

b⁻ quarkonia production at LHCb

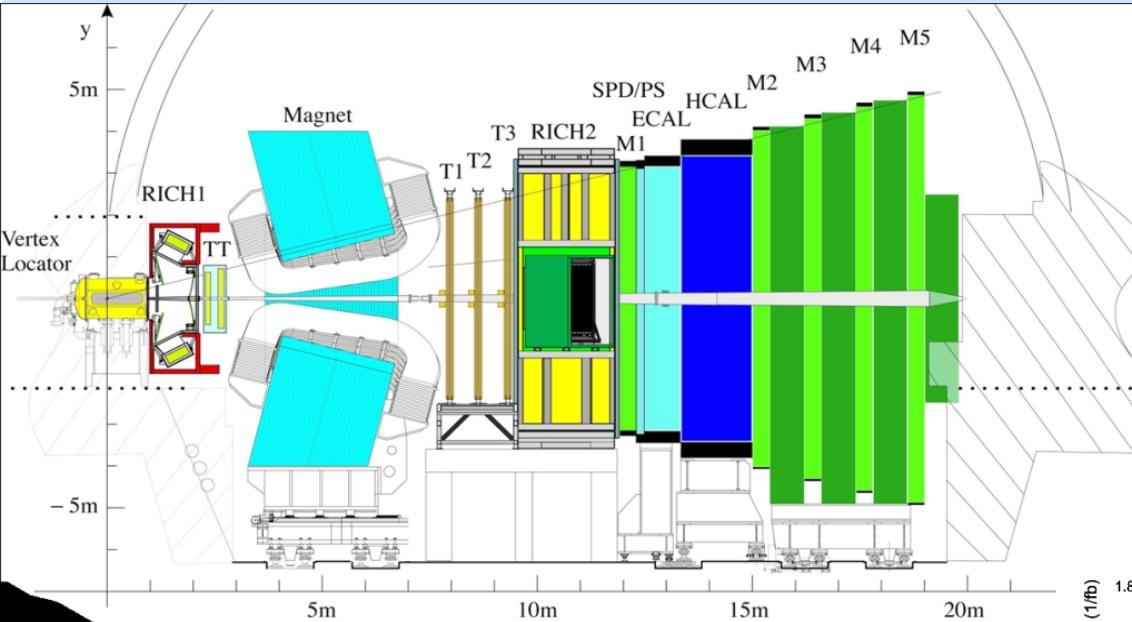
V. Obraztsov, IHEP, Protvino
on behalf of LHCb collaboration

“LHC on the March”, IHEP, Protvino, 16-18 November 2011

The talk layout

- LHCb detector
- $\Upsilon(1S)$, $\Upsilon(2S)$, $\Upsilon(3S)$ cross sections measurement
- $\chi_b(1P)$ observation
- Conclusions

LHC is good for b-physics: the $b\bar{b}$ xsection is large: $\sqrt{s}=7 \text{ TeV}$ $\sigma(b\bar{b}) \sim 300 \mu\text{b}$; B_s , B_c , Λ_b available
 $b\bar{b}$ production peaks at low angles \rightarrow forward spectrometer $2 < \eta < 5$ ($0.9^\circ < \theta < 16^\circ$) $\sim 1/4 \sigma(b\bar{b})$
 Designed to work @ $L=2 \times 10^{32}$ $\rightarrow 2 \text{ fb}^{-1}/\text{year}$ (10^7 sec);
 Features: PV @ IP resolution; particle ID: p, K, π , γ , μ , e



37 pb $^{-1}$ $\sqrt{s} = 7 \text{ TeV}$

in 2010

1.1 fb $^{-1}$ $\sqrt{s} = 7 \text{ TeV}$

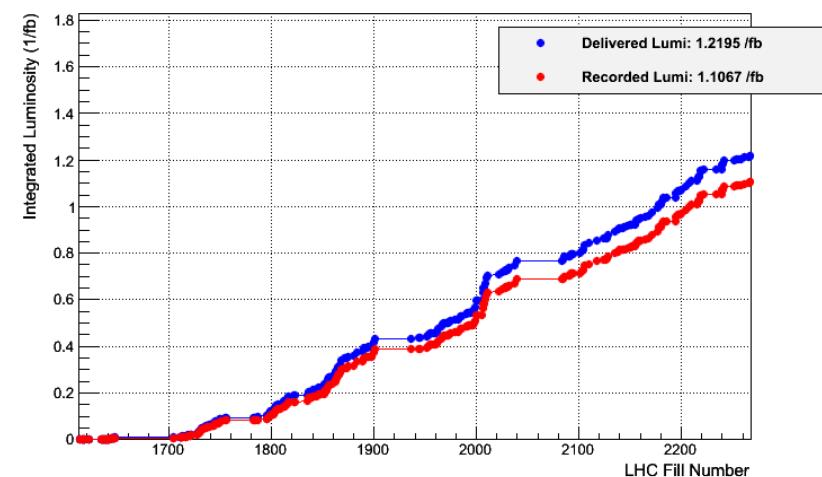
in 2011

Only 2010 data are used for the present analysis

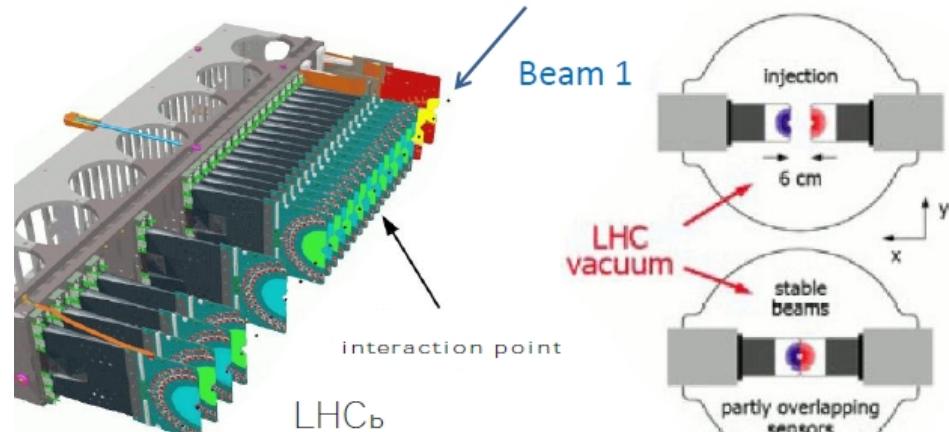
Main parts of LHCb

- ✓ vertex detector
- ✓ Tracking system (+magnet)
- ✓ RICH1,2
- ✓ Calorimeters
- ✓ Muon system
- ✓ Trigger&DAQ

LHCb Integrated Luminosity at 3.5 TeV in 2011



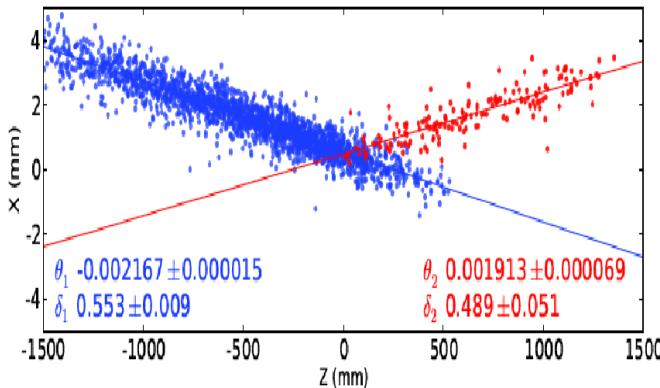
Vertex Locator (VELO)



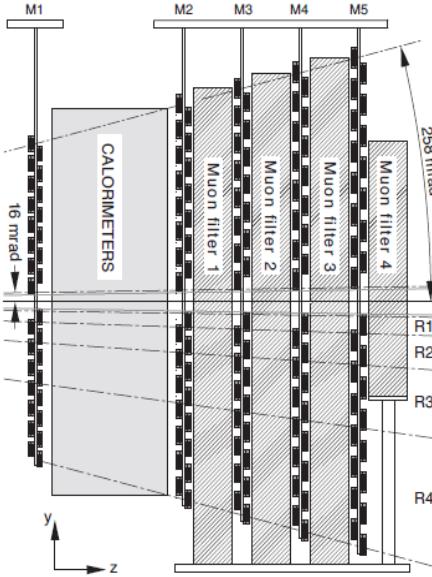
2 retractable halves; 21 stations/half R/ ϕ -strips;
pitch 40-100 μm ; 8 mm from beam (physics);
172K channels

Primary Vertex resolution : $\sigma_{x,y} \sim 13 \mu\text{m}$; $\sigma_z \sim 70 \mu\text{m}$ (> 20 trk)

