

QCD Factorization of Hard Exclusive Processes and Tomographic Image of Proton

Zhite Yu

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In collaboration with Jianwei Qiu

JHEP 08 (2022) 103

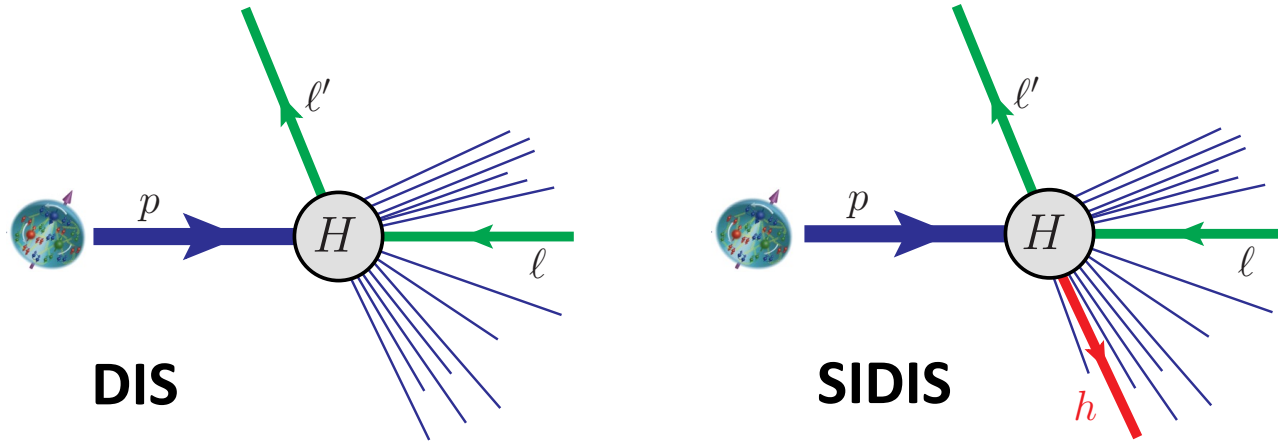
PRD 107 (2023) 014007

PRL 131 (2023) 161902

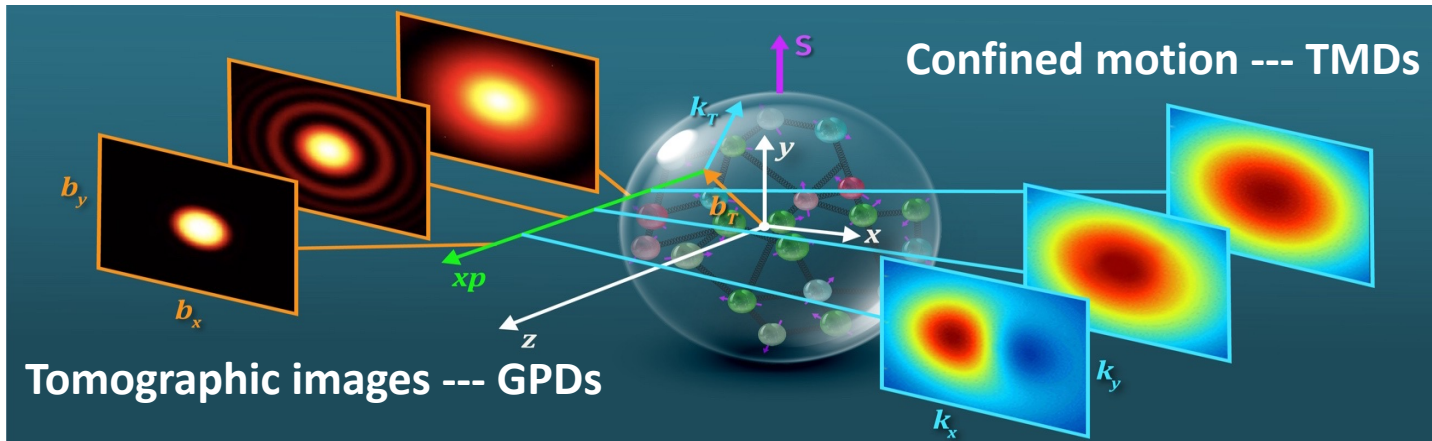
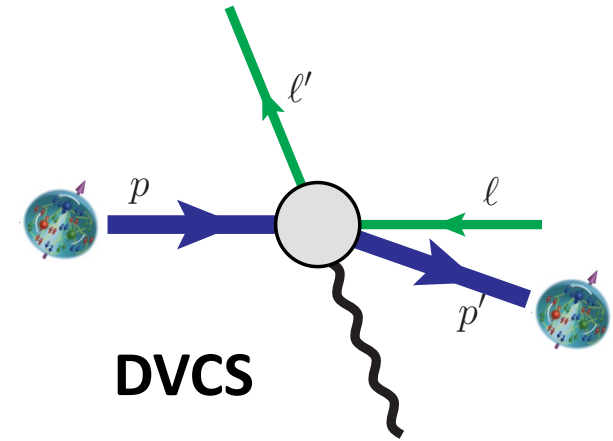
PRD 109 (2024) 074023

Hard probes and proton structure

Inclusive scattering

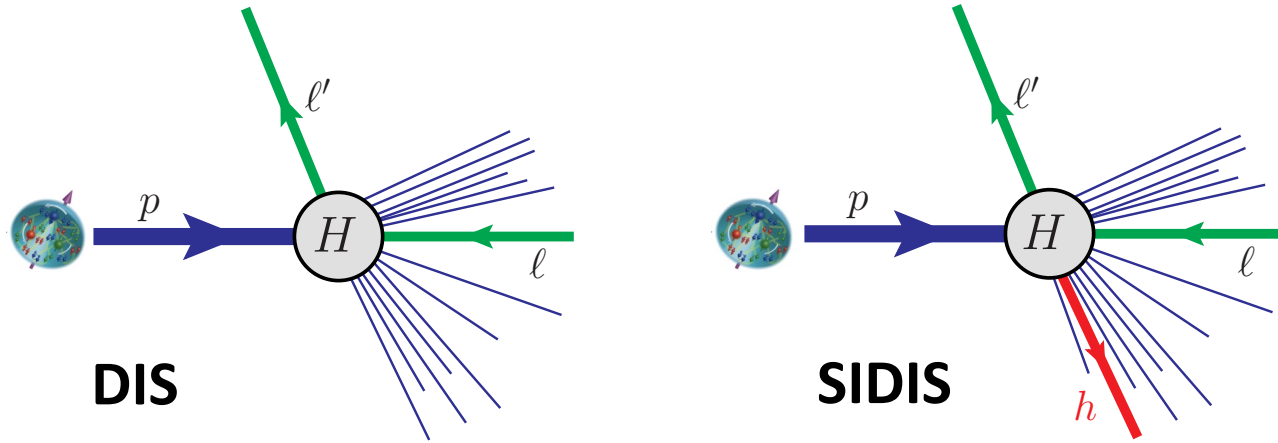


Exclusive diffraction

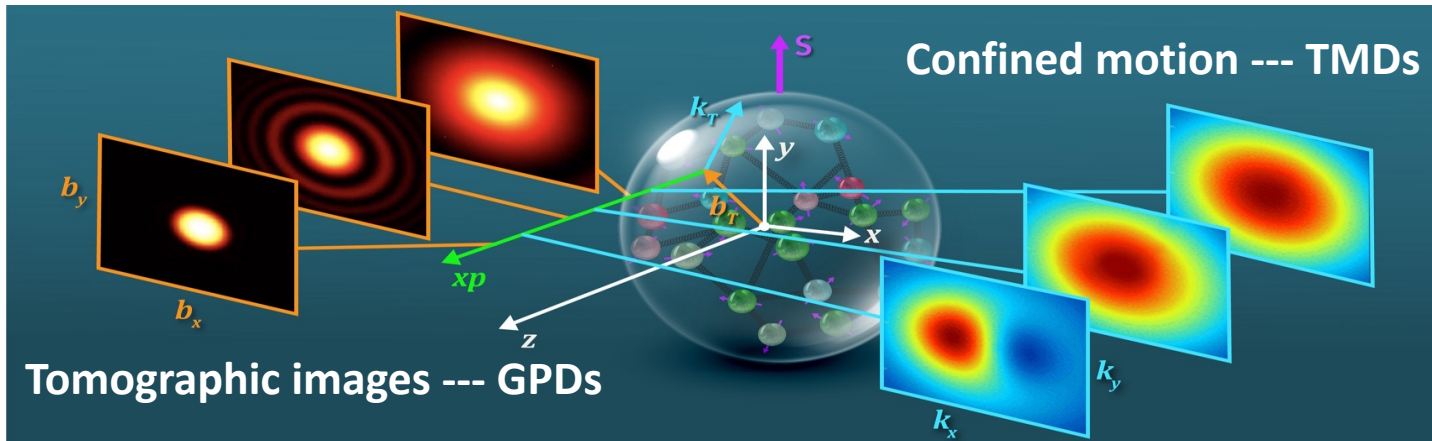
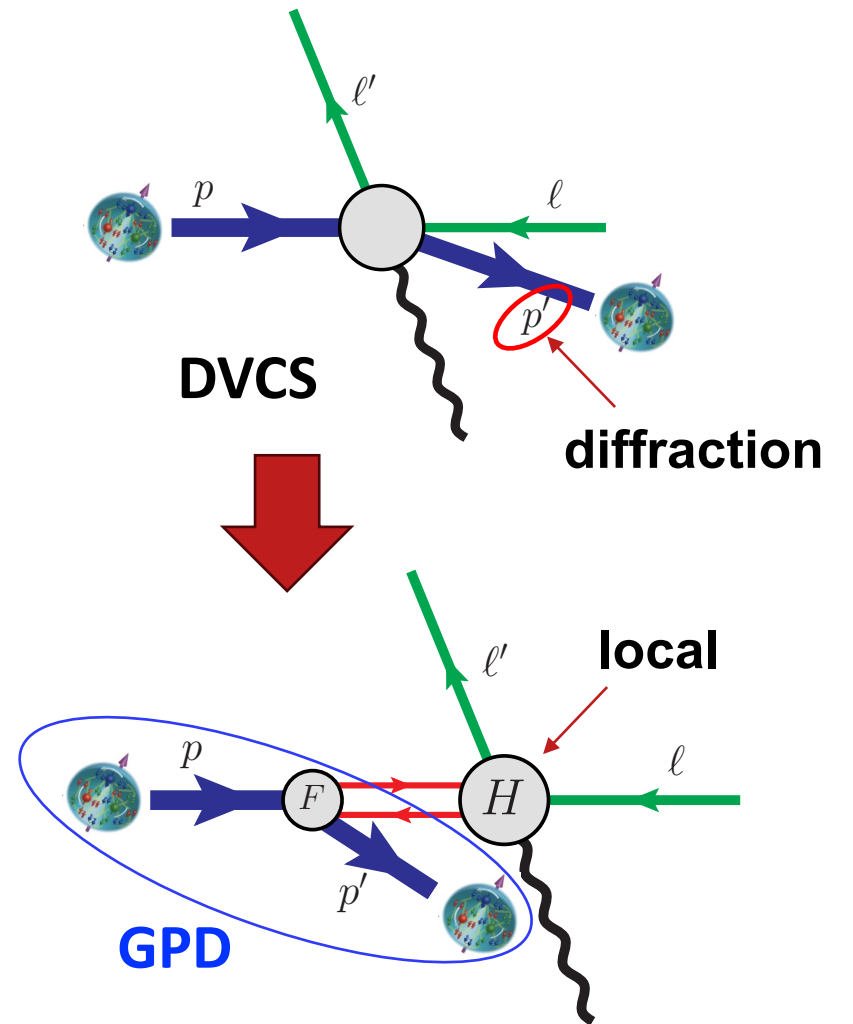


