## Higgs Properties in CMS

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on behalf of CMS collaboration Johns Hopkins University August 31, 2011





- Intro: status of CMS
- Properties of SM Higgs
- Angular distributions
- Discriminating signal/background
- Separating hypotheses
- Measuring parameters
- Conclusions

- LHC continues to perform better than expected
  - Already  $\sim 2.5 \, fb^{-1}$  on tape!
- $\bullet\,$  Can expect to  $\sim\,$  double int. lumi. in 2011
- 2012: 10  $fb^{-1}$  per experiment?
- Allowed parameter space for SM Higgs shrinking quickly



- Only free parameter of SM Higgs is its mass
  - Assuming a given mass, all properties of SM Higgs can, in principle, be calculated
- Given large excess in data what will we know:
  - mass, width, cross section
- What about Higgs specific properties?
  - $J^P = 0^+$
  - full angular correlations of final state particles
  - unique manifestations of  $J^P$
  - has been demonstrated to be good handle for determining spin and parity of resonances



• For scalar resonance decaying into 2 vector bosons, most general amplitude:

$$A(X \rightarrow V_1 V_2) = v^{-1} \epsilon_1^{*\mu} \epsilon_2^{*\nu} (a_1 g_{\mu\nu} M_X^2) + a_2 q_{1\mu} q_{2\nu} + a_3 \epsilon_{\mu\nu\alpha\beta} q_1^{\alpha} q_2^{\beta})$$

- SM Higgs $\rightarrow$ ZZ,WW:  $a_1 \neq 0, a_2 \sim O(10^{-2}), a_3 \sim O(10^{-11})$
- SM Higgs  $\rightarrow \gamma \gamma$ :  $a_1 = -a_2/2 \neq 0$
- BSM psuedo-scalar Higgs  $a_3 \neq 0$
- One can write a general formula for all fermionic final states
- Can be applied to spin 1 & 2 resonances as well
- Including amplitude for production of X and decay of V's and integrating:  $\frac{d\Gamma(\vec{\Omega};a_1,a_2,a_3)}{\Gamma d\vec{\Omega}}$



## Kinematics of Decay

- Kinematics of final state fermions can be separated into three sets of (mostly) uncorrelated variables
  - $P_T^X, Y^X$
  - $m_{f1,f2}, m_{f3,f4}, m_{f1,f2,f3,f4}$
  - $\cos \theta^*, \Phi_1, \cos \theta_1, \cos \theta_2, \Phi$
- cos θ\*, Φ<sub>1</sub> are related to production of the Z's (production angles)
- $\cos \theta_1$ ,  $\cos \theta_2$ ,  $\Phi$  are related to Z decays (helicity angles)
- The production/helicity angular distributions determined by helicity amplitudes  $A_{00} = -\frac{m_X^4}{v}(a_1 x + a_2 \frac{M_Z M_*}{M_X^2}(x^2 - 1)),$   $A_{\pm\pm} = \frac{m_X^2}{v}(a_1 \pm i a_3 \frac{M_Z M_*}{M_H^2} \sqrt{x^2 - 1})$   $x = \frac{M_H^2 - M_Z^2 - M_*^2}{2M_Z M_*}$



# Helicity Angular Distribution $(J_X = 0)$

$$\begin{aligned} d\Gamma(\theta^*, \Phi_1, \theta_1, \theta_2, \Phi) \propto & 4(1 - f_{++} - f_{--})\sin^2\theta_1 \sin^2\theta_2 \\ &+ (f_{++} + f_{--})((1 + \cos^2\theta_1)(1 + \cos^2\theta_2) + 4R_1R_2\cos\theta_1\cos\theta_2) \\ &- 2(f_{++} - f_{--})(R_1\cos\theta_1(1 + \cos^2\theta_2) + R_2(1 + \cos^2\theta_1)\cos\theta_2) \\ &+ 4\sqrt{f_{++}(1 - f_{++} - f_{--})}(R_1 - \cos\theta_1)\sin\theta_1(R_2 - \cos\theta_2)\sin\theta_2\cos(\Phi + \phi_{++}) \\ &+ 4\sqrt{f_{--}(1 - f_{++} - f_{--})}(R_1 + \cos\theta_1)\sin\theta_1(R_2 + \cos\theta_2)\sin\theta_2\cos(\Phi - \phi_{--}) \\ &+ 2\sqrt{f_{++}f_{--}}\sin^2\theta_1\sin^2\theta_2\cos(2\Phi + \phi_{++} - \phi_{--}) \end{aligned}$$

- Flat distribution of production angles, cos θ\*, Φ<sub>1</sub> (background & J > 0 have non-trivial distributions)
- $f_{ij}$  and  $\phi_{ij}$  determined by helicity amplitudes  $\rightarrow$  couplings

$$f_{ij} = |A_{ij}|^2 / \sum_{k,l} |A_{kl}|^2,$$

$$\phi_{ij} = arg(A_{ij}/A_{00})$$

•  $R_{1,2}$  determined by fermion type

# Angular Distributions $(X \rightarrow ZZ \rightarrow 4I)$



‡Note there are no detector effects here.

