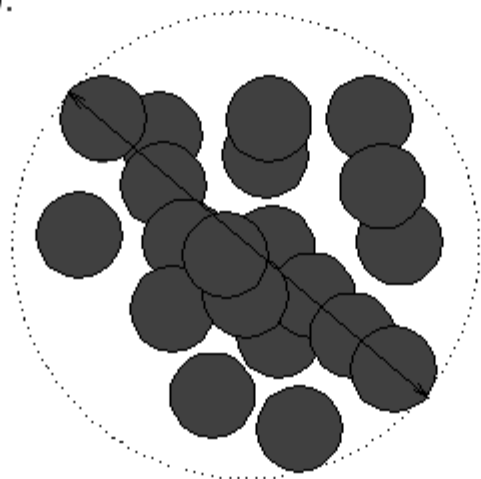




- **Color strings** are stretched between the projectile and target
 - **Strings = Particle sources**: particles are created via sea $q\bar{q}$ production in the field of the string
 - **Color strings = Small areas** in the transverse space filled with color field created by the colliding partons
 - With growing energy and/or atomic number of colliding particles, the number of sources grows
 - So the elementary color sources start to overlap, forming clusters, very much like disk in the 2-dimensional percolation theory
 - In particular, at a certain critical density, a macroscopic cluster appears, which marks the percolation phase transition
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(N. Armesto et al., PRL77 (96); J.Dias de Deus et al., PLB491 (00); M. Nardi and H. Satz(98).

- **How?:** Strings fuse forming clusters. At a certain critical density η_c (central PbPb at SPS, central AgAg at RHIC, central SS at LHC) a macroscopic cluster appears which marks the percolation phase transition (second order, non thermal).



$$\eta = N_{st} \frac{S_1}{S_A}, \quad S_1 = \pi r_0^2, \quad r_0 = 0.2 \text{ fm}, \quad \eta_c = 1.1 \div 1.2.$$

$$\mu_n = \sqrt{\frac{nS_n}{S_1}} \mu_1 ; \langle p_T^2 \rangle_n = \sqrt{\frac{nS_1}{S_n}} \langle p_T^2 \rangle_1$$

Energy-momentum of the cluster is the sum of the energy-momentum of each string.

As the individual color field of the individual string may be oriented in an arbitrary manner respective to one another, $Q_n^2 = nQ_1^2$

■ At high densities

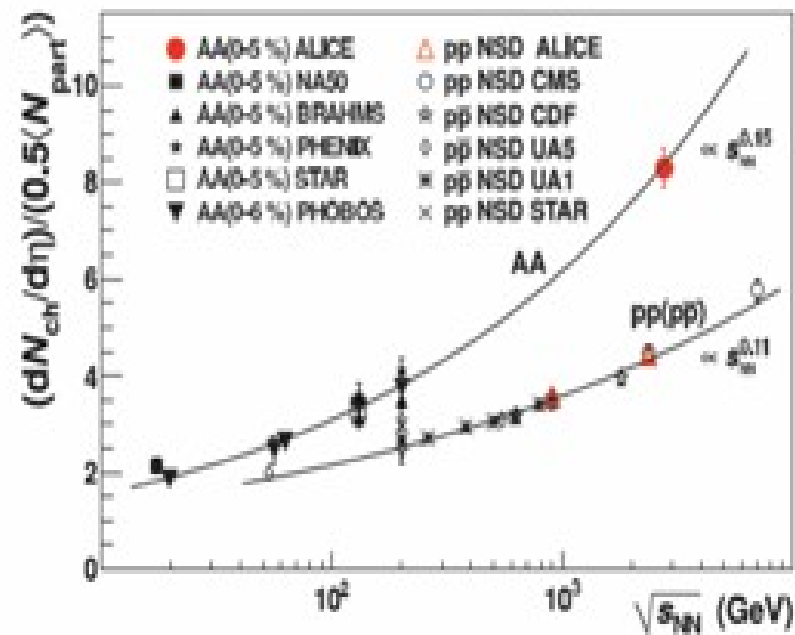
• $\langle \mu \rangle_n = nF(\eta) \langle \mu \rangle_1 \quad \langle p_T^2 \rangle_n = \frac{\langle p_T^2 \rangle_1}{F(\eta)}$

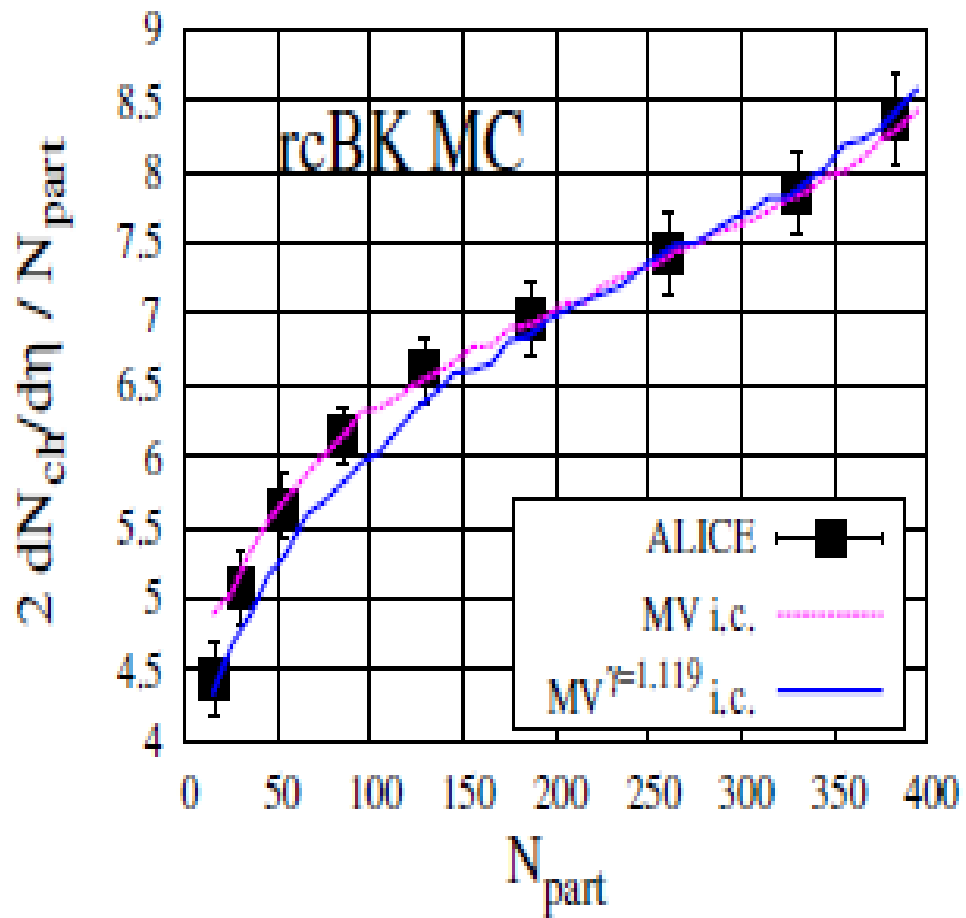
• $F(\eta) = \sqrt{\frac{1-e^{-\eta}}{\eta}}, \quad \eta = N_S \frac{\pi r_0^2}{S_A}$

• r_0 is the transverse size of a single string $\simeq 0.2$ fm.

New data: multiplicities

[ALICE arxiv:1011.3916]





$$\frac{1}{N_A} \frac{dn}{dy} |_{N_A N_A} = \frac{dn}{dy} |_{pp} \left[1 + \gamma(\sqrt{s}) \frac{F(\eta_{N_A})}{F(\eta_p)} (N_A^{1/3} - 1) \right]$$

$$\text{with } \gamma(\sqrt{s}) = 1 - \frac{\sqrt{s_{th}}}{\sqrt{s}}$$

at high energy $\gamma \rightarrow 1$

$$\frac{1}{N_A} \frac{dn}{dy} |_{N_A N_A} = \frac{dn}{dy} |_{pp} \left[1 + \left(\frac{N_A}{A} \right)^{1/6} \left(1 - \frac{1}{N_A^{1/3}} \right) \right]$$

(The shape as a function of N_A is independent of energy)

Two-Particle Angular Correlations

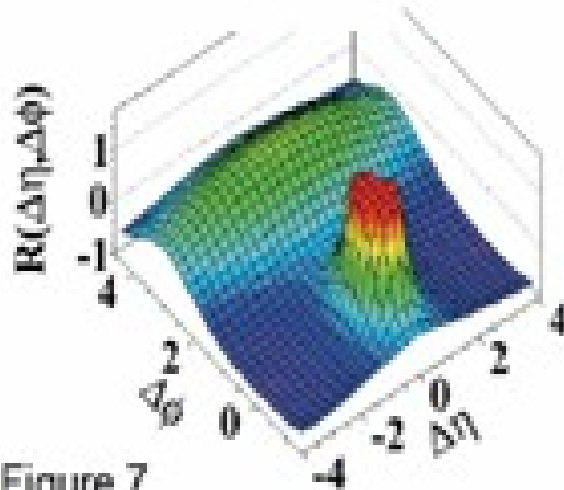
Published in

J. High Energy Phys. 09 (2010) 091

First **surprising** result from the LHC:
Observation of Long-Range Near-Side
Angular Correlations in pp Collisions

