

Overview on the three-body interactions with strangeness

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Three-body dynamics with hyperons

Dynamics of baryons involves formation of hadronic excitations

H.-W. Hammer, S. König, U. van Kolck RMP 92 (2020)





Three-body forces in Effective Field Theories

E. Epelbaum, H.-W. Hammer, U.-G. Meißner, RMP 81, 1773 (2009)



*g*_i constants to be fixed by the **experimental data**



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Ν π π Ν Ν Short-range dynamics (qm) Scattering data ь $\Lambda p \rightarrow \Lambda p$ Sechi-Zorn et al
 Alexander et al. Hauptman et al Piekenbrock 200 Cusp structure: $\Lambda N-\Sigma N$ coupling 100 Talk by L. Serksnyte 0 220 45 135 310 k^* (MeV/c) J. Haidenbauer, U. Meißner, EPJA 56 (2020), 3, 91

Three-body forces in **Effective Field Theories**

E. Epelbaum, H.-W. Hammer, U.-G. Meißner, RMP 81, 1773 (2009)



• Average distances: about 2 fm (Hypertriton: Λ -d about 10 fm)



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The NA and NNA interactions in neutron stars





D. Lonardoni et al., PRL 114 (2019)



Small particle distances can be accessed using femtoscopy!

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Femtoscopy in three-particle system



Correlation function:

$$C(Q_3) = \int S(\rho) |\psi(Q_3, \rho)|^2 \rho^5 d\rho$$

Three-body scattering wave function

Hyper-momentum:

$$Q_3 = 2\sqrt{k_{12}^2 + k_{23}^2 + k_{31}^2}$$

R. Del Grande et al. EPJC 82 (2022) 244 ALICE Coll., EPJ A 59, 145 (2023) Hyper-radius:

$$\rho = 2\sqrt{r_{12}^2 + r_{23}^2 + r_{31}^2}$$

L. E. Marcucci et al., Front. in Phys. 8, 69 (2020).

Extension to three-particle system

- First measurement of the free scattering of three hadrons
- Deviation from unity in p-p-p and p-p- Λ correlation functions



ТЛП

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Source function for three particles in pp collisions

- Source function derived from three independent Gaussian emitters ALICE Coll., Phys. Lett. B 811, 135849 (2020) ALICE Coll., arXiv:2311.14527 (2023)
- In hypersperical coordinates

 $S(\rho) = \frac{1}{\pi^3 \rho_0^6} e^{-\left(\frac{\rho}{\rho_0}\right)^2}$

with $\rho_0 = 2 r_o$ and r_o is two-body source size. A. Kievsky, et al., Phys.Rev.C 109 (2024) 3, 034006

- The value of ρ_0 is determinted from the $m_{\rm T}$ of the pairs in the triplets
- In pp collsions at the LHC small source: $ho_0 =$ 1-3 fm



A. Kievsky, et al., Phys.Rev.C 109 (2024) 3, 034006



- Wave function in hyperspherical harmonics $\Psi(\rho, Q_3) = \sum_{K} R_K(\rho) Y_K(\Omega)$
- First ever full three-body correlation function calculations

three-proton wave function

 $C(Q_3) = \int \rho^5 d\rho \, S(\rho, \rho_0) |\Psi(\rho, Q_3)|^2$ hyperradius

- Interactions:
 - pp strong interaction (AV18)
 - \circ Coulomb
 - \odot No three-body forces

A. Kievsky, et al., Phys.Rev.C 109 (2024) 3, 034006

 p-p-p correlation function: superposition of many partial waves



A. Kievsky, et al., Phys.Rev.C 109 (2024) 3, 034006

Influence of the source size

• Using m_{τ} (source size) differential studies we can probe the interaction with the distances





Comparison Run-2 data

Comparison with the ALICE Run-2 measurement:

 calculations can describe the shape observed in the data

Required improvements:

- source model based on two-body femtoscopy measurements
- feed-down contribution from p-p-Λ must be evaluated
- More precision in the data is required



p-p-p correlation function in Run-3





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- Hadron-nuclei correlations at the LHC can be used to study many-body dynamics







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$p\Lambda$ and $pp\Lambda$ interactions

• AN interaction modelled using Gaussian potentials anchored to scattering parameters

$$V_{p\Lambda}(r) = \sum_{S} V_{S} e^{-(r/r_{s})^{2}} \mathcal{P}_{0,S}$$

NLO19 (600):
$$V_0 = -31.9 \text{ MeV}$$
 $r_0 = 1.4 \text{ fm}$
 $V_1 = -42.1 \text{ MeV}$ $r_1 = 1.1 \text{ fm}$

