PARAMETER ESTIMATION FOR INFLATIONARY GRAVITATIONAL WAVE BACKGROUNDS WITH LISA

BASED ON PROJECT 15 OF THE COSWG

11TH LISA COSWG MEETING JUN 17 2024

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STOCHASTIC GRAVITATIONAL WAVE BACKGROUNDS





DATA & ANALYSIS PIPELINE

Data:

$$\tilde{d}_i(f) = \int_{-T/2}^{T/2} d_i(t) e^{-2\pi i f t} dt$$

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Channel *i*

Data:

$$\tilde{d}_i(f) = \int_{-T/2}^{T/2} d_i(t) e^{-2\pi i f t} dt$$

Observation time

Data:

$$\tilde{d}_i(f) = \int_{-T/2}^{T/2} \frac{d_i(t) e^{-2\pi i f t} dt}{\int}$$

Time domain data stream

Data:

$$\tilde{d}_{i}(f) = \int_{-T/2}^{T/2} d_{i}(t) e^{-2\pi i f t} dt = \sum_{\nu} \tilde{n}_{i}^{\nu}(f) + \sum_{\sigma} \tilde{s}_{i}^{\sigma}(f)$$

<u>Assumption 1</u>: only stochastic components.

Transients, deterministic sources and glitches

are removed

Data:

$$\tilde{d}_i(f) = \int_{-T/2}^{T/2} d_i(t) \, \mathrm{e}^{-2\pi i f t} \, \mathrm{d}t = \sum_{\nu} \tilde{n}_i^{\nu}(f) + \sum_{\sigma} \tilde{s}_i^{\sigma}(f)$$

<u>Assumptions</u>: only stochastic

Data:



Assumptions: only stochastic

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$$\tilde{d}_i(f) = \int_{-T/2}^{T/2} d_i(t) e^{-2\pi i f t} dt = \sum_{\nu} \tilde{n}_i^{\nu}(f) + \sum_{\sigma} \tilde{s}_i^{\sigma}(f)$$

~



TDI variables

 $\mathbf{A} = \frac{\mathbf{Z} - \mathbf{X}}{\sqrt{2}} \; ,$

$$\mathbf{E} = \frac{\mathbf{X} - 2\mathbf{Y} + \mathbf{Z}}{\sqrt{6}} \; ,$$

$$T = \frac{X + Y + Z}{\sqrt{3}}$$

Orthogonal channels

<u>Assumptions</u>: only stochastic

Data:

$$\tilde{d}_i(f) = \int_{-T/2}^{T/2} d_i(t) \, \mathrm{e}^{-2\pi i f t} \, \mathrm{d}t = \sum_{\nu} \tilde{n}_i^{\nu}(f) + \sum_{\sigma} \tilde{s}_i^{\sigma}(f)$$

(Implicit) assumption 2: equilateral geometry and equal noise in each spacecraft. (See Hartwig et al 2303.15929 for a paper relaxing these assumptions)

<u>Assumptions</u>: only stochastic

Data:

Properties:

$$\tilde{d}_i(f) = \int_{-T/2}^{T/2} d_i(t) \, \mathrm{e}^{-2\pi i f t} \, \mathrm{d}t = \sum_{\nu} \tilde{n}_i^{\nu}(f) + \sum_{\sigma} \tilde{s}_i^{\sigma}(f)$$

Assumption 3: noise and signals are

Gaussian and stationary.

The signal is also assumed to be isotropic

and non-chiral.

Assumptions: only stochastic, equilateral geometry

$$\tilde{d}_i(f) = \int_{-T/2}^{T/2} d_i(t) \, \mathrm{e}^{-2\pi i f t} \, \mathrm{d}t = \sum_{\nu} \tilde{n}_i^{\nu}(f) + \sum_{\sigma} \tilde{s}_i^{\sigma}(f)$$

Properties:

$$\langle \tilde{d}_i(f) \rangle = 0 \; ,$$

$$\langle \tilde{d}_i(f)\tilde{d}_j^*(f')\rangle = \frac{\delta(f-f')}{2} \left[\sum_{\nu} P_{N,ij}^{\nu}(f) + \sum_{\sigma} P_{S,ij}^{\sigma}(f)\right]$$

