



Higher-order corrections for VBS and polarised multi-boson processes

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LHC EW WG General Meeting CERN, July 10–12, 2024





Introduction

- Electroweak vector-boson scattering at the LHC
- Olarised vector bosons
- General remarks on results for EW corrections

5 Conclusion







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Motivation

• Vector-boson scattering (VBS)

important process to test the Standard Model (SM)

- sensitive to quartic gauge couplings
- sensitive to Higgs sector and electroweak symmetry breaking (EWSB)
- polarised vector bosons

offer new observables to test SM and mechanism of EWSB

- production of polarised vector boson pairs
- longitudinal VBS scattering

Experiments use general purpose codes like SHERPA, MADGRAPH5_AMC@NLO or POWHEG/PYTHIA including

- NLO QCD corrections including parton-shower matching etc.
- but no NLO EW corrections

Focus of the talk

- EW corrections
- available results for (mostly) leptonic final states
- $\bullet\,$ remarks on incorporation in experimental analyses $\Rightarrow\,$ discussion





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Processes: $pp \rightarrow VV + 2j \rightarrow 4\ell + 2j$

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Vector-boson scattering (VBS) signal (decays not shown)



Irreducible background to VBS



- $\mathcal{O}(\alpha^4)$ for stable Vs, $\mathcal{O}(\alpha^6)$ with decays • EW process:
- QCD-induced process: $\mathcal{O}(\alpha_s^2 \alpha^2)$ for stable Vs, $\mathcal{O}(\alpha_s^2 \alpha^4)$ with decays
- interferences between EW and QCD contributions:
 - $\mathcal{O}(\alpha_{\rm s}\alpha^3)$ for stable Vs, $\mathcal{O}(\alpha_{\rm s}\alpha^5)$ with decays
- gluonic channels for neutral final states

UNIVERSITÄT Expansion in multiple couplings Example: $pp \rightarrow 4\ell jj$ (vector-boson scattering: $pp \rightarrow VV jj$) LO: pure EW diagrams $\mathcal{O}(e^6)$ and diagrams with gluons $\mathcal{O}(e^4g_s^2)$

NLO: EW and QCD corrections to both types of diagrams at level of cross section:



consequences:

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- QCD and EW corrections cannot be separated in general
- consider complete (well-defined) orders $\mathcal{O}(\alpha_s^n \alpha^m)$
- \bullet QCD corrections to leading LO terms, i.e. $\mathcal{O}(\alpha_{\rm s}^3\alpha^4)$, well defined
- QCD corr. to EW LO overlaps with EW corrections to LO interference

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Virtual diagrams mix QCD and EW corrections:

- EW correction to LO QCD amplitude
- QCD correction to LO EW amplitude
- QED and QCD IR singularities

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 \Rightarrow separation into QCD and EW is not well-defined at NLO

real subtraction terms with both gluons and photons needed





Strategy

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- isolate EW process
- subtract QCD-induced process based on theoretical predictions
- \Rightarrow measurement of EW VVjj production

Issues

- QCD and EW corrections must be included in EW signal
- QCD and EW corrections should be included in irreducible QCD background (potentially large!)
- QCD-EW interference should be taken into account (if relevant)



Vector-boson scattering (VBS) topologies: $\mathcal{O}(q^6)$ all t channel (u channel)



irreducible background to VBS:

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EW background $\mathcal{O}(g^6)$, s channel (also t channel)

QCD background $\mathcal{O}(g_s^2 g^4)$ t channel (also s channel)

t channel: incoming quarks/antiquarks connected to outgoing quarks/antiquarks u channel: exchange identical quarks/antiquarks in final state s channel: incoming quark and anti-quark connected, all boson propagators time like

UNIVERSITÄT VBS approximation WÜRZBURG



VBS approximation

- Neglect interferences between t- and u-channel contributions and all s-channel contributions
 ⇒ keep only squares of t- and u-channel contributions
- calculation simplifies considerably (~ 1000 loop diagrams per channel at $\mathcal{O}(\alpha_{s}\alpha^{6})$)
- only applicable to order α^6 and corresponding corrections for VBS cuts (tailored to VBS processes, not applicable to $\alpha_s^2 \alpha^4$)
- EW and QCD corrections to VBS uniquely defined (interferences neglected by definition!)
- VBS approximation works within $\lesssim 1\%$ at LO for $M_{\rm jj} > 500\,{\rm GeV}$ Denner, Hošeková, Kallweit 1209.2389, Ballestrero et al. 1803.07943
- ullet VBS approximation fails for NLO QCD corrections for small $M_{\rm jj}$
 - $\sim 6\%$ for $pp \rightarrow \mu^+ \nu_\mu e^+ \nu_e jj$ and $M_{jj} \sim 500 \, {\rm GeV}$ 1803.07943
 - $\sim 24\%$ for ${\rm pp} \to \mu^+\mu^-{\rm e^+e^-jj}$ and $M_{\rm jj}\gtrsim 100\,{\rm GeV}$ 2009.00411
 - \Rightarrow contribution to theoretical error



