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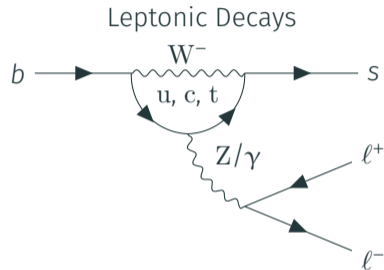
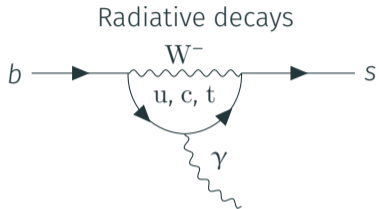
Amplitude and angular measurements in B decays

Implications of LHCb measurements and future prospects

Lakshan Madhan
on behalf of the LHCb collaboration
24-10-2024



FCNC decays

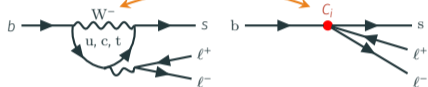


Flavour changing neutral currents (FCNC)
Proceed via loops in SM
Small SM contribution makes them sensitive probes for NP

Studying $b \rightarrow s \ell^+ \ell^-$ transitions

Effective Field Theory approach

Integrate Out



$$\mathcal{H}_{SM} \longrightarrow \mathcal{H}_{eff} = -\frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i (C_i^{SM} + C_i^{NP}) \mathcal{O}_i + \text{chiral flipped}$$

Wilson Coefficients: C_i

- Perturbative, short distance physics
- Describes heavy SM+NP effects

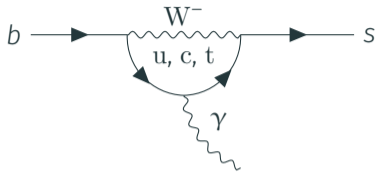
Operators: \mathcal{O}_i

- Non-perturbative, long distance physics
- Strong interactions, difficult to calculate

Operator \mathcal{O}_i	$B_{S(d)} \rightarrow V_{S(d)} \mu^+ \mu^-$	$B_{S(d)} \rightarrow \mu^+ \mu^-$	$B_{S(d)} \rightarrow V_{S(d)} \gamma$
\mathcal{O}_7 EM	✓		✓
\mathcal{O}_9 Vector dilepton	✓		
\mathcal{O}_{10} Axial-vector dilepton	✓	✓	
$\mathcal{O}_{S,P}$ (Pseudo-)Scalar dilepton	(✓)	✓	

Radiative Decays

Radiative decays



Parameters and Observables

- $\mathcal{B}(b \rightarrow s\gamma) \propto |C_7|^2 + |C_7'|^2$
- CP asymmetry sensitive to $2\text{Im}C_7^{\text{eff}} \Delta C_7$
- Photon polarization sensitive to $\frac{C_7'}{C_7}$

Recent LHCb Measurements

Amplitude analysis of $\Lambda_b^0 \rightarrow pK\gamma$
[JHEP 06 (2024) 098]

- Studies the composition of the pK spectrum.
- Feeds into interpretations of $\Lambda_b^0 \rightarrow pK\ell\ell$ decays
- Input to QCD studies

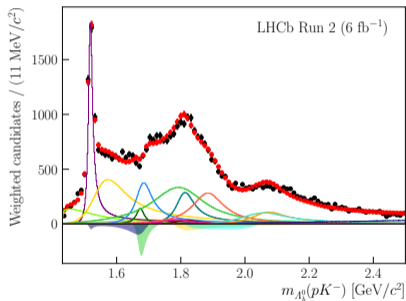
Amplitude analysis of $B_s^0 \rightarrow K^+K^-\gamma$
[JHEP 08 (2024) 093]

- Studies the composition of K^+K^- spectrum
- Measurement of the total tensor contribution $\mathcal{F}_{\{f_2\}} = 16.8 \pm 0.5(\text{stat}) \pm 0.7(\text{syst})\%$
- First observation of $B_s^0 \rightarrow f_2(1525)\gamma$

Recent LHCb Measurements

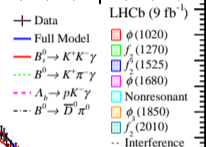
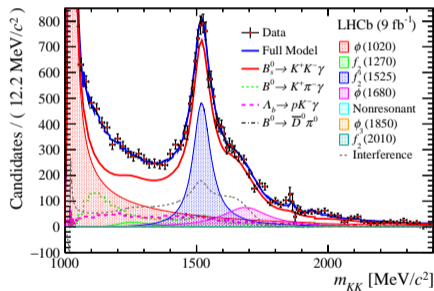
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[JHEP 08 (2024) 093]

