

331 models

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Number of fermion families

The SM has three generations of fermions

$$\begin{pmatrix} u \\ d \end{pmatrix}_L, \quad \begin{pmatrix} c \\ s \end{pmatrix}_L, \quad \begin{pmatrix} t \\ b \end{pmatrix}_L, \quad \begin{pmatrix} t' \\ b' \end{pmatrix}_L \cdots \cdots \\
 \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L, \quad \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L, \quad \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L, \quad \begin{pmatrix} \nu_4 \\ l_4 \end{pmatrix}_L \cdots \cdots$$

SM does not offer any explanation why there is only three generations. **In principle there could be more.**

⇒ 331-models offer explanation to the number of fermion generations.

331-gauge extension

331-models replace the SM gauge group $SU(3)_c \times SU(2)_L \times U(1)_Y$ with $SU(3)_c \times SU(3)_L \times U(1)_X$

\Rightarrow 5 more generators = 5 more gauge bosons

\Rightarrow Scalar sector has to be extended in order to give masses to the new gauge bosons

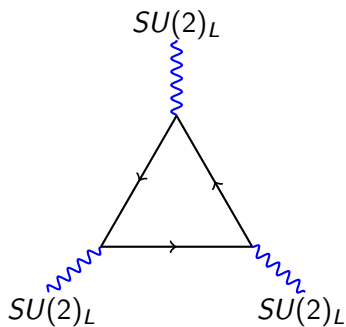
Fundamental representation is now a $SU(3)_L$ -triplet for left-handed fermions. \Rightarrow Fermion sector needs to be extended (U, D_1, D_2, ν'_i)

$$Q_{L,i} = \begin{pmatrix} u_i \\ d_i \\ q_{\text{exotic}} \end{pmatrix}_L, \quad L_{L,i} = \begin{pmatrix} \nu_i \\ e_i \\ l_{\text{exotic}} \end{pmatrix}_L.$$

Anomaly cancellation in 331-model

Standard model = $SU(3)_C \times SU(2)_L \times U(1)_Y$

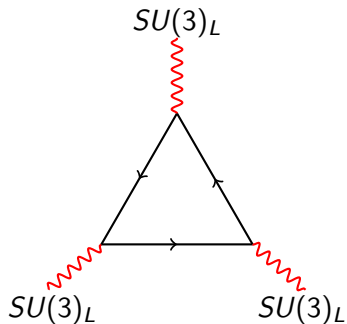
331 - model = $SU(3)_C \times SU(3)_L \times U(1)_X$



$SU(2)_L: \left\{ \frac{\sigma^a}{2}, \frac{\sigma^b}{2} \right\} = \frac{\delta^{ab}}{2} \Rightarrow$ Anomaly cancels automatically

$SU(3)_L: \left\{ \frac{T^a}{2}, \frac{T^b}{2} \right\} \neq \delta^{ab} \Rightarrow$ Anomaly doesn't cancel automatically!

In 331 number of families = 3



Pure $SU(3)_L$ -anomaly cancels only if the number of families is 3

Family discrimination

Pure $SU(3)_L$ anomaly cancels only if number of fermion triplets equals the number of antitriplets. [How to assign them?](#)

6 triplets (quarks have 3 colours):

$$L_{L,i} = \begin{pmatrix} \nu_i \\ e_i \\ \nu'_i \end{pmatrix}_L, \quad Q_{L,1} = \begin{pmatrix} u_1 \\ d_1 \\ U \end{pmatrix}_L$$

6 antitriplets (quarks have 3 colours):

$$Q_{L,2} = \begin{pmatrix} d_2 \\ -u_2 \\ D_1 \end{pmatrix}_L, \quad Q_{L,3} = \begin{pmatrix} d_3 \\ -u_3 \\ D_2 \end{pmatrix}_L$$

One of the quarks is in different representation!

Family discrimination

Pure $SU(3)_L$ anomaly cancels only if number of fermion triplets equals the number of antitriplets. [How to assign them?](#)

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$$L_{L,i} = \begin{pmatrix} \nu_i \\ e_i \\ \nu'_i \end{pmatrix}_L, \quad Q_{L,2} = \begin{pmatrix} u_2 \\ d_2 \\ U \end{pmatrix}_L$$

6 antitriplets (quarks have 3 colours):

$$Q_{L,1} = \begin{pmatrix} d_1 \\ -u_1 \\ D_1 \end{pmatrix}_L, \quad Q_{L,3} = \begin{pmatrix} d_3 \\ -u_3 \\ D_2 \end{pmatrix}_L$$

But which generation to pick?