Impact Parameter resolution : $\sigma_{x,y} \sim 13 \mu\text{m}$ (large P_T)



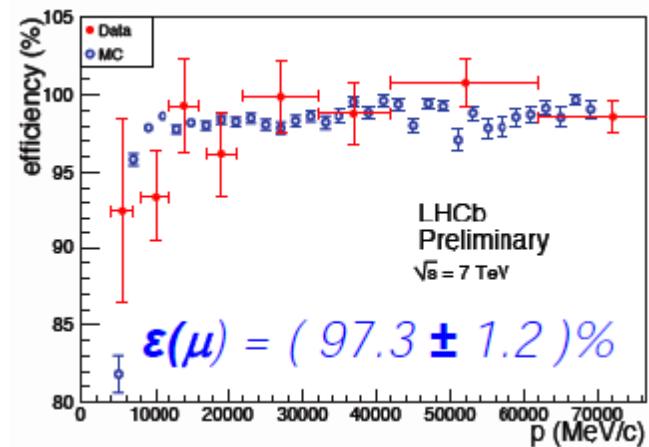
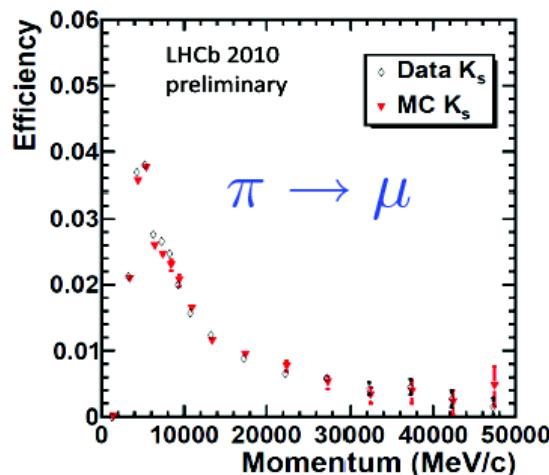
VELO provides new method of luminosity measurement
«beam-gas imaging» method by reconstructing vertexes of
beam-gas interactions.



M1-M5 1380 MWPC's ; 435 m² 3M wires

Ar: CO₂: CF₄ = 40:55:5

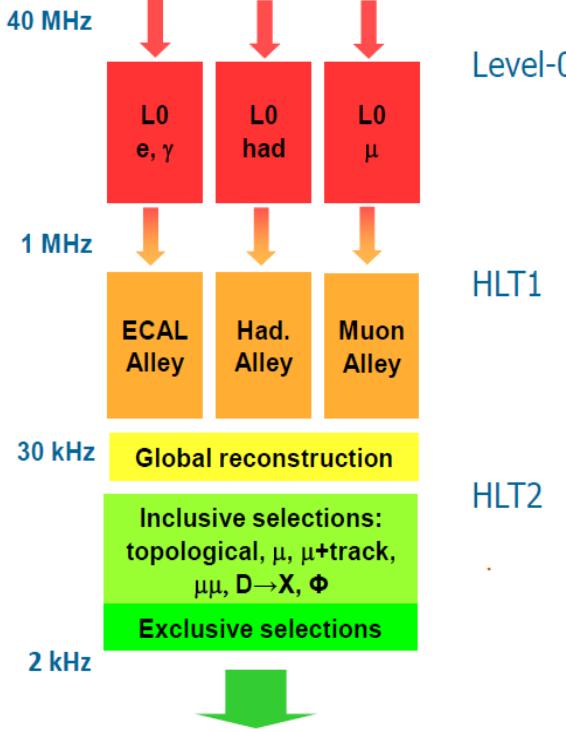
The system provides L0 μ -trigger (with p_T cut) and μ -identification



μ -efficiency and $\pi-\mu$ misidentification determined using $K_s \rightarrow \pi\pi$ and $J/\Psi \rightarrow \mu\mu$ decays

Trigger

LHCb trigger is a challenge:
 $\sigma_{bb}/\sigma_{in} \sim 0.005$



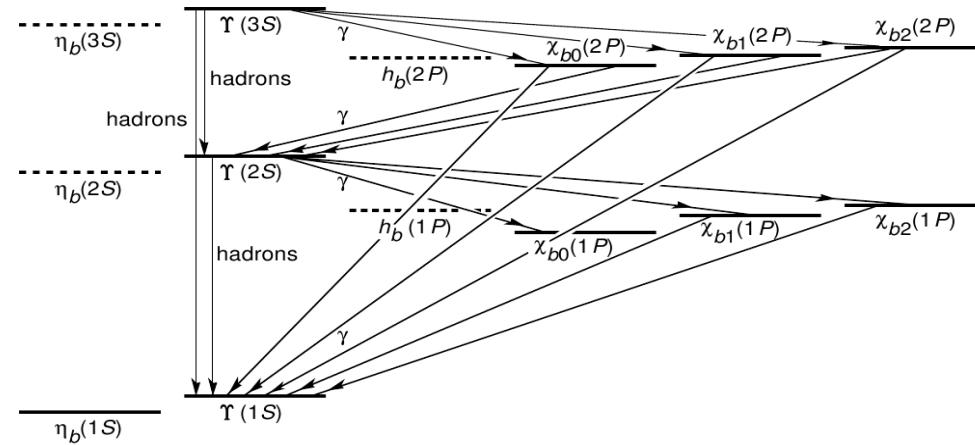
L0 is hardware trigger with fixed latency ($4\mu s$) @ $L=2 \times 10^{32}$:
 $\epsilon_B \sim 90\%$ for $\mu P_T > 1.1 \text{ GeV}$; $\Sigma_{PT} > 1.3 \text{ GeV}$;
 $\sim 50\%$ for $B \rightarrow hh E_T > 4.5 \text{ GeV}$;
 $\sim 70\%$ for e.m. $E_T > 3 \text{ GeV}$ for prompt D's $\epsilon_D \sim 10\%$

HLT runs on event filter farm ~ 2000 nodes
 HLT1 confirms L0 using tracking system, local reconstruction
 HLT2: use full detector information

Trigger lines important for $\Upsilon \rightarrow \mu^+ \mu^-$ events

L0 Trigger	Single Muon	$p_T > 1.4 \text{ GeV}/c$
	Di-Muon	$p_{T1} > 0.56 \text{ GeV}/c, p_{T2} > 0.48 \text{ GeV}/c$
LT1 Trigger	Single Muon	Confirm L0 Single Muon and $p_T > 1.8 \text{ GeV}/c$ (<i>Prescaled</i>)
	Di-Muon	Confirm L0 Di-Muon and $m(\mu^+ \mu^-) > 2.5 \text{ GeV}/c^2$
HLT2 Trigger	Di-Muon	$m(\mu^+ \mu^-) > 2.9 \text{ GeV}/c^2$ or cuts on vertex and track quality

- Quarkonia production has been under study for a long time nevertheless the production mechanism in pp-collision is not fully understood
 - Large cross-section is expected at LHC
 - Several theoretical models are around
 - Color singlet (CSM)
- Under-predicts xsection, no polarization prediction
- Extended to Color octet (COM) mechanisms, (NRQCD)
- better agreement for xsections; predicts transverse polarization, not confirmed by experiments
- NLO CSM better describes xsection and allows longitudinal polarization
 - Other models (Color evaporation (CEM), kt factorization, soft color interaction)
- New data from LHC experiments will help to resolve these issues
 - $b\bar{b}$ system has advantages as compared to $c\bar{c}$ (real NR)



