

Jet Modification and Hard-Soft Correlations (SoftJet 2024)

**Study of jet-induced medium response  
in high-energy heavy-ion collisions  
within the CoLBT-hydro model**

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Tokyo Japan  
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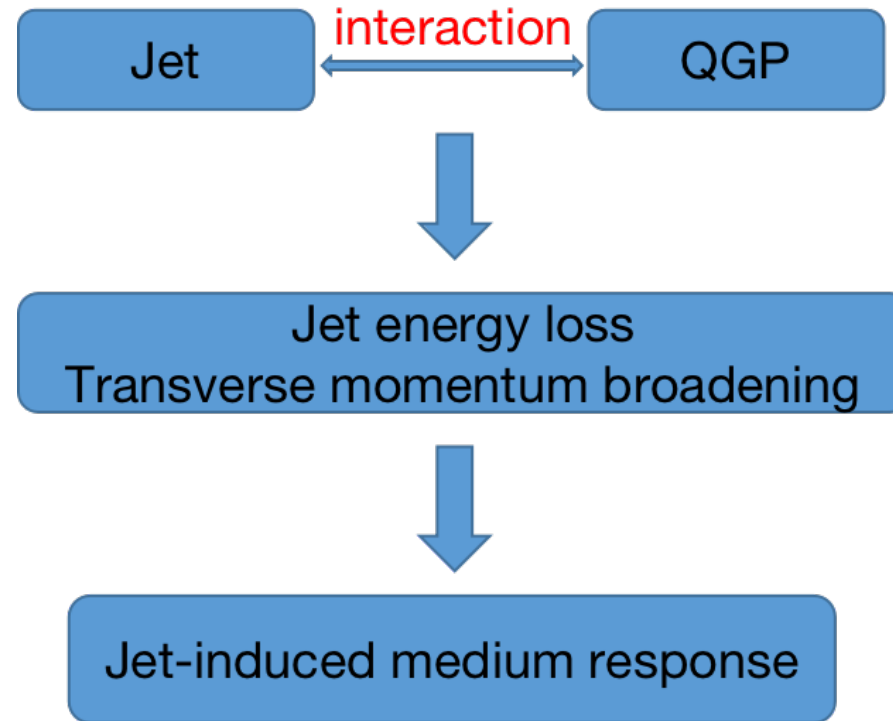
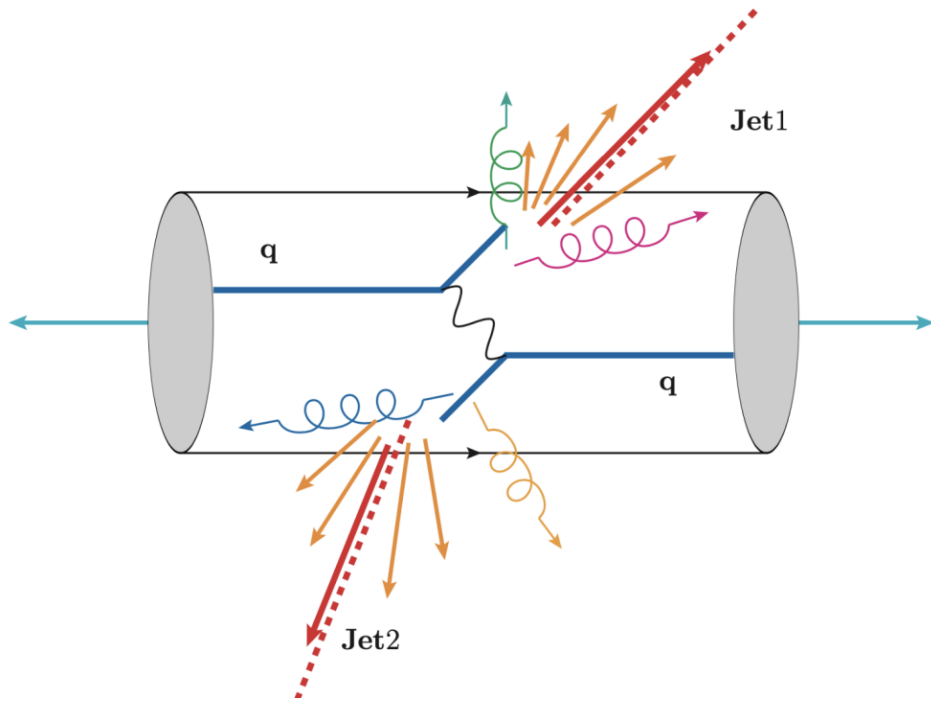


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# Outline

- **Introduction**
- **LBT/CoLBT-hydro model**
- **Jet-induced medium response**
  - **Jet-induced diffusion wake**
  - **3D structure of diffusion wake**
  - **EECs in realistic heavy-ion collisions**
- **Summary and Outlook**

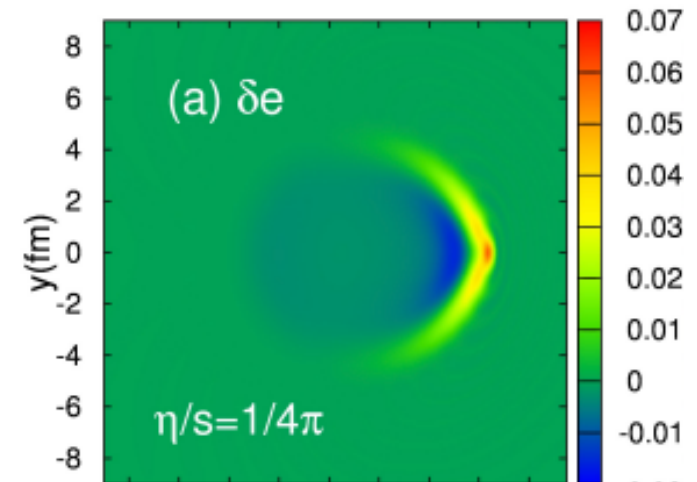
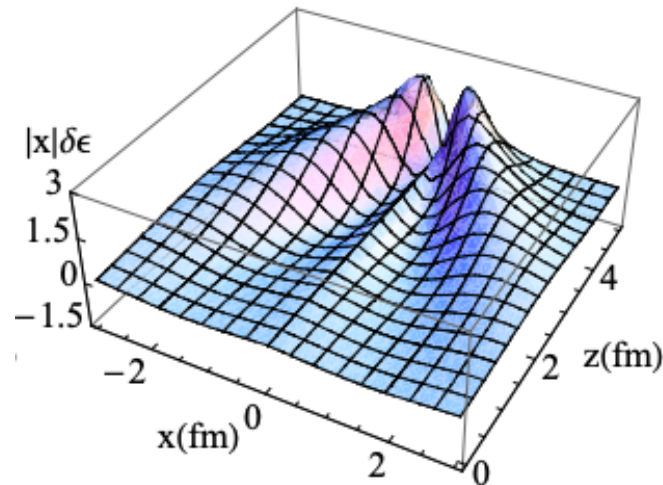
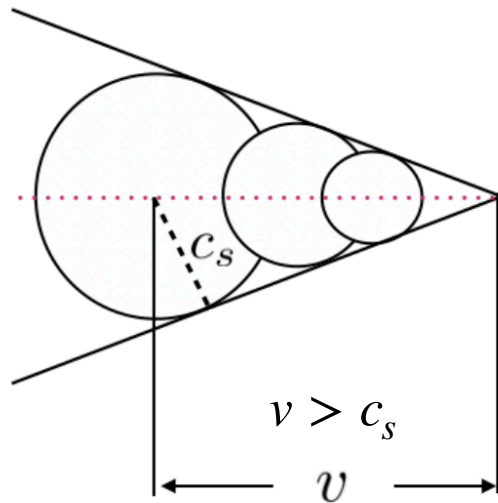
# Jet in heavy-ion collisions



# Jet-induced medium response

## Jet-induced medium response in the form of Mach-cone-like excitation.

[Casalderrey-Solana, Shuryak, Teaney, 2005; Ruppert, Muller, 2005; Gubser, Pufu, 2008; Qin, Majumder, Song, Heinz, 2009; Yan, Jean, Gale, 2017; ...]

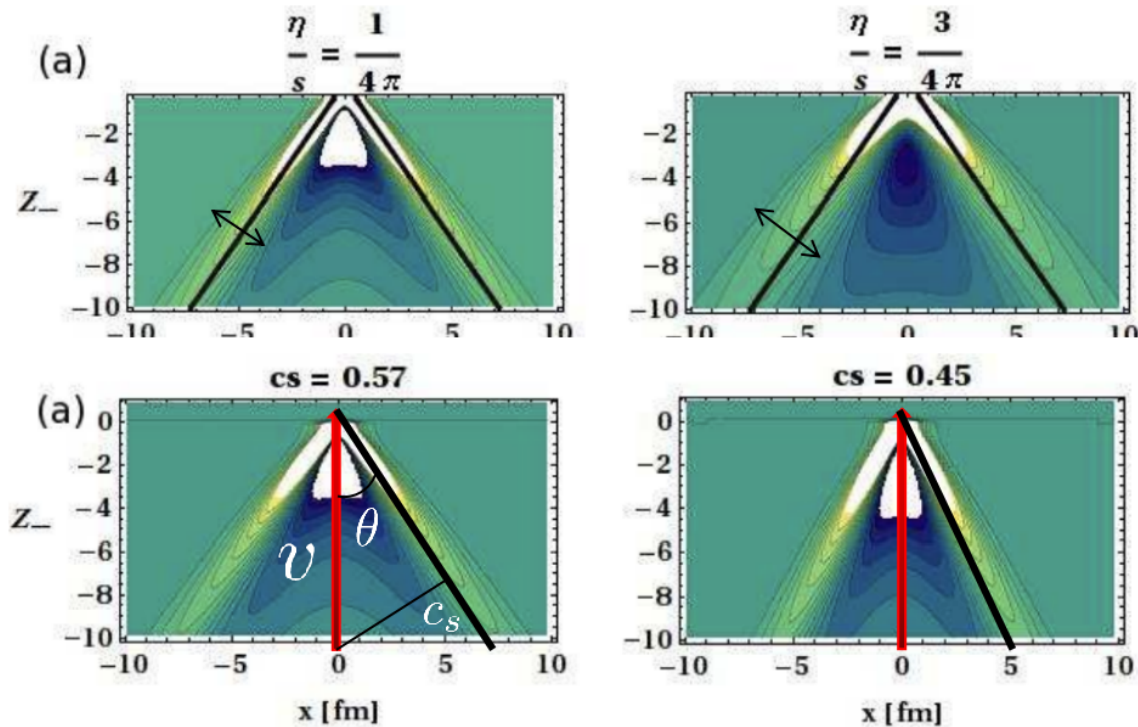


The front wake which is along with the jet direction will modify the particles distribution in jet.



# Jet-induced medium response

Jet-induced Mach-cone could extract the QGP properties



(1) **Width of front wake** of Mach cone is related with viscous properties of QGP medium;

(2) **Mach cone angle** is sensitive to EoS.

$$\sin\theta = \frac{c_s}{v}$$

R.B.Neufeld Phys.Rev.C 79 (2009) 054909

# Linear Boltzmann Transport(LBT) model

$$p_1 \partial f_1 = - \int dp_2 dp_3 dp_4 (f_1 f_2 - f_3 f_4) |M_{12 \rightarrow 34}|^2 (2\pi)^4 \delta^4(\sum_i p^i) + inelastic$$

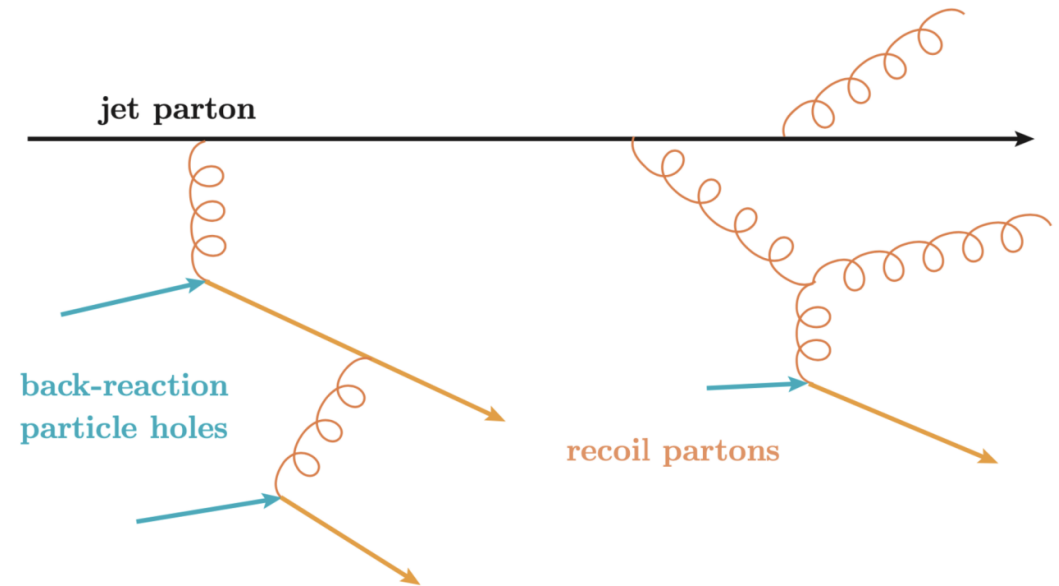
## Medium-induced gluon(High-Twist):

[Wang, Guo, 2001]

$$\frac{dN_g}{dz d^2 k_{\perp} dt} \approx \frac{2C_A \alpha_s}{\pi k_{\perp}^4} P(z) \hat{q} (\hat{p} \cdot u) \sin^2 \frac{k_{\perp}^2 (t - t_0)}{4z(1-z)E}$$

