



High and Low-Multiplicity pp-Collisions and Long-Range Angular Correlations

Report by G.Feofilov at LUHEP seminar May 14, 2024,
Petergof , 12:00—13:00

➤ Яркий сигнал процесса слияния кварк-глюонных струн?

Layout

- Emergence of Long-Range Angular Correlations in Low-Multiplicity Proton-Proton Collisions(PHYSICAL REVIEW LETTERS 132, 172302 (2024))—
кратко
- Percolation string model and old predictions for critical behaviour -- where to search for the critical point
- High and Low-Multiplicity pp-Collisions and Long-Range Angular Correlations
- Plans

Emergence of Long-Range Angular Correlations in Low-Multiplicity Proton-Proton Collisions

S. Acharya et al.*
(ALICE Collaboration)

This Letter presents the measurement of near-side associated per-trigger yields, denoted ridge yields, from the analysis of angular correlations of charged hadrons in proton-proton collisions at $\sqrt{s} = 13$ TeV. Long-range ridge yields are extracted for pairs of charged particles with a pseudorapidity difference of $1.4 < |\Delta\eta| < 1.8$ and a transverse momentum of $1 < p_T < 2$ GeV/ c , as a function of the charged-particle multiplicity measured at midrapidity. This Letter extends the measurements of the ridge yield to the low multiplicity region, where in hadronic collisions it is typically conjectured that a strongly interacting medium is unlikely to be formed. The precision of the new low multiplicity results allows for the first direct quantitative comparison with the results obtained in e^+e^- collisions at $\sqrt{s} = 91$ GeV and $\sqrt{s} = 183$ – 209 GeV, where initial-state effects such as preequilibrium dynamics and collision geometry are not expected to play a role. In the multiplicity range $8 \lesssim \langle N_{\text{ch}} \rangle \lesssim 24$ where the e^+e^- results have good precision, the measured ridge yields in pp collisions are substantially larger than the limits set in e^+e^- annihilations. Consequently, the findings presented in this Letter suggest that the processes involved in e^+e^- annihilations do not contribute significantly to the emergence of long-range correlations in pp collisions.

The observables presented in this analysis are extracted using two-particle angular correlations for pairs of charged particles. The two-particle per-trigger yield is measured as a function of relative azimuthal angle $\Delta\varphi$ and pseudorapidity $\Delta\eta$ of two particles—traditionally called trigger and associated—and is defined as

$$\frac{1}{N_{\text{trig}}} \frac{d^2 N_{\text{pair}}}{d\Delta\eta d\Delta\varphi} = B(0, 0) \frac{S(\Delta\eta, \Delta\varphi)}{B(\Delta\eta, \Delta\varphi)}. \quad (1)$$

$32 < N_{\text{ch}} \leq 37$
 $1 < p_{\text{T, trig/assoc}} < 2 \text{ GeV}/c$

ALICE
 $pp \sqrt{s} = 13 \text{ TeV}$

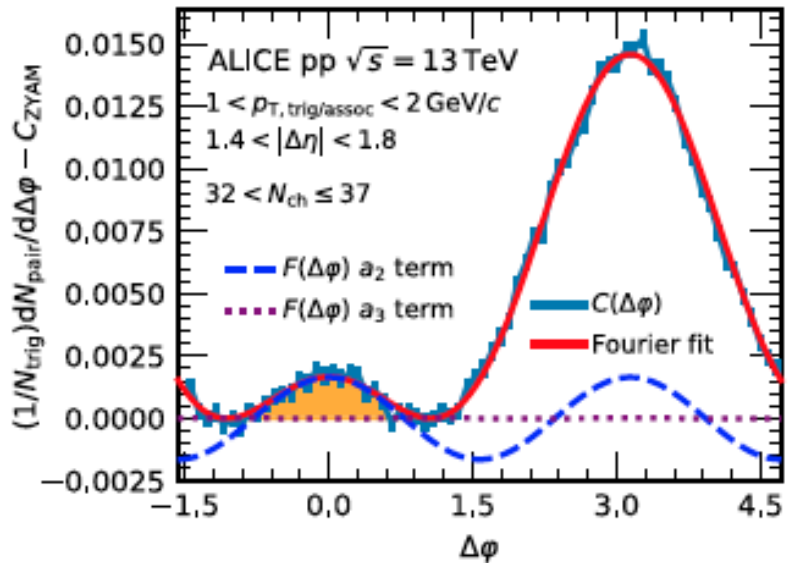
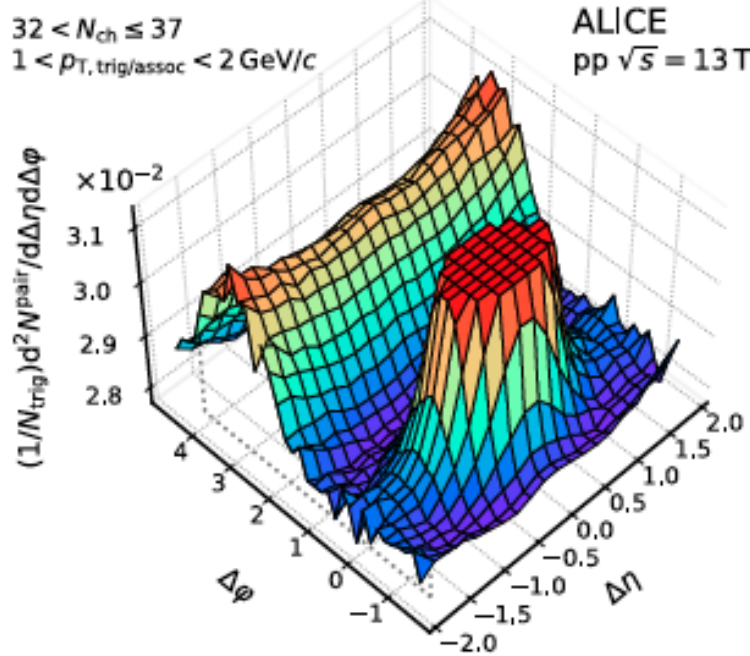


FIG. 1. Two-particle per-trigger yield measured for charged track pairs with $1 < p_{\text{T, trig}} < 2 \text{ GeV}/c$ and $1 < p_{\text{T, assoc}} < 2 \text{ GeV}/c$ within the multiplicity range $32 < N_{\text{ch}} \leq 37$. The jet fragmentation peak has been truncated to ensure a better visibility of the long-range structure. The right panel shows the zero-suppressed projection to $\Delta\phi$ overlaid with $F(\Delta\phi)$ (red line) and the area in which the ridge yield is extracted (shaded area). The blue and purple lines represent the second and third harmonic terms of $F(\Delta\phi)$.

The per-trigger yield distribution as a function of $\Delta\varphi$ is obtained by integrating the two-dimensional two-particle per-trigger yield in the long-range intervals $1.4 < |\Delta\eta| < 1.8$, in order to exclude the region dominated by the jet fragmentation peak

$$\begin{aligned}
 Y(\Delta\varphi) &= \frac{1}{N_{\text{trig}}} \frac{dN_{\text{pair}}}{d\Delta\varphi} \\
 &= \int_{1.4 < |\Delta\eta| < 1.8} \left(\frac{1}{N_{\text{trig}}} \frac{d^2N_{\text{pair}}}{d\Delta\eta d\Delta\varphi} \right) \frac{1}{\delta_{\Delta\eta}} d\Delta\eta, \quad (2)
 \end{aligned}$$

where $\delta_{\Delta\eta} = 0.8$ is the normalization constant for the chosen $\Delta\eta$ range. The ridge yield Y^{ridge} is extracted by integrating the near-side area of the associated per-trigger yield using

$$Y^{\text{ridge}} = \int_{|\Delta\varphi| < |\Delta\varphi_{\text{min}}|} Y(\Delta\varphi) d\Delta\varphi - 2|\Delta\varphi_{\text{min}}| C_{\text{ZYAM}}. \quad (3)$$

Ridge yield as a function of $\langle N_{ch} \rangle$

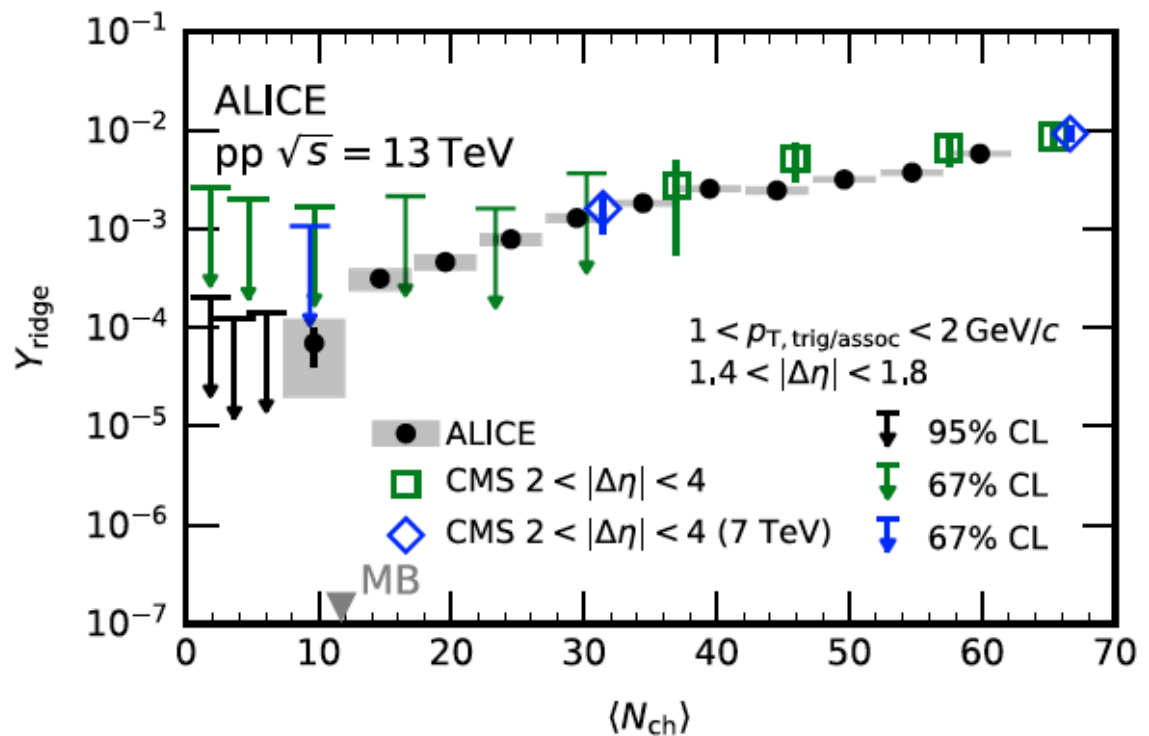


FIG. 2. Ridge yield as a function of multiplicity. The black points correspond to the measurement presented in this Letter, while data from CMS [8,38] are drawn as green and blue markers. Vertical bars denote statistical uncertainties while systematic uncertainty is shown as a shaded area. For both results, at low multiplicity where the lower uncertainty reaches zero, an upper limit is reported, which is drawn as a bar and down arrow. Such points are given at 95% CL for the results from this Letter and at 67% for the results from CMS. The “MB” arrow at $\langle N_{ch} \rangle = 11.3$ indicates the multiplicity averaged over the entire studied multiplicity range.

Comparison to e^+e^- collisions

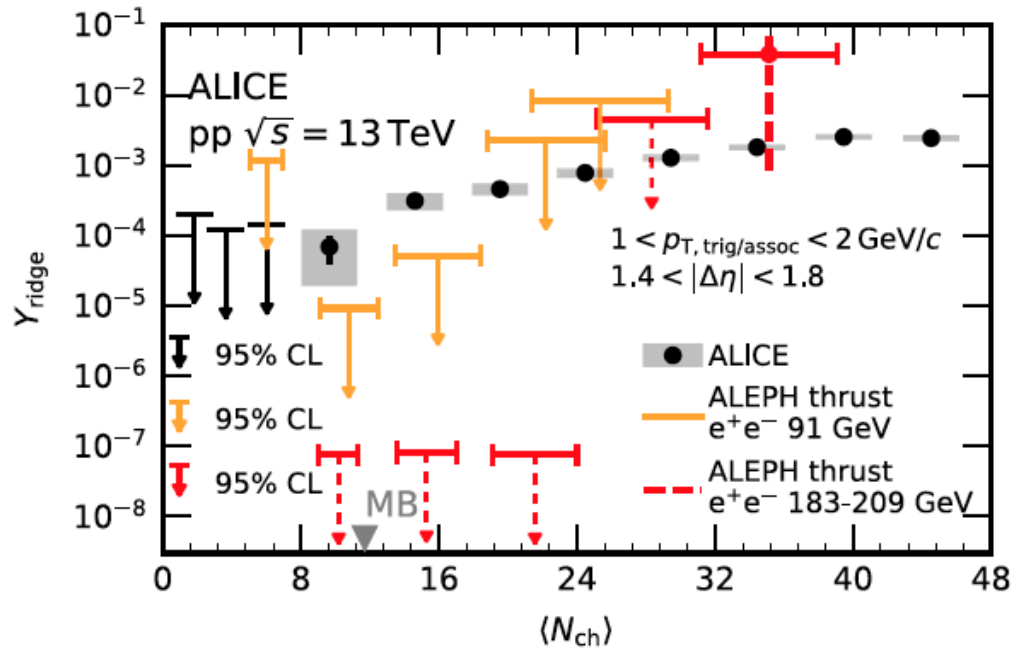


FIG. 3. Ridge yield as a function of multiplicity, compared to the upper limits on the ridge yield in e^+e^- collisions. Vertical bars denote statistical uncertainties while systematic uncertainty is shown as the shaded areas. The orange limits represent the measurement in the thrust-axis reference frame with ALEPH [34]. The horizontal bars in the ALEPH points represent the uncertainty related to the multiplicity conversion from the ALEPH to the ALICE acceptance (see text). All limits are given at 95% CL.

In the multiplicity range 8 to 18 (24) the yields in pp collisions are substantially above the ALEPH limit at $\sqrt{s} = 91$ ($\sqrt{s} = 183\text{--}209$ GeV) while outside this range the limits from e^+e^- collisions are above the pp measurement.

Similarly to the previous figure, in order to translate the ALEPH multiplicity into the ALICE acceptance range, a scaling factor is estimated with PYTHIA8.3 events by counting the resulting particles in the acceptance ranges of both experiments ($|\eta| < 1.738$, $p_T > 0.2$ GeV/ c in case of ALEPH).

