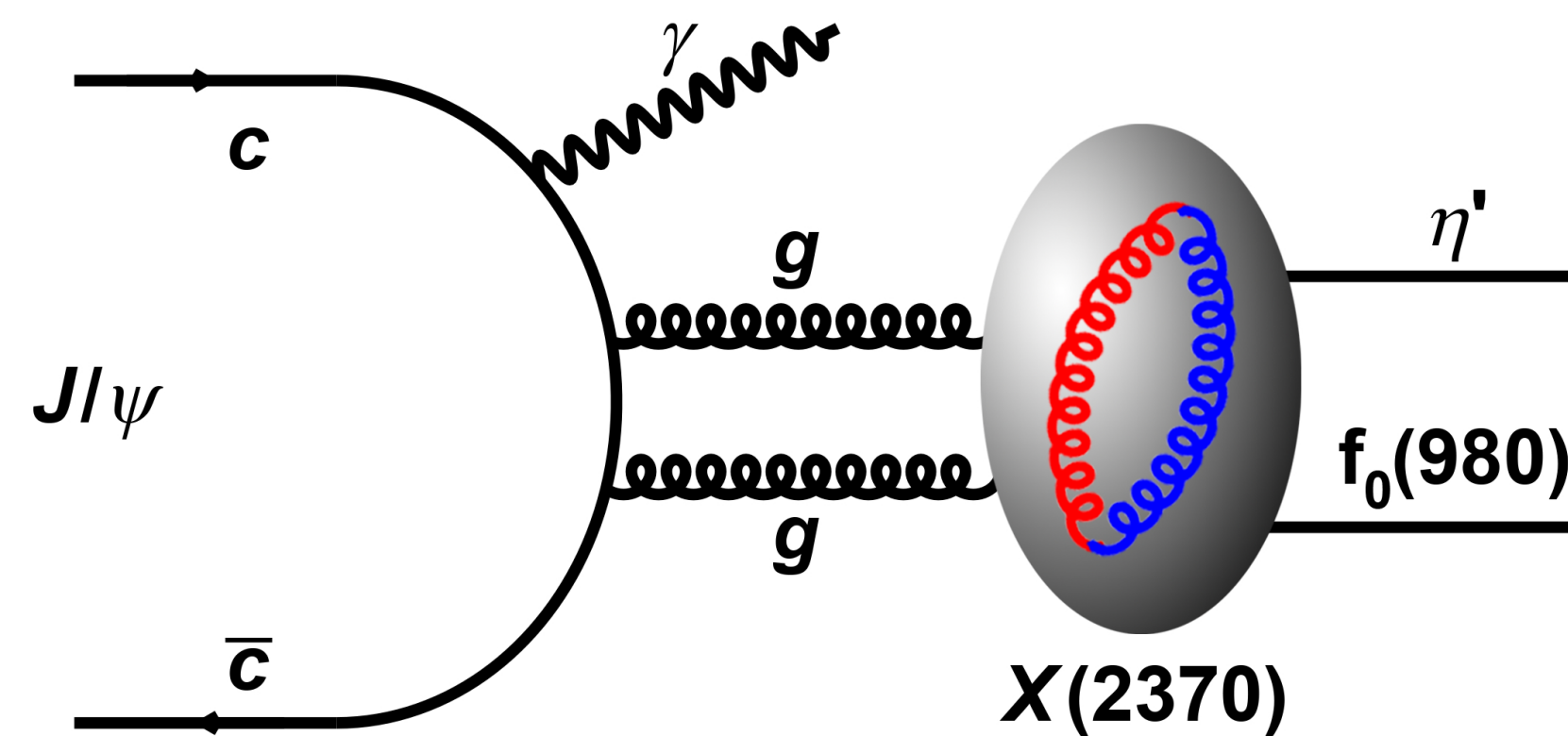


Discovery of a Glueball-like particle $X(2370)$ @ BESIII



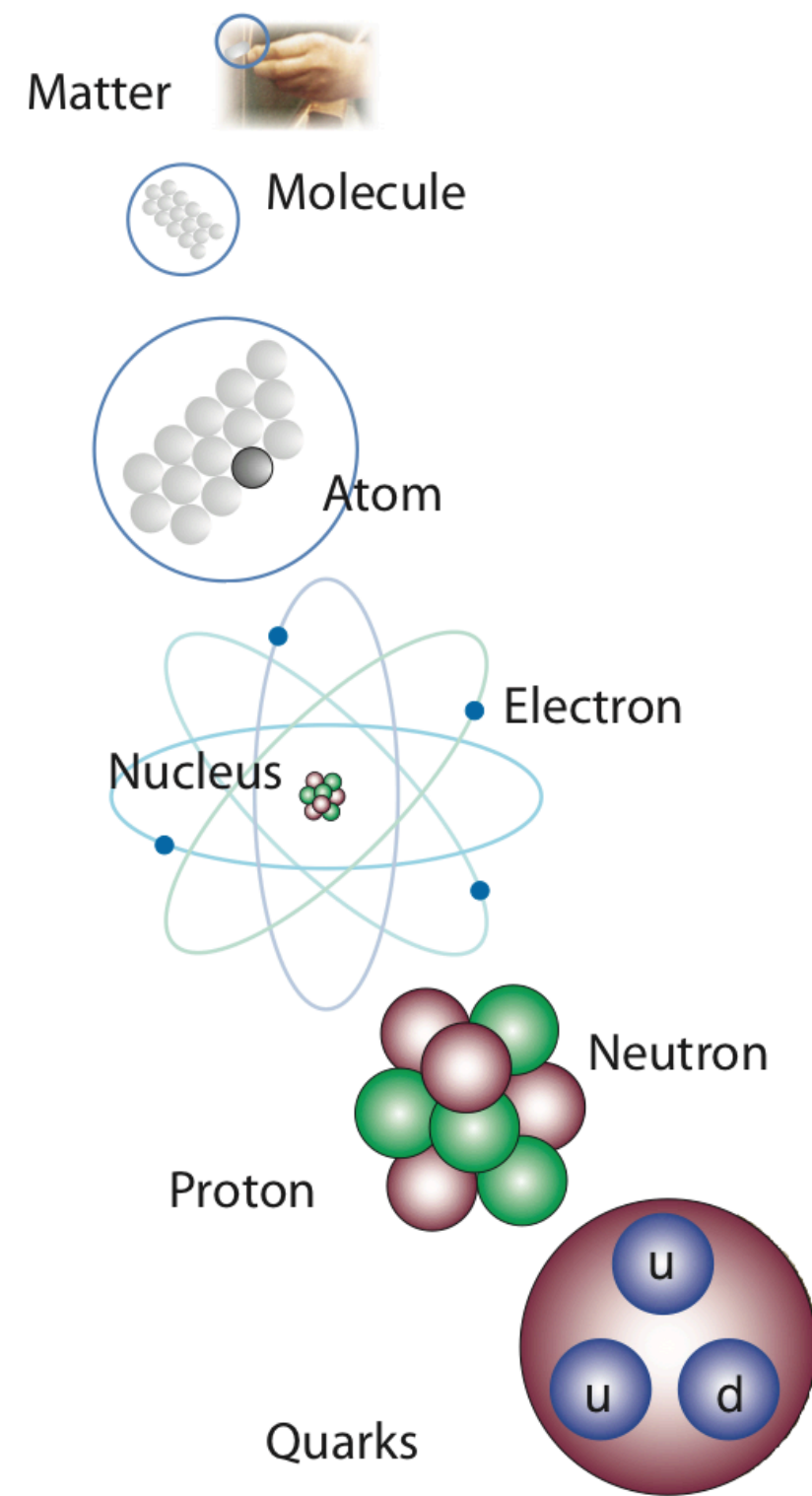
Yanping Huang

Institute of High Energy Physics, CAS

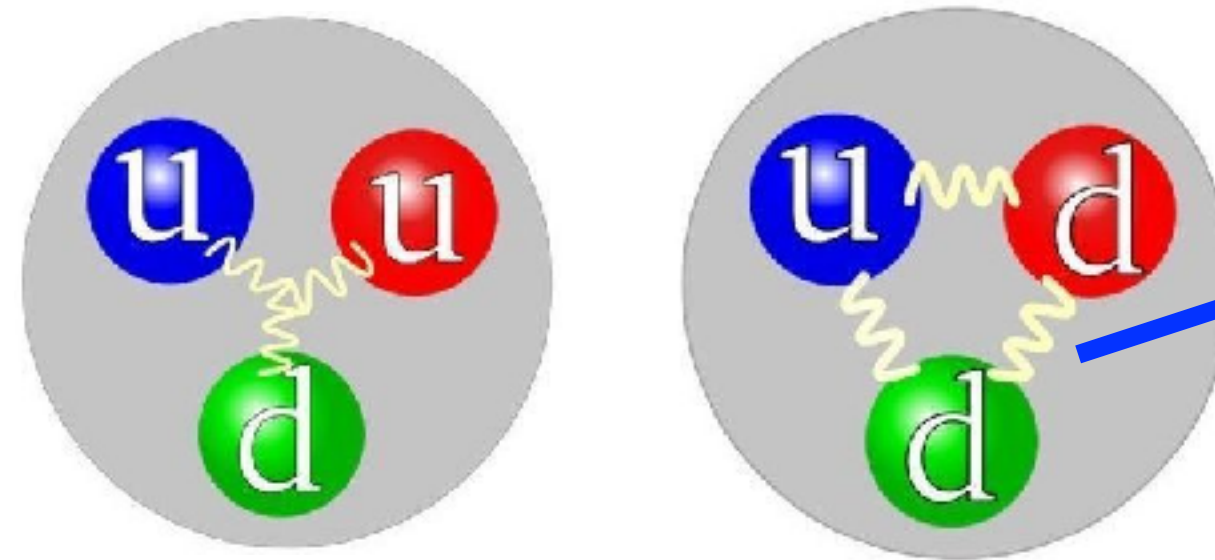
(On behalf of the BESIII Collaboration)

The Standard Model

All matters are made out of quarks and leptons



Proton: uud Neutron: udd



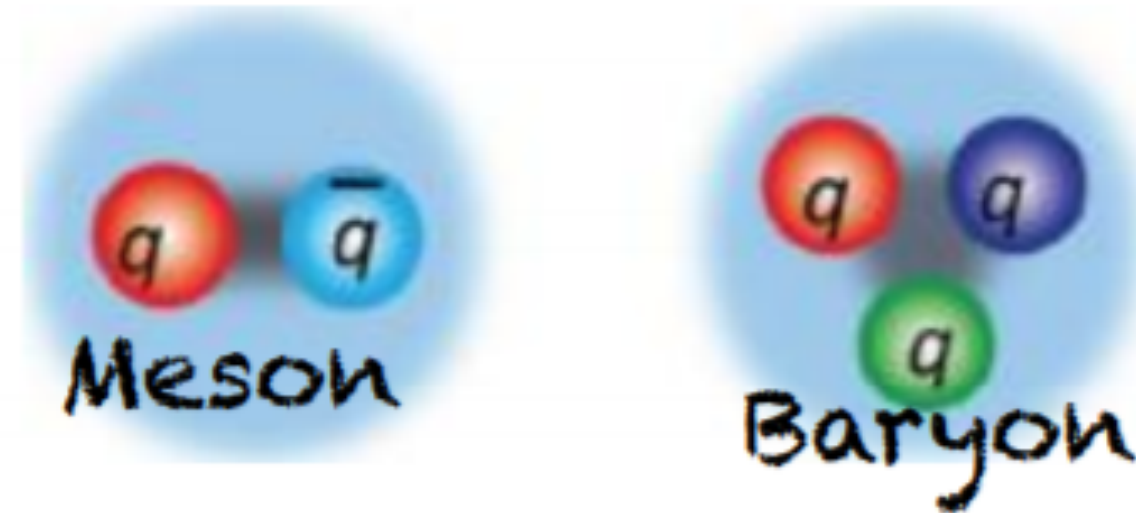
Standard Model of Elementary Particles

		three generations of matter (fermions)			interactions / force carriers (bosons)	
		I	II	III		
QUARKS	mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 125.11 \text{ GeV}/c^2$
	charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
	spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
		u up	c charm	t top	g gluon	H higgs
		d down	s strange	b bottom	γ photon	
LEPTONS	mass	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$	
	charge	-1	-1	-1	0	
	spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
		e electron	μ muon	τ tau	Z Z boson	
		ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	
					GAUGE BOSONS VECTOR BOSONS	
					SCALAR BOSONS	

- ◆ The basic theory for strong interactions is quantum chromodynamics (QCD)
 - ◆ **Gluons are the force carriers of the strong interactions**
 - ◆ **Gluon self-interaction:** prediction of non-Abelian Gauge SU(3) QCD theory

Forms of hadrons

◆ In quark model:



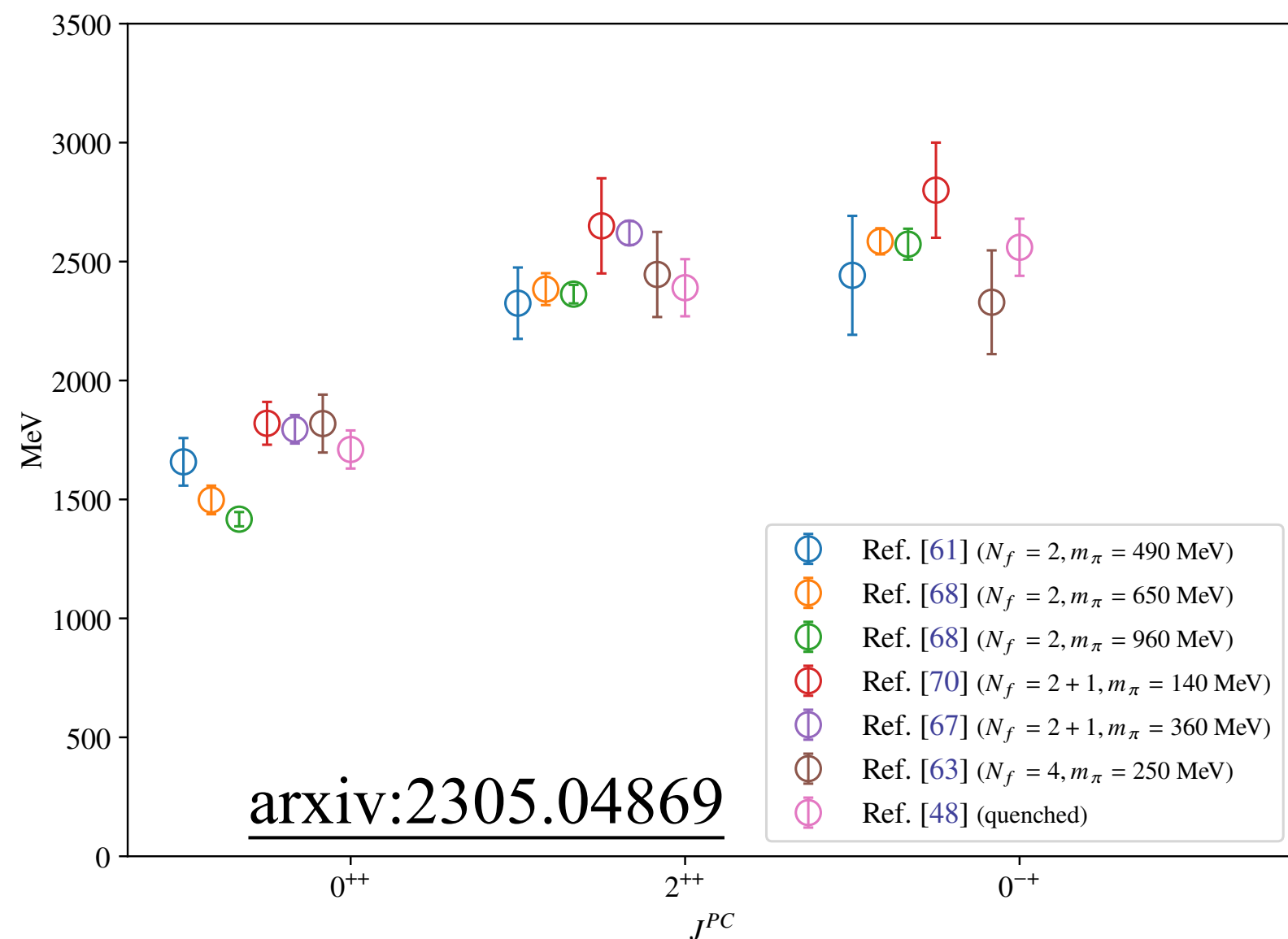
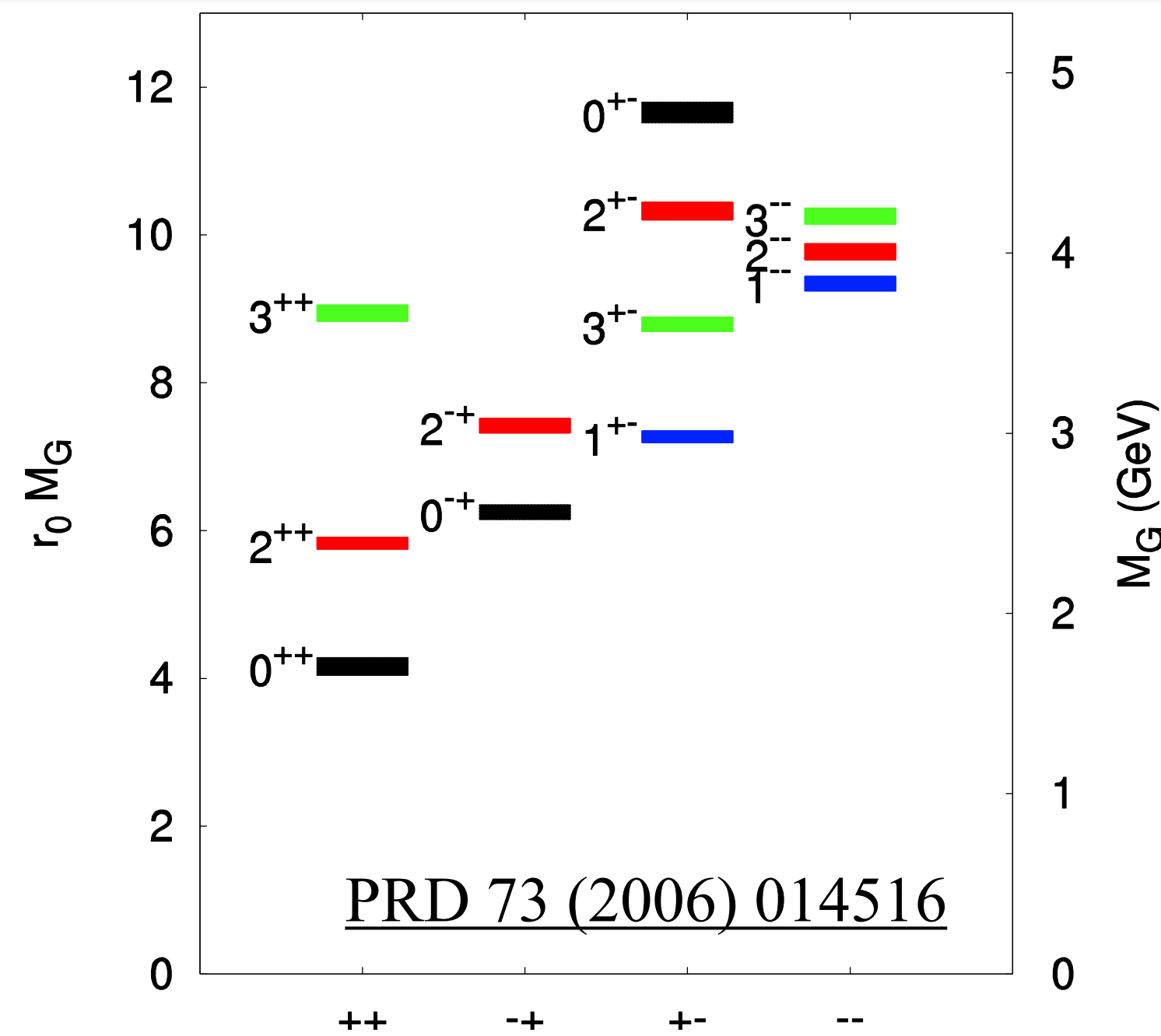
◆ Other forms of hadrons:

- ◆ **Multi-quark:** quark number ≥ 4
- ◆ **Hybrid state:** the mixture of quark and gluon
- ◆ **Glueball:** composed of gluons (**gg, ggg, gggg**)



Glueballs are unique particles formed by gluons (force carriers) due to non-Abelian gauge self-interactions of gluons

Theory Prediction



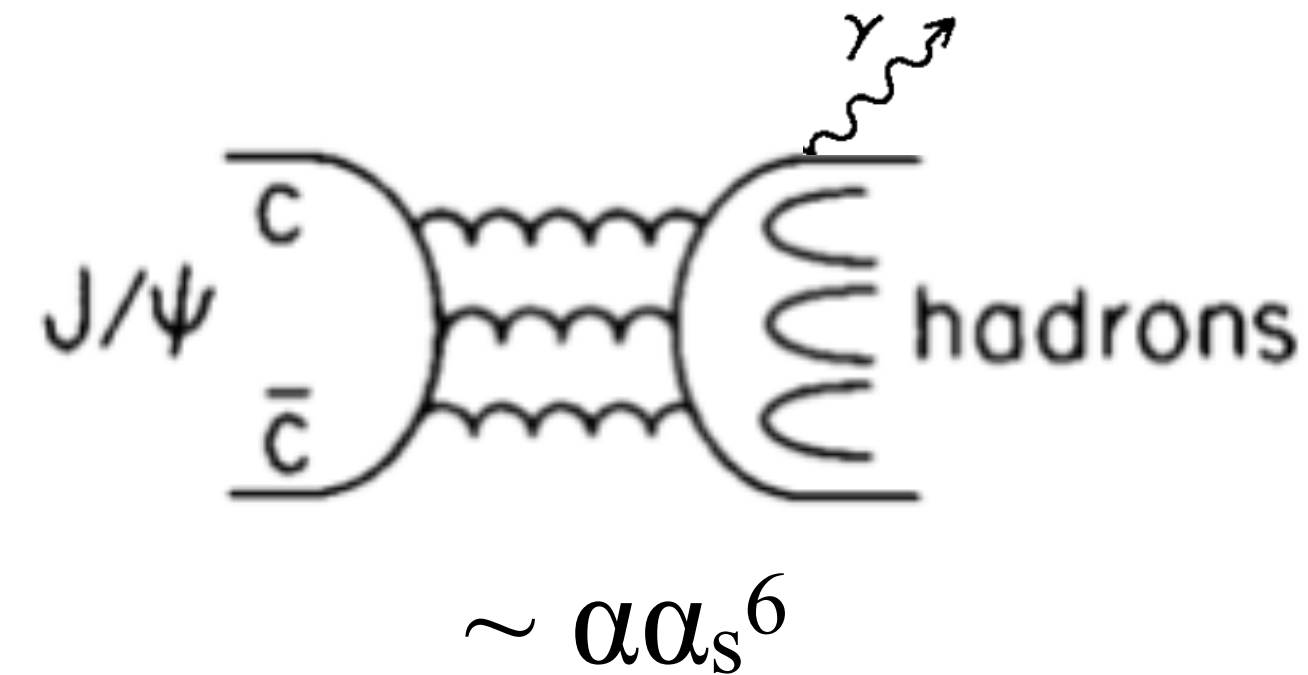
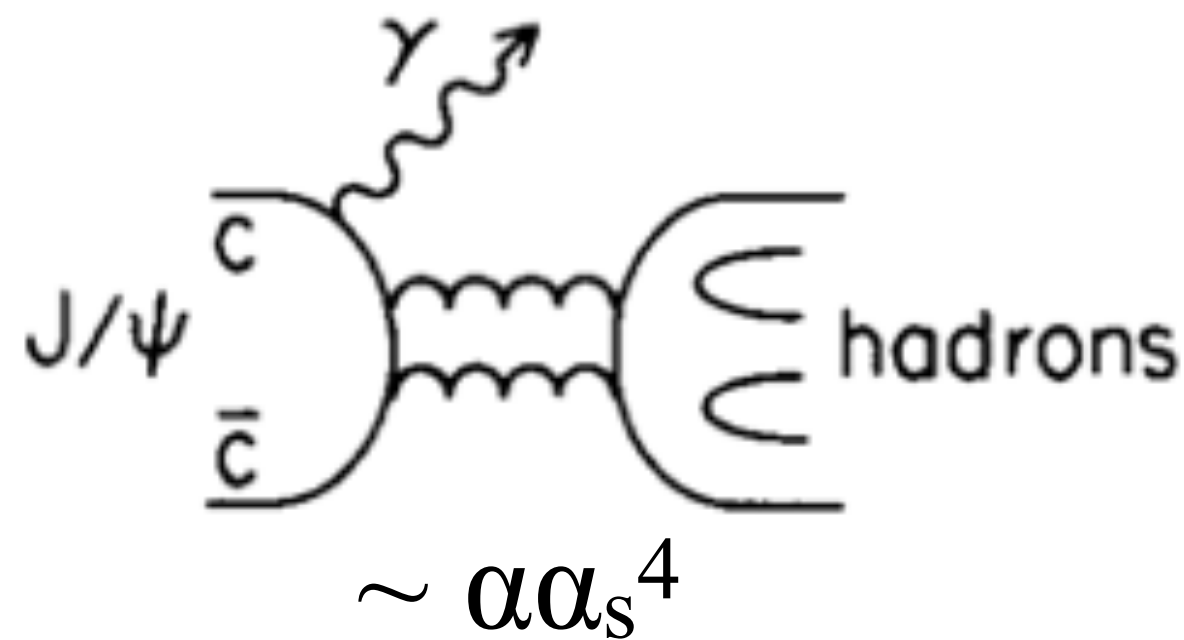
- ◆ **Quenched lattice QCD** is the non-perturbative method for this theory from the first principles.
- ◆ The predictions of **masses and production rates** of pure glueballs are expected to be reliable.
- ◆ Lattice QCD predictions:
 - ◆ **0⁺⁺ ground state: 1.5 - 1.7 GeV/c²**
 - ◆ **2⁺⁺ ground state: 2.3 - 2.4 GeV/c²**
 - ◆ **0⁻⁺ ground state: 2.3 - 2.6 GeV/c²**

Glueball Search

- ◆ **Many experiments searched for glueballs over the past 4 decades, mostly in J/ψ radiative decays from electron-positron collision**
- ◆ The advantage for **glueball searches via J/ψ radiative decays from electron-positron collision**
- ◆ **Glueball production in J/ψ radiative decays**
- ◆ **Glueball decays**
- ◆ **Many historical glueball candidates, but also some difficulties**

J/ψ radiative decays

◆ Gluon rich environment



◆ **Isospin filter:** final states dominated by **I=0** processes

◆ **Spin-parity filter:** **C** parity must be +, so $J^{PC}=0^{-+}, 0^{++}, 1^{++}, 2^{++}, 2^{-+} \dots$

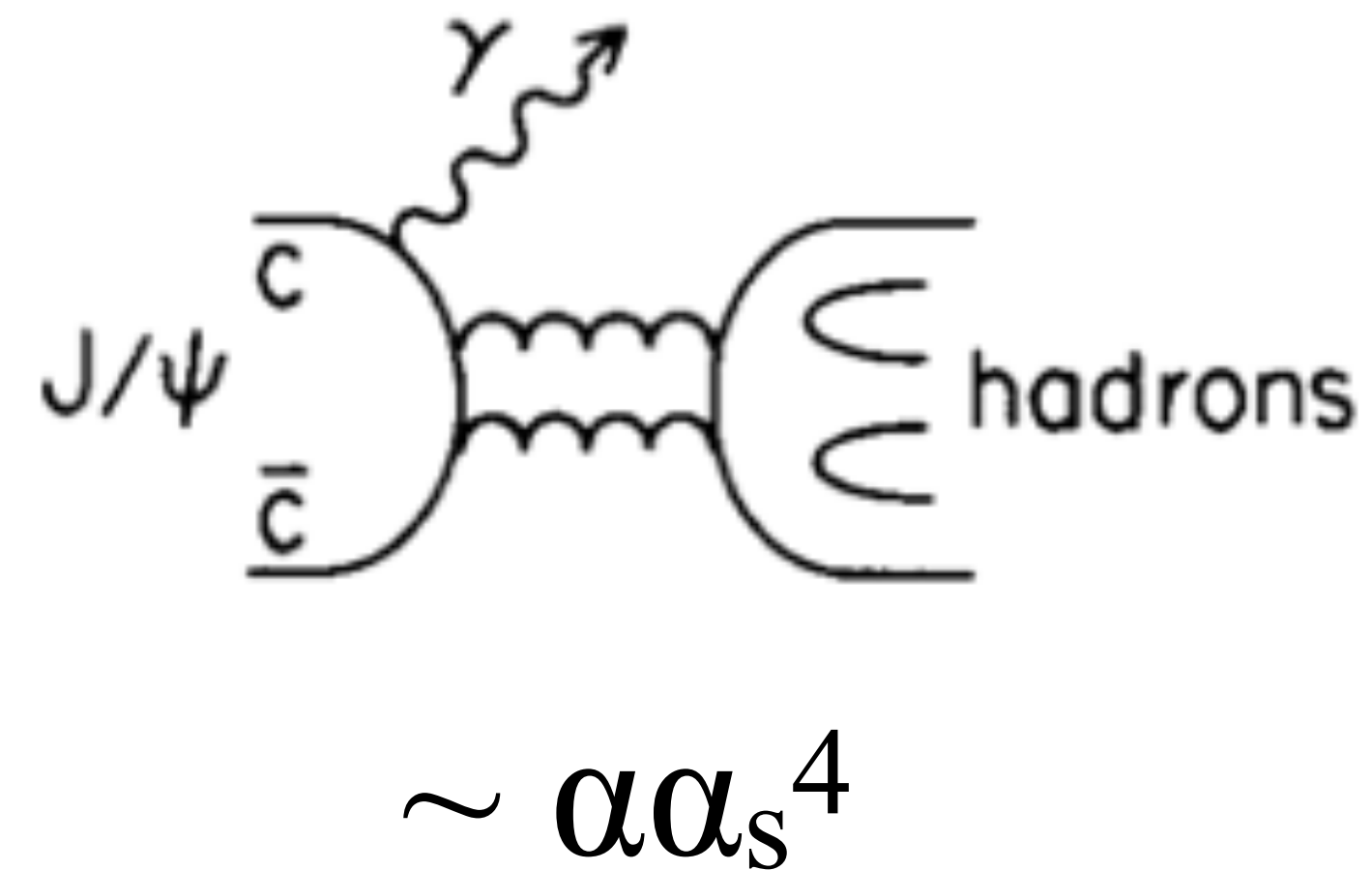
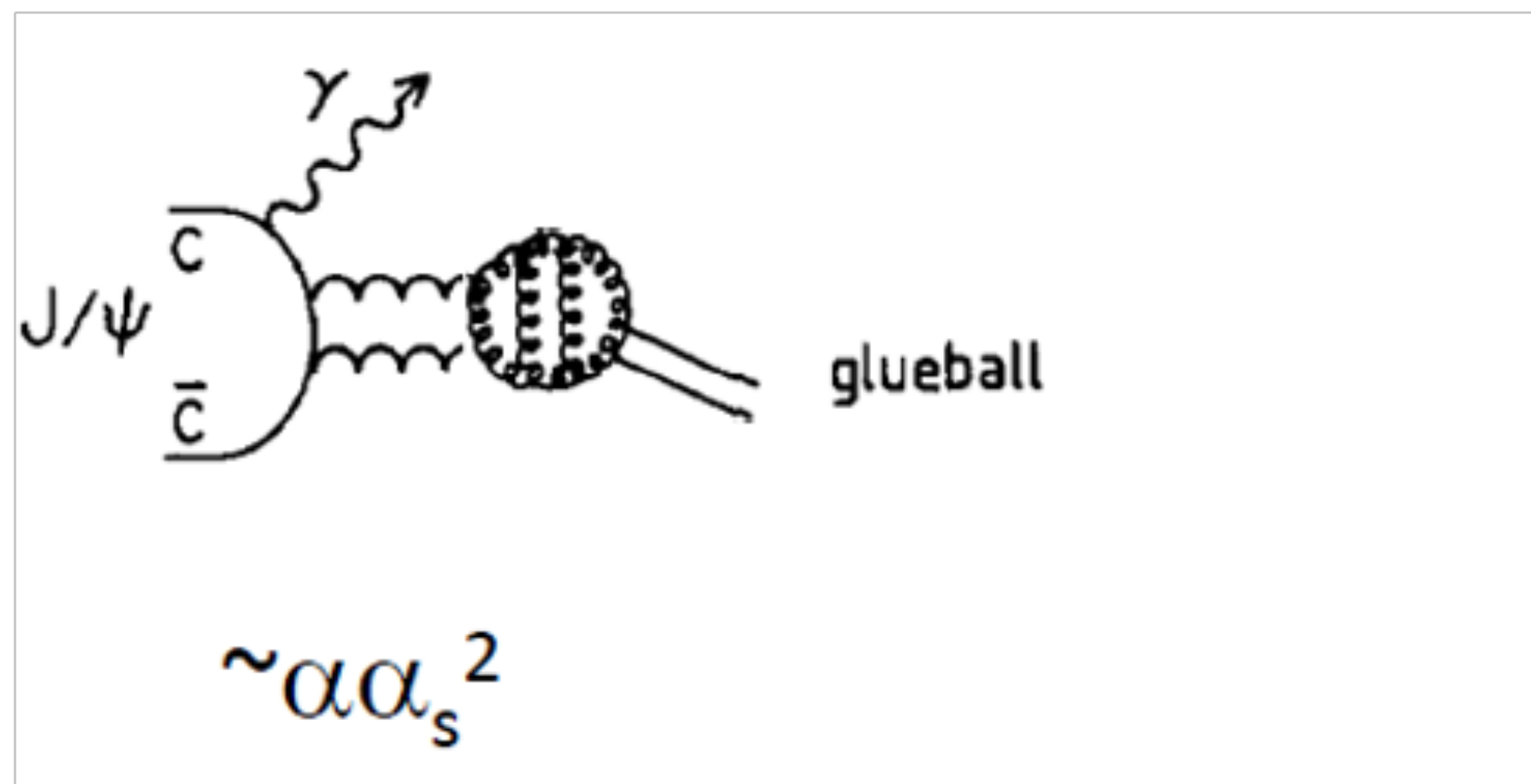
◆ **Clean environment in electron-positron collision:** very different from proton-antiproton collision

➔ **Ideal place to search for glueballs**

Glueball Production in J/ψ radiative decays

◆ Rich production in J/ψ radiative decays:

- ◆ Glueball production rate in J/ψ radiative decays could be higher than normal hadrons



Glueball Decays

- ◆ **Flavor symmetric decays**
- ◆ **No rigorous predictions on decay patterns and their branching ratios**
- ◆ The glueball decays could have **similar decays to the Charmonium families** since both of them can only decay via gluons
 - ◆ e.g. the 0^{-+} glueball could have similar decays of η_c
 - ◆ One of the favorite decay modes of η_c is $\pi\pi\eta'$, so $J/\psi \rightarrow \gamma \pi\pi\eta'$ could be a good place to search for the 0^{-+} glueball

Historical Glueball Candidates — 0^{++} scalar Glueball

- ◆ **$f_0(1710)$** : discovered by MarkII in 1980's as $\theta_2(1720)$ and $J^{PC} = 2^{++}$ from a simple fit to the angular distribution. Lots of studies at MarkII, DM2, BES I, BES II, BES III
- ◆ **J^{PC} was firstly changed to 0^{++}** on a full PWA of $J/\psi \rightarrow \gamma KK$ @ BES II
- ◆ The high production rate of $J/\psi \rightarrow \gamma f_0(1710)$ and the suppression of $f_0(1710) \rightarrow \eta\eta'$ strongly support its interpretation:
 - **A scalar glueball or large glueball content if it is a mixture of glueball and normal meson**
- ◆ **With PS subtracted, $\Gamma(f_0(1710) \rightarrow \pi\pi:KK) = 1:2.43$** , which should be 1:1 for a pure glueball decays
- ◆ **Difficulty:** What causes the **flavor symmetric breaking** needs to be understood from first principle of QCD (not just phenomenological understanding).

$$\begin{aligned} & \bullet B[J/\psi \rightarrow \gamma f_0(1710) \rightarrow \gamma\pi\pi] = (4.01 \pm 1.0) \times 10^{-4} \\ & \bullet \text{BESII: PLB 642 441 (2006)} \\ & \bullet B[J/\psi \rightarrow \gamma f_0(1710) \rightarrow \gamma K_S^0 K_S^0] = (2.00_{-0.02}^{+0.03+0.31}) \times 10^{-4} \\ & \bullet \text{BESIII: PRD 98 072003 (2018)} \end{aligned}$$

$$\begin{aligned} & B[f_0(1500) \rightarrow \eta\eta'] / B[f_0(1500) \rightarrow \pi\pi] = (1.66_{-0.40}^{+0.42}) \times 10^{-1} \\ & B[f_0(1710) \rightarrow \eta\eta'] / B[f_0(1710) \rightarrow \pi\pi] < (2.9_{-0.9}^{+1.1}) \times 10^{-3} \\ & \bullet \text{BESIII: PRD 106 072012 (2022)} \end{aligned}$$

A lot of studies from BES in $J/\psi \rightarrow \gamma KK, \gamma\pi\pi, \gamma\eta\eta, \gamma\eta\eta'$

Historical Glueball Candidates — 0^{++} scalar Glueball

- ◆ **$f_0(1500)$** : discovered by Crystal Barrel in 1990's as a unique 0^{++} candidate since $f_0(1710)$ was f_2 at that time.
 - ◆ Difficulty: compared with $f_0(1710)$, much lower production rate than $f_0(1710)$ disfavors its interpretation as a scalar glueball
- ◆ **Mixing between $f_0(1500)/f_0(1710)$, or even with $f_0(1790)$**
 - ◆ Difficulty: dynamic mixing mechanism needs to be understood from the first principle of QCD (not just phenomenological understanding).

