PDF ISSUES: VIEWPOINT FROM THE EXPERIMENTS

LHC EW WG GENERAL MEETING

JUNE 11th, 2024

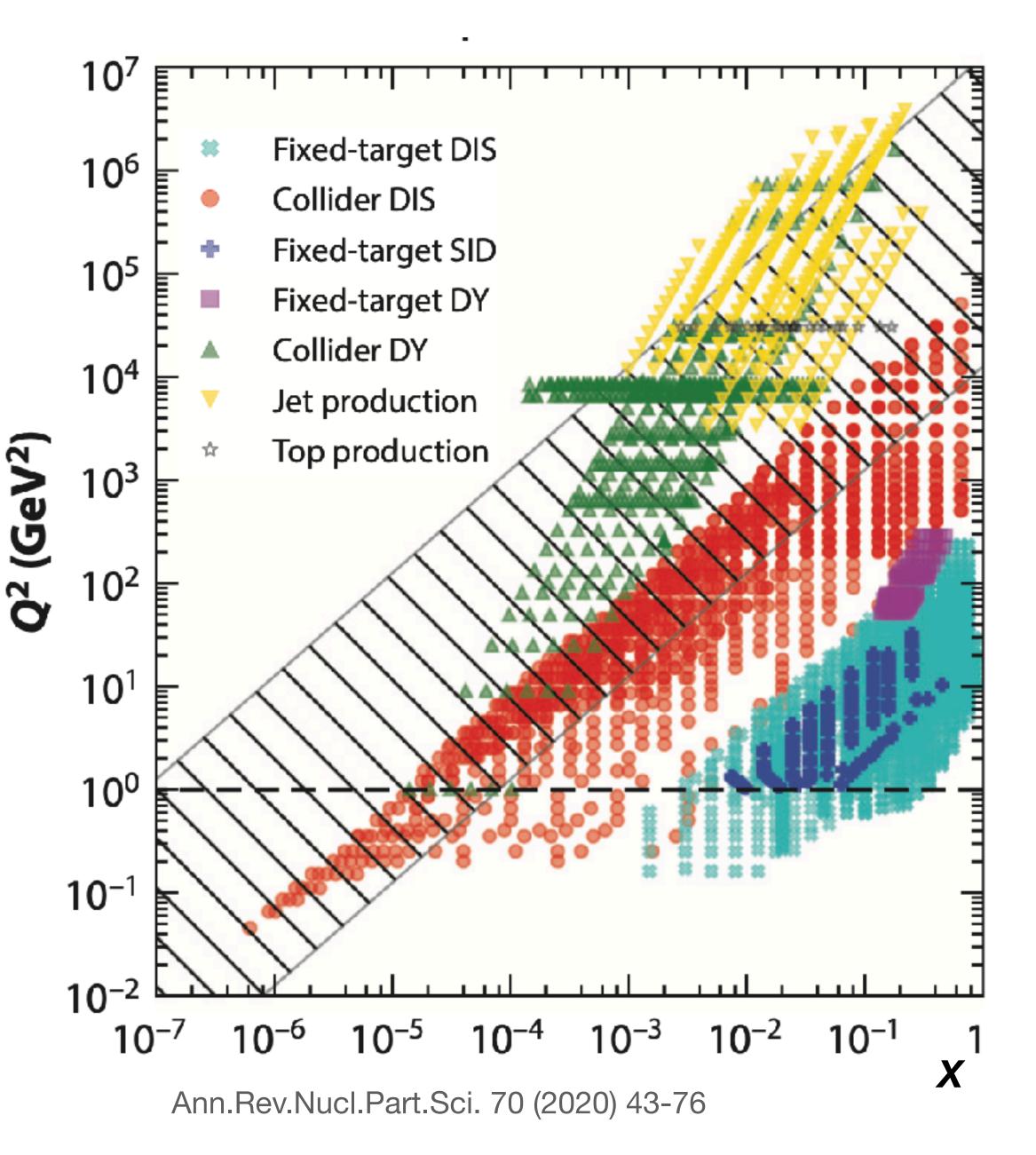
SIMONE AMOROSO (DESY)

PARTON DISTRIBUTION FUNCTIONS

O Predictions at a hadron collider require knowledge of the proton structure

$$\sigma = \sum_{i,j} \int dx_1 dx_2 f_{i/p}(x_1) f_{j/p}(x_2) \,\hat{\sigma}(x_1 x_2 s)$$

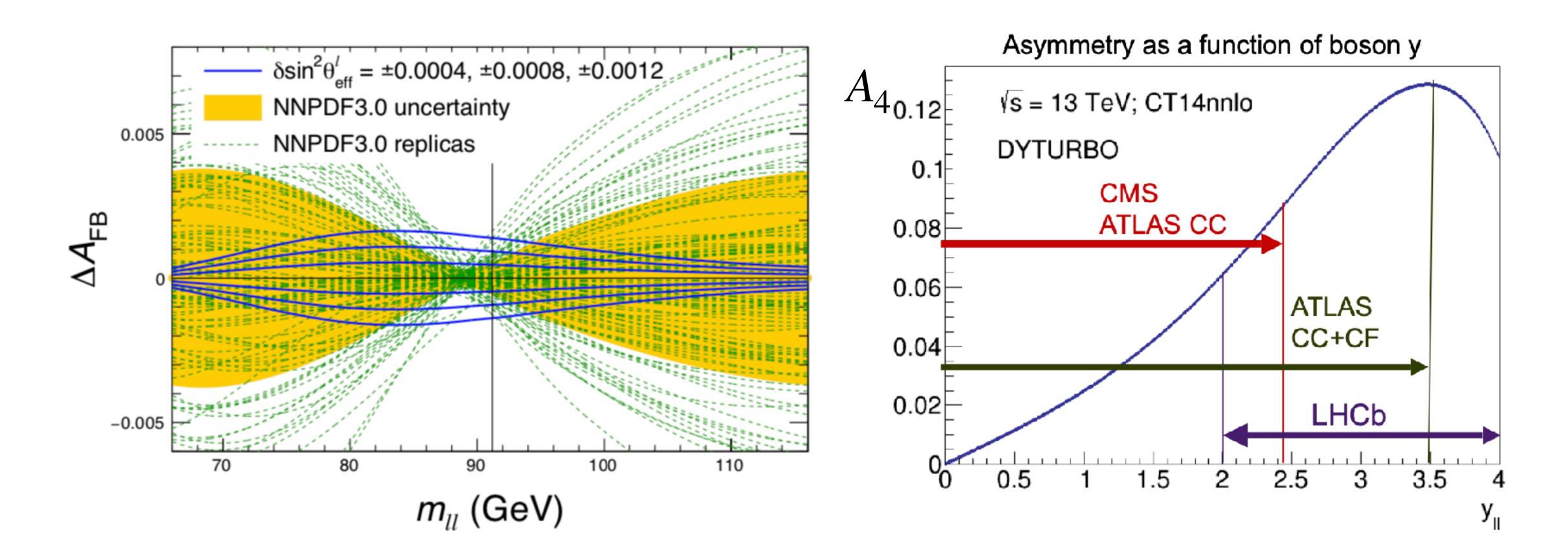
- O Cross-sections calculated as convolution of short-distance cross-sections with Parton Distribution Functions (PDFs)
- O A universal quantity, PDFs are inferred from a given set of measurements and can be used to predict any cross-section





$\sin^2 \theta_{\rm eff}^l$ at the LHC and PDFs

- Ο
- Ο



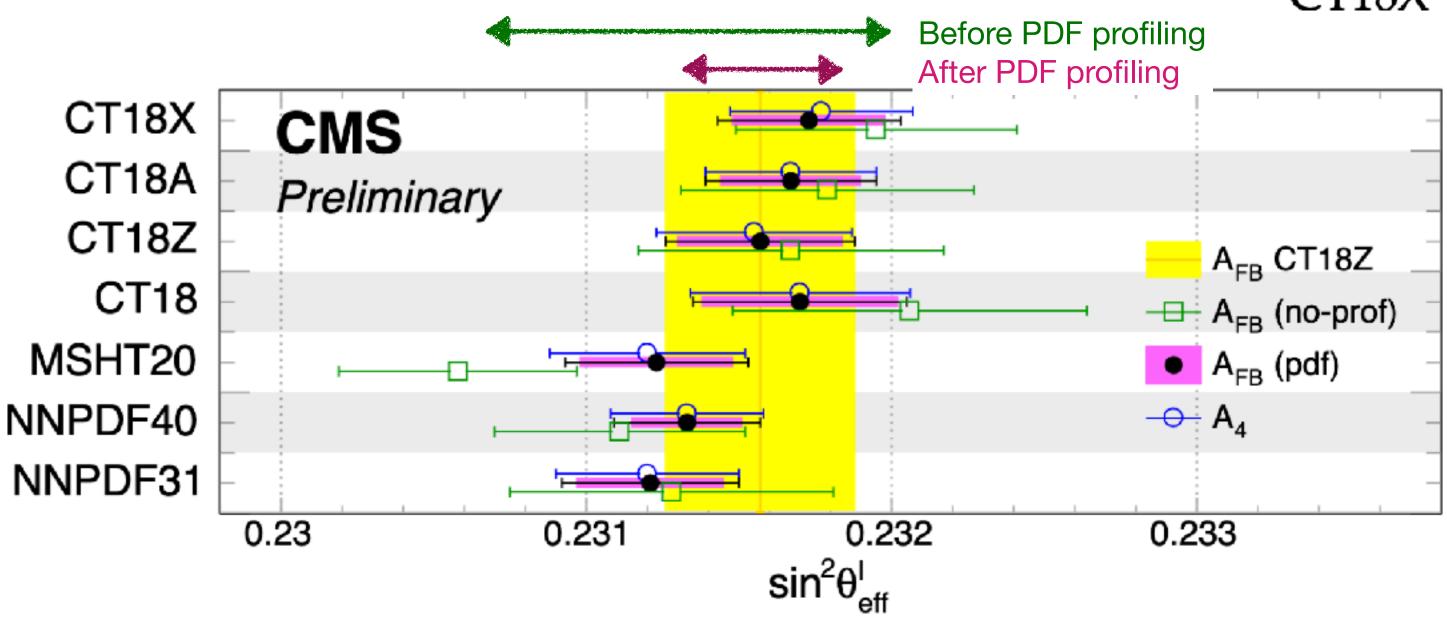
Weak mixing angle extracted from polarization asymmetries in Drell-Yan

Parton-level effect measured at particle level -> dependence on PDFs Direction of incoming quark/anti-quark inferred from Z rapidity sign



PDF uncertainties on $\sin^2 \theta_{\rm eff}^l$ - CMS13 TeV

- All PDF sets provide an equally Ο good description of the data
- PDF spread and uncertainties Ο reduced in the fit



SMP-22-010

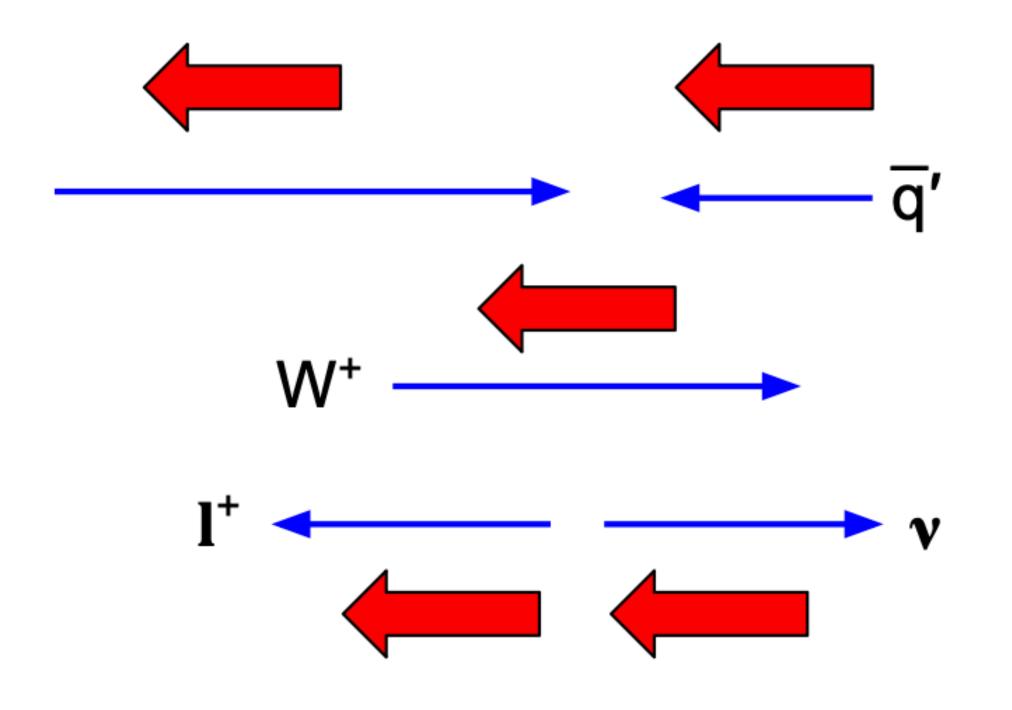
PDF	A _{FB} (816 bins)		A_4 (63 bins)	
	$\chi^2_{ m min}$	$\sin^2 heta_{ m eff}^\ell$	$\chi^2_{ m min}$	$\sin^2 heta_{ m eff}^\ell$
NNPDF31	724.7	23121 ± 29	58.5	23120 ± 30
NNPDF40	730.5	23133 ± 24	62.6	23133 ± 25
MSHT20	735.8	23123 ± 30	71.0	23120 ± 32
CT18	728.4	23170 ± 35	62.2	23170 ± 36
CT18Z	730.7	23157 ± 31	61.3	23155 ± 32
CT18A	730.3	23167 ± 28	63.6	23167 ± 28
CT18X	728.5	23173 ± 30	61.8	23177 ± 30

- **O** $\sin^2 \theta_{\text{eff}}^l$ values with different PDFs consistent at the ~1 sigma level
 - Use CT18Z as covering the central values obtained with the other sets



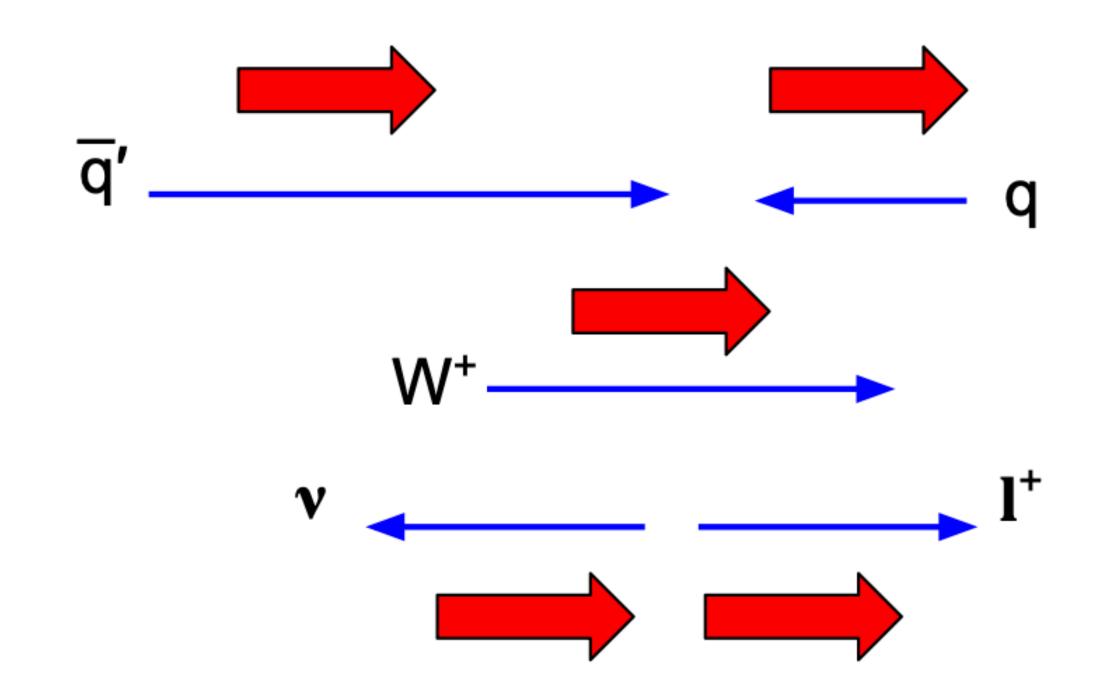
PDFS AND W-HELICITIES

- Ο
- Ο lepton



Effect induced by PDFs, important uncertainty in W-mass determinations \mathbf{O}

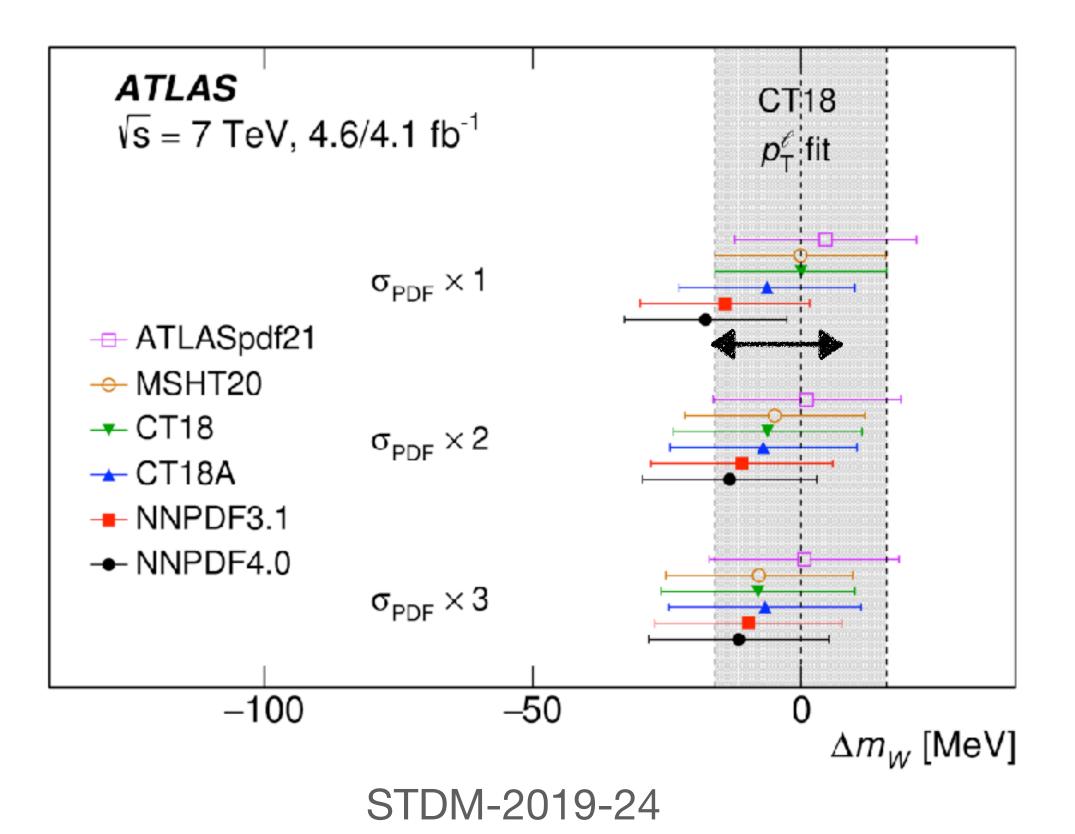
Lepton direction in W-boson decays retains information on the boson polarization Left-handed couplings of the W correlate polarization and rapidity of the boson with the direction of the quark/anti-quark, and hence the direction of the outgoing