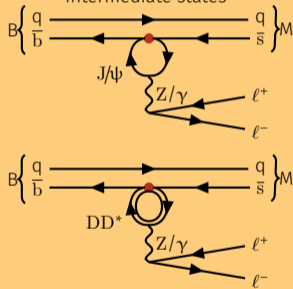


Leptonic Decays

Non-local contributions

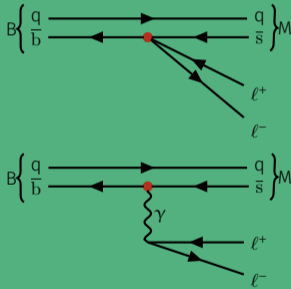
$$(\mathcal{O}_{1q,2q})$$

Intermediate states

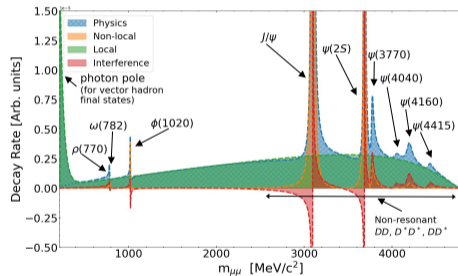


Local contributions

$$(\mathcal{O}_{9,10,7})$$



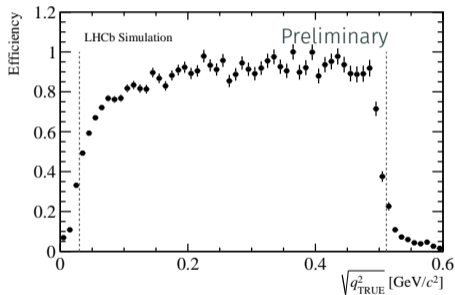
The invariant $m_{\mu\mu}$ spectrum of
 $B \rightarrow M\mu^+\mu^-$



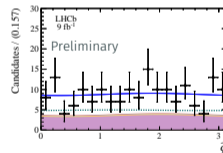
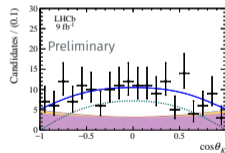
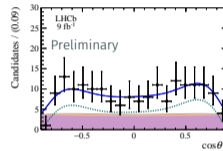
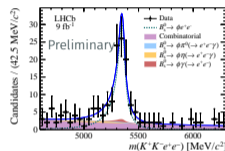
$$q_{\mu\mu}^{22}$$

- Fit to $B_s^0 \rightarrow \phi e^+ e^-$ in effective $m_{ee} \in [30, 511.37] \text{ MeV}/c^2$ using Run1+Run2 (9 fb^{-1}) LHCb data
- Result quoted in effective range due to sharp drops in efficiency at the edges and to allow for theoretical interpretations without simulation input.
 - Follows similar procedure used in the $B^0 \rightarrow K^{*0} e^+ e^-$ analysis [JHEP 04 (2015) 064]
- Using $B_s^0 \rightarrow \phi \gamma$ as control with γ reconstructed electronically and $m_{ee} < 10 \text{ MeV}/c^2$

- $A_T^{(2)}$ and $A_T^{\mathcal{I}m\mathcal{CP}}$ sensitive to photon polarisation
- $A_T^{\mathcal{R}e\mathcal{CP}}$ related to forward-backward asymmetry
- F_L is the longitudinal polarisation



- A cut of $|\cos\theta_l| < 0.9$ imposed to remove $B_s^0 \rightarrow D_s^\mp (\rightarrow \phi e^\mp \nu_e) e^\pm \nu_e$
- Shape of peaking background component of $B_s^0 \rightarrow \phi \gamma$ is extracted from control ($m_{KKe\ell}, \cos\theta_K$) +simulation samples ($\cos\theta_l, \tilde{\varphi}$).
- Peaking Backgrounds $B_s^0 \rightarrow \phi \eta$ and $B_s^0 \rightarrow \phi \pi^0$ modelled from simulation
- Fit with ≈ 100 signal candidates



$$A_T^{(2)} = -0.045 \pm 0.235(\text{stat.}) \pm 0.014(\text{syst.}),$$

$$A_T^{\mathcal{I}m\mathcal{P}} = 0.002 \pm 0.247 \pm 0.016,$$

$$A_T^{\mathcal{R}e\mathcal{P}} = 0.116 \pm 0.155 \pm 0.006,$$

$$F_L < 11.5\% \text{ at } 90\% \text{ CL.}$$

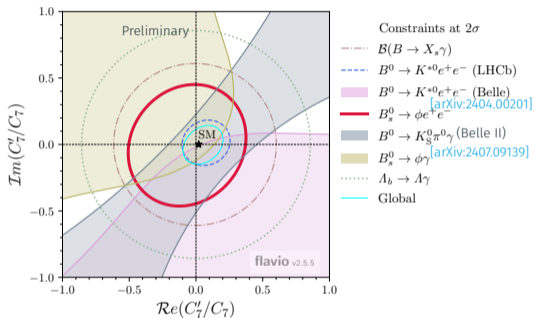
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$$A_T^{\text{ReCP}} = 0.116 \pm 0.155 \pm 0.006,$$

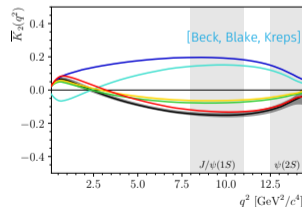
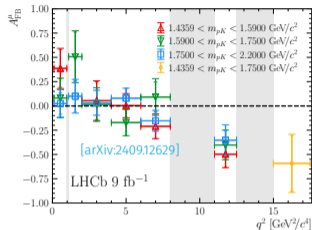
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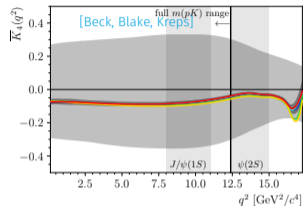
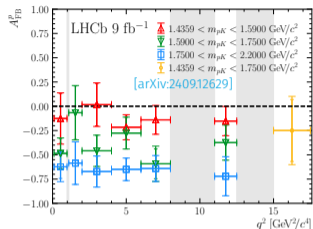
- Using Run1+Run2 LHCb data of 9 fb^{-1}
- Measured in bins of q^2 and m_{pK}^2 and normalised with $\Lambda_b^0 \rightarrow pK^- J/\psi$
- Yields and branching fraction measured from $m_{pK^- \mu^+ \mu^-}$
- Decay rate described by 46 angular moments

$$\frac{d\Gamma^5}{d\Phi} = \frac{3}{8\pi} \sum_{i=0}^{46} K_i(q^2, m_{pK}^2) f(\cos\theta_\mu, \cos\theta_p, \phi)$$

