

# SUSY Higgs and interpretation of LHC results

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Recall: SUSY  $\rightarrow$  at least 2 Higgs doublets  $H_u, H_d$  (+ Singlets ?)

All CP-even Higgs bosons  $H_i$  “share” their reduced couplings  $\xi_i$  (normalized w.r.t. the SM-like coupling) to electroweak gauge bosons:

$$\sum_i \xi_i^2 = 1$$

The couplings to  $b$ -quarks,  $t$ -quarks ( $\rightarrow$  gluon-gluon) can be enhanced or reduced, depending on  $\tan\beta$  and the Higgs mixing angles

**At present:** (mild) excesses of events at ATLAS/CMS  
for  $M_H \sim 140 - 145$  GeV

What if these excesses are confirmed?

## MSSM:

In the MSSM (minimal Higgs sector with two Higgs doublets) the mass  $M_h$  of the Standard Model-like Higgs Boson is bounded from above by

$$\left(M_h^{upper}\right)^2 \simeq M_Z^2 \cos^2 2\beta + \frac{m_t^4}{4\pi^2 v^2} \left( \ln \left( \frac{M_{stop}^2}{m_t^2} \right) + \frac{A_t^2}{M_{stop}^2} \left( 1 - \frac{A_t^2}{12M_{stop}^2} \right) \right) + \dots$$

$$\rightarrow M_h \lesssim 135 \text{ GeV} \text{ if } M_{stop} \lesssim 3 \text{ TeV}$$

The second “heavy” CP-even Higgs boson  $H$  can have a mass in the 140-150 GeV range (with  $M_h \sim 115$  GeV,  $\tan \beta \sim 8 - 10$ ), but then its  $BR$  into  $WW/ZZ$  is small,  $< 10^{-2}/10^{-3}$

(Question: could this be enhanced in the case of CP-violation?)

$\rightarrow$  Split Susy, extra vector-like matter, extra  $U(1)'$ , NMSSM ( $\lambda$ -Susy),...

## Split SUSY:

Supersymmetry does not solve the Hierarchy Problem, but

- a light gaugino (bino) provides a dark matter candidate
- sparticles lead to gauge coupling unification
- scalar sparticles like stops can be heavy

$$\rightarrow M_H \gtrsim 140 \text{ GeV if } M_{stop} \gtrsim 10^6 \text{ GeV}$$

(D. Alves, E. Izaguirre, Jay G. Wacker, 1108.3390)

## Extra Vector-like Matter

- Add extra SU(5) multiplets, keeping perturbative unification at the GUT scale: e.g. one  $10 + \overline{10}$  and a superpotential  
$$W = \dots + H_u 10 10 + M_{10} 10 \overline{10}$$
- Assume that the extra scalars and fermions are sufficiently massive in order to comply with negative direct searches and electroweak precision observables ( $\gtrsim 300 - 400$  GeV)
- Rad. corr. to  $M_H$  of the same type as top-stop loops:

$$\rightarrow M_H \gtrsim 140 \text{ GeV if } M_{\text{scalar}} \gtrsim 5M_{\text{fermion}}$$

(M. Asano, T. Moroi, R. Sato, T. Yanagida, 1108.2402,  
M. Endo, K. Hamaguchi, S. Iwamoto, N. Yokozaki, 1108.3071,  
J. Evans, M. Ibe, T. Yanagida, 1108.3437)

## Extra U(1)'

- If the Higgs doublets carry U(1)' charges:  
Extra quartic Higgs couplings from U(1)'-D-terms,  
extra SM-singlets (see the NMSSM below)
- Larger SM-like Higgs mass possible
- Models must comply with lower bounds on  $M_{Z'} \gtrsim 2$  TeV,  
absence of U(1)' anomalies → enlarged field content

Scenarios with  $M_H \gtrsim 140$  GeV are possible

(E. Ma, 1108.4029)

# NMSSM

Solution of the  $\mu$ -problem of the MSSM by an additional gauge singlet superfield  $S$  (the simplest Susy extension of the SM with a scale invariant superpotential):

$$W_{MSSM} = \dots + \mu H_u H_d \quad \rightarrow \quad W_{NMSSM} = \dots + \lambda S H_u H_d + \frac{\kappa}{3} S^3 \quad (+\dots)$$

→ 3 CP-even Higgs bosons, mixtures of  $H_u$ ,  $H_d$  and  $S$

→ Upper bound on the lightest CP-even state, valid also for the general NMSSM with additional Susy mass/tadpole terms:

$$\left(M_H^{upper}\right)^2 \simeq \left(M_{H, MSSM}^{upper}\right)^2 + \lambda^2 M_Z^2 \sin^2 2\beta + \dots$$

(The second tree-level term decreases as  $\tan^{-2} \beta$ )

