

No unique definition of diffraction

1. Diffraction is elastic (or quasi-elastic) scattering caused, via **s-channel** unitarity, by the absorption of components of the wave functions of the incoming particles

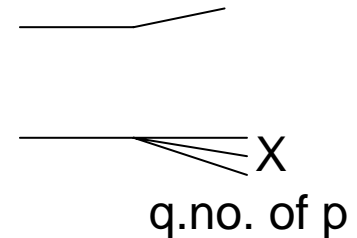
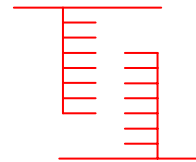
e.g. $pp \rightarrow pp$,

$pp \rightarrow pX$ (single proton dissociation, SD),

$pp \rightarrow XX$ (both protons dissociate, DD)

Good for quasi-elastic proc.

– but not high-mass dissocⁿ



2. A diffractive process is characterized by a large rapidity gap (LRG), which is caused by **t-channel** Pomeron exch. (or, to be more precise, by the exchange corresponding to the rightmost singularity in the complex angular momentum plane with vacuum quantum numbers).

Only good for very LRG events – otherwise

Reggeon/fluctuation contaminations

Why study diffraction ?

Intrinsic interest. The LHC should reach, for the first time, sufficiently HE to distinguish between the different theoretical asymptotic scenarios for HE interactions.

In HE pp collisions about 40% of σ_{tot} comes from diffractive processes, like elastic scatt., SD, DD.
Need to study diffraction to understand the structure of σ_{tot} and the nature of the underlying events which accompany the sought-after rare hard subprocesses.
(Note LHC detectors do not have 4π geometry and do not cover the whole rapidity interval. So minimum-bias events account for only part of total $\sigma_{\text{inelastic}}$.)

Study needed to estimate the survival probabilities of LRG to soft rescattering.

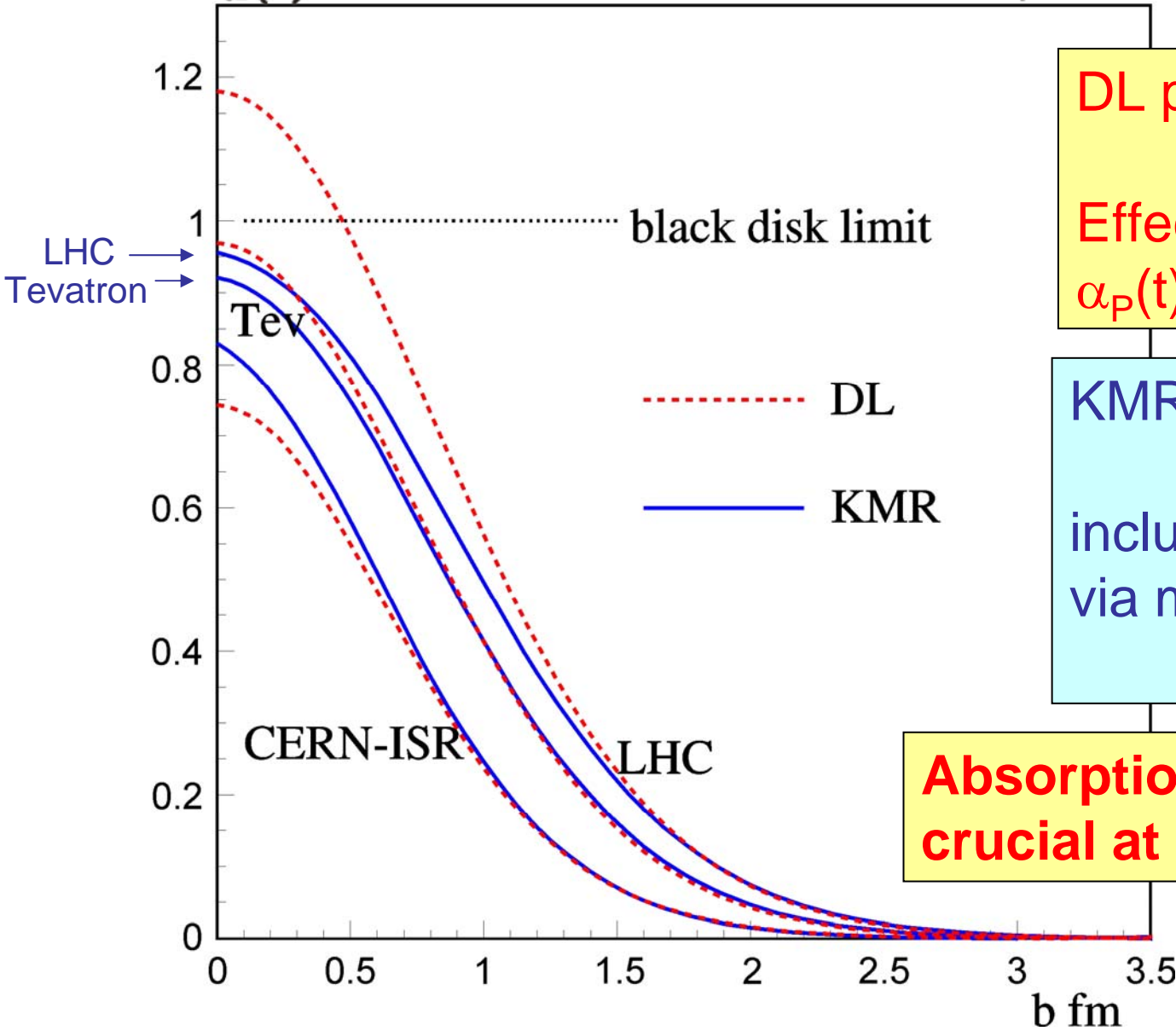
Recall “hard” exclusive diffractive processes (e.g., $pp \rightarrow p + \text{Higgs} + p$) are an excellent means of suppressing the background for New Physics signals

Needed so as to understand the structure of HE cosmic ray phenomena (e.g. Auger experiment).

Finally, the hope is that a study of diffraction may allow the construction of a MC which merges “soft” and “hard” HE hadron interactions in a reliable and consistent way.

$$\text{Im}T_{\text{el}}(b) = \int \sqrt{\frac{d\sigma_{\text{el}}}{dt} \frac{16\pi}{1+\rho^2}} J_0(qb) \frac{qdq}{2\pi}$$

$\text{Im}T_{\text{el}}(b)$

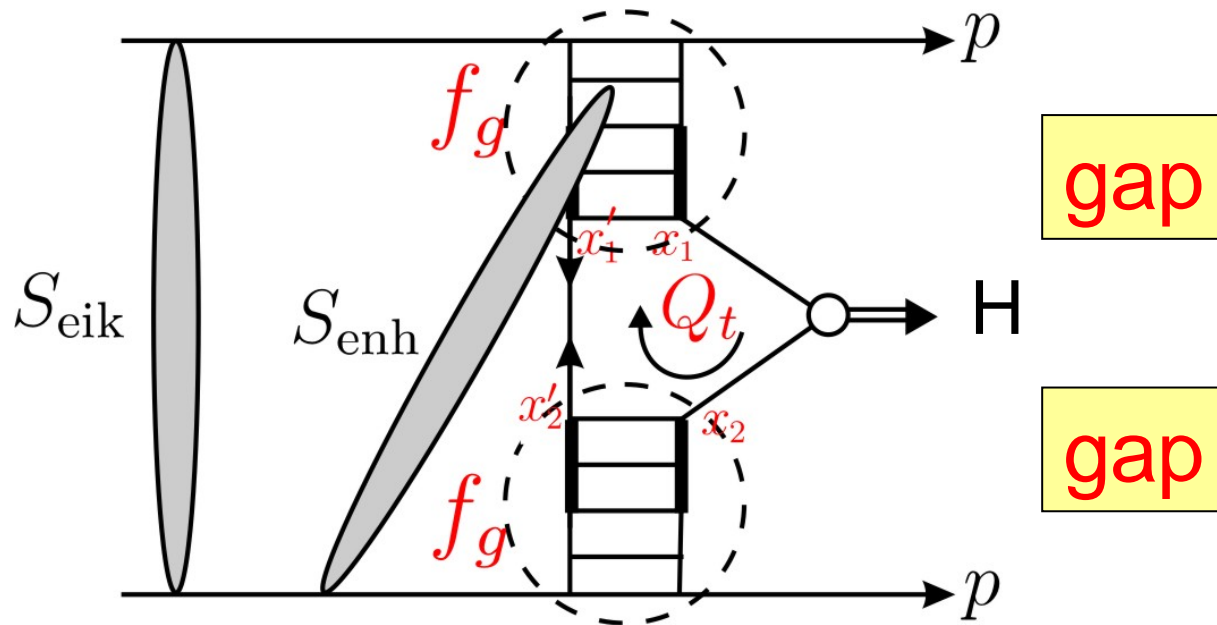


DL parametrization:
Effective Pom. pole
 $\alpha_P(t) = 1.08 + 0.25t$

KMR parametrization
includes absorption
via multi-Pomeron
effects

Absorption/ s-ch unitarity
crucial at small b at LHC

... “soft” scatt. can easily destroy the gaps



soft-hard
factorizⁿ

eikonal rescatt: between protons

← conserved

enhanced rescatt: involving intermediate partons

← broken

→ need a model for soft physics at HE, see later..

