



Results on Inclusive Diffraction From The ZEUS Experiment



Presented by B.Loehr on behalf of ZEUS

Data from the running period 1999-2000.

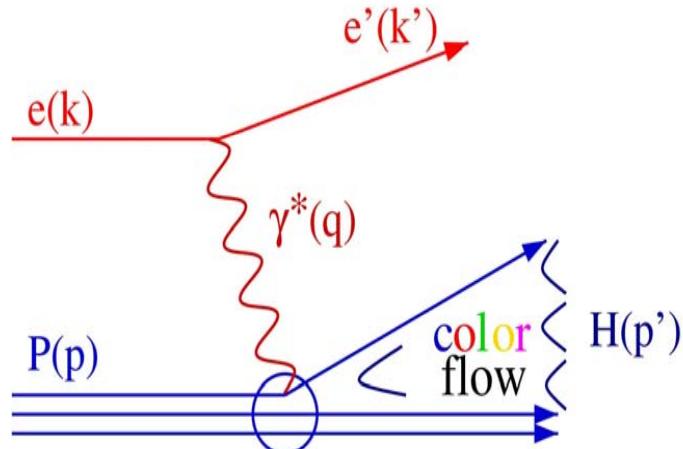
The last period with the ZEUS Forward Plug Calorimeter (FPC) and the Leading Proton Spectrometer (LPS) installed.

Three methods to extract inclusive diffractive events:

- Leading proton spectrometer
- The M_x - method
- Large rapidity gap method

We attempt to get a consistent picture from these three methods using data from the same running period.

Inclusive DIS events :



$$s = (k+p)^2$$

center of mass energy squared

$$Q^2 = -q^2 = -(k-k')^2$$

virtuality, size of the probe

$$W^2 = M_H^2 = (p+q)^2$$

γ^* - proton cms energy squared

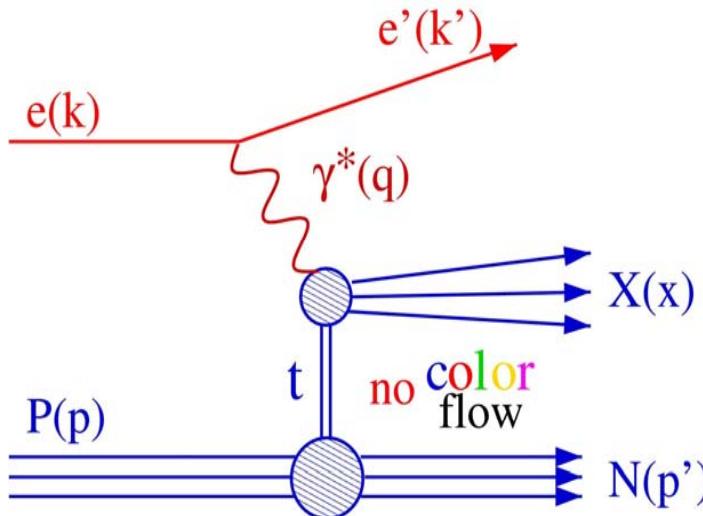
$$x = \frac{Q^2}{2p \cdot q} \quad y = \frac{p \cdot q}{p \cdot k}$$

x : fraction of the proton carried by the struck parton

$$Q^2 = x \cdot y \cdot s$$

y : inelasticity, fraction of the electron momentum carried by the virtual photon

Diffractive DIS events :



For diffractive events in addition 2 variables

$$M_x$$

mass of the diffractive system x

$$t = (p-p')^2$$

four-momentum transfer squared at the proton vertex

$$x_{IP} = \frac{(p-p') \cdot q}{p \cdot q} = \frac{M_x^2 + Q^2}{W^2 + Q^2}$$

momentum fraction of the proton carried by the Pomeron

$$\beta = \frac{Q^2}{2(p-p') \cdot q} = \frac{x}{x_{IP}} = \frac{Q^2}{M_x^2 + Q^2}$$

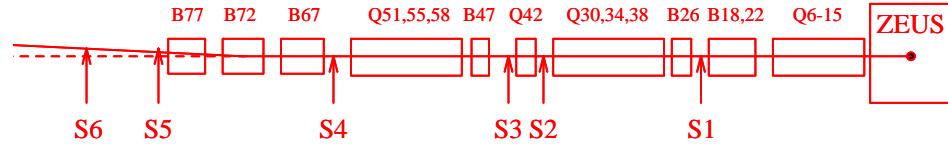
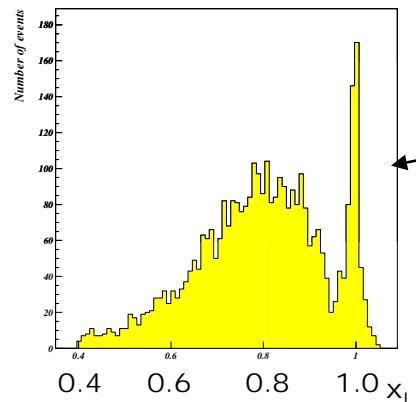
fraction of the Pomeron momentum which enters the hard scattering



Extraction of diffractive events (I)



1.) Forward proton detection

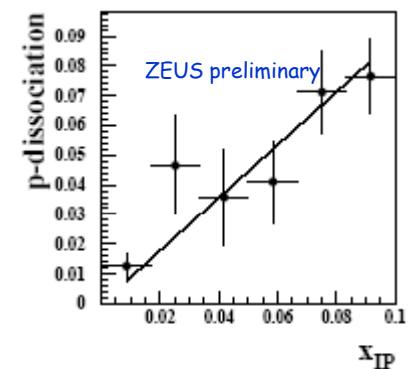


$$x_L \approx 1 \rightarrow \text{diffractively scattered proton} ; \quad x_L \approx 1 - x_{IP}$$

$$t = -\frac{p_T^2}{x_L} - \frac{(1-x_L)^2}{x_L} M_p^2 \quad \text{the only method to measure the } t\text{-distribution}$$

Forward proton tagged events are practically free of proton dissociation background.

They contain, however, contributions from Reggeon exchange at high x_{IP} or low x_L .

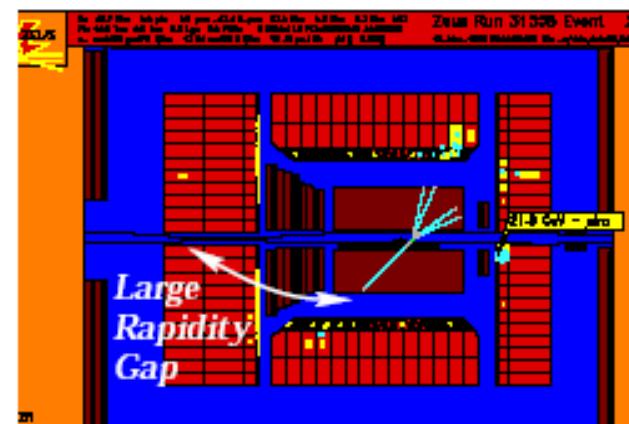


2.) The large (pseudo)rapidity (η_{max}) method

$$\eta_{max} = -\ln \tan(\Theta_{min}/2)$$

No tracks or energy deposits in calorimeter for rapidities greater than η_{max} or at angles less than Θ_{min} .

Events tagged by a large rapidity are dominated by diffraction but they contain contributions from proton dissociation and from Reggeon exchange.

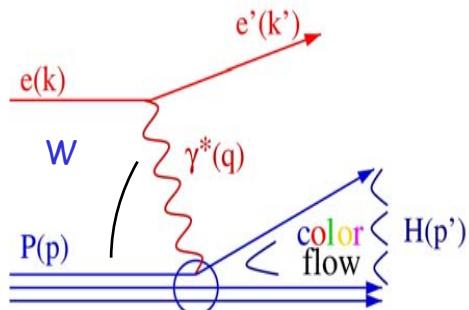


3.) The M_x -method

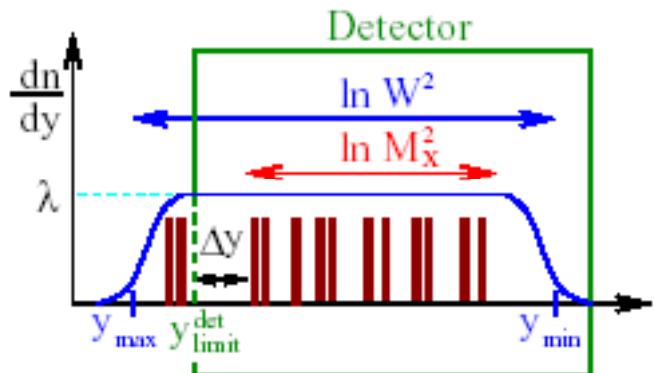
(i) Nondiffractive events :

Rapidity $y = \frac{1}{2} \ln \frac{E+p_z}{E-p_z}$

Property of a produced particle



Uncorrelated particle emission between incoming p-direction and scattered quark.



$$W^2 = c_0 e^{y_{\max} - y_{\min}}$$

$$M_x^2 = c_0 e^{y_{\text{limit}} - y_{\min}}$$

Poisson distr. for Δy in nondiffractive events

$$\frac{dN_{\text{part}}}{dy} = \lambda = \text{const.}$$

$$P(0) = e^{-\lambda \Delta y}$$

$$\frac{dN_{\text{nondiff}}}{d \ln M_x^2} = c \cdot e^{b \cdot \ln M_x^2}$$

(ii) Diffractive events :

$$\frac{dN_{\text{diff}}}{dM_x^2} \propto \frac{1}{(M_x^2)^n}$$

At high energies and
not too low M_x
 $n \approx 1$



$$\frac{dN_{\text{diff}}}{d \ln M_x^2} \approx \text{const.}$$



Extraction of diffractive events (III)



(iii) Nondiffractive + diffractive contributions

$$\frac{dN}{d \ln M_X^2} = D + c \cdot e^{b \cdot \ln M_X^2}$$

D is the diffractive contribution

Two approaches :

1.) take $D = \text{const.}$ for a limited range
in $\ln M_X^2$

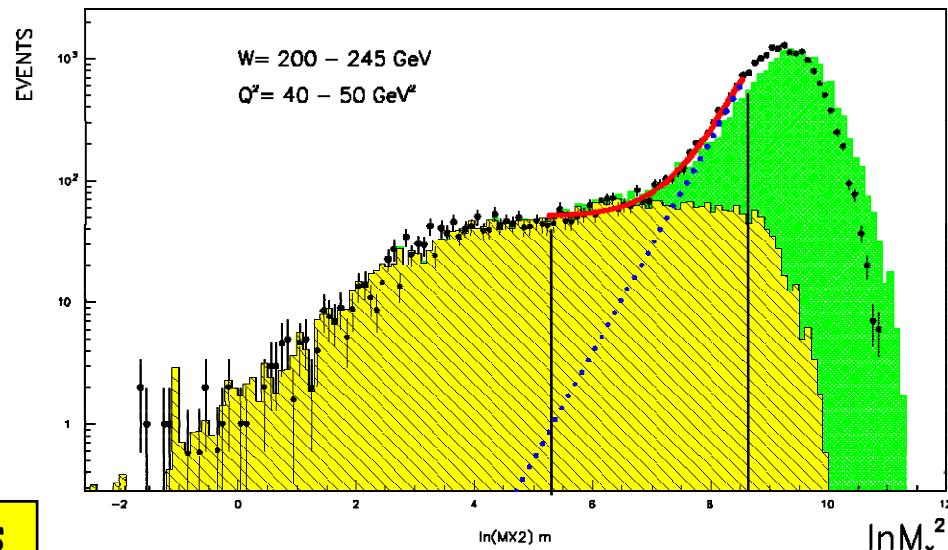
2.) take D from a BEKW-model (*see later*)
parametrization which describes
our measured data.
This is an iterative procedure.

Fit slope b, c and D

for $\ln M_X^2 \leq \ln W^2 - \eta_0$

Determine diffractive events
by subtracting nondiffractive
events from measured data bin
by bin as calculated from
fitted values b and c.