Exclusive processes and GPDs

Inclusive scattering

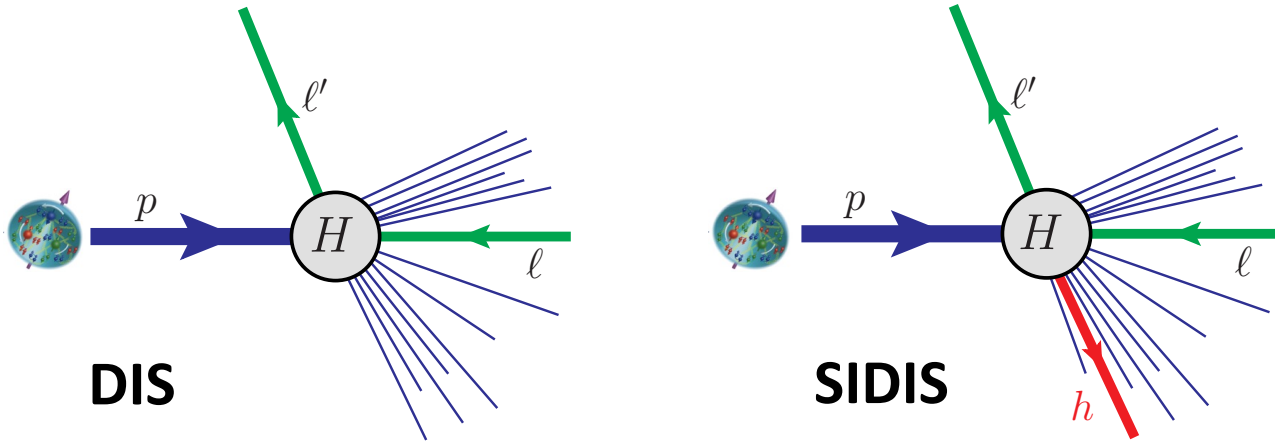


Exclusive diffraction

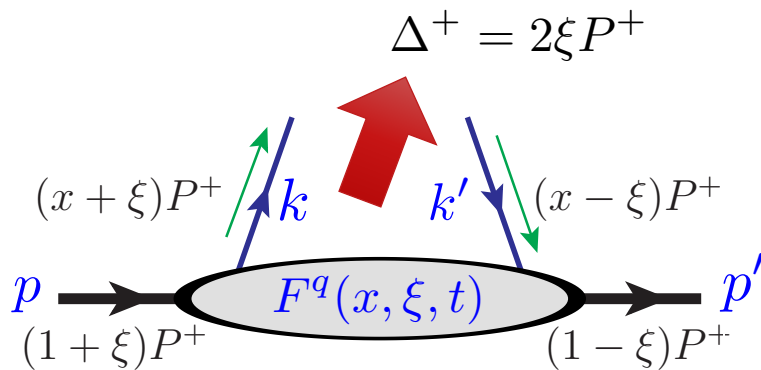
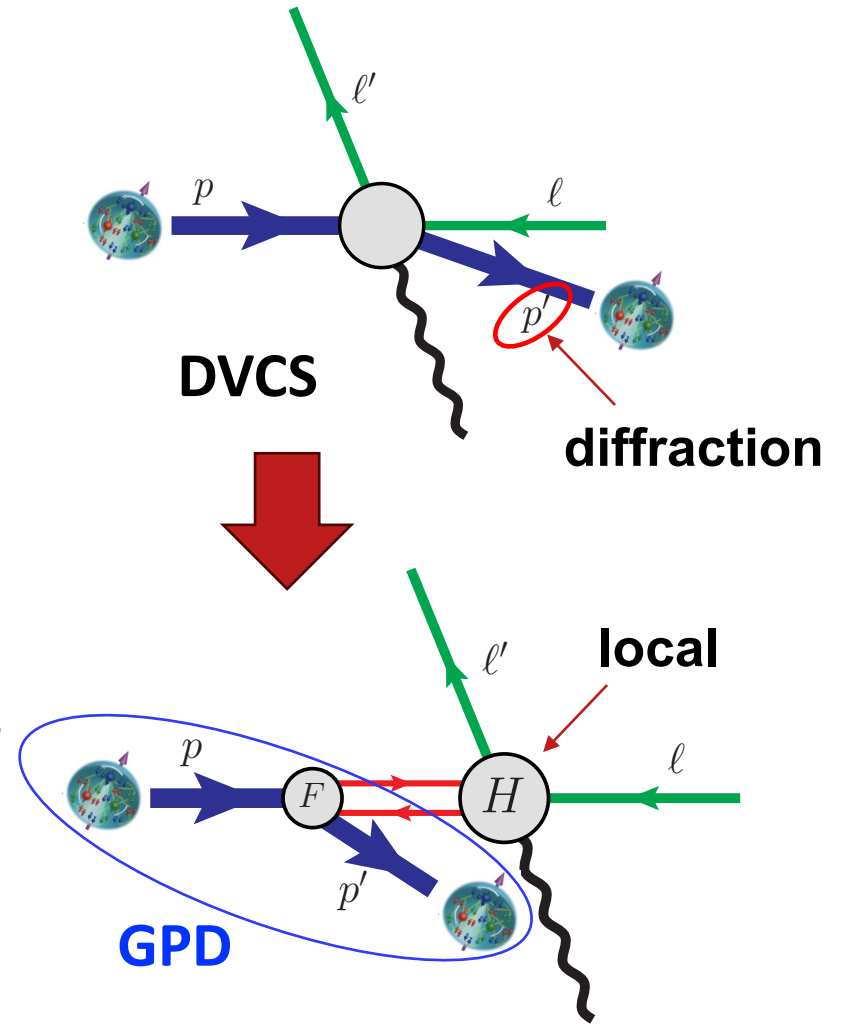


Exclusive processes and GPDs

Inclusive scattering



Exclusive diffraction

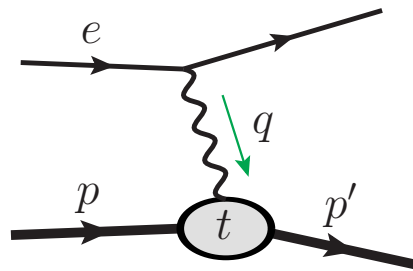


$$F^q(x, \xi, t) = \int \frac{dz^-}{4\pi} e^{-ixP^+z^-} \langle p' | \bar{q}(z^-/2) \gamma^+ q(-z^-/2) | p \rangle$$

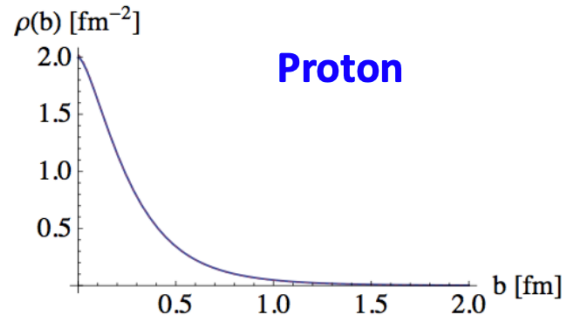
$$= \frac{1}{2P^+} \left[H^q(x, \xi, t) \bar{u}(p') \gamma^+ u(p) - E^q(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\alpha} \Delta_\alpha}{2m} u(p) \right]$$

GPD and 3D tomography

□ Diffraction probes form factors and **spatial density**



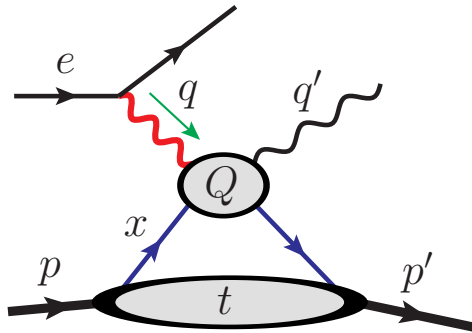
F. T.
 \longrightarrow
 $F_{1,2}(t)$



Proton

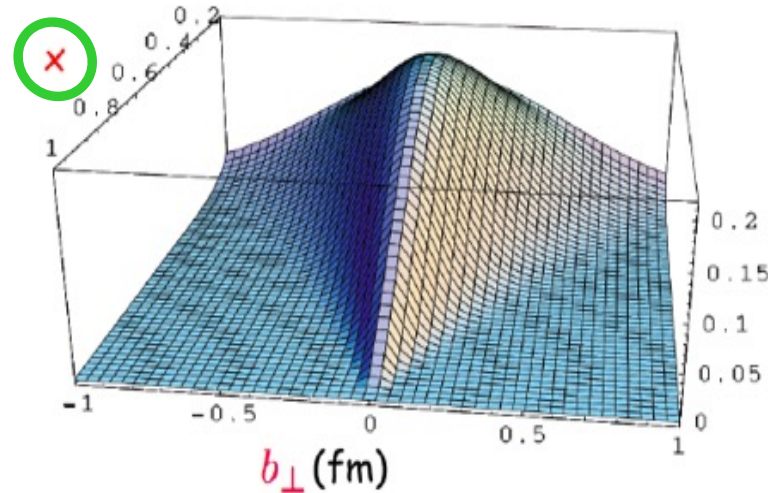
\longrightarrow Electric charge radius

□ **Two-scale** diffraction probes 3D tomography



F. T.
 \longrightarrow

3D image



\longrightarrow Proton radius
in terms of
parton contents

$$f_i(x, \mathbf{b}_T) = \int d^2 \Delta_T e^{i \Delta_T \cdot \mathbf{b}_T} F_i(x, 0, -\Delta_T^2)$$

Parton density in $dx d^2 \mathbf{b}_T$

[M. Burkardt, 2000, 2003]

GPD and Hadronic Property

QCD energy-momentum tensor

$$T^{\mu\nu} = \sum_{i=q,g} T_i^{\mu\nu} \quad \text{with} \quad T_q^{\mu\nu} = \bar{\psi}_q i\gamma^{(\mu} \overleftrightarrow{D}^{\nu)} \psi_q - g^{\mu\nu} \bar{\psi}_q (i\gamma \cdot \overleftrightarrow{D} - m_q) \psi_q \quad \text{and} \quad T_g^{\mu\nu} = F^{a,\mu\eta} F^{a,\eta\nu} + \frac{1}{4} g^{\mu\nu} (F_{\rho\eta}^a)^2$$

Gravitational form factor

$$\langle p' | T_i^{\mu\nu} | p \rangle = \bar{u}(p') \left[A_i(t) \frac{P^\mu P^\nu}{m} + J_i(t) \frac{iP^{(\mu} \sigma^{\nu)\Delta}}{2m} + D_i(t) \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{4m} + m \bar{c}_i(t) g^{\mu\nu} \right] u(p)$$

Connection to GPD moments

$$\int_{-1}^1 dx x F_i(x, \xi, t) \propto \langle p' | T_i^{++} | p \rangle \propto \bar{u}(p') \left[\underbrace{(A_i + \xi^2 D_i)}_{\int_{-1}^1 dx x H_i(x, \xi, t)} \gamma^+ + \underbrace{(B_i - \xi^2 D_i)}_{\int_{-1}^1 dx x E_i(x, \xi, t)} \frac{i\sigma^{+\Delta}}{2m} \right] u(p)$$

Angular momentum sum rule

$$J_i = \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H_i(x, \xi, t) + E_i(x, \xi, t)]$$

$$i = q, g$$

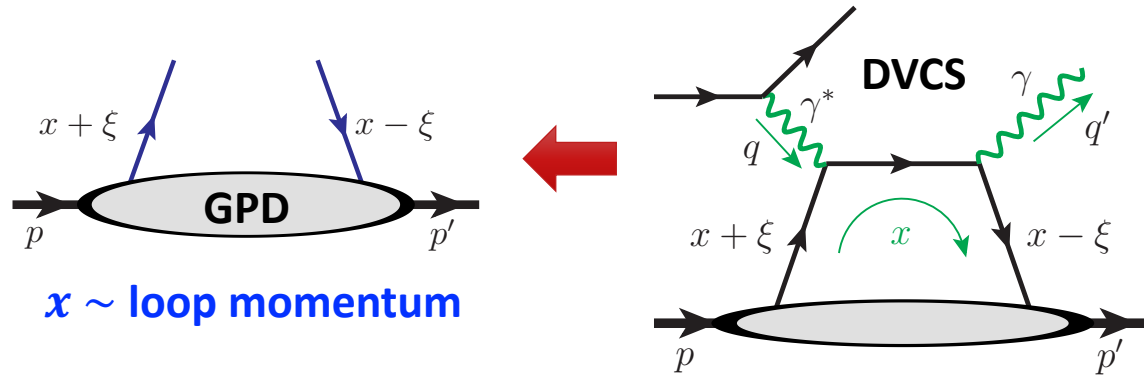
[X.-D. Ji, 1997]

- 3D tomography
- relations to GFF
- angular momentum
- ...

 **x-dependence!**

x -dependence problem for GPD --- why is it so difficult?

□ Amplitude nature: exclusive processes

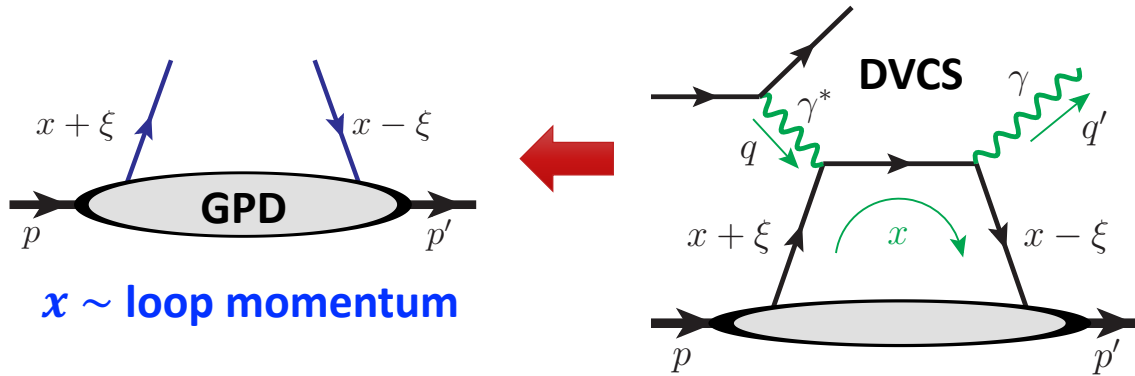


$$i\mathcal{M} \sim \int_{-1}^1 dx F(x, \xi, t) \cdot C(x, \xi; Q/\mu)$$

never pin down to some x

x -dependence problem for GPD --- why is it so difficult?

Amplitude nature: exclusive processes



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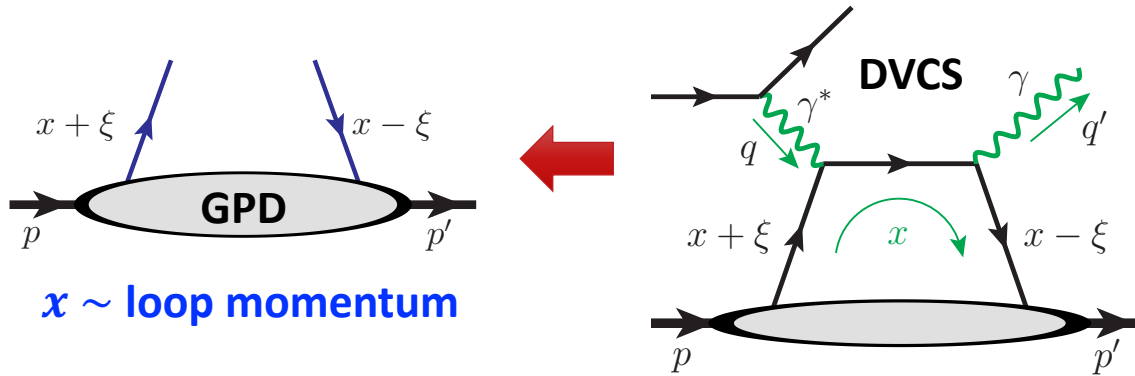
Compare with DIS

cross section: cut diagram

$$\sigma_{\text{DIS}} \simeq \int_{x_B}^1 dx f(x) \hat{\sigma}(x/x_B)$$

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Amplitude nature: exclusive processes