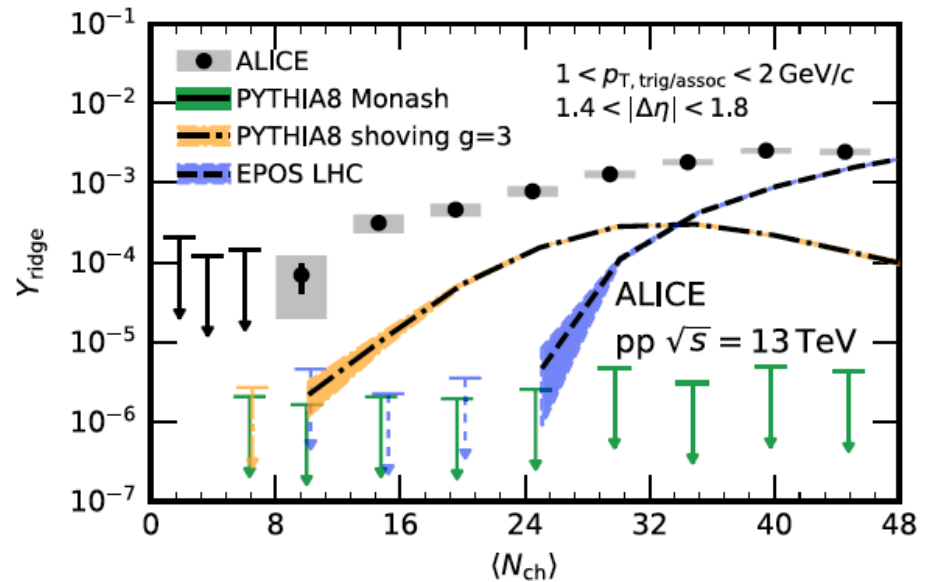
e^+e^-

In Fig. 3, the result is compared to a recent measurement performed in e^+e^- collisions at $\sqrt{s} = 91$ [34] and at $\sqrt{s} = 183\text{--}209$ GeV [35] in the thrust-axis reference frame using ALEPH archived data.

Because of the absence of beam remnants, the thrust axis provides an estimate of the longitudinal color field between the initially created outgoing $q\bar{q}$ pair and is therefore the sensible choice in e^+e^- collisions to search for collective effects.

Compare to MC

This suggests that the mechanisms for ridge yield production in very small hadronic collisions have not been understood and more theoretical work is needed.



The results presented in this Letter suggest that the ridge yield measured from a hadronic system of roughly equivalent multiplicity is nonzero and substantially larger than the limit observed in e^+e^- annihilations. Based on this, one can conclude that additional processes besides those in the e^+e^- annihilations must play a role for the emergence of long-range correlations in pp collisions.

- Очень странный вывод !
- Пытаются сопоставлять эффекты дальних корреляций для e^+e^- и pp , основываясь на такой наблюдаемой как множественность ???
- Некорректный подход??? Разные механизмы адронизации в процессах e^+e^- и pp

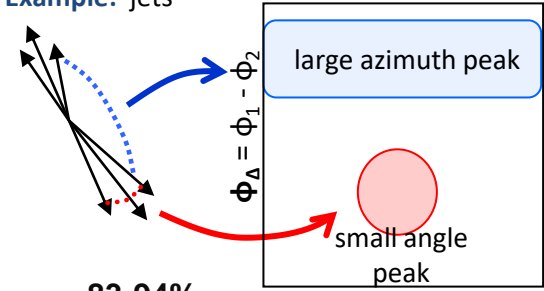
- Emergence of Long-Range Angular Correlations in Low-Multiplicity Proton-Proton Collisions(PHYSICAL REVIEW LETTERS 132, 172302 (2024))—кратко
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O.Kochebina and G.Feofilov, Onset of "ridge phenomenon" in AA and pp collisions and percolation string model arXiv:1012.0173v1

New phenomena in AA collisions observed by STAR [1]

Example: jets

Axial Autocorrelations



Variation of **low- p_T** “ridge” with centrality (Npart).
($p_T > 0.15 \text{ GeV}/c$)

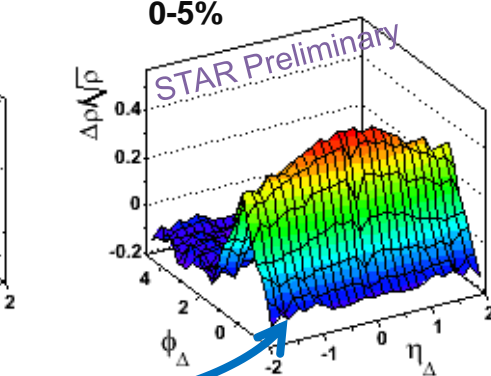
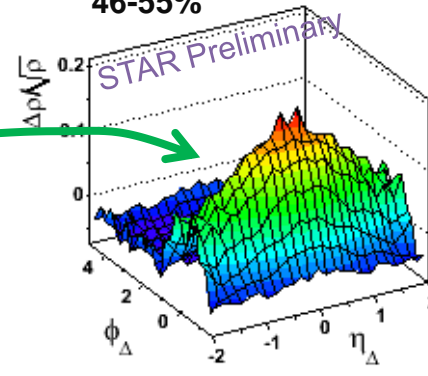
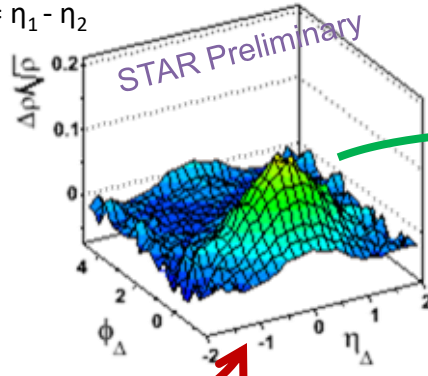
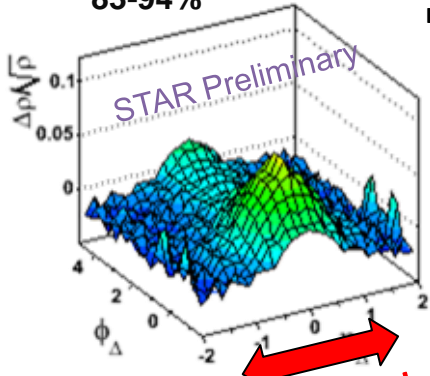
83-94%

$\eta_\Delta = \eta_1 - \eta_2$

55-65%

46-55%

0-5%



η_Δ width

Shape changes little from peripheral to the transition

Large change within $\sim 10\%$ centrality

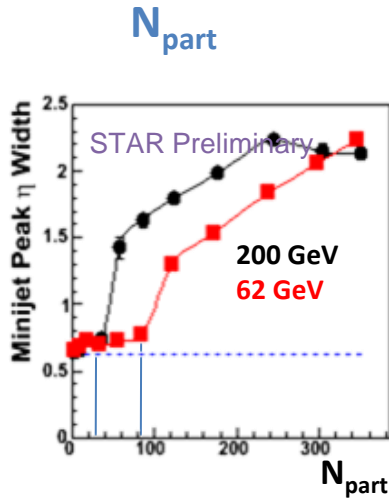
Smaller change from transition to most central

Low- p_T manifestation of the “ridge”

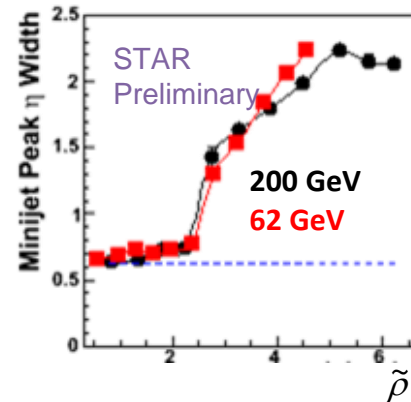
The data showed a sharp transition at some definite energy-dependent centrality: **growing of peak amplitude and stretching of width.**

Variation of low- p_T “ridge” with centrality (Npart).

A sudden increase of the **peak amplitude** and **η width** of *the near-side low p_T ridge* were found at an *energy-dependent* centrality point at definite number of participating nucleons **Npart**.



Transverse Particle Density



$$\tilde{\rho} = \frac{3}{2} \frac{dN_{ch}}{d\eta} / S$$

$$\tilde{\rho}_{crit} = 2,6 \pm 0,2 \text{ fm}^{-2}$$

62 GeV

“Critical value” $N_{part} \approx 90$

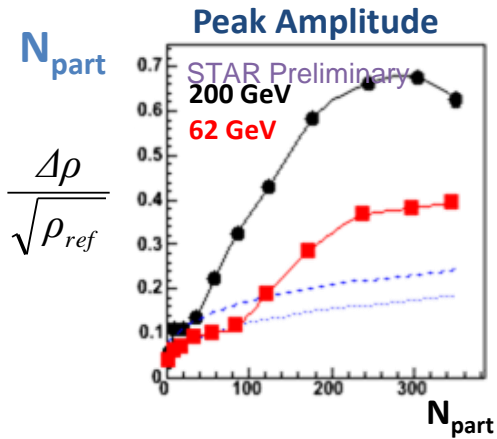
200 GeV

“Critical value” $N_{part} \approx 40$

Same-side gaussian peak width. Points show eleven centrality bins for each energy (84-93%, 74-84%, 65-74%, 55-65%, 46-55%, 37-46%, 28-37%, 19-28%, 9-19%, 5-9%, and 0-5%) transformed to transverse density.

Variation of low- p_T “ridge” with centrality (Npart).

A sudden increase of the **peak amplitude** and **η width** of *the near-side low p_T ridge* were found at an *energy-dependent* centrality point at definite number of participating nucleons **Npart**.

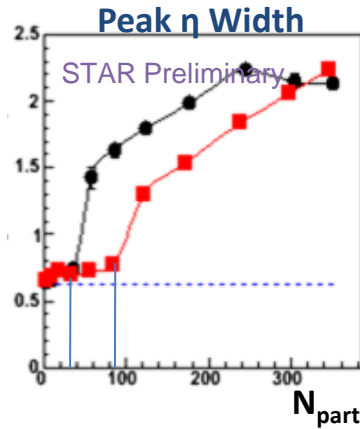


62 GeV

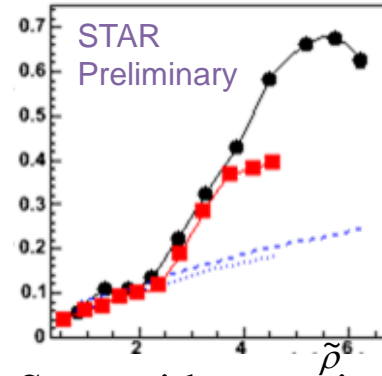
“Critical value” $N_{part} \approx 90$

200 GeV

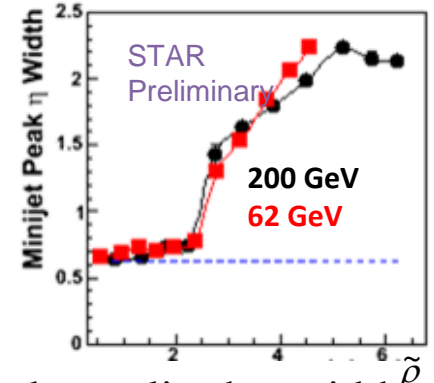
“Critical value” $N_{part} \approx 40$



Transverse Particle Density



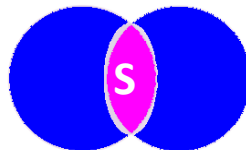
Peak η Width



Same-side gaussian peak amplitude, width. Points show eleven centrality bins for each energy (84-93%, 74-84%, 65-74%, 55-65%, 46-55%, 37-46%, 28-37%, 19-28%, 9-19%, 5-9%, and 0-5%) transformed to transverse density.

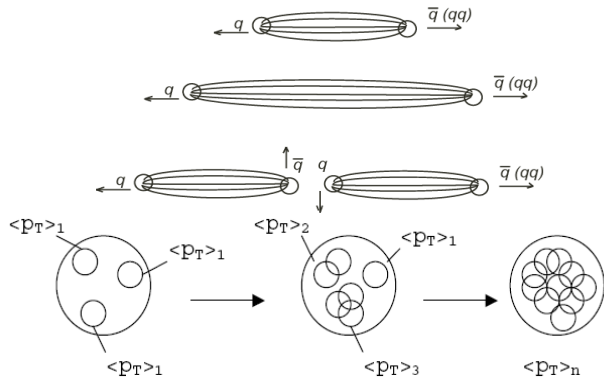
$$\tilde{\rho} = \frac{3}{2} \frac{dN_{ch}}{d\eta} / S$$

$$\tilde{\rho}_{crit} = 2,6 \pm 0,2 \text{ fm}^{-2}$$



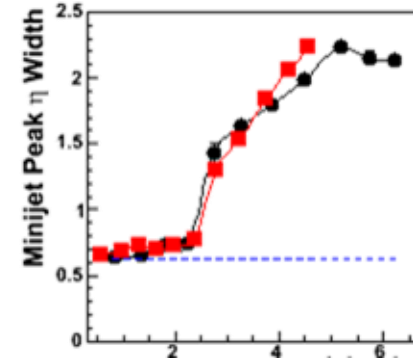
New! Our calculation

Percolation string theory[3]



Percolation model parameters
at the “critical” point

Transverse particle density



$$n = N_{string} \pi r_0^2 / \textcircled{S}$$

$$n_{crit} = 1,12 - 1,175 [4]$$

$$\tilde{\rho}_{crit} = \frac{\tilde{\rho}}{n_{crit}}$$

$$\tilde{\rho} = \frac{3}{2} \frac{dN_{ch}}{d\eta} / \textcircled{S}$$

$$\tilde{\rho}_{crit} = 2,6 \pm 0,2 \text{ fm}^{-2} [1]$$

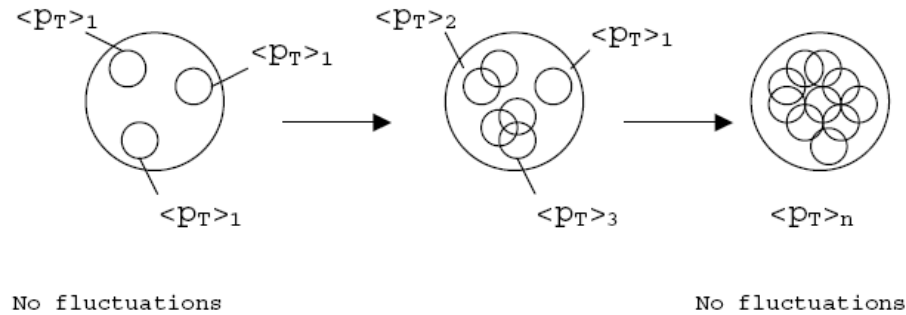
Parameters of the model

[3] C. Pajares // arXiv:hep-ph/0501125v1 14 Jan 2005

[4] N. Armesto, M.A. Braun, E.G. Ferreiro, C. Pajares // Phys. Rev. Lett. 77, 3736 (1996)

[5] V.V. Vechernin, R.S. Kolevatov, 2007, Yad. Fiz., 2007, Vol. 70, No. 10, pp. 1858–1867.

