

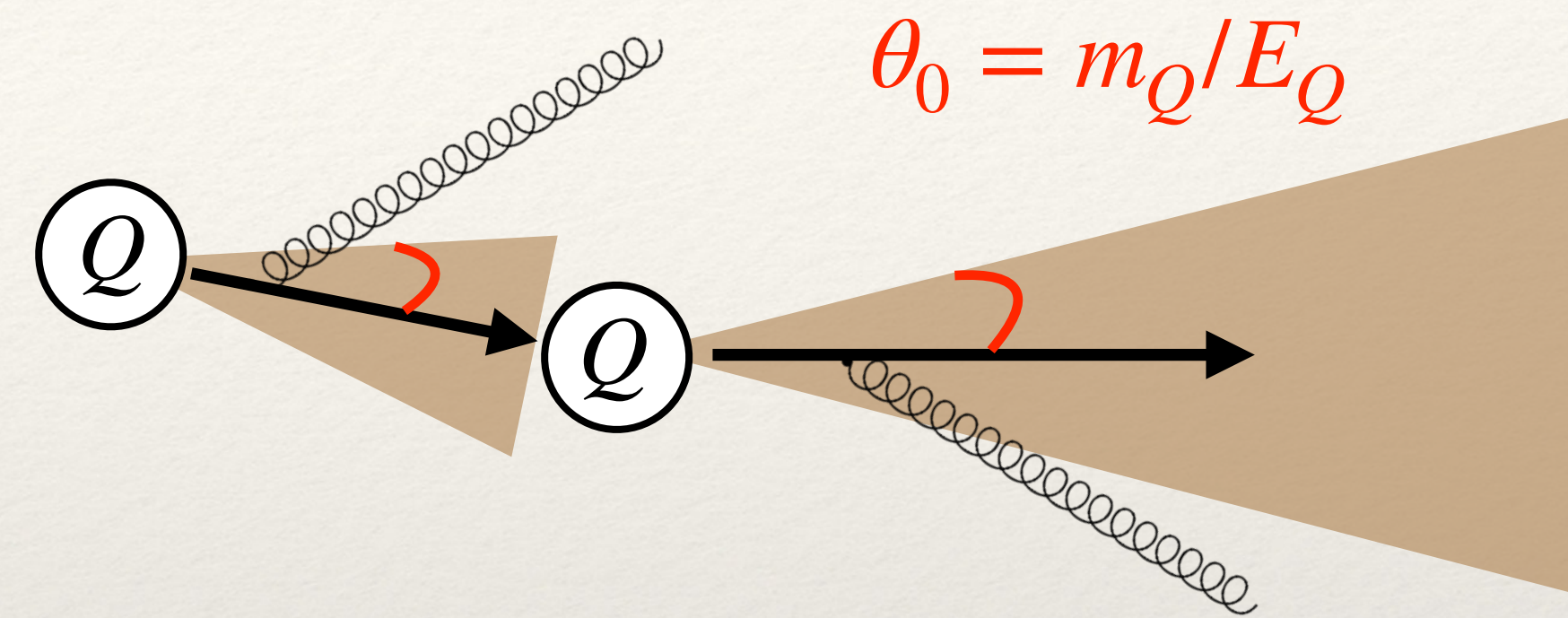
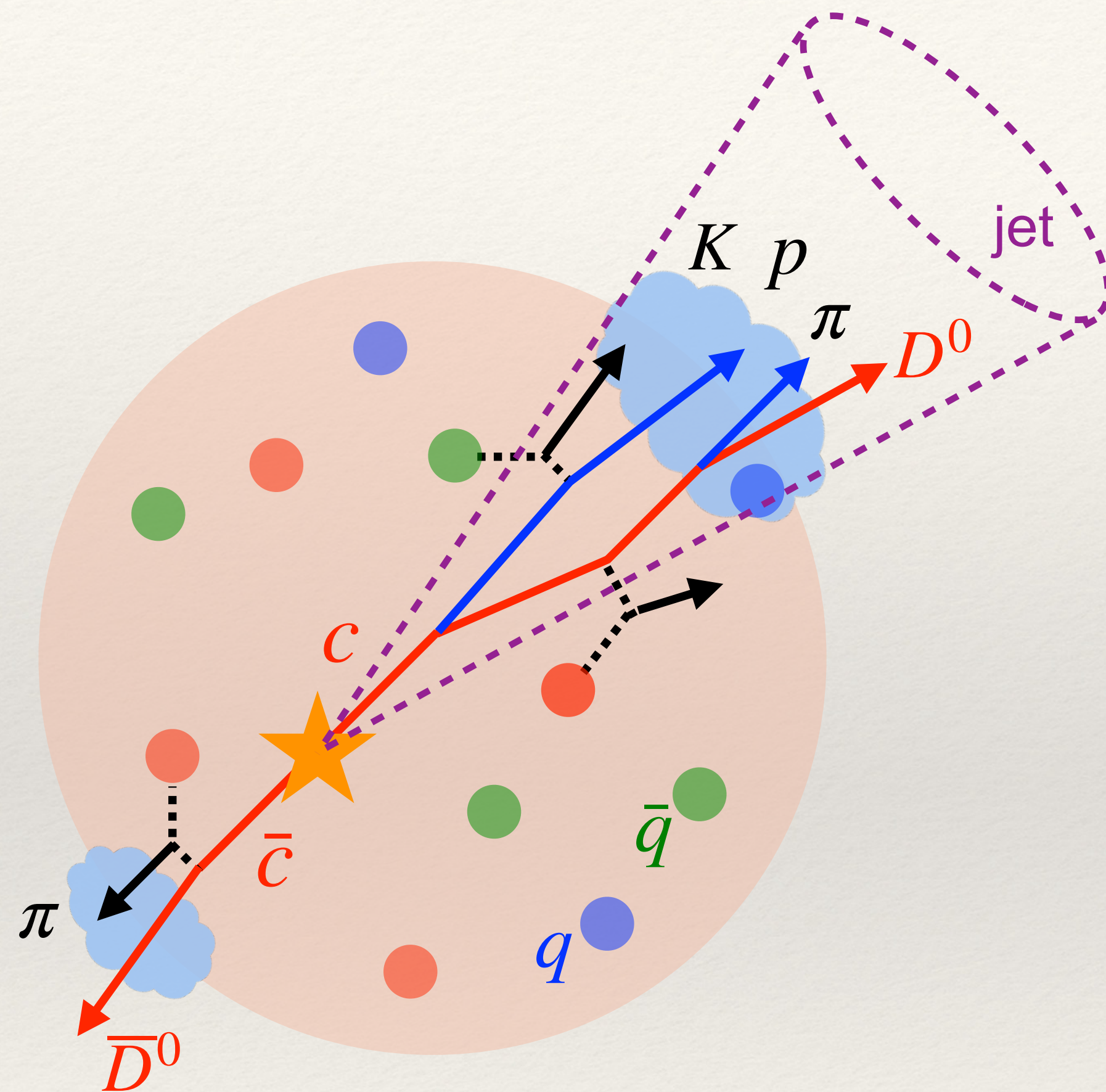
Heavy quark transport: from heavy flavor hadrons to heavy flavor jets

Shanshan Cao
Shandong University

September 29, 2024
SoftJet 2024, Tokyo



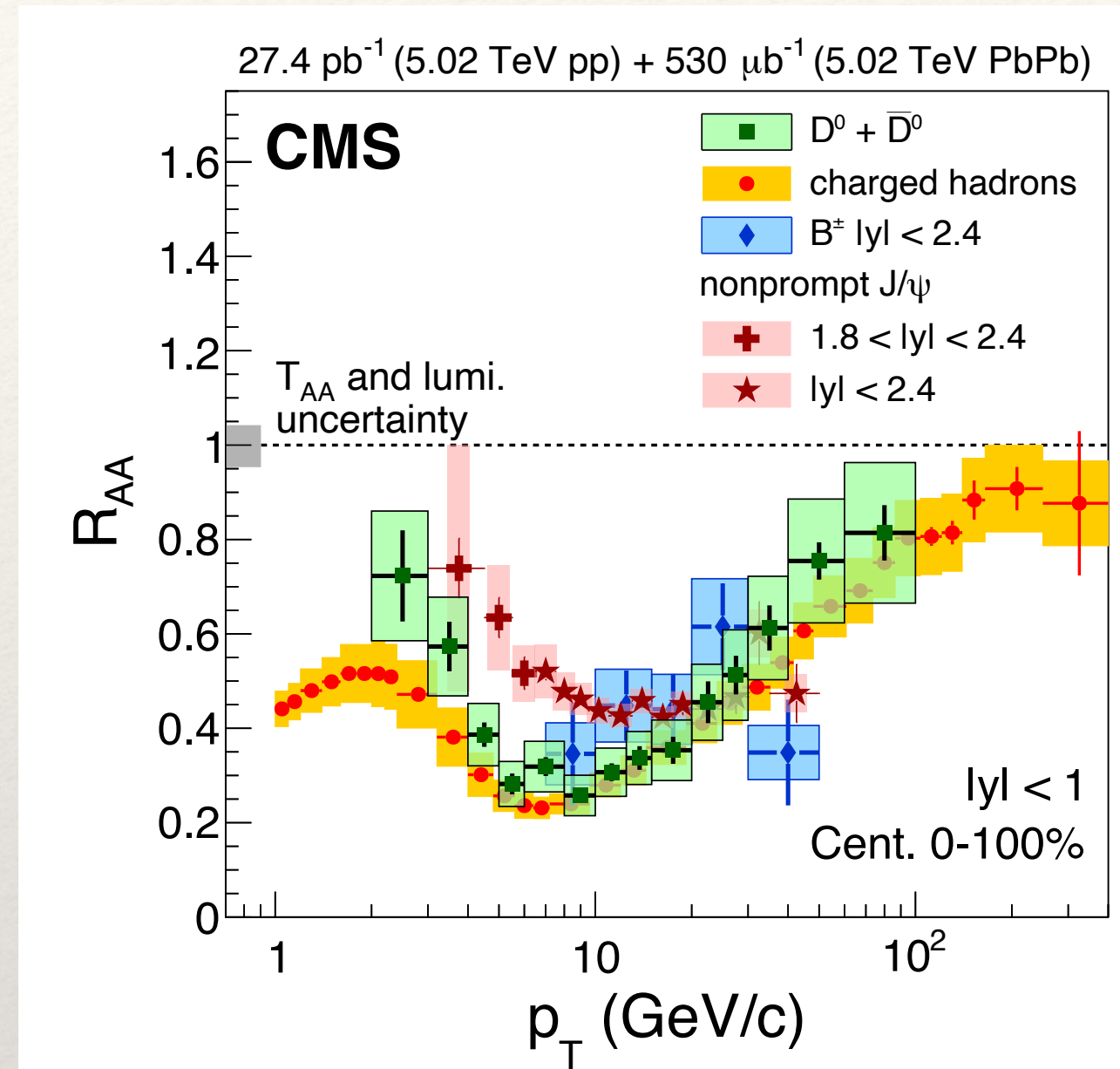
Heavy quarks in high-energy nuclear collisions



- Produced from initial hard scatterings
- Serve as an ideal probe of the QGP properties
- Provide a unique opportunity for studying the flavor dependence of parton splitting (dead cone effect)

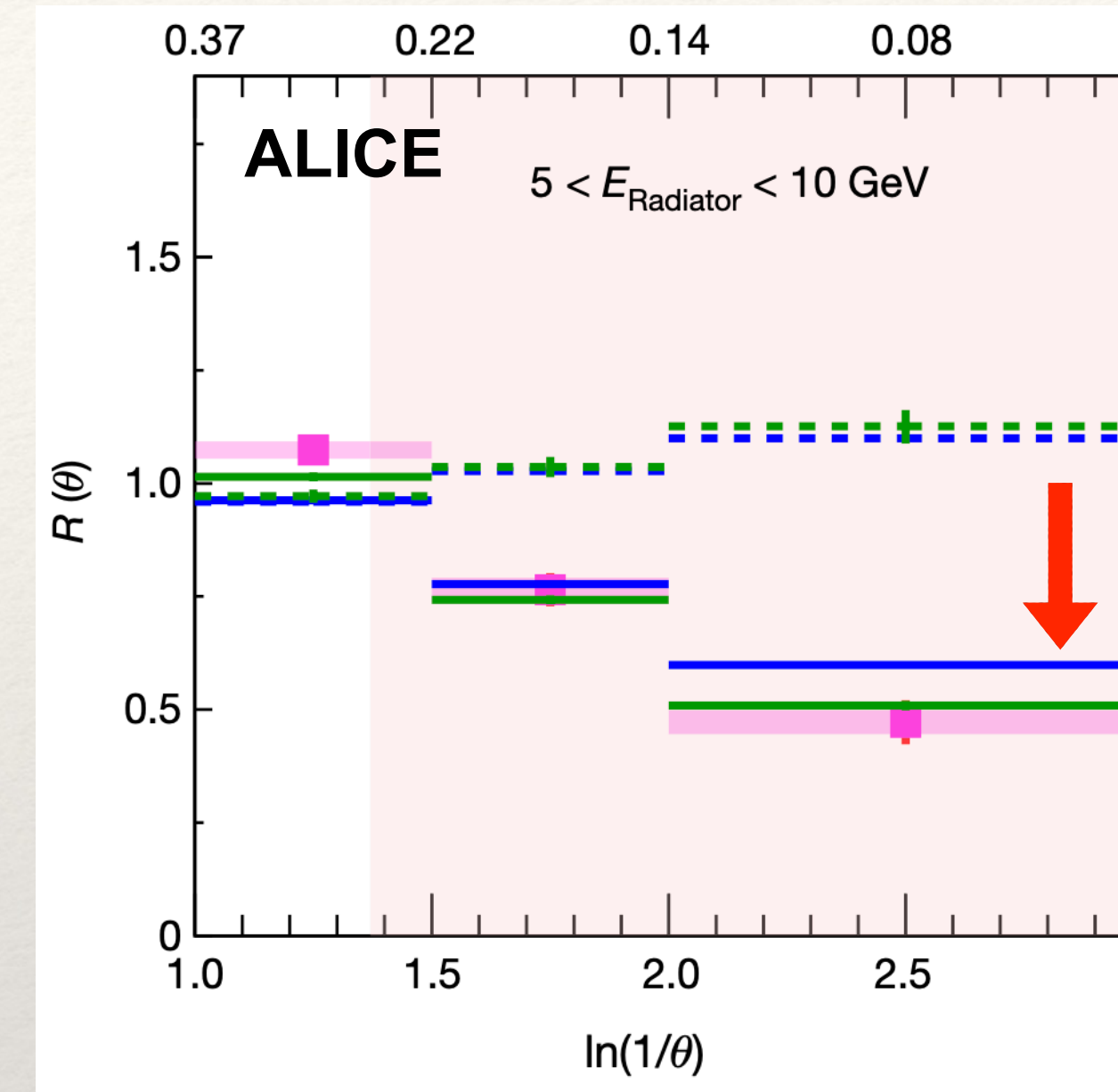
Searches for the flavor dependence of parton splitting