Large PDF dependence: NNPDF4.0 and CT18NNLO differ by 21 MeV Ο

PDF uncertainties smaller and at the 3-9 MeV level Ο



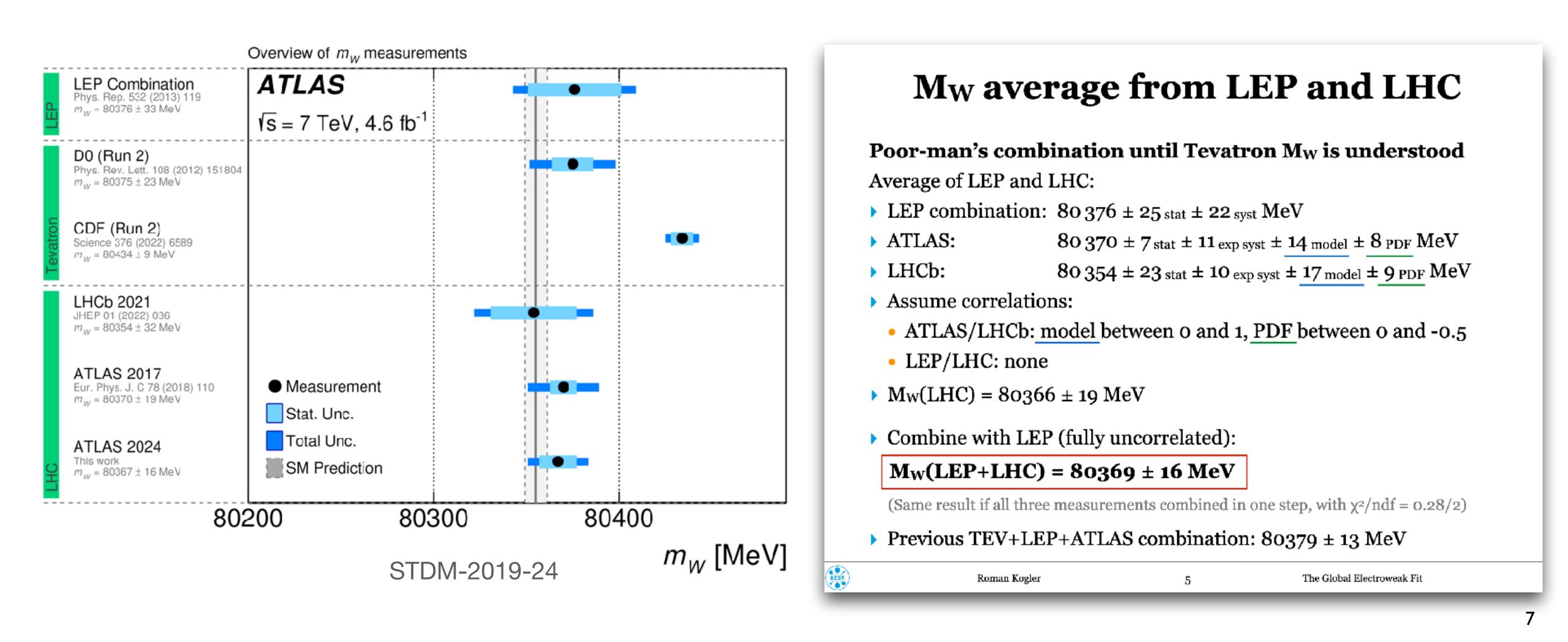
PDFS IN THE ATLAS 7 TEV m_W

PDF set	Combined m _W [MeV]
CT14	80363.6 ± 15.9
CT18	80366.5 ± 15.9
CT18A	80357.2 ± 15.6
MMHT2014	80366.2 ± 15.8
MSHT20	80359.3 ± 14.6
ATLASpdf21	80367.6 ± 16.6
NNPDF3.1	80349.6 ± 15.3
NNPDF4.0	80345.6 ± 14.9



PDFs in m_W combinations

O Different hadron collider measurements are correlated through the PDFsO How much? Measurements all made with very different sets

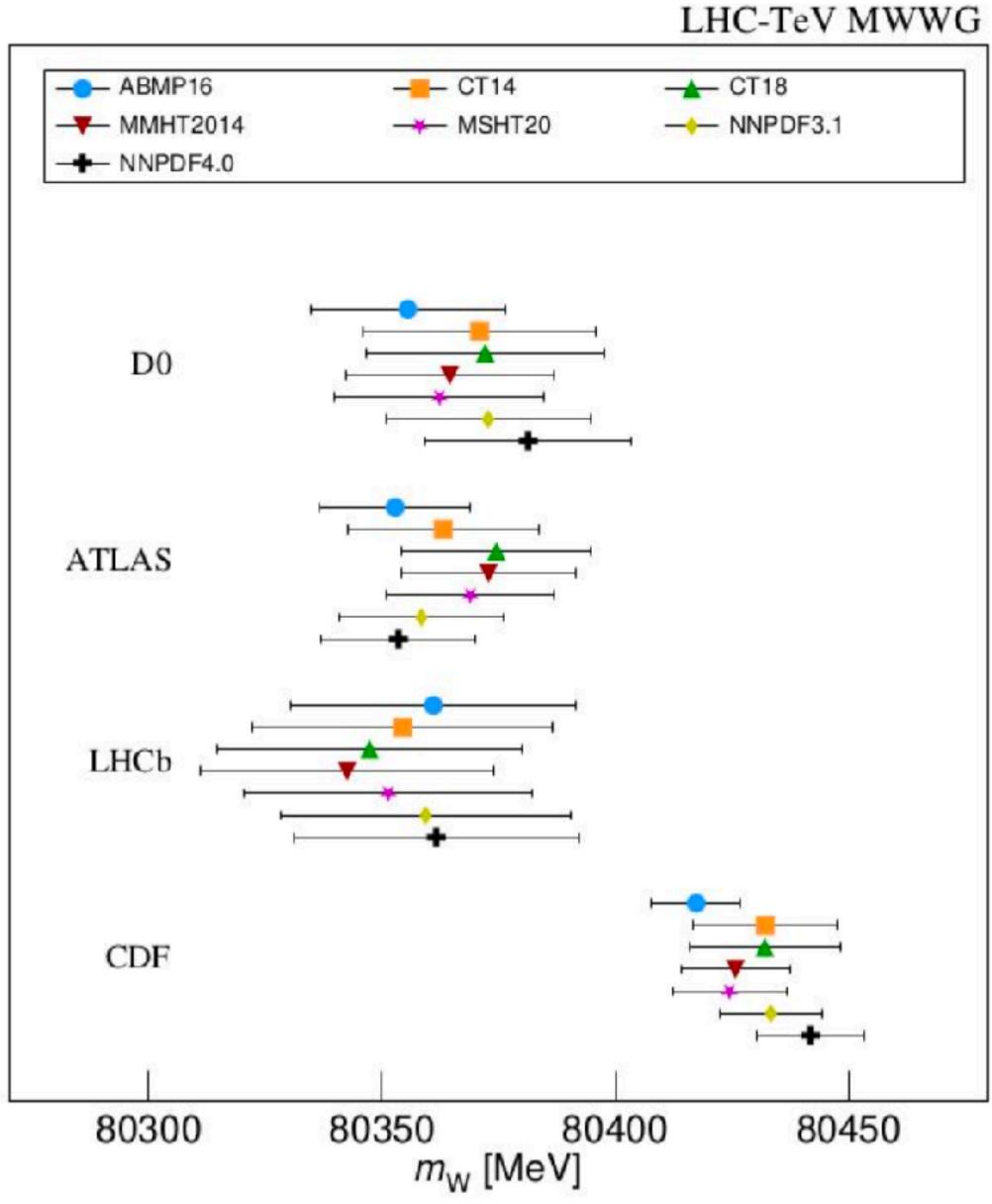


PDFs in the Tevatron/LHC m_W combination

- Combination of hadron collider Ο W-boson mass measurements
- **O** Measurements corrected to the same PDF before averaging
- Different PDF choices evaluated
- Variations larger than the PDF Ο uncertainties and often as large as the total uncertainty

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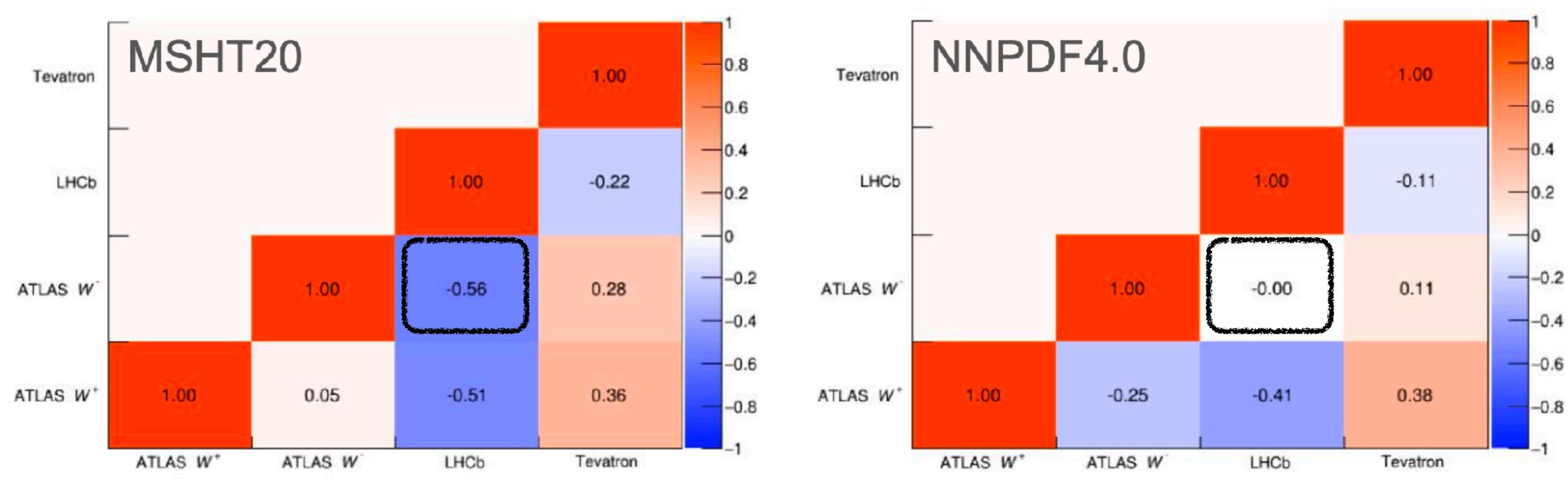






PDFs in the Tevatron/LHC m_W combination

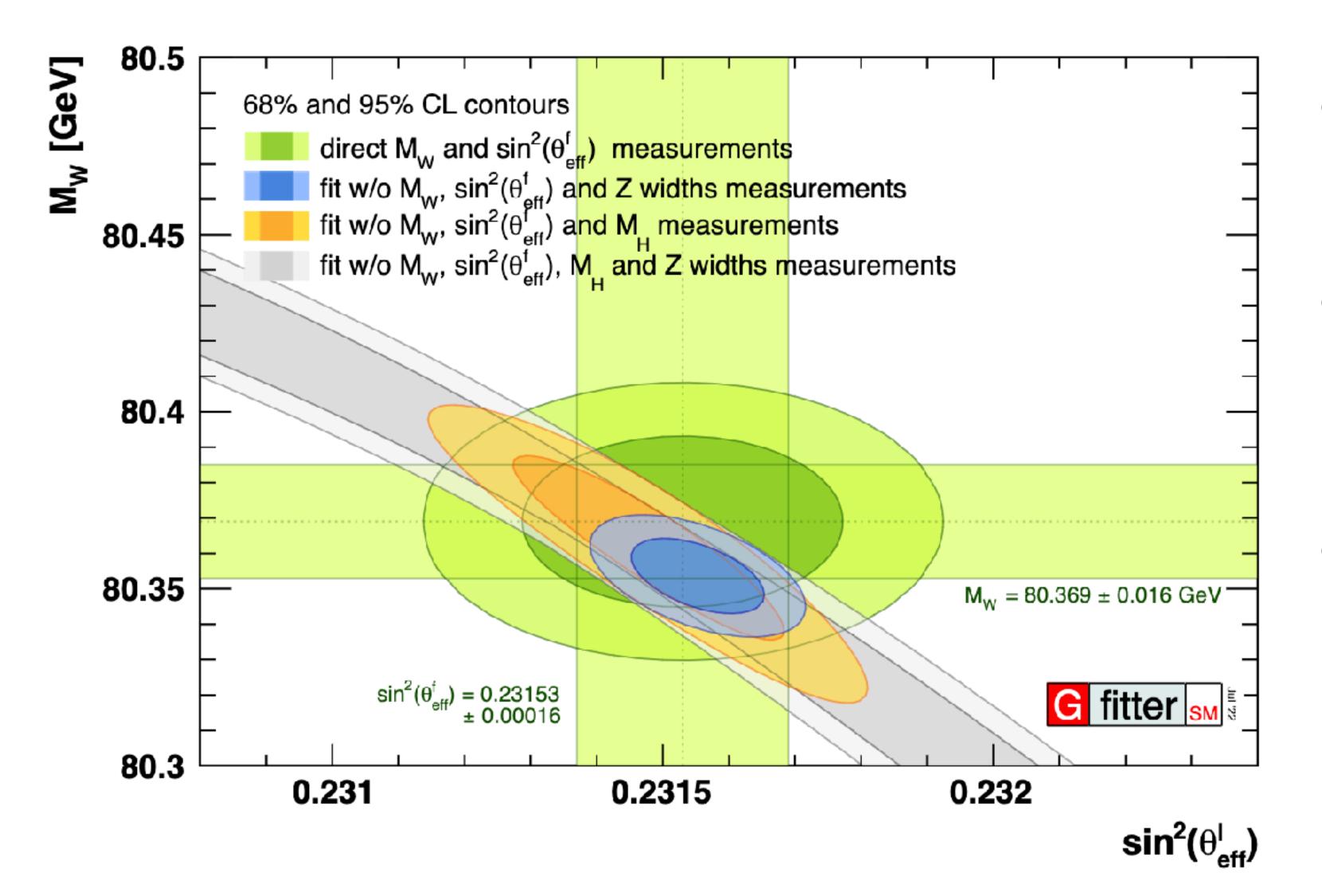
O Different predicted patter of PDF correlations for the different experiments



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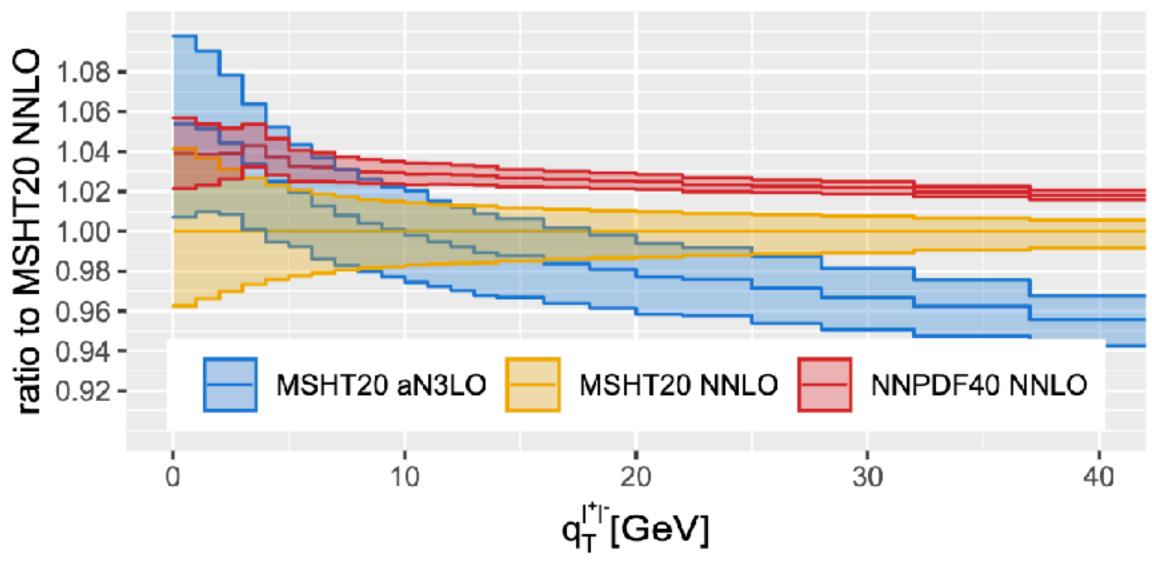
HADRON COLLIDER MEASUREMENTS IN THE EW FIT



- O Different measurements taken to be uncorrelated
- O Not true anymore if
 PDFs become the
 largest uncertainty
- O How much correlated? How much does it depend on the PDF chosen?



T. Neumann, MCFM NNLO Z+jet

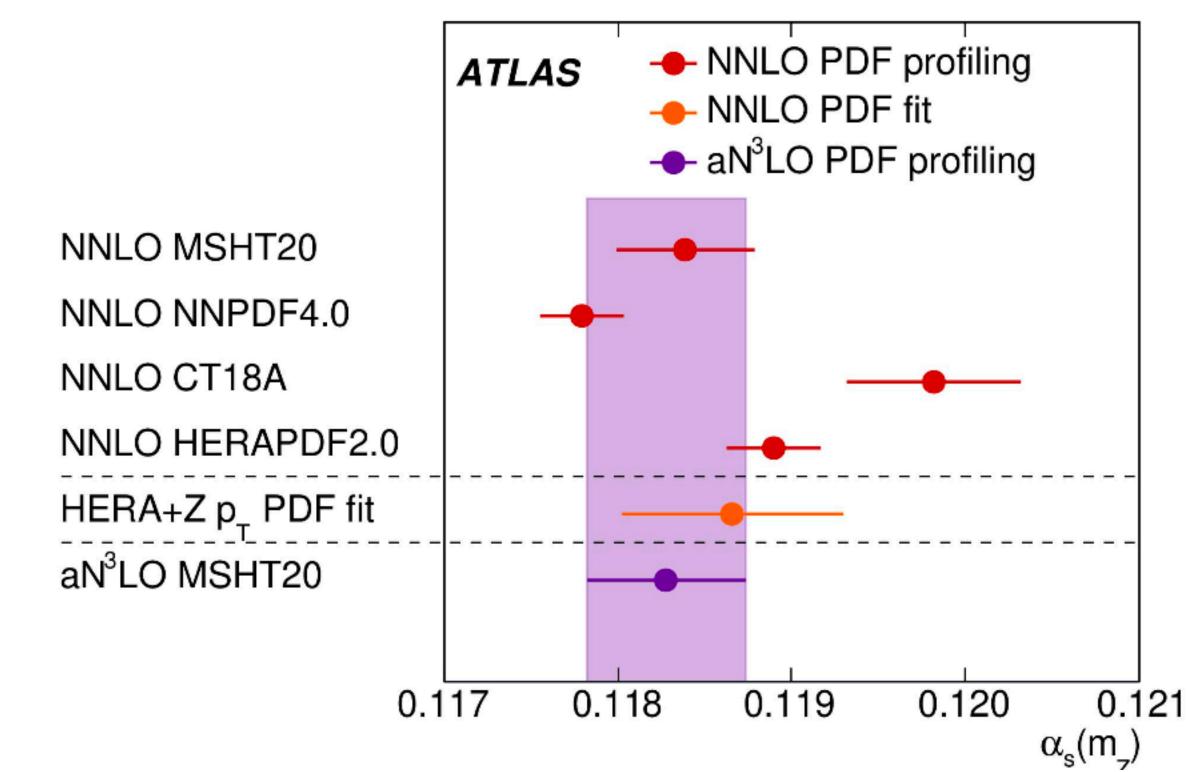


Spread in NNLO PDFs ~0.00102, Þ driven by NNPDF4.0 vs CT18A

- CT18 not compatible with other set within PDF uncertainties
- Often as large as the total D measurement uncertainty