Recall the 2nd definition of diffraction

Diffraction is any process caused by **Pomeron exchange**.

(Old convention was any event with LRG of size $\delta\eta > 3$, since Pomeron exchange gives the major contribution)

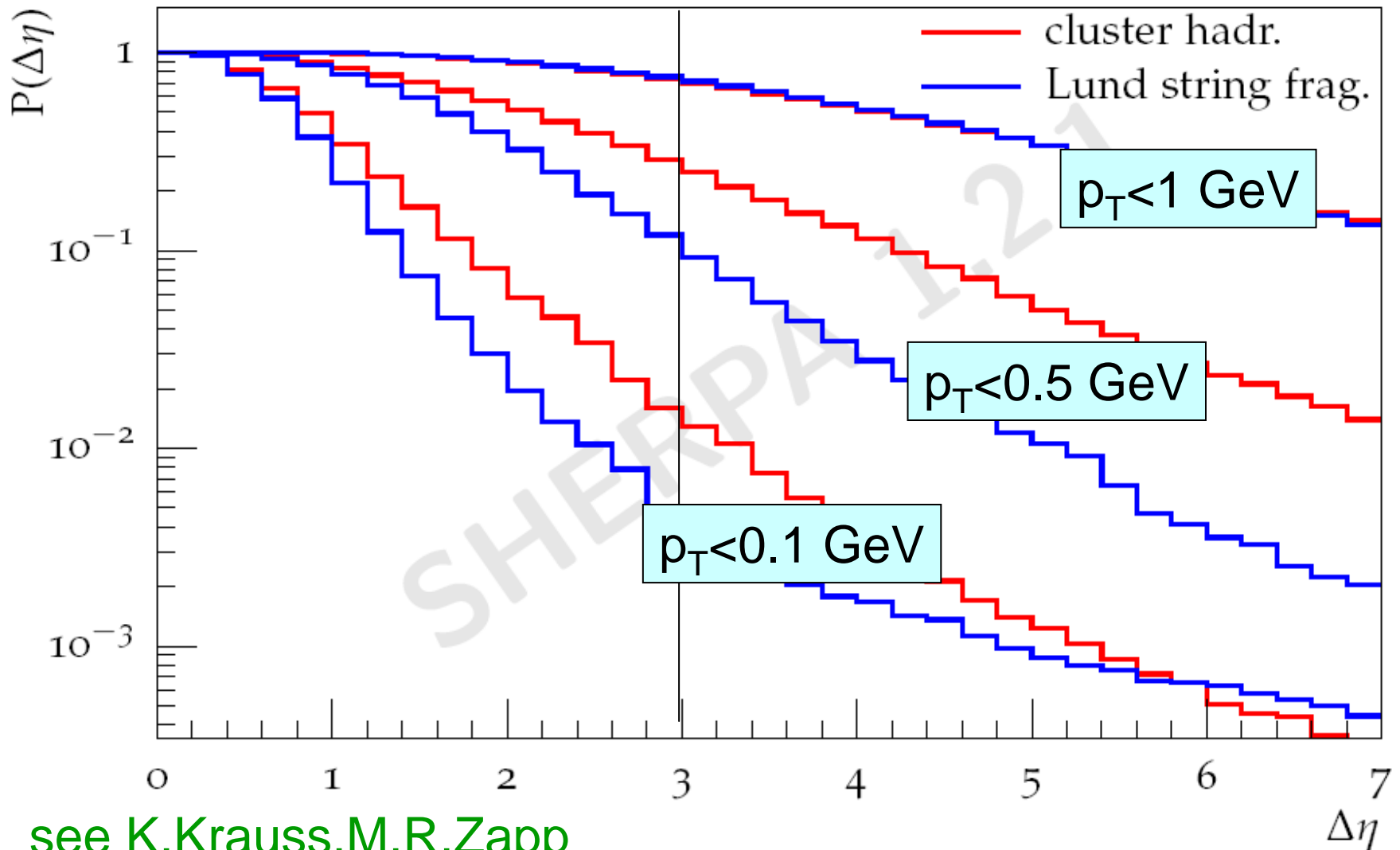
However LRG in the distribution of secondaries can also arise from

- (a) Reggeon exchange
- (b) **fluctuations** during the hadronization process

Indeed, at LHC energies LRG of size $\delta\eta > 3$ do not unambiguously select diffractive events.

Prob. of finding gap larger than $\Delta\eta$ in inclusive event at 7 TeV
due to fluctuations in hadronization

gap anywhere in $-5 < \eta < 5$, different threshold p_T

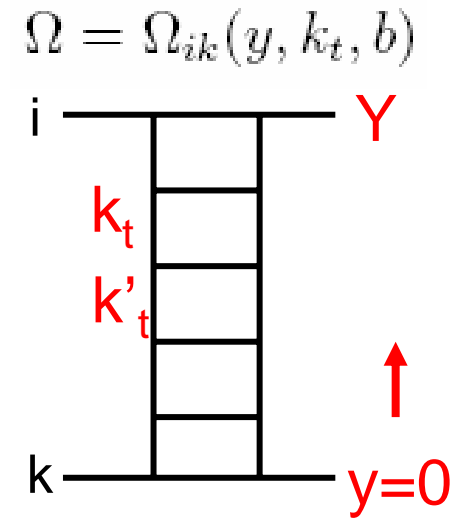


Partonic structure of “bare” Pomeron

BFKL evolⁿ in rapidity generates ladder

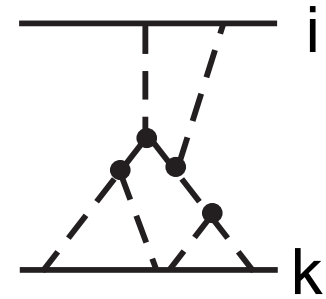
$$\frac{\partial \Omega(y, k_t)}{\partial y} = \bar{\alpha}_s \int d^2 k'_t K(k_t, k'_t) \Omega(y, k'_t)$$

- At each step k_t and b of parton can be changed – so, in principle, we have 3-variable integro-diff. eq. to solve
- **Inclusion of k_t crucial to match soft and hard domains. Moreover, embodies less screening for larger k_t comp^{ts}.**
- We use a simplified form of the kernel K with the main features of BFKL – diffusion in $\log k_t^2$, $\Delta = \alpha_P(0) - 1 \sim 0.3$
- b dependence during the evolution is prop' to the Pomeron slope α' , which is v.small ($\alpha' < 0.05 \text{ GeV}^{-2}$) -- so ignore. Only b dependence comes from the starting evolⁿ distribⁿ
- Evolution gives $\longrightarrow \Omega = \Omega_{ik}(y, k_t, b)$



Multi-Pomeron contributions

Now include rescatt of intermediate partons with the “beam” i and “target” k

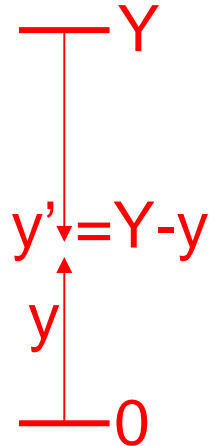


evolve up from $y=0$

$$\frac{\partial \Omega_k(y)}{\partial y} = \bar{\alpha}_s \int d^2 k'_t \exp(-\lambda(\Omega_k(y) + \Omega_i(y'))/2) K(k_t, k'_t) \Omega_k(y)$$

evolve down from $y'=Y-y=0$

$$\frac{\partial \Omega_i(y')}{\partial y'} = \bar{\alpha}_s \int d^2 k'_t \exp(-\lambda(\Omega_i(y') + \Omega_k(y))/2) K(k_t, k'_t) \Omega_i(y')$$



where $\lambda \Omega_{i,k}$ reflects the different opacity of protons felt by intermediate parton, rather the proton-proton opacity $\Omega_{i,k}$ $\lambda \sim 0.2$

solve iteratively for $\Omega_{i,k}(y, k_t, b)$

inclusion of k_t crucial

Note: data prefer $\exp(-\lambda \Omega) \rightarrow [1 - \exp(-\lambda \Omega)] / \lambda \Omega$

Form is consistent with generalisation of AGK cutting rules

In principle, knowledge of $\Omega_{ik}(y, k_t, b)$ allows the description of all soft, semi-hard pp high-energy data:

σ_{tot} , $d\sigma_{\text{el}}/dt$, $d\sigma_{\text{SD}}/dtdM^2$, DD, DPE...