- Forward-backward asymmetry of dimuons sensitive to interference between vector and axial-vector parts. $A_{FB}^\mu = \frac{2}{3} \bar{K}_2$
- Similar pattern to what is seen in $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ decays but consistent with SM (sign change w.r.t mesons due to convention)

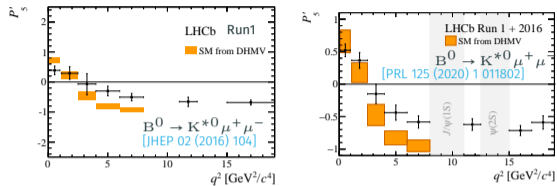


- Standard Model
- $C_9 = -C_9^{\text{SM}}$
- $C_{10} = -C_{10}^{\text{SM}}$
- $C_9^i = C_9^{\text{SM}}$
- $C_{10}^i = C_{10}^{\text{SM}}$
- global fit



$$A_{FB}^{\mu} = \frac{3}{2} \bar{K}_4 - \frac{\sqrt{21}}{8} \bar{K}_{10} + \frac{\sqrt{33}}{16} \bar{K}_{16}$$

- Hadronic spectrum qualitatively similar to $\Lambda_b^0 \rightarrow pK^- \gamma$ and $\Lambda_b^0 \rightarrow pK^- J/\psi$
- Difficult to directly interpret due to unknown interference patterns between the states.
- A more complete understanding of the different contributing states needed.



$$\frac{d^4\Gamma(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{d\vec{\Omega} dq^2} = \sum_i J_i(q^2) f_i(\vec{\Omega})$$

Angular coefficients angular functions

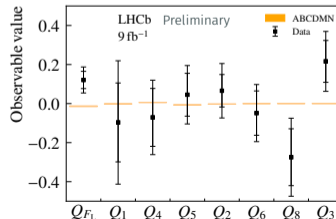
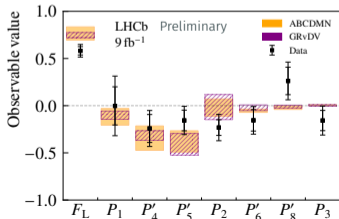
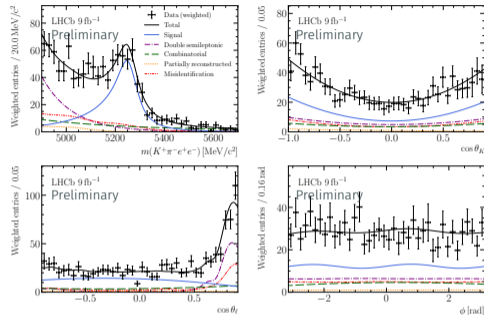
- $J_i(q^2) \propto (\mathcal{A}_\lambda \mathcal{A}_\lambda^*)$, C_i and FF dependent amplitudes.
- Coefficients $P_{1,\dots,8}$ constructed to cancel form factors to leading order
- Additional nuisance terms introduced for non-resonant S-wave configuration.

- Global fit to CP averaged observables show deviation with SM change from 3.0σ to 3.3σ between 3 fb^{-1} and 4.6 fb^{-1}
- Recent preliminary CMS measurement compatible with LHCb [CMS-PAS-BPH-21-002]
- Tensions also seen in $B^+ \rightarrow K^{*+} \mu^+ \mu^-$ and $B_s^0 \rightarrow \phi \mu^+ \mu^-$ decays
- Large dependence on treatment of SM hadronic effects

Could data help us understand the hadronic effects better?

Angular Analysis of $B^0 \rightarrow K^{*0} e^+ e^-$ in central q^2

- Results from fit to $B^0 \rightarrow K^{*0} e^+ e^-$ in central $q^2 \in [1.1, 6.0]$ GeV^2/c^4 using Run1+Run2 (9 fb^{-1}) LHCb data.
- Good agreement seen with SM but similar pattern seen with P'_5 as in the muonic mode.
- Construct LFU observables $Q_i = P_i^\mu - P_i^e$ and refit to data from [PRL 125 (2020) 011802] (consistent S-Wave treatment).
- Results compatible with LFU hypothesis. [\[more on other LFU tests by Sara Celani\]](#)



A closer look using $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ decays

Work in Progress

Method 1 Binned

- Measure angular observables in q^2 bins
- Update using Run1+Run2 LHCb data

Low model dep.

Work in Progress

Method 2

Amplitude ansatz

- Fit in q^2 regions using a Ansatz $\mathcal{A} = \sum_i \alpha_i L_i(q^2)$
- Extract amplitude components with correlated uncertainties.
- Allows one to recompute any observable and generate synthetic data to study any choice of model.

Low model dep.

Discussed Here

Method 3 z Expansion

- Fit to q^2 regions using a polynomial to describe the non-local effects.
- Use theoretical and experimental inputs to constrain them.

Med. model dep.