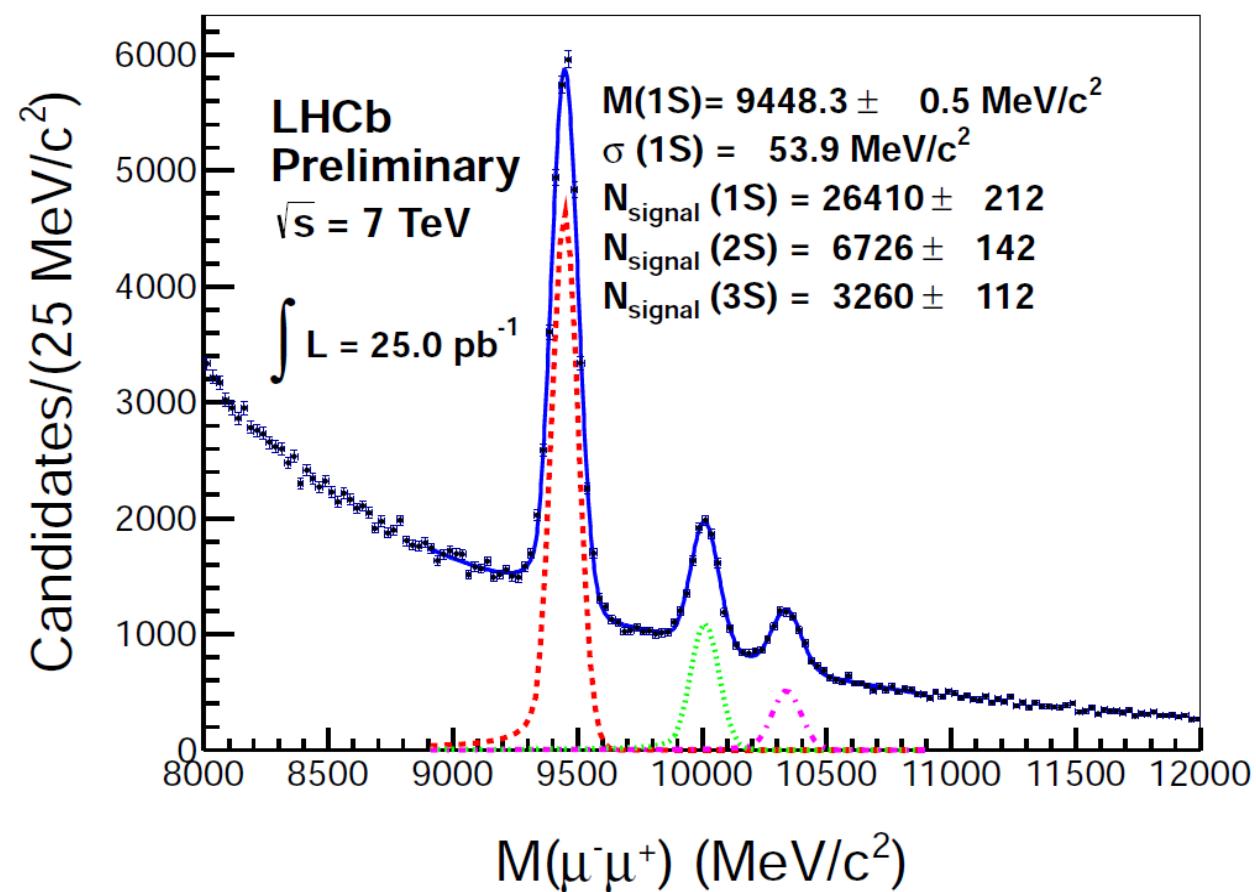
$$JPC = \quad 0-+ \quad 1-- \quad 1+- \quad 0++ \quad 1++ \quad 2++$$

25 pb^{-1} of 2010 data μ -tracks: well reconstructed tracks identified as muons in muon detector, $p_T > 1 \text{ GeV}/c$, Track fit quality, vertex fit quality $\text{Prob}(\chi^2) > 0.5\%$

3 Crystal Ball(CB) +
exp. for background

Parameters α and n
of CB are fixed:
 $\alpha = 2$, $n = 1$

The same function is
used for individual
bin fits



$$\frac{d\sigma(pp \rightarrow Y)}{dp_T dy} = \frac{N^{fit}(p_T, y, \epsilon_{tot})}{Br(Y \rightarrow \mu^+ \mu^-) \cdot \int L dt \cdot \Delta p_T \cdot \Delta y}$$

- N^{fit} – number of candidates in the mass peak in each Δp_T , Δy bin, obtained from the fit and corrected for efficiency
- ϵ_{tot} – total efficiency (*including acceptance*)
- $\Delta p_T, \Delta y$ – bins of p_T and y $0 < p_T < 15 \text{ GeV}/c$; $2 < y < 4.5$
- $\int L dt$ – integrated luminosity
- $\epsilon_{tot} = \epsilon_{Acc} \times \epsilon_{rec} \times \epsilon_{trig}$

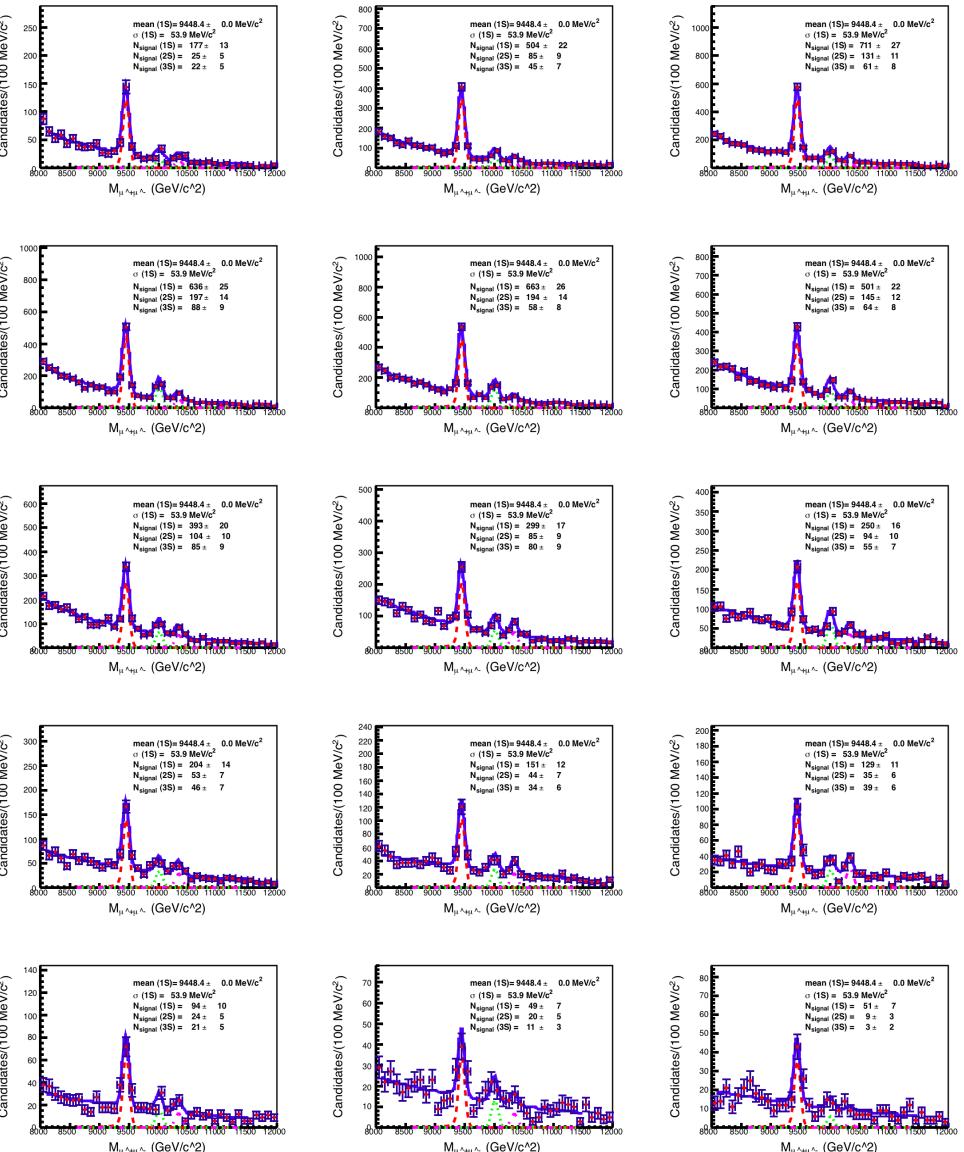
$BR(Y(1S) \rightarrow \mu\mu) = (2.48 \pm 0.05)\%$