PDMS NN СΊ HE

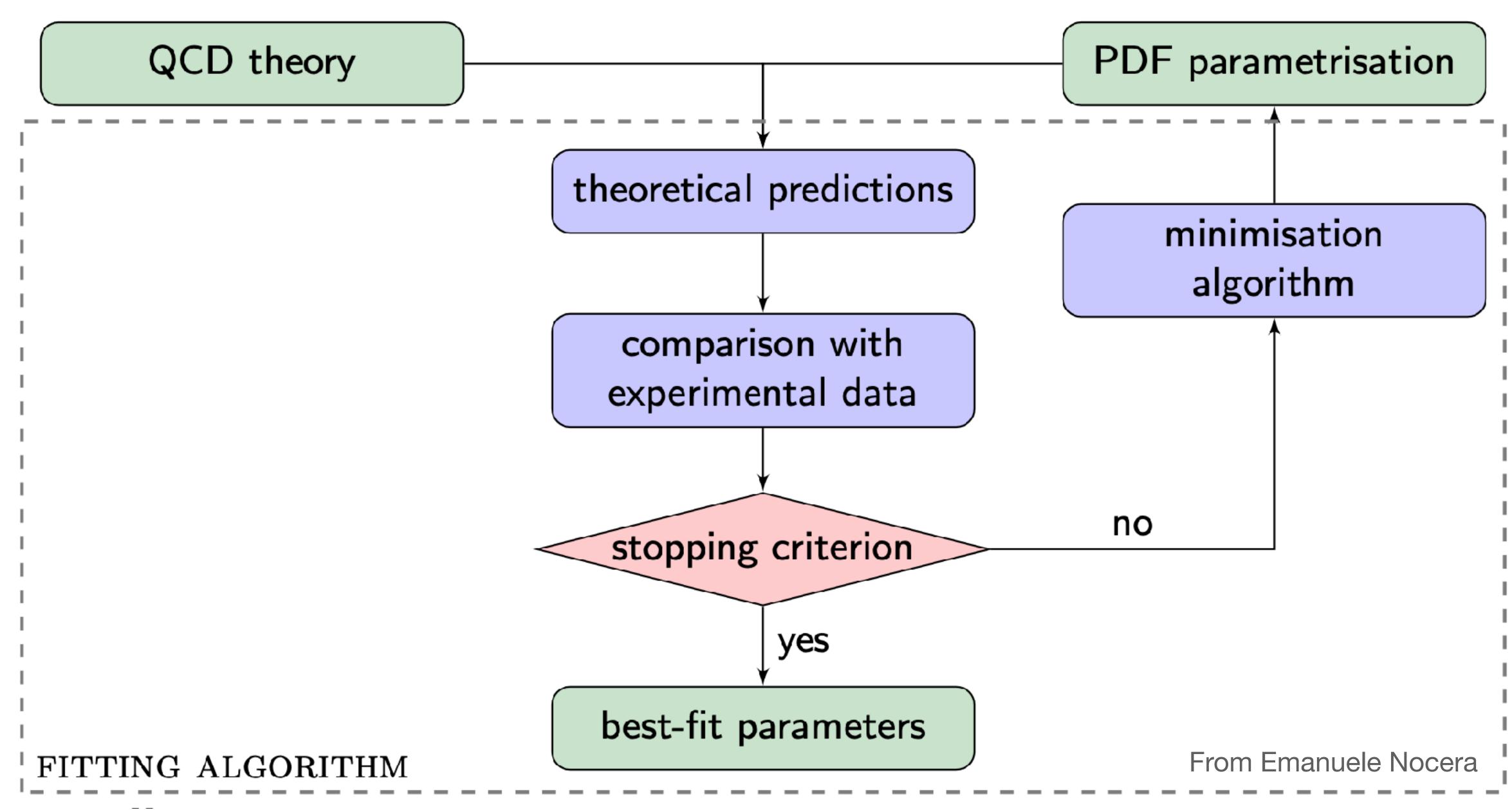
PDFS IN ATLAS $\alpha_S(m_Z)$ FROM Z PT



DF set	$\alpha_{ m s}(m_Z)$	PDF uncertainty
SHT20 [37]	0.11839	0.00040
NPDF4.0 [84]	0.11779	0.00024
$\Gamma 18A$ [29]	0.11982	0.00050
ERAPDF2.0 [65]	0.11890	0.00027



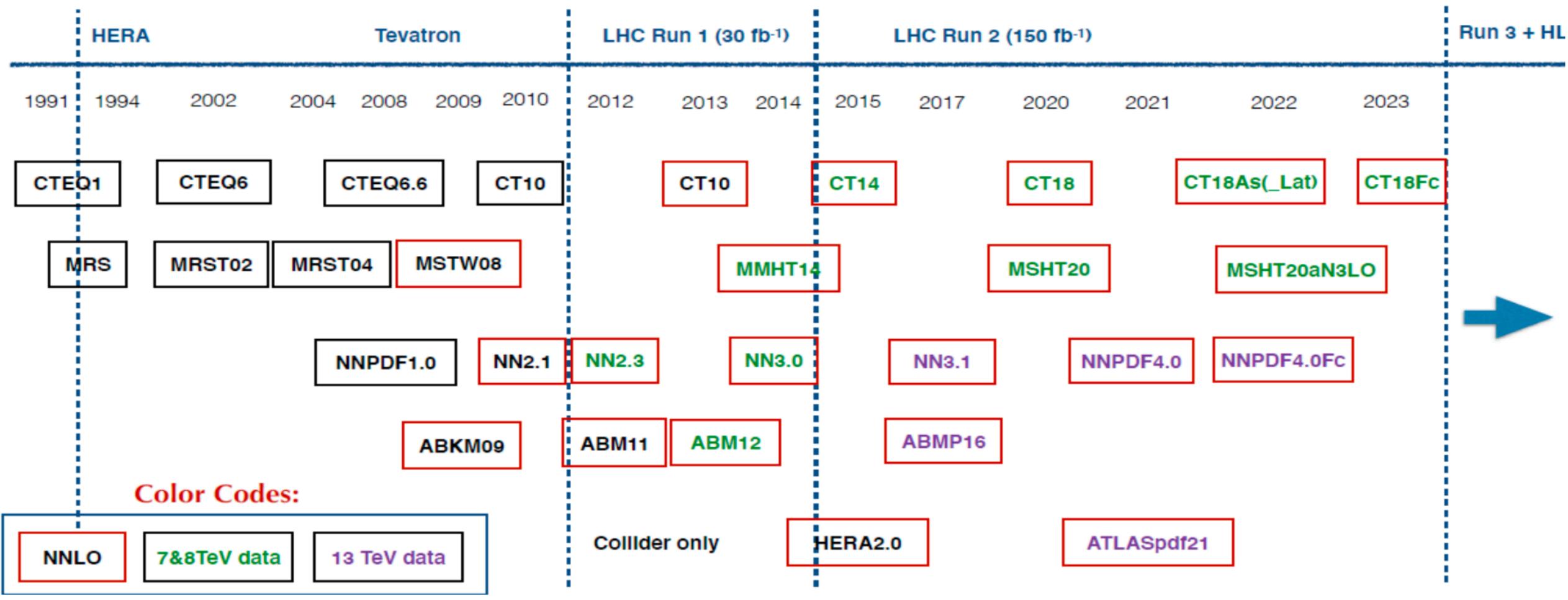




PDFs fitting 101

HISTORY OF PDF DETERMINATIONS

Ο non-perturbative parametrization, heavy-flavor scheme and fit methodology: ABMP, CTEQ-TEA, MSHT, NNPDF



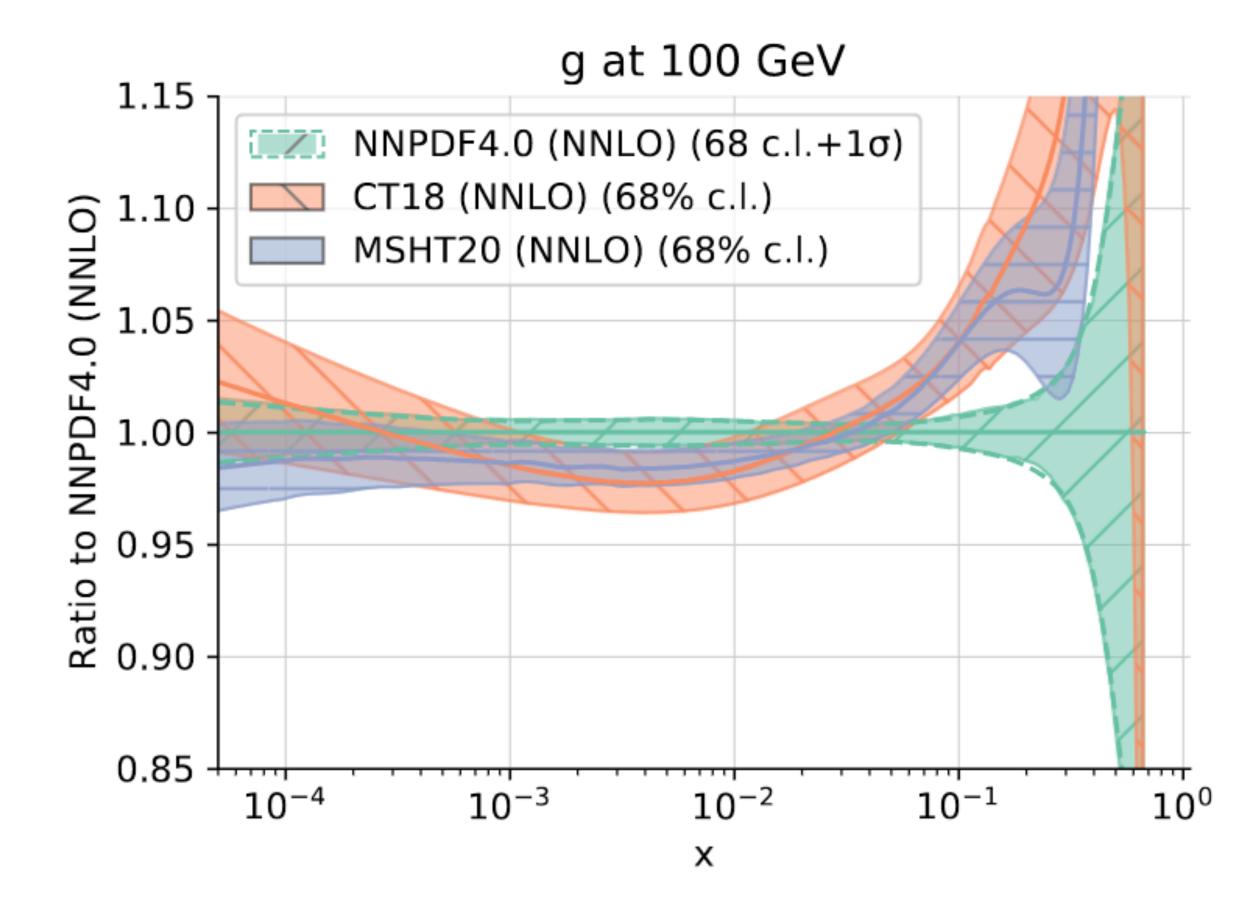
Global PDFs extracted by several groups making different choices in their input data,

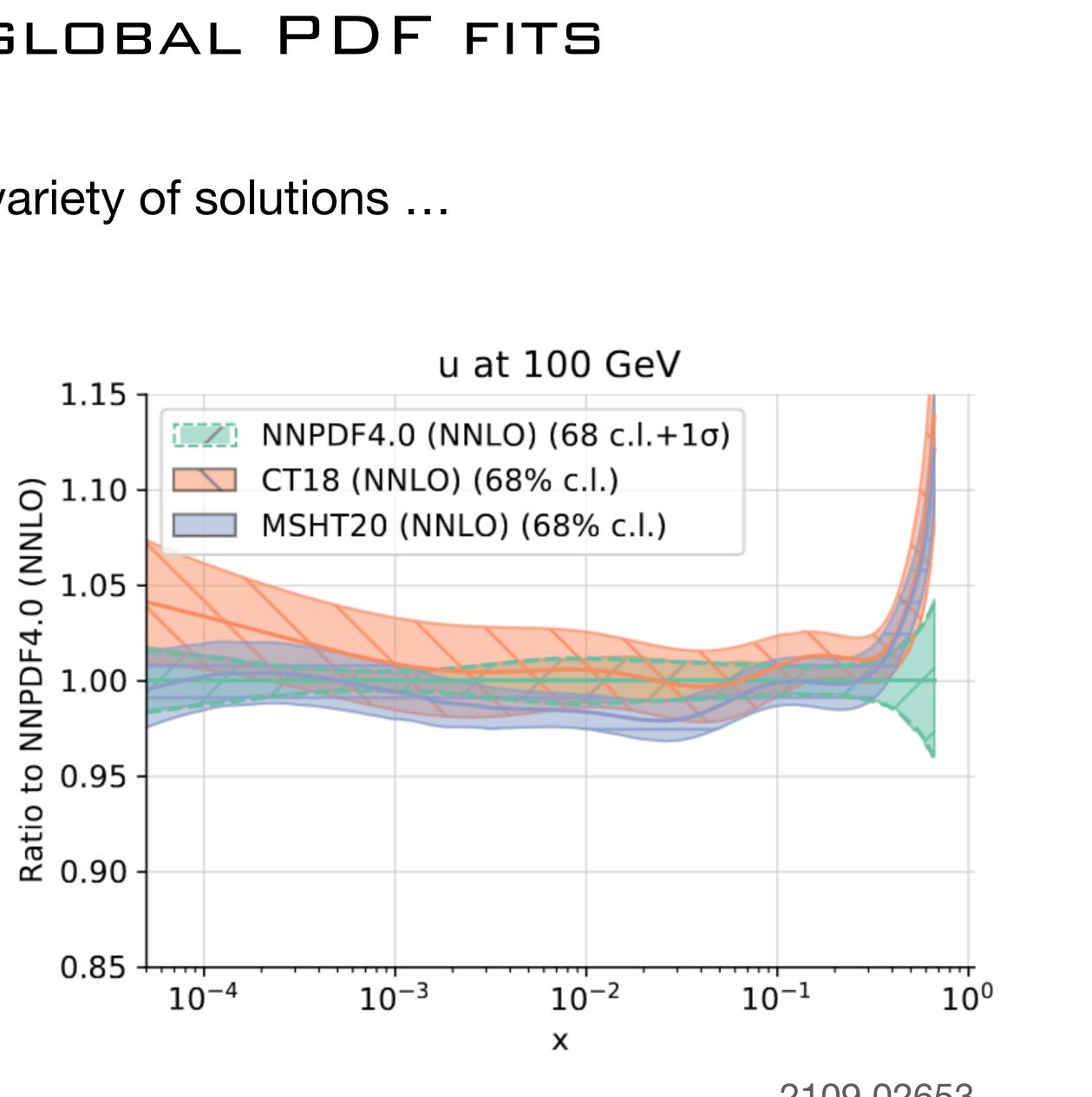
From Jun Gao



THE STATUS OF GLOBAL PDF FITS

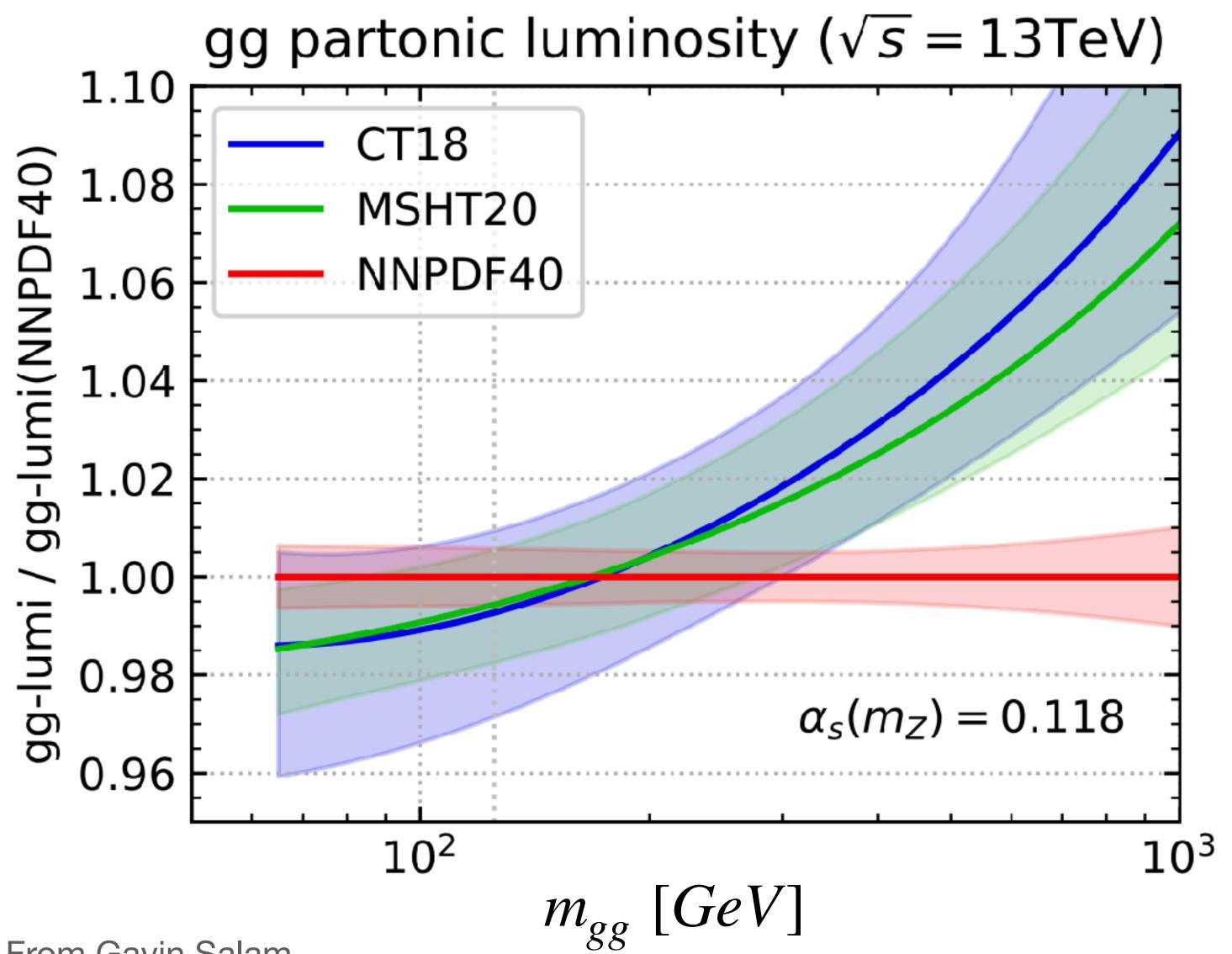
A complex problem that leads to a variety of solutions ... Ο





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From Gavin Salam

THE STATUS OF GLOBAL PDF FITS

gg-lumi, ratio to PDF4LHC15 @ m_H

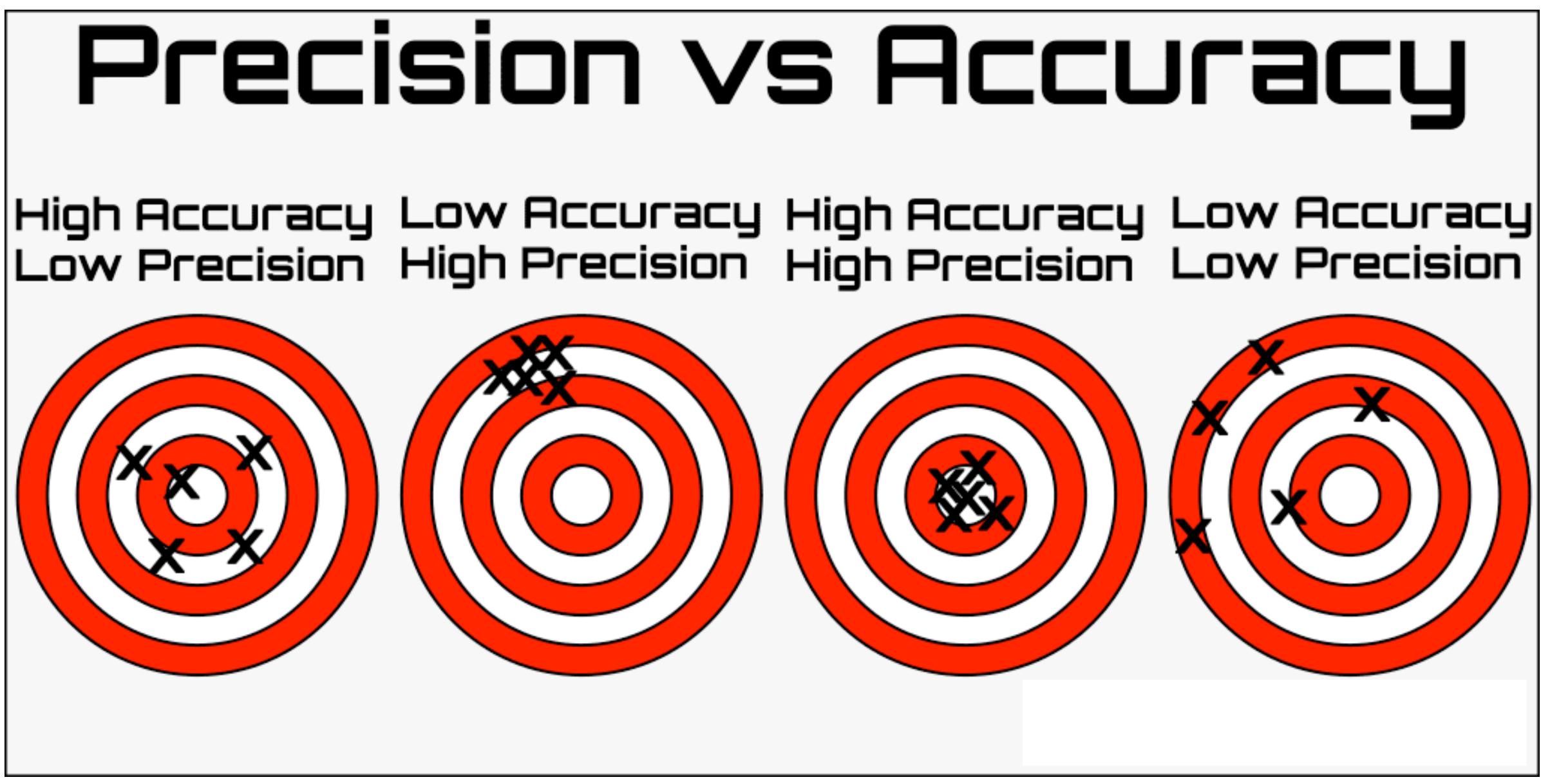
PDF4LHC15	1.0000	±	0.0184
PDF4LHC21	0.9930	±	0.0155
CT18	0.9914	\pm	0.0180
MSHT20	0.9930	\pm	0.0108
NNPDF40	0.9986	\pm	0.0058