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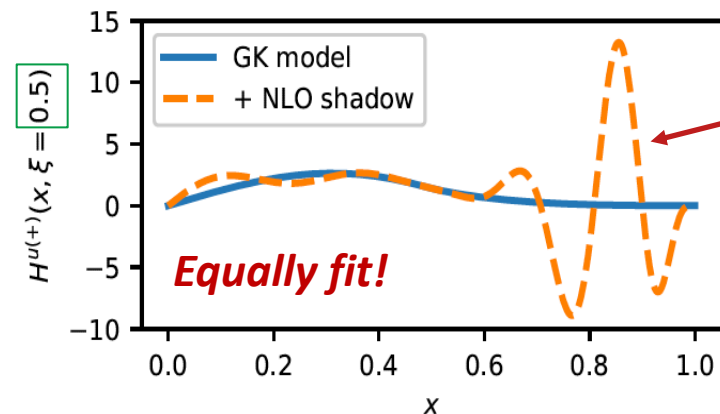
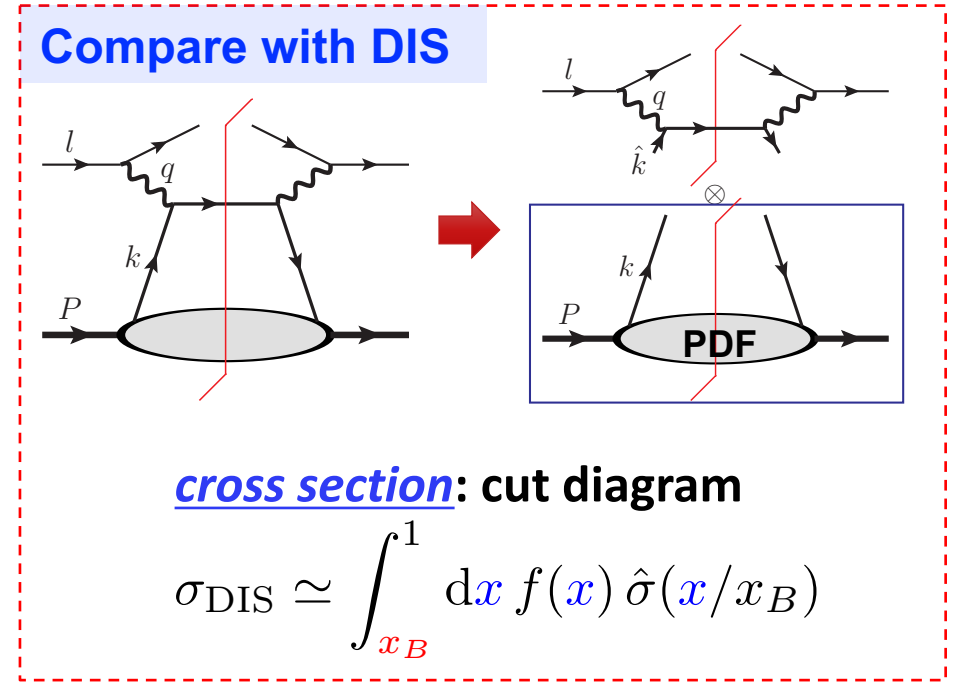
never pin down to some x

Sensitivity to x : comes from $C(x, \xi; Q/\mu)$

$$C(x, \xi; Q/\mu) = T(Q/\mu) \cdot G(x, \xi) \propto \frac{1}{x - \xi + i\epsilon} \dots$$

$$\Rightarrow i\mathcal{M} \propto \int_{-1}^1 dx \frac{F(x, \xi, t)}{x - \xi + i\epsilon} \equiv \text{“}F_0(\xi, t)\text{”} \quad \text{“moment”}$$

“LO scaling”



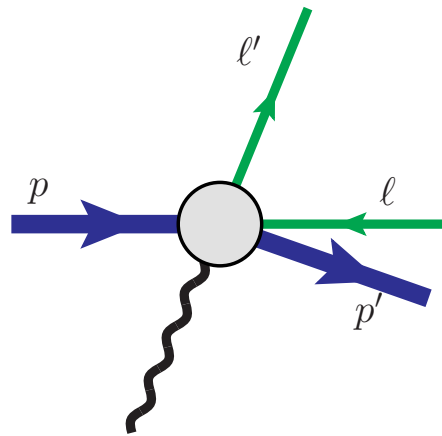
[Bertone et al. PRD '21]

A unified framework of GPD processes

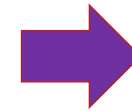
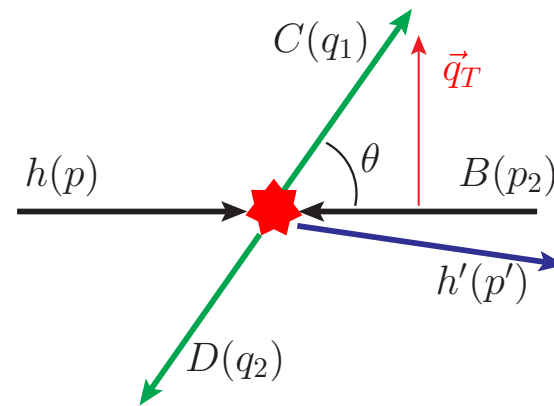
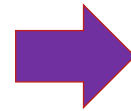
□ Single diffractive hard exclusive process (SDHEP)

[Qiu & Yu, PRD 107 (2023) 014007]

DVCS in lab frame



$$h(p) + B(p_2) \rightarrow h'(p') + C(q_1) + D(q_2)$$



2 → 3: *minimal*
kinematic configuration!

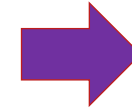
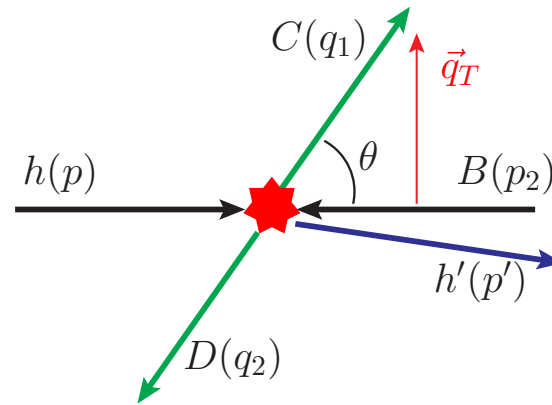
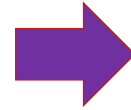
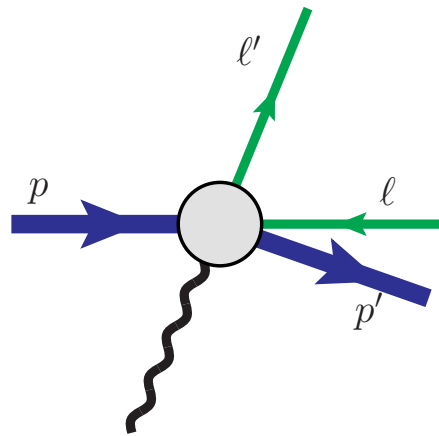
A unified framework of GPD processes

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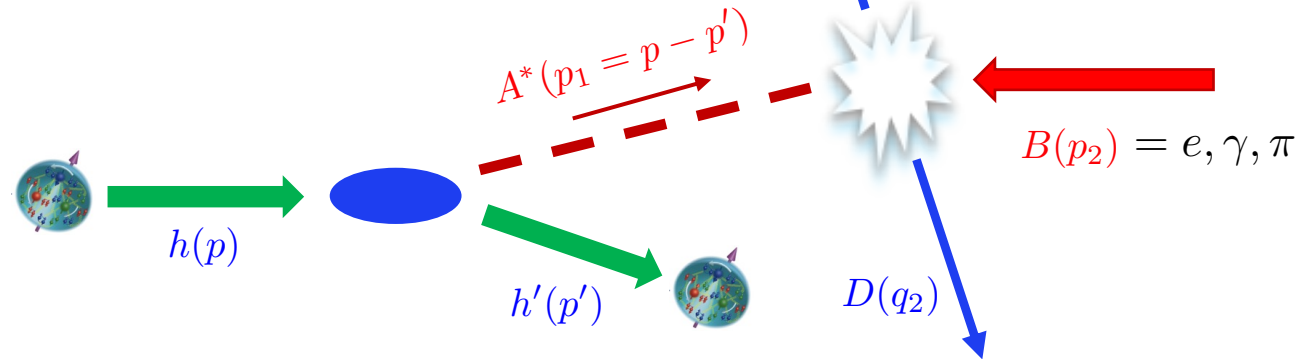
2 → 3: minimal kinematic configuration!

□ Two-stage process paradigm

Single diffractive: $h(p) \rightarrow h'(p') + A^*(p_1 = p - p')$

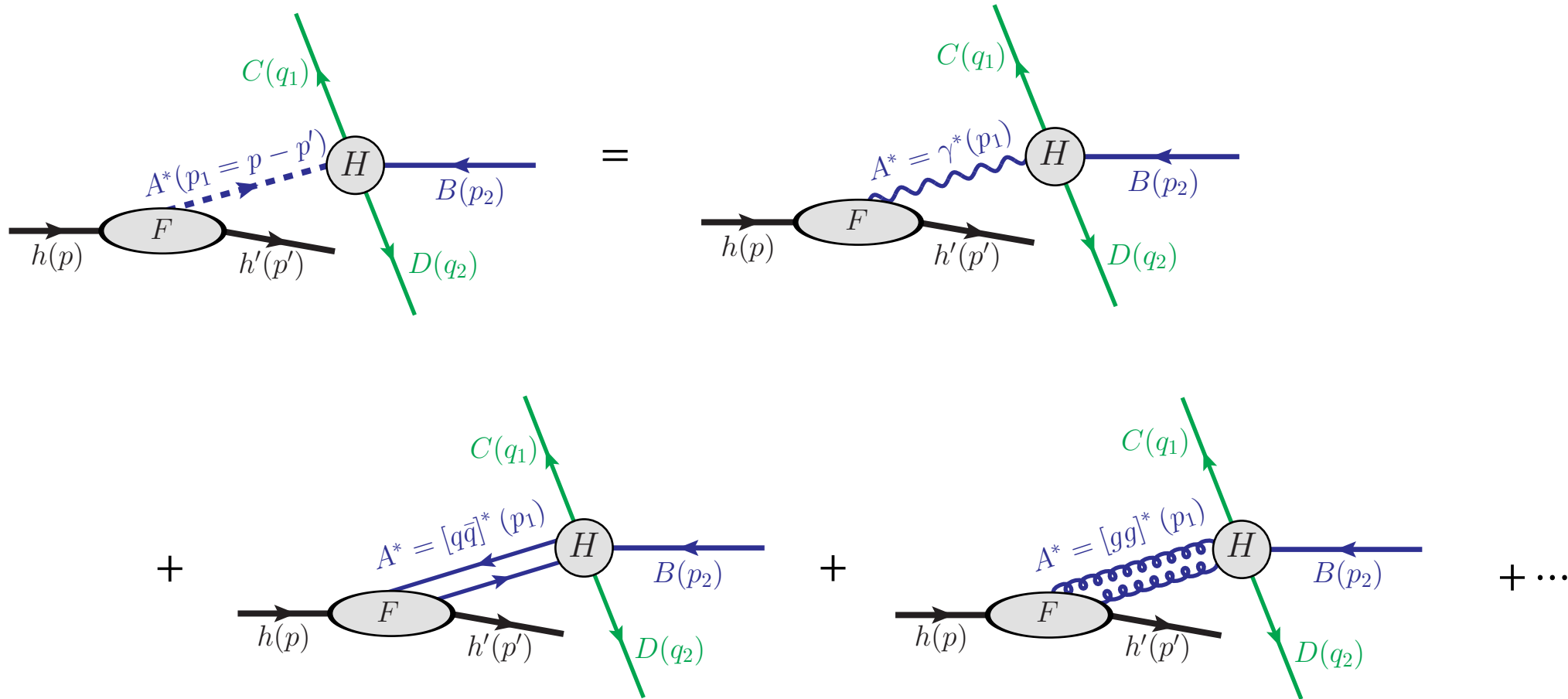
↓ factorize

Hard exclusive: $A^*(p_1) + B(p_2) \rightarrow C(p_3) + D(p_4)$

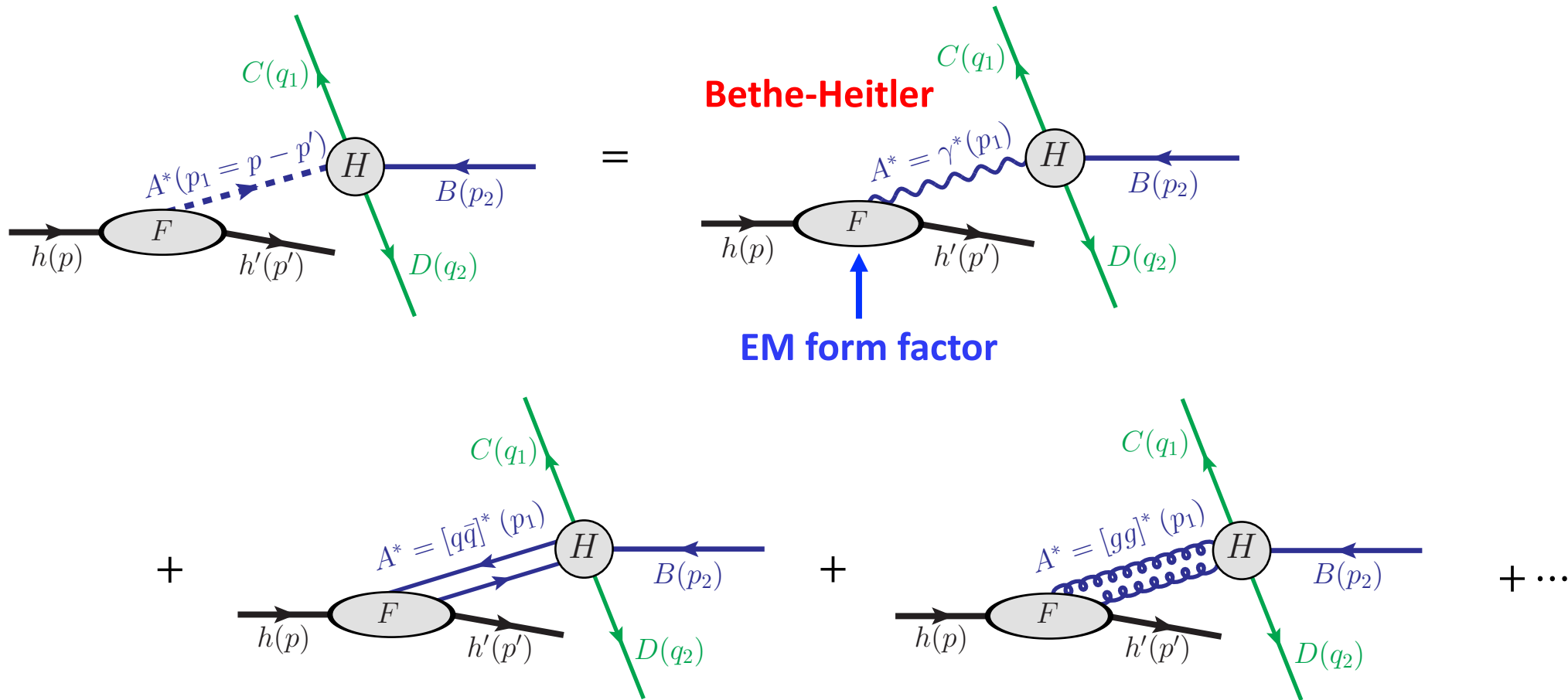


Necessary condition for factorization: $q_T \gg \sqrt{-t} \simeq \Lambda_{\text{QCD}} \quad t = (p - p')^2$