LRG survival factors S^2 (to both eikonal, enhanced rescatt)

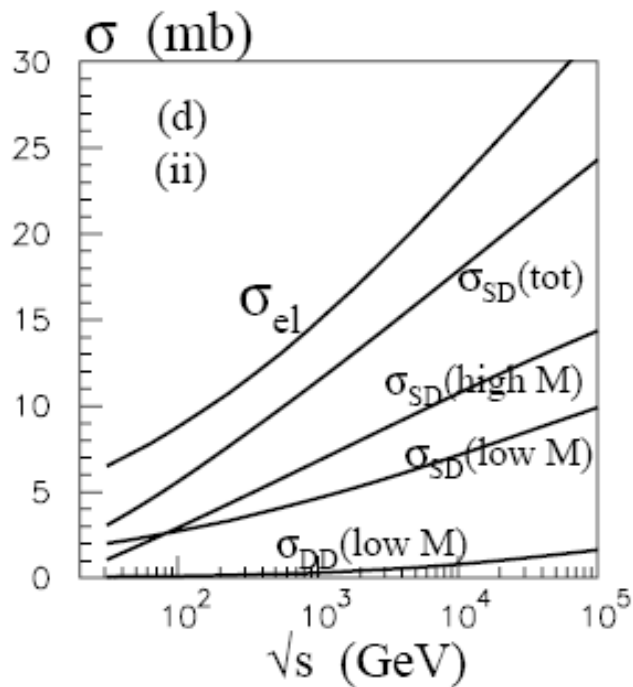
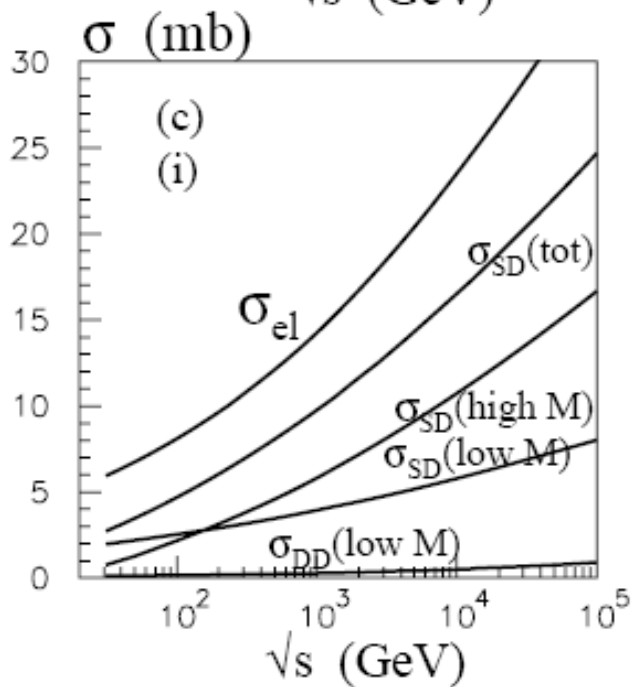
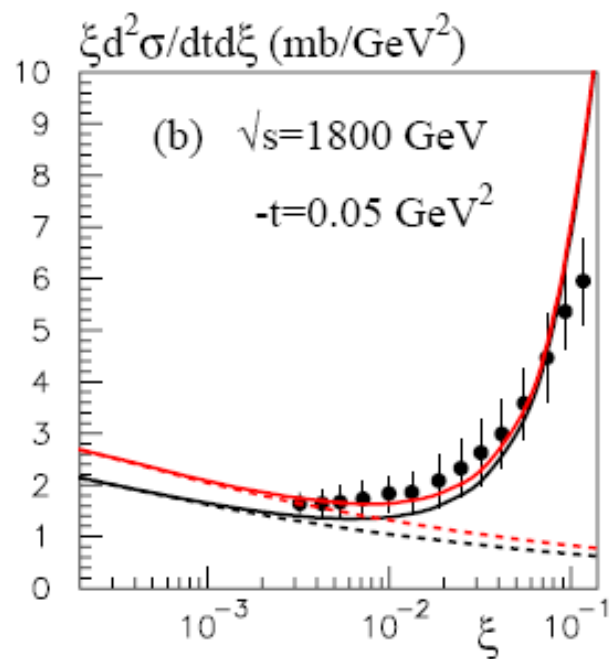
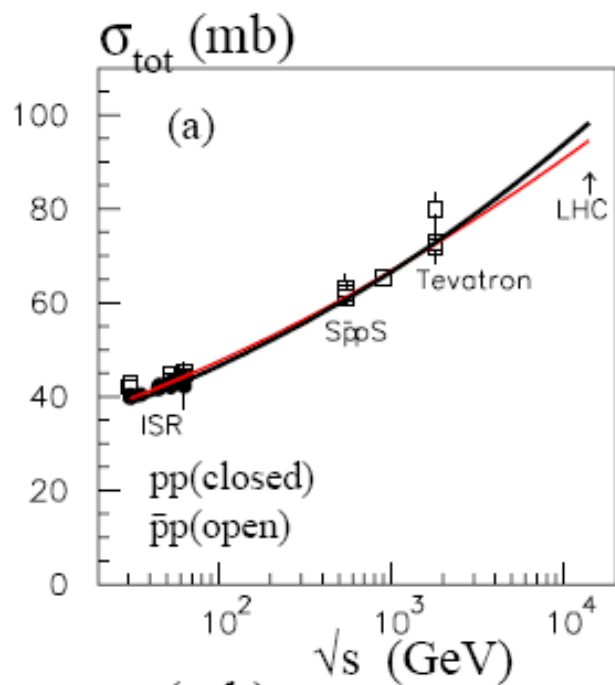
PDFs and diffractive PDFs at low x and low scales

Indeed, such a model can describe the main features of all the data, in a semi-quantitative way, with just a few physically motivated parameters.

energy	σ_{tot}	σ_{el}	$\sigma_{\text{SD}}^{\text{low}M}$	$\sigma_{\text{SD}}^{\text{high}M}$	$\sigma_{\text{SD}}^{\text{tot}}$	$\sigma_{\text{DD}}^{\text{low}M}$
1.8	72.8/72.5	16.3/16.8	4.4/5.2	8.3/11.1	12.7/16.3	0.8/1.7
7	89.0/86.8	21.9/21.6	5.5/6.7	12.0/14.4	17.5/21.1	1.4/2.8
14	98.3/94.6	25.1/24.2	6.1/7.5	14.0/15.9	20.1/23.4	1.8/3.6
100	127.1/117.4	35.2/31.8	8.0/9.9	20.6/20.0	28.6/29.9	3.4/6.2

$d\sigma/dy \sim s^{0.2}$ like the LHC data for 0.9 to 7 TeV

$S^2 = 0.010-0.016$ for gaps in $pp \rightarrow p + H + p$ (120GeV SM Higgs at 14TeV)



Conclusions

- **s-ch unitarity** is important for quasi-elastic scatt or LRG events
- **Multi-Pomeron** exchange diagrams restore unitarity:
(i) **eikonal** pp rescatt. (ii) **enhanced** with intermediate partons
- Altho' $g_{3P} \sim 0.2g_N$, high-mass p dissociation is **enhanced** at the LHC
- Unitarity is restored for LRG by small **survival prob.** S^2 of gaps
e.g. $S^2 \sim 0.015$ for $pp \rightarrow p + H + p$ ($M_H = 120$ GeV at 14 TeV)
- LRG also from **fluct^{ns} in hadⁿ**: study different p_T cuts and Δy
also study long-range rapidity correlations at the LHC
- QCD/BFKL Pom. \rightarrow Pomeron describing soft physics
- Partonic struct. of Pom, with multi-Pom contrib^{ns} can describe
all soft ($\sigma_{tot,el,SD..}$) and semihard (**PDFs, minijets..**) physics - **KMR**
- Forms the basis of **"all purpose" MC** - **Krauss, Hoeth, Zapp**

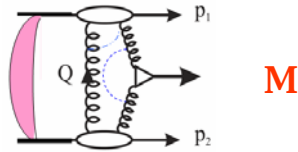
CEP as a way to study old and new heavy resonances.