Hadron R_{AA} (parton energy loss)



Phys. Lett. B 782
(2018) 474-496

Distribution of splitting angles in pp



Nature 605
(2022) 7910

No clear separation between charged hadrons, D , and B , except at very low p_T

Clear suppression of splitting at small θ in D -jets vs. inclusive jets

Goals:

- Understand flavor hierarchies embedded in both hadrons and jets
- Use heavy flavor observables to probe the QGP properties

Transport model for parton-QGP interactions

Linear Boltzmann Transport (LBT)

$$p_a \cdot \partial f_a(x_a, p_a) = E_a (\mathcal{C}_a^{\text{el}} + \mathcal{C}_a^{\text{inel}})$$

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$$p_a \cdot \partial f_a(x_a, p_a) = E_a (\mathcal{C}_a^{\text{el}} + \mathcal{C}_a^{\text{inel}})$$

Elastic energy loss ($ab \rightarrow cd$)

$$\mathcal{C}_a^{\text{el}} = \sum_{b,c,d} \int \prod_{i=b,c,d} \frac{d[p_i]}{2E_a} (\gamma_d f_c f_d - \gamma_b f_a f_b) \cdot (2\pi)^4 \delta^4(p_a + p_b - p_c - p_d) \left| \mathcal{M}_{ab \rightarrow cd} \right|^2$$

2 → 2 scattering matrices

Transport model for parton-QGP interactions

Linear Boltzmann Transport (LBT)

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Elastic energy loss ($ab \rightarrow cd$)

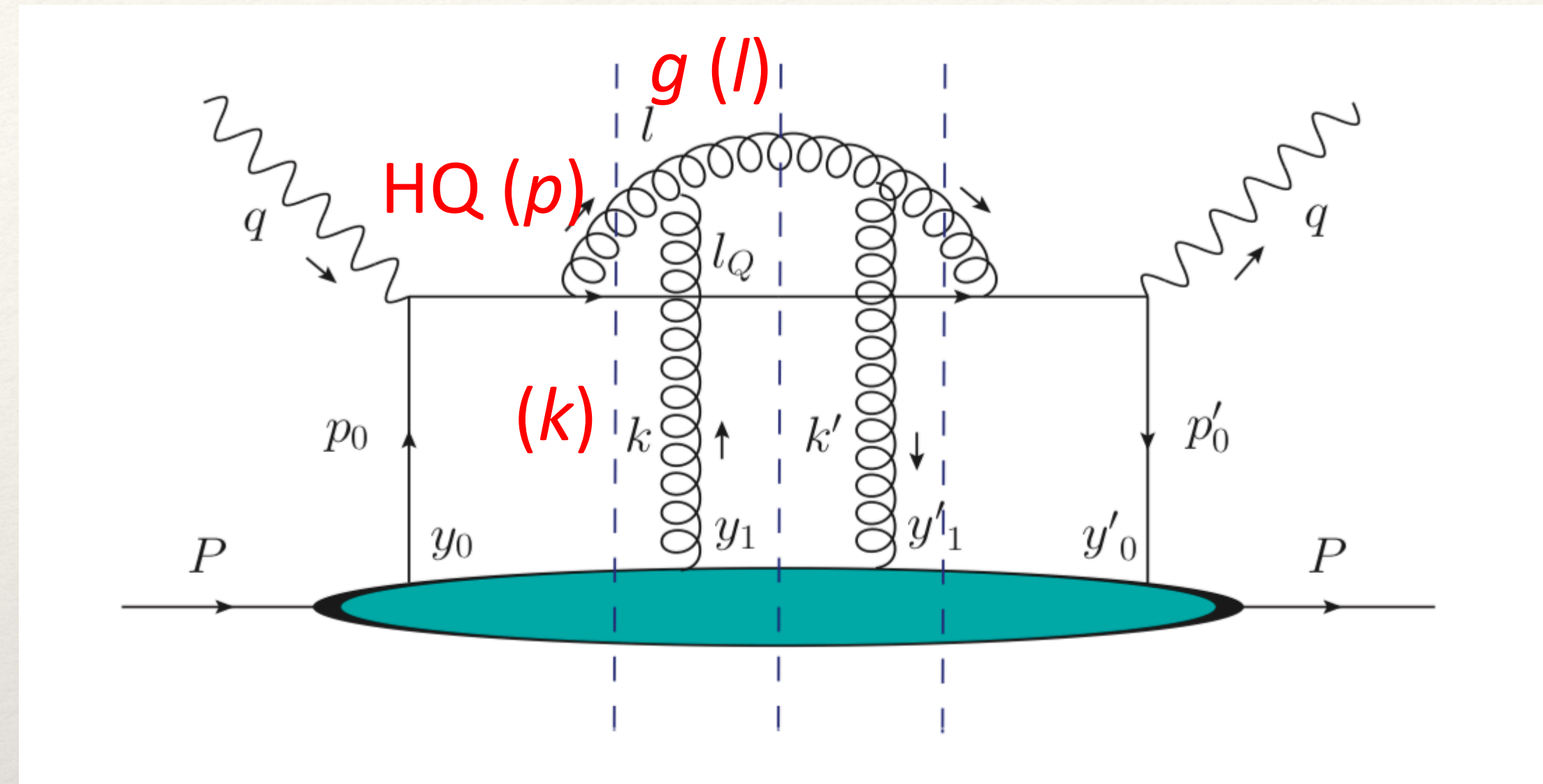
$$\mathcal{C}_a^{\text{el}} = \sum_{b,c,d} \int \prod_{i=b,c,d} \frac{d[p_i]}{2E_a} (\gamma_d f_c f_d - \gamma_b f_a f_b) \cdot (2\pi)^4 \delta^4(p_a + p_b - p_c - p_d) \left| \mathcal{M}_{ab \rightarrow cd} \right|^2$$

2 → 2 scattering matrices

loss term: **scattering rate**
(for Monte-Carlo simulation)

$$\Gamma_a^{\text{el}}(\mathbf{p}_a, T) = \sum_{b,c,d} \frac{\gamma_b}{2E_a} \int \prod_{i=b,c,d} d[p_i] f_b \cdot (2\pi)^4 \delta^{(4)}(p_a + p_b - p_c - p_d) \left| \mathcal{M}_{ab \rightarrow cd} \right|^2$$

Inelastic energy loss



Majumder PRD 85 (2012); Zhang, Wang and Wang, PRL 93 (2004)

- **Higher-twist formalism:** collinear expansion ($\langle k_{\perp}^2 \rangle \ll l_{\perp}^2 \ll Q^2$)

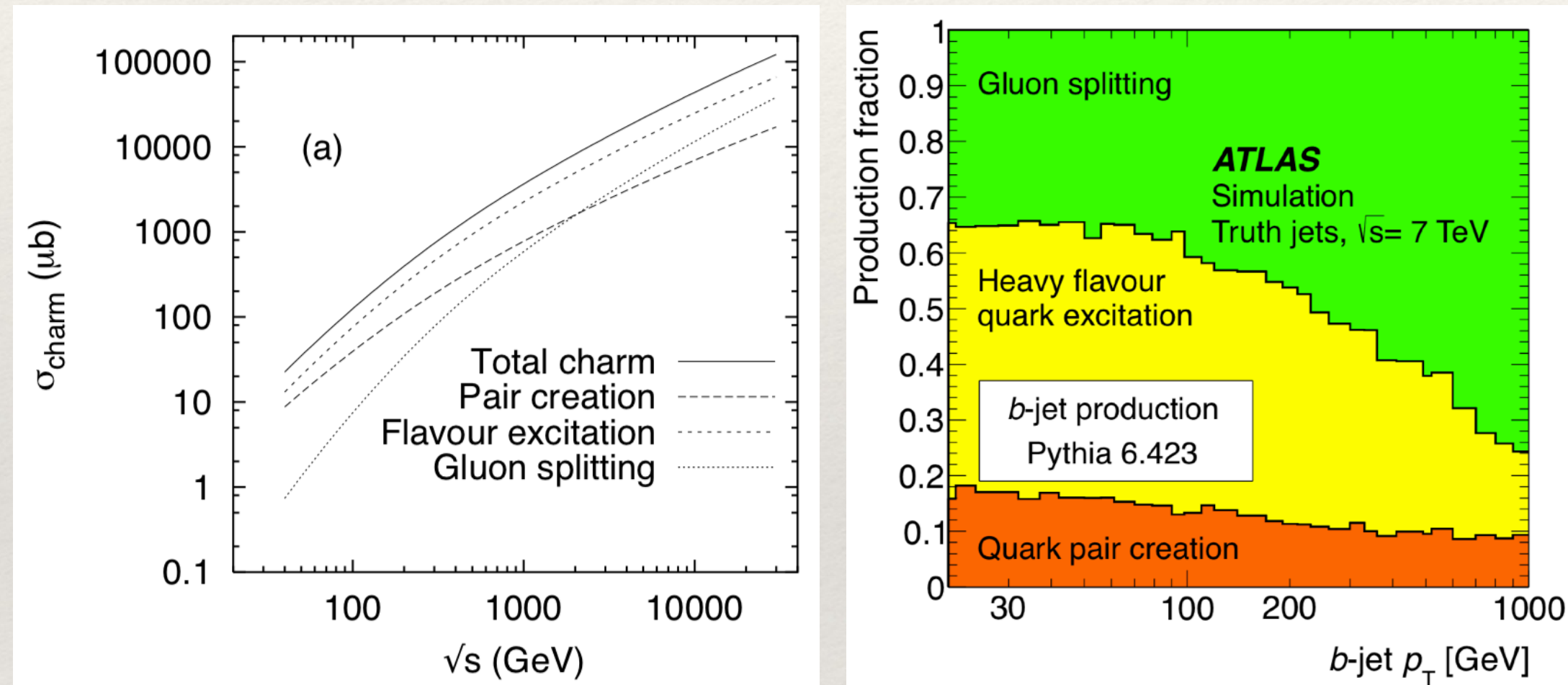
$$\frac{d\Gamma_a^{\text{inel}}}{dz dl_{\perp}^2} = \frac{dN_g}{dz dl_{\perp}^2 dt} = \frac{6\alpha_s P(z) l_{\perp}^4 \hat{q}}{\pi(l_{\perp}^2 + z^2 M^2)^4} \sin^2 \left(\frac{t - t_i}{2\tau_f} \right)$$

- Medium information absorbed in $\hat{q} \equiv d\langle p_{\perp}^2 \rangle / dt$

Flavor hierarchy in hadron suppression

- Hadron production in pp collisions: NLO production + fragmentation

NLO contribution to HQ production in Pythia simulation (gluon splitting)

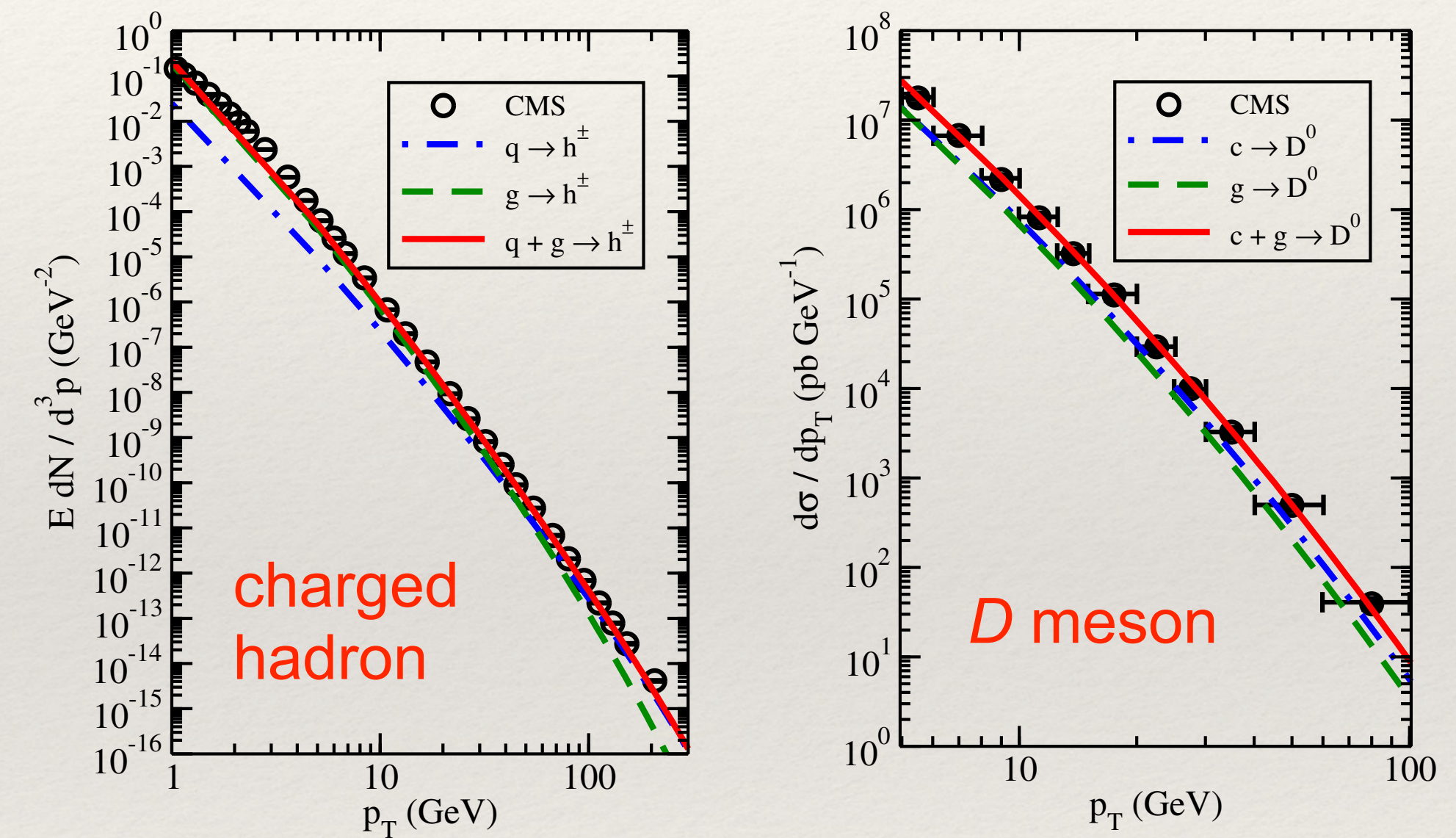


[Norrbin and Sjostrand, EPJC 17 (2000)]