NNPDF and MSHT now Ο at %-level precision

Yet in significant tension \mathbf{O} with each other









PDFS AND THEIR UNCERTAINTIES

- Ο
- Expect PDF uncertainties to be largely correlated (but how much exactly?) Ο



- Dataset choice ()
- Experimental uncertainties Ο
- Statistical + systematics Ο
- Correlations \mathbf{O}
- χ^2 definition
- **Outliers treatment**

Methodological:

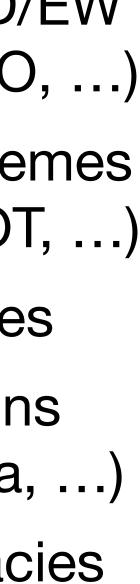
- Parametrization choice \mathbf{O}
- Regularization \mathbf{O}
 - (what is a good PDF?)
- Uncertainties prescription O (i.e. tolerances)

Most of the measurements and theory inputs are common among PDF groups

Data = PDF $\otimes \sigma_H$

Theoretical:

- Perturbative QCD/EW ()order (NLO, NNLO, ...)
- Heavy-flavor schemes \mathbf{O} (FONLL, RT, ACOT, ...)
- Perturbative scales
- Theory calculations \mathbf{O} (q_T-subtr, antenna, ...)
- Stat./Grid accuracies ()



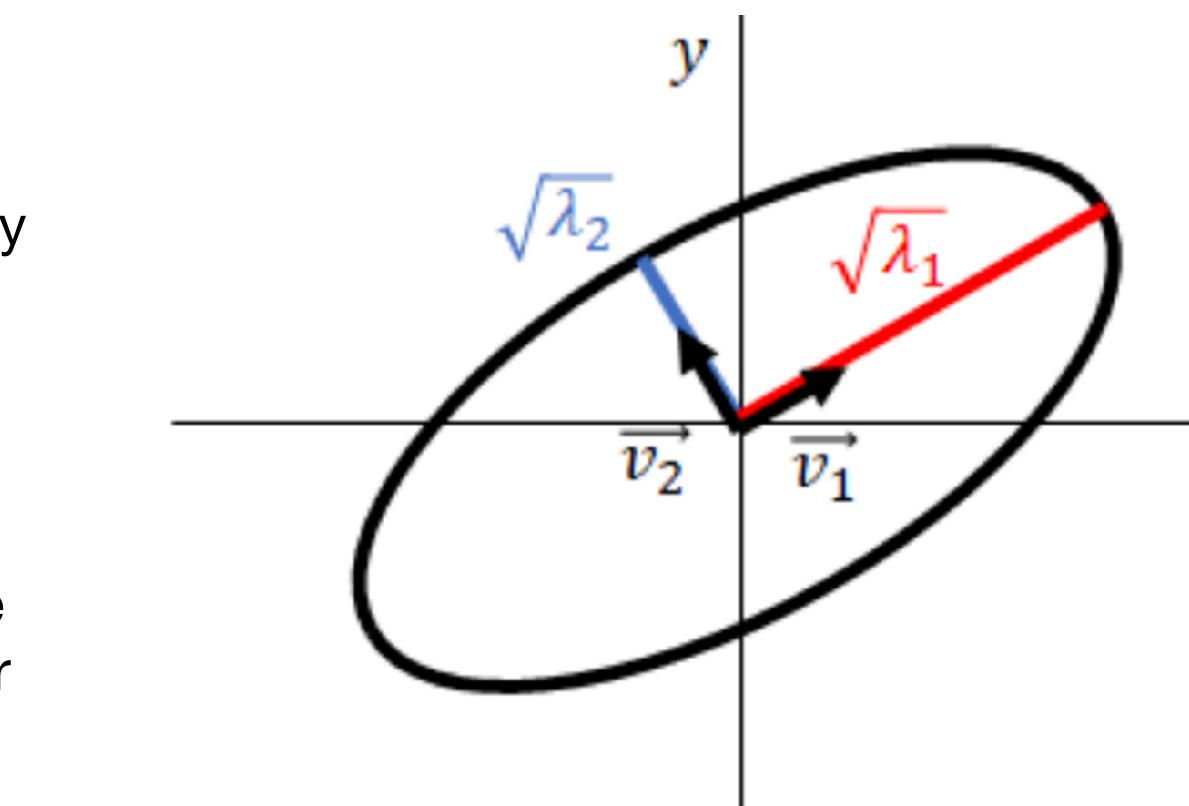
CONSTRAINING PDFS USING NEW DATA

- Ο
- Ο

- Our data carries information on \mathbf{O} the PDFs but not enough to fully determine them
- Combine the likelihood of our measurement with the (approximated) likelihood of the PDF fit to extract our parameter

PDF groups typically provide uncertainties as Hessian eigenvector Or as replicas, but the following considerations would equally apply

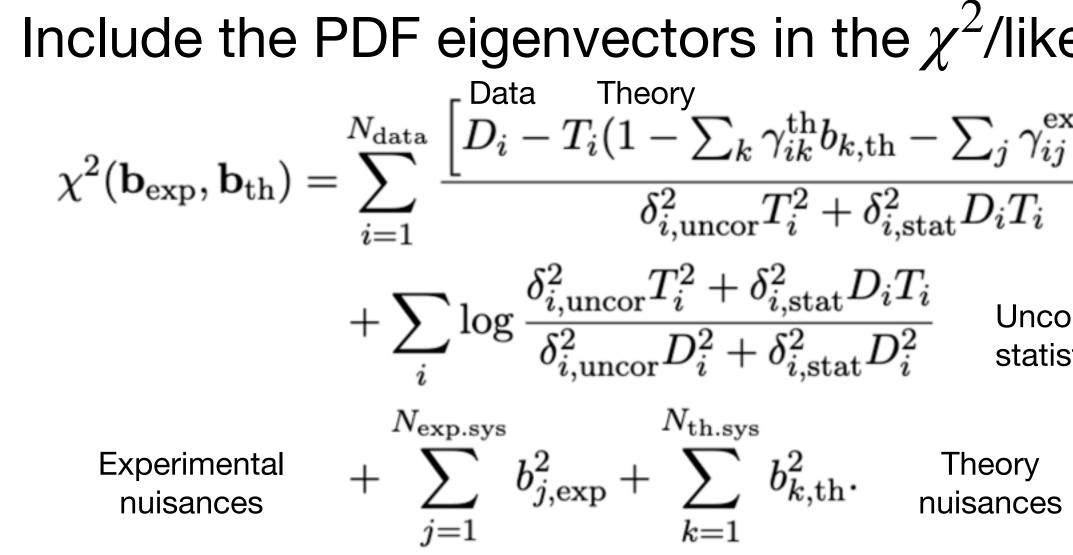
Eigenvectors provide a suitable representation of the PDF likelihood near the fit minimum suitable for propagating them in statistical analyses







HESSIAN PROFILING OF PDFS



O

Ο

The values of the nuisance parameters at the minimum define a \mathbf{O} new profiled PDF with (generally) smaller uncertainties

$$f'_0 = f_0 + \sum_k b_{k,\text{th}}^{\min} \left(\frac{f_k^+ - f_k^-}{2} + b_k^{\text{r}} \right)$$

This reduction in PDF uncertainties happens as long as their covariance is \mathbf{O} included in the fit, even if the nuisance parameters are not explicitly used

Include the PDF eigenvectors in the χ^2 /likelihood as covariance/nuisances

$$\left[b_{k, \text{th}} - \sum_{j} \gamma_{ij}^{\exp} b_{j, \exp}) \right]^2 + \delta_{i, \text{stat}}^2 D_i T_i$$

 $\left[\frac{\text{At}}{\text{At}} \frac{D_i T_i}{D_i^2} \right]$ Uncorrelated and statistical uncertainties

Nuisance parameter impacts

The data pulls and constrains (a linear combination of) the PDF nuisances

 $p_{k,\text{th}}^{\min} \frac{f_k^+ + f_k^- - 2f_0}{2}$ Defines "profiled" PDFs



PDF UNCERTAINTIES AND TOLERANCES

- **O** PDF uncertainties are derived using "Tolerance factors" $\Delta \chi^2 = T^2$
 - Introduced to avoid underestimated uncertainties due to bad goodness-of-fit
 - Corresponds to scaling the errors on ALL input measurements by a factor T^2
- O Crude approach to deal with model deficiencies, and analogous to PDG scaling

- O Different groups use different approaches:
 - CTEQ-TEA:
 Global $T^2 \sim 30$ designed to also cover for different PDF parameterizations
 - MSHT:
 Dynamical $T^2 \sim 10$ different for each eigenvector
 - NNPDF/ABMP: No tolerance applied, T²=1 (does this mean their GoF is good?)



PROFILING AND TOLERANCES IN LHC MEASUREMENTS

- When we propagate PDF eigenvectors we ignore tolerances ($T^2 = 1$) Ο
 - The impact of our data on PDFs is overestimated
 - For CT (MSHT) PDFs we effectively assign a weight of ~30 (3) to the new data
 - Equivalent to taking the PDG average of some quantity, but removing the scaling factor on its uncertainty
- Our result is not equivalent to including our data in the original PDF fit Ο Consistent approach requires rescaling of PDF eigenvectors by T^2 and using $\Delta \chi^2 = T^2$ when deriving the constraints on the parameters
- At the same time, we extract a parameter from our measurement, not the PDFs Ο and we do not make an average of different measurements
- What sense does it make in scaling up our measurement uncertainties due \mathbf{O} to (not-understood) tensions among some other data or with theory?



PROFILING AND TOLERANCES - III

- Need to construct a procedure that allows us to: \mathbf{O}
 - Estimate the impact of our data to the PDFs consistently with how the original PDF uncertainties have been derived (with tolerance)
 - Evaluate the measurement uncertainties on the parameter of interest without inflating them (without tolerances)
- How to have different $\Delta \chi^2$ for different uncertainties ? Ο
 - Perform a profiling of PDF with the new measurement including the proper eigenvectors rescaling and tolerance factors on $\Delta \chi^2$
 - Use the PDFs such obtained to extract the POI for fixed PDFs using $\Delta \chi^2 = 1$
 - Determine the PDF uncertainty on the POI through fixed error propagation of the reduced/rotated eigenvectors obtained in step 1

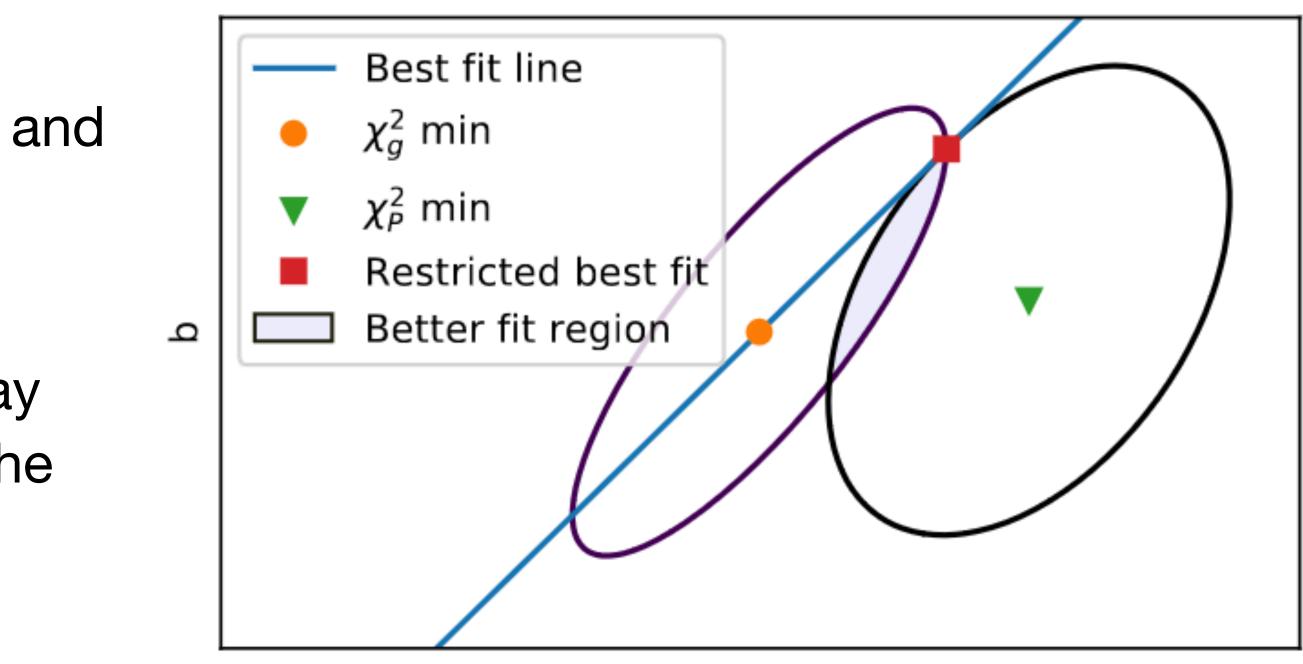
(from T. Cridge)



PDF PROFILING AND $\alpha_S(m_Z)$

- O The situation when we extract the strong coupling is more delicate
 - $\sim \alpha_S(m_Z)$ is obtained performing a χ^2 scan using PDFs obtained at fixed $\alpha_S(m_Z)$
 - PDF uncertainty evaluated at the nominal value of $\alpha_S(m_Z)$

- lgnores correlations between $\alpha_S(m_Z)$ and PDFs
- The result of a full PDF+ $\alpha_S(m_Z)$ fit may well be in a direction not probed by the factorised approximation







PDFs for precise EW measurements

- Ο largest possible set of input measurements
 - Requires compromises in the level of understanding one can achieve
 - For certain measurements we lack detailed experimental informations, or there are known theoretical issues in the interpretation, ...

- Ο
 - Can we foresee "reduced data PDF fits" using a subset of accurate and self-consistent measurements with state-of-the-art theory?
 - Can we have an estimate of the impact of each individual choice \triangleright entering a PDF fit (input data, theory, tolerances, ...)?

Global PDF fits aim to provide the best average description using the

What we need are accurate PDFs with reliable and complete uncertainties (not necessarily the smallest possible) for specific data/processes



THE EW PRECISION LEGACY OF THE LHC

- O Result of LHC EW precision measurements will stay with us for a long time
 - Essential to ensure future-proof measurements that can be updated to newer PDFs and theoretical models and combined across experiments
 - At LEP, this was done through the definition of pseudo-observables
- O Closest analogue at the LHC: simultaneous PDF+EW parameter extractions
 - Full control over input choices and assumptions entering the PDFs
 - Better defined (and possibly smaller) PDF uncertainties
 - Combinations and updates become (in principle) trivial
- **O** Measurements of unfolded (NC/CC) DY cross-sections, asymmetries (A_{FB}, A_i) and boson/lepton p_T would encompass the full program of EW precision measurements (m_W, $\sin^2 \theta_{eff}^l$, $\alpha_S(m_Z)$)



SUMMARY

Ο

- A deeper understanding of PDFs and their uncertainties is a must Ο
 - We are not after the best average description of all measurements,
- Ο
- Ο and reinterpretation of LHC EW precision measurements
 - D of the input data, and the many other prescriptions related to PDFs, ...

Uncertainties propagation for EW measurements at the LHC *almost broken* by PDFs

but for the most accurate PDFs using the most reliable measurements

PDF uncertainties on (most) hadron collider measurements are underestimated Tolerances should be accounted for when including PDFs in fits (but how exactly?) Spread of central values often larger than PDF uncertainties, how to quote it?

Ultimately, simultaneous PDF+EW fits are probably the way to go to allow control over PDFs and their uncertainties and would provide a framework for preservation

But opens up new questions related to the definition, choice and consistency

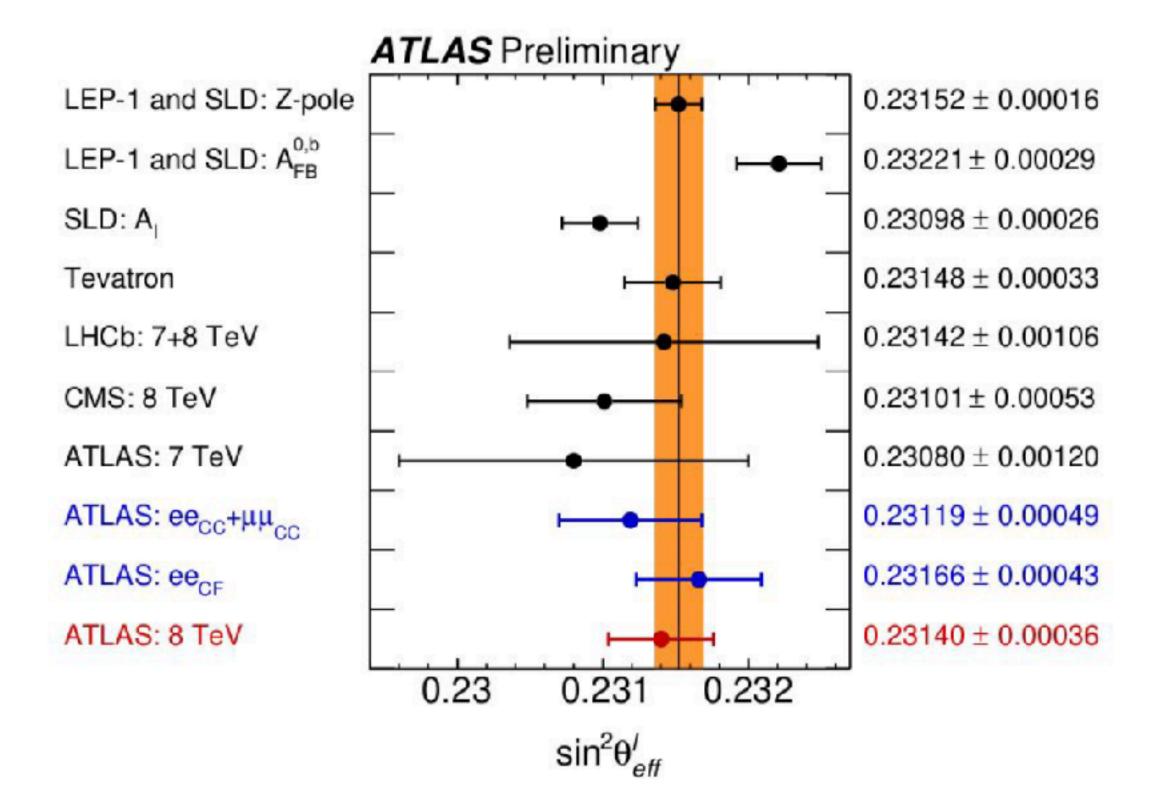


BACKUP



PDFs in $\sin^2 \theta_{\rm eff}^l$ - ATLAS 8 TeV

0.23140 ± 0.00021 (stat.) ± 0.00024 (PDF) ± 0.00016 (syst.)

