LHC EFT WG General Meeting - CERN, 3rd December 2024

Istituto Nazionale di Fisica Nucleare

DI PERUGIA

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Dim-6 and dim-8 EFT studies with VBS at CMS

on behalf of the CMS Collaboration

Why is VBS so charming?

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i.e. Wilson coefficients

- and EWSB mechanism
- **Anomalies modeling with EFT approach** independent approaches like EFT.

Why is VBS so charming?

- **Investigation of both SM and (possible) BSM effects** VBS processes are strictly related to **unitarity preservation** in SM and precise measurements can probe the nature of the Higgs sector and EWSB mechanism
- **Anomalies modeling with EFT approach** VBS processes exhibit both QGC and TGC vertices and can be sensitive to deviations from the SM, parameterizable via modelindependent approaches like EFT.

* c_i and f_i are EFT couplings, i.e. Wilson coefficients

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Overview of VBS same-sign W with one hadronic tau study

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH (CERN)

CERN-EP-2024-234 2024/10/08

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EWK production of same-sign W boson pairs with a hadronically decaying τ in the final state (@13 TeV, full RunII data):

$$
qq' \rightarrow W^{\pm}W^{\pm}q''q''' \rightarrow l^{\pm}\tau_{h}jj \nu_{l}\nu_{\tau} \quad (l = e, \mu)
$$

FIRST

TAU!

CMS-SMP-22-008

Study of same-sign W boson scattering and anomalous couplings in events with one tau lepton from pp collisions at $\sqrt{s} = 13 \text{ TeV}$

The CMS Collaboration[®]

Abstract

A first measurement is presented of the cross section for the scattering of same-sign W boson pairs via the detection of a τ lepton. The data from proton-proton collisions at the center-of-mass energy of 13 TeV were collected by the CMS detector at the LHC, and correspond to an integrated luminosity of 138 fb^{-1} . Events were selected that contain two jets with large pseudorapidity and large invariant mass, one τ lepton, one light lepton (e or μ), and significant missing transverse momentum. The measured cross section for electroweak same-sign WW scattering is $1.44^{+0.63}_{-0.56}$ times the standard model prediction. In addition, a search is presented for the indirect effects of processes beyond the standard model via the effective field theory framework, in terms of dimension-6 and dimension-8 operators.

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FIRST

TAU!

Dileptonic tt production (1%)

Both signal and background events are fully processed through detector reconstruction \rightarrow Full simulation analysis

Main backgrounds sources in SR are:

- events containing nonprompt leptons from QCD-mediated multijet, W+jets, hadr. and semi-leptonic tt (**95%)**
- Z/γ^* + jets (2%)

SM measurements

Measurement of purely EW ssWW signal strength and EWK+QCD ssWW signal strength:

EFT sensitivity study

Likelihood scan performed using a dedicated transverse mass as the discriminating variable.

- 1D scans, both for dim-6 and dim-8 operators (first time in VBS)
- 2D simultaneous scans of relevant operator pairs, first time ever with different dimension.

Focus of today: dim-6 vs dim-8 fits.

Cross Section [fb] 28.7 51.0

Deliverables from the analysis

SM measurements

Measurement of purely EW ssWW signal strength and EWK+QCD ssWW signal strength:

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Deliverables from the analysis

List of operators included

 $c_i^{(6)}$

$$
\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \left(\sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)}\right)
$$

4

Dim-6 operators Dim-8 operators

$$
\mathcal{O}_{ll} = \delta_{pr} \delta_{st} (\bar{l}_p \gamma_\mu l_r) (\bar{l}_s \gamma^\mu l_t)
$$

\n
$$
\mathcal{O}'_{ll} = \delta_{pt} \delta_{sr} (\bar{l}_p \gamma_\mu l_r) (\bar{l}_s \gamma^\mu l_t)
$$

\n
$$
\mathcal{O}^{(1)}_{qq} = \delta_{pr} \delta_{st} (\bar{q}_p \gamma_\mu q_r) (\bar{q}_s \gamma^\mu q_t)
$$

\n
$$
\mathcal{O}^{(1)'}_{qq} = \delta_{pt} \delta_{sr} (\bar{q}_p \gamma_\mu q_r) (\bar{q}_s \gamma^\mu q_t)
$$

\n
$$
\mathcal{O}^{(3)}_{qq} = \delta_{pr} \delta_{st} (\bar{q}_p \gamma_\mu \tau^i q_r) (\bar{q}_s \gamma^\mu \tau^i q_t)
$$