[ATLAS, EPJC 73 (2013)]

- NLO contribution increases with \sqrt{s}
- NLO contribution increases with b -jet p_T

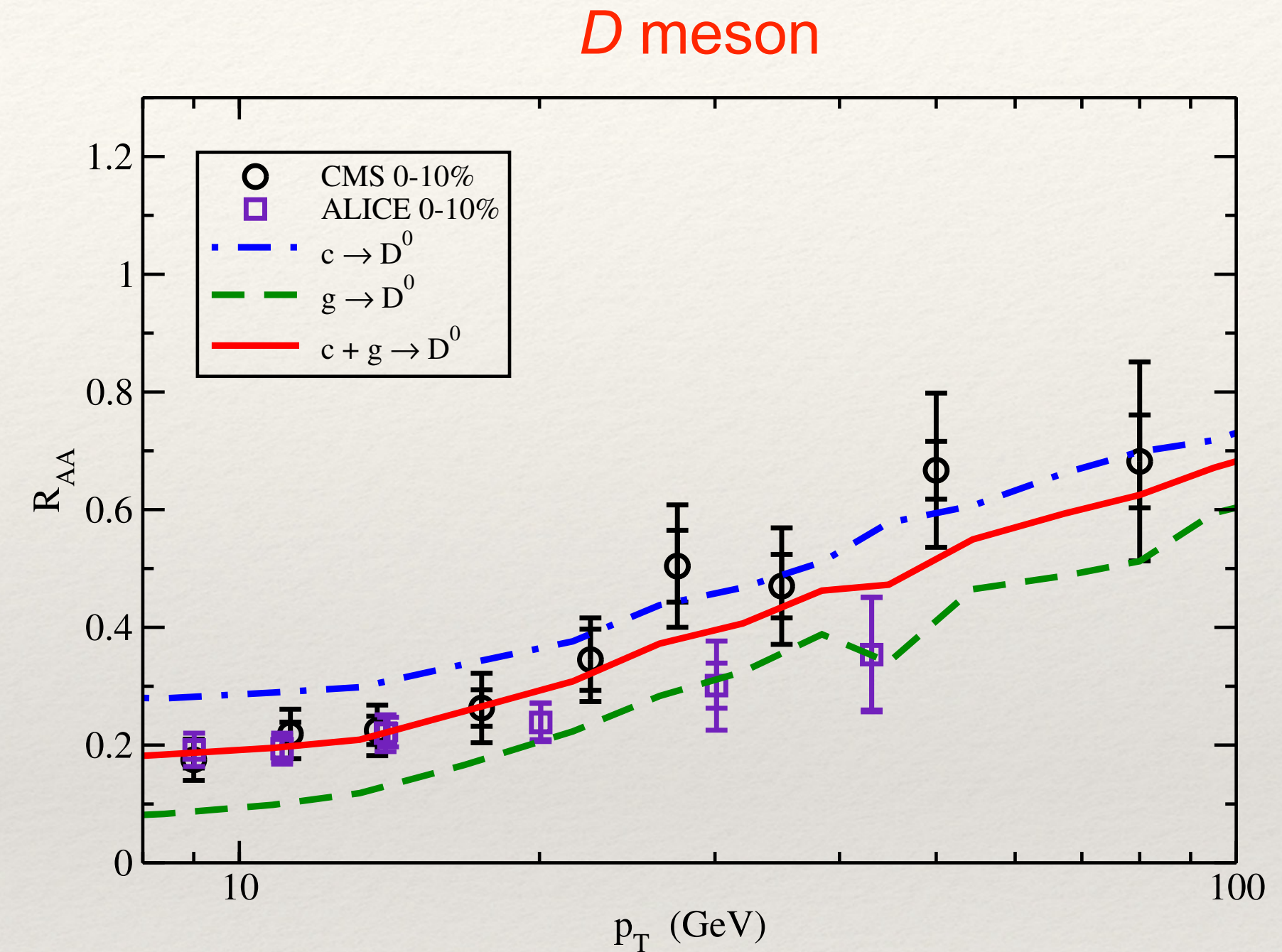
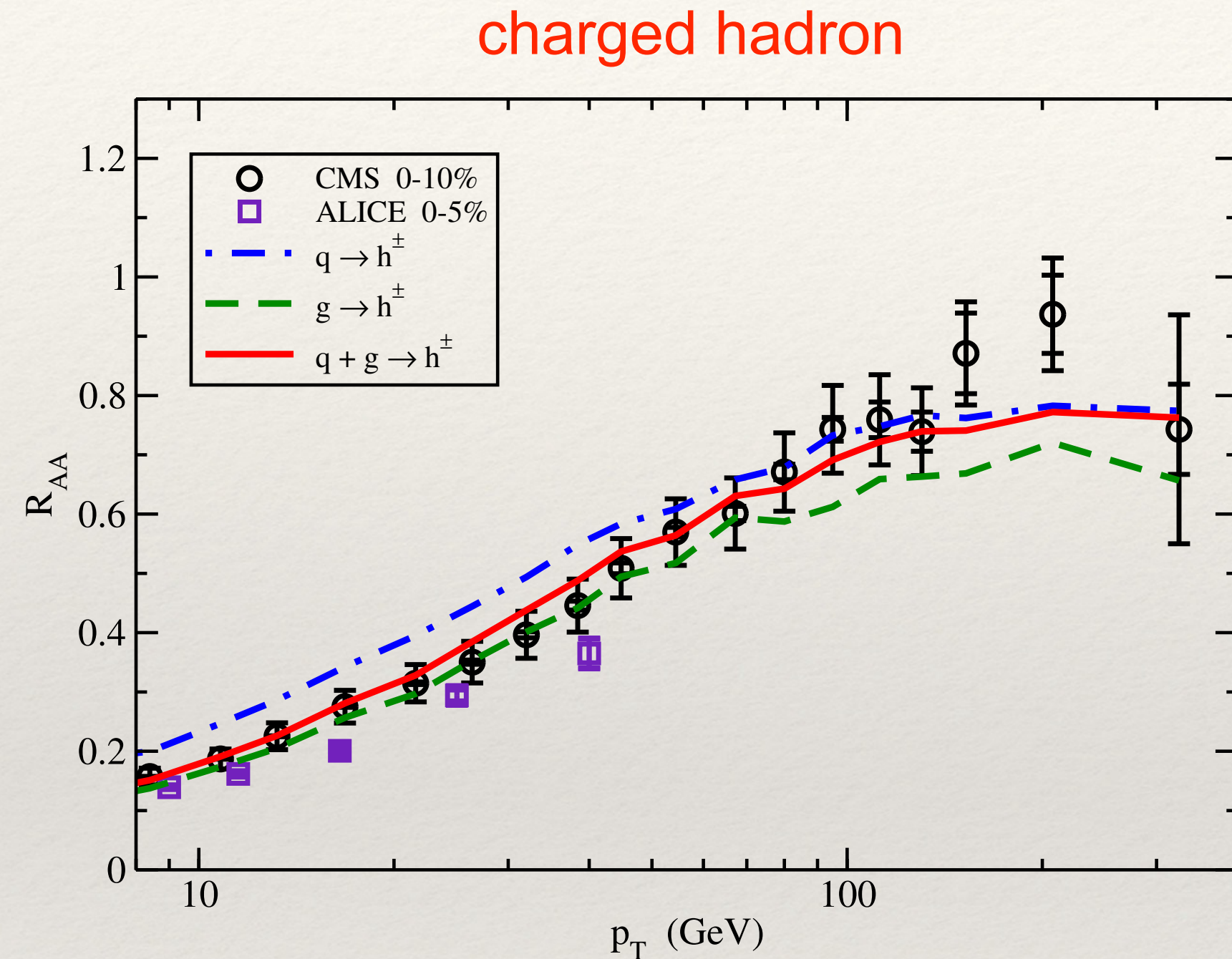
Different NLO contributions to light and heavy flavor hadrons



- dominates h^\pm production up to 50 GeV
- contributes to over 40% D up to 100 GeV

Flavor hierarchy in hadron suppression

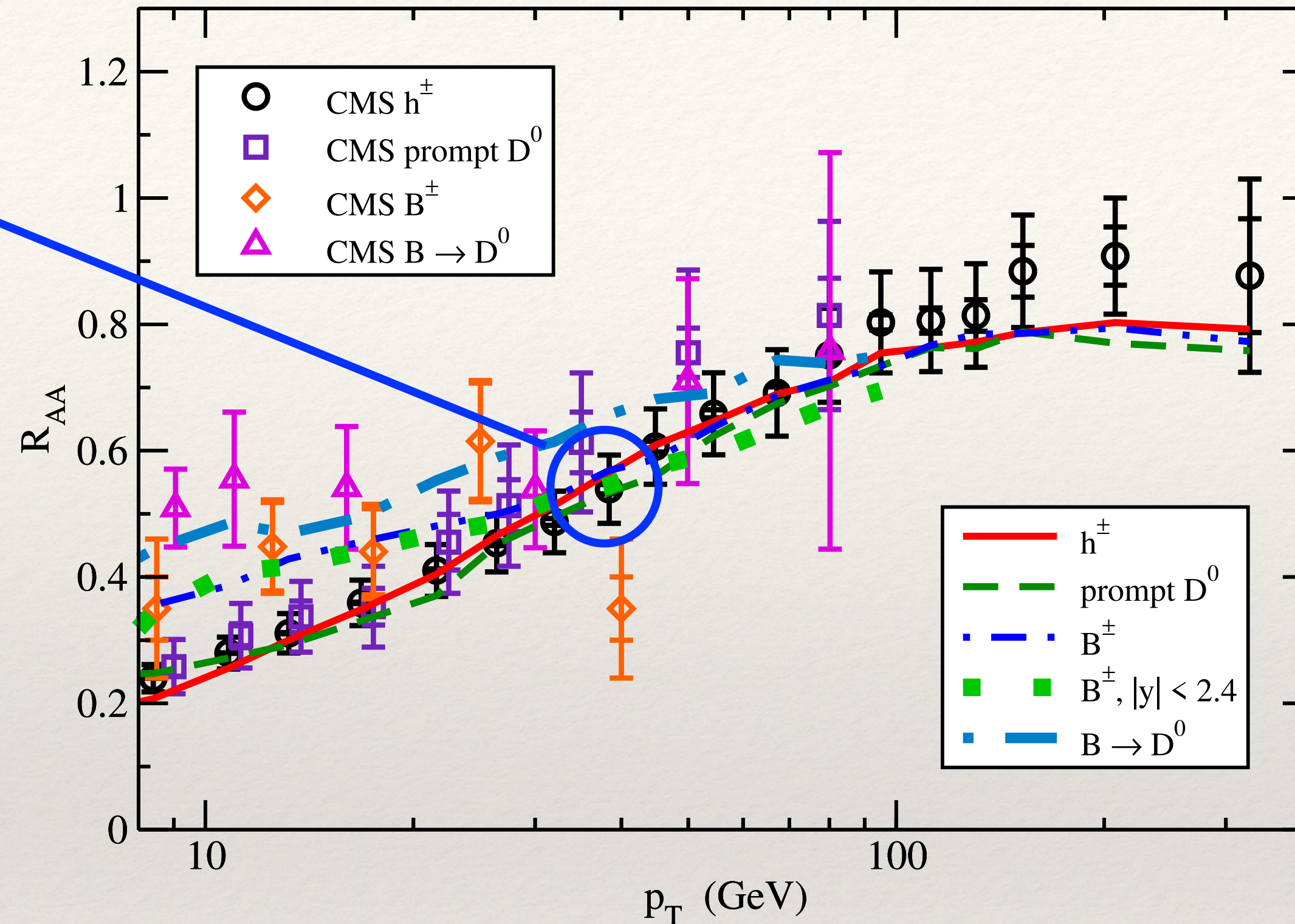
NLO initial production and fragmentation + Boltzmann transport + hydrodynamic medium for QGP



- g -initiated h & D $R_{AA} < q$ -initiated h & D R_{AA} [$\Delta E_g > \Delta E_{q/c}$]
- $R_{AA}(c \rightarrow D) > R_{AA}(q \rightarrow h)$ [$\Delta E_q > \Delta E_c$], $R_{AA}(g \rightarrow D) < R_{AA}(g \rightarrow h)$ [different FFs] $\Rightarrow R_{AA}(h) \approx R_{AA}(D)$
- Signature of flavor hierarchy of parton ΔE offset by NLO production/fragmentation in hadron R_{AA}

Flavor hierarchy in hadron suppression

Merging of D and B
 R_{AA} at $p_T \sim 40$ GeV



Xing, SC, Qin and Xing, Phys. Lett. B
805 (2020) 135424

- A simultaneous description of charged hadron, D meson, B meson, B -decay D meson R_{AA} 's starting from $p_T \sim 8$ GeV
- Predict R_{AA} separation between B and h / D below 40 GeV, but similar values above – **wait for confirmation from future precision measurement**

Extraction of parton energy loss from hadron R_{AA}

NLO initial production and fragmentation + **Parametrized** parton energy loss inside the QGP

• **Mean p_T loss:** $\langle \Delta p_T^j \rangle = C_j \beta_g p_T^\gamma \log(p_T)$