 $P_{N,ij}(f) \equiv \sum_{\nu} P_{N,ij}^{\nu}(f) = \left[T_{ij,lk}^{\text{TM}}(f) S_{lk}^{\text{TM}}(f) + T_{ij,lk}^{\text{OMS}}(f) S_{lk}^{\text{OMS}}(f) \right]$



 $P_{N,ij}(f) \equiv \sum_{\nu} P_{N,ij}^{\nu}(f) = \left[T_{ij,lk}^{\text{TM}}(f) S_{lk}^{\text{TM}}(f) + T_{ij,lk}^{\text{OMS}}(f) S_{lk}^{\text{OMS}}(f) \right]$

Test mass:
$$S_{lk}^{\text{TM}}(f) = A_{lk}^2 \left(1 + \left(\frac{0.4 \text{mHz}}{f}\right)^2\right) \left(1 + \left(\frac{f}{8 \text{mHz}}\right)^4\right) \left(\frac{1}{2\pi fc}\right)^2 \left(\frac{\text{fm}^2}{\text{s}^3}\right)$$

$$\begin{array}{ll} \underline{\textbf{Optical metrology}}\\ \underline{\textbf{system}} \end{array} & S_{lk}^{\text{OMS}}(f) = P_{lk}^2 \ \left(1 + \left(\frac{2 \times 10^{-3} \text{Hz}}{f}\right)^4\right) \times \left(\frac{2\pi f}{c}\right)^2 \ \times \left(\frac{\text{pm}^2}{\text{Hz}}\right)^2 \end{array} \right) \\ \end{array}$$

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$$\underline{\text{Test mass}}: \qquad S_{lk}^{\text{TM}}(f) = A_{lk}^2 \left(1 + \left(\frac{0.4 \text{mHz}}{f}\right)^2 \right) \left(1 + \left(\frac{f}{8\text{mHz}}\right)^4 \right) \left(\frac{1}{2\pi fc}\right)^2 \left(\frac{\text{fm}^2}{\text{s}^3}\right)$$

$$\begin{array}{ll} \underline{\textbf{Optical metrology}}\\ \underline{\textbf{system}} \end{array} \quad S_{lk}^{\text{OMS}}(f) = P_{lk}^2 \left(1 + \left(\frac{2 \times 10^{-3} \text{Hz}}{f}\right)^4 \right) \times \left(\frac{2\pi f}{c}\right)^2 \\ \times \left(\frac{\text{pm}^2}{\text{Hz}}\right) \end{array}$$

SIGNAL

$$h^2 \Omega_{\rm GW}(f) \equiv \frac{4\pi^2 f^3}{3H_0^2} S_h(f)$$

$$P_{S,ij}(f) \equiv \sum_{\sigma} P_{S,ij}^{\sigma}(f) = \mathcal{R}_{ij}(f) \left[S_{\text{Gal}}(f) + S_{\text{Ext}}(f) + S_{\text{Cosmo}}(f) \right]$$

Galactic:
$$S_{\text{Gal}}(f) = A_{\text{Gal}} \left(\frac{f}{1 \text{ Hz}}\right)^{-\frac{7}{3}} \times e^{-(f/f_1)^{\alpha}} \times \frac{1}{2} \left[1 + \tanh \frac{f_{\text{knee}} - f}{f_2}\right]$$

Extragalactic:
$$h^2 \Omega_{\text{Ext}} = 10^{\log_{10}(h^2 \Omega_{\text{Ext}})} \left(\frac{f}{0.001 \text{Hz}}\right)^{2/3}$$

SIGNAL AND NOISE IN THE LISA BAND



Divide full data stream into N_d segments, each lasting $T_{\rm obs}/N_d$. $T_{\rm obs}/N_d \sim 11.5 {\rm ~days} \mapsto \Delta f \sim 10^{-6} {\rm ~Hz}$

<u>Assumption 4</u>: $T_{obs} = 4$ years

Divide full data stream into N_d segments, each lasting T_{obs}/N_d .

Generate N_d Gaussian realizations of noise and data.

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Average over realizations $\bar{D}_{ij}^{k} \equiv \sum_{s=1}^{N_d} \tilde{d}_i^s(f_k) \tilde{d}_j^s(f_k) / N_d$

Divide full data stream into N_d segments, each lasting T_{obs}/N_d .