Family discrimination

Pure $SU(3)_L$ anomaly cancels only if number of fermion triplets equals the number of antitriplets. [How to assign them?](#)

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$$L_{L,i} = \begin{pmatrix} \nu_i \\ e_i \\ \nu'_i \end{pmatrix}_L, \quad Q_{L,3} = \begin{pmatrix} u_3 \\ d_3 \\ U \end{pmatrix}_L$$

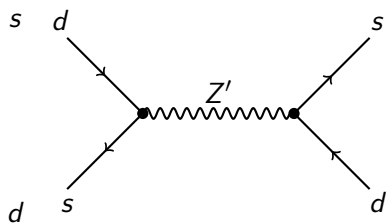
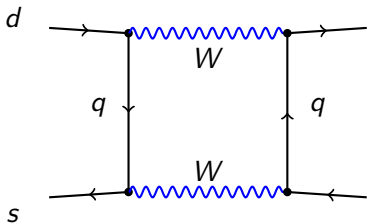
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No one knows!

Flavour changing neutral currents

- In order to cancel the anomalies, one quark family has to be in a different representation than the other two. (And no one knows which!)
- \Rightarrow Scalar mediated flavour changing neutral currents (FCNC) at tree-level!
- \Rightarrow Neutral meson mixing at tree-level!
- Different quark generation assignments give different flavour violating couplings \Rightarrow Can be tested \Rightarrow New physics scale ~ 10 TeV



331-model

The $SU(3)_L$ -gauge group has one additional diagonal generator compared to the $SU(2)_L$:

⇒ Freedom in electric charge definition: $Q = T_3 + \beta T_8 + X$ (compare to SM: $Q = T_3 + Y/2$).

Two types of models:

- $\beta = \pm\sqrt{3}$:

F. Pisano and V. Pleitez, Phys. Rev. D **46**, 410 (1992).

P. H. Frampton, Phys. Rev. Lett. **69**, 2889 (1992).

- $\beta = \pm\frac{1}{\sqrt{3}}$:

M. Singer, J. W. F. Valle and J. Schechter, Phys. Rev. D **22**, 738 (1980).

$$\beta = \begin{pmatrix} + \\ - \end{pmatrix} \frac{1}{\sqrt{3}}$$

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Fermion representations:

$$\text{Triplets: } L_{L,i} = \begin{pmatrix} \nu_i \\ e_i \\ N_i \end{pmatrix}_L \sim (1, 3, -\frac{1}{3}), \quad Q_{L,1} = \begin{pmatrix} u \\ d \\ U \end{pmatrix}_L \sim (3, 3, \frac{1}{3})$$

$$\text{Antitriplets: } Q_{L,2} = \begin{pmatrix} s \\ c \\ D_1 \end{pmatrix}_L \sim (3, 3^*, 0), \quad Q_{L,3} = \begin{pmatrix} b \\ t \\ D_2 \end{pmatrix}_L \sim (3, 3^*, 0)$$

$N_{L,i}$ = new neutrino-like fields, U = new up-type quark, D_1, D_2 = new down-type quarks.

Minimal Scalar Sector:

$$\eta = \begin{pmatrix} \eta^+ \\ \eta^0 \\ \eta'^+ \end{pmatrix} \sim (1, 3, \frac{2}{3}), \quad \rho = \begin{pmatrix} \rho^0 \\ \rho^- \\ \rho'^0 \end{pmatrix}, \quad \chi = \begin{pmatrix} \chi^0 \\ \chi^- \\ \chi'^0 \end{pmatrix} \sim (1, 3, -\frac{1}{3})$$