Assume  $\lambda \lesssim 0.75$  ( $\lambda(M_{GUT}) < \infty$ ),  $M_{stop} \lesssim 1$  TeV,  $m_{top} \sim 173$  GeV, incl. full electroweak 1-loop and  $h_t/\alpha_s$ -2-loop corrections (Degrassi, Slavich):

$$M_H^{upper} \simeq 138 \text{ GeV}$$

**Moreover:**  $M_H = M_H^{upper}$  only if  $H - S$ -mixing terms vanish!

Ways out:

a)  $\lambda > 0.75$  ( $\lambda$ -Susy): either

— additional matter/gauge interactions such that  $\lambda(Q^2)$  does not increase (strongly) with  $Q^2$ , or

— accept  $\lambda(Q'^2) \rightarrow \infty$ , new strong interactions at  $Q' < M_{GUT}$   
( $S =$  composite field, fat Higgs...)

→  $M_H \sim 200$  GeV possible

b) Drop the (unnecessary) assumption that  $H_{SM} \equiv H_1$ :

IF  $H_{SM} = H_2$  ( $H_1$  mostly singlet, with a mass possibly below 114 GeV):

Typically  $M_{H_{SM}} > M_H^{upper}$  due to  $H - S$ -mixing!

**BUT:** The reduced coupling to electroweak gauge bosons  $\xi_2$  of  $H_2$  becomes  $\xi_2 < 1$  due to  $H - S$ -mixing; also: slightly reduced couplings (by 70-90%) to  $b$ -quarks,  $t$ -quarks/gluon-gluon (U.E., 1108.0157):

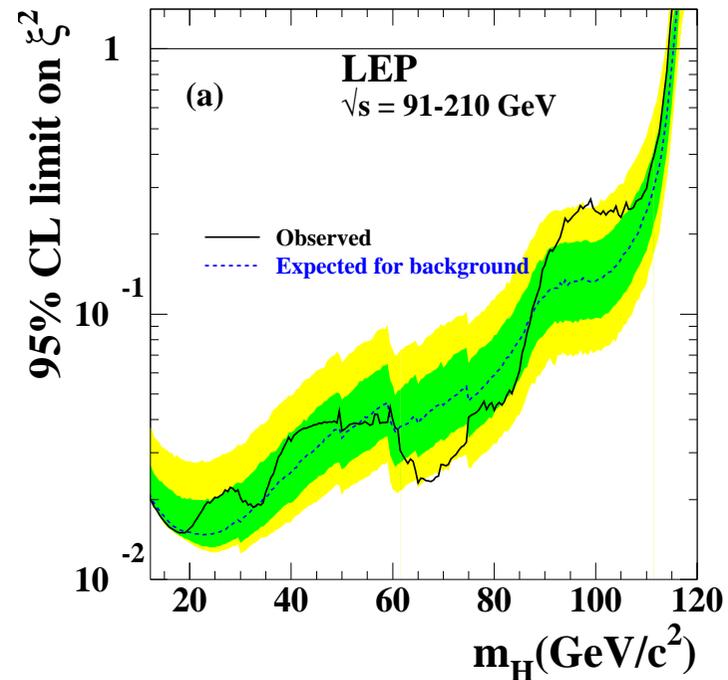
Examples ( $\lambda \sim 0.7$ ,  $\tan \beta \sim 2.7$ ):

$M_{H_2}$	140 GeV	145 GeV	150 GeV
$\xi_2$	0.92	0.86	0.73
$M_{H_1}$	91 GeV	97 GeV	115 GeV
$\xi_1$	0.40	0.51	0.68

Note:  $\xi_1^2 + \xi_2^2 \simeq 1$ ,

$M_{H_1}$  and  $\xi_1$  are consistent with LEP constraints, which are particularly weak for  $M_{H_1} \sim 95$  GeV (Dermisek, Gunion)

→  $M_{H_{SM}} \gtrsim 140$  GeV is possible, but slightly reduced signal rates!



## Note:

- $H_1$  could be visible at the Tevatron (if CDF/D0 look for  $M_H$  as light as  $\sim 100$  GeV)
- The  $H_1$  production cross section  $\times BR(H_1 \rightarrow \gamma\gamma)$  could be enhanced by a factor 6 w.r.t. the SM, if the  $H_d$  component of  $H_1$  happens to be small (U.E., 1012.1201)
  - the reduced width for  $H_1 \rightarrow b\bar{b}$  enhances the  $BR(H_1 \rightarrow \gamma\gamma)$  by a factor up to 10

**Moreover:** possibly *less*  $E_T^{miss}$  in sparticle decay cascades in the NMSSM compared to the MSSM from additional bino  $\rightarrow$  singlino transitions

→ An attractive scenario within an attractive model, but:

Finally, after  $\sim 35$  years of Susy extensions of the Standard Model, DATA will decide and NOT “naturalness arguments” of theorists (including mine...)