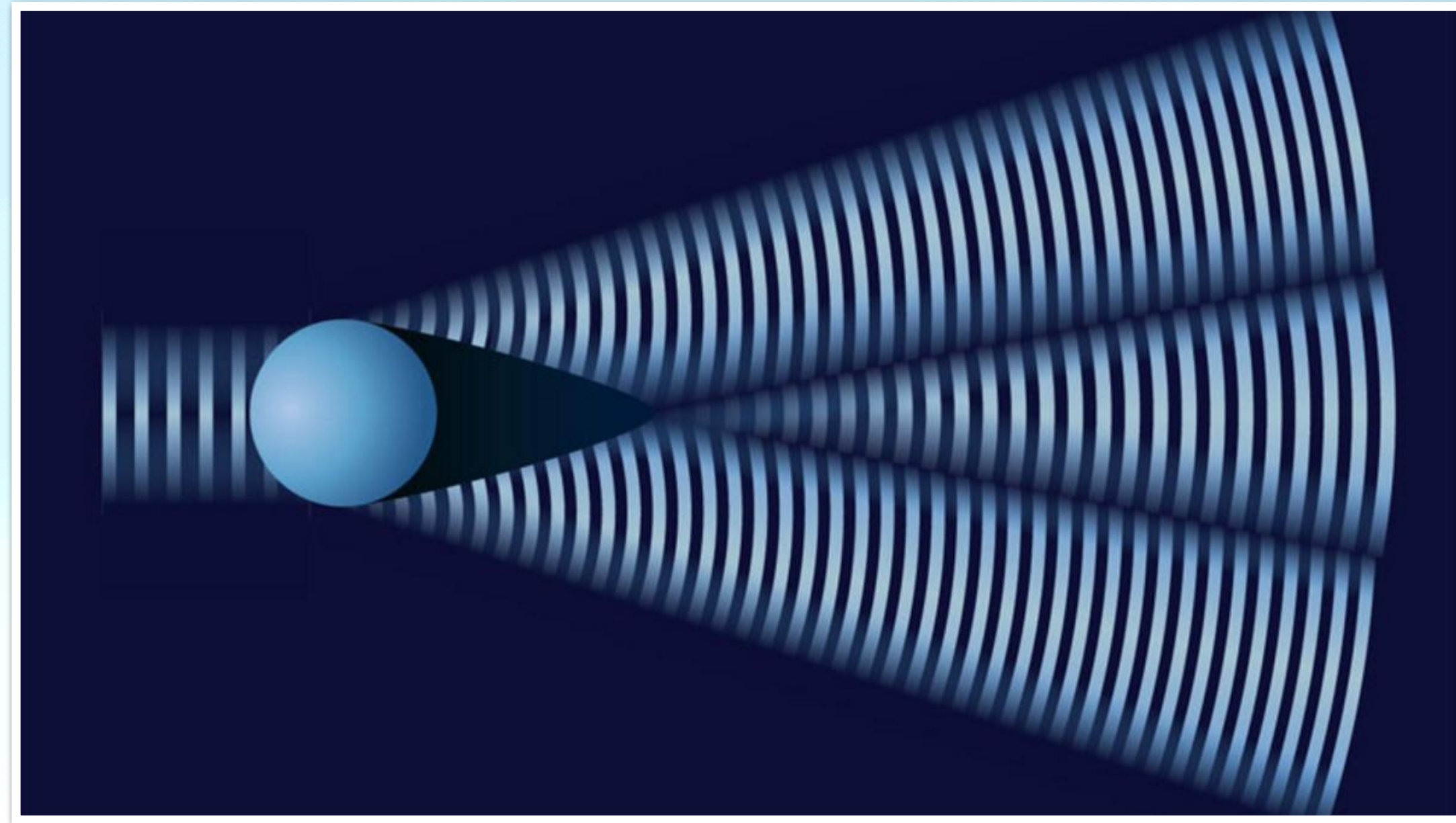


Proper time path integrals for GWs

An improved wave optics framework

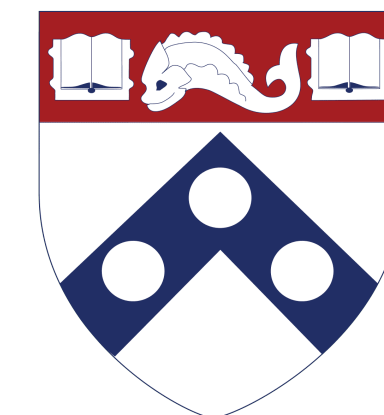


Alice Garoffolo

aligar@sas.upenn.edu

Based on: 2405.20208

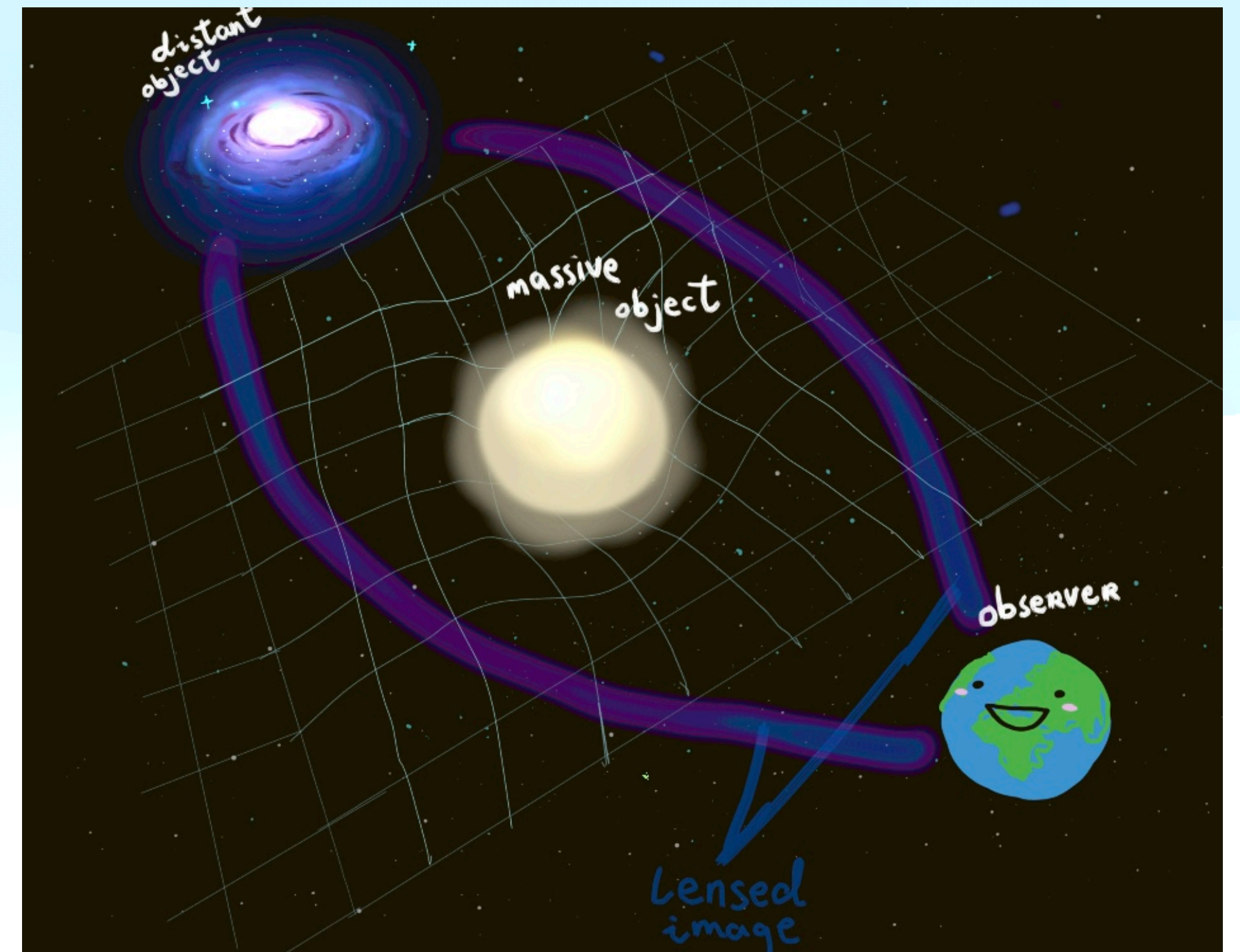
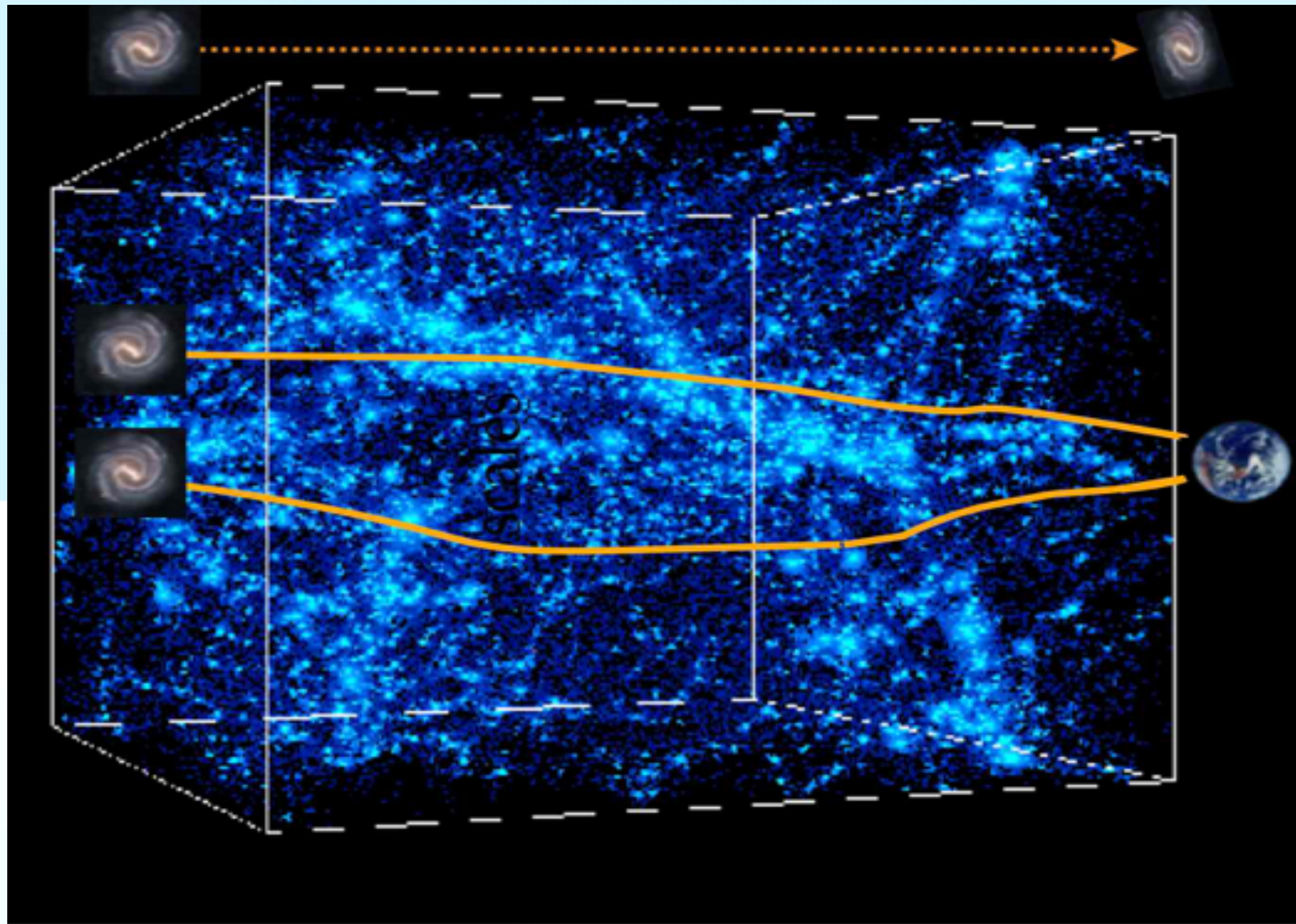
G. Braga, AG, A. Ricciardone, N. Bartolo, S. Matarrese



