



# Adjoint Approach to Optimization and Sensitivity Analysis of RF Sources\*

## <u>Alexander N. Vlasov</u><sup>1</sup>, Thomas M. Antonsen, Jr.<sup>2</sup>, David Chernin<sup>2</sup> Igor A. Chernyavskiy<sup>1</sup>,

<sup>1</sup> U.S. Naval Research Laboratory, Washington, DC USA

<sup>2</sup>Leidos Inc., Billerica, MA

\*This work is supported by the U.S. Office of Naval Research Distribution Statement A. Approved for public release. Distribution unlimited.

# Motivation: Optimization of VE RF Sources



- Design of advanced VE RF sources requires optimization with respect to many parameters
- > Accurate, geometry-driven modeling of VE devices requires substantial computational resources
- Current state-of-the-art of VE devices optimization is limited by the technical approaches and computational complexity

Key Elements of CAD Design and Optimization of RF Sources:

Accurate, Computationally Efficient Codes

U.S. NAVAL RESEARCH LABORATORY

Efficient Multi-Parameter Optimization Algorithms



# NRL/Leidos Design Codes for RF Sources

- 1. MAGY Gyro-devices
- 2. Christine 1D
  - Christine Helix
  - Christine CC-TWT
  - Christine-FW
  - Christine KL
- 3. Christine 2.5D Helix
- 4. TESLA 2.5D
  - TESLA (Klystron)
  - TESLA CC-TWT
  - TESLA FW
  - TESLA MBK

5. TESLA-Z & Christine-Z

6. 3D PIC Code Neptune



Computationally efficient, fast runs, optimization study, parametric sweep, search for optimal design, promising for stability study



Very accurate design tool, suitable to resolve 2D fields in beam tunnel and 3D electron motion, promising for stability study



External Information on Circuit Properties as Model Parameters or Z-matrix is Needed from 3D EM Codes:

- ANSYS HFSS
- Cadence ANALYST-MP

TESLA FW and TESLA –Z (Hybrid 2.5 D Codes)

# **Geometry Driven Z-Matrix Approach**

## **Network of generalized ports:** Case with Input/Output ports & multiple-gaps

U.S. NAVAL RESEARCH



Example based on structure with 29 generalized ports (27 gaps + 2 i/o ports)

$$\begin{pmatrix} V_{Pin} \\ V_{B1} \\ V_{B2} \\ \dots \\ V_{BN} \\ V_{Pout} \end{pmatrix} = \begin{pmatrix} Z_{PinPin} Z_{PinB_1} & Z_{PinB_2} & \dots & Z_{PinB_N} & Z_{PinPout} \\ Z_{B_1Pin} & Z_{B_1B_1} & Z_{B_1B_2} & \dots & Z_{B_NB_N} & Z_{B_1Pout} \\ Z_{B_2Pin} & Z_{B_2B_1} & Z_{B_2B_2} & \dots & Z_{B_2B_N} & Z_{B_2Pout} \\ \dots & Z_{B_2Pin} & Z_{B_2B_1} & Z_{B_2B_2} & \dots & Z_{B_2B_N} & Z_{B_2Pout} \\ \dots & Z_{B_NPin} & Z_{B_NB_1} & Z_{B_NB_2} & \dots & Z_{B_NB_N} & Z_{B_NPout} \\ Z_{PoutPin} Z_{PoutB_1} & Z_{PoutB_2} & \dots & Z_{PoutB_N} & Z_{PoutPout} \\ \end{pmatrix} \\ \times \begin{pmatrix} I_{Pin} \\ I_{B_1} \\ I_{B_2} \\ \dots \\ I_{B_N} \\ I_{Pout} \end{pmatrix} \quad Z_{BP} = V_P / I_P \\ Z_{PP} = V_P / I_P \end{pmatrix}$$

Complex matrix and vectors with voltages and currents on all ports of given structure:

 $\vec{V}(f) = \hat{Z}(f) \cdot \vec{I}(f)$  Z-matrix: 29 x 29



TESLA-Z & CHRISTINE-Z Modeling with generalized frequency dependent response:

Large-signal codes TESLA-Z & CHRISTTINE-Z use of generalized frequency dependent impedance Z matrix<sup>1</sup>

$$\begin{pmatrix} V_{\sigma} \\ V_{\mu} \end{pmatrix} = \hat{Z} \begin{pmatrix} I_{\sigma} \\ I_{\mu} \end{pmatrix} \longrightarrow \begin{pmatrix} I_{\sigma} \\ 2I_{\mu}^{inc} \end{pmatrix} = \hat{\mathbf{Y}} \begin{pmatrix} V_{\sigma} \\ V_{\mu} \end{pmatrix} \qquad \hat{\mathbf{Y}} = \hat{\mathbf{Z}}^{-1} + \hat{\mathbf{Y}}_{\mu} \\ y_{p} \vec{u}_{p} = \hat{\mathbf{Y}} \cdot \vec{u}_{p}$$

<sup>1</sup>I.A. Chernyavskiy, T.M. Antonsen, Jr., J.C. Rodgers, A.N. Vlasov, D. Chernin and, B. Levush, "Modeling Vacuum Electronic Devices Using Generalized Impedance Matrices", IEEE TED, Vol.64, No.2, pp.536-542, Feb. 2017.

leidos

**Output Wave Port** 



# **3D GPU Based PIC Code Neptune**



# Neptune Fundamentals: Time Dependent Maxwell's Equations



Boundary conditions:

- Second order accuracy cut-cell BC on curved boundary
- Advanced port matching conditions

S.J. Cooke et all, ICOPS 2017

Helix TWTA with dielectric support rods

**U.S.NAVAL** 

ABORATORY





- We often need to know the sensitivity of an amplifier figure of merit *F* (e.g., gain, bandwidth, output power, ...) to small changes in the values of VED design parameters  $\vec{p}$  (e.g., circuit pitch profile, beam voltage, current, radius, attenuation profile ...).
- **These sensitivities are expressed as derivatives,**  $\nabla_p F$ . They are needed for:
  - **1.** Manufacturing tolerance analysis
  - 2. Design optimization, using a derivative based algorithm like steepest descent.

We often have many (10's or 100's) design parameters, but just a few figures of merit or performance metrics (gain, Pout, efficiency ...) that we care about:

$$\delta F = \delta \overrightarrow{p} \cdot \nabla_p F$$

The RHS may have many terms, one for each design parameter. We need an efficient way of computing the gradient  $\nabla_p F$ .



# **Optimization Approaches at Large Scale**

**Optimization Problem** 

- Figure of Merit (FoM) or goal function (g)
- Parameters of Optimization



• Gradient Based Approaches

combination of parameters

Evaluate FoM

Select optimum

**Derivative Free Approaches** 

Runs codes to solve full size problem for many

- ✓ Calculate derivatives of *FoM* with respect to parameters
- ✓ Move in parameter space toward optimum

Gradient Approaches are extremely efficient for local optimization with respect to many parameters. Key Challenge: Efficient calculation of derivatives

Adjoint approach is an example of gradient based approaches with efficient way of multidimensional derivative calculations. Adjoint approach demonstrated superiority with respect to other optimization approaches for multi-parameter problems, in particular for shape optimization.

## U.S.NAVAL LRESEARCH



# Direct Approach to Calculation of $\nabla_p F$

- Calculation of a figure of merit F usually involves a computer simulation. Some examples are:
  - Circuit dispersion (HFSS or ANALYST)
  - Beam radius at gun exit (MICHELLE)
  - Gain\*bandwidth (CHRISTINE or TESLA or NEPTUNE)
- The most obvious way to compute  $\nabla_p F$  is to run the computer simulation for each parameter of interest and approximate the derivative as a finite difference:
  - Run 1: p=p0 => F(p0)
  - Run 2: p=p0+dp => F(p0+dp)
  - Approximate: dF/dp ≈ (F(p0+dp)-F(p0))/dp
- This 'Direct Approach' becomes time consuming if the number of parameters N is large:
  - Total number of simulations = N+1 (ideally)
  - However, the Direct Approach may be inaccurate if dp is:
    - Too large (non-linearities)
    - Too small (roundoff)
  - As a practical matter, therefore, the need to test at least several values of dp means that the number of simulations required may be ~ m\*N+1, where m  $\gtrsim$  3-5.



# Alternative: Adjoint Approach



## **Circuit/network theory**

 ✓ S.W. Director and R.A. Rohrer, IEEE Trans. Circuit Theory (1969)

## **Electromagnetic analysis**

- ✓ N.K. Nikolova, et al. IEEE Trans. MTT (2004)
- ✓ HFSS computation of adjoint derivatives

## **Fluid dynamics**

- ✓ A. Jameson, Computational Fluid Dynamics Review (1995)
- ✓ C. Othmer, Journal of Mathematics in Industry (2014)

## Meteorology

✓ R.M. Errico, Bulletin Am Met. Soc. (1997) [review]

## **Plasma Physics**

✓ T. M. Antonsen, Jr. and K.R. Chu, PoF 25, (1982)

## **Beam Optics**

✓ T. M. Antonsen, Jr. and J.J. Petillo PoP 26, (2019)

Example: Calculating the Sensitivity Functions Over an Entire Surface C. Othmer, Journal of Mathematics in Industry (2014)

# Surface Sensitivities

Adjoint methods have been applied in many fields and demonstrated superiority with respect to other optimization approaches for multi-parameter problems, in particular for shape optimization.