\n
$$
\mathcal{O}^{(3)'}_{qq} = \delta_{pt} \delta_{sr} (\bar{q}_p \gamma_\mu \tau^i q_r) (\bar{q}_s \gamma^\mu \tau^i q_t)
$$

Warsaw basis [arXiv:1008.4884](https://arxiv.org/abs/1008.4884)

SMEFTsim@LO CP conservation U(3)5 flavor symmetry

$$
\mathcal{O}_{\varphi l}^{(1)} = (\varphi^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} \varphi)(\bar{l}_{p} \gamma^{\mu} l_{r}) \qquad \mathcal{O}_{W} = \epsilon^{ijk} W_{\mu}^{\nu i} W_{\nu}^{\rho j} W_{\rho}^{\mu k}
$$
\n
$$
\mathcal{O}_{\varphi l}^{(3)} = (\varphi^{\dagger} i \overset{\leftrightarrow}{D^{i}}_{\mu} \varphi)(\bar{l}_{p} \tau^{i} \gamma^{\mu} l_{r}) \qquad \mathcal{O}_{\varphi \Box} = (\varphi^{\dagger} \varphi) \Box (\varphi^{\dagger} \varphi)
$$
\n
$$
\mathcal{O}_{\varphi q}^{(1)} = (\varphi^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} \varphi)(\bar{q}_{p} \gamma^{\mu} q_{r}) \qquad \mathcal{O}_{\varphi W} = \varphi^{\dagger} \varphi W_{\mu\nu}^{i} W^{\mu\nu i}
$$
\n
$$
\mathcal{O}_{\varphi q}^{(3)} = (\varphi^{\dagger} i \overset{\leftrightarrow}{D^{i}}_{\mu} \varphi)(\bar{q}_{p} \tau^{i} \gamma^{\mu} q_{r}) \qquad \mathcal{O}_{\varphi W B} = \varphi^{\dagger} \tau^{i} \varphi W_{\mu\nu}^{i} B^{\mu\nu}
$$

Eboli basis [arXiv:1604.03555v1](https://arxiv.org/abs/1604.03555) generating genuine **anomalous-Quartic-Couplings dedicated UFO model** U(3)5 flavor symmetry

 $\frac{f_i^{(8)}}{\Lambda^4}$ $\mathcal{O}_i^{(8)}$

 ${\cal O}_{S,0}=\left[(D_\mu\Phi)^\dagger D_\nu\Phi\right]\times \left[(D^\mu\Phi)^\dagger D^\nu\Phi\right]$ ${\cal O}_{S,1}=\left[(D_\mu\Phi)^\dagger D^\mu\Phi\right]\times \left[(D_\nu\Phi)^\dagger D^\nu\Phi\right]$ ${\cal O}_{S,2} = \left[(D_\mu \Phi)^\dagger D_\nu \Phi \right] \times \left[(D^\nu \Phi)^\dagger D^\mu \Phi \right]$

$$
\begin{aligned} \mathcal{O}_{M,0} &= \text{Tr}\bigl[\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}\bigr] \times \bigl[\bigl(D_\beta \Phi\bigr)^\dagger D^\beta \Phi\bigr] \\ \mathcal{O}_{M,1} &= \text{Tr}\bigl[\hat{W}_{\mu\nu} \hat{W}^{\nu\beta}\bigr] \times \bigl[\bigl(D_\beta \Phi\bigr)^\dagger D^\mu \Phi\bigr] \\ \mathcal{O}_{M,7} &= \bigl[\bigl(D_\mu \Phi\bigr)^\dagger \hat{W}_{\beta\nu} D^\nu \Phi\bigr] \end{aligned}
$$

$$
\begin{aligned} \mathcal{O}_{T,0} &= \text{Tr}\big[\hat{W}_{\mu\nu}\hat{W}^{\mu\nu}\big] \times \text{Tr}\big[\hat{W}_{\alpha\beta}\hat{W}^{\alpha\beta}\big] \\ \mathcal{O}_{T,1} &= \text{Tr}\big[\hat{W}_{\alpha\nu}\hat{W}^{\mu\beta}\big] \times \text{Tr}\big[\hat{W}_{\mu\beta}\hat{W}^{\alpha\nu}\big] \\ \mathcal{O}_{T,2} &= \text{Tr}\big[\hat{W}_{\alpha\mu}\hat{W}^{\mu\beta}\big] \times \text{Tr}\big[\hat{W}_{\beta\nu}\hat{W}^{\nu\alpha}\big] \end{aligned}
$$

Simulation @LO with MADGRAPH5 aMC@NLO v2.6.5

 \mathbf{b}

Combined EFT operators study

• CMS analysis framework breaks down in presence of negative histograms (coming from EFT linear terms). Since SM \cdot dim8 interference can give negative contributions, dim 8^2 term is included.