Penn
UNIVERSITY of PENNSYLVANIA

GWs through the perturbed Universe

Probe of large scale structures and compact objects

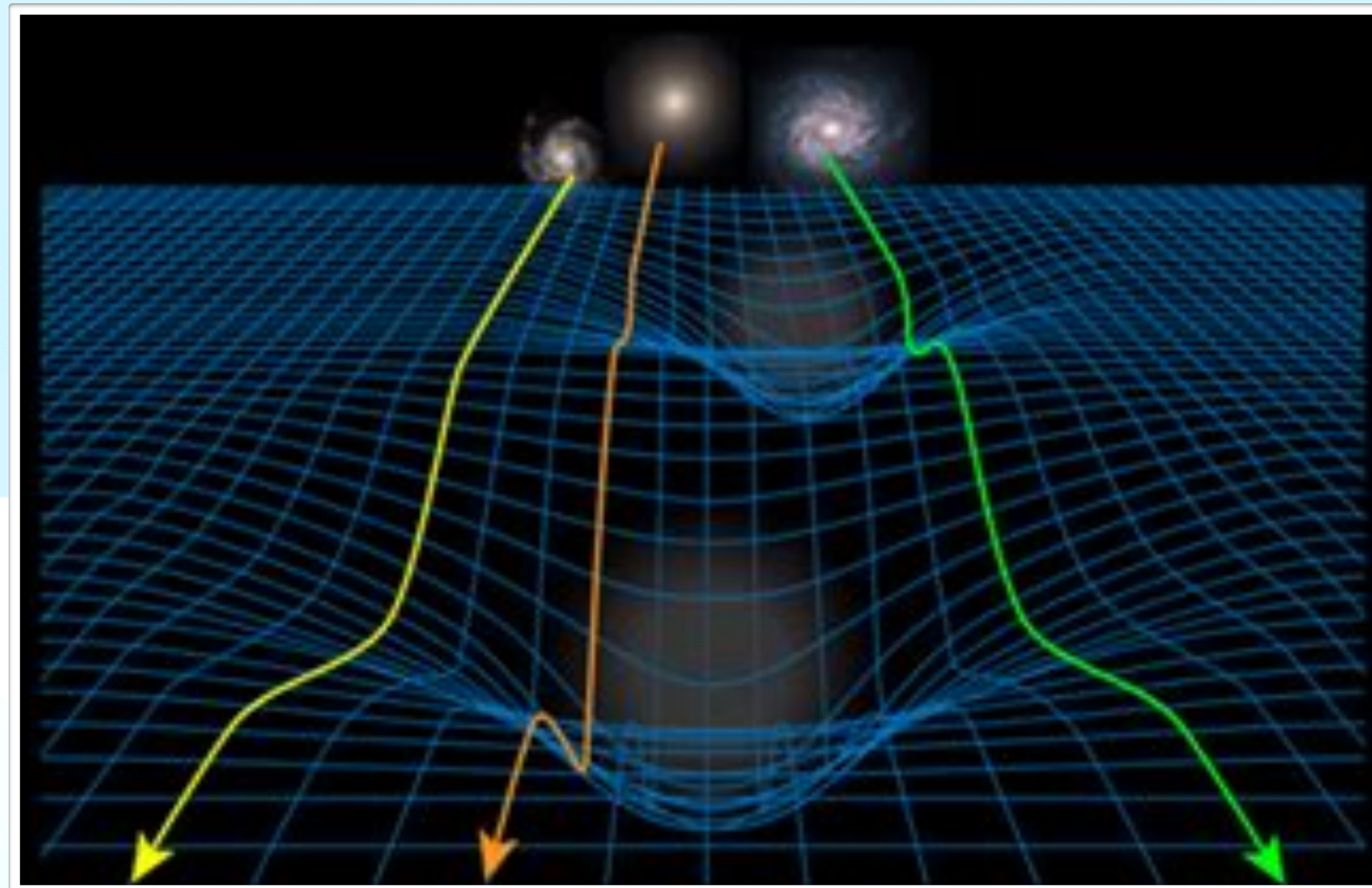


Propagation effects carry cosmological and astrophysical information

Optical regimes

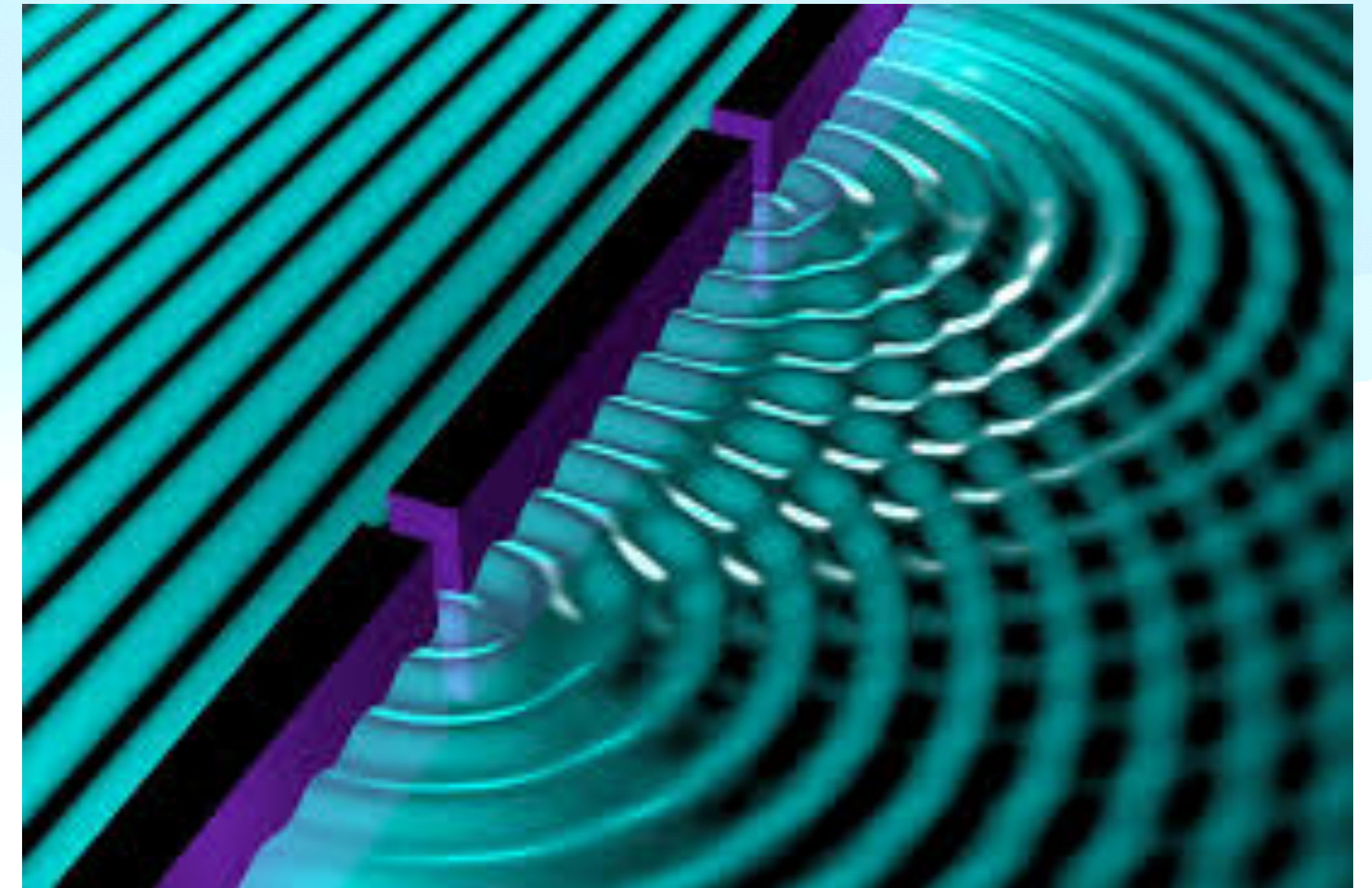
Geometric vs Wave optics

High Frequency: $\omega R_S \gg 1$



Ray description

Low Frequency: $\omega R_S \lesssim 1$



Wave effects

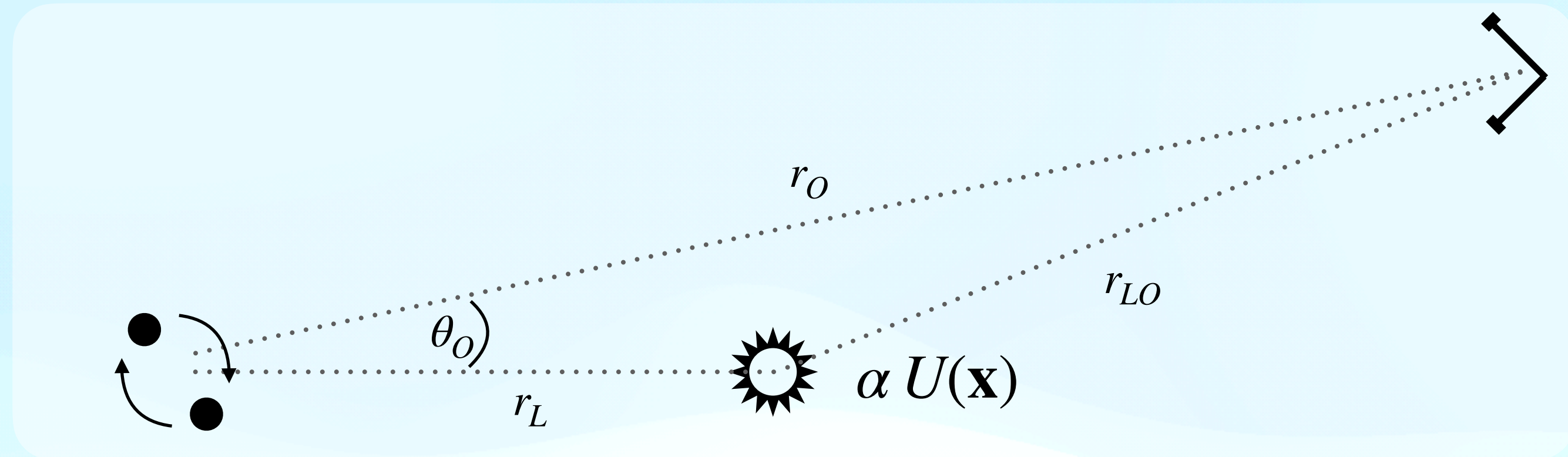
3.4 Wave effects

LISA CosGW: 2204.05434

GWs can be emitted at low frequencies ($\omega \lesssim 1$), allowing the observation of wave diffractive phenomena. For typical LISA sources, wave optics as in Eq. (15) needs to be considered for lenses with masses $M_L \sim 10^6 - 10^9 M_\odot$, cf. Eq. (13).

Diffraction integral for a scalar wave

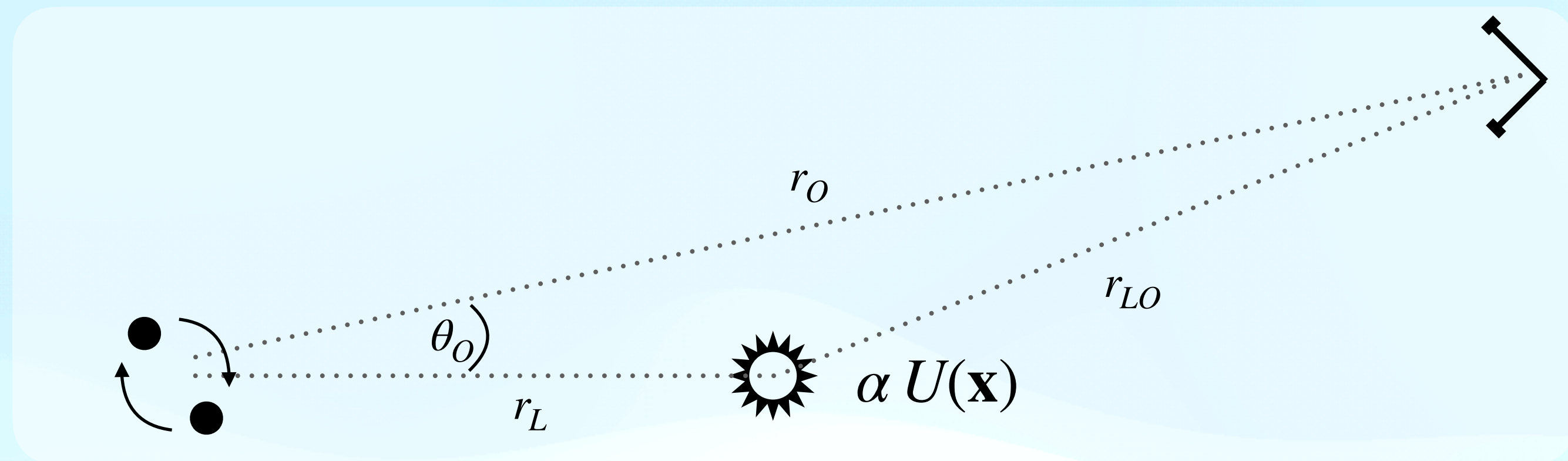
Nakamura&Deguchi 1999



Diffraction integral for a scalar wave

Nakamura&Deguchi 1999

1. Klein-Gordon Eq.: $[\nabla^2 + \omega^2(1 - 4\alpha U)] \tilde{\Psi}_\omega(\mathbf{x}) = 0$

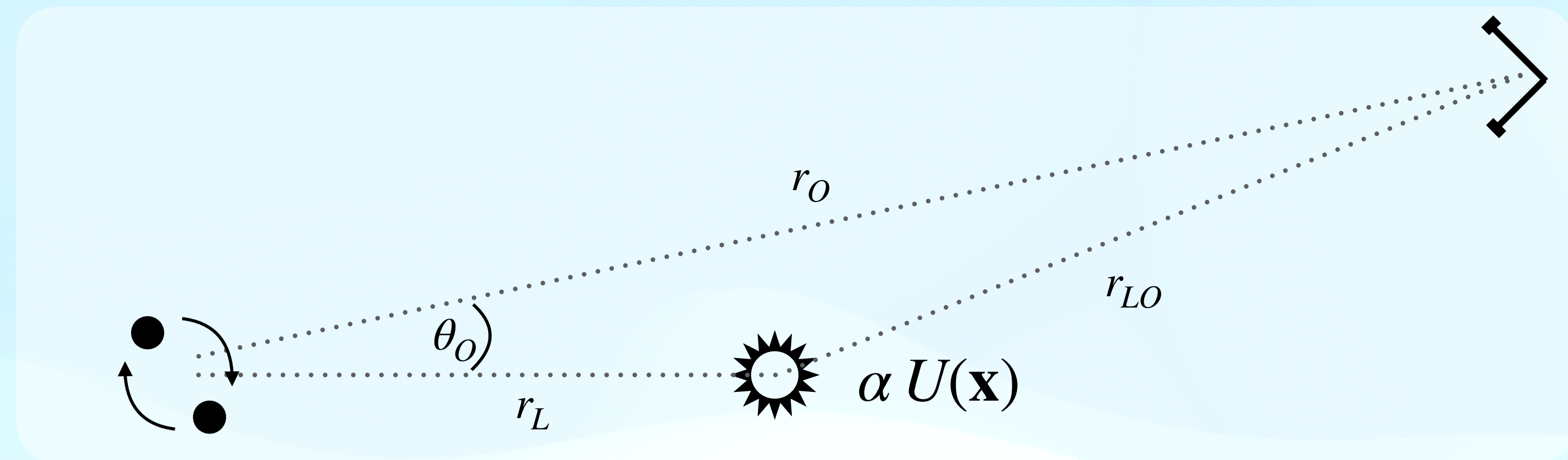


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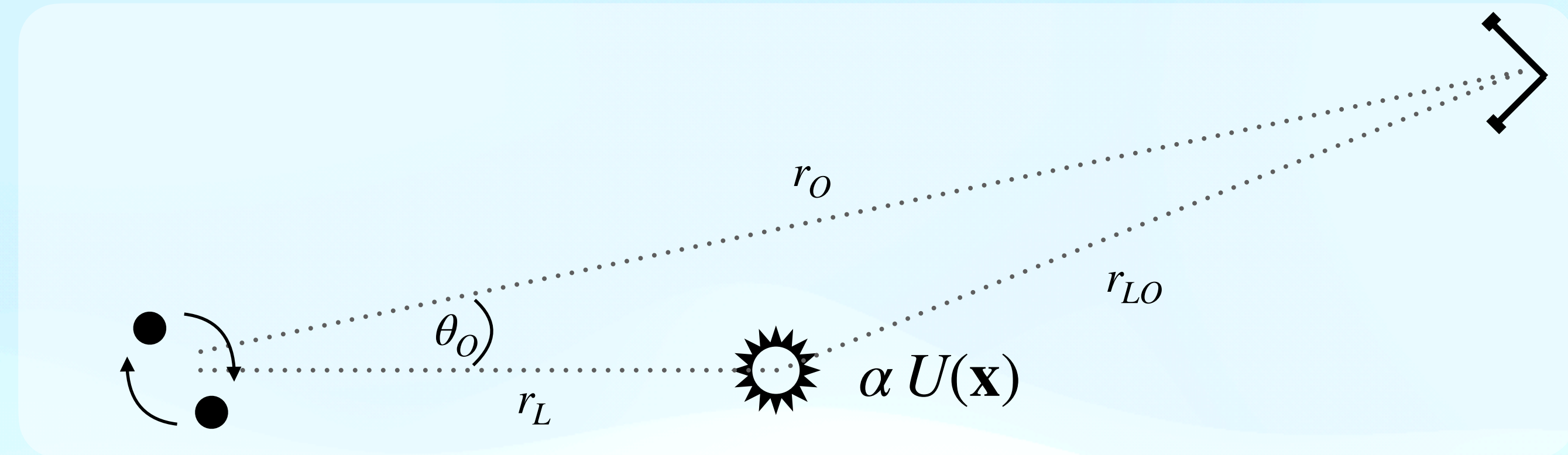
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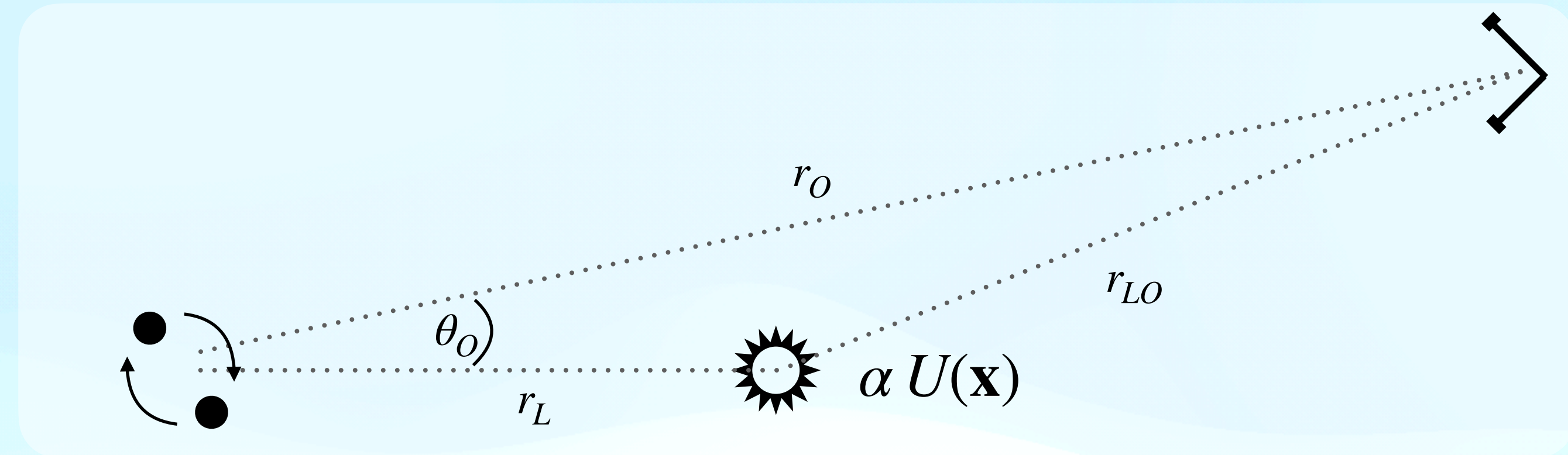
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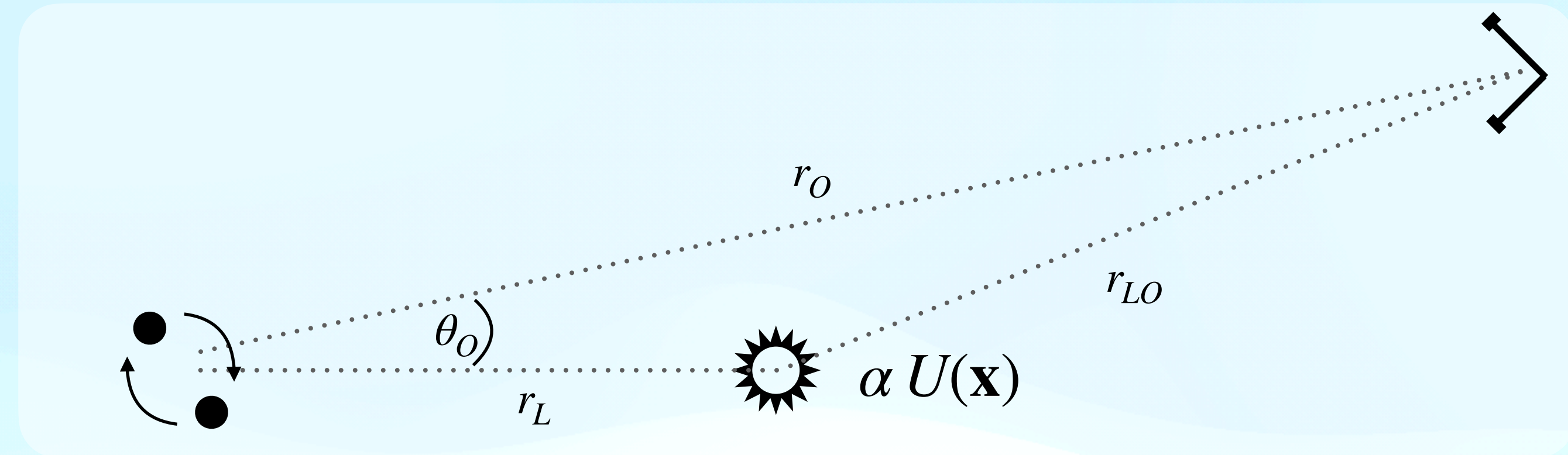
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Diffraction integral:

$$F(\vec{r}_O) = \int \mathcal{D}\theta(r) \exp \left\{ i\omega \int_0^{r_O} dr \left[\frac{r^2}{2} |\dot{\theta}|^2 - 2\alpha U(r, \theta) \right] \right\}$$

