

Angle- and Polarization-dependent Synchrotron and SSC blazar model

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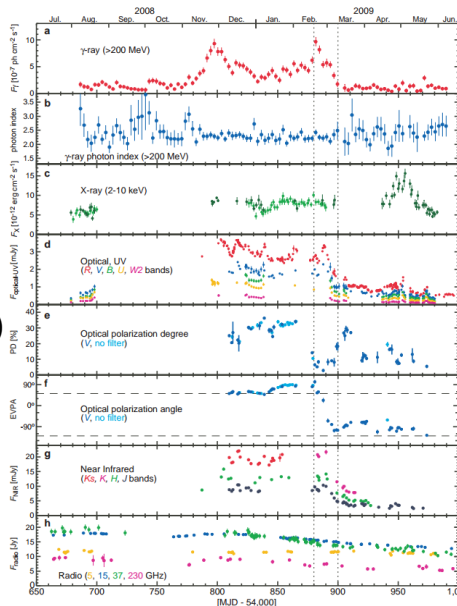
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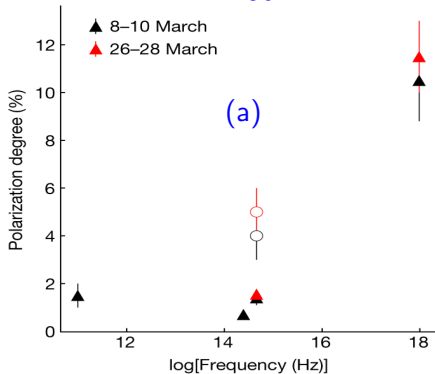
Multiwavelength Polarization of Blazars

- Blazars exhibit polarization in radio-optical bands and X-rays (in some cases).
- Multi-band polarimetry effectively probes the jet magnetic field and emitting regions.
- Blazar polarization (PD and PA) can change with time (see the right plot, from [Abdo et al., 2010 - 3C 279](#)).
- Optical-radio polarimetry effectively probes less active jet regions.



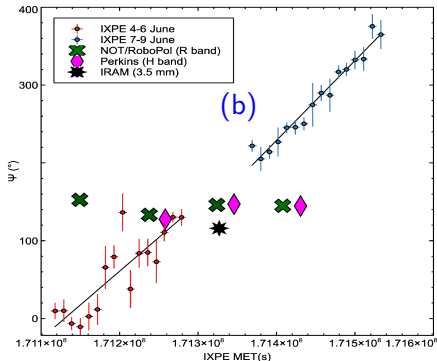
IXPE Observation Results

Mrk 501



Liodakis et al. (2022)

Mrk 421



Di Gesu et al. (2023)

Related Conducted Work

- Lyutikov et al. (2005) reviewed synchrotron polarization properties with purely helical B -fields.
- Jamil & Böttcher (2012) studied the angle-dependent S+SSC blazar emission with partially ordered B -fields.
- Zhang & Böttcher (2013) analyzed the expected SSC polarization signatures of blazars in a perfectly ordered B -field.
- Joshi et al. (2020) explored the effects of the magnetic field orientation on the SEDs of a typical blazar.

So, the fully angle- and polarization-dependent S+SSC blazar emission has never been explored.

- To develop a comprehensive, fully angle- and polarization-dependent synchrotron and SSC blazar model.
- Explore the SED and linear PD variations when the magnetic field is oriented perpendicular, obliquely, and toroidally relative to the jet axis.
- Conducted in the emission frame moving along the jet \rightarrow no Doppler boosting effects into the observer's frame.
- LoS, jet axis \parallel z-axis.

Model (Synchrotron Radiation)

Synchrotron polarization and emissivity (Rybicki & Lightman, 1979):

$$\Pi_{\nu}^{\text{syn}}(\varphi) = \frac{\langle G(x) \rangle}{\langle F(x) \rangle} = \frac{\int n_e(\gamma) G(x) d\gamma}{\int n_e(\gamma) F(x) d\gamma}, \quad (1)$$

$$j_{\nu}^{\text{syn}}(\varphi) \propto \sin(\varphi) \int n_e(\gamma) F(x) d\gamma, \quad (2)$$

where

$$F_{\text{exact}}(x) = x \int_x^{\infty} K_{5/3}(\xi) d\xi, \quad F_{\text{approx}}(x) = 1.42 \sqrt{\frac{\pi}{2}} x^{3/10} e^{-x} \quad (3)$$

$$G(x) = x K_{2/3}(x), \quad (4)$$

$$x = \frac{\nu}{\nu_{\varphi}(\gamma)}, \quad (5)$$

$$\nu_{\varphi}(\gamma) \approx 4.2 \times 10^6 B \gamma^2 \sin \varphi \text{ [Hz]}, \quad (6)$$

$$n_e(\gamma) = n_0 \gamma^{-p} \text{ for } (\gamma_1 \leq \gamma \leq \gamma_2). \quad (7)$$

Model (Synchrotron Self-Compton Process)

SSC emissivity (Böttcher et al., 2012):

$$j^{\text{SSC}}(\epsilon_s, \Omega_s) = \frac{hc\epsilon_s}{4\pi} \int_1^\infty n_e(\gamma) d\gamma \int_{-1}^{+1} d\eta_{\text{ph}} \int_0^{2\pi} d\phi_{\text{ph}} \int_0^\infty n_{\text{ph}}(\epsilon, \Omega_{\text{ph}}) (1 - \beta\mu) \frac{d\sigma_C}{d\epsilon_s} d\epsilon. \quad (8)$$