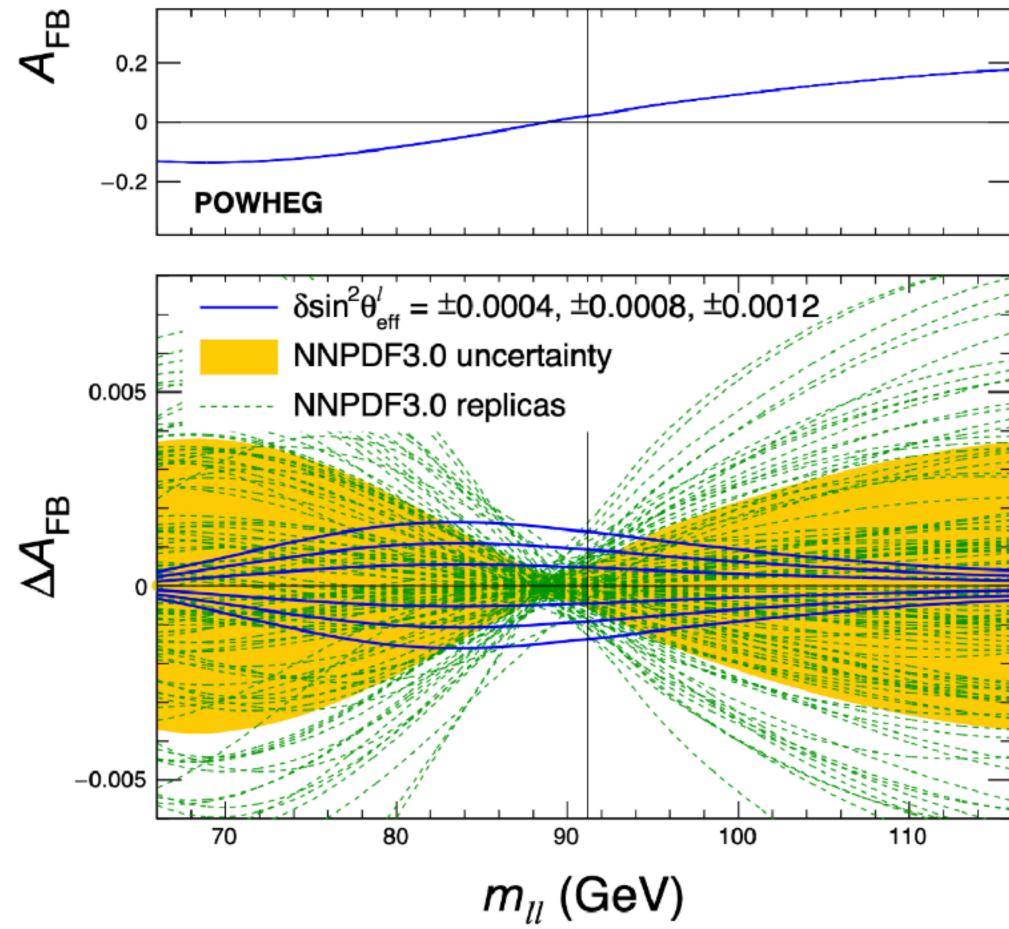


	CT10	CT14	MMHT14	NNPDF31
$\sin^2 \theta_{\text{eff}}^{\ell}$	0.23118	0.23141	0.23140	0.23146
	Uncertainties in measurements			
Total	39	37	36	38
Stat.	21	21	21	21
Syst.	32	31	29	31

Large uncertainty from envelope of PDFs, $3 \cdot 10^{-3}$, but using old PDF sets



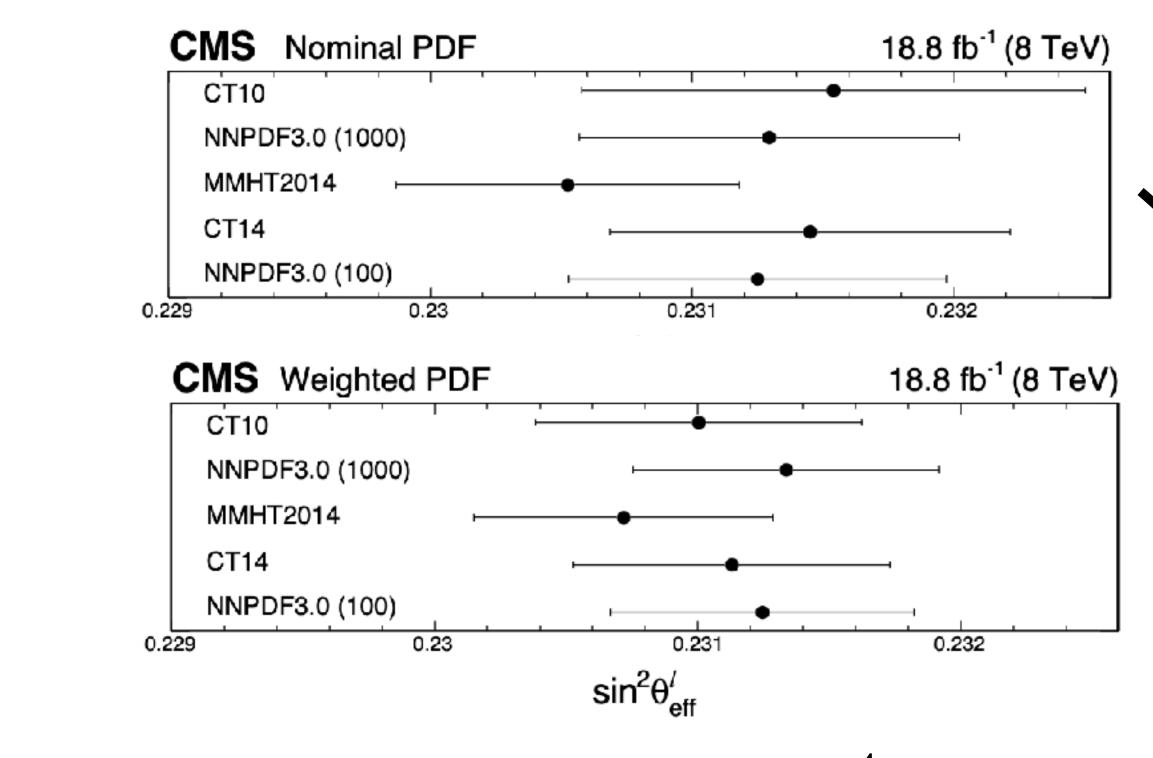
impact on the weak mixing angle





LHC measurements rely on the correlation pattern in the PDFs to reduce their

 $\sin^2 \theta_{eff}^{\ell} = 0.23101 \pm 0.00036$ (stat) ± 0.00018 (syst) ± 0.00016 (theo) ± 0.00031 (PDF).



PDF uncertainty of $3 \cdot 10^{-4}$ vs MSHT14/NNPDF30 spread of $6 \cdot 10^{-4}$





CONSTRAINING PDFS USING NEW DATA

- Ο
- Ο
- Ο result of including it in a new fit

Most of the times our data carries useful information about the PDFs that we would want to exploit, but not enough to fully determine them

Ideally would include the new data in a simultaneous fit of the parameter of interest and the PDF (requires more data and is complicated, see later)

Approximations allow to propagate PDF uncertainties in statistical analyses in a way that, under certain assumptions, would reproduce the

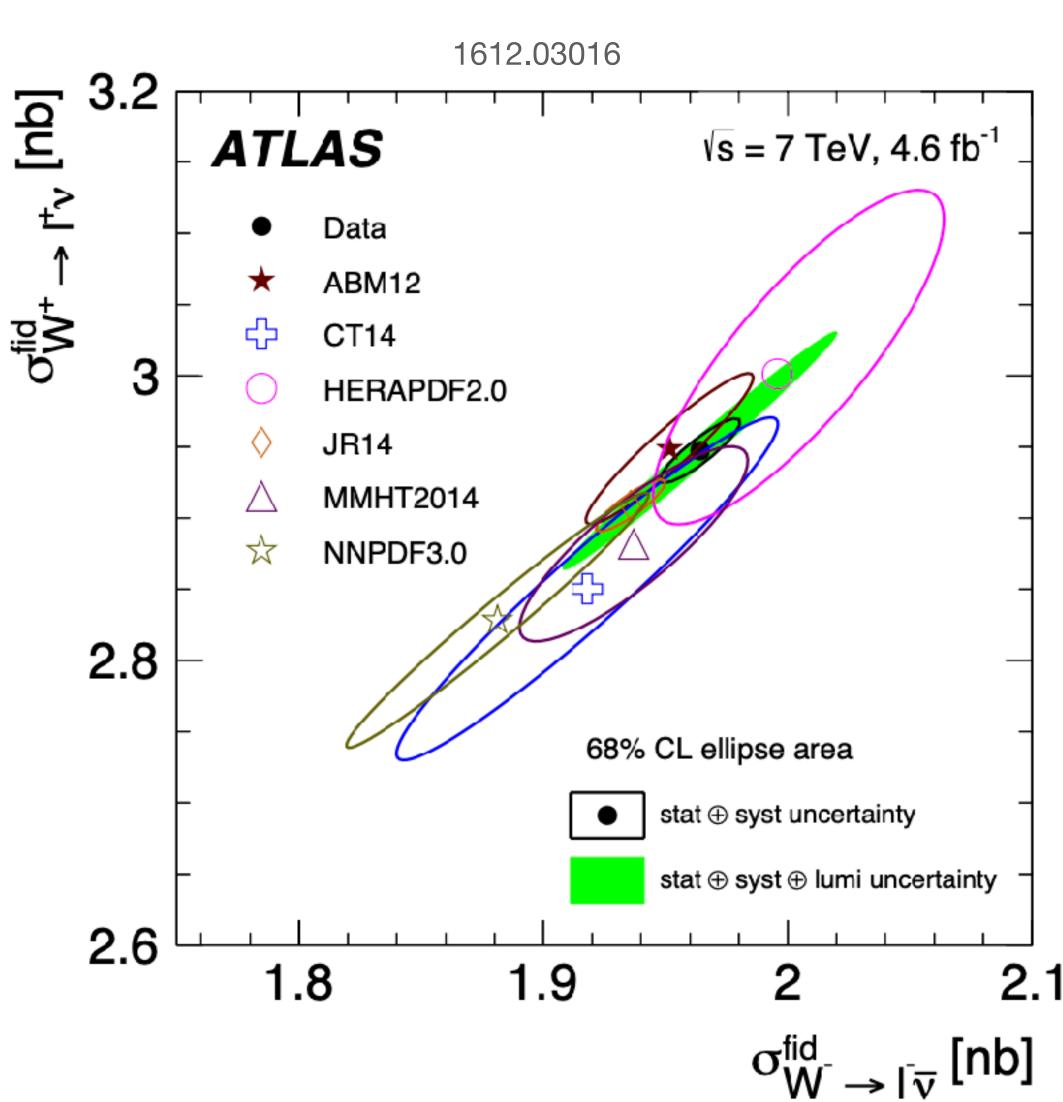


FORCED ERROR PROPAGATION

- PDF uncertainty on the measured parameter
- In this case we would only like to propagate the input PDF uncertainties to the parameter of interest
 - Forced error propagation (also called externalized or offset uncertainties)
 - Repeat the measurement for each O eigenvector/replica
 - Evaluate the difference in the fitted Ο parameter

$$\delta_{\theta} = \hat{\mu}(\hat{\theta} + \sigma_{\theta}) - \hat{\mu}(\hat{\theta})$$

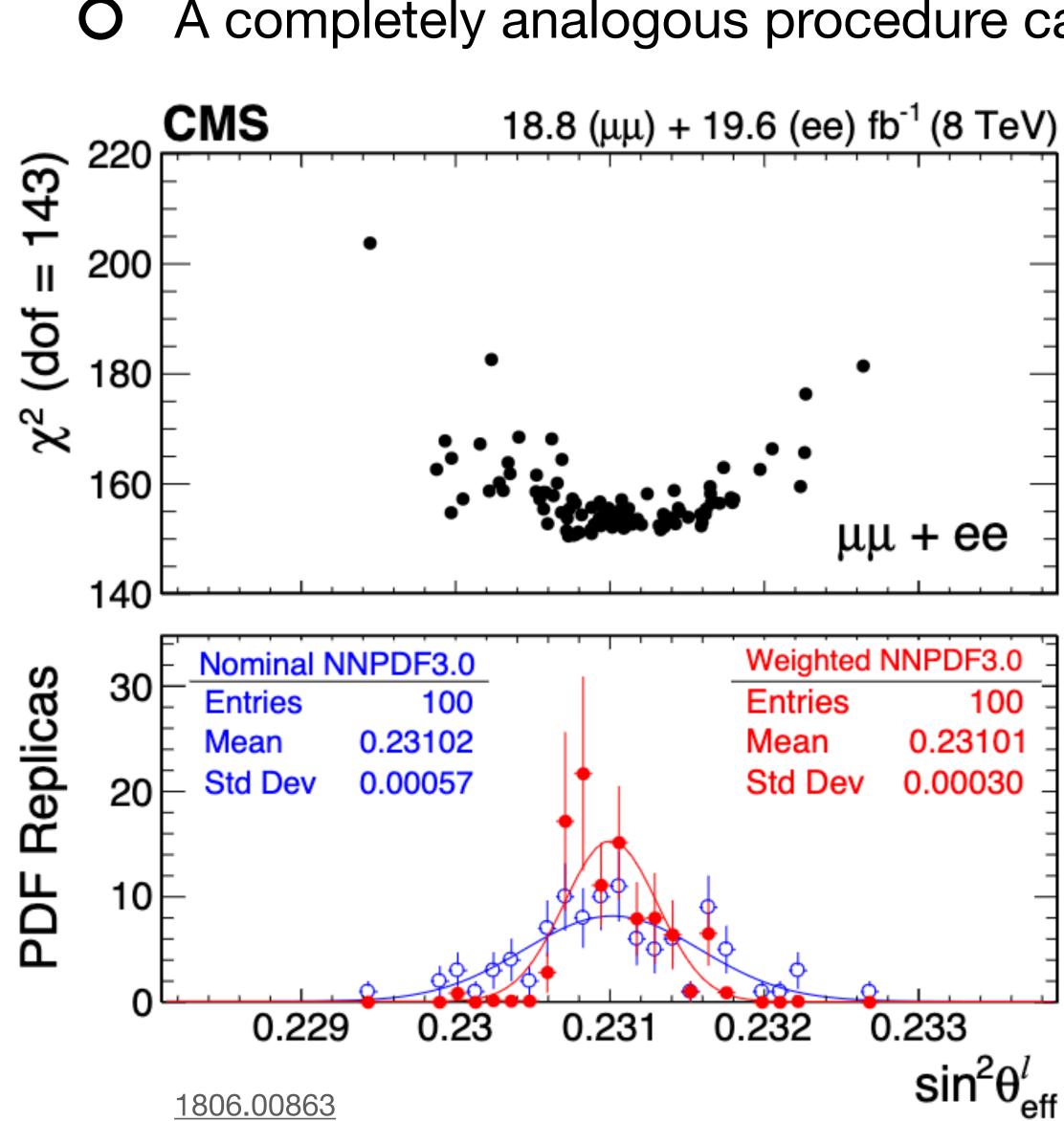
Sometimes our data is not very sensitive to PDF or we simply want to report the original







BAYESIAN REWEIGHING



A completely analogous procedure can be defined for replica uncertainties

Downweight replicas far from Ο the data according to:

$$w_{k} = \frac{(\chi_{k}^{2})^{\frac{1}{2}(N_{data}-1)} \exp^{-\frac{1}{2}\chi_{k}^{2}}}{\frac{1}{N_{rep}} \sum_{k=1}^{N_{rep}} (\chi_{k}^{2})^{\frac{1}{2}(N_{data}-1)} \exp^{-\frac{1}{2}\chi_{k}^{2}}}$$

Recompute the uncertainty on your Ο observable with the reweighted replicas

$$\langle O^{\text{new}}(\text{PDF}) \rangle = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} w_k O(\text{PDF})$$

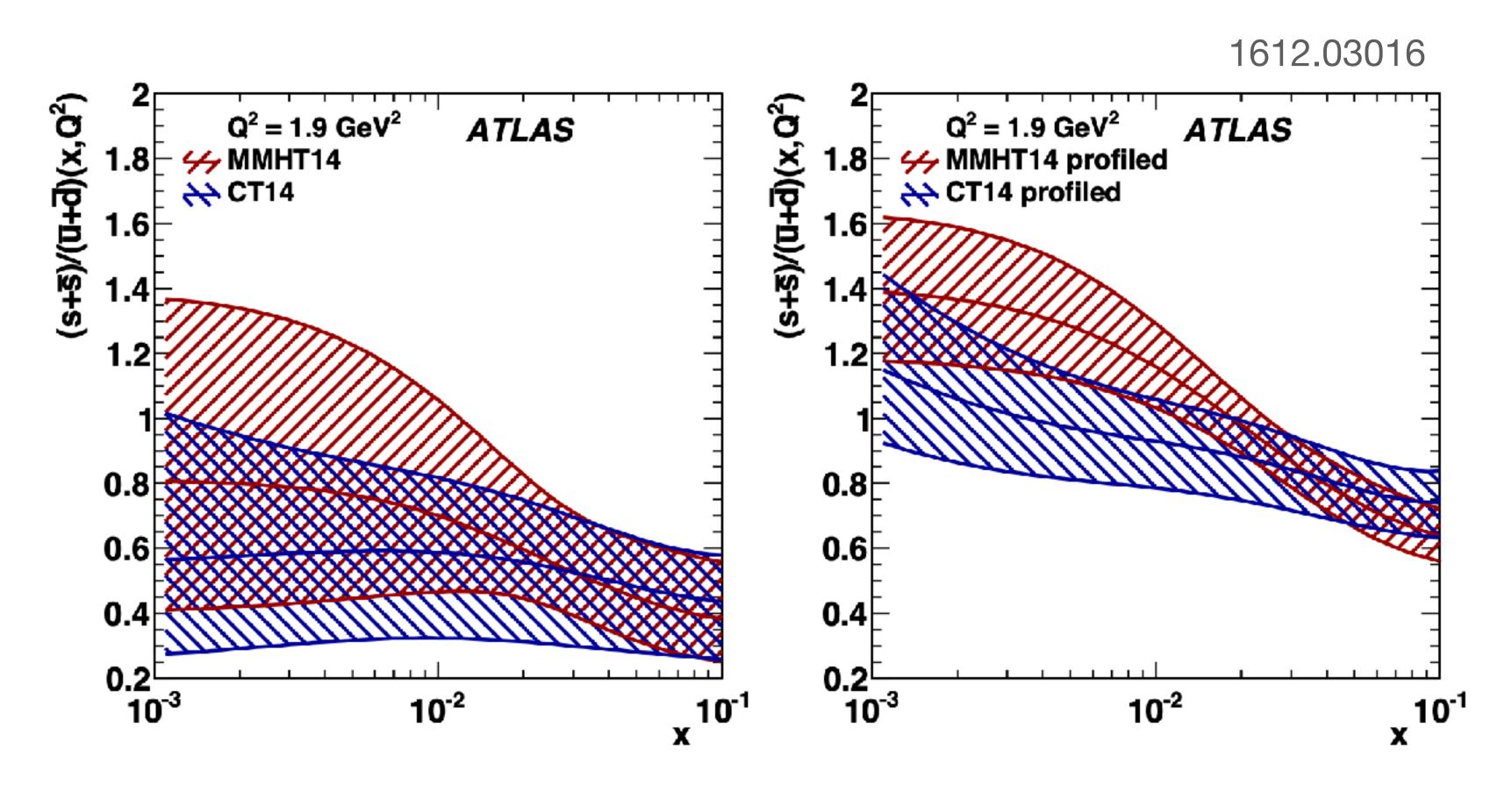






PROFILING AND TOLERANCES - I

- Hessian approximation is only valid around the minimum
- the original PDFs
- When are the profiling results reliable?