### Tracked partons:

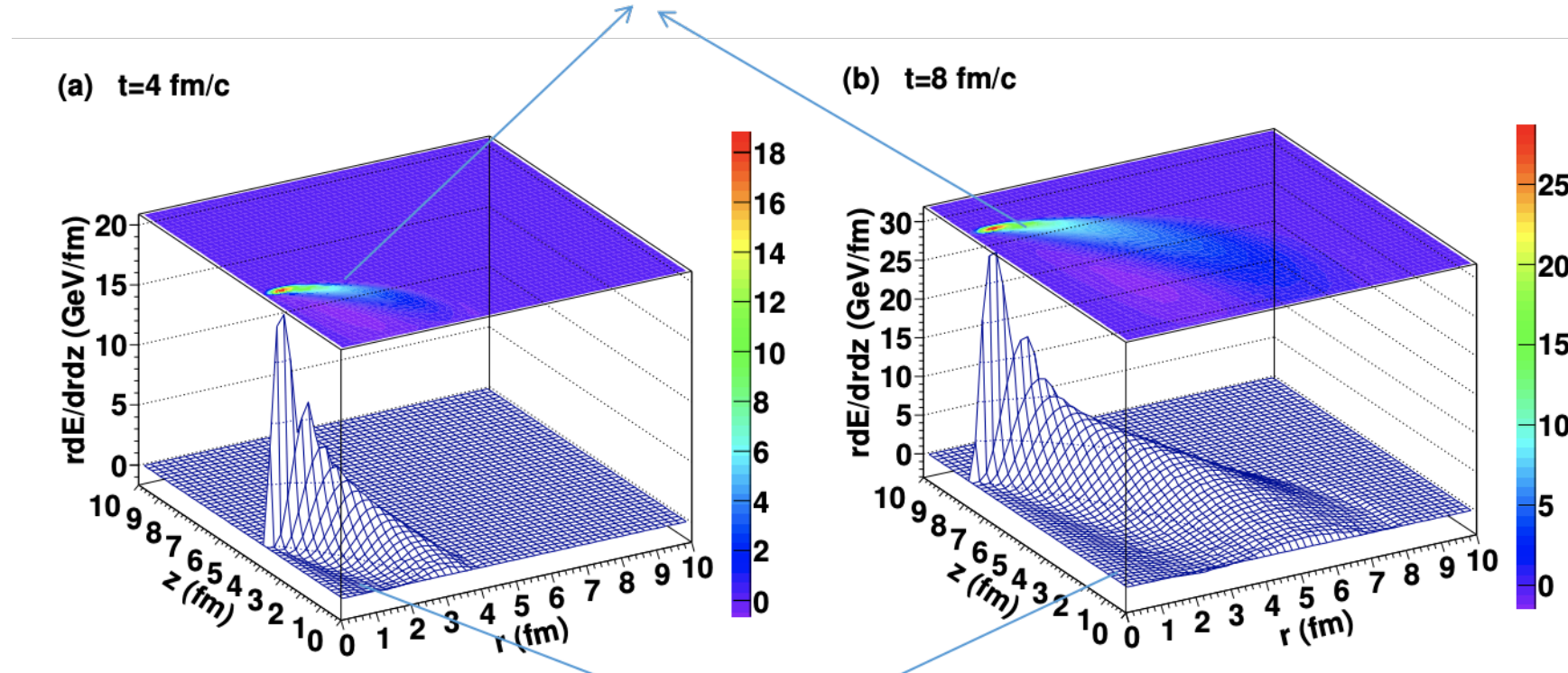
- Jet shower partons
- Thermal recoil partons
- Radiated gluons
- Negative partons



LBT: Pure pQCD description of parton transport

# LBT: Jet-induced medium response

Shock wave: propagation of recoil particles



He, Luo, Wang & Zhu, Phys.Rev.C 91 (2015) 054908

Diffusion wake: propagation of negative partons

**Medium response: recoil and negative particles**

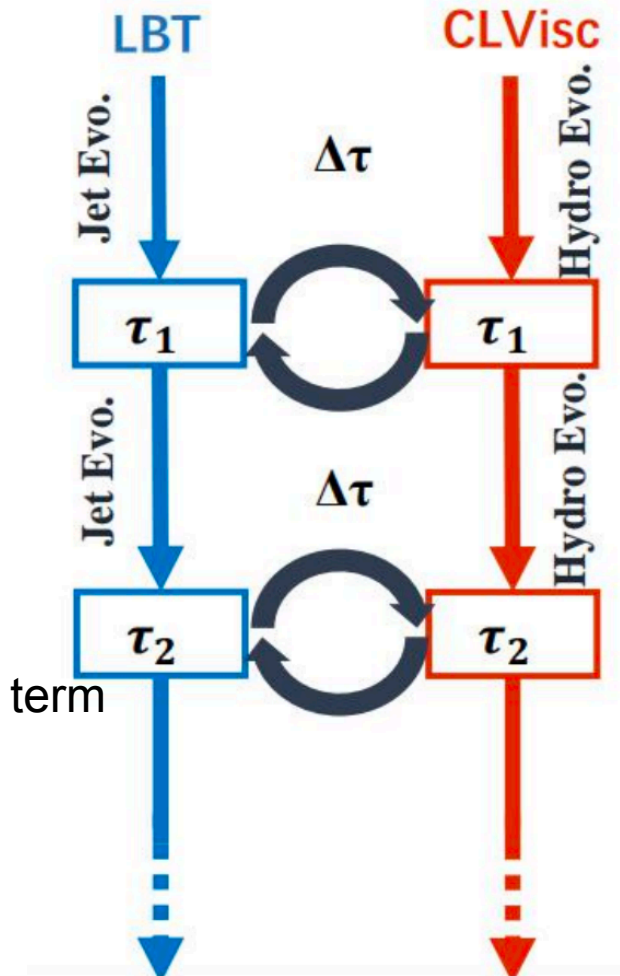
# CoLBT-hydro model

1. LBT for energetic partons(jet shower and recoil)
2. Hydrodynamic model for bulk and soft particles: CLVisc
3. Sorting partons according to a cut-off parameter  $p_{cut}^0$  (2 GeV)  
 Hard partons:  $p \partial f(p) = -C(p) \quad (p \cdot u > p_{cut}^0)$   
 Soft and negative partons:

$$j^\nu = \sum_i p_i^\nu \delta^{(4)}(x - x_i) \theta(p_{cut}^0 - p \cdot u)$$

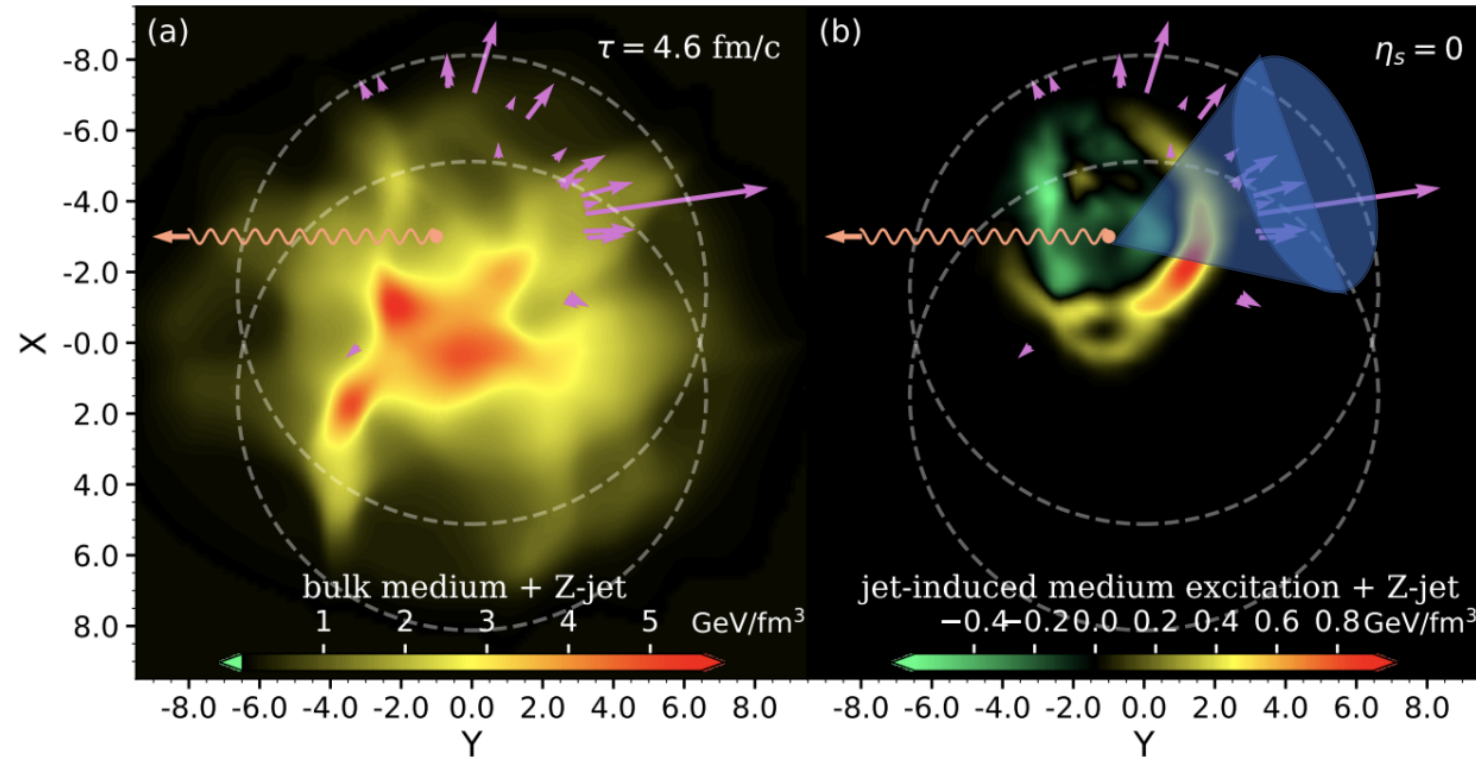
4. Updating medium information by solving the hydrodynamics equation with source term  

$$\partial_\mu T^{\mu\nu} = j^\nu$$
5. The final hadron spectra:  
 (1) hadronization of hard partons within a parton recombination model  
 (2) jet-induced hydro response via Cooper-Frye freeze-out



Wei Chen

# CoLBT-hydro: Jet-induced medium response

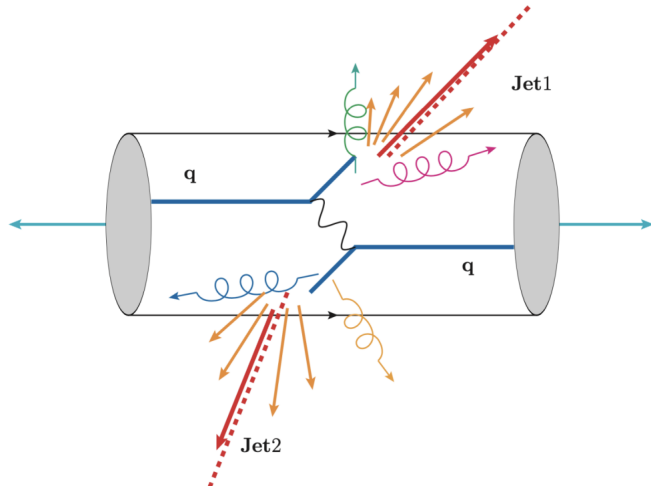


We run model twice with and without jet to subtract hydro background

Chen, Yang, He, Ke, Pang and Wang, PRL 127 (2021) 8, 082301

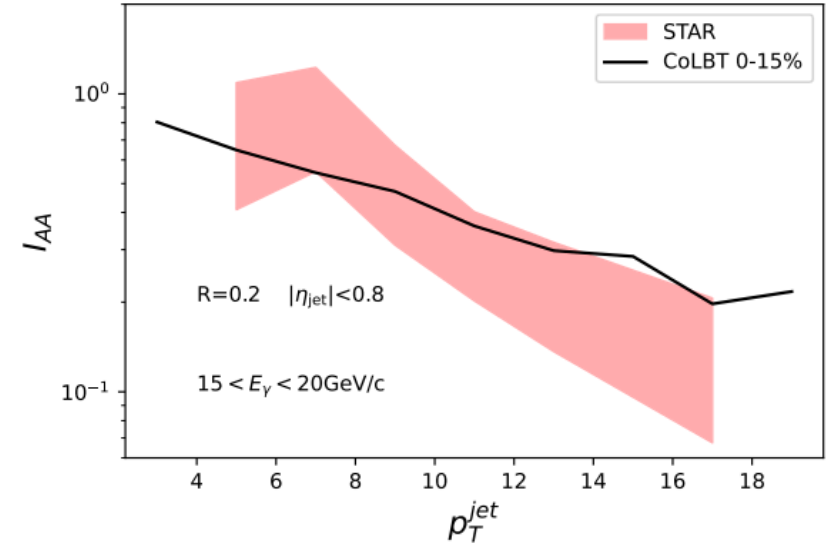
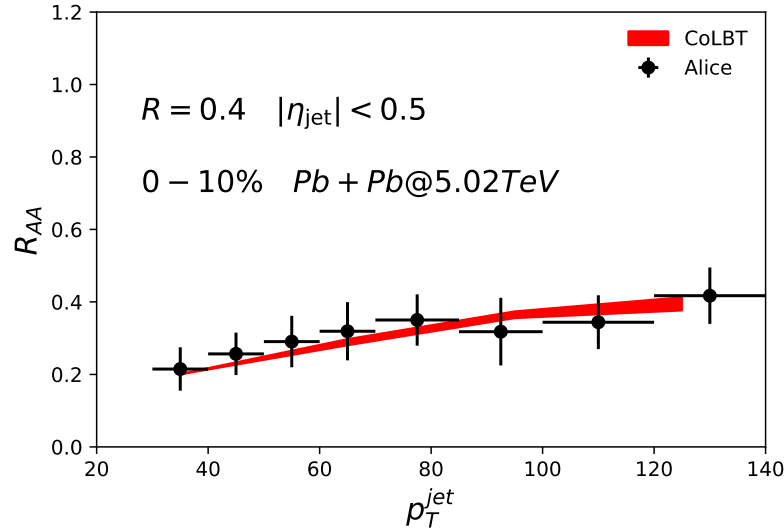
The Mach-cone-like jet-induced medium response including the diffusion wake is clearly seen in the right panel.