MinBias

(b) MinBias, $1.00\text{GeV} < p_T < 3.00\text{GeV}$



high multiplicity ($N > 110$)

(d) $N > 110$, $1.00\text{GeV} < p_T < 3.00\text{GeV}$

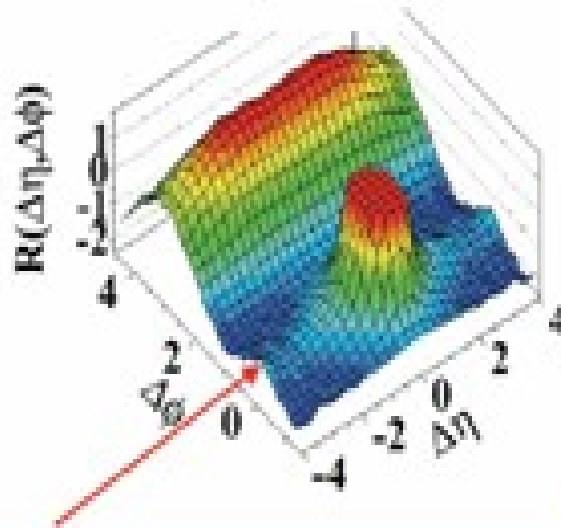


Figure 7

Why Protons?

In String Percolation...

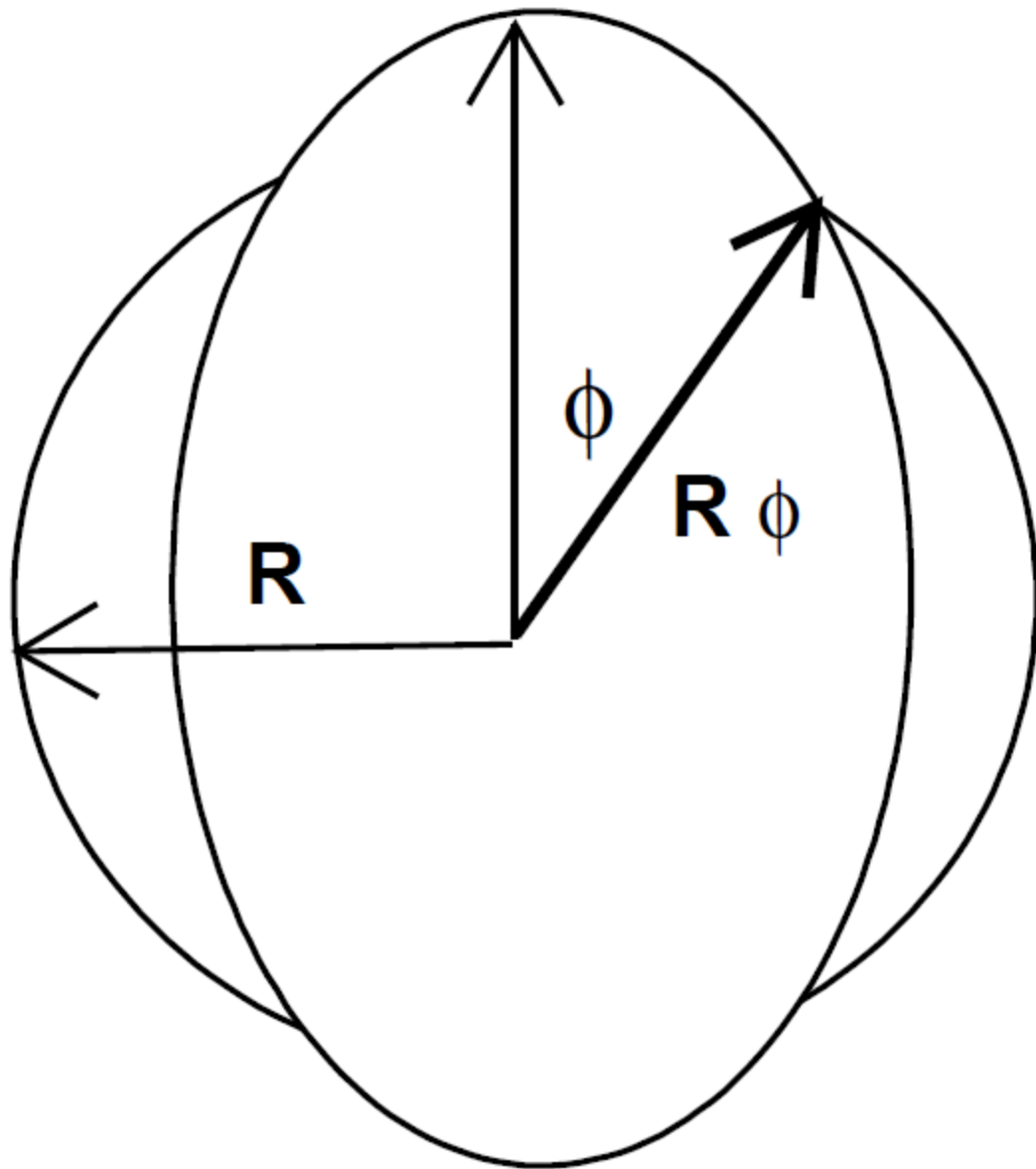
$$\eta_{AA} = \left(\frac{r}{R}\right)^2 \bar{N}^s \cong \frac{N_A^{4/3}}{N_A^{2/3}} \left(\frac{r}{R_p}\right)^2 \bar{N}_p^s$$

$$\eta_{AA}(s) = N_A^{2/3} \eta_{pp}(s) \quad \text{and} \quad \bar{N} \sim s^{2/7}$$

$$\eta_c \approx 1.15 \begin{cases} \swarrow \eta_{PbPb}(\sqrt{s}) \cong 20 GeV \\ \searrow \eta_{PP}(\sqrt{s}) \cong 6 TeV \leftarrow \text{LHC} \end{cases}$$

- As the string density in Au-Au peripheral collisions at 200Gev is the same as in pp high multiplicity events(three times m.b)
If there is a ridge structure in Au-Au at RHIC
It should be seen the same structure in pp as it was seen CMS collaboration

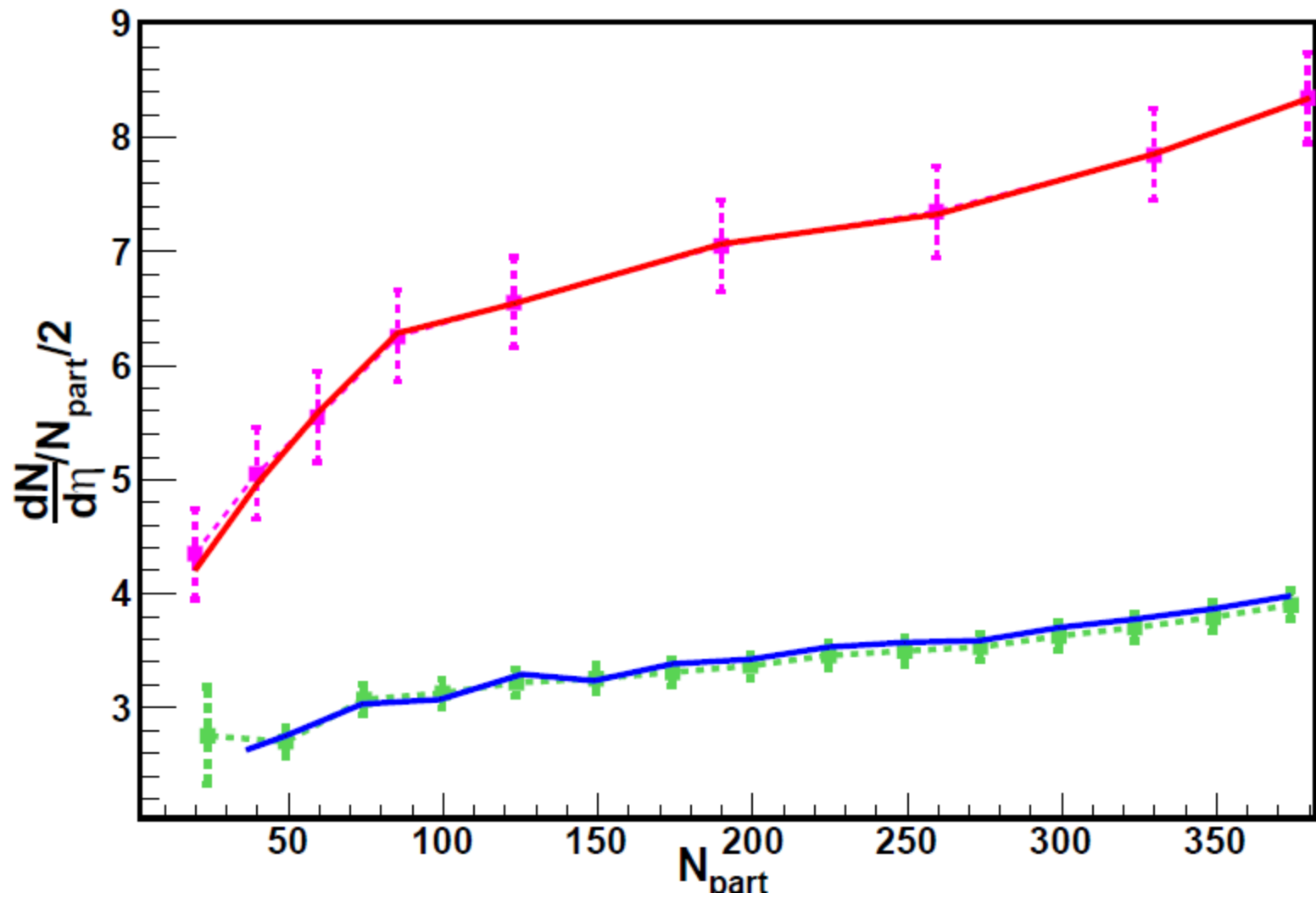
L.Cunqueiro, J Dias de Deus and CP
Eur Phys J C 65 423 (2010)