Both approaches give the same results

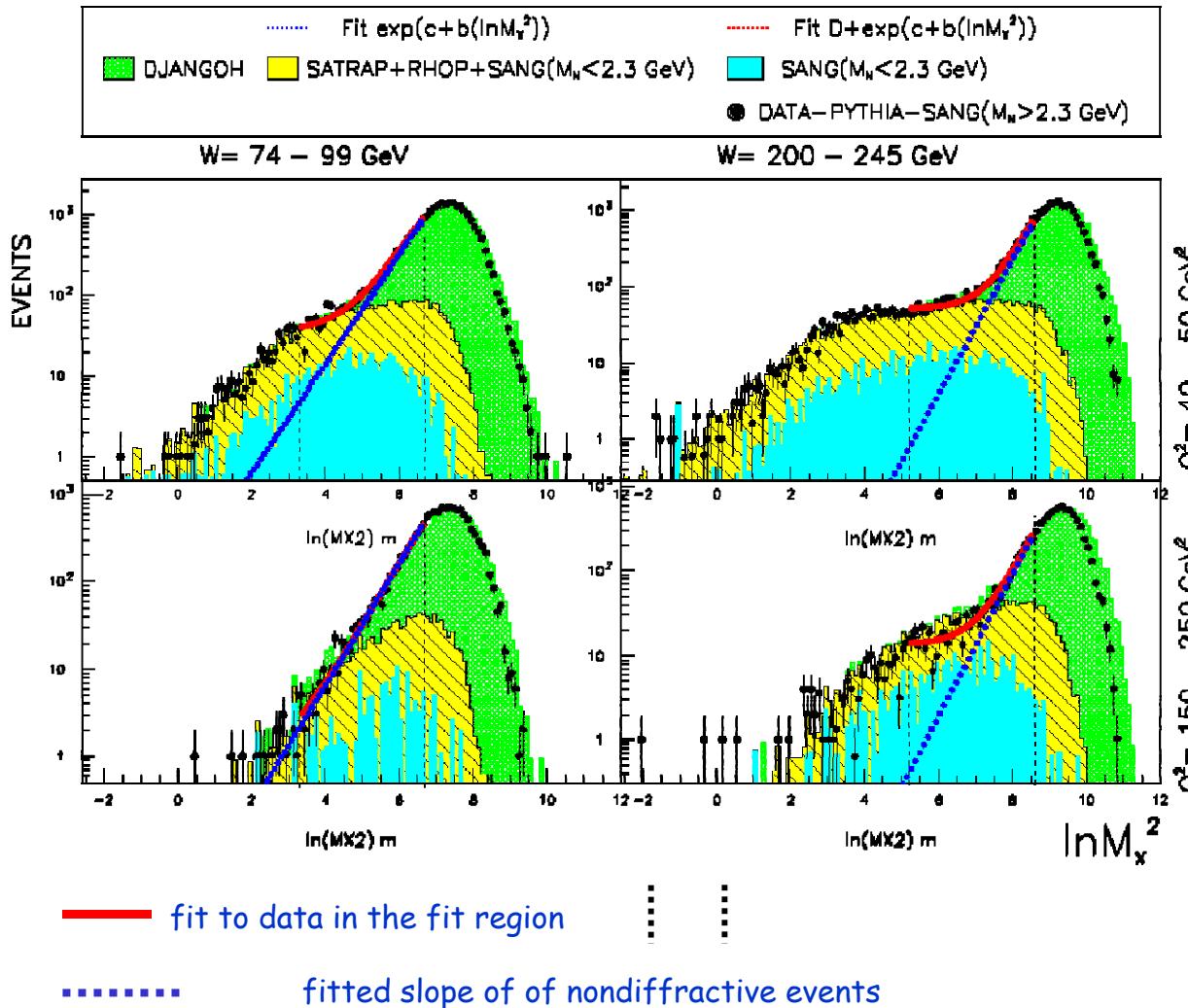




The ZEUS M_X -Analysis (I)



Example of $\ln M_X$ -distributions for four kinematical bins :



Diffractive data selected by the M_X -method contain proton dissociative events but no contributions from Regge exchange

MC-simulation :

nondiffractive : DJANGOH

diffractive : SATRAP

proton diss.: SANG

SANG adjusted to fit data which are dominated by proton dissociation

Proton dissociation can be reliably calculated for $M_N > 2.3$ GeV and has been subtracted from data

The ZEUS M_X -results contain contributions from proton dissociation for masses $M_N < 2.3$ GeV.



ZEUS M_X - data from 1998 - 2000 (II)

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ZEUS-99-PR01-00-prep(9.)

Mx 98-99 : *

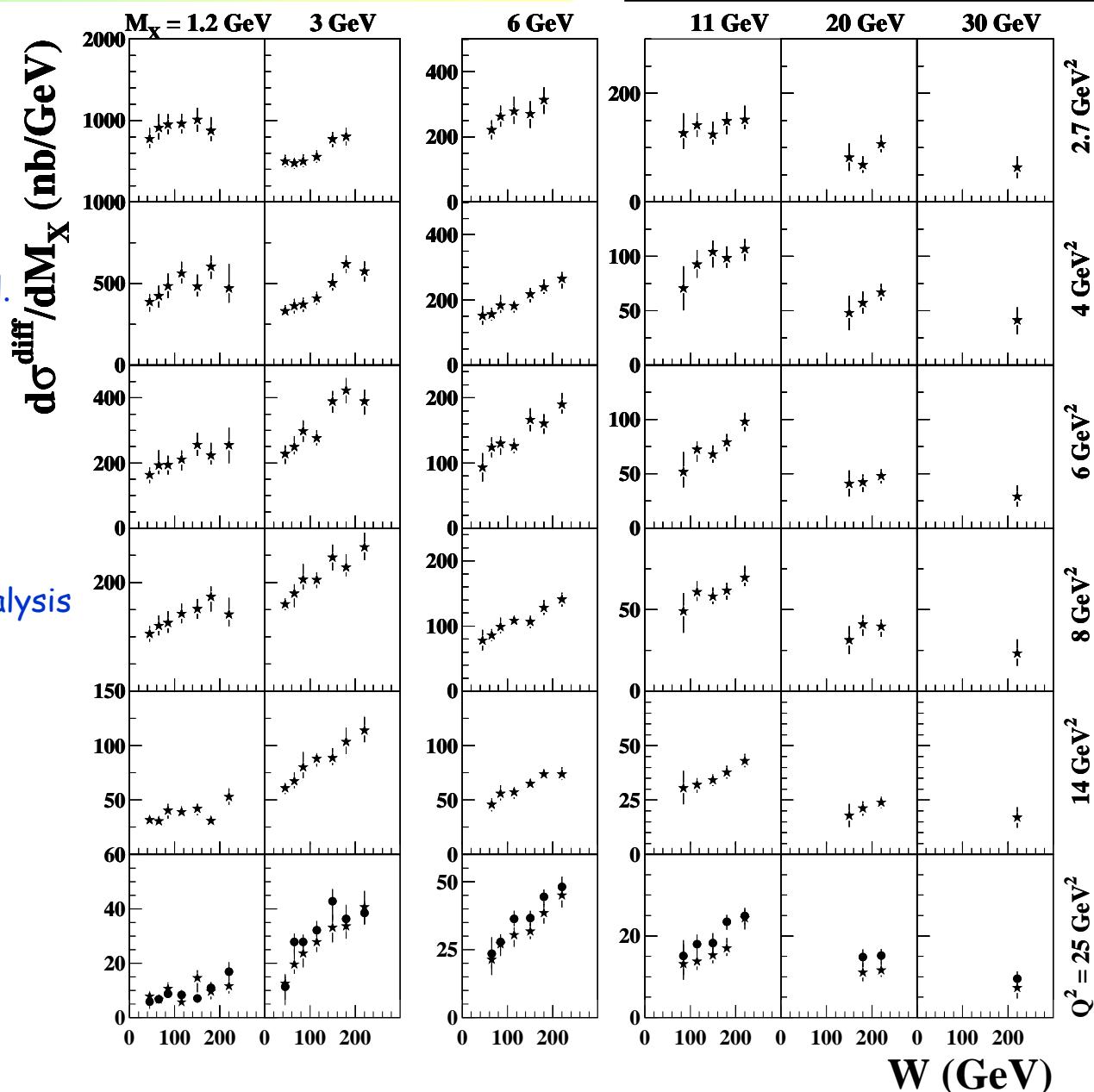
Published data from
1998-1999 period

(ZEUS Coll., S.Chekanov et al.
Nucl. Phys B 713, 3 (2005))

Prel. Mx 99-00: ●

Preliminary results from
1999-2000 period.
Extension of Mx 98-99 analysis
to higher Q^2 .

Mx 98-99 and Mx 99-00
analyses have common bin
at $Q^2 = 25 \text{ GeV}^2$





ZEUS M_X - data from 1998 - 2000 (III)



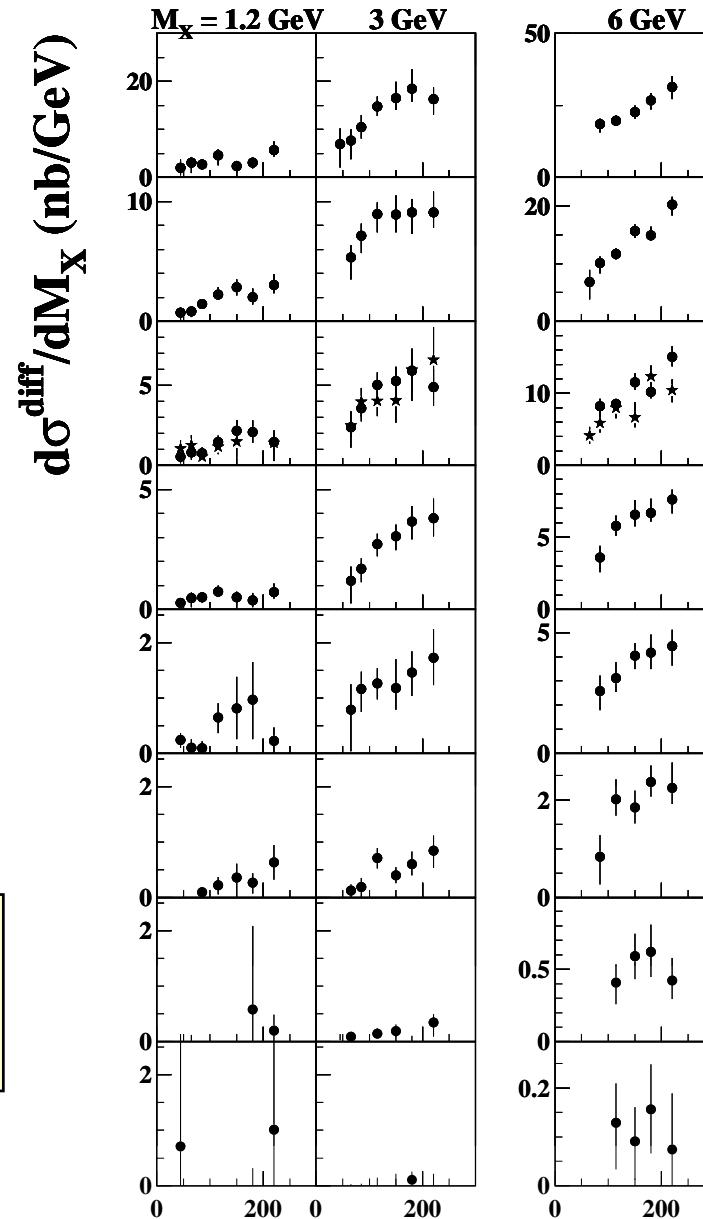
M_X 98-99: *

Prel.

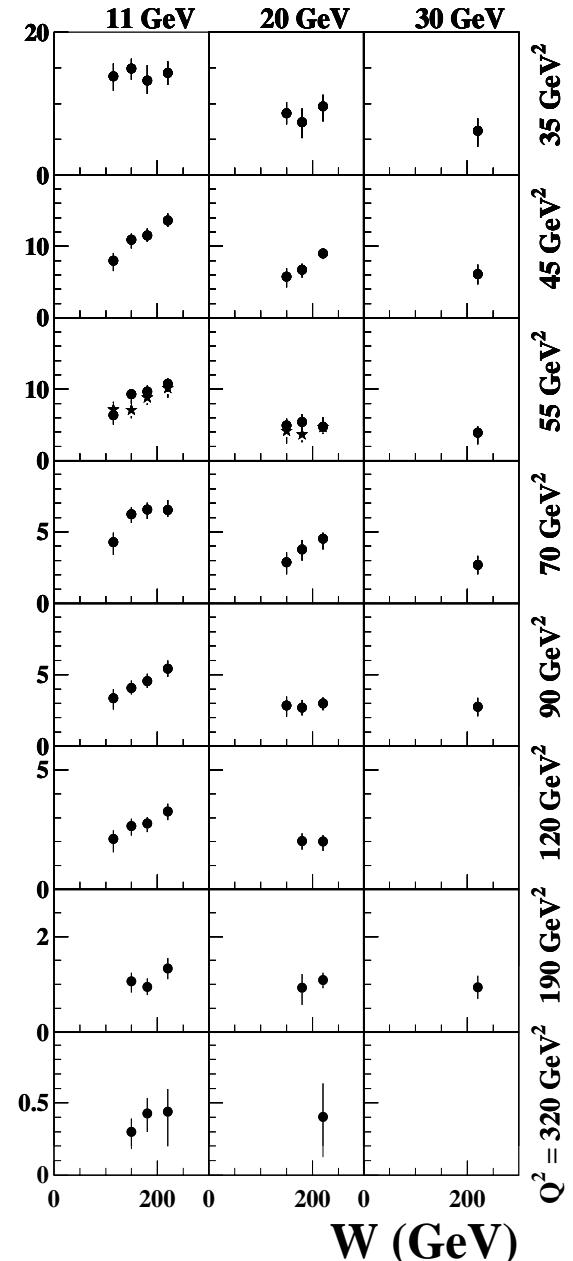
M_X 99-00: ●

M_X 98-99 and M_X 99-00
analyses have common bin
at $Q^2 = 55 \text{ GeV}^2$

Within syst. errors good
agreement between
 M_X 98-99 and
 M_X 99-00 results



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ZEUS M_X 98-99, ZEUS M_X 00 (prel.)