In most cases including the new data strongly constraints and modifies



PROFILING AND TOLERANCES: CTEQ VIEW

Augmented likelihood for PDFs with global tolerance

- 2. Write the **augmented** likelihood density for this definition:
 - $P(D_i|T_i) \propto e^{-\chi^2/(2T^2)}$
- 3. When profiling 1 new experiment with the prior imposed on PDF nuisance parameters $\lambda_{\alpha,th}$:

$$\chi^{2}(\vec{\lambda}_{exp},\vec{\lambda}_{th}) = \sum_{i=1}^{N_{pt}} \frac{\left[D_{i} + \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,exp} - T_{i} - \sum_{\alpha} \beta_{i,\alpha}^{th} \lambda_{\alpha,th}\right]^{2}}{s_{i}^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} T^{2} \lambda_{\alpha,th}^{2}. \qquad \beta_{i,\alpha}^{th} = \frac{T_{i}(f_{\alpha}^{+}) - T_{i}(f_{\alpha}^{+})}{2}$$
new experiment priors on expt. systematics and PDF params
4. Alternatively, we can reparametrize $\chi^{2'} \equiv \chi^{2}/T^{2}$, so that 68% c.l. $\Leftrightarrow \Delta \chi^{2'} = 1$. We have
$$\frac{P(D_{i}|T_{i}) \propto e^{-\chi^{2'}/2}}{s_{i}^{2}T^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,th}^{2}.$$
consistent redefinition
$$\chi^{2'}(\vec{\lambda}_{exp}, \vec{\lambda}_{th}) = \sum_{i=1}^{N_{pt}} \frac{\left[D_{i} + \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,exp} - T_{i} - \sum_{\alpha} \beta_{i,\alpha}^{th} \lambda_{\alpha,th}\right]^{2}}{s_{i}^{2}T^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,exp}^{2}.$$
for $P(D_{i}|T_{i}) \propto e^{-\chi^{2'}/2}$

$$\chi^{2'}(\vec{\lambda}_{exp}, \vec{\lambda}_{th}) = \sum_{i=1}^{N_{pt}} \frac{\left[D_{i} + \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,exp} - T_{i} - \sum_{\alpha} \beta_{i,\alpha}^{th} \lambda_{\alpha,th}\right]^{2}}{s_{i}^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,exp}^{2}.$$
and $P(D_{i}|T_{i}) \propto e^{-\chi^{2'}/2}$
or $P(D_{i}|T_{i}) \propto e^{-\chi^{2'}/2}$

$$\chi^{2'}(\vec{\lambda}_{exp}, \vec{\lambda}_{th}) = \sum_{i=1}^{N_{pt}} \frac{\left[D_{i} + \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,exp} - T_{i} - \sum_{\alpha} \beta_{i,\alpha}^{th} \lambda_{\alpha,th}\right]^{2}}{s_{i}^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,exp}^{2}.$$

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$$\chi^{2'}(\vec{\lambda}_{exp}, \vec{\lambda}_{th}) = \sum_{i=1}^{N_{pt}} \frac{\left[D_{i} + \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,exp} - T_{i} - \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,th}\right]^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2}$$

$$\chi^{2}(\vec{\lambda}_{exp},\vec{\lambda}_{th}) = \sum_{i=1}^{N_{pt}} \frac{\left[D_{i} + \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,exp} - T_{i} - \sum_{\alpha} \beta_{i,\alpha}^{th} \lambda_{\alpha,th}\right]^{2}}{s_{i}^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} T^{2} \lambda_{\alpha,th}^{2}. \qquad \beta_{i,\alpha}^{th} = \frac{T_{i}(f_{\alpha}^{+}) - T_{i}(f_{\alpha}^{-})}{2}$$
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$$P(D_{i}|T_{i}) \propto e^{-\chi^{2'}/2}$$
consistent redefinition
$$\chi^{2'}(\vec{\lambda}_{exp},\vec{\lambda}_{th}) = \sum_{i=1}^{N_{pt}} \frac{\left[D_{i} + \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,exp} - T_{i} - \sum_{\alpha} \beta_{i,\alpha}^{th} \lambda_{\alpha,th}\right]^{2}}{s_{i}^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,th}^{2}.$$
5. Inconsistent redefinitions:
$$\chi^{2'}(\vec{\lambda}_{exp},\vec{\lambda}_{th}) = \sum_{i=1}^{N_{pt}} \frac{\left[D_{i} + \sum_{\alpha} \beta_{i,\alpha}^{exp} \lambda_{\alpha,exp} - T_{i} - \sum_{\alpha} \beta_{i,\alpha}^{th} \lambda_{\alpha,th}\right]^{2}}{s_{i}^{2}} + \sum_{\alpha} \lambda_{\alpha,exp}^{2} + \sum_{\alpha} \lambda_{\alpha,th}^{2}.$$
and $P(D_{i}|T_{i}) \propto e^{-\chi^{2'}/2}$
required to $s_{i} \rightarrow s_{i}/T$ or $\lambda_{\alpha,th} \rightarrow \lambda_{\alpha,th}T$ without $\beta_{i,\alpha,th} \rightarrow \beta_{i,\alpha,th}/T$]

2023-11-17

1. Start by defining the correspondence between $\Delta \chi^2$ and cumulative probability level: 68% c.l. $\Leftrightarrow \Delta \chi^2 = T^2$.

P. Nadolsky, PDF4LHC meeting



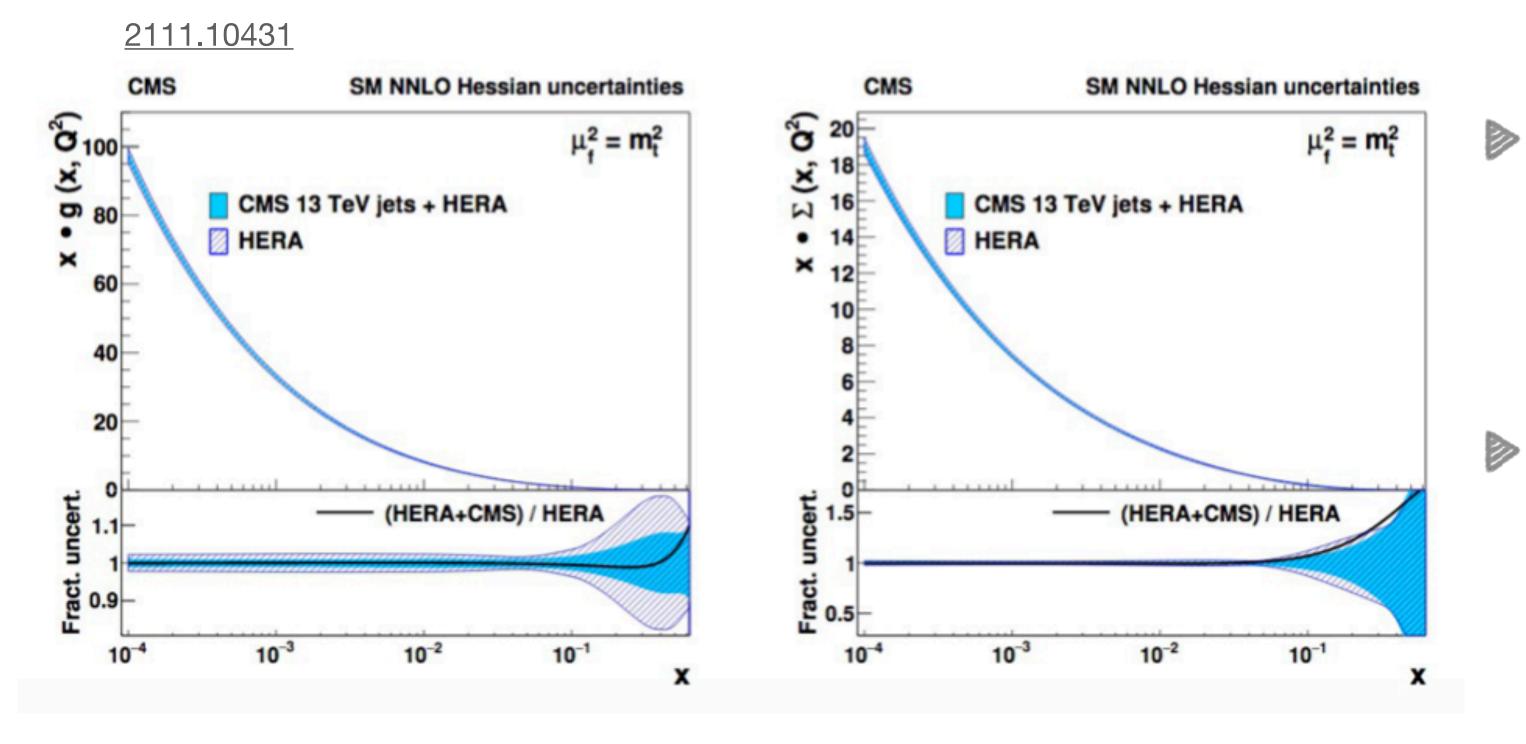


REFITTING PDFS

Ο

Or our measurement is already included in PDF sets, using it to extract other SM parameters would give biased results

In those cases perform a new PDF fit including minimal relevant data Ο



Sometimes data is so sensitive to PDFs that we know profiling would not work

- - Often done in the context of alphas extractions from jet measurements
 - CMS 13 teV inclusive jets fitted with HERA DIS data to obtain:

 $\alpha_{\rm S}(m_{\rm Z}) = 0.1166 \pm 0.0014 \,({\rm fit}) \pm 0.0007 \,({\rm model}) \pm 0.0004 \,({\rm scale}) \pm 0.0001 \,({\rm param.})$



LHC MEASUREMENTS AND PDFS

- Ο
- Ο

Drell-Yan	Flavour dec
W+charm	Strange PD
Jets	High-x gluor
Photon	Medium-x g
Top pair	Medium- an

PDFs most precisely determined from DIS data, but not all combinations probed • d_v is less precisely determined than u_v , no flavour decomposition of the light sea

LHC data cannot replace DIS, but can provide complementary information and help resolve tensions and disagreements which happen in global fits (or among them)

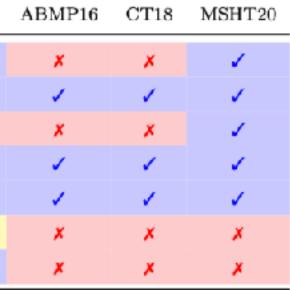
> composition of the sea, u_{i} , d_{j} , γ PDF F n PDF Iluon PDF าd high-x gluon PDF



LHC DATA IN PDF FITS

Data set	Ref.	NNPDF3.1	NNPDF4.0	ABMP16	CT18	MSHT20	Data set	Ref.	NNPDF3.1	NNPDF4.0	ABMP16	CT18	
ATLAS W, Z 7 TeV ($\mathcal{L} = 35 \text{ pb}^{-1}$)	[51]	1	1	1	1	1	CMS W asym. 7 TeV ($\mathcal{L} = 36 \text{ pb}^{-1}$)	[267]	×	×	X	x	
ATLAS W, Z 7 TeV ($\mathcal{L} = 4.6 \text{ fb}^{-1}$)	[52]	1	1	x	(✓)	1	CMS Z 7 TeV ($\mathcal{L} = 36 \text{ pb}^{-1}$)	[268]	×	×	×	×	
ATLAS low-mass DY 7 TeV	[53]	1	1	×	()	×	CMS W electron asymmetry 7 TeV	[55]	1	1	×	1	
ATLAS high-mass DY 7 TeV	[54]	1	1	×	(•)	1	CMS W muon asymmetry 7 TeV	[56]	1	1	1	1	
ATLAS W 8 TeV	[79]	×	(✔)	×	×	1	CMS Drell-Yan 2D 7 TeV	[57]	1	1	×	(✔)	
ATLAS DY 2D 8 TeV	[78]	×	1	×	×	1	CMS Drell-Yan 2D 8 TeV	[269]	(🗸)	×	×	×	
ATLAS high-mass DY 2D 8 TeV	[77]	×	1	×	()	1	CMS W rapidity 8 TeV	[58]	1.1	1.1	1	1	
ATLAS $\sigma_{W,Z}$ 13 TeV	[81]	×	1	1	×	×	CMS $W, Z p_T 8$ TeV ($\mathcal{L} = 18.4 \text{ fb}^{-1}$)	[270]	×	×	×	(✔)	
ATLAS W+jet 8 TeV	[93]	×	1	×	×	1	CMS $Z p_T 8$ TeV	[64]	1	1	×	(✔)	
ATLAS $Z p_T$ 7 TeV	[259]	(1)	×	×	(✔)	×	${\rm CMS}\;W+c\;7\;{\rm TeV}$	[76]	1	1	×	(🖍)	
ATLAS $Z p_T$ 8 TeV	[63]	1	1	×	1	1	CMS $W + c$ 13 TeV	[84]	×	1	×	×	
TLAS $W + c$ 7 TeV	[83]	×	1	×	()	×	CMS single-inclusive jets 2.76 TeV	[75]	1	×	×	×	
ATLAS σ_{tt}^{tot} 7, 8 TeV	[65]	1	- 2	1	×	x	CMS single-inclusive jets 7 TeV	[147]	1	(•)	×	1	
ATLAS σ_{tt}^{tot} 7, 8 TeV	[260-265]	×	×	-	-	×	CMS dijets 7 TeV	[74]	×	1	×	×	
ATLAS σ_{tt}^{tot} 13 TeV ($\mathcal{L} = 3.2 \text{ fb}^{-1}$)	[66]		×		÷.	1	CMS single-inclusive jets 8 TeV	[87]	×	1	×	1	
	[134]	×	- 2	×	Ŷ.	Ĵ,	CMS 3D dijets 8 TeV	[149]	×	(✔)	×	×	
ATLAS σ_{tt}^{tot} 13 TeV ($\mathcal{L} = 139 \text{ fb}^{-1}$)						×	CMS σ_{tt}^{tot} 5 TeV	[88]	×	1	*	×	
ATLAS σ_{tt}^{tot} and Z ratios	[266]	×	×	×	×	(🗸)	CMS σ_{tt}^{tot} 7, 8 TeV	[146]	1	1	×	×	
ATLAS $t\bar{t}$ lepton+jets 8 TeV	[67]	×	· ·	<u>^</u>	1	- *	CMS σ_{tt}^{tot} 8 TeV	[271]	×	×	×	×	
ATLAS $t\bar{t}$ dilepton 8 TeV	[89]	×		×	×		CMS σ_{tt}^{tot} 5, 7, 8, 13 TeV	[68, 272 - 280]	×	×	1	×	
ATLAS single-inclusive jets 7 TeV, R=0.6	[73]	· · ·	(✔)	×	1	· ·	CMS σ_{tt}^{tot} 13 TeV	[69]	1	1	1	×	
ATLAS single-inclusive jets 8 TeV, R=0.6	[86]	×	1	×	*	×	CMS $t\bar{t}$ lepton+jets 8 TeV	[70]	1	1	×	×	
ATLAS dijets 7 TeV, R=0.6	[148]	×	1	×	×	×	CMS $t\bar{t}$ 2D dilepton 8 TeV	[90]	×	1	×	1	
ATLAS direct photon production 8 TeV	[100]	×	(✔)	×	×	×	CMS $t\bar{t}$ lepton+jet 13 TeV	[91]	×	1	×	×	
ATLAS direct photon production 13 TeV	[101]	×	1	×	×	×	CMS $t\bar{t}$ dilepton 13 TeV	[92]	×	1	X	×	
ATLAS single top R_t 7, 8, 13 TeV	[94, 96, 98]	×	1	1	×	×	CMS single top $\sigma_t + \sigma_{\bar{t}}$ 7 TeV	[95]	×	1	1 - A	×	
ATLAS single top diff. 7 TeV	[94]	×	1	×	×	×	CMS single top R_t 8, 13 TeV	[97, 99]	×	1	×	×	
ATLAS single top diff. 8 TeV	[96]	×	1	×	×	×	CMS single top 13 TeV	[281, 282]	×	×	×	×	

Data set	Ref.	NNPDF3.1	NNPDF4.0
LHCb Z 7 TeV ($\mathcal{L} = 940 \text{ pb}^{-1}$)	[59]	1	1
LHCb $Z \rightarrow ee$ 8 TeV ($\mathcal{L} = 2 \text{ fb}^{-1}$)	[61]	1	1
LHCb W 7 TeV ($\mathcal{L} = 37 \text{ pb}^{-1}$)	[283]	×	×
LHCb $W,Z \to \mu$ 7 TeV	[60]	1	1
LHCb $W,Z \to \mu$ 8 TeV	[62]	1	1
LHCb $W \rightarrow e \ 8 \ { m TeV}$	[80]	×	(✔)
LHCb $Z \rightarrow \mu \mu, ee~13$ TeV	[82]	×	1

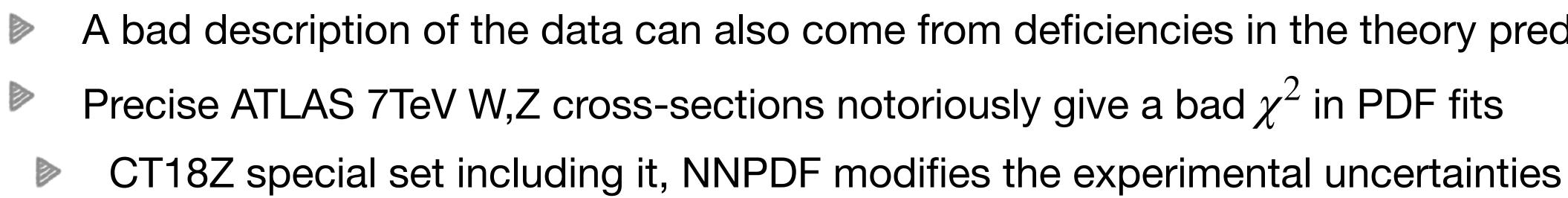


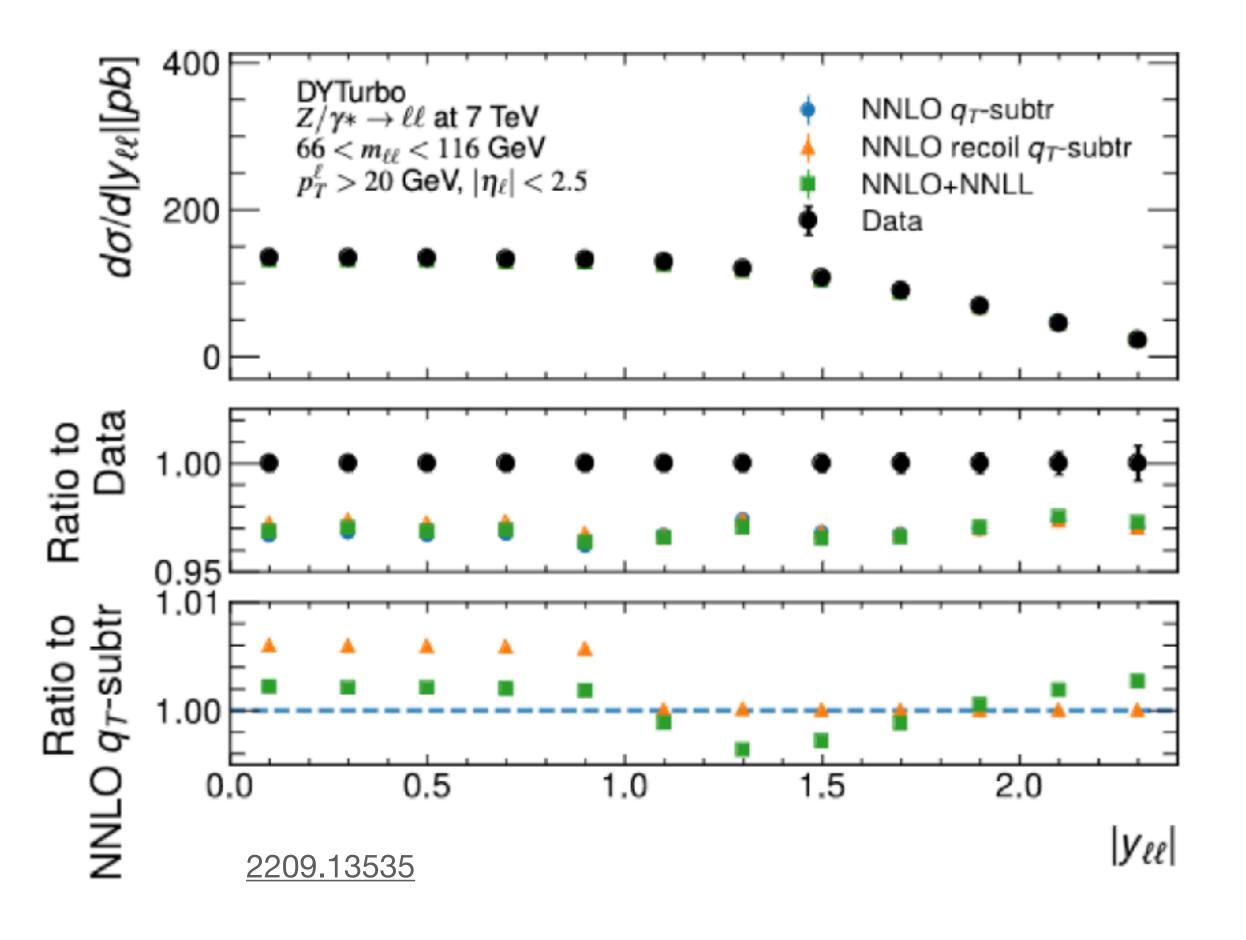
in baseline dataset not considered impact assessed but excluded from baseline

2109.02653



THEORY DEFICIENCIES





A bad description of the data can also come from deficiencies in the theory predictions

Impact of fiducial cuts on theory neglected until a few years ago

Large χ^2 improvement for all PDFs

_		Total χ^2 (ndf=61)		
	PDF set	NNLO	NNLO	NNLO+1
-		$q_T ext{ subtr.}$	recoil q_T -subtr	
_	* •	1		
	CT18ANNLO68	96	84	74
	MSHT20nnlo	111	87	79
-	NNPDF31	91	84	71
	NNPDF40nnlo	89	83	69

In how many other cases are we D blaming data for theory issues?