• Full LO predictions: Ballestrero, Franzosi, Maina '10 (PHANTOM)

NLO QCD separately for EW ($\mathcal{O}(\alpha^6)$) and QCD-induced production ($\mathcal{O}(\alpha_s^2 \alpha^4)$)

- NLO QCD corrections to EW production in VBS approximation: Jäger, Oleari, Zeppenfeld (+ Bozzi) '06, '07, '09 (VBFNLO); Denner, Hošeková, Kallweit '12
 PS matching: Jäger, Zanderighi '11, '13 + Karlberg '14 (W⁺W[±], ZZ) Rauch, Plätzer '16 (W⁺W⁻), Jäger, Karlberg, Scheller '18, '24 (WZ)
- NLO QCD corrections to QCD production: Melia, Melnikov, Röntsch, Zanderighi '10, '11 (W⁺W⁺); Greiner et al. '12 (W⁺W⁻); Campanario, Kerner, Ninh, Zeppenfeld '13, '14 (VBFNLO) (W⁺W⁺, WZ, ZZ)
 PS matching: Melia, Nason, Röntsch, Zanderighi '11 (W⁺W[±], WZ, ZZ)
- $\bullet~{\rm EW}$ corrections for complete processes ${\rm pp} \rightarrow 4f + 2{\rm j}$
 - NLO EW and QCD corrections for VBS into $W^+W^\pm,\,W^+Z,\,ZZ$ Biedermann et al.'16; Denner et al.'19, '20, '22
 - full NLO corrections to VBS into W⁺W⁺ and ZZ Biedermann, Denner, Pellen '17; Denner, Franken, Pellen, Schmidt '21
 - NLO EW matched to EW PS and interfaced to QCD PS for W[±]W[±] within POWHEG-BOX-RES Chiesa, Denner, Lang, Pellen '19

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Existing NLO calculations – state of the art



Calculations for massive VBS within the SM

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- all processes known at NLO QCD accuracy matched to QCD PS
 - for both QCD-/EW-induced process
 - $\bullet\,$ all available in $\rm VBFNLO$
 - $\bullet\,$ all available in ${\rm POWHEG\text{-}BOX} \Rightarrow \mathsf{parton\text{-}shower}$ (PS) matching
 - often in VBS approximation (no int., s channel sometimes included)
 - \bullet possible to generate in $\rm MG5_AMC@NLO$ or $\rm Sherpa$
- NLO EW corrections known for $W^{\pm}W^{\pm}$, WZ, ZZ, and $W^{+}W^{-}$ with leptonic decays NLO EW matched to EW PS and interfaced to QCD PS for $W^{\pm}W^{\pm}$ in POWHEG-BOX-RES
- $\bullet\,$ full NLO computation only for $\mathrm{W^+W^+}$ and ZZ with leptonic decays
- no NLO results for hadronically decaying vector bosons
- no NLO results for polarised vector bosons
- no NNLO results known



Large NLO EW corrections to VBS processes

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process	$\sigma_{ m LO}^{{\cal O}(lpha^6)}$ [fb]	$\Delta \sigma_{ m NLO, EW}^{{\cal O}(lpha^7)}$ [fb]	$\delta_{\rm EW}$ [%]
Biedermann et al. 1708.00268	(Dittmaier et al.	2308.16716)	
$pp \rightarrow \mu^+ \nu_\mu e^+ \nu_e jj (W^+W^+)$	1.4178(2)	-0.2169(3)	-15.3
Denner et al. 1904.0088 pp $\rightarrow \mu^+ \mu^- e^+ \nu_e jj (ZW^+)$	0.25511(1)	-0.04091(2)	-16.0
Denner et al. 2009.00411 pp $\rightarrow \mu^+ \mu^- e^+ e^- jj$ (ZZ)	0.097681(2)	-0.015573(5)	-15.9
Denner et al. 2202.10844 pp $\rightarrow \mu^+ \mu^- e^+ e^- jj (W^+ W^-)$	2.6988(3)	-0.307(1)	-11.4

- EW corrections similar for all processes and rather independent of cuts ⇒ intrinsic feature of VBS process
- smaller corrections to W⁺W⁻ due to Higgs resonance in fiducial phase space (Higgs contribution about 25%, corresponding EW corrections -6.5%)
- NLO EW corrections to fiducial cross section well described by simple logarithmic approximation (Sudakov approximation $s, |t|, |u| \gg M_W^2$)
- NLO EW corrections to distributions not well described by Sudakov approximation

Julius-Maximilians-**UNIVERSITÄT** Distributions for pp $\rightarrow \mu^+ \mu^- e^+ \nu_{ejj}$ (ZW⁺jj)