Method 4 Dispersion

- Unbinned fit to the full q^2 spectrum.
- Use dispersion relations to explicitly model non-local states

High model dep.

$$\mathcal{A}_\lambda^{L,R} = \mathcal{N}_\lambda \left\{ \left[(C_9 \pm C'_9) \mp (C_{10} \pm C'_{10}) \right] \mathcal{F}_\lambda(q^2) + \frac{2m_b M_B}{q^2} \left[(C_7 \pm C'_7) \mathcal{F}_\lambda^T(q^2) - 16\pi^2 \frac{M_B}{m_b} \mathcal{H}_\lambda(q^2) \right] \right\}$$

Wilson Coefficients, C_i

- Real part of C_9 , C'_9 , C_{10} and C'_{10} treated as free fit parameters. Imag. fixed to zero

Local Form Factors, $\mathcal{F}_\lambda(q^2)$

- Constrained to LCSR + Lattice inputs
[Gubernari, Kokulu & van Dyk] + [Horgan, Liu, Meinel & Wingate]

Non-local hadronic matrix elements (charm-loop), $\mathcal{H}_\lambda(q^2)$

- Exploit analytic properties of hadronic matrix elements,
- Following [Gubernari, Reboud, van Dyk & Virto],

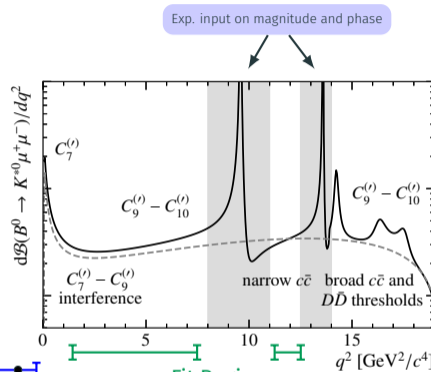
$$q^2 \rightarrow z(q^2) \equiv \frac{\sqrt{t_+ - q^2} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - q^2} + \sqrt{t_+ - t_0}}$$

Polarisation
 $\lambda \in (\parallel, \perp, 0)$

$$\mathcal{H}_\lambda(z) = \frac{1 - zZ_{J/\psi}}{z - Z_{J/\psi}} \frac{1 - zZ_{\Psi(2S)}}{z - Z_{\Psi(2S)}} \phi_\lambda^{-1}(z) \sum_k a_{\lambda,k} z^k,$$

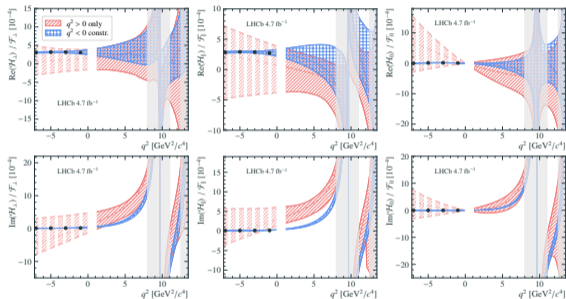
- Experimental inputs for magnitudes and phases of resonances [PRD 90 (2014) 112009], [PRD 76 (2007) 031102], [PRD 88 (2013) 074026], [PRD 88 (2016) 052002], [EPJC 72 (2012) 2118]

- Using 4.7 fb^{-1} of LHCb data (Run1+2016)
- Unbinned fit for FFs, WCs, and non-local polynomial coefficients $a_{\lambda,k}$
- 6D fit to 3 angles, q^2 , $m^2(K\pi)$ and $m(K\pi\mu\mu)$
 - Amplitudes modified to incl. dependence on $m_{K\pi}$ to model P-Wave as relativistic Breit-Wigner and S-Wave with LASS
 - $m(K\pi\mu\mu)$ modelled with double Crystal Ball
- BF measured relative to $B^0 \rightarrow K^+ \pi^- J/\psi$



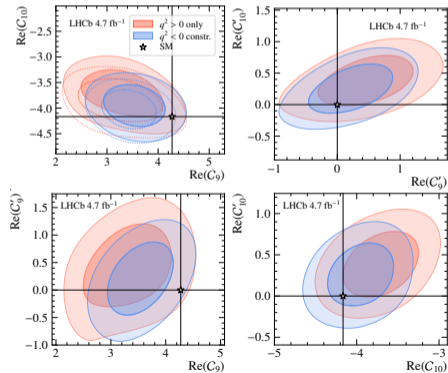
$q^2 < 0$ theory input
 [Gubernari, Reboud, Dyk & Virto]

Two versions of fit:
Without and **with** theory input on non-local component from $q^2 < 0$



Good agreement seen between both versions,
 without and with $q^2 < 0$ theory input

Deviation from SM	ΔC_9	ΔC_{10}	$\Delta C'_9$	$\Delta C'_{10}$
$q^2 > 0$ only	1.9σ	1.5σ	0.9σ	1.5σ
$q^2 < 0$ prior	1.8σ	0.9σ	0.5σ	1.0σ



Non-local effects described as corrections to C_9 .

[Cornella,Isidori,König,Liechti,Owen,Serra]

$$C_9 \rightarrow C_9^{\text{eff},\lambda} = C_9 + Y_{c\bar{c}}^{(0),\lambda} + \Delta Y_{light}^{1P,\lambda}(q^2) + \Delta Y_{c\bar{c}}^{1P,\lambda}(q^2) + \Delta Y_{c\bar{c}}^{2P,\lambda}(q^2) + Y_{\tau\tau}(q^2)$$

$$C_7^{\text{eff},\lambda} = C_7 + \epsilon^\lambda e^{i\omega^0}$$

Non-local effects described as corrections to C_9 .