String model. Estimate of string percolation parameter.



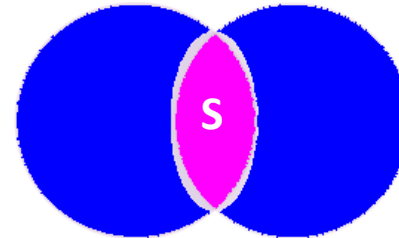
With growing **energy** and/or **atomic number** of colliding particles, the number of strings grows and they start to **overlap**, forming clusters, **new type of particle emitting source**.

At a critical density a macroscopic cluster appears that marks the percolation phase transition. [3]

Percolation parameter:

$$\eta(b) = \frac{N_{str}(b) \pi r_0^2}{S(b)}$$

$N_{str} - ?$
 $S - ?$



$$\eta_c = 1,12-1,175 \text{ ([4])}$$

$$N_{str} = N_{Sea} + N_{Valent}$$

$$\eta(N_{part}) = N_{str}(N_{part}) \pi r_0^2 / S(N_{part})$$

$$N_{str} = x N_{str}^{(pp)} N_{coll} + (1-x) N_{part} \text{ [5]}$$

N_{Str} - number of strings, πr_0^2 string transverse area, S overlap area of two nucleons.

$r_0 = 0,2-0,3 \text{ fm}$ – change of string radius value results in different percolation parameter

[3] C. Pajares // arXiv:hep-ph/0501125v1 14 Jan 2005

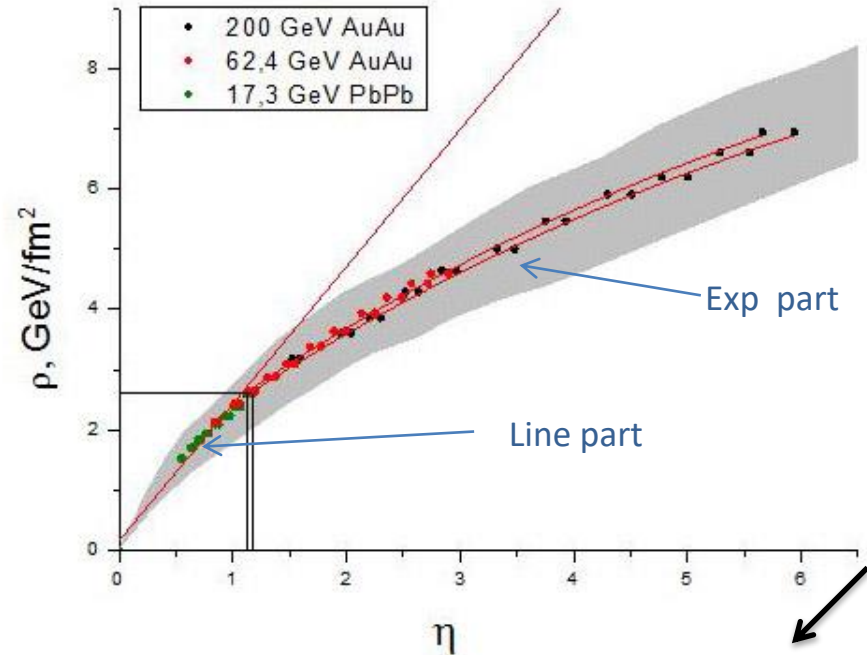
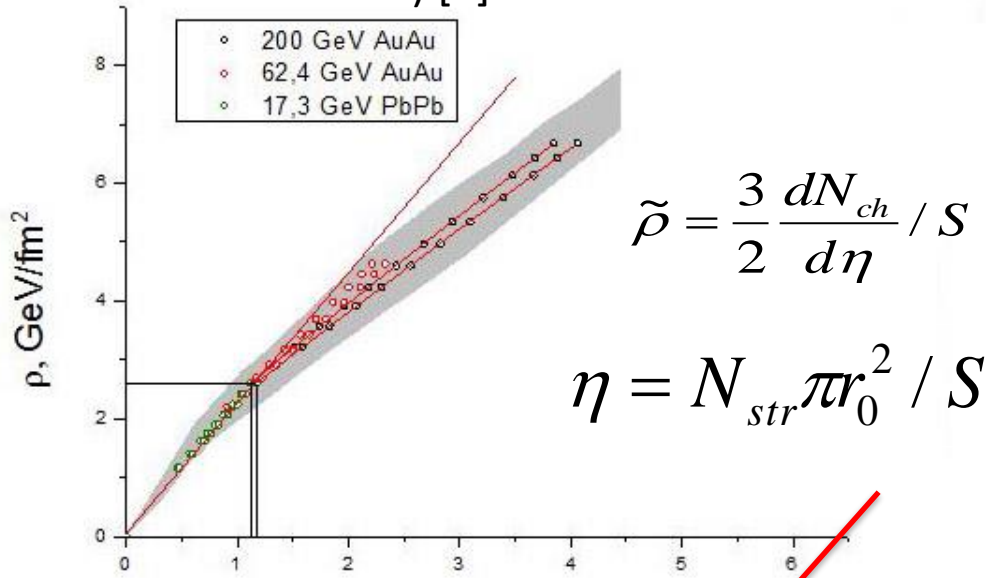
[4] N. Armesto, M.A. Braun, E.G. Ferreiro, C. Pajares // Phys. Rev. Lett. 77, 3736 (1996)

[5] V.V. Vechernin, R.S. Kolevatov, 2007, published in Yadernaya Fizika, 2007, Vol. 70, No. 10, pp. 1858–

Transverse particle density vs. percolation parameter

Modified Glouber Model (taken in account momentum loss) [6]

Classical Glouber Model



Estimates for pp (for $S \sim 1 \text{ fm}^2$):

vs, GeV	dNch/dy[7]	ρ , GeV/fm ²	η	Nstr
7000	6,02	9,0 ± 0,5	5,8 ± 0,5	24 ± 2
2360	4,68	7,0 ± 0,4	4,3 ± 0,4	18 ± 2
900	3,78	5,6 ± 0,4	3,3 ± 0,3	13 ± 1
200	2,30	3,4 ± 0,3	1,7 ± 0,2	7,1 ± 0,7
130	2,03	3,0 ± 0,2	1,5 ± 0,1	6,0 ± 0,6
62,4	1,64	2,5 ± 0,6	1,1 ± 0,1	4,4 ± 0,5
19,4	1,17	1,8 ± 0,5	0,8 ± 0,1	3,1 ± 0,3
17,3	1,14	1,7 ± 0,3	0,7 ± 0,1	3,0 ± 0,3

vs, GeV	dNch/dy[7]	ρ , GeV/fm ²	η	Nstr
7000	6,02	9,0 ± 0,5	10,4 ± 0,9	42 ± 4
2360	4,68	7,0 ± 0,4	6,1 ± 0,5	25 ± 3
900	3,78	5,6 ± 0,4	4,2 ± 0,4	17 ± 2
200	2,30	3,4 ± 0,2	1,9 ± 0,2	7,6 ± 0,8
130	2,03	3,0 ± 0,2	1,5 ± 0,1	6,1 ± 0,6
62,4	1,64	2,5 ± 0,1	1,05 ± 0,09	4,3 ± 0,4
19,4	1,17	1,75 ± 0,09	0,73 ± 0,07	3,0 ± 0,3
17,3	1,14	1,70 ± 0,08	0,70 ± 0,06	2,9 ± 0,3

[6] G.Feofilov, A.Ivanov // Journal of Physics: Conference Series 5 (2005) 230–237

[7] ALICE Collaboration: arXiv:1004.3034[hep-ex], arXiv: 1004.3514[hep-ex]

➤ For pp collision the threshold would appear at energy 70-85 GeV.

Ridge in pp

String density in pp collisions

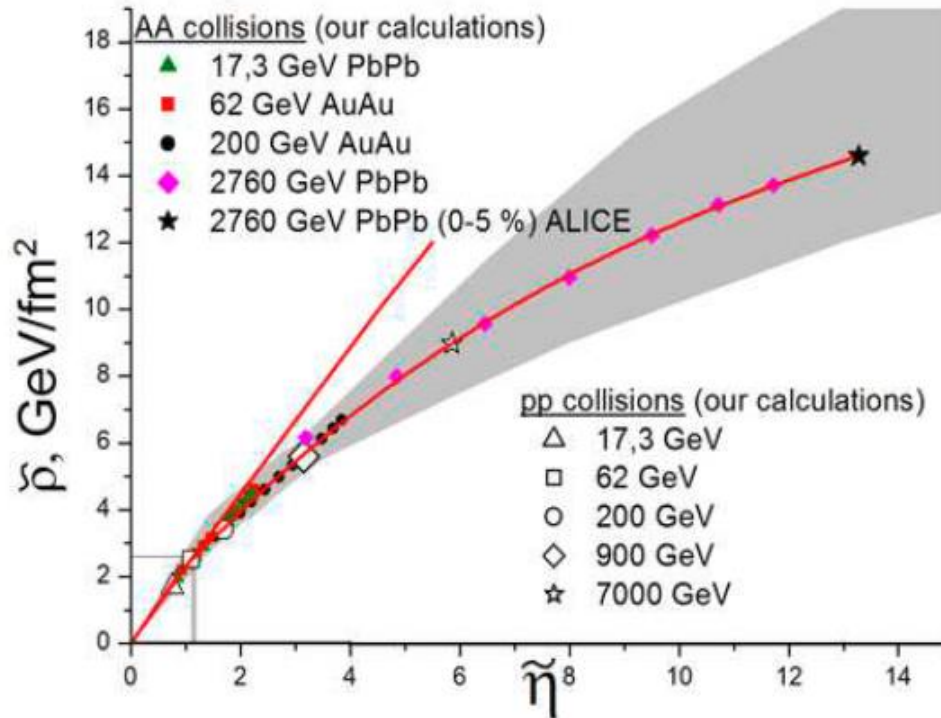


Figure 4: Transverse particle density $\tilde{\rho}$ vs. string density $\tilde{\eta}$ in AA collisions for energy $\sqrt{s} = 17.3$ GeV, $\sqrt{s} = 62$, $\sqrt{s} = 200$ GeV and $\sqrt{s} = 2.76$ TeV. Points are our calculations. Lines are fits to the points. Points for pp collisions are obtained from the experimental data in the range of energies from $\sqrt{s} = 17.3$ GeV to $\sqrt{s} = 7$ TeV (see text).

O.Kochebina and G.Feofilov, Onset of "ridge phenomenon" in AA and pp collisions and percolation string model, arXiv:1012.0173v1



Azimuthal flow in hadron collisions from quark-gluon string repulsion

Igor Altsybeev, Grigory Feofilov
St.Petersburg State University

Quarks-2016
19th International Seminar on High Energy Physics

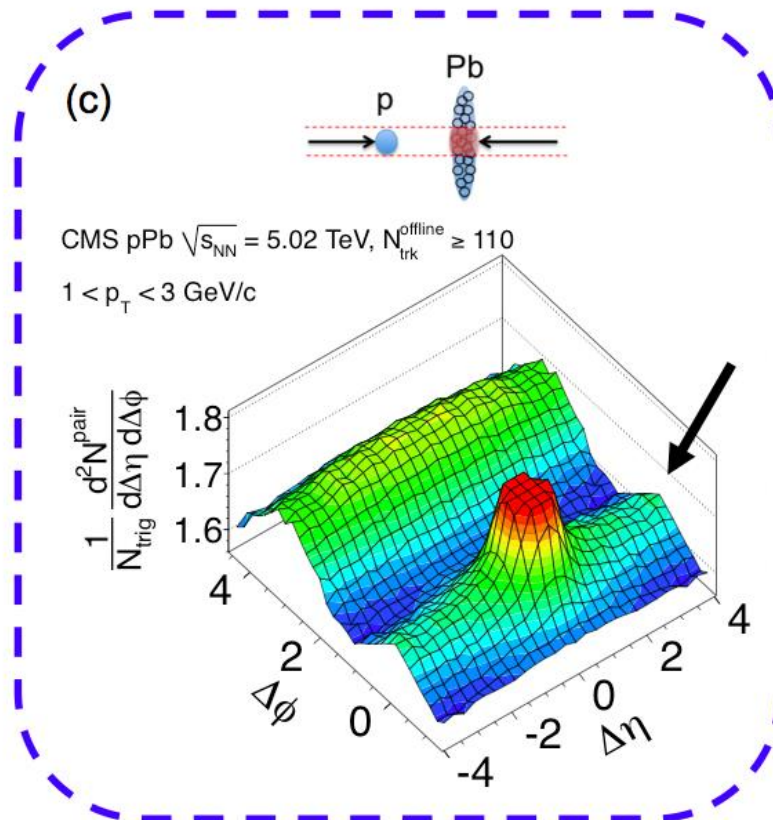
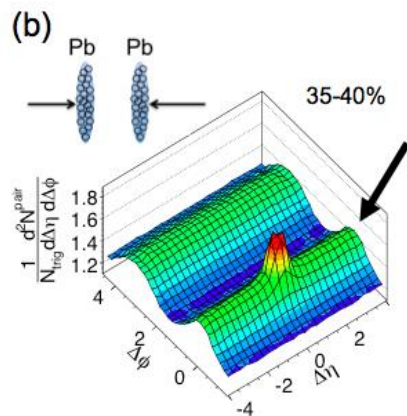
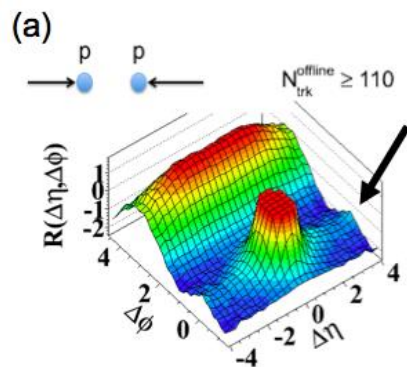
May 31, 2016

Unexplained long-range correlations

– “Ridge” by CMS in p-Pb at LHC

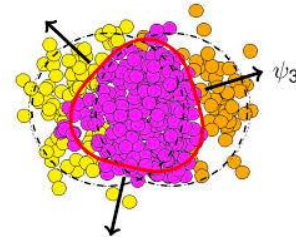
– similar to pp and Pb-Pb !