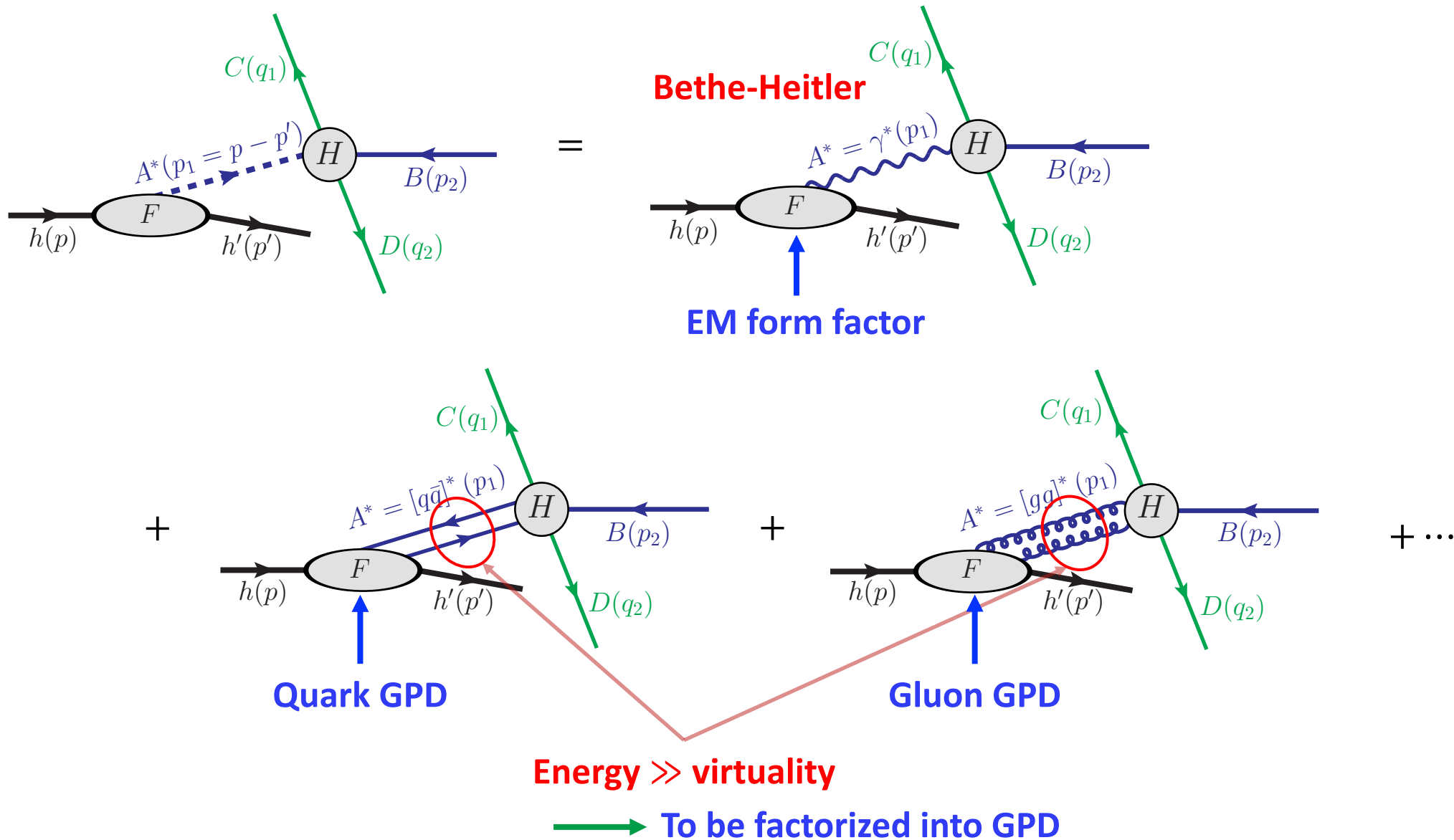
SDHEP: Two-stage paradigm and channel expansion



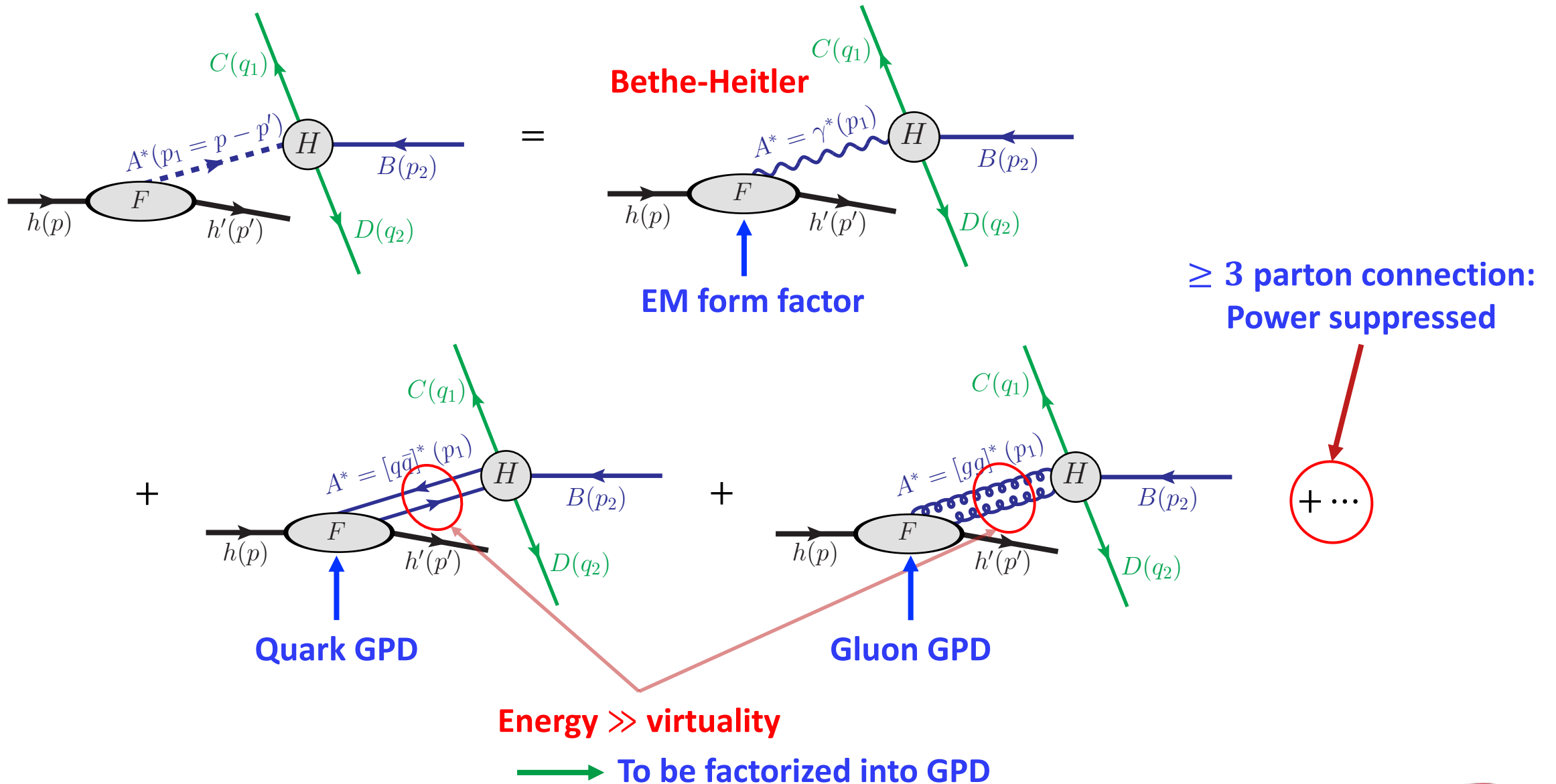
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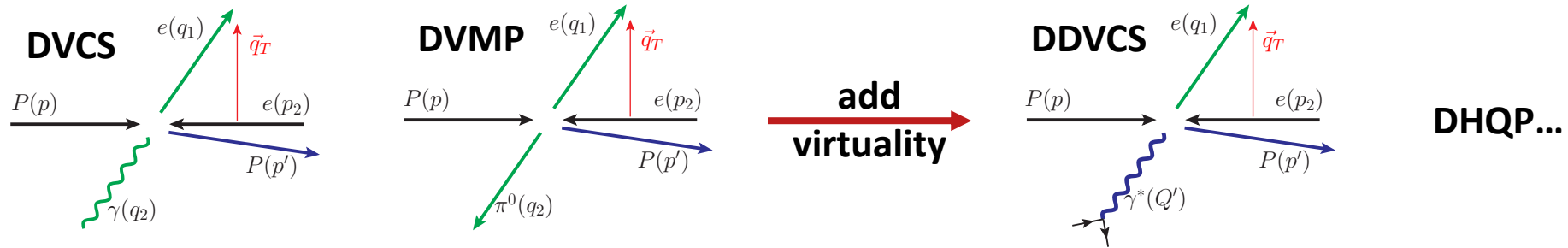


SDHEP: Two-stage paradigm and channel expansion (**twist expansion**)

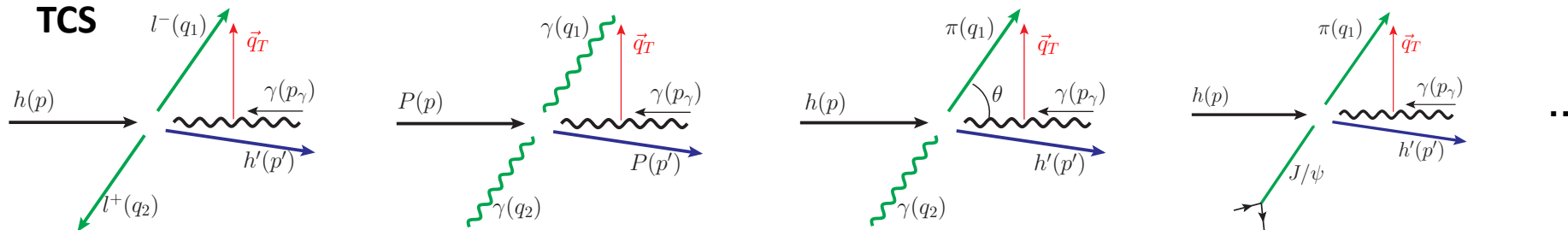


Classification of SDHEPs

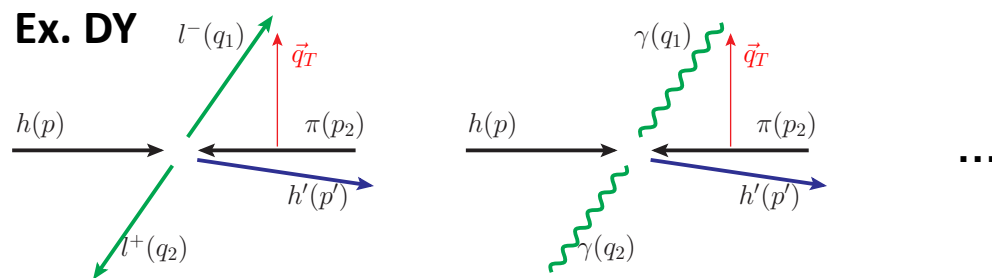
□ Electro-production (JLab, EIC, ...)



□ Photo-production (JLab, EIC, ...)



□ Meso-production (AMBER, J-PARC, ...)

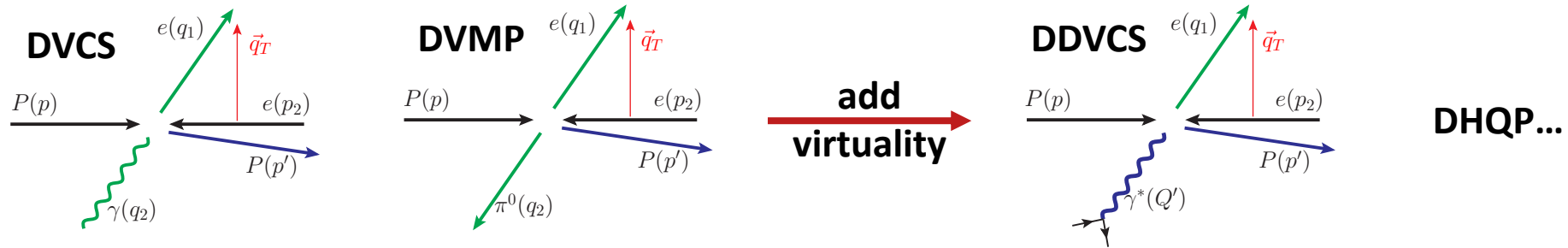


Generic discussion

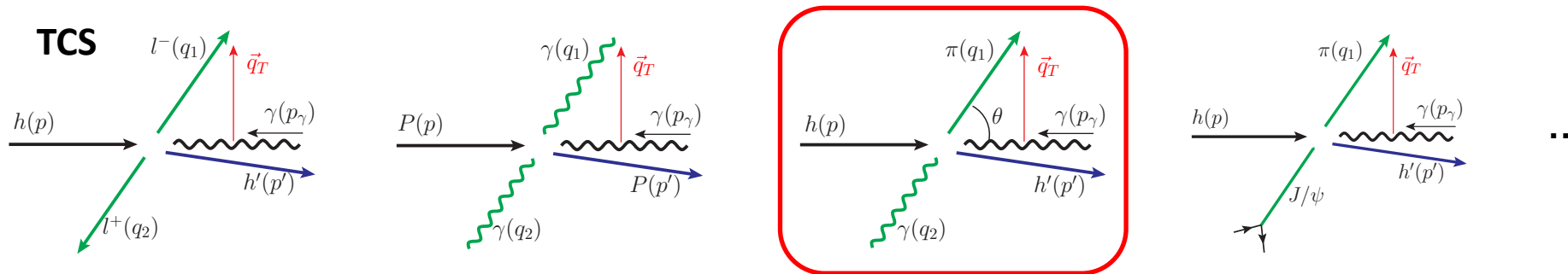
[Qiu, Yu, PRD 107 (2023), 014007]