# Jet energy loss simulated by CoLBT-hydro



## Jet-medium interaction:

the suppression of large transverse momentum jets and hadrons in AA collisions



$$R_{AA} = \frac{d\sigma_{\text{jet}}^{AA}}{\langle N_{bin} \rangle d\sigma_{\text{jet}}^{pp}}$$

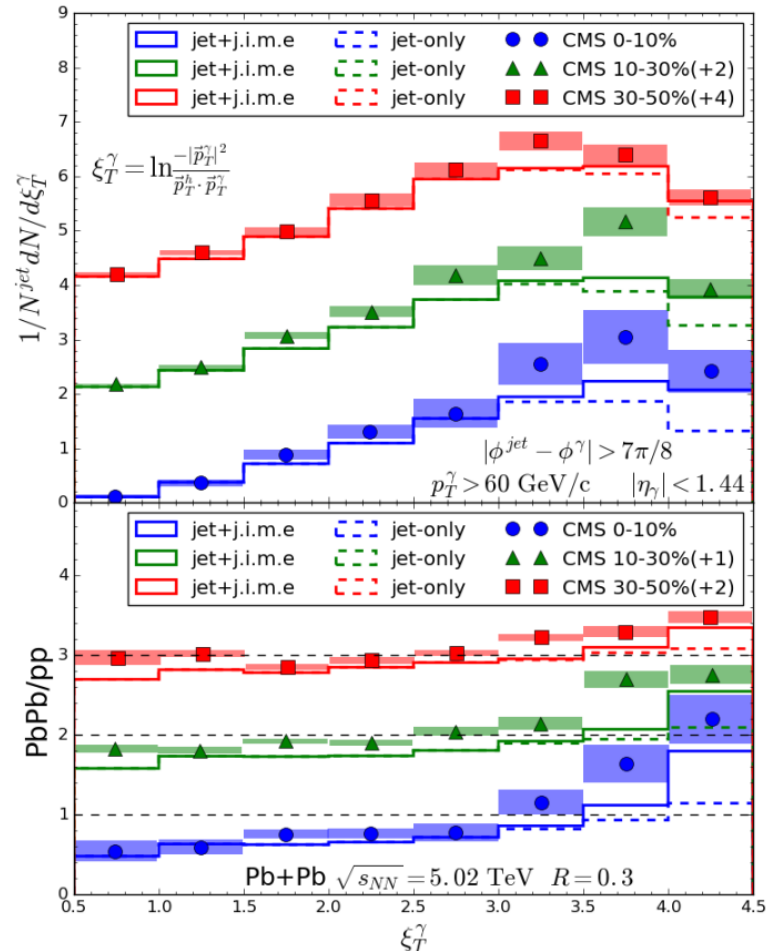
$$I_{AA} = \frac{dN_{AA}}{dp_T^{\text{jet}}} / \frac{dN_{pp}}{dp_T^{\text{jet}}}$$

**CoLBT-hydro model is an effective model to describe jet energy loss in QGP (RHIC, LHC, single jet and trigger-jet)**



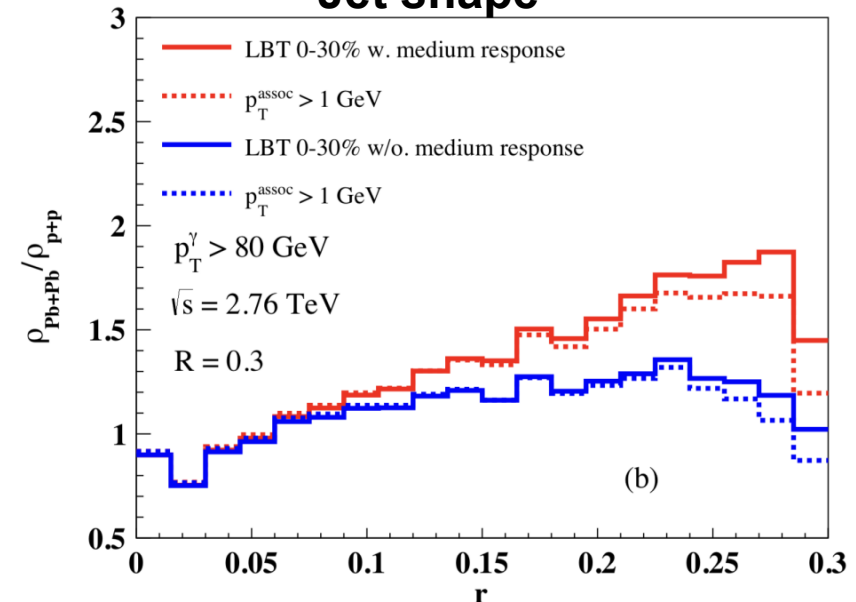
# Studying of jet-induced medium response

## Jet fragmentation function



Chen, Cao, Luo, Pang & Wang PLB 810 (2020) 135783

## Jet shape



Luo, Cao, He & Wang, PLB 782 (2018) 707-716

$$\rho(r) = \frac{1}{\delta r} \frac{\sum_{\text{jets}} \sum_{r < r_{\text{trk}} < r + \delta r} (p_T^{\text{trk}} / p_T^{\text{jet}})}{\sum_{\text{jets}} \sum_{r_{\text{trk}} < R} (p_T^{\text{trk}} / p_T^{\text{jet}})}$$

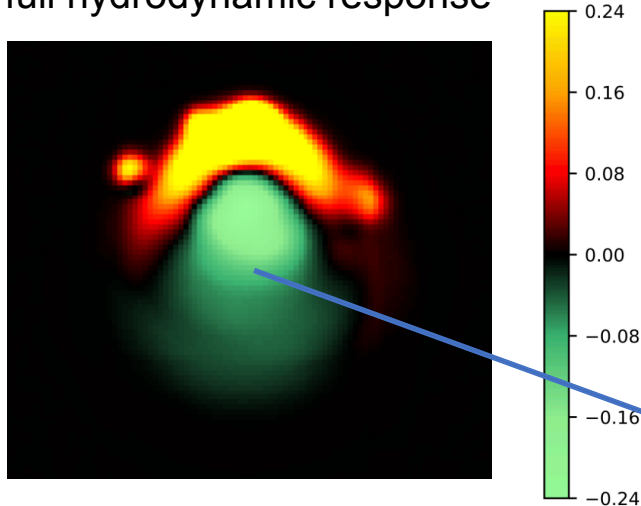
**Jet-induced medium response lead to enhancement of soft hadrons at large angle inside jet**

# Medium response and soft gluon radiation

**Medium response** leads to enhancement of soft hadrons along the direction of jet.

**Medium-induced gluon radiation** has the similar effect.

full hydrodynamic response



Zhong Y, et al. arXiv:2206.02393

Medium response:  $\delta f(p) \sim e^{-p \cdot u/T}$

Medium induced gluon radiations:  $\omega \approx \lambda^2 \hat{q}/2 \sim T$

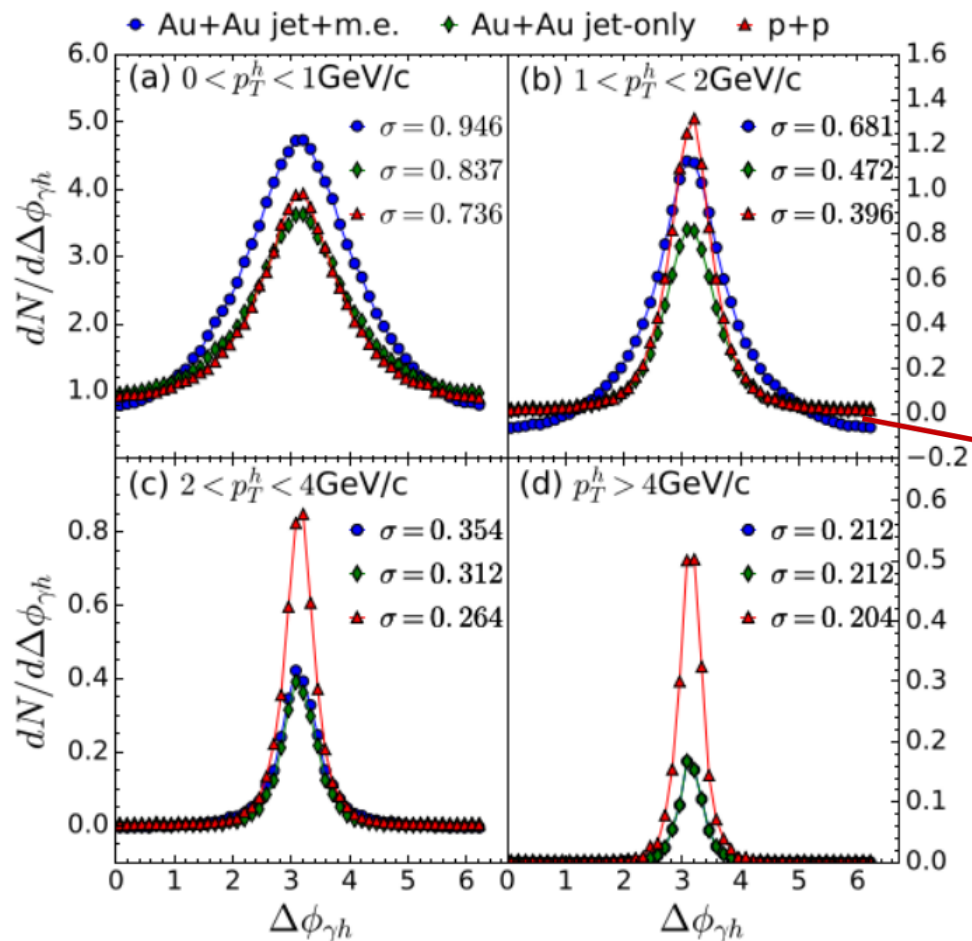
Formation time:  $\tau_f = \frac{2\omega}{k_T^2}$   $k_T^2 = \hat{q}\tau_f$   $\tau_f \approx \sqrt{2\omega/\hat{q}}$

Mean-free-path limits the formation time:  $\tau_f \leq \lambda \sim 1/T$   $\hat{q} \sim T^3$

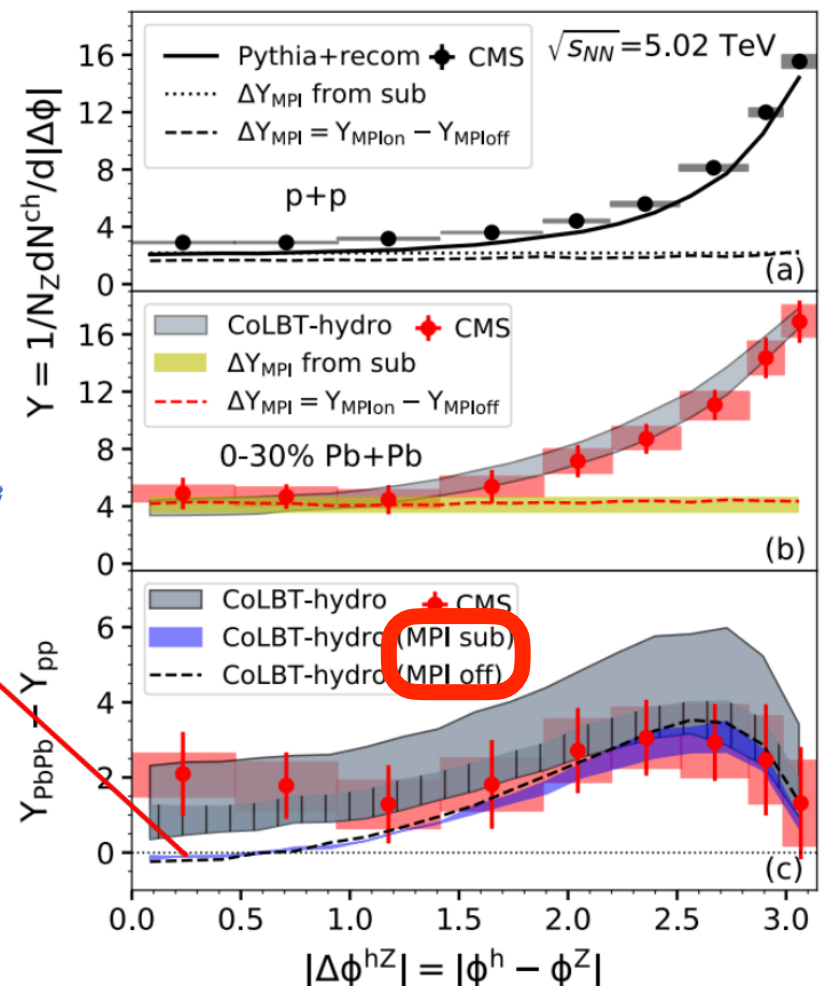
**Diffusion wake:** an unambiguous part of the jet-induced medium response. It can lead to depletion of soft hadrons in the opposite direction of the jet.