BESIII:

$$B(J/\psi \rightarrow \gamma f_0(1500)) \sim 0.29 \times 10^{-3}$$

$$B(J/\psi \rightarrow \gamma f_0(1710)) \sim 2.2 \times 10^{-3}$$

Historical Glueball Candidates — 2^{++} Tensor Glueball

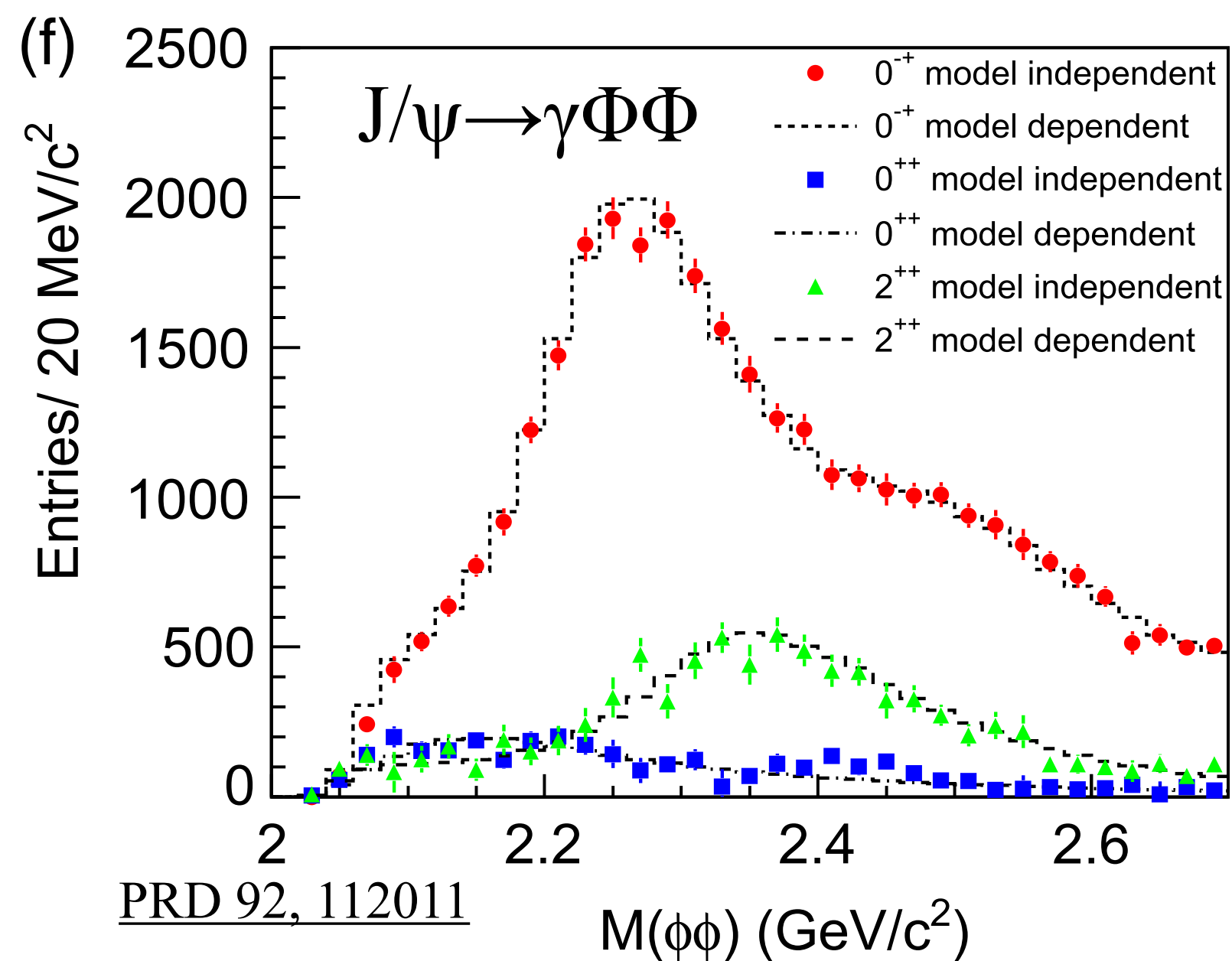
◆ $\xi(2230)$

- ◆ First observed by MarkIII is $J/\psi \rightarrow \gamma KK$ in 1980's, then by BES I in 1990's in $J/\psi \rightarrow \gamma KK, \gamma \pi\pi, \gamma p\bar{p}$ with very narrow mass peak.
- ◆ It was a tensor glueball candidate due to good flavor symmetric decay property.
- ◆ Difficulty: it was not confirmed by BES II, nor BES III with much higher statistics.

Historical Glueball Candidates — 2^{++} Tensor Glueball

◆ $f_2(2340)$

- ◆ **Its large production rate** in $J/\psi \rightarrow \gamma(KK/\eta\eta/\eta'\eta'/\phi\phi)$ favors its interpretation as a **tensor glueball**.
- ◆ **Many wide f_2 mesons in the mass region of 2.3 GeV** of 2^{++} glueball mass from the LQCD predictions
- ◆ **Difficulty: no clear mass peak of these f_2 mesons** due to large overlaps among various wide resonances. (PWA components)
- **More PWA studies are needed to check the consistency among various decays modes.** However, due to large overlaps again, no independent mass and width scan can be performed in PWA.



Resonance	M (MeV/c ²)	Γ (MeV/c ²)	B.F. ($\times 10^{-4}$)	Sig.
$\eta(2225)$	2216^{+4+21}_{-5-11}	185^{+12+43}_{-14-17}	$(2.40 \pm 0.10^{+2.47}_{-0.18})$	28σ
$\eta(2100)$	2050^{+30+75}_{-24-26}	$250^{+36+181}_{-30-164}$	$(3.30 \pm 0.09^{+0.18}_{-3.04})$	22σ
$X(2500)$	$2470^{+15+101}_{-19-23}$	230^{+64+56}_{-35-33}	$(0.17 \pm 0.02^{+0.02}_{-0.08})$	8.8σ
$f_0(2100)$	2101	224	$(0.43 \pm 0.04^{+0.24}_{-0.03})$	24σ
$f_2(2010)$	2011	202	$(0.35 \pm 0.05^{+0.28}_{-0.15})$	9.5σ
$f_2(2300)$	2297	149	$(0.44 \pm 0.07^{+0.09}_{-0.15})$	6.4σ
$f_2(2340)$	2339	319	$(1.91 \pm 0.14^{+0.72}_{-0.73})$	11σ
0^{++} PHSP			$(2.74 \pm 0.15^{+0.16}_{-1.48})$	6.8σ

Historical Glueball Candidates — 0^{-+} Pseudoscalar Glueball

◆ $\eta(1405)$

- ◆ first discovered by MarkII in 1980's, named as $\eta(1440)$ with complicated structures. Lots of studies at MarkII, MarkIII, DM2 and BES.
- ◆ Believed as the first glueball candidate due to its large production rate in J/ψ radiative decays and **lack of reliable LQCD predictions in 1980's**
- ◆ No longer to be 0^{-+} glueball candidate due to its large different mass from LQCD prediction.

Historical Difficulties in Glueball Searches

◆ Theoretically:

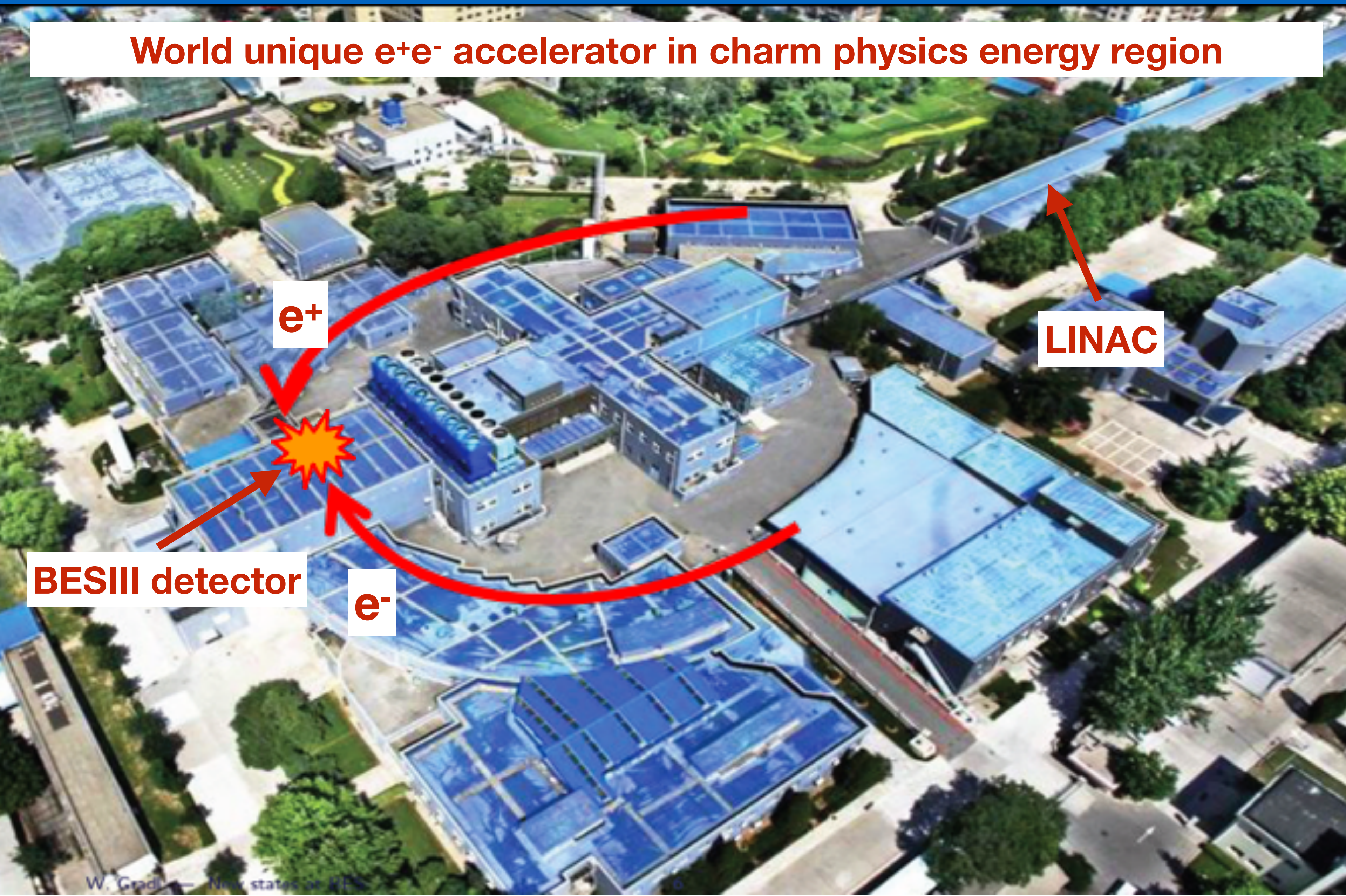
- ◆ **No prediction on the decay branch ratios** so far (even the order)
- ◆ **Very rare prediction on the glueball production rate** $\Gamma(J/\psi \rightarrow \gamma G)$
- ◆ **Mix** with qqbar mesons or even with 4q, qqg, mesons? **Mixing dynamics?**

◆ Experimentally:

- ◆ Data sample was not big enough
- ◆ No good way **modeling background** in many cases.
- ◆ **Interference among mesons** makes the analysis more complicated:
 - **PWA is a must, but it is complicated and takes a quite long time.**

Beijing Electron Positron Collider (BEPCII)

World unique e^+e^- accelerator in charm physics energy region



2004: Construction

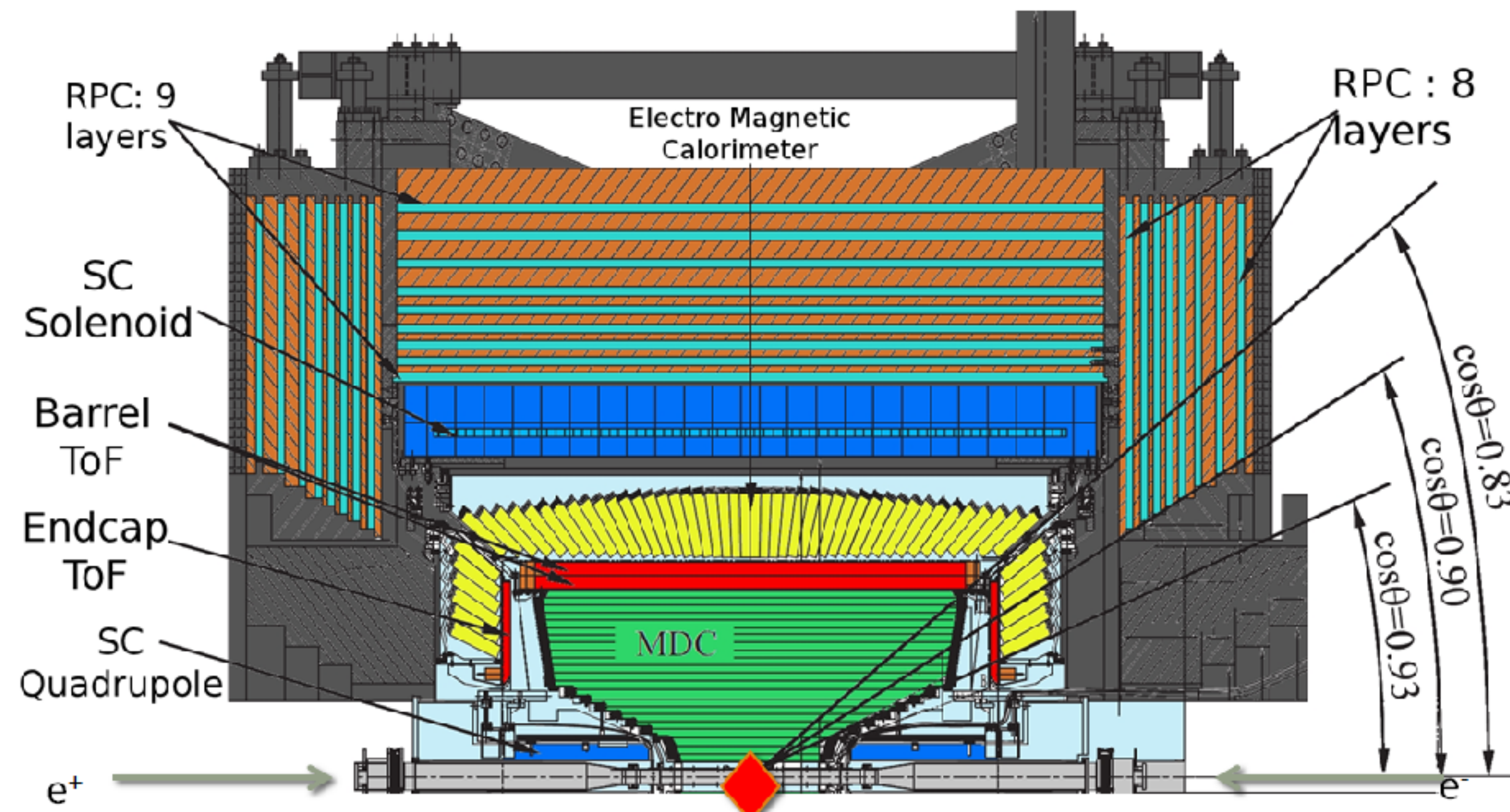
- Double rings
- Beam energy:
1.0 - 2.3 (2.45) GeV
- Designed luminosity:
 $1 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$

2008: test run

2009-now: BESIII physics runs

BESIII detector

Designed for neutral and charged particle with excellent resolution, PID, and large coverage



Total weight 730 ton,
~40,000 readout channel
Data rate: 5kHz, 50Mb/s

- ◆ Magnet: 1T Super conducting
- ◆ MDC: small cell & He gas
 - $\sigma_{xy} = 130 \mu\text{m}$
 - $\sigma_p/p = 0.5\% @ 1\text{GeV}$
 - $dE/dx = 6\%$
- ◆ TOF: plastic scintillator/MRPC
 - $\sigma_T = 80 \text{ ps}$ Barrel
 - $\sigma_T = 110 (60) \text{ ps}$ Endcap
- ◆ EMC: CsI crystals
 - $\Delta E/E = 2.5\% @ 1\text{GeV}$ - Barrel
 - $\Delta E/E = 5\% @ 1\text{GeV}$ - Endcap
- ◆ Muon ID: 9 layer RPC

Has been in full operation since 2008, all sub-detectors are in very good status!

BESIII Collaboration

Political Map of the World, November 2011



BESIII

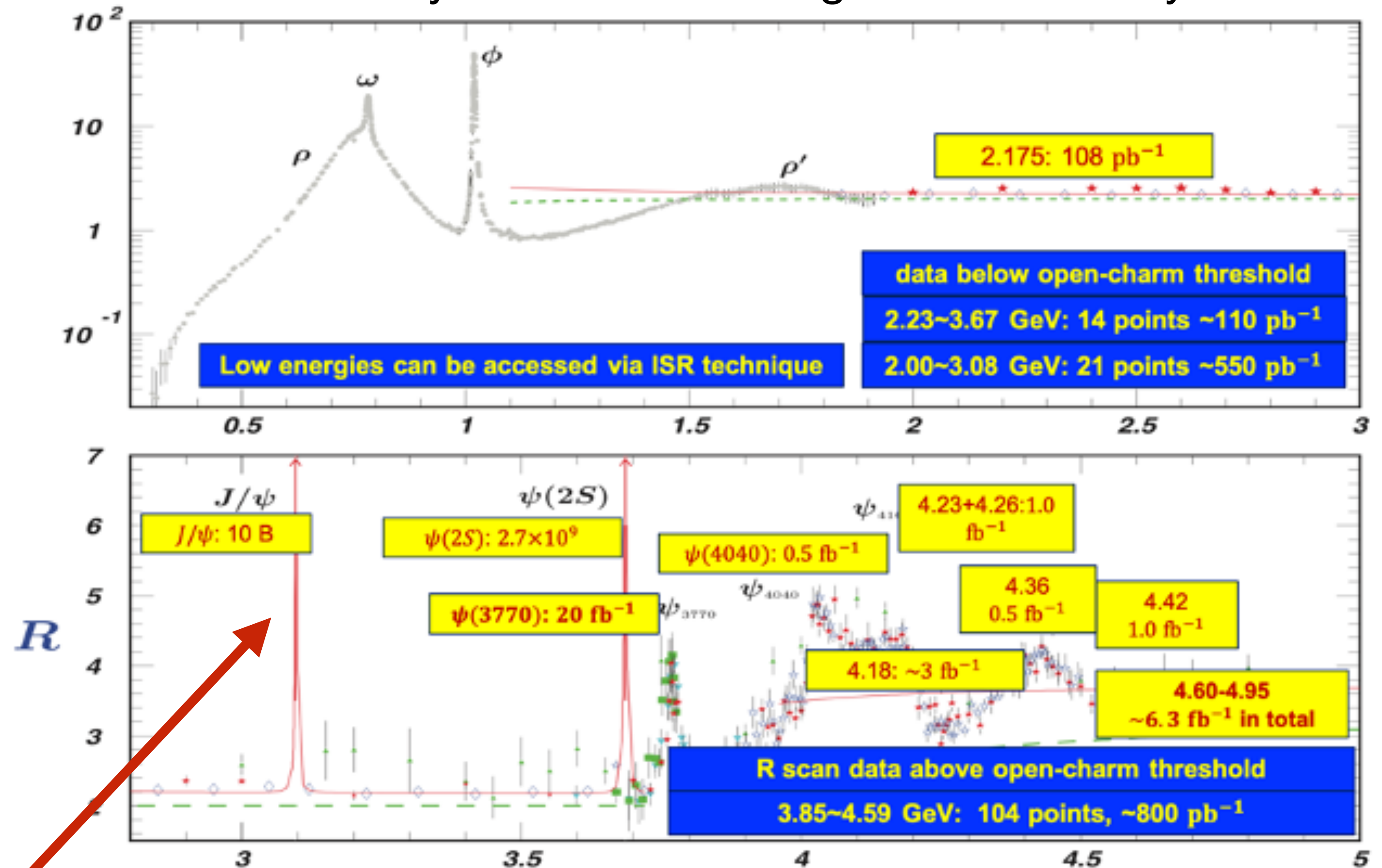
~600 members
(more than 130 from outside of China)
From 89 institutions in 17 countries