Heavy Quarkonia



Zoo of charmonium -like XYZ states



- $\chi_{c,b}$ production is of special interest: (star reactions!)
 - Heavy quarkonium production can shed light on the physics of bound states (lattice, NRQCD. . .).
 - Potential to produce different J^P states, which exhibit characteristic features (e.g. angular distributions of forward protons).
 - Could shed light on the nature of the various 'exotic' charmonium states observed recently.

(X,Y,Z) charmonium-like states.

Spin-Parity Analyzer

(KKMR-2003)

CDF & new LHCb measurements are all in good agreement (factor "few") with the Durham group predictions.

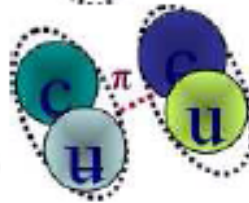


Zoo of charmonium -like XYZ states

Tetraquark:
four tightly bound quarks



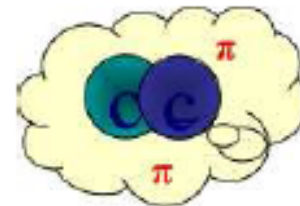
Molecular state:
two loosely bound mesons



Hybrid: states with
excited gluonic degrees of freedom



Hadrocharmonium: charmonium state,
"coated" by excited light-hadron matter



- X(3872) -
- XYZ(3940) & X(3915) -
- Y(4140)/Y(4280) & X(4350)

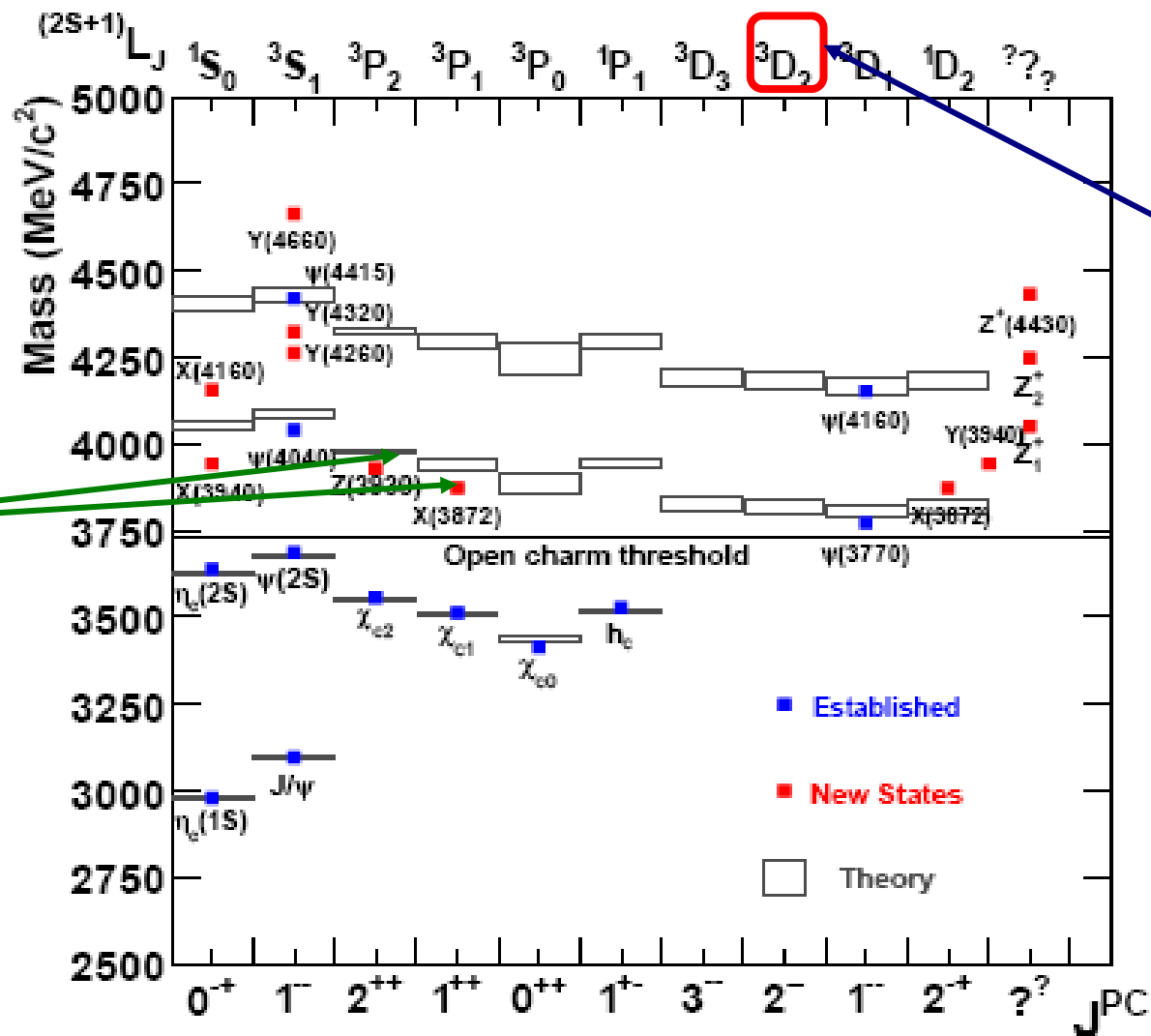


Figure 1: The mass versus the quantum numbers (J^{PC}) for the charmonium-like states. The boxes represent the predictions; blue boxes show the established states, and the red boxes indicate the new states discovered at the B -factories.



Overview of exotic/charmonium spectroscopy (2)

State	M , MeV	Γ , MeV	J^{PC}	Process
$X(3872)$	3871.52 ± 0.20	1.3 ± 0.6 (< 2.2)	$1^{++}/2^{-+}$	$B \rightarrow K(\pi^+\pi^- J/\psi)$ $p\bar{p} \rightarrow (\pi^+\pi^- J/\psi) + \dots$ $B \rightarrow K(\omega J/\psi)$ $B \rightarrow K(D^{*0}D^0)$ $B \rightarrow K(\gamma J/\psi)$ $B \rightarrow K(\gamma\psi(2S))$
$X(3915)$	3915.6 ± 3.1	28 ± 10	$0/2^{?+}$	$B \rightarrow K(\omega J/\psi)$ $\gamma\gamma \rightarrow (\omega J/\psi)$
$X(3940)$	3942_{-8}^{+9}	37_{-17}^{+27}	$?^{?+}$	$e^+e^- \rightarrow J/\psi(D\bar{D}^*)$ $e^+e^- \rightarrow J/\psi(\dots)$
$Y(4008)$	4008_{-49}^{+121}	226 ± 97	1^{--}	$e^+e^- \rightarrow \gamma(\pi^+\pi^- J/\psi)$
$Z_1(4050)^+$	4051_{-43}^{+24}	82_{-55}^{+51}	$?$	$B \rightarrow K(\pi^+\chi_{c1}(1P))$
$Y(4140)$	4143.4 ± 3.0	15_{-7}^{+11}	$?^{?+}$	$B \rightarrow K(\phi J/\psi)$
$X(4160)$	4156_{-25}^{+29}	139_{-65}^{+113}	$?^{?+}$	$e^+e^- \rightarrow J/\psi(D\bar{D}^*)$
$Z_2(4250)^+$	4248_{-45}^{+185}	177_{-72}^{+321}	$?$	$B \rightarrow K(\pi^+\chi_{c1}(1P))$
$Y(4260)$	4263 ± 5	108 ± 14	1^{--}	$e^+e^- \rightarrow \gamma(\pi^+\pi^- J/\psi)$ $e^+e^- \rightarrow (\pi^+\pi^- J/\psi)$ $e^+e^- \rightarrow (\pi^0\pi^0 J/\psi)$
$Y(4360)$	4353 ± 11	96 ± 42	1^{--}	$e^+e^- \rightarrow \gamma(\pi^+\pi^- \psi')$
$Z(4430)^+$	4443_{-18}^{+24}	107_{-71}^{+113}	$?$	$B \rightarrow K(\pi^+\psi(2S))$
$X(4630)$	4634_{-11}^{+9}	92_{-32}^{+41}	1^{--}	$e^+e^- \rightarrow \gamma(\Lambda_c^+\Lambda_c^-)$
$Y(4660)$	4664 ± 12	48 ± 15	1^{--}	$e^+e^- \rightarrow \gamma(\pi^+\pi^-\psi(2S))$

● **X(3872)** first and most puzzling state
(observed in 2003 at Belle)



- Discovered by BELLE in 2003, confirmed by BaBar, CDF, D0
- Possible spin-parity assignment: 1^{++} or 2^{-+}
- May well be of exotic nature : loosely bound molecule, diquark-antidiquark, hybrid,..... but a conventional 2 P-wave charmonium interpretation is still on the table (recent renewal of interest).
- BaBar (2010) seems to favour 2^{-+} though various theory groups find this assignment highly problematic.
- According to PDG $\Gamma(\pi^+ \pi^- J/\psi(1S))/\Gamma_{\text{total}} > 2.6\%$; $\Gamma(\gamma\psi(2S))/\Gamma_{\text{total}} > 3.0\%$, $\Gamma < 2.3$ MeV.
(maybe two different states X(3872), X(3875))

CEP as a spin-parity analyzer could help to resolve the X(3872) puzzle.