The number of expected events is proportional to square module of scattering amplitude.

$$
|\mathcal{A}_{\text{TOT}}|^2 = |\mathcal{A}_{SM}|^2 \left| + \sum_{i}^{N_{dim6}} \left[\frac{c_i}{\Lambda^2} 2 \Re \left(\mathcal{A}_{\mathcal{O}_i^{(6)}} \mathcal{A}_{\mathcal{S}N}^* \right) \right. \\ \left. + \sum_{k}^{N_{dim8}} \left[\frac{f_k}{\Lambda^4} 2 \Re \left(\mathcal{A}_{\mathcal{O}_k^{(8)}} \mathcal{A}_{\mathcal{S}N}^* \right) \right] \right|
$$

- EFT operator contributions up to Λ^4 are considered;
-

Which pairs are we considering?

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We want to avoid looking at the statistical effects coming with multiple EFT contributions. The coupling criteria are the following:

-
-

• Operators acting on the same vertex $(W^{\pm}W^{\pm} \rightarrow W^{\pm}W^{\pm}$ scattering vertex); • Operators giving similar contributions to the process (ratio between linear and quadratic terms). **Any effect arising from such a combination can be interpreted in terms of physics effects.**

2D scans: dim6 VS dim8 pairs

7

cW cHW cHbox

L = $138 \text{ fb}^{-1} (13 \text{ TeV})$

 c_{Hbox}

 c_{Hbox}

 c_{Hbox}

FS0

FT0

FM0

7

2D scans: dim6 VS dim8 pairs

cW cHW cHbox

Conclusions

In EFT, we cannot always assume that dim-6 operators are negligible compared to dimension-8 ones.

- Overestimation of **dim-8 experimental sensitivity** occurs if **dim-6 operators** are not included in the analysis.
- The interplay between dimension-6 and dimension-8 effects cannot be determined a priori without explicit consideration in the analysis -> potential bias on NP

- SmeftFR model [\(arXiv:2302.01353](https://arxiv.org/abs/2302.01353)) includes operators of both dimensions:
	- Warsaw basis for dim-6;
- **Murphy's basis*** for dim-8 bosonic operators.
- Possibility to perform dim-6+dim-8 simulations including effects of dim8 on TGCs.
	- ***** In Eboli we have dim8 operators affecting QGCs only and with a limited number of derivatives. Murphy's one is a full basis of dim8 operators. It's possible to map operators from one basis to another.

Common model for dim-6 and dim-8 operators: SmeftFR model

Smeft FR v3 – Feynman rules generator for the Standard Model **Effective Field Theory**

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Abstract

We present version 3 of SmeftFR, a Mathematica package designed to generate the Feynman rules for the Standard Model Effective Field Theory (SMEFT) including the complete set of gauge invariant operators up to dimension-6 and the complete set of bosonic operators of dimension-8. Feynman rules are generated with the use of FeynRules package, directly in the physical (mass eigenstates) basis for all fields. The complete set of interaction vertices can be derived, including all or any chosen subset of SMEFT operators. As an option, the user can also choose preferred gauge fixing, generating Feynman rules in unitary or R_{ξ} -gauges. The novel feature in version-3 of SmeftFR is its ability to calculate SMEFT interactions consistently up to dimension-8 in EFT expansion (including quadratic dimension-6 terms) and express the vertices directly in terms of user-defined set of input-parameters. The derived Lagrangian in the mass basis can be exported in various formats supported by FeynRules, such as UFO, FeynArts, etc. Initialisation of numerical values of Wilson coefficients of higher dimension operators is interfaced to WCxf format. The package also includes a dedicated Latex generator allowing to print the result in clear human-readable form. The SmeftFR v3 is publicly available at www.fuw.edu.pl/smeft.

Keywords: Standard Model Effective Field Theory, Feynman rules, unitary and R_{ξ} -gauges

HIGH NUMBER OF OPERATORS!

PROS:

Effects of interplay between operators of different dimensions on TGCs

CONS:

High number of operators \rightarrow increase of degrees of freedom in the fits (leads to instability of fits due to limited statistics)

Common model for dim-6 and dim-8 operators: SmeftFR model

Summary

-
-

In this analysis we demonstrated that there is an interplay between different dimensions (even though we are not considering the dim6-dim8 interference term!)