Diffraction integral for a scalar wave

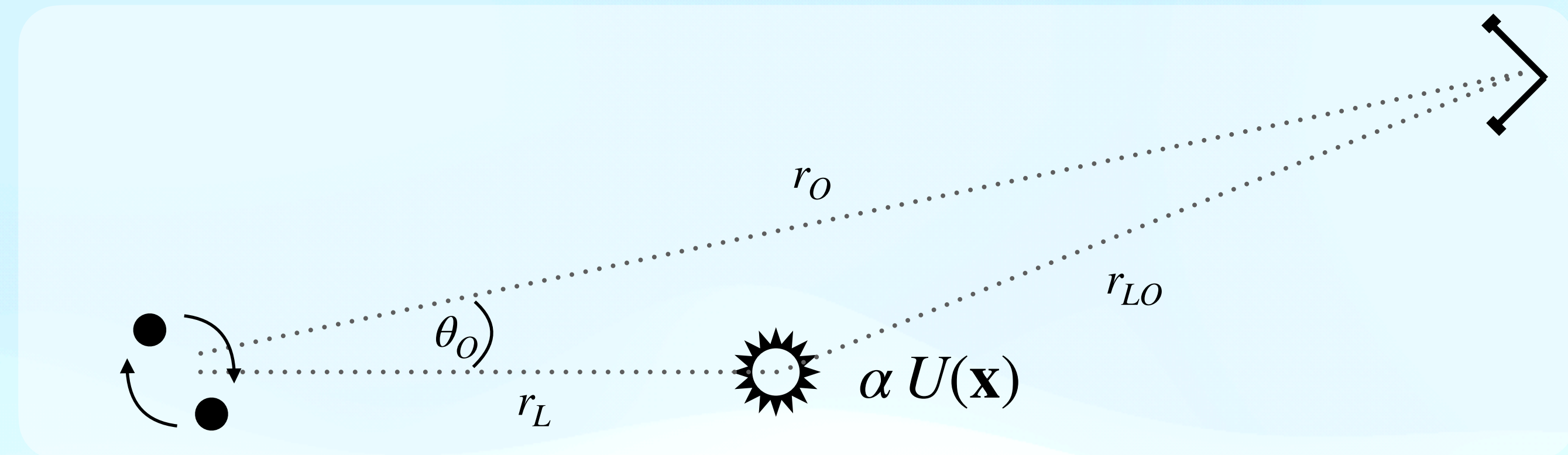
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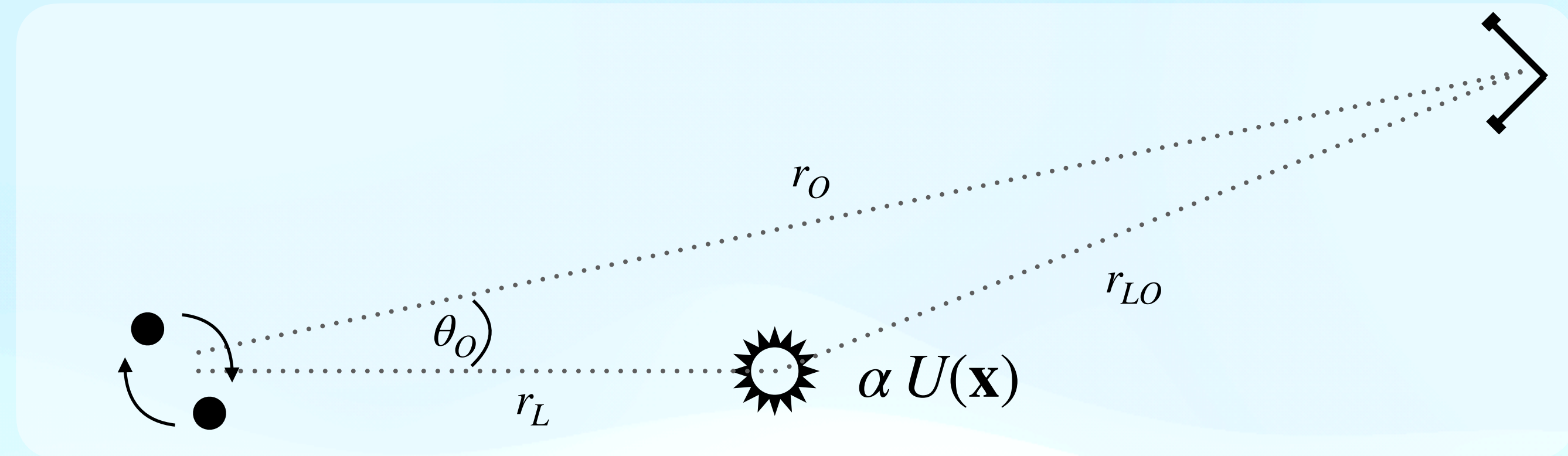
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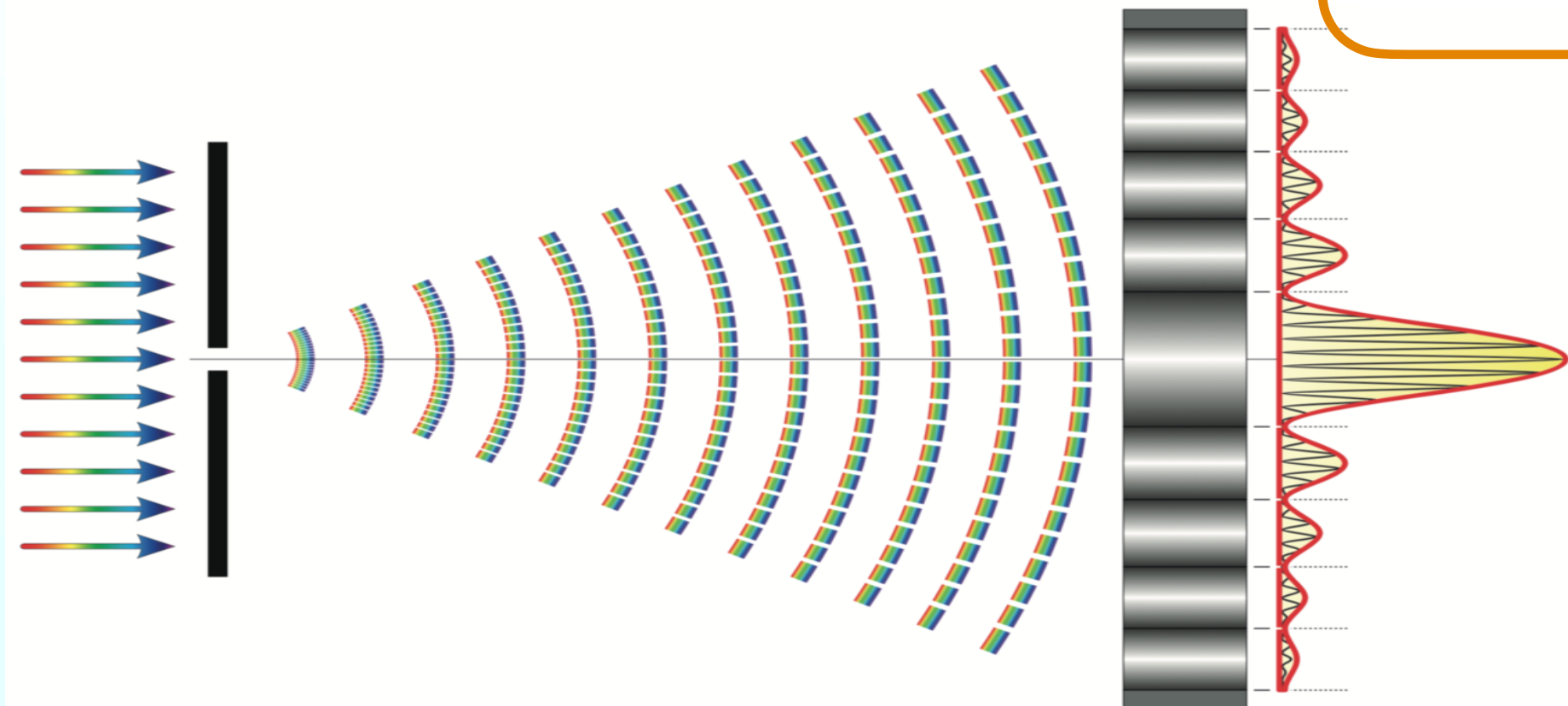
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Analogy between wave and quantum effects:
interference between all paths.
Geometric optics = classical limit

Diffraction integral: Pros and Cons

PROs

CONs

Diffraction integral: Pros and Cons

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1. Wave optics effects are frequency dependent

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Diffraction integral: Pros and Cons

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Citation per year of
Nakamura&Deguchi 1999



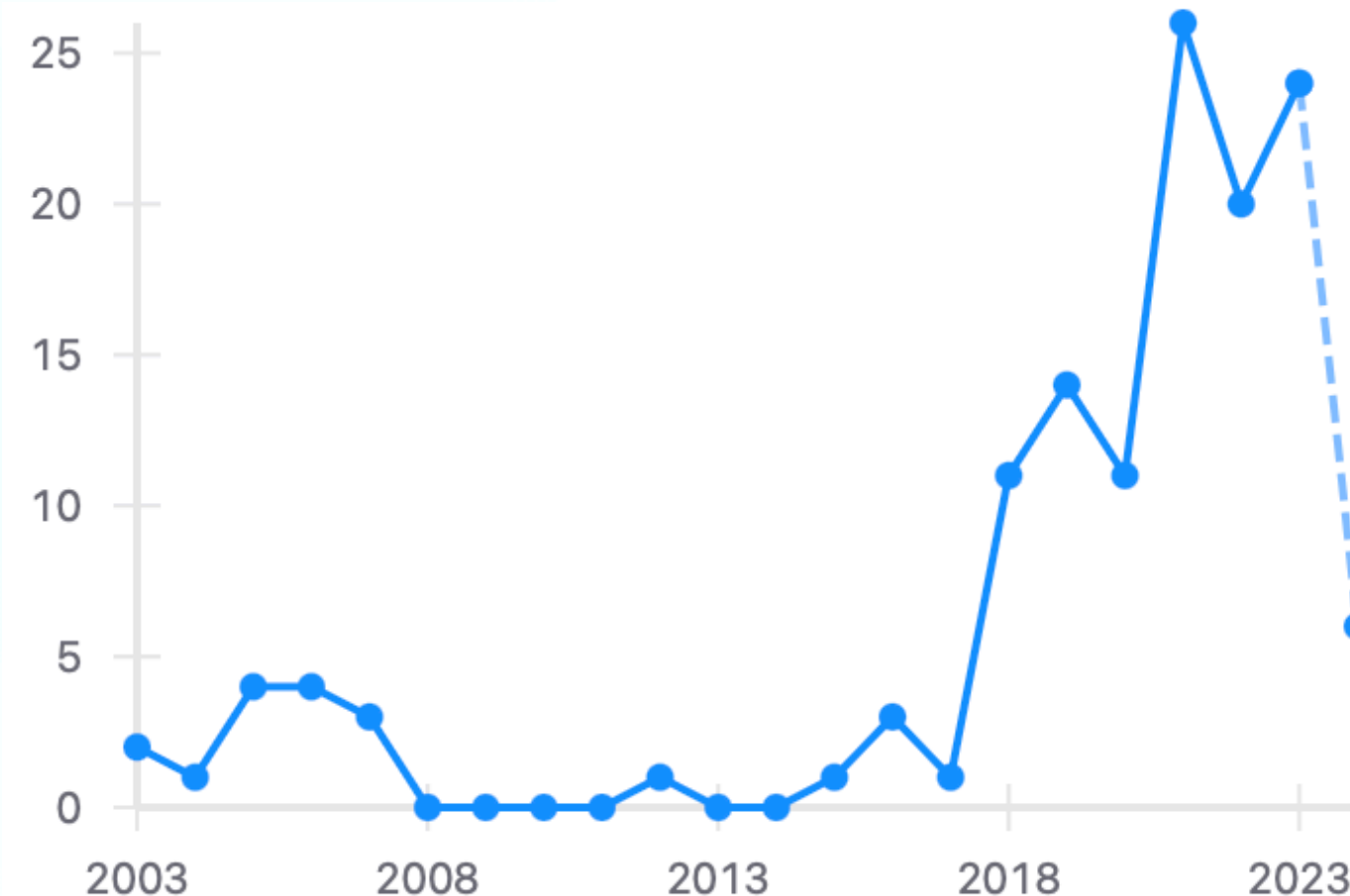
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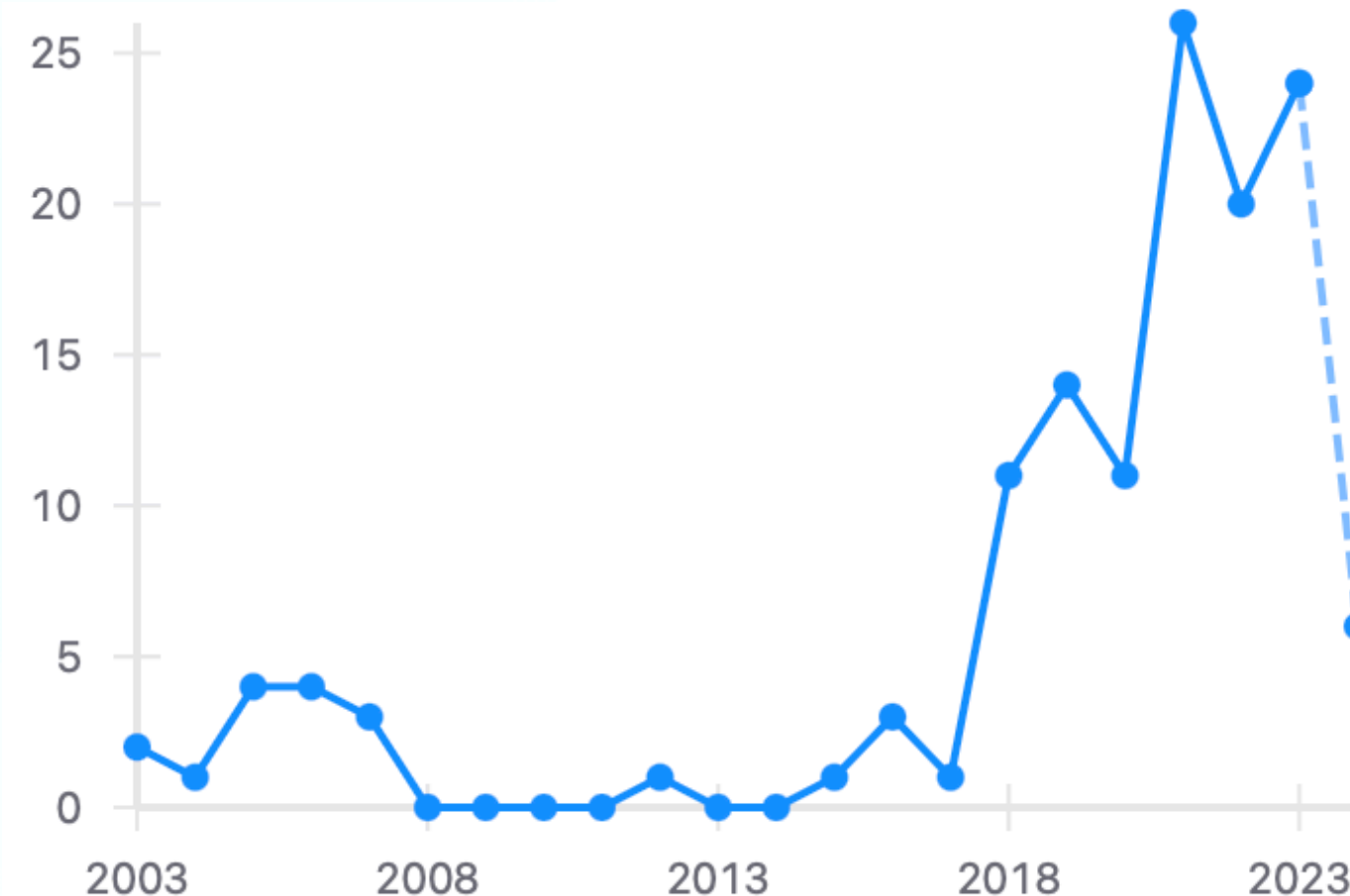
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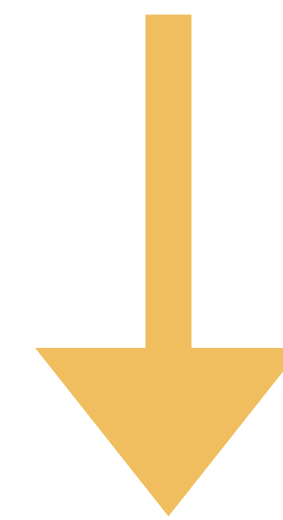
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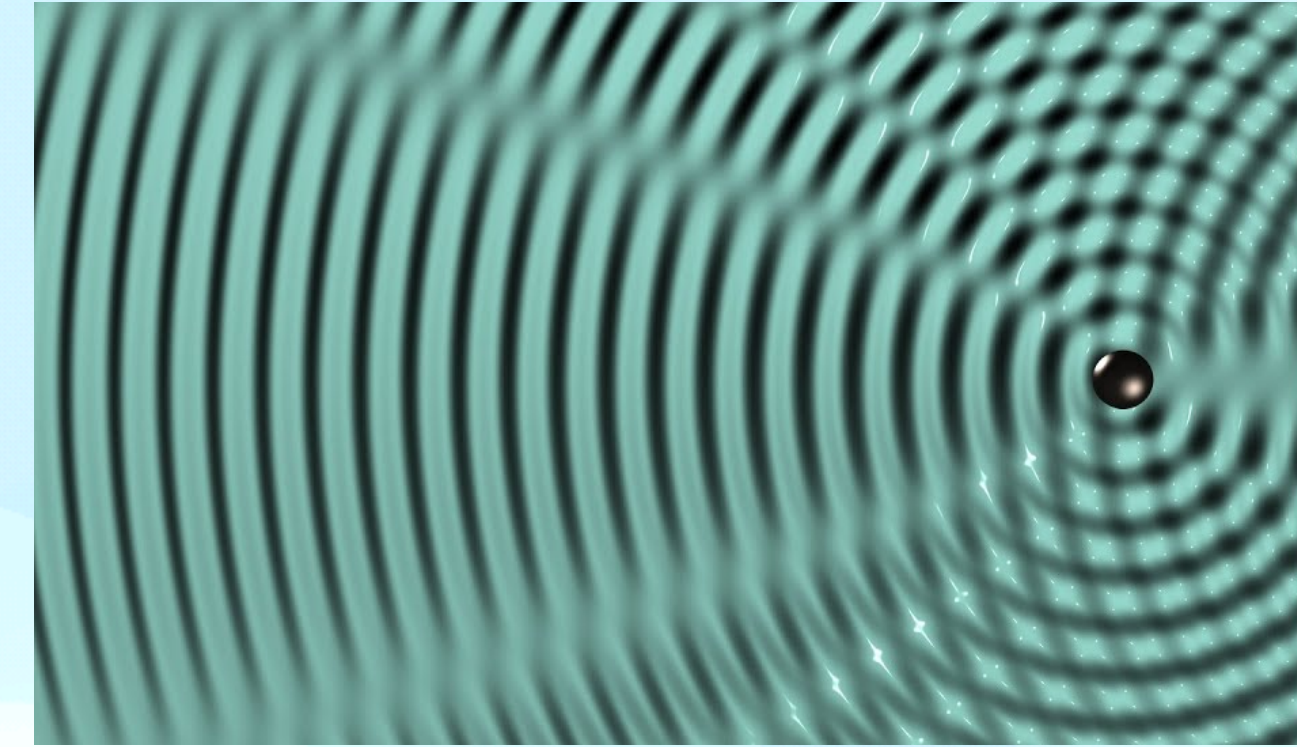


Proper time path integrals

2405.20208

Proper time path integral

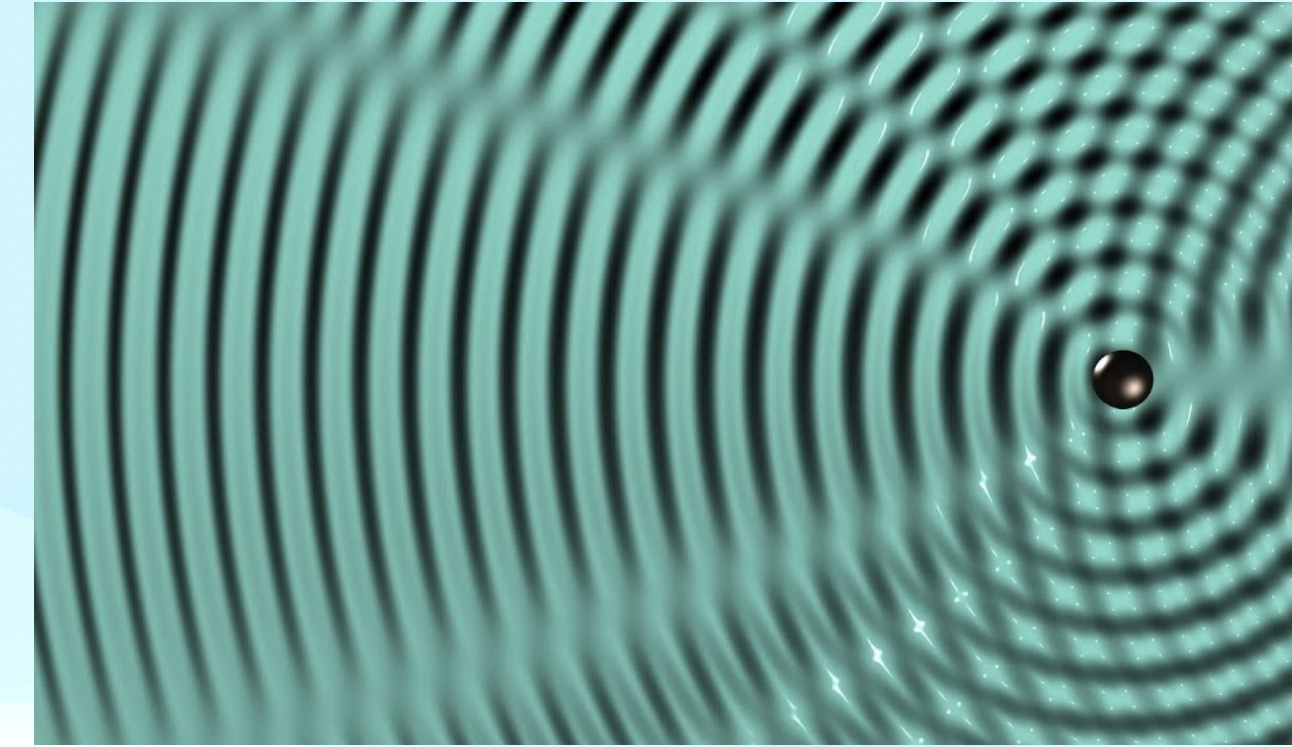
Finding a Schrödinger equation without Eikonal approximation