# Azimuthal distribution of soft hadrons



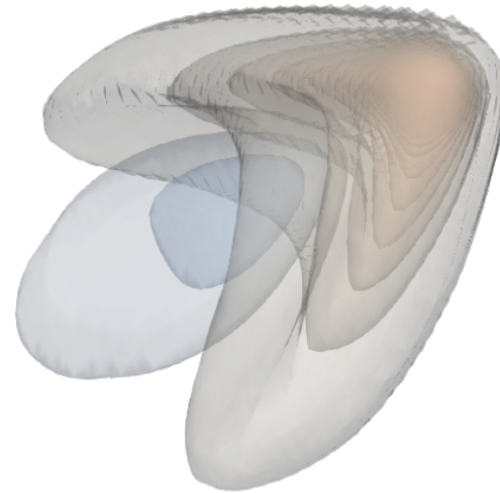
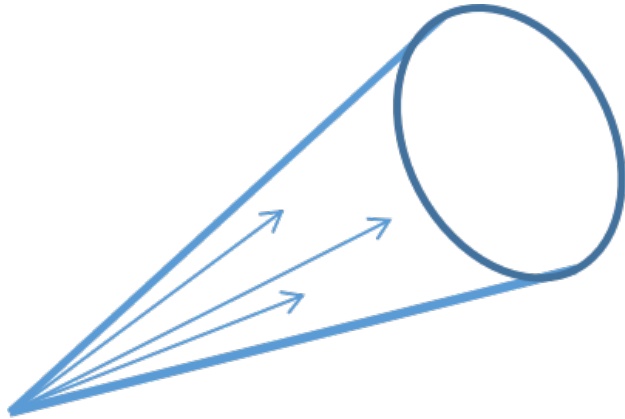
**Diffusion wake**



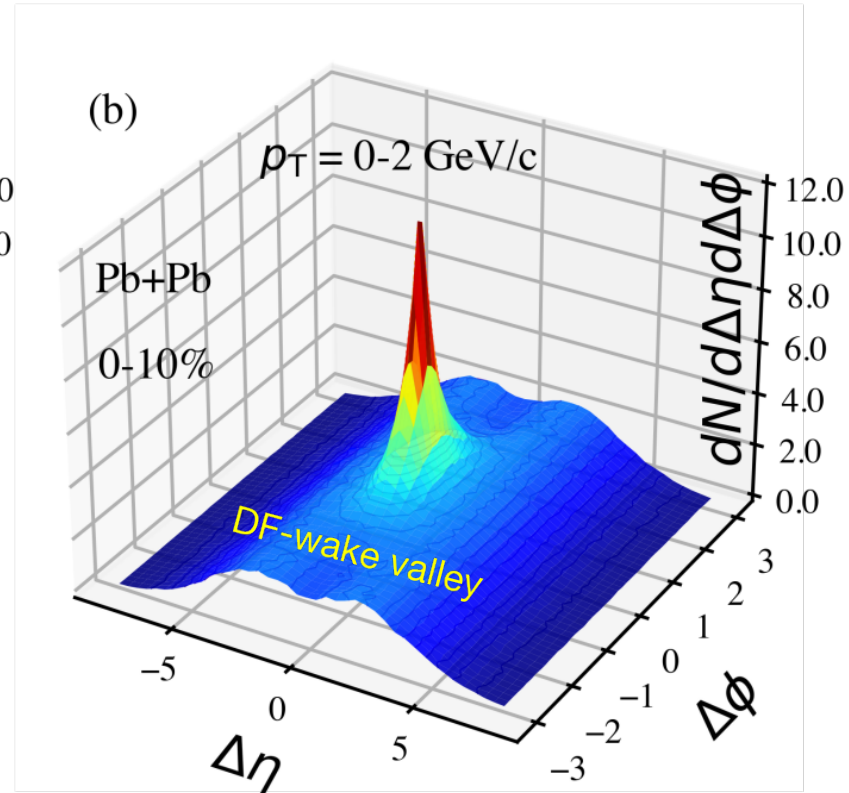
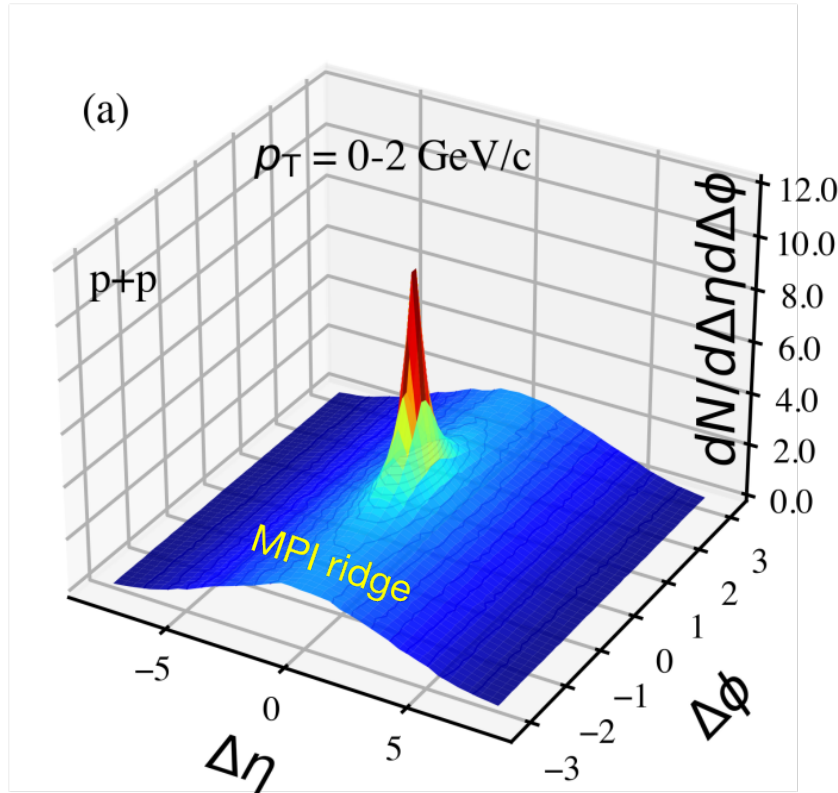
# Motivation to study 3D structure of DW

**(1) The previous studies of diffusion wake focus on the azimuthal angle.**

**(2) The jet is a 3D observable, thus the diffusion wake should also have a 3D structure.**



# 3D structure of diffusion wake



$$\Delta\eta = \eta_h - \eta_{\text{jet}}$$

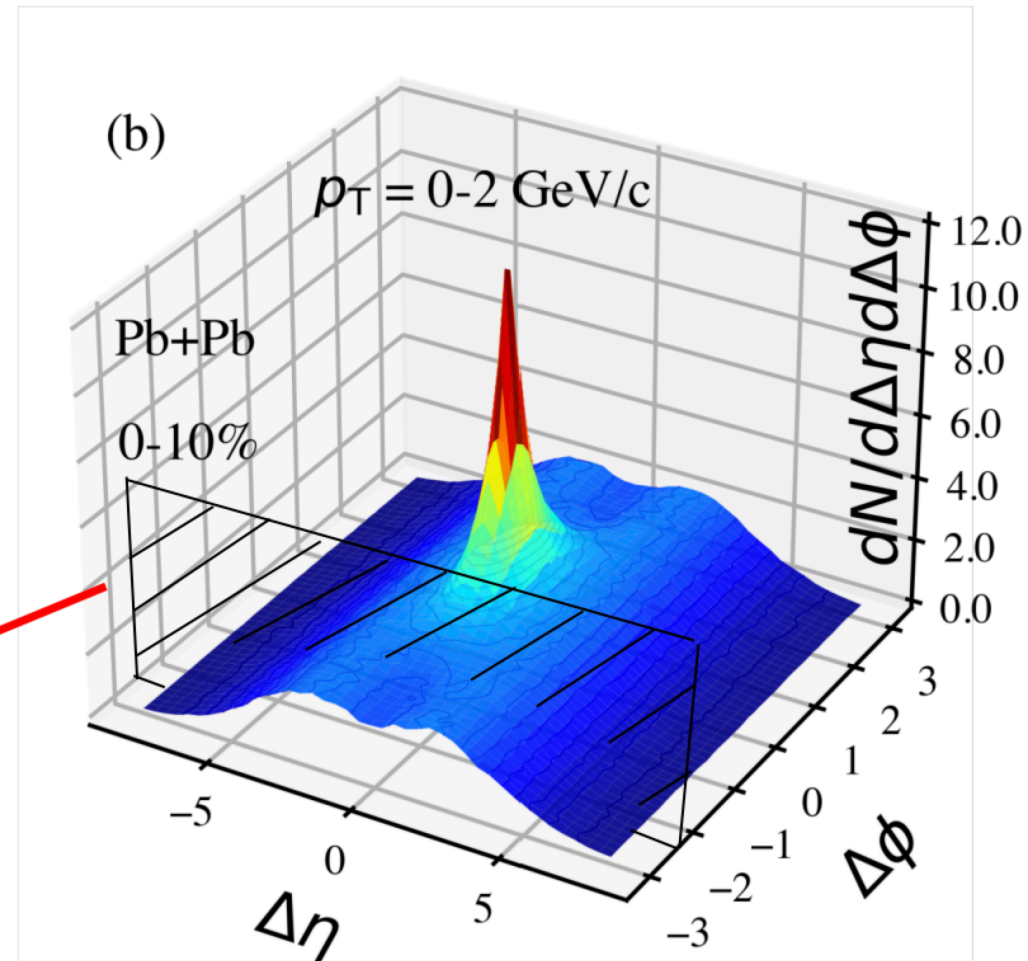
$$\Delta\phi = \phi_h - \phi_{\text{jet}}$$

**Diffusion wake valley(DF-wake valley):** a valley is formed on top of the MPI ridge due to the depletion of soft hadrons by jet-induced diffusion wake.

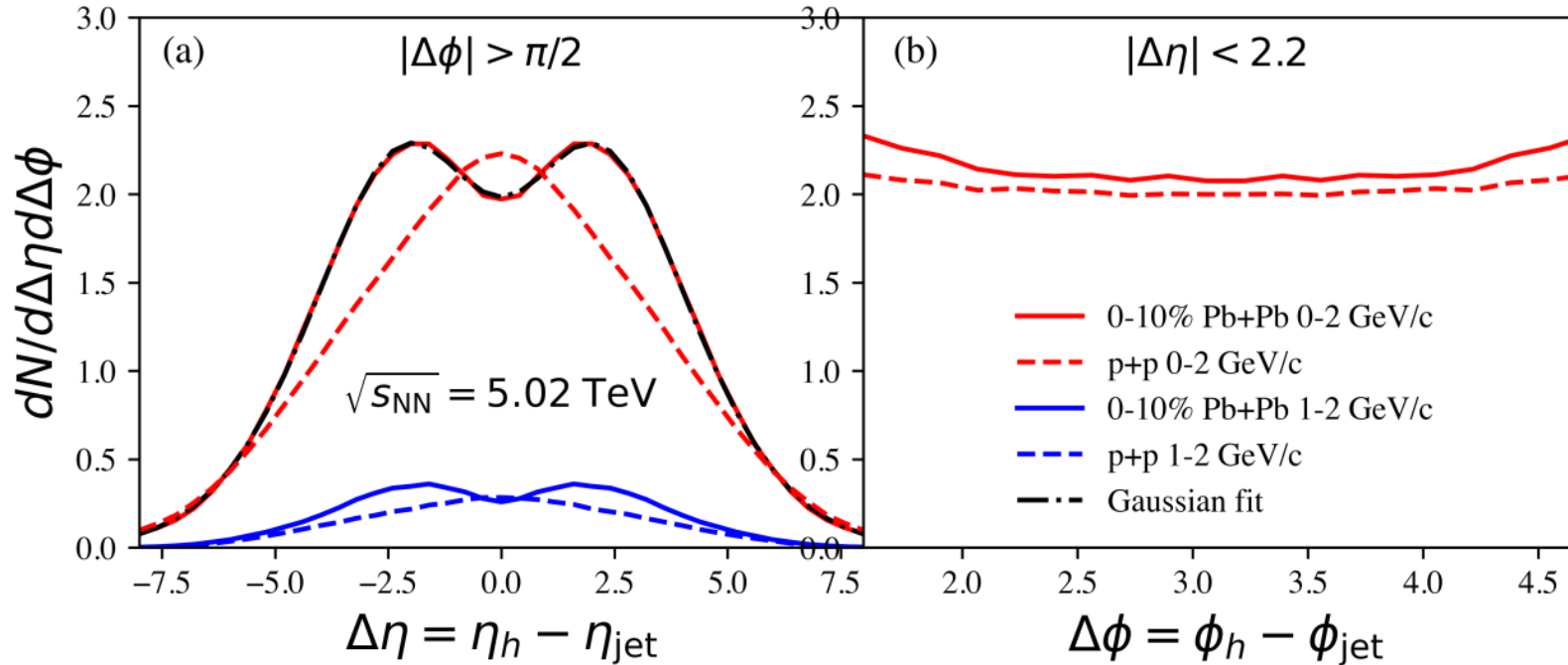
# 3D structure of diffusion wake

**Gamma side**

$$|\Delta\phi| > \pi/2$$



# 3D structure of diffusion wake



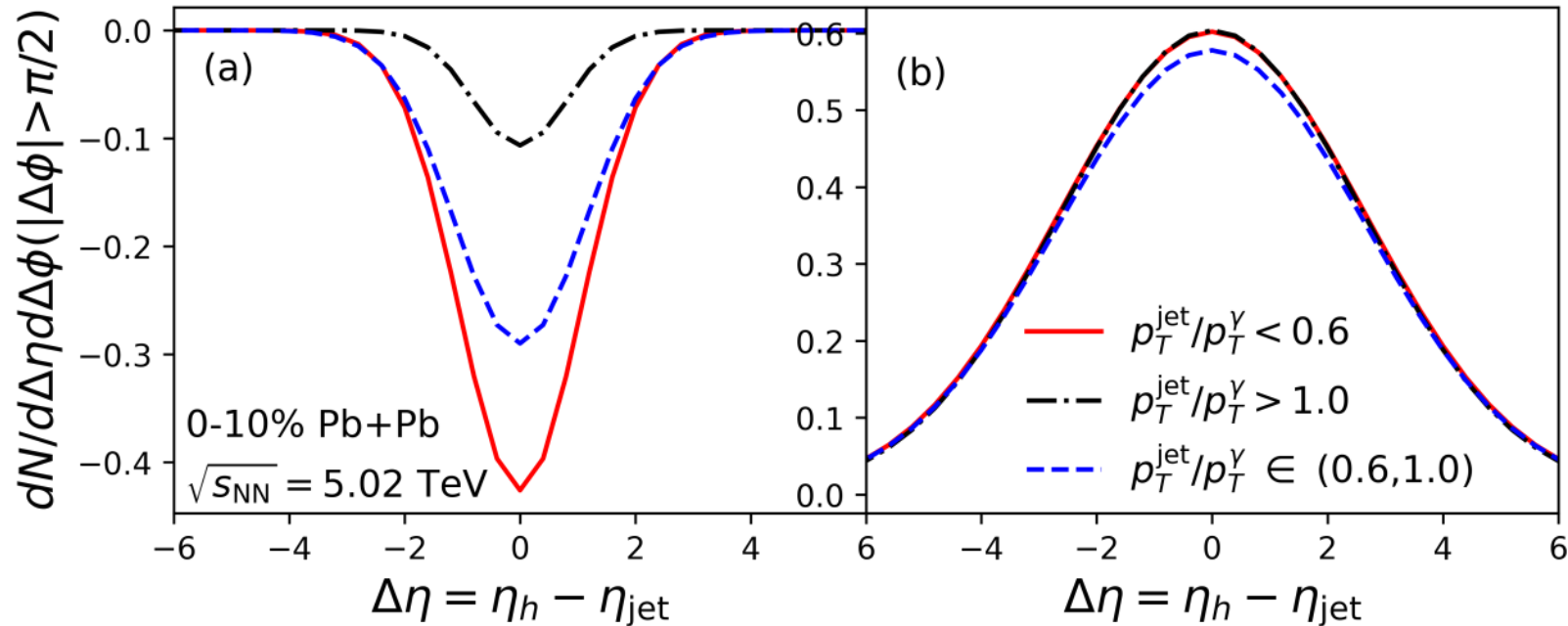
**2-Gaussian fitting:**

$$F(\Delta\eta) = \int_{\eta_{j1}}^{\eta_{j2}} d\eta_j F_3(\eta_j) (F_2(\Delta\eta, \eta_j) + F_1(\Delta\eta))$$