Classification of SDHEPs

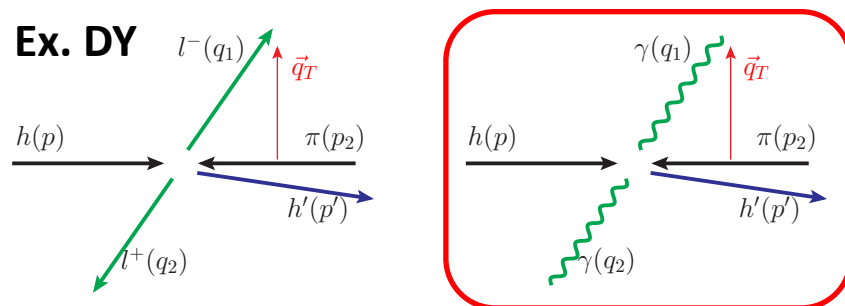
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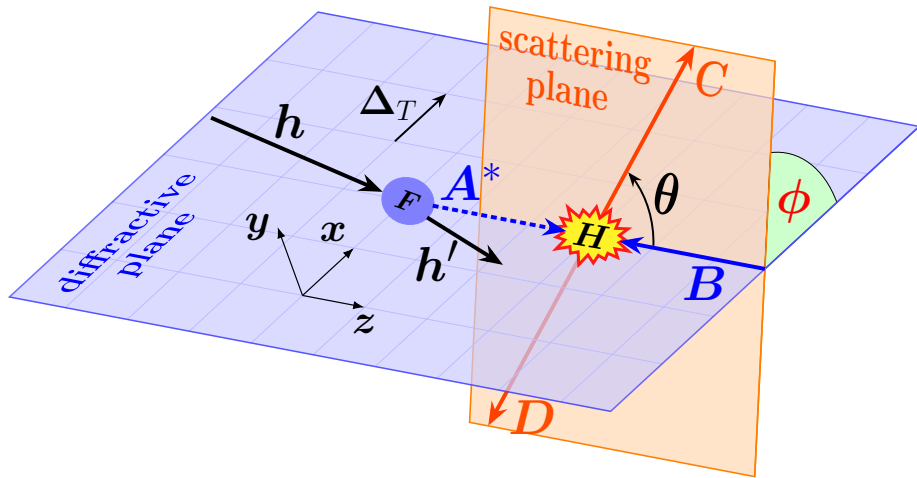
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Generic discussion

[Qiu, Yu, PRD 107 (2023), 014007]

Where does the **x -sensitivity** come from?



□ **x -sensitivity $\Leftrightarrow 2 \rightarrow 2$ hard scattering**

Kinematics:

1. $\hat{s} = 2 \xi s / (1 + \xi)$ ← ξ

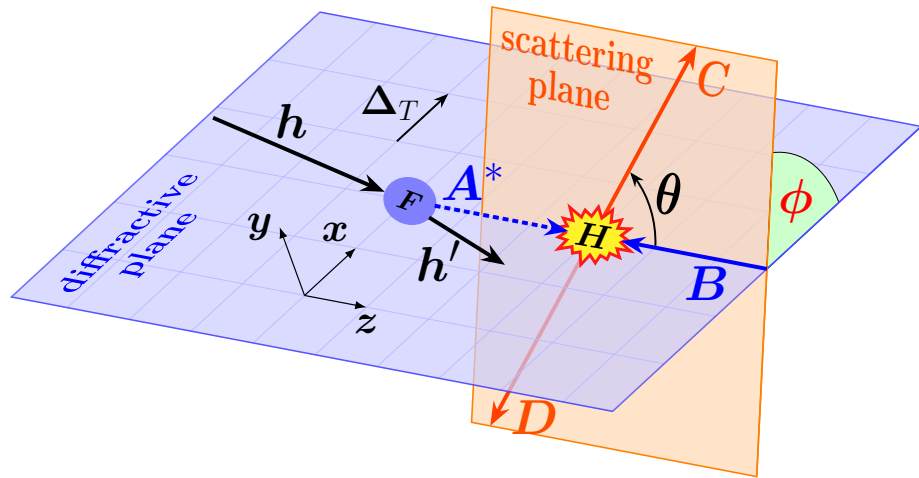
2. θ or $q_T = (\sqrt{\hat{s}}/2) \sin\theta$ ↔ x

3. ϕ ← (A^*B) spin states

$$\mathcal{M}(Q, \phi) \simeq \sum_A e^{i(\lambda_A - \lambda_B)\phi} \cdot \int_{-1}^1 dx F_A(x) C_A(x; Q) \quad (Q = \theta \text{ or } q_T)$$

[suppressing t and ξ dependence]

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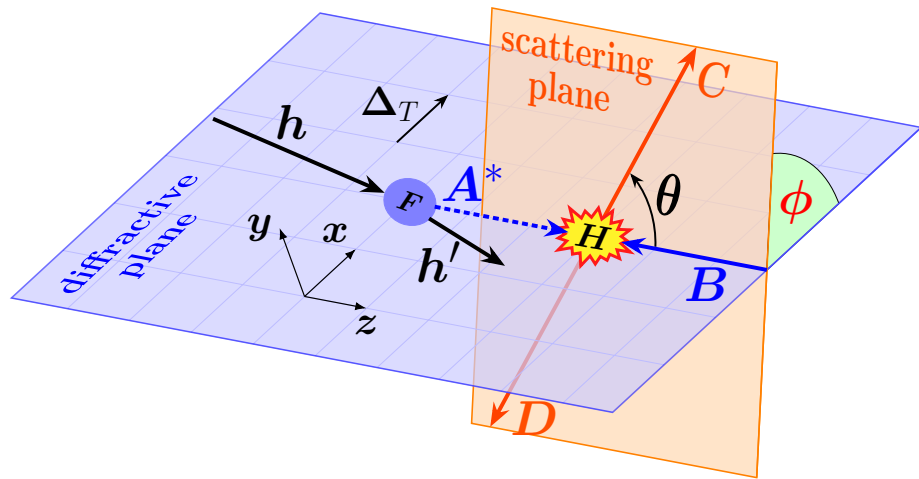
[suppressing t and ξ dependence]

➤ **Moment-type sensitivity** $C(x; Q) = G(x) \cdot T(Q)$ → $F_G = \int_{-1}^1 dx G(x) F(x, \xi, t)$ **Independent of Q .**
Scaling for F_G .

➔ **Inversion problem: shadow GPD**

$$S_G = \int_{-1}^1 dx G(x) S(x, \xi) = 0 \quad [\text{Bertone et al. PRD '21}]$$

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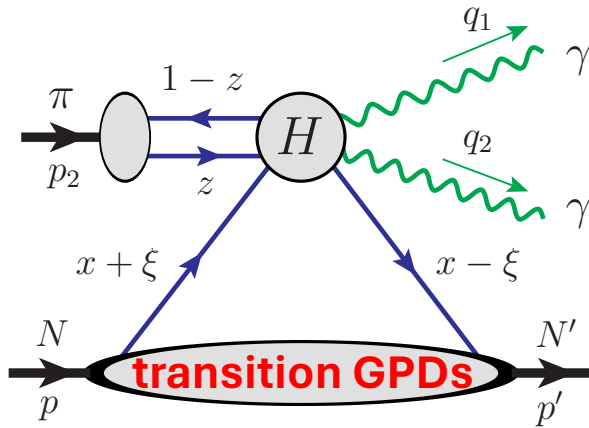
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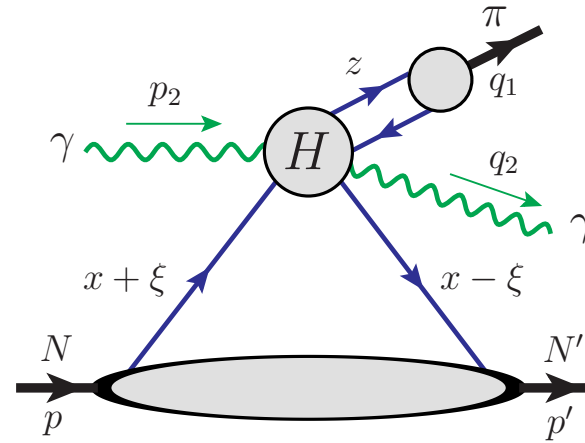
➤ **Enhanced sensitivity** $C(x; Q) \neq G(x) \cdot T(Q)$ → $d\sigma/dQ \sim |C(x; Q) \otimes_x F(x, \xi, t)|^2$

Two new example processes with enhanced x -sensitivity



J-PARC, AMBER

Qiu & Yu, JHEP 08 (2022) 103
 Qiu & Yu, PRD 109 (2024) 074023

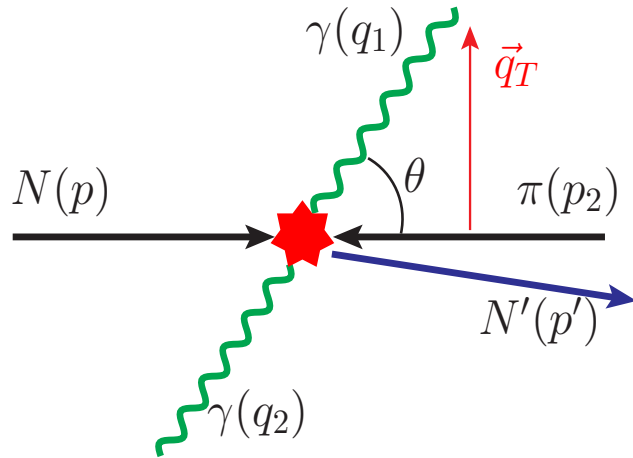


JLab Hall D

G. Duplancic et al., JHEP 11 (2018) 179
 G. Duplancic et al., JHEP 03 (2023) 241
 G. Duplancic et al., PRD 107 (2023), 094023
 Qiu & Yu, PRD 107 (2023), 014007
 Qiu & Yu, PRL 131 (2023), 161902

Enhanced x -sensitivity: (1) diphoton mesoproduction

[Qiu & Yu, JHEP 08 (2022) 103;
PRD 109 (2024) 074023]



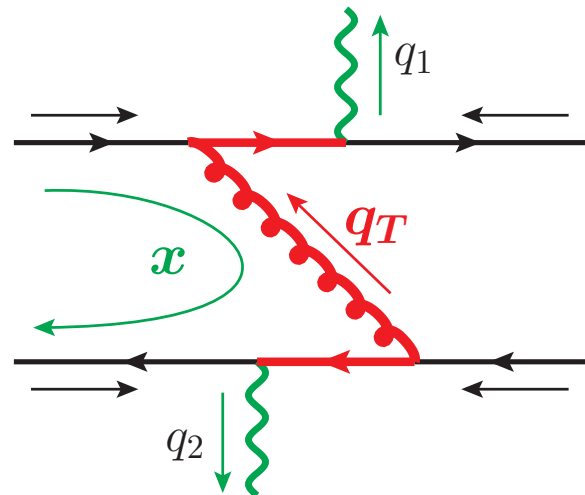
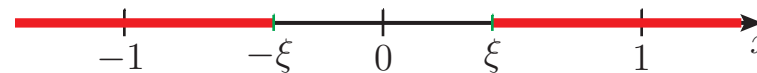
In addition to

$$F_0(\xi, t) = \int_{-1}^1 \frac{dx F(x, \xi, t)}{x - \xi + i\epsilon}$$

$i\mathcal{M}$ also contains

$$I(t, \xi; z, \theta) = \int_{-1}^1 \frac{dx F(x, \xi, t)}{x - \rho(z; \theta) + i\epsilon \operatorname{sgn}[\cos^2(\theta/2) - z]}$$

$$\rho(z; \theta) = \xi \cdot \left[\frac{1 - z + \tan^2(\theta/2) z}{1 - z - \tan^2(\theta/2) z} \right] \in (-\infty, -\xi] \cup [\xi, \infty)$$



Enhanced x -sensitivity: (1) diphoton mesoproduction

[Qiu & Yu, PRD 109 (2024) 074023]

□ **Diphoton process:** $N\pi \rightarrow N'\gamma\gamma$: (1) $p\pi^- \rightarrow n\gamma\gamma$; (2) $n\pi^+ \rightarrow p\gamma\gamma$

$$\frac{d\sigma}{d|t|d\xi d\cos\theta} = 2\pi \left(\alpha_e \alpha_s \frac{C_F}{N_c} \right)^2 \frac{1}{\xi^2 s^3} \cdot \left[(1 - \xi^2) \sum_{\alpha=\pm} \left(|\mathcal{M}_\alpha^{[\tilde{H}]}|^2 + |\tilde{\mathcal{M}}_\alpha^{[H]}|^2 \right) - \left(\xi^2 + \frac{t}{4m^2} \right) \sum_{\alpha=\pm} |\tilde{\mathcal{M}}_\alpha^{[E]}|^2 - \frac{\xi^2 t}{4m^2} \sum_{\alpha=\pm} |\mathcal{M}_\alpha^{[\tilde{E}]}|^2 - 2\xi^2 \sum_{\alpha=\pm} \text{Re} \left(\tilde{\mathcal{M}}_\alpha^{[H]} \tilde{\mathcal{M}}_\alpha^{[E]*} + \mathcal{M}_\alpha^{[\tilde{H}]} \mathcal{M}_\alpha^{[\tilde{E}]*} \right) \right]$$