Inclusive DIS:

For small x , F_2 rises rapidly as $x \rightarrow 0$

$$F_2 = c \cdot x^{-\lambda} \quad W \propto \frac{1}{x}$$

$$\lambda = \alpha_{IP}(0) - 1$$

Inclusive diffractive DIS:

$$\frac{d\sigma_{\gamma^* p \rightarrow XN}}{dM_X} = h \cdot \left(\frac{W}{W_0} \right)^{\alpha_{IP}^{diff}}$$

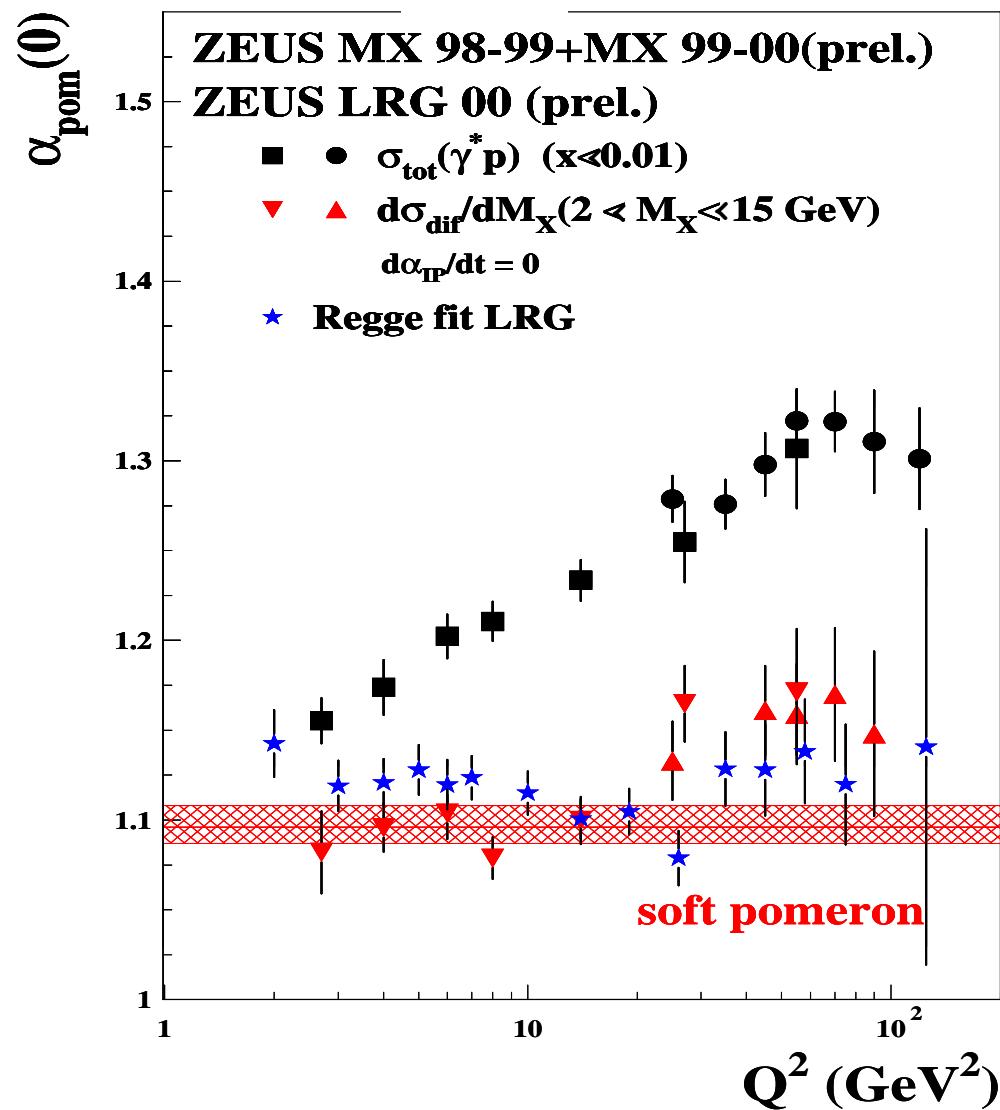
$$\bar{\alpha}_{IP} = 1 + \frac{\alpha_{IP}^{diff}}{4} \text{ averaged over } t$$

$$\alpha_{IP}(t) = \alpha_{IP}(0) + \alpha'_{IP} \cdot t$$

$$\frac{d\sigma}{dt} = f(t) \cdot e^{2(\alpha_{IP}(t)-1) \cdot \ln\left(\frac{W}{W_0}\right)^2}$$

$$\frac{d\sigma}{dt} \propto e^{A \cdot t} \text{ for small } t.$$

take $A = 7.9 \pm 0.5(\text{stat.})^{+0.9}_{-0.5}(\text{syst.}) \text{ GeV}^2$
as measured by ZEUS LPS



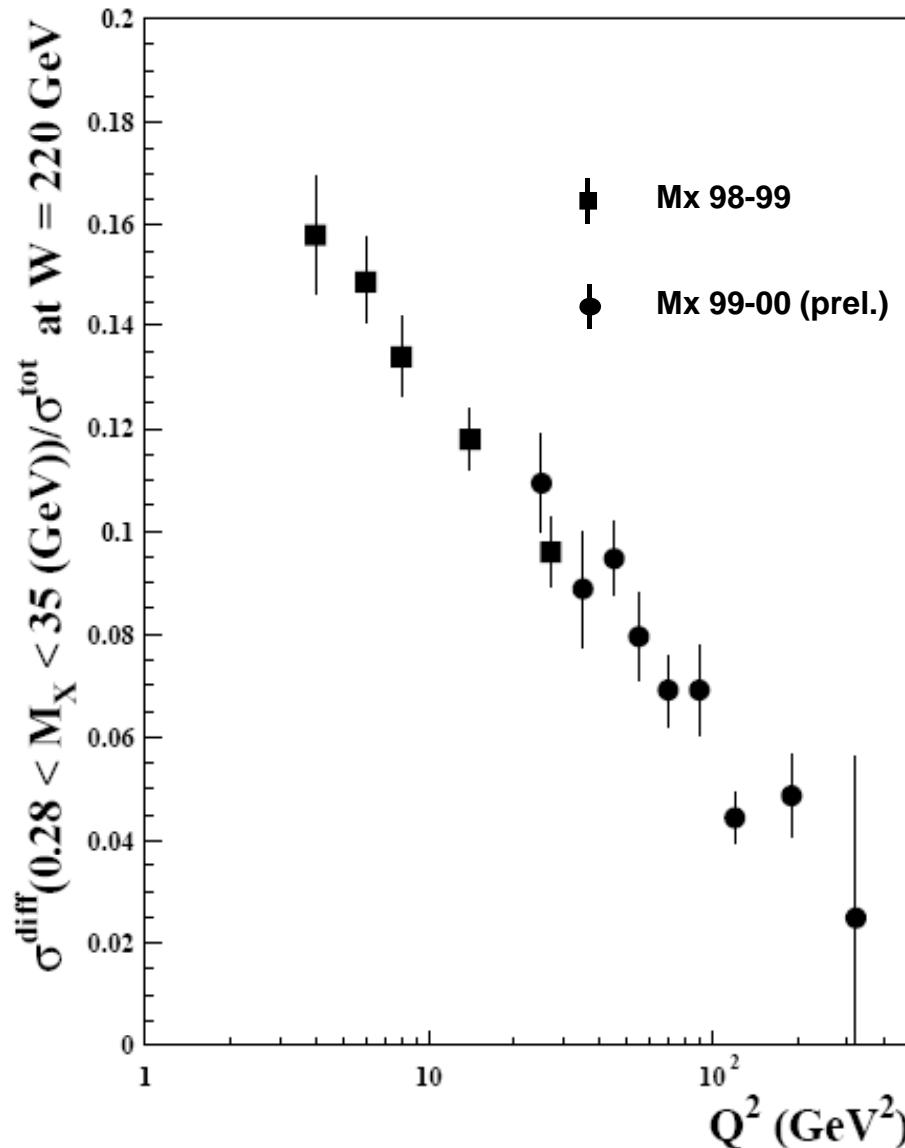
Inclusive DIS and inclusive diffractive DIS are not described by the same 'Pomeron'.



Ratio of total diffractive cross-section to total DIS cross-section



Ratio plotted at $W=220 \text{ GeV}$ because only there the full M_X range is covered by measurements



$$r = \sigma^{\text{diff}}(0.28 < M_X < 35 \text{ GeV}) / \sigma^{\text{tot}}$$

Within the errors of the measurements
r is independent of W.

At $W=220 \text{ GeV}$, r can be fitted by

$$r = 0.22 - 0.034 \cdot \ln(1+Q^2)$$

This logarithmic dependence of the ratio
of total diffractive cross-section to the
total DIS cross section indicates that
diffraction is a leading twist process
for not too low Q^2 .



Diffractive Cross-Section and Diffractive Structure Functions



$$\frac{d^4\sigma}{dQ^2 dt dx_{IP} d\beta} = \frac{2\pi\alpha_{em}}{\beta Q^2} [1 - (1-y)^2] \cdot F_2^{D(4)}(Q^2, t, x_{IP}, \beta)$$

ZEUS neglects the contribution from longitudinal structure function

H1 defines : sizable only at high y

$$\sigma_r^D = F_2^D - \overbrace{\frac{y^2}{1+(1-y)^2} F_L^D}$$

If t is not measured :
(LRG and M_x -method)

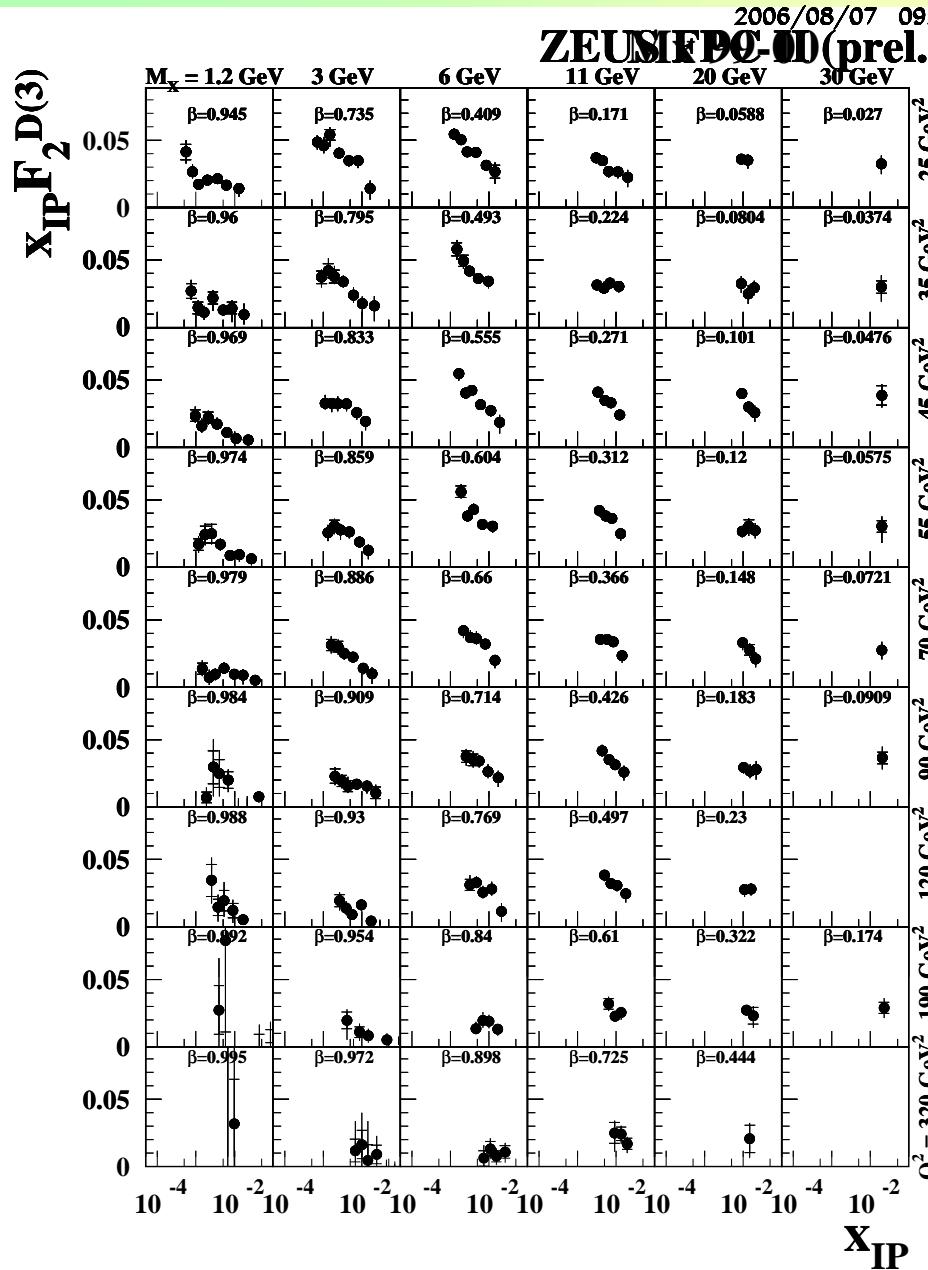
$$\frac{d^3\sigma_{\gamma^* p \rightarrow XN}^{diff}}{dQ^2 d\beta dx_{IP}} = \frac{2\pi\alpha^2}{\beta Q^4} [1 + (1-y)^2] \cdot F_2^{D(3)}(\beta, x_{IP}, Q^2)$$

$$\frac{1}{2M_X} \frac{d\sigma_{\gamma^* p \rightarrow XN}^{diff}(M_X, W, Q^2)}{dM_X} = \frac{4\pi^2\alpha}{Q^2(Q^2 + M_X^2)} x_{IP} F_2^{D(3)}(\beta, x_{IP}, Q^2)$$

If $F_2^{D(3)}(\beta, x_{IP}, Q^2)$ is interpreted in terms of quark densities, it specifies the probability to find in a proton which undergoes a diffractive interaction a quark carrying a fraction $x = \beta x_{IP}$ of the proton momentum.



