

# Search for the critical point in NA61/SHINE



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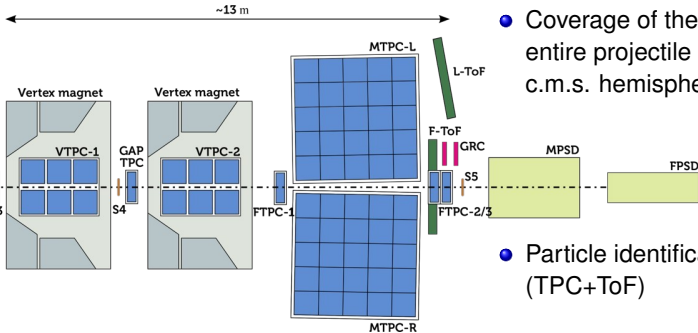
- 1 The NA61/SHINE experiment
- 2 Net electric charge fluctuations
- 3 Bose-Einstein (HBT) correlations (femtoscscopy)
- 4 Intermittency analysis of scaled factorial moments
- 5 Intermittency methodology improvements
- 6 Summary & Outlook

# The NA61/SHINE experiment

NA61/SHINE  
at CERN SPS



Multipurpose fixed target spectrometer  
with unique capabilities



- Coverage of the entire projectile c.m.s. hemisphere

- Particle identification (TPC+ToF)

- **Strangeness in quark matter:**  
 $K^+$ ,  $K^-$ ,  $K_s^0$ ,  $K^*$ ,  
 $\Lambda$ ,  $\phi$

- **Heavy quarks:**  
 $D^0$  and  $\bar{D}^0$

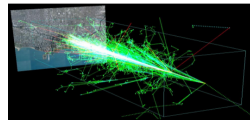
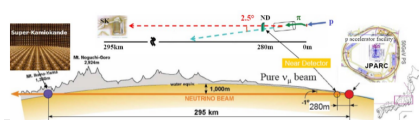
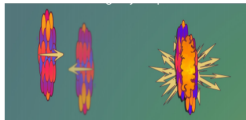
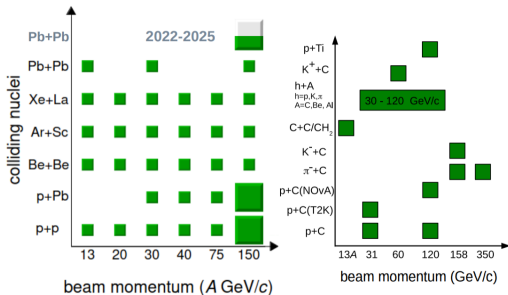
- **Correlations, fluctuations, HBT, intermittency...**

## Strong interactions

- study the onset of deconfinement
- **search for the critical point**
- measurement of open charm

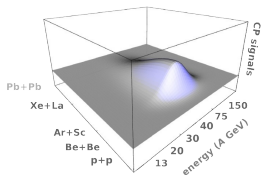
## Neutrino and cosmic ray physics

- measurements for neutrino programs (J-PARC, Fermilab)
- measurements for cosmic-ray physics (Pierre-Auger, KASCADE, satellite experiments)

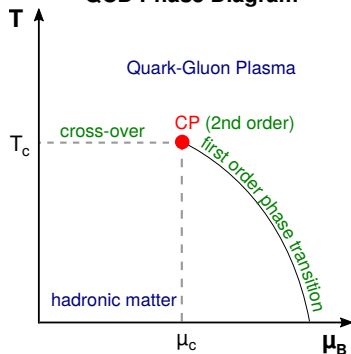


# The NA61/SHINE critical point search

- **Critical point (CP)** — a hypothetical end point of **first order** phase transition line (QGP-HM) that has properties of **second order** phase transition
- **2<sup>nd</sup> order** phase transition → **scale invariance** → **power-law form of correlation function**
- Expectations for **enhanced fluctuations and correlations** in **configuration space** → projected to **momentum space** via quantum statistics and/or collective flow

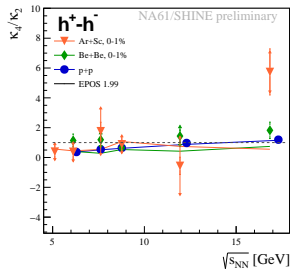
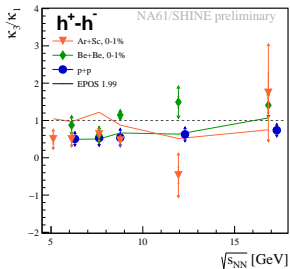
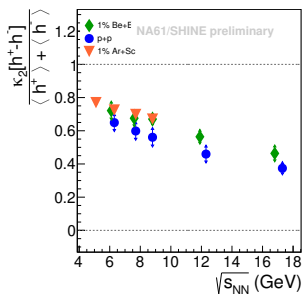