Distribution in transverse momentum of the leading jet Denner et al. 1904.00882



- $\mathcal{O}(\alpha^7) \sim -30\%$ at $p_{T,i_1} = 800 \, \text{GeV}$ (Sudakov logarithms) dominant correction larger than QCD scale uncertainty
- $\mathcal{O}(\alpha_{\rm s}\alpha^6) \lesssim 10\%$ for $p_{T,i_1} > 100 \,\mathrm{GeV}$ small QCD scale uncertainty owing to suitable dynamical scale $\mu = \sqrt{p_{\mathrm{T},j_1} p_{\mathrm{T},j_2}}$
- large correction for small p_{T,i_1} due to phase-space suppression at LO (all jets have small $p_{\rm T}$) redistribution of events at NLO



Corrections of order $\mathcal{O}(\alpha_s^2 \alpha^5)$ normalised to $\mathcal{O}(\alpha_s^2 \alpha^4)$ ("EW corrections to LO QCD contribution")

• pp $\rightarrow \mu^+ \nu_\mu e^+ \nu_e jj$ (W⁺W⁺):

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- negligible: 0.2% for $M_{\rm j_1j_2}>500\,{\rm GeV}$, result of cancellations
- EW corrections to LO QCD $\mathcal{O}(\alpha_{\rm s}^2 \alpha^4)$: -12.3%
- QCD corrections to LO interference $\mathcal{O}(\alpha_{\rm s}\alpha^5)$: 12.5%
- LO interference reaches 30% of LO QCD contribution
- $\bullet\,$ LO QCD amounts to 12% of LO cross section
- pp $\rightarrow \mu^+ \mu^- e^+ e^- jj$ (ZZ):
 - sizeable: -10% (-8%) for $M_{j_1j_2} > 500 \,{
 m GeV}$ ($100 \,{
 m GeV}$)
 - EW corrections to LO QCD $\mathcal{O}\bigl(\alpha_{\rm s}^2\alpha^4\bigr):~-11.6\%~(-8.2\%)$
 - QCD corrections to LO interference $\mathcal{O}(\alpha_{\rm s}\alpha^5)$: 1.3% (0.3%)
 - LO interference only 0.8% (4%) of LO QCD contribution
 - LO QCD amounts to 60% (82%) of LO cross section

 $\Rightarrow \mathsf{EW} \ \mathcal{O}(\alpha_{\rm s}^2 \alpha^5) \text{ corrections generically large} (also seen in individual partonic processes)$





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Observables with polarised massive vector bosons

- are important probes of Standard Model gauge and Higgs sectors,
- may provide discrimination power between SM and beyond-SM physics.
- Longitudinal polarisation mode of vector bosons is
 - a consequence of the EW Symmetry Breaking
 - very sensitive to deviations from SM: unitarity of cross sections with longitudinally polarised vector bosons realised in SM via cancellation of different contributions.

Challenges and problems

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- Unstable massive vector bosons appear only as virtual particles \Rightarrow
 - no unique definition of vector-boson polarisations for off-shell bosons
 - diagrams without resonant vector bosons contribute to physical final state



vector bosons are massive ⇒
 definition of polarisation depends on reference frame



Idea: use pole expansion to extract resonant (vector-boson) contributions in gauge-invariant way Ballestrero, Maina, Pelliccioli 1710.09339, 1907.04722

formulation for NLO developed by Denner, Pelliccioli 2006.14867

 not all diagrams involve required resonances resonant diagrams
 non-resonant diagrams



• split full matrix element into resonant part and non-resonant part using pole expansion (gauge-invariant)

$$\begin{aligned} \mathcal{A} &= \frac{R(k^2)}{k^2 - M^2 + iM\Gamma} + N(k^2) \\ &= \frac{R(M^2)}{k^2 - M^2 + iM\Gamma} + \frac{R(k^2) - R(M^2)}{k^2 - M^2} + N(k^2) = \mathcal{A}_{\text{res}} + \mathcal{A}_{\text{nonres}} \end{aligned}$$

• consider non-resonant part as irreducible background: no resonance

• define polarisation for on-shell residue $R(M^2)$ (gauge invariant)

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Definition of polarisation based on pole approximation II