- There is technical complexity in considering both dimensions, but it's a more correct approach to preserve model independence in NP searches;
- A validated model that allows us to perform a complete dim6+dim8 study on TGCs is available.

We may need a full model including dim6-dim8 interference terms to see correlations between two

different dimensions.

Backup

1D results

Uncertainties

Uncertainty source Theory (PDF, scales, ISR, FS. Nonprompt background est $t\bar{t}$ normalization Trigger mistiming Luminosity b tagging and mistagging Jet energy scale, resolution, Pileup LO-to-NLO VBS corrections Unclustered energy τ_h energy scale and identific Charge misidentification Lepton reconstruction, ident Background statistical Total systematic Data statistical Total

Details on Control and Signal Regions

Generation details

Reweighting

define $p = g u c d s b u^2 c^2 d^2 s^2 b^2$ define $j = p$ define $l+ = e+ mu+ ta+$ $define 1- = e- mu- ta$ define $v1 = ve$ vm vt define $v1 - v$ = $ve - v$ m $v1 - v$ generate p p > 1+ $v1$ 1+ $v1$ j j SMHLOOP=0 QCD=0 NP=1 add process p p > 1- $v1$ ~ 1- $v1$ ~ j j SMHLOOP=0 QCD=0 NP=1 output WWjjTolnulnu_SS_ewk_dim6

```
import model QCKM_5_Aug21v2
define p = g u c d s b u~ c~ d~ s~ s~ v~define j = pdefine l+ = e+ mu+ ta+define 1- = e- mu- ta-define v1 = ve vw vtdefine v1 - v = ve- vm- vt-
generate p p > w+ w+ j j QED=4 QCD=0 NP=1, w+ > 1+ v1 @ 1
add process p \, p > w - w - j j QED=4 QCD=0 NP=1, w - > 1 - v1 \sim 0 2
output WWjj_SS_dim8_ewk
```
import model SMEFTsim_U35_MwScheme_UFO-cW_cHWB_cHDD_cHbox_cHW_cHl1_cHl3_cHq1_cHq3_cqq1_cqq11_cqq31_cqq3_cll_cll1_massless

Combine formula for negative histograms

Combine can handle only non-negative defined histograms, then the simple model has

to be re-written, such that only non-negative defined inputs are provided.

 $A = A$

 $N = SM + k$

 $N = SM + k \cdot Lin + k^2 \cdot Quad$ $= SM + k \cdot (SM + Lin + Qu)$ $= SM \cdot (1-k) + k \cdot (SM +$

$$
A_{SM} + k * Op
$$

$$
k * Lin + k2 * Quad
$$

$$
Can be negative
$$

$$
uad) - k \cdot SM - k \cdot Quad + k^2 \cdot Quad
$$

Lin + Quad) + $(k^2 - k) \cdot Quad$

SmeftFR - dim8

Dim-5 + dim6 operators from Warsaw basis Dim8 bosonic operators from Murphy's basis Doesn't consider dim8 fermionic opertors, nor dim7.

Same input schemes of SMEFTsim (AEM and GF)

 $W^{J\mu\nu}$ $W^{J\mu\nu}$

 $2\mu\nu$ $2\mu\nu$

 $\overline{W^I_{\mu\nu}} \ \widetilde{W}^I_{\mu\nu} \ \nu \varphi) W^K_{\mu\nu} \ \nu \varphi) \widetilde{W}^K_{\mu\nu}$ $\mu\nu$

 $\epsilon_{\mu\nu}$

SmeftFR - validation

 Δ

EFT insertions in Murphy

Murphy Cphi6Box, Cphi6D2, CW2phi4n*, CW3phi2n*, CWphi4D2n*, CB2phi4n*, CBphi4D2n*, CW2Bphi2n*, CWBphi4n*

Effects on UV matching

Benefits from including dimensions consistently with power counting:

- More consistent results in terms of experimental findings
- Preserving model independence in EFT approach
- Unlock reliable matching with BSM models (better hints on where NP hides)

Example with Z2RSE model [arXiv:2311.16897](https://arxiv.org/abs/2311.16897)

$$
\phi = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}} (v + h_1 + iG_0) \\ H \end{pmatrix}, \quad S = \frac{v_s + h_2}{\sqrt{2}}
$$

$$
\begin{pmatrix} h \\ H \end{pmatrix} = \begin{pmatrix} c_{\chi} & -s_{\chi} \\ s_{\chi} & c_{\chi} \end{pmatrix} \begin{pmatrix} h_1 \\ h_2 \end{pmatrix}
$$