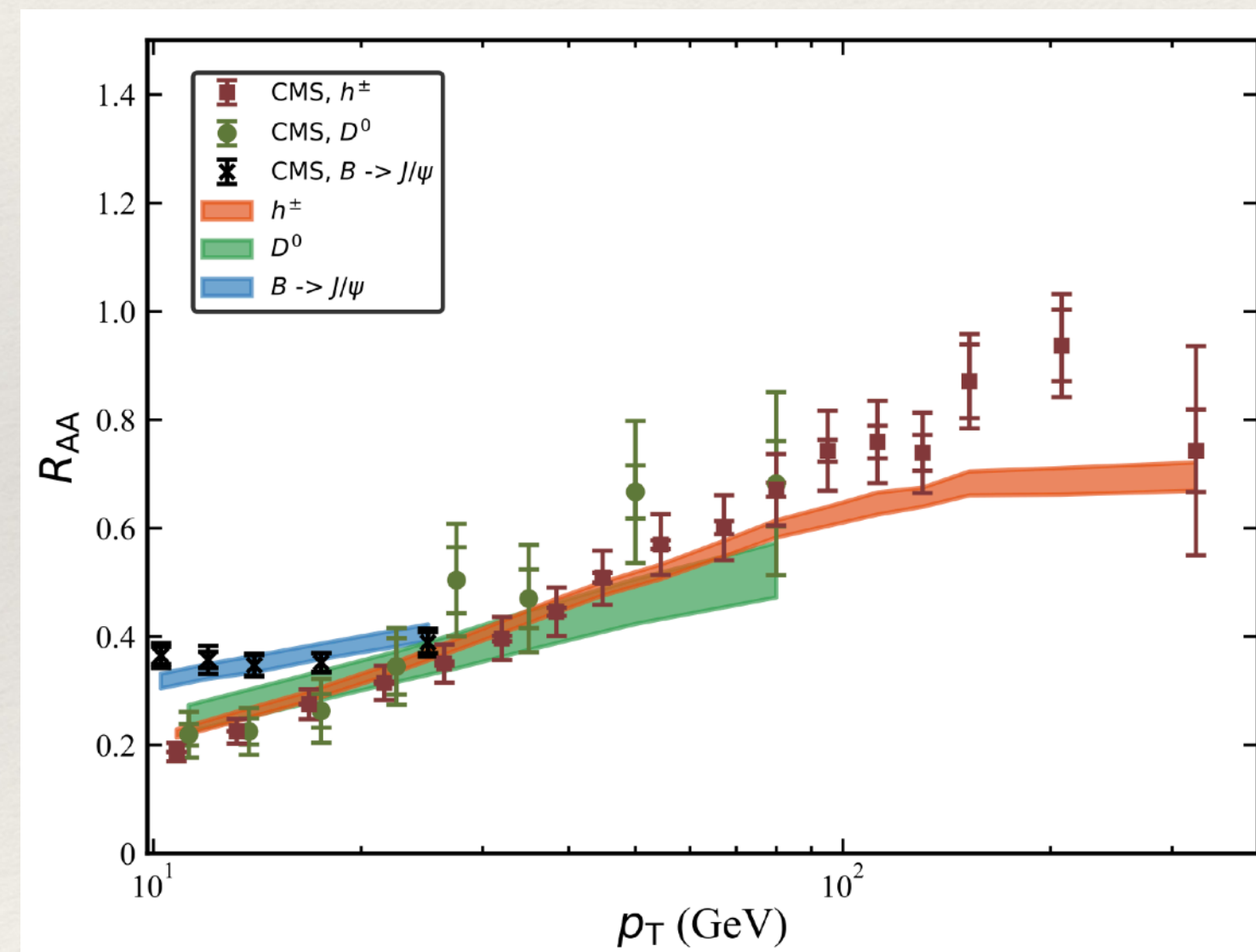
[Xing, SC, Qin, Phys. Lett. B 850 (2024) 138523]