$BR(Y(2S) \rightarrow \mu\mu) = (1.93 \pm 0.17)\%$

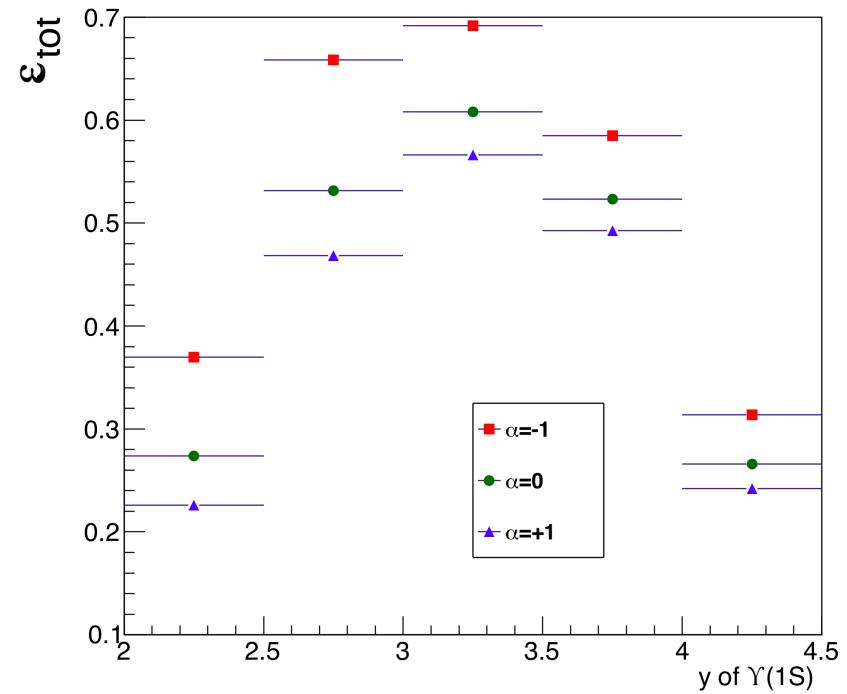
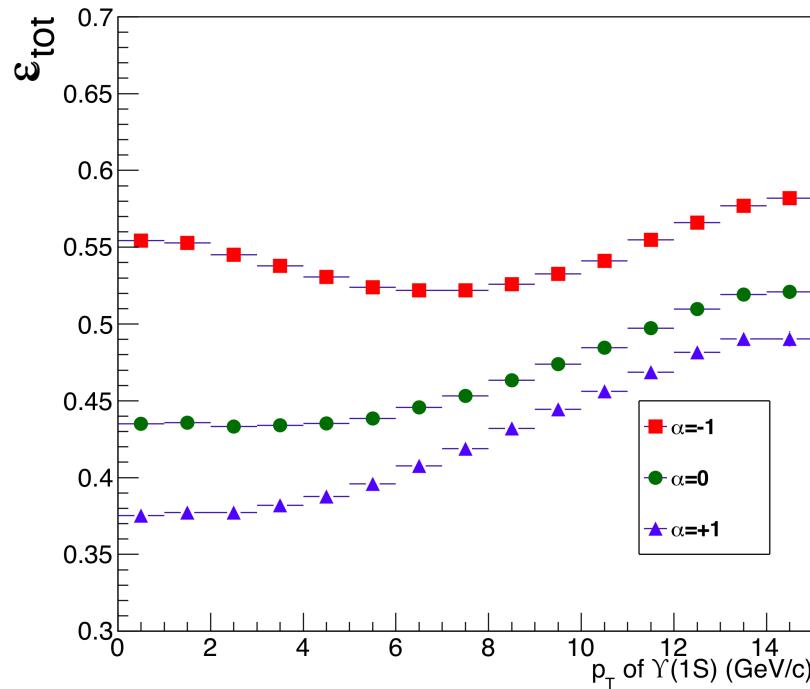
$BR(Y(3S) \rightarrow \mu\mu) = (2.18 \pm 0.21)\%$

N^r in bins of p_T

- Numbers of Υ candidates are extracted from Crystal Ball (CB) part of the fit with 3·CB+exponential.
- Rapidity interval
 $2.0 < y < 4.5$



Total efficiency : $\Upsilon(1S)$

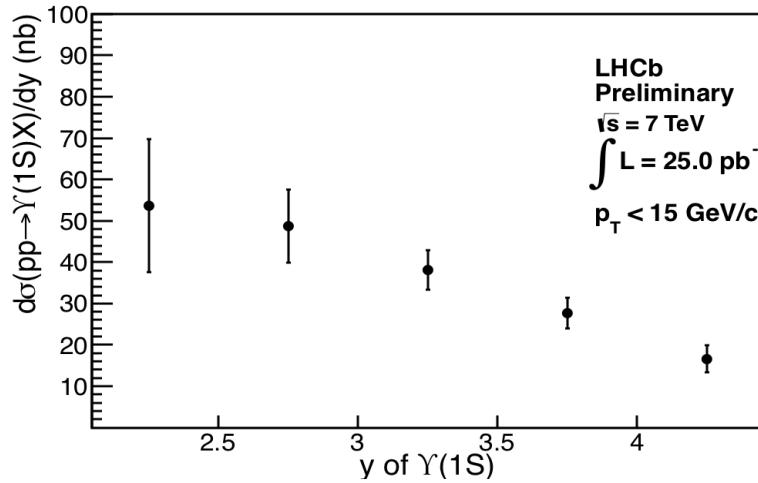
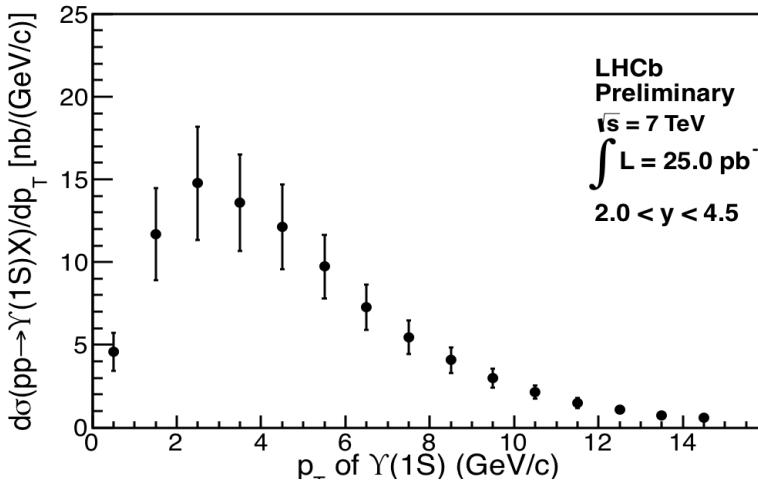
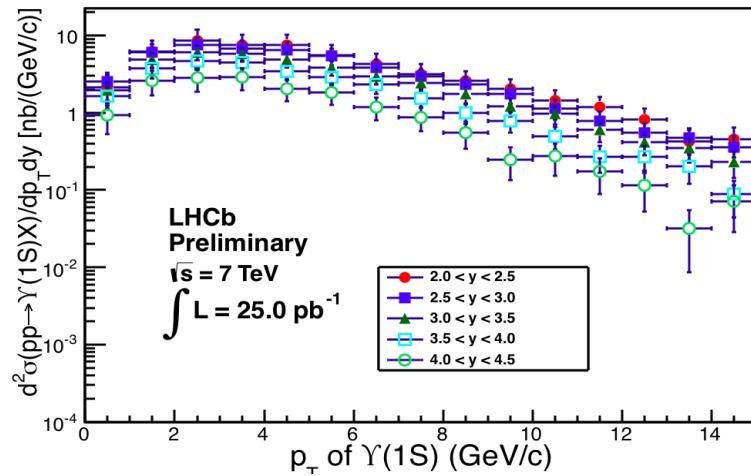
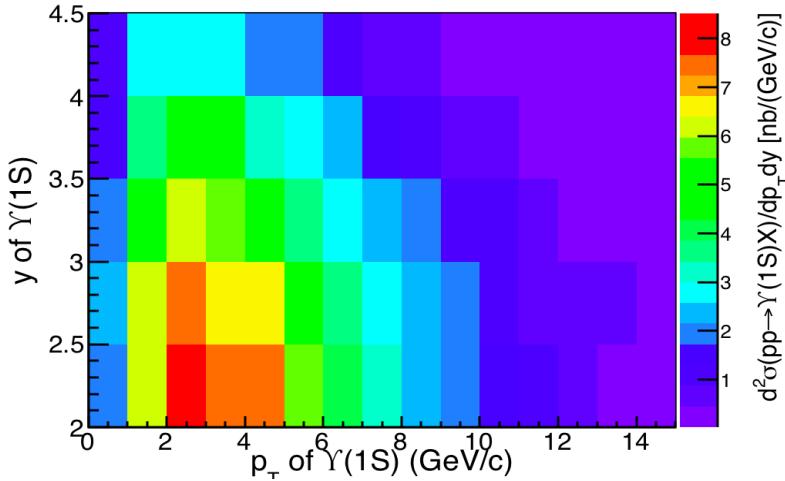


$\alpha = -1$	– full longitudinal polarization
$\alpha = +1$	– full transverse polarization
$\alpha = 0$	– no polarization