[Cornella,Isidori,König,Liechti,Owen,Serra]

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$$C_7^{\text{eff},\lambda} = C_7 + \epsilon^\lambda e^{i\omega^0}$$

Local Contributions

Wilson Coefficients

Real part of C_9 , C_9' , C_{10} and C_{10}' treated as free fit parameters.

Imaginary part set to 0, implicitly assumes no CPV in $B \rightarrow K^* \mu^+ \mu^-$ decays (left for a future measurement)

Polarisation dependent shift

Included as a fit parameter to allow for helicity dependent complex phase

Non-local effects described as corrections to C_9 .

[Cornella, Isidori, König, Liechti, Owen, Serra]

$$C_9 \rightarrow C_9^{\text{eff},\lambda} = C_9 + Y_{c\bar{c}}^{(0),\lambda} + \Delta Y_{\text{light}}^{1P,\lambda}(q^2) + \Delta Y_{c\bar{c}}^{1P,\lambda}(q^2) + \Delta Y_{c\bar{c}}^{2P,\lambda}(q^2) + Y_{\tau\tau}(q^2)$$

$$C_7^{\text{eff},\lambda} = C_7 + \epsilon^\lambda e^{i\omega^0}$$

Non-Local Contributions

Subtraction term

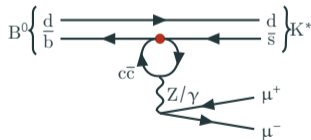
Theoretically calculated at negative q^2

[Asatrian, Greub, Vito]

Introduced to ensure convergence of dispersion relation.

Negligible impact for light resonances

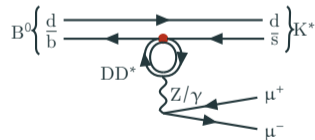
1-particle contributions



Includes vector resonances:
 light $\rightarrow \rho(770), \omega(782), \phi(1020)$,
 $c\bar{c} \rightarrow J/\psi, \psi(2S), \psi(3770), \psi(4040), \psi(4160)$

Mag (except J/ψ) and phase are fit parameters

2-particle contributions



Includes non-resonant:
 DD, DD^*, D^*D^*
 Real and Imag. parts are fit parameters

Non-local effects described as corrections to C_9 .

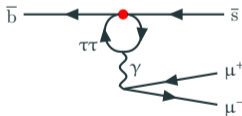
[Cornella,Isidori,König,Liechti,Owen,Serra]

$$C_9 \rightarrow C_9^{\text{eff},\lambda} = C_9 + Y_{c\bar{c}}^{(0),\lambda} + \Delta Y_{\text{light}}^{1P,\lambda}(q^2) + \Delta Y_{c\bar{c}}^{1P,\lambda}(q^2) + \Delta Y_{c\bar{c}}^{2P,\lambda}(q^2) + Y_{\tau\tau}(q^2)$$

$$C_7^{\text{eff},\lambda} = C_7 + \epsilon^\lambda e^{i\omega^0}$$

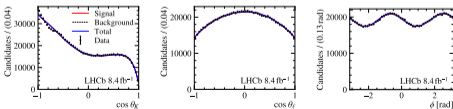
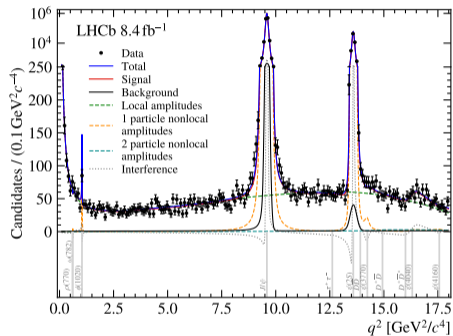
Non-Local Contributions

Tau-scattering contribution



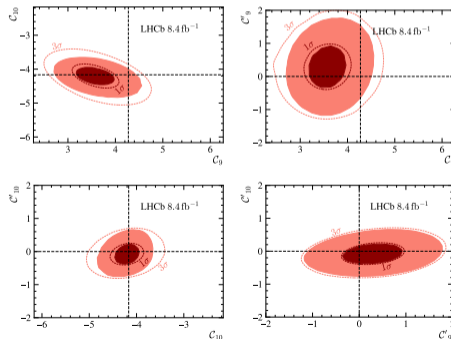
Wilson C_9^τ is fit parameter
gives indirect access to $\text{BF}(B^0 \rightarrow K^* \tau^+ \tau^-)$

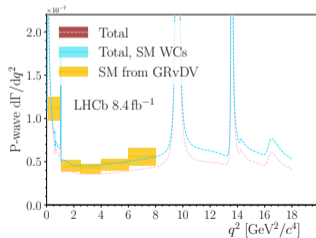
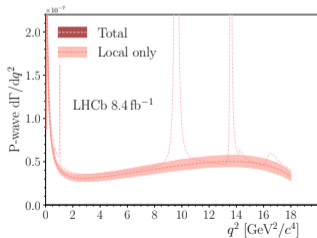
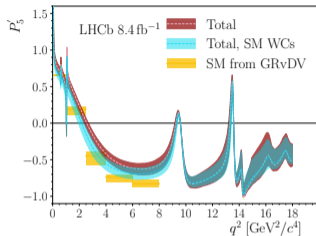
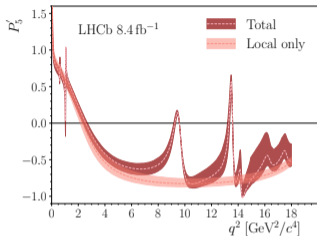
Unbinned fit to full q^2 using 8.4 fb^{-1} (2011+2012, 2016+2017+2018) LHCb data