## QCD Phase Diagram



- **Scan** in the **experimentally controlled parameters** (collision energy, nuclear mass number, centrality). Conjecture is that, by varying them, we vary **freeze-out conditions** ( $T$ ,  $\mu_B$ )

# Net electric charge fluctuations



$$\kappa_1 = \langle N \rangle$$

$$\kappa_2 = \langle (\delta N)^2 \rangle$$

$$\kappa_3 = \langle (\delta N)^3 \rangle$$

$$\kappa_4 = \langle (\delta N)^4 \rangle$$

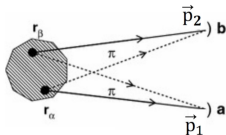
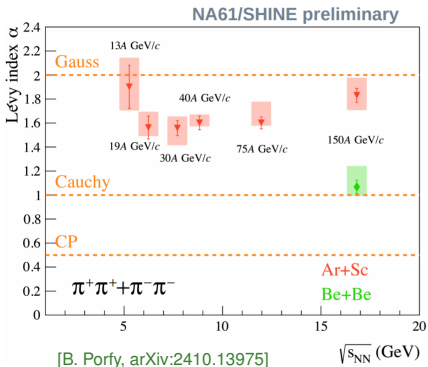
$$\langle N_2 \rangle \sim \xi^2$$

$$\langle N_4 \rangle \sim \xi^7$$

No significant non-monotonic signal observed

[NA61/SHINE, PoS(PANIC2021)238]  
[NA61/SHINE, Status Report 2022]

# Bose-Einstein (HBT) correlations (femtoscology)



Correlation function from Lévy source:

$$C(q) = 1 + \lambda e^{-|qR|^\alpha}$$

$$q = |\vec{p}_1 - \vec{p}_2|$$

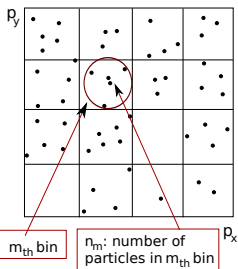
- Bose-Einstein correlations (femtoscology) reveal the space-time structure of hadron production
- The Lévy parameter  $\alpha$  describes the shape of the source and is sensitive to the system freezing out at the CP
- The new Ar+Sc results are close to Gaussian, and **far from the CP**

[Csörgő, Hegyi, Novák, Zajc, AIP Conf. Proc. 828 (2006) 525]

Ar+Sc, 0-10% central, NA61/SHINE preliminary

Be+Be, 0-20% central, NA61/SHINE, EPJC 83 (2023) 919

# Proton intermittency – scaled factorial moments $F_r(M)$



- When the system **freezes out at CP**, the **scaled factorial moments  $F_r(M)$**  are expected to follow a **power-law** behaviour:

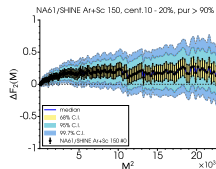
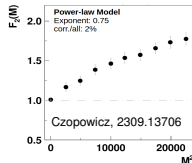
$$F_r(M) \sim (M^2)^{\phi_r}$$

- For **protons** and  $r = 2$ ,  $\phi_2 = 5/6$  is expected
- Either **correlated<sup>1</sup>** or **statistically independent data points<sup>2</sup>** can be used

- Cumulative variables<sup>3</sup>** or **mixed-event moment subtraction<sup>4</sup>** handle **baseline correlations**

$$F_r(M) \equiv \frac{\left\langle \frac{1}{M^2} \sum_{m=1}^{M^2} n_m(n_m - 1) \dots (n_m - r + 1) \right\rangle}{\left\langle \frac{1}{M^2} \sum_{m=1}^{M^2} n_m \right\rangle^r}$$

$M^2$  – number of bins;  $\langle \dots \rangle$  – averaging over events



[Białas, Peschanski, NPB 273 (1986) 703]

[Wosiek, APPB 19 (1988) 863]

[Asakawa, Yazaki, NPA 504 (1989) 668]

[Barducci et al., PLB 231 (1989) 463]

[Satz, NPB 326 (1989) 613]

[Antoniou et al., PRL 97 (2006) 032002]

<sup>1</sup>[NA61/SHINE, APP.Supp. 13 (2020) 637]

<sup>2</sup>[NA61/SHINE, EPJC 83 (2023) 881]