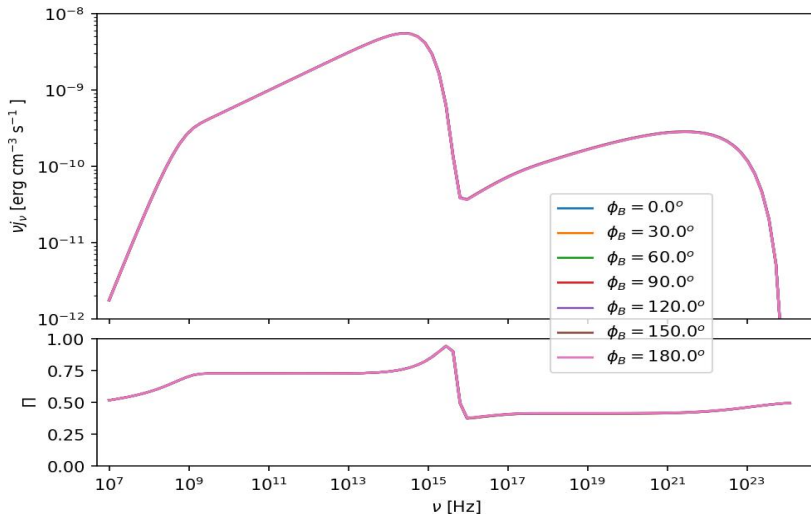
SSC Polarization (Bonometto & Saggion, 1973):

$$\Pi_\nu^{\text{SSC}} = \frac{P_\perp^{\text{SSC}} - P_\parallel^{\text{SSC}}}{P_\perp^{\text{SSC}} + P_\parallel^{\text{SSC}}} \quad (9)$$

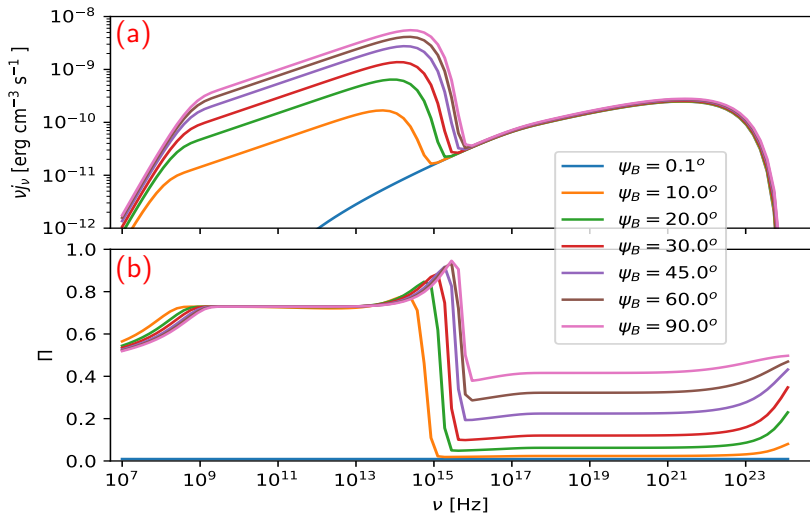
Total Polarization:

$$\Pi_\nu^{\text{Total}} = \frac{\Pi_\nu^{\text{Syn}} \cdot j_\nu^{\text{Syn}} + \Pi_\nu^{\text{SSC}} \cdot j_\nu^{\text{SSC}}}{j_\nu^{\text{Syn}} + j_\nu^{\text{SSC}}} \quad (10)$$

Results (SED and Total Polarization)

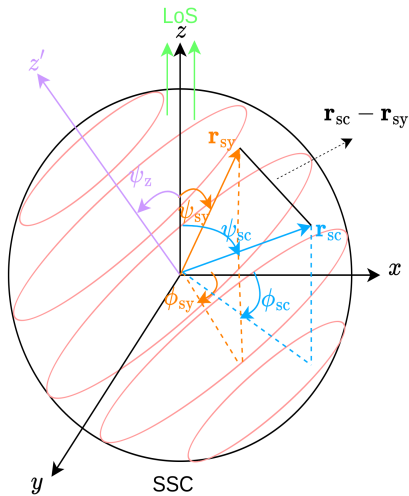
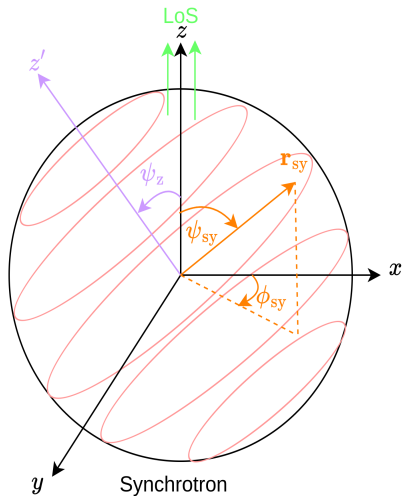


Results (cont...)



Currently \rightarrow Toroidal Field

- Exploring the toroidal (and eventually, helical) magnetic field structures.



Summary and Outlook

- We developed an angle- and polarization-dependent S+SSC blazar model.
- New aspect was to treat the anisotropy of the radiation produced and the polarization in the treatment of the synchrotron and SSC mechanisms.
- Synchrotron SED component and SSC PD depend strongly on the B -field polar angle compared to the synchrotron PD and SSC SED component.
- First step in the development of a comprehensive, fully angle- and polarization-dependent blazar radiation transfer code.
- Future work: We will also consider more complex B -field geometries (such as helical) and the effects of SSA and gamma-ray absorption.