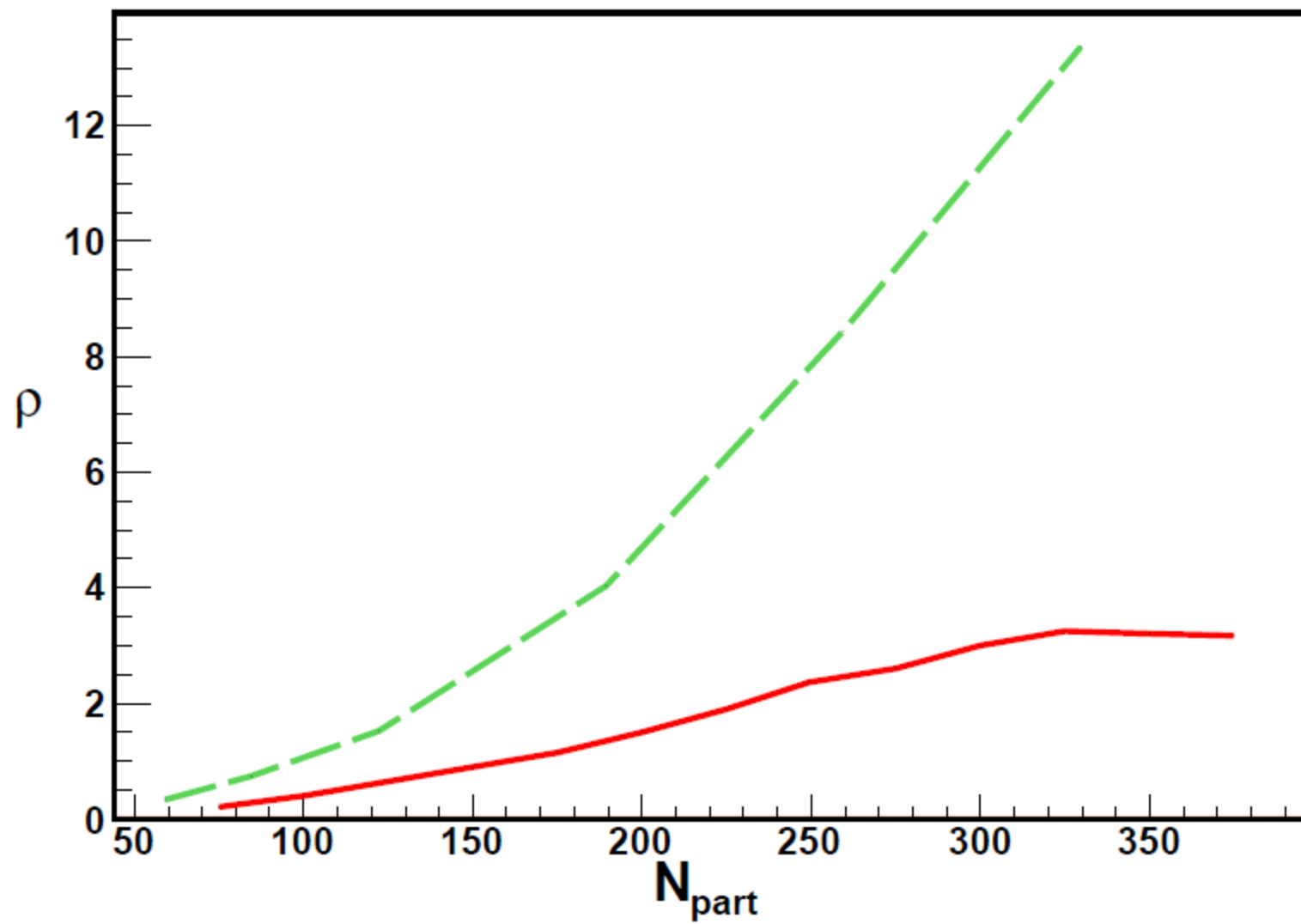


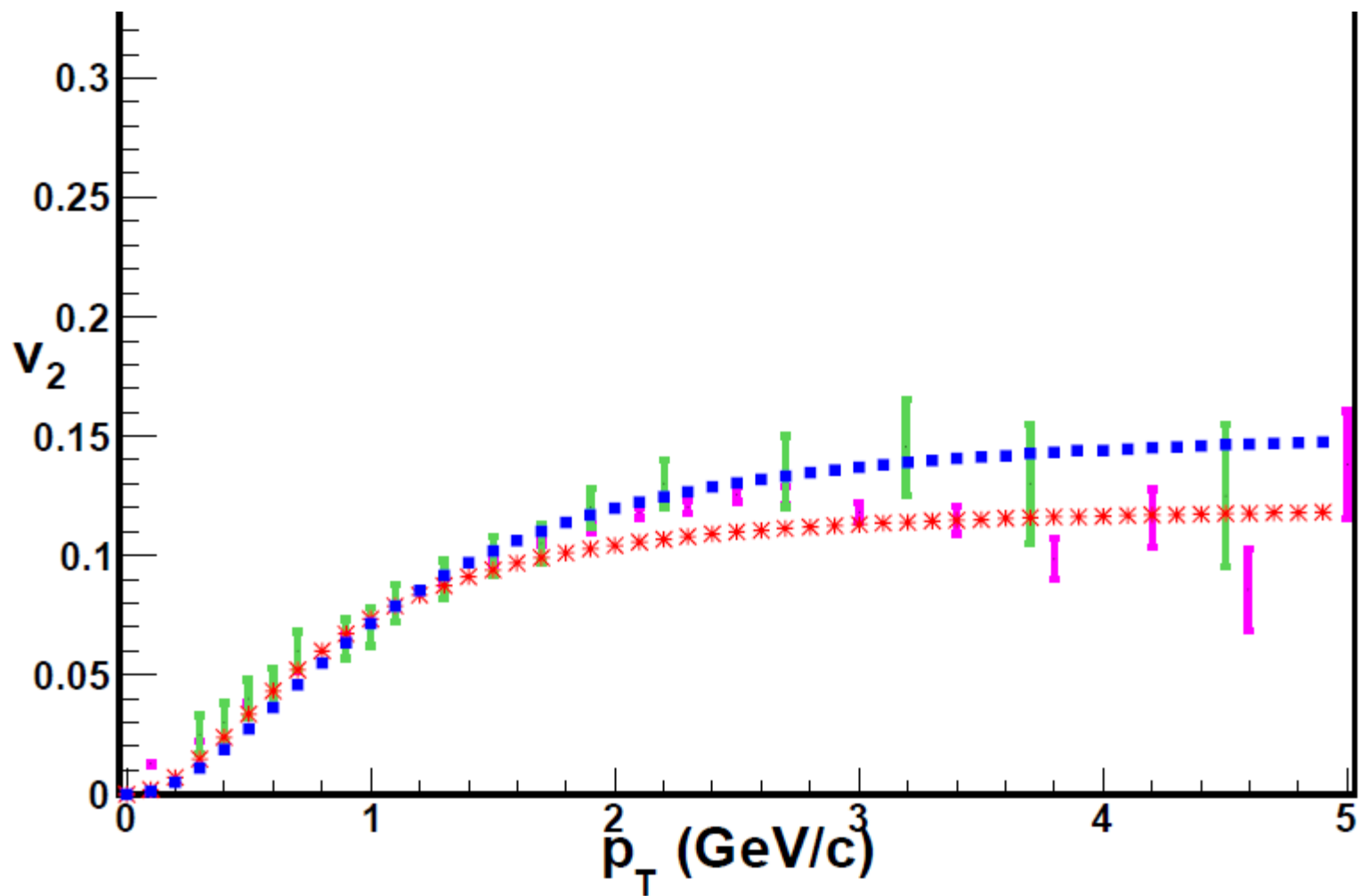
$$\eta_\varphi = \eta \left(\frac{R}{R_\varphi} \right)^2$$