Proper time path integral

Finding a Schrödinger equation without Eikonal approximation

1. Proper time:
$$\tilde{\Psi}_\omega(\mathbf{x}) = -\frac{i}{\omega} \int_0^{+\infty} d\tau e^{i\omega\tau} \psi_\omega(\tau, \mathbf{x})$$

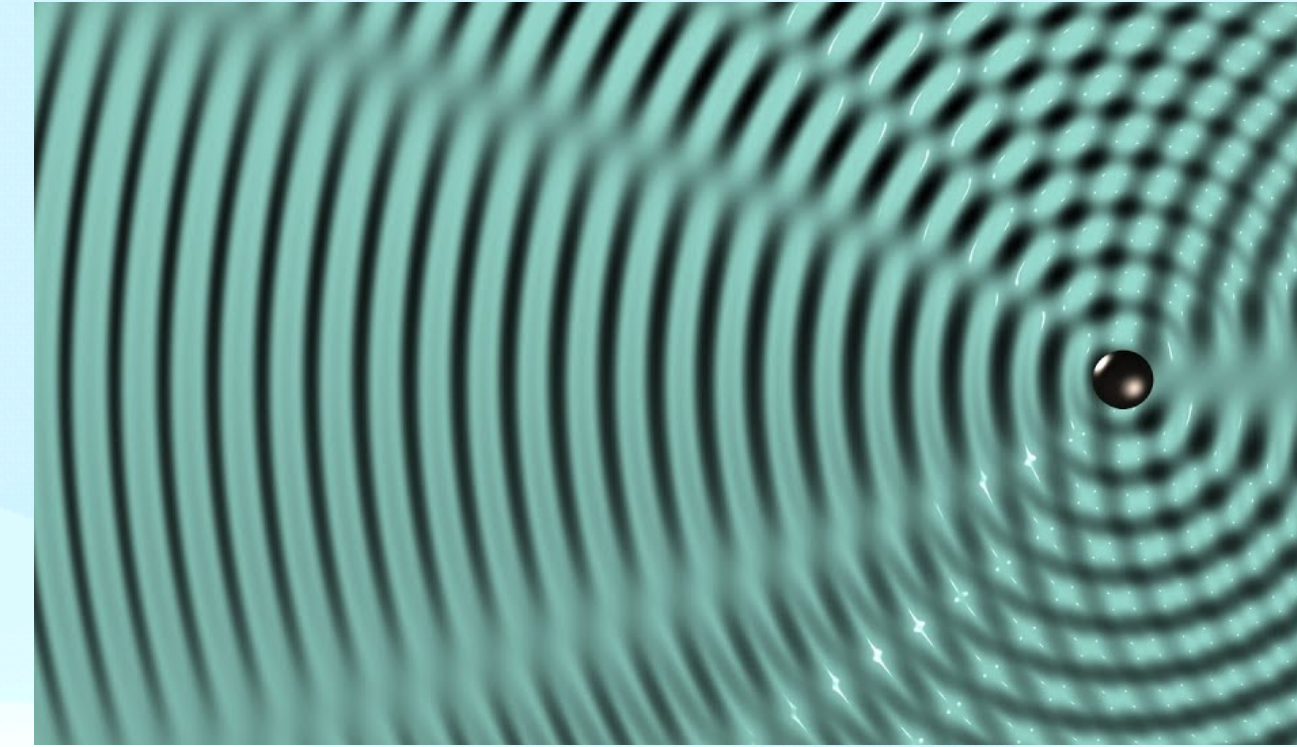


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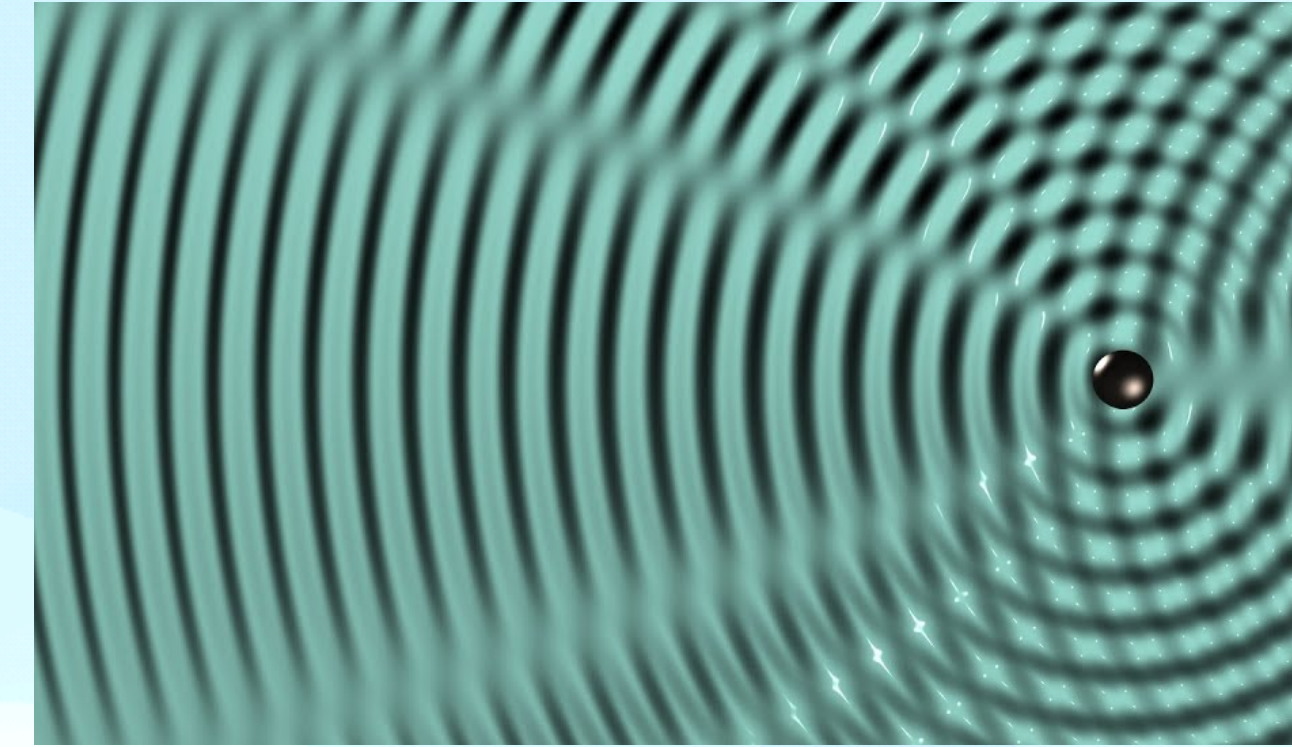


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Proper time path
integral:

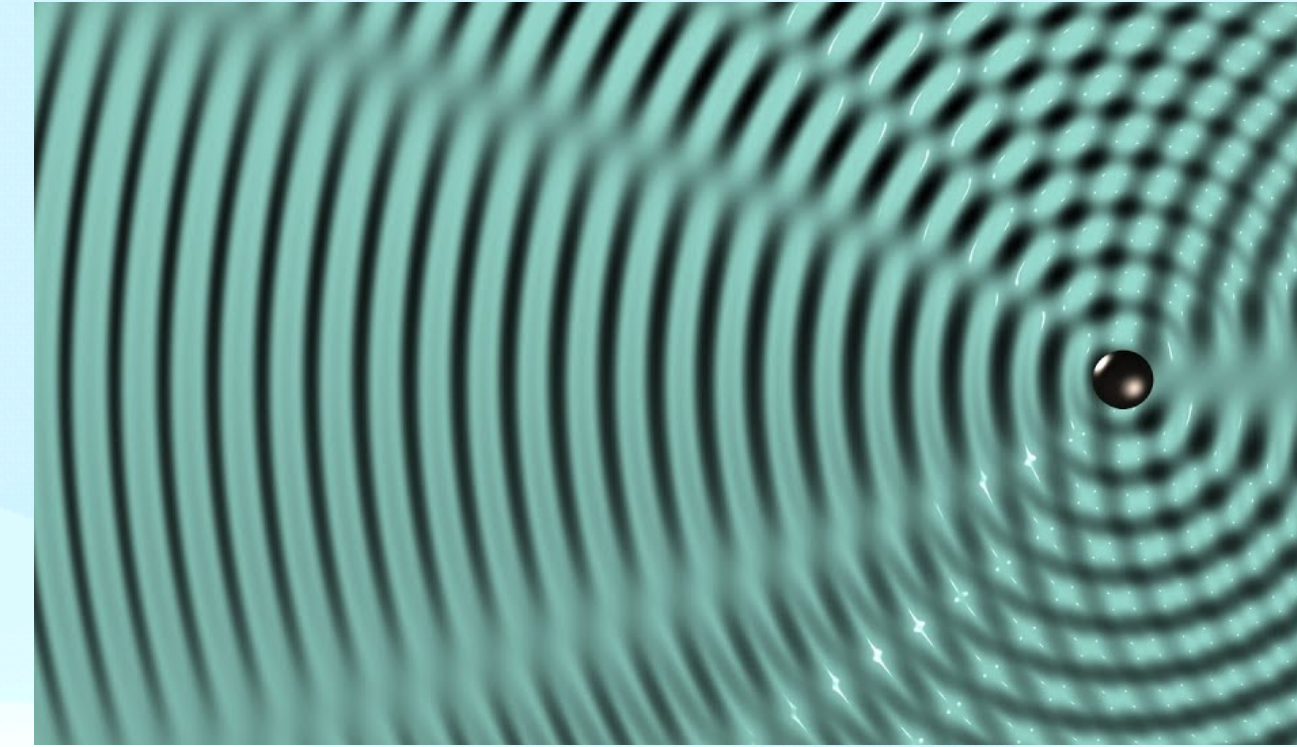
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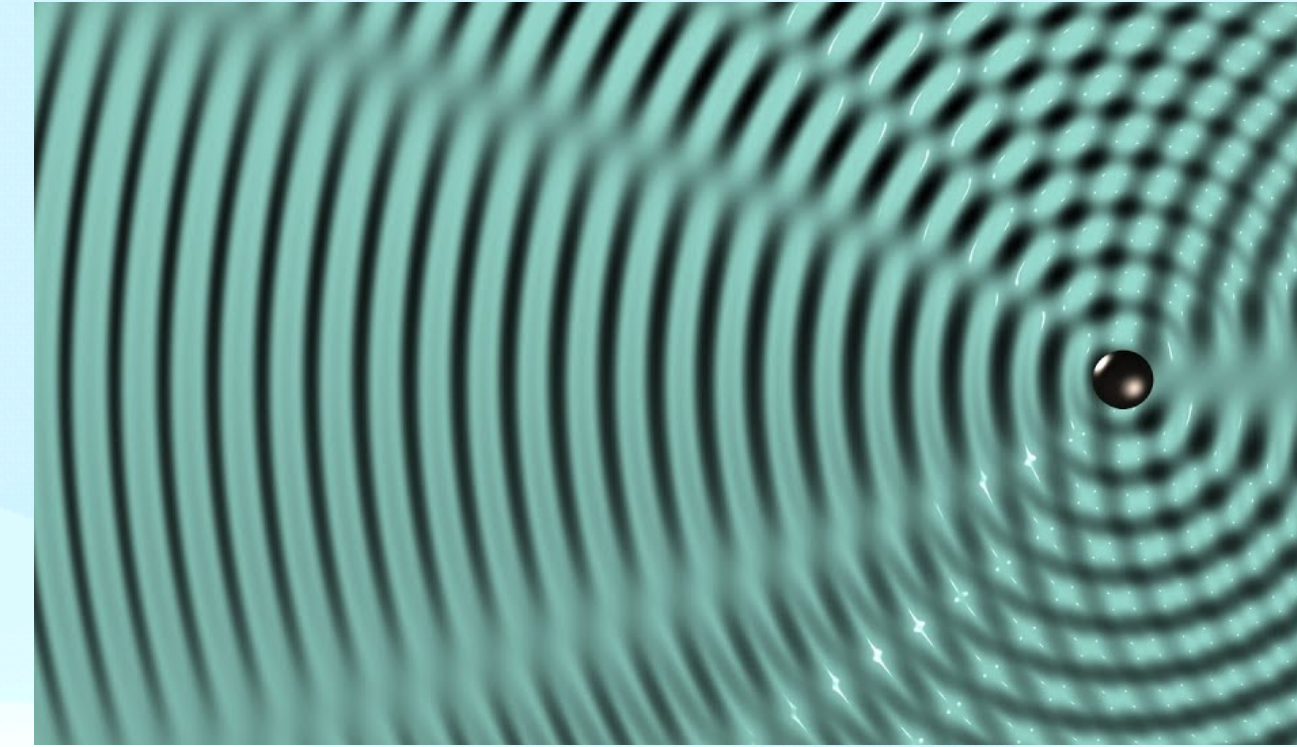
Sum over paths

Proper time path integral

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Particle action

Proper time path
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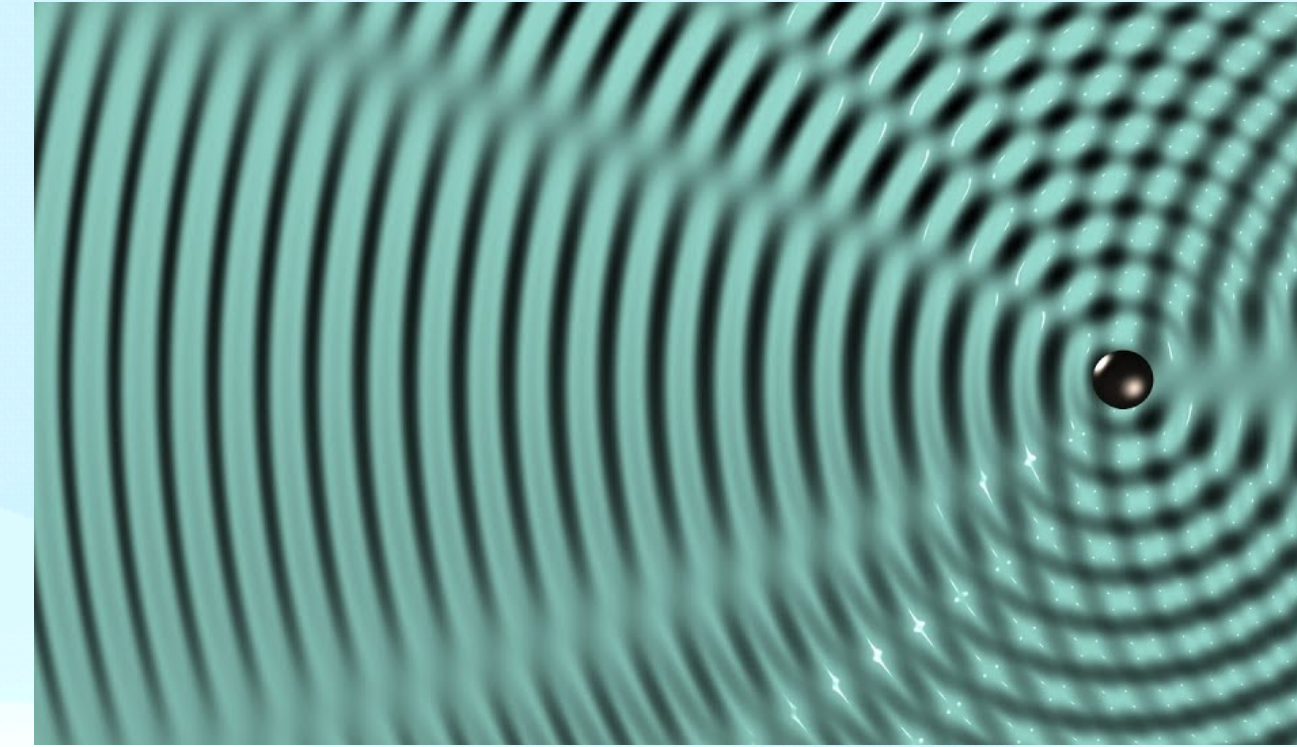
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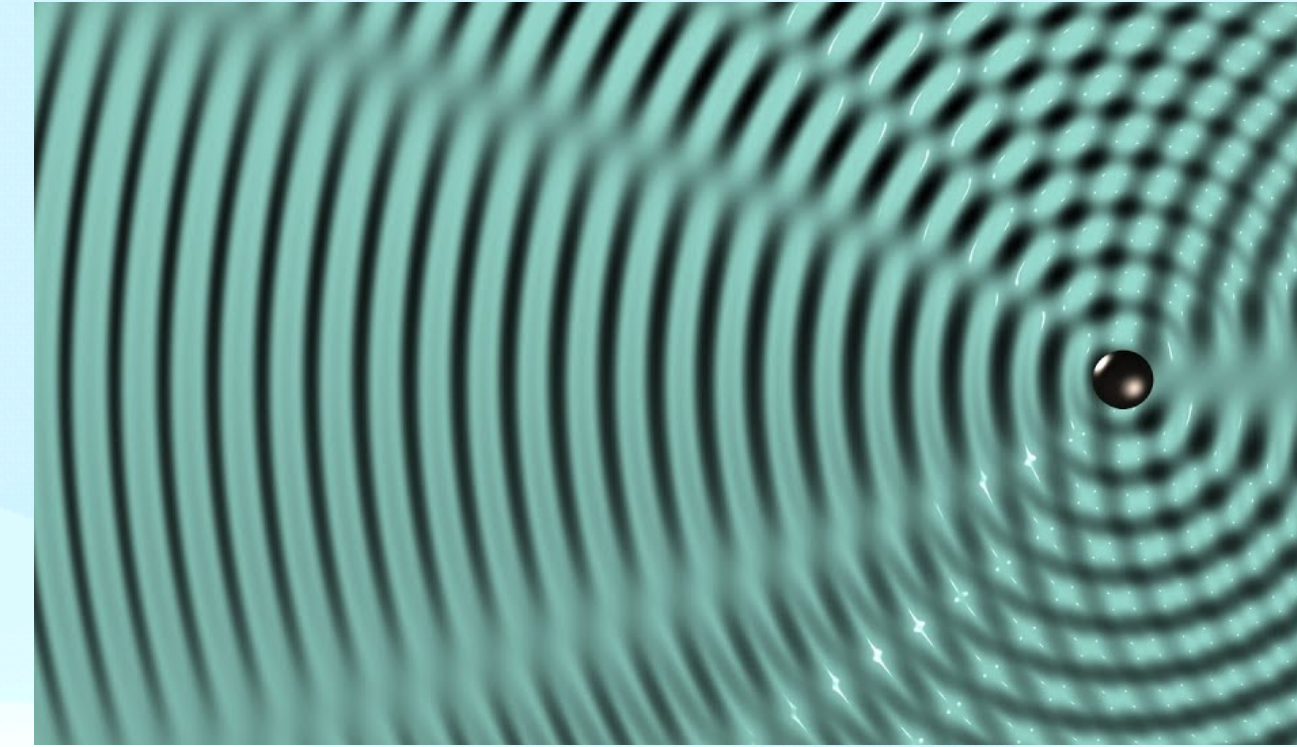
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Integrate away τ