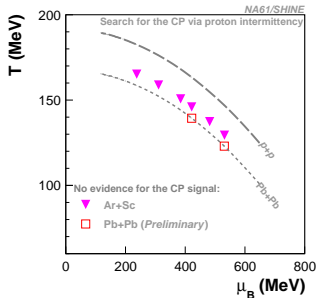
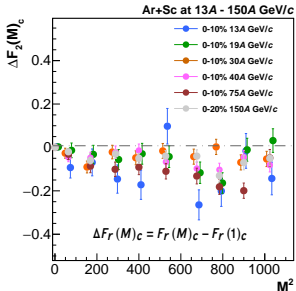
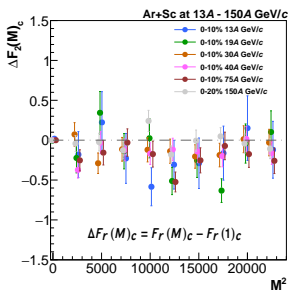
<sup>3</sup>[Białas, Gazdzicki, PLB 252 (1990) 483]

<sup>4</sup>[NA49, EPJC 75 (2015) 587]



# SHINE $^{40}\text{Ar} + ^{45}\text{Sc}$ independent bin proton intermittency

- **No signal** indicating the **critical point** in **cumulative independent bin analysis**



$$1^2 \leq M^2 \leq 150^2$$

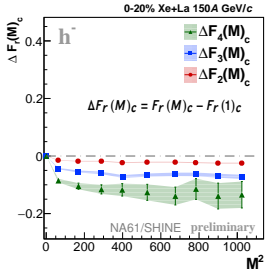
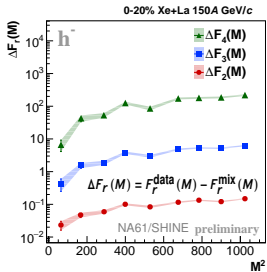
$$1^2 \leq M^2 \leq 32^2$$



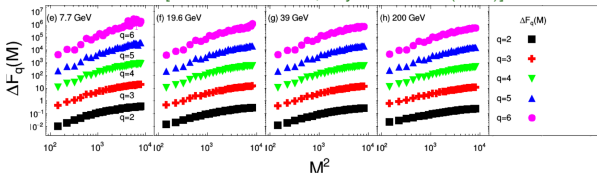
number of subdivisions in  
cumulative transverse momentum space

[NA61/SHINE, EPJC 83 (2023) 881]  
[NA61/SHINE, EPJC 84 (2024) 7, 741]

# Intermittency of negatively charged hadrons

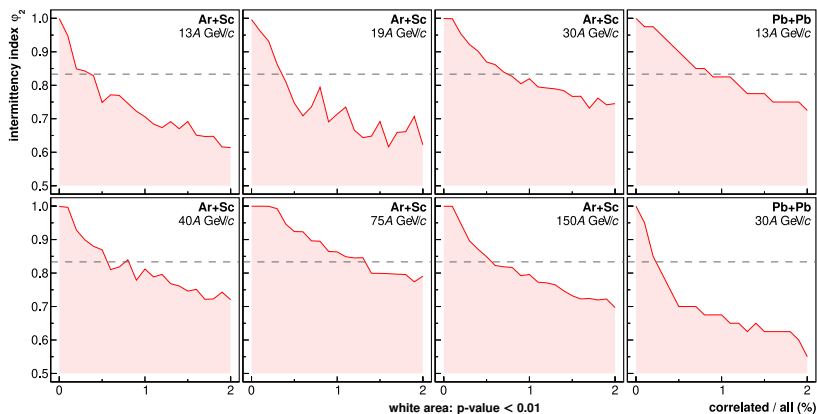


[STAR collaboration, Phys.Lett.B 845 (2023)]



- **STAR collaboration:** intermittency results of  $\Delta F_q(M) \equiv F_q^{\text{data}}(M) - F_q^{\text{mixed}}(M)$ ,  $q = 2 - 6$  of all charged hadrons in **0-5% Au+Au collisions**
- **NA61/SHINE:**  $\Delta F_r(M)$ ,  $r = 2 - 4$  for **non-cumulative and cumulative transformed  $p_T$  binning**, of **negatively charged ( $h^-$ ) hadrons in 0-20% Xe + La collisions**
- $\Delta F_q(M)$  increases with  $M^2$  up to  $M^2 \sim 4000$  in STAR results; interpretation of this increase was **unclear**, with **no specific theoretical prediction** given for  $h^\pm$  critical scaling
- **Cumulative transform and/or short-range correlation cut removes** corresponding effect in **SHINE  $h^-$  analysis**, indicating a **systematic effect** as the origin of observed STAR scaling