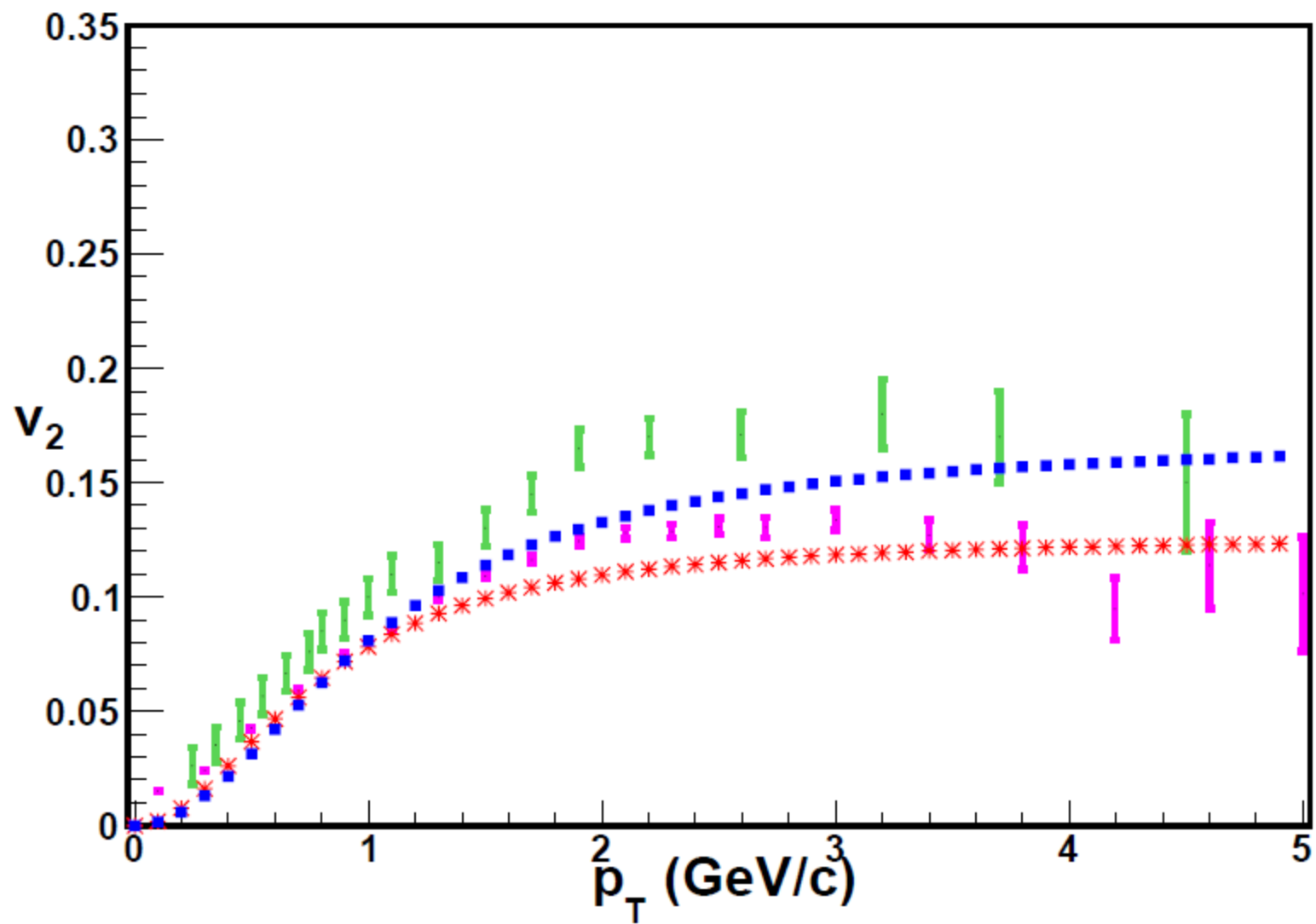
$$\begin{aligned} v_2(p_T^2, y) &= \frac{2}{\pi} \int_0^{\pi/2} d\varphi \cos(2\varphi) \left[1 + \frac{\partial \ln f(p_T^2, \eta, y)}{\partial R^2} (R_\varphi^2 - R^2) \right] \\ &= \frac{2}{\pi} \int_0^{\pi/2} d\varphi \cos 2\varphi \left(\frac{R_\varphi}{R} \right)^2 \left(\frac{e^{-\eta} - F(\eta)^2}{2F(\eta)^2} \right) \frac{F(\eta) p_T^2 / \langle p_T^2 \rangle_1}{(1 + F(\eta) p_T^2 / \langle p_T^2 \rangle_1)} \end{aligned}$$

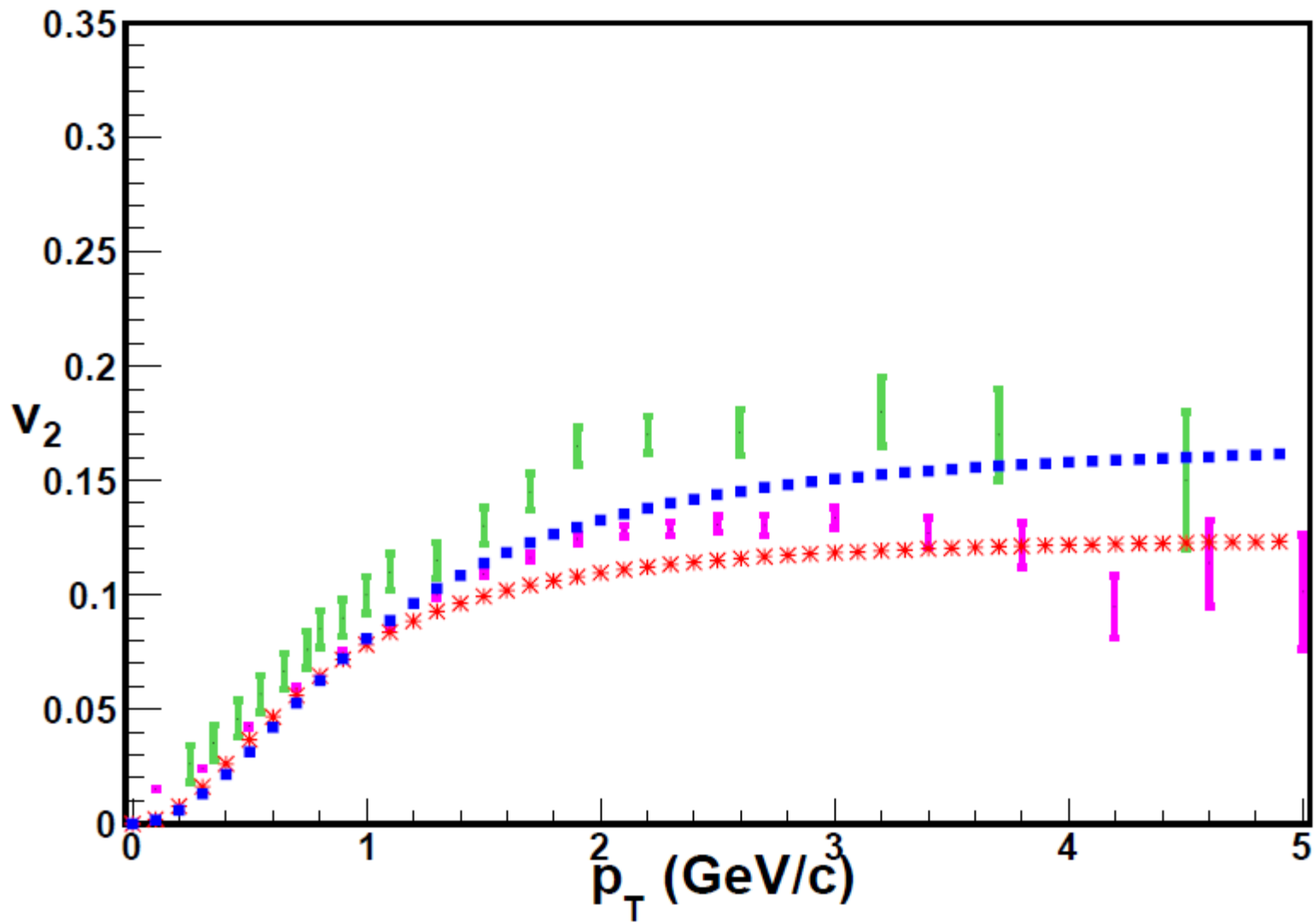
$$v_2 = \frac{2}{\pi} \int_0^{\pi/2} d\varphi \cos(2\varphi) \left(\frac{R_\varphi}{R} \right) \left(\frac{e^{-\eta} - F(\eta)^2}{2F(\eta)^3} \right) \frac{R}{R-1}$$