LHC

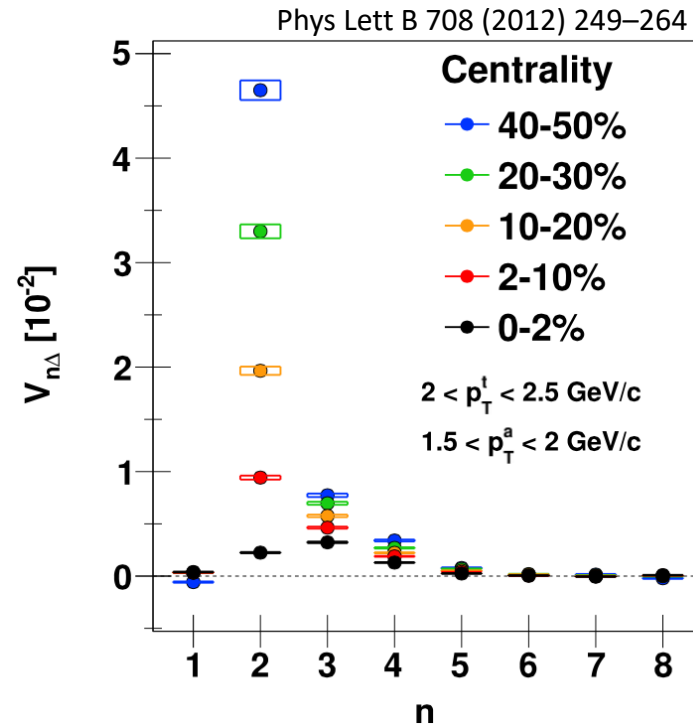
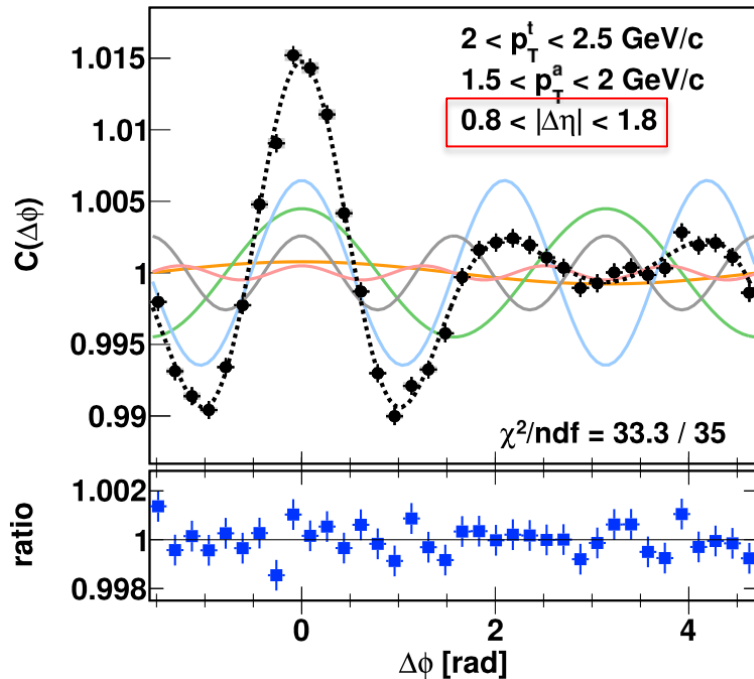


Experiment: Harmonic decomposition of long-range two-particle angular correlations

- shape of the overlap region \rightarrow large elliptic flow v_2 , non-zero v_4
- fluctuations of the initial energy density profile \rightarrow affect v_2 , give v_3, \dots



Most central (0-2%) Pb-Pb collisions at 2.76 TeV:



“Double hump” structure – due to larger v_3 relative to v_2 in most central collisions.

Our approach:

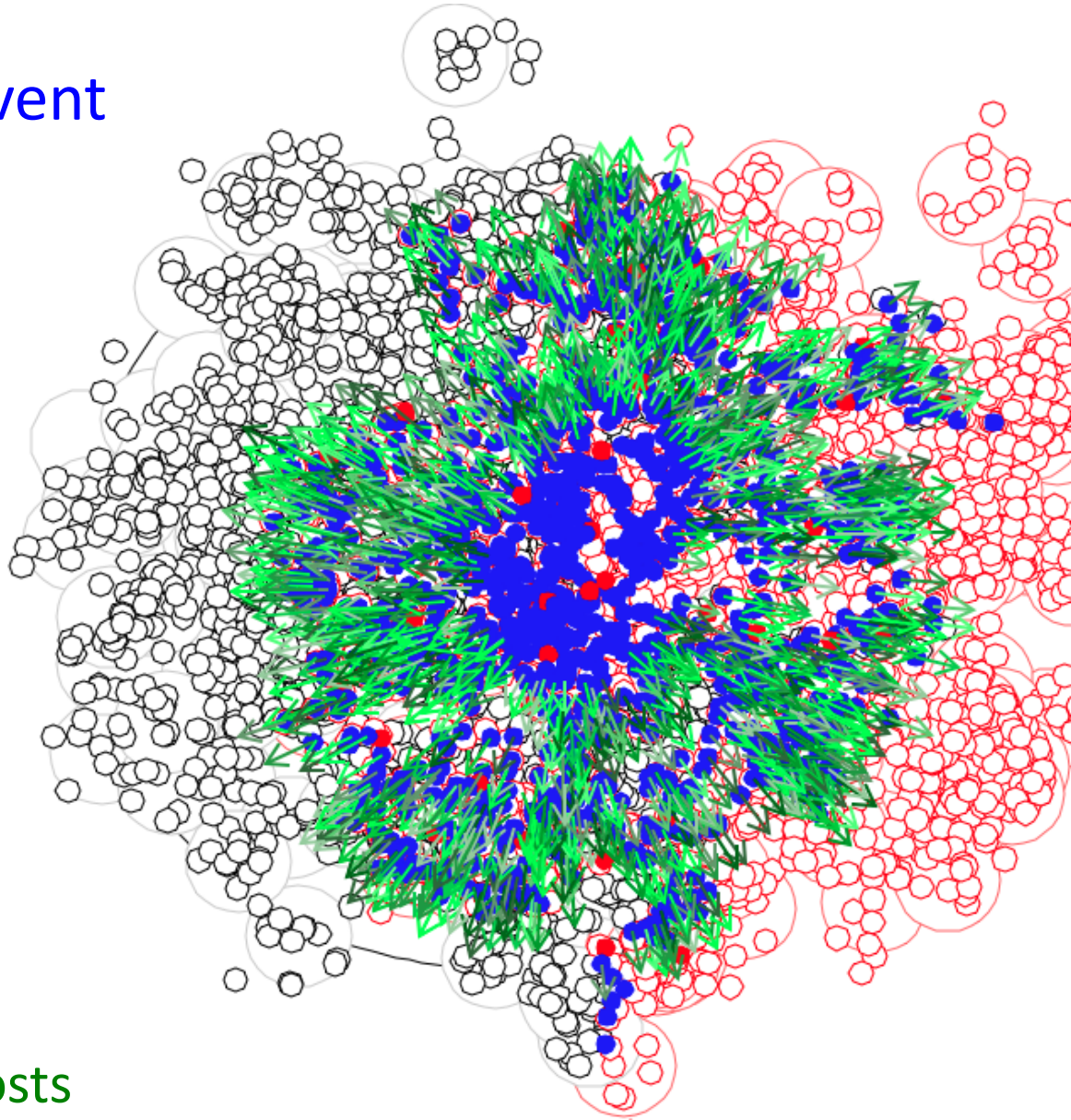
MC model of repulsive strings*

Assumptions:

- Pair of partons from colliding nucleons can form a string (a color flux tube)
- Strings are repulsing each other within some distance (efficient radius – R_{eff})
- Thus each string gets a summarized momentum kick
- In its rest frame, string decays into particles isotropically in azimuth, but in the lab frame the particles are boosted due to string momentum kick

* see report by I.Altsybeev "Mean transverse momenta correlations in hadron-hadron collisions in MC toy model with repulsing strings" at **Quark Confinement and the Hadron Spectrum XI** 7-12 September 2014 St. Petersburg, arXiv:1502.03608, AIP Conf.Proc. 1701 (2016)

Visualization of a toy event



Pb-Pb collision at $b=4$ fm

blue circles – strings

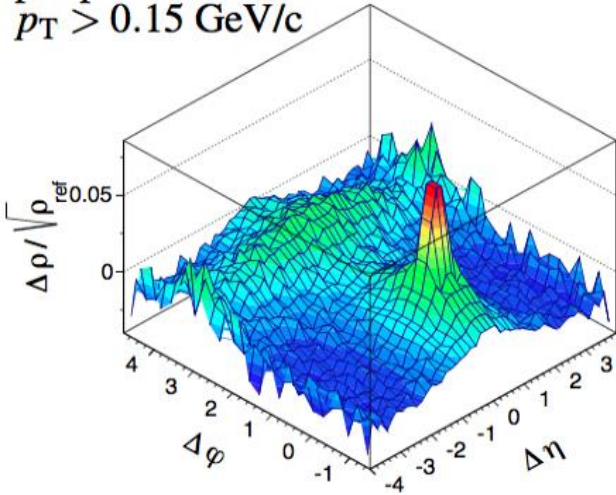
green arrows – string boosts

Two-particle correlations in MC toy model

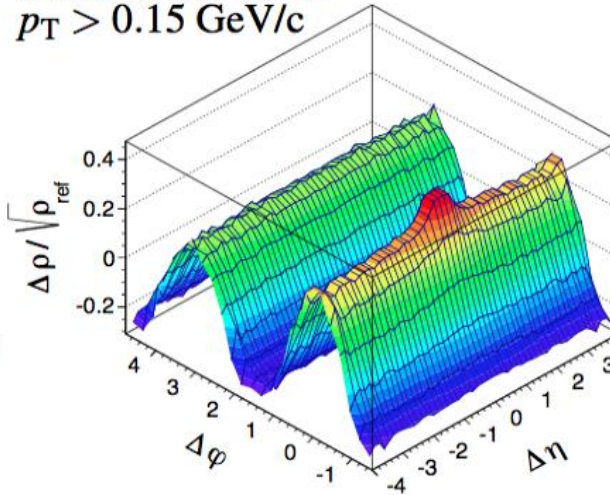
MC model results for A-A events at high energies:

(string effective interaction radius is 2 fm)

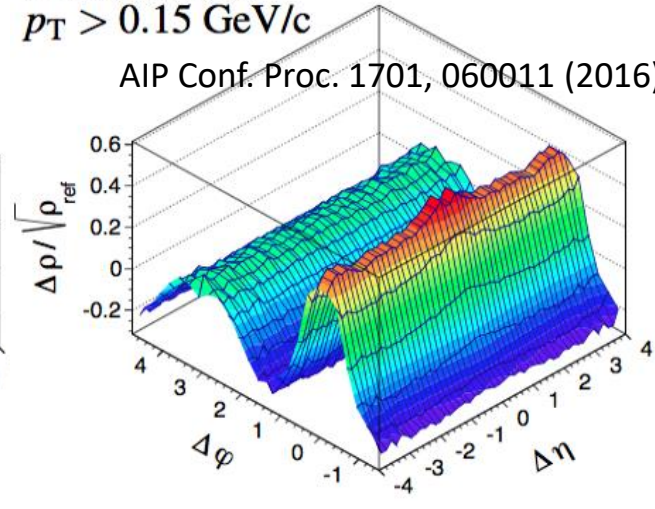
peripheral events
 $p_T > 0.15$ GeV/c



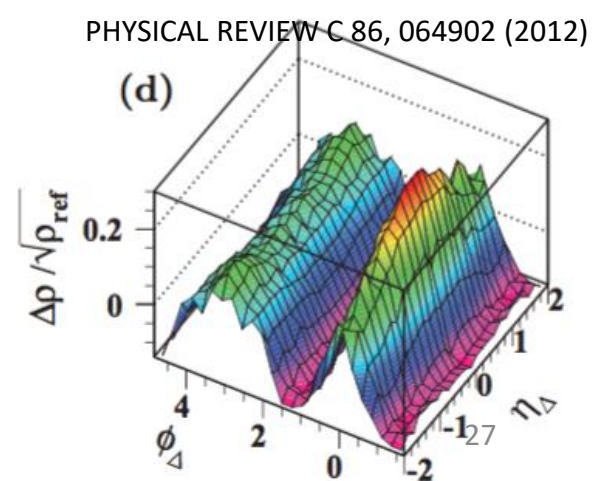
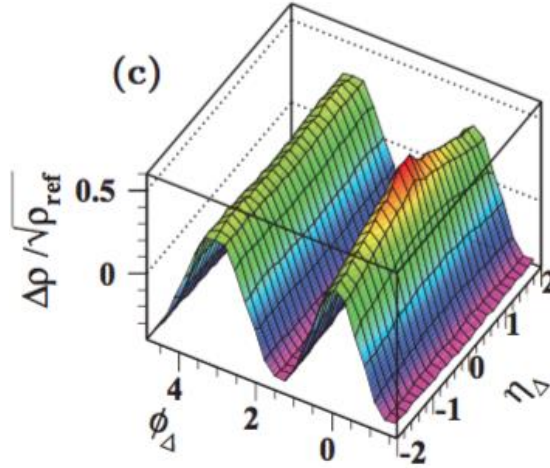
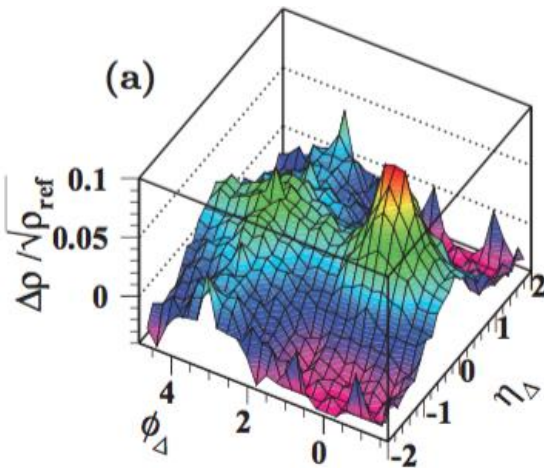
semicentral events
 $p_T > 0.15$ GeV/c



central events
 $p_T > 0.15$ GeV/c



Compare with STAR data (Au-Au):