Nucleon transition GPDs

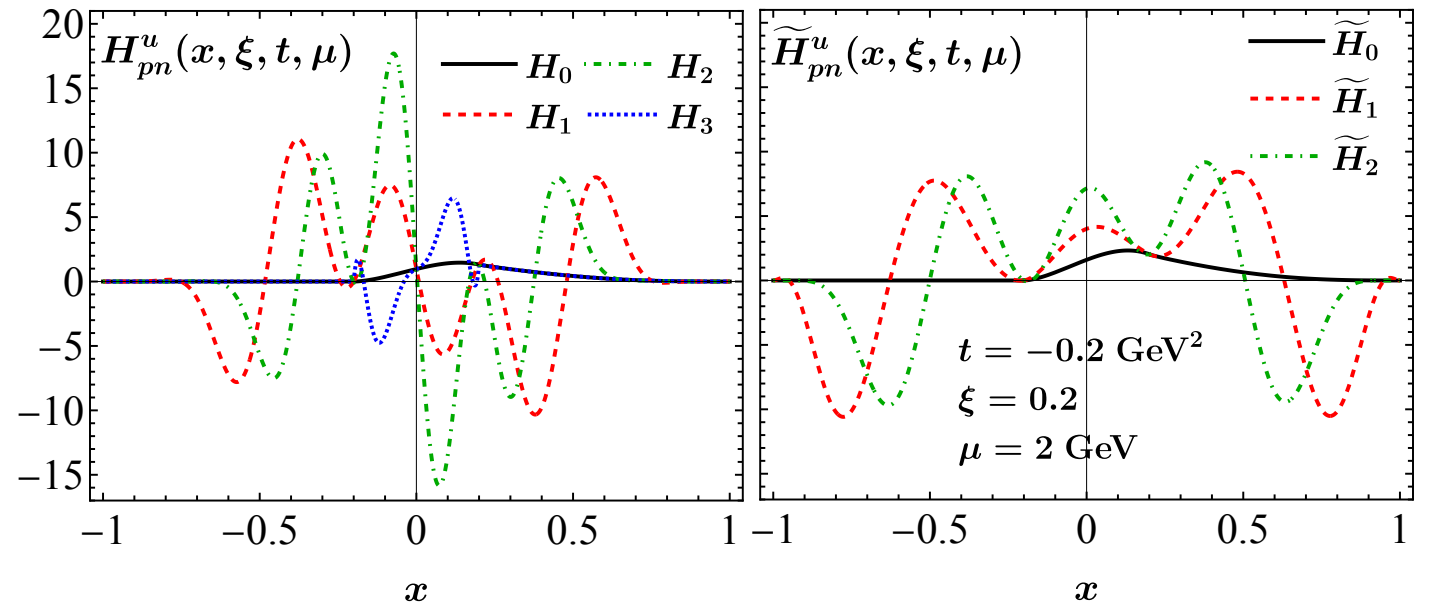
$$H_{pn}^u = H_p^u - H_p^d, \text{ etc.}$$

GPD models = GK model + shadow GPDs

Goloskokov & Kroll, '05, '07, '09

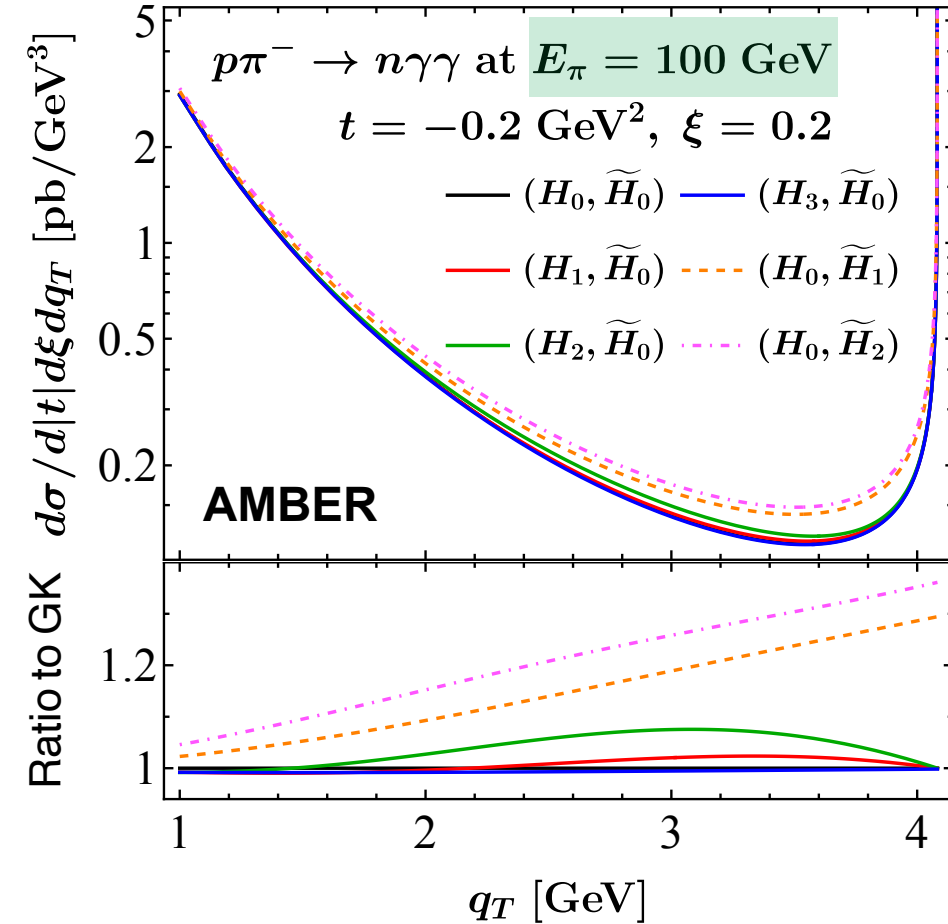
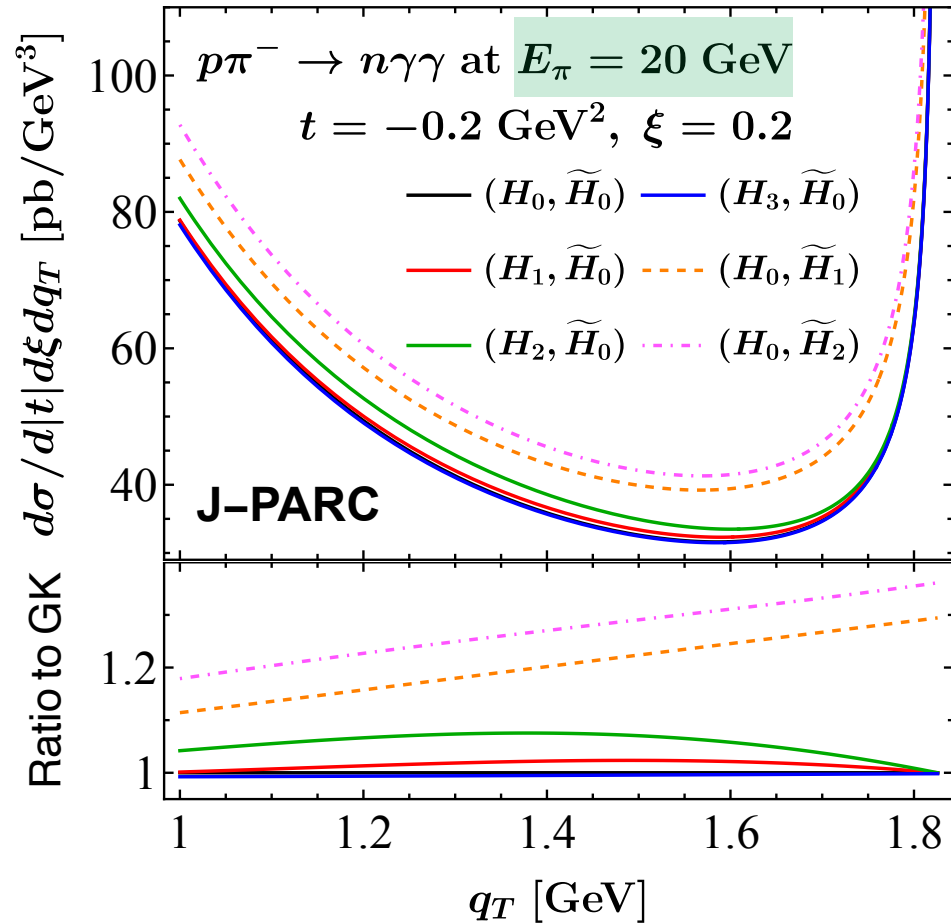
Bertone et al. '21
Moffat et al. '23

$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$



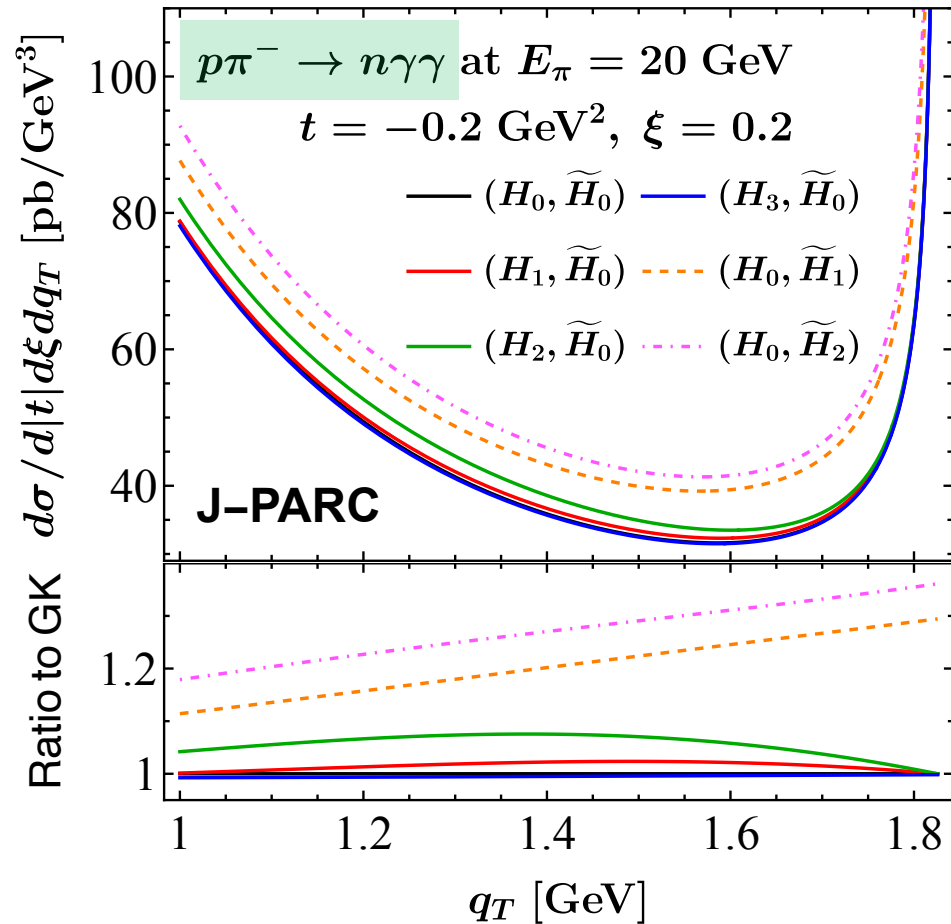
Enhanced x -sensitivity: (1) diphoton mesoproduction (at J-PARC or AMBER)

[Qiu & Yu, PRD 109 (2024) 074023]



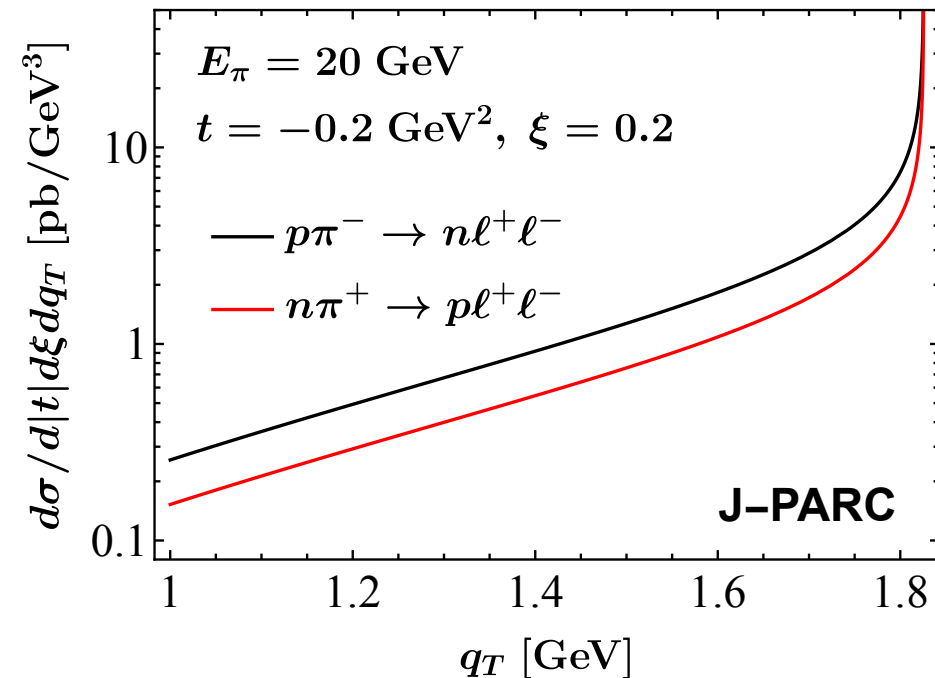
Enhanced x -sensitivity: (1) diphoton mesoproduction (at J-PARC or AMBER)

[Qiu & Yu, PRD 109 (2024) 074023]



Exclusive Drell-Yan dilepton mesoproduction

$$N + \pi \rightarrow N' + \gamma^* [\rightarrow \ell^+ + \ell^-]$$



- **Blind to shadow GPDs**
- **Lower rate**

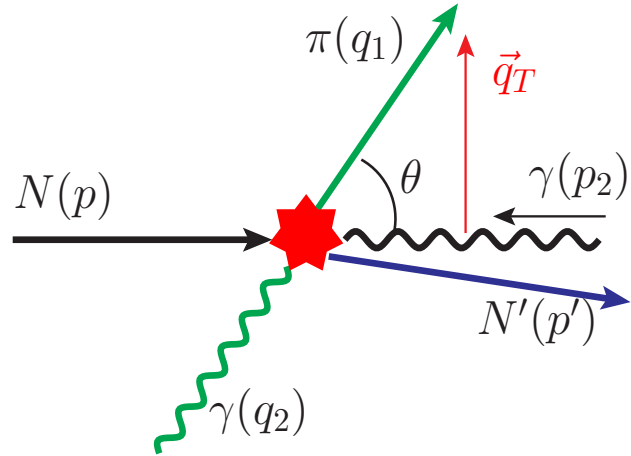
Enhanced x -sensitivity: (2) γ - π pair photoproduction