$x_{IP}F_2^D(3)$ Results from the Mx 99-00 Analysis



ZEUS Mx 99-00 (prel.)

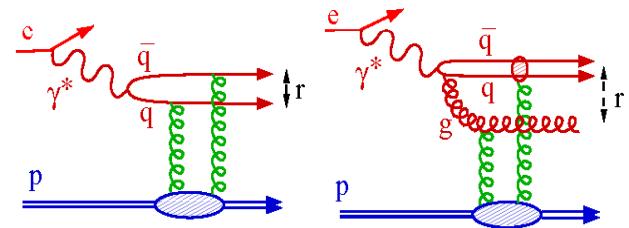


ZEUS modified BEKW Fit



Fit with BEKW model

(Bartels, Ellis, Kowalski and Wüsthoff, 1998)



- $x_{IP} F_2^{D(3)} = c_T \cdot F_{q\bar{q}}^T + c_L \cdot F_{q\bar{q}}^L + c_g \cdot F_{q\bar{q}g}^T$

$$F_{q\bar{q}}^T = \left(\frac{x_0}{x_{IP}}\right)^{n_T(Q^2)} \cdot \beta(1 - \beta),$$

$$F_{q\bar{q}}^L = \left(\frac{x_0}{x_{IP}}\right)^{n_L(Q^2)} \cdot \frac{Q_0^2}{Q^2 + Q_0^2} \cdot [\ln(\frac{7}{4} + \frac{Q^2}{4\beta Q_0^2})]^2 \cdot \beta^3 (1 - 2\beta)^2,$$

$$F_{q\bar{q}g}^T = \left(\frac{x_0}{x_{IP}}\right)^{n_g(Q^2)} \cdot \ln(1 + \frac{Q^2}{Q_0^2}) \cdot (1 - \beta)$$

assume $n_T(Q^2) = c_4 + c_7 \ln(1 + \frac{Q^2}{Q_0^2})$, $n_L(Q^2) = c_5 + c_8 \ln(1 + \frac{Q^2}{Q_0^2})$,

$$n_g(Q^2) = c_6 + c_9 \ln(1 + \frac{Q^2}{Q_0^2})$$

The ZEUS data support taking $n_T(Q^2) = n_g(Q^2) = n_1 \ln(1 + Q^2/Q_0^2)$ and $n_L = 0$

Taking $x_0 = 0.01$ and $Q_0^2 = 0.4 \text{ GeV}^2$ results in the modified BEKW model (BEKW(mod)) with the 5 free parameters :

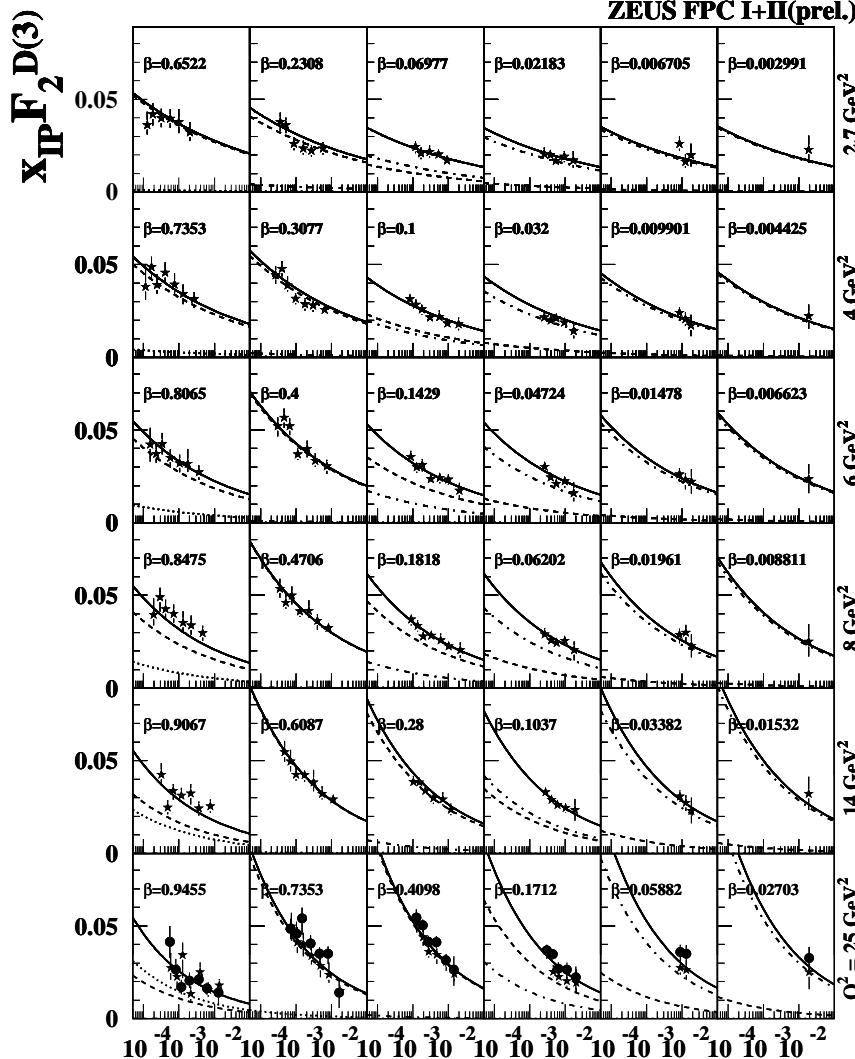
$c_T, c_L, c_g, n_1^{T,g}, \gamma$



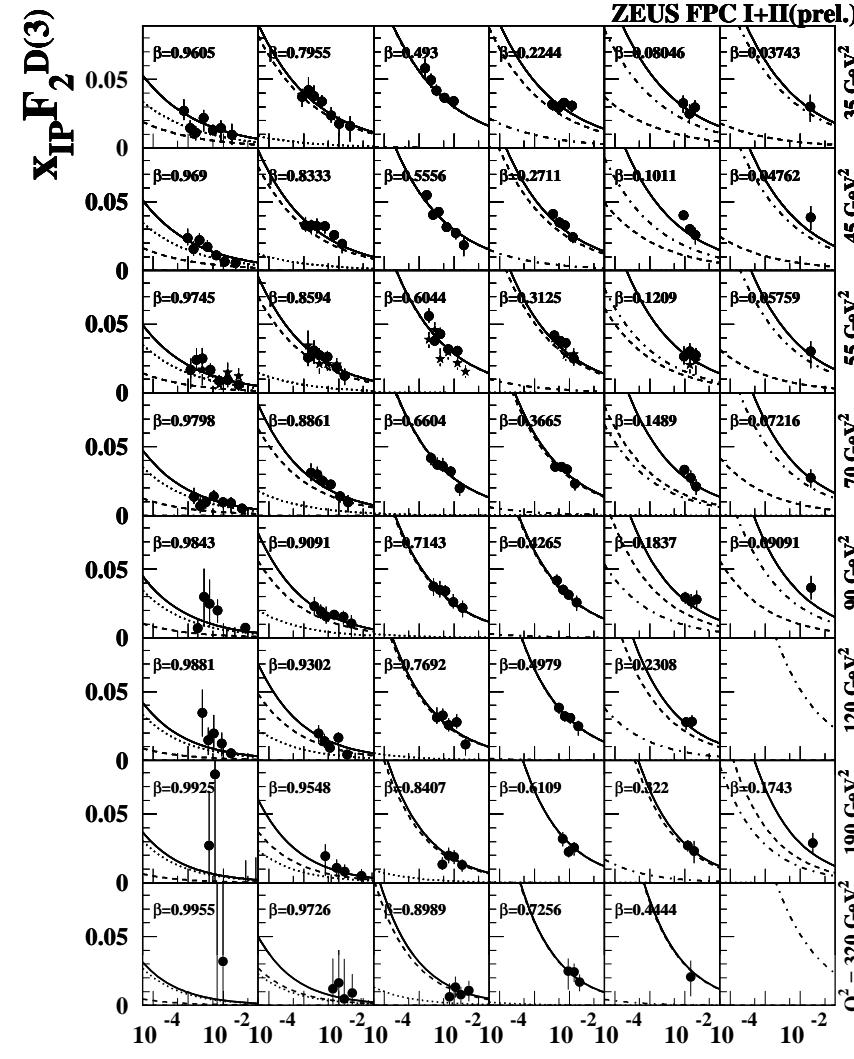
$x_{IP} F_2 D(3)$ Results from the Mx 98-99 and Mx 99-00 Analyses with BEKW(mod) Fit (I)

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BEKW-fit (prel.):

> 400 points, 5 parameters
 $\chi^2/n_D = 0.71$, total errors



sum of all contributions

— - - transverse qq contribution

..... longitudinal qq contribution

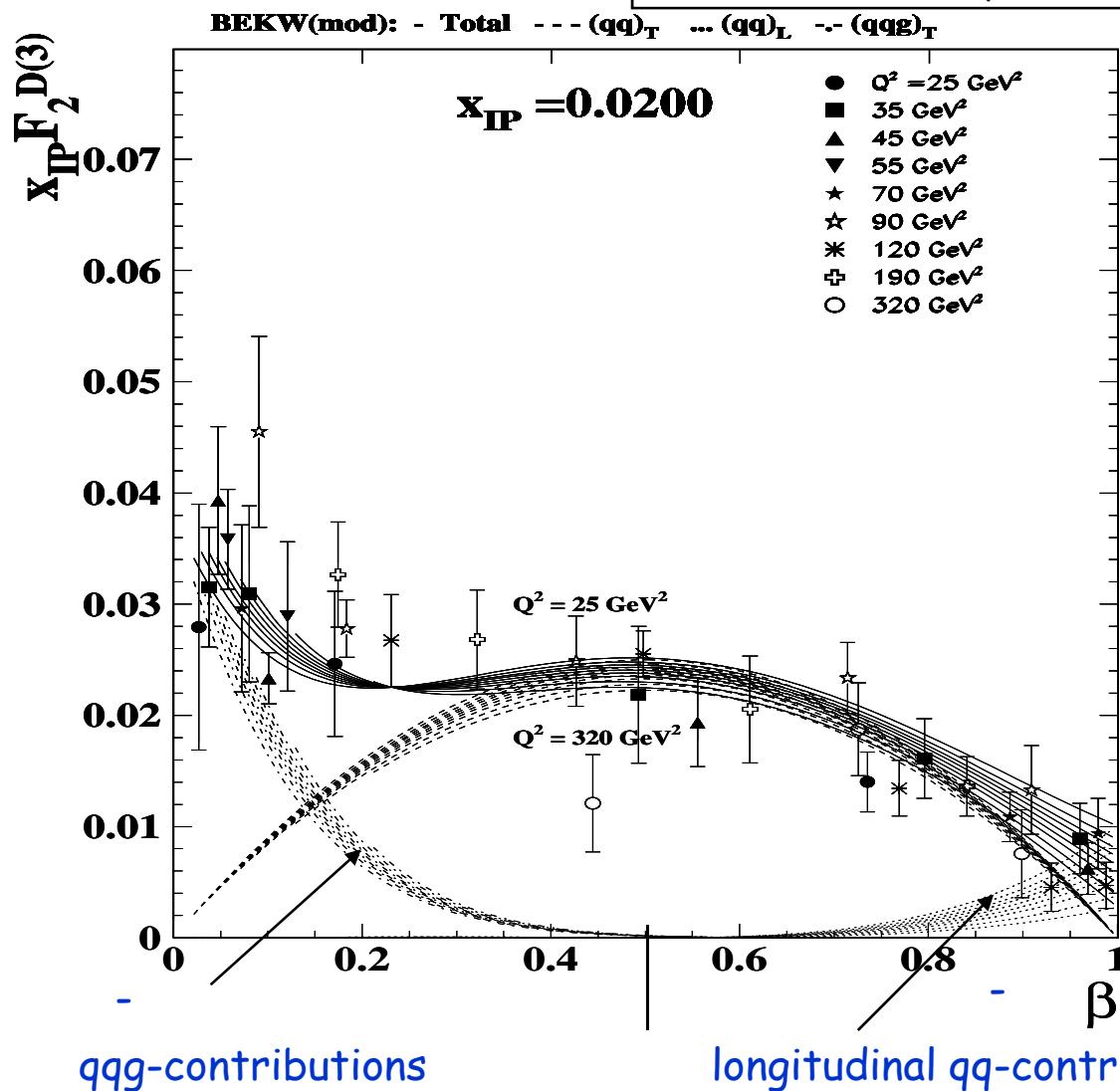
- · - transverse qgg contribution



$x_{IP}F_2D(3)$ Results from the Mx 99-00 Analysis with BEKW(mod) Fit (II)



ZEUS EPC-II (prel.)