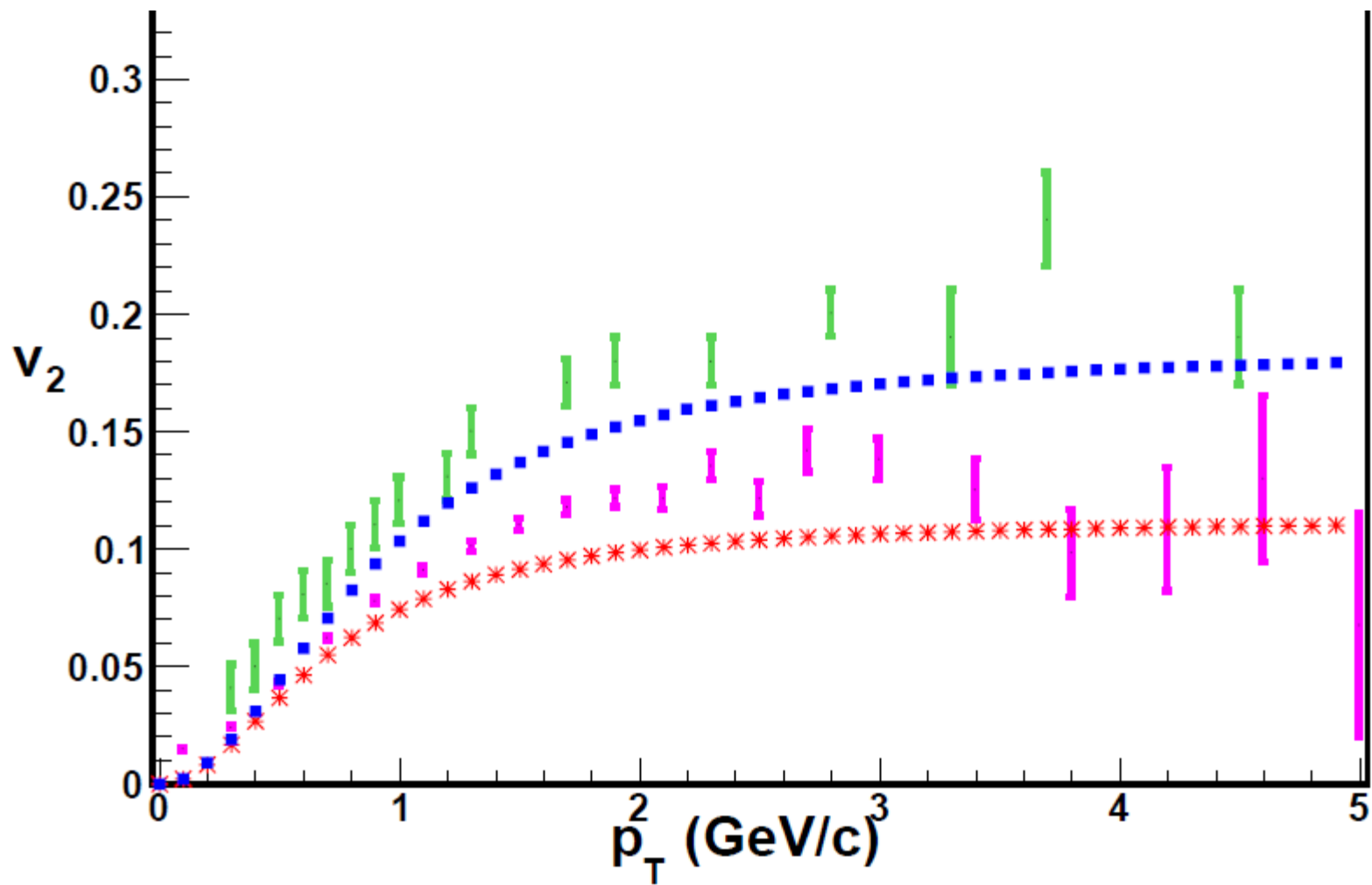


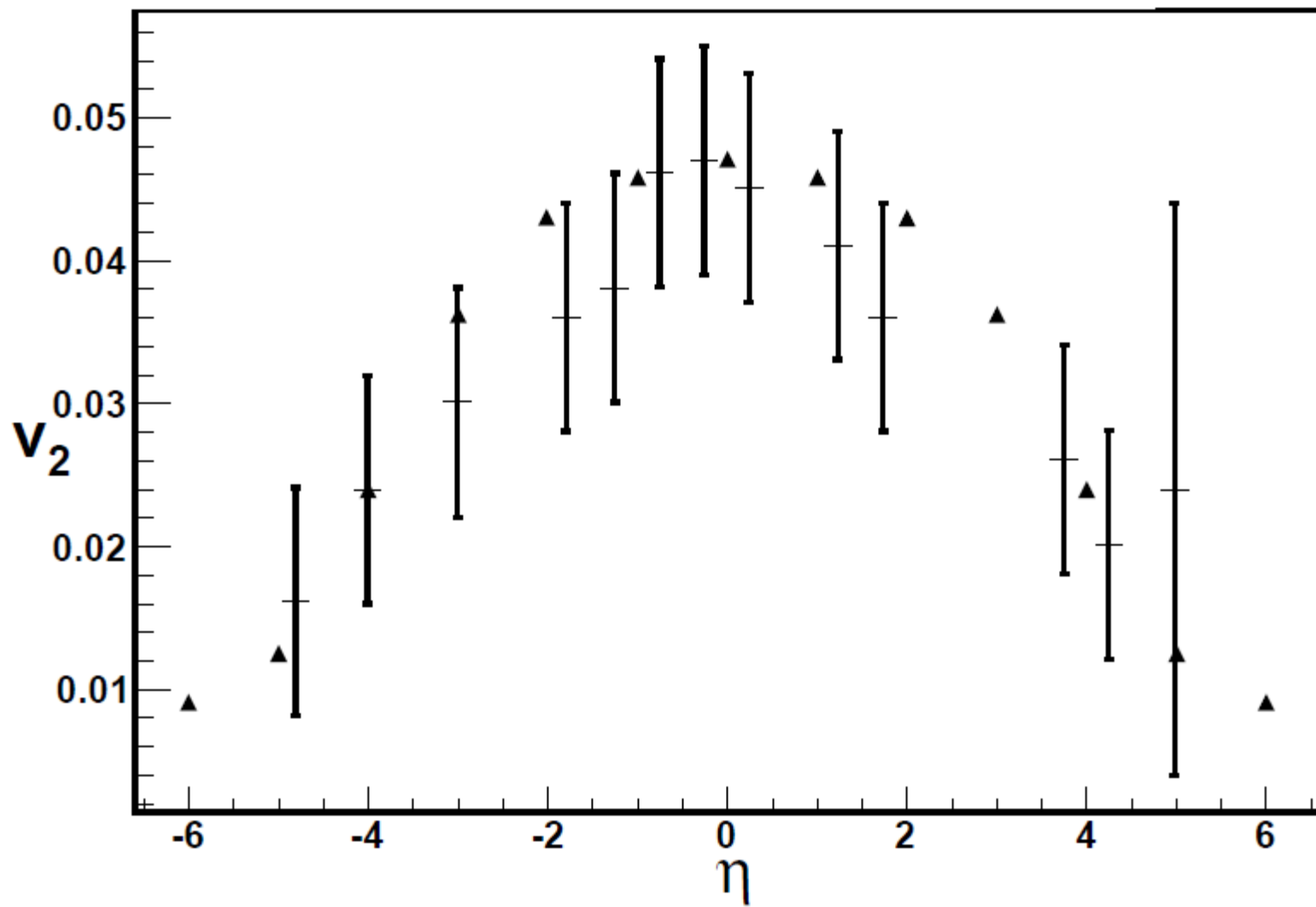


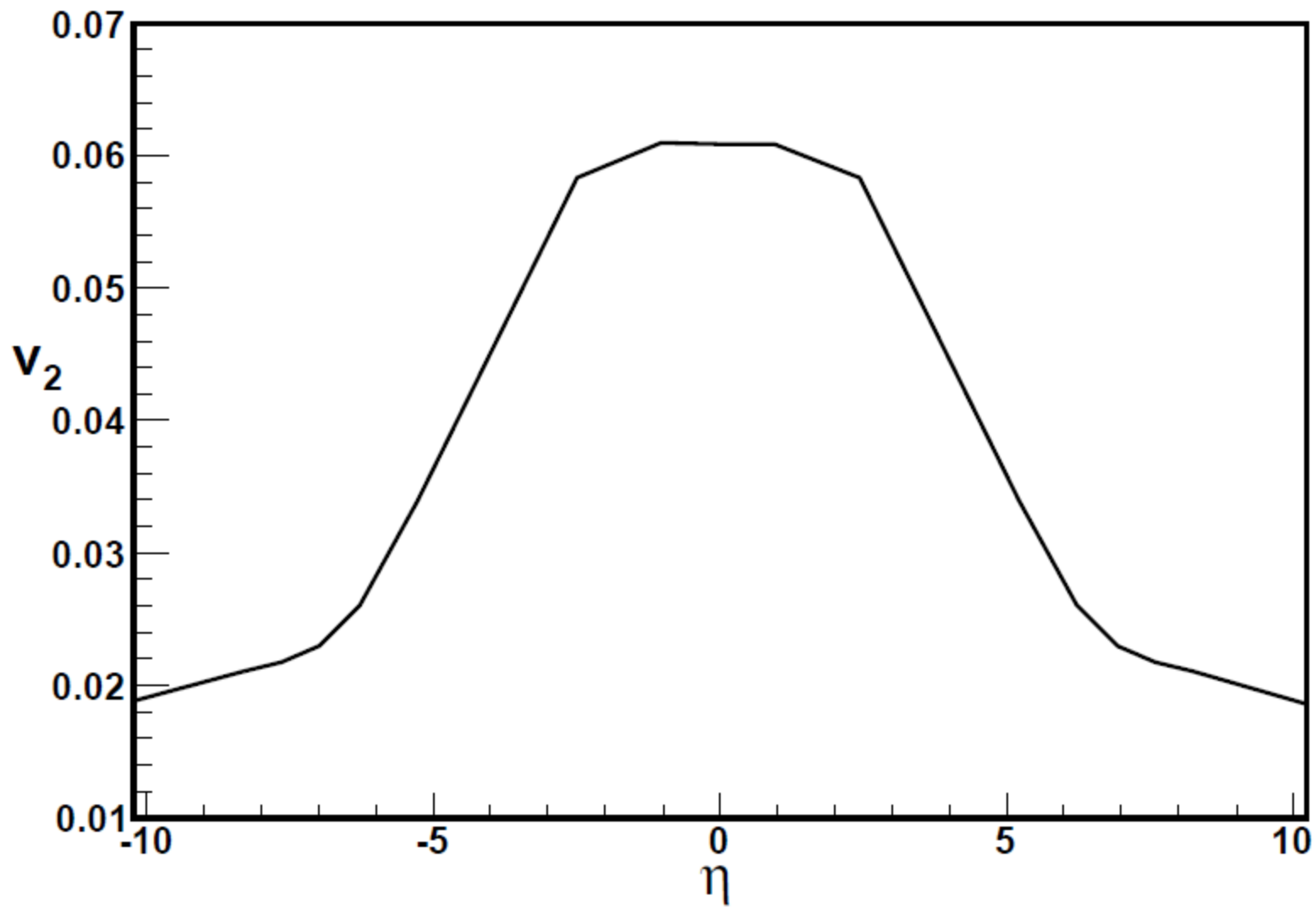


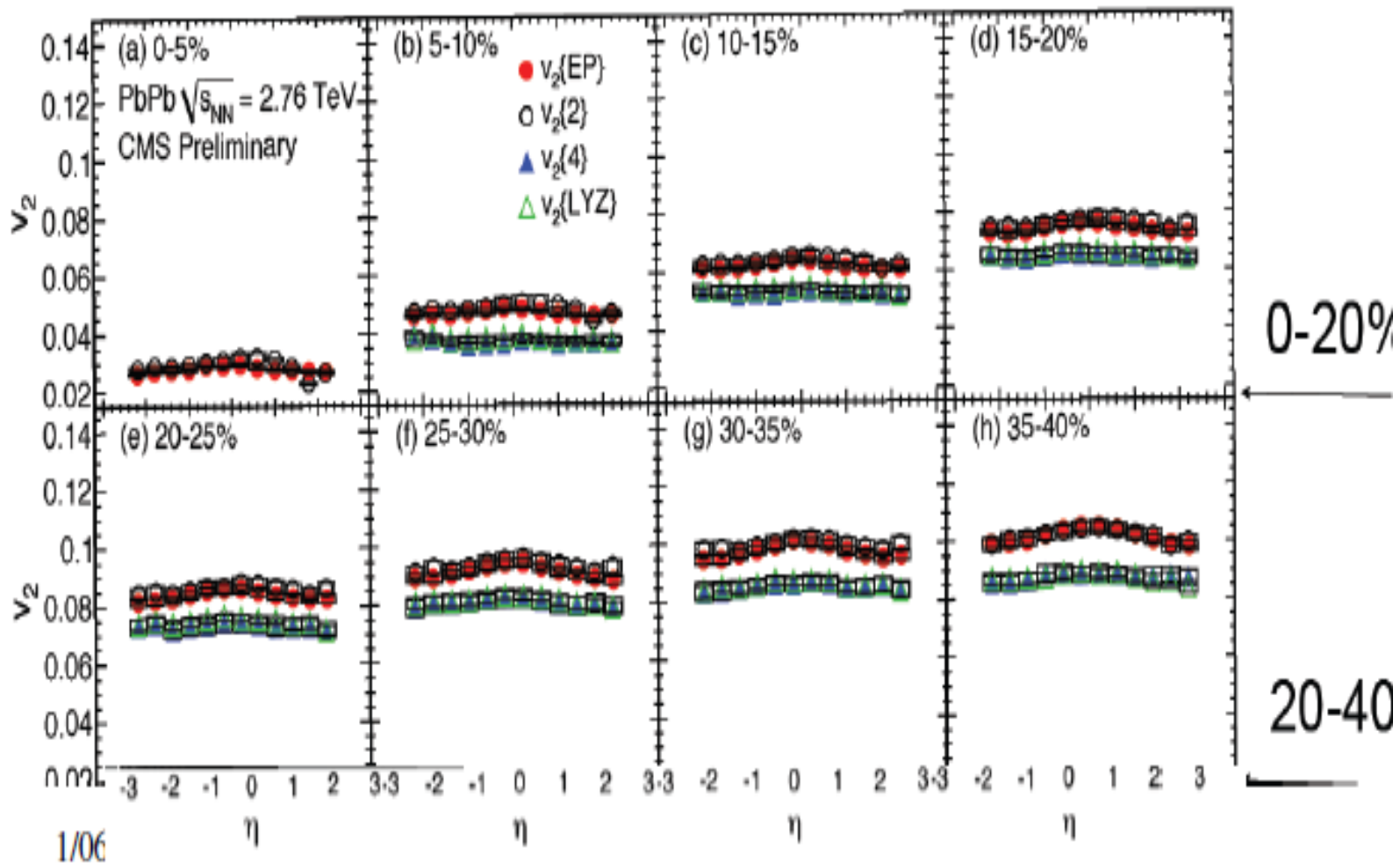


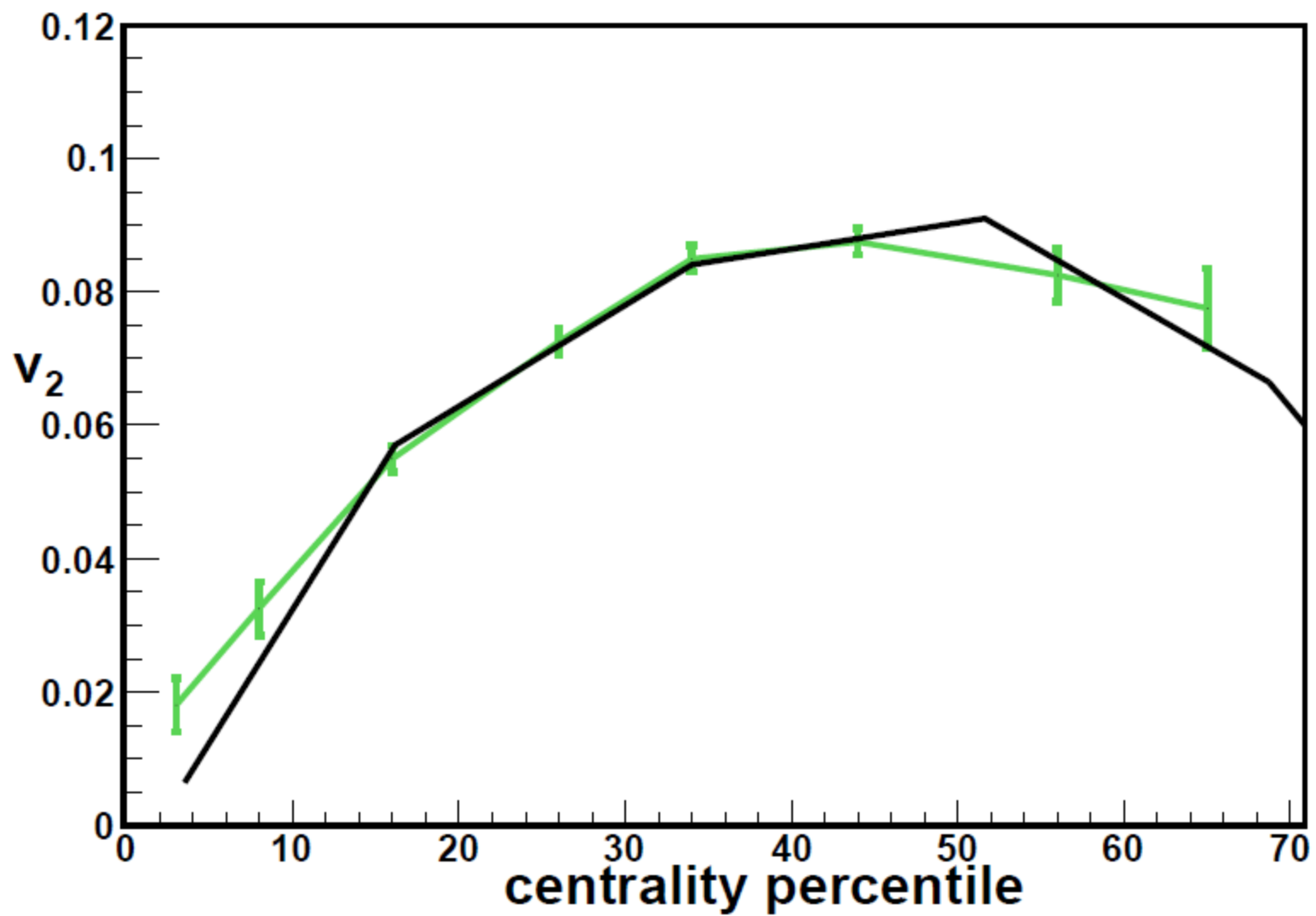