Application of adjoint techniques to particular problem requires unique theoretical work on formulation of adjoint problem.



- <u>Shot noise on gyrotron beams</u>, **T.M. Antonsen**, **Jr**., W. Manheimer and A. Fliflet, PoP (2001).
- <u>Beam optics sensitivity function</u>, **T.M. Antonsen**, **Jr**., D. Chernin, J. Petillo, Phys. Plasmas 26, 013109 (2019);
- <u>Stellarator Optimization and Sensitivity</u>, E. Paul, M. Landreman, T.M. Antonsen, Jr., *J. Plasma Phys.* (2019), vol. 85, 905850207, *J. Plasma Phys.* (2021), vol. 87, 905870214
- <u>Optimization of Flat to Round Transformers in Particle</u> <u>Accelerators</u>, L. Dovlatyan, B. Beaudoin, S. Bernal, I. Haber, D. Sutter and **T.M. Antonsen**, Jr., Phys Rev Accel and Beams V25, 044002 (2022).
- <u>Adjoint Equations for Beam-Wave Interaction and</u>
  <u>Optimization of TWT Design</u>, A. Vlasov, **T.M. Antonsen**, Jr., D.
  Chernin and I. Chernyavskiy, IEEE Trans. Plasma Sci. V. March (2022).

# Adjoint Approach in VE Related Fields

An Adjoint Method is a technique based on a reciprocity principle, which facilitates the computation of the effects of small changes of design parameters on selected figures of merit.

U.S.NAVA



Example of reciprocity: Antenna receiving pattern = Antenna transmitting pattern

- We have previously developed adjoint methods for the evaluation of the effects of (1) small changes in geometry in an electron gun an small changes in beam voltage, current, gap spacing profiles in a FWG TWT(2), Helix TWT(3).
  - 1) T.M. Antonsen, Jr., D. Chernin, and J.J. Petillo, "Adjoint approach to beam optics sensitivity based on Hamiltonian particle dynamics," Phys. Plasmas 26, 013109 (2019).
  - 2) A.N. Vlasov, T.M. Antonsen, Jr., D. Chernin, and I.A. Chernyavskiy, "Adjoint Equations for Beam-Wave Interaction and Optimization of TWT Design," IEEE Trans. Plasma Sci 50 2568-2577 (2022).
  - **3)** D. Chernin, A.N. Vlasov, T.M. Antonsen, Jr., and I.A. Chernyavskiy, "Adjoint Equations and Their Application to Helix TWT Design," IEEE Trans. Electron Devices (Early Access) (June 2024).

Application of the adjoint method requires only 3 runs of a simulation code t compute all N partial derivatives, compared with (at least) N+1 runs for the 'Direct Method' !!



# Fundamentals of Adjoint Approach to Beam-Wave Interaction

leidos

## Symplectic Area Conservation Law for particles interacting with EM fields:

- Works explicitly for Hamilton's Equations
- Works implicitly for other forms of particles- field equations

Symplectic Area conserved for any choice of perturbed trajectories 1 and 2

$$\frac{d}{dt} \left( \mathcal{O}\mathbf{p}_1 \times \mathcal{O}\mathbf{q}_2 - \mathcal{O}\mathbf{p}_2 \times \mathcal{O}\mathbf{q}_1 \right) = 0$$

(q(t), p(t)) $(dq_2(t), dp_2(t))$ Coordinates Canonical

 $\frac{d\mathbf{q}}{dt}$   $\frac{d\mathbf{q}}{dt}$   $\frac{d\mathbf{q}}{dt}$   $\frac{d\mathbf{q}}{dt}$   $\frac{d\mathbf{q}}{dt}$   $\frac{d\mathbf{q}}{dt}$   $\frac{d\mathbf{q}}{dt}$ 

Hamilton's Equations  $d\mathbf{q} \quad \P H$ 

¶p

= ....

= -...