	ΔC_9	ΔC_{10}	$\Delta C'_9$	$\Delta C'_{10}$	C_9^τ
Deviation from SM	2.1σ	0.6σ	0.7σ	0.4σ	0.4σ

$|C_9^\tau| < 500$ at 90% C.L., $C_9^{\tau, SM} \sim 4$
 Current best limits $|C_9^\tau| < 680$ (600) at 90% C.L. $C_{10}^\tau = SM$ ($C_{10}^\tau = -C_9^\tau$)
 [PRD 108, L01102 (2023)], flavio



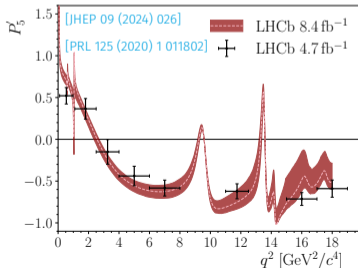
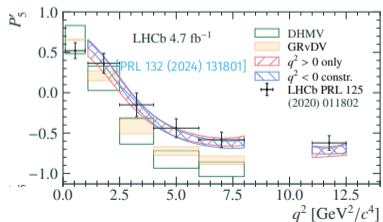
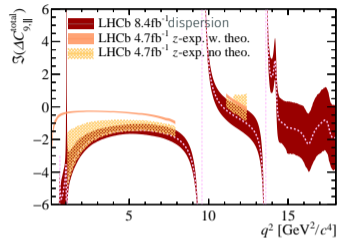
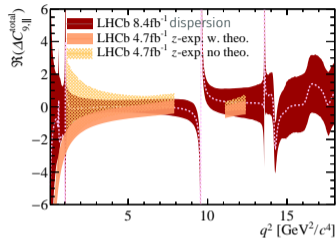


- Modification of differential observable from non-local state
- Cyan band on right calculated by fixing WC to SM values and non-local from data.
- Data seem to prefer non-local values larger than SM prediction
- Nevertheless, Wilson C_9 still prefers a shifted value w.r.t SM (indicated by the Total band)
- Missing other non-local components eg. $B \rightarrow D^* D_s^+ \rightarrow K^* \mu \mu$? [Ciuchini et al.]

A similar measurement with $B^+ \rightarrow K^+ \mu^+ \mu^-$ in progress

Compatibility of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ Analyses

- Polarisation dependent amplitude of non-local states (eg. \parallel)
- Good agreement seen between the two analyses.



- Comparison of P_5' with binned analysis (black points)
 - Good agreement between the different analyses
- Data consistently prefers a deviation from SM

Conclusion

- The study of FCNC B decays is an important research topic at LHCb.
- Amplitude and angular measurements of such decays offer a wide and complimentary information.
- Many measurements statistically dominated, however for our most sensitive probes the systematics are becoming more and more important.
- Currently some dominant systematics are external (normalisation mode and form factor).
- Theory inputs on local and non-local form factors crucial for measurements and interpretations.
[\[see talk by Arianna and Martin\]](#)

What can be done from experiment side to help extract C_9^{NP} accurately?

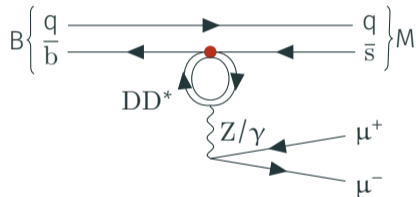
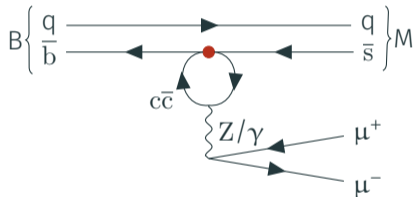
Interpretations of old measurements becomes difficult with updates to model
How could we work together to address this issue?

Thanks for your attention!

Backup

Non-local charm-loop effects

Experimentally can mimic NP and difficult to estimate theoretically



Can we measure these non-local effects from data?

$$\begin{aligned} \frac{1}{d(\Gamma + \bar{\Gamma})/dq^2} \frac{d^3(\Gamma + \bar{\Gamma})}{d \cos \theta_L d \cos \theta_K d \tilde{\varphi}} &= \frac{9}{32\pi} \left\{ \frac{3}{4} (1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K \right. \\ &+ \left[\frac{1}{4} (1 - F_L) \sin^2 \theta_K - F_L \cos^2 \theta_K \right] \cos 2\theta_L \\ &+ \frac{1}{2} (1 - F_L) A_T^{(2)} \sin^2 \theta_K \sin^2 \theta_L \cos 2\tilde{\varphi} \\ &+ (1 - F_L) A_T^{\mathcal{R}e\mathcal{CP}} \sin^2 \theta_K \cos \theta_L \\ &\left. + \frac{1}{2} (1 - F_L) A_T^{\mathcal{I}m\mathcal{CP}} \sin^2 \theta_K \sin^2 \theta_L \sin 2\tilde{\varphi} \right\}. \end{aligned}$$