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#### Angular Distributions cont.



# Helicity Angles as a Background Discriminant

- Simplest application of angular distributions
- Within the  $H \rightarrow ZZ$  decay channel this was applied to two final states, 4l and 2l2j
- Has been shown to increase sensitivity in 4l final state by  $\sim 20\%$  ([1], arxiv.org:1001.3396)
- Model angular distributions as seen in detector
- Signal: (Ideal)×(uncorrelated acceptance)

 $P_{\textit{sig}} = P_{\textit{IDEAL}}(\theta^*, \theta_1, \theta_2, \Phi, \Phi_1; \vec{\xi}) A_{\theta^*}(\theta^*) A_{\theta_1}(\theta_1) A_{\theta_2}(\theta_2) A_{\Phi}(\Phi) A_{\Phi_1}(\Phi_1)$ 

- $\xi = (f_{ij}, \phi_{ij})$  fixed to SM Higgs values
- Background: product of 1D, uncorrelated functions

$$P_{bkg} = D_{\theta^*}(\theta^*) D_{\theta_1}(\theta_1) D_{\theta_2}(\theta_2) D_{\Phi}(\Phi) D_{\Phi_1}(\Phi_1)$$

• Define discriminant as:

$$D=rac{P_{sig}}{P_{sig}+P_{bkg}}$$

•  $D\epsilon[0,1]$ ; cuts applied to D (e.g. D > .7 are signal-like)

# Angular Distributions $(ZZ \rightarrow 4I)$



# Angular Distributions $(ZZ \rightarrow 2/2j)$



# Helicity Likelihood Discriminant $(ZZ \rightarrow 2I2j)$



‡ For more details, see twiki/PAS here - https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsHIG

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#### **Background Parameterization**

- 4I final state: SM ZZ production is the major background
- Use MC, fit helicity amplitudes to

 $q\bar{q} \rightarrow ZZ \rightarrow 4I,$  $gg \rightarrow ZZ \rightarrow 4I$ 

 Using helicity amplitudes as basis for fits can recover correlations in background

- Example of helicity amplitude fit to SM ZZ events near 250 GeV
- Can use to measure fraction of gg vs qq initiated events in data

parameter	$qar{q}  ightarrow ZZ$	gg  ightarrow ZZ
f <sub>00</sub>	0.025	0.398
$f_{++}$	0.206	0.430
f	0.005	0.012
$f_{+0}$	0.007	0.047
<i>f</i> <sub>0-</sub>	0.147	0.007
<i>f</i> <sub>+-</sub>	0.228	0.026

- $\bullet~$  Cutting on  $D{\rightarrow}~$  lose information
- Instead, maximum likelihood (ML) fit would be better

$$L = \exp(-n_{sig} - n_{bkg}) \prod_{i}^{N} (n_{sig} \times P_{sig}(\vec{\Theta_i}, m_{ZZ}; \vec{\xi}) + n_{bkg} \times P_{bkg}(\vec{\Theta_i}, m_{ZZ}))$$

• Fixing  $\vec{\xi}$  and floating  $n_{bkg}$ ,  $n_{sig}$  one can calculate upper limit, significance, etc. of  $n_{sig}$  for a given resonance hypothesis

# Separating Signal Hypotheses $(ZZ \rightarrow 4I)$

- Using 5D likelihood for a given model (SM Higgs, pseudo-scalar, RS graviton, SM ZZ...)
  - evaluate  $-2ln(L_1/L_2)$  for data and two choice models (e.g. SM Higgs, pseudo-scalar)
  - using MC psuedo-experiments, separation significance can be calculated
- Example: resonance with

$$m = 250, n_{sig} = 30, n_{bkg} = 24 (\sim 5 fb^{-1} @ \sqrt{s} = 14 TeV)$$
  
model 1:  $J^P = 0^+$ , model 2:  $J^P = 0^-$  (A)  
model 1:  $J^P = 0^+$ , model 2:  $J^P = 2^+_m$  (B)



#### Separating Signal Hypotheses $(ZZ \rightarrow 4I)$ contd

 Separation significance, S, has been calculated for a number of hypothetical models (S - # of widths between peaks)

• all using a resonace of 250 GeV,  

$$n_{sig} = 30, n_{bkg} = 24 (\sim 5 \text{ fb}^{-1} \text{ @ } \sqrt{s} = 14 \text{ TeV})$$