# Independent bin proton intermittency – exclusion plots



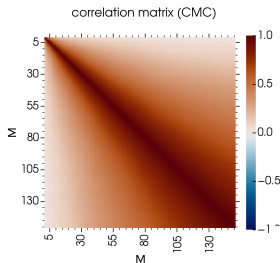
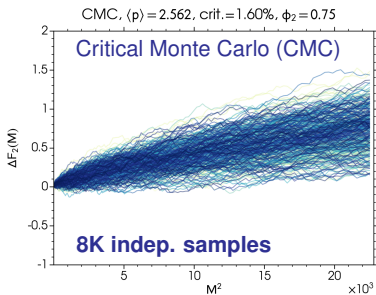
Exclusion plots for parameters of simple power-law model:

- power-law exponent  $\phi$  in  $|\Delta\vec{p}_T|$  correlation function  $\rho(|\Delta\vec{p}_T|) = |\Delta\vec{p}_T|^{-\phi}$ ,  $\varphi_2 = (\phi + 1)/2$
- fraction of correlated particles

Expected intermittency index:  $\varphi_2 = 5/6$  (3D Ising universality class)

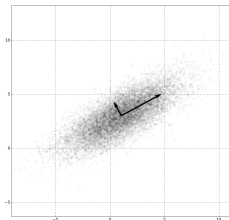
# Handling bin correlations through PCA

- $F_2(M)$  values for different **M-bin sizes** are **correlated**, if the **same events** are used to calculate **different bins**; this **invalidates** fitting & model comparison
- **Independent points** can be used, but **point uncertainties increase!**



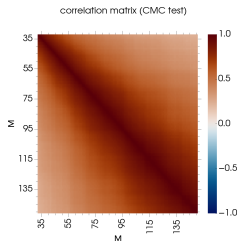
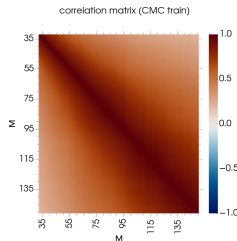
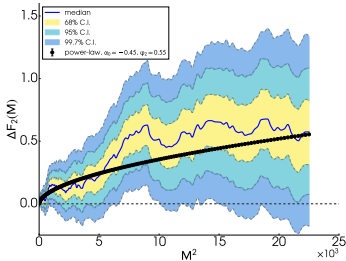
- An alternative is to **untangle correlations**

- We can do this via **Principal Component Analysis (PCA)**: center and scale sample points in  **$M$ -space**, then **rotate** the axes to make **independent linear combinations of  $M$ -bins**. Finally, keep **only few significant components** [N. Davis, arXiv:2409.14185]

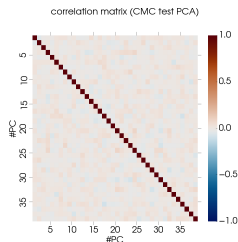
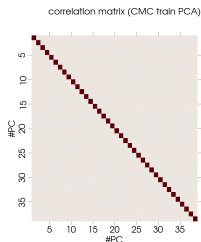
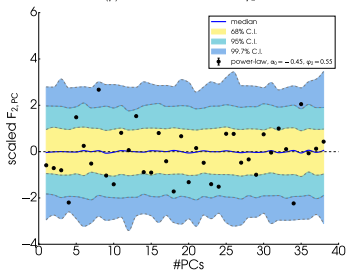


# Performing a scan in power-laws with PCA (CMC “data”)

CMC,  $\langle p \rangle = 2.562$ , crit.=1.00%,  $\phi_2 = 0.825$

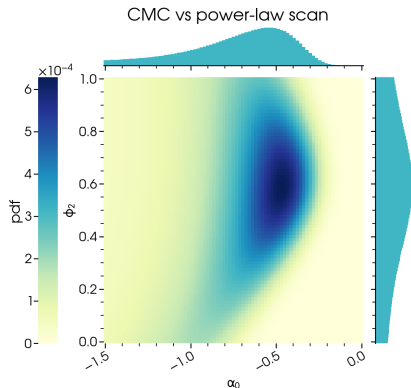
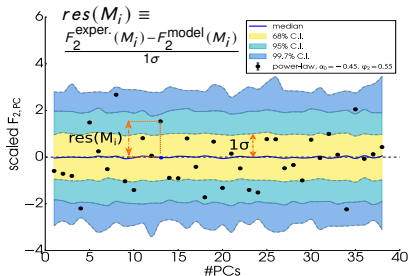


CMC,  $\langle p \rangle = 2.562$ , crit.=1.00%,  $\phi_2 = 0.825$



- **Original sample becomes PC baseline; all power-laws compared to it.**

# Estimating power-law model likelihoods



[N. Davis, arXiv:2409.14185]