BESIII Data samples

Totally about 50fb⁻¹ integrated luminosity

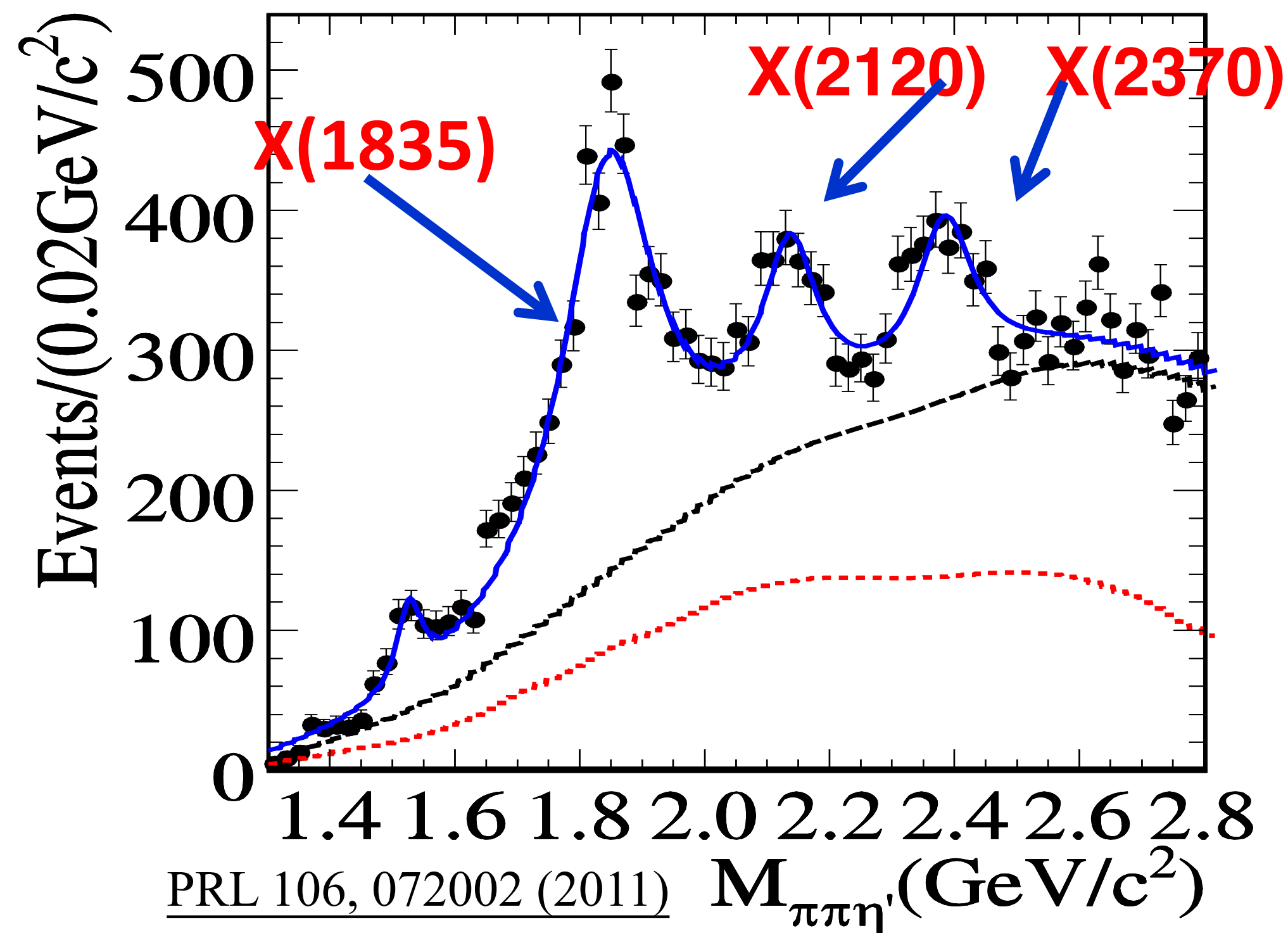
Data sets collected so far include

- ◆ 10×10⁹ J/ψ events
- ◆ 2.7×10⁹ ψ(2S) events
- ◆ 20 fb⁻¹ ψ(3770)
- ◆ Scan data between 2.0 and 3.08 GeV, and above 3.74GeV
- ◆ Large datasets for XYZ studies:
Scan with >500pb⁻¹ per energy point space 10-20MeV apart



World largest J/ψ data sample : ~10 billion

Observation of the X(2370) in 2011

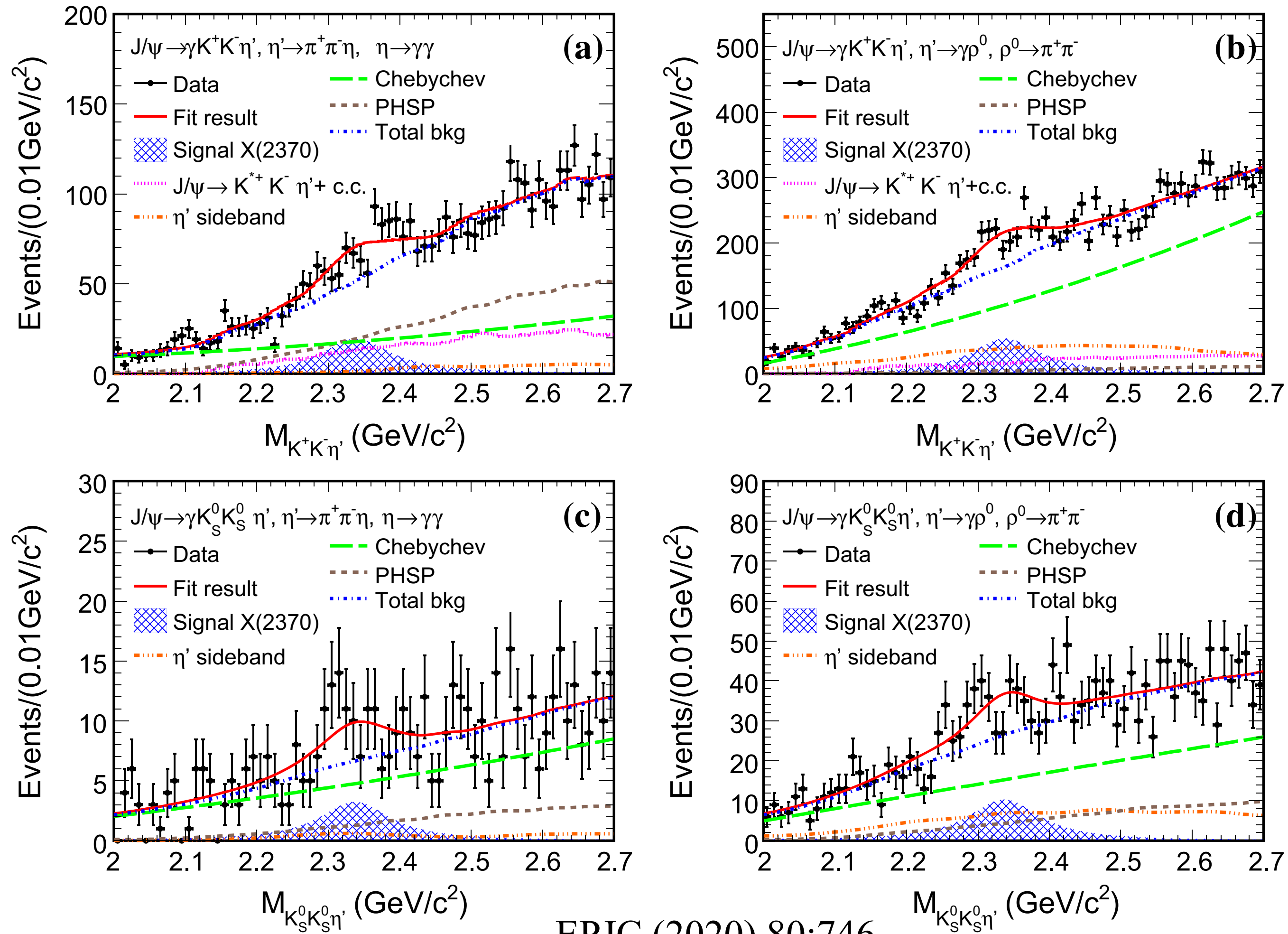


$J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$ With $\sim 225\text{M}$ J/ψ events

	$M(\text{MeV}/c^2)$	$\Gamma(\text{MeV}/c^2)$	Sig.
X(1835)	$1836.5 \pm 3.0^{+5.6}_{-2.1}$	$190.1 \pm 9.0^{+38}_{-36}$	$>20\sigma$
X(2120)	$2122.4 \pm 6.7^{+4.7}_{-2.7}$	$83 \pm 16^{+31}_{-11}$	7.2σ
X(2370)	$2376.3 \pm 8.7^{+3.2}_{-4.3}$	$83 \pm 17^{+44}_{-6}$	6.4σ

- ◆ **Discovery of X(2370) in $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$ with the statistic significance of 6.4σ**
- ◆ **First observation of one particle: a good candidate for 0^{++} glueball**
- ◆ **Mass, production and decay property are consistent with the LQCD prediction**

Confirmation of the X(2370) in $J/\psi \rightarrow \gamma K K \eta'$



EPJC (2020) 80:746

◆ Combination with 1.31×10^9 J/ψ events

- $J/\psi \rightarrow \gamma K^+ K^- \eta'$ and $J/\psi \rightarrow \gamma K_S^0 K_S^0 \eta'$
- $\eta' \rightarrow \gamma \pi \pi$ and $\eta' \rightarrow \pi \pi \eta$

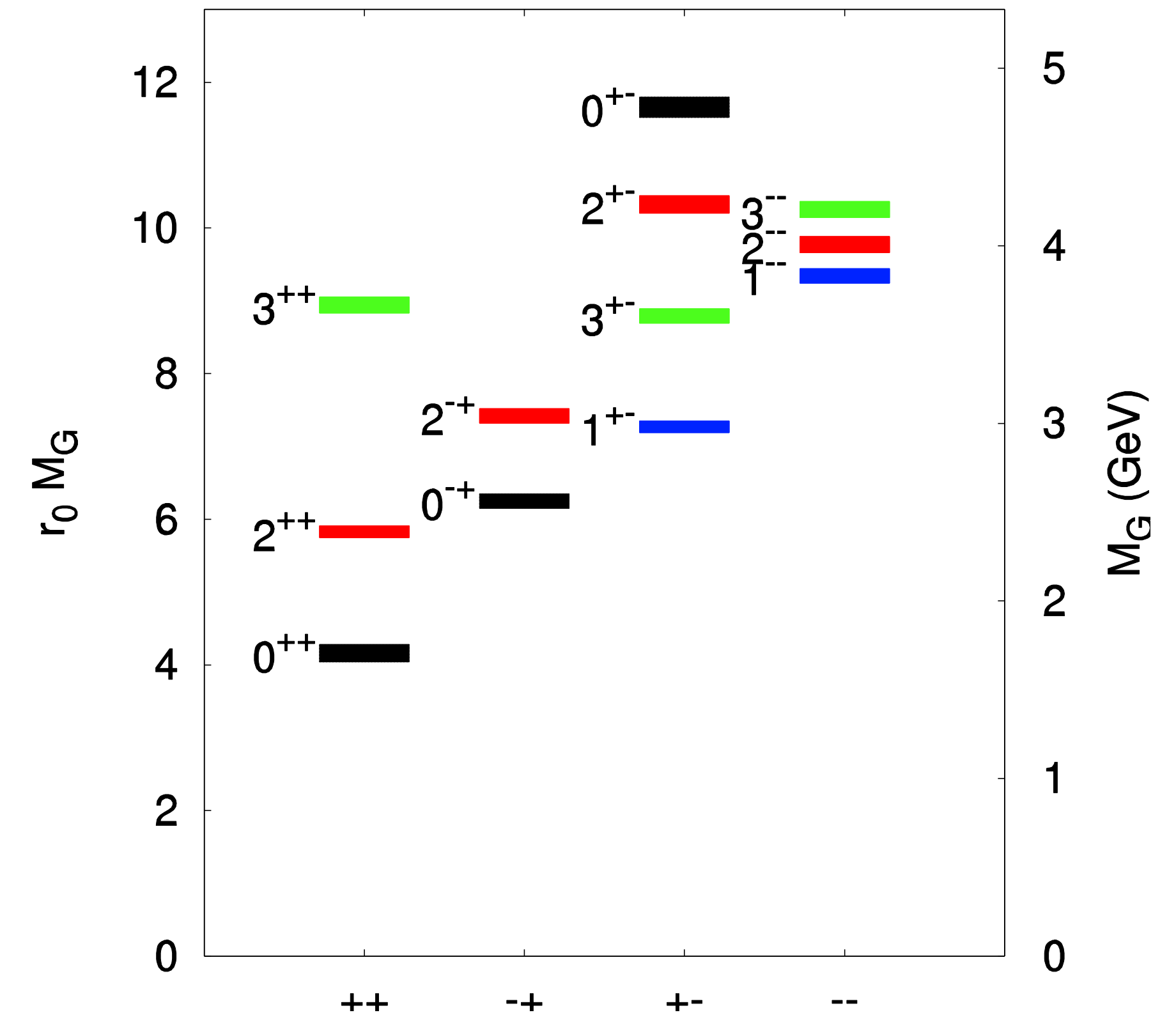
◆ Confirmation of the X(2370) with 8.3σ

- $M = 2341.6 \pm 6.5(\text{stat.}) \pm 5.7(\text{syst.}) \text{ MeV}$
- $\Gamma = 117 \pm 10(\text{stat.}) \pm 8(\text{syst.}) \text{ MeV}$
- $\text{Br}(J/\psi \rightarrow \gamma X(2370) \rightarrow \gamma K^+ K^- \eta') = (1.79 \pm 0.23 \pm 0.65) \times 10^{-5}$
- $\text{Br}(J/\psi \rightarrow \gamma X(2370) \rightarrow \gamma K_S^0 K_S^0 \eta') = (1.18 \pm 0.32 \pm 0.39) \times 10^{-5}$

Observation: X(2370) new decay mode of $KK\eta'$

X(2370) - good candidate of 0^{-+} glueball

- ◆ Its mass is consistent with LQCD prediction on the 0^{-+} glueball
- ◆ Observed in flavor symmetric decay modes of $\pi^+\pi^-\eta'$ and $KK\eta'$ — favorite decay modes of 0^{-+} glueball
- ◆ We need to know its spin-parity



Improved situation in Glueball Searches

◆ Theoretically:

◆ Guidance from η_c decays

◆ Now we have prediction on glueball production rate from LQCD: $B(J/\psi \rightarrow \gamma G_{0^{++}}) = 2.31 \pm 0.80 \times 10^{-4}$

◆ Luckily, for the X(2370), there is **no other 0^{++} resonance nearby** (in $\sim 200\text{MeV}$ range) in J/ψ radiative decays

◆ Experimentally:

◆ World largest J/ψ data sample : ~ 10 billion

◆ Physics channels with few background

◆ GPU technique helps to speed up PWA [J.Phys.Conf. Ser. 219, 042031]

◆ It takes a long time in PWA for the complicated interference and comprehensive test of different combinations

Spin-Parity determination of the $X(2370)$ in $J/\psi \rightarrow \gamma K_s^0 K_s^0 \eta'$

Make use of three advantages:

- ◆ **Clean $J/\psi \rightarrow \gamma K_s^0 K_s^0 \eta'$ process**

- ◆ **Almost no background:** possible dominant background processes of $J/\psi \rightarrow \pi^0 K_s^0 K_s^0 \eta'$ and $J/\psi \rightarrow K_s^0 K_s^0 \eta'$ are forbidden by **exchange symmetry and C-parity conservation.**

- ◆ **~10B clean J/ψ events**

- ◆ **High-precision efficiency and resolution of charged particles and photons:**
good reconstruction for K_s^0/η'