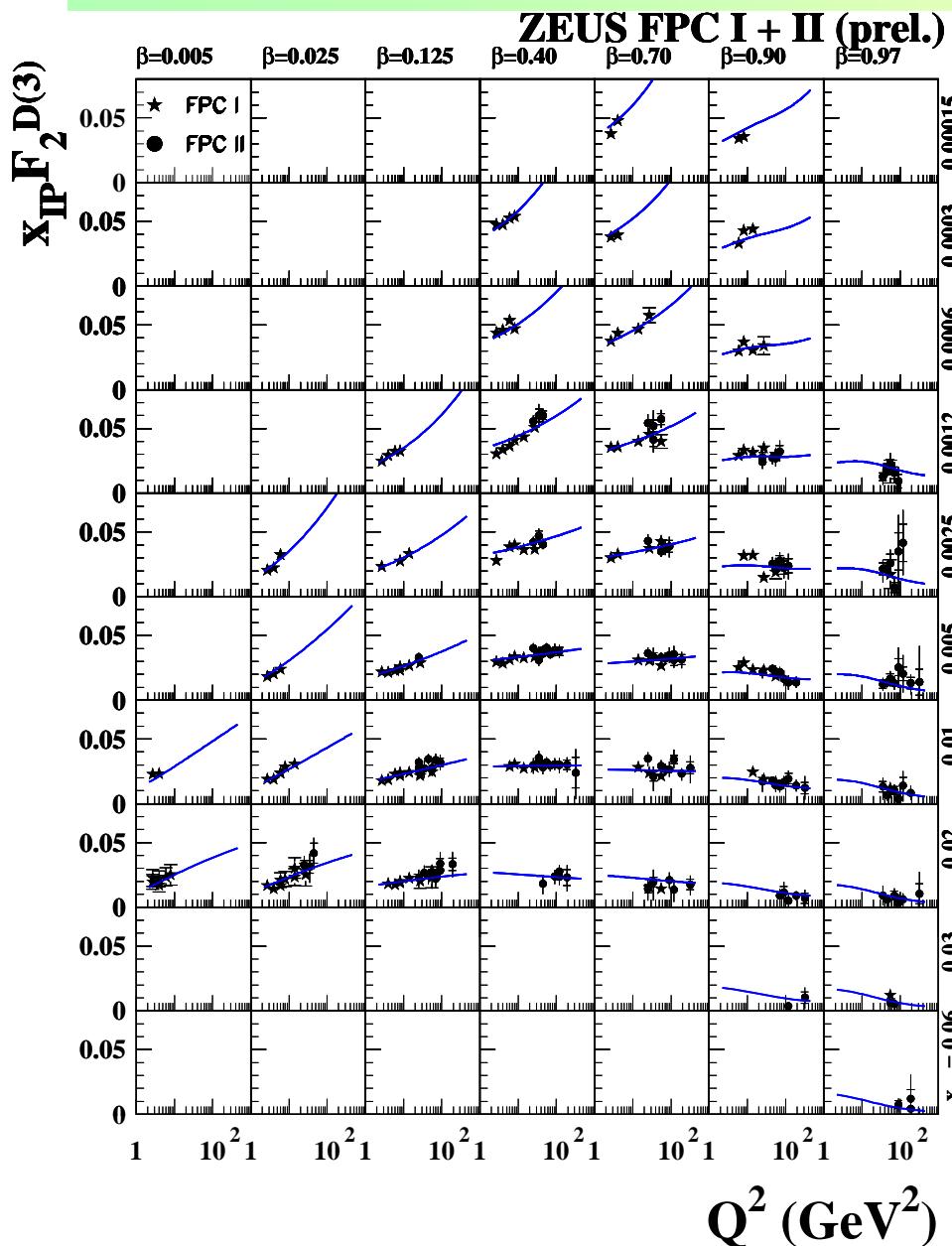
Fixed $x_{IP} = 0.02$

$25 < Q^2 < 320 \text{ GeV}^2$
in one plot

The 3 contributions
from BEKW(mod)
fit for the above
 Q^2 values plotted



$x_{IP}F_2D(3)$ Results from the Mx 98-99 and Mx 99-00 Analyses with BEKW(mod) Fit: III



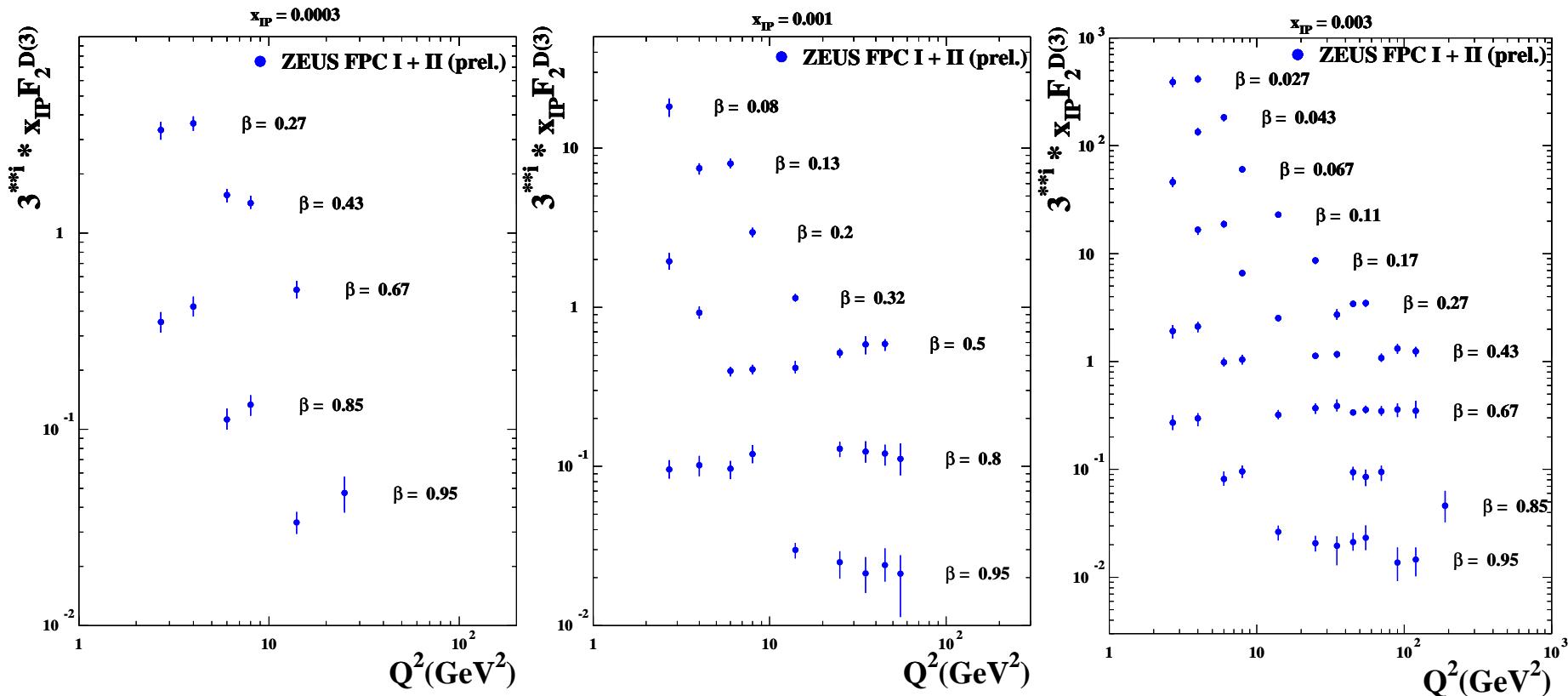
— Result of the BEKW(mod) fit

$x_{IP}F_2D(3)$ shows considerable
scaling violations:

from positive scaling violations over
near constancy to
negative scaling violations.

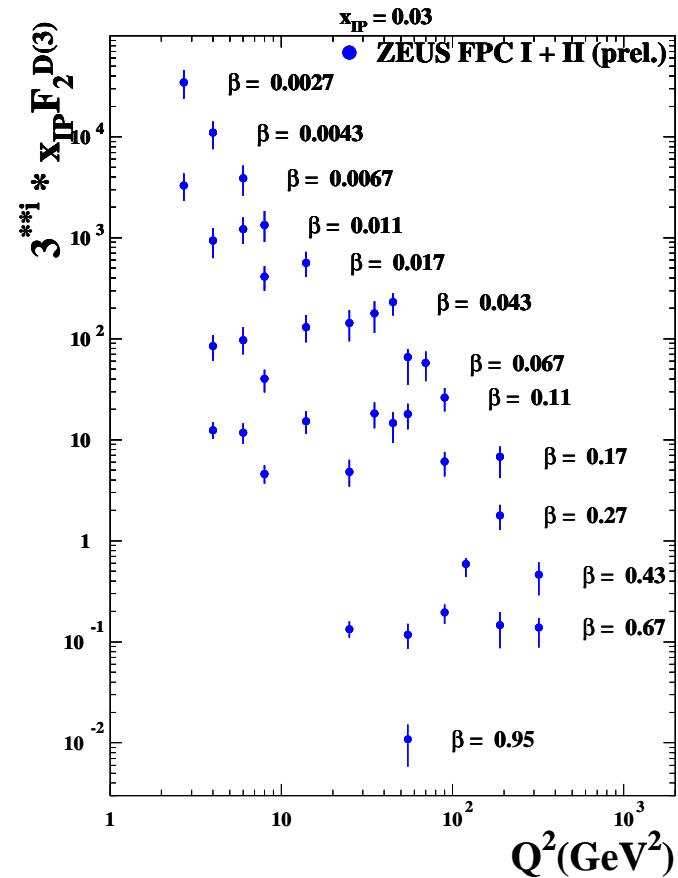
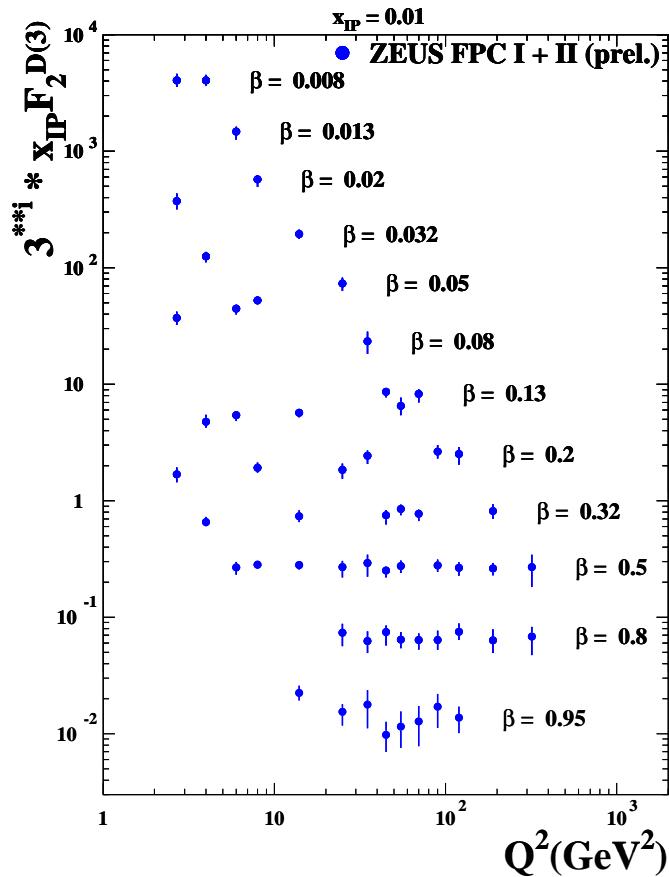


$x_{IP} F_2 D(3)$ Results from the Mx 98-99 and Mx 99-00 Analyses : Q²-dependence (I)





$x_{IP}F_2D(3)$ Results from the Mx 98-99 and Mx 99-00 Analyses : Q²-dependence (II)