Separate polarisation modes of resonant amplitude

split propagator numerator of resonant particle



$$\begin{split} \mathcal{A}_{\rm res} &= \mathcal{P}_{\mu} \, \frac{-g^{\mu\nu}}{k^2 - M_{\rm W}^2 + \mathrm{i}\Gamma_{\rm W}M_{\rm W}} \, \mathcal{D}_{\nu} = \mathcal{P}_{\mu} \, \frac{\sum_{\lambda} \varepsilon_{\lambda}^{\mu*}(k)\varepsilon_{\lambda}^{\nu}(k)}{k^2 - M_{\rm W}^2 + \mathrm{i}\Gamma_{\rm W}M_{\rm W}} \, \mathcal{D}_{\nu} \\ &= \sum_{\lambda=\mathrm{L},\pm} \, \frac{\mathcal{M}_{\lambda}^{\mathrm{prod}}\,\mathcal{M}_{\lambda}^{\mathrm{dec}}}{k^2 - M_{\rm W}^2 + \mathrm{i}\Gamma_{\rm W}M_{\rm W}} =: \sum_{\lambda=\mathrm{L},\pm} \mathcal{A}_{\lambda} \,, \\ \mathcal{A}_{\mathrm{res}} \Big|^2 &= \sum_{\lambda} \left| \mathcal{A}_{\lambda} \right|^2 + \sum_{\lambda\neq\lambda'} \mathcal{A}_{\lambda}^* \, \mathcal{A}_{\lambda'} \end{split}$$

• incoherent sum $\sum_{\lambda} |A_{\lambda}|^2$: $|A_{\lambda}|^2 \propto$ "polarised cross sections", "polarisation fractions": $f_{\lambda} = \frac{|A_{\lambda}|^2}{\sum_{\lambda} |A_{\lambda}|^2}$

• interferences $\sum_{\lambda \neq \lambda'} A_{\lambda}^* A_{\lambda'}$ vanish for observables fully inclusive in decay products in general to be considered as extra background

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Method

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- applicable to arbitrary processes and multiple resonances at LO, NLO QCD and EW and beyond
- needs pole approximation/ pole expansion for all NLO contributions including subtraction terms ⇒ technical complication
- implementation for processes with two polarised vector bosons complete in double-pole approximation (DPA) for NLO QCD and EW
- use of spin-correlated narrow-width approximation (NWA) also possible (simpler but potentially less precise)

Method allows to separate

- polarised cross sections in arbitrary frames
- interference contributions between polarisations
- irreducible background.

All contributions that are not small need to be taken into account!



Fixed-order results at (N)NLO

- results at LO for VBS for ss-WW, WZ, ZZ, os-WW Ballestrero, Maina, Pelliccioli '17, '19, '20 [PHANTOM]
- results at NLO QCD for

• pp $\rightarrow jj\ell^+\ell^-$ (W⁺Z)

- pp $\rightarrow \mu^+ \nu_\mu e^+ \nu_e (W^+ W^-)$ Denner, Pelliccioli 2006,14867 • pp $\rightarrow \mu^+ \mu^- e^+ \nu_e$ (W⁺Z)
 - Denner, Pelliccioli 2010.07149
 - Denner, Haitz, Pelliccioli '22
- results at NLO EW for (diboson production)
 - pp $\rightarrow \mu^+ \mu^- e^+ e^-$ (ZZ) Denner, Pelliccioli 2107.06579
 - pp $\rightarrow \mu^+ \mu^- e^+ \nu_e$ (W⁺Z) Baglio, Dao, Le 2203.01470, 2208.09232
 - pp $\rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu$ (WW) Denner, Pelliccioli 2311.16031, Dao, Le 2311.17027
- results at NNLO QCD for
 - pp $\rightarrow \mu^+ \nu_\mu e^+ \nu_e$ (W⁺W⁻) (DPA and NWA) Poncelet, Popescu 2102.13583

Implementation in Monte Carlo generators

- MADGRAPH5_AMC@NLO: spin-correlated narrow-width approximation (NWA), LO Franzosi, Mattelaer, Ruiz, Shil 1912.01725
- SHERPA: approximate NLO QCD (NWA) Hoppe, Schönherr, Siegert 2310.14803
- POWHEG-BOX-RES: for diboson processes at NLO QCD

Pelliccioli, Zanderighi 2311.05220



Massive diboson production with leptonic decays within the SM

- \bullet NLO QCD corrections for all processes $\mathrm{W^+W^-},\,\mathrm{WZ},\,\mathrm{ZZ}$
- \bullet PS matching available via $\operatorname{POWHEG-BOX-RES}$
- $\bullet\,$ NNLO QCD results for $\rm W^+W^-$
- \bullet NLO EW corrections for all processes $\mathrm{W^+W^-},\,\mathrm{WZ},\,\mathrm{ZZ}$

Massive diboson production with semi-leptonic decays within the SM

 \bullet NLO QCD results for ${\rm ZW} \to \ell^+ \ell^- j j$

Massive VBS within the SM

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- LO results exist within PHANTOM
- first NLO QCD+EW results being calculated

Implementation in Monte Carlo generators

- $Madgraph5_AMC@NLO:$ spin-correlated NWA at LO
- SHERPA: approximate NLO QCD in NWA