- β_g : overall magnitude for g
- C_j : flavor dependence
- γ : p_T dependence

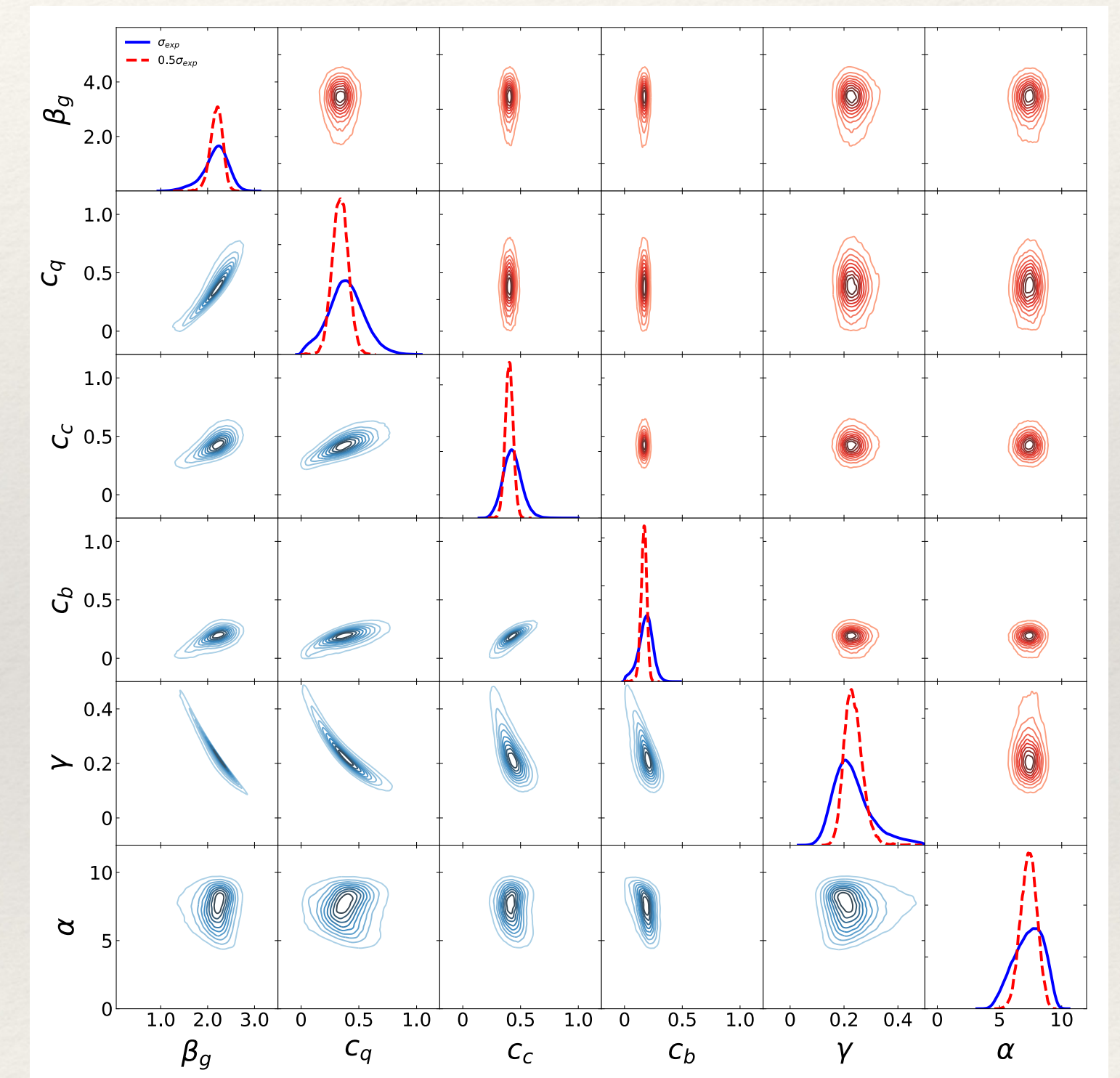
• **p_T loss distribution:**

$$W_{AA}(x) = \frac{\alpha^\alpha x^{\alpha-1} e^{-\alpha x}}{\Gamma(\alpha)}$$

$$x \equiv \Delta p_T / \langle \Delta p_T \rangle$$



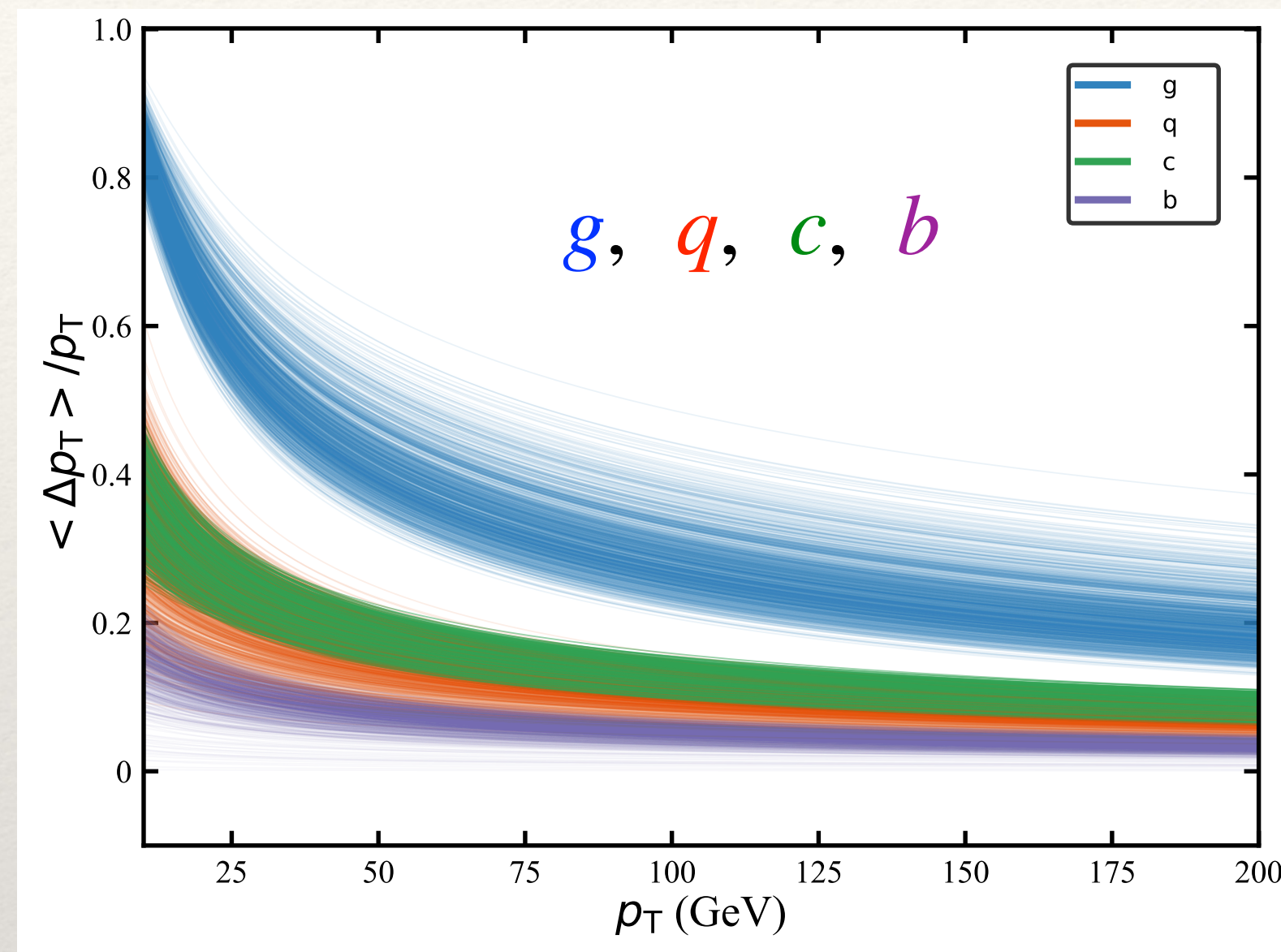
Bayesian calibration to data



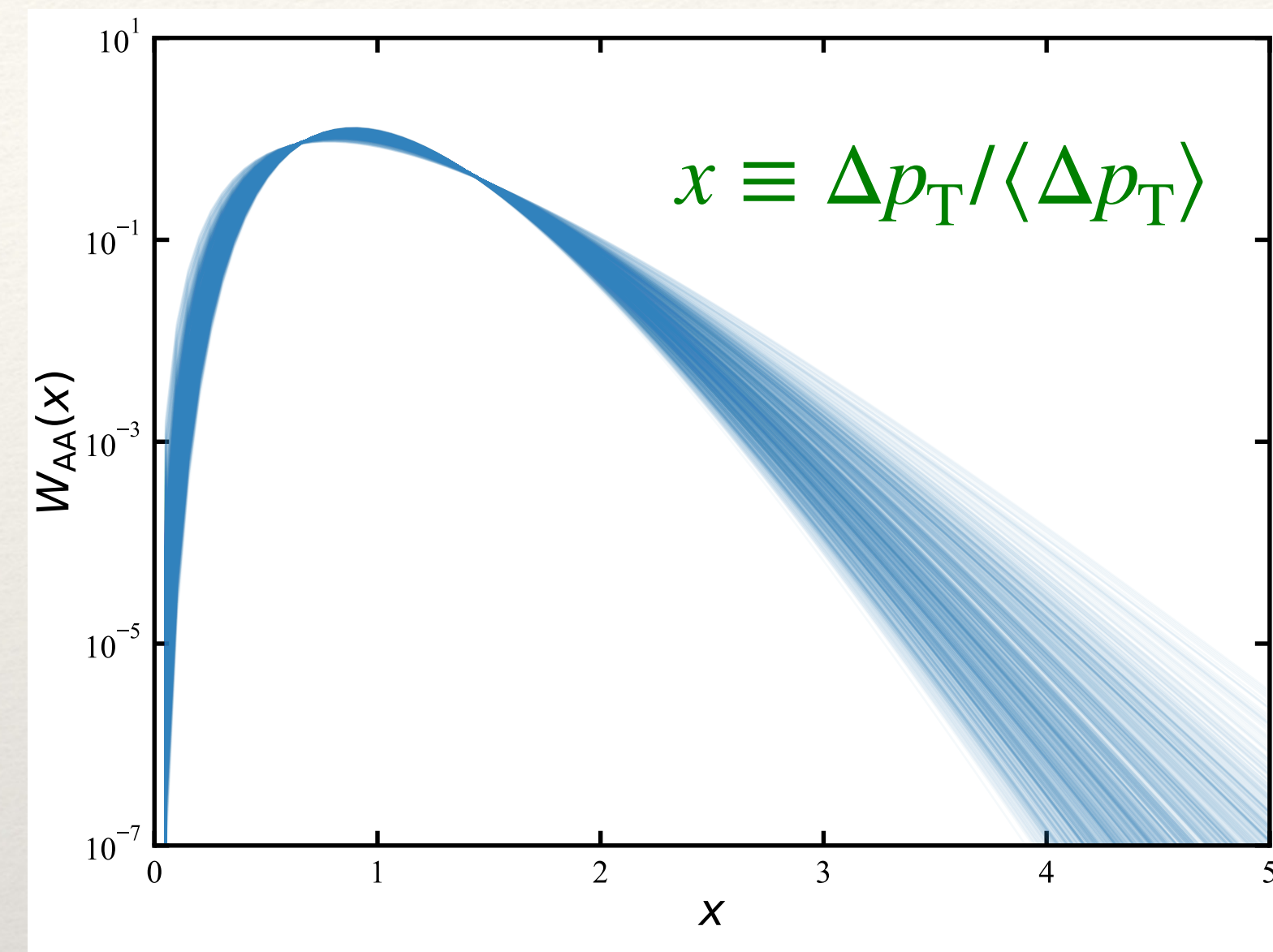
Constraints on parameters

Extraction of parton energy loss from hadron R_{AA}

Average energy loss



Energy loss distribution

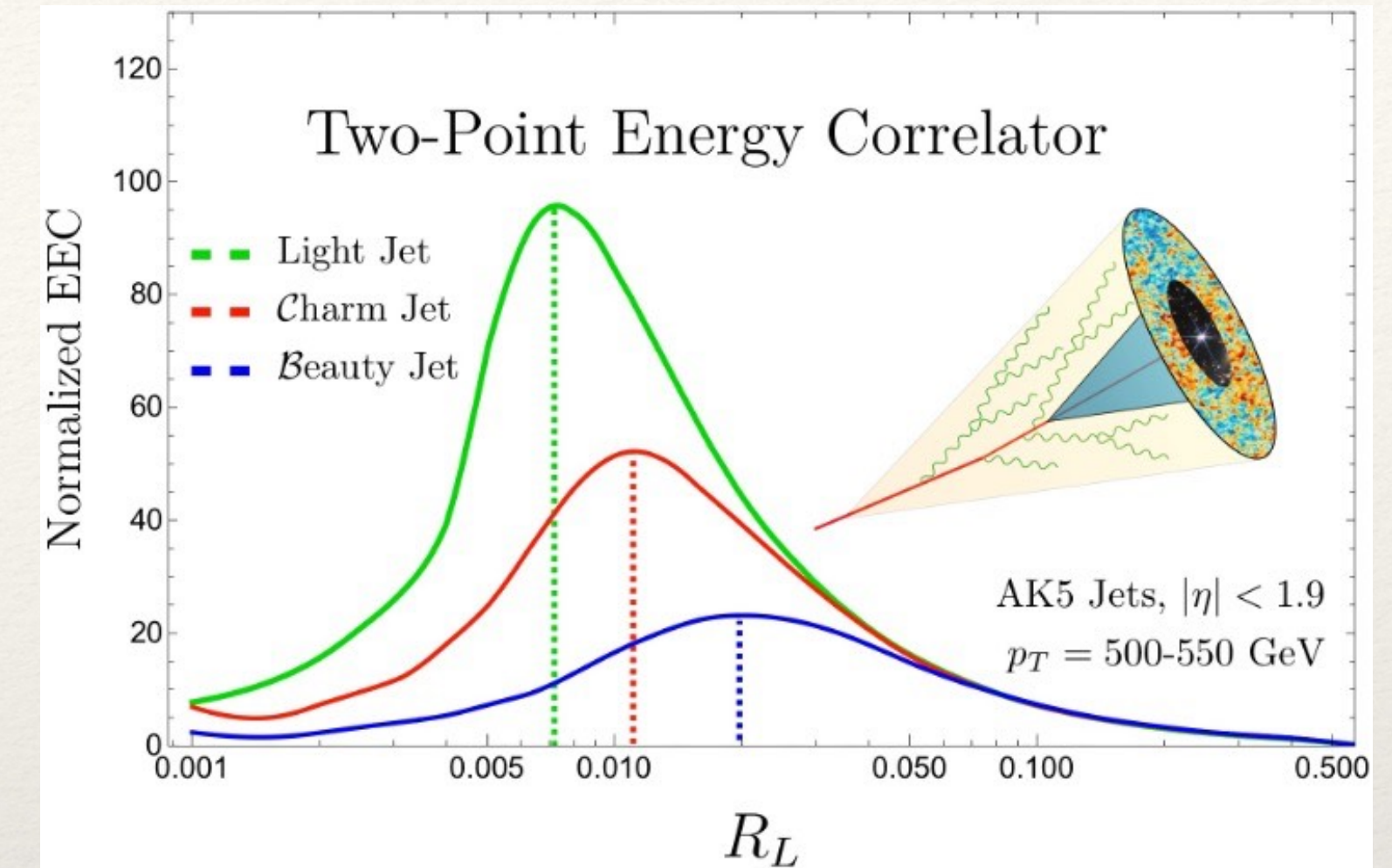
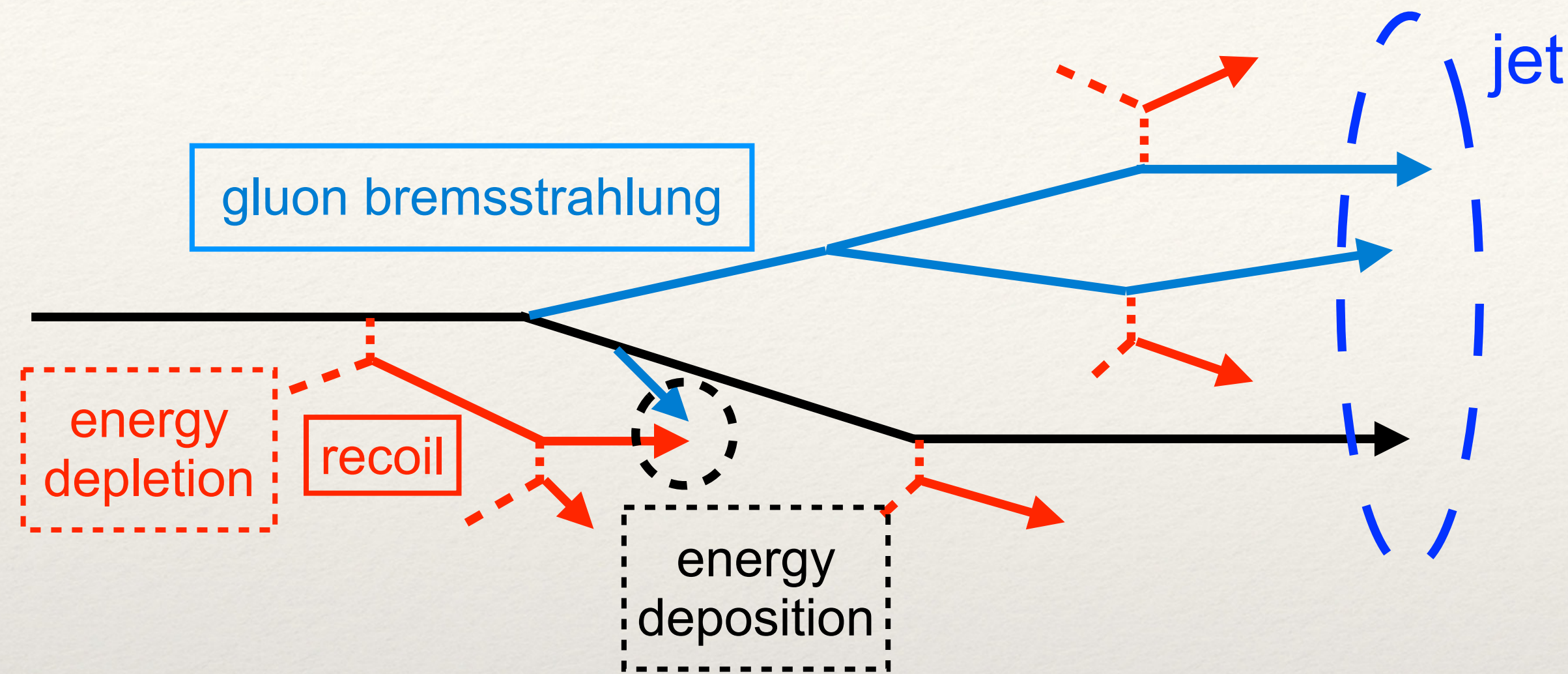


- $\Delta E_g > \Delta E_q \sim \Delta E_c > \Delta E_b$

- More stringent test on QCD calculation

- Flavor hierarchy of parton energy loss is encoded in the hadron R_{AA} data
- No obvious hierarchy for the hadron R_{AA} itself, due to the interplay between energy loss and NLO production and fragmentation

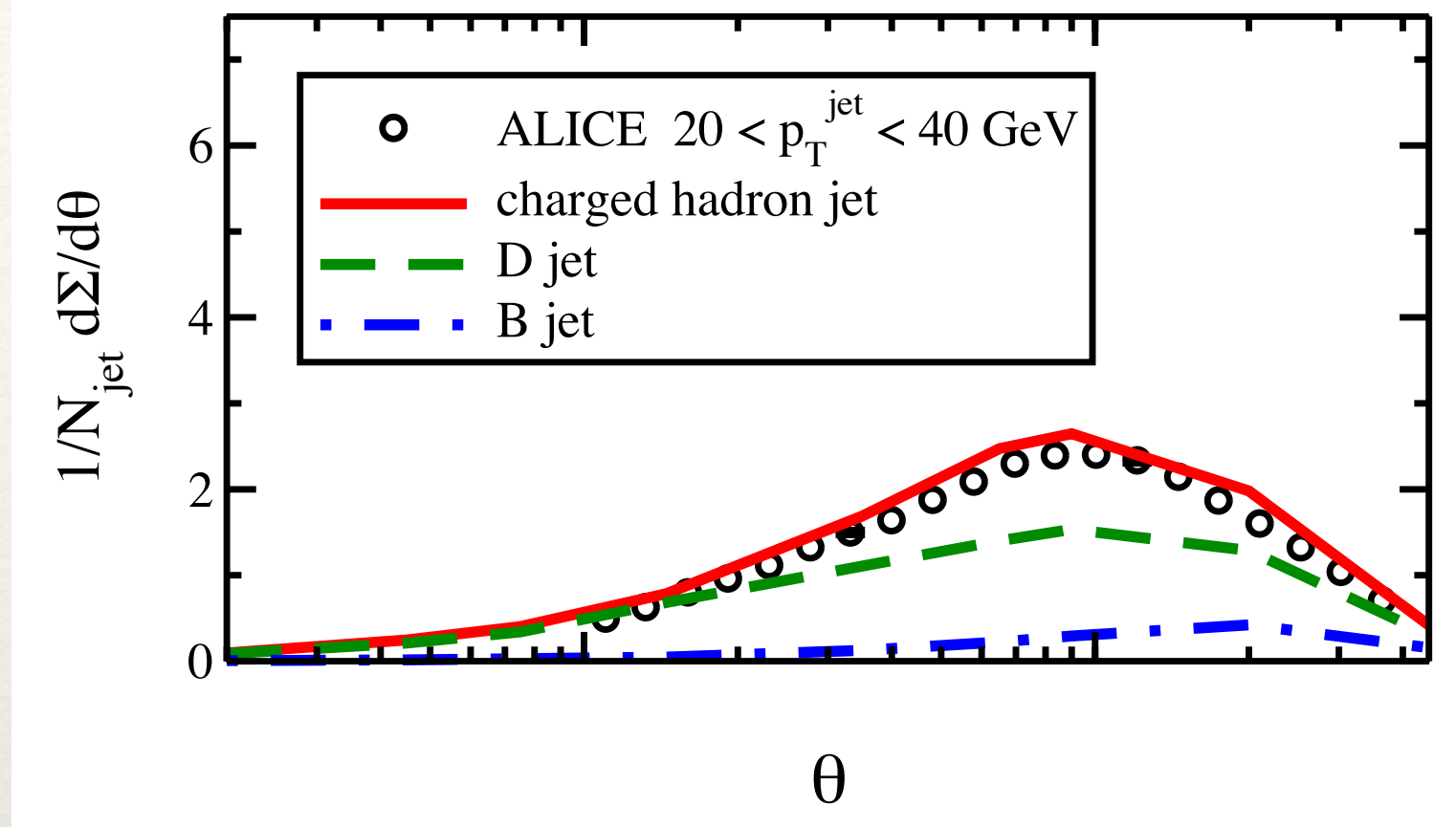
From heavy flavor hadrons to heavy flavor jets



[Craft et. al., arXiv:2210.09311]

- Description of full jets requires both **medium modification on jets** and **jet-induced medium excitation (recoil + energy depletion)**
- Energy-energy correlator (EEC) of jets presents a clear angular scale separation between perturbative and non-perturbative (e.g. hadronization) regions
- EEC can also reveal the flavor dependence of splitting angles of partons in pp collisions
- Implement a first realistic calculation on light and heavy flavor jet EEC in AA collisions