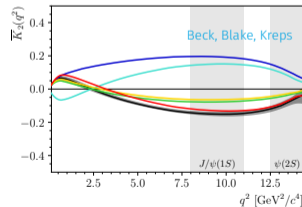
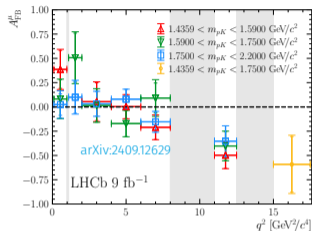
$$\lim_{q^2 \rightarrow 0} A_T^{(2)}(q^2) = \frac{2 \left[\mathcal{R}e[C_7] \mathcal{R}e[C_7'] + \mathcal{I}m[C_7] \mathcal{I}m[C_7'] + \frac{y}{2} [(\mathcal{R}e[C_7])^2 - (\mathcal{I}m[C_7])^2] \right]}{(\mathcal{R}e[C_7])^2 + (\mathcal{I}m[C_7])^2}, \quad (1)$$

$$\lim_{q^2 \rightarrow 0} A_T^{\mathcal{I}m\mathcal{CP}}(q^2) = \frac{2 \left[\mathcal{R}e[C_7] \mathcal{I}m[C_7'] - \mathcal{I}m[C_7] \mathcal{R}e[C_7'] - y \mathcal{R}e[C_7] \mathcal{I}m[C_7] \right]}{(\mathcal{R}e[C_7])^2 + (\mathcal{I}m[C_7])^2}, \quad (2)$$

Source of systematic	$A_T^{(2)}$	$A_T^{Im\mathcal{CP}}$	$A_T^{Re\mathcal{CP}}$	F_L
$\Delta\Gamma_s/\Gamma_s$	0.008	< 0.001	< 0.001	< 0.001
Corrections to simulation	0.002	<0.001	<0.001	0.010
Acceptance function modelling	<0.001	<0.001	0.001	0.002
Simulation sample size for acceptance	0.006	0.008	0.005	0.002
Background contamination	0.009	0.014	0.004	0.006
Angles resolution	-0.005	< 0.001	-	-
Total systematic uncertainty	0.014	0.016	0.006	0.012
Statistical uncertainty	0.235	0.247	0.155	+0.056

- Yields and branching fraction measured from $m_{pK^- \mu^+ \mu^-}$
- The sPlot technique is used derive weights to separate signal and background.
- With further efficiency corrections moments derived from weighted sum over basis functions.

$$\bar{K}_i = \frac{1}{N} \sum_{\text{event } n} w(\vec{\Phi}_n) f_i(\vec{\Omega}_n)$$

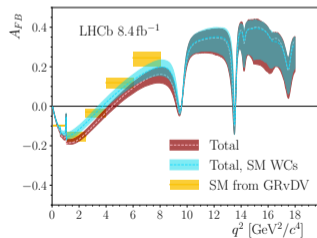
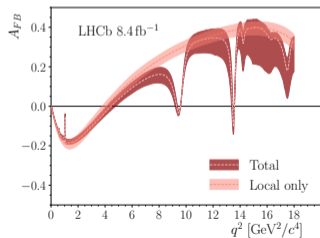
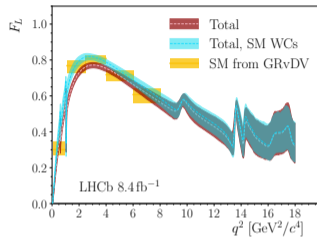
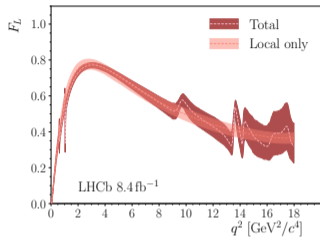


- Standard Model
- $C_9 = -C_9^{\text{SM}}$
- $C_{10} = -C_{10}^{\text{SM}}$
- $C_9^L = C_9^{\text{SM}}$
- $C_{10}^L = C_{10}^{\text{SM}}$
- global fit

$1.1 < q^2 < 6.0 \text{ GeV}^2/c^4$				$1.1 < q^2 < 6.0 \text{ GeV}^2/c^4$	
F_L	$0.58 \pm 0.05 \pm 0.05$	F_L	$0.58 \pm 0.04 \pm 0.05$	Q_{FL}	$0.12 \pm 0.05 \pm 0.05$
S_3	$-0.00 \pm 0.04 \pm 0.02$	P_1	$-0.00 \pm 0.20 \pm 0.24$	Q_1	$-0.10 \pm 0.26 \pm 0.24$
S_4	$-0.12 \pm 0.07 \pm 0.04$	P'_4	$-0.24 \pm 0.15 \pm 0.12$	Q_4	$-0.07 \pm 0.17 \pm 0.12$
S_5	$-0.08 \pm 0.05 \pm 0.03$	P'_5	$-0.16 \pm 0.11 \pm 0.10$	Q_5	$0.05 \pm 0.13 \pm 0.10$
A_{FB}	$-0.15 \pm 0.05 \pm 0.04$	P_2	$-0.23 \pm 0.08 \pm 0.11$	Q_2	$0.07 \pm 0.10 \pm 0.11$
S_7	$-0.08 \pm 0.06 \pm 0.04$	P'_6	$-0.16 \pm 0.11 \pm 0.09$	Q_6	$-0.05 \pm 0.14 \pm 0.09$
S_8	$0.13 \pm 0.07 \pm 0.06$	P'_8	$0.26 \pm 0.15 \pm 0.14$	Q_8	$-0.28 \pm 0.17 \pm 0.14$
S_9	$0.07 \pm 0.05 \pm 0.02$	P_3	$-0.16 \pm 0.11 \pm 0.11$	Q_3	$0.22 \pm 0.14 \pm 0.11$