$\mathrm{pp} \rightarrow \mathrm{e}$	$e^+ \nu_e \mu^-$	$\bar{\nu}_{\mu}$	(WW	'):
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Denner, Haitz, Pelliccioli 2311.16031

state	$\sigma_{ m LO}$ [fb]	$\sigma_{ m NLOEW}$ [fb]	$\delta_{\rm EW}[\%]$	$f_{\rm NLOEW}[\%]$		
${ m b}ar{ m b}$ included, $\gamma{ m b},\gammaar{ m b}$ excluded						
full	259.02(2)	253.95(9)	-1.96	103.4		
unp.	249.97(2)	245.49(2)	-1.79	100.0		
LL	21.007(2)	20.663(2)	-1.64	8.4		
LT	33.190(3)	33.115(3)	-0.23	13.5		
TL	34.352(5)	34.230(5)	-0.35	13.9		
ΤT	182.56(2)	178.21(3)	-2.38	72.6		
int.	-21.14(5)	-20.6(2)	-2.45	-8.4		

- \bullet irreducible background (3.4%) consistent with DPA accuracy
- ullet sizeable interferences (-8.4%) from $p_{\rm T}$ cuts on charged leptons
- NLO EW corrections differ for various polarised and unpolarised cross sections
- expect larger differences for VBS





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Existing results for EW corrections

- no public codes except for $W^{\pm}W^{\pm}$ within POWHEG-BOS-RES NLO EW matched to EW PS and interfaced to QCD PS Chiesa, Denner, Lang, Pellen 1906.01863
- fiducial cross sections and various distributions for specific setups (one-dimensional histograms)
- for other setups obtainable with some numerical efforts
- can be used to reweight Monte Carlo events (as has been done in the past)

Long-term goal (not available soon)

- Les Houches events including EW corrections
- suitable for matching to (e.g. PYTHIA) parton shower
- public codes





QCD corrections

- Standard procedure: 7-point scale variations to estimate uncertainty owing to missing higher-order corrections
- conventional, simple, easy to use, but not reliable

EW corrections

- no standard method to estimate missing higher-order corrections
- available approaches
 - differences between renormalisation schemes, e.g. $\alpha(M_{\rm Z})$ versus $\alpha_{G_{\mu}}$
 - square of relative NLO correction as uncertainty
 - resummation in Sudakov approximation as uncertainty estimate
 - counting of powers of couplings and logarithms
 - use process-specific combination of estimates
- sensible estimate needs agreement of experts ⇒ task for LHC EW WG?

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Other sources of errors

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- use of approximations
 - VBS approximation
 - pole approximation
 - NWA
- QED final-state radiation
- photon-induced contributions

Estimates from comparison of results of different codes

- \bullet done for unpolarised W^+W^+ scattering at NLO QCD within VBSCAN COST network Ballestrero et al. 1803.07943
- ongoing for polarised ZZ production within COMETA WG1 https://foswiki.web.cern.ch/COMETA/WkG1





• check of our results by other codes

done for W⁺W⁺ scattering by Freiburg group Dittmaier et al. 2308.16716 bug with minor impact found $[0.005 \,\text{fb} \text{ in } \mathcal{O}(\alpha_s^2 \alpha^5)]$

- NLO corrections to polarised VBS \Rightarrow in the making for W^+W^+ needed for other channels?
- estimate of theoretical uncertainty \Rightarrow see above
- resummation of EW logarithms \Rightarrow needed for LHC?

more relevant for higher energies

matching of EW corrections to parton showers (including EW effects)
 ⇒ work ongoing by different groups

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Higher-order (EW) corrections in VBS

- NLO QCD matched to QCD PS available in Monte Carlo generators
- results for EW corrections exist in form of one-dimensional histograms
- EW corrections typically at the level of 10–20% similarly large for irreducible background
- estimate of theoretical uncertainties to be developed

Polarised multi-boson processes

- MADGRAPH5_AMC@NLO: spin-correlated NWA at LO
- SHERPA: approximate NLO QCD in NWA
- diboson production with leptonic decays
 - $\bullet~$ NLO QCD and matching to PS via $\operatorname{Powheg-Box-Res}$
 - EW corrections in form of one-dimensional histograms
 - NLO (EW) corrections depend on boson polarisations
- VBS
 - LO results exist within PHANTOM
 - NLO QCD+EW results being calculated.