## Power-law model

$$\Delta F_2(M) \equiv 10^{\alpha_0} \left( \frac{M^2}{10^4} \right)^{\varphi_2}$$

$\alpha_0$  : power-law strength,  
 $\varphi_2$  : power-law exponent

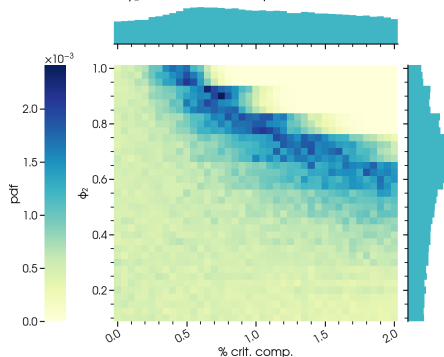
$$\chi^2 = \sum_i res^2(M_i) \Rightarrow \text{Model Weight} \sim e^{-\frac{\chi^2}{2}}$$

- Scan in power-law parameters  $\Rightarrow$  best-fitting power laws to the data
- Critical component can be estimated by power-law strength  $\alpha_0$
- PCA transformation ensures valid model weighting

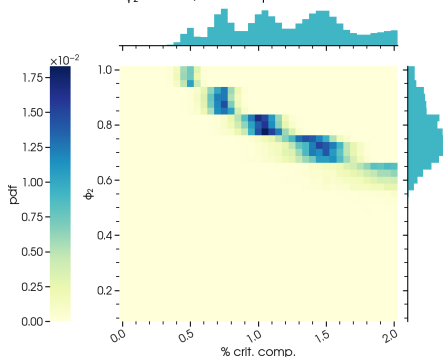
# The role of event statistics

- **Event statistics** (number of analyzed events) **greatly affects predictive power** of model scans!

CMC signal vs itself, PCA  
 $\phi_2 = 0.825$ , crit. comp. = 1.00%



CMC signal vs itself, PCA 10x  
 $\phi_2 = 0.825$ , crit. comp. = 1.00%



- **Critical Monte Carlo (CMC)** model simulations indicate that a  $\sim 10\times$  increase in event statistics for **SHINE** could detect **as weak as a 1%** critical component signal! The **upgraded NA61/SHINE detector** is expected to provide **sufficient data**

[M. Maćkowiak-Pawłowska *et al.* [NA61/SHINE], CERN-SPSC-2023-022]

# Summary

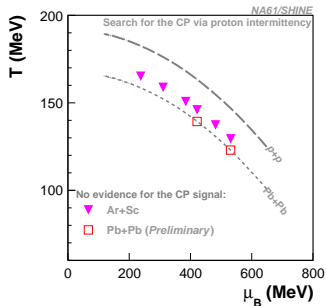
- Results on **net-charge fluctuations** in p+p, Be+Be and Ar+Sc energy scans show **no non-monotonic signal**
- Obtained exponents from the Lévy-shaped source fit in the **HBT analyses** of pions produced in Be+Be at  $\sqrt{s_{NN}} \approx 17$  GeV and Ar+Sc energy scan are **far from the values predicted for the critical point**
- Results on the dependence of **proton scaled factorial moments** of multiplicity distribution on cumulative momentum bin size, analyzed using independent data points for:

- protons in Pb+Pb at  $\sqrt{s_{NN}} \approx 5$  GeV
- protons in Pb+Pb at  $\sqrt{s_{NN}} \approx 7.5$  GeV
- protons in Ar+Sc at  $\sqrt{s_{NN}} \approx 5 - 17$  GeV

show **no indication of a power-law increase**

- **No indication of a power-law increase** in **negatively charged hadron factorial moments** in Xe+La ( $\sqrt{s_{NN}} \approx 17$  GeV) when **cumulative  $p_T$  bins** are used
- Please also see **F. Diakonov talk** for a discussion of the **STAR  $h^{+-}$  intermittency** result!

