

Solutions of the homeworks

HW/1

HW1: (a)
$$P_\mu P^\mu - m^2 c^2 = (\beta^K P_K - mc)(\gamma^\lambda P_\lambda + mc) = 0$$

$$= \beta^K \gamma^\lambda P_K P_\lambda + (\beta^K - \gamma^K) P_K mc - m^2 c^2$$

$$P_\mu P_\nu g^{\mu\nu} - \underbrace{\beta^K \gamma^\lambda P_K P_\lambda}_{\beta^K \gamma^K P_K P_K} - mc(\beta^K - \gamma^K) P_K = 0$$

$$(g^{\mu\nu} - \beta^K \gamma^K) P_\mu P_\nu - mc(\beta^K - \gamma^K) P_K = 0$$

$$= 0 \rightarrow \boxed{\beta^K = \gamma^K, \forall K}$$

(b)
$$(g^{\mu\nu} - \gamma^\mu \gamma^\nu) P_\mu P_\nu = 0$$

symmetric in μ, ν

$$\frac{1}{2} \underbrace{\{\gamma^\mu, \gamma^\nu\}}_{S^{\mu\nu}} + \frac{1}{2} \underbrace{[\gamma^\mu, \gamma^\nu]}_{A^{\mu\nu}}$$

$$S^{\mu\nu} = S^{\nu\mu} \text{ \& \ } A^{\mu\nu} = -A^{\nu\mu}$$

$$(g^{\mu\nu} - S^{\mu\nu} - A^{\mu\nu}) P_\mu P_\nu = 0$$

= 0 since symmetric x Antisymmetric

$$(g^{\mu\nu} - S^{\mu\nu}) P_\mu P_\nu = 0$$

$$= 0 \rightarrow S^{\mu\nu} = g^{\mu\nu} \rightarrow \boxed{\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}}$$

(c)
$$\left. \begin{aligned} \mu = \nu = 0, \quad 2(\gamma^0)^2 &= 2g^{00} \rightarrow (\gamma^0)^2 = 1 \\ \mu = \nu = i, \quad (\gamma^i)^2 &= -1 \\ \mu = 0, \nu = i \quad 2\gamma^0 \gamma^i &= 0 \end{aligned} \right\}$$

choose $\gamma^0 = +1, \gamma^i = i$
 $(\gamma^0)^2 = 1, (\gamma^i)^2 = -1, \gamma^0 \gamma^i = 0$

→ No numbers could satisfy the algebra

(If γ^0 & γ^i are some commuting numbers

$$(d) (\alpha) \quad \mu=0, \nu=i \rightarrow \gamma^0 \gamma^i + \gamma^i \gamma^0 = 2g^{0i} = 0$$

$$\gamma^0 / \gamma^0 \gamma^i = -\gamma^i \gamma^0$$

$$(\gamma^0)^2 \gamma^i = -\gamma^0 \gamma^i \gamma^0$$

$$\stackrel{=1}{\gamma^i} = -\gamma^0 \gamma^i \gamma^0$$

$$\text{Tr}[\gamma^i] = -\text{Tr}[\gamma^0 \gamma^i \gamma^0] \quad (\text{Tr}[ABC] = \text{Tr}[BCA] = \dots)$$

$$= -\text{Tr}[\gamma^i \underbrace{\gamma^0 \gamma^0}_{=1}] = -\text{Tr}[\gamma^i]$$

$$\rightarrow 2\text{Tr}[\gamma^i] = 0 \rightarrow \boxed{\text{Tr}[\gamma^i] = 0}$$

same holds for γ^0 ; $\text{Tr}[\gamma^0] = 0$.

$$(b) \text{ For } \gamma^0: \text{ since } (\gamma^0)^2 = 1, (\gamma^0)^2 - 1 = 0 \rightarrow (\gamma^0 - 1)(\gamma^0 + 1) = 0$$

$$\downarrow$$

$$\gamma^0 = \pm 1$$

$$\text{For } \gamma^i: \text{ since } (\gamma^i)^2 = -1 \rightarrow (\gamma^i)^2 + 1 = 0$$

$$\downarrow$$

$$\gamma^i = \pm i$$

\exists a basis where γ^0 is diagonal with ± 1 eign.