Our results

G.Feofilov, I.Altysbeev,O.Kochebina, *PoS (Baldin ISHEPP XXII) 067 (2015)*

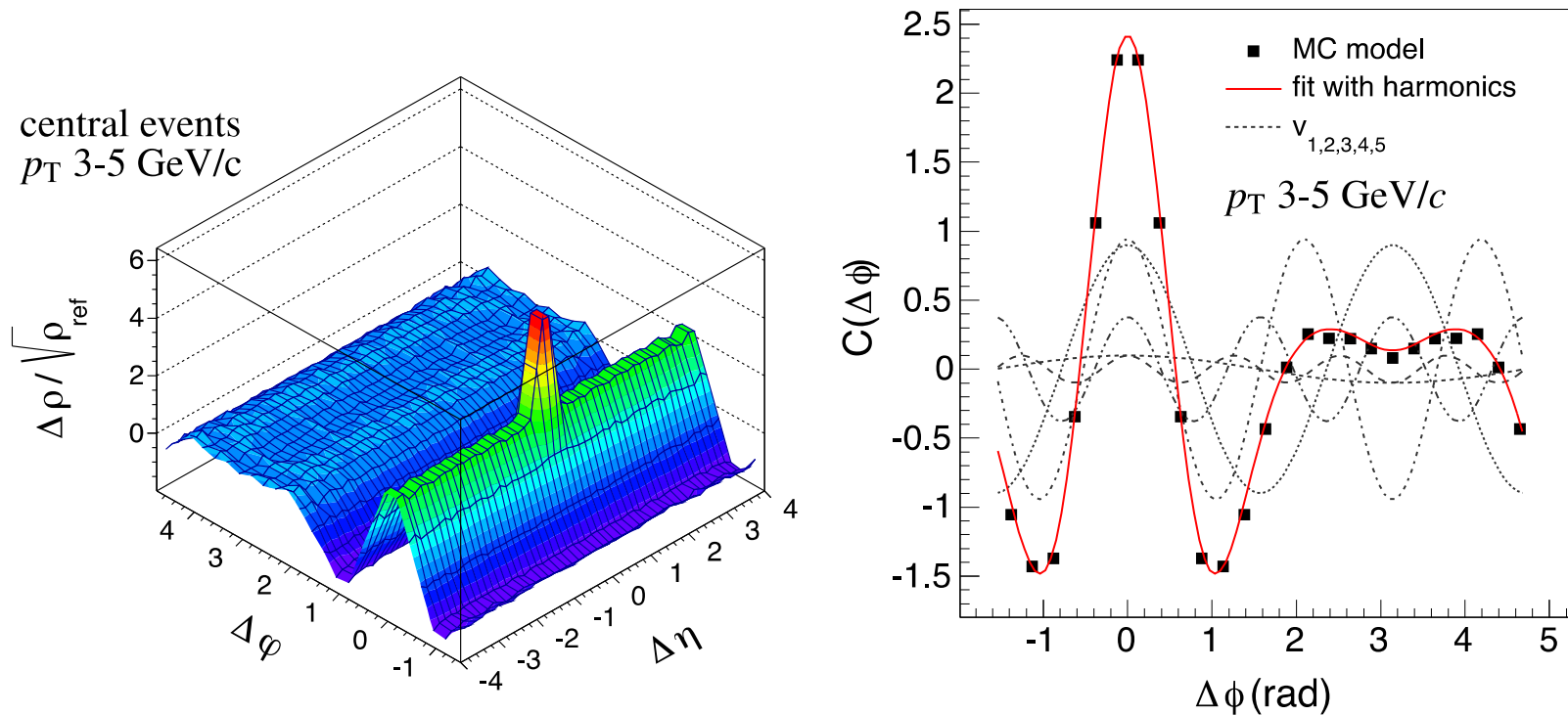
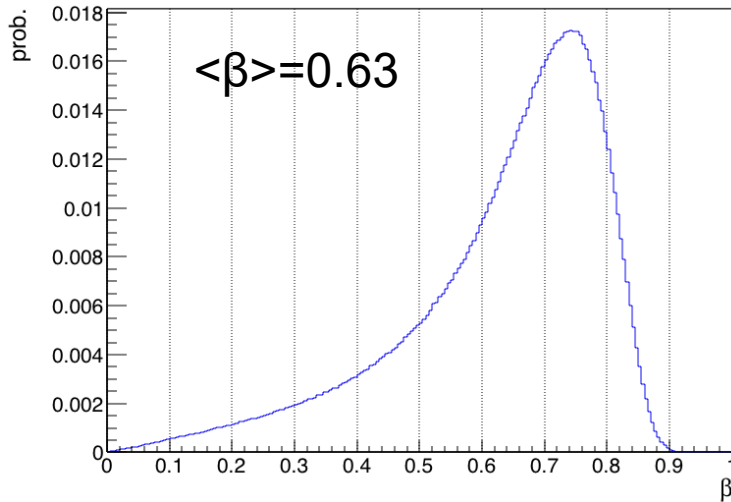
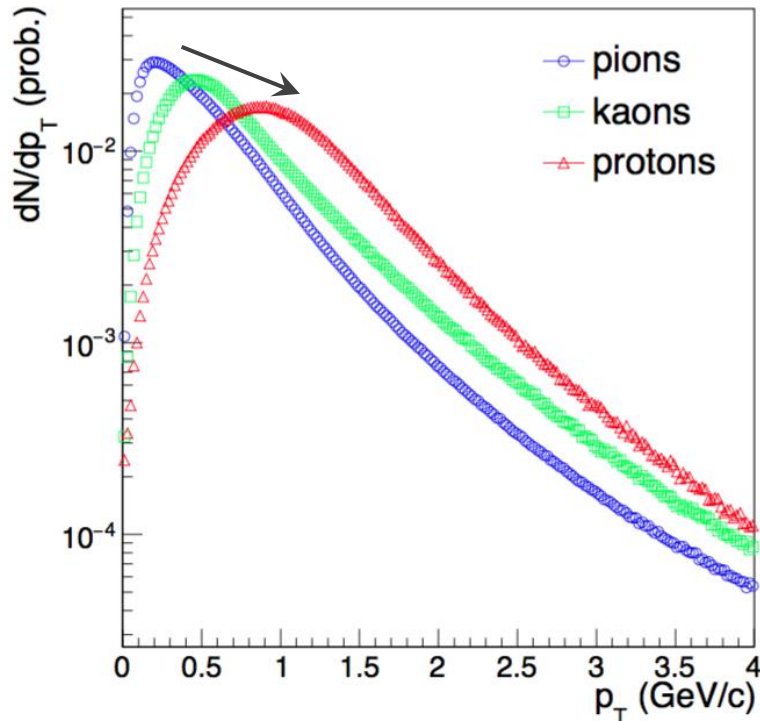
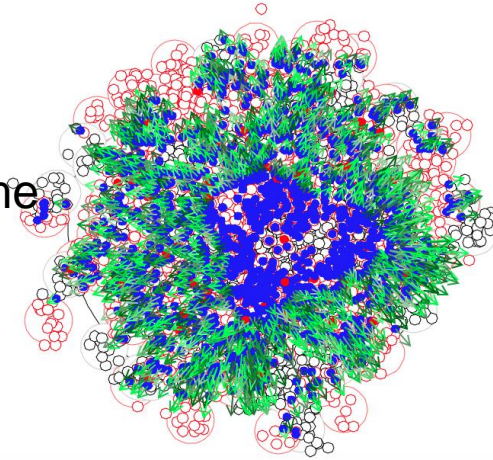


Figure 5: Left: Toy model [19] two-particle correlation function obtained in the most central events for charged particles with $p_T \in [3,5]$ GeV/c. String-string interaction radius R_{int} is 2 fm. Right: The harmonic decomposition of the azimuthal profile of the correlation function shown in the left pad. The solid red line

Radial boost of the strings: influence on p_T spectra



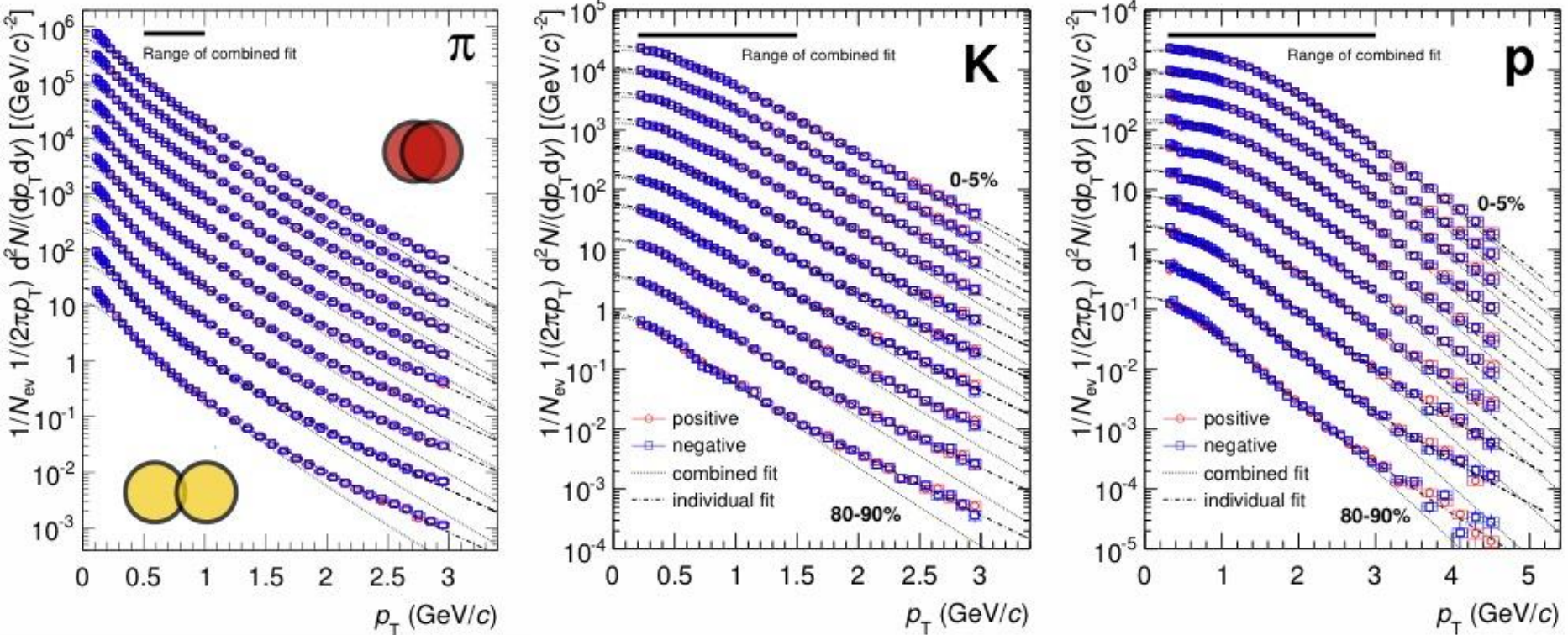
String repulsion leads to some distribution of transverse boosts of the strings.



- Particles from strings are boosted in transverse direction.
- More massive particles gain more p_T shift. - hardening of spectrum for protons may be observed

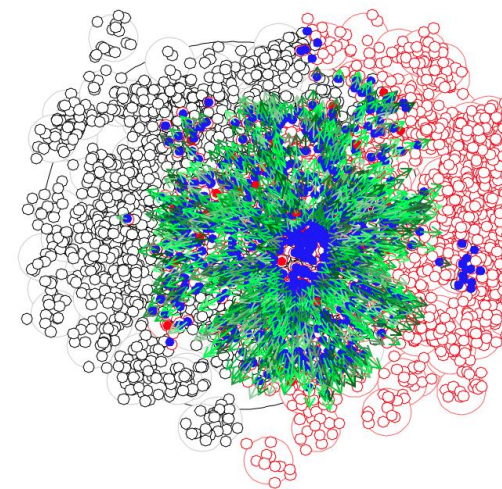
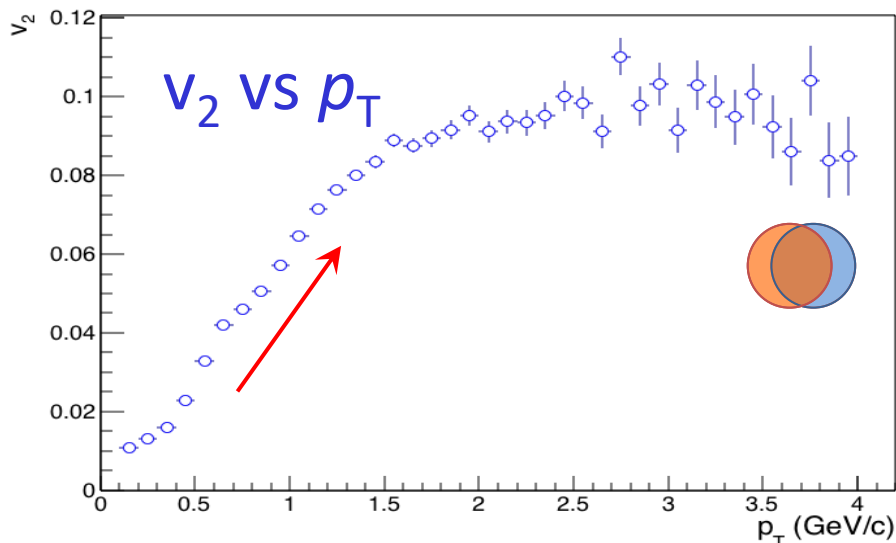
compare to ALICE data →