Light vs. heavy flavor jet EEC in pp collisions



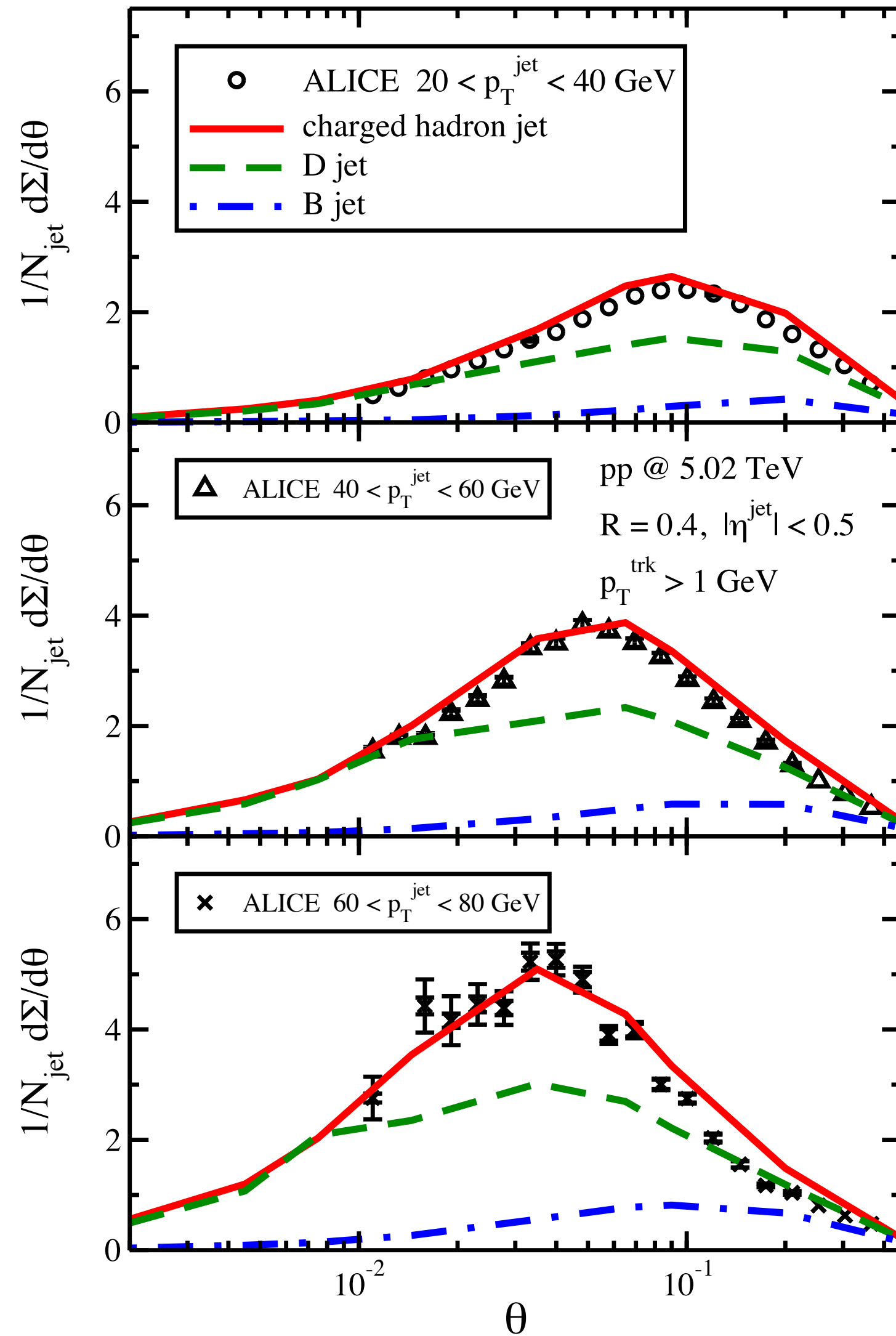
- Jet in pp : Pythia 8 simulation
- EEC of a jet (constituents denoted by i, j)

$$\frac{d\Sigma(\theta)}{d\theta} = \frac{1}{\Delta\theta} \sum_{|\theta_{ij}-\theta|<\theta/2} \frac{p_{T,i}(\vec{n}_i) p_{T,j}(\vec{n}_j)}{p_{T,\text{jet}}^2}$$

- Flavor (mass) dependence:
 - Overall magnitude: charged jet $>$ D -jet $>$ B -jet
 - Typical (peak) angle: charged jet $<$ D -jet $<$ B -jet

Suppression of splitting within $\theta_0 \sim m_Q/E_Q$ in vacuum

Light vs. heavy flavor jet EEC in pp collisions



- Jet in pp : Pythia 8 simulation
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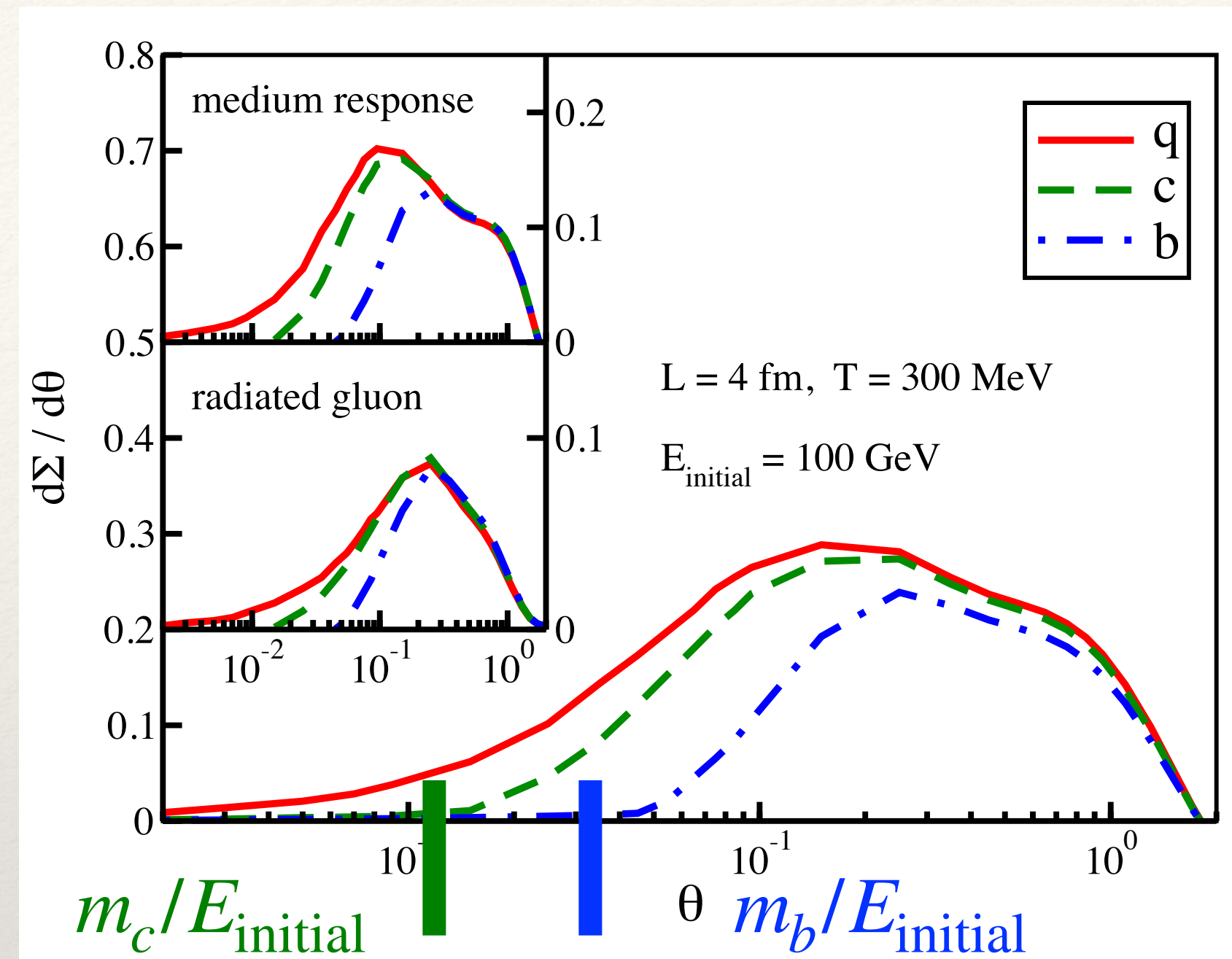
- Flavor (mass) dependence:
 - Overall magnitude: charged jet $>$ D -jet $>$ B -jet
 - Typical (peak) angle: charged jet $<$ D -jet $<$ B -jet

Suppression of splitting within $\theta_0 \sim m_Q/E_Q$ in vacuum

- Jet energy dependence
 - Higher $p_T \rightarrow \Sigma$ peaks at smaller θ

$p_T \theta_{\text{peak}} \sim$ transition scale between pert. and non-pert.

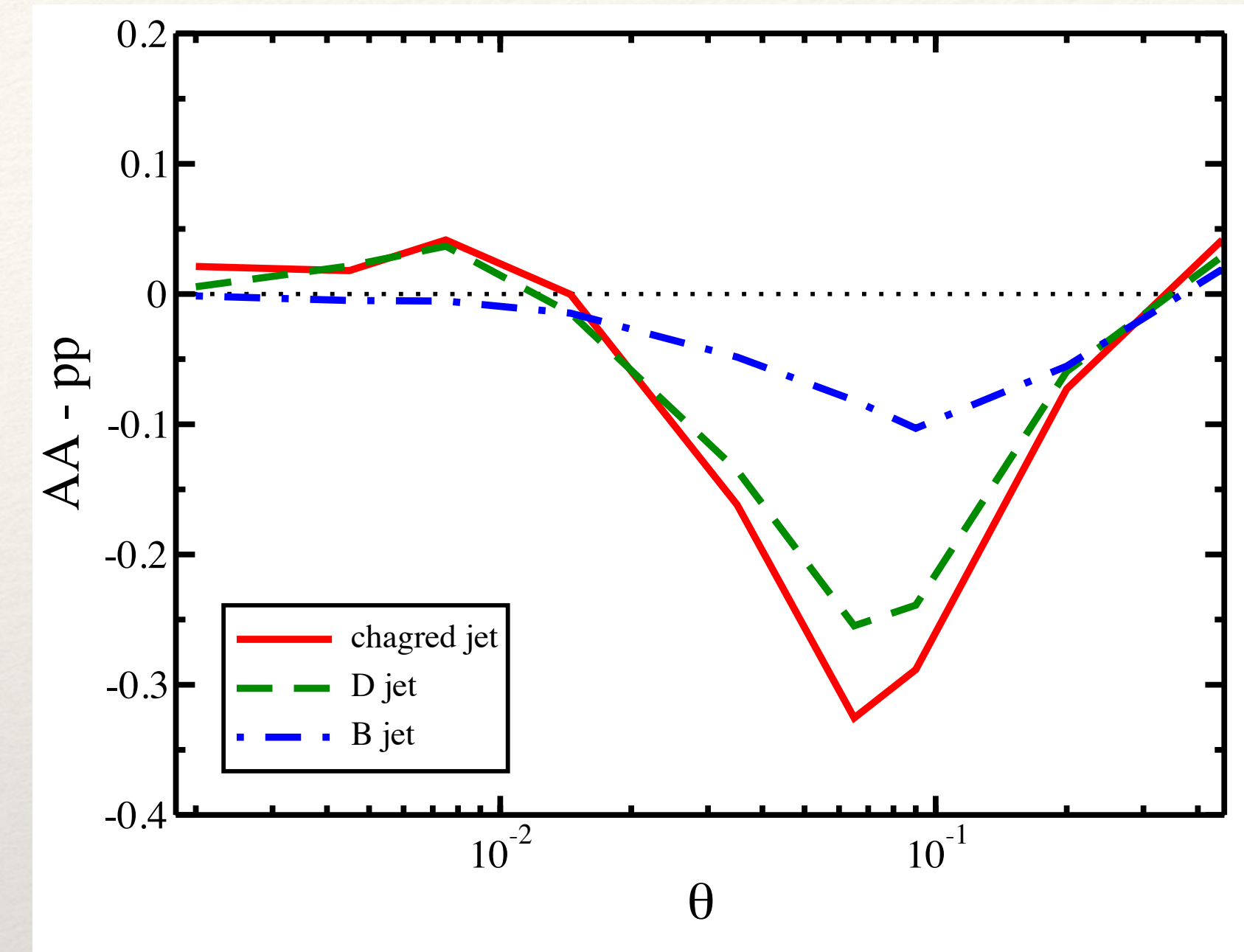
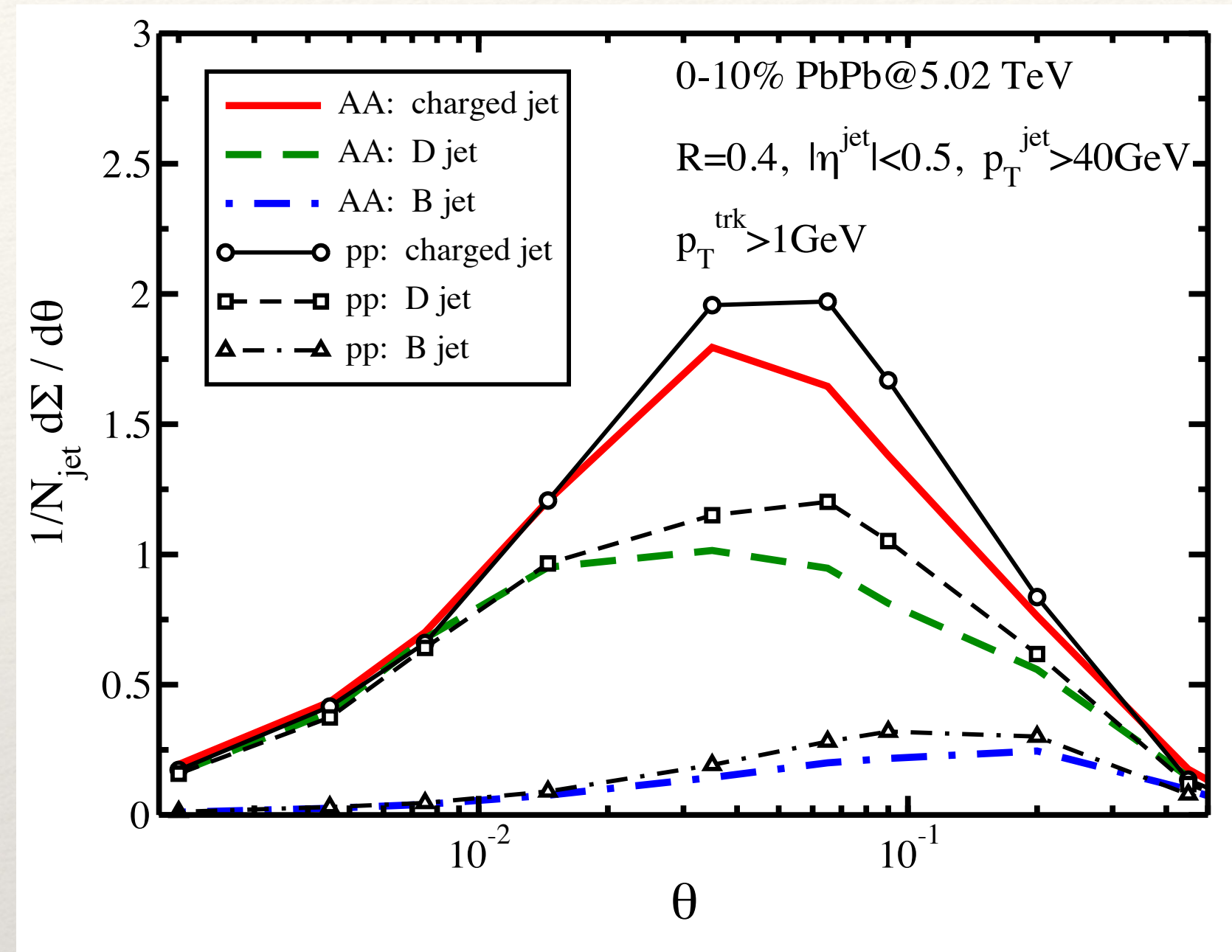
EEC of partons developed from a single quark