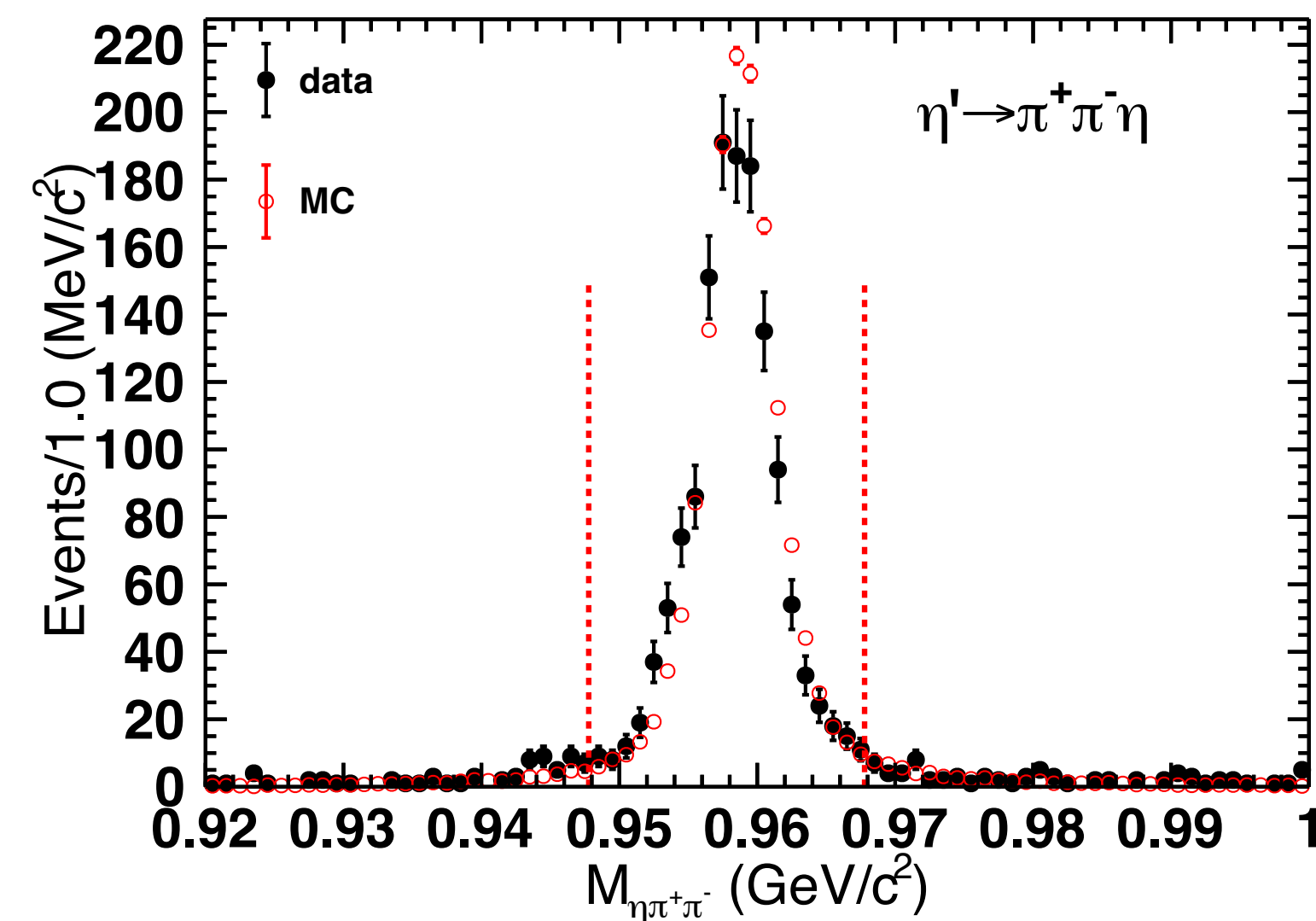
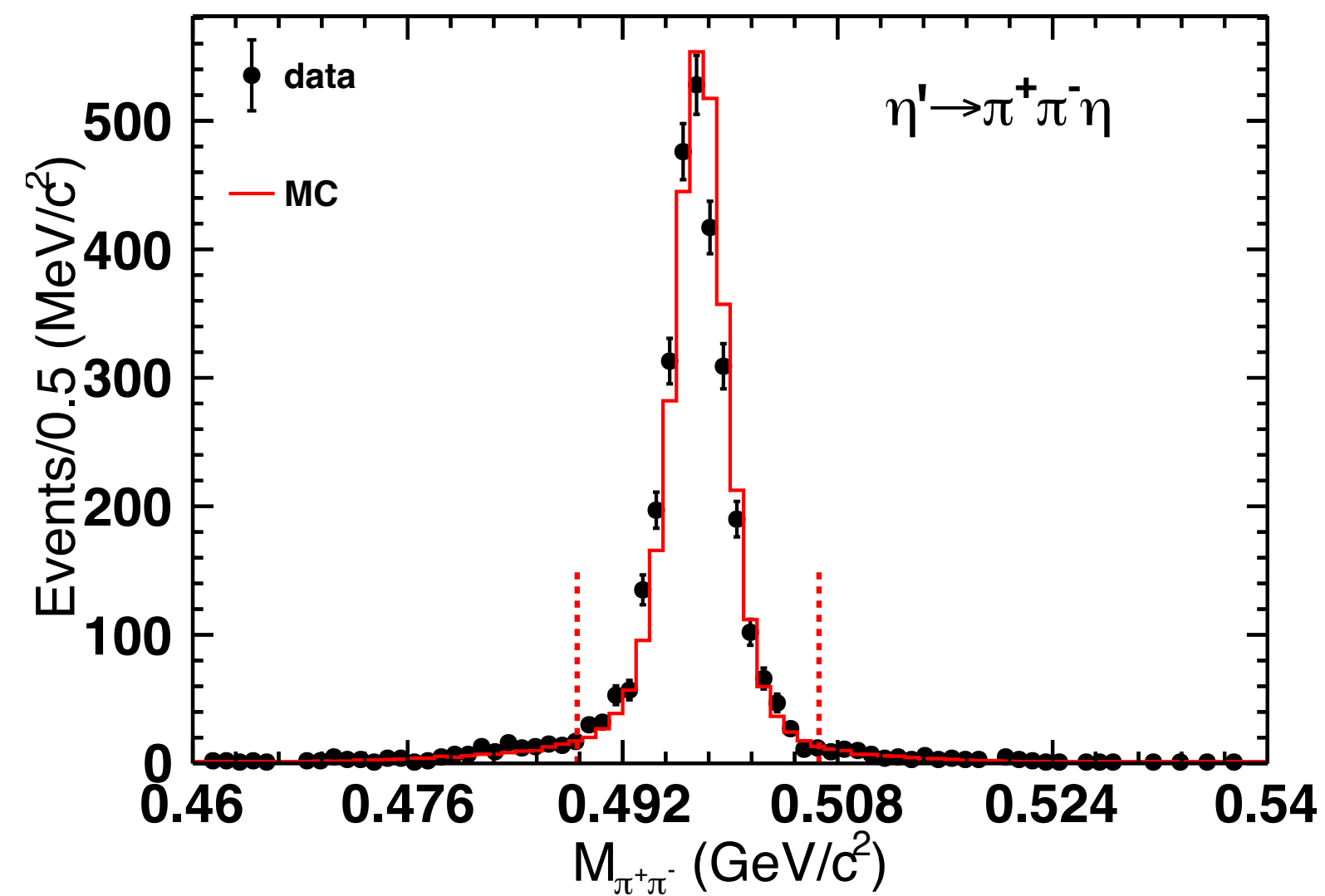
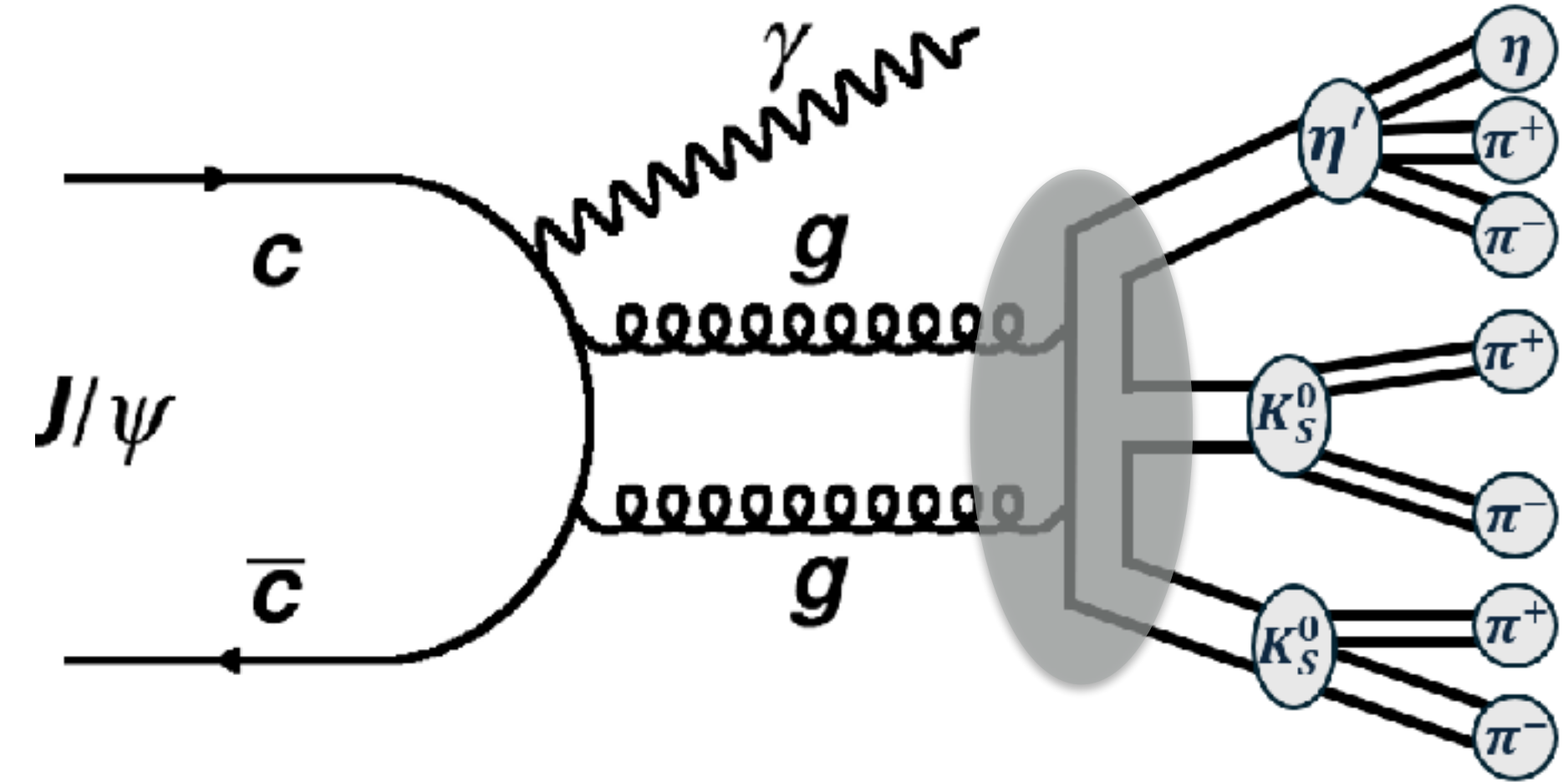
Selection for $J/\psi \rightarrow \gamma K^0_s K^0_s \eta'$, $\eta' \rightarrow \pi^+ \pi^- \eta$

Signal selection:

- At least 3 charged pairs + 3 photons
- Constraint kinematic fit with energy-momentum conservation
- K^0_s reconstruction: $|M_{\pi\pi} - m_{K^0_s}| < 9 \text{ MeV}/c^2$
- η' reconstruction: $|M_{\pi\pi\eta} - m_{\eta'}| < 10 \text{ MeV}/c^2$

Background veto:

- π^0 veto: $|M_{\gamma\gamma} - m_{\pi^0}| > 20 \text{ MeV}/c^2$



Clean K^0_s and η' Signal

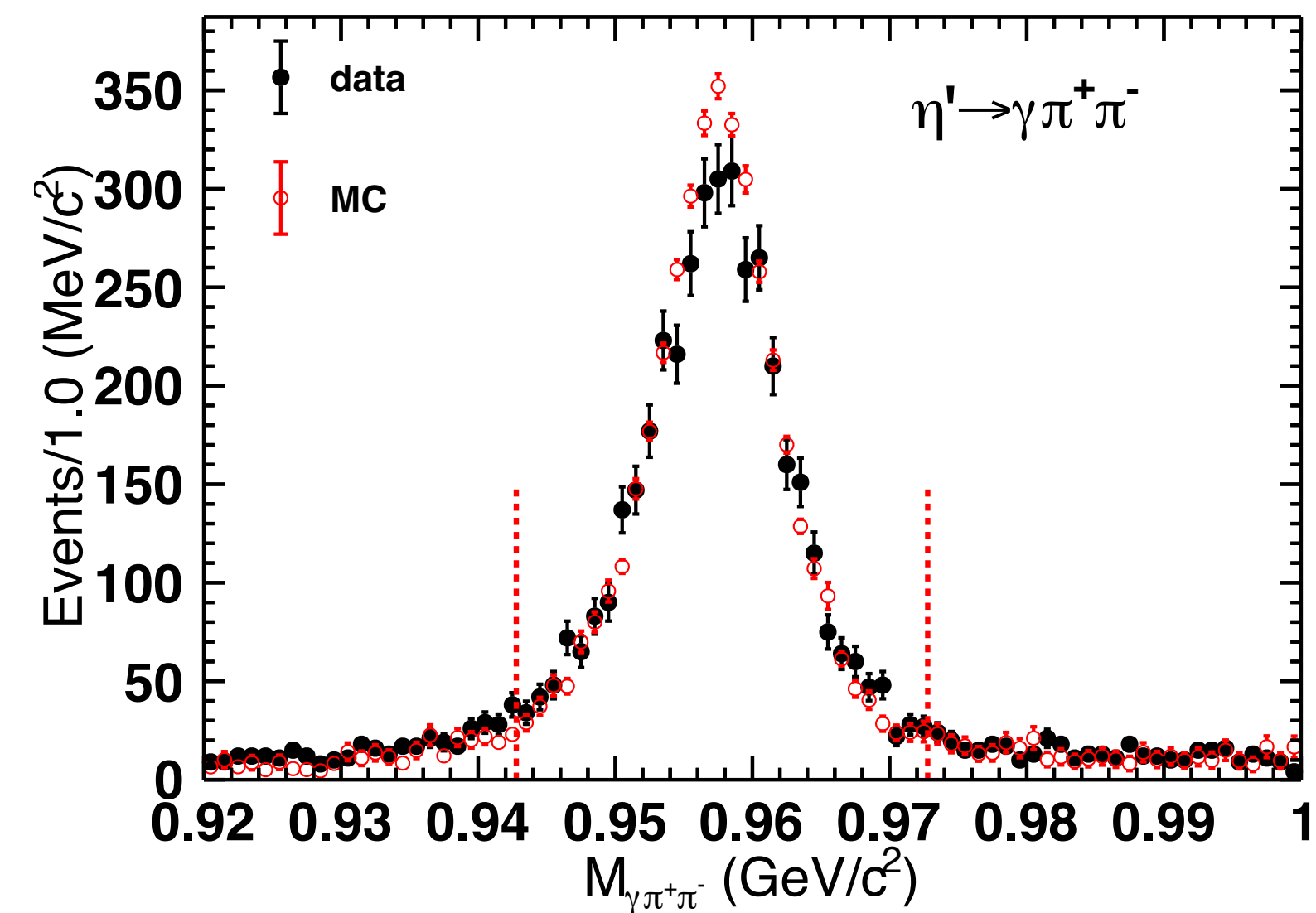
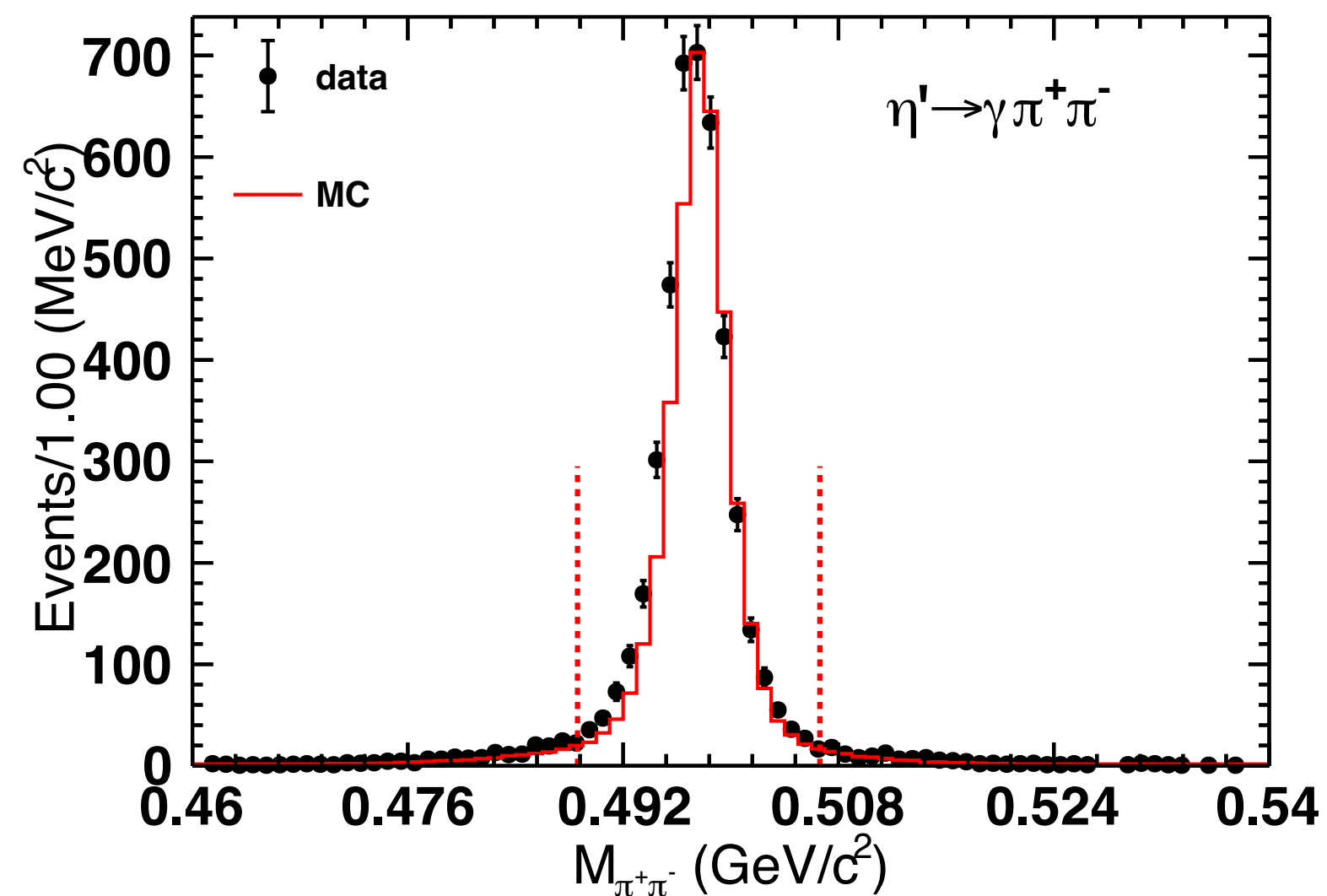
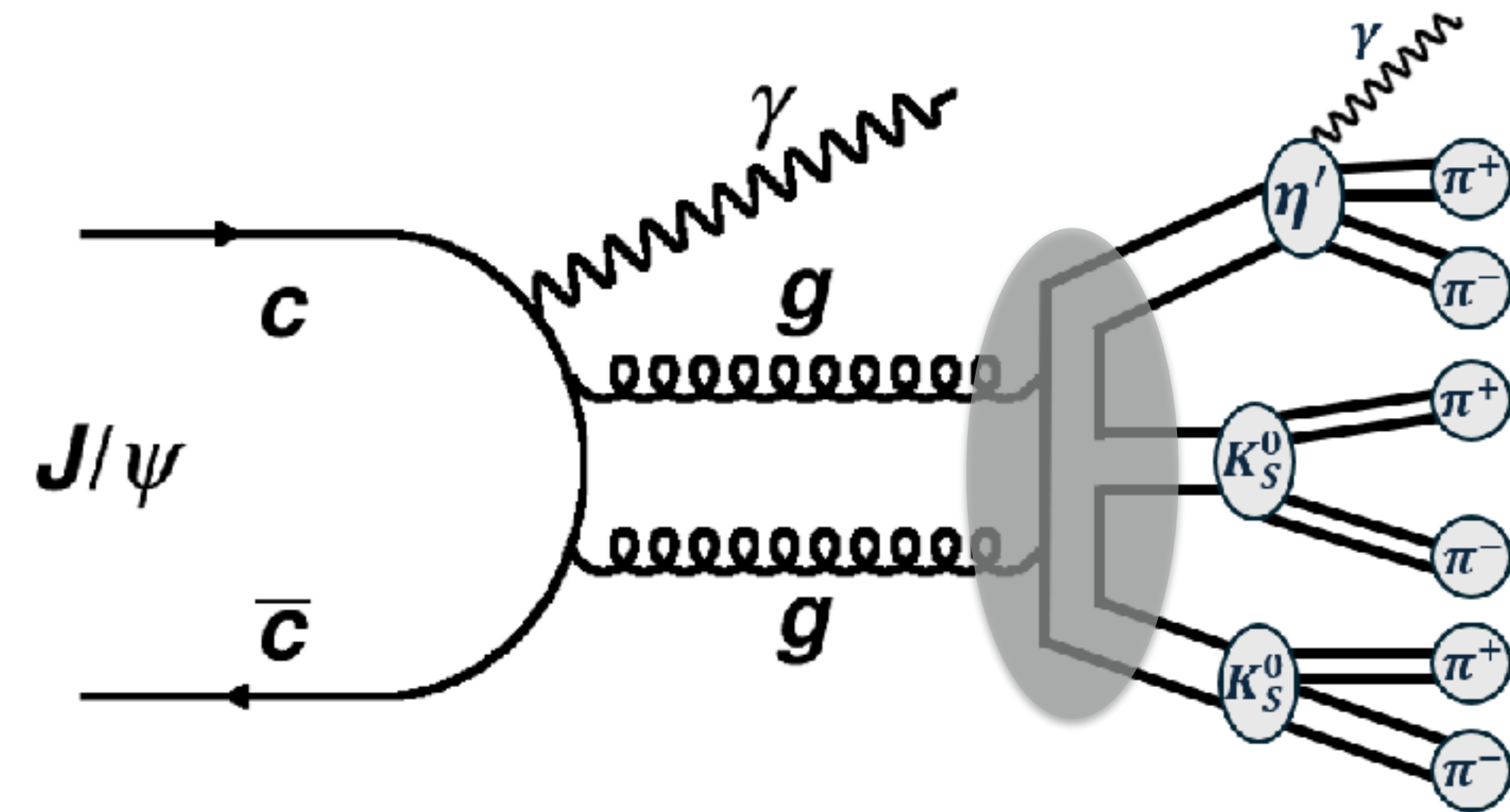
Selection for $J/\psi \rightarrow \gamma K_s^0 K_s^0 \eta'$, $\eta' \rightarrow \gamma \pi^+ \pi^-$

Signal selection:

- At least 3 charged pairs + 2 photons
- Constraint kinematic fit with energy-momentum conservation
- K_s^0 reconstruction: $|M_{\pi\pi} - m_{K_s}| < 9 \text{ MeV}/c^2$
- η' reconstruction: $|M_{\pi\pi\eta} - m_{\eta'}| < 15 \text{ MeV}/c^2$

Background veto:

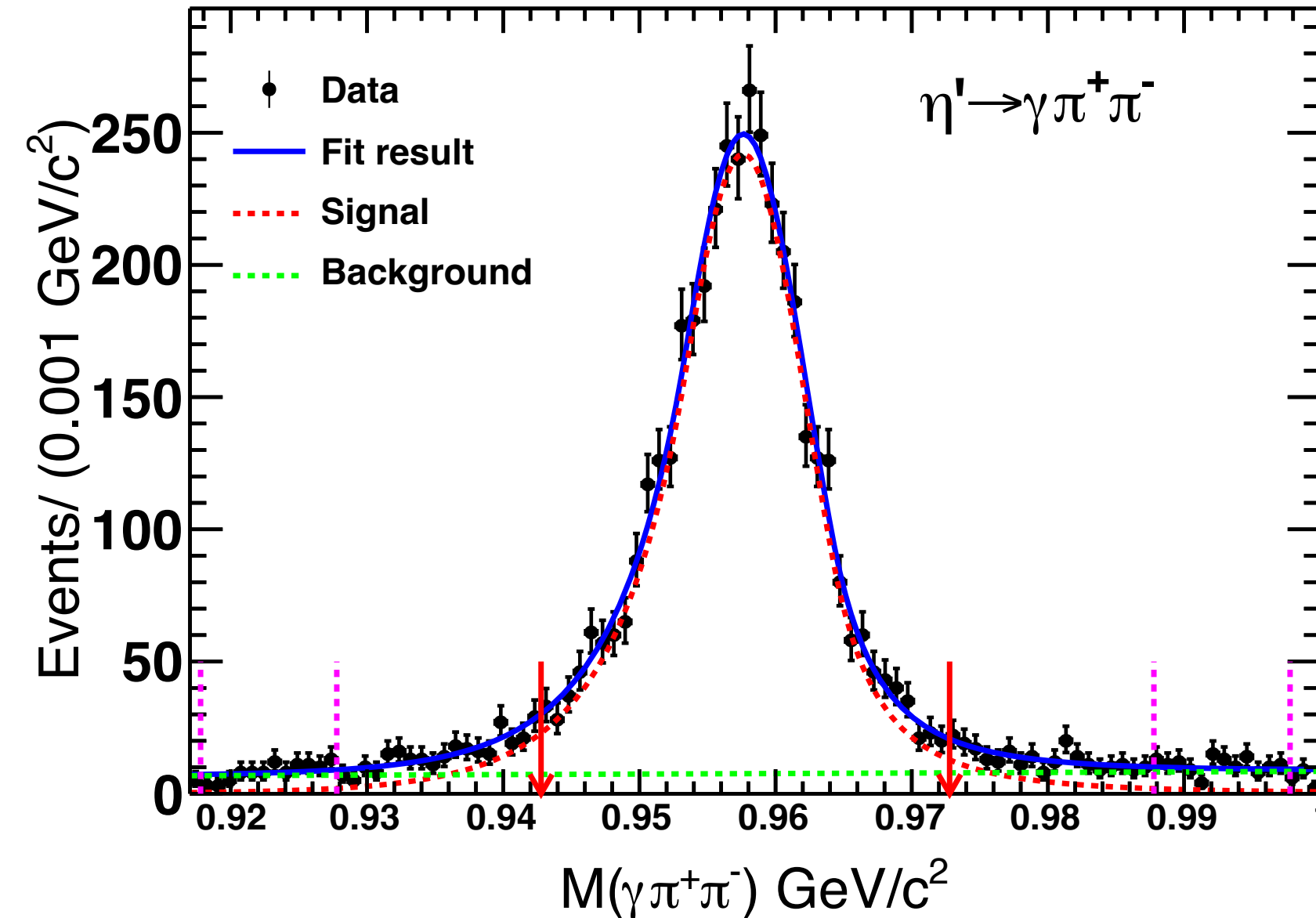
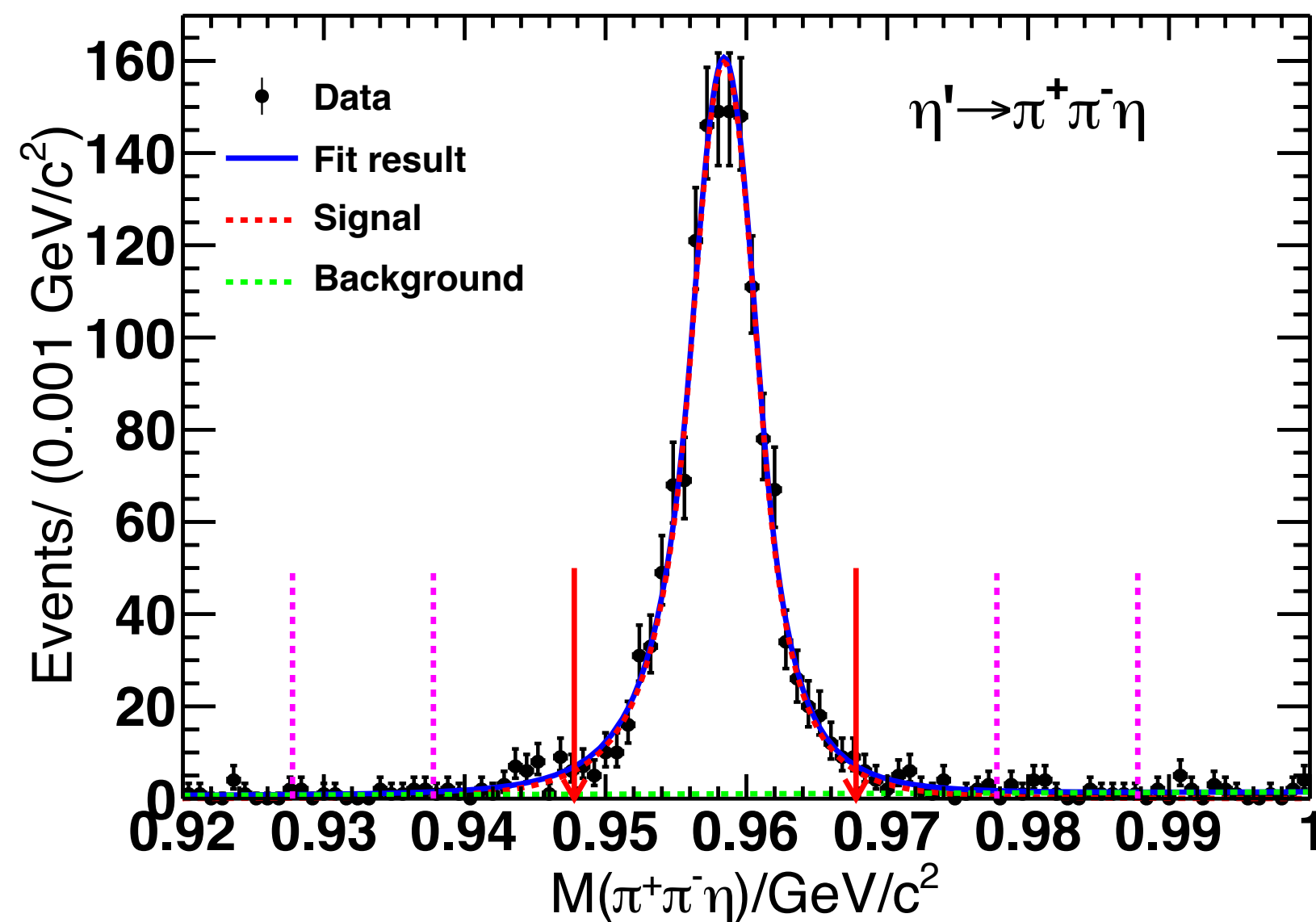
- π^0/η veto: $|M_{\gamma\gamma} - m_{\pi^0}| > 20 \text{ MeV}/c^2$, $|M_{\gamma\gamma} - m_{\eta}| > 30 \text{ MeV}/c^2$



Clean K_s^0 and η' Signal

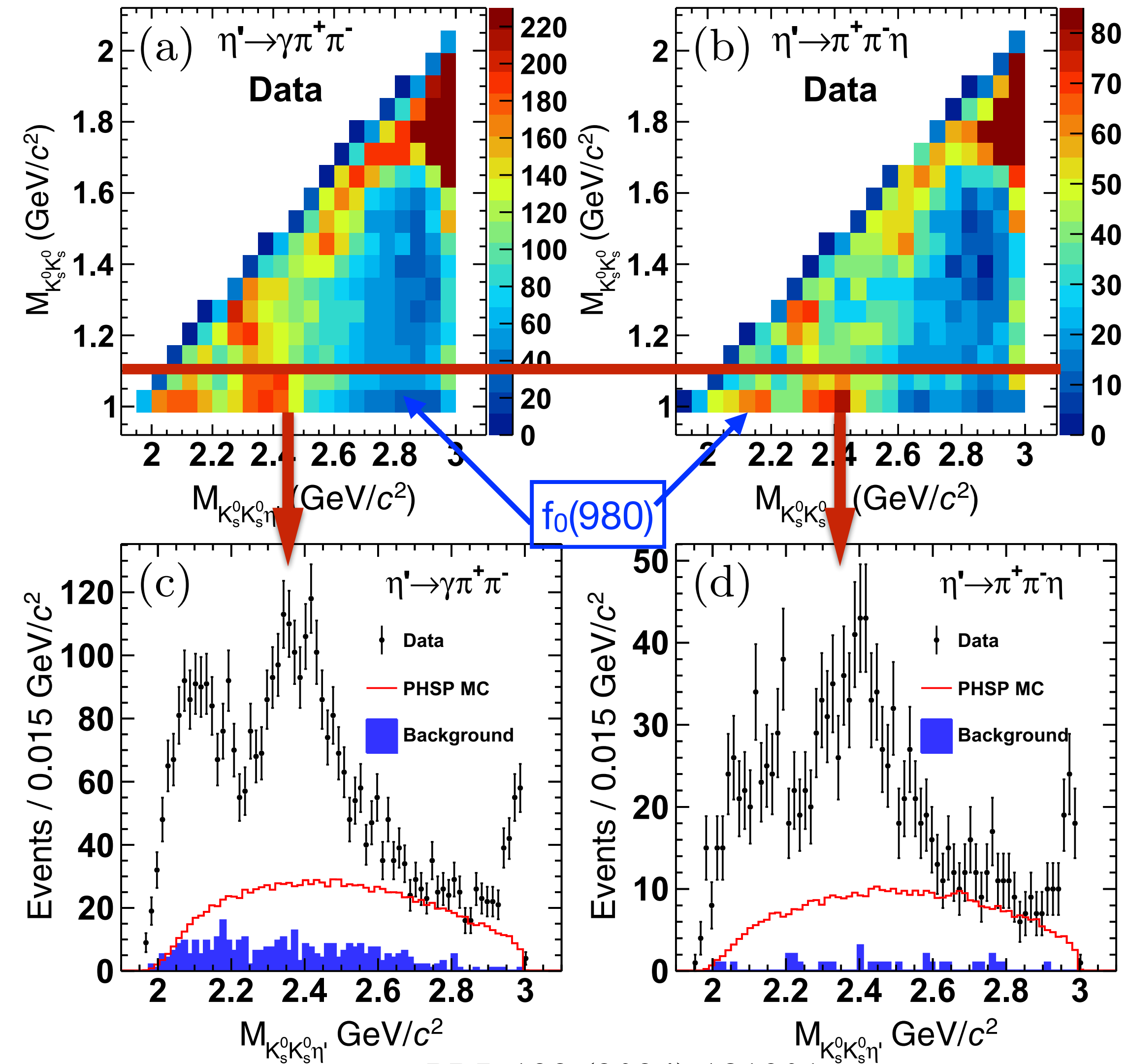
Background estimation

- ◆ **Negligible mis-combination for K^0_s reconstruction ($<0.1\%$)**
- ◆ **No background from $J/\psi \rightarrow \pi^0 K^0_s K^0_s \eta'$** : further validation directly from data
- ◆ **Little background from non- η' processes**: estimated directly from η' mass sideband region:
 - ◆ No peaking background
 - ◆ **Non- η' background fraction: 1.8% for $\eta' \rightarrow \pi^+ \pi^- \eta$ 6.8% for $\eta' \rightarrow \gamma \pi^+ \pi^-$**



The process with almost no background is good for the PWA

Mass spectrum after final selection



PRL 132 (2024) 181901

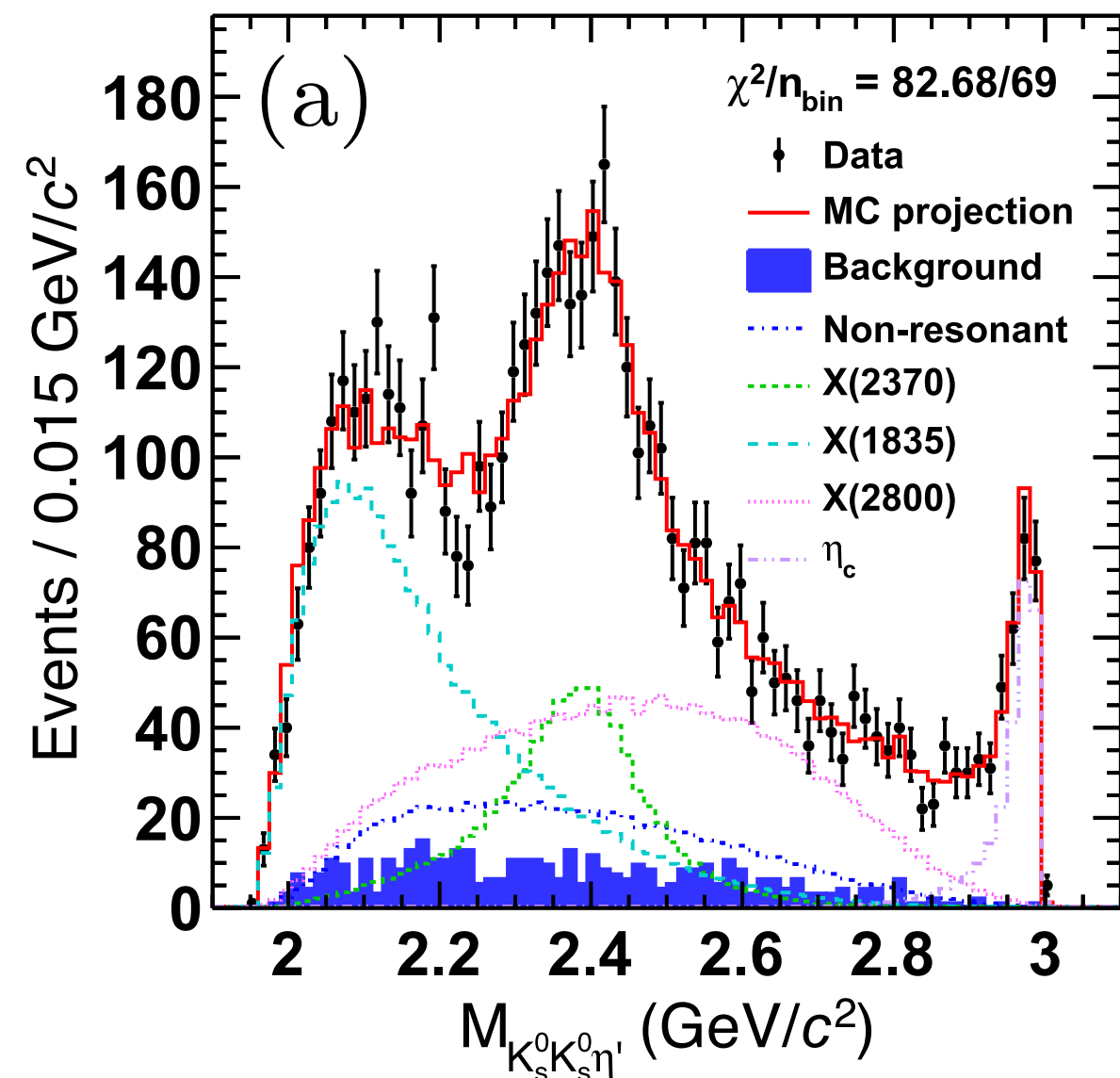
- ◆ **Similar structures in $\eta' \rightarrow \pi^+ \pi^- \eta$ / $\gamma \pi^+ \pi^-$ modes:**
 - ◆ Evident $f_0(980)$ in $K_s^0 K_s^0$ mass threshold
 - ◆ **A clear connection between the $f_0(980)$ and X(2370)**
- ◆ **$f_0(980)$ selection with $M(K_s^0 K_s^0) < 1.1 \text{ GeV}/c^2$**
 - ◆ Clear signal of the X(2370) and η_c
 - ◆ Reduce PWA complexities from additional intermediate processes

Partial wave analysis

- ◆ It is necessary to perform partial wave analysis:
 - ◆ To determine quantum numbers and interferences
- ◆ Amplitude construction with **covariant tensor formalism** [EPJA 16 (2003) 537]
- ◆ **Parametrization with quasi-sequential two-body decays (only spin $J(X) < 3$ states):**
 - ◆ $J/\psi \rightarrow \gamma X, X \rightarrow Y \eta'$: Y represent $K^0_s K^0_s$
 - ◆ $J/\psi \rightarrow \gamma X, X \rightarrow Z K^0_s$: Z represent $K^0_s \eta'$
- ◆ **An un-binned maximum likelihood fit** on the combination of $\eta' \rightarrow \pi^+ \pi^- \eta$ / $\gamma \pi^+ \pi^-$ modes:
 - ◆ Non- η' background subtraction with the NLL values of events from η' sideband region

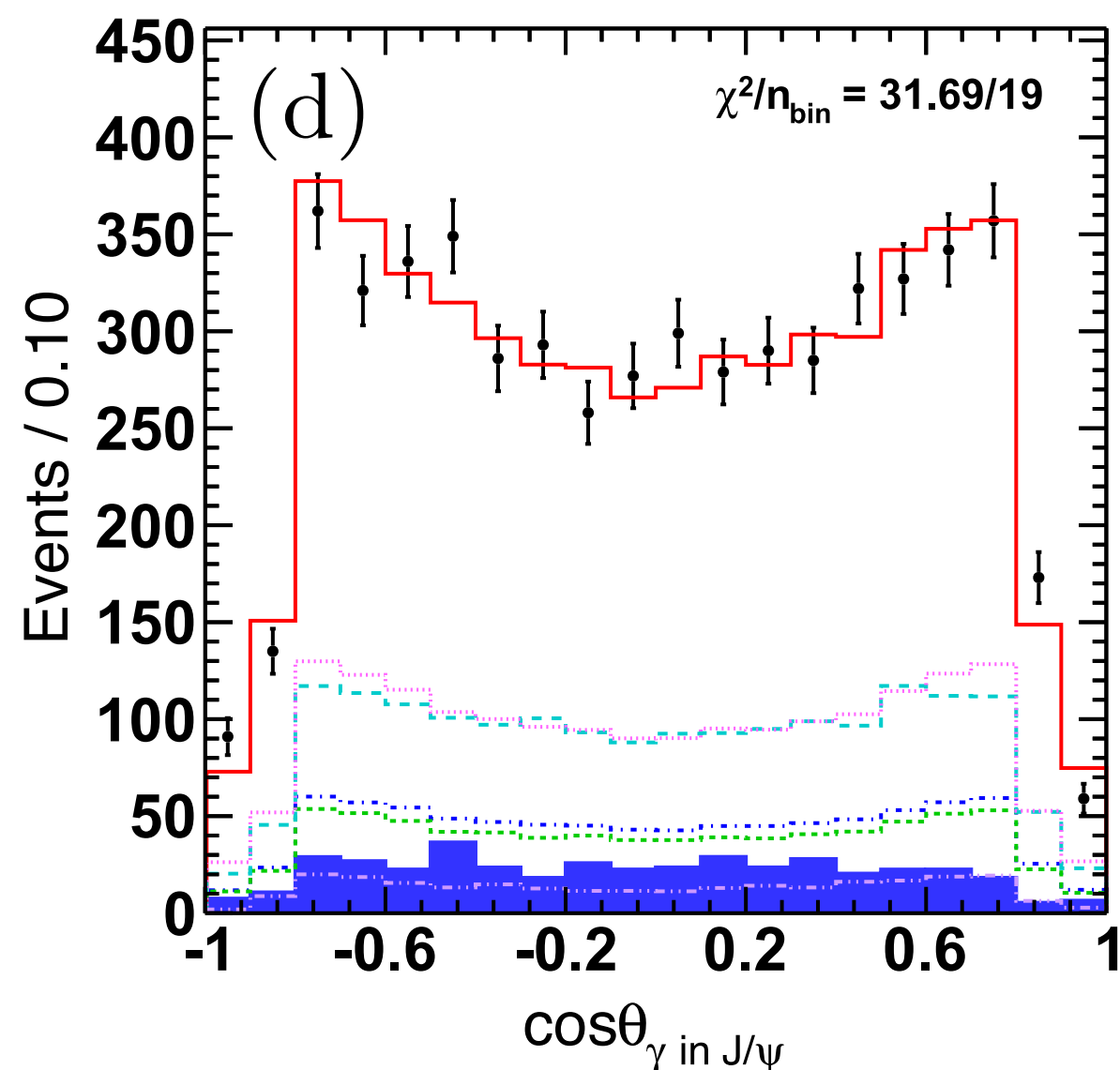
$$S = (-\ln \mathcal{L})_{data} - (-\ln \mathcal{L})_{bg} = - \left(\sum_{i=1}^{N_{data}} \ln \frac{\omega(\xi_i)}{\sigma} \right) + f_{norm} \cdot \left(\sum_{i=1}^{N_{sideband}} \ln \frac{\omega(\xi_i)}{\sigma} \right)$$