Sum over paths

$$\omega = 1/\hbar$$

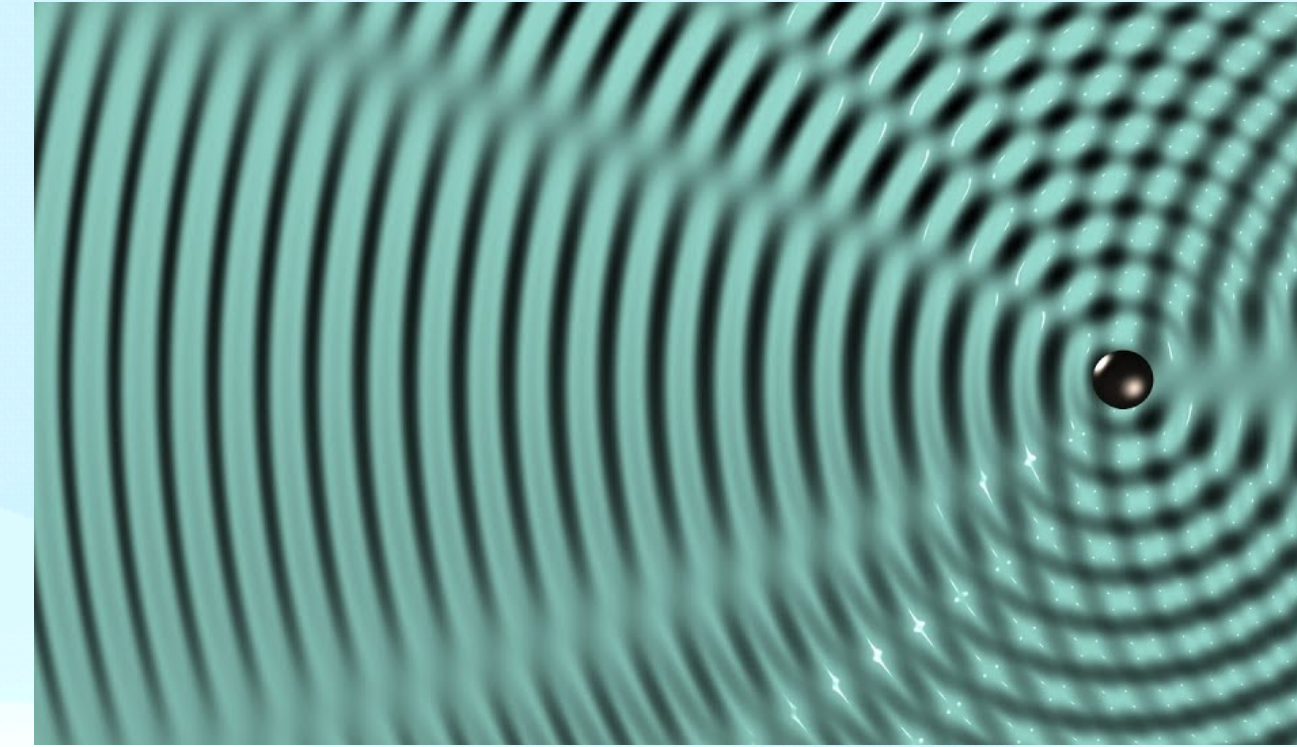
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Particle action

Exact particle-like solution WITHOUT the need of Eikonal approximation

What you can find in 2405.20208:

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1. $\omega \rightarrow \infty$ limit to recover geometric optics

$$(\delta\hat{S}/\delta\tau = 0 \text{ and } \delta\hat{S}/\delta\mathbf{x} = 0)$$

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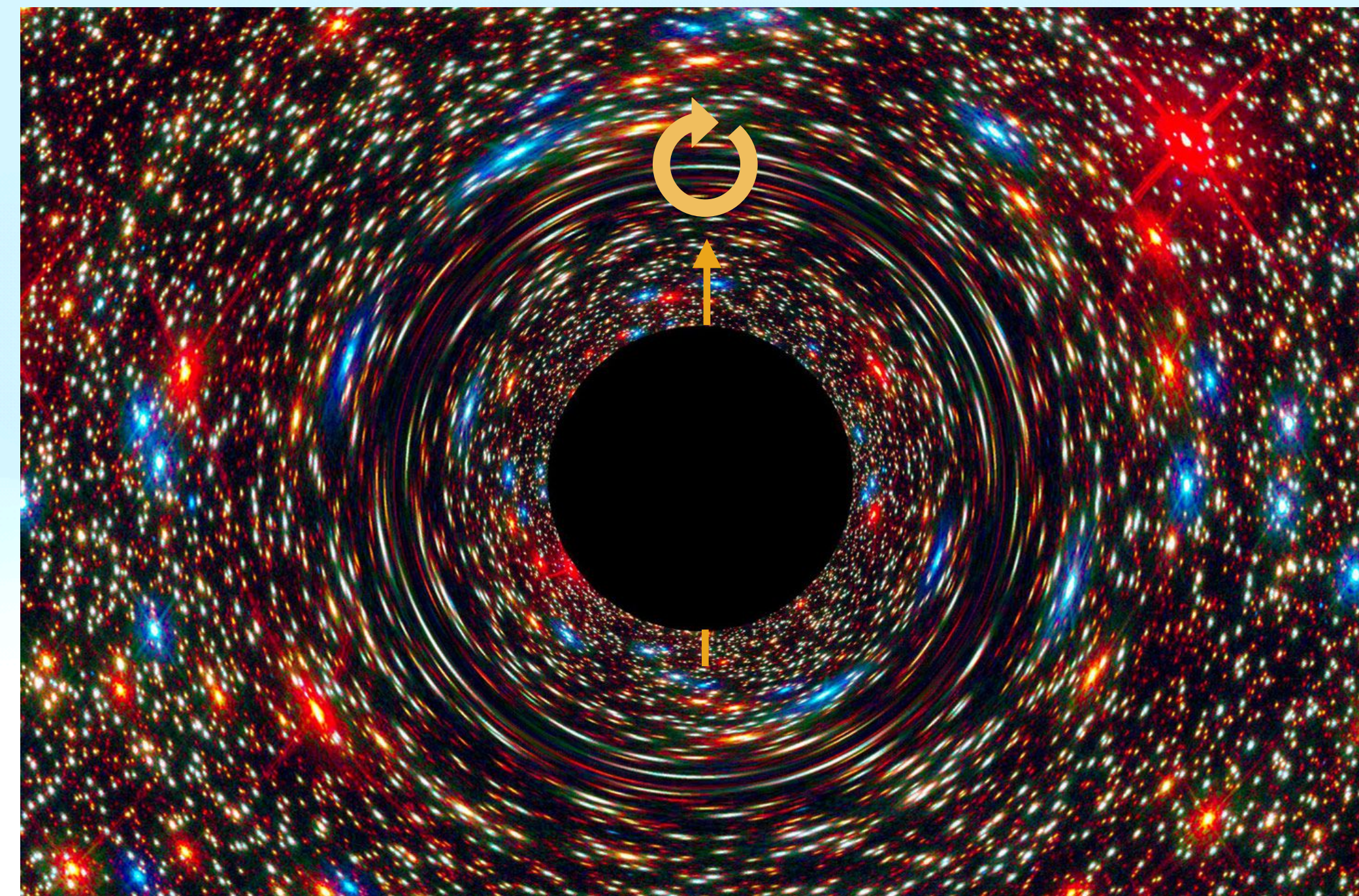
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6. Polarization effects

Polarization effects on a Kerr background

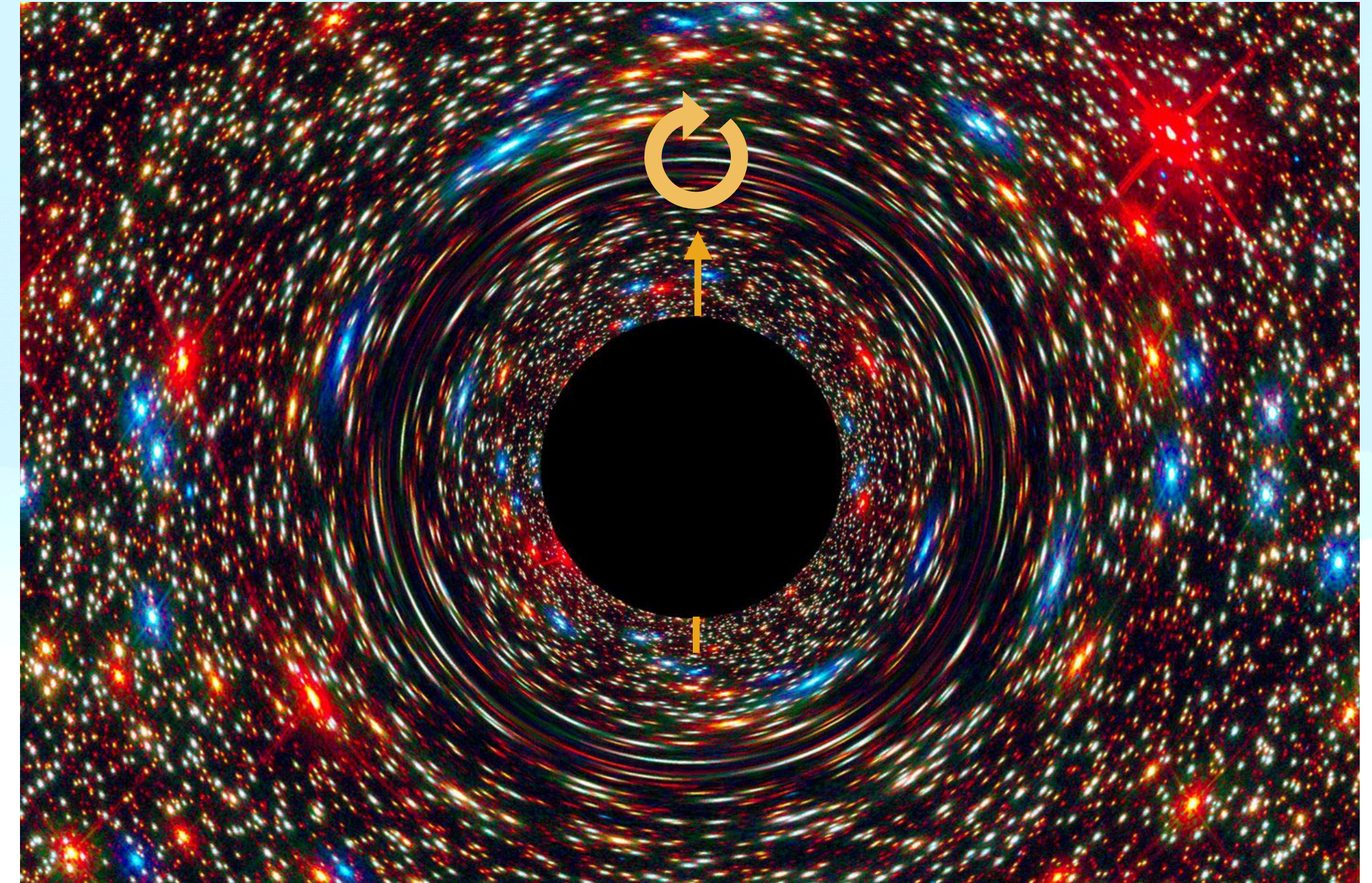
S. Teukolsky (1973)



Polarization effects on a Kerr background

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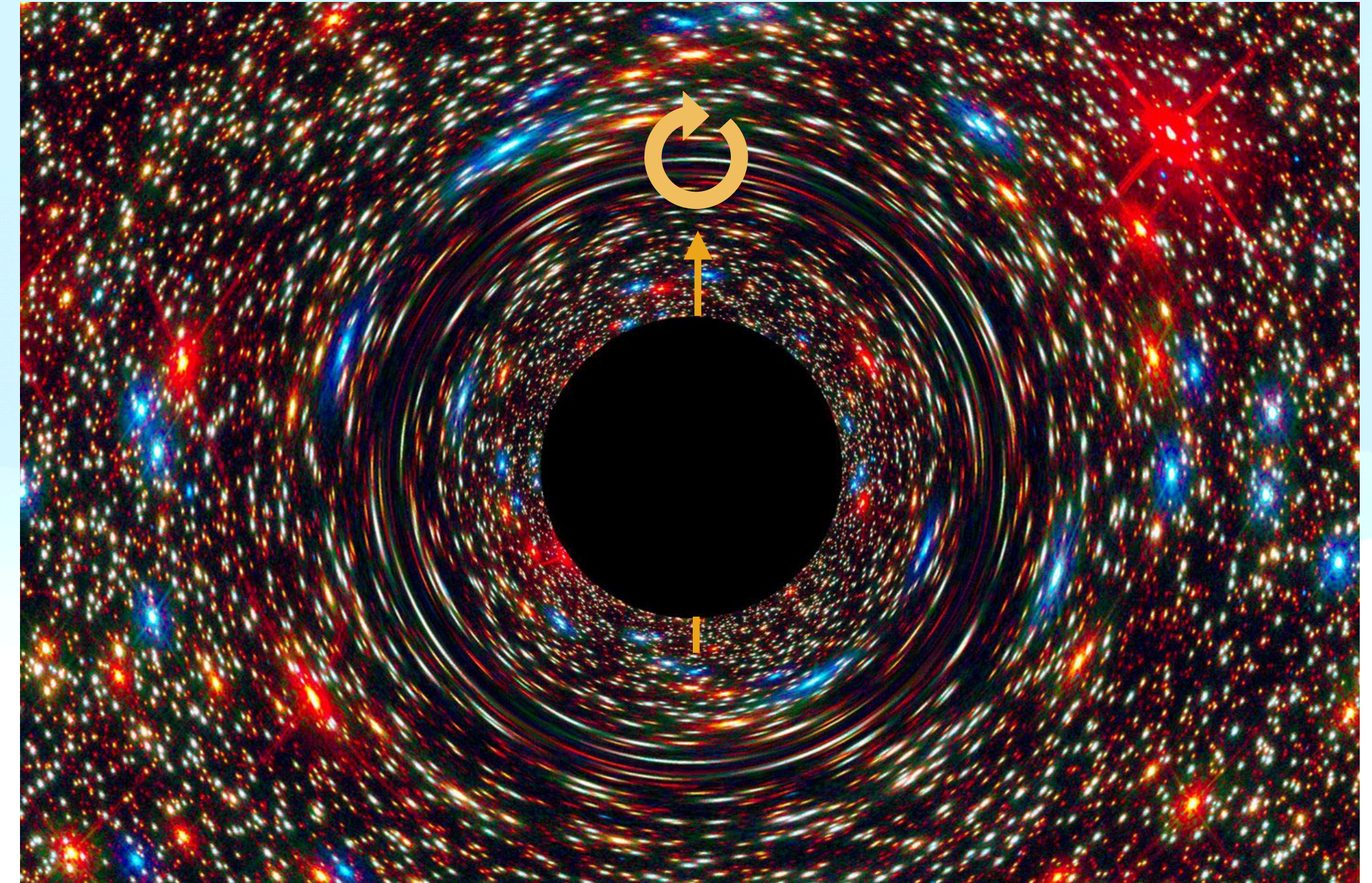
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1. Lens = Kerr BH
2. Use BH perturbation theory long-standing results

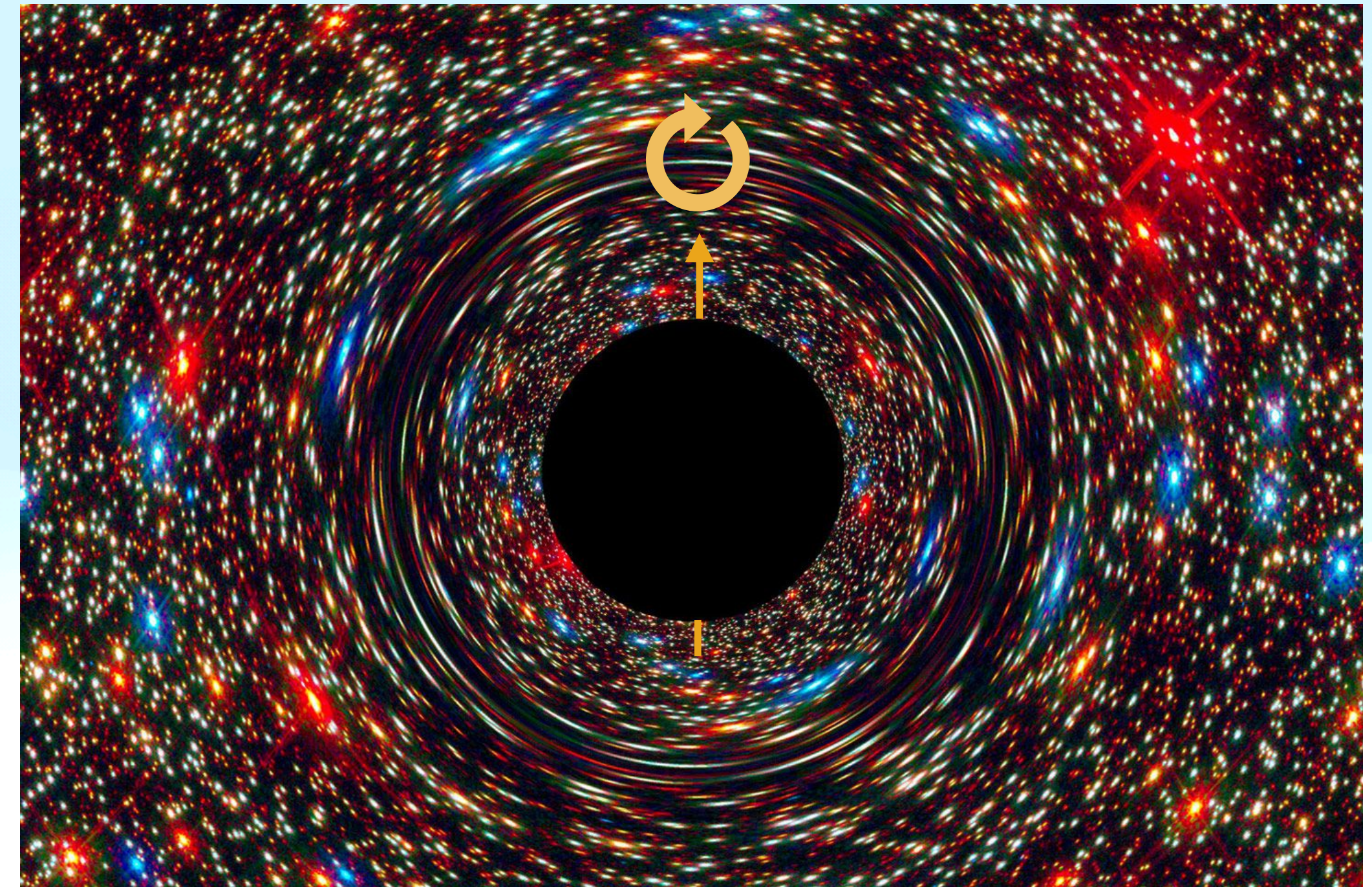


Polarization effects on a Kerr background

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$$\hat{O}[M, a, \omega, s] \psi_{\omega}^s(r, \theta, \varphi) = 0$$