[Qiu & Yu, PRL 131 (2023) 161902]

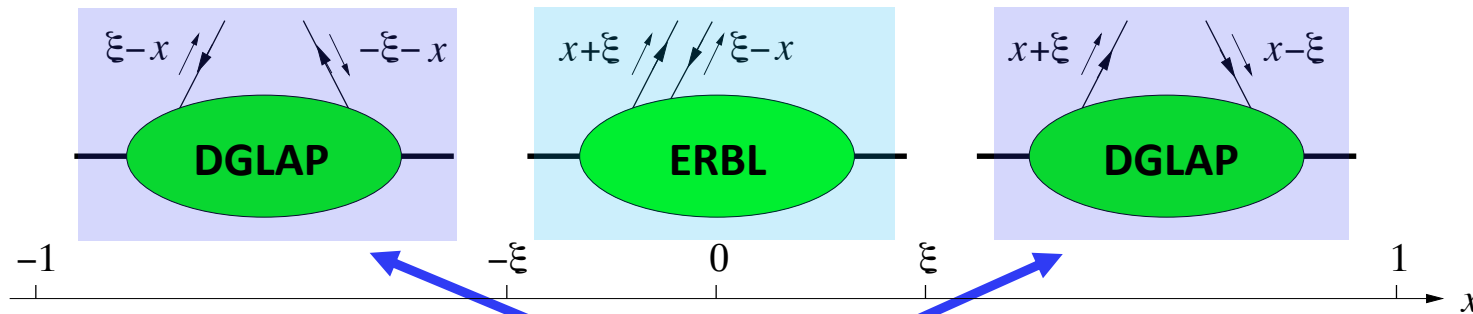
$i\mathcal{M}$ also contains the special integral

$$I'(t, \xi; z, \theta) = \int_{-1}^1 \frac{dx F(x, \xi, t)}{x - \rho'(z; \theta) + i\epsilon}$$

$$\rho'(z; \theta) = \xi \cdot \left[\frac{\cos^2(\theta/2) (1 - z) - z}{\cos^2(\theta/2) (1 - z) + z} \right] \in [-\xi, \xi]$$



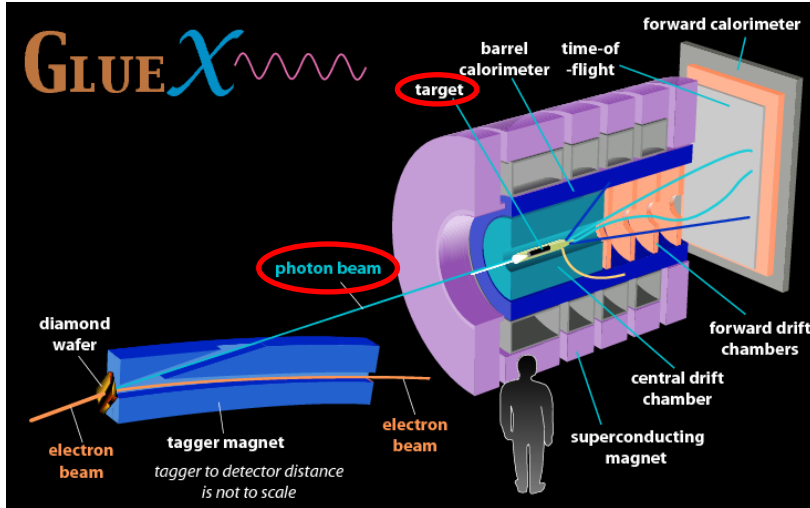
Complementary sensitivity



$$N \pi \rightarrow N' \gamma \gamma$$

Enhanced x -sensitivity: (2) γ - π pair photoproduction (at JLab Hall D)

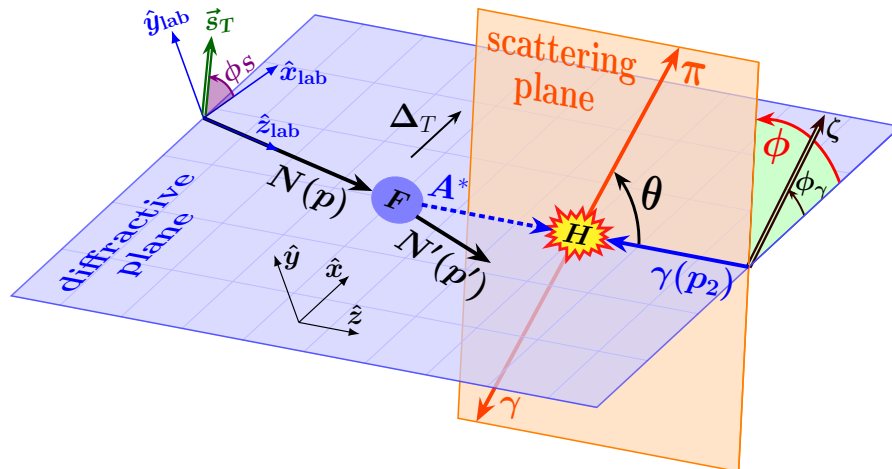
[Qiu & Yu, PRL 131 (2023) 161902]



Polarization asymmetries

$$\frac{d\sigma}{d|t| d\xi d \cos \theta d\phi} = \frac{1}{2\pi} \frac{d\sigma}{d|t| d\xi d \cos \theta} \cdot [1 + \lambda_N \lambda_\gamma A_{LL} + \zeta A_{UT} \cos 2(\phi - \phi_\gamma) + \lambda_N \zeta A_{LT} \sin 2(\phi - \phi_\gamma)]$$

$$\frac{d\sigma}{d|t| d\xi d \cos \theta} = \pi (\alpha_e \alpha_s)^2 \left(\frac{C_F}{N_c} \right)^2 \frac{1 - \xi^2}{\xi^2 s^3} \Sigma_{UU}$$



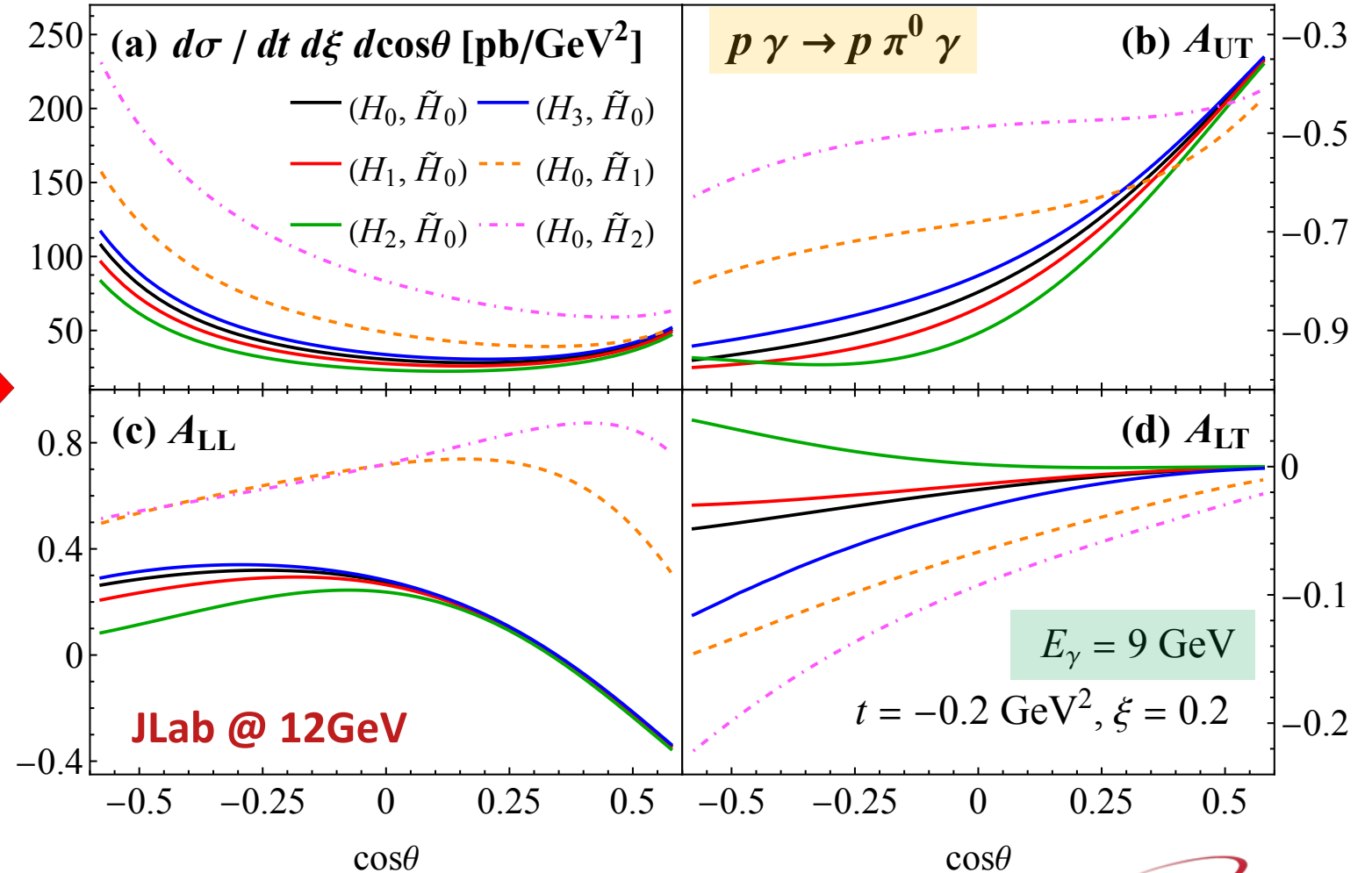
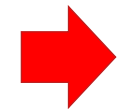
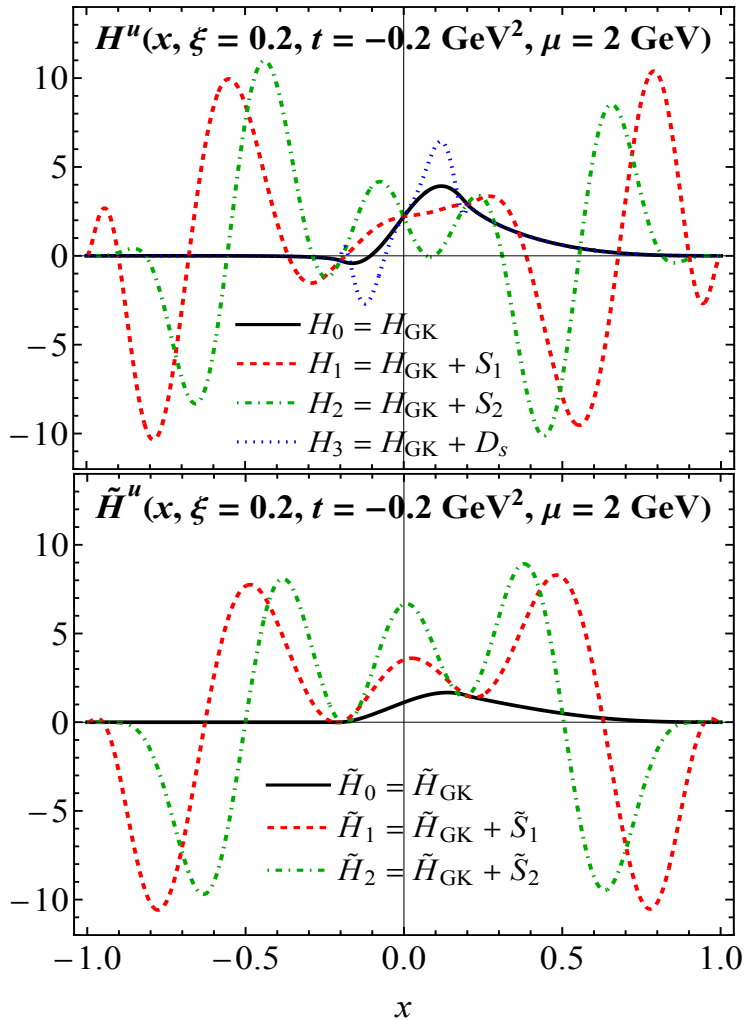
$$\begin{aligned} \Sigma_{UU} &= |\mathcal{M}_+^{[\tilde{H}]}|^2 + |\mathcal{M}_-^{[\tilde{H}]}|^2 + |\tilde{\mathcal{M}}_+^{[H]}|^2 + |\tilde{\mathcal{M}}_-^{[H]}|^2, \\ A_{LL} &= 2 \Sigma_{UU}^{-1} \text{Re} \left[\mathcal{M}_+^{[\tilde{H}]} \tilde{\mathcal{M}}_+^{[H]*} + \mathcal{M}_-^{[\tilde{H}]} \tilde{\mathcal{M}}_-^{[H]*} \right], \\ A_{UT} &= 2 \Sigma_{UU}^{-1} \text{Re} \left[\tilde{\mathcal{M}}_+^{[H]} \tilde{\mathcal{M}}_-^{[H]*} - \mathcal{M}_+^{[\tilde{H}]} \mathcal{M}_-^{[\tilde{H}]*} \right], \\ A_{LT} &= 2 \Sigma_{UU}^{-1} \text{Im} \left[\mathcal{M}_+^{[\tilde{H}]} \tilde{\mathcal{M}}_-^{[H]*} + \mathcal{M}_-^{[\tilde{H}]} \tilde{\mathcal{M}}_+^{[H]*} \right]. \end{aligned}$$

Enhanced x -sensitivity: (2) γ - π pair photoproduction (at JLab Hall D)