Generate N_d Gaussian realizations of noise and data.

Average over realizations
$$ar{D}_{ij}^{
m k}\equiv\sum_{s=1}^{N_d} ilde{d}_i^s(f_{
m k}) ilde{d}_j^s(f_{
m k})/N_d$$

Down sample data by coarse-graning $\bar{D}_{ii}^k \mapsto D_{ii}^k$

$$\ln \mathcal{L}(\vec{\theta}) = \frac{1}{3} \ln \mathcal{L}_{G}(\vec{\theta}|D_{ij}^{k}) + \frac{2}{3} \ln \mathcal{L}_{LN}(\vec{\theta}|D_{ij}^{k})$$

$$\ln \mathcal{L}_{G}(\vec{\theta}|D_{ij}^{k}) = -\frac{N_{d}}{2} \sum_{k} \sum_{i,j} w_{ij}^{k} \left[1 - D_{ij}^{k}/D_{ij}^{Th}(f_{ij}^{k},\vec{\theta})\right]^{2}$$
$$\ln \mathcal{L}_{LN}(\vec{\theta}|D_{ij}^{k}) = -\frac{N_{d}}{2} \sum_{k} \sum_{i,j} w_{ij}^{k} \ln^{2} \left[D_{ij}^{Th}(f_{ij}^{k},\vec{\theta})/D_{ij}^{k}\right]$$

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 $\vec{\theta} = \{A, P, \Omega_{\text{gal}}, \Omega_{\text{ext}}, \vec{\theta}_{\text{Inf}}\}$

$$\begin{split} &\ln \mathcal{L}(\vec{\theta}) = \frac{1}{3} \ln \mathcal{L}_{\mathrm{G}}(\vec{\theta}|D_{ij}^{k}) + \frac{2}{3} \ln \mathcal{L}_{\mathrm{LN}}(\vec{\theta}|D_{ij}^{k}) \\ &\ln \mathcal{L}_{\mathrm{G}}(\vec{\theta}|D_{ij}^{k}) = -\frac{N_{d}}{2} \sum_{k} \sum_{i,j} w_{ij}^{k} \left[1 - D_{ij}^{k} / D_{ij}^{\mathrm{Th}}(f_{ij}^{k}, \vec{\theta}) \right]^{2} \\ &\ln \mathcal{L}_{\mathrm{L}N}(\vec{\theta}|D_{ij}^{k}) = -\frac{N_{d}}{2} \sum_{k} \sum_{i,j} w_{ij}^{k} \ln^{2} \left[D_{ij}^{\mathrm{Th}}(f_{ij}^{k}, \vec{\theta}) / D_{ij}^{k} \right] \\ &\vec{\theta} = \{A, P, \Omega_{\mathrm{gal}}, \Omega_{\mathrm{ext}}, \vec{\theta}_{\mathrm{Inf}} \} \end{split}$$

and foregrounds

TEMPLATE BANK FOR INFLATIONARY MODELS & FORECASTS

TEMPLATE BANK FOR INFLATIONARY MODELS & FORECASTS

- 1. POWER LAW
- 2. LOGNORMAL BUMP
- **3. BROKEN POWER LAW**
- 4. DOUBLE PEAK
- **5. EXCITED STATES**
- 6. LINEAR OSCILLATIONS
- 7. LOGARITHMIC OSCILLATIONS

POWER LAW

TEMPLATE DEFINITION

$$h^2 \Omega_{\text{GW}}^{\text{PL}}(f, \vec{p}) = 10^{\alpha_*} \left(\frac{f}{f_*}\right)^{n_t} \qquad \vec{p} = \{\alpha_*, f_*, n_t\}$$



INFLATIONARY MODELS PRODUCING THE SIGNAL

Axion inflation:

Barnaby & Peloso 1011.1500, Sorbo 1101.1525

The inflaton is an axion coupled to a gauge field through an axial interaction. The rolling axion strongly amplifies the gauge field, which in turn produces a strong SGWB.