Systematic uncertainties

SOURCE	METHOD	VALUE $\Upsilon(1S)$	VALUE $\Upsilon(2S)$	VALUE $\Upsilon(3S)$
statistics MC samples		0-2.5%	0-6.8%	0-5.9%
luminosity	see section 3.2.2		3.5%	
ϵ^{trig} measurement	difference data-MC		3%	
polarisation on A	extreme scenarios	0-33%	0-37%	0-37%
polarisation on ϵ^{rec}	extreme scenarios	0-18%	0-18%	0-17%
polarisation on ϵ^{trig}	extreme scenarios	0-2.5%	0-2.2%	0-2.0%
choice of fit function	different function	1%	1%	1.5%
unknown p_T spectrum	p_T spectrum distribution		1%	
GEC	statistical uncertainty of data		0.6%	
$\epsilon^{trackquality}$	difference data-MC		0.5% per track	
$\epsilon^{track-finding}$	difference data-MC		2.4%	
$\epsilon^{track-finding}$	difference data-MC		0.7% per track	
vertexing	difference data-MC		1%	
ϵ^{muonID}	tag and probe [21]		1.1%	
$BR(\Upsilon \rightarrow \mu^+ \mu^-)$	[7]	2%	8.8%	9.6%
statistics of data		0.02-0.5%	0.06-0.5%	0.08-0.6%

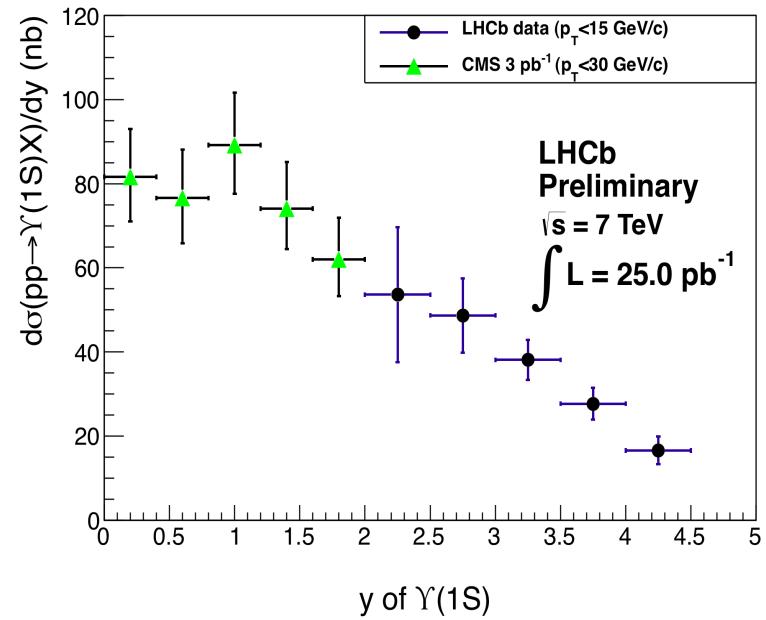
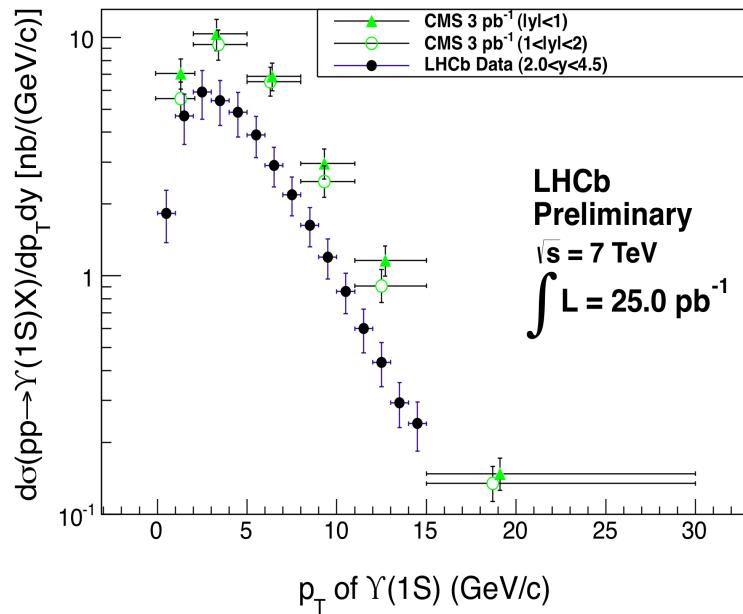
$\Upsilon(1S)$ cross-section



$$92.3 \pm 2.9(\text{stat}) \stackrel{+15.8}{-8.3}(\text{pola}) \pm 4.8(\text{osys}) \pm 3.2(\text{lumi}) \text{ nb} = 92.3 \pm 2.9(\text{stat}) \stackrel{+16.9}{-10.1}(\text{syst}) \text{ nb}$$

LHCb-CMS comparison

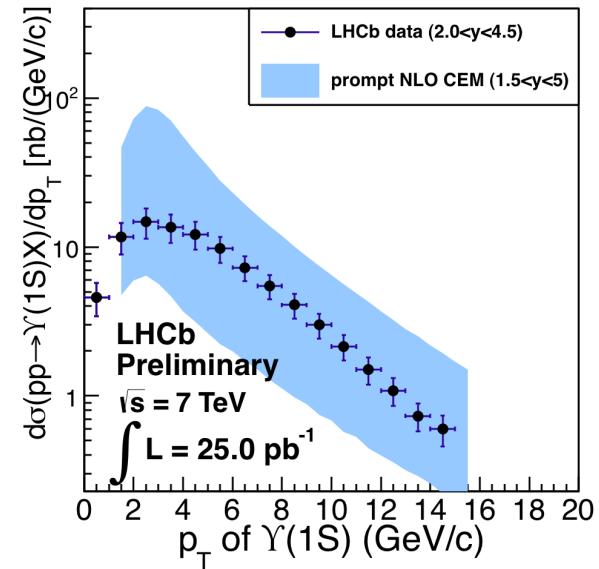
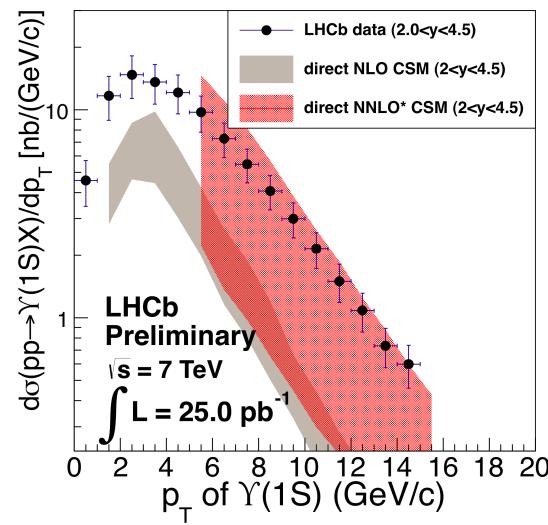
CMS publication: [Phys. Rev. **D83** (2011) 112004; arXiv: hep-ex/1012.5545v1]



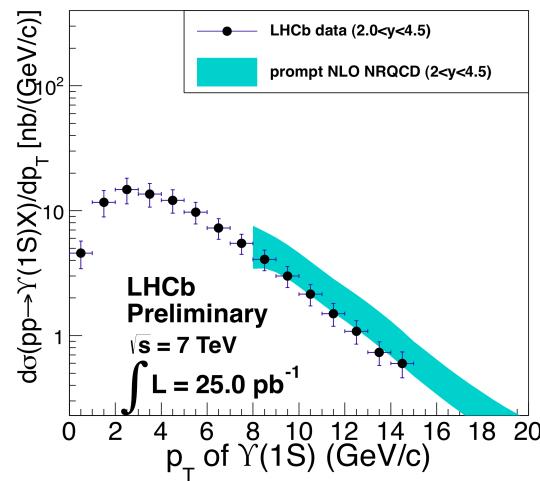
- CSM measurement as a function of p_T is higher, but it should be noted that CMS measures cross-section for $|y| < 2$, while LHCb for $2.0 < y < 4.5$.
- These two measurements supplement each other and are in a good agreement.