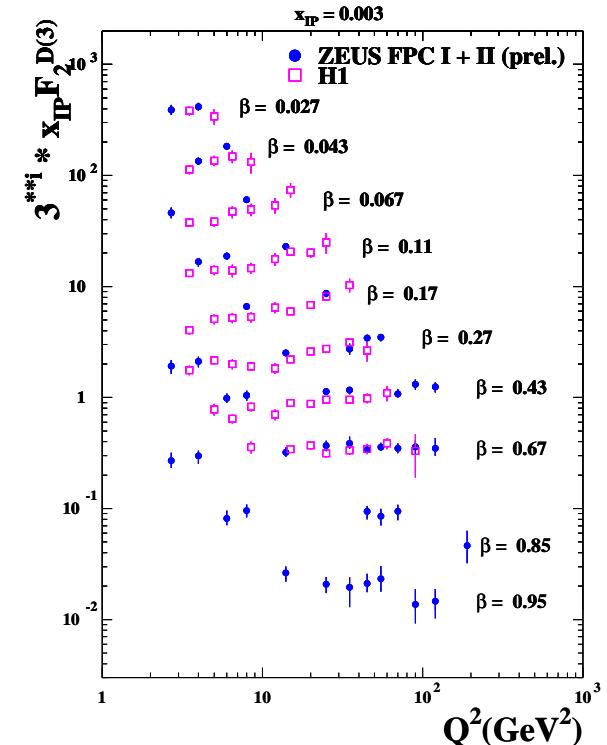
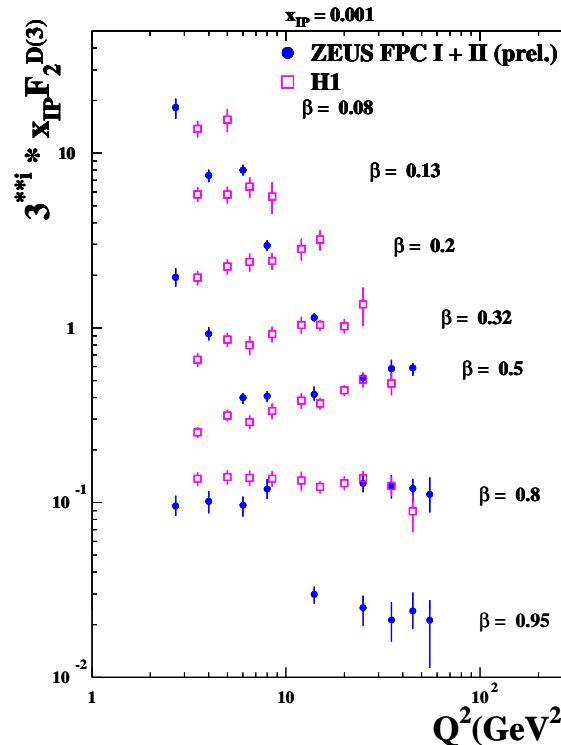
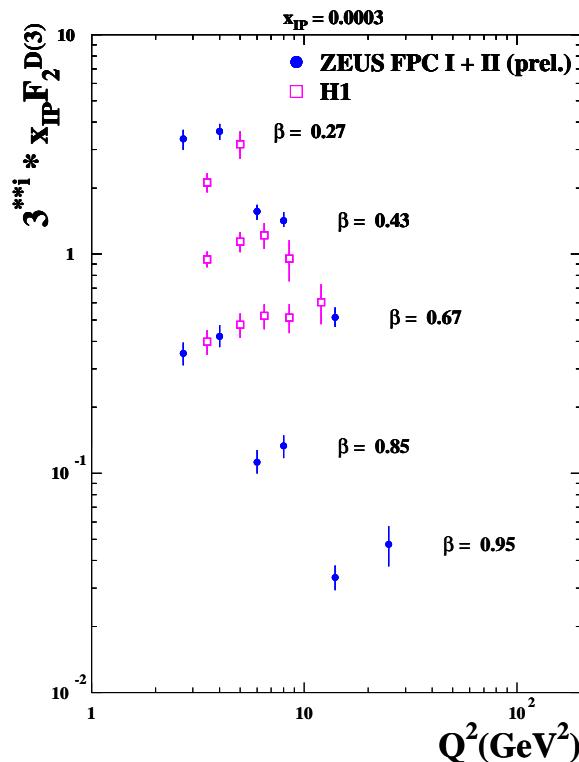


Q^2 dependence of $x_{IP}F_2D(3)$

The region $x_{IP}\beta = x < 6 \cdot 10^{-4}$ is dominated by positive scaling violations.
For $0.002 < x < 0.02$ constancy is observed



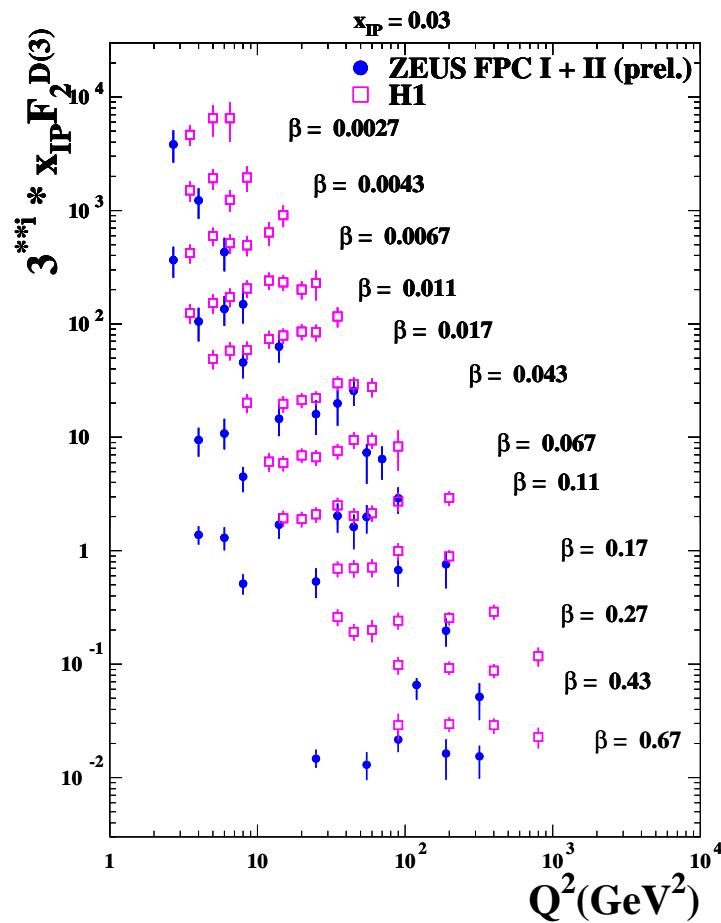
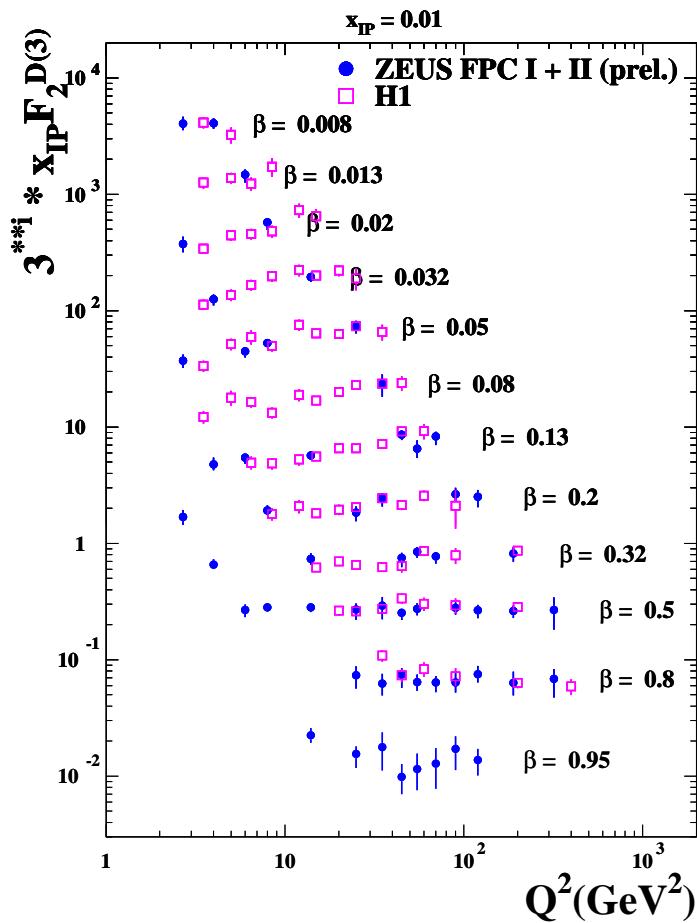
$x_{IP} F_2 D(3)$ Results from the Mx 98-99 and Mx 99-00 Analyses Comparison with H1 Results (I)



Note: ZEUS results contain contributions from p-dissociation with masses $M_{p\text{-diss}} < 2.3 \text{ GeV}$,
H1 results contain contributions with masses $M_{p\text{-diss}} < 1.6 \text{ GeV}$.
ZEUS results do not contain contributions from Reggeon-exchanges,
H1 results may contain such contributions for higher x_{IP} .



$x_{IP} F_2 D(3)$ Results from the Mx 98-99 and Mx 99-00 Analysis Comparison with H1 Results (II)



Comparison to H1 data

Fair agreement,
except maybe for a few (x_{IP}, Q^2) bins

Note: ZEUS points are shifted to
H1 bins using BFKL parametrization.
Only those ZEUS point are shown for
which the shift was <30%.



ZEUS Results from the LPS I

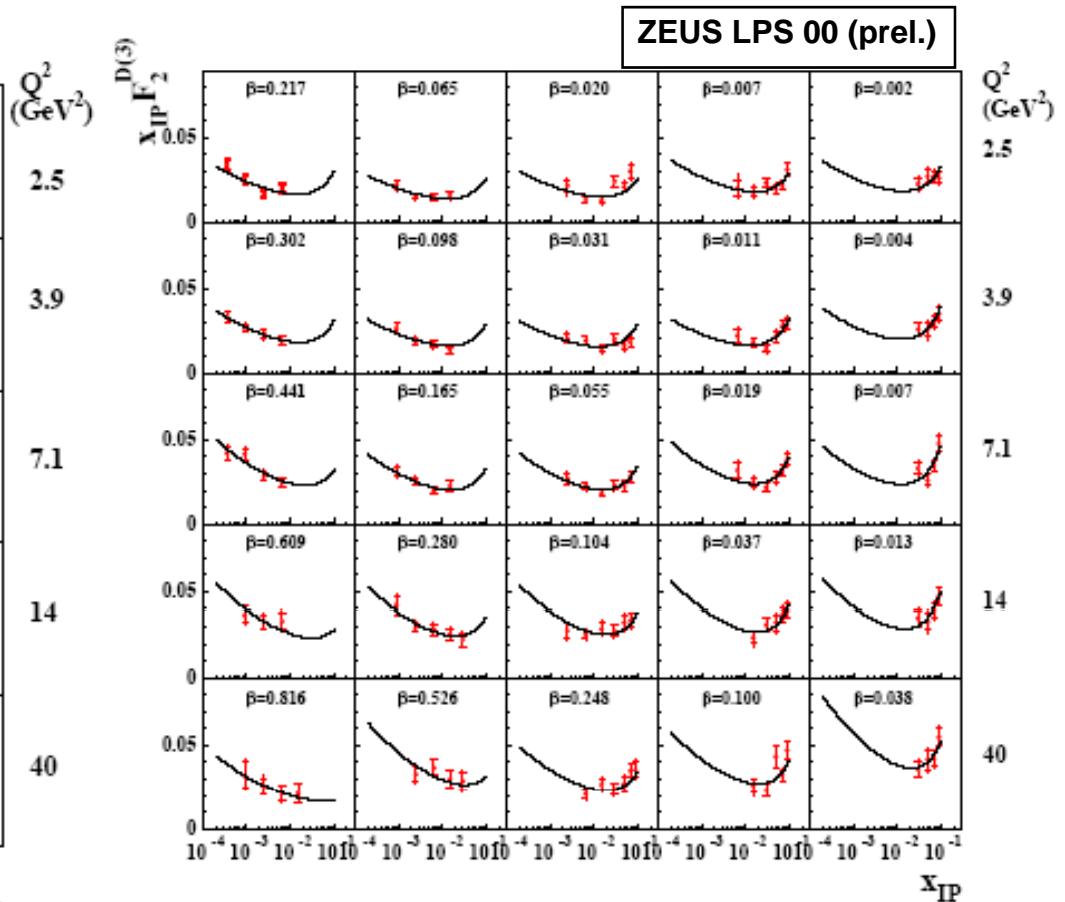
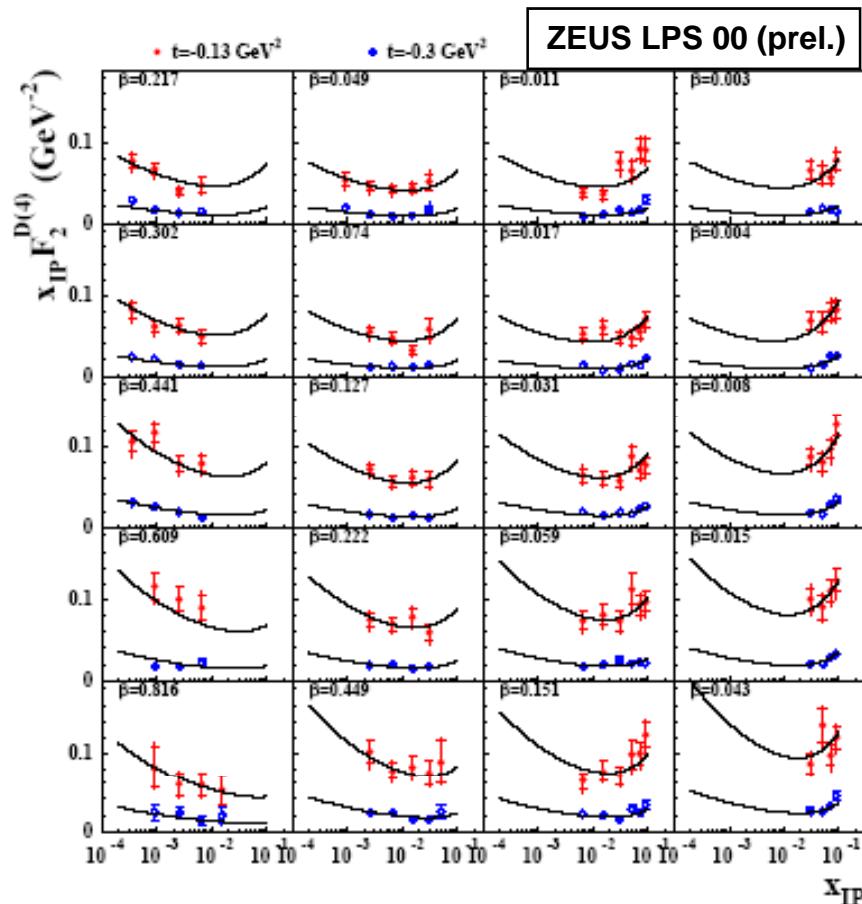
Diffractive structure functions from ZEUS LPS mesurements