$$F_1(\Delta\eta) = A_1 e^{-\Delta\eta^2/\sigma_1^2}$$

$$F_2(\Delta\eta, \eta_j) = A_2 e^{-(\Delta\eta + \eta_j)^2/\sigma_2^2}$$

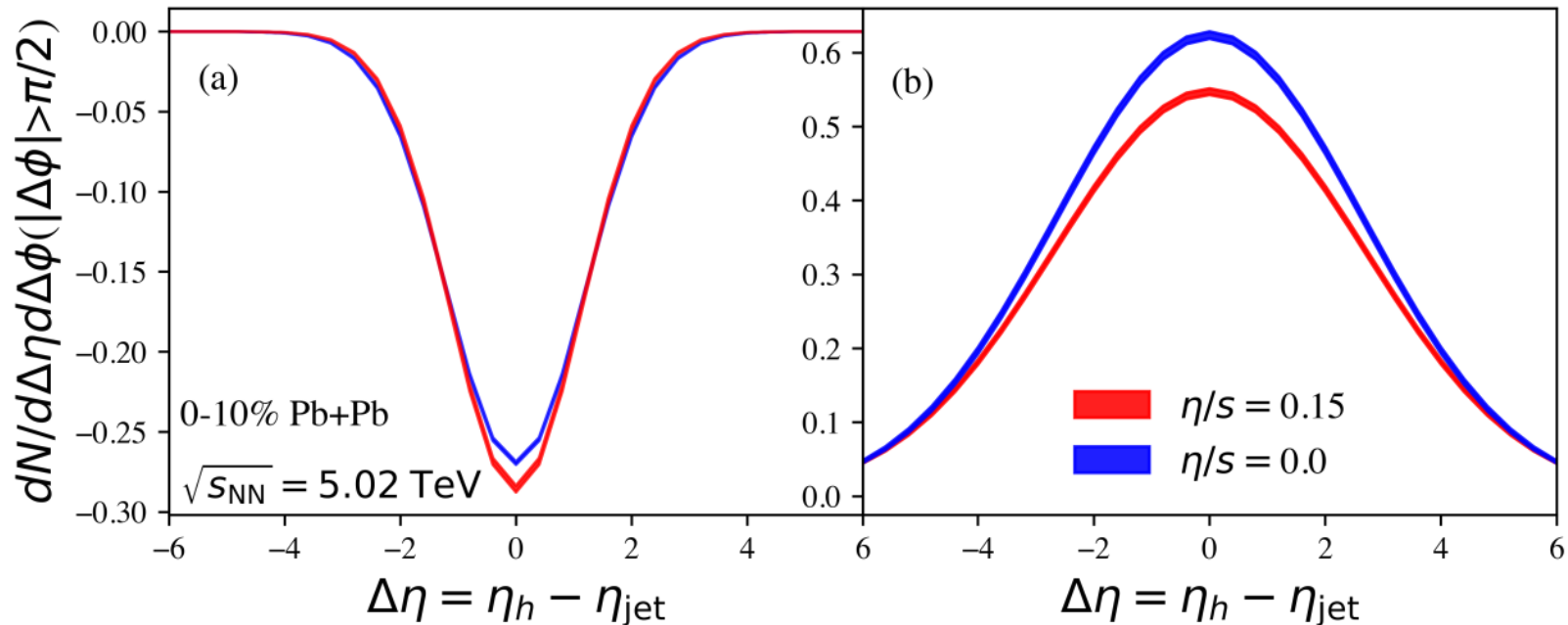
# Sensitivity to jet energy loss



Longer propagation length and larger jet energy loss leads to deeper DF-W valley.

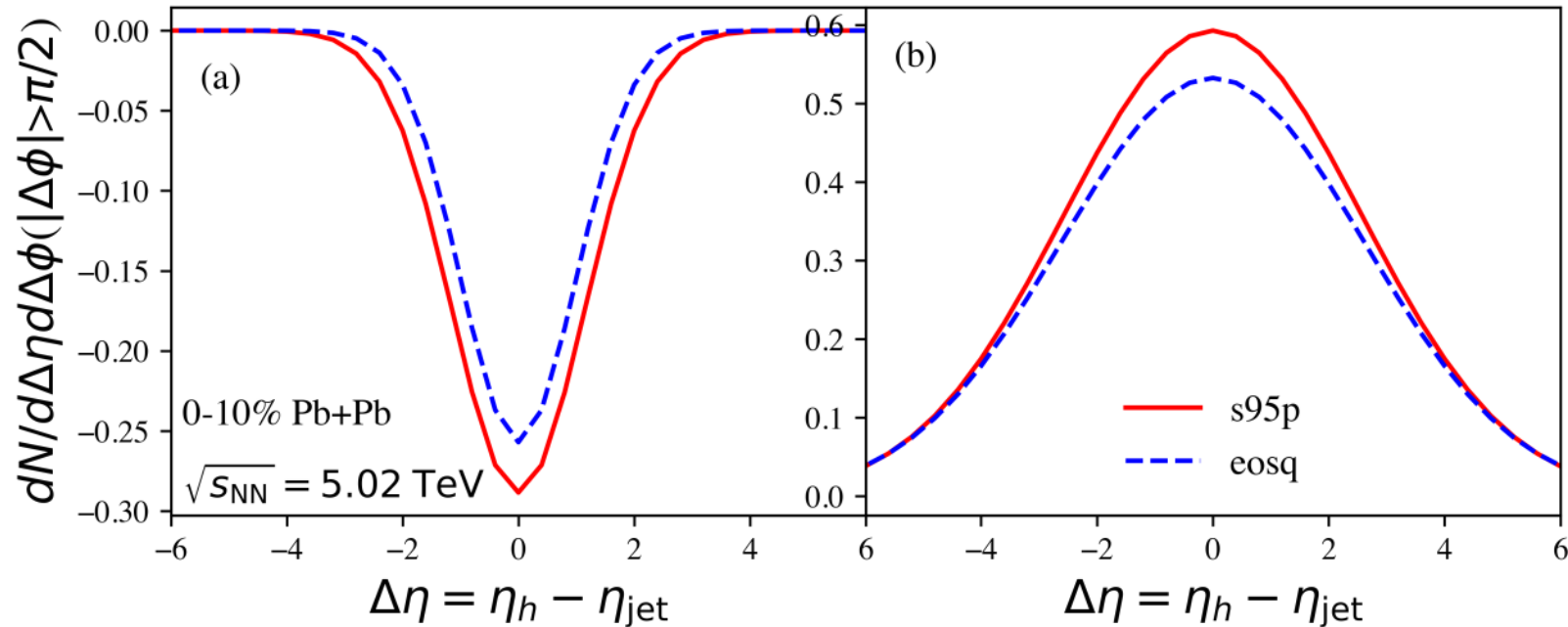
The MPI ridge has a very weak and non-monotonic dependence on  $x_{j\gamma}$  due to the non-monotonic dependence of the propagation length on  $x_{j\gamma}$  for mini-jets from MPI.

# Sensitivity to shear viscosity

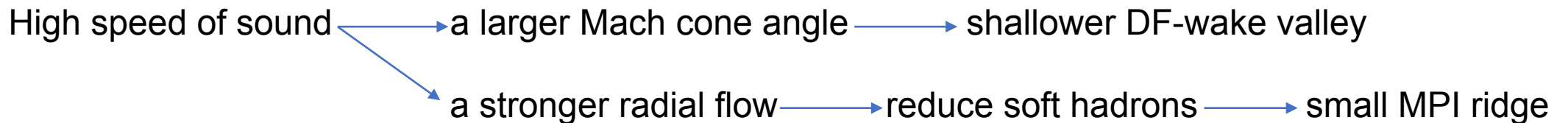


Competition between **increased radial flow** and **negative longitudinal pressure** in the shear correction of the energy momentum tensor leads to a slightly smaller MPI ridge and a deeper DF-wake valley in viscous hydro than in an ideal hydro.

# Sensitivity to equation of state



The effective speed of sound is higher in eosq than s95.



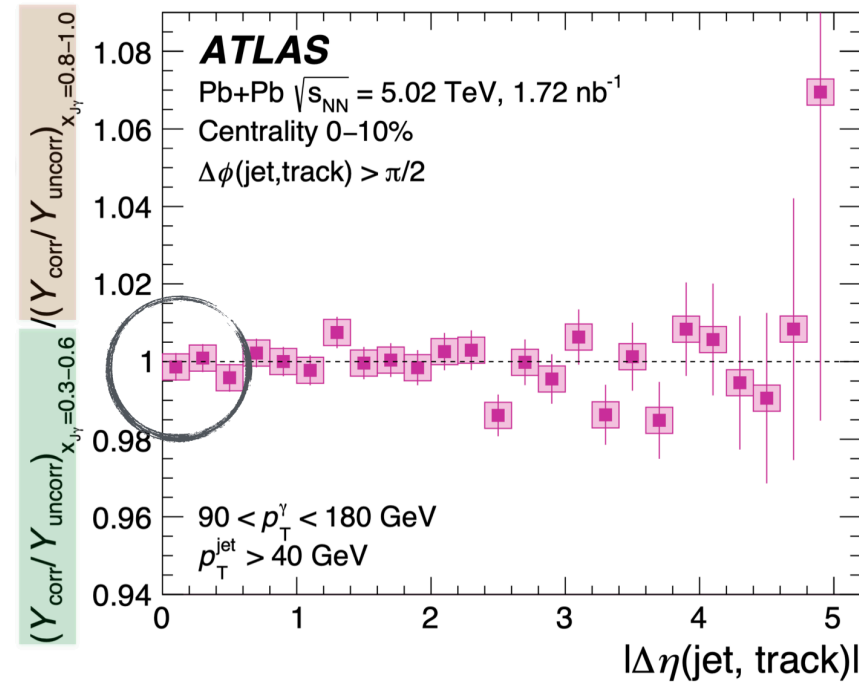
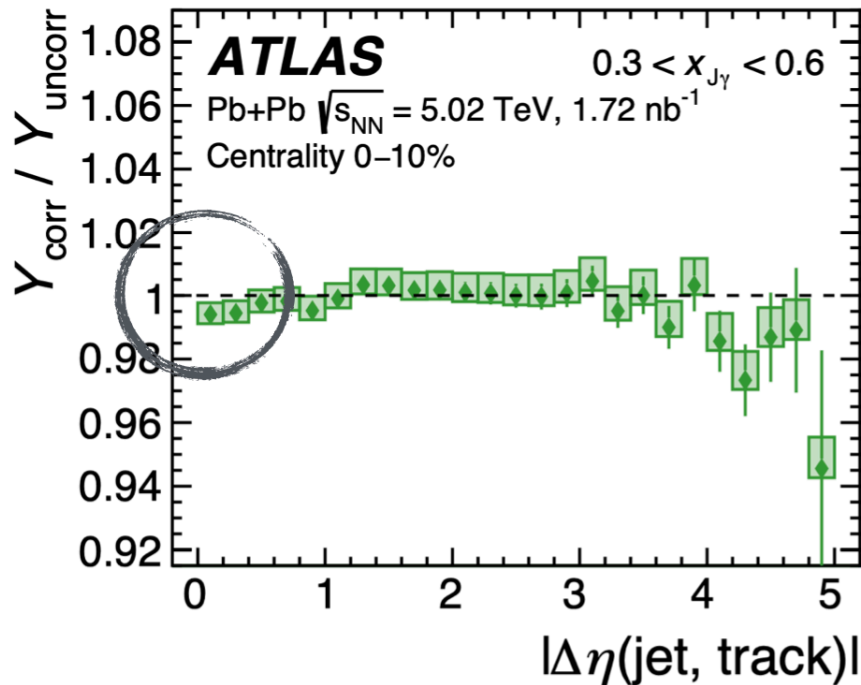


# Measurement of 3D structure at LHC

**ATLAS**

Jet-hadron correlation

arXiv:2408.08599



$$Y_{corr} = \frac{1}{N_{\gamma\text{-jet}}} \frac{d^2 N^{\text{jet-track}}}{d\Delta\eta d\Delta\phi}$$

The ratio of diffusion wake compared to hydro background is a few. No significant  $x_{j\gamma}$ -dependence of the diffusion wake is found.