$$\beta = (+)\sqrt{3}$$

F. Pisano and V. Pleitez, Phys. Rev. D **46**, 410 (1992).

P. H. Frampton, Phys. Rev. Lett. **69**, 2889 (1992).

$$\text{Triplets: } L_{L,i} = \begin{pmatrix} \nu_{L,i} \\ e_{L,i} \\ (e_{R,i})^c \end{pmatrix} \sim (1, 3, 0), \quad Q_{L,1} = \begin{pmatrix} u \\ d \\ q_1^{+5/3} \end{pmatrix} \sim (1, 3, \frac{2}{3}),$$

$$\text{Antitriplets: } Q_{L,2} = \begin{pmatrix} c \\ s \\ q_2^{-4/3} \end{pmatrix}, \quad Q_{L,3} = \begin{pmatrix} t \\ b \\ q_3^{-4/3} \end{pmatrix} \sim (1, 3^*, -\frac{1}{3}).$$

$q_1^{+5/3}$, $q_2^{-4/3}$ and $q_3^{-4/3}$ are new quarks with exotic electric charges!

$\beta = (+_-)\sqrt{3}$, scalar sector

$$\mathcal{L}_m = \epsilon_{\alpha\beta\gamma} G_{ij} \bar{L}_{L,i}^\alpha (L_{L,j}^\beta)^c \langle \eta^* \rangle^\gamma = \left(\epsilon_{\alpha\beta\gamma} G_{ij} \langle \eta^{0*} \rangle \right) \bar{e}_{L,i} e_{R,j}$$

⇒ Antisymmetric charged lepton mass matrix! Need to add scalar sextet.

Minimal Scalar Sector:

$$\eta = \begin{pmatrix} \eta^0 \\ \eta^- \\ \eta'^+ \end{pmatrix} \sim (1, 3, 0), \quad \rho = \begin{pmatrix} \rho^+ \\ \rho^0 \\ \rho^{++} \end{pmatrix} \sim (1, 3, 1),$$

$$\chi = \begin{pmatrix} \chi^- \\ \chi^{--} \\ \chi^0 \end{pmatrix} \sim (1, 3, -1), \quad S = \begin{pmatrix} \sigma_1^0 & h_2^- & h_1^+ \\ h_2^- & H_1^{--} & \sigma_2^0 \\ h_1^+ & \sigma_2^0 & H_2^{++} \end{pmatrix} \sim (1, 6, 0).$$

$\beta = (+)\sqrt{3}$, Gauge bosons

$$D_\mu = \partial_\mu - ig_3 \sum_{a=1}^8 T_a W_{a\mu} - ig_x X B_\mu,$$

$$\sum_{a=1}^8 T_a W_{a\mu} = \frac{1}{\sqrt{2}} \begin{pmatrix} \frac{1}{\sqrt{2}} W_{3\mu} + \frac{1}{\sqrt{6}} W_{8\mu} & W'_\mu{}^+ & V'_\mu{}^- \\ W'_\mu{}^- & -\frac{1}{\sqrt{2}} W_{3\mu} + \frac{1}{\sqrt{6}} W_{8\mu} & U_\mu{}^{--} \\ V'_\mu{}^+ & U_\mu{}^{++} & -\frac{2}{\sqrt{6}} W_{8\mu} \end{pmatrix},$$

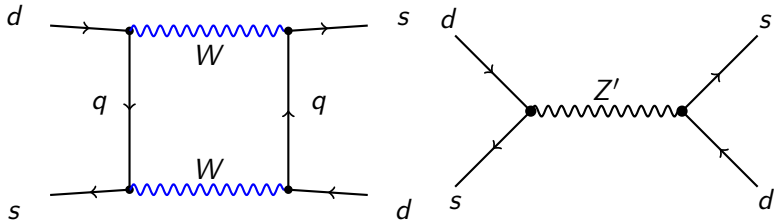
The neutral mass eigenstates are: **photon**, Z_μ and Z'_μ . The charged mass eigenstates are W_μ^\pm , V_μ^\pm and $U^{\pm\pm}$.