Broken space diffeomorphisms:

Ricciardone & Tasinato 1611.04516, 1711.02635, Fujita et al 1808.02381

The breaking of space diffeomorphisms can give rise to a massive graviton during the inflationary epoch which tilts the SGWB spectrum towards the blue

$$h^2 \Omega_* \simeq 1.5 \times 10^{-13} \frac{H_*^4}{M_{\rm Pl}^4} \frac{e^{4\pi\xi_*}}{\xi_*^6}$$

 $n_t \simeq -4\epsilon_* + (4\pi\xi_* - 6)(\epsilon_* - \eta_*)$

$$h^2\Omega_* =$$
 Model dependent

$$n_t \simeq \frac{2}{3} \frac{m_h^2}{H_*^2} > 0$$

FORECASTS (FISHER)



FORECASTS (NESTED SAMPLING)






LOGNORMAL BUMP

TEMPLATE DEFINITION

$$h^2 \Omega_{\rm GW}^{\rm LBp}(f,\vec{p}) = 10^{\alpha_*} \exp\left[-\frac{1}{2\rho^2} \log_{10}^2 \left(\frac{f}{f_*}\right)\right] \quad \vec{p} = \{\alpha_*, f_*, \rho\}$$



INFLATIONARY MODELS PRODUCING THE SIGNAL

Axion spectator:

Namba et al 1509.07521

The axion coupled to a gauge field is now a spectator field rolling only for a short time ΔN .

$$h^2 \Omega_* \simeq 1.5 \times 10^{-13} \frac{H_*^4}{M_{\rm Pl}^4} \frac{\mathrm{e}^{4\pi\xi_*}}{\xi_*^6}$$

 $\rho \propto \Delta N$
 $f_* \propto a(N_*)H(N_*)$

FORECASTS (NESTED SAMPLING)



INTERMEZZO: SCALAR-INDUCED Gravitational waves

Large scalar perturbations source gravitational waves at 2nd order in perturbation theory when they re-enter the horizon during radiation era

$$\Omega_{\text{ind}}(k) = 0.387 \,\Omega_{\text{R}}\left(\frac{g_{*,s}^{4}g_{*}^{-3}}{106.75}\right)^{-\frac{1}{3}} \frac{1}{6} \int_{-1}^{1} dx \int_{1}^{\infty} dy \,\mathscr{P}\left(\frac{y-x}{2}k\right) \,\mathscr{P}\left(\frac{x+y}{2}k\right) F(x,y)$$

see Robert's talk and Domenech 2109.01398 for a review

SIGW FROM A LOGNORMAL $\mathcal{P}(k)$

Bumps in $\mathscr{P}(k)$ are often modeled using a Lognormal template $\mathscr{P}(k) = \frac{A_{\zeta}}{\sqrt{2\pi}\Delta} \exp\left[-\frac{\ln^2 k/k_*}{2\Delta^2}\right] \qquad 10^{-1}$



See Pi & Sasaki 2005.12306 for analytical solutions for $\Omega_{
m GW}(f)$

BROKEN POWER LAW

TEMPLATE DEFINITION

$$h^{2}\Omega_{\rm GW}^{\rm BPL}(f,\vec{p}) = 10^{\alpha_{*}} \frac{\left(\frac{f}{f_{*}}\right)^{n_{t,1}}}{\left\{\frac{1}{2}\left[1 + \left(\frac{f}{f_{*}}\right)^{1/\delta}\right]\right\}^{(n_{t,1}-n_{t,2})\delta}} \quad \vec{p} = \{\alpha_{*}, f_{*}, n_{t,1}, n_{t,2}, \delta\}$$



INFLATIONARY MODELS PRODUCING THE SIGNAL. MODEL 1

Second slow-roll stage:

Franciolini & Urbano 2207.10056

An ultra-slow-roll phase amplifies primordial scalar perturbations and is followed by a second slow-roll regime generating a plateau.







Amplitude: ratio H^2/ϵ during the second SR stage, abundance of Primordial Black Holes $f_{\rm PBH}$.



Amplitude: ratio H^2/ϵ during the second SR stage, abundance of Primordial Black Holes f_{PBH} . Frequency of the turn: time of the onset of the second SR stage, mass of Primordial Black Holes M/M_{\odot}



Amplitude: ratio H^2/ϵ during the second SR stage, abundance of Primordial Black Holes f_{PBH} . Frequency of the turn: time of the onset of the second SR stage, mass of Primordial Black Holes M/M_{\odot} . IR spectral index: related to the scalar IR spectral index.