Status of NA61/SHINE CP search via proton intermittency



Points indicate analyzed reactions with no evidence for CP. They are placed at  $T - \mu_B$  values calculated from Becattini, Manninen, Gazdzicki, Phys. Rev. C73 (2006)



- Long-standing **bin-by-bin correlation problem** now **effectively solved**; **Principal Component Analysis (PCA)** methodology allows **direct handling** of **factorial moment bin correlations** using **the full event statistics**
- **Critical Monte Carlo (CMC)** model simulations indicate that a  **$\sim 10\times$  increase** in event statistics for **SHINE** could detect **as weak as a 1%** critical component signal. The **upgraded NA61/SHINE detector** is expected to provide **sufficient data**



*Thank You!*

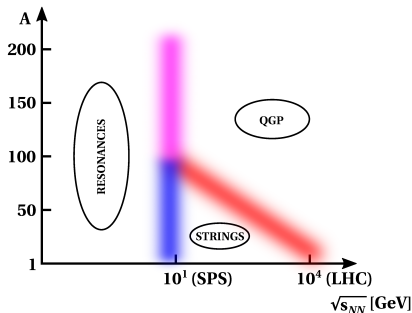
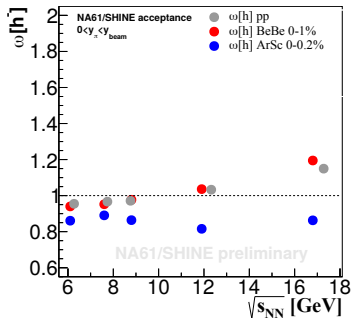


# Backup Slides

# Backup Slides Outline

- 7 Multiplicity & multiplicity- $p_T$  fluctuations
- 8 The bootstrap
- 9 Critical Monte Carlo Simulations
- 10 Independent bin analysis with cumulative variables
- 11  $h^-$  intermittency

# Multiplicity fluctuations



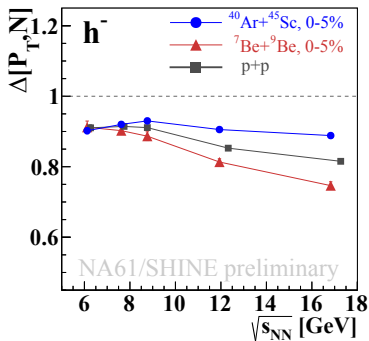
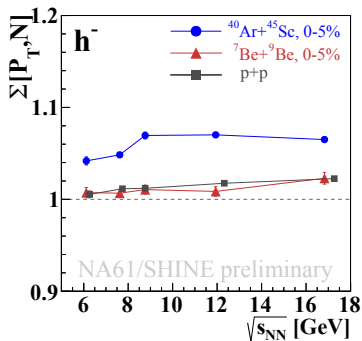
$$\omega[N] = \frac{\text{Var}[N]}{\langle N \rangle}$$

Be+Be similar to p+p, Ar+Sc different → onset of fireball (?).

No collision energy dependence that could be related to the critical point observed in Ar+Sc

[NA61/SHINE, PoS CPOD2017 (2018) 012]  
 [Andronov, Kuich, Gazdzicki, Universe 9 (2023) 2, 106]

# Multiplicity-transverse momentum fluctuations



$$\Sigma[A, B] = \frac{1}{C_\Sigma} \left[ \langle B \rangle \omega[A] + \langle A \rangle \omega[B] - 2 \left( \langle AB \rangle - \langle A \rangle \langle B \rangle \right) \right]$$

$$P_T = \sum_{i=1}^N p_{T_i}, \quad C_\Delta = C_\Sigma = \langle N \rangle \omega[p_T]$$

$$\Delta[A, B] = \frac{1}{C_\Delta} \left[ \langle B \rangle \omega[A] - \langle A \rangle \omega[B] \right]$$

Be+Be similar to p+p, Ar+Sc different → onset of fireball (?).

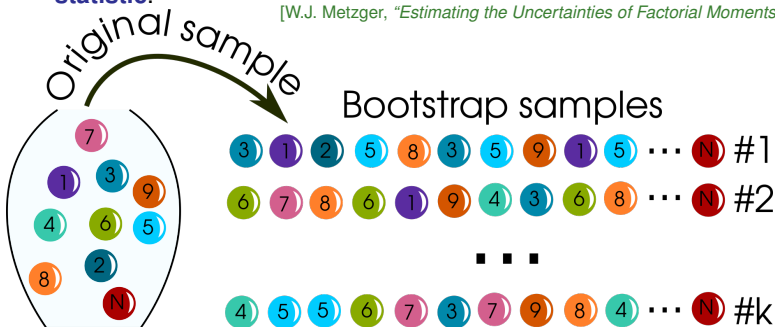
No collision energy dependence that could be related to the critical point observed in Ar+Sc

# Intermittency analysis tools: the bootstrap

- Random **sampling** of events, **with replacement**, from the original set of events;
- $k$  bootstrap samples ( $k \sim 1000$ ) of the **same number of events** as the original sample;
- Each **statistic** ( $\Delta F_2(M)$ ,  $\phi_2$ ) **calculated for bootstrap** samples as for the **original**; [B. Efron, *The Annals of Statistics* 7,1 (1979)]
- **Variance of bootstrap values** estimates **standard error of statistic**.