Coordinates Canonical Momenta  $(Qq_1(t), Qp_1(t))$ 

 $\frac{d\mathbf{p}}{dt} = -\frac{\P H}{\P \mathbf{q}}$ 

## Perturbed orbit #1

Perturbed orbit #2

$$\frac{d\mathcal{O}\mathbf{q}_{1}}{dt} = \frac{\P^{2}H}{\P\mathbf{p}\P\mathbf{q}} \times \mathcal{O}\mathbf{q}_{1} + \frac{\P^{2}H}{\P\mathbf{p}\P\mathbf{p}} \times \mathcal{O}\mathbf{p}_{1} \qquad \qquad \frac{d\mathcal{O}\mathbf{q}_{2}}{dt}$$
$$\frac{d\mathcal{O}\mathbf{p}_{1}}{dt} = -\frac{\P^{2}H}{\P\mathbf{q}\P\mathbf{q}} \times \mathcal{O}\mathbf{q}_{1} - \frac{\P^{2}H}{\P\mathbf{q}\P\mathbf{p}} \times \mathcal{O}\mathbf{p}_{1} \qquad \qquad \frac{d\mathcal{O}\mathbf{p}_{2}}{dt}$$

### U.S. NAVAL RESEARCH LABORATORY Beam-Wave Interaction in 1D Approximation

Folded Waveguide TWT



Implemented in Christine-Z Code

- RF fields at input frequency + temporal harmonics
- 1D Equations of Electron Motion + Disk Model for Space Charge Fields

• Bunched Beam Effect

• Use Z-matrices for circuit field calculation

 $I_n = I_b \int dz \, e_n^*(z - z_n) \, \langle e^{i\omega t_i(z)} \rangle$  $V_n = \sum_{n'=1,N} Z_{nn'} I_{n'} + 2 \sum_{n'=N+1,P} Z_{nn'} I_{n'}^+$ 

**1D Equations of Motion can be expressed in Hamiltonian Form** 

$$\begin{split} P(E_i, t_i; z) &= m\gamma v_z(E_i) - \sum_n \left[ \frac{iqV_n e_n(z)}{\omega} e^{-i\omega t_i} + c.c. \right] - \left[ \frac{qE_0}{\omega} e^{-i\omega t_i} \langle e^{i\omega t_i} \rangle + c.c. \right] \\ \frac{dE_i}{dz} &= -\frac{\partial}{\partial t_i} P(E_i, t_i; z). \\ \end{split}$$

- 1. A.N. Vlasov, T.M. Antonsen, D. Chernin, I.A. Chernyavskiy, "Adjoint Equations for Beam-Wave Interaction and Optimization of TWT Design, " IEEE Trans. on Plasma Science, IEEE Trans. Plasma Sci 50 2568-2577 (2022).
- I.A. Chernyavskiy, T.M. Antonsen, Jr., J.C. Rodgers, A.N. Vlasov, D. Chernin and B. Levush," Modeling Vacuum Electronic Devices Using Generalized Impedance Matrices," IEEE Trans. Electr. Dev., vol 64, pp. 536-542, Feb 2017,

## U.S. NAVAL RESEARCH Adjoint Approach for 1D Beam-Wave Interaction



wave interaction problem:

- Solution of base problem (known, one run of the code was needed)
- Solution of Adjoint problem (known, two runs of the code was needed)
  - Derivatives for variations of our interest now can be found as sums or integrals without repetitive code runs

14



**Circuit Parameters Variations Term in Adjoint Relationship** 

$$\sum_{\substack{n=1,P\\n'=1,P}} \left[ \delta V_n^{Y^*} \delta Y_{n,n'} V_{n'} - c.c. \right]$$

3D EM calculations

 $\delta Y$  matrix structure



## $\delta$ Z matrix structure



**Evaluation of Local Multi Parameter Variations By N Cells Transmission Line Model** 



N cells based on line-shunt model to calculate  $\delta Y$  matrix (properties of each individual cell might be varied)

Unit cell parameters are fitted using calculated dispersion curve



- We are developing a general purpose design optimization capability for TWTs, based on the adjoint method.
  - By "optimization" we mean finding a maximum or minimum of a scalar figure of merit as a function  $F(\vec{p})$  of a set of design parameters  $\vec{p}$ .
  - $F(\vec{p})$  is usually computed by running a large signal TWT code like CHRISTINE-Z.
- The adjoint method may be used to compute the partial derivatives of  $F(\vec{p})$ , with respect to the design parameters.
  - The adjoint method requires just 2 specially designed runs of a large signal code like CHRISTINE-Z, to compute  $\nabla F(\vec{p})$ , independent of the dimension of  $\vec{p}$  !
  - These partial derivatives may be used in a derivative-based optimizer, e.g., steepest descent

TWT with uniform 54 gaps hybrid serpentine structure connected to input/output waveguides: 20 kV beam

A.N. Vlasov, T.M. Antonsen, D. Chernin, I.A. Chernyavskiy, "Adjoint Equations for Beam-Wave Interaction and Optimization of TWT Design, " IEEE Trans. on Plasma Science, v. 50, pp.2568-2577 June 2022.