<u>Amplitude</u>: ratio H^2/ϵ during the second SR stage, abundance of Primordial Black Holes f_{PBH} . Frequency of the turn: time of the onset of the second SR stage, mass of Primordial Black Holes M/M_{\odot} . **IR spectral index:** related to the scalar IR spectral index. **<u>UV spectral index:</u>** predicted to be

flat in this model



<u>Amplitude</u>: ratio H^2/ϵ during the second SR stage, abundance of Primordial Black Holes f_{PBH} . Frequency of the turn: time of the onset of the second SR stage, mass of Primordial Black Holes M/M_{\odot} . **IR spectral index:** related to the scalar IR spectral index. **<u>UV spectral index:</u>** predicted to be flat in this model. **Smoothing parameter:** related to the

sharpness of the USR-SR transition.

INFLATIONARY MODELS PRODUCING THE SIGNAL. MODEL 2

Hybrid inflation:

Braglia, Linde, Kallosh, Finelli 2211.14262

Primordial scalar perturbations are sourced by tachyonically amplified isocurvature perturbations during a waterfall phase producing a very broad bump.









The signal is **well reconstructed**, but affected by **strong degeneracies**. The parameters are not well constrained.

Suboptimal parameterization.

LISA cannot measure

moderately **asymmetries** in very broad bumps, due to its limited bandwidth.

DOUBLE PEAK

$$TEMPLATE DEFINITION \quad \vec{p} = \{\alpha_*, f_*, \beta, \kappa_1, \kappa_2, \rho, \gamma\}$$

$$h^2 \Omega_{\rm GW}^{\rm DP}(f, \vec{p}) = 10^{\alpha_*} \left[\beta \left(\frac{f}{\kappa_1 f_*} \right)^{n_p} \left[\frac{c_1 - f/f_*}{c_1 - \kappa_1} \right]^{\frac{n_p}{\kappa_1}(c_1 - \kappa_1)} \Theta\left(c_1 - \frac{f}{f_*}\right) + \exp\left[-\frac{1}{2\rho^2} \log_{10}^2 \left(\frac{f}{\kappa_2 f_*} \right) \right] \left\{ 1 + \operatorname{erf}\left[-\gamma \log_{10} \left(\frac{f}{\kappa_2 f_*} \right) \right] \right\} \right]$$



TEMPLATE DEFINITION
$$\vec{p} = \{\alpha_*, f_*, \beta, \kappa_1, \kappa_2, \rho, \gamma\}$$

 $h^2 \Omega_{\rm GW}^{\rm DP}(f, \vec{p}) = 10^{\alpha_*} \left[\beta \left(\frac{f}{\kappa_1 f_*} \right)^{n_p} \left[\frac{c_1 - f/f_*}{c_1 - \kappa_1} \right]^{\frac{n_p}{\kappa_1}(c_1 - \kappa_1)} \Theta\left(c_1 - \frac{f}{f_*}\right) + \exp\left[-\frac{1}{2\rho^2} \log_{10}^2 \left(\frac{f}{\kappa_2 f_*} \right) \right] \left\{ 1 + \operatorname{erf} \left[-\gamma \log_{10} \left(\frac{f}{\kappa_2 f_*} \right) \right] \right\} \right]$



INFLATIONARY MODELS PRODUCING THE SIGNAL

Broken power law $\mathcal{P}_{\underline{\zeta}}(k)$:

$$\mathcal{P}_{\zeta}^{\text{bpl}}(k) = \frac{\mathcal{A}_s(p_1 + p_2)}{\left[p_2\left(\frac{k}{k_*}\right)^{-p_1} + p_1\left(\frac{k}{k_*}\right)^{p_2}\right]}$$

Lognormal $\mathscr{P}_{\zeta}(k)$:

$$\mathcal{P}_{\zeta}^{\ln}(k) = \mathcal{A}_s \exp\left[-\frac{1}{2\Delta^2}\ln^2\left(\frac{k}{k_*}\right)\right]$$

See Ozsoy & Tasinato 2301.03600 for a review of models producing a peak in the scalar power spectrum

Prior on the 7 model parameters chosen so as to reproduce the GW background from these models.