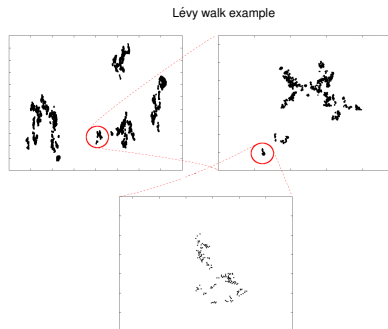
[W.J. Metzger, "Estimating the Uncertainties of Factorial Moments", HEN-455 (2004).]



# Critical Monte Carlo (CMC) algorithm for baryons

- Simplified version of CMC\* code:

- Only **protons** produced;
- **One cluster** per event, produced by sampling random Lévy walk of **adjustable dimension  $d_F$** , e.g.:  
 $d_F^B = 1/3 \Rightarrow \phi_2 = 1 - d_F^B/2 = 5/6$
- **Lower / upper bounds** of Lévy walks  $p_{min,max}$  plugged in;
- Cluster center **adjustable** to **experimental set mean proton  $p_T$**  per event;
- **Poissonian** proton multiplicity distribution.



## Input parameters (example)

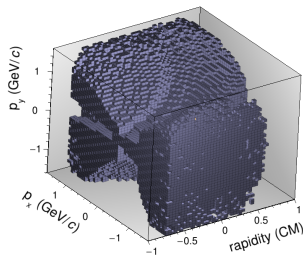
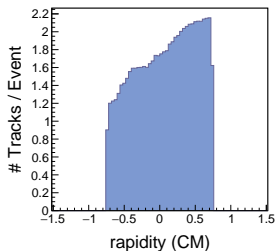
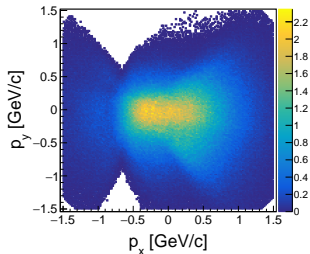
Parameter	$p_{\min}$ (MeV)	$p_{\max}$ (MeV)	$\lambda_{\text{Poisson}}$
Value	0.1 $\rightarrow$ 1	800 $\rightarrow$ 1200	$\langle p \rangle_{\text{non-empty}}$

\*[Antoniou, Diakonou, Kapoyannis and Kousouris, Phys. Rev. Lett. 97, 032002 (2006).]



# CMC – background simulation & detector effects

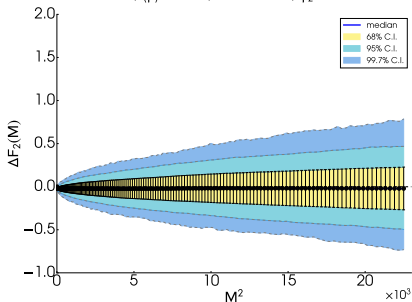
- **Non-critical background simulation: replace critical tracks by uncorrelated (random) tracks, with fixed probability:  $\mathcal{P}_{track} = 1 - \mathcal{P}_{crit}$ , where  $\mathcal{P}_{crit}$  is the percentage of critical component;**
- **$p_T$  distribution of background tracks plugged in to match experimental data;**
- **$y_{CM}$  rapidity value generated orthogonal to  $p_T$ , matching experimental distribution;**
- **$p_T, y_{CM},$  quality & acceptance cuts applied, same as in NA61/SHINE data;**



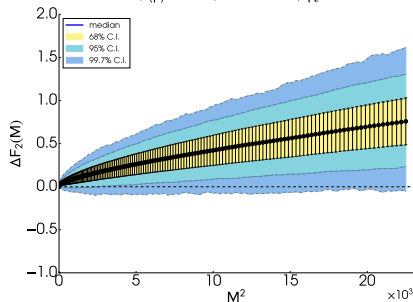
# CMC scan $\Delta F_2(M)$ – examples

- Results shown for **CMC  $\Delta F_2(M)$** , with  $\langle p \rangle = 2.562$ , corresponding to **NA61/SHINE Ar+Sc @ 150A GeV/c, cent.10-20%**;
- 2 settings:**
  - $\phi_2 = 0.125$ , crit.% = 1.60%;
  - $\phi_2 = 0.750$ , crit.% = 1.60%;
- For each setting,  **$\sim 8K$  independent samples** of  **$\sim 400K$  events** are generated; event statistics selected to **match NA61/SHINE data**.