Polarization effects on a Kerr background

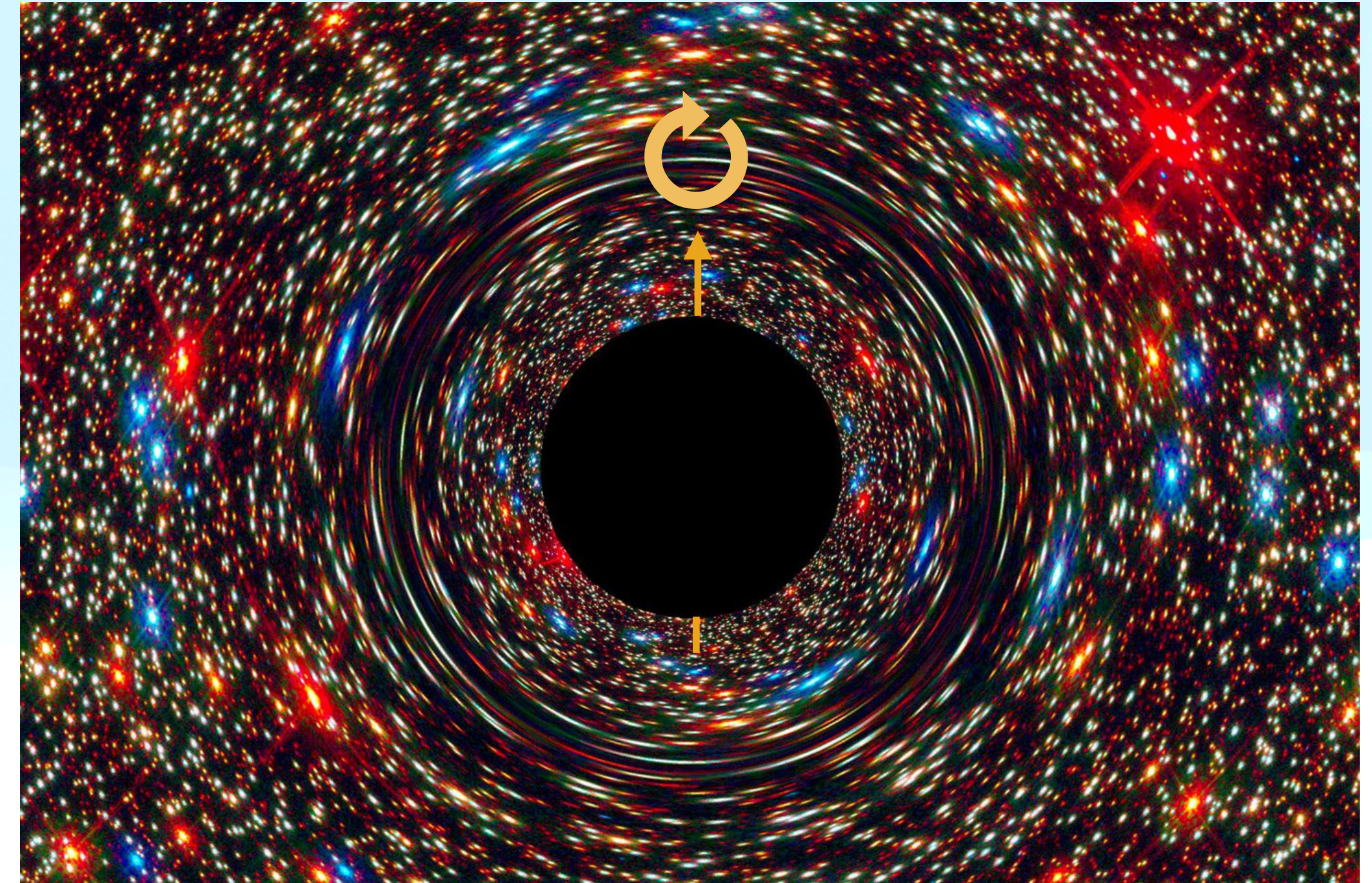
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Differential
operator



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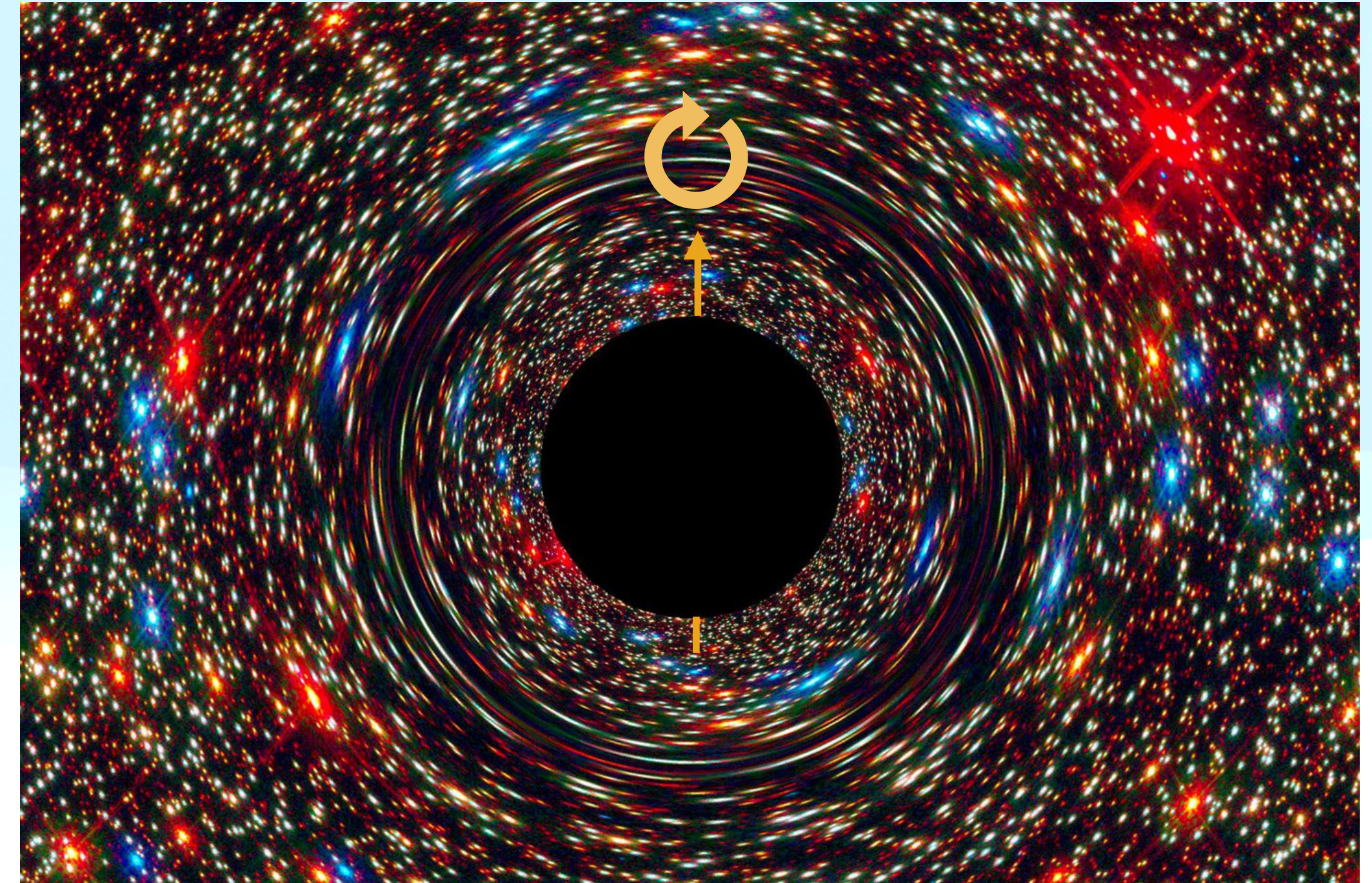
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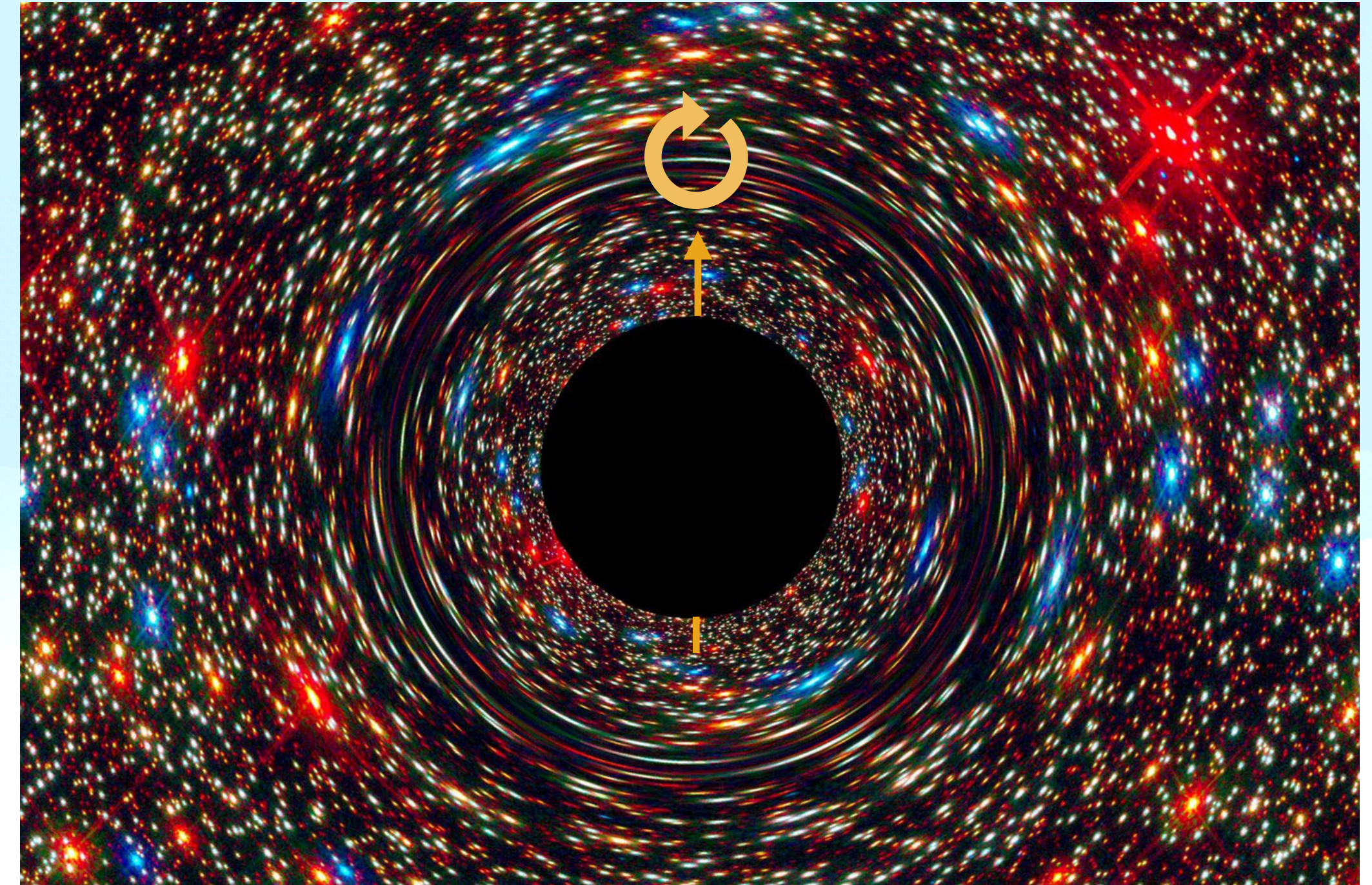
↓
Differential
operator

↓
Newman-Penrose scalar, e.g.:
 $\psi_{\omega}^{s=2} \supset \{\ddot{h}_+ \pm i \ddot{h}_x\}$



Polarization effects on a Kerr background

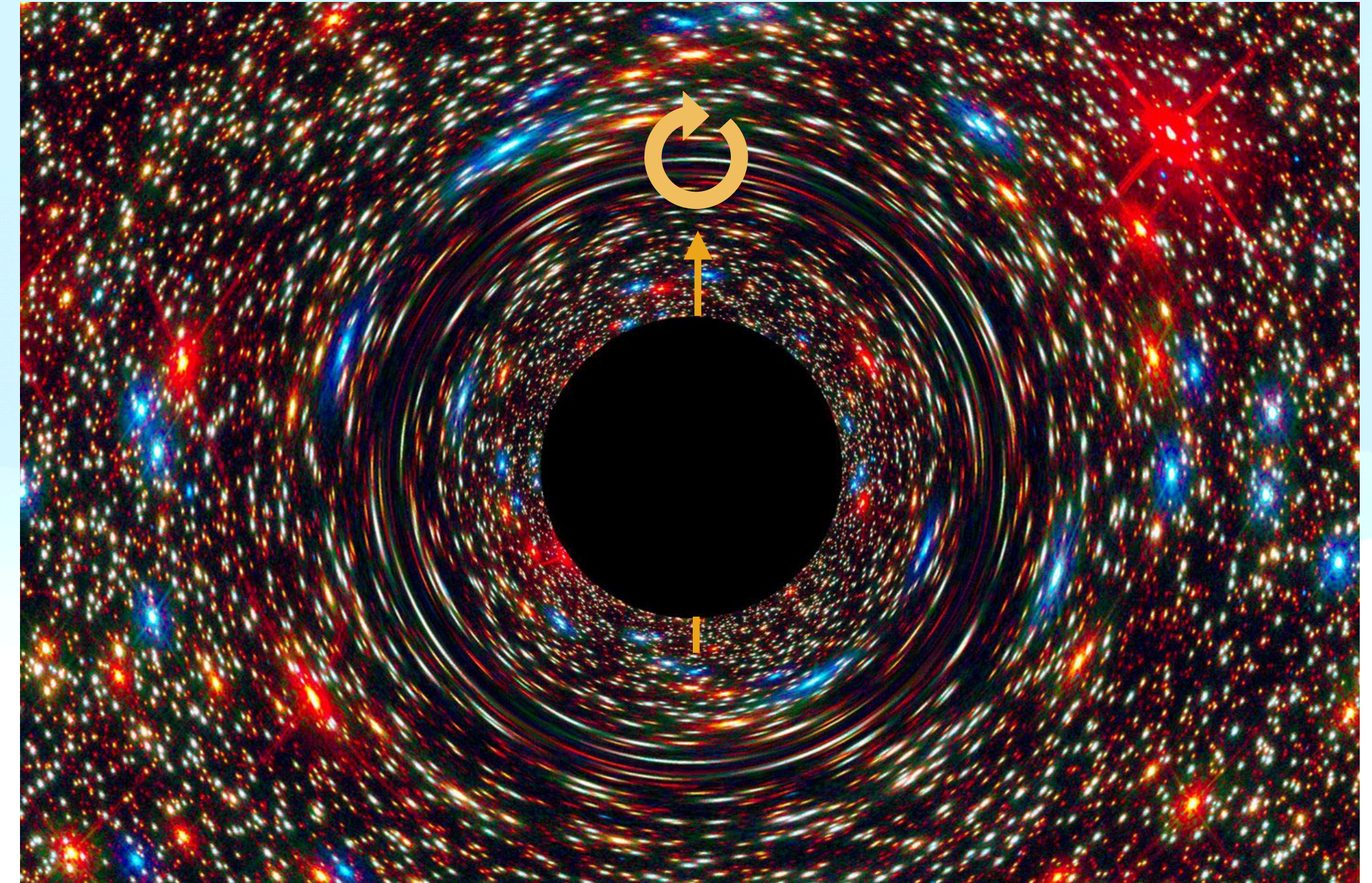
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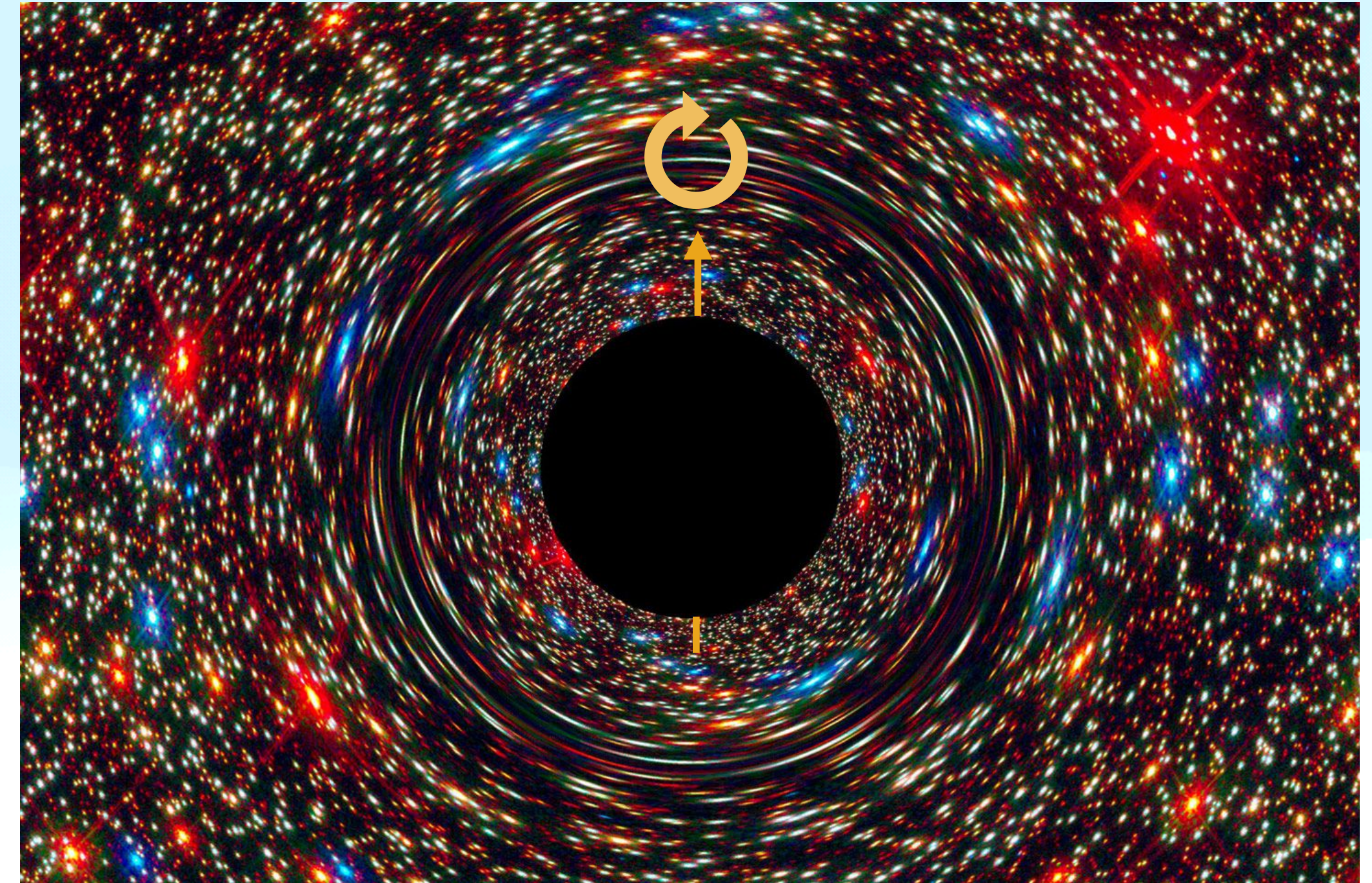