Summary

- In 331 models anomaly cancellation fixes number of generations to be 3
- The new physics scale is $\gtrsim 10$ TeV
- Probable at the LHC

Neutral meson mixing

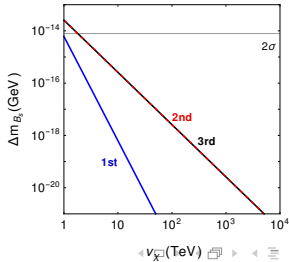
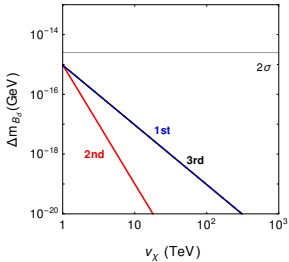
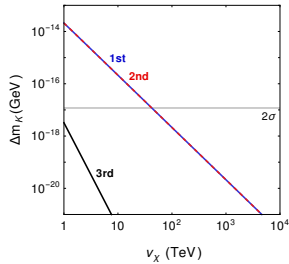
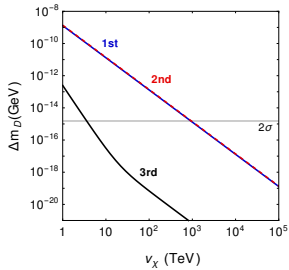
Different quark generation assignments give different flavour violating couplings \Rightarrow can be tested



Neutral mesons:

$$\begin{array}{ll}
 D^0 = c\bar{u} & \Delta m_D = (6.56 \pm 0.763) \times 10^{-15} \text{ GeV} \\
 K^0 = d\bar{s} & \Delta m_K = (3.48 \pm 0.00592) \times 10^{-15} \text{ GeV} \\
 B_d^0 = d\bar{b} & \Delta m_{B_d} = (3.33 \pm 0.0125) \times 10^{-13} \text{ GeV} \\
 B_s^0 = s\bar{b} & \Delta m_{B_s} = (1.17 \pm 0.000395) \times 10^{-11} \text{ GeV}
 \end{array}$$

Plot



Benchmarks for LHC

	BP1	BP2	BP3
m_U/TeV	1.362	2.368	1.829
m_{D1}/TeV	2.150	6.931	9.076
m_{D2}/TeV	16.007	36.369	26.922

How to make predictions?

- The different quark generation assignments lead to different flavour violating structure \Rightarrow Can be used to test which generation is discriminated
- One needs to make assumptions about the structure of quark mass matrices
- We assume that CKM, $V_{\text{CKM}} = U_L^u U_L^{d\dagger}$, is generated without accidental cancellations between elements of U_L^u and $U_L^{d\dagger}$. This is achieved when ($\epsilon = 0.23$, Cabibbo angle)

$$V_{\text{CKM}} \sim \begin{pmatrix} 1 & \epsilon & \epsilon^3 \\ \epsilon & 1 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix}, \quad U_L^u, U_L^d \sim \begin{pmatrix} 1 & \epsilon & \epsilon^3 \\ \epsilon & 1 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix}$$

$$m_u \sim \frac{v}{\sqrt{2}} \begin{pmatrix} \epsilon^{L_1+R_1} & \epsilon^{L_1+R_2} & \epsilon^{L_1+R_3} \\ \epsilon^{L_2+R_1} & \epsilon^{L_2+R_2} & \epsilon^{L_2+R_3} \\ \epsilon^{L_3+R_1} & \epsilon^{L_3+R_2} & \epsilon^{L_3+R_3} \end{pmatrix}, \quad \begin{matrix} L_1 = 3 \\ L_2 = 2 \\ L_3 = 0 \end{matrix} \Rightarrow \text{CKM}$$

331 mass matrix

- In 331 exotic quarks mix with SM quarks
- We take it into account and do not ignore it!
- But mixing is suppressed by $SU(3)_L$ -breaking scale u . $\delta = \frac{v_{\text{SM}}}{u}$

$$m_u = \frac{v_{\text{SM}}}{\sqrt{2}} \begin{pmatrix} y_{11}^u & y_{12}^u & y_{13}^u & y_{14}^u \\ y_{21}^u & y_{22}^u & y_{23}^u & y_{24}^u \\ y_{31}^u & y_{32}^u & y_{33}^u & y_{34}^u \\ \frac{u}{v_{\text{SM}}} y_{41}^u & \frac{u}{v_{\text{SM}}} y_{42}^u & \frac{u}{v_{\text{SM}}} y_{43}^u & \frac{u}{v_{\text{SM}}} y_{44}^u \end{pmatrix}$$