Comparison with theory for $\Upsilon(1S)$

NLO & NNLO CSM
 J.-P.Lansberg,Eur.Phys.J.
 C 61 (2009)693

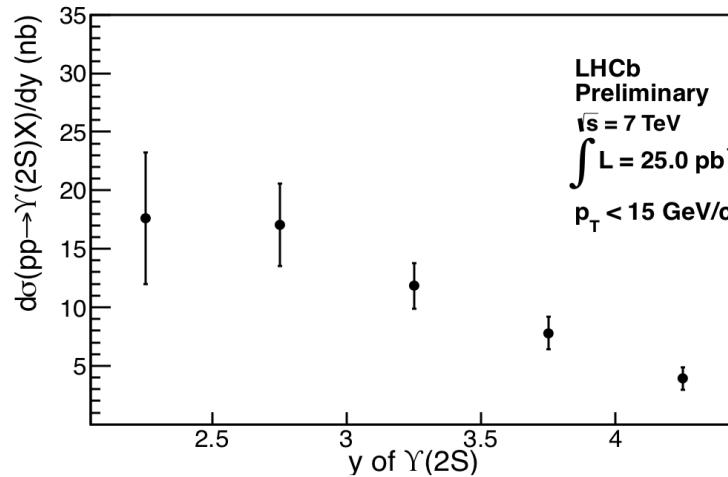
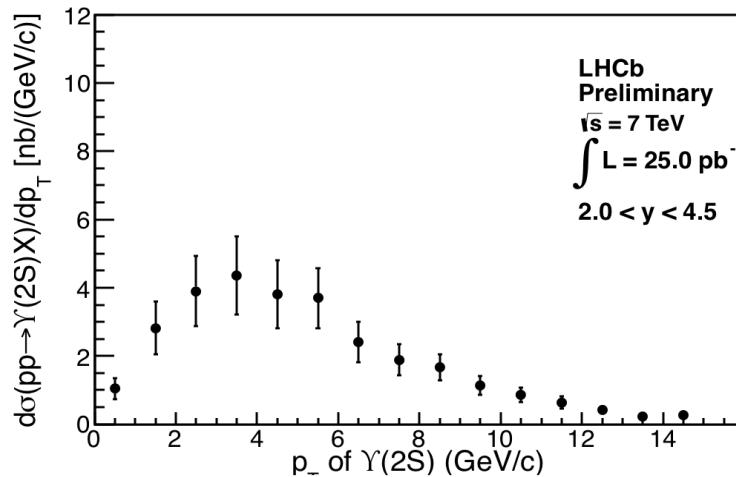
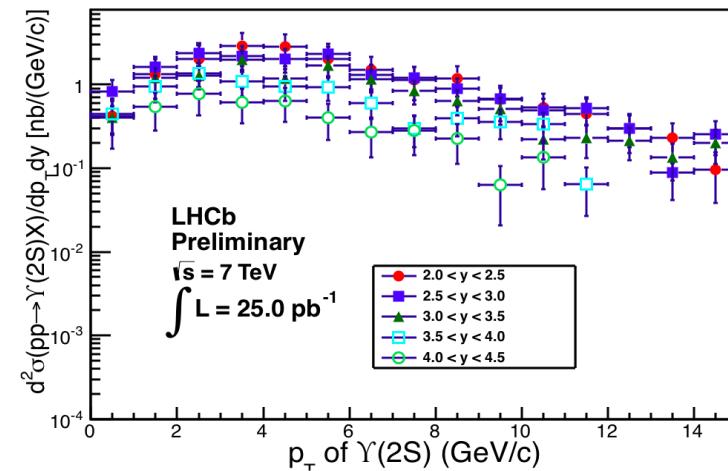
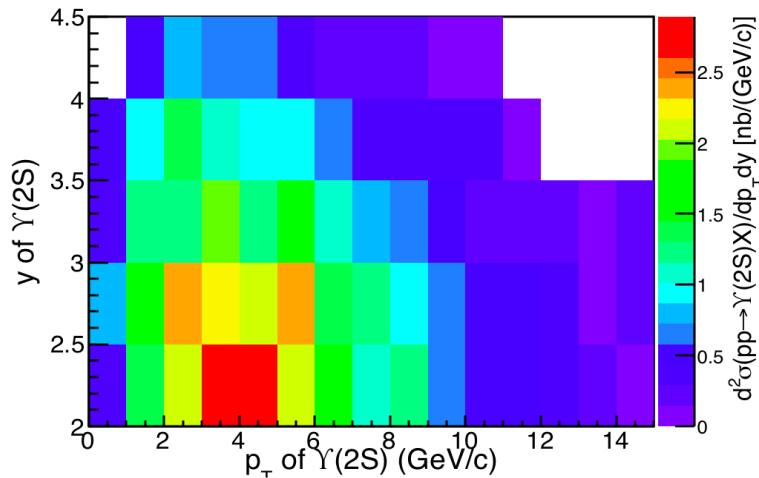


NLO NRQCD
 Y.Q.Ma, et al.,Rev.Lett.
 106 (2011) 042002.



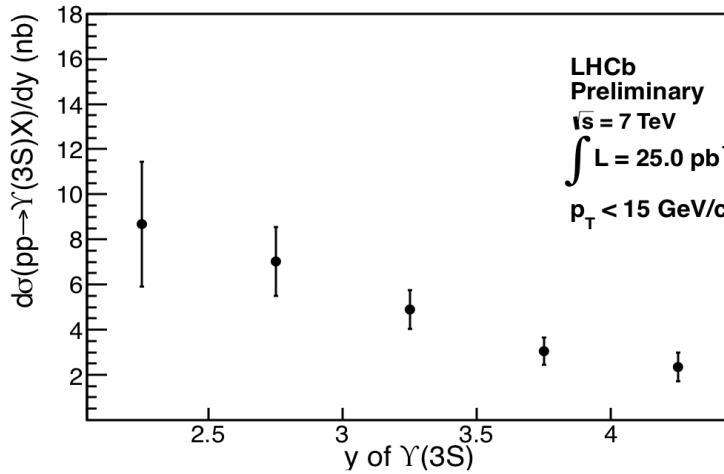
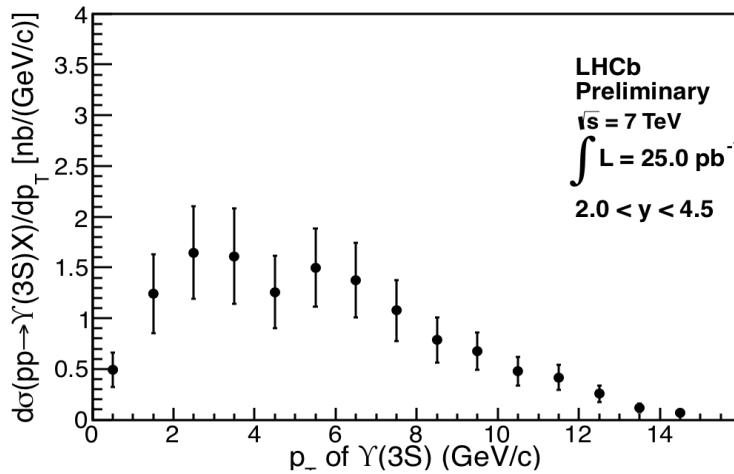
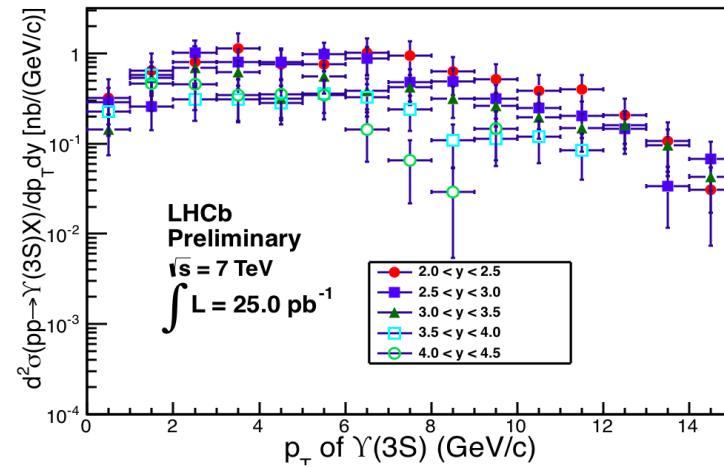
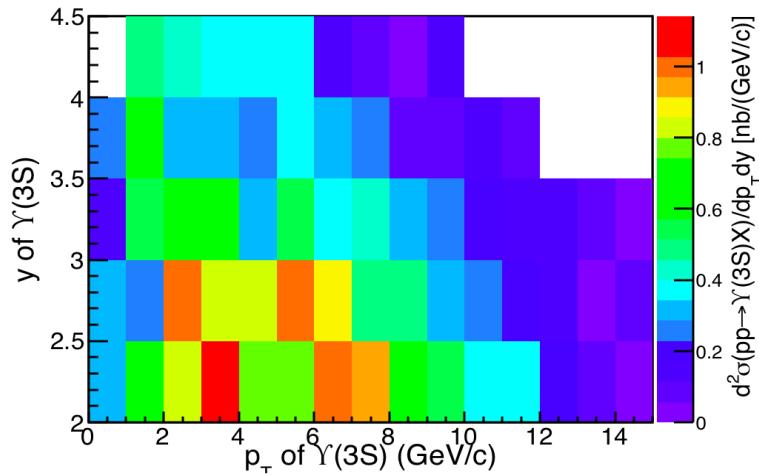
NLO CEM
 A.D.Frawley, et al.,
 Phys. Rep.462 (2008)
 125.