Xing, SC, Qin, Wang,
arXiv:2409.12843

- Single quark \rightarrow LBT + static medium \rightarrow EEC of daughter partons
- Flavor (mass) hierarchy of EEC:
 - Magnitude: charged $> D > B$ -jet; peak position: charged $< D < B$ -jet (similar to vacuum jets)
 - Clear strong suppression of Σ below $\theta_0 \sim m_Q/E_{\text{initial}}$
- Contributions from medium response and gluon emission show similar hierarchies

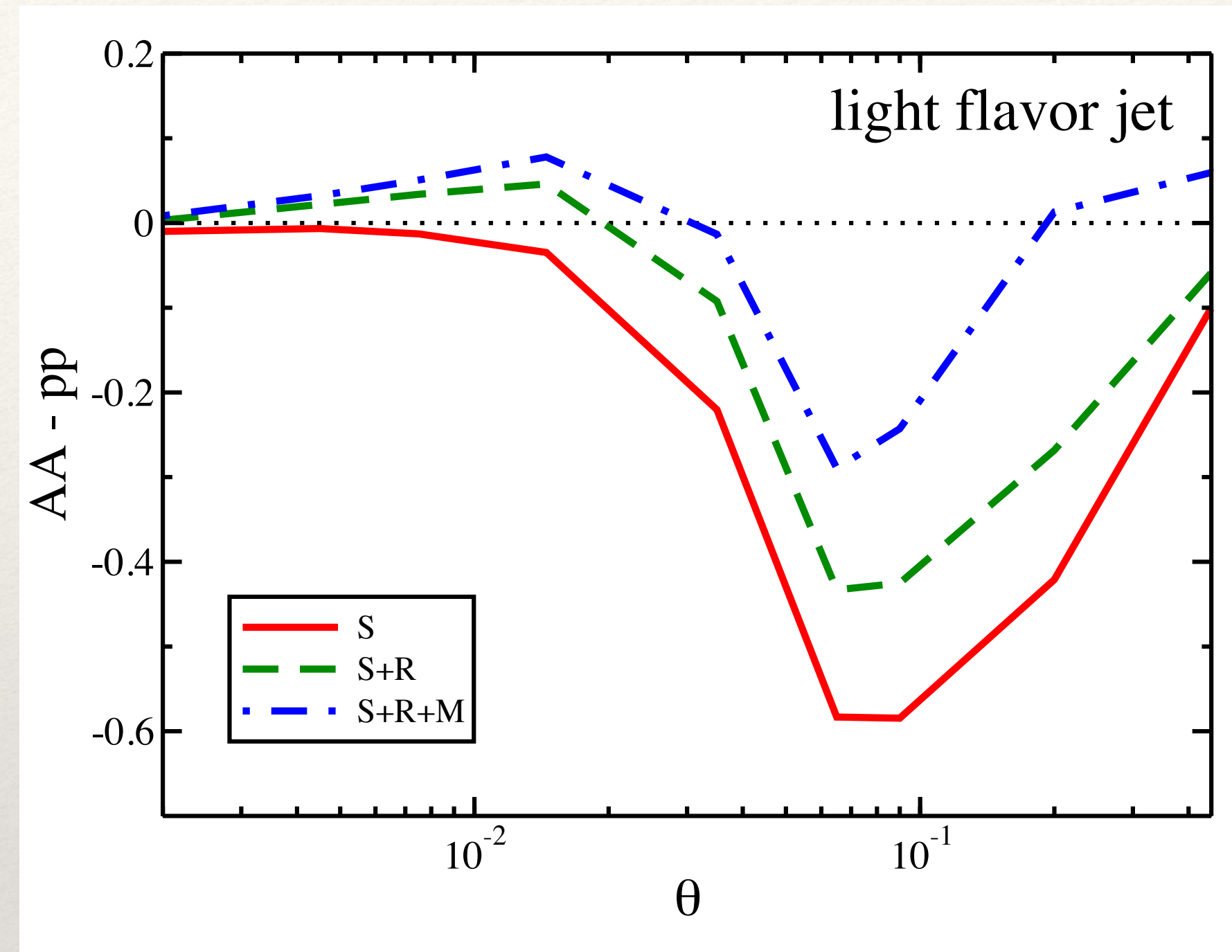
Light vs. heavy flavor jet EEC in central PbPb collisions



Nuclear modification ($AA - pp$) — Pythia + LBT in hydro

- General features: suppression at intermediate θ , enhancement at small θ (except for B -jet) and large θ
- Flavor hierarchy: weaker nuclear modification (both suppression and enhancement) for jets tagged with heavier mesons

Different contributions to medium modification on EEC



S: shower partons inherited from Pythia
S+R: add medium-induced gluons
S+R+M: further add medium response

- Jet energy loss causes suppression over the entire θ region
- Medium-induced gluon emission enhances EEC at small θ
- Medium response enhances EEC at large θ

Probing the EoS of QGP using heavy quarks

[F.-L. Liu, X.-Y. Wu, SC, G.-Y. Qin, X.-N. Wang, Phys. Lett. B 848 (2024) 138355]

Transport

$$p_a \cdot \partial f_a(x_a, p_a) = E_a (\mathcal{C}_a^{\text{el}} + \mathcal{C}_a^{\text{inel}})$$

Strong coupling strength

$$g(E, T)$$

Thermal mass of partons

$$m_g^2 = \frac{1}{6} g^2 \left[(N_c + \frac{1}{2} n_f) T^2 + \frac{N_c}{2\pi^2} \sum_q \mu_q^2 \right]$$

$$m_{u,d}^2 = \frac{N_c^2 - 1}{8N_c} g^2 \left[T^2 + \frac{\mu_{u,d}^2}{\pi^2} \right]$$

$$m_s^2 - m_{0s}^2 = \frac{N_c^2 - 1}{8N_c} g^2 \left[T^2 + \frac{\mu_s^2}{\pi^2} \right]$$

Equation of state

$$\begin{aligned} P_{qp}(m_u, m_d, \dots, T) &= \sum_{i=u,d,s,g} d_i \int \frac{d^3 p}{(2\pi)^3} \frac{|\vec{p}^2|}{3E_i(p)} f_i(p) - B(T) \\ &= \sum_i P_{kin}^i(m_i, T) - B(T) \end{aligned}$$

$$\epsilon = T dP(T)/dT - P(T), \quad s = (\epsilon + P)/T$$

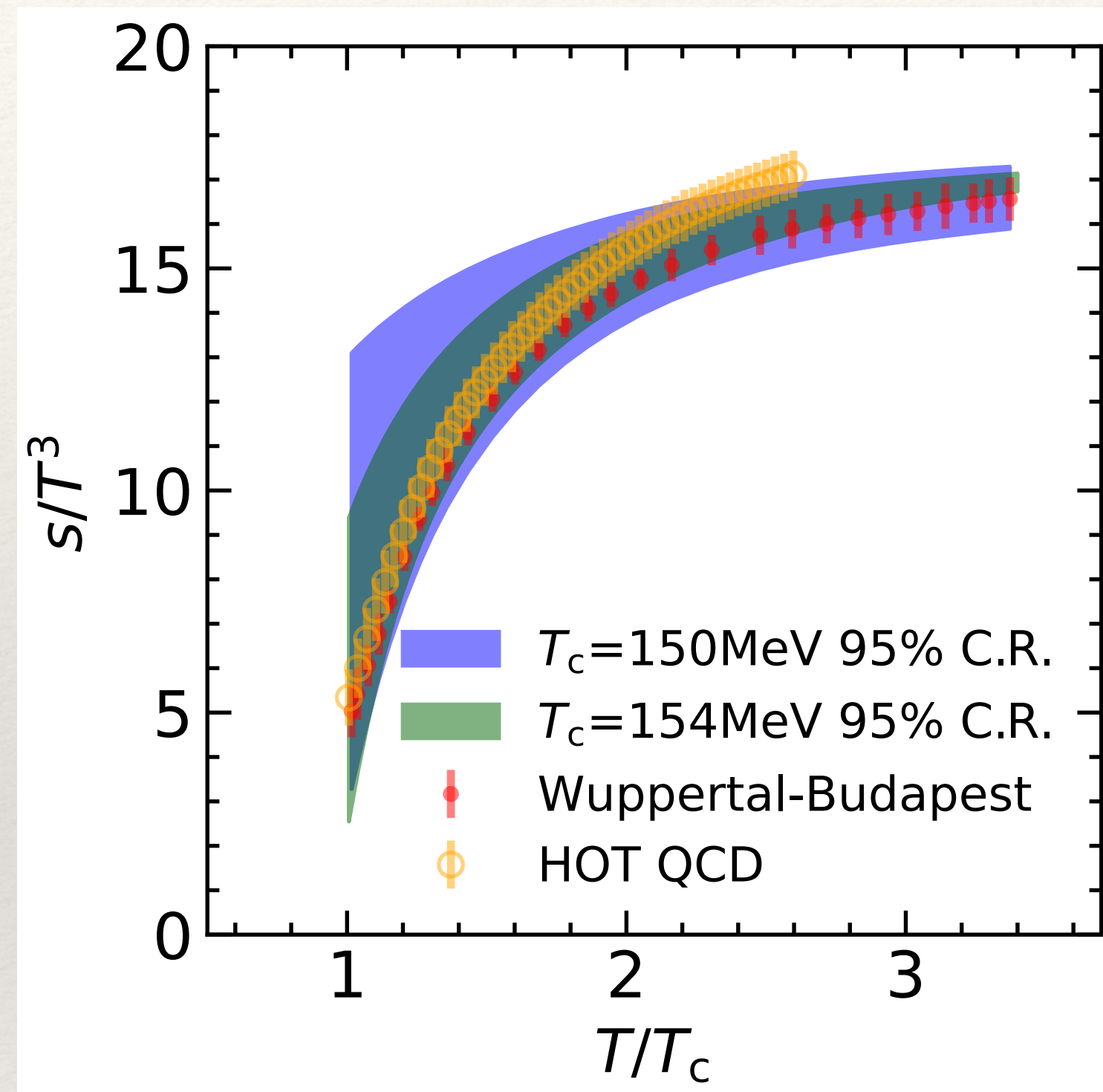
Strategy:

Fit g from comparing
transport model to data

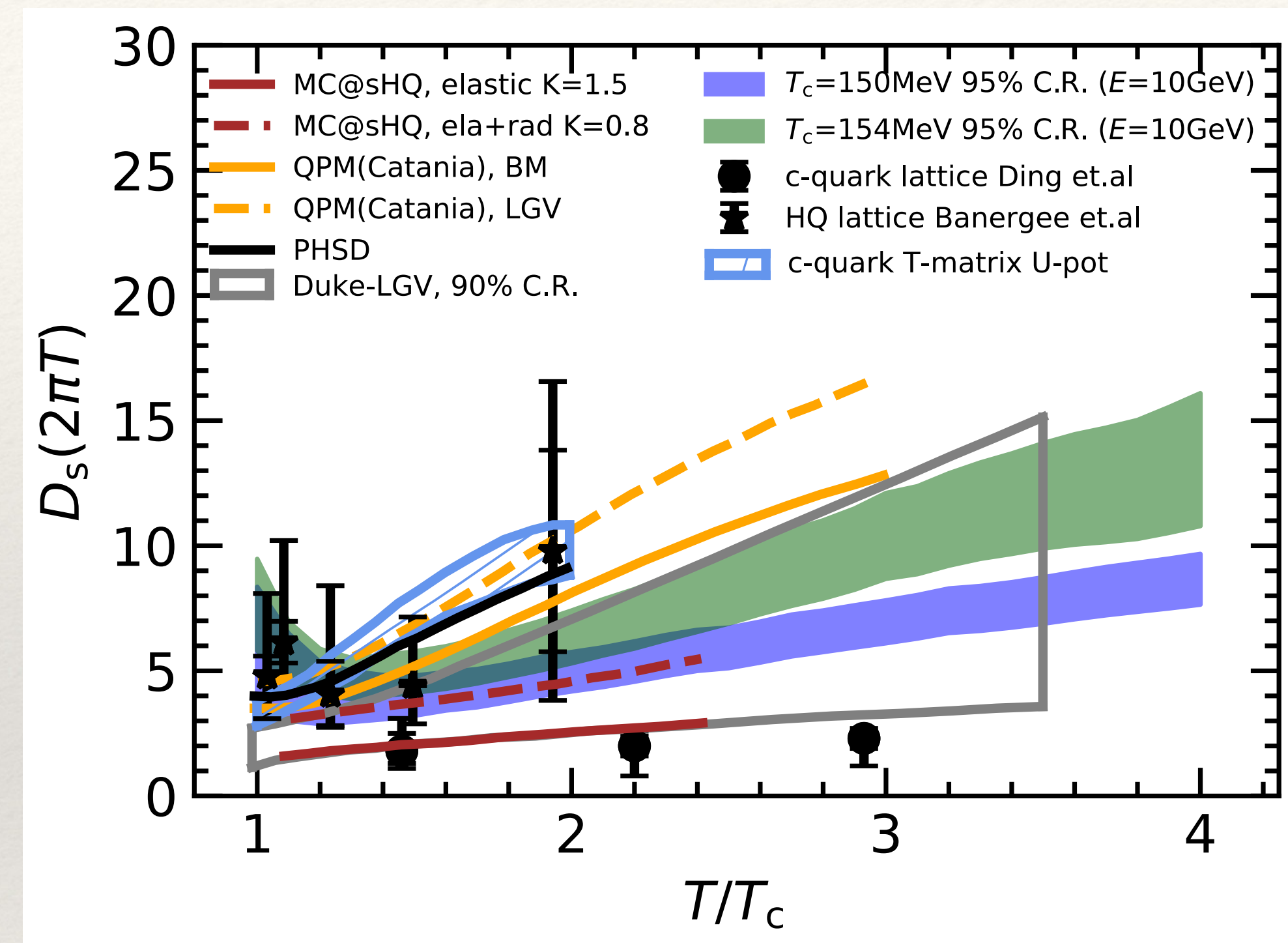
Calculate EoS from g

EoS of QGP and diffusion coefficient of heavy quarks

Equation of state of the QGP



Diffusion coefficient of heavy quarks



- Agreement with the lattice data
- Simultaneous constraint on QGP properties and transport properties of hard probes