Write Yukawas in suggestive form:

$$y_{ij} = \epsilon^{L_i+R_j} \Rightarrow (U_L)_{ij} \sim \epsilon^{|L_i-L_j|}, (U_R)_{ij} \sim \epsilon^{|R_i-R_j|}$$

- $L_1 = 3, L_2 = 2, L_3 = 0$ gives CKM and $L_4 \equiv L_U$ affects mixing with exotic quarks
- R_i determine the right-handed rotation matrix (not relevant) $\Rightarrow R_i$ are chosen to produce correct SM masses

331 quark rotation matrices

$$U_L^u \sim \begin{pmatrix} 1 & \epsilon & \epsilon^3 & \delta\epsilon^{3-L_U} \\ \epsilon & 1 & \epsilon^2 & \delta\epsilon^{2-L_U} \\ \epsilon^3 & \epsilon^2 & 1 & \delta\epsilon^{-L_U} \\ \delta\epsilon^{3-L_U} & \delta\epsilon^{2-L_U} & \delta\epsilon^{-L_U} & 1 \end{pmatrix}, \quad \delta = \frac{v_{\text{sm}}}{u}$$

$$U_L^d \sim \begin{pmatrix} 1 & \epsilon & \epsilon^3 & \delta\epsilon^{3-L_{D_1}} & \delta\epsilon^{3-L_{D_2}} \\ \epsilon & 1 & \epsilon^2 & \delta\epsilon^{2-L_{D_1}} & \delta\epsilon^{2-L_{D_2}} \\ \epsilon^3 & \epsilon^2 & 1 & \delta\epsilon^{-L_{D_1}} & \delta\epsilon^{-L_{D_2}} \\ \delta\epsilon^{3-L_{D_1}} & \delta\epsilon^{2-L_{D_1}} & \delta\epsilon^{-L_{D_1}} & 1 & \epsilon^{L_{D_1}-L_{D_2}} \\ \delta\epsilon^{3-L_{D_2}} & \delta\epsilon^{2-L_{D_2}} & \delta\epsilon^{-L_{D_2}} & \delta\epsilon^{L_{D_1}-L_{D_2}} & 1 \end{pmatrix}$$

$$V_{\text{CKM}}^{331} = U_L^u \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix} U_L^{d\dagger} \quad (1)$$

- Quark rotation matrices determine the flavor violating couplings
- Each generation assignment gives different flavour violating couplings
 \Rightarrow difference can be seen in neutral meson mixing!

Example: $\mathcal{L} = \bar{u}_L U_L^u (\lambda_{Z'}^u) U_L^{u\dagger} \gamma^\mu u_L Z'_\mu + \bar{d}_L U_L^d (\lambda_{Z'}^d) U_L^{d\dagger} \gamma^\mu d_L Z'_\mu$

$$\lambda_{1st}^u = \begin{pmatrix} a+b & 0 & 0 & 0 \\ 0 & a & 0 & 0 \\ 0 & 0 & a & 0 \\ 0 & 0 & 0 & a+c \end{pmatrix}, \lambda_{2nd}^u = \begin{pmatrix} a & 0 & 0 & 0 \\ 0 & a+b & 0 & 0 \\ 0 & 0 & a & 0 \\ 0 & 0 & 0 & a+c \end{pmatrix}, \lambda_{3rd}^u = \begin{pmatrix} a & 0 & 0 & 0 \\ 0 & a & 0 & 0 \\ 0 & 0 & a+b & 0 \\ 0 & 0 & 0 & a+c \end{pmatrix}$$

Neutral mesons: $K^0 = d\bar{s}$, $B_d^0 = d\bar{b}$, $B_s^0 = s\bar{b}$, $D^0 = c\bar{u}$

$$\begin{array}{ll} D^0 = c\bar{u} & \Delta m_D = (6.56 \pm 0.763) \times 10^{-15} \text{ GeV} \\ K^0 = d\bar{s} & \Delta m_K = (3.48 \pm 0.00592) \times 10^{-15} \text{ GeV} \\ B_d^0 = d\bar{b} & \Delta m_{B_d} = (3.33 \pm 0.0125) \times 10^{-13} \text{ GeV} \\ B_s^0 = s\bar{b} & \Delta m_{B_s} = (1.17 \pm 0.000395) \times 10^{-11} \text{ GeV} \end{array}$$