PWA Fit



◆ Best fit can well describe the data including resonances ($>5\sigma$):
X(1835), X(2370), X(2800), η_c

- ◆ **Spin-parity of the X(2370) is determined to be 0^{-+}** with significance larger than 9.8σ w.r.t. other J^{PC} assumptions
- ◆ X(2800): a broad structure for the effective contributions from possible high mass resonances and the tail of η_c lineshape



state	J^{PC}	Decay mode	Mass (MeV/c^2)	Width (MeV/c^2)	Significance
X(2370)	0^{-+}	$f_0(980)\eta'$	2395^{+11}_{-11}	188^{+18}_{-17}	14.9σ
X(1835)	0^{-+}	$f_0(980)\eta'$	1844	192	22.0σ
X(2800)	0^{-+}	$f_0(980)\eta'$	2799^{+52}_{-48}	660^{+180}_{-116}	16.4σ
η_c	0^{-+}	$f_0(980)\eta'$	2983.9	32.0	$> 20.0\sigma$
PHSP	0^{-+}	$\eta'(K_S^0 K_S^0)_{S-wave}$	---	---	9.0σ
		$\eta'(K_S^0 K_S^0)_{D-wave}$	---	---	16.3σ

PWA Validations

◆ Additional decay modes: significance $<3\sigma$ and impact is ignored

- ◆ **J^{PC} and decay modes for each components:** $f_0(1500)\eta'$, $f_2(1270)\eta'$, $K^*(1410)K_s^0$, $K_0^*(1430)K_s^0$, $K_0^*(1430)K_s^0$, $K_2^*(1430)K_s^0$, $K_0^*(1680)K_s^0$, $(K_s^0K_s^0)_s\eta'$, $(K_s^0K_s^0)_D\eta'$, $(K_s^0\eta')_PK_s^0$, $(K_s^0\eta')_DK_s^0$

◆ Additional resonance checks: significance $<5\sigma$

- ◆ No evidence of the **X(2120)** in the K_sK_s mass threshold region for $J/\psi \rightarrow \gamma K_s K_s \eta'$ only
- ◆ The significance of **X(2600) \rightarrow $f_0(980)\eta'$** is 4.2σ
- ◆ Impact from the X(2120) and X(2600) is taken into account as systematic uncertainty

◆ The **X(2800)** with a mass of 2799 MeV and width of 660 MeV:

- ◆ Used to described effective contributions from high mass region
- ◆ **Strongly reply on the description of η_c lineshape:** different variations are included into the systematic uncertainty
- ◆ Statistical uncertainties of the X(2800) mass and width are included in the systematic uncertainties on the X(2370) measurements

Final results

X(2370) measurements:

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$J^{PC} = 0^{-+}$ with significance $>9.8\sigma$

$M = 2395 \pm 11^{+26}_{-94}$ MeV

$\Gamma = 188^{+18}_{-17}{}^{+124}_{-33}$ MeV

**$B(J/\psi \rightarrow \gamma X(2370))B(X(2370) \rightarrow f_0(980)\eta')$
 $B(f_0(980) \rightarrow K^0_s K^0_s)$
 $= 1.31 \pm 0.22^{+2.85}_{-0.84} \times 10^{-5}$**

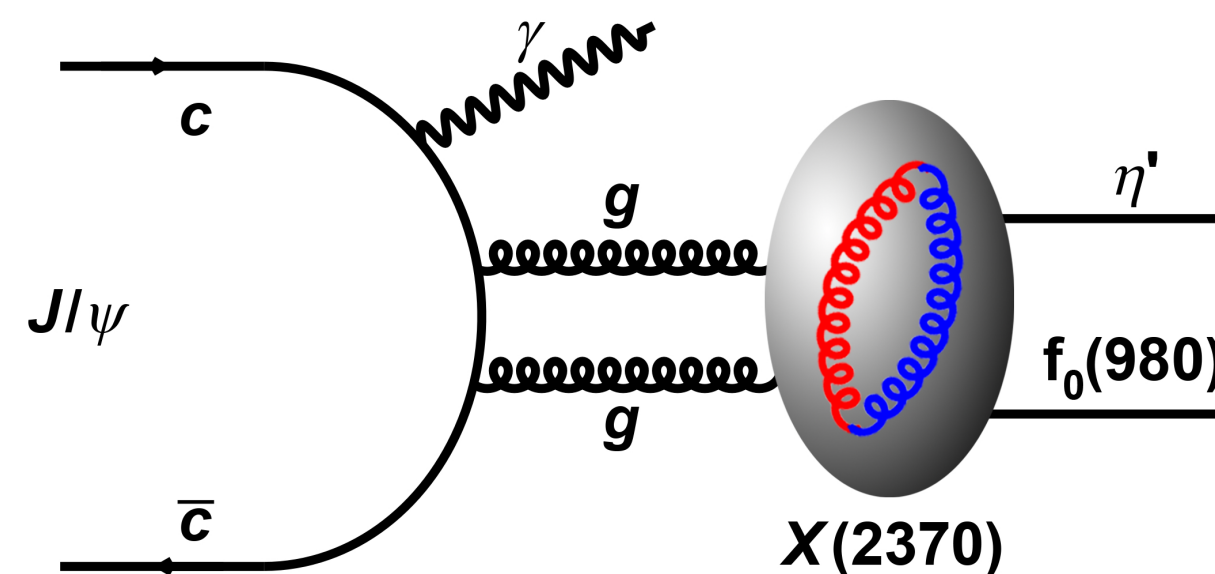
LQCD prediction on lightest pseudoscalar glueball:

$J^{PC} = 0^{-+}$

$M = 2395 \pm 14$ MeV

$B(J/\psi \rightarrow \gamma G_{0^{-+}}) = 2.31 \pm 0.80 \times 10^{-4}$

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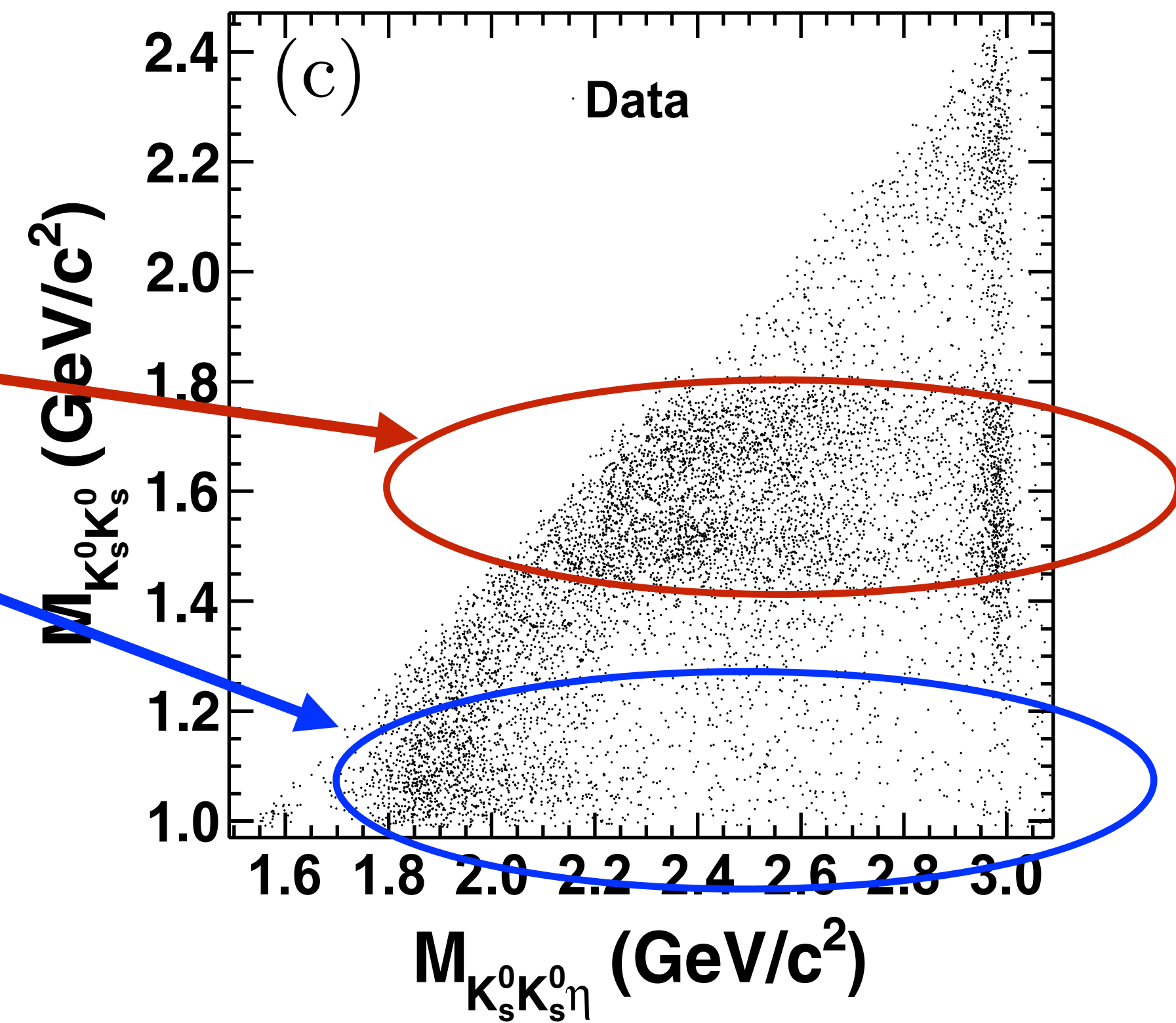


- ◆ The measurements are in a agreement with the predictions on **lightest pseudoscalar glueball**
- ◆ **The spin-parity of the X(2370) is determined to be 0^{-+} for the first time**
- ◆ **Mass is in a good agreement with LQCD predictions**
- ◆ The estimation on $B(J/\psi \rightarrow \gamma X(2370))$ and prediction on $B(J/\psi \rightarrow \gamma G_{0^{-+}})$ are consistent within errors (assuming $\sim 5\%$ decay rate, $B(J/\psi \rightarrow \gamma X(2370)) = 10.7^{+22.8}_{-7} \times 10^{-4}$)

Study in $J/\psi \rightarrow \gamma K_s^0 K_s^0 \eta$

Observation and Spin-Parity Determination of the $X(1835)$ in $J/\psi \rightarrow \gamma K_s^0 K_s^0 \eta$

- ◆ Qualitatively, we can clearly observe:
 - ◆ In the upper M_{KK} mass band of 1.5-1.7 GeV range, clear signals of both $X(2370)$ and η_c
 - ◆ In the lower M_{KK} mass band of $f_0(980)$, no $X(2370)$, nor η_c .
- ◆ Such high similarity between $X(2370)$ and η_c decay modes also strongly supports the glueball interpretation of the $X(2370)$.



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Summary

- ◆ **Glueballs are important predictions from LQCD:**
 - ◆ **Unique particles formed** by gluons (force carriers) due to non-Abelian Gauge self-interactions of gluons
- ◆ **The X(2370) is the first particle that matches the theoretical expectations for a glueball**
 - ◆ **Spin-parity quantum numbers** are determined to be $J^{PC} = 0^{-+}$
 - ◆ Measurements and predictions on **mass and production rate** are consistent within errors
 - ◆ **production and decay properties:** the X(2370) is observed in J/ψ radiative decay and flavor symmetric decay modes (favorite decay modes of 0^{-+} glueball)
 - **Glueball-like particle, X(2370) is discovered by BESIII**

Many thanks to the efficient work:

The BESIII detector maintenance and offline software teams, computing center

The BEPCII accelerator operation team which provide stable detector operation

Prospects

- ◆ **More decay modes of the $X(2370)$ will be studied at BESIII**
 - ◆ Including $KK\pi$, $\pi\pi\eta$
 - ◆ To check their similarities with η_c , and to understand the decay pattern of this glueball-like particle
- ◆ **Improve the measurements on the mass, width, branching ratio and production rates of the $X(2370)$**
 - ◆ Need to have better ways to understand and control the interferences in PWA.
- ◆ **Looking forward to more reliable LQCD studies on the glueball properties**

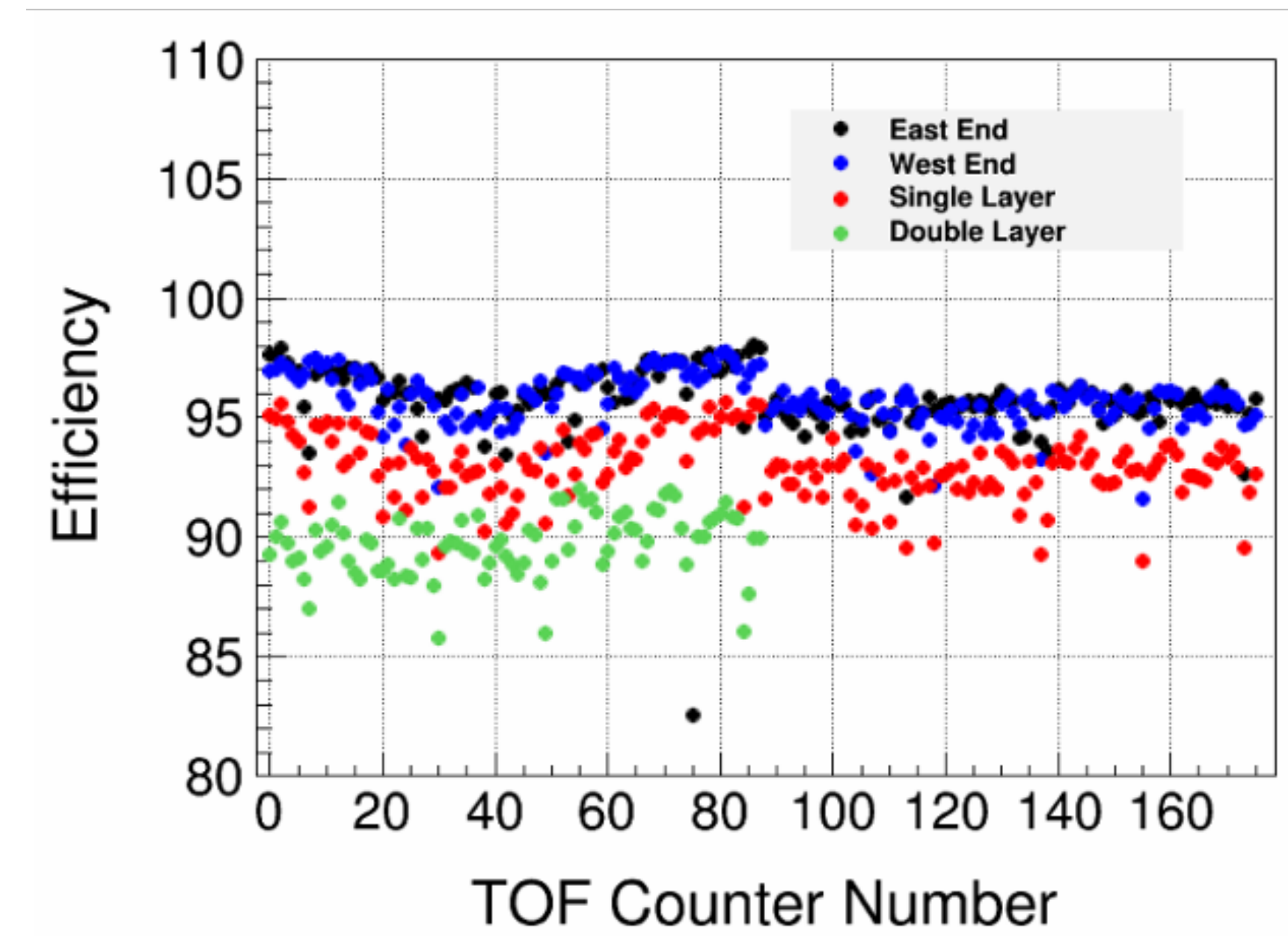
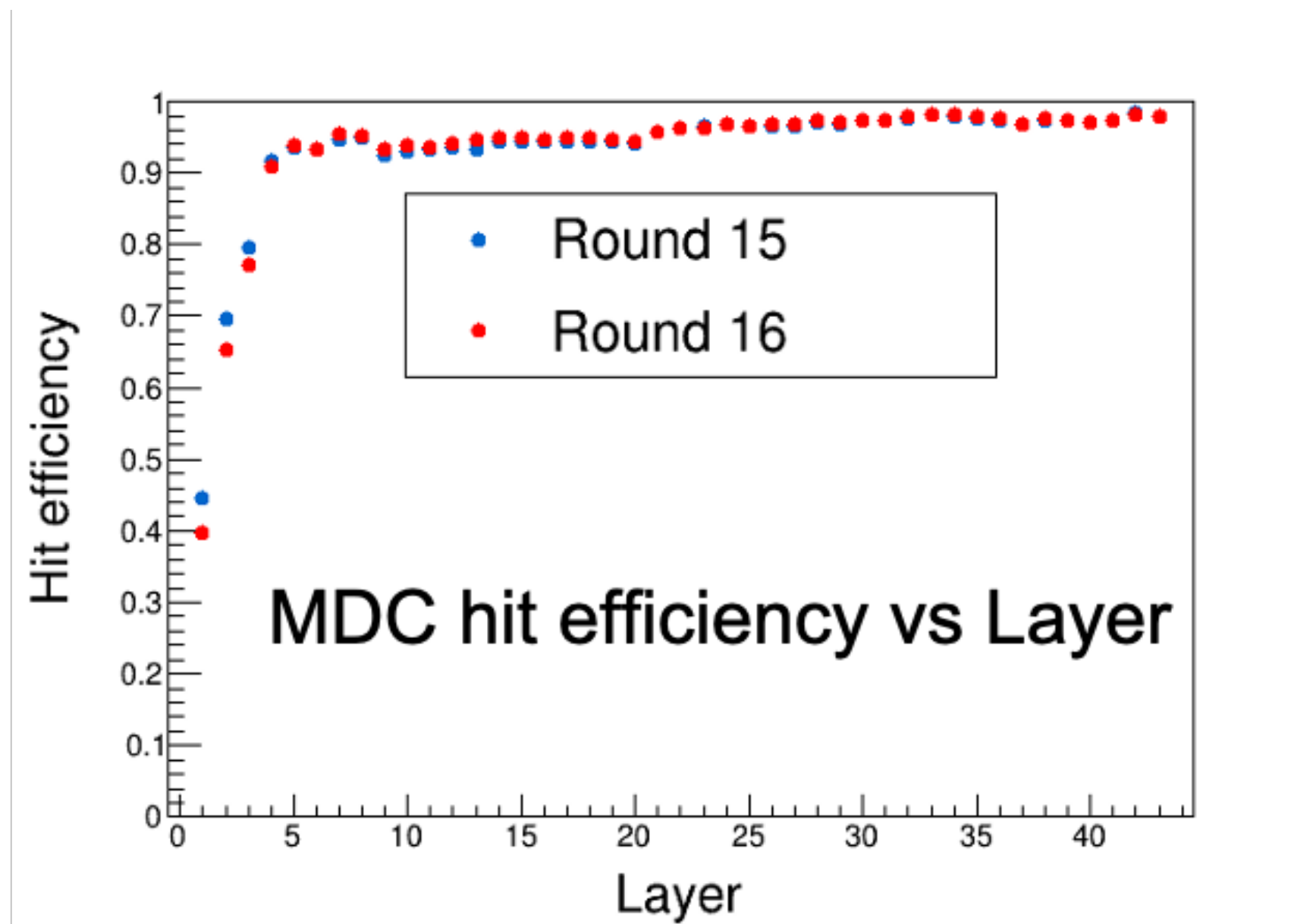
Back-up