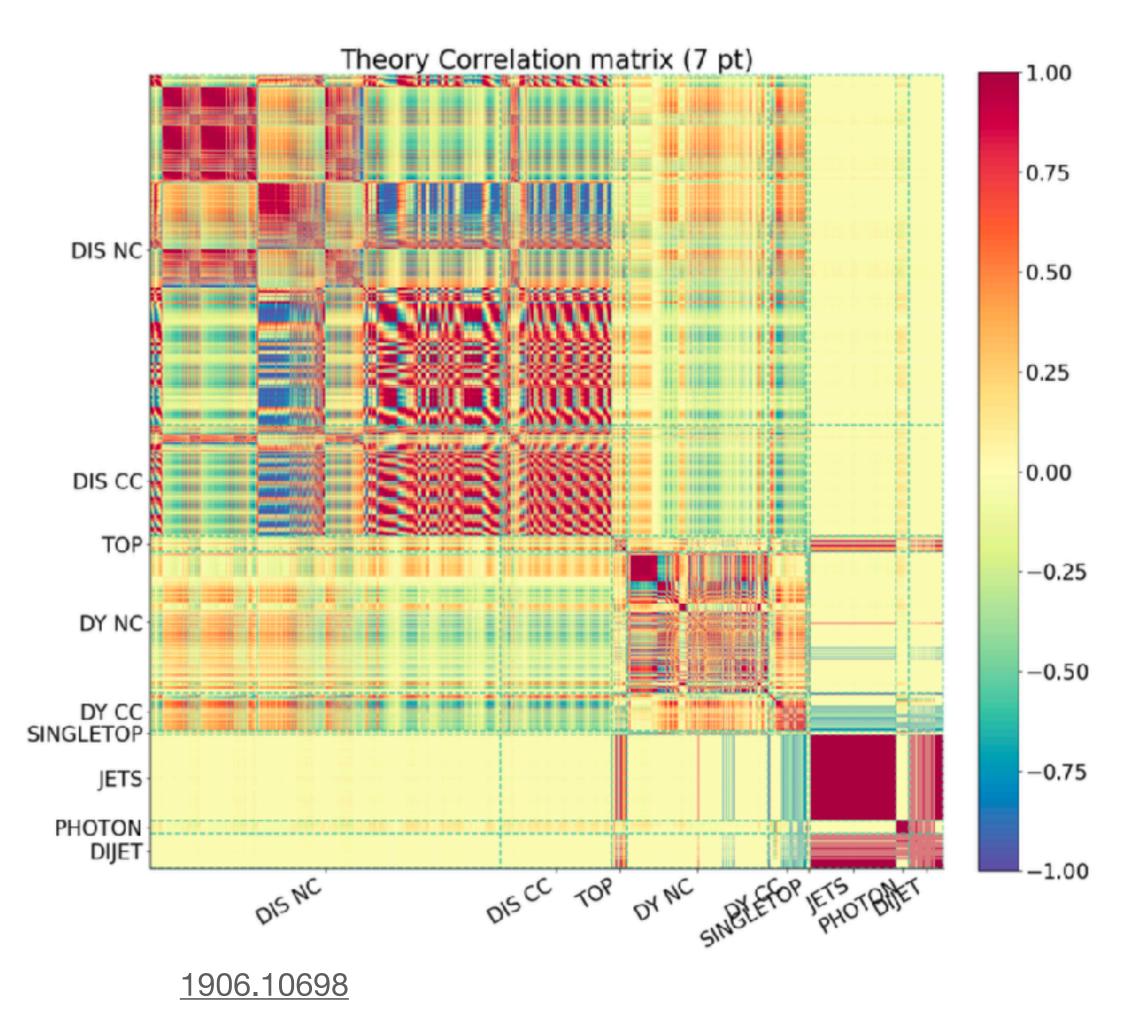






PDFS AND THEORY UNCERTAINTIES

PDF fits historically performed without any theoretical uncertainty The precision of our measurements is clearly challenging this choice



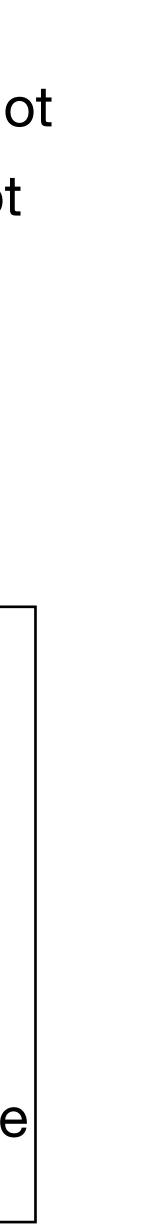
- Attempts by NNPDF and MSHT (at n3lo) to include uncertainties for QCD missing orders
- Shown to give more flexibility to the PDFs and slightly alleviate tensions
- However move the problem to converting scale variations into nuisance parameters (see F.Tackmann, A.Huss, ...)





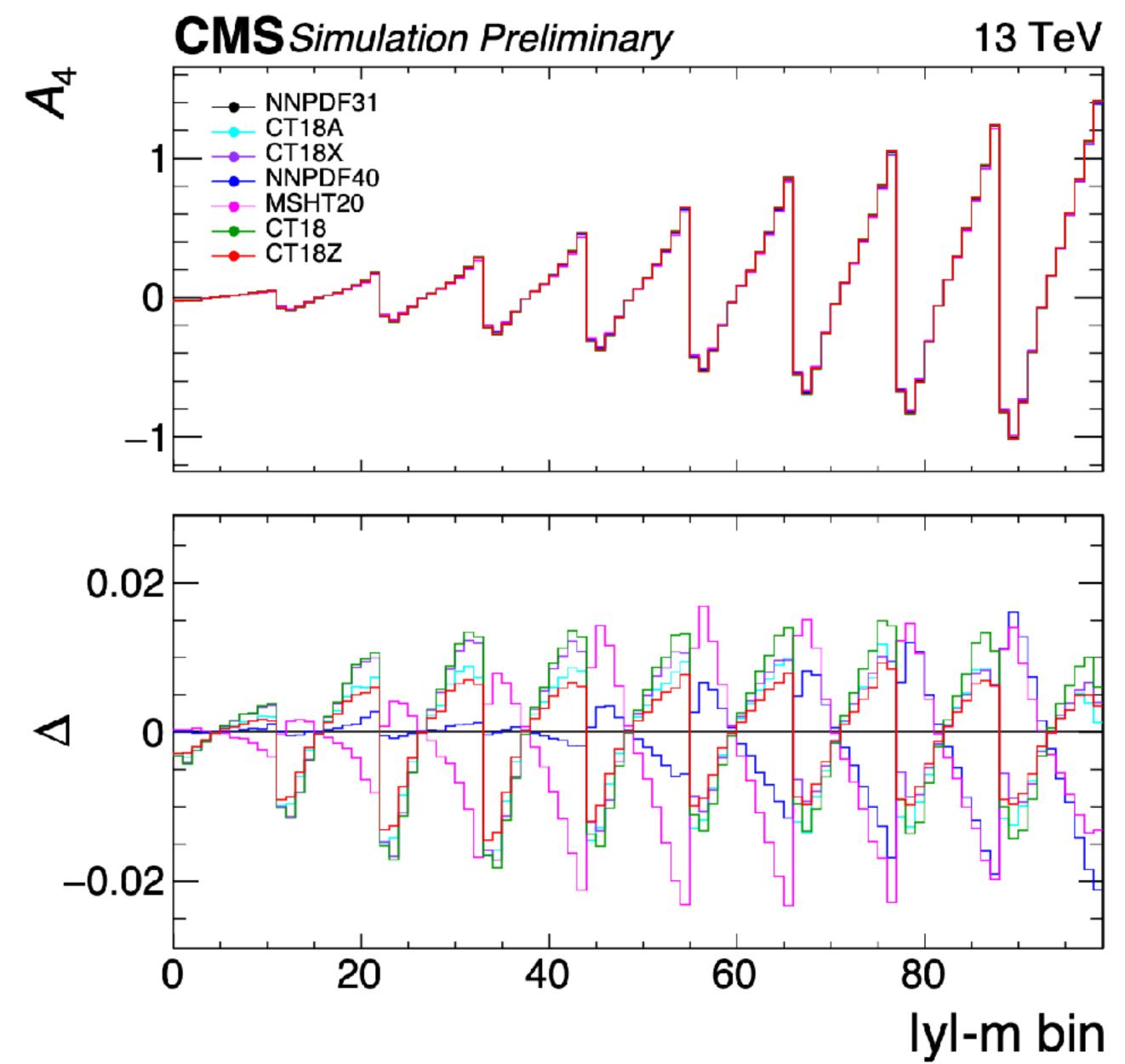
NOT ALL MEASUREMENTS ARE BORN EQUAL

- Certain measurements provide enough of information for reinterpretation, others do not Certain measurements have been cross-checked across multiple channels, others not
- D D
- Certain measurements can be shown to agree with theory, others not Ø
- Certain observables are direct measurements, others extrapolations using theory D (stable tops, parton-level jets, Born-leptons, ...)
- Suggest to identify a subset of precise and self-consistent measurements which we believe to be well described by theory to be used for "reduced data PDF fits"
- Involving both PDF fitting groups and experimental collaborations
- Similar to PDF4LHC benchmarking, but aimed at a deeper understanding of differences in PDFs and alleviate the need for tolerances
- Could consider a "PDF challenge" in which we provide you with pseudo data generated under a known probability distribution (including tensions) and we compare the PDF+uncertainties returned by the various PDF fitting approaches



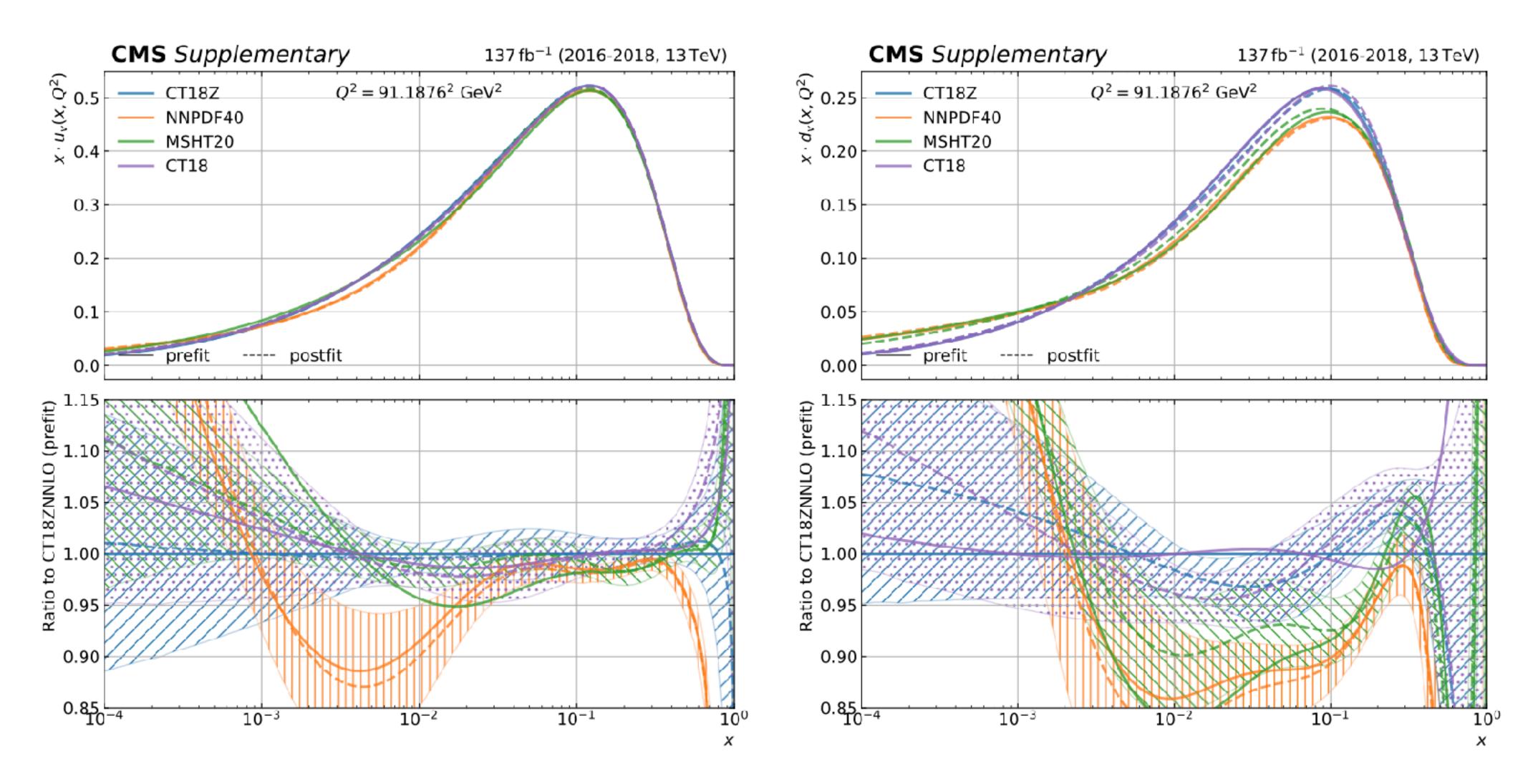


PARTON DISTRIBUTION FUNCTIONS DEPENDENCE OF A4





PRE- AND POST-FIT PDFS COMPARISON



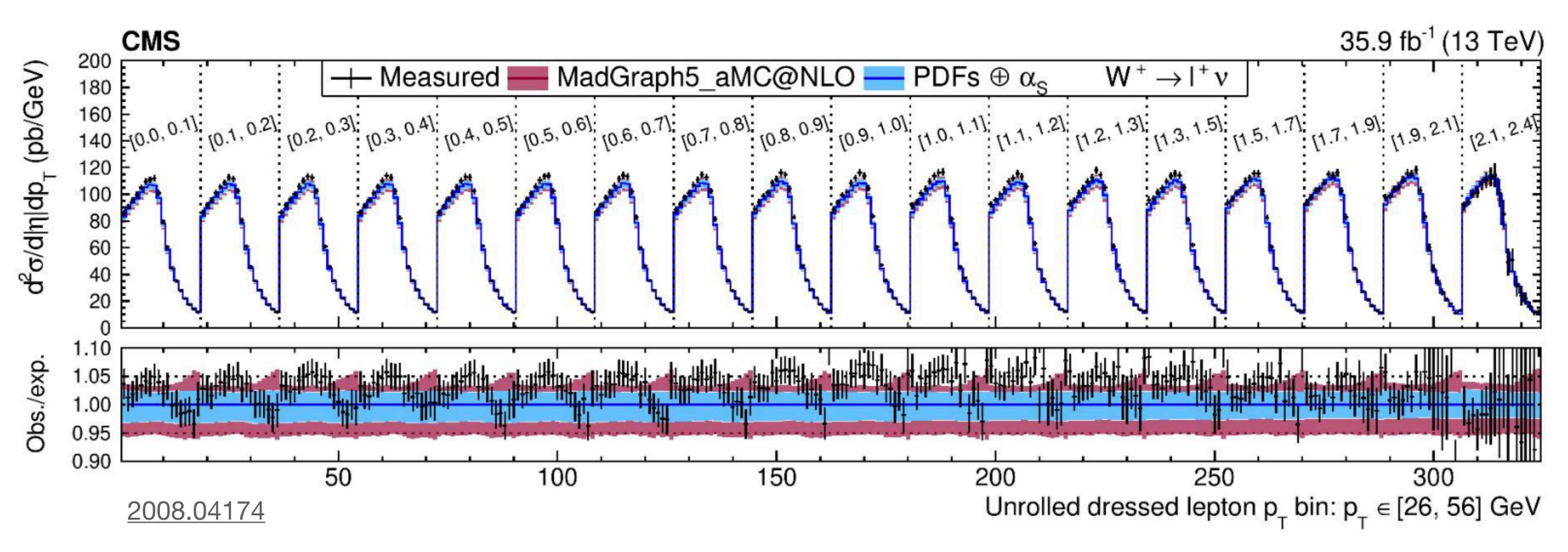
()Important validation of the results of PDF profiling \mathbf{O}

The profiled PDFs are pulled by less than one sigma wrt the original ones



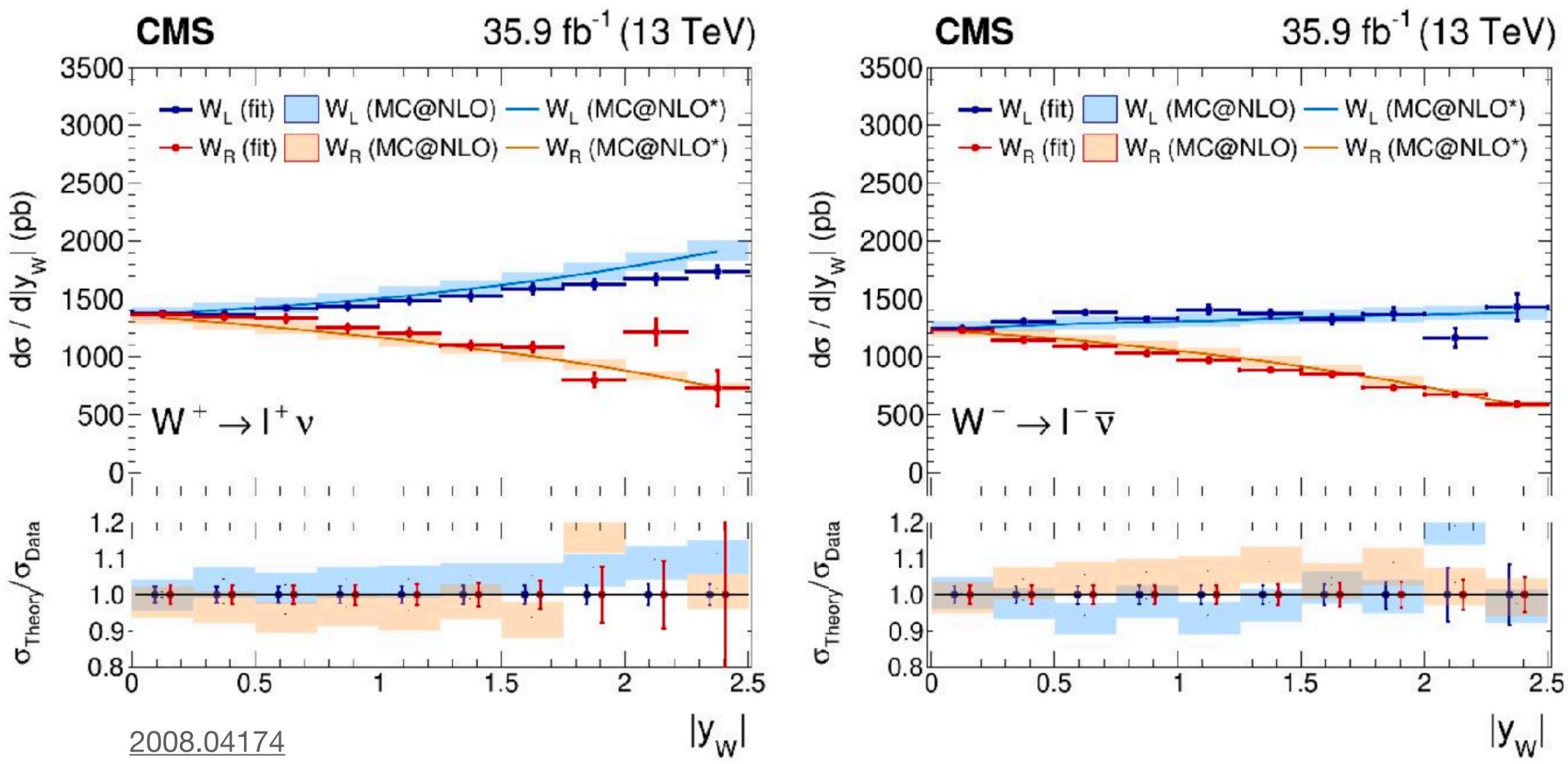
LEPTON PT/ETA CROSS-SECTIONS

- W-boson rapidity and helicity can be inferred statistically from $p_{\rm T}^{\rm lep}$ $\eta^{\rm lep}$ Need predictions with q_T-resummation to describe p_{T}^{lep}