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NLO QCD and EW corrections to ZZ production in VBS





- Loose VBS cut: $M_{j_1j_2} > 100 \,\text{GeV}$ based on 1708.02812 (CMS)
- 24% NLO QCD corrections to fiducial cross section
- *s*-channel NLO contribution involving tri-boson production



Less suppression at NLO owing to extra gluon jet

⇒ include tri-boson contribution (s-channel) and interferences for loose VBS cuts

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Source of large EW corrections

AP2

Double-pole approximation (DPA) for outgoing W bosons effective vector-boson approximation (EVBA) for incoming W bosons

- DPA and EVBA reduce discussion to $V_1V_2 \rightarrow V_3V_4$
- $\bullet\,$ DPA accurate for cross section within 1%
- EVBA crude approximation ($\sim 50\%$) Kuss, Spiesberger '96, Dittmaier et al. '23 sufficient to understand dominant effects



high-energy, logarithmic approximation for $V_1V_2 \rightarrow V_3V_4$

Denner, Pozzorini '00

$$d\sigma_{\rm LL} = d\sigma_{\rm LO} \left[1 - \frac{\alpha}{4\pi} 4C_{\rm W}^{\rm EW} \log^2 \left(\frac{Q^2}{M_{\rm W}^2} \right) + \frac{\alpha}{4\pi} 2b_{\rm W}^{\rm EW} \log \left(\frac{Q^2}{M_{\rm W}^2} \right) \right]$$
$$C_{\rm W}^{\rm EW} = \frac{2}{s_{\rm w}^2}, \quad b_{\rm W}^{\rm EW} = \frac{19}{6s_{\rm w}^2} \quad \text{for transverse W bosons,} \quad Q \to M_{4\ell}$$

(double EW logs, collinear single EW logs, and single logs from parameter renormalisation included) (angular-dependent logarithms omitted, $\log \frac{t}{u} \log \frac{Q}{M_W}$)

large NLO EW corrections intrinsic feature of VBS

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AP

Simple formula for total cross section

$$d\sigma_{\rm LL} = d\sigma_{\rm LO} \left[1 - \frac{\alpha}{4\pi} 4 C_{\rm W}^{\rm EW} \log^2 \left(\frac{Q^2}{M_{\rm W}^2} \right) + \frac{\alpha}{4\pi} 2 b_{\rm W}^{\rm EW} \log \left(\frac{Q^2}{M_{\rm W}^2} \right) \right]$$

process	$\delta_{\rm EW}$ [%]	$\delta_{\rm EW}^{\rm log,int}$ [%]	$\delta_{\rm EW}^{\rm log, diff}$ [%]	$\langle M_{4\ell} \rangle$ [GeV]
$pp \rightarrow \mu^+ \nu_\mu e^+ \nu_e jj$	-16.0	-16.1	-15.0	390
$pp \rightarrow \mu^+ \mu^- e^+ \nu_e jj$	-16.0	-17.5	-16.4	413
$pp \rightarrow \mu^+ \mu^- e^+ e^- jj$	-15.9	-15.8	-14.8	385

- surprisingly good agreement with complete calculation
- large EW corrections are due to large gauge couplings of vector bosons ($C^{\rm EW}$) and large scale $Q \sim \langle M_{4\ell} \rangle \sim 400 \, {\rm GeV}$
- angular-dependent logarithms different for different processes $\sim 1{-}2\%$ owing to cancellations

large NLO EW corrections intrinsic feature of VBS

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Process $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu jj$

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Denner, Franken, Schmidt, Schwan 2202.10844

- EW corrections smaller than for other VBS processes
- fiducial phase space contains Higgs resonance

setup	$\sigma_{ m LO}^{{\cal O}(lpha^6)}$ [fb]	$\Delta \sigma_{ m NLO, EW}^{{\cal O}(lpha^7)}$ [fb]	$\delta_{\rm EW}$ [%]
VBS setup	2.6988(3)	-0.307(1)	-11.4
VBS setup, Higgs cut out	1.6117(2)	-0.239(2)	-14.8
		0.060(0)	10.0
VBS setup, Higgs cut out	1.9750(2)	-0.260(2)	-13.2
VBS setup, Higgs contribution	0.7238(2)	-0.047(2)	-6.5
Higgs setup	1.5322(2)	-0.103(1)	-6.7

- $\bullet~{\rm EW}$ corrections to generic VBS contributions $\sim -15\%$
- $\bullet~{\rm EW}$ corrections to Higgs resonance contribution $\sim -6.5\%$
- Higgs cut: $|M_{4\ell} M_{\rm H}| > 20\Gamma_{\rm H} \approx 80 \,{\rm MeV}$, removes 98.4% of resonance
- Higgs contribution: $|M_{4\ell} M_{\rm H}| < 20\Gamma_{\rm H}$
- Higgs setup: cuts inspired by CMS Higgs search CMS 1806.05246



Biedermann, Denner, Pellen 1708.00268

PDFs

NNPDF3.0QED Ball et al. '13, '14 factorisation and renormalisation scales: $\mu_{\rm F} = \mu_{\rm R} = \sqrt{p_{\rm T,j_1}p_{\rm T,j_2}}$

Recombination / jet clustering

Anti- $k_{\rm T}$ algorithm with R = 0.4 Cacciari, Salam, Soyez '08 recombination of photons with charged partons with R = 0.1