CMC,  $\langle p \rangle = 2.562$ , crit.=1.60%,  $\phi_2 = 0.125$



CMC,  $\langle p \rangle = 2.562$ , crit.=1.60%,  $\phi_2 = 0.75$



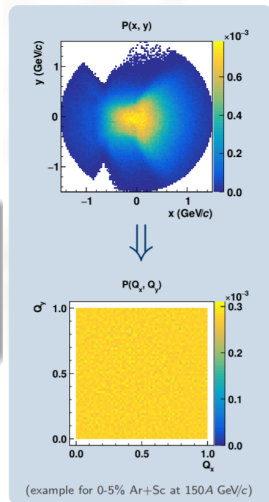
# Independent bin analysis with cumulative variables

- **M-bin correlations** complicate uncertainties estimations for  $\Delta F_2(M)$  &  $\phi_2$ ; one way around this problem is to use **independent bins** – a **different subset** of events is used to calculate  $F_2(M)$  for **each M**;
- **Advantage:** correlations are no longer a problem; **Disadvantage:** we **break up statistics**, and can only calculate  $F_2(M)$  for a **handful of bins**.
- Furthermore, instead of  $p_x$  and  $p_y$ , one can use **cumulative quantities**: [Bialas, Gazdzicki, PLB 252 (1990) 483]

$$Q_x(x) = \int_{min}^x P(x) dx \Bigg| \int_{min}^{max} P(x) dx;$$

$$Q_y(x, y) = \int_{ymin}^y P(x, y) dy \Bigg| P(x)$$

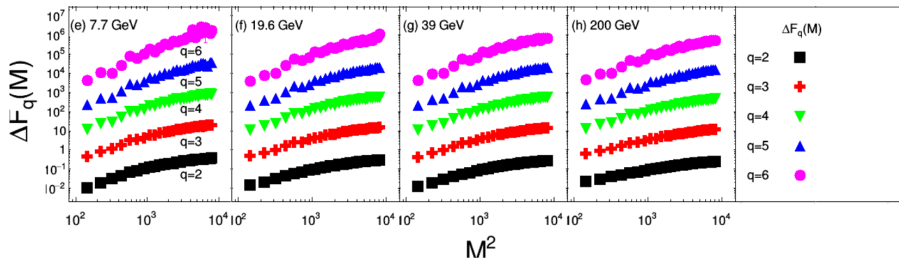
- transform any distribution into **uniform** one (0, 1);
- **remove the dependence** of  $F_2$  on the shape of the **single-particle distribution**;
- approximately **preserves ideal power-law** correlation function. [Antoniou, Diakonou, <https://indico.cern.ch/event/818624/>]



# STAR $h^\pm$ intermittency analysis

- In March 2023, the **STAR collaboration** published intermittency results of  $\Delta F_q$  of charged hadrons in **0-5% Au+Au collisions** at four example energies;

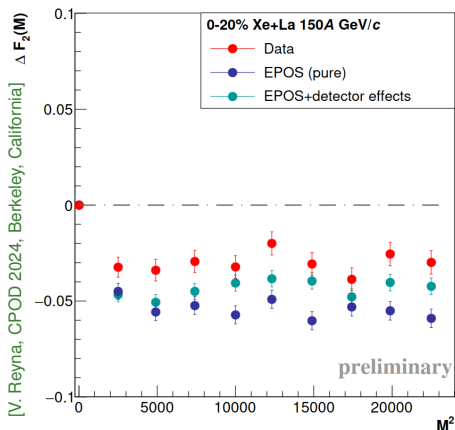
[STAR collaboration, Phys.Lett.B 845 (2023)]



- Plots:  $\Delta F_q(M) = F_q^{\text{data}}(M) - F_q^{\text{mixed}}(M)$  ( $q = 2 - 6$ ), in double-logarithmic scale;
- STAR reported that  $\Delta F_q(M)$  **increases** with  $M^2$  and **saturates** when  $M^2$  is larger than  $M^2 > 4000$ ;
- Interpretation of the **source** of this increase was **unclear**; **no specific theoretical prediction** is given for  $h^\pm$  **critical scaling**.

# SHINE Xe + La negatively charged hadrons intermittency

- Intermittency analysis performed on **negatively charged hadrons ( $h^-$ )** in **SHINE Xe + La collisions @ 150A GeV/c**; motivated by **corresponding STAR analysis**; [STAR collaboration, Phys.Lett.B 845 (2023)]



- Results after **cumulative transform** and **short-range correlation  $\Delta p_T$  cut** ( $\Delta p_T < 100 \text{ MeV}/c$  removed) **do not show any signal** indicating the **critical point**;
- Could the results of **STAR** (reported **increase of  $\Delta F_2$  with  $M$** ) also be interpreted as due to **short-range correlations**?