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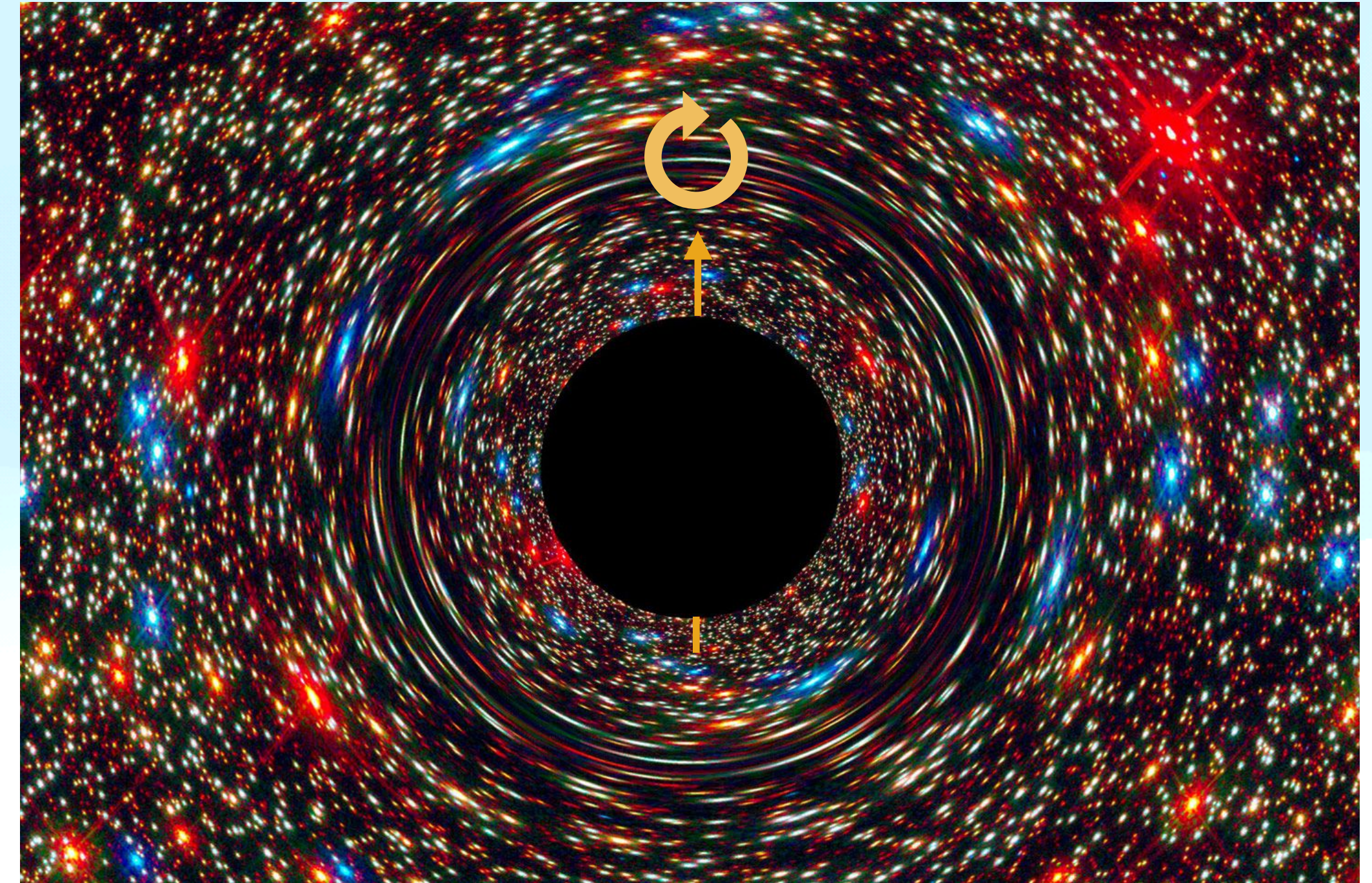
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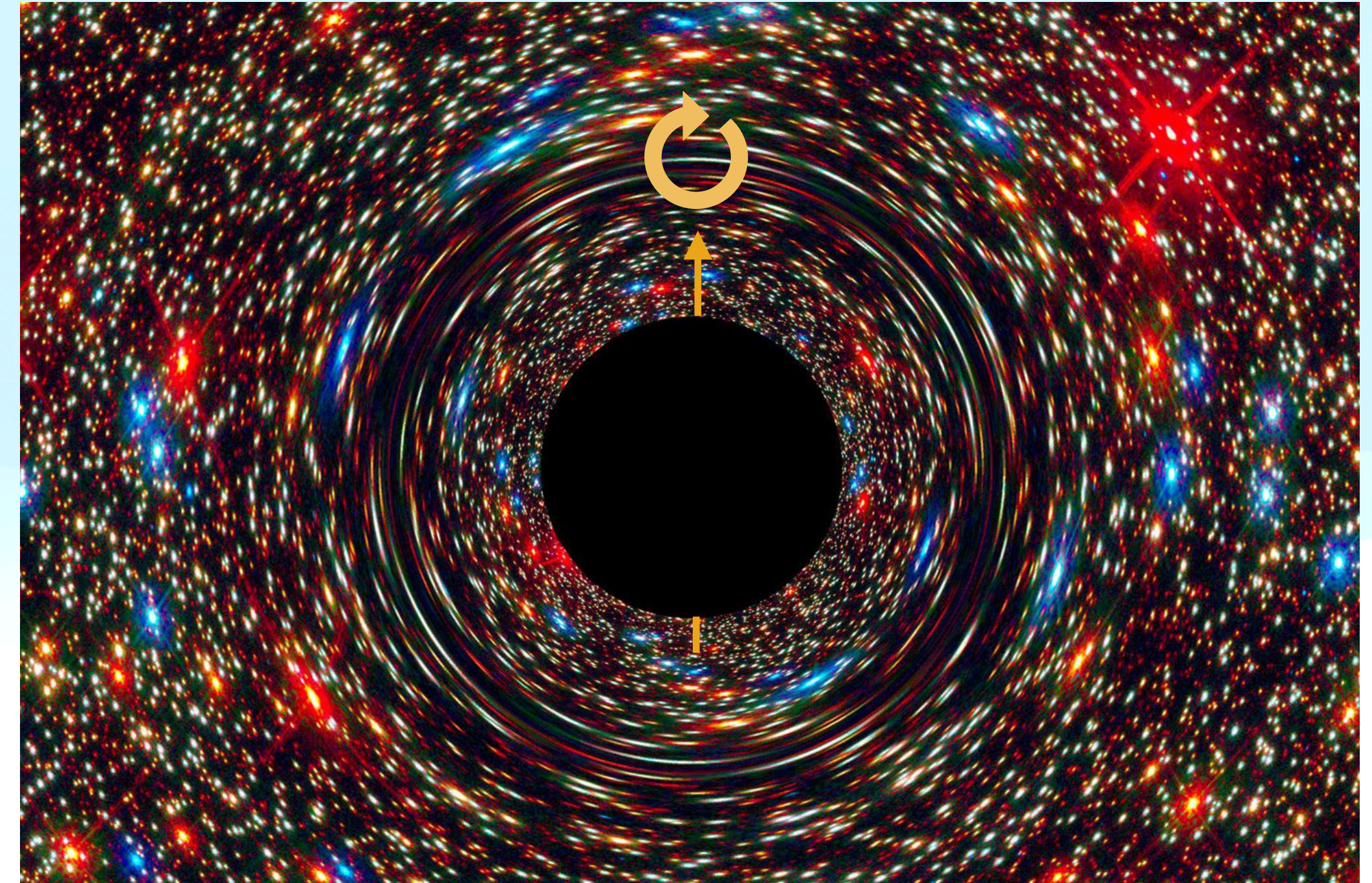
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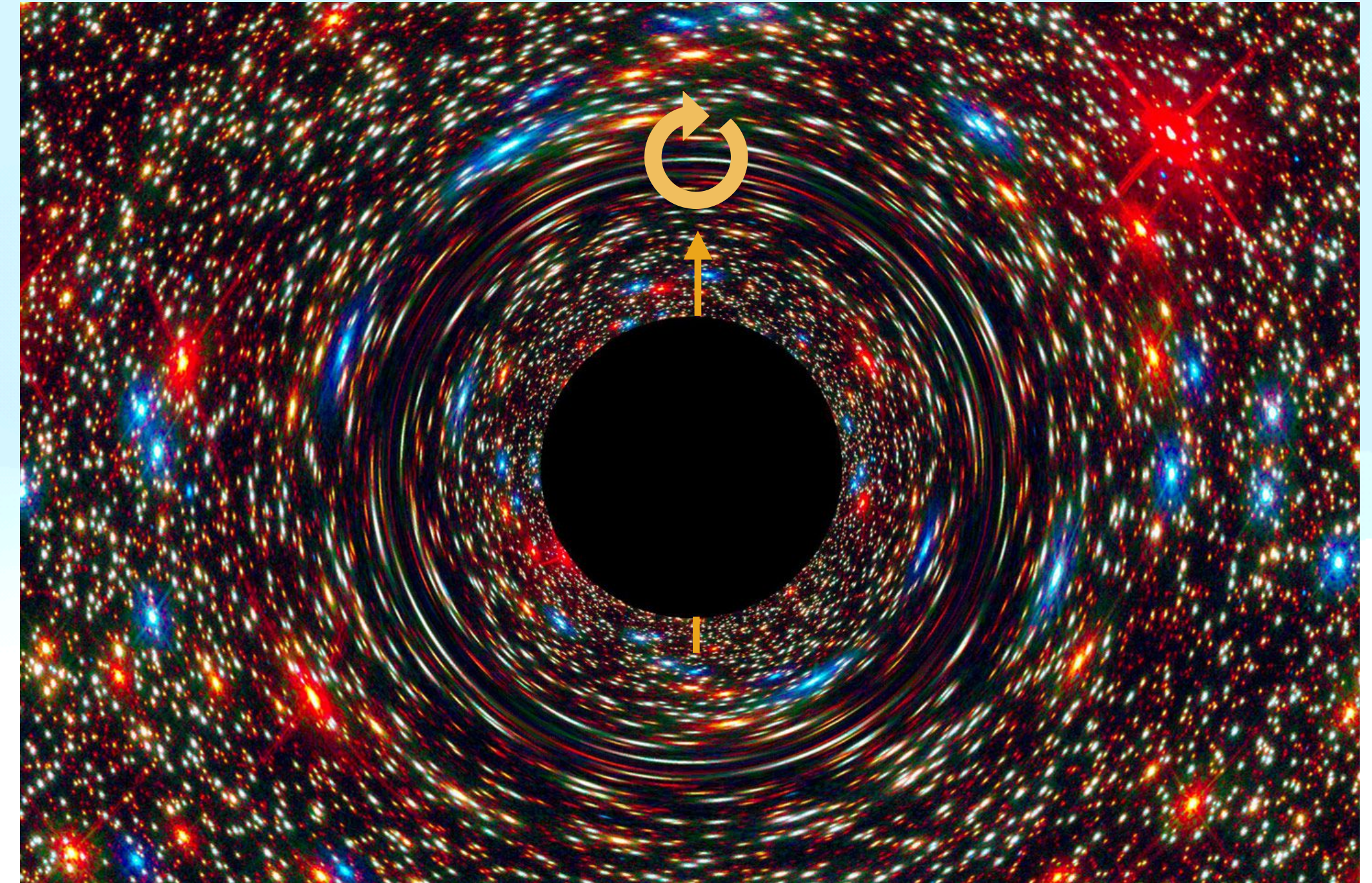
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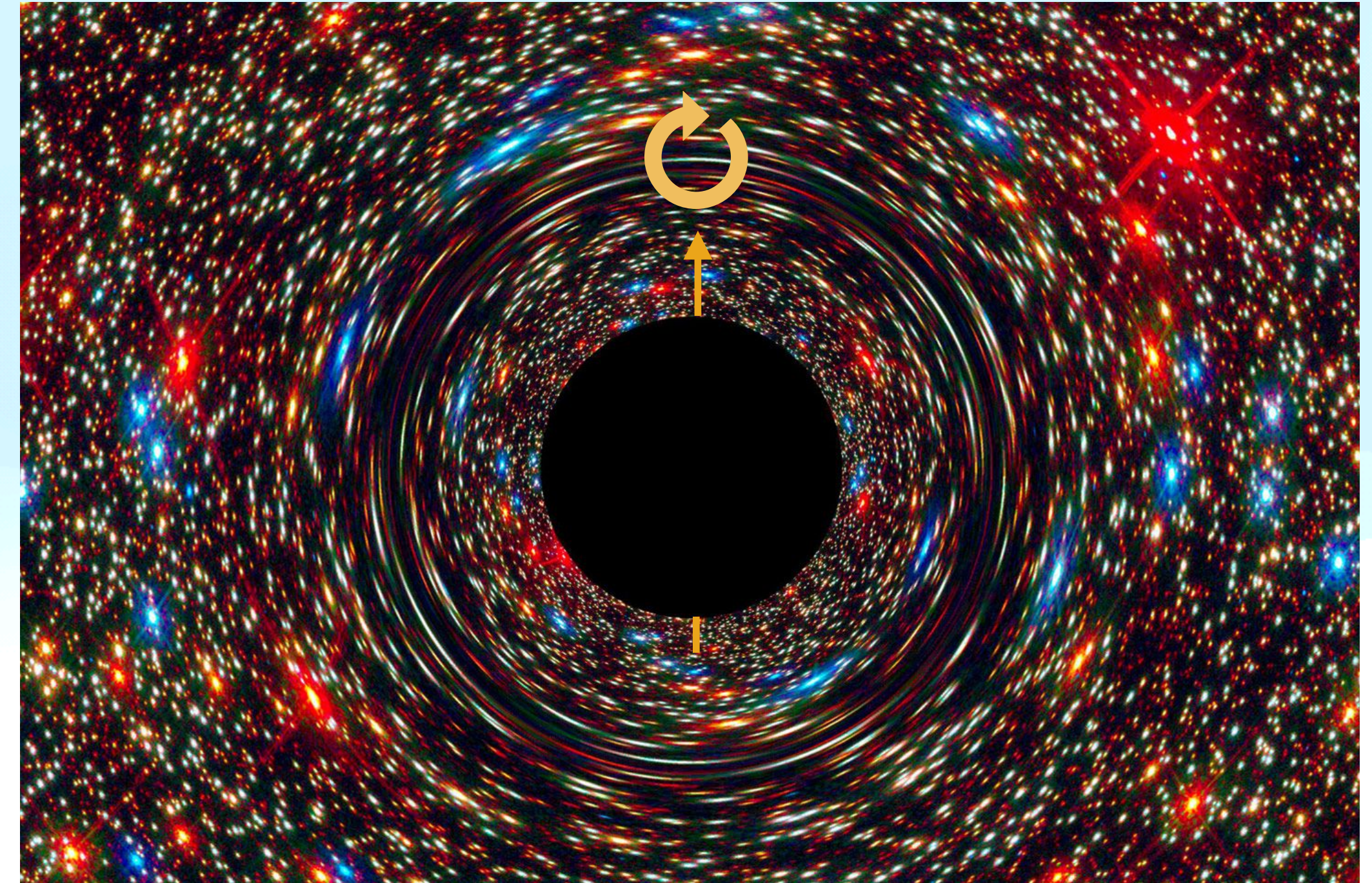
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$$4\tilde{U}_{\ell m}^s(\omega, r) \approx \underbrace{-4\frac{M}{r}}_{\text{Same as diffraction integral}} + \underbrace{\frac{\ell(\ell+1) + s(s+1)}{\omega^2 r^2}}_{\text{Angular momentum (decomposition)}} - \frac{2is}{\omega r}$$

Same as
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Angular momentum
(decomposition)

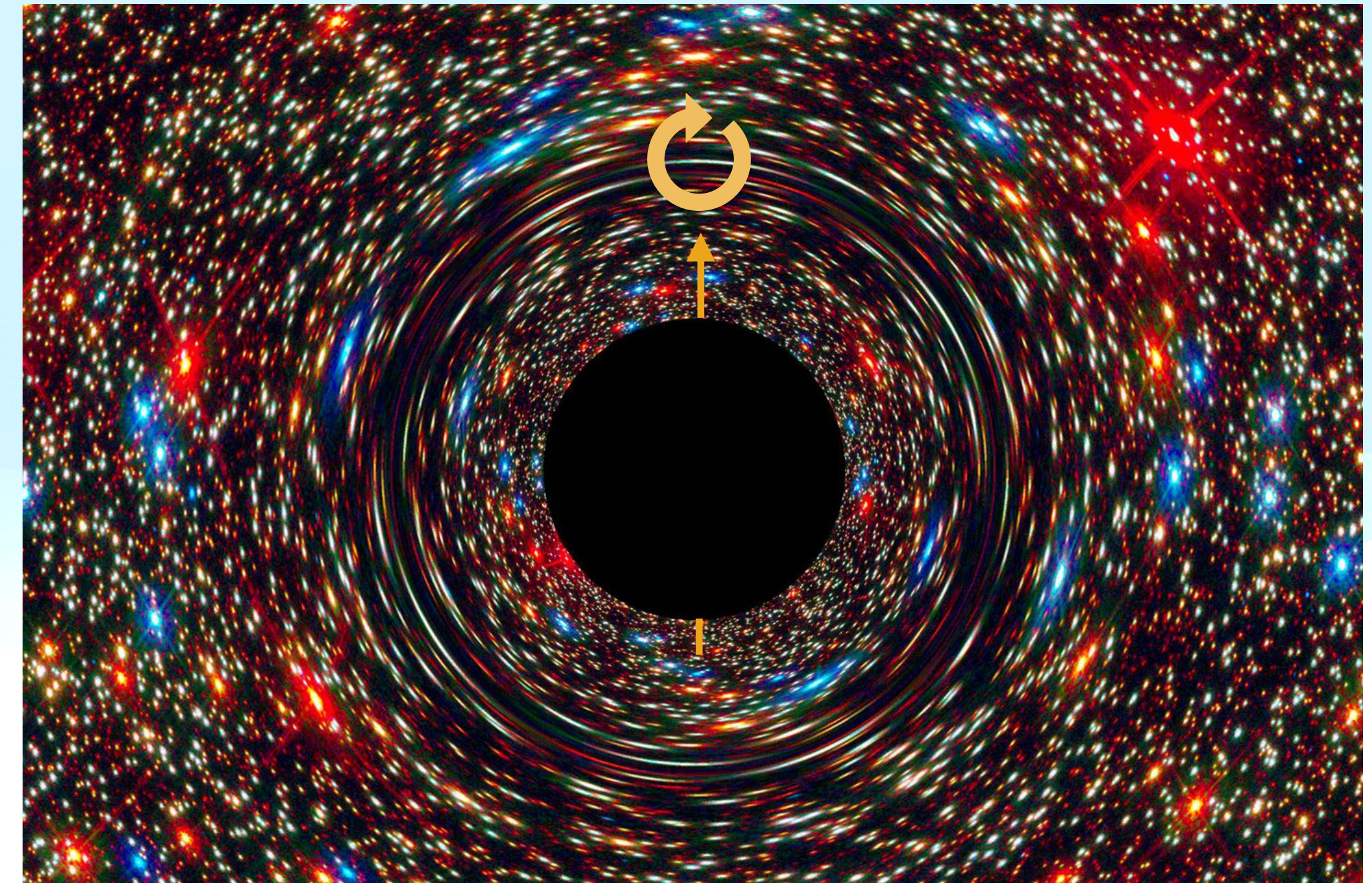


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Spin dependent terms:
Negligible in $\omega \gg 1$
limit