Bulk particle production in Pb-Pb

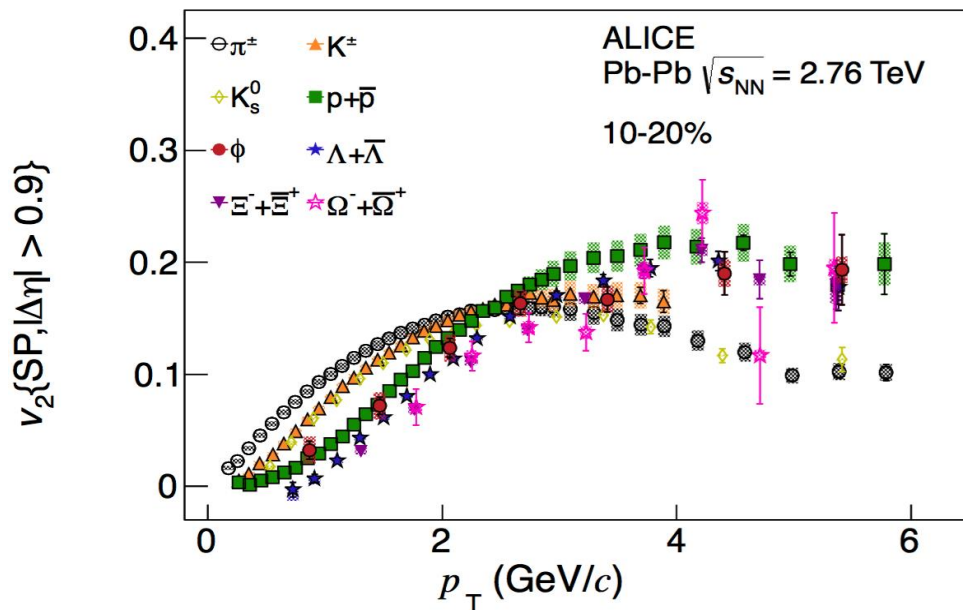


clear evolution of particle spectra \rightarrow hardening with centrality
 more pronounced for protons than for pions
mass ordering as expected from collective hydro expansion

Differential flow of charged particles

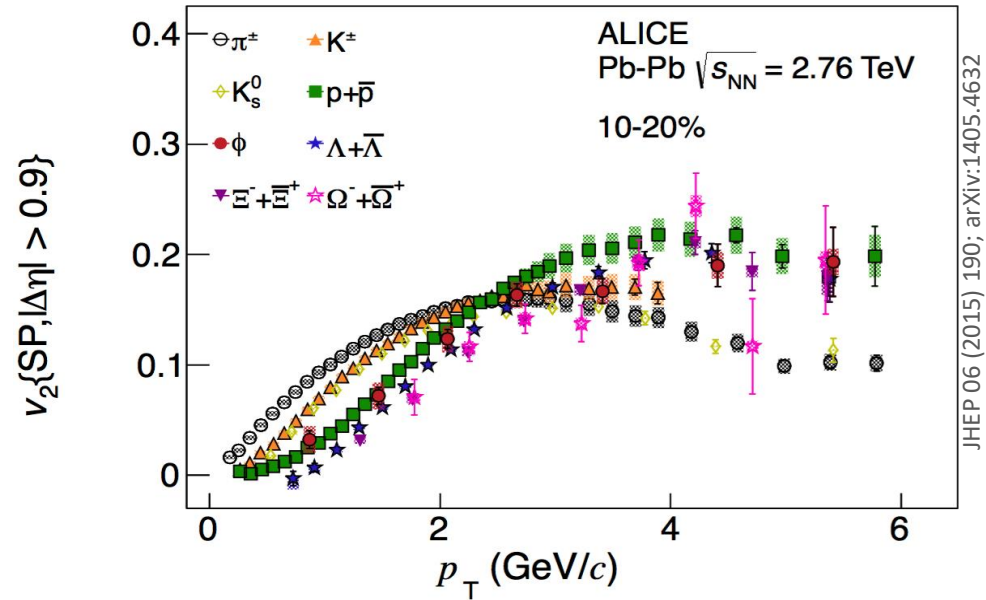
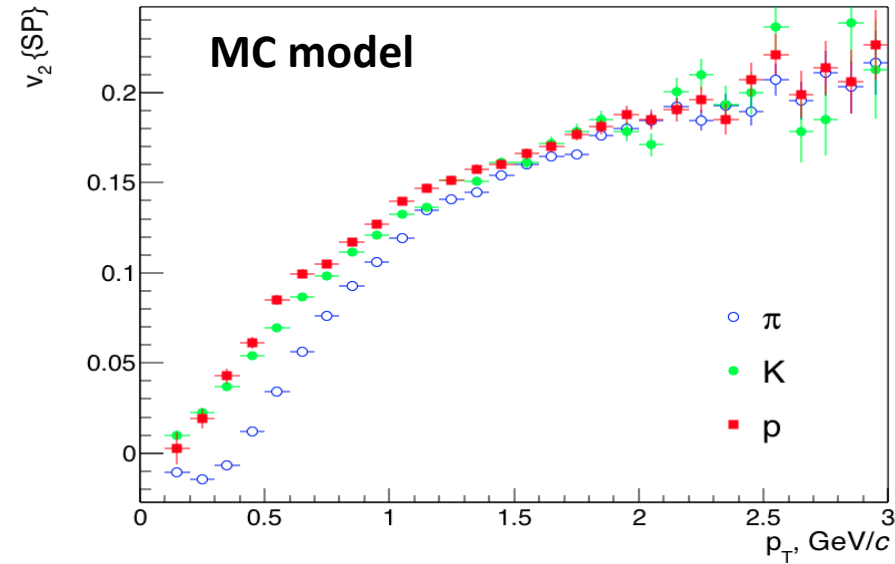


- Qualitative behavior of $v_2(p_T)$ is reproduced in MC model
- “Mass ordering” – to be checked



JHEP 06 (2015) 190; arXiv:1405.4632

“Mass ordering” in MC model with repulsive strings



- Qualitative behavior of $v_2(p_T)$ is reproduced in the MC model
- However, “mass ordering” is NOT properly reproduced (compare with ALICE results)

→ modifications of the model are needed

Possible modifications:

- “dynamical motion” of strings
- “quenching” of hadrons in medium

Conclusions

- **High density string medium** could be formed in hadronic collisions.
- It is shown in Monte Carlo toy-model that collective effects appear due to (rather weak) **repulsive type of string interaction between color flux tubes** (strings): the particle-emitting sources (strings) are boosted by the combined repulsion by neighbours.
- Thus the initial space anisotropy is converted into the harmonics of azimuthal momenta distributions of charged particles.
- Similar results are expected at sufficiently high string density collisions (as is expected for **p-Pb** and **pp** collisions **PoS (Baldin ISHEPP XXII) 067 (2015)**).
- ...More detailed quantitative analysis including mass ordering of flows, in the case of proton-proton and proton-nucleus collisions, will follow...

Layout

- Emergence of Long-Range Angular Correlations in Low-Multiplicity Proton-Proton Collisions(PHYSICAL REVIEW LETTERS 132, 172302 (2024))—
кратко
- Percolation string model and old predictions for critical behaviour -- where to search for the critical point
- **High and Low-Multiplicity pp-Collisions and Long-Range Angular Correlations**
- Plans

Вопросы:

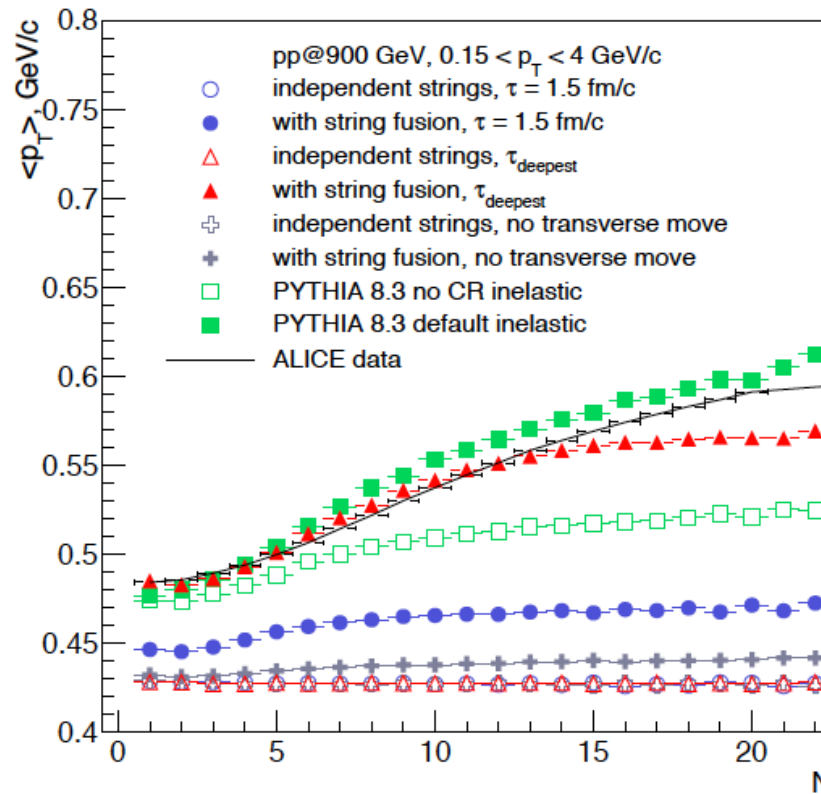
- 1) Количество струн в pp-столкновениях большой множественности
– струны какого типа?

Доминирование морских струн [1].

- 2) Плотность морских струн в самых центральных pp-столкновениях и дальние корреляции?
- 3) $\langle p_T \rangle - N_{ch}$ - корреляции
- 2) $\langle p_T \rangle - N_{ch}$ -дальние корреляции

[1] O.Kochebina and G.Feofilov, Onset of "ridge phenomenon" in AA and pp collisions and percolation string model, arXiv:1012.0173v1

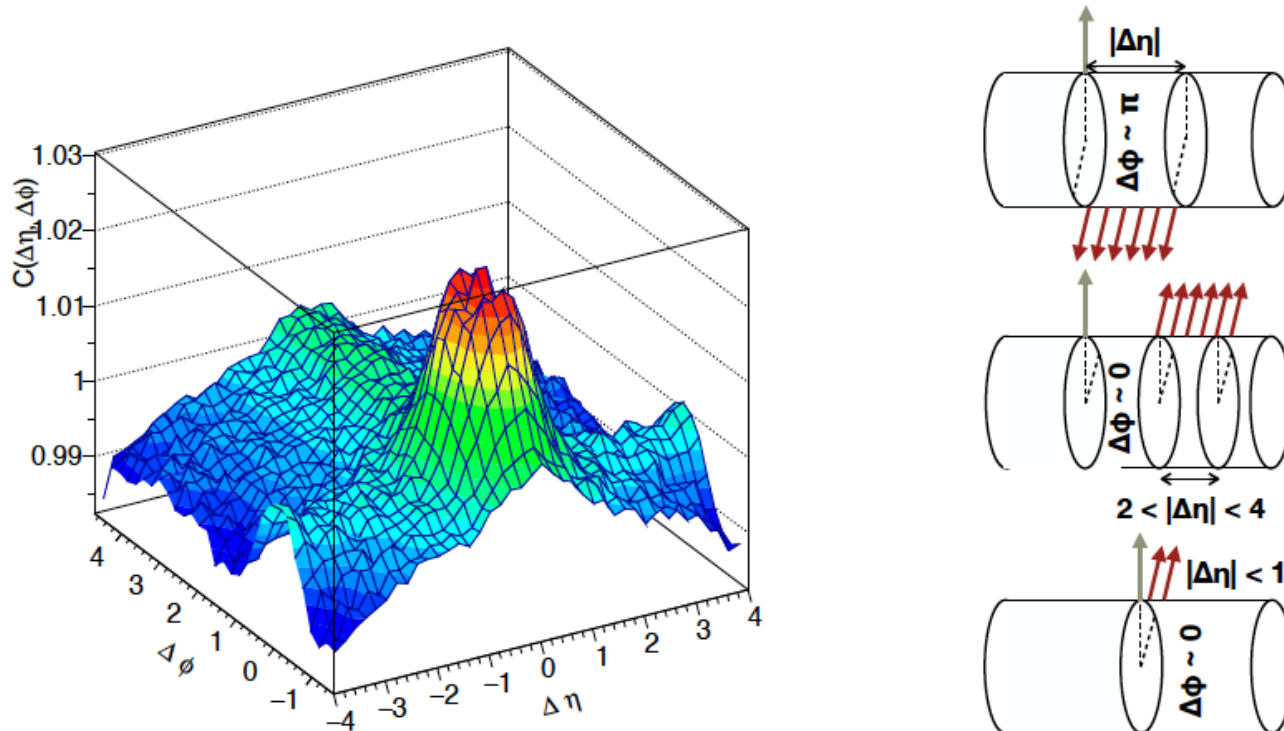
Модель взаимодействующих струн, конечных по быстроте, учитывающая поперечную и продольную динамику (кандидатская Дарьи Прохоровой)



$\langle p_T \rangle - N_{ch}$ - корреляции
в центральной области быстрой

Рис. 2.8: $\langle p_T \rangle - N$ корреляционная функция, вычисленная в аксептансе $|\eta| < 0.8$, $0.15 < p_T < 4$ ГэВ для неупругих $p + p$ взаимодействий при $\sqrt{s} = 900$ ГэВ. Результаты модели для струнной системы, эволюционирующей до $\tau_{deepest}$, представлены полными красными треугольниками, настроены в соответствии с данными ALICE [129] (черная линия: точки данных соединены, чтобы улучшить восприятие). Подробности см. в тексте.

Модель взаимодействующих струн, конечных по быстроте, учитывающая поперечную и продольную динамику (кандидатская Дарьи Прохоровой)

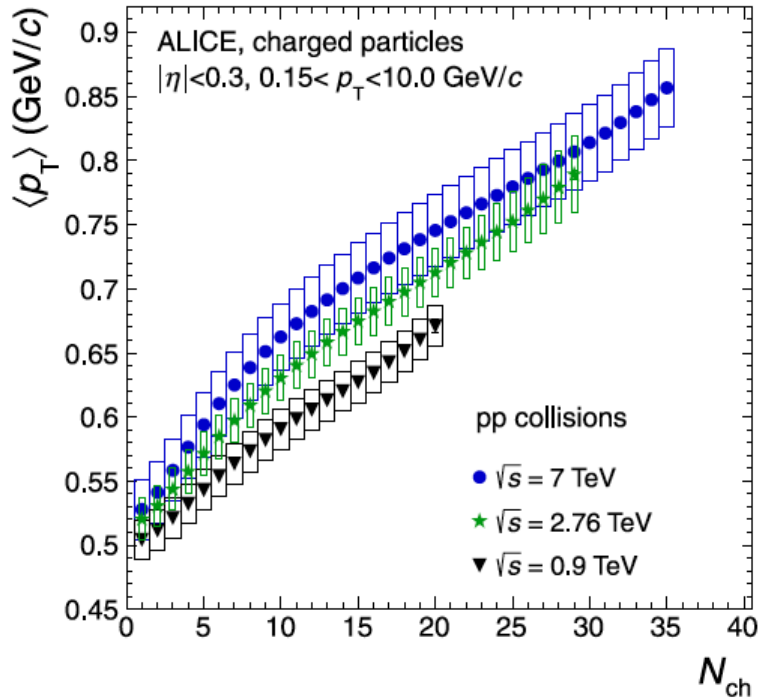


ыстрот

Рис. 3.7: Слева: Модельный результат для двухчастичной корреляционной функции $C(\Delta\eta, \Delta\phi)$, рассчитанной для частиц с $|\eta| < 2.5$ и $0.3 < p_T < 3.0$ ГэВ и представленной для класса событий 0 – 10% на основе множественности заряженных частиц N_{ch}^{sel} с $|\eta| < 2.5$ и $p_T > 0.6$ ГэВ для неупругих $p + p$ взаимодействий при $\sqrt{s} = 13$ ТэВ. Справа: схематическое объяснение наблюдаемых $\Delta\eta - \Delta\phi$ корреляций.