$$(e) \quad R \gamma^0 R^T = \gamma^0_{\text{Diag}} = \begin{pmatrix} +1 & & & \\ & -1 & & \\ & & +1 & \\ & & & \ddots \end{pmatrix}$$

Since γ^0_{Diag} is traceless, \exists $(+1, -1)$ pairs in the diagonal

\downarrow
 $N \equiv \text{Dim of } \gamma^0 \text{ is even.}$
 Same for γ^i .

$$N = 2, 4, 6, \dots$$

$N=2$ case, Pauli matrices σ^i 's & $I_{2 \times 2}$

$$\text{Take } \gamma^0 = I_{2 \times 2}, \gamma^i = i\sigma^i$$

\downarrow
 Not satisfying properties

one must go $N=4$.

(4) If $\gamma_{Dirac}^0 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$, $\gamma_{Dirac}^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix}$

One can show that $(\gamma_{Dirac}^0)^2 = 1$, $(\gamma_{Dirac}^i)^2 = -1$ & $\gamma_{Dirac}^0 \gamma_{Dirac}^i - \gamma_{Dirac}^i \gamma_{Dirac}^0 = 0$
 $\rightarrow \xi_0^{\pm} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

Alternative representation

$\gamma_{Dirac}^0 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\gamma_{Dirac}^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix} \rightarrow \xi_{Dirac}^{\pm} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

HW2: Show that $H_0 = \frac{c}{i} \vec{\alpha} \cdot \vec{p} + mc^2 \beta$ for a free Dirac particle

$(\not{p} - mc)\psi = 0$, $\not{p} = \gamma^{\mu} p_{\mu} = i\hbar \gamma^{\mu} \partial_{\mu}$

$(i\hbar \not{\partial} - mc)\psi = 0$

$\cancel{p} / i\hbar \cancel{p} \psi + i\hbar \cancel{p}^2 \psi - mc\psi = 0$

$c / i\hbar (\cancel{p})^2 \psi + i\hbar \cancel{p}^2 \psi - mc\cancel{p} \psi = 0$

$i\hbar \frac{\cancel{p}}{\cancel{p}} = -i\hbar c \underbrace{\sum_{i=1}^3 \sigma_i}_{\vec{\alpha}} \psi + mc^2 \beta \psi$
 $= (-i\hbar c \vec{\alpha} \cdot \vec{p} + mc^2 \beta) \psi$
 $\equiv H_0$

HW3: (a) $\gamma^{\mu T} = \gamma^0 \gamma^{\mu} \gamma^0$

From HD which is hermitian, both γ^0 & $\gamma^0 \gamma^i$ must be hermitian as well.

$$\gamma^{0T} = \gamma^0$$

$$(\gamma^0 \gamma^i)^{\dagger} = \gamma^0 \gamma^i \rightarrow \gamma^{iT} \gamma^{0T} = \gamma^0 \gamma^i / \gamma^0$$

$$\gamma^{iT} \gamma^0 \gamma^0 = \gamma^0 \gamma^i \gamma^0 \rightarrow \gamma^{iT} = \gamma^0 \gamma^i \gamma^0$$

$$\text{Also, } \gamma^{0T} = \gamma^0 \gamma^0 \gamma^0 = \gamma^0 \gamma^0 \gamma^0 \checkmark$$

(b) $(S^{\alpha\beta})^{\dagger} = -\gamma^0 S^{\alpha\beta} \gamma^0$

$$S^{\alpha\beta} = \frac{1}{4} [\gamma^{\alpha}, \gamma^{\beta}] \rightarrow (S^{\alpha\beta})^{\dagger} = \frac{1}{4} [\gamma^{\alpha}, \gamma^{\beta}]^{\dagger}$$

$$= \frac{1}{4} [\gamma^{\beta T}, \gamma^{\alpha T}]$$

$$= \frac{1}{4} [\gamma^0 \gamma^{\beta} \gamma^0, \gamma^0 \gamma^{\alpha} \gamma^0]$$