$\Upsilon(2S)$ Cross section New result!



$$29.1 \pm 1.3(\text{stat}) \stackrel{+5.0}{-2.6}(\text{pola}) \pm 2.9(\text{osys}) \pm 1.0(\text{lumi}) \text{ nb} = 29.1 \pm 1.3(\text{stat}) \stackrel{+5.9}{-4.1}(\text{syst}) \text{ nb}$$

$\Upsilon(3S)$ Cross section New !

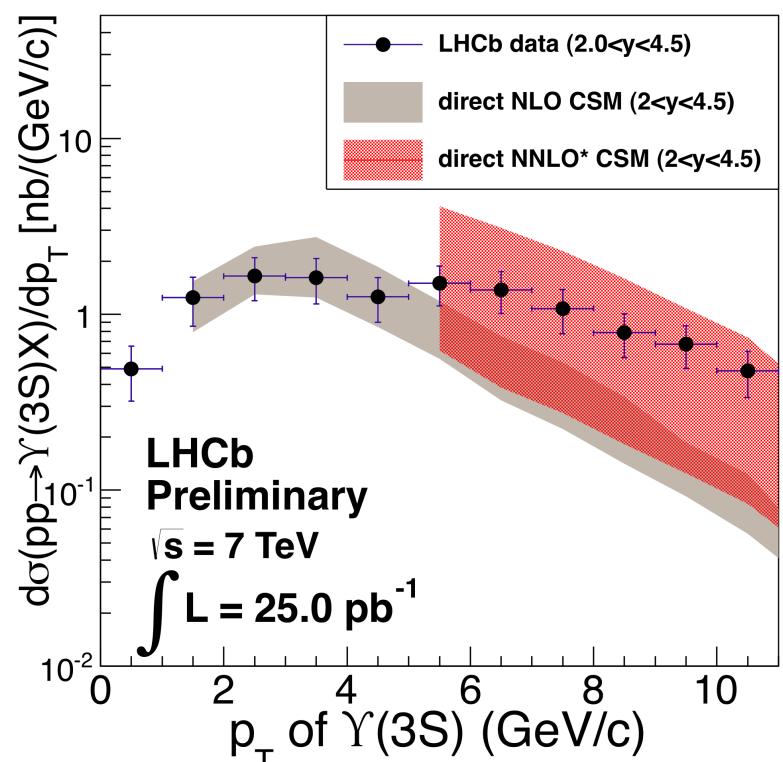
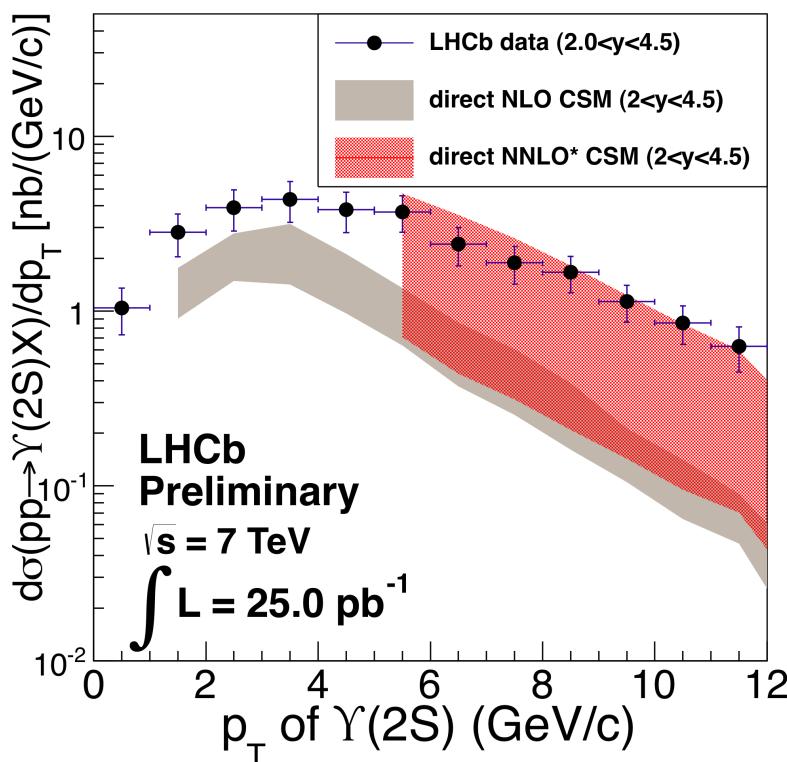


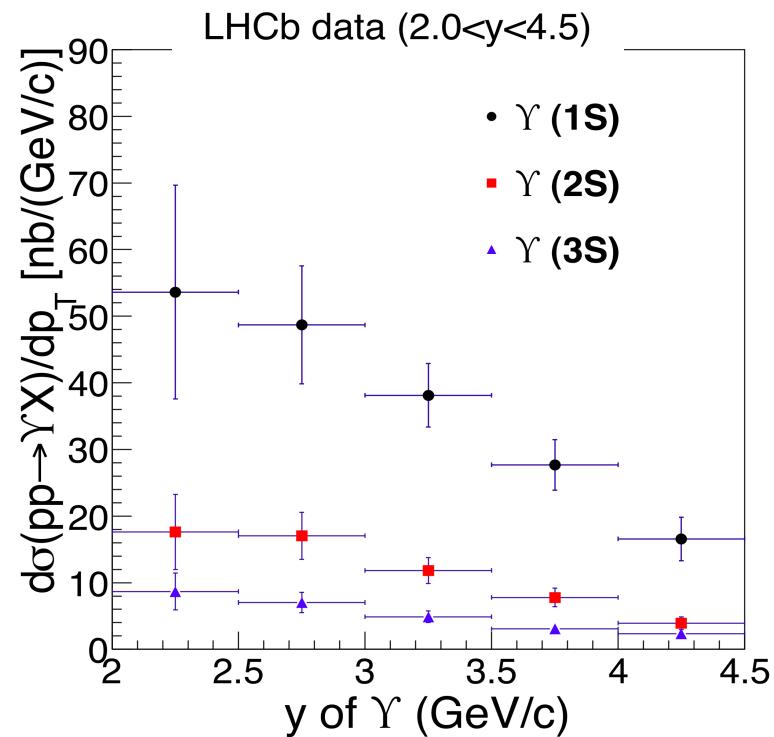
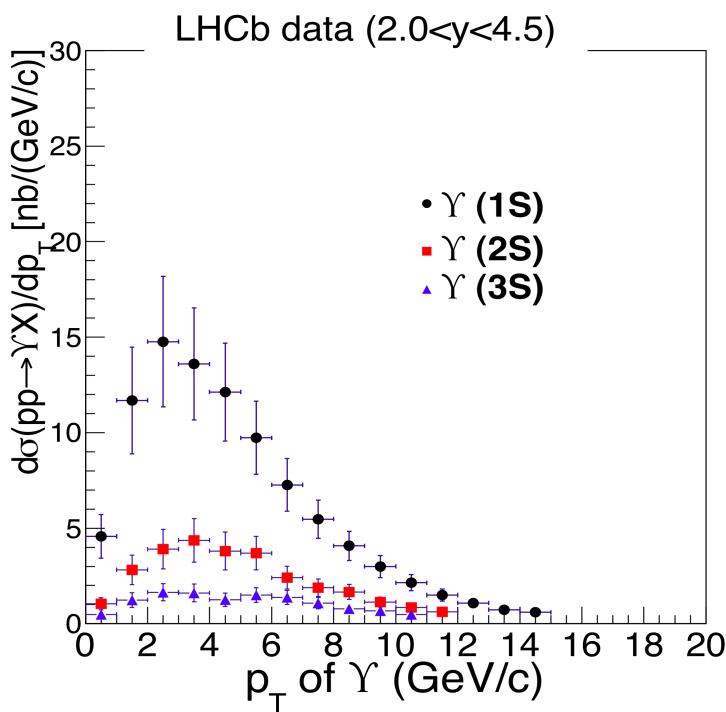
$$13.0 \pm 0.7(\text{stat})^{+2.3}_{-1.2}(\text{pola}) \pm 1.4(\text{osys}) \pm 0.5(\text{lumi}) \text{ nb} = 13.0 \pm 0.7(\text{stat})^{+2.8}_{-1.9}(\text{syst}) \text{ nb}$$