Multiplicity dependence of the average transverse momentum in pp, p–Pb, and Pb–Pb collisions at the LHC

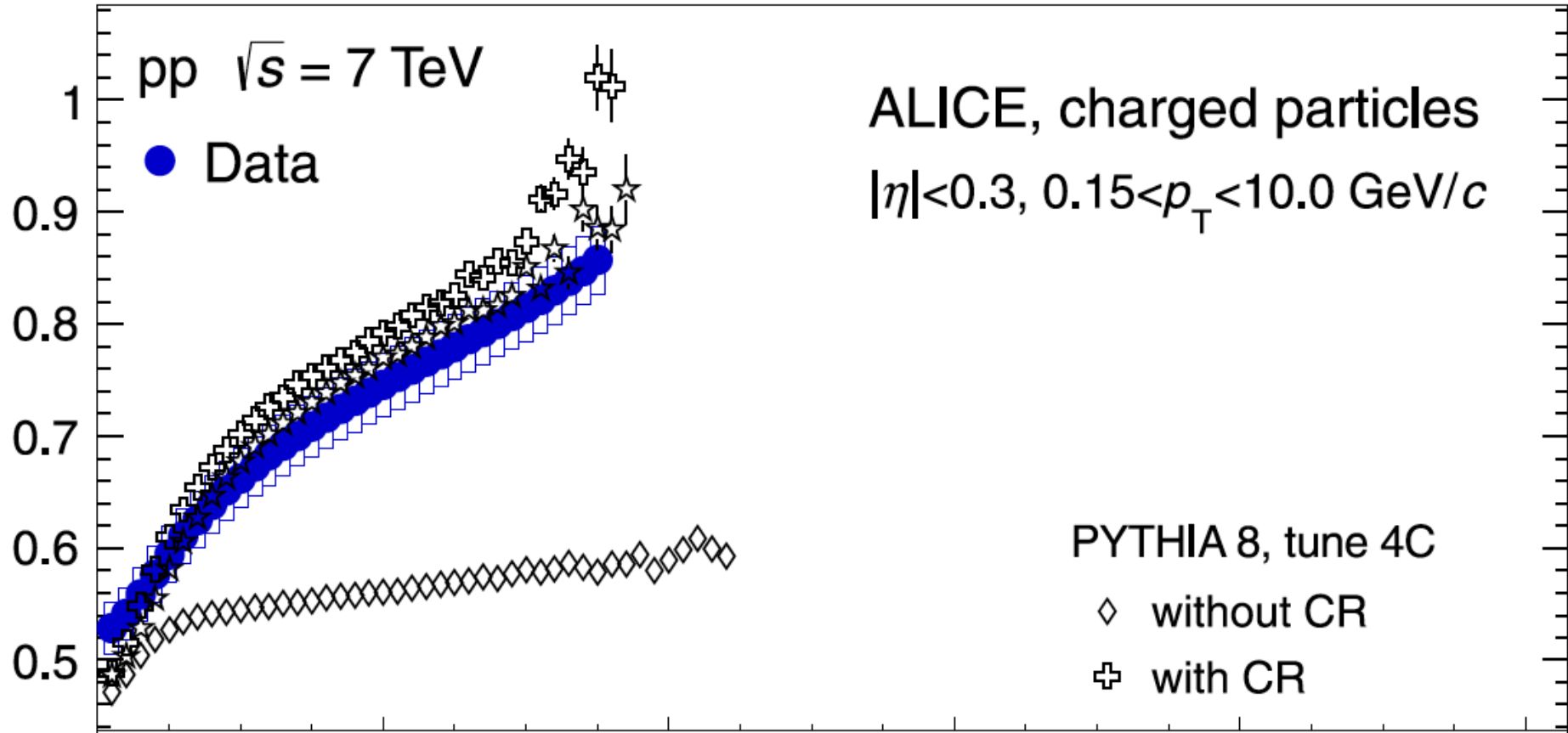
Physics Letters B, Volume 727, Issues 4–5, 18
ALICE Collaboration
December 2013, Pages 371-380



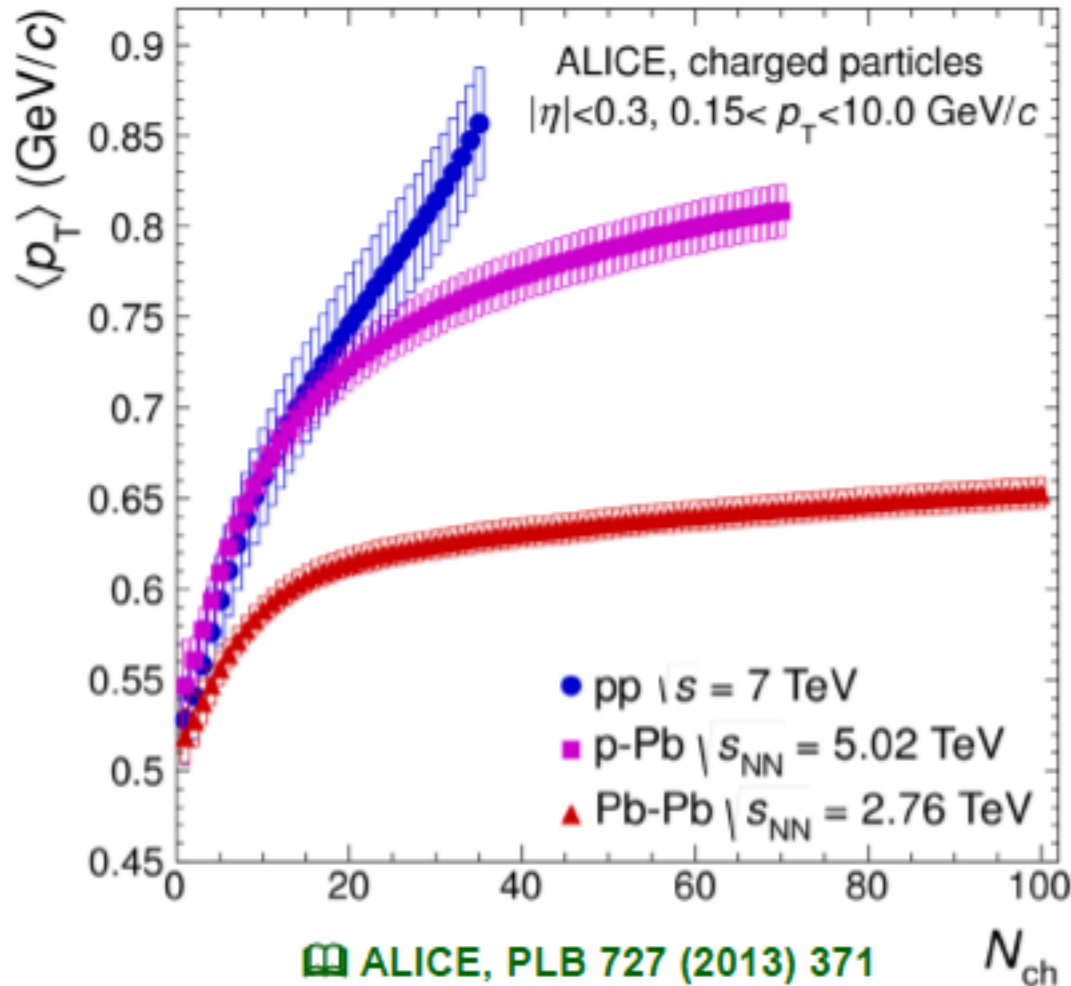
➤ $\langle p_T \rangle - N_{ch}$ - корреляции
в центральной области быстрот
--- а дальние корреляции?

Fig. 1. Average transverse momentum $\langle p_T \rangle$ in the range $0.15 < p_T < 10.0 \text{ GeV}/c$ as a function of charged-particle multiplicity N_{ch} in pp collisions at $\sqrt{s} = 0.9, 2.76,$ and 7 TeV , for $|\eta| < 0.3$. The boxes represent the systematic uncertainties on $\langle p_T \rangle$. The statistical errors are negligible.

PYTHIA



Mean p_T in pp , p -Pb and Pb -Pb



- Three different \sqrt{s} for pp , p -Pb and Pb -Pb
 - ⇒ but \sqrt{s} dependence expected to be weak
- Much stronger increase of $\langle p_T \rangle$ in p -Pb than in Pb -Pb
- p -Pb follows pp up to $N_{ch} \sim 14-15$
- $N_{ch} > 14$ corresponds to
 - ⇒ $\sim 10\%$ of pp x-section:
 - ✓ *pp already highly biased*
 - ⇒ 50% of p -Pb x-section
 - ✓ *only centrality bias*

ЕРОХИН Андрей Александрович

Выпускная квалификационная работа

Исследование энергетической зависимости дальних корреляций множественности заряженных частиц в протон-протонных столкновениях в эксперименте ALICE на LHC

Уровень образования: аспирантура

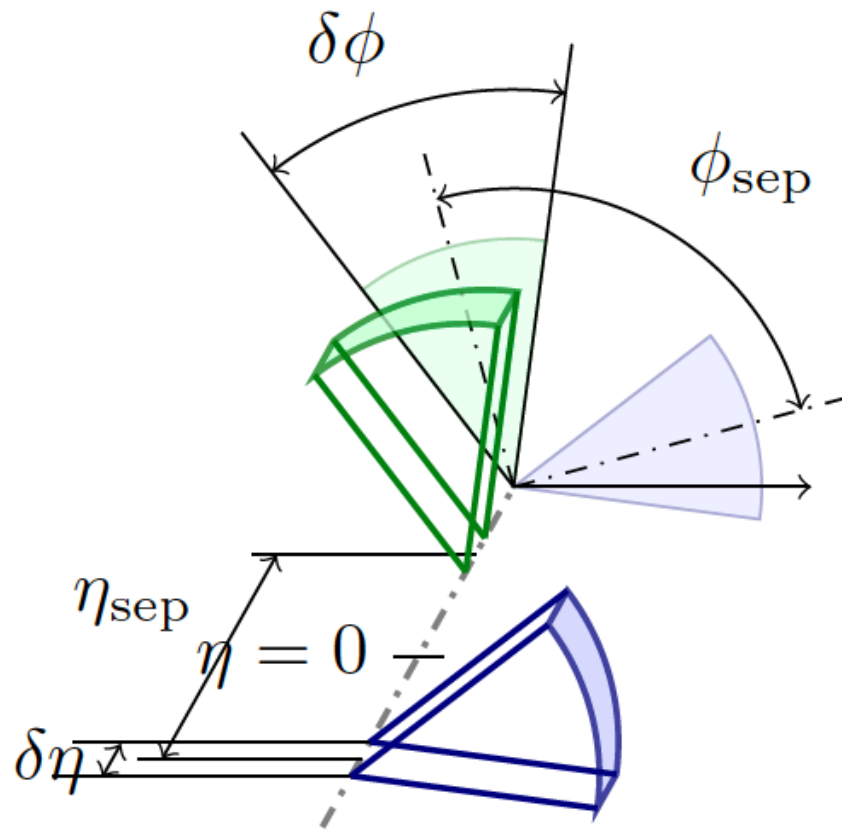


Рис. 6: Схематическое изображение окон, в которых производились измерения корреляций множественности заряженных частиц.