● **Z(3930) $\equiv \chi_{c2}(2P)$**

■ Above DD threshold .

■ Vertex detection at LHCb & RHIC→

■ **exclusive open charm: $D^+ D^- , D^0 \bar{D}^0$,**

■ Roughly the same expectations for CEP as for χ_{c2}

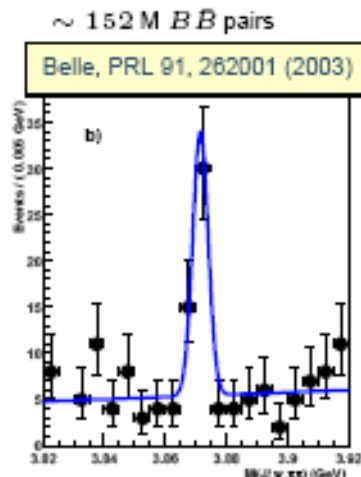
Triggering on J/ψ: $M \rightarrow J/\psi + \gamma, J/\psi + \rho \dots$



$X(3872)$ properties

- ◆ B-factories
- ◆ Spectroscopy overview
- ◆ η_c & $\eta_c(2S)$ properties
- ◆ **$X(3872)$**
- ◆ $\omega J/\psi$
- ◆ Charmonium in $\Upsilon(1S)$ decays
- ◆ Summary
- ◆ Backup

Discovered by Belle in $B^\pm \rightarrow J/\psi \pi^+ \pi^- K^\pm$ decays. Confirmed by D0, CDF, BaBar. What do we know about it?



$$M(X) - (m_{D^0} + m_{D^{*0}}) = -0.32 \pm 0.35 \text{ MeV}$$

$$X(3872) \rightarrow \pi^+ \pi^- J/\psi$$

$M(\pi^+ \pi^-)$ is consistent with $\rho \Rightarrow$
 large isospin violation;
 $C = +1.$

$$X(3872) \rightarrow \pi^+ \pi^- \pi^0 J/\psi \Rightarrow$$

$M(\pi^+ \pi^- \pi^0)$ is consistent with ω ,
 different isospin then $\rho J/\psi.$

Angular analysis: $J^{PC} = 1^{++}$ or 2^{-+}

S -wave $D^0 \bar{D}^{*0}$ molecular state?

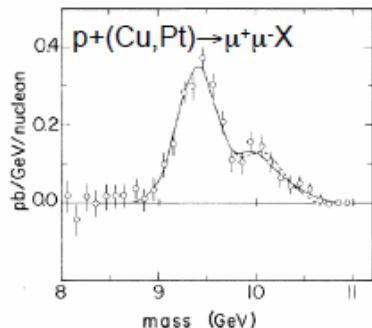
E.S.Swanson, PLB 588, 189 (2004)

P-wave Bottomonia

Bottomonium history started 30 years ago

(PRL 39, 242 (1977) and PRL 39,1240 (1977))

30 years later....



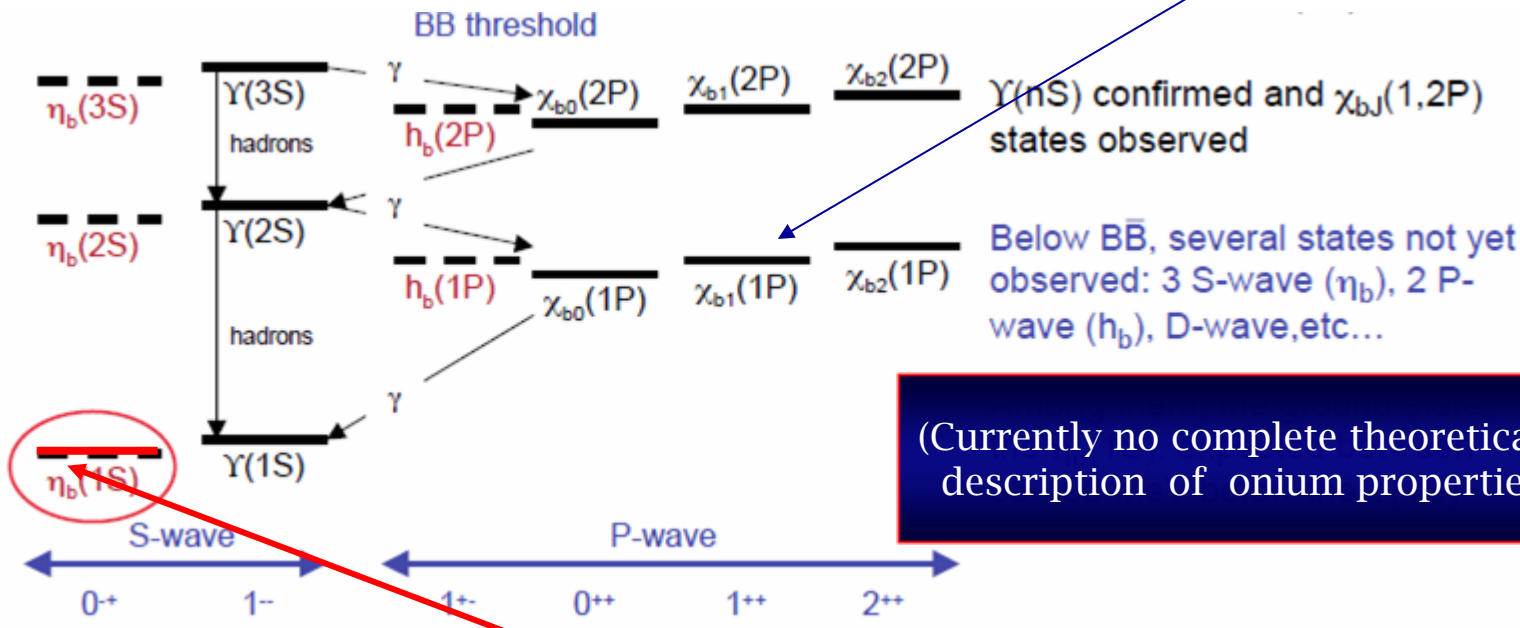
FNAL, E288

$M(\Upsilon) = 9.40 \pm 0.013$

$M(\Upsilon') = 10.00 \pm 0.04$

$M(\Upsilon'') = 10.43 \pm 0.12$

(spins- still unconfirmed)



(Currently no complete theoretical description of onium properties.)

(BABAR (2008))

(Still some puzzles)



The heaviest and most compact quark-antiquark bound state in nature

$\chi_{b0}(1P)$ ^[d]

$I^G(J^{PC}) = 0^+(0^{++})$

J needs confirmation.

Mass $m = 9859.44 \pm 0.42 \pm 0.31$ MeV

$(1.73 \pm 0.35)\%$
CLEO-2011

$\chi_{b0}(1P)$ DECAY MODES	Fraction (Γ_i/Γ)	Confidence level	p (MeV/c)
$\gamma \Upsilon(1S)$	<6 %	90%	391

(Crystal Ball-1986)

$\chi_{b1}(1P)$ ^[d]

$I^G(J^{PC}) = 0^+(1^{++})$

J needs confirmation.

$\gamma \Upsilon(1S)$	$(35 \pm 8) \%$		423
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$\chi_{b2}(1P)$ ^[d]

$I^G(J^{PC}) = 0^+(2^{++})$

J needs confirmation.