Summary

What's next

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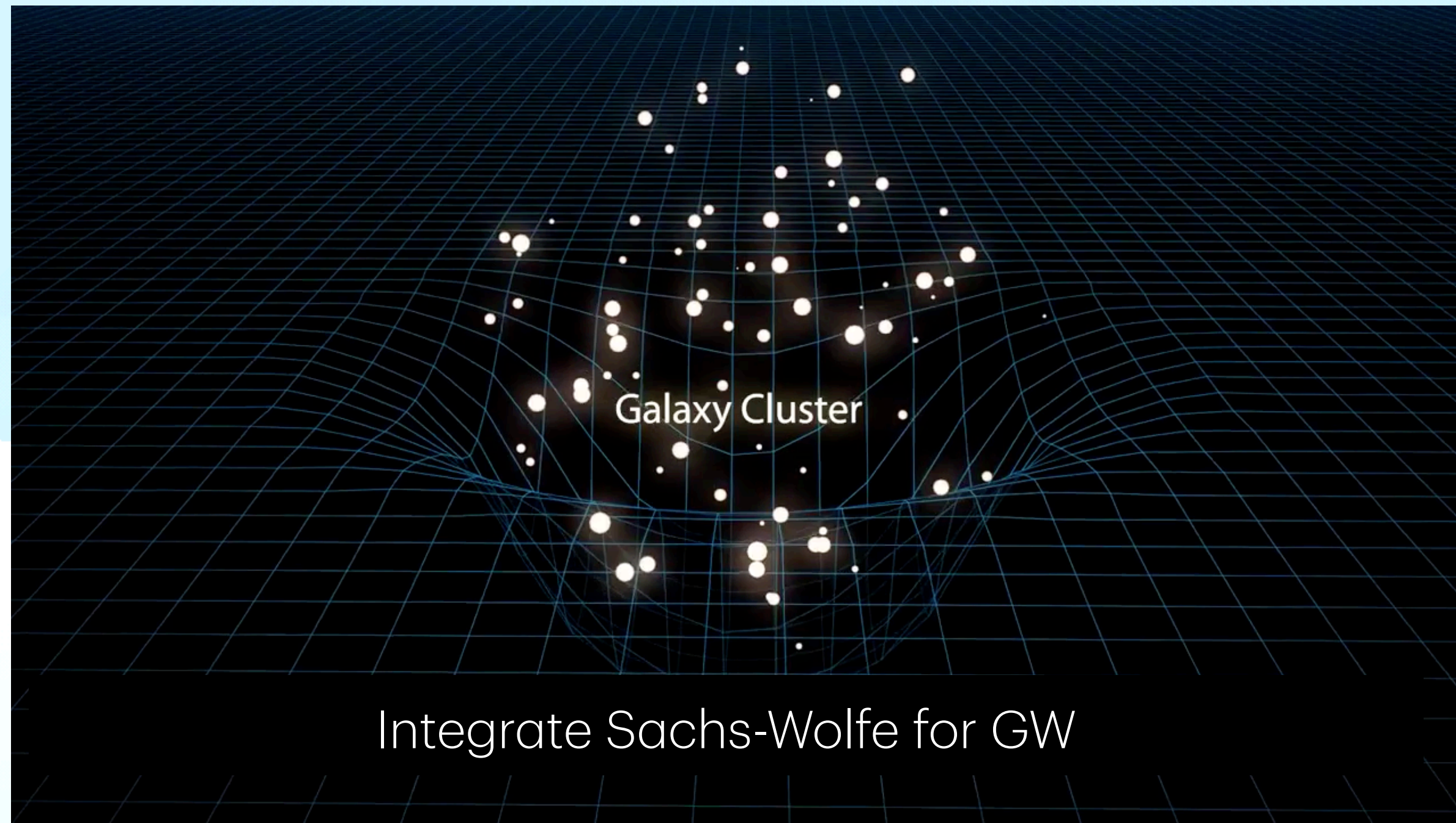
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Thank you!

Optical regimes

Geometric vs Wave optics

High Frequency: $\omega R_S \gg 1$



Ray description

Low Frequency: $\omega R_S \lesssim 1$

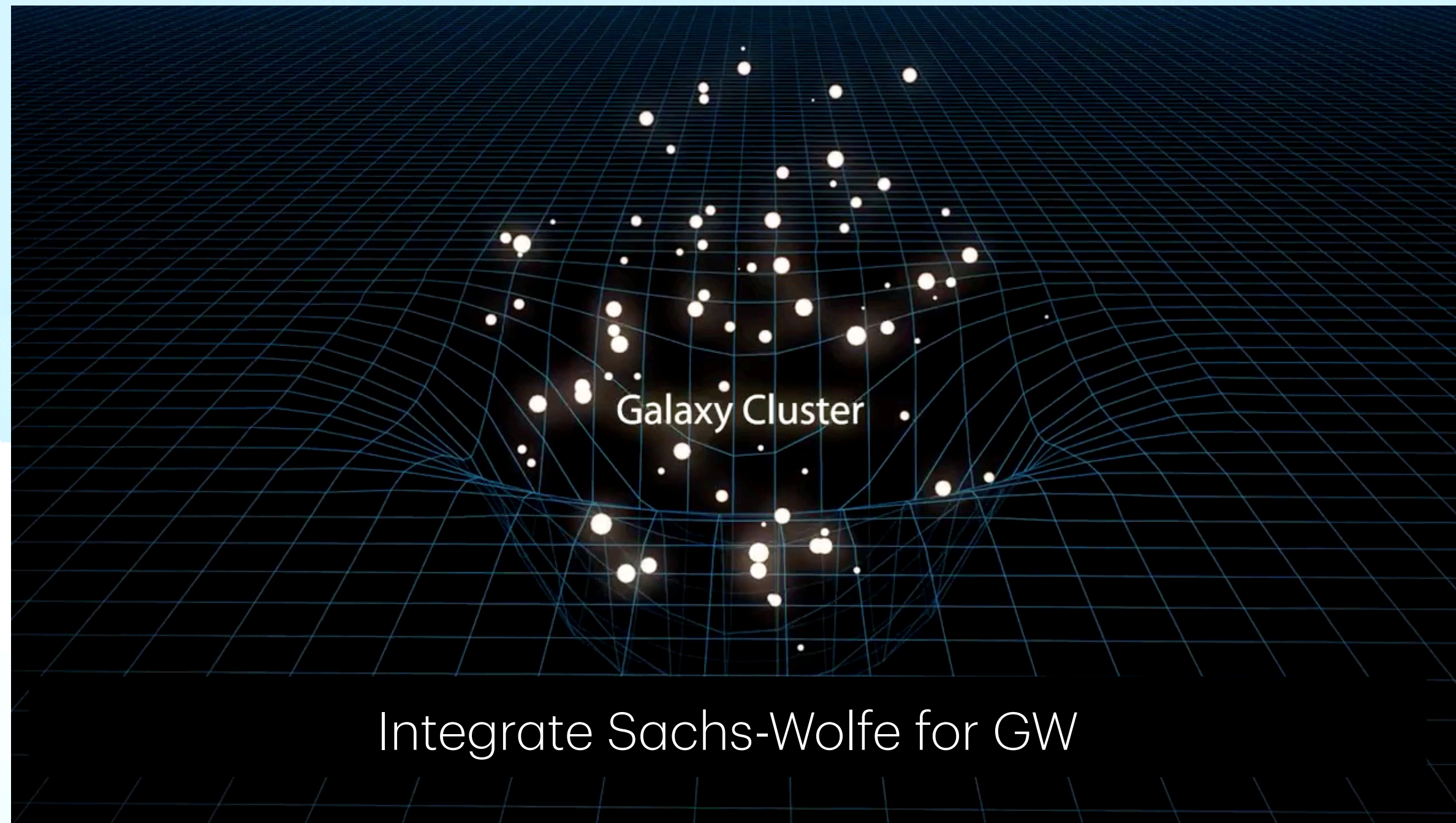


Wave effects

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Wave effects

3.4 Wave effects

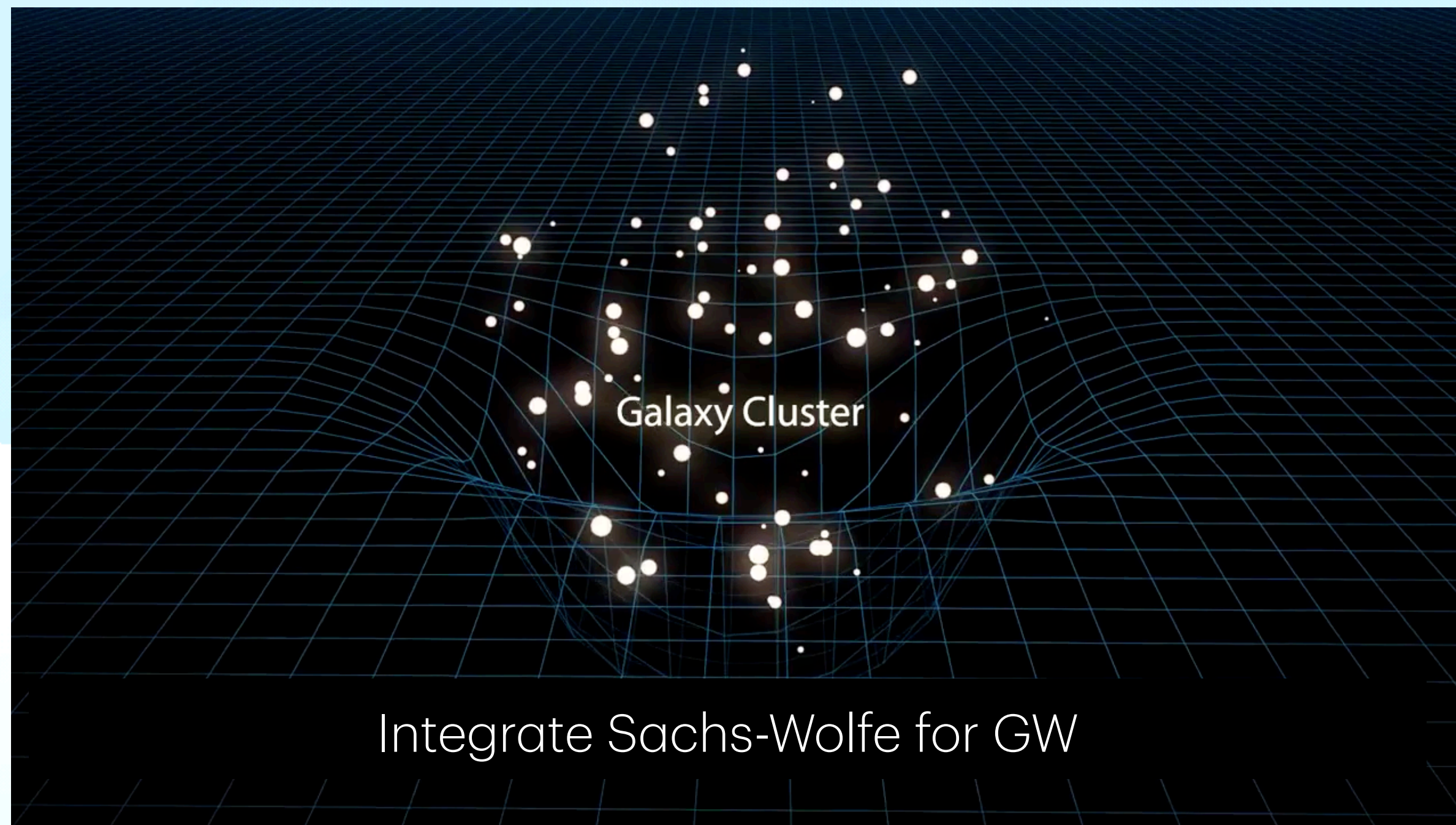
LISA CosGW: 2204.05434

GWs can be emitted at low frequencies ($\omega \lesssim 1$), allowing the observation of wave diffractive phenomena. For typical LISA sources, wave optics as in Eq. (15) needs to be considered for lenses with masses $M_L \sim 10^6 - 10^9 M_\odot$, cf. Eq. (13).

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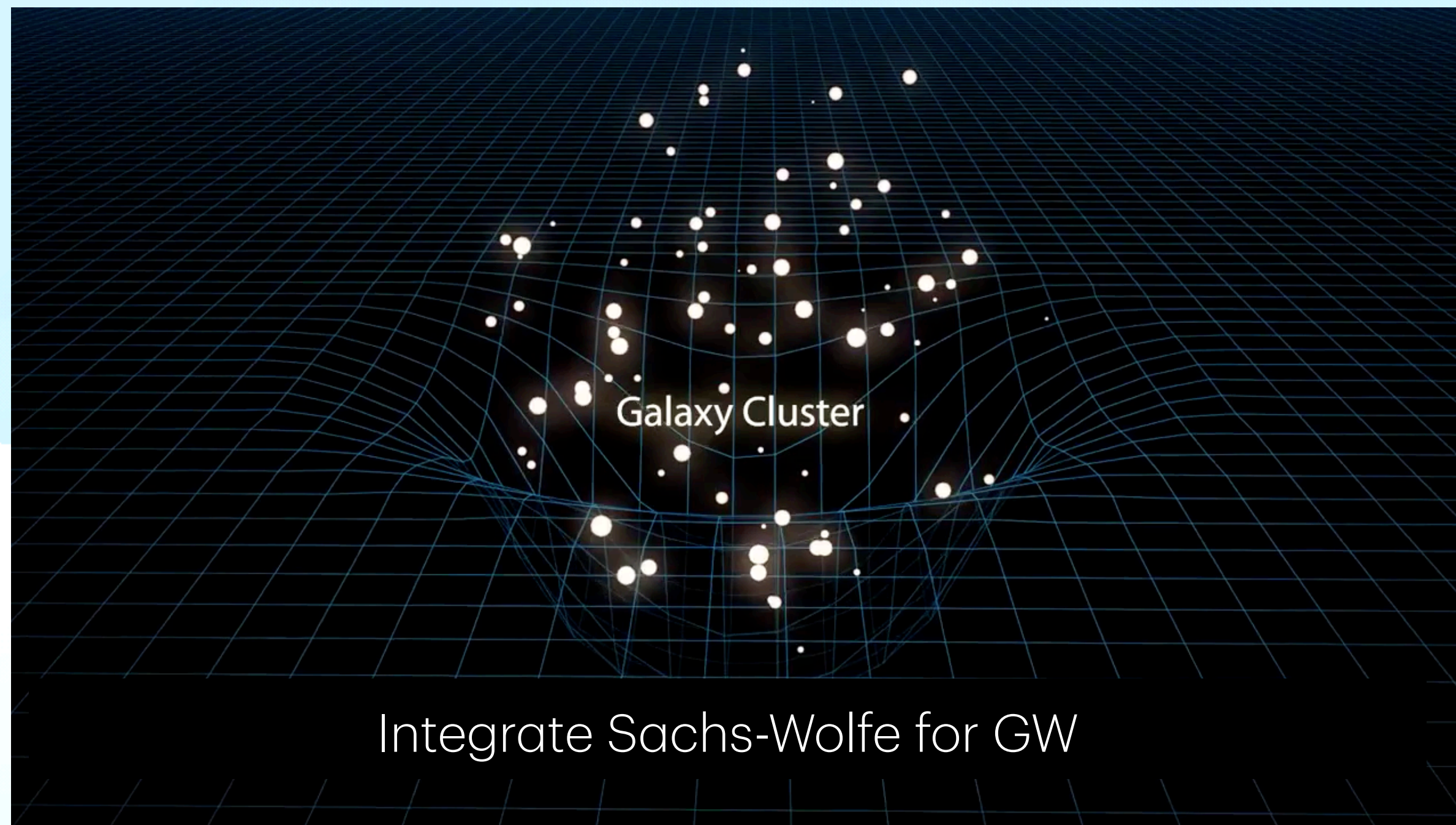
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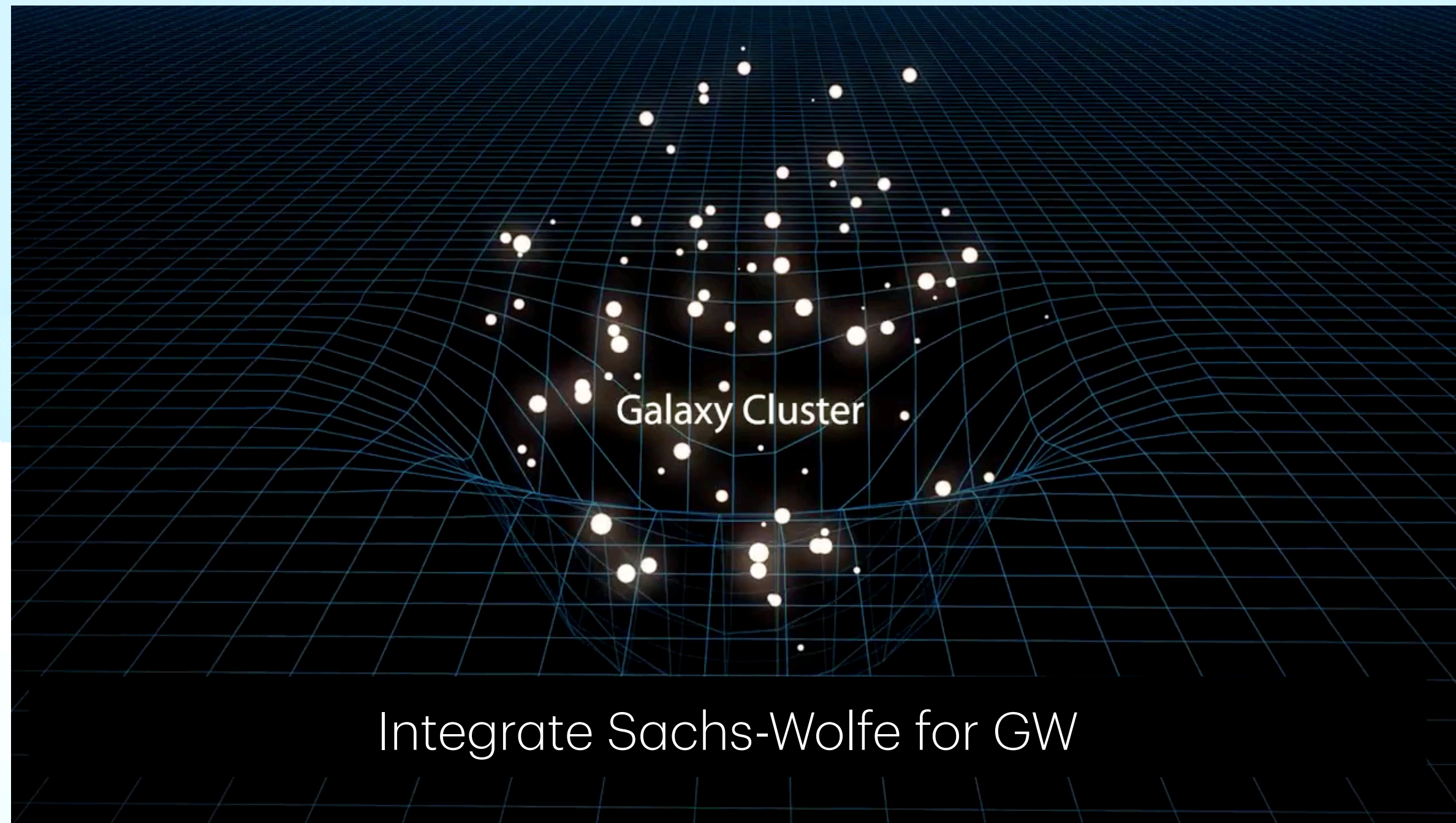
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