More efficient approach: start from the parameters for $\mathscr{P}_{\zeta}(k)$ instead. See Robert's talk.

FORECAST. BNK 1



Very loud signal:

tight constraints on all

parameters.

FORECAST. BNK 2



Moderately loud signal:

some parameters are not constrained, but signal is reconstructed quite well.

FORECAST. BNK 3



Faint signal: only certain features of the signal are constrained.

EXCITED STATES

TEMPLATE DEFINITION

$$h^2 \Omega_{\rm GW}^{\rm ES}(f,\vec{p}) = \frac{10^{\alpha_*}}{0.052} \frac{1}{x^3} \left[1 - \frac{x^2}{4\gamma_{\rm ES}^2} \right]^2 \left[\sin(x) - 2\frac{1 - \cos(x)}{x} \right]^2 \Theta(x_{\rm cut} - x)$$

$$x \equiv (f \omega_{\rm ES})/2$$
 $\vec{p} = \{\alpha_*, \gamma_{\rm ES}, \omega_{\rm ES}\}$



INFLATIONARY MODELS PRODUCING THE SIGNAL

Scalar- induced GWs during

inflation:

Inomata 2109.03972, Fumagalli et al 2111.14664

A sharp feature induces a temporary large amplification of sub horizon scalar modes. GWs are amplified by loop effects.



<section-header>FORECASTS



ES-BNK_3

-9.5

LINEAR OSCILLATIONS

TEMPLATE DEFINITION

 $h^2 \Omega_{\rm GW}^{\rm LO}(f,\vec{p}) = \left[1 + \mathcal{A}_{\rm lin} \cos\left(\omega_{\rm lin} f + \theta_{\rm lin}\right)\right] h^2 \Omega_{\rm GW}^{\rm env}(f,\vec{p}_{\rm env})$

 $\vec{p} = \{ \vec{p}_{env}, \mathcal{A}_{lin}, \omega_{lin}, \theta_{lin} \}$



INFLATIONARY MODELS PRODUCING THE SIGNAL. MODEL 1

Sharp features during inflation:

Fumagalli et al 2012.02861, Braglia, Chen, Hazra 2012.05821

Scalar induced signal from sharp features during inflation produced at horizon re-entry.





FORECASTS



Because of their particular spectral shape oscillations can be <u>well constrained</u> by LISA.

RESONANT (LOGARITHMIC) OSCILLATIONS

TEMPLATE DEFINITION

 $h^{2}\Omega_{\rm GW}^{\rm RO}(f,\vec{p}) = \left\{ 1 + \mathcal{A}_{1}(A_{\rm log},\omega_{\rm log})\cos\left[\omega_{\rm log}\ln(f/{\rm Hz}) + \theta_{\rm log,1}\right] + \mathcal{A}_{2}(A_{\rm log},\omega_{\rm log})\cos\left[2\omega_{\rm log}\ln(f/{\rm Hz}) + \theta_{\rm log,2}\right] \right\} h^{2}\Omega_{\rm GW}^{\rm env}(f,\vec{p}_{\rm env})$ $\vec{p}_{\rm RO} = \left\{\alpha_{*},A_{\rm log},\omega_{\rm log},\phi_{\rm log}\right\}$



INFLATIONARY MODELS PRODUCING THE SIGNAL. MODEL 1

<u>Resonant features during inflation:</u> Fumagalli et al 2105.06481

Currently a theoretical proposal based on analogies with resonant mechanisms in axion monodromy or primordial standard clock models. Models going in this direction include:

Battacharya & Zavala 2205.06065, Mavromatos et al 2206.07963, Calcagni & Kuroyanagi 2308.05904



FORECASTS


SUMMARY

- We initiated the collection of a template bank for $\Omega_{Gw}(f)$ based on motivated inflationary scenarios.
- Within the LISA CosWG ongoing efforts for Cosmic strings and phase transitions
- Our results showcase the potential of LISA to constrain Inflation based on the reconstruction of the spectral shape of the GW background
- Other observables or correlation with other experiments may be needed to underpin the inflationary origin of the background if the spectral shape is degenerate with other models