Regge fit :

$$F_2^{D(4)} = f_{IP}(x_{IP}, t) \cdot F_2^{IP}(\beta, Q^2) + n_{IR} \cdot f_{IR}(x_{IP}, t) \cdot F_2^{IR}(\beta, Q^2)$$

- $t = 0.13 \text{ GeV}^2$
- $t = 0.3 \text{ GeV}^2$



Fit parameters: $\alpha_{IP}(0) = 1.10 \pm 0.02(\text{stat.})^{+0.01}_{-0.02}(\text{syst.}) + 0.02(\text{model})$ $\alpha'_{IP} = -0.03 \pm 0.07(\text{stat.})^{+0.04}_{-0.08} \text{ GeV}^{-2}$

$$B_{IP} = 7.2 \pm 0.7(\text{stat.})^{+1.4}_{-0.7}(\text{syst.}) \text{ GeV}^{-2}$$

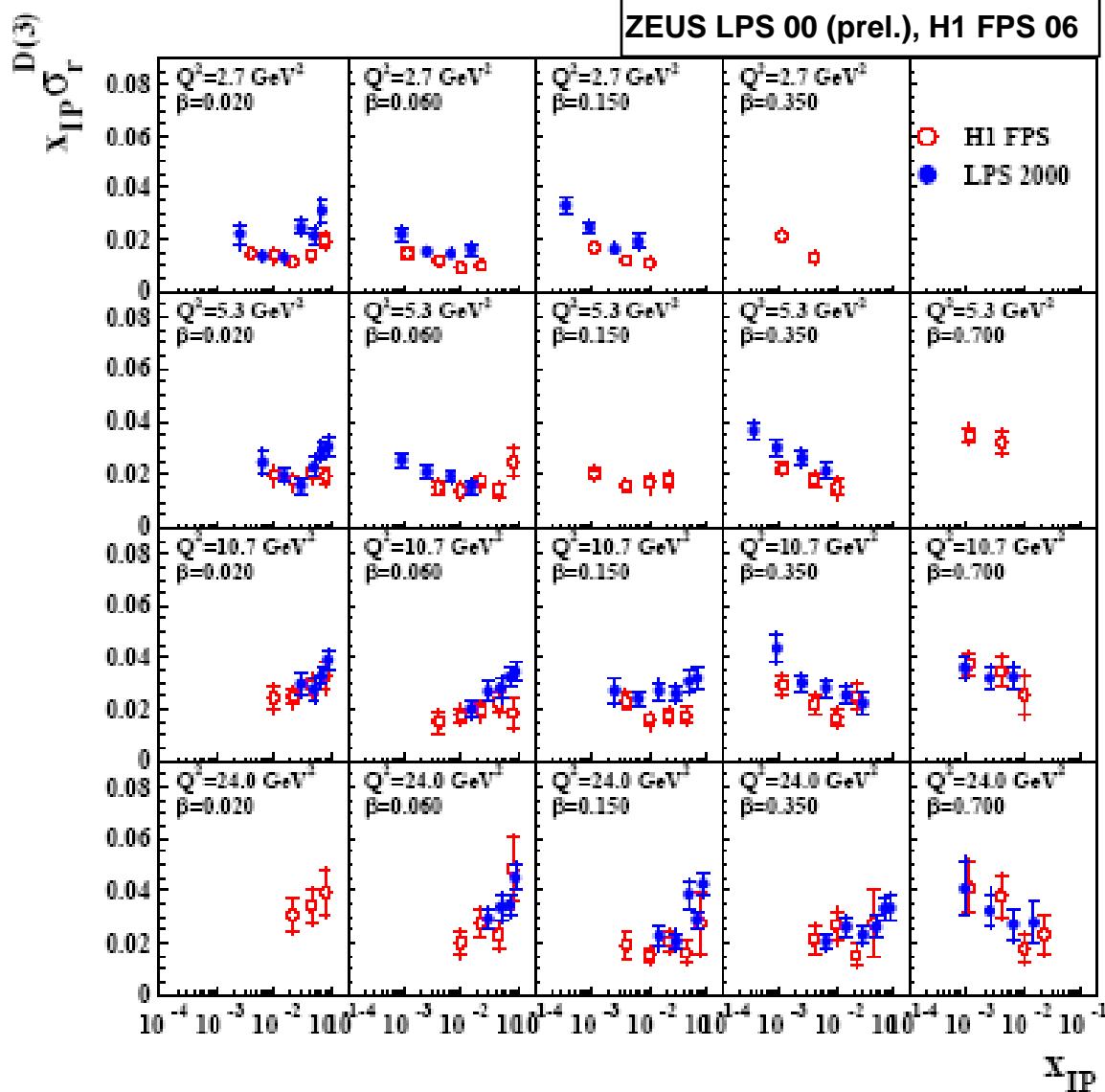
$$\alpha_{IR}(0) = 0.75 \pm 0.07(\text{stat.})^{+0.02}_{-0.04}(\text{syst.}) \pm 0.05(\text{model})$$



ZEUS Results from the LPS (II)



Comparison of LPS results with recent H1 FPS results



$$R^D = \sigma_L^{\gamma^* p \rightarrow pX} / \sigma_T^{\gamma^* p \rightarrow pX}$$

$$R^D = 0 \rightarrow x_{\text{IP}} F_2^{D(3)} = x_{\text{IP}} \sigma_r^{D(3)}$$

Not shown are the normalization
Uncertainties of +12/-10 % for
the ZEUS LPS data
and +/-10% for the H1 FPS data.

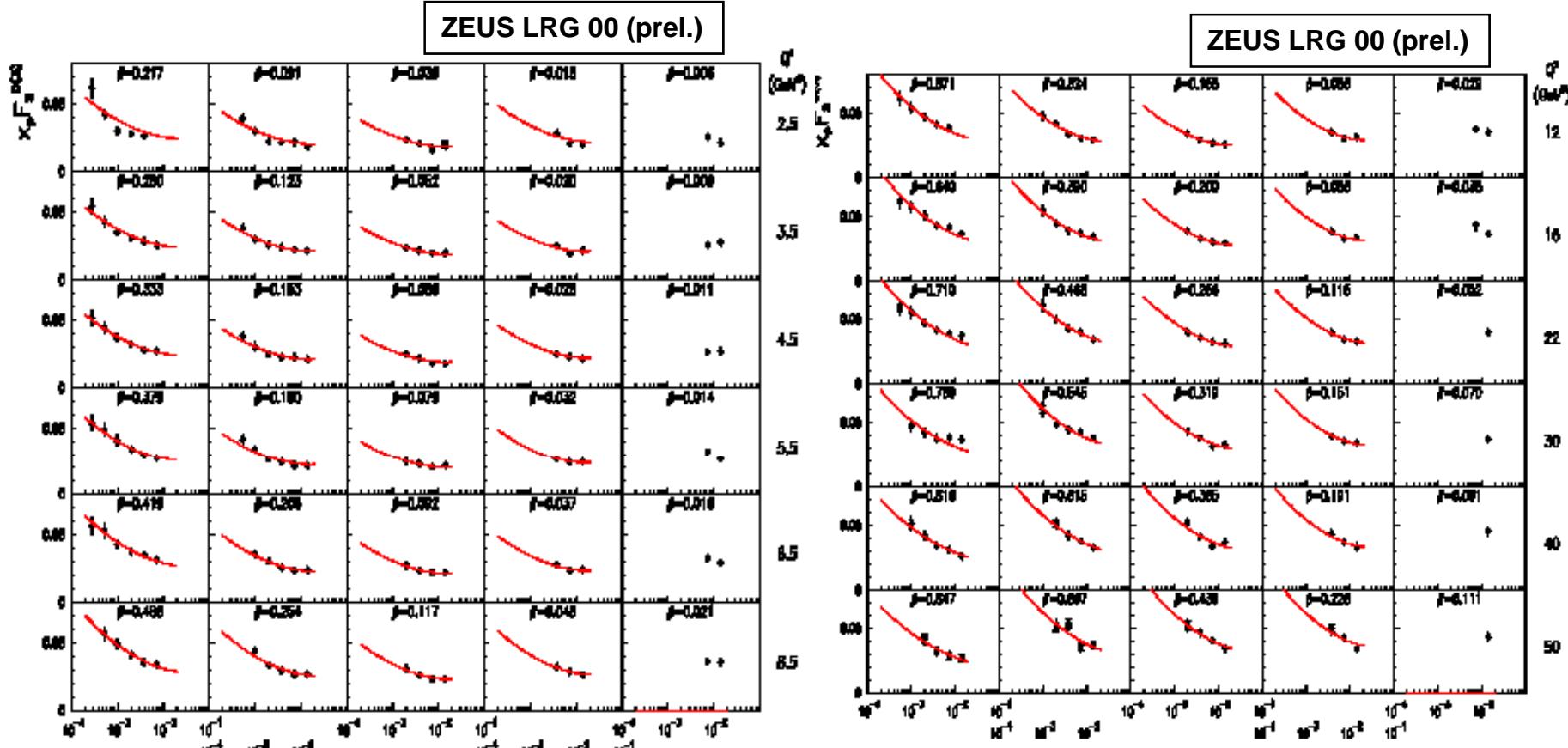
The agreement is good



ZEUS Results from the LRG Method (I)



Events selected by : $\eta_{\text{max}} < 3.0$ and energy in the Forward Plug Calorimeter (FPC) $< 1 \text{ GeV}$



$$F_2^{D(3)} = f_{IP}(x_{IP}) \cdot F_2^{IP}(\beta, Q^2) + n_{IR} \cdot f_{IR}(x_{IP}) \cdot F_2^{IR}(\beta, Q_2) \quad \text{with} \quad f_{IP,IR}(x_{IP}) = \int \frac{e^{B_{IP,IR}t}}{x_{IP,IR}^{2\alpha_{IP,IR}(t)-1}} dt$$