Cuts: based on ATLAS 1405.6241, 1611.02428 and CMS 1410.6315

 $\begin{array}{ll} p_{\mathrm{T},j} > 30 \, \mathrm{GeV}, & |y_{j}| < 4.5, \quad \Delta R_{j\ell} > 0.3 \\ p_{\mathrm{T},\ell} > 20 \, \mathrm{GeV}, & |y_{\ell}| < 2.5, \quad \Delta R_{\ell\ell} > 0.3, \quad \Delta R_{ij} = \sqrt{(\Delta y_{ij})^{2} + (\Delta \phi_{ij})^{2}} \\ E_{\mathrm{T,miss}} > 40 \, \mathrm{GeV} \\ M_{jj} > 500 \, \mathrm{GeV}, & |\Delta y_{jj}| > 2.5 \quad (\mathsf{VBF cuts}) \\ \mathrm{require} \geq 2 \text{ jets}, 2 \text{ same-sign leptons and missing energy} \end{array}$





Energy: 13 TeV (14 TeV)

Denner et al. 1904.00882

PDFs

NNPDF3.1QED Ball et al. '14, Bertone et al. '17 factorisation and renormalisation scales: $\mu_{\rm F} = \mu_{\rm R} = \sqrt{p_{\rm T,j_1} p_{\rm T,j_2}}$

Recombination / jet clustering

Anti- $k_{\rm T}$ algorithm with R=0.4 $\,$ Cacciari, Salam, Soyez '08 recombination of photons with charged partons with R=0.4

Cuts: loose fiducial region of CMS 1901.04060

 $\begin{array}{ll} p_{\mathrm{T},j} > 30 \, \mathrm{GeV}, & |y_{j}| < 4.7, \quad \Delta R_{\mathrm{j}\ell} > 0.4 & \Delta R_{ij} = \sqrt{(\Delta y_{ij})^{2} + (\Delta \phi_{ij})^{2}} \\ p_{\mathrm{T},\ell} > 20 \, \mathrm{GeV}, & |y_{\ell}| < 2.5, \quad M_{3\ell} > 100 \, \mathrm{GeV}, \quad M_{\ell\ell} > 4 \, \mathrm{GeV} \\ |M_{\mu^{+}\mu^{-}} - M_{\mathrm{Z}}| < 15 \, \mathrm{GeV} & \\ M_{\mathrm{j}j} > 500 \, \mathrm{GeV}, & |\Delta y_{\mathrm{j}j}| > 2.5 & (\mathsf{VBF cuts}) \\ \mathrm{require} \geq 2 \, \mathrm{jets}, \, 3 \, \mathrm{leptons} \end{array}$





Denner, Franken, Pellen, Schmidt 2009.00411

PDFs

NNPDF3.1QED Ball et al. '14, Bertone et al. '17 factorisation and renormalisation scales: $\mu_{\rm F} = \mu_{\rm R} = \sqrt{p_{\rm T,j_1} p_{\rm T,j_2}}$

Recombination / jet clustering

Cuts: inspired by CMS 1708.02812

 $\begin{array}{ll} p_{\mathrm{T},j} > 30 \, \mathrm{GeV}, & |y_{j}| < 4.7, \quad \Delta R_{j\ell} > 0.4 \qquad \Delta R_{ij} = \sqrt{(\Delta y_{ij})^{2} + (\Delta \phi_{ij})^{2}} \\ p_{\mathrm{T},\ell} > 20 \, \mathrm{GeV}, & |y_{\ell}| < 2.5, \quad \Delta R_{\ell\ell'} > 0.05, \quad M_{\ell^{+}\ell'^{-}} > 4 \, \mathrm{GeV} \\ 60 \, \mathrm{GeV} < M_{\ell^{+}\ell^{-}} < 120 \, \mathrm{GeV} \\ \text{inclusive setup: } M_{jj} > 100 \, \mathrm{GeV}, \qquad \mathsf{VBS \ setup} \ M_{jj} > 500 \, \mathrm{GeV} \\ \text{require} \geq 2 \ \text{jets}, \ 4 \ \text{leptons} \end{array}$





Denner, Franken, Schmidt, Schwan 2202.10844

PDFs

NNPDF3.1QED Ball et al. '14, Bertone et al. '17 factorisation and renormalisation scales: $\mu_{\rm F} = \mu_{\rm R} = \sqrt{p_{\rm T,j_1}p_{\rm T,j_2}}$

