

CONNECTING SCIENCES

Inflation

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$\mathsf{KBFl}\bullet \underline{\mathsf{History}}$ of the Universe \bullet

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European Space Apency

The cosmic microwave background (CMB) is a radiation bathing the Earth from all directions, with the form of a black-body radiation with $T \approx 2.725K$. Furthermore, the T coming from different parts of the sky is astonishingly uniform.

Figure: The COBE (1989-1993) measurement of the CMB anisotropy. The oval is a map of the sky showing the dipole anisotropy $\Delta\, T/T \sim 10^{-3}.$

$\frac{1}{2}$ KBFI • Comoving coordinates •

CMB uniformity

CMB anisotropy

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CMB uniformity \rightarrow Cosmological Principle

"The Universe, at large scales, is homogeneous and isotropic"

CMB anisotropy

● Comoving coordinates ●

CMB uniformity \rightarrow Cosmological Principle

"The Universe, at large scales, is homogeneous and isotropic"

Think about a grid which expands with time: galaxies are fixed in the comoving system, but they move in the physical one.

The evolution of the Universe is described by the Friedmann eqs.

$$
\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} \qquad M_P = (8\pi G)^{-1/2}, \quad c = 1
$$

$$
\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho + 3\rho)
$$
 $\left(k < 0 \text{ open}\right)$

 \bullet k describes the geometry of the Universe $\left\{\n \begin{array}{l}\n k < 0 \quad \text{ope} \\
k = 0 \quad \text{flat} \\
k > 0 \quad \text{clos}\n\end{array}\n\right.$ $k > 0$ closed

- \bullet ρ : energy density of the Universe
- \bullet p: pressure of the Universe

• critical density:
$$
\rho_c = \frac{3(\dot{a}/a)^2}{8\pi G} \Rightarrow k = 0
$$
 $\rho_c(t)!!!!$

Now is $\rho_c(t_0) = 1.88 h^2 \times 10^{-26}$ kg m⁻³, h = 0.674 ± 0.005

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• density parameter: $\Omega = \frac{\rho}{\rho}$ ρ_c $\Omega = 1 \Leftrightarrow \rho = \rho_c$ • observations: $|\Omega - 1| \ll 1$ in different eras \Rightarrow simplest but unnatural choice $\rightarrow \Omega = 1 \Rightarrow k = 0$ • theory: the gravitational pull of matter should slow down the expansion of the Universe \Rightarrow $\dot{a}(t)$ is a decreasing function of time Friedmann eq. $\Rightarrow |\Omega - 1| = \frac{|k|}{\lambda^2}$ \dot{a}^2

 $\Rightarrow |\Omega - 1|$ is an increasing function of time

- \Rightarrow Ω will inevitably depart from 1
- \Rightarrow flatness problem

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$$

- $\Rightarrow |\Omega 1|$ is an increasing function of time
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- ⇒ flatness problem

inflation: $\left| \vec{a} \right| > 0$

● Inflation ●

Looking at the acceleration equation

$$
\frac{\ddot{a}}{a}=-\frac{4\pi G}{3}(\rho+3\rho)
$$

we see immediately that inflation implies $\rho + 3p < 0$. Since we always assume $\rho > 0$, this requires

Alan Guth proposed a solution (PRD 23(1981) 347-356):

$$
p<-\frac{\rho}{3}
$$

How can we get a $p < 0$? 1st idea: **Cosmological Constant (CC)**

● Cosmological constant ●

CC is a fluid with a constant energy density: Λ $8\pi G$

Performing all the computations of the case:

- negative pressure! $p_\Lambda = -\rho_\Lambda$
- $a(t)$ increases exponentially with time

$$
a(t) = a_{t=0} \exp\left(\sqrt{\frac{\Lambda}{3}}t\right) \Rightarrow |\Omega - 1| \propto \exp\left(-\sqrt{\frac{4\Lambda}{3}}t\right)
$$

 $\Omega \rightarrow 1$ very fast. We want to get so close that what happens after inflation cannot move it away again.

How much inflation is needed? The result is expressed in N_e

$$
\frac{a(t_{\text{final}})}{a(t_{\text{initial}})} = \exp[N_e]
$$

More accurate model building will provide $N_e \in [50, 60]$

- Inflation via a CC never stops \rightarrow PROBLEM
- solution \rightarrow a temporary CQ
- a scalar particle (ϕ , inflaton): $\mathcal{L}_m = \mathcal{T} \mathcal{V}$

- kinetic energy density : $T = \frac{1}{2}\dot{\phi}^2$
- potential energy density : $V(\phi)$
- If $T \ll V$, then $\rho \simeq V \simeq -p \approx CC$ behaviour