$$= \frac{1}{4} (\underbrace{\gamma^0 \gamma^{\beta} \gamma^0}_{\mathbb{1}} \underbrace{\gamma^0 \gamma^{\alpha} \gamma^0}_{\mathbb{1}} - \underbrace{\gamma^0 \gamma^{\alpha} \gamma^0}_{\mathbb{1}} \underbrace{\gamma^0 \gamma^{\beta} \gamma^0}_{\mathbb{1}})$$

$$= \frac{1}{4} (\gamma^0 (\gamma^{\beta} \gamma^{\alpha} - \gamma^{\alpha} \gamma^{\beta}) \gamma^0)$$

$$[\gamma^{\beta}, \gamma^{\alpha}] = -[\gamma^{\alpha}, \gamma^{\beta}]$$

$$= -4 S^{\alpha\beta}$$

$$= -\gamma^0 S^{\alpha\beta} \gamma^0$$

(c) $S[\Lambda]^{\dagger} = \gamma^0 S[\Lambda]^{-1} \gamma^0$

$$S[\Lambda] = e^{\frac{1}{2} \Omega_{\alpha\beta} S^{\alpha\beta}} \rightarrow S[\Lambda]^{\dagger} = e^{\frac{1}{2} \Omega_{\alpha\beta} (S^{\alpha\beta})^{\dagger}}$$

$$= e^{-\frac{1}{2} \Omega_{\alpha\beta} \gamma^0 S^{\alpha\beta} \gamma^0}$$

$$\begin{aligned}
 S[\Lambda]^\dagger &= e^{-\frac{1}{2} \gamma^0 \Omega \cdot S \gamma^0} \quad (e^x = 1 + x + \frac{x^2}{2!} + \dots) \\
 &= \underbrace{1}_{\gamma^0 \gamma^0} - \frac{1}{2} \gamma^0 \Omega \cdot S \gamma^0 + \frac{1}{2!} \frac{1}{4} (\gamma^0 \Omega \cdot S \gamma^0)^2 + \dots \\
 &= \gamma^0 \Omega \cdot S \underbrace{\gamma^0 \gamma^0}_{=1} \Omega \cdot S \gamma^0 \\
 &= \gamma^0 (\Omega \cdot S)^2 \gamma^0 \\
 &= \gamma^0 \left(1 - \frac{1}{2} \Omega \cdot S + \frac{1}{2!} \left(\frac{1}{2} \Omega \cdot S \right)^2 + \dots \right) \gamma^0 \\
 &\quad \underbrace{e^{\frac{1}{2} \Omega \cdot S}}_{= S[\Lambda]^{-1}}
 \end{aligned}$$

$$\underline{S[\Lambda]^\dagger = \gamma^0 S[\Lambda]^{-1} \gamma^0}$$

(d) $\bar{\psi}\psi$ is a Lorentz scalar.

$$\begin{aligned}
 \psi(x) &\xrightarrow{L} S[\Lambda] \psi(\Lambda^{-1}x) \\
 \bar{\psi}(x) &\xrightarrow{L} \overline{S[\Lambda] \psi(\Lambda^{-1}x)} = (S[\Lambda] \psi(\Lambda^{-1}x))^\dagger \gamma^0 \\
 &= \psi^\dagger(\Lambda^{-1}x) \underbrace{S[\Lambda]^\dagger}_{\gamma^0 \gamma^0} \gamma^0 \\
 &= \bar{\psi}(\Lambda^{-1}x) \underbrace{\gamma^0 S[\Lambda]^\dagger \gamma^0}_{S[\Lambda]^{-1} \text{ (from (c))}} \\
 &= \bar{\psi}(\Lambda^{-1}x) S[\Lambda]^{-1}
 \end{aligned}$$

$$\begin{aligned}
 \bar{\psi}(x) \psi(x) &\rightarrow \bar{\psi}(\Lambda^{-1}x) \underbrace{S[\Lambda]^{-1} S[\Lambda]}_{=1} \psi(\Lambda^{-1}x) \\
 &= \bar{\psi}(\Lambda^{-1}x) \psi(\Lambda^{-1}x) \quad \checkmark
 \end{aligned}$$