Recombination / jet clustering

Cuts: similar to CMS 2205.05711

 $\begin{array}{ll} p_{\mathrm{T},j} > 30 \, \mathrm{GeV}, & |y_{\mathrm{j}}| < 4.5, \quad \Delta R_{\mathrm{j}\ell} > 0.4 \quad \Delta R_{ij} = \sqrt{(\Delta y_{ij})^2 + (\Delta \phi_{ij})^2} \\ p_{\mathrm{T},\ell} > 25 \, \mathrm{GeV}, & |y_\ell| < 2.4, \quad p_{\mathrm{T},\ell^+\ell^-} > 30 \, \mathrm{GeV}, \quad M_{\ell^+\ell^-} > 20 \, \mathrm{GeV} \\ p_{\mathrm{T},\mathrm{miss}} > 20 \, \mathrm{GeV} & \\ M_{\mathrm{j}j} > 500 \, \mathrm{GeV}, \quad \Delta y_{\mathrm{j}j} > 2.5 \\ \mathrm{require} \ge 2 \, \mathrm{jets}, \, 2 \, \mathrm{leptons} \end{array}$





Denner, Franken, Schmidt, Schwan 2202.10844

PDFs

NNPDF3.1QED Ball et al. '14, Bertone et al. '17 factorisation and renormalisation scales: $\mu_{\rm F} = \mu_{\rm R} = \sqrt{p_{\rm T,j_1} p_{\rm T,j_2}}$

Recombination / jet clustering

Anti- $k_{\rm T}$ algorithm with R=0.4 $\,$ Cacciari, Salam, Soyez '08 recombination of photons with charged partons with R=0.4

Cuts: following CMS 1806.05246 (Higgs search)

$$\begin{array}{ll} p_{\mathrm{T},j} > 30 \, \mathrm{GeV}, & |y_j| < 4.7, \quad \Delta R_{j\ell} > 0.4 \qquad \Delta R_{ij} = \sqrt{(\Delta y_{ij})^2 + (\Delta \phi_{ij})^2} \\ p_{\mathrm{T},\ell}^{\mathrm{lead}} > 25 \, \mathrm{GeV}, & p_{\mathrm{T},\ell}^{\mathrm{trail}} > 10 \, \mathrm{GeV}, \quad |y_\ell| < 2.4, \\ p_{\mathrm{T},\ell+\ell^-} > 30 \, \mathrm{GeV}, & M_{\ell+\ell^-} > 12 \, \mathrm{GeV}, \quad \Delta R_{\ell+\ell^-} > 0.4 \\ p_{\mathrm{T},\mathrm{miss}} > 20 \, \mathrm{GeV}, & 60 \, \mathrm{GeV} < M_{\mathrm{T},\ell+\ell^-,\mathrm{miss}} < 125 \, \mathrm{GeV} \\ M_{j_1j_2} > 400 \, \mathrm{GeV}, & \Delta y_{j_1j_2} > 3.5, \quad p_{\mathrm{T},j_3} < 30 \, \mathrm{GeV} \text{ (jet veto on 3rd jet)} \\ |z_{\ell j_1j_2}| < 0.5, \quad z_{\ell j_1j_2} = \frac{2y_\ell - y_{j_1} - y_{j_2}}{2|y_{j_1} - y_{j_2}|}, \quad \mathrm{require 2 \, jets, 2 \, leptons} \end{array}$$



$$pp \rightarrow e^+ e^- \mu^+ \mu^-$$
 (ZZ):

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Denner, Pelliccioli 2107.06579

mode	$\sigma_{ m LO}$ [fb]	$\delta_{ m QCD}$	δ_{EW}	$\delta_{ m gg}$	$\sigma_{ m NLO_+}$ [fb]
full	$11.1143(5)^{+5.6\%}_{-6.8\%}$	+34.9%	-11.0%	+15.6%	$15.505(6)^{+5.7\%}_{-4.4\%}$
unpol.	$11.0214(5)^{+5.6\%}_{-6.8\%}$	+35.0%	-10.9%	+15.7%	$15.416(5)^{+5.7\%}_{-4.4\%}$
$\rm Z_L \rm Z_L$	$0.64302(5)^{+6.8\%}_{-8.1\%}$	+35.7%	-10.2%	+14.5%	$0.9002(6)^{+5.5\%}_{-4.3\%}$
$\rm Z_L \rm Z_T$	$1.30468(9)^{+6.5\%}_{-7.7\%}$	+45.3%	-9.9%	+2.8%	$1.8016(9)^{+4.3\%}_{-3.5\%}$
$Z_{\rm T} Z_{\rm L}$	$1.30854(9)^{+6.5\%}_{-7.7\%}$	+44.3%	-9.9%	+2.8%	$1.7933(9)^{+4.3\%}_{-3.4\%}$
$Z_{\rm T} Z_{\rm T}$	$7.6425(3)^{+5.2\%}_{-6.4\%}$	+31.2%	-11.2%	+20.5%	$10.739(4)^{+6.2\%}_{-4.7\%}$

- small irreducible background (0.5%) and interferences (1.2%)
- sizeable QCD and EW corrections
- substantial contribution from loop-induced gg fusion for LL and TT
- polarisation fractions roughly conserved by NLO corrections owing to cancellations