GPD models = GK model + shadow GPDs

$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$

Goloskokov, Kroll, '05, '07, '09
 Bertone et al. '21
 Moffat et al. '23

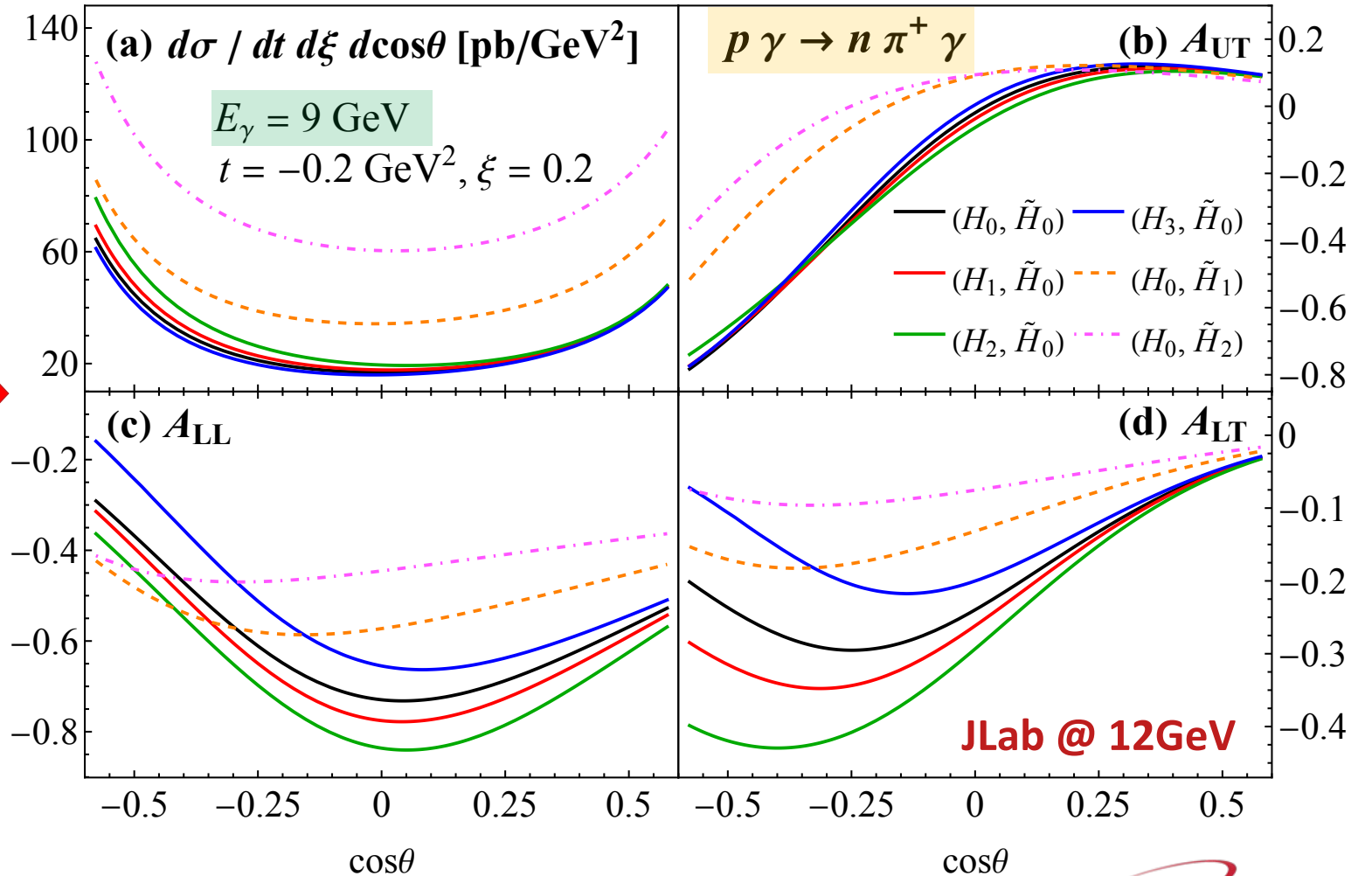
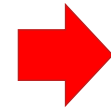
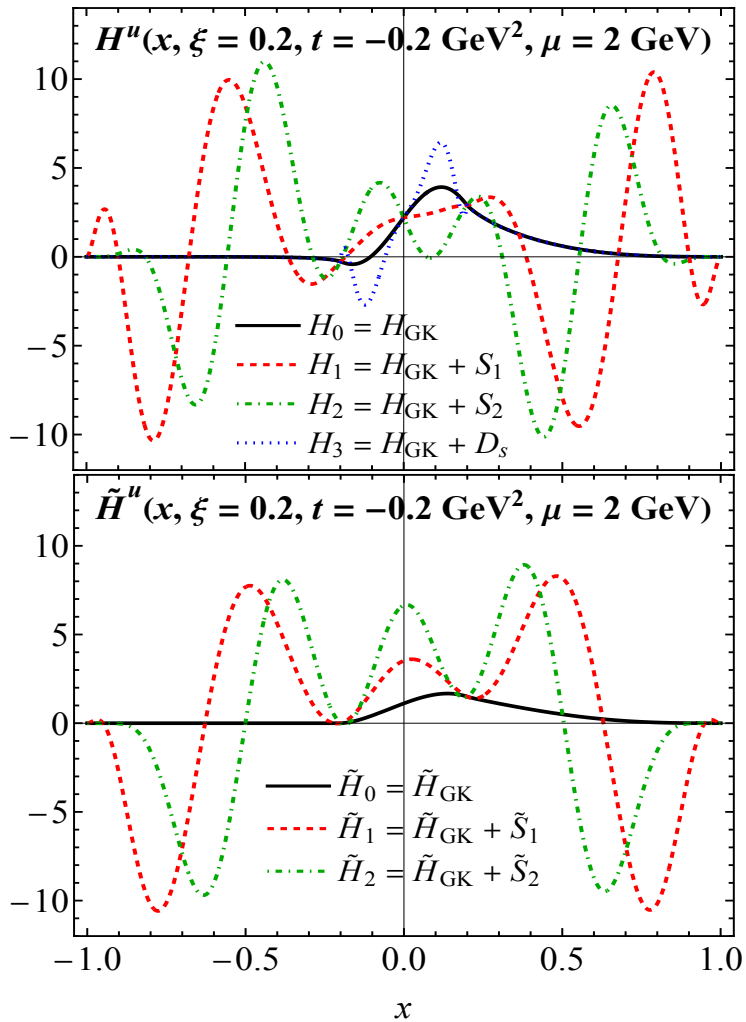


Enhanced x -sensitivity: (2) γ - π pair photoproduction (at JLab Hall D)

GPD models = GK model + shadow GPDs

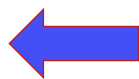
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Goloskokov, Kroll, '05, '07, '09
 Bertone et al. '21
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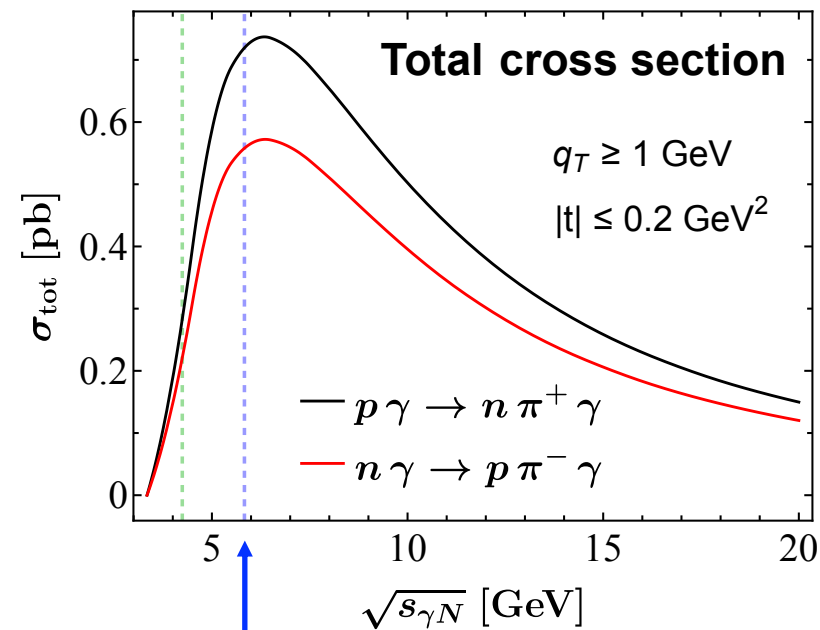
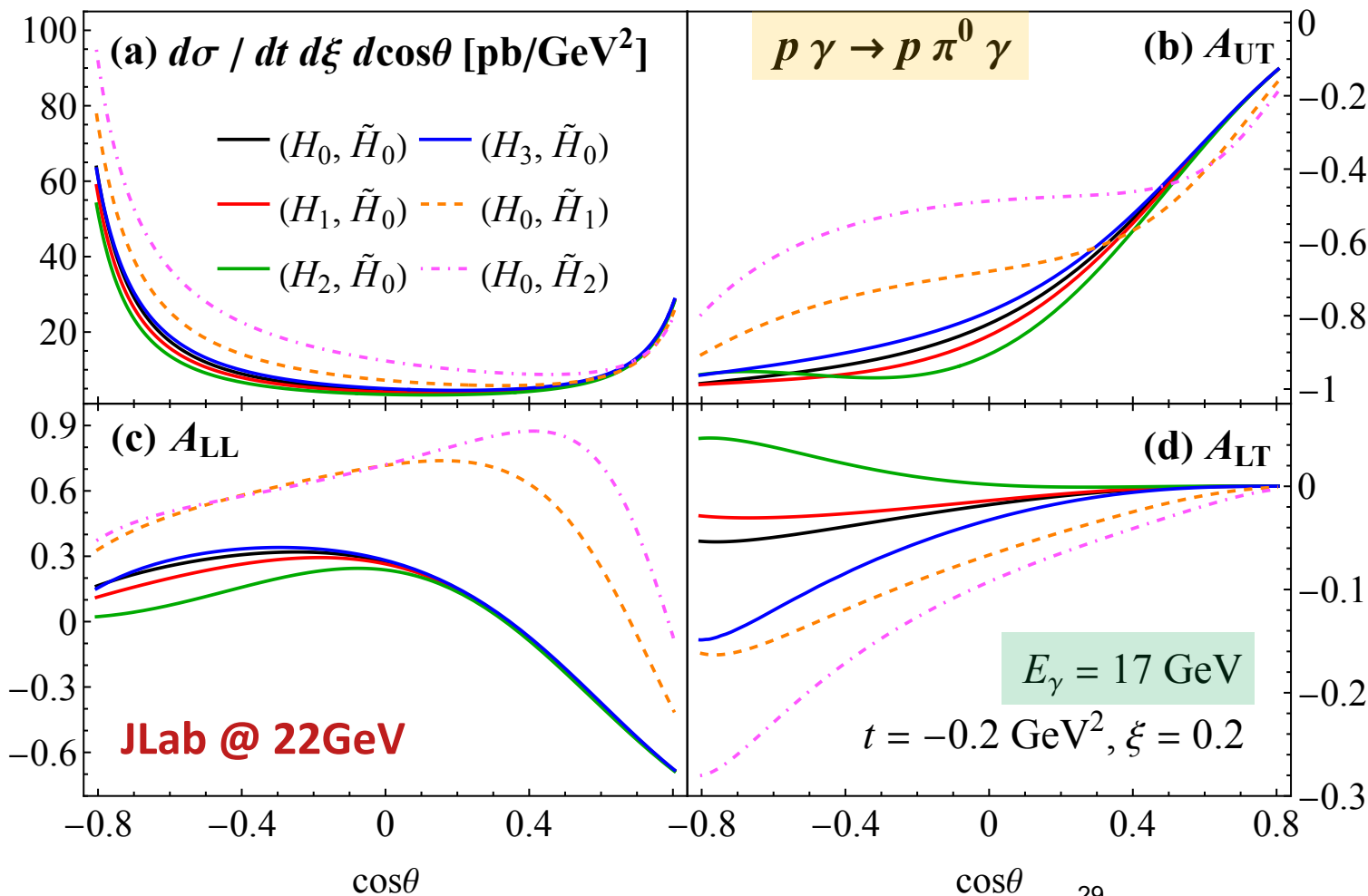
Enhanced x -sensitivity: (2) γ - π pair photoproduction (at upgraded JLab energy)

GPD models = GK model + shadow GPDs



$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$

Goloskokov, Kroll, '05, '07, '09
 Bertone et al. '21
 Moffat et al. '23



JLab @ 22GeV [arXiv:2306.09360]

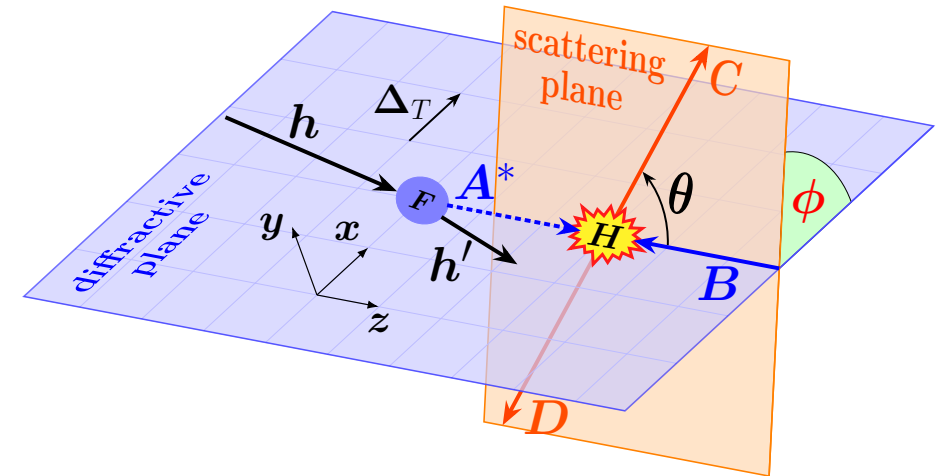
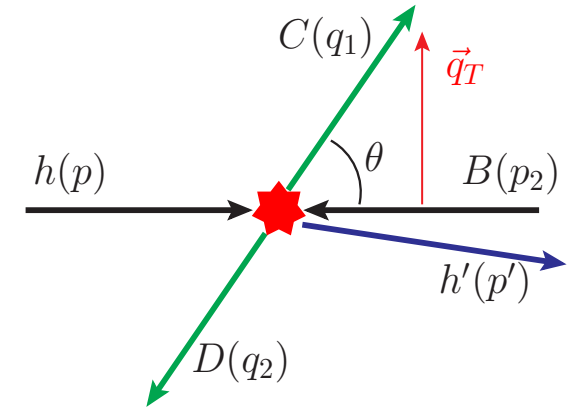
Summary

- ❑ A new perspective to frame GPD processes
 - Single diffractive hard exclusive processes
 - Unified description: kinematics and factorization
 - Two new processes to give **enhanced x -sensitivity**

- ❑ Still not there yet

- Need more processes, more observables
- Need input from lattice QCD
- Need global fit

--- A long challenging but exciting way to go!



Thank you!