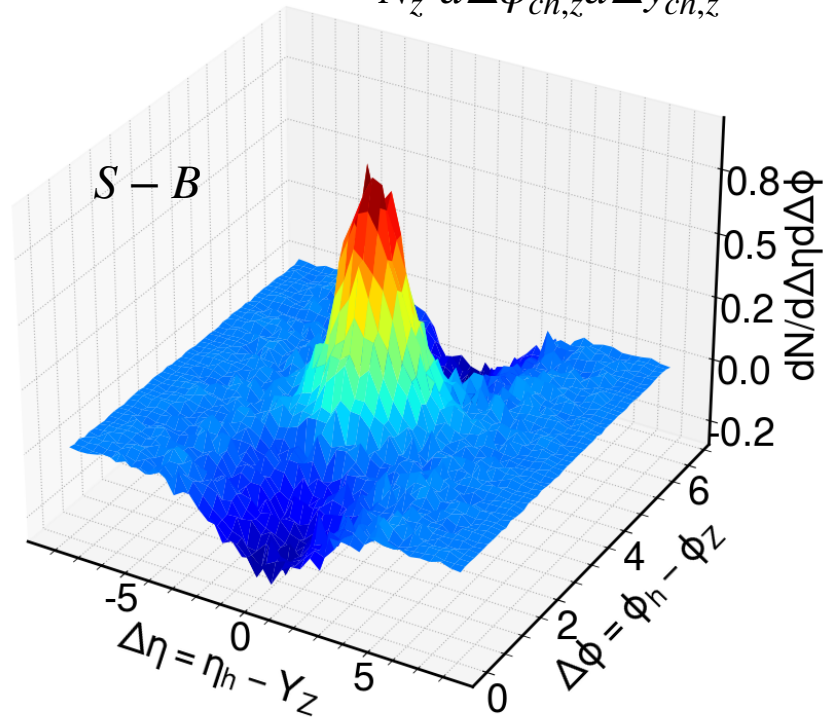
# Measurement of 3D structure at LHC

**CMS**

Z-hadron correlation

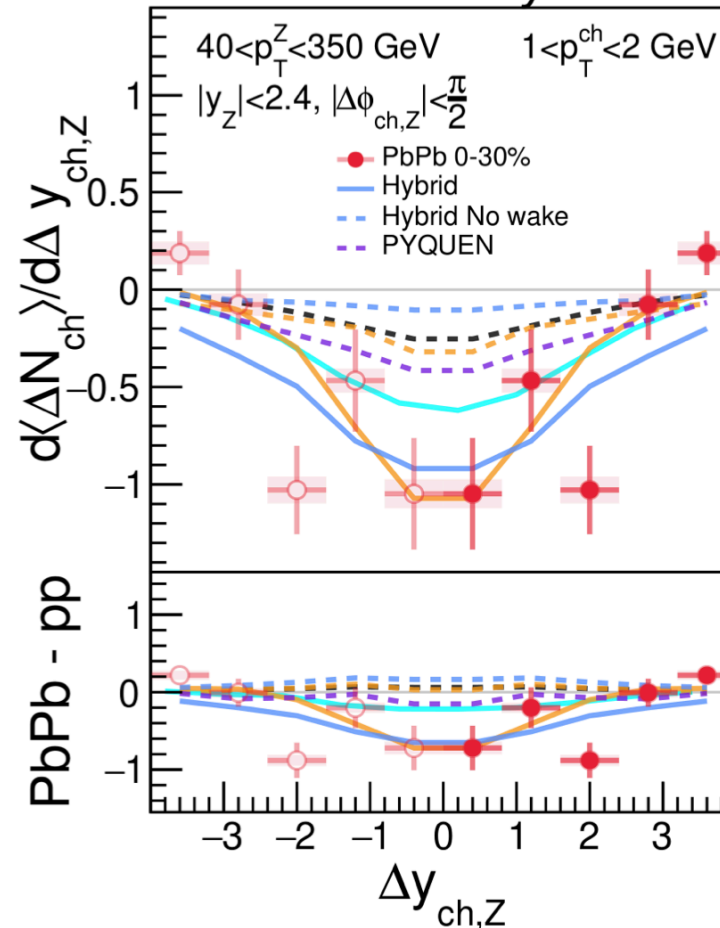
$$S(\Delta\phi_{ch,z}, \Delta y_{ch,z}) = \frac{1}{N_z} \frac{d^2 N^{same}}{d\Delta\phi_{ch,z} d\Delta y_{ch,z}}$$

$$B(\Delta\phi_{ch,z}, \Delta y_{ch,z}) = \frac{1}{N_z} \frac{d^2 N^{mix}}{d\Delta\phi_{ch,z} d\Delta y_{ch,z}}$$



**CMS Preliminary**

Yen-Jie's talk at HP2024



**The first direct evidence  
of medium response in  
QGP!!!**

# Energy-energy correlators

Energy-energy correlators (EEC) have recently emerged as excellent jet substructure observables for studying the space-time structure of the jet shower.

$$\langle \varepsilon^{(n)}(\vec{n}_1) \dots \varepsilon^{(n)}(\vec{n}_k) \rangle$$

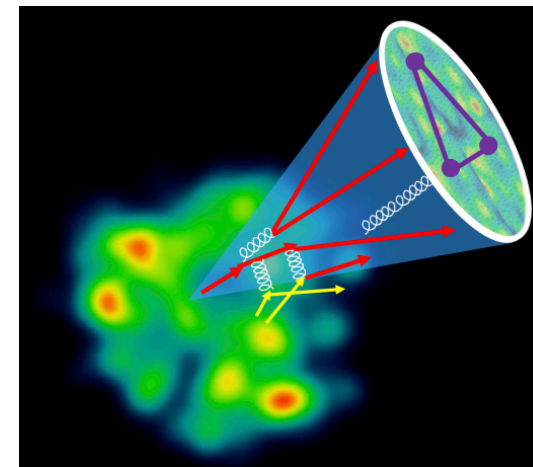
$\varepsilon^{(n)}(\vec{n}_1)$  measures the asymptotic energy flux in the direction  $\vec{n}_1$

$$\varepsilon^{(n)}(\vec{n}_1) = \lim_{r \rightarrow \infty} \int dt r^2 n_1^i T_{0i}(t, r\vec{n}_1)$$

The  $n$ -th weighted normalized two-point correlation:

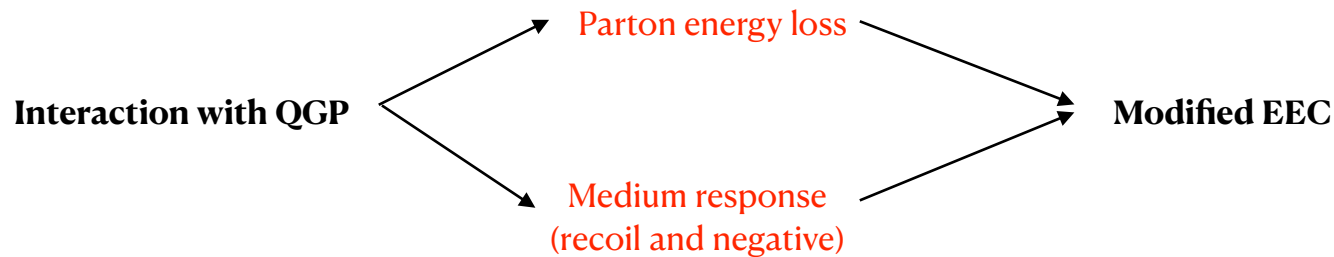
$$\frac{\langle \varepsilon^{(n)}(\vec{n}_1) \varepsilon^{(n)}(\vec{n}_2) \rangle}{Q^{2n}} = \frac{1}{\sigma} \sum_{ij} \frac{d\sigma_{ij}}{d\vec{n}_i d\vec{n}_j} \frac{E_i^n E_j^n}{Q^{2n}} \delta^{(2)}(\vec{n}_i - \vec{n}_1) \delta^{(2)}(\vec{n}_j - \vec{n}_2) \quad n = 1$$

$$\frac{d\Sigma^{(n)}}{d\theta} = \int dn_{1,2} \frac{\langle \varepsilon^{(n)}(\vec{n}_1) \varepsilon^{(n)}(\vec{n}_2) \rangle}{Q^{2n}} \delta(\vec{n}_1 \cdot \vec{n}_2 - \cos\theta) \quad \cos\theta = \vec{n}_1 \cdot \vec{n}_2$$



# EECs from a quark going through the QGP brick

## EECs from a **quark** going through the QGP brick

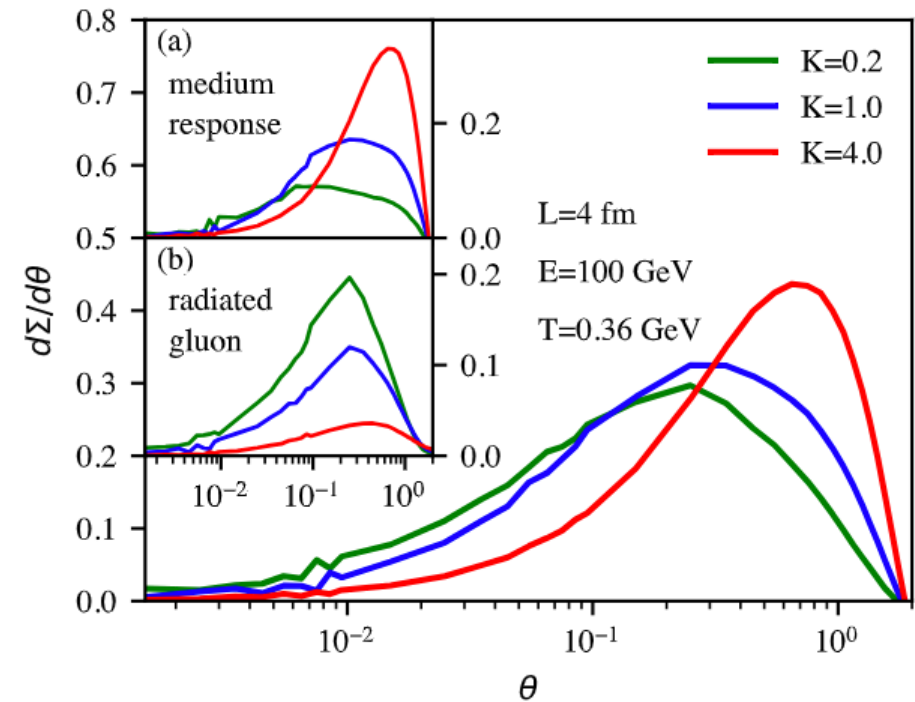


LBT model has a cut-off in the transverse momentum transfer in terms of the Debye screening mass.

$$\mu_D^2 = \frac{3}{2} K g^2 T^2 \quad K = 1(\text{default}), 0.2, 4.0$$

It determines the typical momentum and angular scale of the in-medium interaction.

The EEC distribution from the medium response shifts to a larger angle with an enhanced magnitude if  $\mu_D$  increases. While the quark-radiated-gluon correlator decreases with  $\mu_D$  and peak shifts slightly to large angles.



Transverse momentum transfer:  $q_{\perp} \sim \mu_D$

Energy transfer to the medium:  $\delta E \sim \mu_D^2/T$

# EECs from a parton shower going through the QGP brick

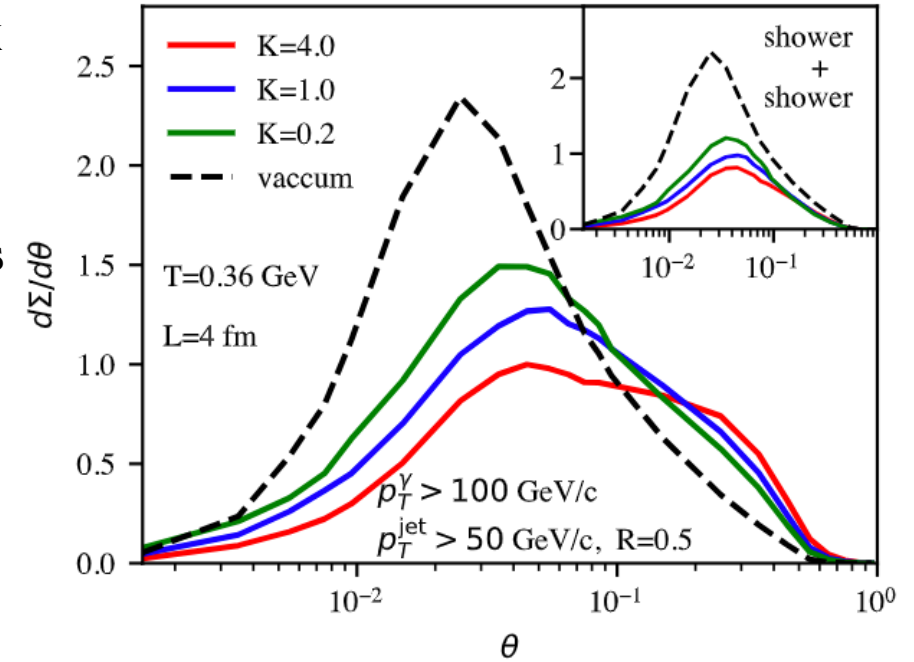
EECs from a **parton shower** going through the QGP brick

$$R = 0.5 \quad p_T^\gamma \geq 100 \text{ GeV}/c \quad p_T^{\text{jet}} \geq 50 \text{ GeV}/c$$

Jet  $\longrightarrow$  Parton showers  $\longrightarrow$  Multiple elastic and inelastic scatterings