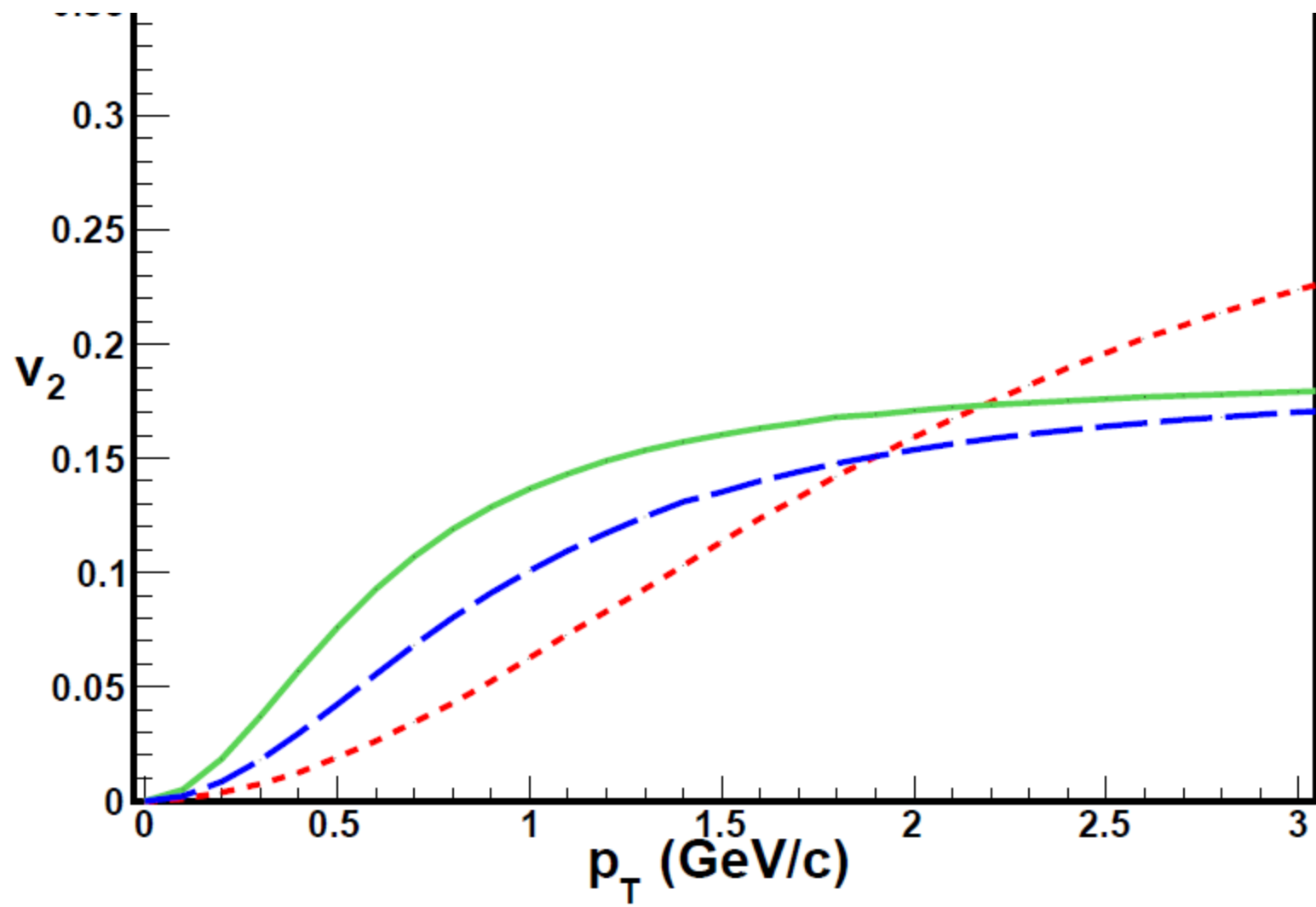


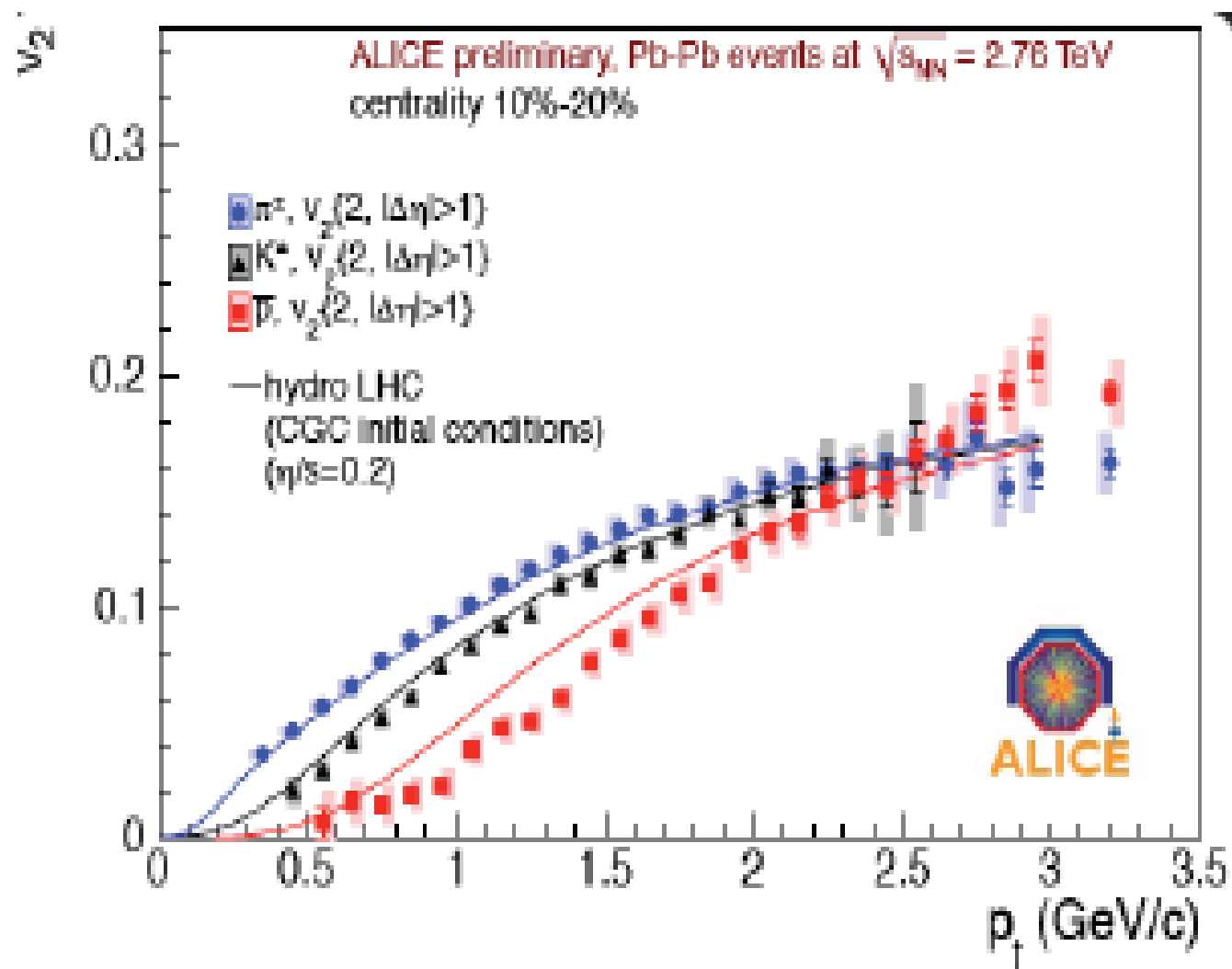


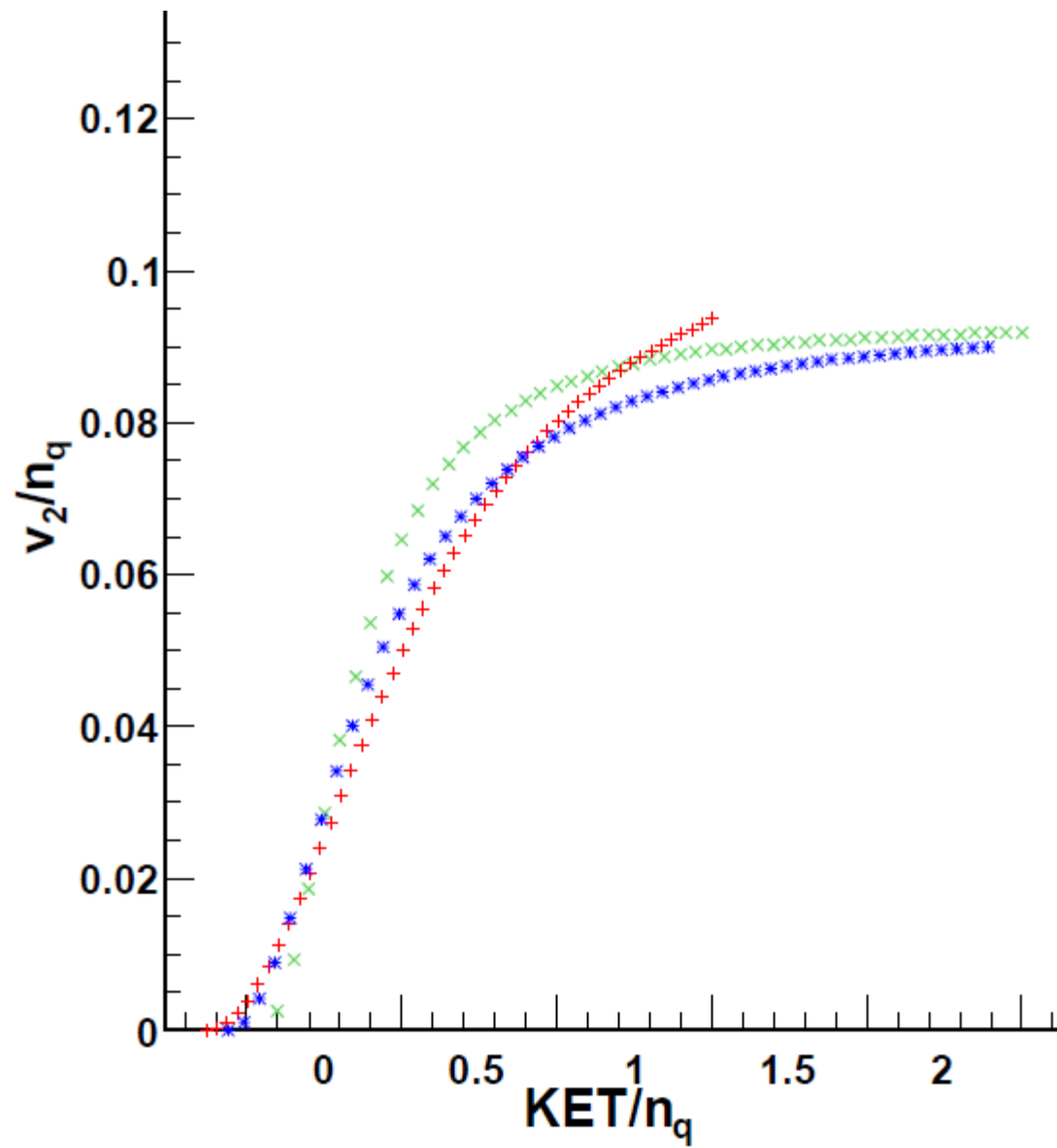




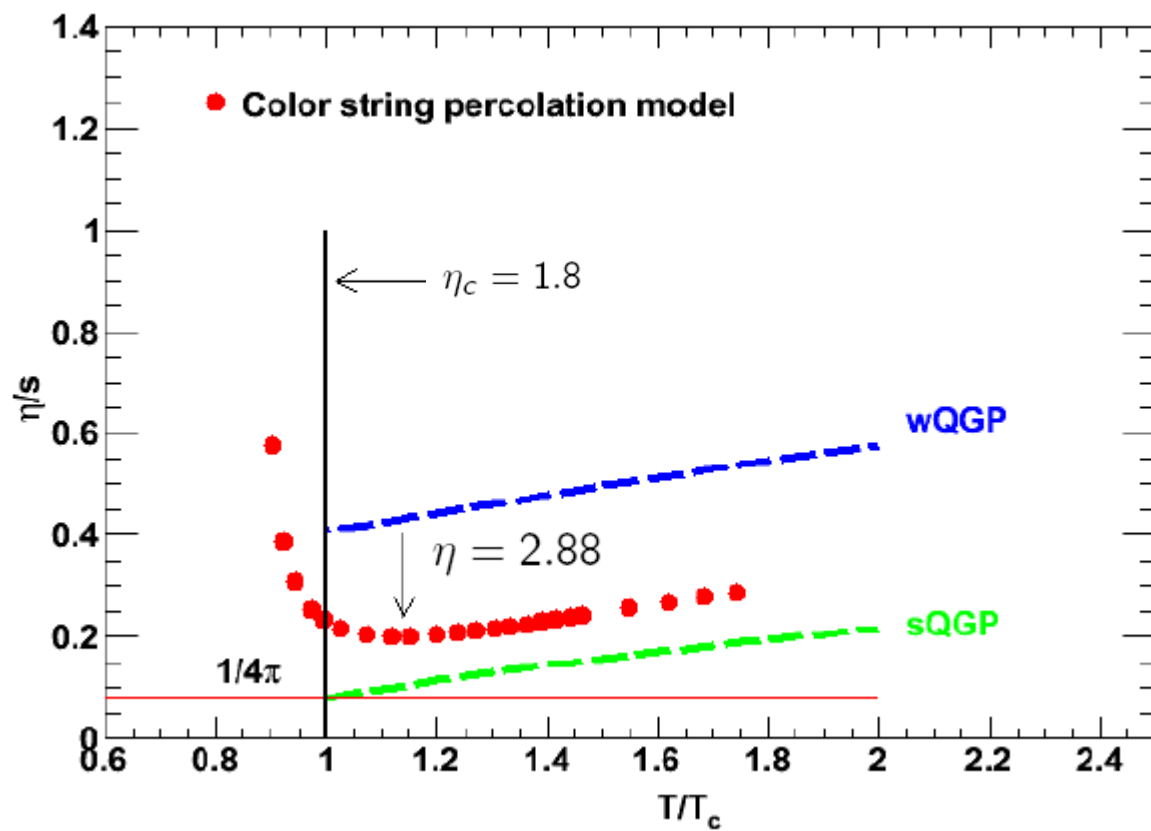








$$\frac{\eta}{s} = \frac{1}{5\sqrt{2}} \frac{\langle p_T \rangle_1 \eta^{1/4}}{(1-e^{-\eta})^{5/4}} L$$



Conclusions

- A good agreement with RHIC and LHC data, for dN/dy dependences on energy and centrality.
- It was predicted rapidity long range correlations and ridge structure in pp at high multiplicity
- Good description of v_2 at RHIC and LHC including the rapidity dependences
- Low ratio shear viscosity/entropy density in the whole energy range RHIC-LHC