Mass $m = 9912.21 \pm 0.26 \pm 0.31$ MeV

$\chi_{b2}(1P)$ DECAY MODES	Fraction (Γ_i/Γ)	p (MeV/c)
$\gamma \Upsilon(1S)$	$(22 \pm 4) \%$	442



Towards Full Acceptance Detector (bj- 1992)



CMS (& ATLAS) currently blind between $\eta = 6.4$ (CASTOR) and beam rapidity ($y_p = 8.9 @ 7 \text{ TeV}$) except ZDC (neutrals). Cannot distinguish most diffractive/non-diffractive events.

IS THERE A WAY OUT ?

Yes, an addition of **Forward Shower Counters** around beam pipes at CMS!

(8 FSC per side see showers from particles with $|\eta| = 7-9$)



(Alice is installing such counters, ongoing studies for LHCb)

(FSC → at least a good foot in the door)



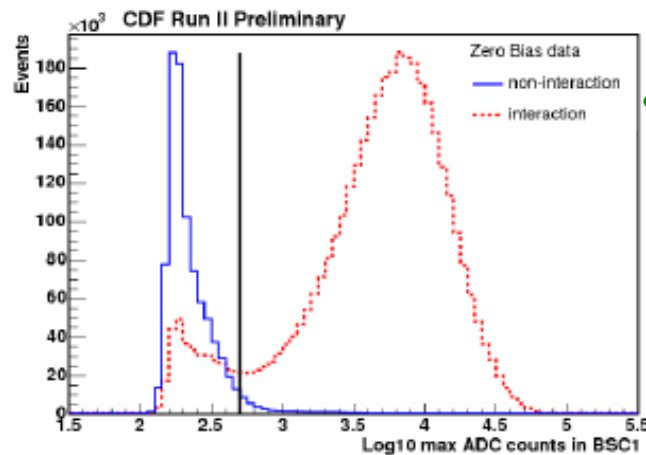
**BSC very important as rap gap detectors.
All LHC experiments should have them!**

FORWARD PHYSICS WITH RAPIDITY GAPS AT THE LHC

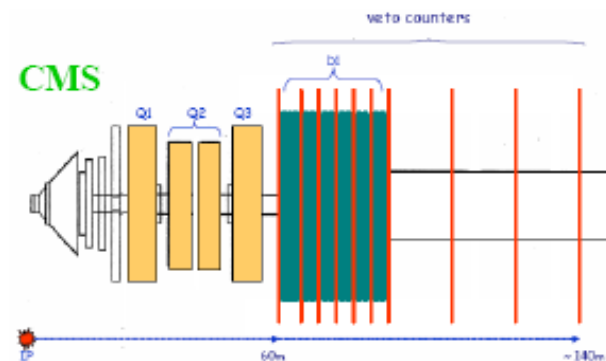
arXiv:0811.0120

Michael Albrow¹, Albert De Roeck², Valery Khoze³, Jerry Lämsä^{4,5}, E. Norbeck⁶,
Y. Onel⁶, Risto Orava⁵, and M.G. Ryskin⁷
Sunday, November 09, 2008

JINST-09



Warm accessible vacuum pipe (circular – elliptical)
Do not see primary particles, but showers in pipe ++
Simple scintillator paddles: **Gap detectors in no P-U events**



Take 0-bias events (Essential!)

{1} = prob no interaction

{2} = prob ≥ 1 interaction

Take hottest PMT of 8 BSC1

Plot log max ADC for {1} and {2}

Separates empty / not empty

Repeat for all detectors





July 19, 2010

Physics and Beam Monitoring with Forward Shower Counters (FSC) in CMS

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Oleg Grachov, Michael Murray

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Andrei Sobol, Vladimir Samoylenko

IHEP, Protvino, Russia

Aldo Penzo



CMS NOTE-2010/015

**Approved by CMS MB
for Jan-Feb 2011 installation.**

“Limited approval” :
Go ahead without detracting from
necessary shutdown work.

Most value is 2011 running
& when $\langle n/x \rangle < \sim 5$
(Do not expect to use > 2012)

Station 3 (114m) Installed on both sides.
March Techn. Stop (28-31.03.11).
Stations 1&2- installed during May
Techn. stop. Commissioning with beam.

The FSC- these are for real !

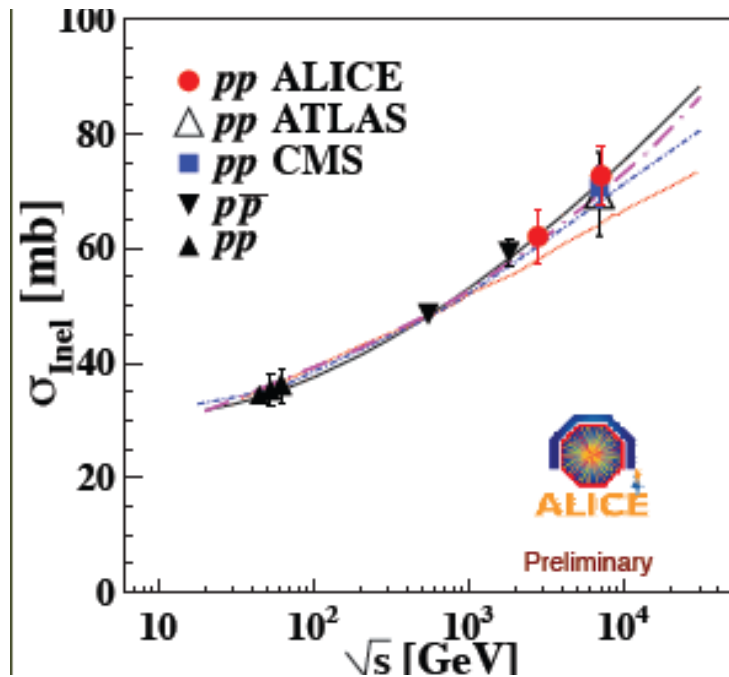
- The installation and commissioning phase of FSC during the March Technical Stop.
- Main concern- lumi per bunch crossing might be too high.

What about precise measurement of SD?

Don't hold your breath, Valery. This certainly needs all the counters and some low lumi run, or at least bunches. (Mike Albrow)



Recent LHC results



ATLAS $\sigma_{inel}(\xi > m_p^2/s) = 69.4 \pm 2.4(\text{exp.}) \pm 6.9(\text{extr.}) \text{ mb}$

CMS $66.8 \leq \sigma_t^{inel} \leq 74.8 \text{ mb.}$

ALICE $\sigma_{Inel}(\sqrt{s} = 7 \text{ TeV}) = 72.7 \pm 1.1(\text{model}) \pm 5.1(\text{lum.}) \text{ mb}$

$|\eta| < 2. \quad -3.7 < \eta < -1.7 \text{ and } 2.8 < \eta < 5.1.$

Gostman et al., arXiv:1010.5323, EPJ. C74, 1553 (2011)
 Kaidalov et al., arXiv:0909.5156, EPJ. C67, 397 (2010)
 Ostapchenko, arXiv:1010.1869, PR D83 114018 (2011)
 Khoze et al., EPJ. C60 249 (2009), C71 1617 (2011)



Can we measure σ_{inel} , σ_{SD} , σ_{DD} with high accuracy?

Achilles' Heel of 'inelastic' measurements : low mass SD,DD

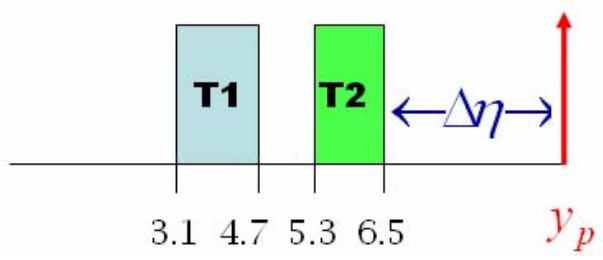
Uninstrumented regions: Totem-CMS :
(at least)

$$M_X \leq 2.5 - 3.5 GeV$$

Atlas: $M_X \leq 7 GeV$ (<15.7 GeV)



(Castor)

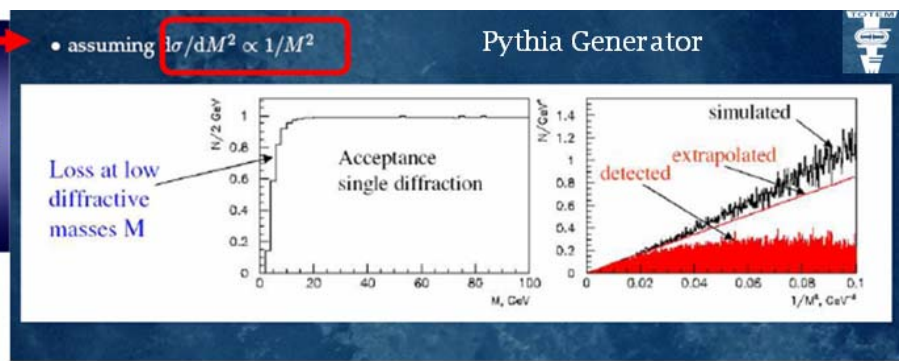


$$\eta = -\ln \tan \frac{\vartheta}{2}$$

$$y_p = \ln(\sqrt{s} / m_p), \Delta\eta \approx (2.4 - 3.1)$$

Can we extrapolate from HM SD ?