Transverse momentum broadening and energy loss



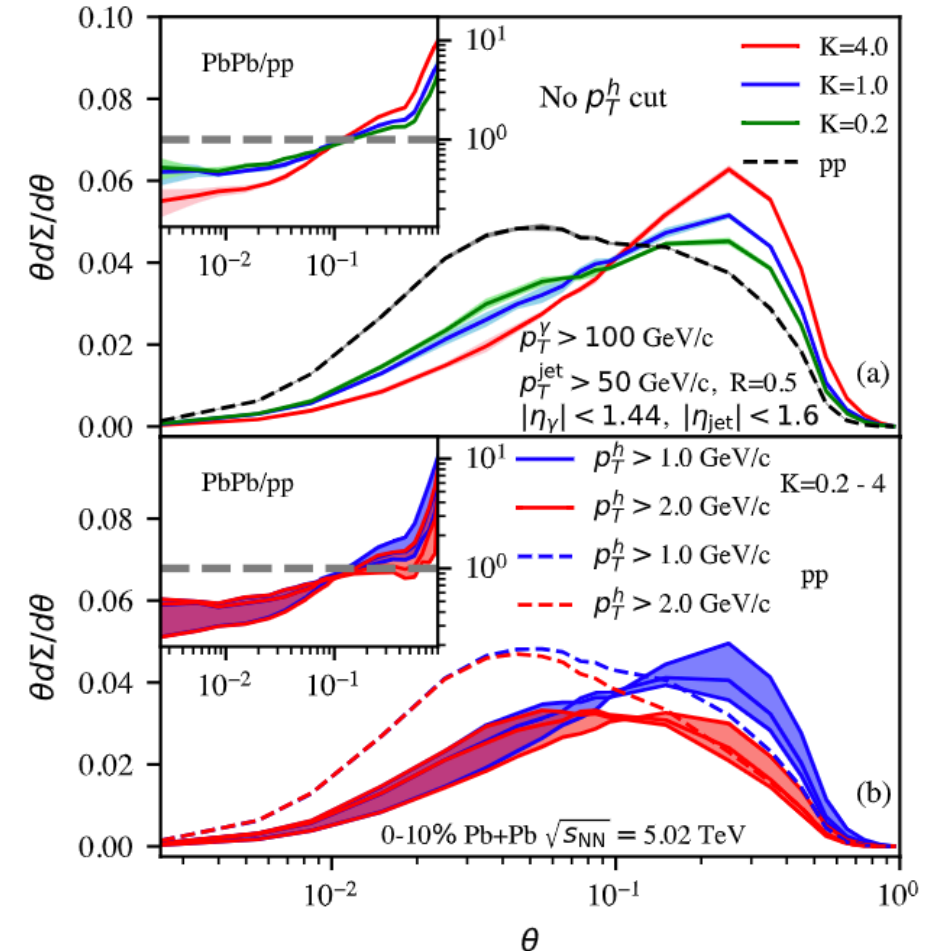
EEC distributions from correlation between shower partons **suppressed at both small and large angles** relative to the vacuum EEC (dashed).

The total correlator of all partons (shower, medium-response and radiated gluons) inside the modified jet **enhances at large angles due to correlations involving medium response or/and radiated gluons.**

# EECs in Pb+Pb collisions at LHC

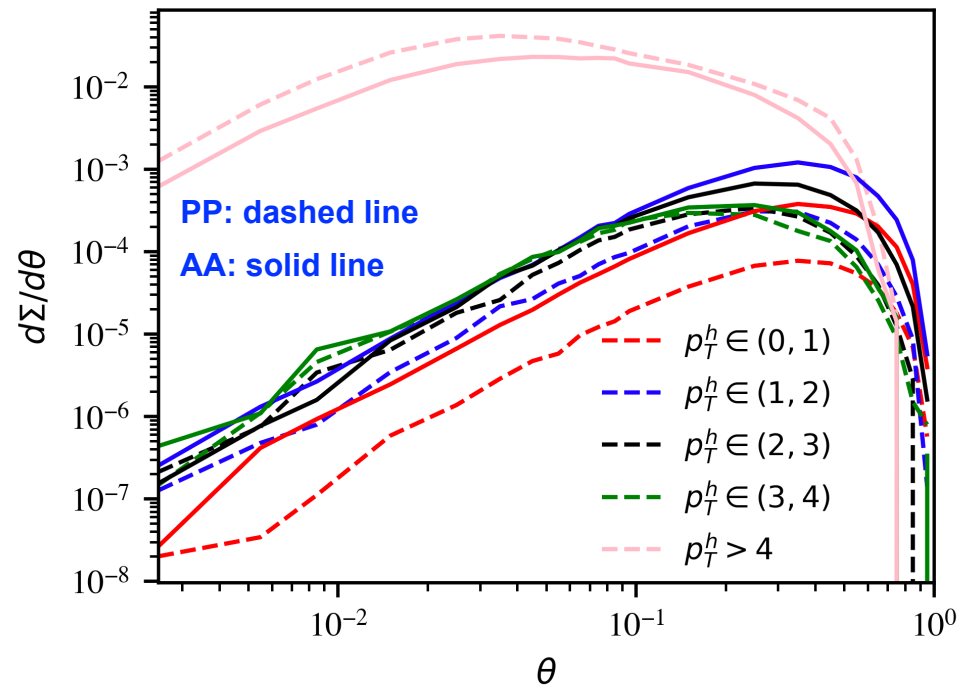
## EECs of $\gamma$ -jets in Heavy-Ion Collisions.

1. Similar to the case of a QGP brick, the EEC's in Pb+Pb collisions are suppressed at small angles due to energy loss, while they are enhanced at large angles.
2. This modification is sensitive to the Debye mass,  $\mu_D$ , which determines the angular scales of each jet-medium scattering and characterizes the structure of the QGP medium in the CoLBT simulations.
3. The enhancement at large angles is reduced but still survives if a  $p_T > 1 \text{ GeV}/c$  cut is imposed on the final hadrons for the purpose of reducing the background in experimental analysis. If  $p_T > 2 \text{ GeV}/c$  cut is used, the medium enhancement at large angles is mostly gone except for the case of  $K = 4$ .



# Transverse momentum dependence of EECs

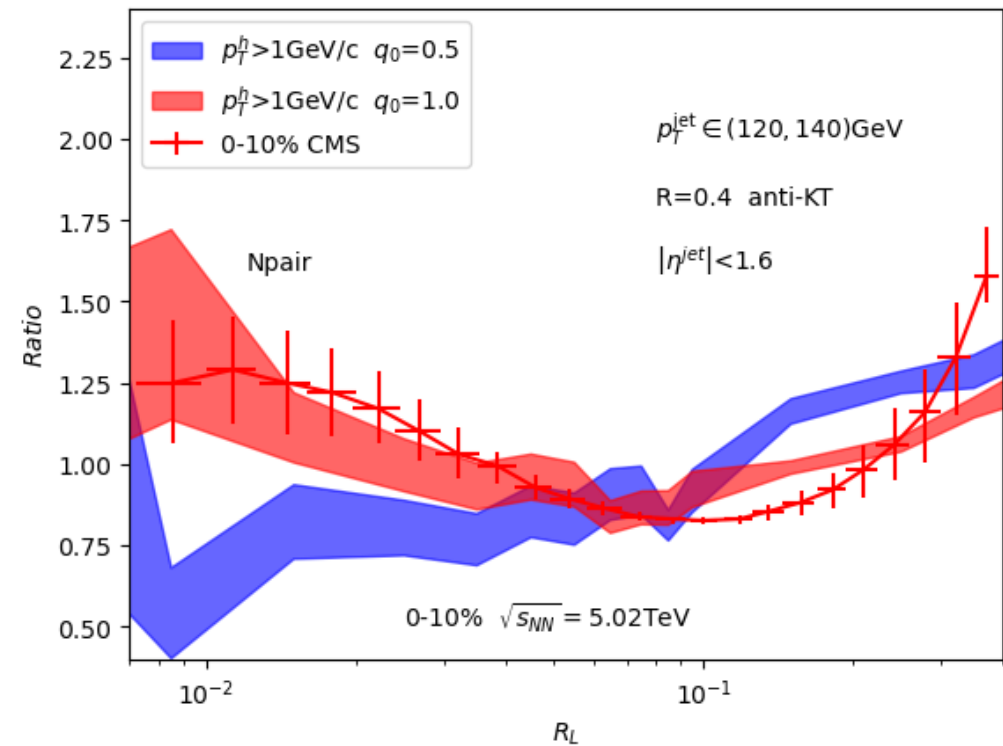
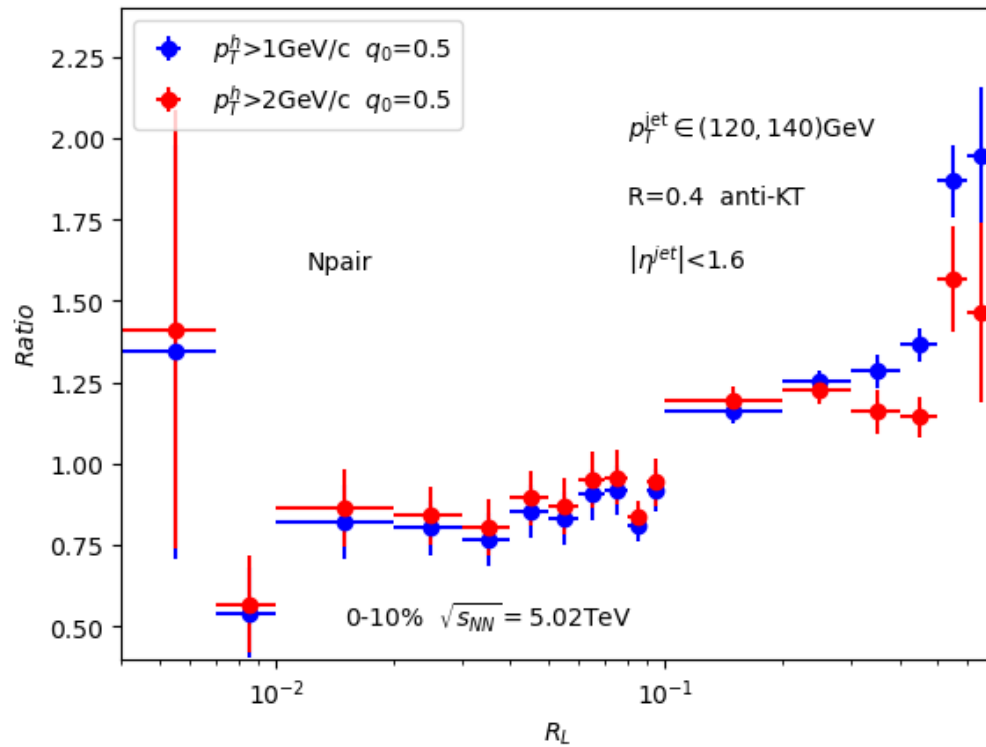
EECs of  $\gamma$ -jets in Heavy-Ion Collisions.



For high  $p_T$  hadrons, PP result is greater than AA result. But, AA results becomes greater when you decrease  $p_T^h$ .

# EECs of single inclusive jets

## EECs of single inclusive jets in Heavy-Ion Collisions.



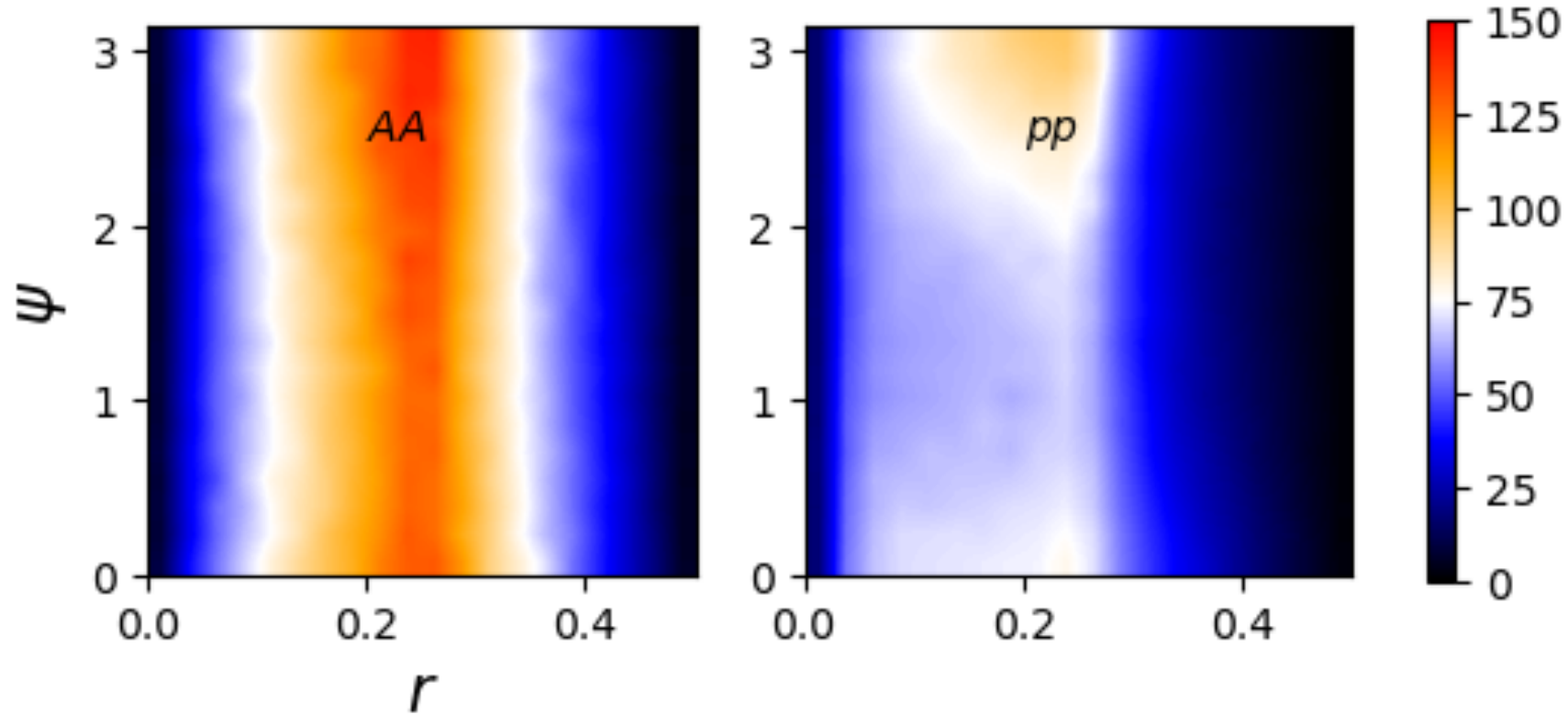
$q_0$  is the minimum value for vacuum splitting in Pythia model



# Summary

1. Studying of jet-induced medium response can help us understand QGP properties.
2. CoLBT-hydro model is an effective model to study jet quenching and jet-induced medium response.
3. Jet-induced diffusion wake is an unambiguous part of medium response which leads to depletion of hadrons in the opposite direction of jet.
4. The double-peak structure of jet-hadron correlation in rapidity direction is a unique signal of diffusion wake, and it's sensitive to energy loss, shear viscosity and EoS.
5. Jet-medium interaction will modify the EECs inside jet ( $\gamma$ -jet and single inclusive jet). And this modification of EECs shows a clear sensitivity to Debye screening mass. The coming experimental result can help us constrain this value of models.