→ BE($^{3}_{\Lambda}$ H) = 2.904 MeV exp: 2.39 MeV

Binding energy from: https://hypernuclei.kph.uni-mainz.de

• ANN interaction modelled using Gaussian potentials

$$W(r_{13}, r_{23}) = W_3 e^{-(r_{13}^2 + r_{23}^2)/\rho_3^2}$$

 $W_3 = 11.8 \text{ MeV}$ (anchored to ${}^3_{\Lambda}\text{H}$ binding energy) $\rho_3 = 2.0 \text{ fm}$ (anchored to four-body hypernuclei, ${}^4_{\Lambda}\text{H}$ and ${}^4_{\Lambda}\text{He}$)

$p-p-\Lambda$ correlation function

- Reference calculations with NLO19 (600)
- ANN interaction gives 50% effect: only one partial wave (K=0) significantly contributes



Results from recent paper: E. Garrido et al., arXiv: 2408.01750 (2024)

$p-p-\Lambda$ correlation function

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Experimental corrections applied to the theory (feed-down from resonances and

- Gauss NLO19 (600): 50% effect of three-body interactions
- Run-2 data: one data point in the region of the maximum

p-p-Λ correlation function







misidentifications)

Results from recent paper: E. Garrido et al., arXiv: 2408.01750 (2024)



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p∧ **interaction**: scattering + femto



Results from recent paper: E. Garrido et al., arXiv: 2408.01750 (2024)



*Private comunication A. Kievsky and E. Garrido

Conclusions and Outlook

- Exciting results from femtoscopy:
 - o pp collisions at the LHC provides access to the interactions at short distances
 - p-p-p correlation function:
 - ✓ overlap of many partial waves
 - ✓ negligible effect of three-body forces (< 1%)</p>
 - p-d correlation function:
 - ✓ Many-body dynamics can be studied using hadron-deuteron correlations
 - \circ p-p- Λ correlations:
 - \checkmark only one contributing wave in the signal region
 - ✓ 50 % effect of three-body forces
- On-going Run 3 and future Run 4
 - Access to precise data on three-particle correlations
 - Sensitivity to the effect of three-body forces in the correlation functions
- Future: combined analysis of femtoscopy + scattering data + hypernuclei

Backup



Effect of the three-body interaction

• Effect of the three-body interaction for different interaction models: 30% (Usmani) - 80% (NLO13)





The $p\Lambda$ interaction in the femtoscopy era



• Cumulant method provides first hint of effects beyond two-body correlations



- A deviation of $n\sigma = 6.7$ from lower-order contributions
- Theoretical predictions necessary to understand the origin of the deviation further



• Cumulant method provides first hint of effects beyond two-body correlations

R. Kubo, J. Phys. Soc. Jpn. 17, 1100-1120 (1962)



Compatible with lower-order contributions (nσ = 0.8)



Hyperspherical Harmonics formalism

• The Jacobi coordinates:



$$\left\{ egin{array}{l} m{x} = m{r}_2 - m{r}_1 \ m{y} = \sqrt{rac{4}{(1+2m/M)}} \; (m{r}_3 - rac{m{r}_1 + m{r}_2}{2}) \end{array}
ight.$$

 \circ $\,$ We introduce the hyperradius and hyperangle:

$$\rho = \sqrt{x^2 + y^2}$$
 $\phi = \arctan\left(\frac{y}{x}\right)$

 $\,\circ\,\,$ Now the 6 variables are: $(\rho,\phi,\,\hat{\mathbf{x}}\,\,,\,\hat{\mathbf{y}}\,\,)\,\,$ $\,$ 1 radius, 5 angles

Hyperspherical Harmonics formalism

Defining the wave function as:

$$\psi = \sum_{[K]} \rho^{-5/2} u_{[K]}(\rho) Y_{[K]}(\Omega)$$

Schroedinger equation with the interaction:

$$\left(\frac{\partial^2 u_{[K]}(\rho)}{\partial \rho^2} - \frac{(K+3/2)(K+5/2)}{\rho^2} u_{[K]}(\rho)\right) + \sum_{[K']} U_{[K][K']}(\rho) u_{[K']}(\rho) = Q^2 u_{[K]}(\rho)$$

Where the hypercentral potential is obtained as

$$U_{[K][K']}(\rho) = \int d\Omega Y_{[K]}^*(\Omega) [V_{12} + V_{23} + V_{31} + V_{123}] Y_{[K']}(\Omega)$$



Kaon/Proton-deuteron correlation

- Effective two-body system
 - Coulomb + Strong interactions via Lednický model; only s-wave
 - Anchored to scattering experiments
 - Emission source: from m_T scaling

System	Spin averaged		S = 1/2		S = 3/2	
	$a_0(\mathrm{fm})$	$d_0(\mathrm{fm})$	$a_0(\mathrm{fm})$	$d_0(\text{fm})$	$a_0(\mathrm{fm})$	$d_0(\mathrm{fm})$
p-d			$1.30^{+0.20}_{-0.20}$	2 2	$11.40^{+1.80}_{-1.20}$	$2.05^{+0.25}_{-0.25}$
			$2.73_{-0.10}^{+0.10}$	$2.27^{+0.12}_{-0.12}$	$11.88_{+0.40}^{-0.10}$	$2.63_{-0.02}^{+0.01}$
			4.0		11.1	<u></u>
			0.024		13.8	
	-		$-0.13^{+0.04}_{-0.04}$	2 	$14.70^{+2.30}_{-2.30}$	3-32
K ⁺ -d	-0.470	1.75				
	-0.540	0.0				