- Theoretically unjustified
- Currently **NO** theoretically solid way to extrapolate HM to LM single diffraction



(UA4-experience \times factor of 2 for $M < 4 GeV$)

A diffractive process is characterized by a large rapidity gap (LRG), which is caused by **t-channel** Pomeron exch.

$$\sigma_{\text{total}} = \sum_{\mathbf{X}} \left| \text{Diagram}(\mathbf{X}) \right|^2 = \text{Im} \left[\text{Diagram}(\text{cut}) \right] = \alpha_{IP}(0)$$

High mass diffractive dissociation

$d^2\sigma/dM^2 dt|_{t=0} \sim \frac{s^{2\epsilon}}{(M^2)^{1+\epsilon}} S^2 \sim 1/M^2$

PPP-diagram

Screening is very important. (semi) enhanced absorption

... (t-dependence !?)

Low mass diffractive dissociation

introduce diff^{ve} estates ϕ_i, ϕ_k (comb^{ns} of p, p*, ...) which **only** undergo "elastic" scattering (Good-Walker)

dual to

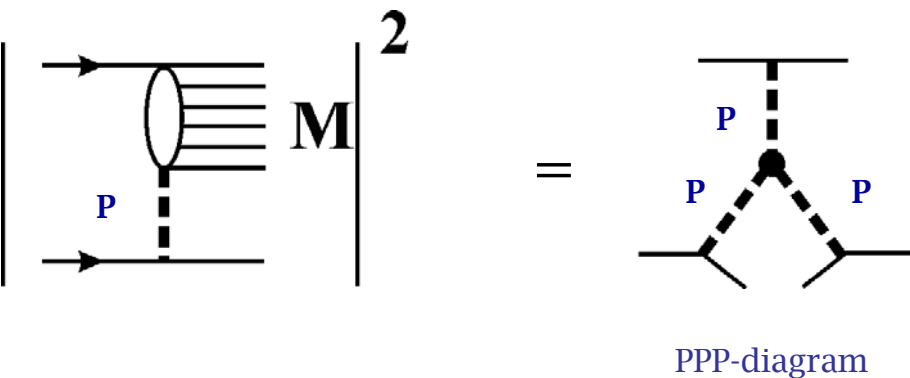
$\sim \frac{s^{2\epsilon}}{(M^2)^{1.5+2\epsilon}} S^2 \sim 1/M^3$

PPR-diagram

A diffractive process is characterized by a large rapidity gap (LRG), which is caused by **t-channel** Pomeron exch.

$$\sigma_{\text{total}} = \sum_{\mathbf{X}} \left| \text{Diagram}(\mathbf{X}) \right|^2 = \text{Im} \left[\text{Diagram}(\text{cut}) \right] = \alpha_{IP}(0)$$

High mass diffractive dissociation



$$d^2\sigma/dM^2 dt|_{t=0} \sim \frac{s^{2\epsilon}}{(M^2)^{1+\epsilon}} S^2 \sim 1/M^2$$

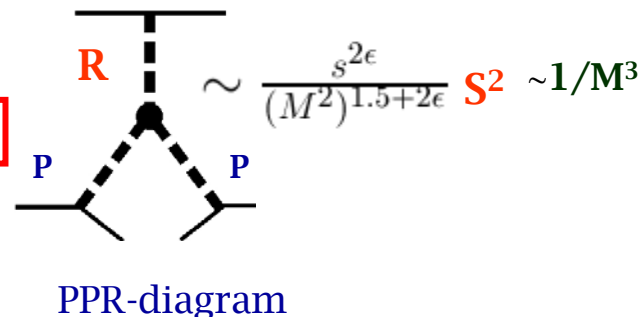
Screening is very important. (semi) enhanced absorption

... (t-dependence !?)

Low mass diffractive dissociation

introduce diff^{ve} estates ϕ_i, ϕ_k (comb^{ns} of p, p*, ...) which **only** undergo "elastic" scattering (Good-Walker)

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$$\sim \frac{s^{2\epsilon}}{(M^2)^{1.5+2\epsilon}} S^2 \sim 1/M^3$$

Model expectations for total inelastic cross-section

- Strong dependence of the longitudinal development of air showers on

$$\sigma_{inel}$$

Various MC generators are used by the CR community (some with full resummation of multi-Pomeron graphs)

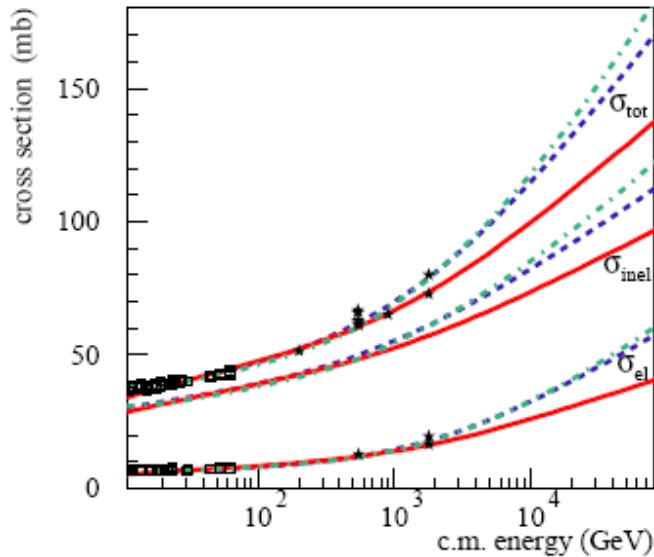


Figure 1: Model predictions for total, elastic, and inelastic proton-proton cross sections: QGSJET-II-4 - solid, QGSJET-II-3 - dashed, and SIBYLL - dot-dashed. The compilation of data is from Ref. [17].

S.Ostapchenko, ArXiv:1103.5684

$\sqrt{s} = 7 \text{ TeV}$

(in mb)	σ_{inel}	$\sigma_{SD}^{LM} + \sigma_{DD}^{LM}$
QGSJET II-04	69.7	7.1
QGSJET II-03	77.5	3.3
SIBYLL	79.6	0
PYTHIA	71.5	0

KMR-11 **65.2/67.1** **6/7.4**

Dino-11 **71 ± 6**

MPS-11 **73.4**

KP-10 **71.6**

For illustration purposes only

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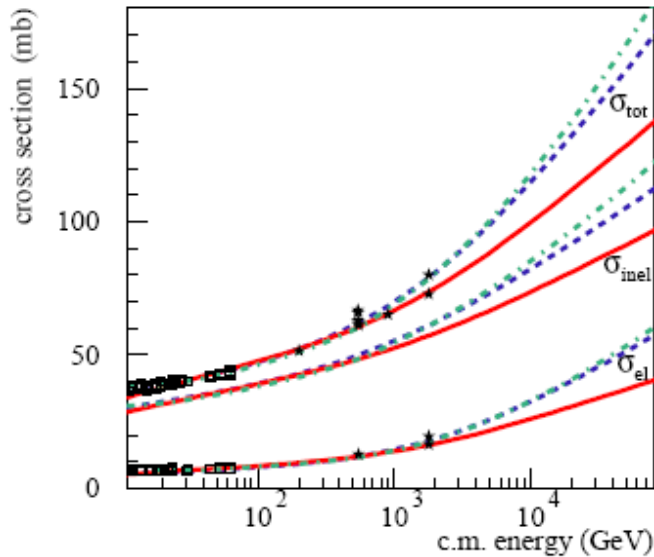


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For illustration purposes only

What about current theoretical uncertainties ?

$$\sqrt{s} = 14 \text{ TeV.}$$

	σ^{tot}	σ^{el}	σ^{SD}	σ^{DD}	$\sigma_{\text{LM}}^{\text{SD}}$	$\sigma_{\text{HM}}^{\text{SD}}$	$\sigma_{\text{LM}}^{\text{DD}}$	$\sigma_{\text{HM}}^{\text{DD}}$
Set (A)	128	37.5	12.1	4.61	8.48	3.92 (3.54)	1.15	2.06
Set (B)	126	37.3	12.4	5.18	8.22	4.24 (4.14)	1.08	2.50
Set (C)	114	33.0	11.0	4.83	5.05	5.22 (5.12)	0.47	3.15
KMR-08	91.7	21.5	19.0		4.9	14.1		
GLMM-08	92.1	20.9	11.8	6.08	10.5	1.28		