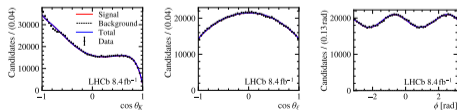
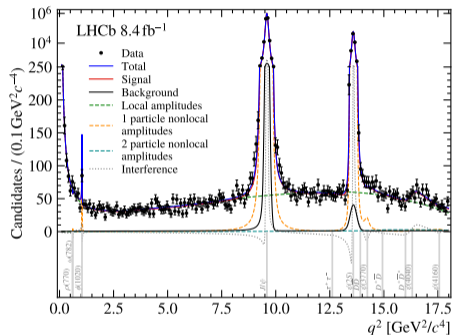
Preliminary results. First error statistical

	Systematic uncertainties			
	\mathcal{C}_9	\mathcal{C}_{10}	\mathcal{C}'_9	\mathcal{C}'_{10}
Amplitude model				
S-wave form factors	< 0.01	< 0.01	< 0.01	< 0.01
S-wave non-local hadronic	0.02	0.02	0.14	0.04
S-wave $m_{K\pi}$				
[i] LASS parameters	< 0.01	< 0.01	< 0.01	< 0.01
[ii] $m_{K\pi}$ model	< 0.01	< 0.01	0.05	0.03
Exotic states on hadronic inputs	0.01	< 0.01	0.02	< 0.01
Total	0.02	0.02	0.15	0.05
External inputs on BR				
$\mathcal{B}(B \rightarrow J/\psi K^+ \pi^-)$	0.05	0.08	0.02	0.01
Fraction of $B \rightarrow J/\psi K^+ \pi^-$ in $m_{K\pi} \in 100$ MeV	0.03	0.03	0.01	< 0.01
Others	0.03	0.04	0.03	0.01
Total	0.07	0.09	0.04	0.01
Background model				
Chebyshev polynomial order	0.01	0.01	0.01	< 0.01
Combinatorial shape on $m_{K\pi}$	0.02	< 0.01	0.02	< 0.01
Background factorisation	0.01	0.01	0.01	0.01
Peaking background	0.01	< 0.01	0.02	0.01
Total	0.03	0.02	0.03	0.01
Experimental effects				
Acceptance parametrisation	< 0.01	< 0.01	< 0.01	< 0.01
Statistical uncertainty on acceptance	0.02	< 0.01	0.02	< 0.01
Total	0.02	< 0.01	0.02	< 0.01
Total systematic uncertainty	0.08	0.10	0.16	0.05

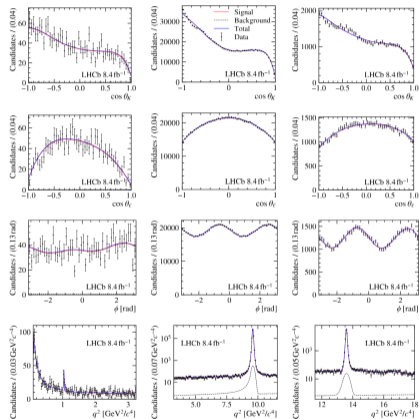


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- Cyan band on right calculated by fixing WC to SM values and non-local from data.
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- Nevertheless, Wilson C_9 still prefers a shifted value w.r.t SM (indicated by the Total band)

Unbinned fit to full q^2 using 8.4 fb^{-1} (2011+2012, 2016+2017+2018) LHCb data

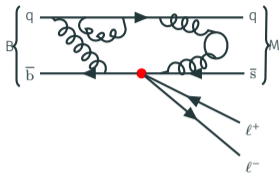


- Total 150 fit parameters
- From simulation:
 - Acceptance model
- From data:
 - Resolution
 - S-Wave Parameters
 - Background model
 - Non-local states
- From theory
 - Local $B \rightarrow K^*$ form factors gaussian constrained [Gubernari, Reboud, van Dyk & Virto]
 - Subtraction point for $c\bar{c}$ [Asatrian, Greub & Virto]



Variable	Branch.	Exotic		Model		
	Frac.	$\psi(2S)$	J/ψ	$D\bar{D}$	$\psi(4160)$	$\psi(3770)$
$ C_g $	45.37%	8.45%	10.71%	24.53%	2.55%	3.77%
$ C_{10} $	86.20%	3.78%	9.63%	4.46%	0.99%	1.51%
$Re(C'_g)$	3.85%	13.32%	20.32%	7.10%	5.07%	6.39%
$Re(C'_{10})$	6.81%	17.63%	22.11%	7.65%	0.80%	2.29%
$Re(C'_9)$	4.93%	15.70%	12.29%	30.56%	1.50%	9.42%

Local Hadronic effects



The form factors are difficult to calculate and leading source of theoretical uncertainty.

Two complimentary approaches,

- Lattice QCD in high q^2 region (low hadronic recoil) [HPQCD], [FNAL/MILC]
- Light Cone Sum Rules in low or -ve q^2 region [Gubernari et al.], [Khodjamirian and Rusov], [Bharucha, Bharucha and Zwicky], [Khodjamirian et al.]

Interpolate to full q^2