## leidos Practically Important FoM for TWT Design U.S. NAVAL Optimization ABORATORY

Find minimum of  $F(\vec{p})$  using steepest Descent Algorithm:



Each calculation of  $F(\vec{p})$  requires 1 runs of CHRISTINE-Z Each calculation of  $\nabla F(\vec{p}_n)$  requires 2 runs of CHRISTINE-Z

A.N. Vlasov, T.M. Antonsen, D. Chernin, I.A. Chernyavskiy, "Adjoint Equations for Beam-Wave Interaction and Optimization of TWT Design, " IEEE Trans. on Plasma Science, v. 50, pp.2568-2577 June 2022.

# Constrained Optimization: Optimize Circuit Pitch



Avoids backward wave and band edge oscillations.

**U.S.NAVAL** 

ABORATORY





# **Optimization of Two Section TWT**



**Optimization of Small Signal Gain** 

Distance between gaps (two sections). 3 different goal functions 2 optimization parameters

$$\square \quad p=L_{g1}, \ \ L_{g2}$$

• Maximize: 
$$F_1(p) = \frac{1}{f_2 - f_1} \int_{f_1}^{f_2} G(f) df$$

Minimize: 
$$F_2(p) = \frac{1}{f_2 - f_1} \int_{f_1}^{f_2} (G(f) - \overline{G})^2 df$$

**Maximize:** 
$$F_3(p) = \Gamma_p \equiv \frac{\langle g \rangle}{\sigma_q^p}$$



# LABORATORY Example: Find Circuit Pitch Profile to Maximizidos Output Power at a Single Frequency (91 GHz)



- 3 parameter optimization
  - flat pitch section (1 parameter) + linear ramp section (2 parameters)
  - Pin = 3W fixed
- Base case: Flat pitch profile, Pout(91 GHz) = 77.9 W
- Result of optimization: Pout(91 GHz) = 101.1 W (+1.1 dB)
  - Peak output power increased from 79.3 W to 136 W (+2.3 dB)

# Example of Klystron Optimization Using Adjoint Approach

CPI 4-cavity X-band klystron\*

USNAVA

ABORATORY

- Fixed cavity properties (Z-matrix)
- Variable beam voltage, cavity spacings
- Base case parameters:
  - $-V_{b} = 115 \text{ kV}$
  - $-I_{b} = 78.1 \text{ A}$
  - $-r_{\rm b} = 0.152$  cm
  - Cavity spacings:
    - » 1-2: 2.776 cm
    - » 2-3: 4.265 cm
    - » 3-4: 1.676 cm
- Example: Maximize average small signal gain across band
  - 6 equally spaced frequencies in [9.8-10.1] GHz

\*Courtesy of R. Begum, CPI





## **Klystron Optimization**





**U.S.NAVAL** 

Result: Increase in average gain by ~0.2 dB by increasing beam voltage by ~ 6.4 kV (5.6%).

135

from optved.fom

130

125

# Klystron Optimization of Spacing Between Cavities

U.S.NAVA

ABORATORY



Result: Increase in average gain by ~1.7 dB by increasing cavity 3-4 spacing by ~ 1.4 cm (84%).



## X-Band 4 Cavity Klystron Optimize Cavities 2-3 and 3-4 Spacing









- Adjoint methods are a powerful way to evaluate parameter dependences in many RF Sources based on self-consistent beam-wave interaction.
- An adjoint method for the evaluation of the sensitivity of the performance of a FW TWT and Klystron to small changes in various design parameters has been formulated and implemented in the CHRISTINE-1D large signal code.
- The method has been successfully tested against direct calculation of changes in gain and output power due to changes in various design parameters, including beam voltage, beam current, circuit phase velocity and circuit impedance.
- Our recent and planned work include development and implementation of the adjoint method for:
  - Multi-parameter optimization for large signal case for klystron
  - Algorithms for Case with large number of design parameters
  - 2D large signal simulations (TESLA-Z code)