	0-	$1^{+}$	1-	$2_{m}^{+}$	$2_{L}^{+}$	2-
0+	4.1	2.3	2.6	2.8	2.6	3.3
$0^{-}$		3.1	3.0	2.4	4.8	2.9
$1^+$			2.2	2.6	3.6	2.9
$1^{-}$				1.8	3.8	3.4
$2_{m}^{+}$					3.8	3.2
$2_{L}^{+}$						4.3

• Most values are  $\gtrsim$  3 and almost all are > 2

#### Measuring Helicity Amplitudes

- floating  $\vec{\xi}$  one could use the ML to measure helicity amplitudes of a given spin hypothesis
- Example study:
  - for  $ZZ \rightarrow 4I$  final state
  - $n_{sig} = 150, \ n_{bkg} = 120 \ (\sim 25 \ fb^{-1} \ \mathbb{Q} \ \sqrt{s} = 14 \ TeV)$
  - Generate MC for  $J^p = 0^+, 0^-$  resonance at 250 GeV (A), (B)

(A)	generated	$m_X = 250 \text{ GeV}$ without detector	with detector	- 150	)		 1 <b>4</b>			
nsig	150	$150 \pm 13$	$153\pm15$	-	-					-
$(f_{++} + f_{})$	0.208	$0.21\pm0.07$	$0.23\pm0.08$	Its	Į.		1	•		-
$(f_{++} - f_{})$	0.000	$0.01\pm0.13$	$0.01\pm0.14$	້ອີ 100	)-		/` 🔺	X		-
$(\phi_{++} + \phi_{})$	$2\pi$	$6.30 \pm 1.46$	$6.39 \pm 1.54$	<u>=</u> .	Ē.	/	′ <b>T</b>	1		-
$(\phi_{++}-\phi_{})$	0	$0.00 \pm 1.06$	$0.01 \pm 1.09$	ber	F	<b>/</b>		•		-
(B)	generated	$m_X = 250 \text{ GeV}$ without detector	with detector	- <u>ũ</u> 50	) - -			**		
n <sub>sig</sub>	150	$150 \pm 13$	$151 \pm 15$	-	- <b>-</b>					
$(f_{++} + f_{})$	1.000	$1.00\pm0.05$	$1.00\pm0.06$	, i	0	0.1	0.2	0.3	0.4	0.5
$(f_{++} - f_{})$	0.000	$0.00 \pm 0.35$	$0.00\pm0.40$				f	⊥f		
$(\phi_{++} + \phi_{})$	N/A	free	free				'+	+ + 1		
$(\phi_{++} - \phi_{})$	$\pi$	$3.15\pm0.31$	$3.14\pm0.41$							
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## $ZZ \rightarrow 4I$ below threshold

- All of the above is valid below threshold also
- Angular distributions require an additional parameter:  $m_{Z^*}$



• Separation significance between  $0^+/0^-$ : S=3.3  $n_{sig} = 20, n_{bkg} = 30, m = 140 \text{ GeV} (\sim 10 \text{ fb}^{-1} @ \sqrt{s} = 7 \text{ TeV})$ 

### Conclusions

- Angular distributions have been very beneficial to Higgs searches thus far
  - have been exploited for signal/background discrimination
  - has been shown to improve sensitivity in the  $ZZ \to 4I$  channel by  $\sim 20\%$
  - can ultimately help to discover new resonances
- Angular variables are physically motivated
  - have been parameterized in terms of helicity amplitude (coupling constants)
  - can be used to measure properties of new resonances
- Methods described above are already being implemented in analyses
  - Already implemented in  $ZZ \rightarrow 2/2j$  analysis
  - Will be implemented in  $ZZ \rightarrow 4I$  analysis

# **BACKUP SLIDES**

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#### CMS documentation

[1] Y.Gao, A.Gritsan, Z.Guo, K.Melnikov, M.Schulze, N.Tran, "Spin determination of single-produced resonances at hadron colliders"

http://arxiv.org/abs/1001.3396

ATLAS documentation

[2] C.P.Buszello, I.Fleck, P.Marquard and J.J. van der Bij, "Prospective Analysis of Spin- and CP-sensitive Variables in  $H \rightarrow ZZ \rightarrow I^+I^-I^+I^-$  with ATLAS", (Eur Phys J C32,209,2004SN-ATLAS-2003-025) • JHU generator is intended for generating resonances with the following decay topologies:

$$ab \rightarrow X \rightarrow ZZ \rightarrow 4I$$
,

$$ab 
ightarrow X 
ightarrow ZZ 
ightarrow 2/2j$$

- Proper angular correlations are computed
- Resonances can be spin 0,1,2 with arbitrary couplings
- Output is a standard LHE file
- Code and further documentation can be found here: http://www.pha.jhu.edu/spin/