KP-10

108

29.5

14.3

For illustration purposes only

(A,B,C) S. Ostapchenko, Phys.Rev.D81:114028,2010.
 KMR-08: KMR, EPJ C54,199(2008); ibid C60,249 (2009).
 GLMM-08: GLMM,EPJ C57,689 (2008).
 KP-10 A.B. Kaidalov, M.Poghosyan

Large variation of $\sigma_{\text{LM}}^{\text{SD}}$ in the range 5- 10.5 mb

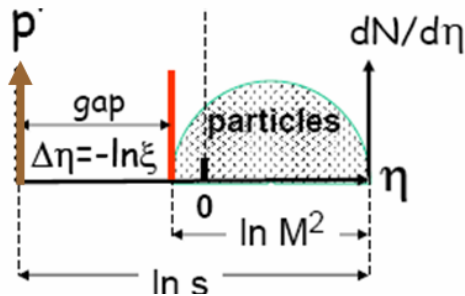


5. A flavour of diffraction in the first LHC runs.

First measurement of $\sigma(\xi > 5 \times 10^{-6})$ at 7 TeV.



(arXiv:1104.0326 [hep-ex] , 2 Apr. 2011)



$$\xi = M_X^2/s > 5 \times 10^{-6}$$

$$M_X > 15.7 \text{ GeV for } \sqrt{s} = 7 \text{ TeV}$$

$\sigma(\xi > 5 \times 10^{-6})$ [mb]	
ATLAS Data 2010	$60.33 \pm 2.10(\text{exp.})$
Schuler and Sjöstrand	66.4
PHOJET	74.2
Ryskin <i>et al.</i>	51.8 / . 56.2
$\sigma(\xi > m_p^2/s)$ [mb]	
ATLAS Data 2010	$69.4 \pm 2.4(\text{exp.}) \pm 6.9(\text{extr.})$
Schuler and Sjöstrand	71.5
PHOJET	77.3
Block and Halzen	69
Ryskin <i>et al.</i>	65.2 / . 67.1
Gotsman <i>et al.</i>	68
Achilli <i>et al.</i>	60 – 75

(model dependence in the definition of ξ .)



$$\Delta\eta = \ln 1/\xi + \ln \langle p_{\perp} \rangle / m_p$$



Can we measure $\frac{d\sigma_{el}}{dt}$ or σ_{el} with σ_{tot} good accuracy?

- $\frac{d\sigma_{el}}{dt} = \frac{\pi}{s p^2} |F_{el}(t)|^2$
- optical theorem: $\sigma_{tot} = \frac{4\pi}{p\sqrt{s}} \text{Im} F_{el}(s, t=0)$

With known lumi (3.5% VdM)



- $L\sigma_{tot} = N_{el} + N_{inel}$
 - Need to separate the Coulomb and hadron scattering
- $$\sigma_{tot} = \frac{16\pi}{1 + \rho^2} \frac{\left. \frac{dN_{el}}{dt} \right|_{t=0}}{N_{el} + N_{inel}}; \quad L = \frac{1 + \rho^2}{16\pi} \frac{(N_{el} + N_{inel})^2}{\left. \frac{dN_{el}}{dt} \right|_{t=0}}$$

(Lumi independent)

Model	ρ
Islam et al.	0.123
Petrov et al. 2P	0.0968
Petrov et al. 3P	0.111
BSW	0.121
Block-Halzen	0.114
COMPETE	0.1316

● measure $\frac{dN_{el}}{dt}$ and extrapolate it to $t = 0 \rightarrow$ needs RP acceptance at small t
 small beam divergence \rightarrow high β^* (parallel to point focusing)

● ALFA- measurement of elastic scattering in the Coulomb interference region

● Will require special LHC runs at high β^* and low \mathcal{L}_{int} : 90m (2011), 2km (2013+)

How Large is Large ?



Diffraction is any process caused by **Pomeron exchange**.

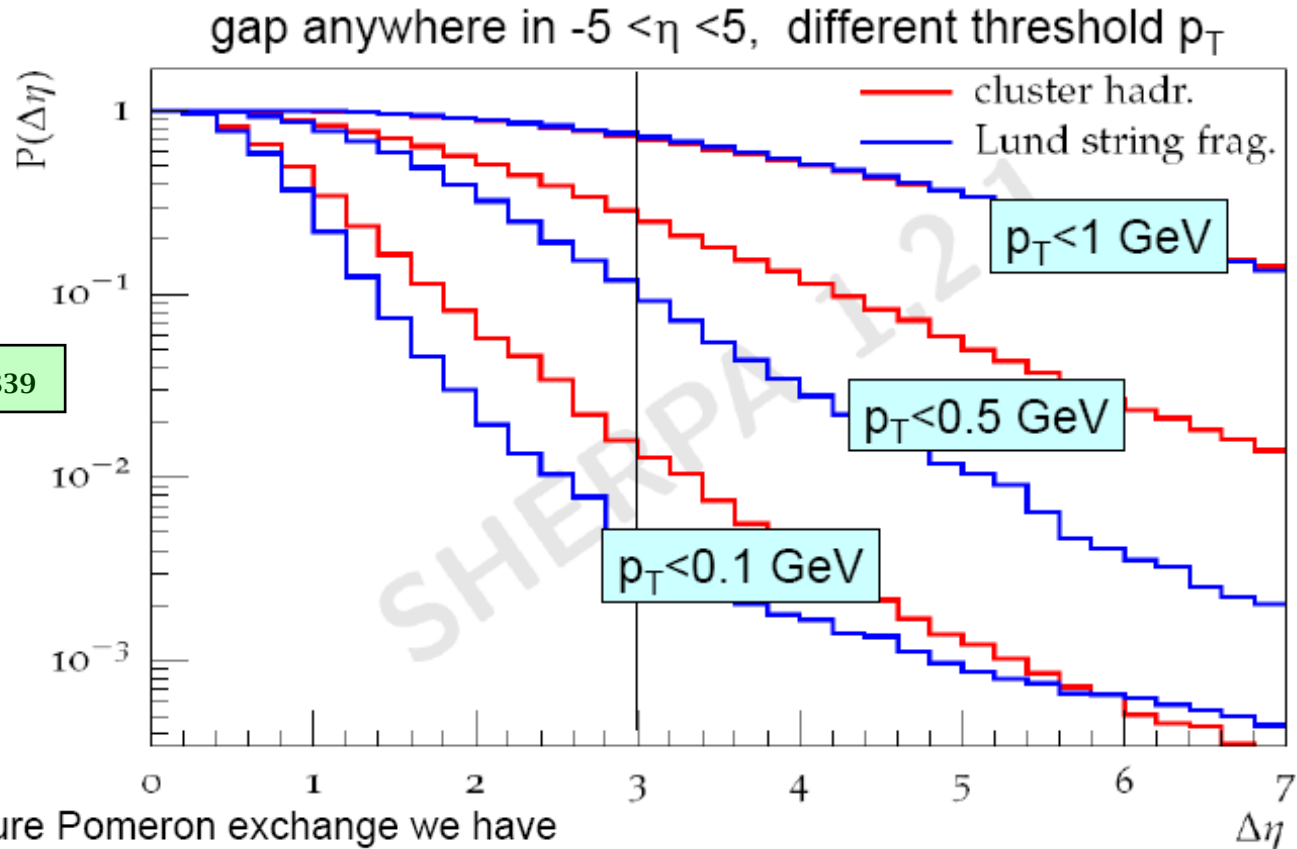
(Old convention was any event with LRG of size $\delta\eta > 3$, since Pomeron exchange gives the major contribution)

However LRG in the distribution of secondaries can also arise from

- (a) Reggeon exchange
- (b) **fluctuations** during the hadronization process

Indeed, at LHC energies LRG of size $\delta\eta > 3$ do not unambiguously select diffractive events.

Prob. of finding gap larger than $\Delta\eta$ in inclusive event at 7 TeV
due to fluctuations in hadronization



KKMRZ, arXiv:1005.4839

So to study pure Pomeron exchange we have

either to select much larger gaps

or to study the Δy dependence of the data, fitting so as to subtract the part caused by Reggeon and/or fluctuations.