**R. Lednicky and V. L. Lyuboshits Sov. J. Nucl. Phys. 35 (1982)

$$C(k^*) = 1 + \sum_{S} \rho_S \left[\frac{1}{2} \left| \frac{f(k^*)^S}{r_0} \right|^2 \left(1 - \frac{d_0^S}{2\sqrt{\pi}r_0} \right) + \frac{2\Re f(k^*)^S}{\sqrt{\pi}r_0} F_1(2k^*r_0) - \frac{2If(k^*)^S}{\sqrt{\pi}r_0} F_2(2k^*r_0) \right]$$

R. Lednický, Phys. Part. Nucl. 40, 307(2009)

W. T. H. Van Oers, & K. W. Brockman Jr, NPA 561 (1967); J. Arvieux et al., NPA 221 (1973); E. Huttel et al., NPA 406 (1983);

A. Kievsky et al., PLB 406 (1997); T. C. Black et al., PLB 471 (1999);



Kaon/Proton-deuteron correlation



It works very well for k-d since this interaction is only repulsive and there are no features of the interaction that appears only at short distances. The asymptotic description is sufficient

Proton-deuteron correlation

- The picture of two point-like particles does not work for p-d
 - the deuteron is a composite object
 - Pauli blocking at work for p-(pn) at short distances
 - The asymptotic interaction is different from the short distance one
 - One need a full-fledged three-body calculation





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Pisa model: p-d as three-body system

- Starting with the p-p-n state that goes into p-d state:
 - Nucleons with the Gaussian sources distributions

Single-particle Gaussian emission source

$$(A_d) C_{pd}(k) = \frac{1}{6} \sum_{m_2, m_1} \int d^3 r_1 d^3 r_2 d^3 r_3 S_1(r_1) S_1(r_2) S_1(r_3) |\Psi_{m_2, m_1}|^2 ,$$

- $\Psi_{m_{(,m)}}(x, y)$ three-nucleon wave function asymptotically behaves as p–d state



Calculation done by PISA theory group: Michele Viviani, Alejandro Kievsky and Laura Marcucci

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- $\Psi_{m_i,m_j}(x,y)$ three-nucleon wave function asymptotically behaves as p-d state
- A_d is the deuteron formation probability using deuteron wavefunction



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$$A_{d}C_{pd}(k) = \frac{1}{6} \sum_{m_{2},m_{1}} \int d^{3}r_{1}d^{3}r_{2}d^{3}r_{3} \underbrace{S_{1}(r_{1})S_{1}(r_{2})S_{1}(r_{3})} |\Psi_{m_{2},m_{1}}|^{2},$$

- $\Psi_{m_{(,m)}}(x, y)$ three-nucleon wave function asymptotically behaves as p-d state
- A_d is the deuteron formation probability using deuteron wavefunction
- Final definition of the correlation with p-p source size R_M :

$$A_d C_{pd}(k) = \frac{1}{6} \sum_{m_2, m_1} \int \rho^5 d\rho d\Omega \, \frac{e^{-\rho^2/4R_M^2}}{(4\pi R_M^2)^3} |\Psi_{m_2, m_1}|^2 \, .$$



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• Point-like particle models anchored to scattering experiments

W. T. H. Van Oers et al., NPA 561 (1967); J. Arvieux et al., NPA 221 (1973); E. Huttel et al., NPA 406 (1983); A. Kievsky et al., PLB 406 (1997); T. C. Black et al., PLB 471 (1999);

- Coulomb + strong interaction using Lednický model Lednický, R. Phys. Part. Nuclei 40, 307–352 (2009)
- Only s-wave interaction
- Source radius evaluated using the universal m_{τ} scaling

Point-like particle description doesn't work for p-d



ALICE Coll. arXiv:2308.16120 (2023), accepted by PRX



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Talk by Laura Serksnyte/Anton Riedel 4 Jun, 17:30

p-d correlation function: d as composite object

The three body wave function with proper treatment of 2N and 3N interaction at very short distances goes to a p-d state.

• Three–body wavefunction for p–d: $\Psi_{m_2,m_1}(x, y)$ describing three-body dynamics,

anchored to p-d scattering observables.

- x = distance of p-n system within the deuteron
- y = p-d distance
- m_2 and m_1 deuteron and proton spin

• $\Psi_{m_2,m_1}(x,y)$ three-nucleon wave function asymptotically behaves as p-d state:









Mrówczyński et al Eur. Phys. J. Special Topics 229, 3559 (2020)