$\Upsilon(2S), \Upsilon(3S)$ comparison with theory

NLO & NNLO CSM

J.-P.Lansberg, Eur.Phys.J. C 61 (2009)693





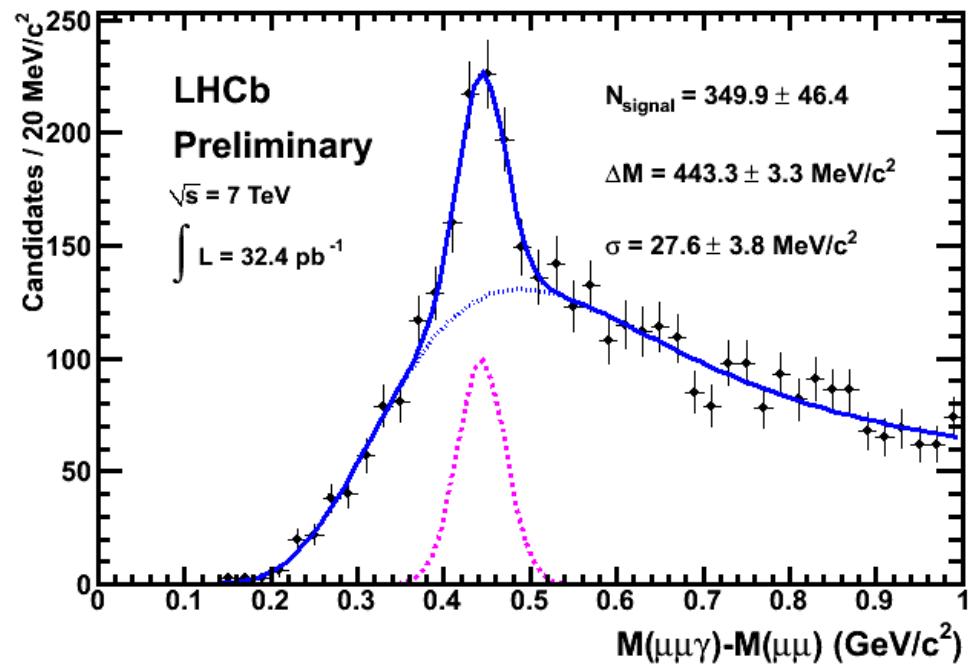
χ_b selection criteria:

$\Upsilon(1S)$ mass band	$9.36 \div 9.56 \text{ GeV}/c^2$
Polar angle of μ^+ in Υ rest frame	$ \cos\theta_{\mu}^* < 0.7$
p_T of photon	$p_T(\gamma) > 700 \text{ MeV}/c$
Polar angle of γ in χ_b rest frame	$\cos\theta_{\gamma}^* > 0$
p_T of χ_b candidate	$p_T(\chi_b) > 7 \text{ GeV}/c$

Gaussian (for signal) + smooth background function is used for fitting to obtain the number of signal events.

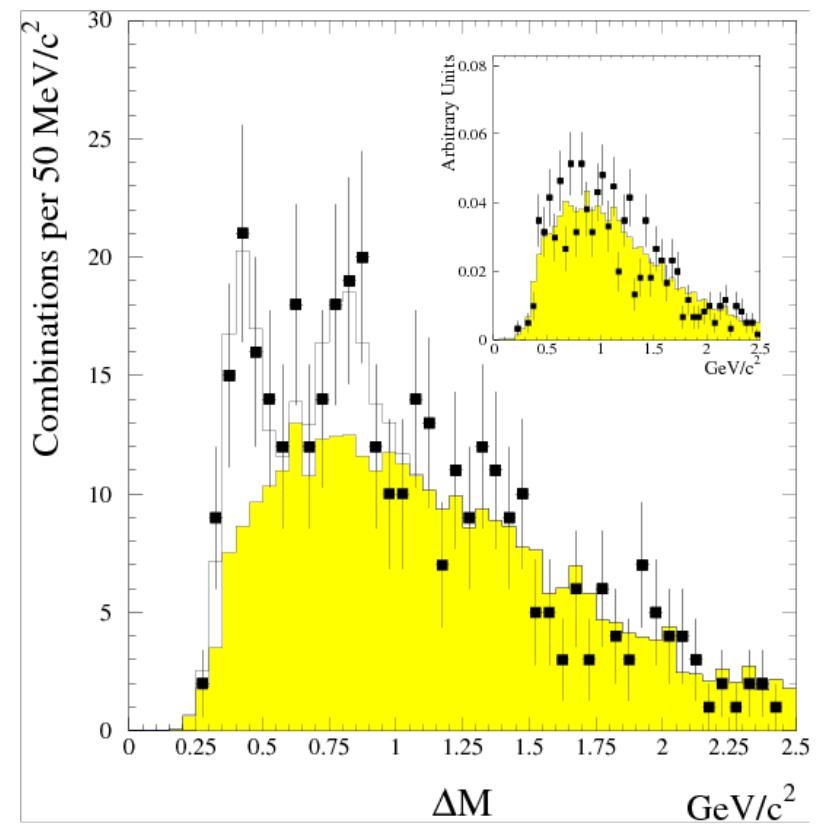
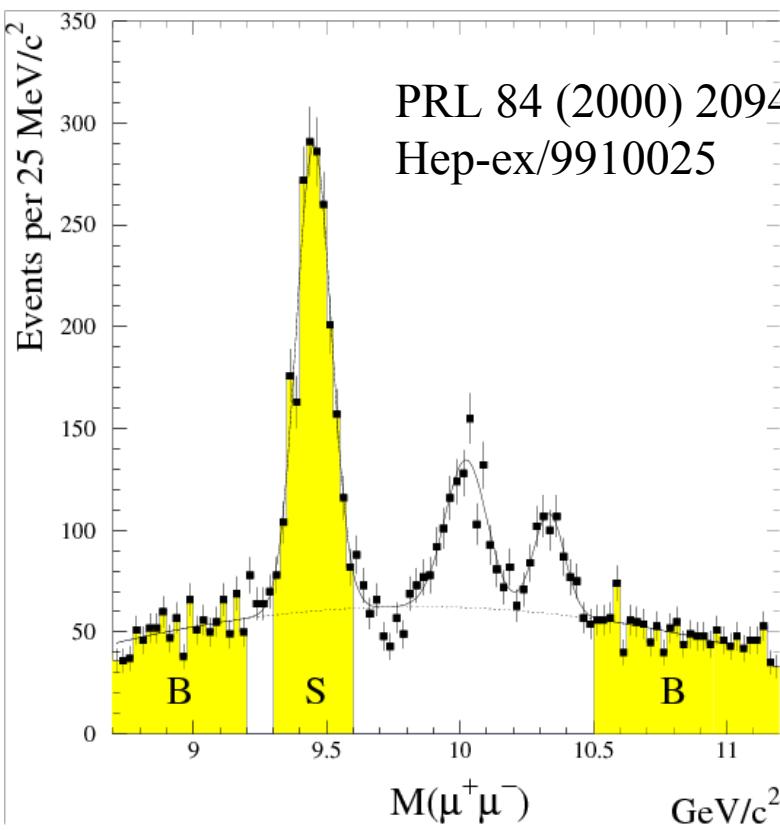
- clear $\chi_b(1P)$ signal is seen.
- No hint for $\chi_b(2P)$ signal ($\Delta M \approx 800 \text{ MeV}/c^2$).
- Cannot resolve states $\chi_{b0}(1P)$, $\chi_{b1}(1P)$, $\chi_{b2}(1P)$.

$$\frac{dN}{dx} = A1 \cdot \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(x-\Delta M)^2}{2\sigma^2}} + A2 \cdot (X - X_0)^\alpha \cdot e^{-(a1 \cdot x + a2 \cdot x^2 + a3 \cdot x^3)}$$



CDF results on $\chi_b(1P)$

- CDF Collaboration with 30 times less statistics of $\Upsilon(1S)$ measure the fractions of $\Upsilon(1S)$ from $\chi_b(1P)$ and $\chi_b(2P)$ for $p_T(\Upsilon) > 8 \text{ GeV}/c$:
 - $(27.1 \pm 6.9(\text{stat.}) \pm 4.4(\text{sys}))\%$ from $\chi_b(1P)$
 - $(10.5 \pm 4.4(\text{stat.}) \pm 1.4(\text{sys}))\%$ from $\chi_b(2P)$



Conclusions

- $\Upsilon(1S)$, $\Upsilon(2S)$, $\Upsilon(3S)$ xsections measured in bins of y and p_T for $2 < y < 4.5$; $0 < p_T < 15$ GeV
- Results agree with other experiments and reasonably agree with theory
- Clear $\chi_b(1P)$ signal has been observed, the measurement of the fraction of $\Upsilon(1S)$ from $\chi_b(1P) \rightarrow \Upsilon(1S)\gamma$ is in progress
- With 1 fb^{-1} of 2011 detailed studies are possible: polarization, χ_{b0} , χ_{b1} , χ_{b2} separation etc.