Like a ball slowly rolling down the slope from the top of a hill

Enough to know $V(\phi) \rightarrow$ Potential Slow Roll Parameters (PSRP)

$$
\epsilon_{V}(\phi) = \frac{M_{P}^{2}}{2} \left(\frac{V_{,\phi}(\phi)}{V(\phi)} \right)^{2}
$$

$$
\eta_{V}(\phi) = M_{P}^{2} \frac{V_{,\phi\phi}(\phi)}{V(\phi)}
$$

Potential Slow Roll Approximation (PSRA): $\epsilon \sqrt{|\eta|} \ll 1$

• end of inflation:
$$
\epsilon_V(\phi_{end}) = 1
$$

• number of e-folds:
$$
N_e \simeq \frac{1}{M_P^2} \int_{\phi_{\text{end}}}^{\phi_N} \frac{V}{V_{,\phi}} d\phi
$$

where
$$
\left\{\begin{array}{c} \phi_N \text{ is taken at the beginning of} \\ \phi_{\text{end}} \text{ is taken at the end of} \right.
$$

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tiny fluctuations in the CMB \Leftarrow pertubations via quantum effects

 $\phi(t, \mathbf{x}) = \overline{\phi}(t) + \delta\phi(t, \mathbf{x}), \qquad g_{\mu\nu}(t, \mathbf{x}) = \overline{g}_{\mu\nu}(t) + \delta g_{\mu\nu}(t, \mathbf{x})$

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$\frac{1}{2}$ KBFI • SR & relevant parameters •

- All observables can expressed in terms of $V(\phi)$, ϵ_v and η_v amplitude of P_S : $A_s = \frac{1}{24s}$ $24\pi^2$ V ∗ M_P^4 1 ϵ_v^* \rightarrow scale of inflation scalar spectral index: $\left| n_{\rm s} = 1 + 2 \eta_{\rm v}^{\star} - 6 \epsilon_{\rm v}^{\star} \right| \rightarrow$ tilt of $V(\phi)$ tensor-to-scalar ratio: $r = \frac{A_T}{A}$ $\frac{AT}{A_s} = 16\epsilon_v^* \rightarrow$ quantum gravity where * stands for evaluated at ϕ = $\overline{\phi}_N$ • we can take any $V(\phi)$ and start computing
	- From the Planck 2018 constraints (arXiv:1807.06211)

 $\ln (10^{10} A_s^{\text{exp}}) = 3.044 \pm 0.014$ ⇒ $A_s^{\text{exp}} \approx 2.1 \times 10^{-9}$

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- minimal Higgs inflation: 0.26 ≲ r ≲ 0.31 is ruled out
- concave potentials are strongly FAVORED!!!
- $V(\phi_N) = \frac{3\pi^2 A_s}{2} r M_P^4 \lesssim (1.6 \times 10^{16} \text{GeV})^4$ \rightarrow inflation a GUT scale phenomenon or just accident?
- non-minimal Higgs inflation still in the game \rightarrow C. Dioguardi's talk tomorrow

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EXAL • Summary & Conclusions •

- inflation is a stage of accelerated expansion of the Universe
- it solves the flatness problem of the Universe
- a scalar particle (inflaton) can drive inflation
- the Higgs boson (in non-minimal setups) can be the inflaton

● Summary & Conclusions ●

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• Expansion of the Universe •

- $H_0 = (67.4 \pm 0.5)$ km/s/Mpc is known as Hubble's constant
- not exact, but average behaviour
- everything is flying away from everything else
	- \rightarrow reverse time \rightarrow initial singularity: Big Bang

● Geometry of the Universe ●

Friedmann equation: $\left(\frac{\dot{a}}{a}\right)$ \overline{a} 2 = 8π G $\frac{\pi G}{3} \rho - \frac{k}{a^2}$ $a²$

 \bullet k describes the geometry of the Universe:

 $k = 0$ Euclidean (aka flat) geometry $k > 0$ spherical (aka close) geometry $k < 0$ hyperbolic (aka open) geometry

Euclidean geometry is based on a set of simple axioms e.g.:

- a straight line is the shortest distance between 2 points
- parallel straight lines remain a fixed distance apart

which lead to the following conclusions:

- The angles of a triangle add up to 180° .
- The circumference of a circle of radius r is $2\pi r$

Such a geometry might well apply to our own Universe \Rightarrow the Universe must be infinite in extent, otherwise the edges would clearly violate homogeneity and isotropy. A Universe with this geometry is often called a flat Universe.

● Spherical geometry ●

• circumference of a circle is less than $2\pi r$

A Universe with $k > 0$ is also called **closed**, because of its finite size.

N.B. The picture shows a 2D spherical surface. Our Universe would eventually be a 3D spherical surface!

Figure 4.2 A sketch of a saddle surface, representing the hyperbolic geometry obtained when k is negative. A rather exaggerated triangle is shown with its sum of angles well below 180 $^{\circ}$.

Main properties

- The angles of a triangle add up to less than 180°
- circumference of a circle is more than $2\pi r$

A Universe with $k < 0$ is also called **open**, because of its finite size.

N.B. The picture shows a 2D hyperbolic surface. Our Universe would eventually be a 3D hyperbolic surface!