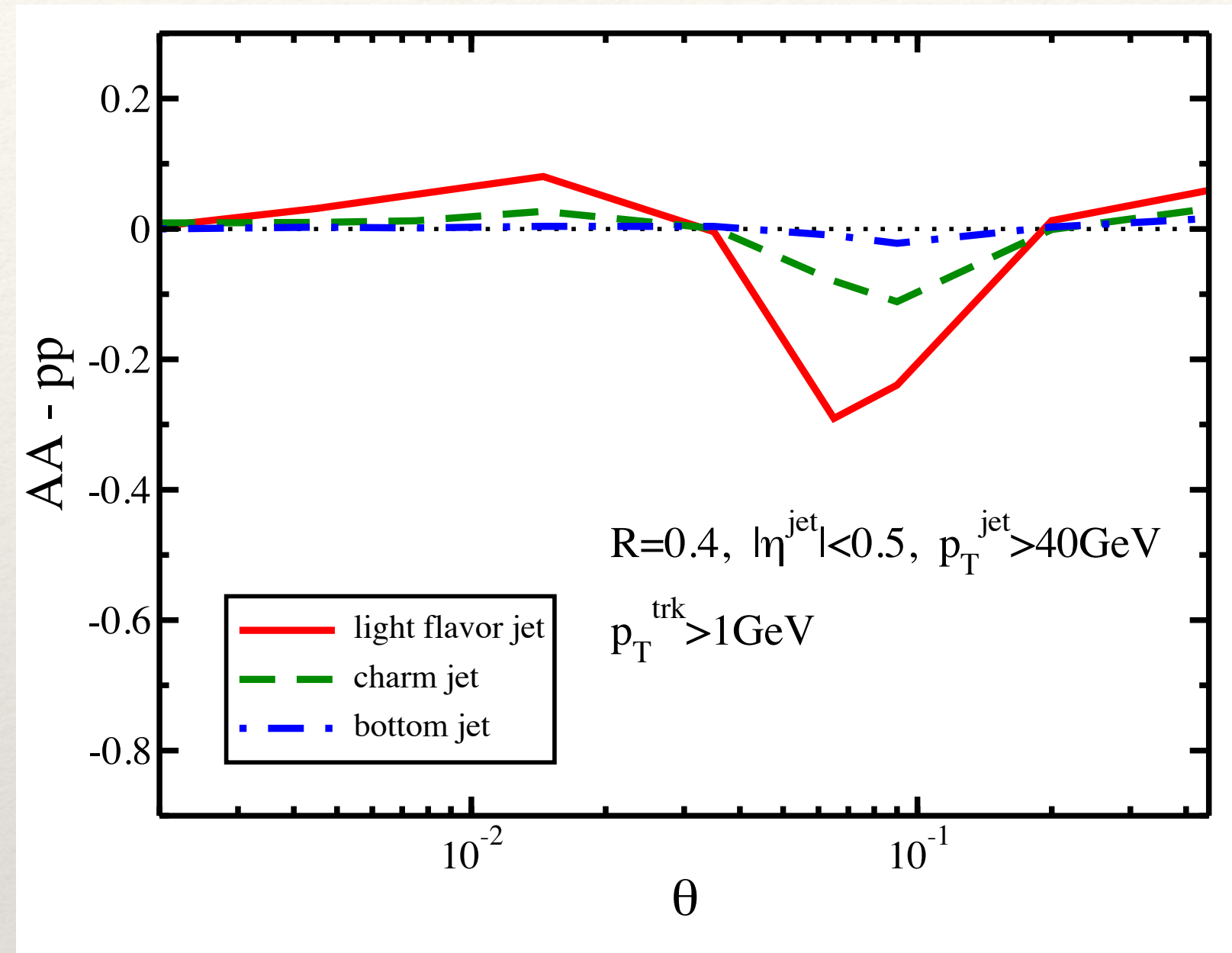
Summary

- Developed a transport model for studying both heavy flavor hadrons and jets
- Flavor hierarchy of parton energy loss is encoded in the hadron R_{AA} , though not explicit due to the interplay between energy loss and NLO contributions
- The jet EEC is an excellent observable to study the dead cone effect on parton splitting (magnitude and peak position of EEC) in both pp and AA collisions
- Heavy flavor observables can be used to constrain the QGP properties (EoS)

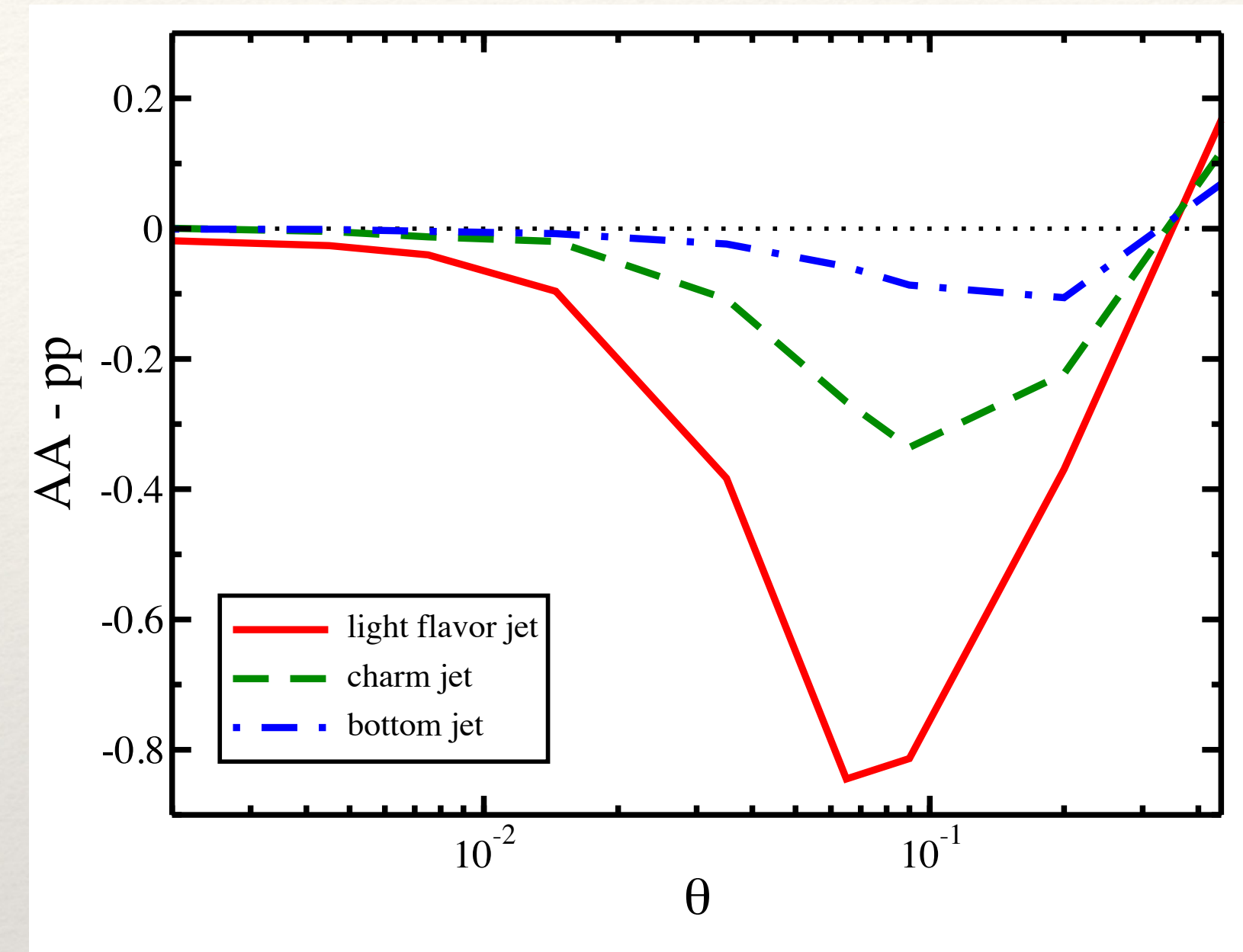
Thank. you!

Effects of trigger bias on the jet EEC

p_T trigger in both pp and AA



p_T trigger only in pp (no trigger bias in AA)



- AA jets with trigger bias originate from pp jets with higher p_T and initial virtuality scale
 - Stronger but narrower vacuum splittings
 - Enhances EEC at small θ , reduces the suppression/enhancement at intermediate/large θ
- Can be tested using γ -jets

Parametrization and Bayesian analysis

Strong coupling strength

Interaction between thermal partons (thermal scale):

$$g^2(T) = \frac{48\pi^2}{(11N_c - 2N_f) \ln \left[\frac{(aT/T_c + b)^2}{1 + ce^{-d(T/T_c)^2}} \right]}$$

Interaction with hard partons (parton energy scale):

$$g^2(E) = \frac{48\pi^2}{(11N_c - 2N_f) \ln [(AE/T_c + B)^2]}$$

Parameters: $\theta = (a, b, c, d, A, B)$

Bayes Theorem

$$P(\theta|\text{data}) \propto P(\text{data}|\theta)P(\theta)$$

posterior distribution

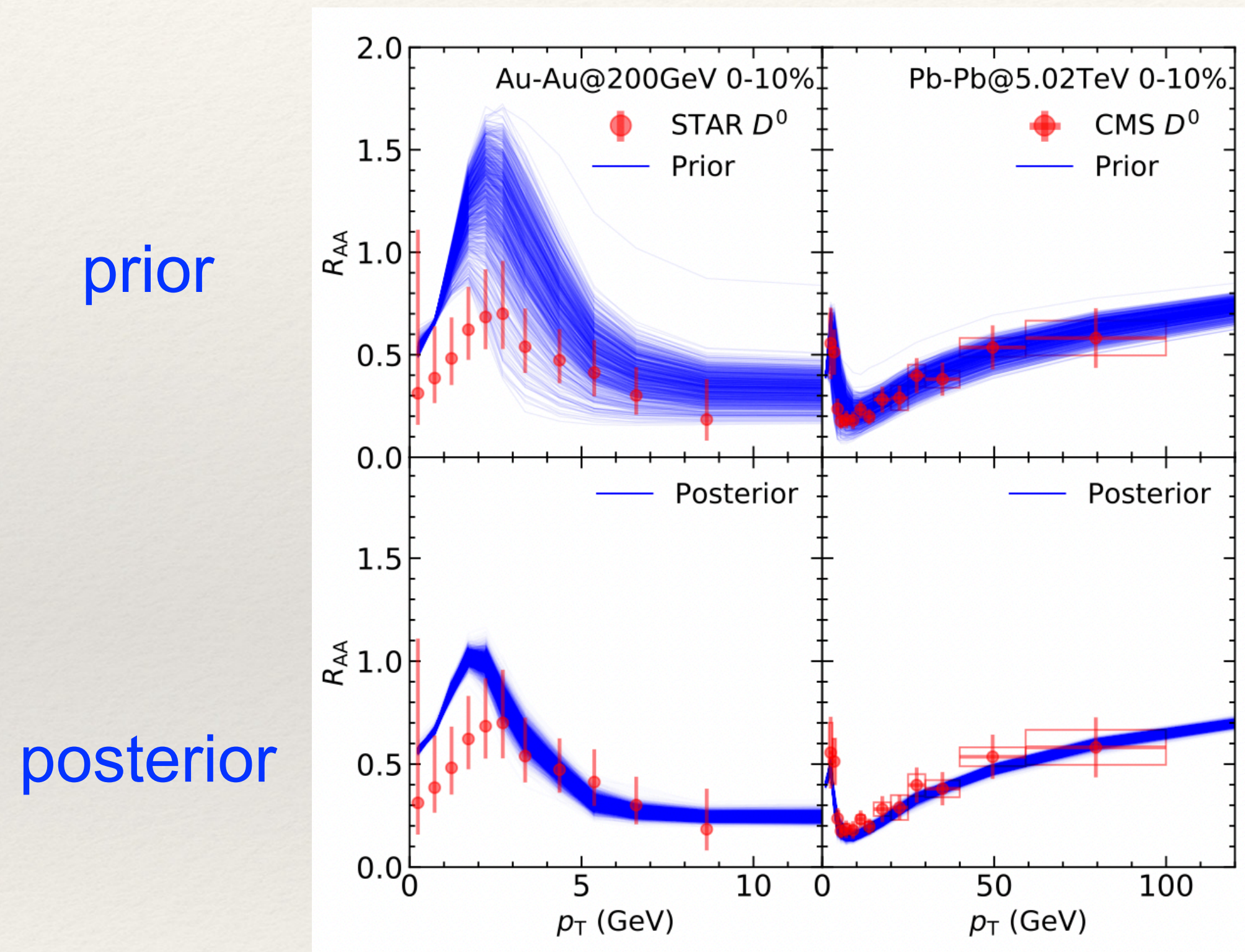
prior distribution

model-to-data comparison

$$P(\text{data}|\theta) = \prod_i \frac{1}{\sqrt{2\pi}\sigma_i} e^{-\frac{[y_i(\theta) - y_i^{\text{exp}}]^2}{2\sigma_i^2}}$$

Model calibration and parameter extraction

Calibration against observables



(Two examples from many observables)

Extraction of model parameters for $g(E, T)$