Results of a **Regge-fit**

Fit results :

$$\alpha_{IP}(0) = 1.117 \pm 0.005 \pm 0.007$$

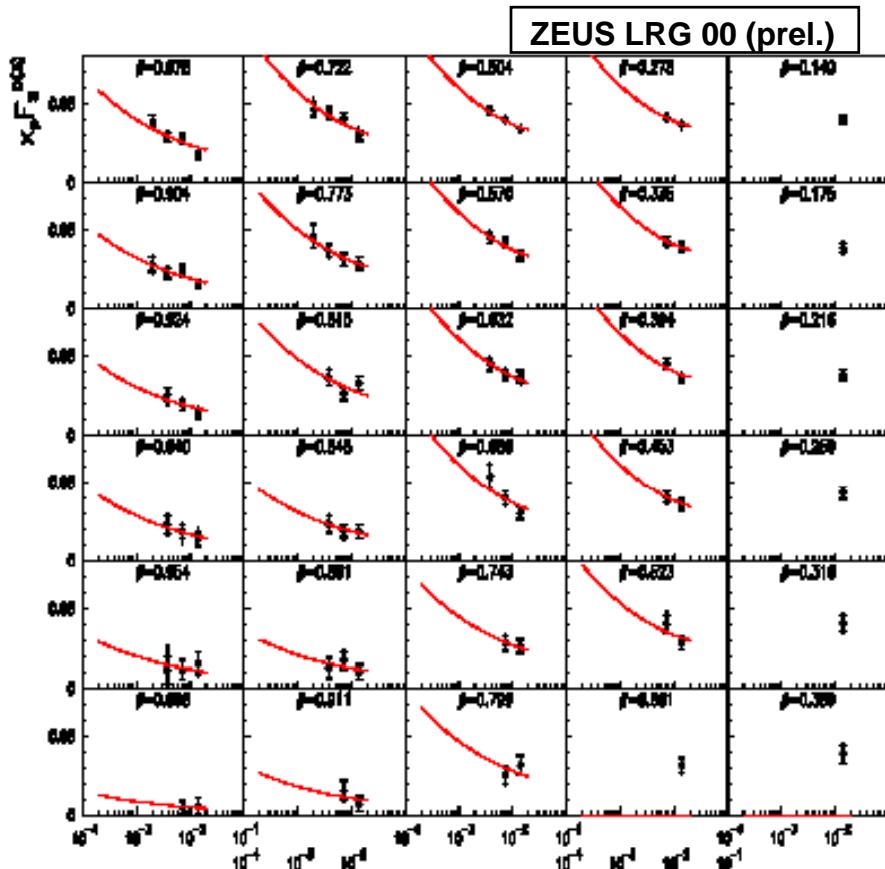
Input parameters to the Regge -fit:

$$\alpha_{IR}(0) = 0.75, \quad B_{IR} = 2.0 \text{ GeV}^{-2}$$

$$\alpha'_{IP} = 0.0 \text{ GeV}^{-2}, \quad B_{IP} = 7.2 \text{ GeV}^{-2}$$

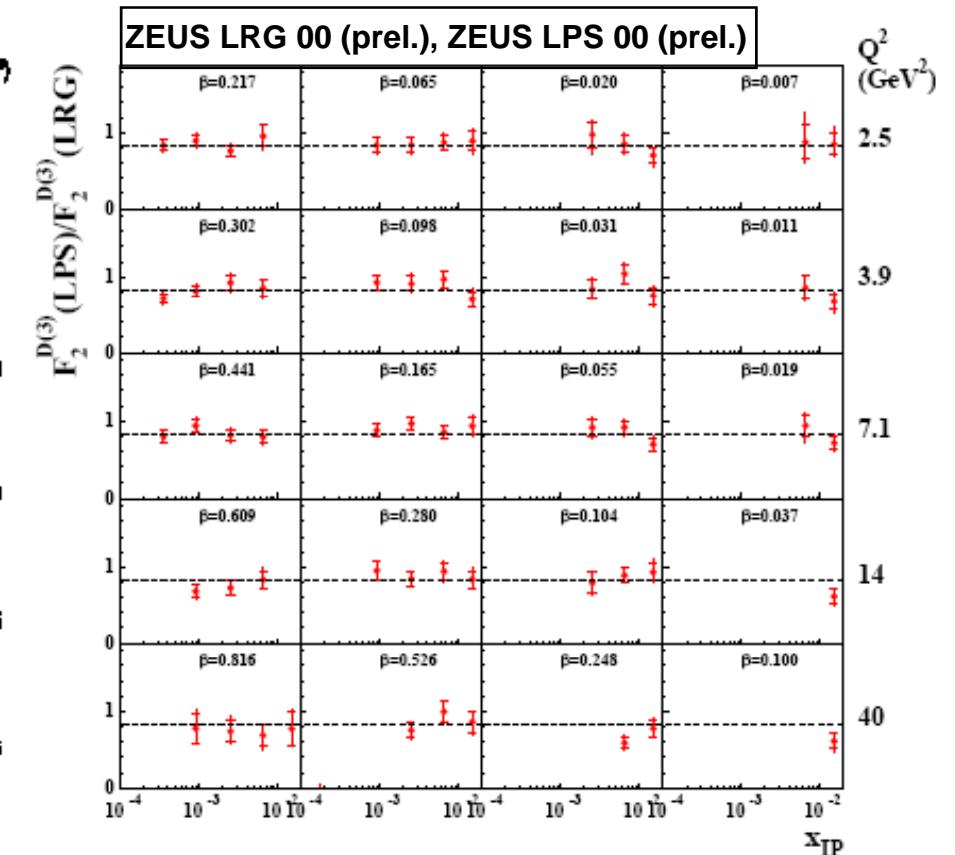


ZEUS Results from the LRG Method (II)



The Regge-fit gives a good description of the ZEUS LRG data with $\chi^2/\text{ndf} = 159/185$

Comparison of ZEUS LRG data with LPS data

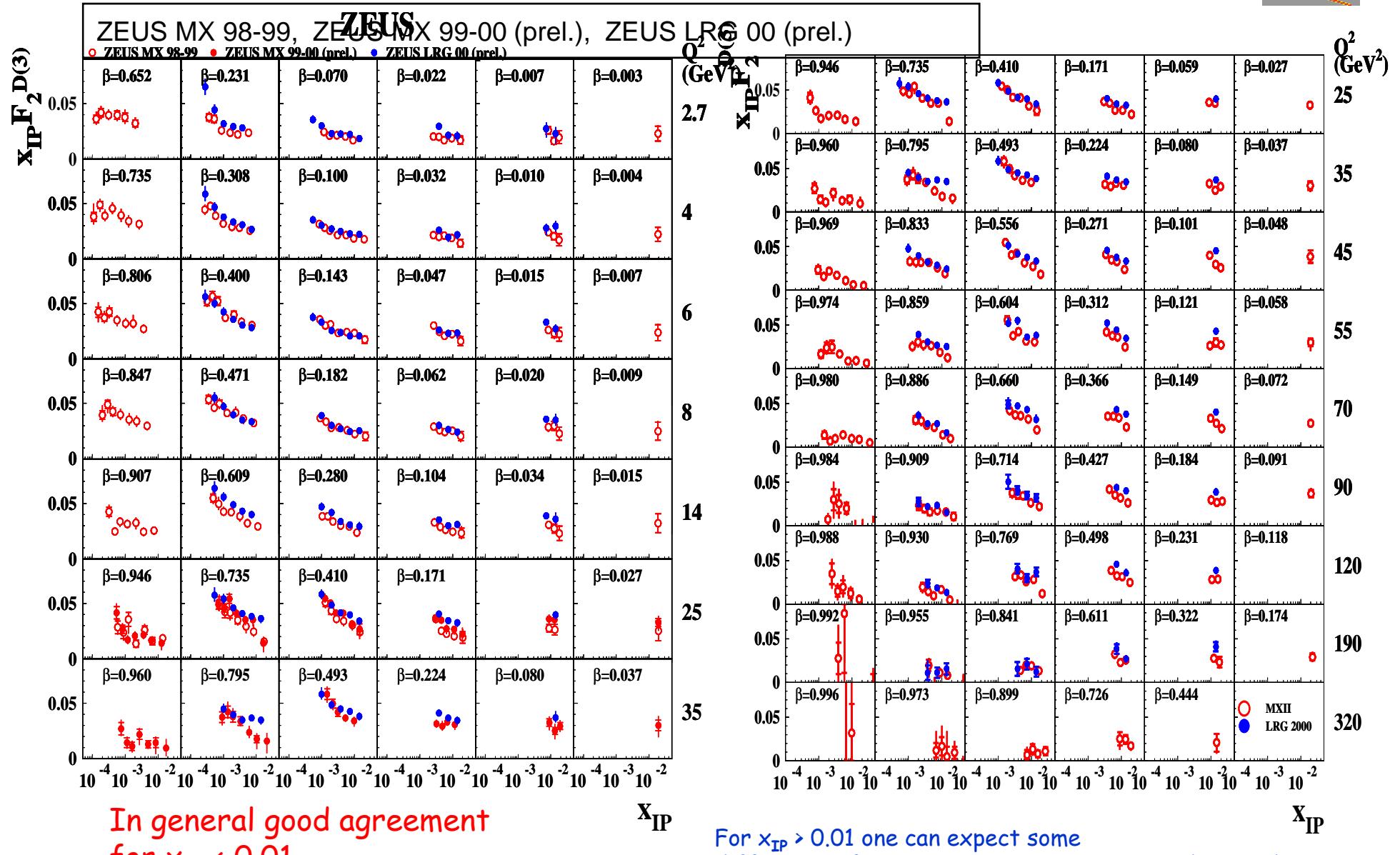


• The ratio $\text{LPS/LRG} = 0.82 +/- 0.01(\text{stat.}) +/- 0.03(\text{syst.})$
it is independent of Q^2 and β

Not shown is the normalization uncertainty
of the LPS measurement of about 10%.

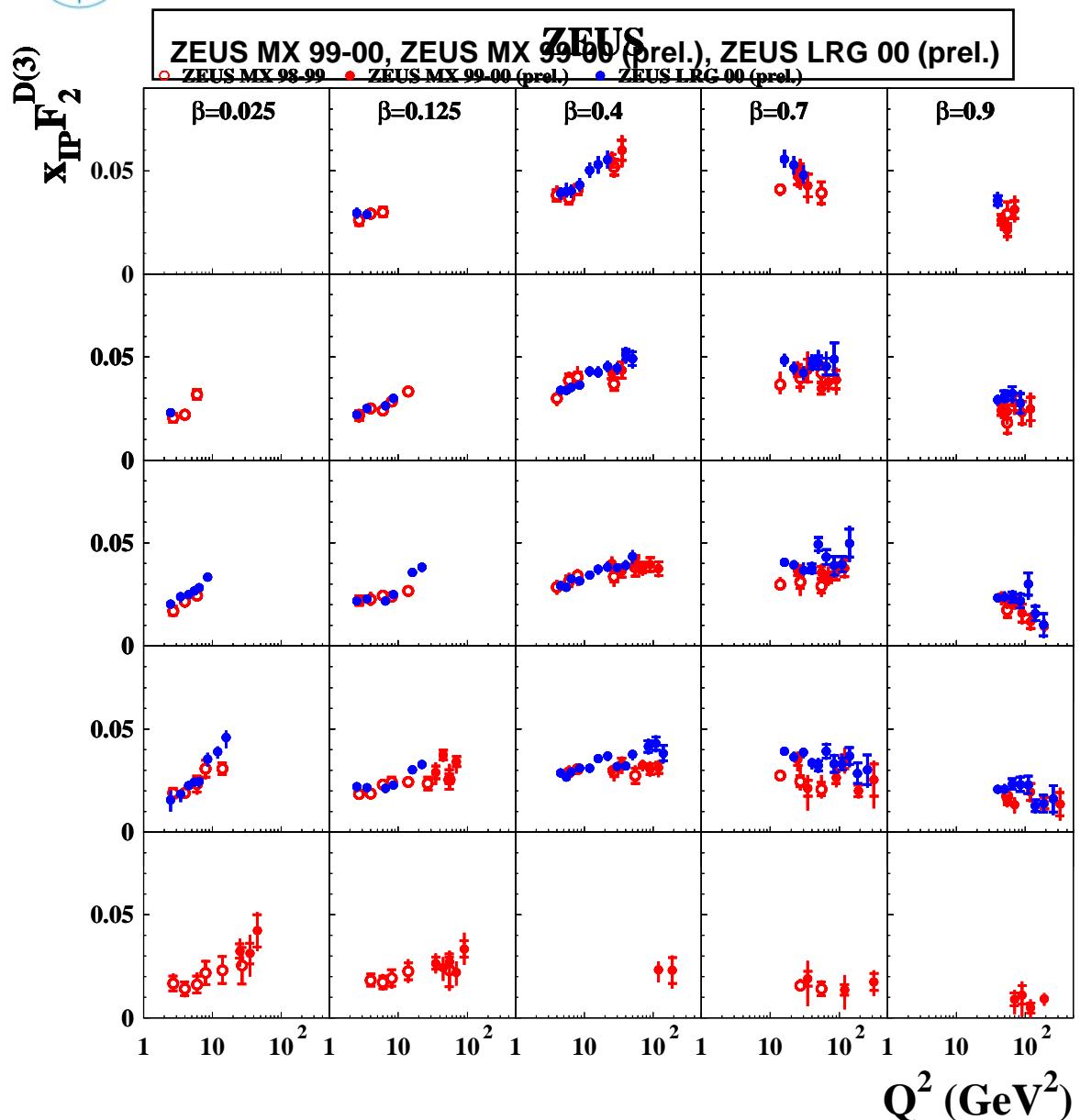


ZEUS : Comparison of Results from the M_x -, and LRG- Method (I)





ZEUS : Comparison of Results from the M_x -, and LRG- Method (II)



Comparison of
ZEUS LRG with ZEUS M_x results:

$x_{IP} F_2^{D(3)}$
as a function of Q^2

- ZEUS Mx 90-99
- ZEUS Mx 99-00 (prel.)
- ZEUS LRG 00 (prel.)

Reasonable agreement,
maybe there is a normalization
difference.

Work is continuing to
understand remaining
differences



Summary



- ZEUS presented **preliminary results** on inclusive diffraction from 3 different methods for the extraction of inclusive diffractive events.
- Results from all 3 methods are derived from **data taken during the same time**.
- The results span a wide range of the kinematic region **up to high Q^2** .
- There is **good to reasonable agreement** for the results from all 3 methods.
- There is **good to reasonable agreement** for the Q^2 -dependence of the structure function between the M_x -method, the LRG-method and the H1 data.
- There is also good **agreement compared to results from H1** for the FPS method.
- **Work continues** to understand some remaining minor differences, in particular with respect to the relative normalisations.
- We try to get a consistent picture out of the results from these three methods.