Exps.	MDC Wire reso.	MDC dE/dx Reso.	EMC Energy reso.
CLEO	110 μm	5%	2.2-2.4 %
Babar	125 μm	7%	2.67 %
Belle	130 μm	5.6%	2.2 %
BESIII	115 μm	<5% (Bhabha)	2.3%

Exps.	TOF time reso.
CDFII	100 ps
Belle	90 ps
BESIII	68 ps (Barrel) 98 ps (Endcap) 60 ps(Endcap, MRPC)

MUC: Eff. \sim 96%

Noise: $< 0.04 \text{ Hz/cm}^2$ (Barrel), $< 0.1 \text{ Hz/cm}^2$ (Endcap)



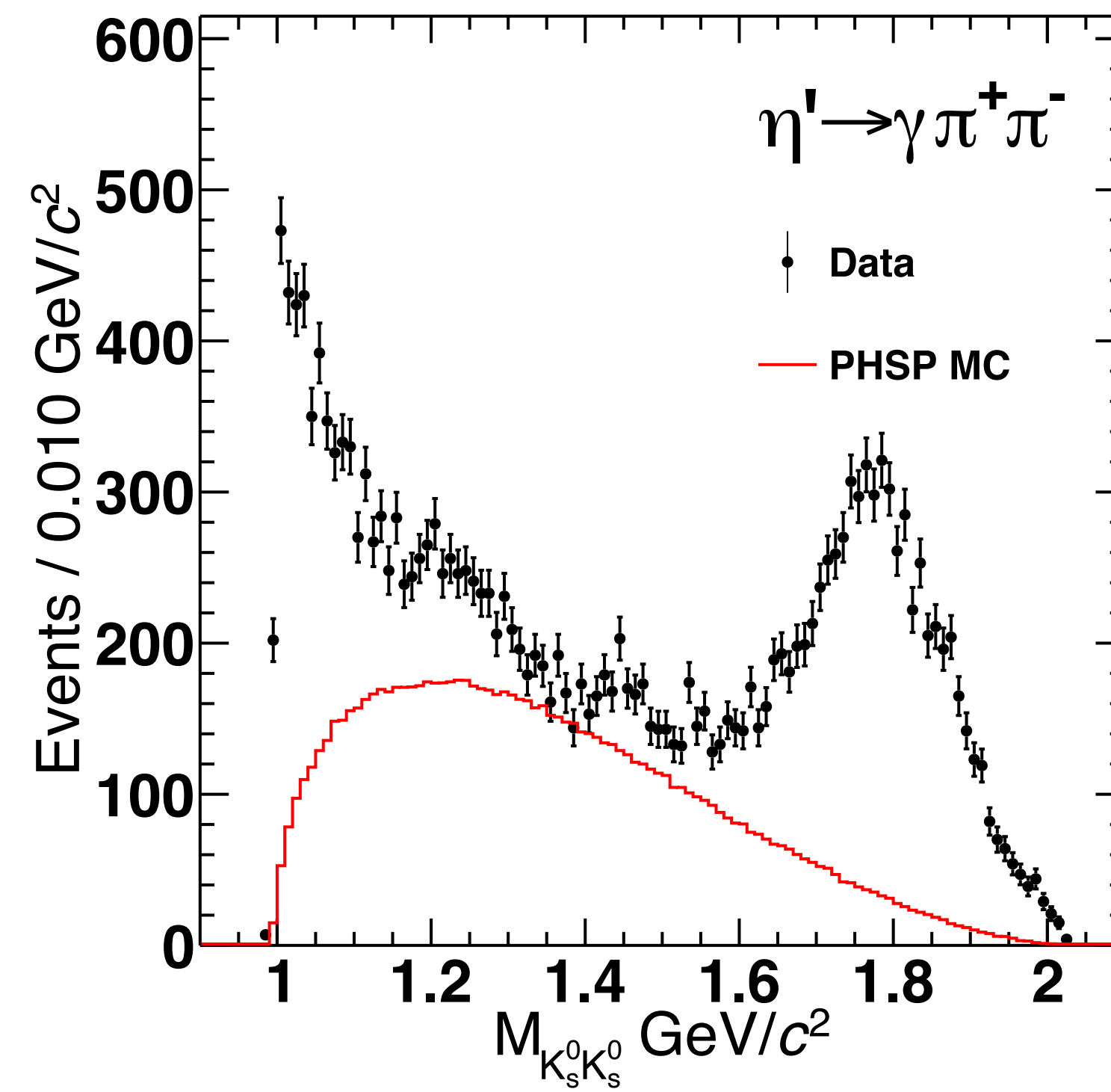
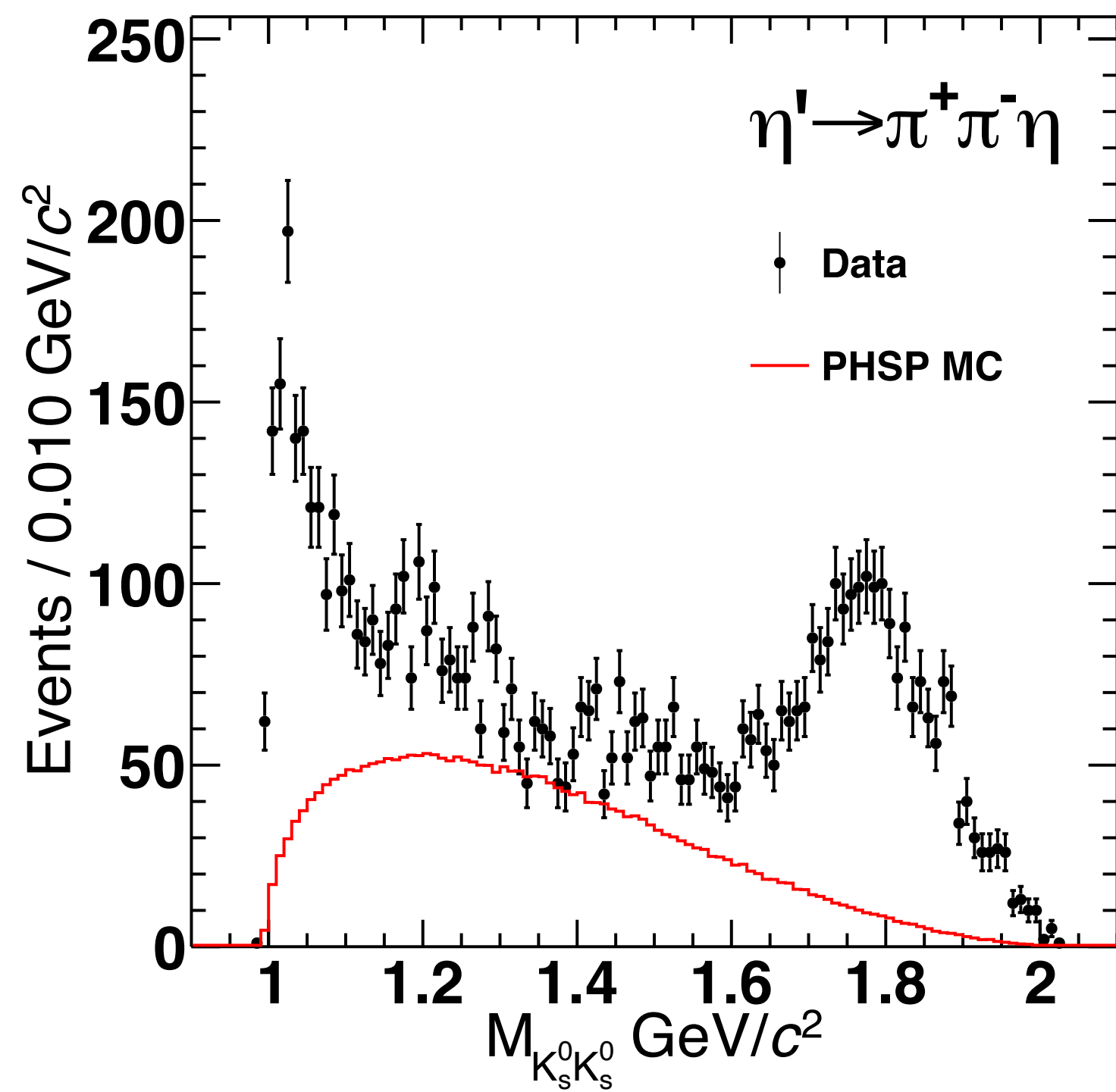
LQCD prediction

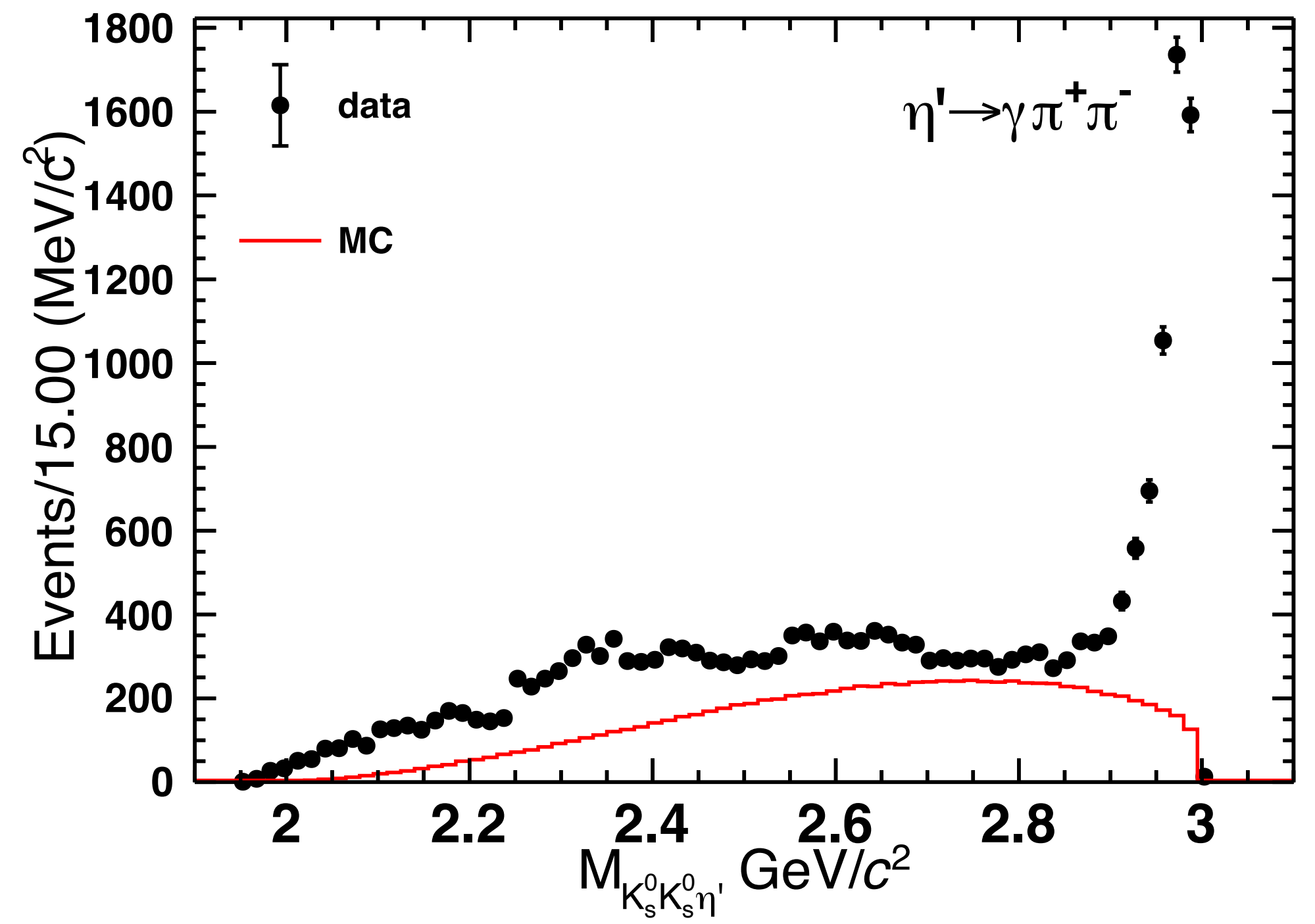
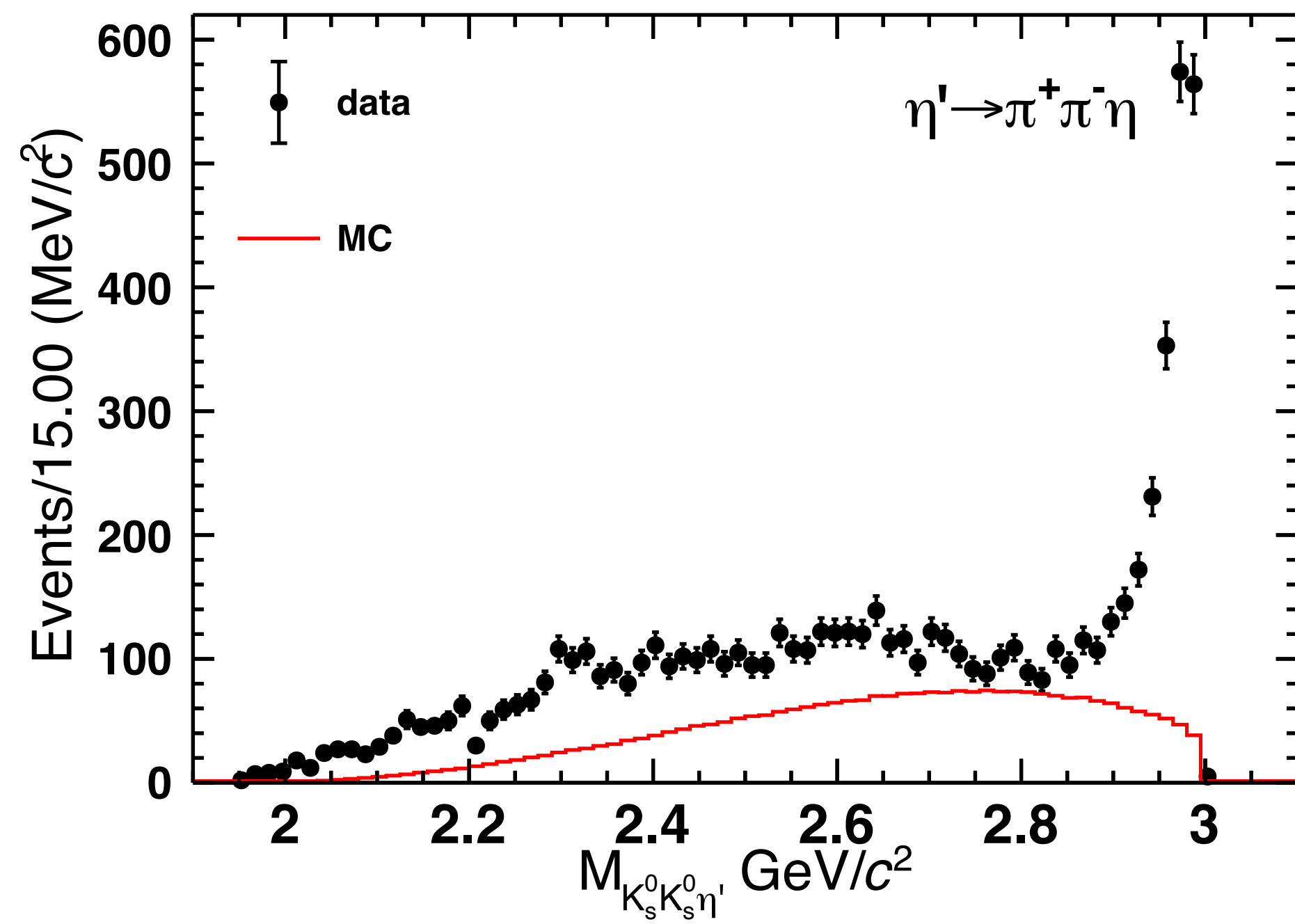
By applying the variational method to a large operator set, we obtain an optimal operator which couples predominantly to the ground state pseudoscalar glueball G . In this work, m_G is determined to be 2.395(14) GeV, and the on-shell form factor of $J/\psi \rightarrow \gamma G$ is derived as $\hat{V}(0) = 0.0246(43)$, in the continuum limit, from which we obtain the following partial decay width and the production rate

$$\begin{aligned}\Gamma(J/\psi \rightarrow \gamma G_{0+}) &= 0.0215(74) \text{ keV} \\ \text{Br}(J/\psi \rightarrow \gamma G_{0+}) &= 2.31(80) \times 10^{-4}.\end{aligned}\quad (32)$$

We introduce an effective coupling g_X to describe the interaction between the pseudoscalar X and the gluonic intermediate states in the processes $J/\psi \rightarrow \gamma X$, as defined in Eq. (30). It is interesting to see that all the g_X 's are comparable for the pseudoscalar glueball and the non-flavored $q\bar{q}$ pseudoscalars (η states). We tentatively attribute the large production rates of the η states to the QCD $U_A(1)$ anomaly which is totally a nonperturbative effect.

Even though this study is performed in the quenched approximation and the uncertainty in the presence of dynamical quarks is not controlled, we hope our result can provide useful theoretical information for experiments to unravel the properties of the possible pseudoscalar glueball.





Sources	ΔM (MeV/ c^2)	$\Delta\Gamma$ (MeV)	$\Delta\mathcal{B}/\mathcal{B}(\%)$
Event selection	± 4.8
Background estimation	+2	+4 -4	+3.7 -5.1
$f_0(980)$ parametrization	-6	+7	± 5.3
X(1835) parametrization	+15 -12	+24 -11	+20.2 -8.3
η_c parametrization	-13	-8	-14.5
Breit-Wigner formula	-1	+6	-8.3
Broad 0^{-+} structure	-88	+111 -21	+211.8 -56.5
Additional resonances	+22 -25	+48 -21	+41.9 -20.8
Total	+26 -94	+124 -33	+217.0 -63.7

component	$X(2370) \rightarrow f_0(980)\eta'$	$X(1835) \rightarrow f_0(980)\eta'$	$X(2800) \rightarrow f_0(980)\eta'$	$\eta_c \rightarrow f_0(980)\eta'$	$PHSP \rightarrow \eta'(K_S^0 K_S^0)_{S-wave}$	$PHSP \rightarrow \eta'(K_S^0 K_S^0)_{D-wave}$
$X(2370) \rightarrow f_0(980)\eta'$	1.00	1.12	1.15	0.08	-1.23	0.02
$X(1835) \rightarrow f_0(980)\eta'$		1.00	0.32	-0.07	-1.41	0.01
$X(2800) \rightarrow f_0(980)\eta'$			1.00	0.33	-1.18	0.02
$\eta_c \rightarrow f_0(980)\eta'$				1.00	0.11	0.00
$PHSP \rightarrow \eta'(K_S^0 K_S^0)_{S-wave}$					1.00	-0.02
$PHSP \rightarrow \eta'(K_S^0 K_S^0)_{D-wave}$						1.00