# Outlook



**Coming 3-particles correlation, sensitive to medium properties...**

**Thank You**

# EECs in vacuum and medium-induced emissions

We focus on the normalized two-point energy correlates

For a quark with energy  $E$  and initial virtuality  $Q=E$ , the vacuum splitting  $q \rightarrow q + g$  at small angles and leading order (LO) in pQCD leads to the angular distribution of the energy correlators.

$$\frac{d\Sigma_q^{\text{vac}}}{d\theta} \approx \frac{\alpha_s}{2\pi} C_F \int_0^1 dz z(1-z) P_{qg}(z) \int_{\mu^2}^{Q^2} \frac{d\mathbf{k}_\perp^2}{\mathbf{k}_\perp^2} \delta\left(\theta - \frac{|\mathbf{k}_\perp|}{z(1-z)E}\right)$$

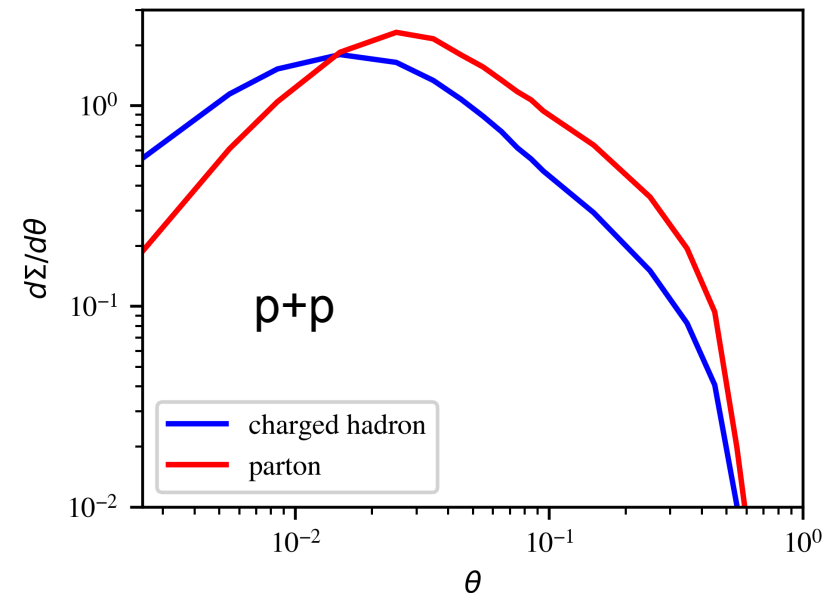
$$P_{qg}(z) = \frac{1 + (1-z)^2}{z} \quad \text{Splitting function}$$

$$\frac{d\Sigma_q^{\text{vac}}}{d\theta} \approx \frac{\alpha_s}{2\pi} \frac{C_F}{2\theta} \left(3 - \frac{2\mu}{E\theta}\right) \sqrt{1 - \frac{4\mu}{E\theta}}$$

$$\theta > 4\mu/E : \quad d\Sigma_q^{\text{vac}}/d\theta \sim 1/\theta$$

$\theta \rightarrow 4\mu/E$  : non-perturbative effects take over and its behavior will be influenced by hadronization processes.

$\mu \ll Q$  the collinear cut-off scale below which non-perturbative effects become dominant.



# EECs in vacuum and medium-induced emissions

The medium-induced gluon radiation is modeled by hist-twist approach

For a massless parton, the formation time for radiated gluon is

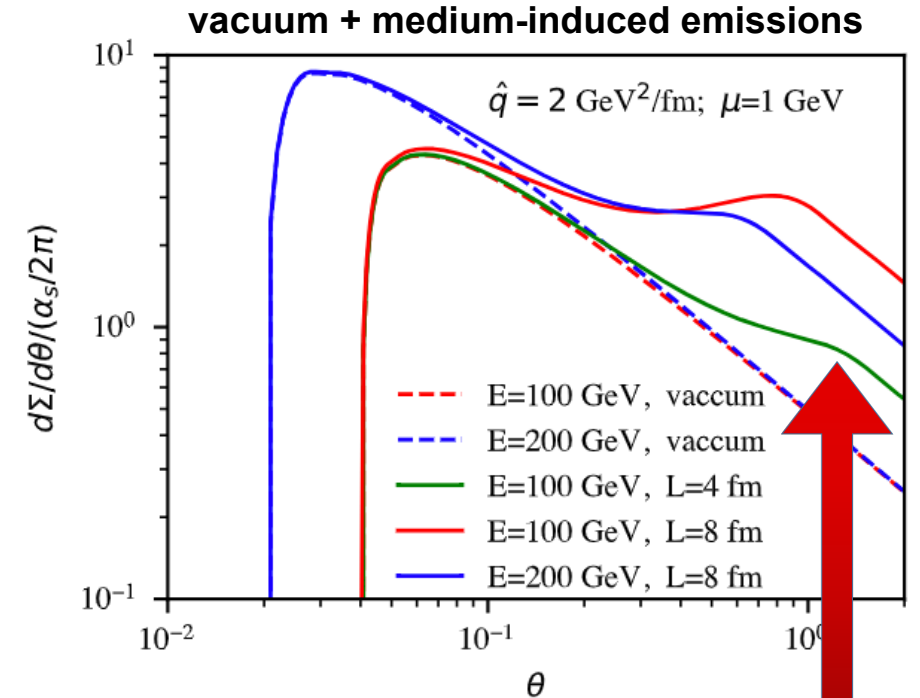
$$\theta_{12} = \frac{2\ell_{\perp}}{Ez(1-z)} \quad \tau_f = \frac{2Ez(1-z)}{\ell_{\perp}^2} = \frac{8}{\theta_{12}^2 z(1-z)E}$$

The corresponding angular contribution to EEC is,

$$\begin{aligned} \frac{d\Sigma_q^{\text{med}}}{d\theta} &= \frac{16\alpha_s C_A}{\pi E^2 \theta^3} \int dx dz \frac{\hat{q} P_{qg}(z)}{z(1-z)} \sin^2\left(\frac{x}{2\tau_f}\right) \\ &= \frac{L^{5/2} \hat{q}}{\pi \sqrt{E}} \frac{8\alpha_s C_A}{(\sqrt{EL}\theta)^3} \int dz \frac{P_{qg}(z)}{z(1-z)} \times \left[ 1 - \frac{\sin ELz(1-z)\theta^2/8}{ELz(1-z)\theta^2/8} \right] \end{aligned}$$

$$\theta < \sqrt{8\pi/EL} : \quad d\Sigma_q^{\text{med}}/d\theta \approx L^3 \hat{q} \alpha_s C_A \theta / (64\pi) \sim \theta$$

$$\theta > \sqrt{8\pi/EL} : \quad \frac{d\Sigma_q^{\text{med}}}{d\theta} \approx \frac{L^2 \hat{q}}{2E} \frac{\alpha_s C_A}{\theta} \left[ 1 + \mathcal{O}\left(\frac{1}{EL\theta^2}\right) \right] \sim 1/\theta$$

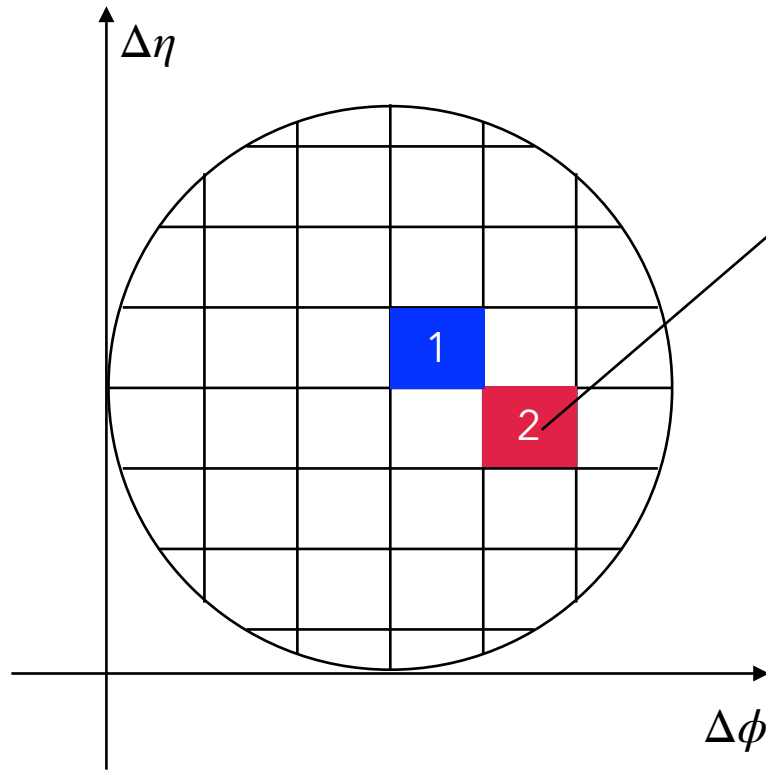


$$\theta_{\text{peak}}^{\text{med}} \sim \sqrt{8\pi/EL}$$

$$\Sigma_{\text{peak}}^{\text{med}} \sim \alpha_s \hat{q} L^{5/2} / \sqrt{E}$$

# How to calculate EEC in LBT model

How to deal with the negative parton in LBT model



The energy deposited in this cell equals  $E_{pos} - E_{neg}$

Therefore, energy correlation between different cells is

$$(E_{pos}^1 - E_{neg}^1)(E_{pos}^2 - E_{neg}^2) = E_{pos}^1 E_{pos}^2 - E_{pos}^1 E_{neg}^2 - E_{neg}^1 E_{pos}^2 + E_{neg}^1 E_{neg}^2$$

**Sign**

+

-

+

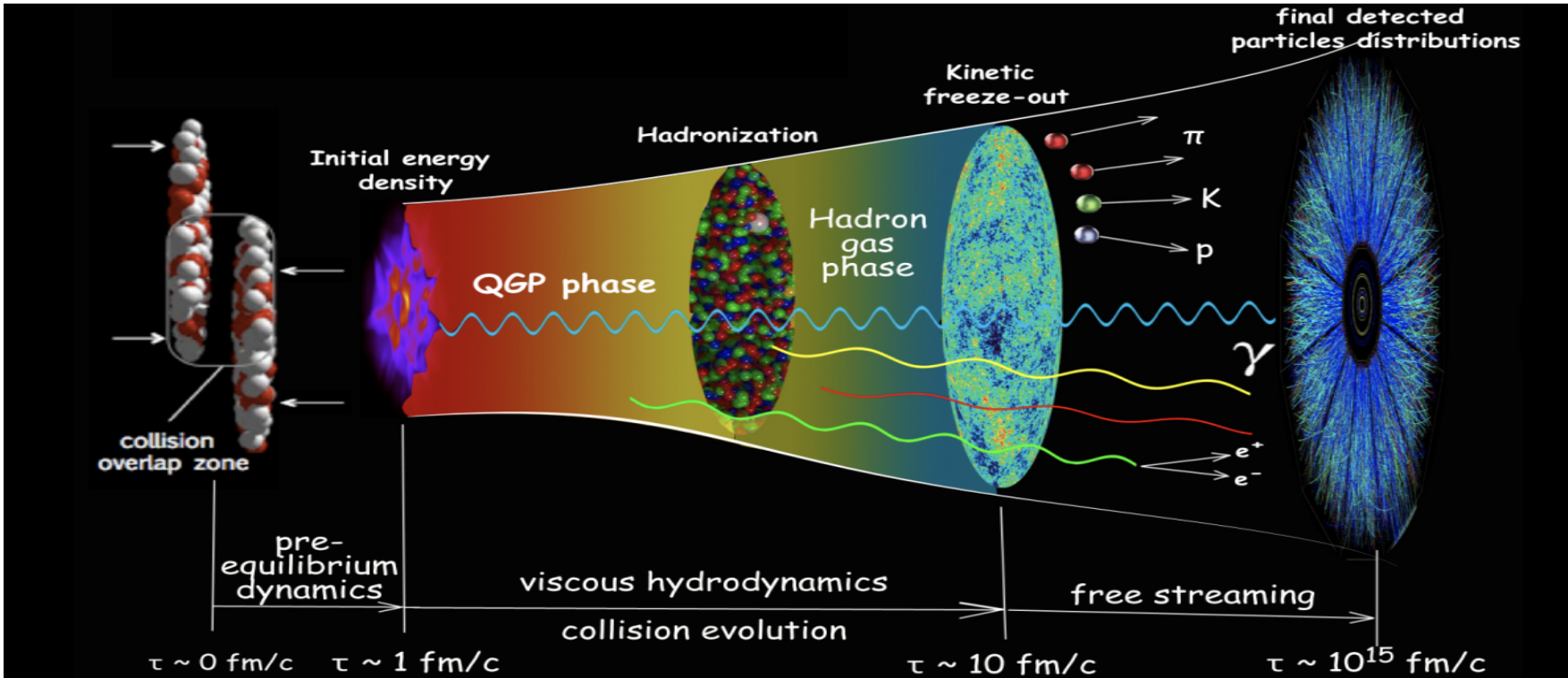
**Pair**

pos+pos

pos+neg

neg+neg

# Introduction



Chun Shen

Looking for and studying QGP are the main programs in high-energy heavy-ion collisions