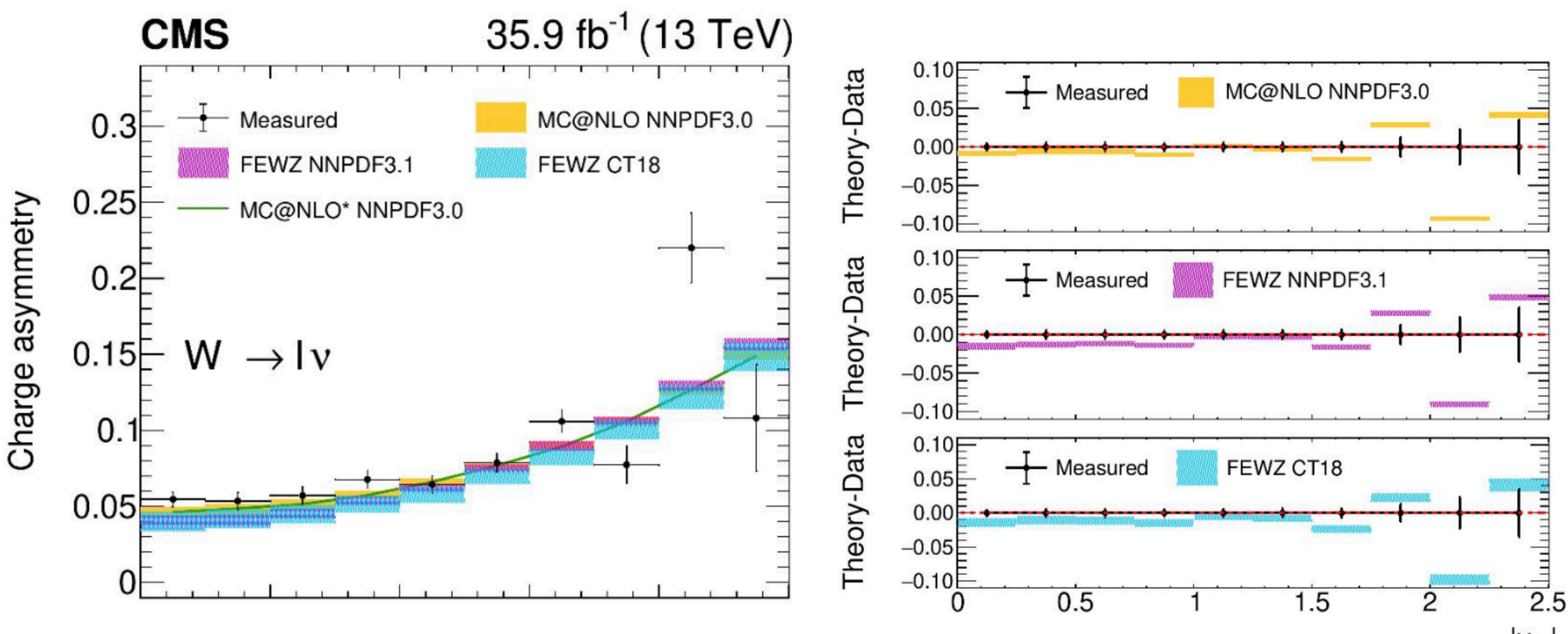
POLARIZED W CROSS-SECTIONS





W CHARGE ASYMMETRY MEASUREMENT

- Ο
- Ο



2008.04174

Helicities in W integrated results measured without assumptions on underlying polarization Avoids circularity in PDF uncertainties in e.g. Tevatron W-asymmetry measurements

y_wl

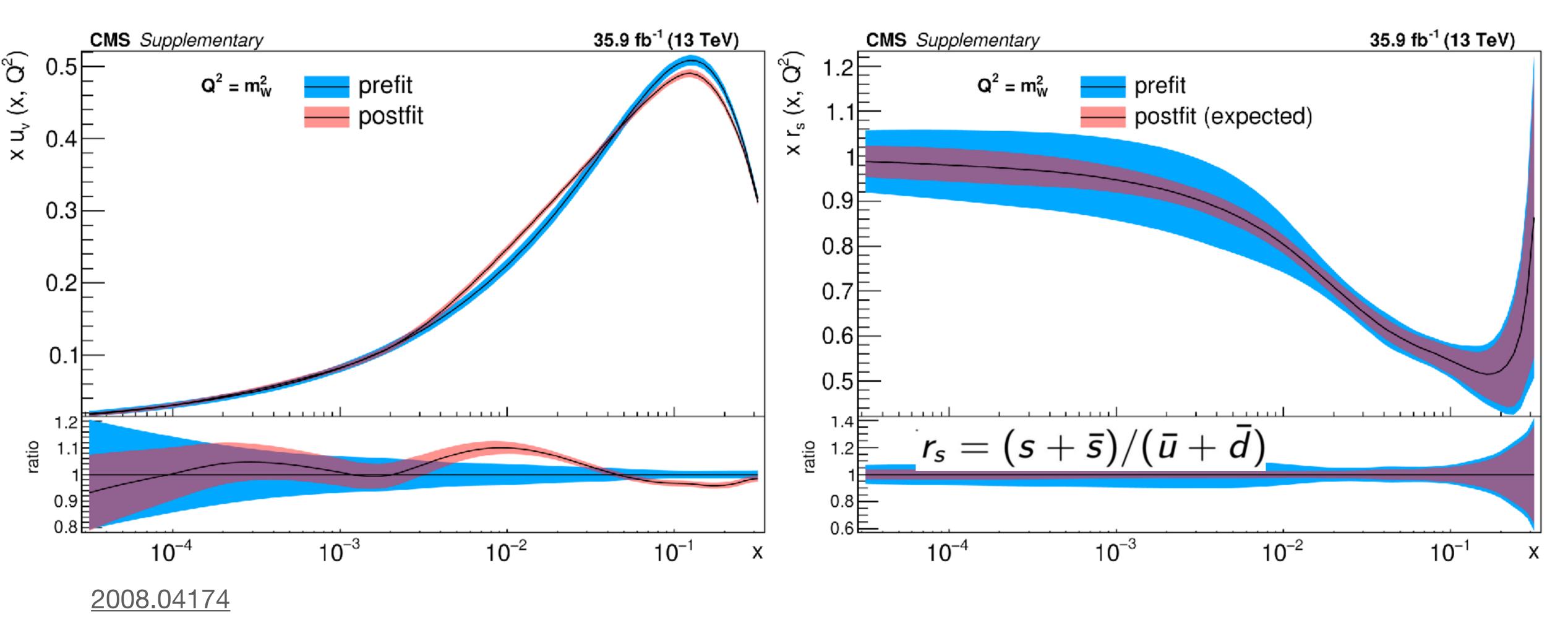






PDF CONSTRAINTS

O Sensitivity to PDFs evaluated using aMC@NLO+Pythia predictions and NNPDF30O Large reduction in uncertainties for valence and strange PDFs

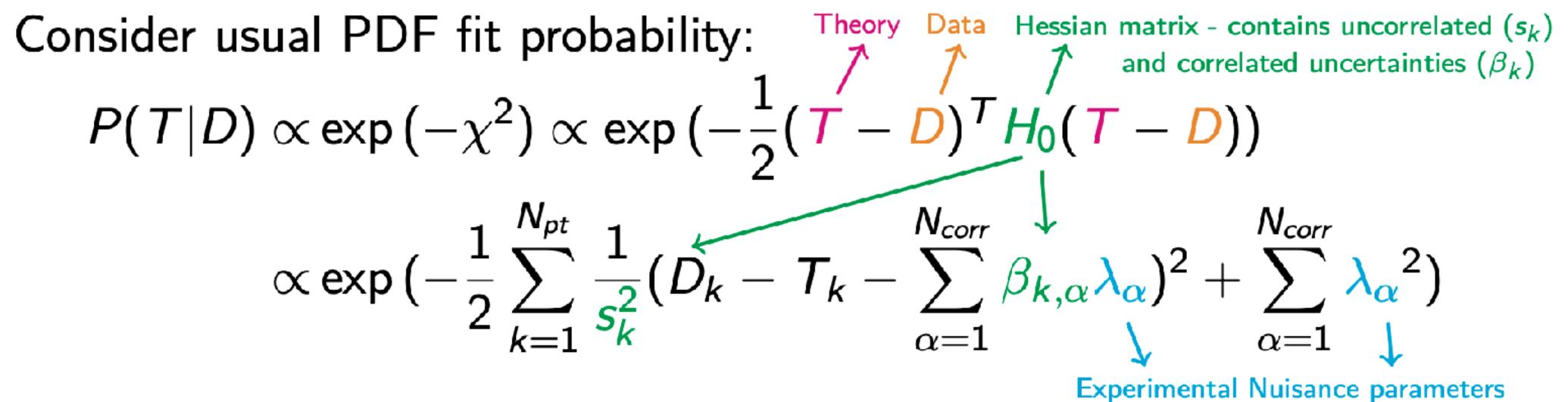




GOODNESS-OF-FIT MEASURE

A /

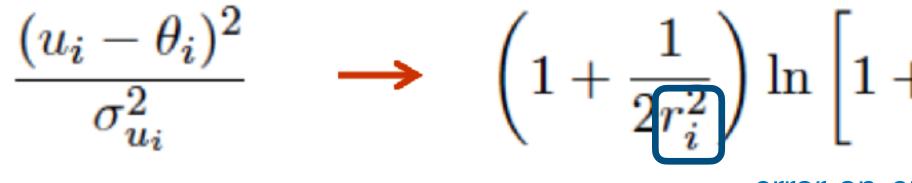
$$\propto \exp\left(-\frac{1}{2}\sum_{k=1}^{N_{pt}}\frac{1}{s_k^2}\right)$$



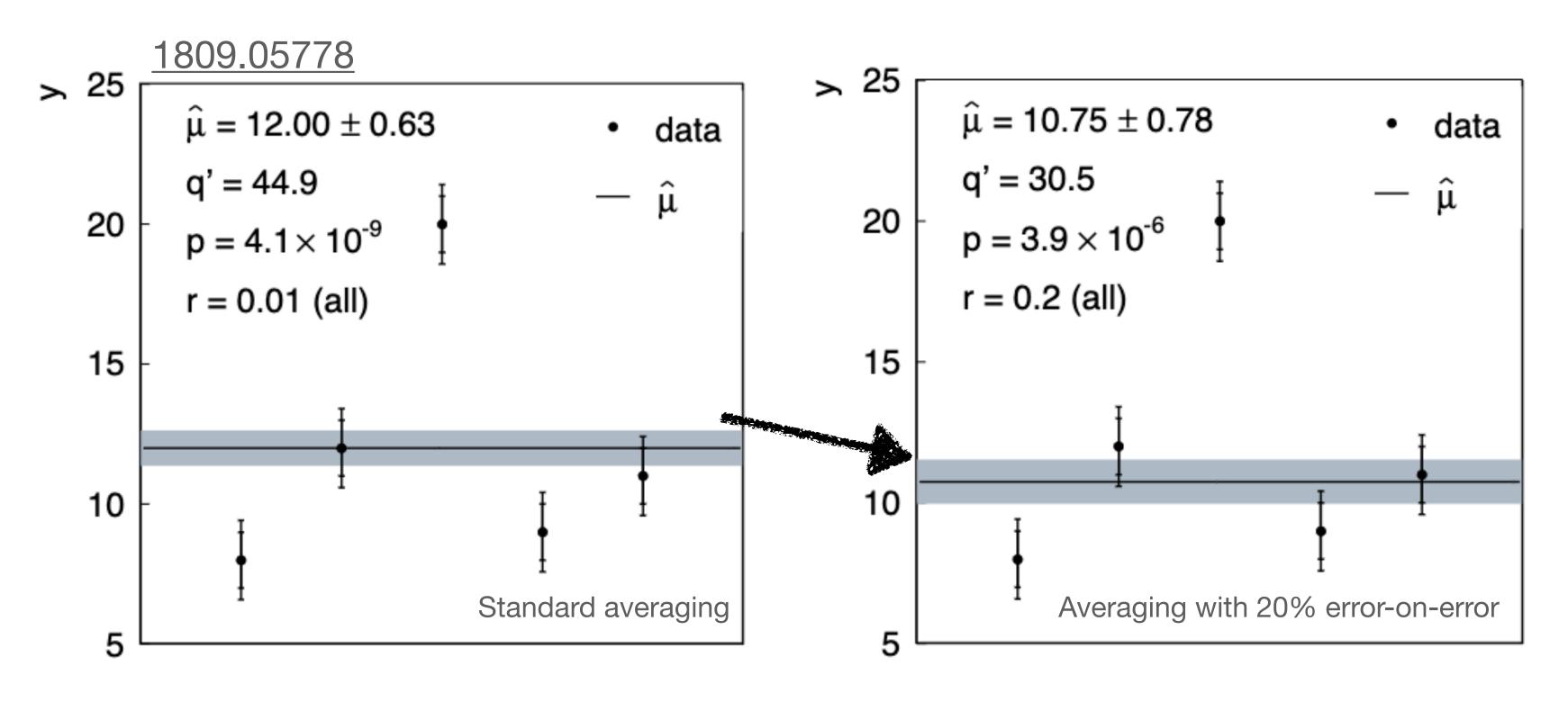


ERRORS ON ERRORS

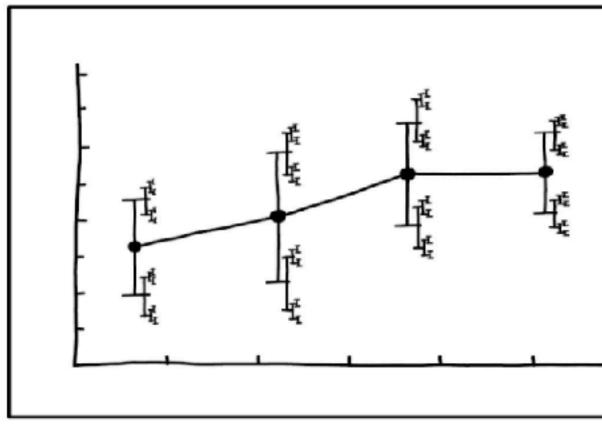
Include an uncertainty on nuisance parameters variances



Replaces Gaussians with Student t distributions with fatter tails



$$+2r_i^2 \frac{(u_i - \theta_i)^2}{v_i}$$



https://xkcd.com/2110/

error-on-error

Small error-on-error (r=20%) sufficient to improve consistency

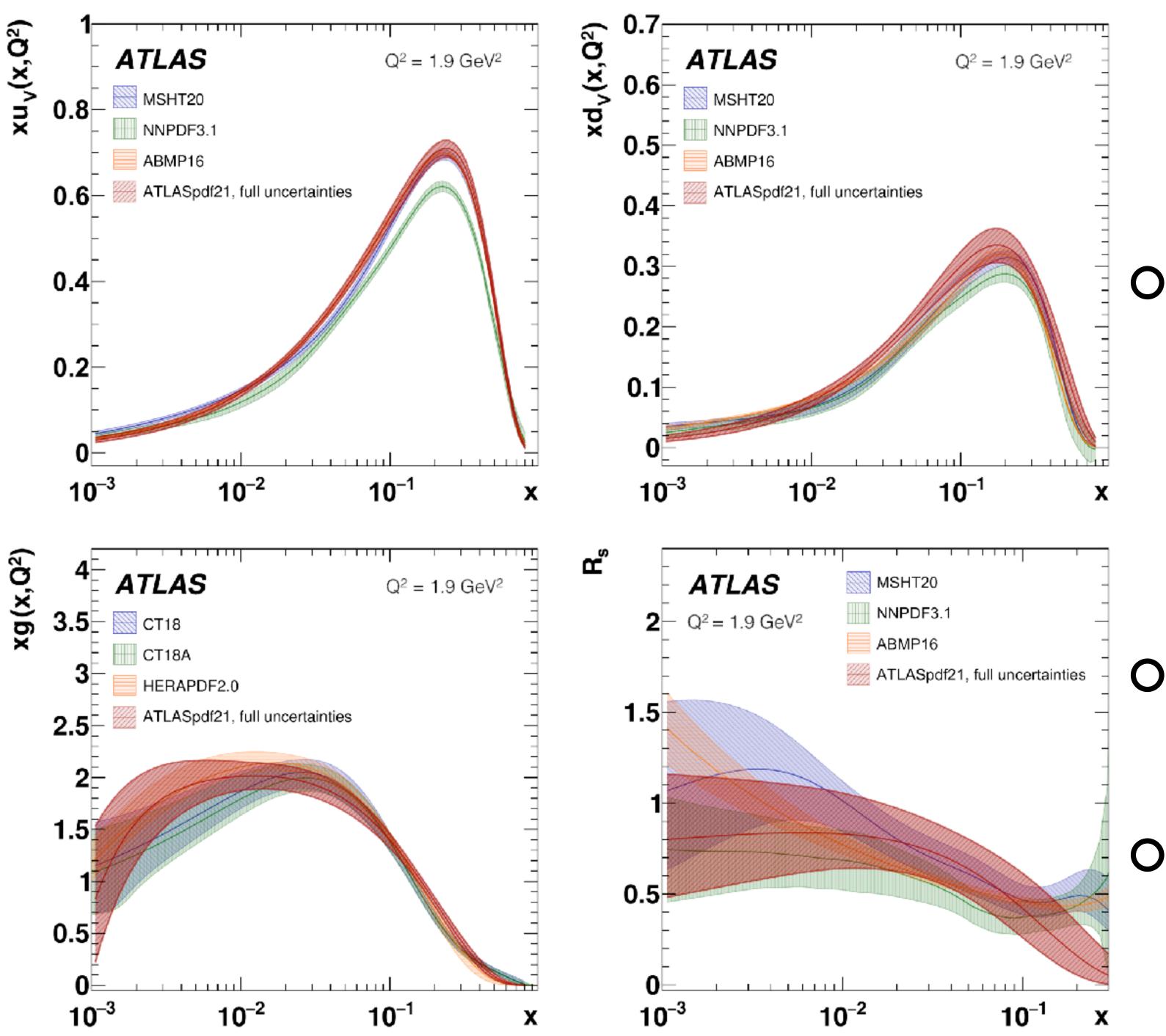
And reduces sensitivity to outlier measurements

How easily could this be tested in a real PDF fit?









ATLASPDF21 RESULTS

Good description of the fitted data

ATLASpdf21	CT18	CT18A			
2010/1620 (1.24)	2135/1641 (1.30)	2133/1641 (1.30)			
MSHT20	HERAPDF2.0	NNPDF3.1			

Significant impact of the ATLAS data on the valence distributions

PDFs in agreement with global fits