$\sqrt{s_{CM}}$, ГэВ	Период pass/AOD	Число отобранных событий
900	LHC10c pass4/AOD221	4 563 735
2760	LHC13g pass1/AOD155	11 349 195
5020	LHC15n pass4/AOD208	99 612 843
7000	LHC10b pass4/AOD221	15 642 255
7000	LHC10d pass4/AOD221	65 100 283
13000	LHC17k pass1/AOD234	85 846 276

Таблица 3: Наборы данных, использованные в анализе

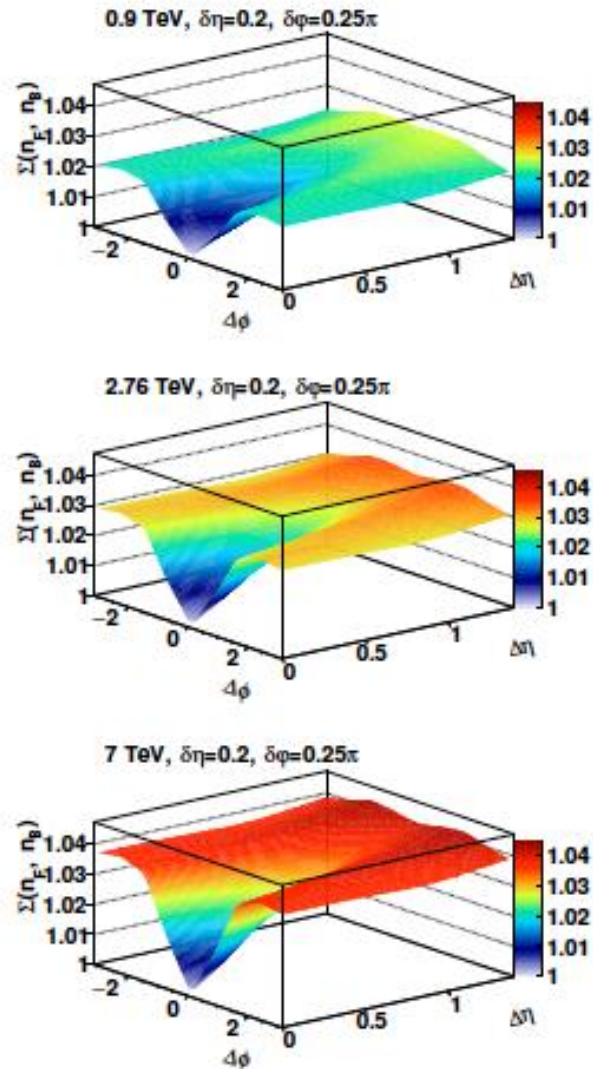
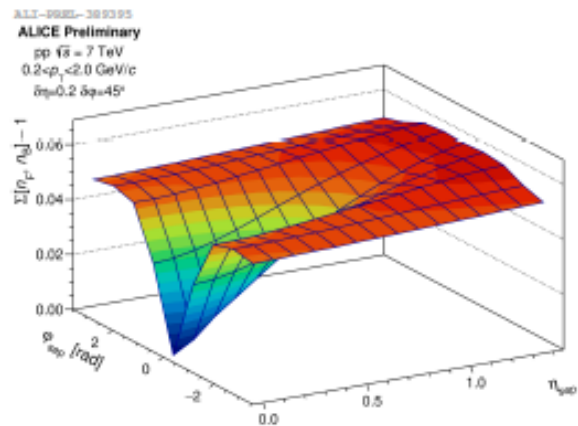
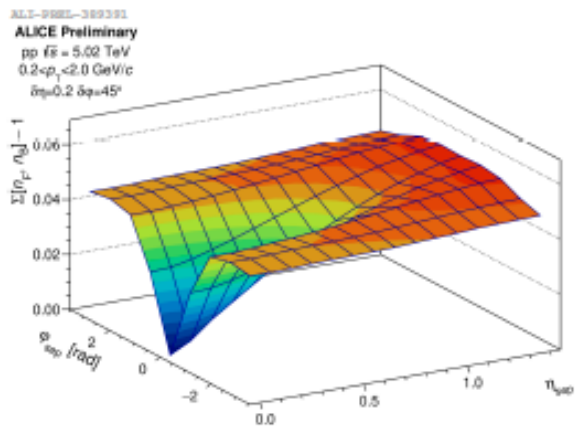
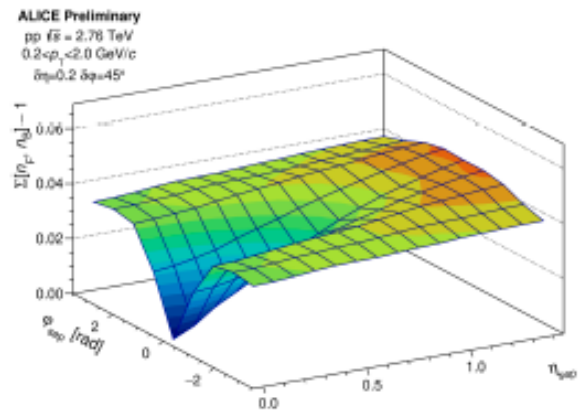
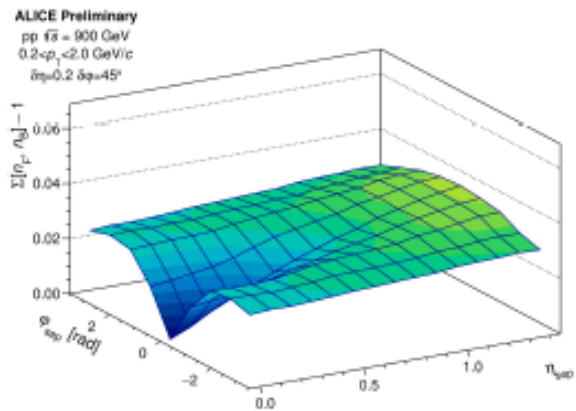
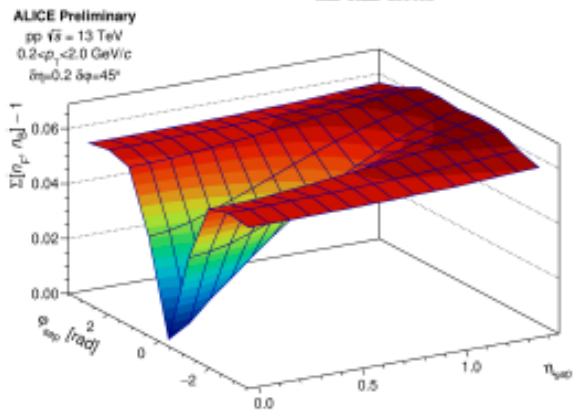
Σ 

Рис. 2: Сильно интенсивная наблюдаемая $\Sigma(n_F, n_B)$ между множественностями в двух окнах шириной $\delta\eta = 0.2$ и $\delta\varphi = \pi/4$ как функция расстояния между центрами окон $\Delta\eta$ по быстрой и $\Delta\varphi$ по азимуту, рассчитанный в модели с независимыми идентичными строками в статье [31] с использованием двухчастичная корреляционная функция струны $\Lambda(\Delta\eta, \Delta\varphi)$ (6) получено путем фитирования [29] экспериментальных данных pp ALICE [8] по FB корреляциям между множественностями при трех начальных энергиях: 0,9, 2,76 и 7 ТэВ в аксептансе псевдобыстрот ALICE TPC. (Рисунок из статьи [31].)



ALI-900EL-389381

ALI-900EL-389395



ALI-900EL-389407

Рис. 9: $\Sigma[n_F, n_B] - 1$ в зависимости расстояния между центрами двух окон размера $\delta\eta = 0.2$ $\delta\varphi = 45^\circ$ для 5 энергий pp столкновений¹.

Поиск нелинейных эффектов в адронных столкновениях на БАК

Магистерская диссертация студента

_____ **Ерохина Андрея Александровича**

Научный руководитель:

_____ к. ф.-м. н., доц. **Г. А. Феофилов**

Рецензент:

_____ к. ф.-м. н., снс. **Е. Л. Крышень**

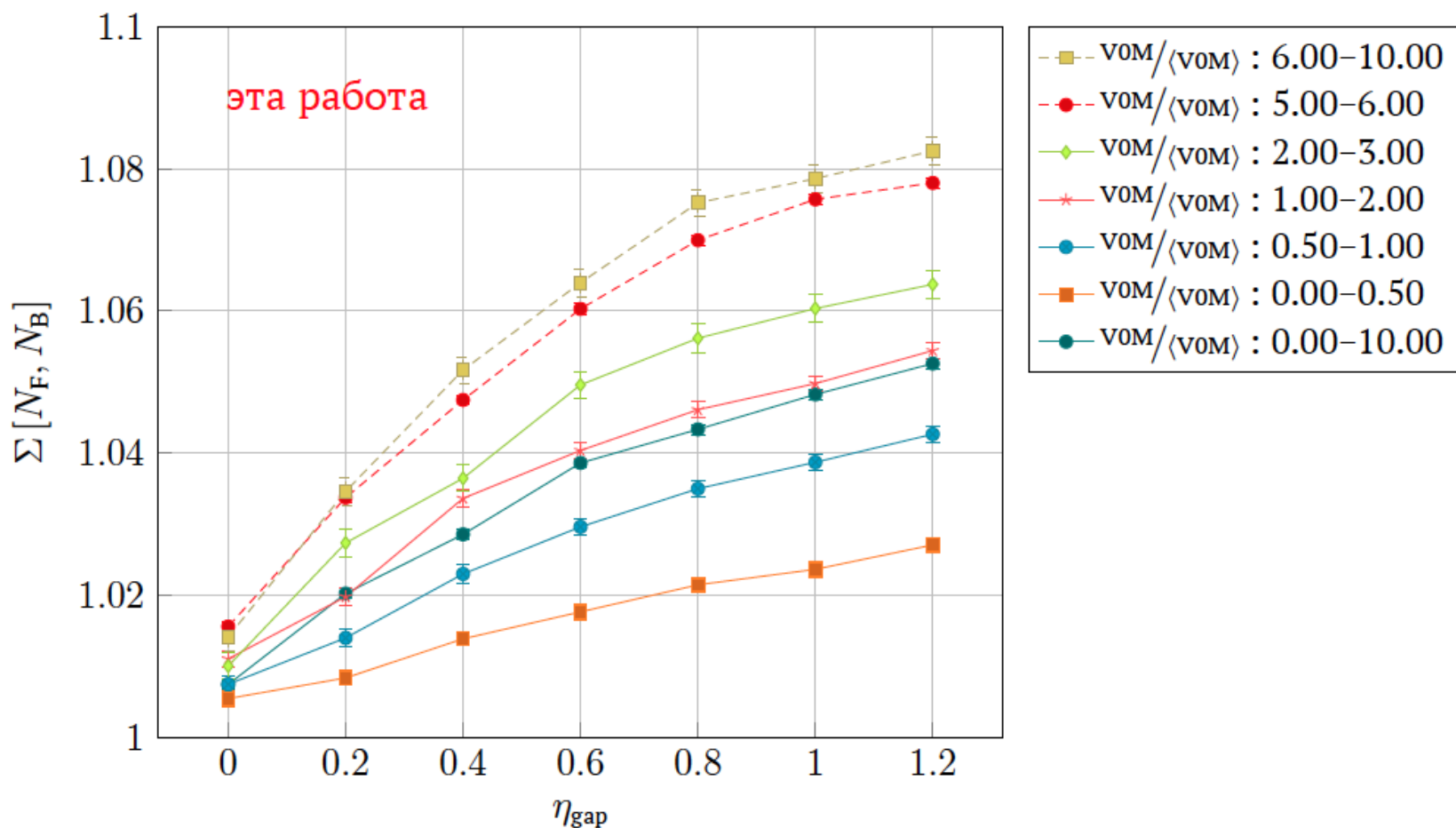


Рисунок 11 – Сечение двумерного графика $\Sigma [N_F, N_B]$ в $\varphi_{\text{sep}} = 0$ для разных классов по множественности. $0.2 \text{ ГэВ} < p_T < 2.0 \text{ ГэВ}$

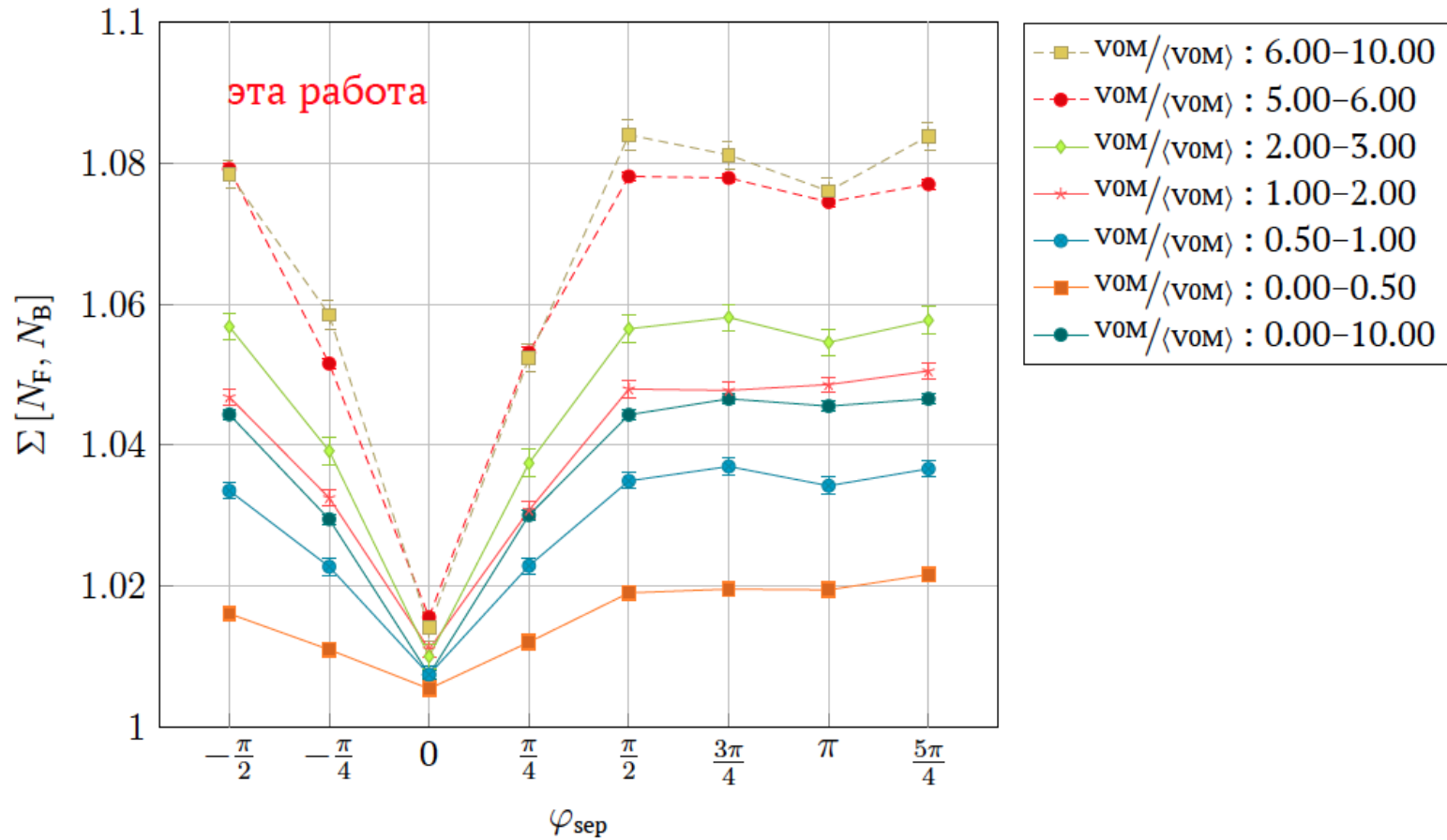


Рисунок 13 – Сечение двумерного графика $\Sigma [N_F, N_B]$ в $\eta_{\text{gap}} = 0$ для разных классов по множественности. $0.2 \text{ ГэВ} < p_T < 2.0 \text{ ГэВ}$

Обнаружен нелинейный плато величины $\Sigma [N_F, N_B]$ в зависимости от энергии и рост плато для событий с большой множественностью. Для событий с большой множественностью ($5 < v_{0M} < 6$) обнаружена особенность топологии — протяжённый вдоль η_{gap} провал для $\varphi_{\text{sep}} = 0$, который соответствует ближнему риджу в двухчастичных корреляционных функциях.

